

Status of the intermediate-mass region: First steps into *terra incognita*

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Nikhef

NWO
Netherlands Organisation
for Scientific Research



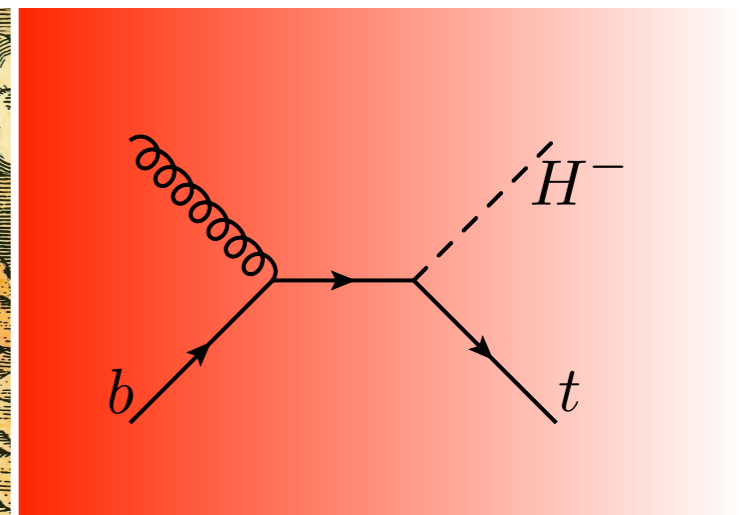
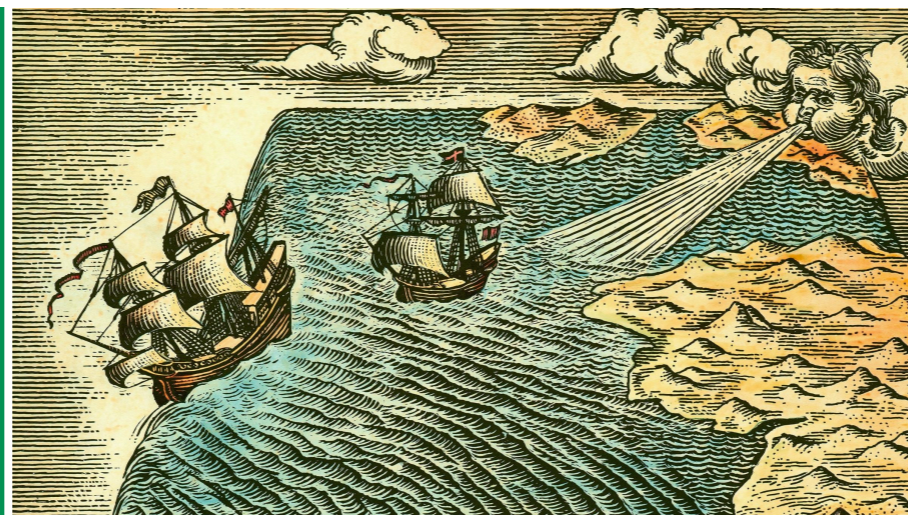
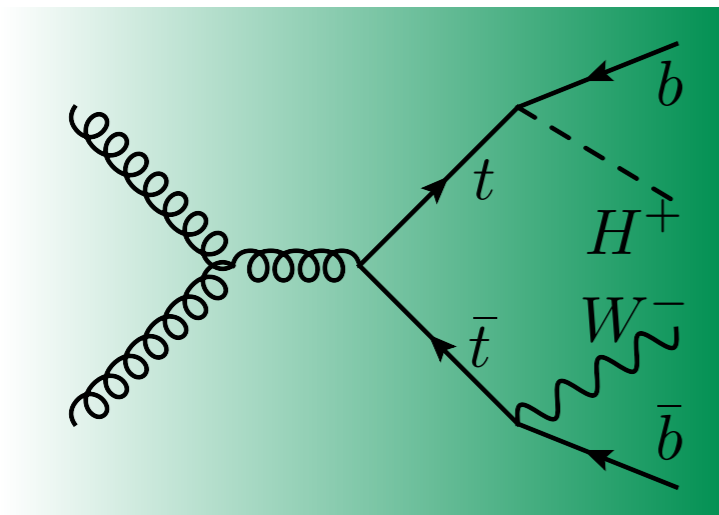
Charged Higgs production in 2HDMs

- In the 2HDM, the dominant production channel depends on the Charged Higgs mass

Light Higgs

Intermediate region

Heavy Higgs



$$M_{H^\pm} < M_t$$

$$M_{H^\pm} \simeq M_t$$

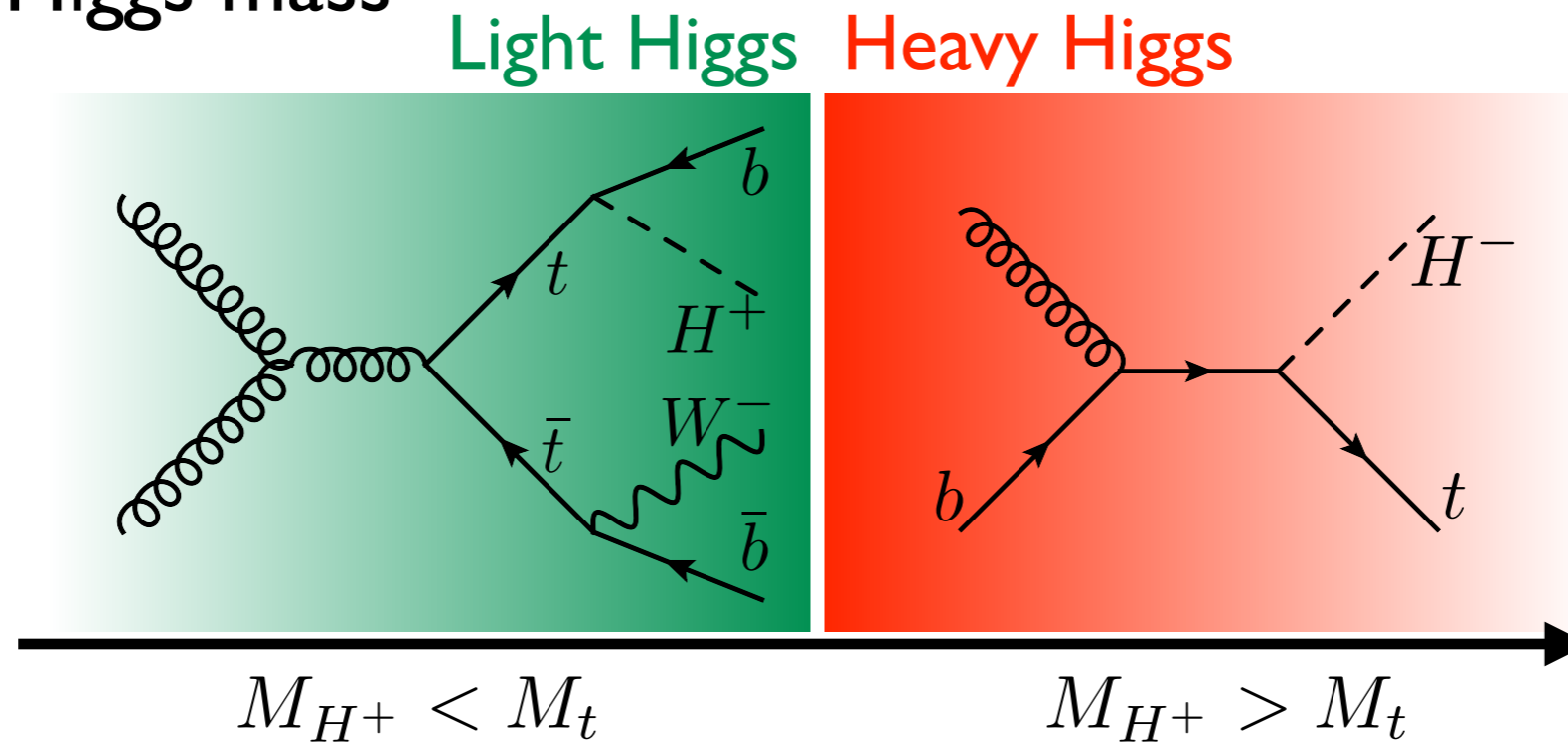
$$M_{H^\pm} > M_t$$

H^\pm mostly produced in $t\bar{t}$ events. The full $pp \rightarrow H^\pm W^\mp b\bar{b}$ process has to be simulated.

H^\pm mostly produced in association with a top quark

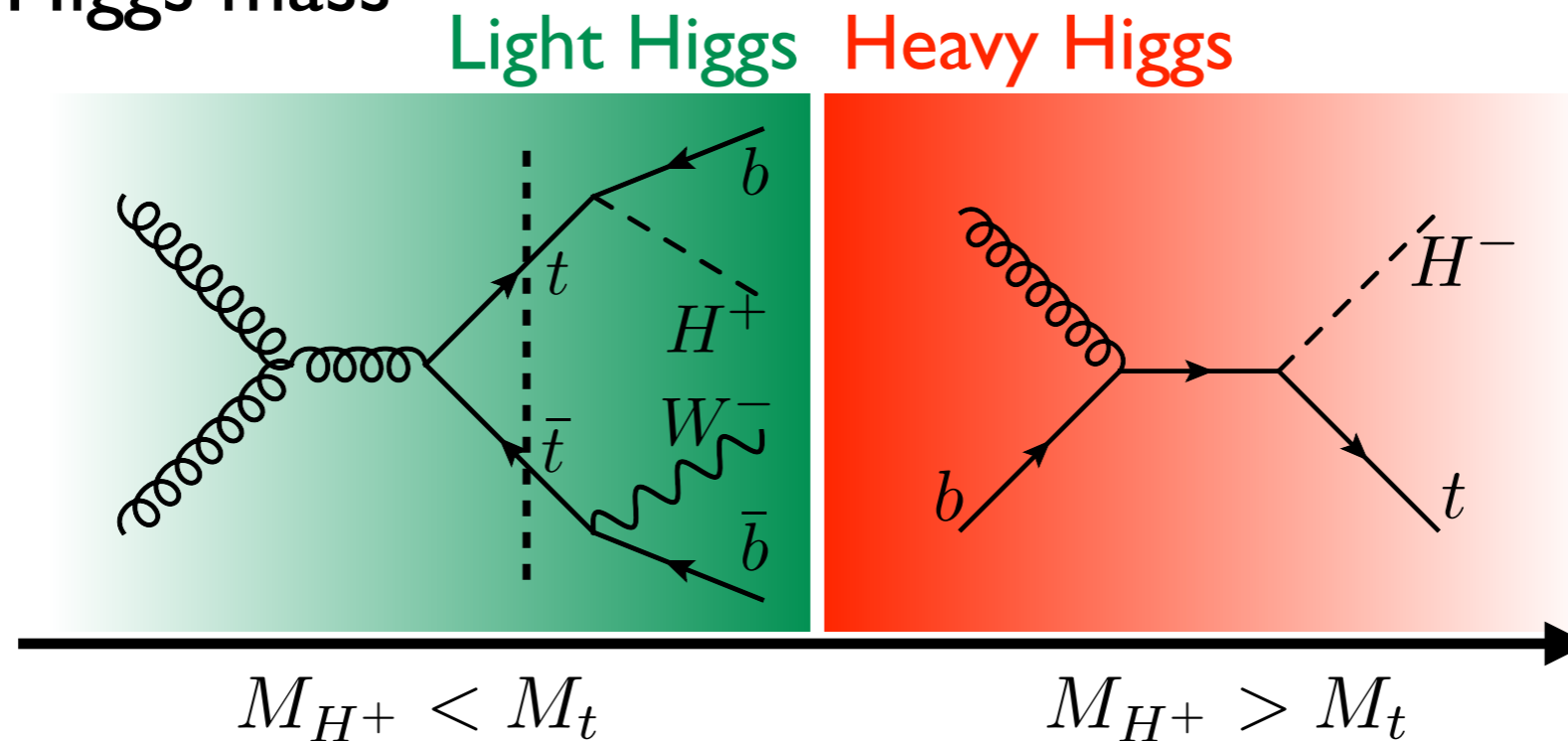
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Can take the $\Gamma_t \rightarrow 0$ limit

Charged Higgs production in 2HDMs

- In the 2HDM, the dominant production channel depends on the Charged Higgs mass

Light Higgs Heavy Higgs

Top cross-section at NNLO QCD: Czakon et al, arXiv:1303.6254, arXiv:1511.00549, ...
 EW corrections to top cross-section: Beenakker et al., Nu.Ph.B.411 (1994), Hollik et al., arXiv:0708.1697, ...
 NNLO QCD corrections to $\Gamma(t \rightarrow W^+b)$: Czarnecki et al, hep-ph/9806244, ...
 NLO (SUSY-)QCD corrections to $\Gamma(t \rightarrow H^+b)$: Reid et al, Z.Phys.C (1990), Li et al, Phys.Rev.D (1990), Czarnecki et al., Phys.Rev.D (1993), ..., Heynemeyer et al., hep-ph/9812320
 ...

NLO (SUSY-)QCD corrections: Zhu, hep-ph/0112109, Plehn, hep-ph/0206121, Berger et al, hep-ph/0312286 (5FS); Dittmaier et al, arXiv:0906.2648 (4FS); Flechl et al, arXiv:1409.5615 (Santander-Martched)
 EW corrections: Beccaria et al, arXiv:0908.1332 (5FS); Nhung et al, arXiv:1210.4087 (4FS)
 Threshold resummation: Kidonakis, arXiv:1005.4451 (5FS)
 Fully differential NLO+PS: Weydert et al, arXiv:0912.3430, Klasen et al, arXiv:1203.1341 (5FS); Degrande et al, arXiv:1507.02549 (4FS)

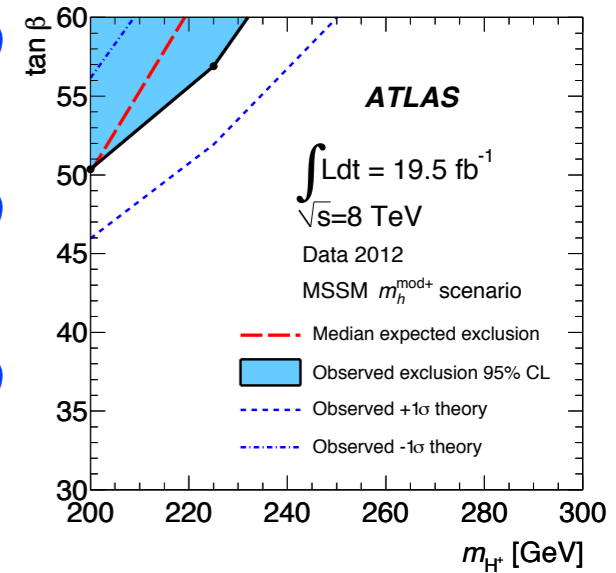
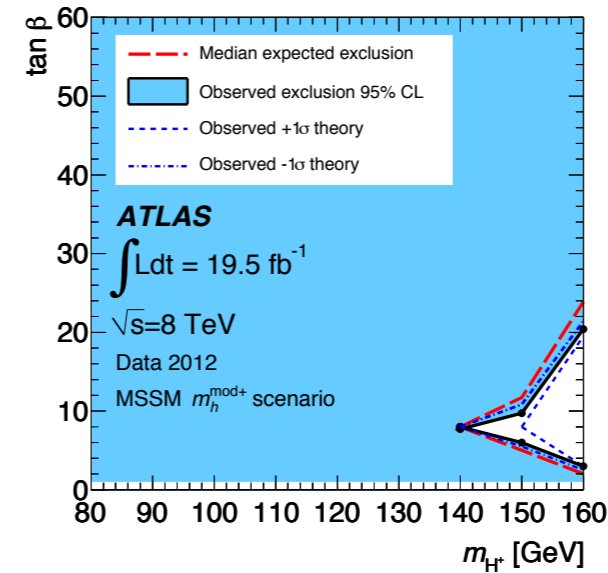
Can take the $\Gamma_t \rightarrow 0$ limit

Both cases are known at or beyond NLO QCD

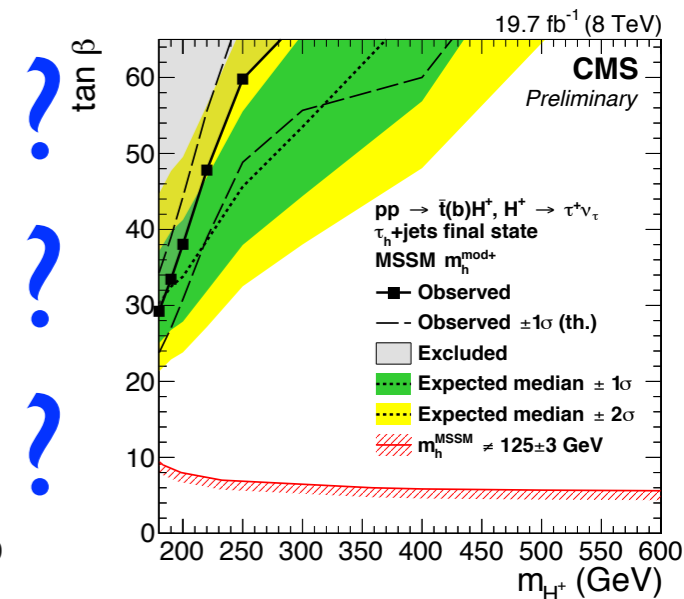
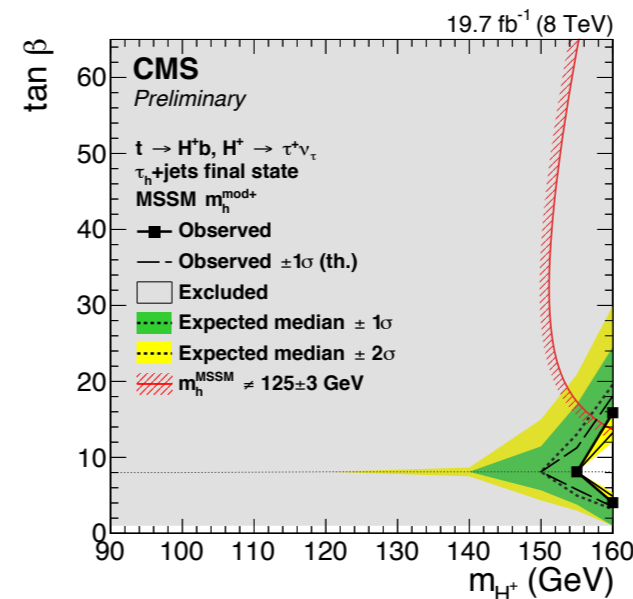
How we did searches at the LHC

- LHC experiments tend to exclude a light charged Higgs
- For a heavy charged Higgs, only very large values of $\tan\beta$ are excluded
- Missing mass window due to non-existence of predictions for the intermediate range beyond LO

ATLAS, arXiv:1412.6663

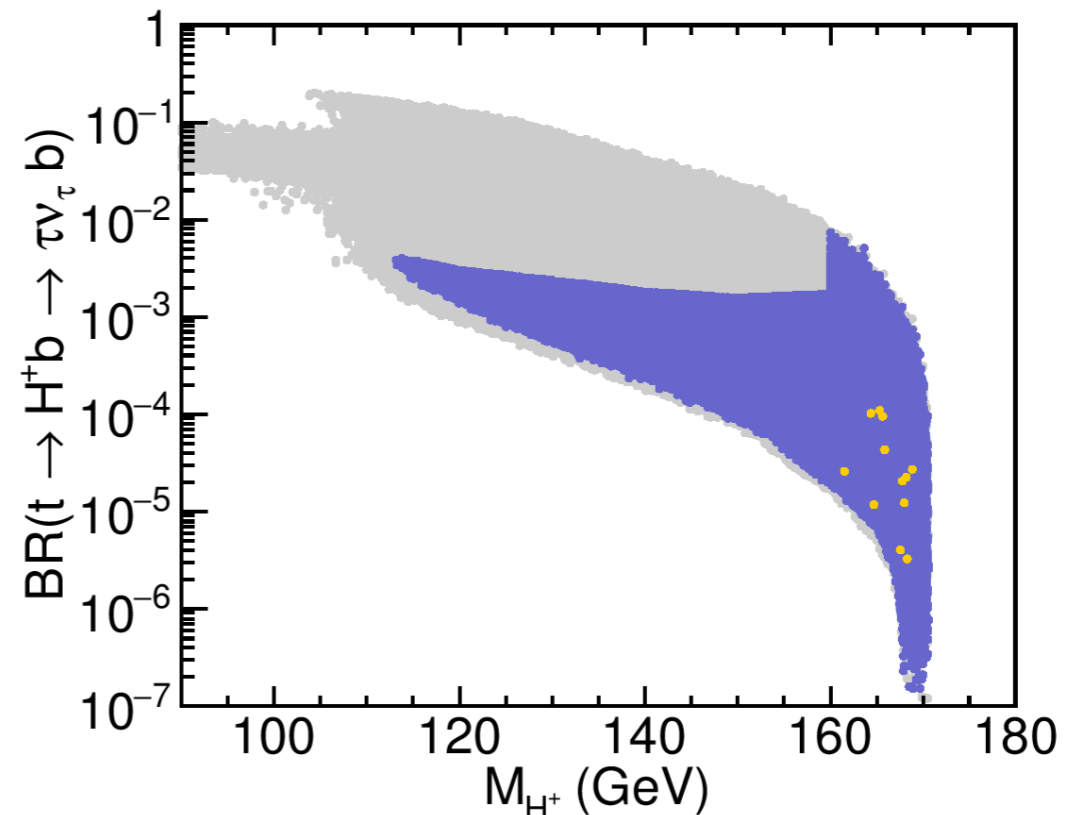
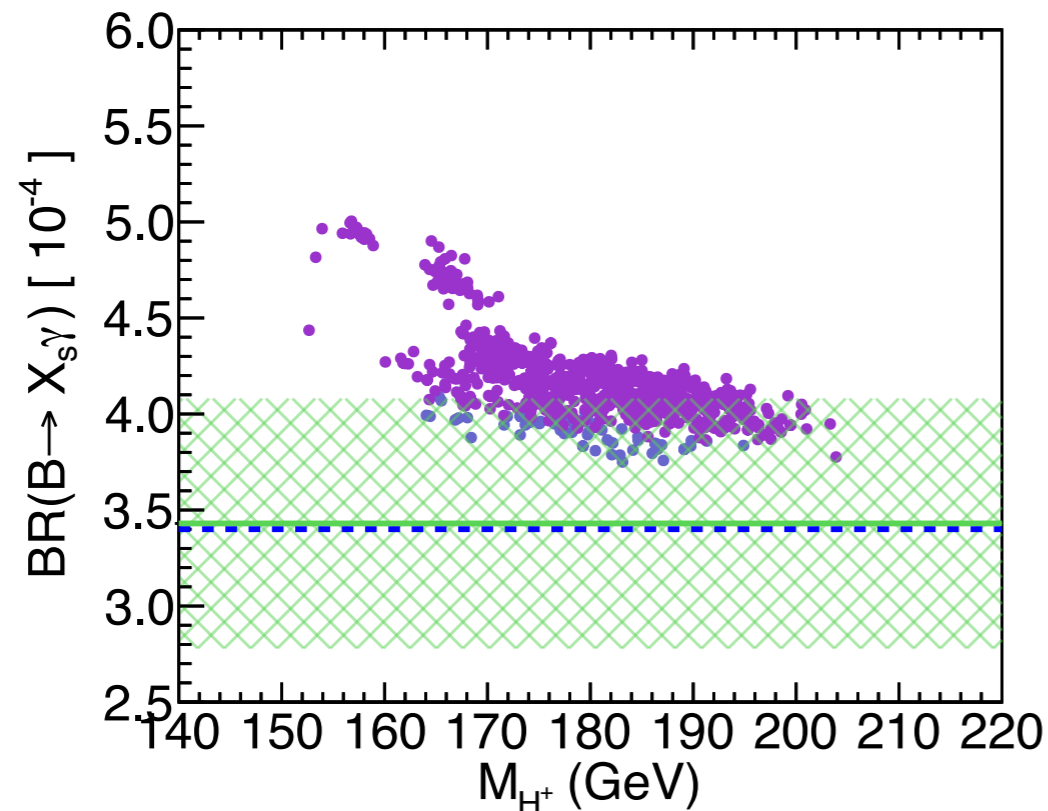


CMS, PAS HIG-14-020



Why to care about the intermediate-mass region?

- Because it is there
- SUSY models where the 125 GeV Higgs is the heavy H boson can have the charged Higgs in the intermediate-mass region
[Bechtle, Haber, Heinemeier, Stal, Stefaniak, Weiglein, Zeune, arXiv:1608.00638](#)
- In some of these models the light Higgs can act as a mediator to DM
[Profumo, Stefaniak, arXiv:1608.06945](#)



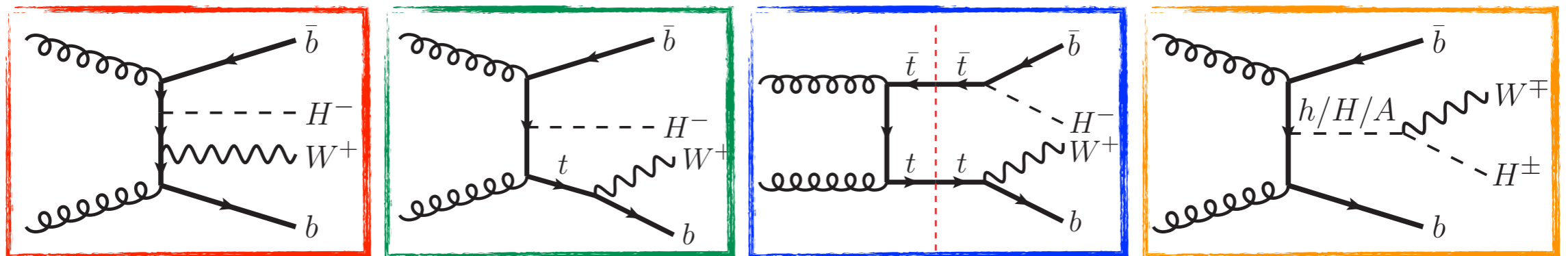
Getting ready to sail



Getting accurate predictions for the intermediate-mass region

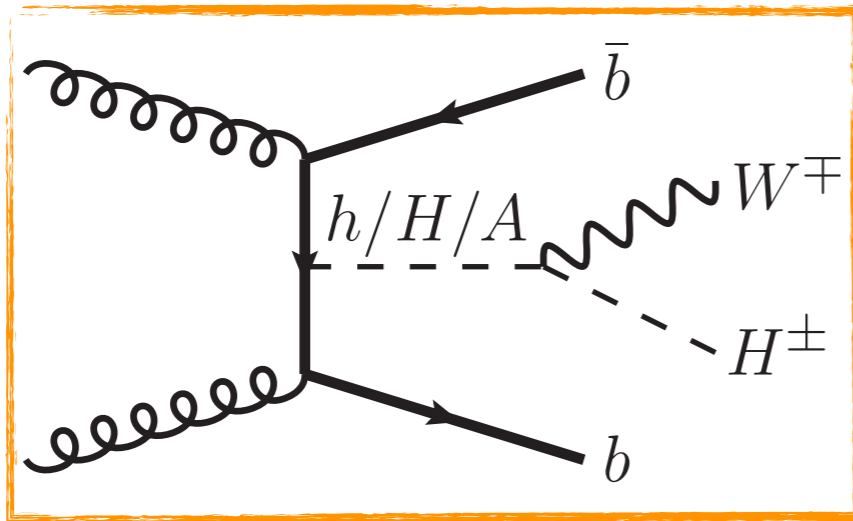
Degrande, Frederix, Hirschi, Ubiali, Wieseemann, MZ, arXiv:1607.05291

- The full process $pp \rightarrow H^\pm W^\mp b\bar{b}$ has to be simulated, consistently including the top quark width. $\Gamma_t = \Gamma_t(m_{H^\pm}, \tan\beta)$
- Diagrams with **0**, **1** and **2** resonant tops contribute to the total cross-section, as well as diagrams with **neutral Higgs bosons**

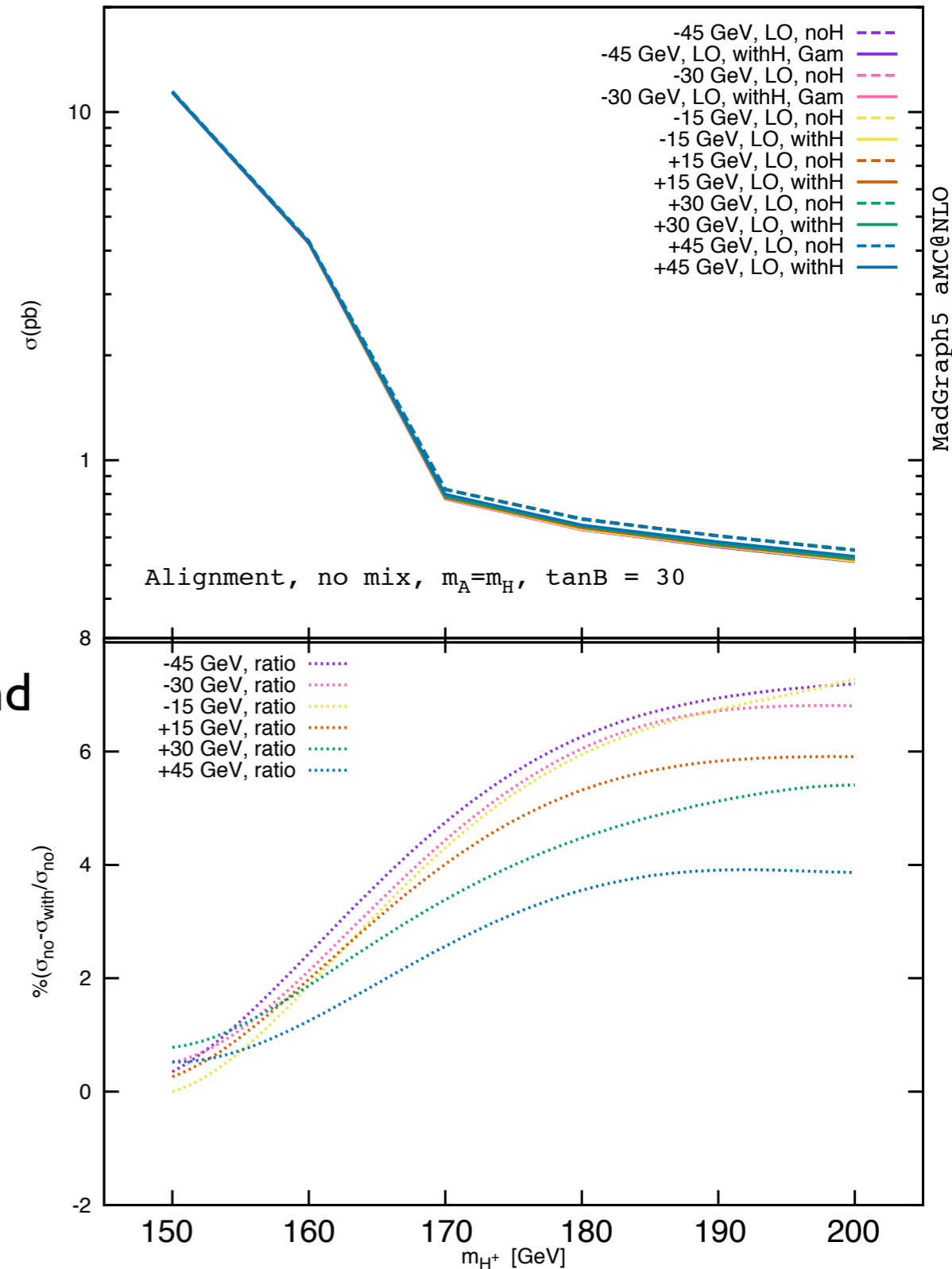


- Cross-section for $m_{H^\pm} > m_t$ ($m_{H^\pm} < m_t$) will get the dominant contribution from **single-** (**double**)-resonant diagrams
- LO total cross section has large (30-50%) theoretical errors. For accurate predictions one needs to compute NLO corrections

Effect of neutral Higgs bosons

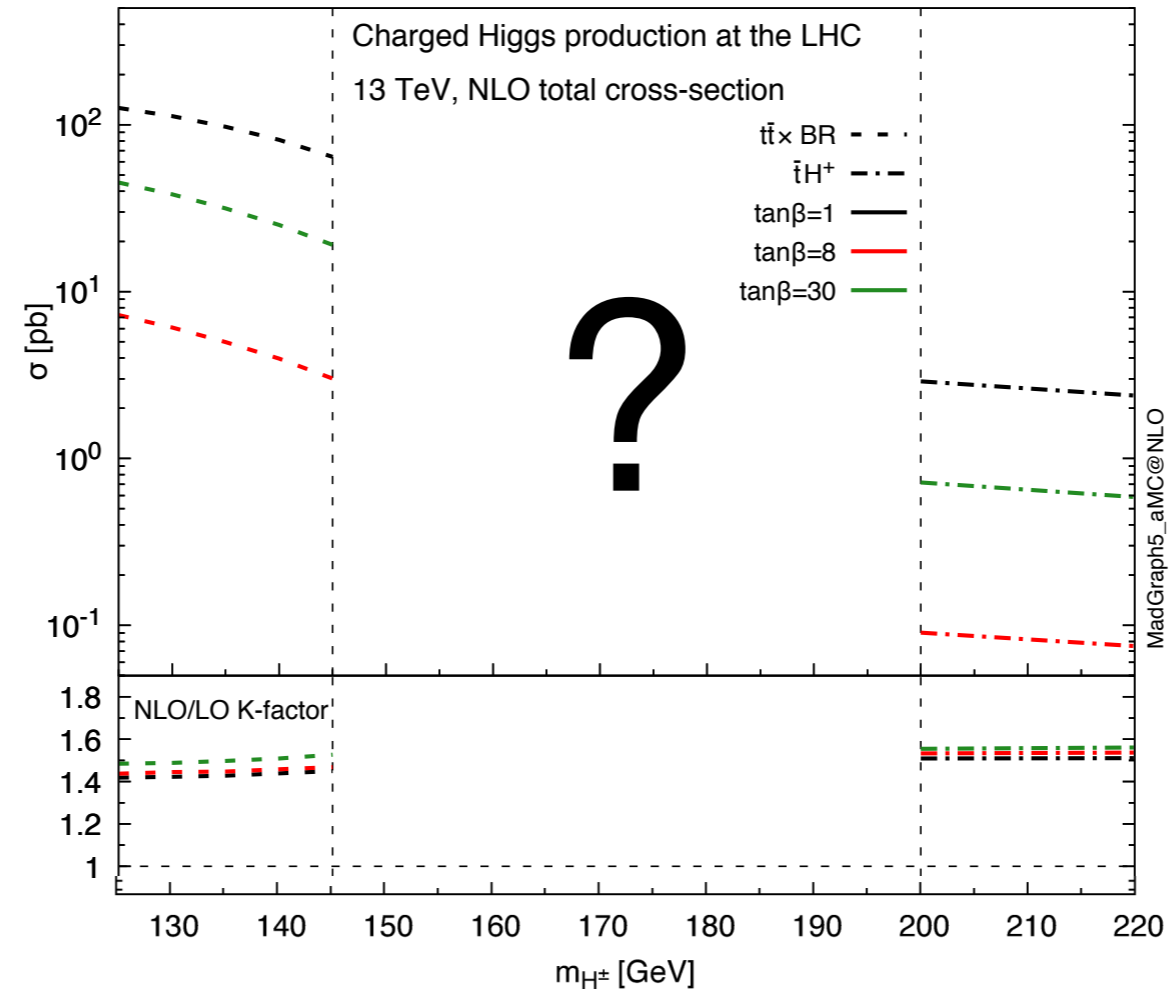


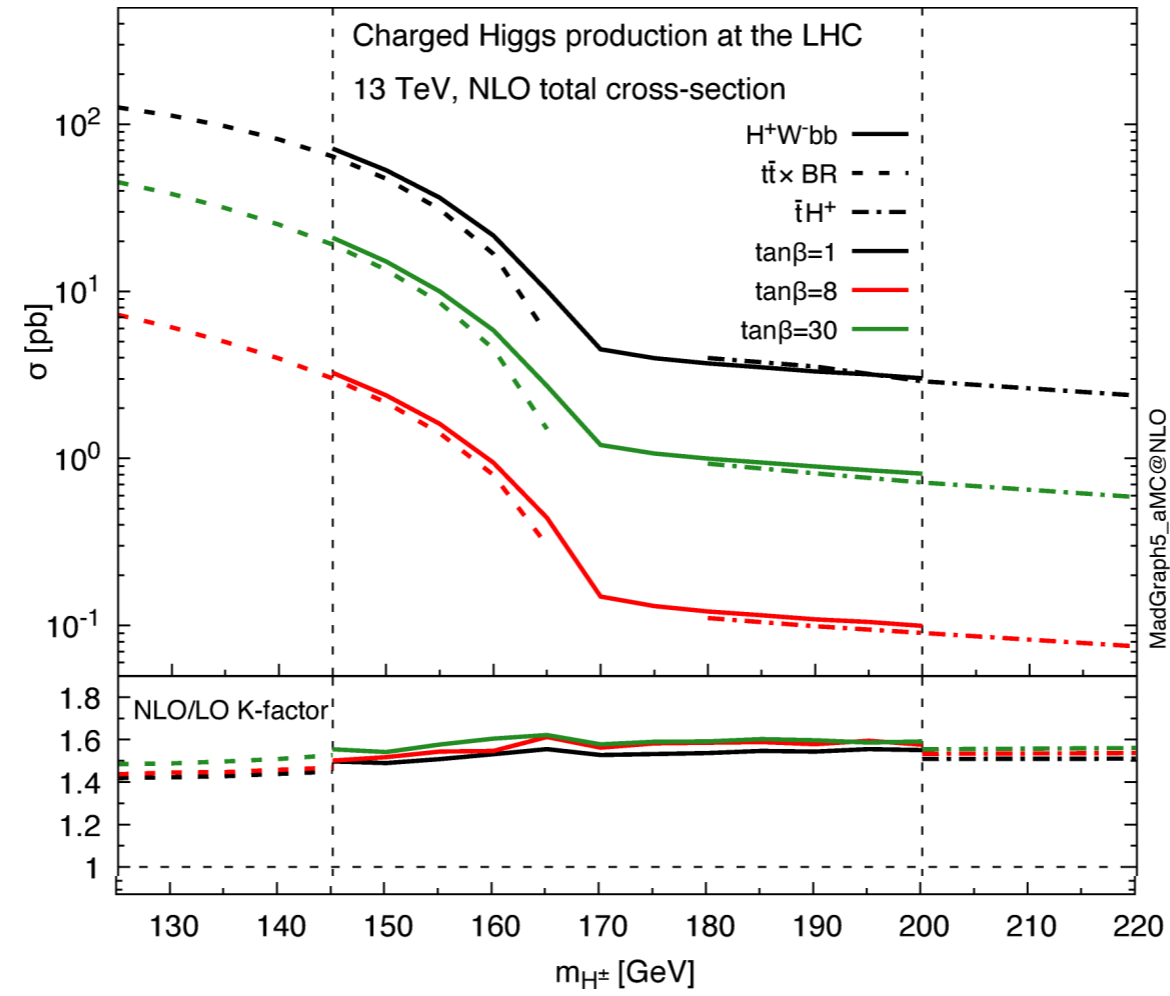
- Diagrams with neutral Higgs bosons introduce additional dependence on the $h/H/A$ masses and on the neutral Higgs mixing α .
- Assuming h to be the SM Higgs ($m_h=125$ GeV and $\cos(\beta-\alpha)\approx 0$), for non-resonant configurations ($m_{h/H/A} < m_W + m_{H^\pm}$) the contribution to the total cross section is small ($\approx 7\%$)
- In practice, these diagrams will be neglected
- Cross section will just depend on m_{H^\pm} and $\tan\beta$ (same as the heavy/light case)



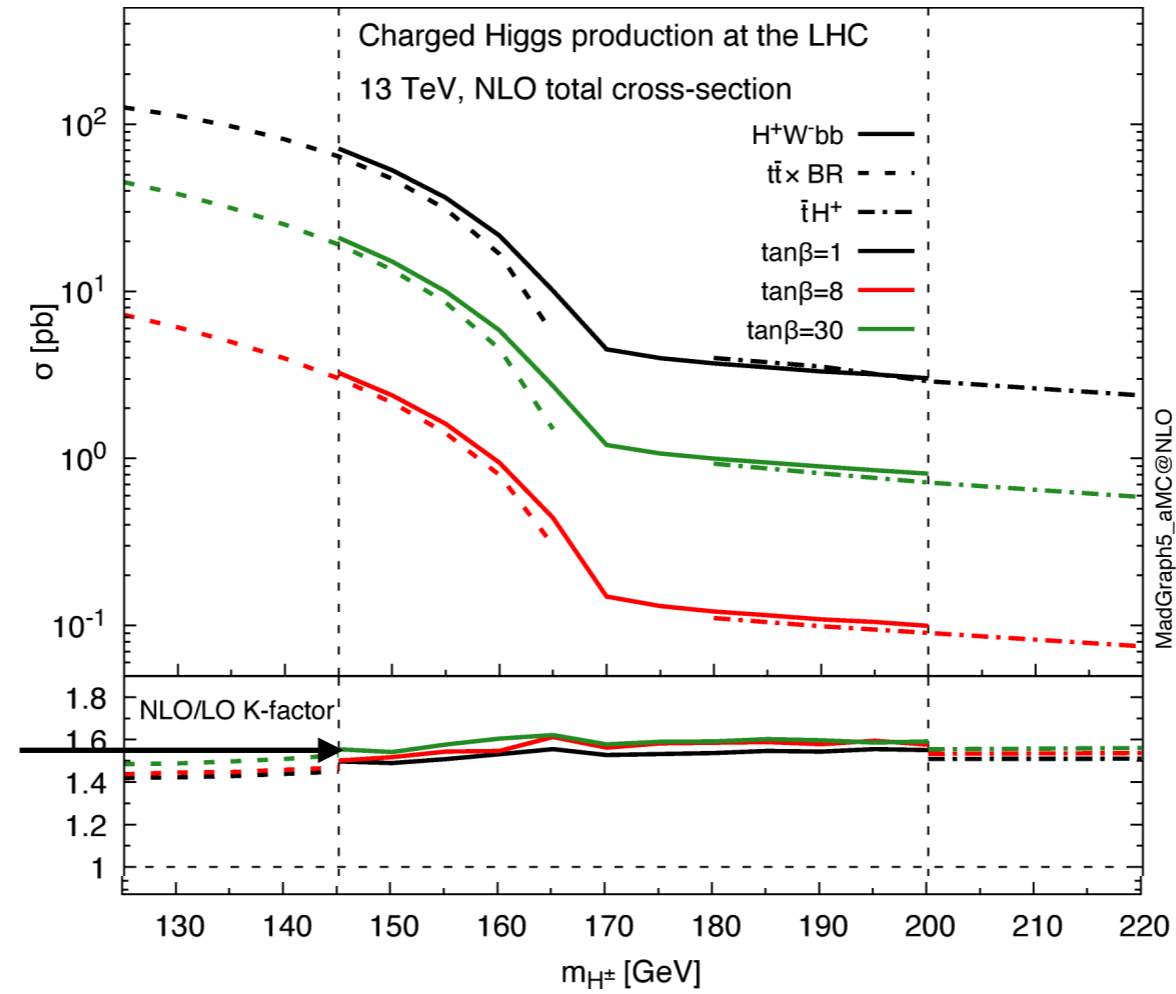
Calculation setup

- Computation carried out with `MADGRAPH5_AMC@NLO`, improved with resonance-aware FKS subtraction [Frederix et al. arXiv:1603.01178](#)
- Focus on type-II 2HDM (extension to MSSM and other 2HDMs will be discussed)
- Use massive bottom quarks (4FS). Use `PDF4LHC_nlo_nf4` PDFs
- Complex top-mass (and Yukawa) scheme to include the top width in a gauge-invariant way. Γ_t computed at NLO for every $(m_{H^\pm}, \tan\beta)$ point
- Use a fixed central scale, $\mu_{R/F}=125$ GeV
 - Matches scales used in the light- and heavy-Higgs regions
 - Scale uncertainties obtained by varying independently up/down of a factor 2
- Use the $\overline{\text{MS}}$ scheme for y_b renormalisation (introduces extra μ_R dependence)
- Scan $145 \text{ GeV} < m_{H^\pm} < 200 \text{ GeV}$
- Three values of $\tan\beta$ will be considered ($\tan\beta = 1, 8, 30$)
- Other input parameters follow the recommendation of the LHC HXSWG



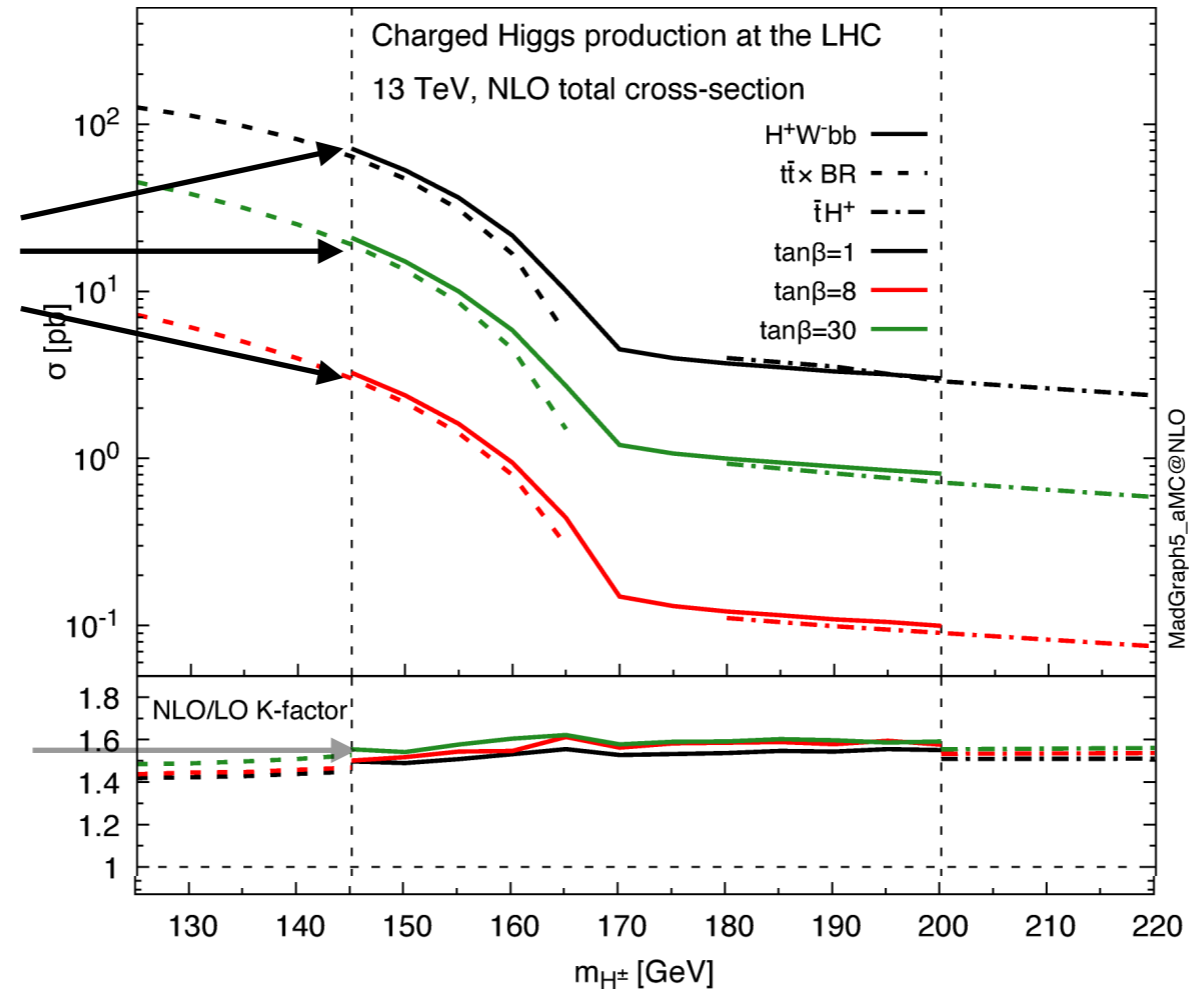


Rather constant K-factor $\sim 1.5-1.6$, with small $\tan\beta$ dependence



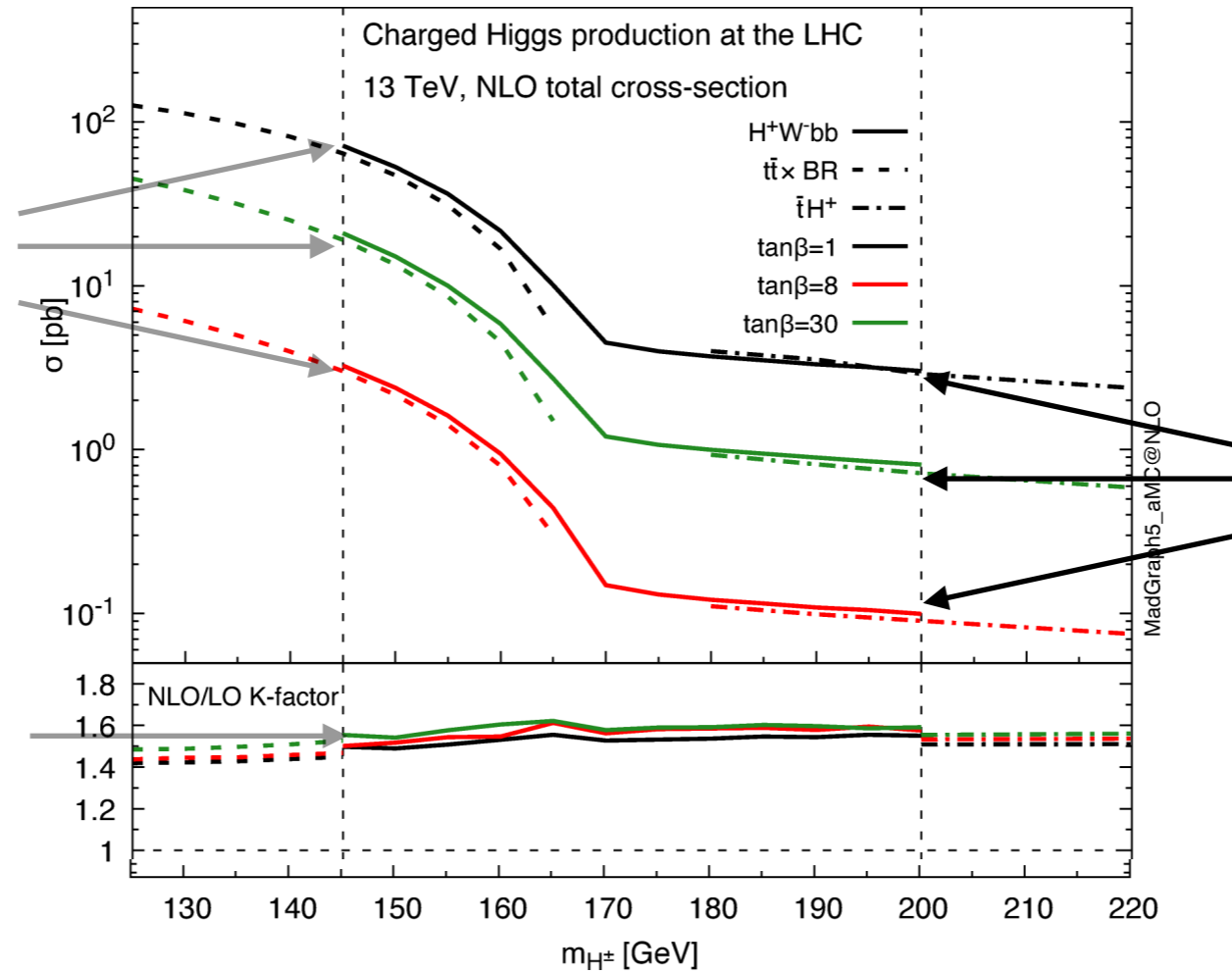
Discontinuity due to single-resonant contributions (tW^\pm and tH^\pm)

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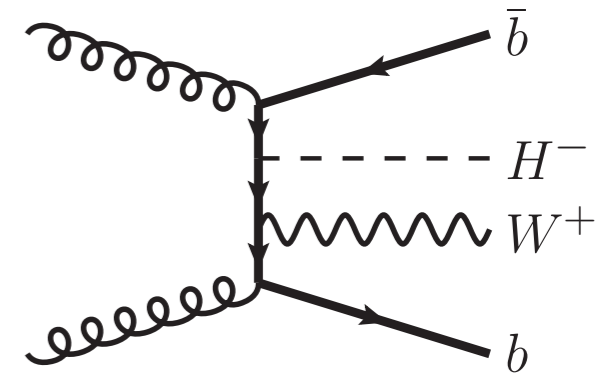


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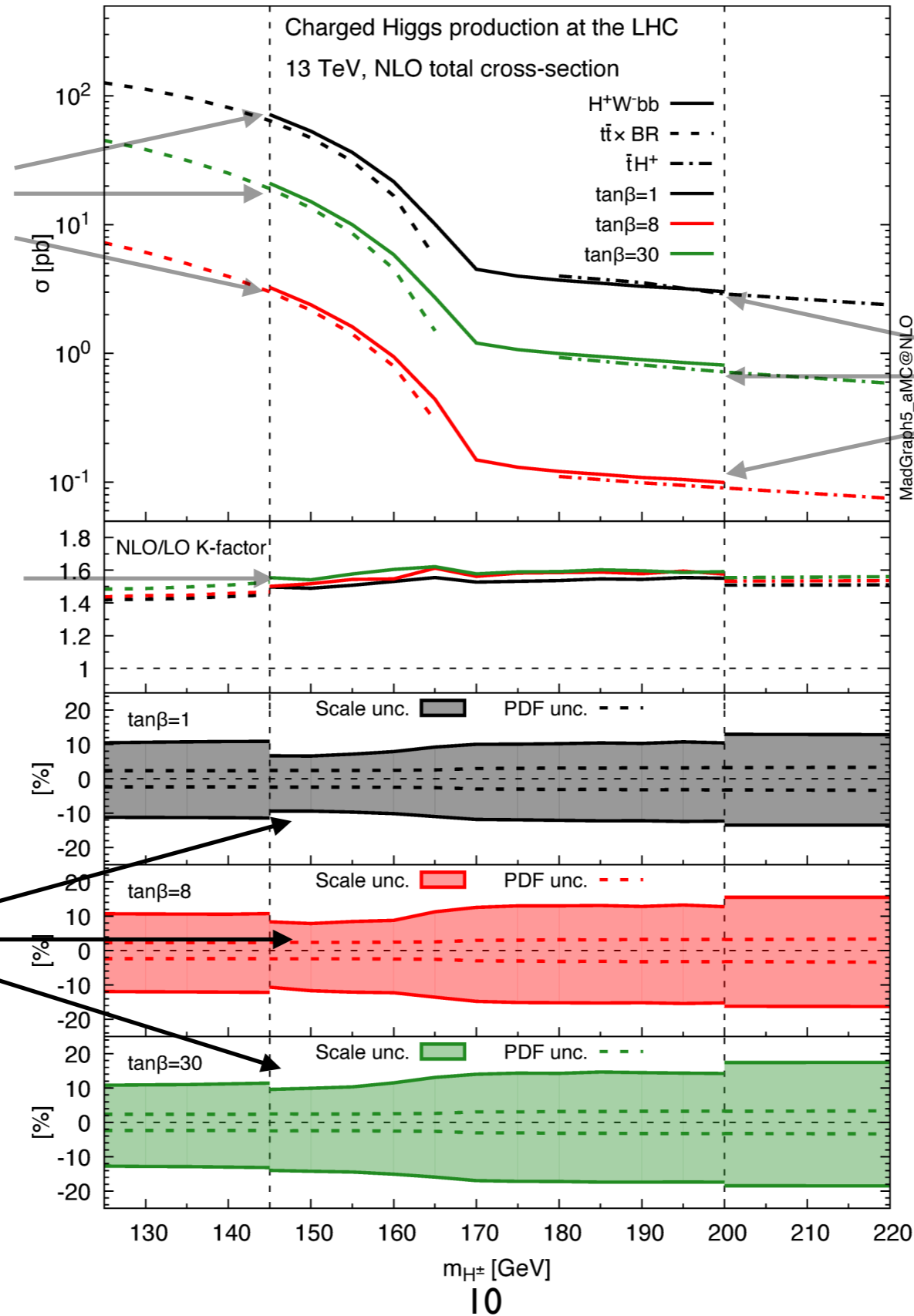
Discontinuity due to non-resonant contributions. Size of discontinuity is $\tan\beta$ dependent because of chiral couplings



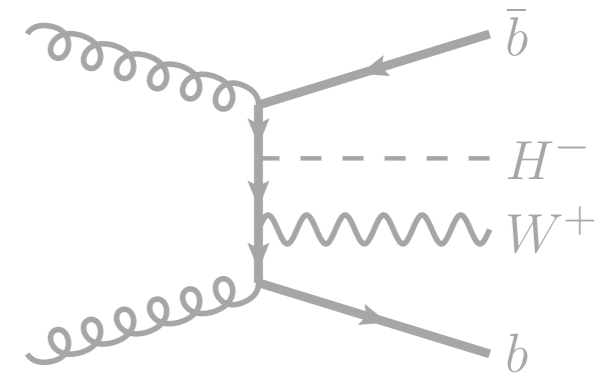
Discontinuity due to single-resonant contributions (tW^\pm and tH^\pm)

Rather constant K-factor $\sim 1.5-1.6$, with small $\tan\beta$ dependence

Scale uncertainties reduced to 10-20% (larger for large $\tan\beta$ because of extra dependence in y_b)



Discontinuity due to non-resonant contributions. Size of discontinuity is $\tan\beta$ dependent because of chiral couplings



Extension to other values of $\tan\beta$ or different 2HDMs

- The charged Higgs cross section receives contributions proportional to y_b^2 ($\sim \tan\beta^2$), y_t^2 ($\sim 1/\tan\beta^2$) and $y_b y_t$ (constant).
- **MADGRAPH5_AMC@NLO** has been extended in order to return the three individual contributions

- The cross section at any value of $\tan\beta$ can be computed as

$$\sigma(\tan\beta') = \left[\left(\frac{\tan\beta'}{\tan\beta} \right)^2 \sigma_{y_b^2}(\tan\beta) + \sigma_{y_b y_t}(\tan\beta) + \left(\frac{\tan\beta}{\tan\beta'} \right)^2 \sigma_{y_t^2}(\tan\beta) \right] \times \left(\frac{\Gamma_t(\tan\beta)}{\Gamma_t(\tan\beta')} \right)^2$$

- Cross-checked by recomputing our results at $\tan\beta=1$ and $\tan\beta=30$, starting from $\tan\beta=8$. Agreement below 1% was found
- Can be extended to obtain the charged Higgs cross section in other 2HDM scenarios (e.g. type I or MSSM with Δ_b corrections)

$$\sigma^{t-I}(\tan\beta') = \frac{\sigma^{t-I}(\tan\beta = 1)}{(\tan\beta')^2} \times \left(\frac{\Gamma_t(\tan\beta)}{\Gamma_t(\tan\beta')} \right)^2$$

- $(m_{H^\pm}, \tan\beta)$ grids for type-I and II 2HDM available at

<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCHXSWGMSMCharged>

2. SUSY QCD NLO corrections



- The leading SUSY QCD corrections can be added as usual for the heavy H^+ case
 - Δb corrections
 - Same simple arithmetic trick to fold them in without knowledge of the three individual terms [from the [LHC H XS WG twiki](#)]
 - Find the Δb value corresponding to t_b
 - Calculate $t_{\text{eff}} = t_b / \sqrt{1 + \Delta b}$
 - Using the cross sections without SUSY-QCD NLO corrections, get the cross section which corresponds to t_{eff} (!)
 - Multiply the result from the previous bullet with $1/(1 + \Delta b) \Rightarrow$ this is your cross section
- Note that for $\tan \beta < 10$, non-factorizable corrections can become significant – $O(10\%)$

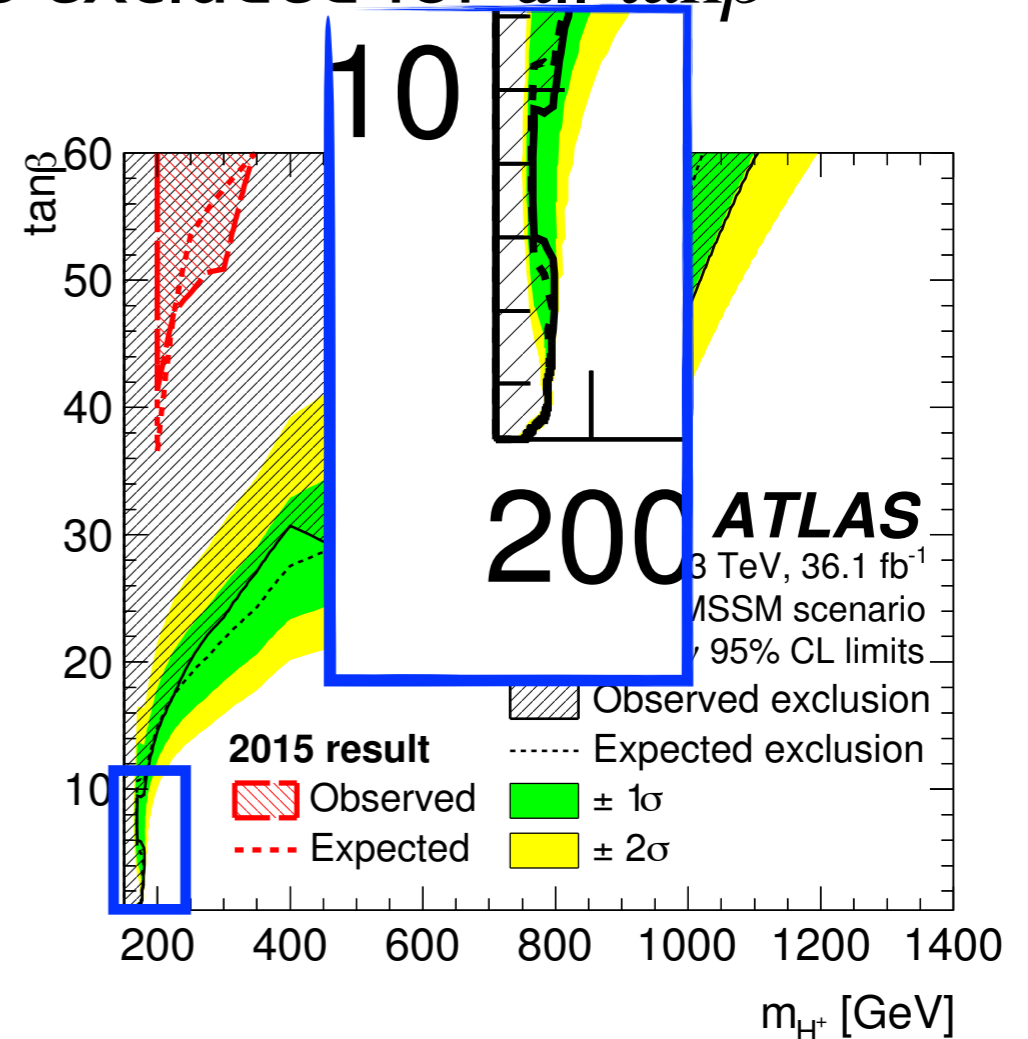
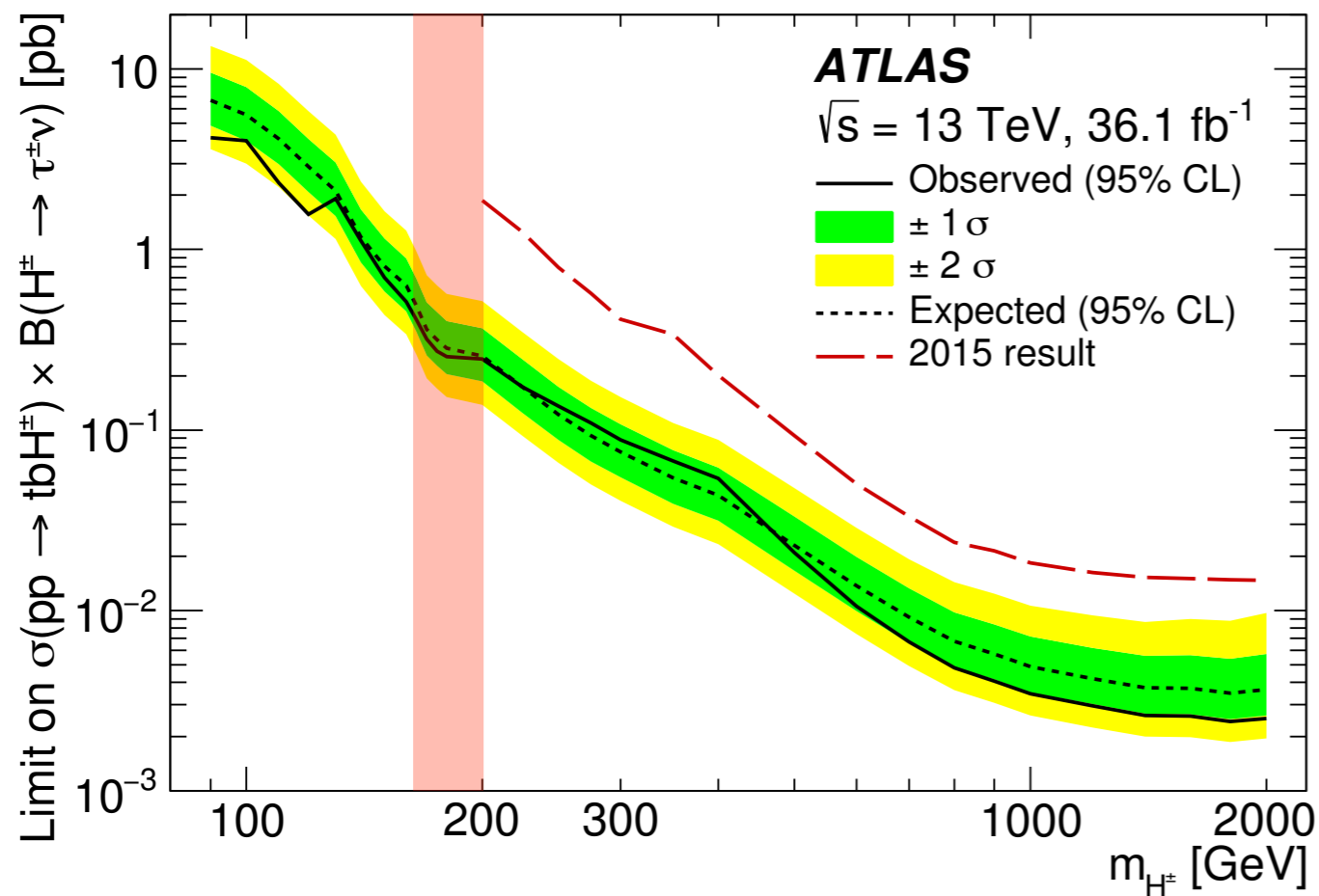
Recommendations for signal simulation

<https://cp3.irmp.ucl.ac.be/projects/madgraph/wiki/chargedHiggs>

- Since a fully-differential NLO computation does not exist, current recommendations for signal simulation rely on LO events, normalised to NLO cross section
- This makes it possible also to add contributions with the extra scalars (h/H/A), which can be evaluated at LO
- More refined approaches can also be performed (reweight jet multiplicity, ...)
- Questions for our experimental colleagues:
 - Is this enough?
 - What is on your wish-list?
 - Any other obstacle that us (theorists) should help removing?

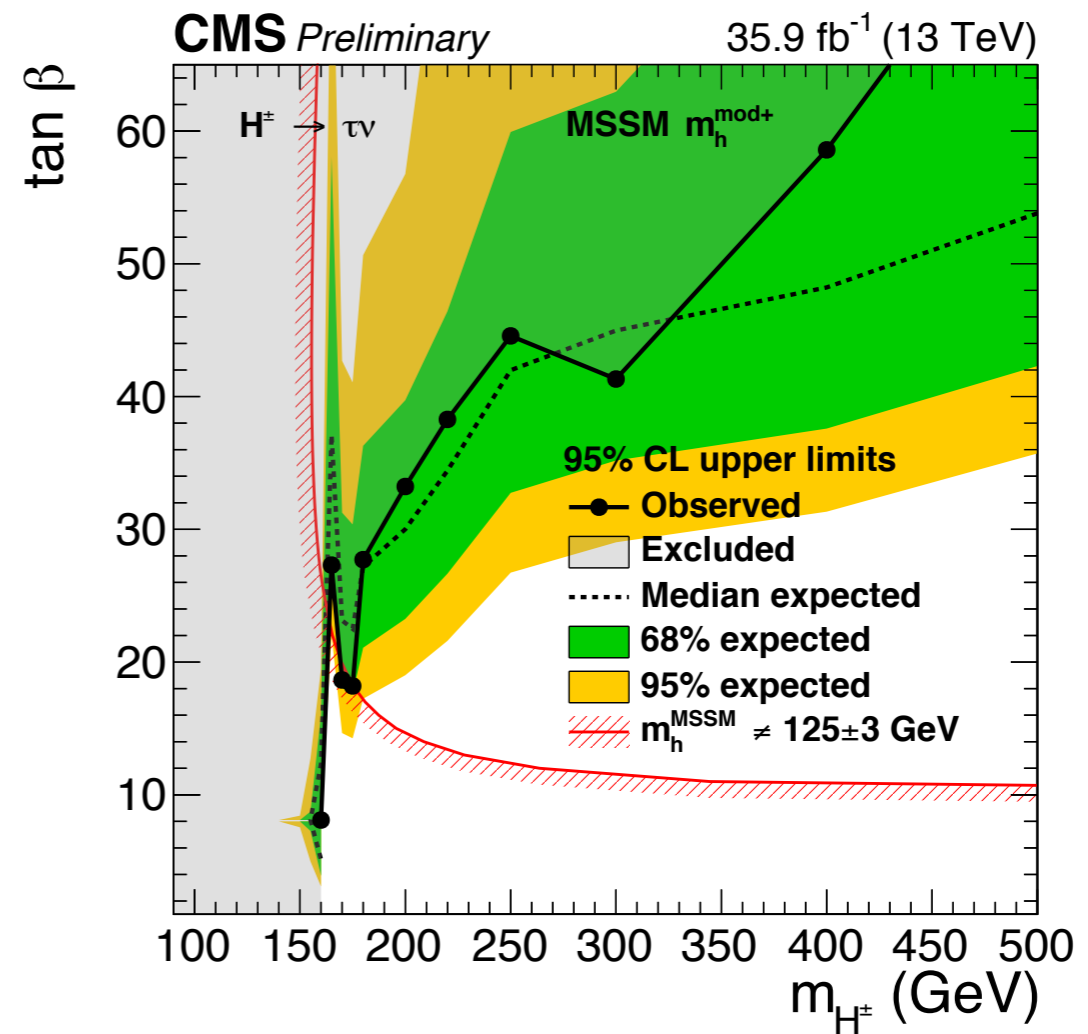
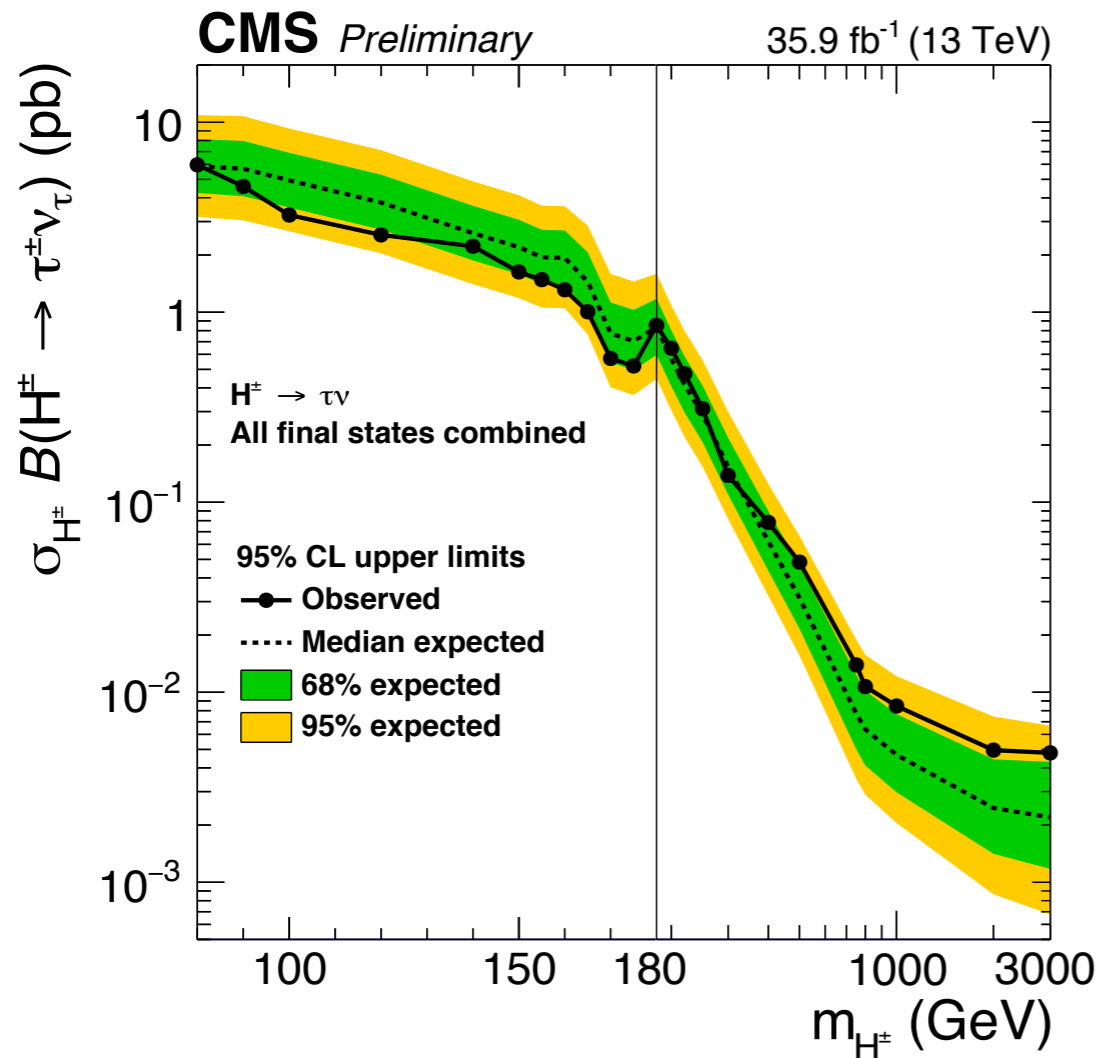
First steps in a new land

- Very recently, ATLAS published updated searches for $H^\pm \rightarrow \tau\nu$ including the intermediate-mass range
(arXiv:1807.07915, see Blake Oliver Burghgrave's talk this morning)
- Charged Higgses below 160 GeV are excluded for all $\tan\beta$



First steps in a new land

- CMS has followed yesterday!
(CMS-PAS-HIG-18-014, see Alexandros Attikis and Jan Eyserman's talks this morning)
- All $\tan\beta$ values excluded up to 150 GeV



Conclusions & Outlook

- The discovery of a charged Higgs boson will be a clear sign of BSM physics
- NLO predictions for the charged Higgs cross-section have been made available for $m_{H^\pm} \sim m_t$, making accurate predictions available in all the mass range
- The full $pp \rightarrow H^\pm W^\mp b \bar{b}$ process interpolates well between the low- and high-mass region
- NLO K-factor is 1.5-1.6, with little $m_{H^\pm}, \tan\beta$ dependence and scale uncertainties reduced to 10-20%
- Full $\tan\beta$ scan (for type-II and type-I 2HDM) available on <https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCHXSWGMSMCharged>
- First analyses in the intermediate-mass region published by **ATLAS** and **CMS**, with 36 fb^{-1} . Masses up **160 GeV** / **150 GeV** excluded in the **hMSSM** / **MSSM** $m_H^{\text{mod}+}$