

## 1 Introduction

**Energy conditions (EC)** constrain the energy momentum tensor  $T_{\mu\nu}$ . Example:

$$\text{Null EC : } T_{\mu\nu} k^\mu k^\nu \geq 0 \quad \forall k_\mu k^\mu = 0$$

- Physical motivation: energy/energy fluxes expected to be non-negative (classically)
- Mathematical consequence: convexity condition for focussing-/singularity theorems
- Problem: all classical ECs violated by reasonable quantum fields**

Therefore, need quantum EC. Example:

$$\text{Averaged NEC : } \int dx^\lambda k_\lambda \langle T_{\mu\nu} k^\mu k^\nu \rangle \geq 0$$

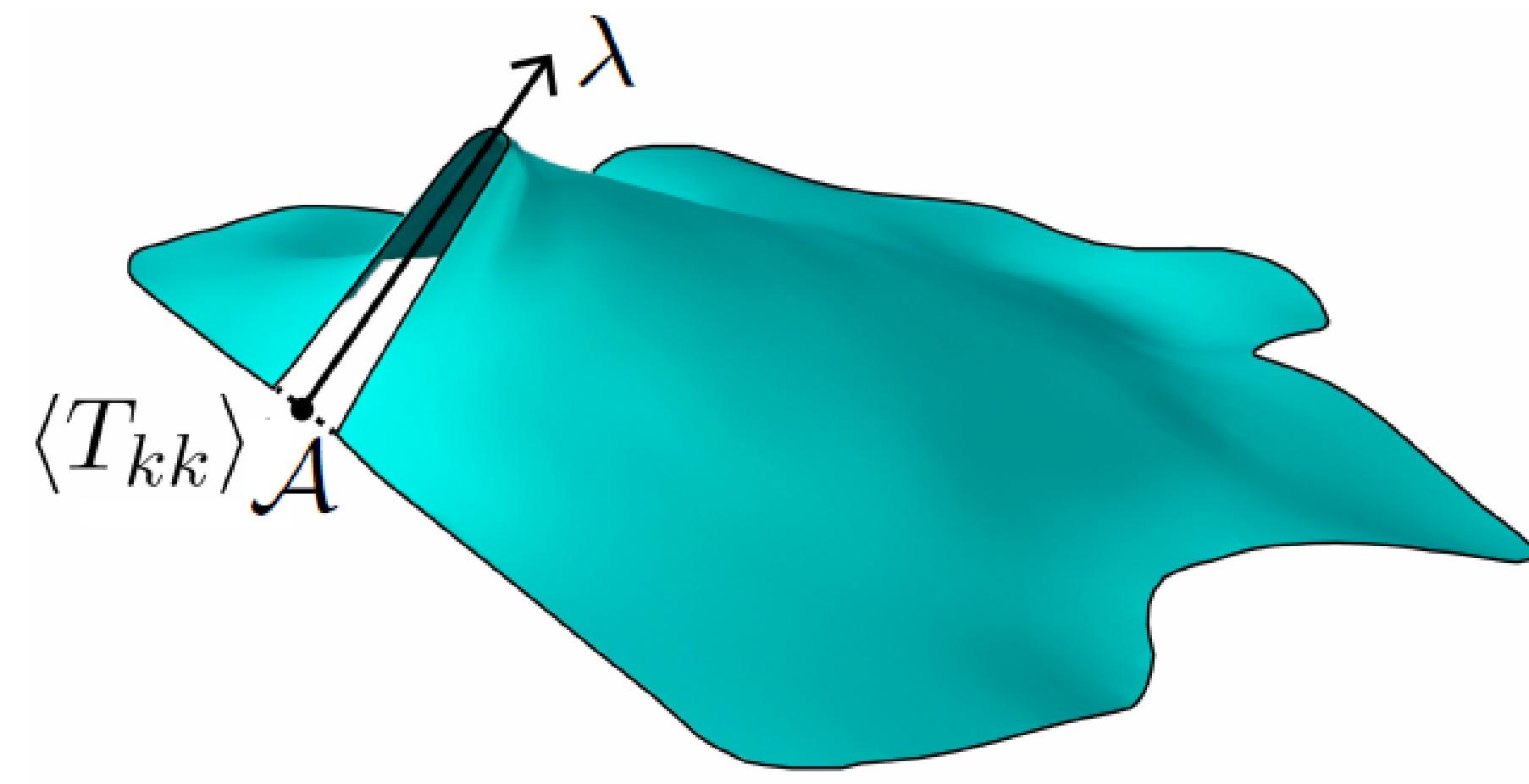
- Proofs of ANEC exist
- Sufficient for proving focussing theorems
- Physical relation to quantum interest conjecture
- Problem: ANEC non-local**

Is there a local quantum energy condition?

## 2 QNEC

- QNEC**, proposed by Bousso, Fisher, Leichenauer and Wall [1] is the following inequality:

$$\langle T_{kk} \rangle \geq \frac{\hbar}{2\pi\sqrt{\gamma}} S''$$



$$T_{kk} = T_{\mu\nu} k^\mu k^\nu \quad \forall k_\mu k^\mu = 0.$$

$\gamma$ : induced metric at the boundary of entangling region  
 $S''$ : second variation of entanglement entropy (EE)  
w.r.t. surface deformations along  $k^\mu$

- QNEC proven for free QFTs [2] holographic CFTs [3] and general CFTs [4]
- QNEC stronger in 2D:

$$\langle T_{kk} \rangle \geq \frac{\hbar}{2\pi} \left( S'' + \frac{6}{c} S'^2 \right)$$

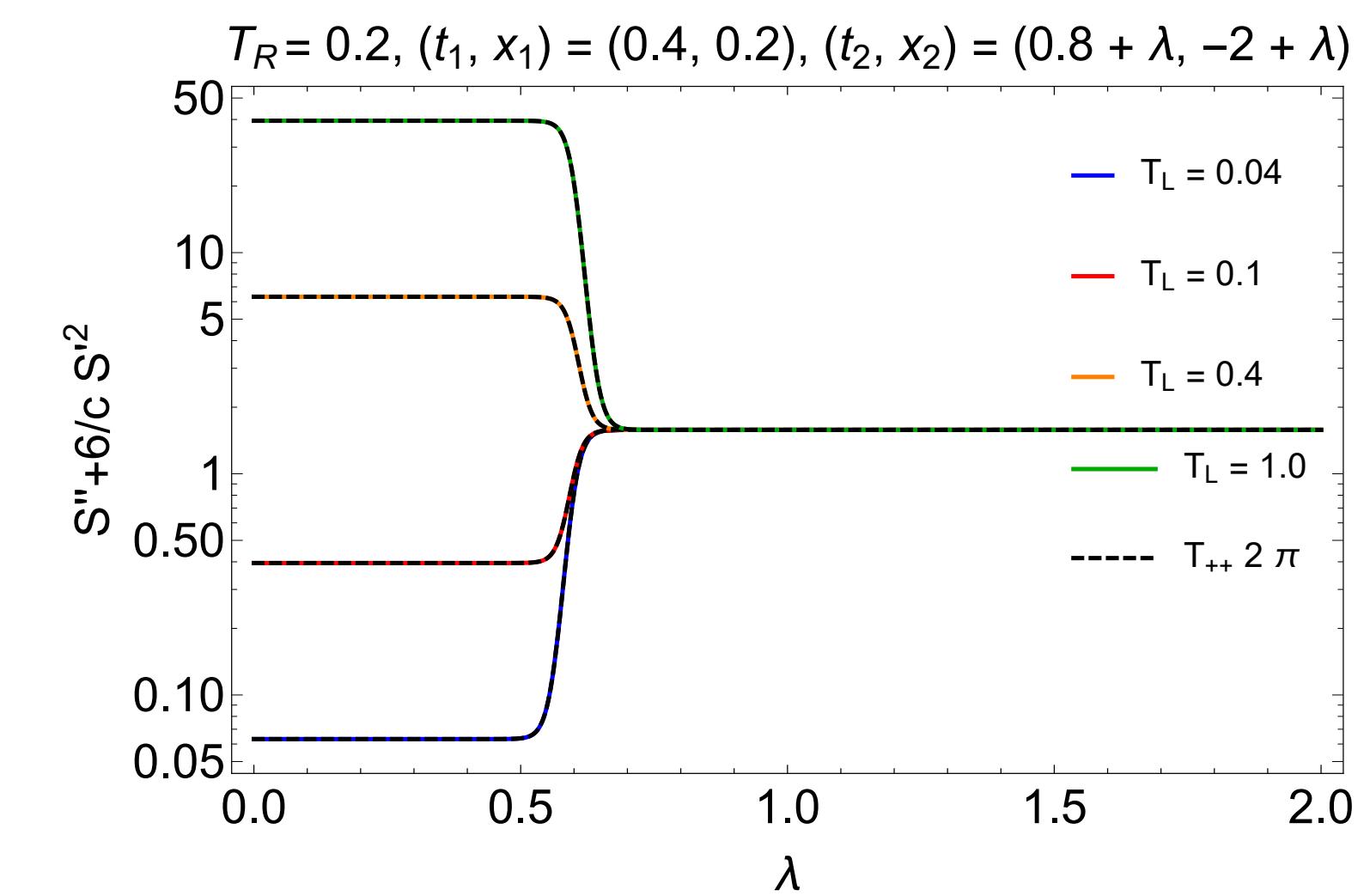
$c$ : central charge

## 3 Analytic Calculation

- $\text{AdS}_3/\text{CFT}_2$ : holographic proof of QNEC saturation for all states dual to Bañados geometries (vacuum, particles in  $\text{AdS}_3$ , BTZ black holes and their Virasoro descendants)
- QNEC saturation even for far-from equilibrium transport in strongly coupled CFT like the hot-cold bath critical quantum system studied in [5]

Sketch of proof:

EE transforms like anomalous weight  $(0, 0)$  operator.  
Define vertex operator:  $V := \exp(\frac{6}{c} S)$ .  
QNEC saturation equivalent to vertex operator solving Hill's equation:  $V'' - \mathcal{L}V = 0$ .  
Conformal symmetries show that for states dual to Bañados geometries vertex operator always solves Hill's equation.



QNEC saturation for hot-cold bath.

## 4 Perturbative Calculation

- Perturbative QNEC calculation for  $\text{AdS}_5$  Schwarzschild black brane for small and large entangling regions; numerically in between [6].
- Finite central charge corrections in  $\text{AdS}_3$  backreacted by massive scalar field with conformal weights  $(h, h)$  (work in progress [9]).

Requires corrections to RT-surface and from bulk EE [8]:

$$S = \frac{\mathcal{A}}{4G_N} + \frac{\delta\mathcal{A}}{4G_N} + S_{\text{bulk}}$$

QNEC satisfied for all conformal weights  $h \geq 1/2$ .

Small interval:

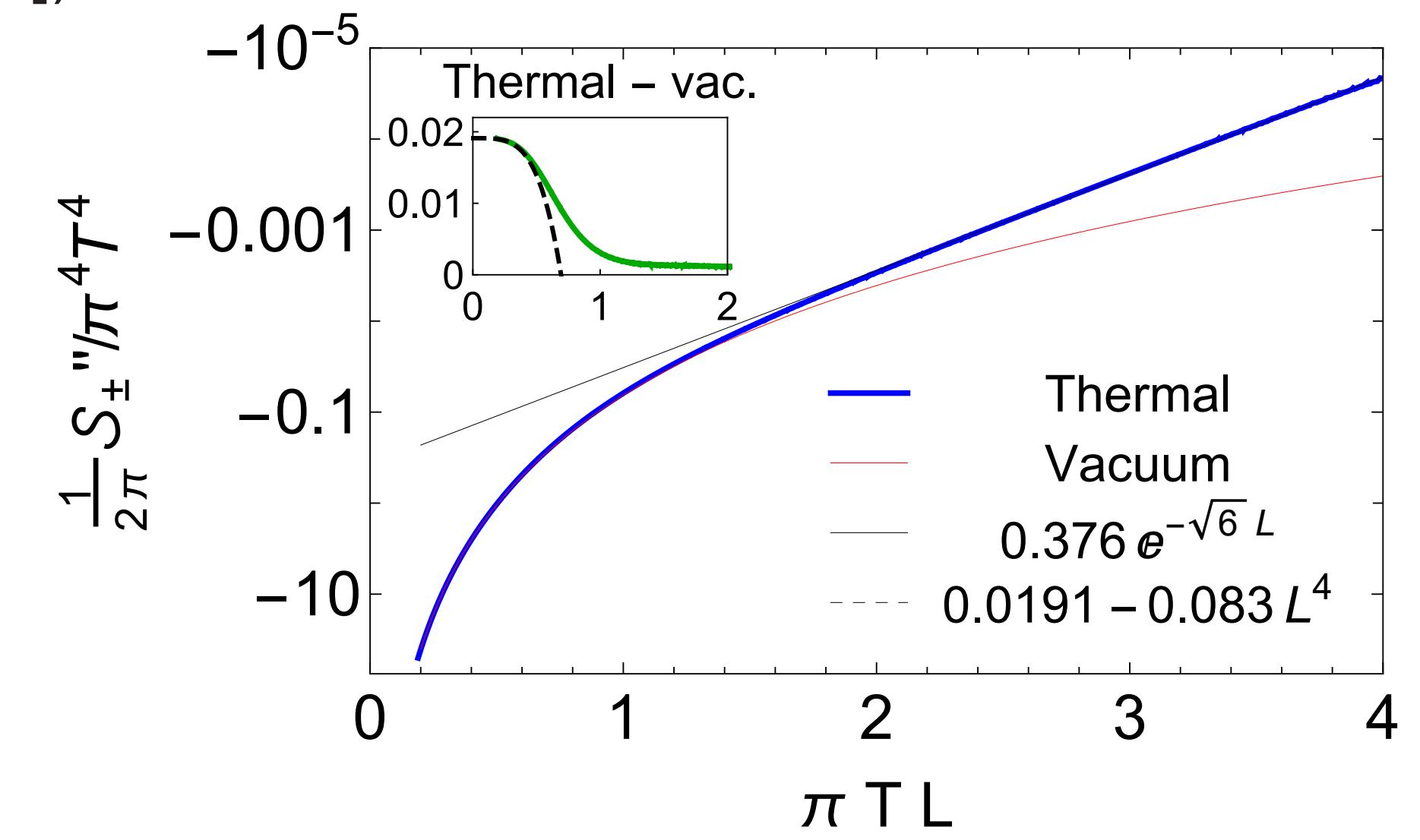
QNEC saturation up to small corrections:

$$2\pi \langle T_{\pm\pm} \rangle - S'' - \frac{6}{c} (S')^2 = +\mathcal{O}(\Delta\varphi^{4h})$$

Half-interval:

QNEC satisfied but gapped by  $h/4$  in large- $h$  limit:

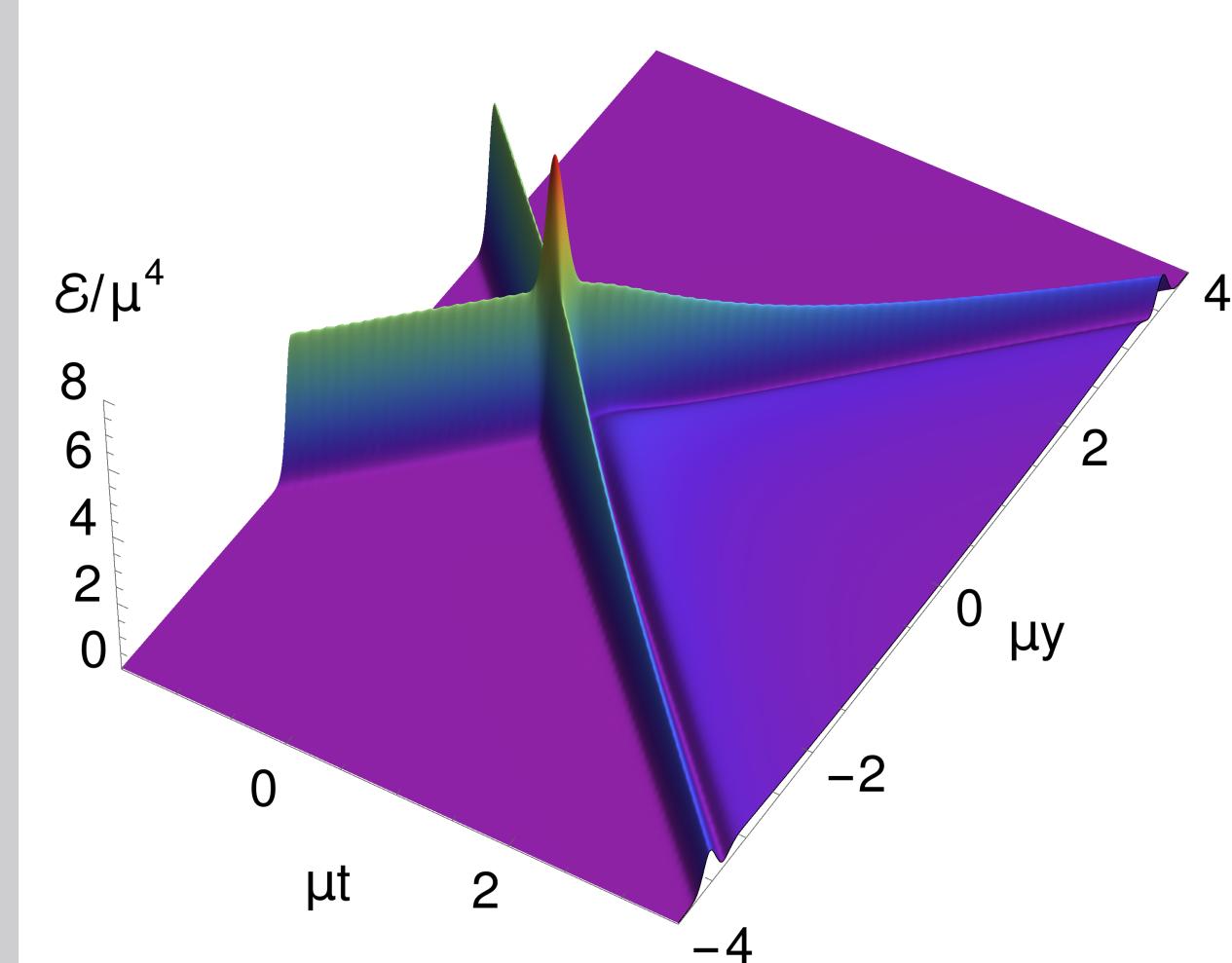
$$2\pi \langle T_{\pm\pm} \rangle - S'' - \frac{6}{c} (S')^2 \Big|_{\Delta\varphi=\pi} = \frac{h}{4} + \dots$$



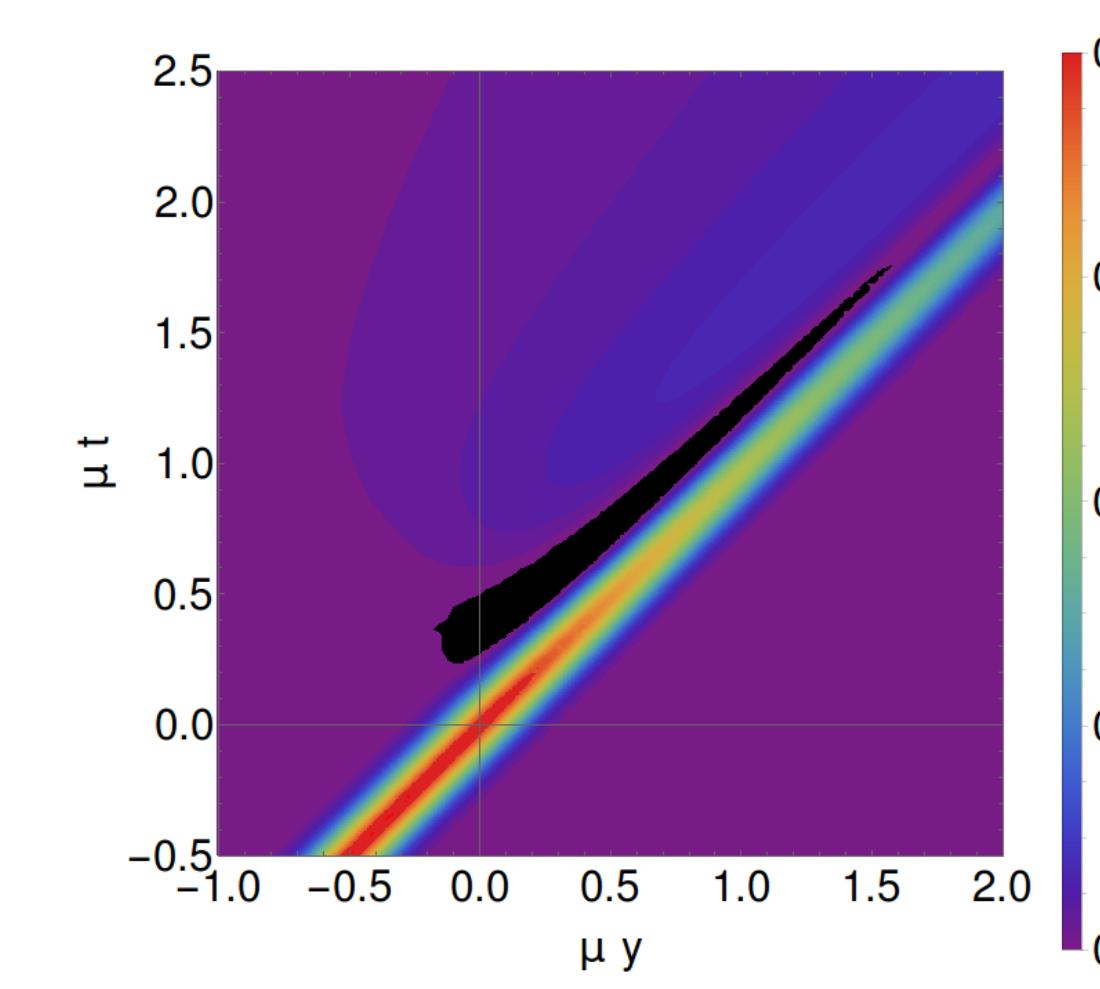
QNEC as function of size of entangling region.  
Black lines are perturbative, colored lines numerical results.

## 5 Numerical Calculation

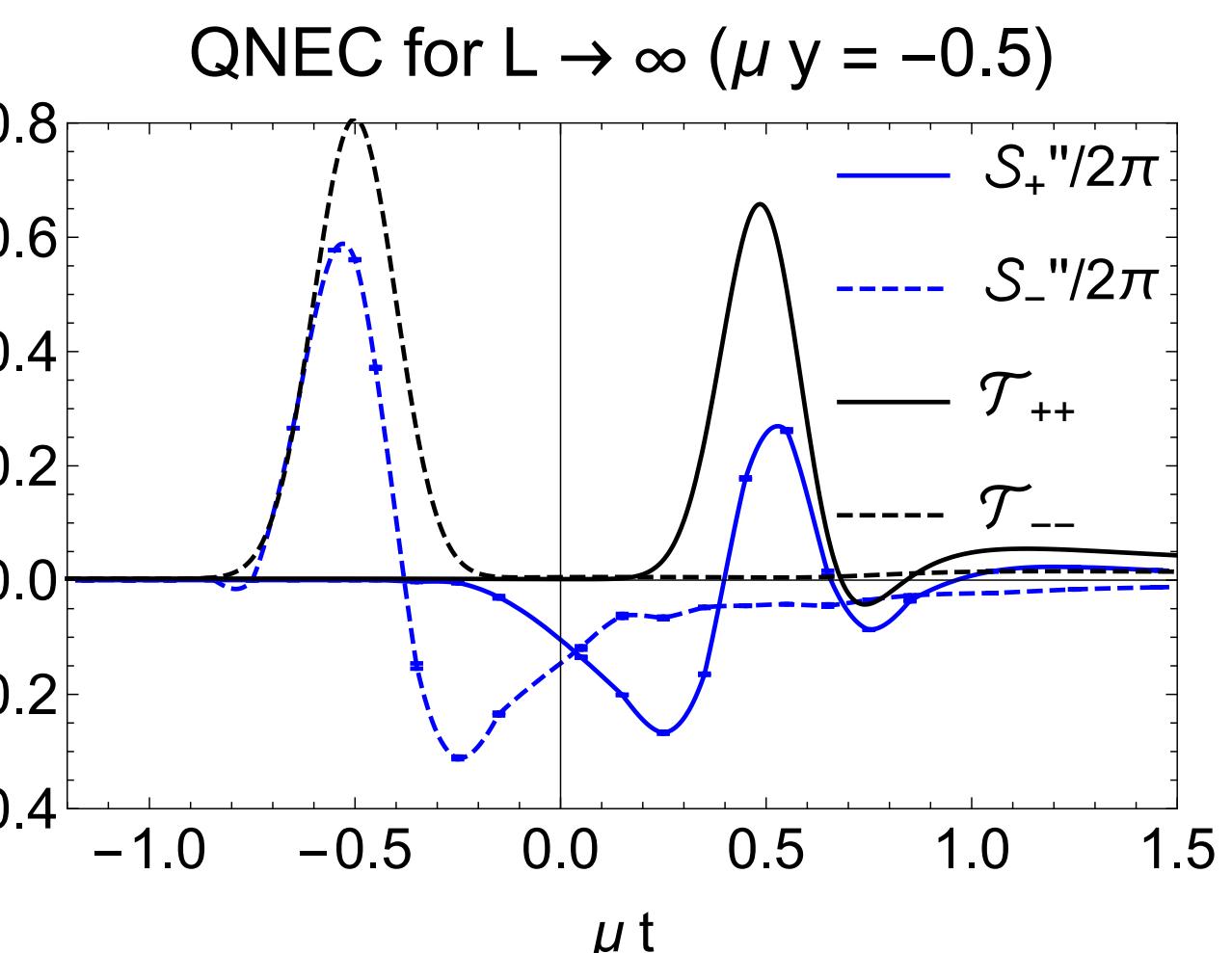
- Holographic model for non-abelian plasma formation (toy model for heavy ion collisions): colliding gravitational shock waves in  $\text{AdS}_5$ .
- Null EC violated, but QNEC satisfied and can saturate in far from equilibrium state [7].



Energy density of colliding shock waves.



NEC violation in shock wave scenario.



Time evolution of QNEC.

## 6 Summary & References

- QNEC only known local EC [1], proven for large class of unitary QFTs [2, 3, 4].
- Provided first numerical studies of QNEC in  $\text{AdS}_5/\text{CFT}_4$  [7].
- Gravitational shockwaves violate Null EC but satisfy QNEC [7].
- QNEC saturation possible in far from equilibrium regime - "quantum equilibrium"? [7, 9].
- Finite central charge corrections to QNEC in  $\text{AdS}_3/\text{CFT}_2$ : saturation at small interval, gap at large interval [9].

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[4] S. Balakrishnan et al., hep-th/1706.09432

[5] M.J. Bhaseen et al., Nature Phys. **11**, 5 (2015)

[6] C. Ecker et al., Phys. Rev. **D97** (2018) no.12, 126016

[7] C. Ecker et al., JHEP **11** (2016) 054

[8] T. Faulkner et al., JHEP **11** (2013) 074.

[9] C. Ecker, D. Grumiller, M.M. Sheikh-Jabbari, P. Stanzer and Wilke van der Schee, in preparation