

# Particle multiplicities in the central region of high-energy collisions from running coupling $k_T$ – factorization

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in collaboration with

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# Motivation

CGC effective theory entered the next-to-leading order era → several processes calculated @ NLO (and their numerical implementation are underway)

- **Regarding the Balitsky-Kovchegov (BK) evolution eq.:**

**Running coupling corrections to the kernel of BK eq. with the solution being able to describe several observables @ HERA, RHIC, LHC**

Balitsky, PRD 75, 014001 (2007); Balitsky and Chirilli, PRD 77, 014019 (2008); Kovchegov and Weigert, NPA 784, 188 (2007), NPA 789, 260 (2007); Kovchegov, Kuokkanen, Rummukainen and Weigert, NPA 823, 47 (2009).

**Large single and double transverse logarithms resummed to all orders in NLO BK eq. → resulting evolution eq. is stable & generates a physically meaningful evolution of the dipole amplitude**

Iancu, Madrigal, Mueller, Soyez and Triantafyllopoulos, PLB 750, 643 (2015); Lappi and Mäntysaari, PRD 91, no. 7, 074016 (2015), PRD 93, no. 9, 094004 (2016).

- **Also, JIMWLK evolution eq. @ NLO;**

Kovner, Lublinsky and Mulian, PRD 89, no. 6, 061704 (2014), JHEP 1408, 114 (2014); Lublinsky and Mulian, arXiv:1610.03453.

# Motivation

- Regarding the “hybrid formalism”:

**NLO corrections calculated and implemented numerically → better agreement with experimental data @ RHIC/LHC energies for forward hadron production;**

Chirilli, Xiao and Yuan, PRL 108, 122301 (2012); Stasto, Xiao and Zaslavsky, PRL 112, no. 1, 012302 (2014); Altinoluk, Armesto, Beuf, Kovner and Lublinsky, PRD 91, no. 9, 094016 (2015); Watanabe, Xiao, Yuan and Zaslavsky, PRD 92, no. 3, 034026 (2015).

- Regarding the  $k_T$  – factorization:

**$k_T$  – factorization formula for inclusive gluon production @ small-x beyond LO with running coupling corrections conjectured**

Horowitz and Kovchegov, Nucl. Phys. A 849, 72 (2011)

**So far, only a qualitative study done in**

Durães, A.V.G., Gonçalves and Navarra, **PRD 94, 054023 (2016)**

**KLN UGD + Local Parton-Hadron Duality + minimum bias collisions**

**Motivates a more robust calculation**

# $k_T$ -factorization: multiplicity in $A+B \rightarrow g+X$ @ low-x

fixed by data; includes “K-factors” due to high order corrections + Frag. Functions

$$\frac{d^3\sigma}{d^2k_T dy} = \frac{N}{C_F} \frac{2}{\mathbf{k}^2} \int d^2b d^2b' d^2q \alpha_s \phi_{h_1}(\mathbf{q}, \mathbf{b}, x_1) \phi_{h_2}(\mathbf{k} - \mathbf{q}, \mathbf{b} - \mathbf{b}', x_2)$$

convolution of the projectile's & target's unintegrated gluon distribution (UGD)

$$\phi(\mathbf{k}, \mathbf{b}, y) = \frac{C_F}{\alpha_s (2\pi)^3} \int d^2r e^{-i\mathbf{k}\cdot\mathbf{r}} \nabla_r^2 \mathcal{N}_{\mathcal{A}}(\mathbf{r}, \mathbf{b}, y) \quad \mathbf{k} = (k_x, k_y)$$

UGD

2-D Fourier Transform of the gluon dipole scattering amplitude

$$x_{1,2} = k_T / \sqrt{s} \exp(\pm y) \quad \text{momentum fraction of the proj./targ. gluon}$$

Originally derived in the fixed coupling (FC) approx.:  $\alpha_s = \text{const.}$

(the impact parameter dependence will be omitted for sake of simplicity)

# $k_T$ -factorization: multiplicity in $A+B \rightarrow g+X$ @ low-x

$$\frac{d^3\sigma}{d^2k_T dy} = \frac{N}{C_F} \frac{2}{\mathbf{k}^2} \int d^2b d^2b' d^2q \alpha_s \phi_{h_1}(\mathbf{q}, \mathbf{b}, x_1) \phi_{h_2}(\mathbf{k} - \mathbf{q}, \mathbf{b} - \mathbf{b}', x_2)$$

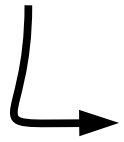
Originally derived in the fixed coupling (FC) approx.:  $\alpha_s = \text{const.}$

Later,  $\alpha_s \rightarrow \alpha_s(Q^2)$  in FC formula: better agreement of theory & data;

$Q^2$  fixed by hand! Distinct choices for  $Q^2 \rightarrow$  similar results

# The running coupling $k_T$ – fact. formula

$$\frac{d^3\sigma}{d^2k_T dy} = \frac{2C_F}{\pi^2} \frac{1}{k^2} \int d^2q \bar{\phi}_{h_1}(\mathbf{q}, x_1) \bar{\phi}_{h_2}(\mathbf{k} - \mathbf{q}, x_2) \frac{\alpha_s(\Lambda_{\text{coll}}^2 e^{-5/3})}{\alpha_s(Q^2 e^{-5/3}) \alpha_s(Q^{*2} e^{-5/3})}$$



Result of resummation of relevant 1-loop corrections into the running coupling

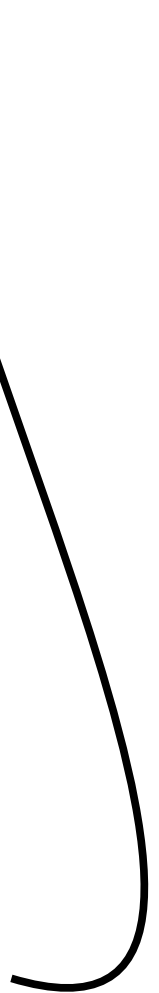
Horowitz and Kovchegov, NPA 849, 72 (2011)

**$Q^2$  from a formal calculation!**

$$\bar{\phi}(\mathbf{k}, \mathbf{b}, y) = \alpha_s \phi(\mathbf{k}, \mathbf{b}, y)$$

$$\Lambda_{\text{coll}} \sim k_T \quad \text{Kovchegov and Weigert, NPA 807, 158 (2008)}$$

$\alpha_s$ -factors appear explicitly in the expression



# The running coupling $k_T$ – fact. formula

$$\frac{d^3\sigma}{d^2k_T dy} = \frac{2C_F}{\pi^2} \frac{1}{\mathbf{k}^2} \int d^2q \bar{\phi}_{h_1}(\mathbf{q}, x_1) \bar{\phi}_{h_2}(\mathbf{k} - \mathbf{q}, x_2) \frac{\alpha_s(\Lambda_{\text{coll}}^2 e^{-5/3})}{\alpha_s(Q^2 e^{-5/3}) \alpha_s(Q^{*2} e^{-5/3})}$$

Horowitz and Kovchegov, NPA 849, 72 (2011)

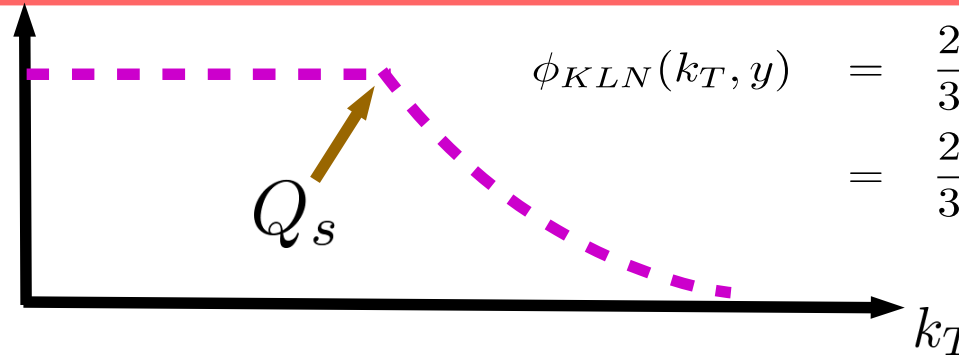
$Q^2$  given by:

$$\begin{aligned} \ln \frac{Q^2}{\mu_{\overline{\text{MS}}}^2} = & \frac{1}{2} \ln \frac{q^2 (\mathbf{k} - \mathbf{q})^2}{\mu_{\overline{\text{MS}}}^4} - \frac{1}{4 q^2 (\mathbf{k} - \mathbf{q})^2 [(\mathbf{k} - \mathbf{q})^2 - q^2]^6} \left\{ k^2 [(\mathbf{k} - \mathbf{q})^2 - q^2]^3 \right. \\ & \times \left\{ [ [(\mathbf{k} - \mathbf{q})^2]^2 - (q^2)^2 ] [ (k^2)^2 + ((\mathbf{k} - \mathbf{q})^2 - q^2)^2 ] + 2 k^2 [ (q^2)^3 - [(\mathbf{k} - \mathbf{q})^2]^3 ] \right. \\ & \left. \left. - q^2 (\mathbf{k} - \mathbf{q})^2 [ 2 (k^2)^2 + 3 [(\mathbf{k} - \mathbf{q})^2 - q^2]^2 - 3 k^2 [(\mathbf{k} - \mathbf{q})^2 + q^2] ] \ln \left( \frac{(\mathbf{k} - \mathbf{q})^2}{q^2} \right) \right\} \right. \\ & + i [(\mathbf{k} - \mathbf{q})^2 - q^2]^3 \left\{ k^2 [(\mathbf{k} - \mathbf{q})^2 - q^2] [ k^2 [(\mathbf{k} - \mathbf{q})^2 + q^2] - (q^2)^2 - [(\mathbf{k} - \mathbf{q})^2]^2 ] \right. \\ & \left. \left. + q^2 (\mathbf{k} - \mathbf{q})^2 \left( k^2 [(\mathbf{k} - \mathbf{q})^2 + q^2] - 2 (k^2)^2 - 2 [(\mathbf{k} - \mathbf{q})^2 - q^2]^2 \right) \ln \left( \frac{(\mathbf{k} - \mathbf{q})^2}{q^2} \right) \right\} \right. \\ & \left. \times \sqrt{2 q^2 (\mathbf{k} - \mathbf{q})^2 + 2 k^2 (\mathbf{k} - \mathbf{q})^2 + 2 q^2 k^2 - (k^2)^2 - (q^2)^2 - [(\mathbf{k} - \mathbf{q})^2]^2} \right\}, \end{aligned}$$

# Going quantitative:

Instead of

$\phi_{KLN}$



$$\begin{aligned}\phi_{KLN}(k_T, y) &= \frac{2C_F}{3\pi^2}, \quad k_T \leq Q_s \\ &= \frac{2C_F}{3\pi^2} \frac{Q_s^2}{k_T^2}, \quad k_T > Q_s\end{aligned}$$

... get UGD from rcBK evolution eq.:

$$\frac{\partial \mathcal{N}(r, Y)}{\partial Y} = \int d^2 r_1 K(r, r_1, r_2) [\mathcal{N}(r_1, Y) + \mathcal{N}(r_2, Y) - \mathcal{N}(r, Y) - \mathcal{N}(r_1, Y) \mathcal{N}(r_2, Y)]$$

$$\mathcal{N}_F(r, Y) \equiv \mathcal{N}(r, Y) \quad ; \quad \mathcal{N}_A = 2\mathcal{N}_F - \mathcal{N}_F^2 \quad ; \quad Y = \ln(x_0/x) \quad ; \quad x_0 = 0.01$$

rcBK provides small-x evolution given an initial condition (I.C.)!

AAMQS I.C.:

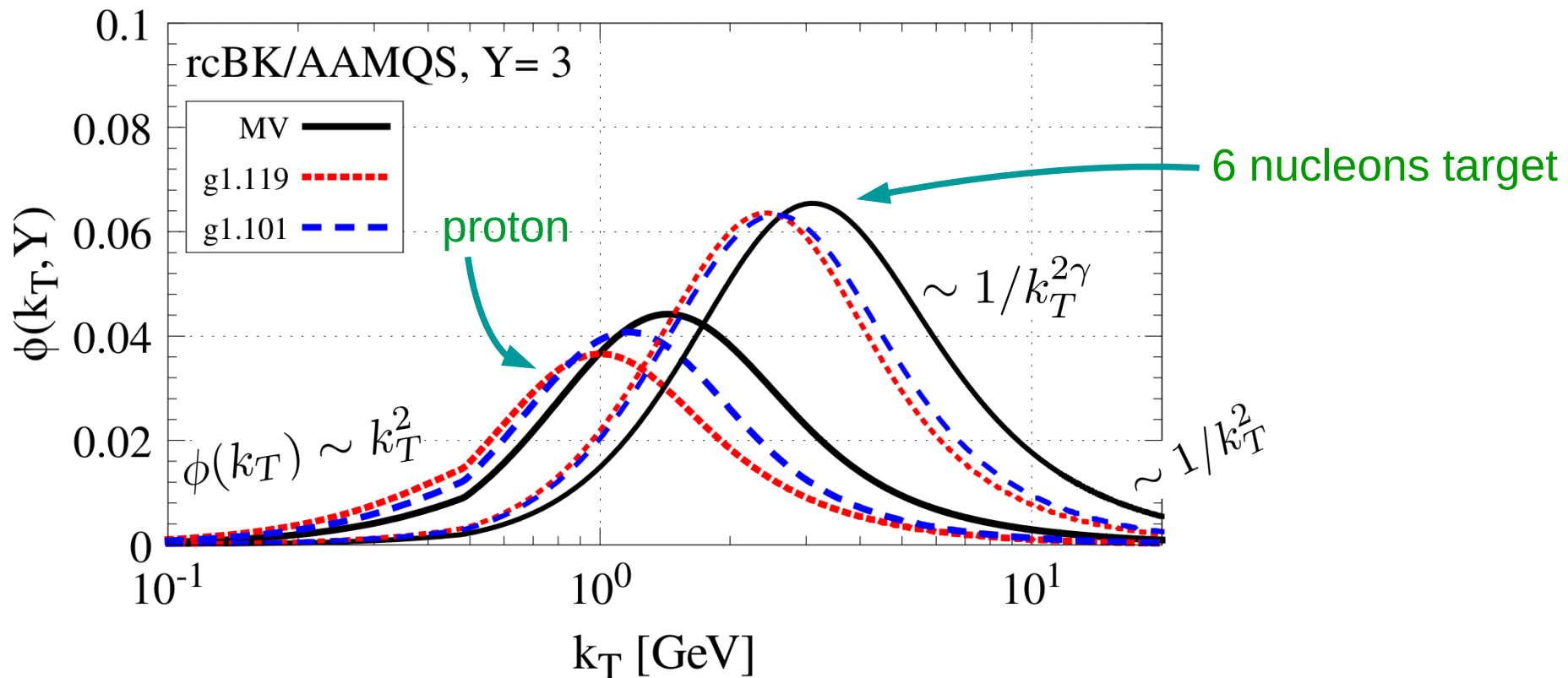
$$\mathcal{N}_F(r, x_0) = 1 - \exp \left[ - \frac{(r^2 Q_{s0, \text{proton}}^2)^\gamma}{4} \ln \left( \frac{1}{\Lambda r} + e \right) \right]$$



$Q_{s0,proton}^2$  = proton's sat. scale at the initial scale  $x_0$  } fitted to data!  
 $\gamma$  = controls steepness of the UGD tail for  $k_T > Q_{s0,proton}^2$  }

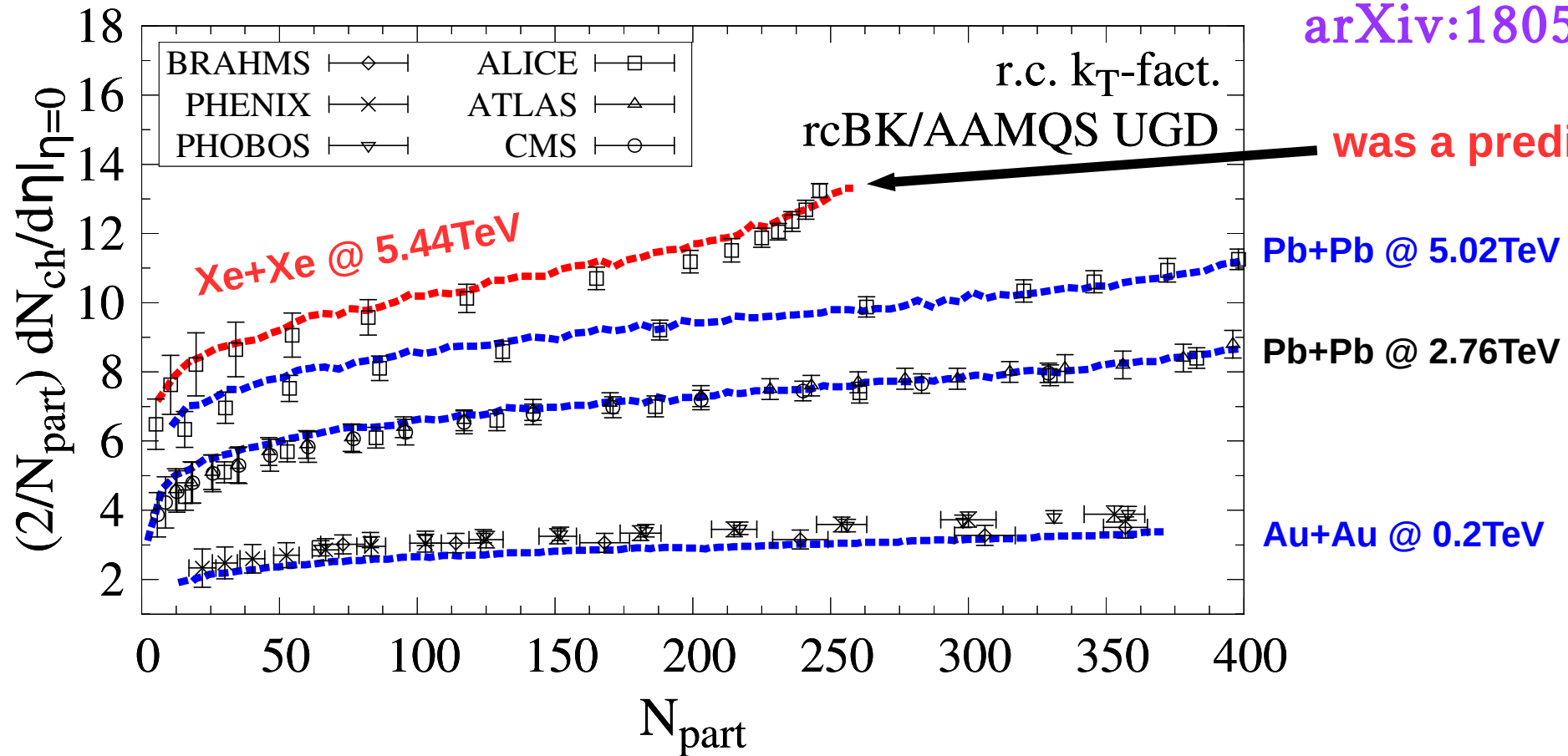
$\gamma = 1$  McLerran-Venugopalan (MV) model as I.C.

For proton: I.C. with  $\gamma > 1$  lead to best fit of HERA e+p data



# Multiplicity vs $N_{\text{part}}$ : A + A, MV I.C.

arXiv:1805.02702

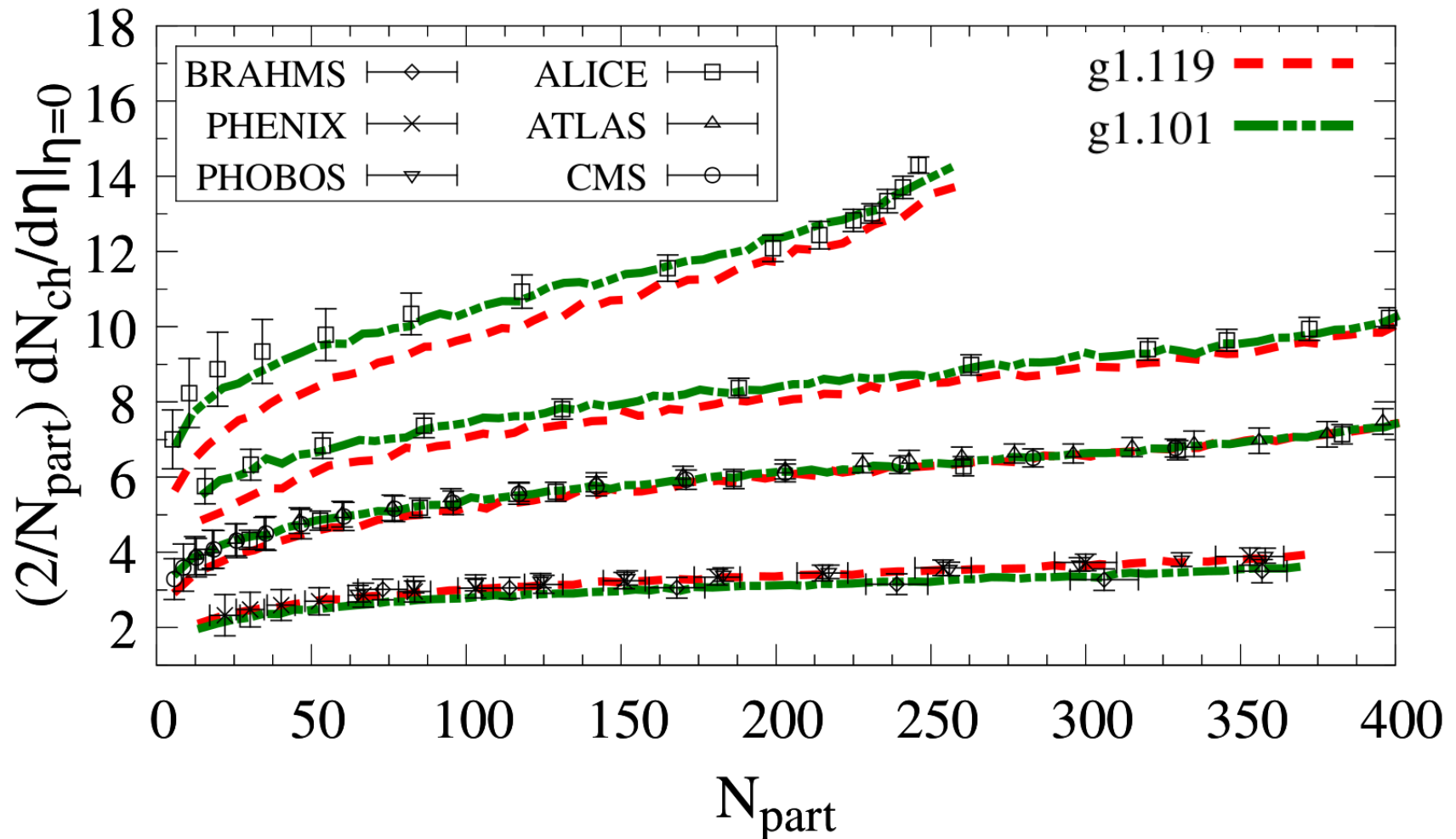


Normalization fixed by Pb+Pb data @ 2.76 TeV; **not changed later!**

**Nice agreement with exp. data from 0.2 TeV to 5.02 TeV & also with new Xe+Xe data @ 5.44 TeV !**

Also, good agreement regarding energy evolution! (Backup slides)

# Multiplicity vs $N_{\text{part}}$ : A + A, $\gamma > 1$ I.C.

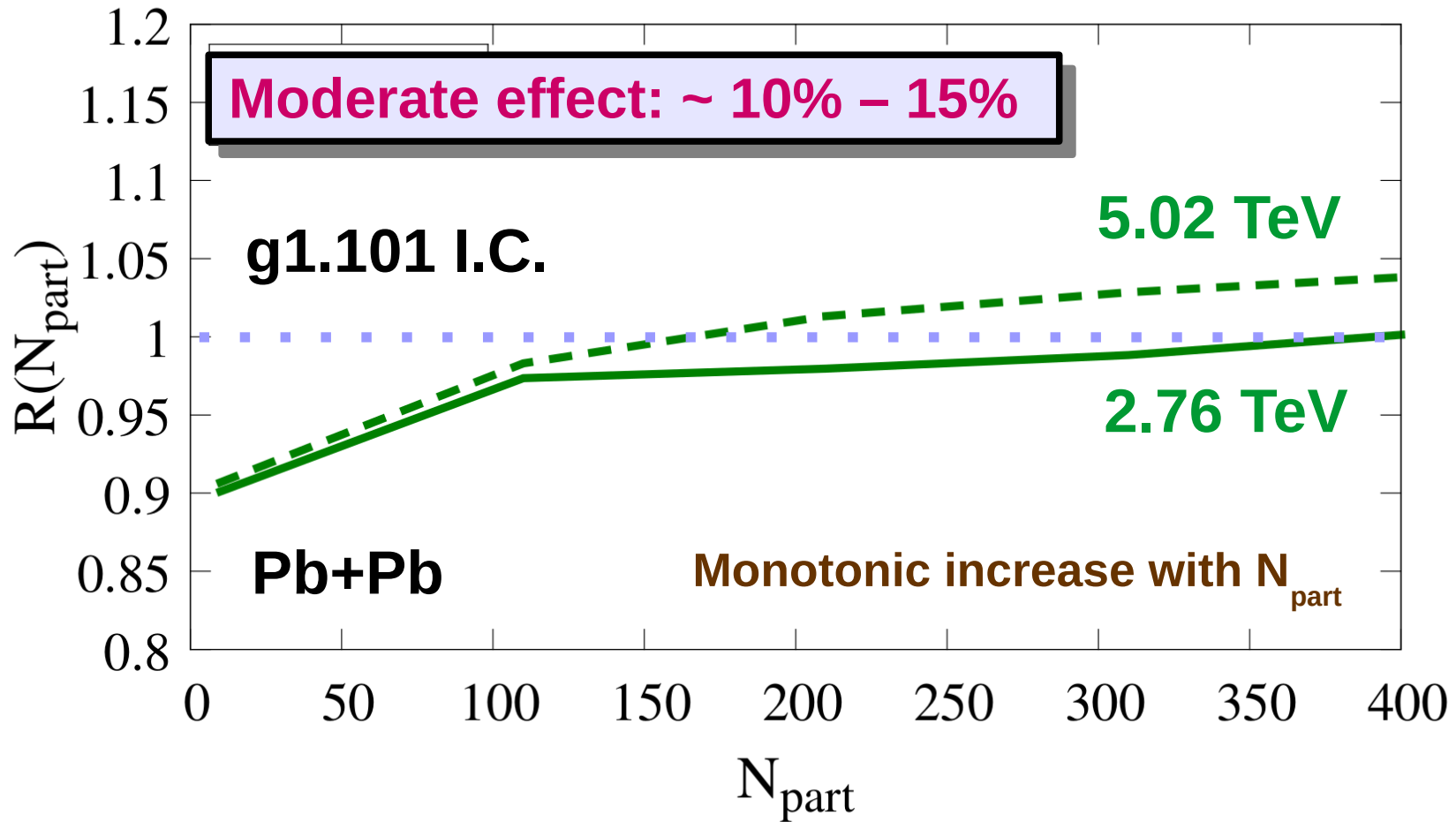


$\gamma = 1.119$  I.C. : poor agreement with data @ highest energies

$\gamma = 1.101$  I.C. : similar results as MV I.C.

# Quantifying the running coupling effects

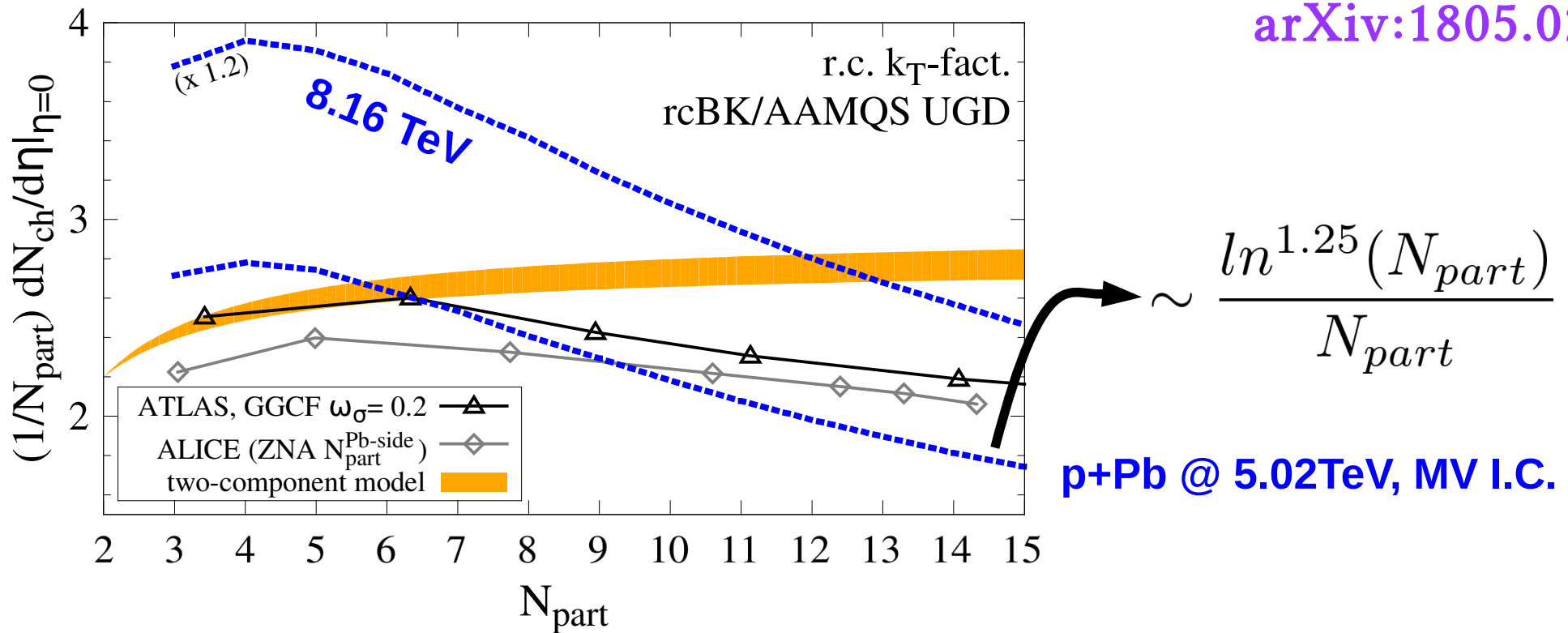
$$R(N_{\text{part}}) = \frac{\text{Fixed coupling expression} + \text{running coupling effects by hand}}{\text{Running coupling expression}}$$



Similar results for all I.C. & p+Pb @ 5.02, 8.16 TeV

# Multiplicity vs $N_{part}$ : p + A, MV I.C.

arXiv:1805.02702



The convolution of two UGDs grows slower than  $N_{part}$ !

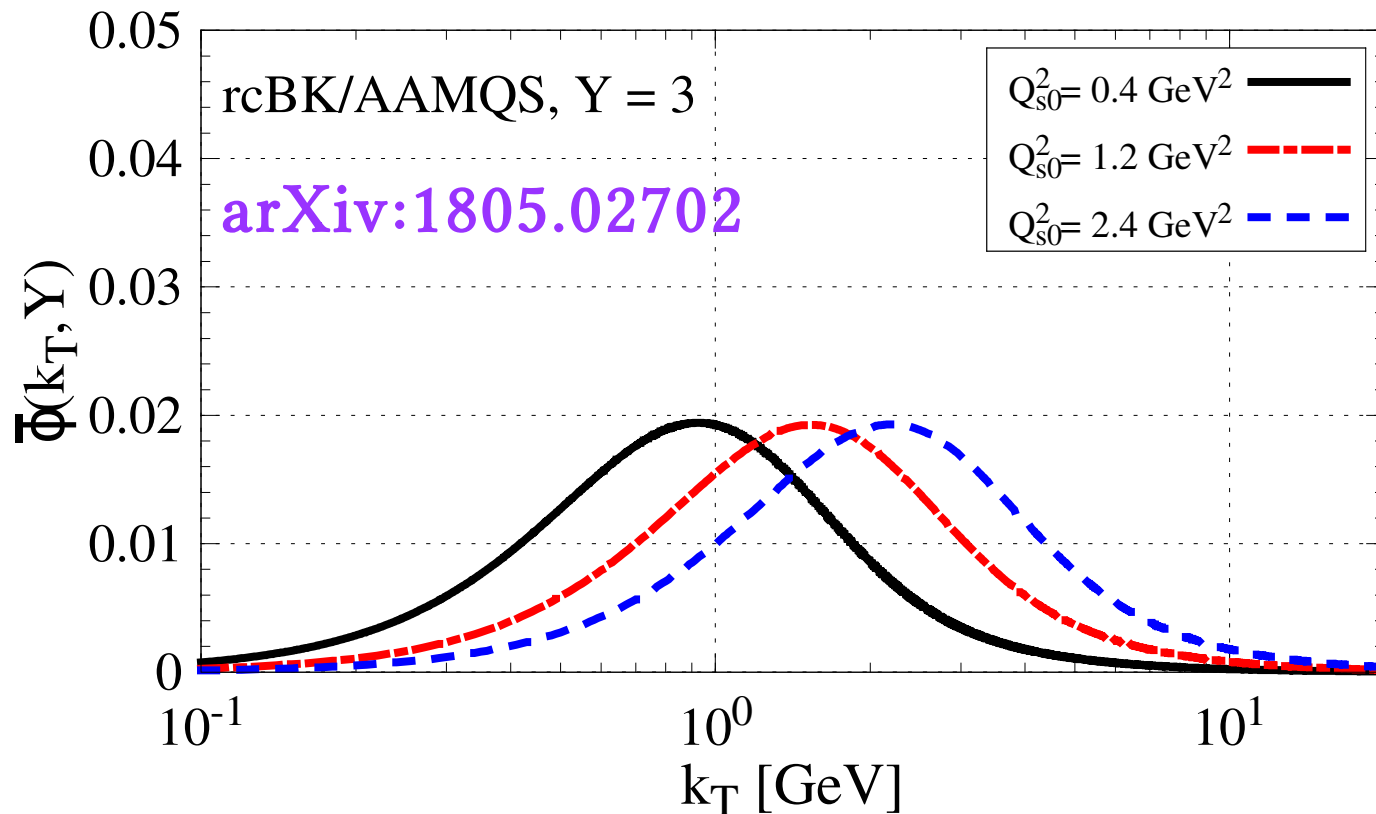
Qualitative agreement with ATLAS & ALICE data @ 5.02 TeV !

$[dE_T/d\eta] / [dN_{ch}/d\eta]$  increases with  $N_{part}$ ! Backup slide

# Multiplicity vs $N_{\text{part}}$ : p + A, MV I.C.

**A+A collisions become more symmetric as  $N_{\text{part}}$  increases**  
(proj. and targ. have  $k_t$  near  $Q_s$ )

**while p+A collisions become more asymmetric!**  
(proton's UGD stay put but target's UGD moves to higher  $k_t$ )

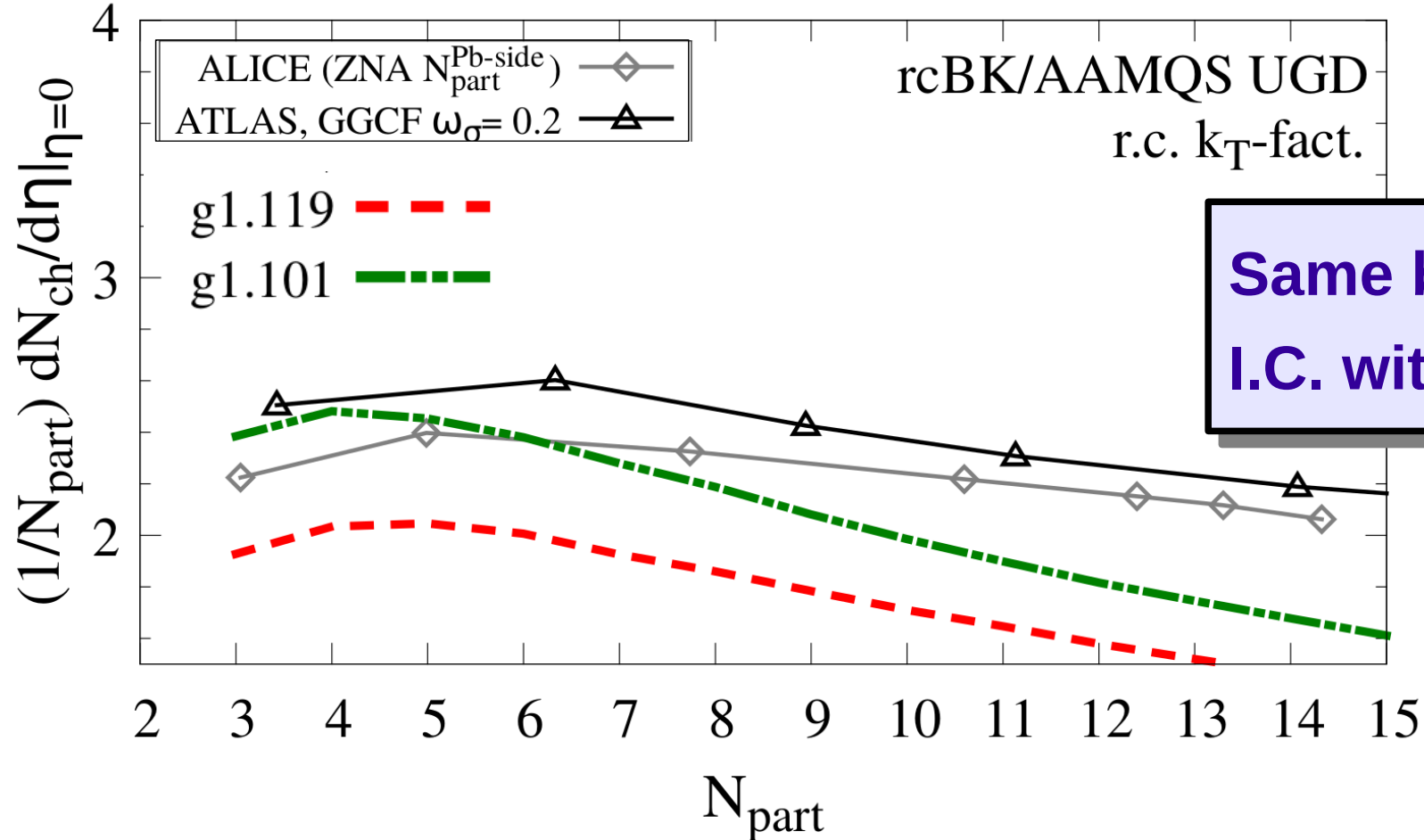


**1 nucleon target (proton)**

**3 nucleons target**

**6 nucleons target**

# Multiplicity vs $N_{\text{part}}$ : $p + A$ , $\gamma > 1$ I.C.



**However: exp. data is flatter than our result!**

**Lack of realistic  $b$ -dependence on proton's UGD?**

**Bias introduced by experimental centrality selection?**

# Multiplicity vs $N_{\text{part}}$ : two-component model, A + A

Two energy dependent shares controlled by  $f$

$$\frac{dN_{AB}}{d\eta} = \left[ \frac{1-f}{2} N_{\text{part}} + f N_{\text{coll}} \right] \frac{dN_{pp}}{d\eta}$$

[“Soft” + “Hard”]  
component

from ALICE: arXiv:1412.6828

From fit of peripheral region [ $N_{\text{part}} < 34$ ] of Pb+Pb & Xe+Xe data:

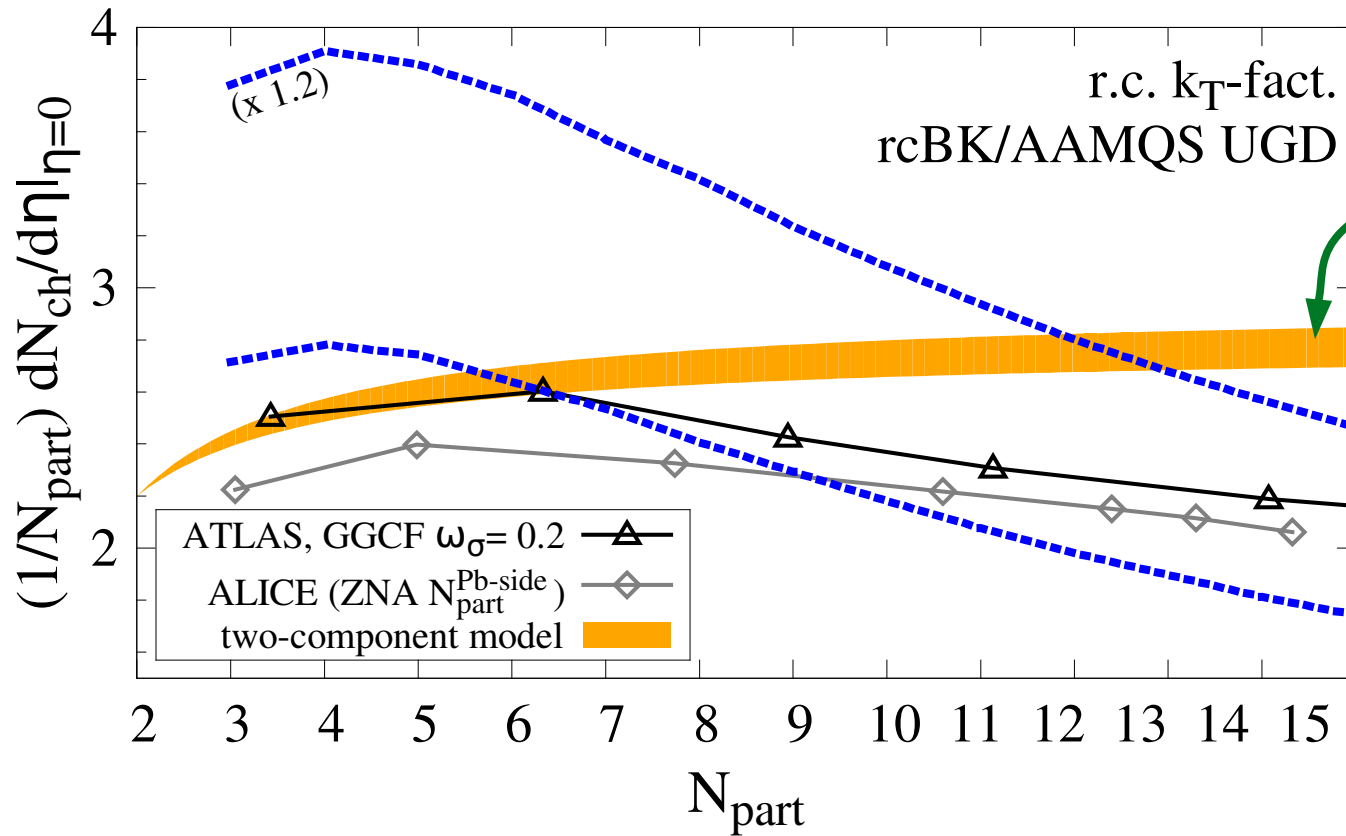
$$f = 0.26 - 0.34$$

Going back to p+Pb...



# Multiplicity vs $N_{part}$ : p + A, two-component model

arXiv:1805.02702



$$f = 0.26 - 0.34$$

**always increasing  
with  $N_{part}$ !**

**For p+A:  $N_{coll} = N_{part} - 1$**

from ALICE: arXiv:1412.6828

$$\frac{1}{N_{part}} \frac{dN_{pA}}{d\eta} = \left[ \frac{1+f}{2} - \frac{f}{N_{part}} \right] \frac{dN_{pp}}{d\eta}$$

# Conclusions

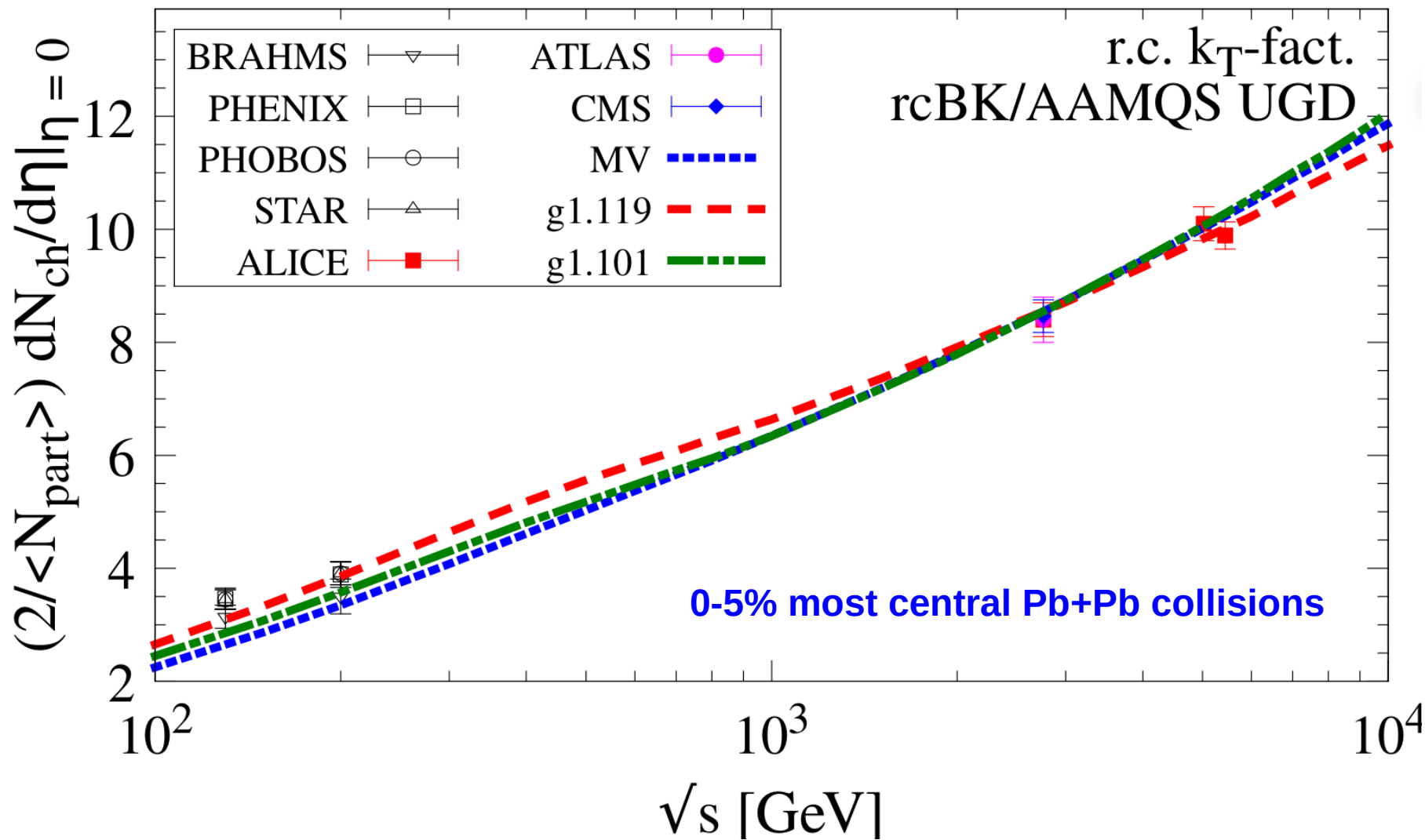
- 1st quantitative comparison of the r.c. kt-fact. with the centrality and energy dependence of particle multiplicities at midrapidity in high-energy p+A and A+A collisions;
- Overall agreement with these observables by adjusting **only one parameter!**
- The CGC framework is in qualitative agreement with the decreasing of the multiplicity per participant with  $N_{part}$  in p+Pb collisions;
- However, exp. data is flatter than our result... Need to know proton UGD better;
- This data is far different from a simple “2-component model” (for  $N_{part} \gtrsim 6$ )
- For p+A: coherent effects from CGC make the multiplicity per participant decrease while the transverse energy per charged particle increases with  $N_{part}$ .

**Thank you for the attention!**

**[and the organizers for the  
opportunity to be here]**

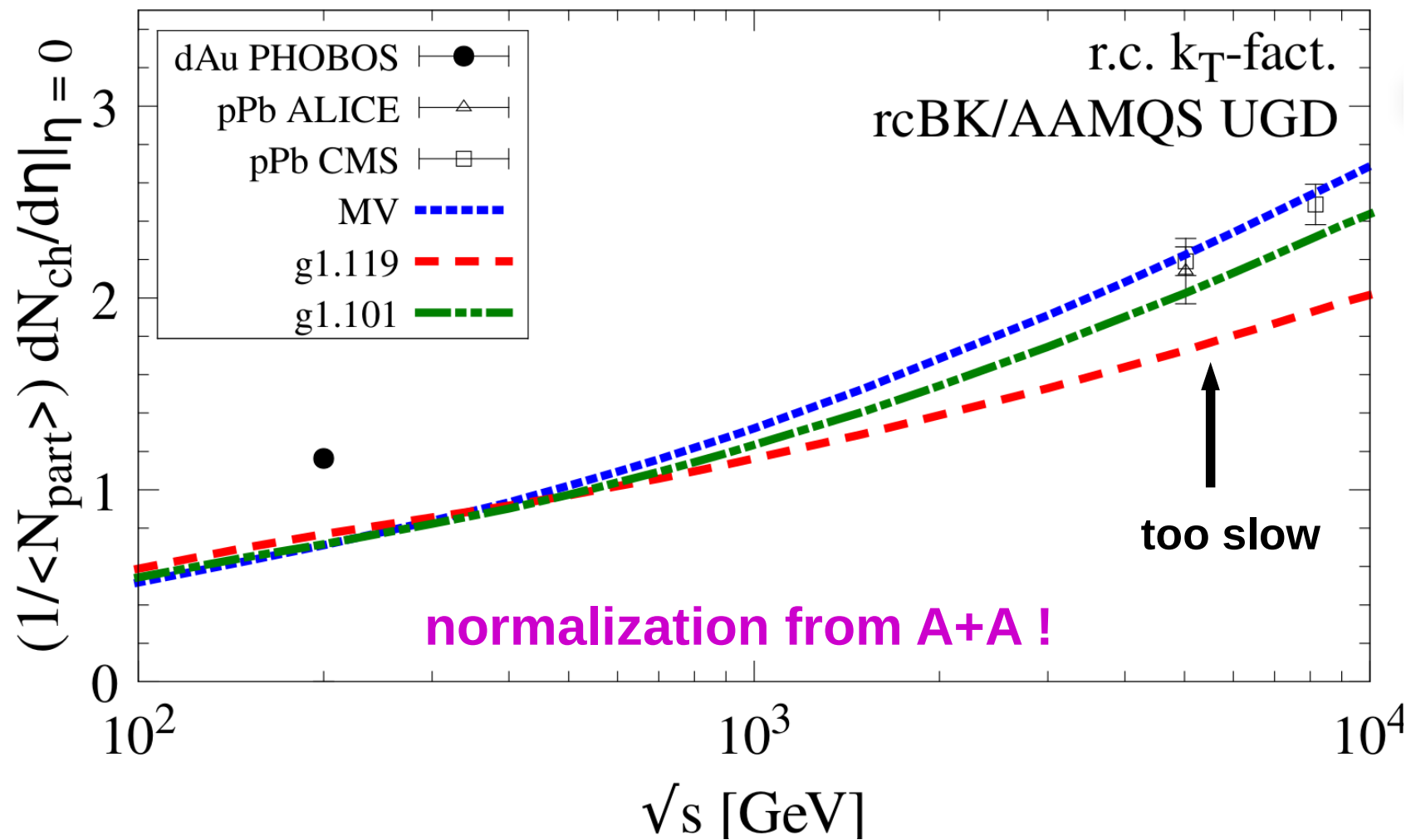
# Backup slides

# Multiplicity vs energy: A+A



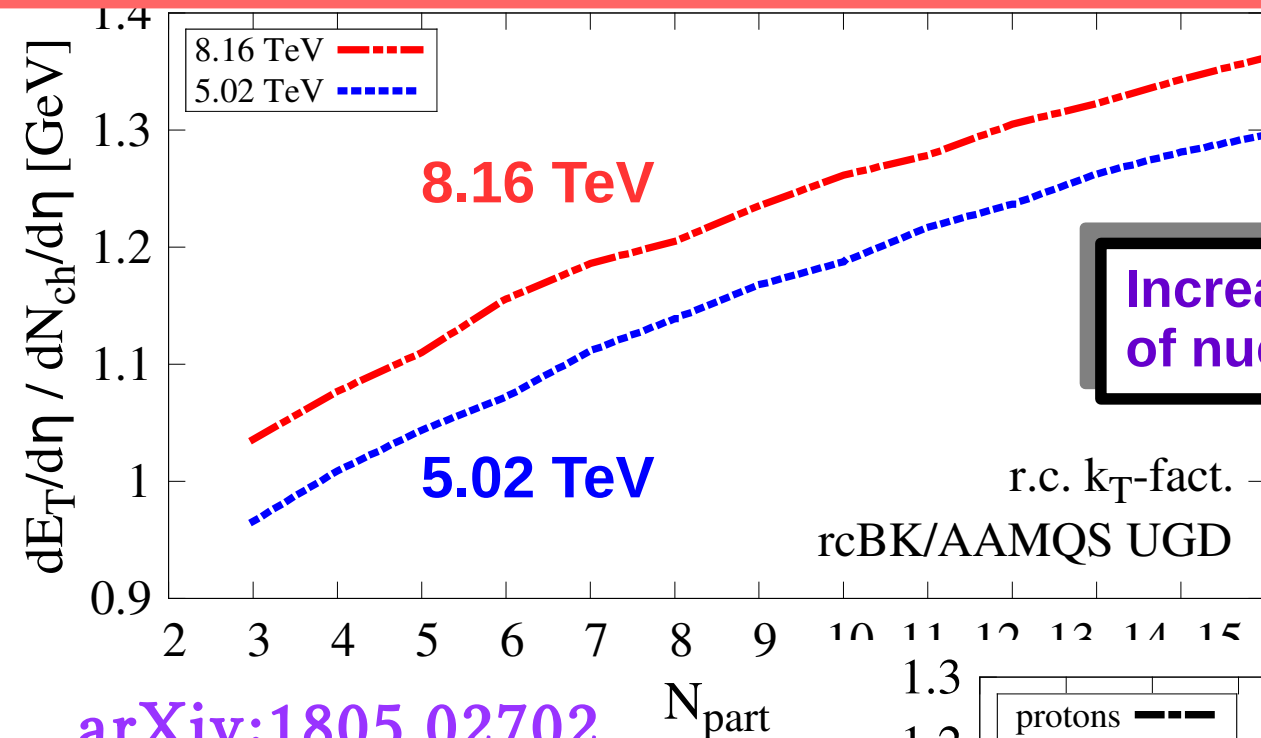
Good description of the energy dependence for all I.C.

# Multiplicity vs energy: p+A



At 200 GeV:  $x \sim 0.01$  and the calculation is most sensitive to the rcBK initial condition rather than the small- $x$  evolution!

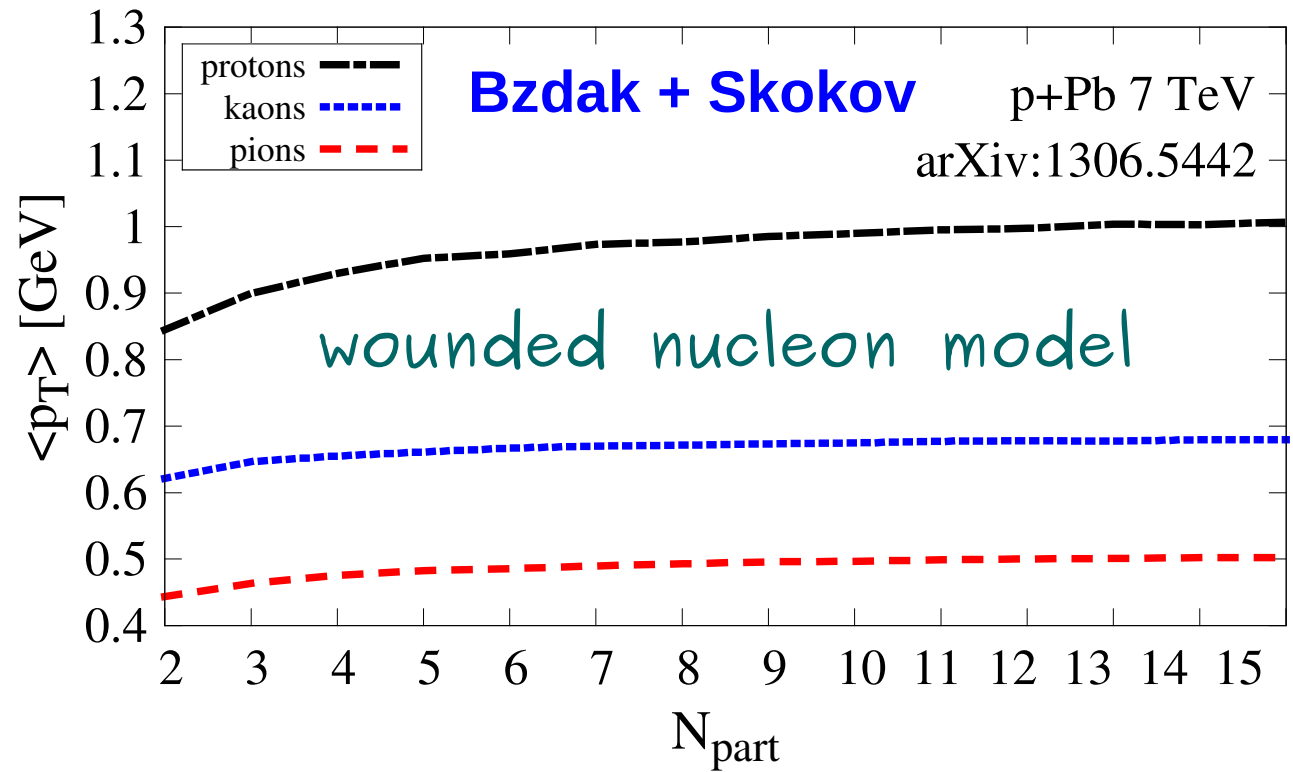
# $[dE_T/d\eta] / [dN_{ch}/d\eta]$ vs $N_{part}$ : p+A



Increase is expected due to increase of nucleus's saturation scale

arXiv:1805.02702

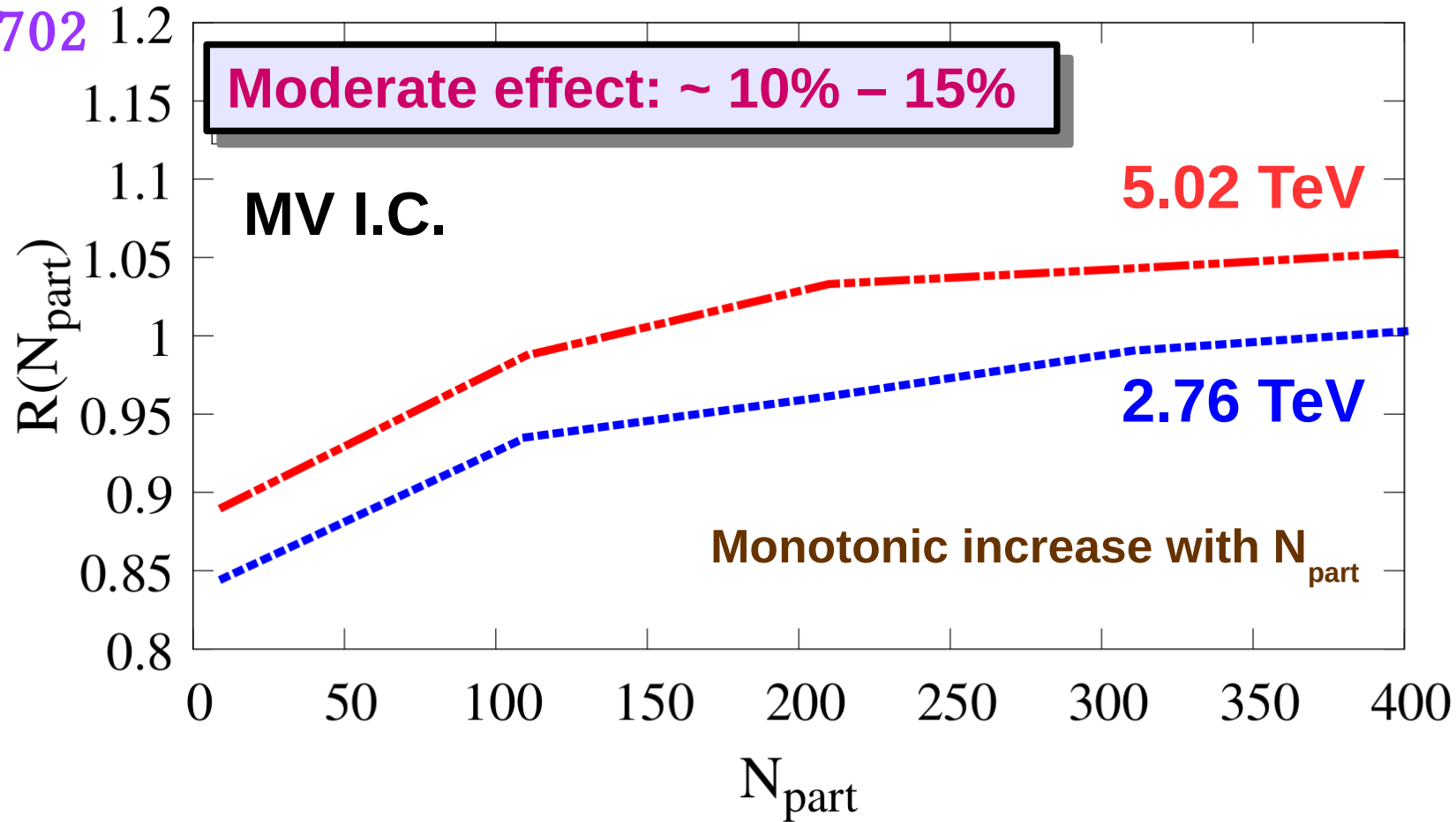
Incoherent model  
 $pA = N_{part} \times pp$



# Quantifying the running coupling effects

$$R(N_{\text{part}}) = \frac{\text{Fixed coupling expression} + \text{running coupling effects by hand}}{\text{Running coupling expression}}$$

arXiv:1805.02702



Similar results for all I.C. & p+Pb @ 5.02, 8.16 TeV