

Hadron Level Event Generation at NLO Accuracy with SHERPA

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30/07/2010



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see also talk by F. Siegert at ICHEP 2010

LO Matrix Elements and Parton Showers

$$d\sigma_{\text{LO}} = d\Phi_B B(\Phi_B) \left[\Delta^{\mathcal{K}}(t_0, t_{\min}) + \int_{t_{\min}}^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \Delta^{\mathcal{K}}(t_0, t) \right]$$

inclusive $2 \rightarrow N$ process
kin. according to $B(\Phi_B)$

$$R^{(\text{PS})}(\Phi_R) = B(\Phi_B) \mathcal{K}(\Phi_{R|B}) = B(\Phi_B) \sum_{ab} P_{ab}(\Phi_{R|B})$$

$$\Delta^{\mathcal{K}}(t_0, t) = \exp \left(- \int_t^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \right)$$

Lo Matrix Elements and Parton Showers

$$d\sigma_{\text{LO}} = d\Phi_B B(\Phi_B) \underbrace{\left[\Delta^{\mathcal{K}}(t_0, t_{\min}) + \int_{t_{\min}}^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \Delta^{\mathcal{K}}(t_0, t) \right]}_{=1}$$

→ cross section unchanged

$$R^{(\text{PS})}(\Phi_R) = B(\Phi_B) \mathcal{K}(\Phi_{R|B}) = B(\Phi_B) \sum_{ab} P_{ab}(\Phi_{R|B})$$

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LO Matrix Elements and Parton Showers

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 exclusive $2 \rightarrow N$ process
 kin. according to $B(\Phi_B)$

approximated real emission ME

$$R^{(\text{PS})}(\Phi_R) = B(\Phi_B) \mathcal{K}(\Phi_{R|B}) = B(\Phi_B) \sum_{ab} P_{ab}(\Phi_{R|B})$$

Sudakov form factor

$$\Delta^{\mathcal{K}}(t_0, t) = \exp \left(- \int_t^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \right)$$

LO Matrix Elements and Parton Showers

$$d\sigma_{\text{LO}} = d\Phi_B B(\Phi_B) \left[\Delta^{\mathcal{K}}(t_0, t_{\min}) + \int_{t_{\min}}^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \Delta^{\mathcal{K}}(t_0, t) \right]$$

↑
 inclusive $2 \rightarrow N+1$ process
 kin. according to $R^{(\text{PS})}(\Phi_R)$

approximated real emission ME

$$R^{(\text{PS})}(\Phi_R) = B(\Phi_B) \mathcal{K}(\Phi_{R|B}) = B(\Phi_B) \sum_{ab} P_{ab}(\Phi_{R|B})$$

Sudakov form factor

$$\Delta^{\mathcal{K}}(t_0, t) = \exp \left(- \int_t^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \right)$$

NLO Matrix Elements and Parton Showers

use POWHEG method → shower independent

$$d\sigma_{\text{NLO}} = d\Phi_B [B(\Phi_B) + V(\Phi_B)] + d\Phi_R R(\Phi_R)$$

$$\bar{B}(\Phi_B) = B(\Phi_B) + V(\Phi_B) + \int d\Phi_{R|B} R(\Phi_R)$$

$$\Delta^{RB}(t_0, t) = \exp \left(- \int_t^{t_0} d\Phi_{R|B} \frac{R(\Phi_R)}{B(\Phi_B)} \right)$$

- NLO accuracy of the ME preserved
- (N)LL accuracy of the PS preserved

NLO Matrix Elements and Parton Showers

use POWHEG method → shower independent

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NLO Matrix Elements and Parton Showers

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$$d\sigma_{\text{NLO}} = d\Phi_B \bar{B}(\Phi_B) \left[\Delta^{RB}(t_0, t_{\min}) + \int_{t_{\min}}^{t_0} d\Phi_{R|B} \frac{R(\Phi_R)}{B(\Phi_B)} \Delta^{RB}(t_0, t) \right]$$

$$\bar{B}(\Phi_B) = B(\Phi_B) + V(\Phi_B) + \int d\Phi_{R|B} R(\Phi_R)$$

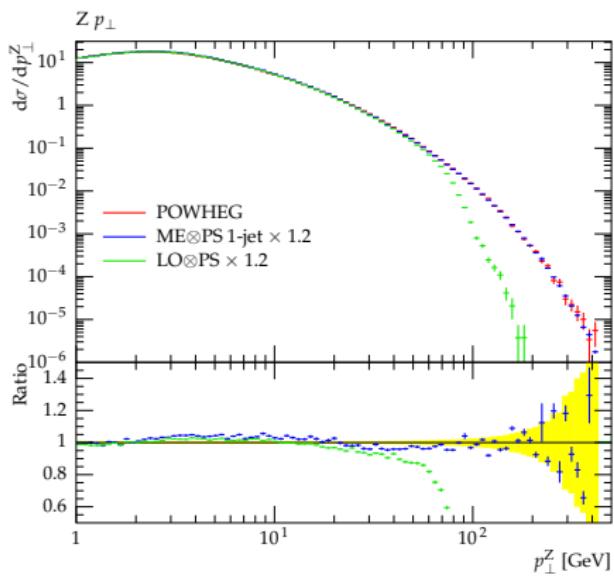
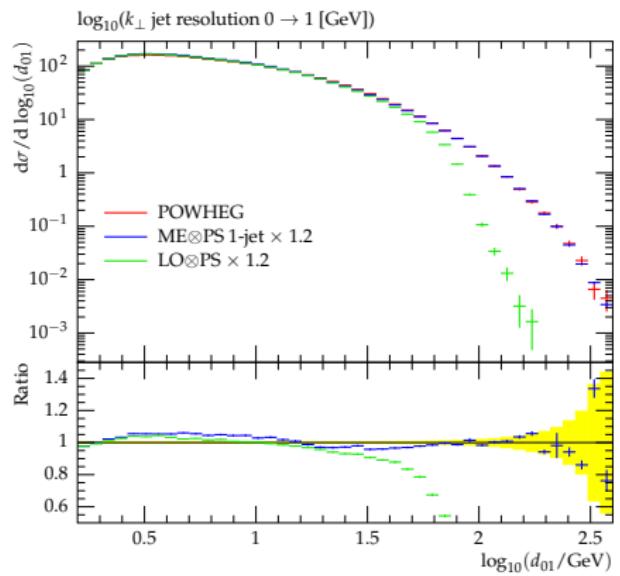
$$\Delta^{RB}(t_0, t) = \exp \left(- \int_t^{t_0} d\Phi_{R|B} \frac{R(\Phi_R)}{B(\Phi_B)} \right)$$

- NLO accuracy of the ME preserved
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Implementation in SHERPA

- automised tree-level ME generators AMEGIC++ and COMIX
[JHEP02\(2002\)044](#), [JHEP12\(2008\)039](#)
- automised Catani-Seymour subtraction term generation
[Eur.Phys.J. C53 \(2008\) 501-523](#)
- Binoth Les Houches Accord Interface for Loop ME
[Comput.Phys.Commun.181:1612-1622,2010](#)
- automised tree-level multi-channel \otimes factorised phase space generator
- PS based on CS-dipole terms CSSOWER++
[JHEP03\(2008\)038](#)

Preliminary Results – $pp \rightarrow Z + X$



Multijet Merging

Multijet merging at Lo

JHEP05(2009)053

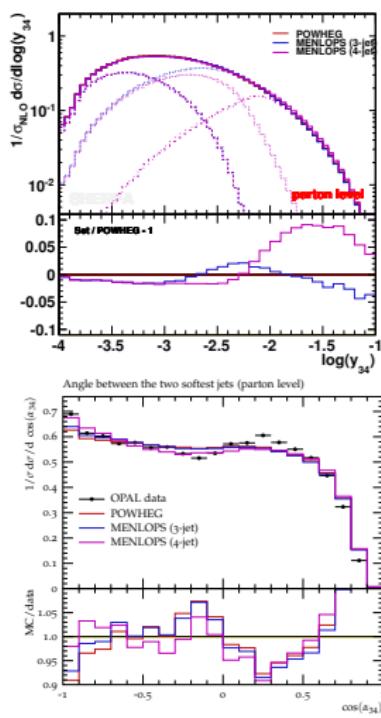
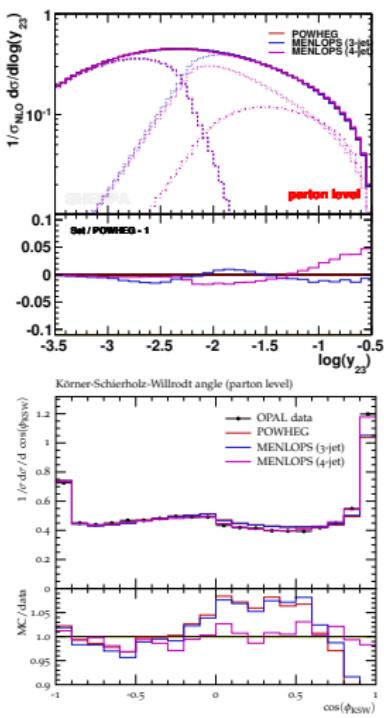
- separate into ME and PS regimes using IR-safe jet criterion Q_{cut}
- reproduces NLO shapes
- Lo total cross section → needs K -factor
- Lo scale uncertainties on multijet observables

Multijet merging at NLO

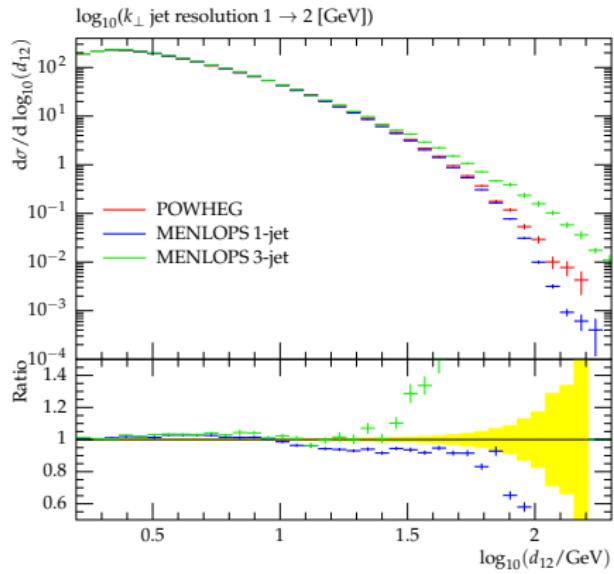
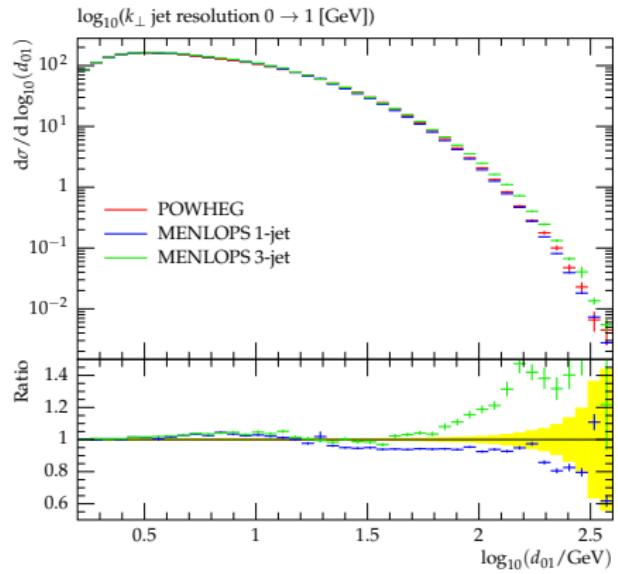
JHEP06(2010)039

- separate into ME and PS regimes as above
- replace $\text{Lo} \otimes \text{PS}$ by $\text{NLO} \otimes \text{PS}$ for lowest multiplicity
- rescale all higher multiplicities by local K -factor $\bar{B}(\Phi_B)/B(\Phi_B)$
- NLO accuracy on total cross section and inclusive observables

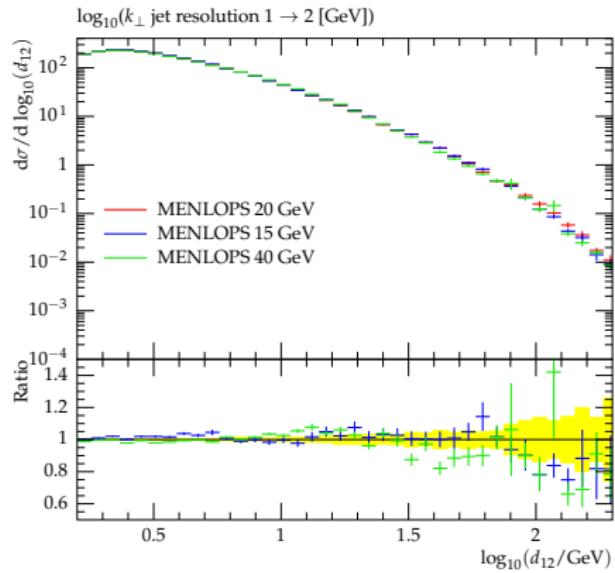
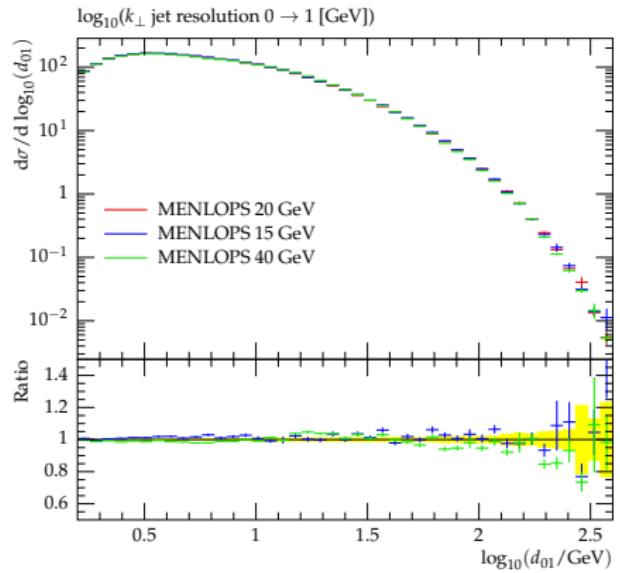
Preliminary Results – $e^+e^- \rightarrow \text{hadrons}$



Preliminary Results – $pp \rightarrow Z + X$



Preliminary Results – $pp \rightarrow Z + X$ – Systematics



Conclusions

- Tree-level merging reproduces shapes well, but needs K-factor
- POWHEG reproduces full NLO cross section and shape of first emission but fails for additional hard radiation
- Combination of full NLO and higher order tree-level MEs with shower achieves both of the above
- Automisation within reach
- Momentarily NLO accuracy only in core process' observables, not in higher jet multiplicities ...
- ... yet

Conclusions

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Multijet Merging at NLO

- factorisability of the Sudakov form factor

$$\Delta(t_1, t_2) = \Delta(t_1, t') \cdot \Delta(t', t_2)$$

- multijet merging at LO
- multijet merging at NLO

$$\begin{aligned} d\sigma_{\text{NLO}} = \bar{B}(\Phi_B) & \left[\Delta^{RB}(t_0, t_{\min}) + \int_{t_{\min}}^{t_0} d\Phi_{R|B} \frac{R(\Phi_R)}{B(\Phi_B)} \Delta^{RB}(t_0, t) \Theta(Q_{\text{cut}} - Q) \right. \\ & \quad \left. + \int_{t_{\min}}^{t_0} d\Phi_{R|B} \frac{R(\Phi_R)}{B(\Phi_B)} \Delta^{\mathcal{K}}(t_0, t) \Theta(Q - Q_{\text{cut}}) \right] \end{aligned}$$

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- multijet merging at LO

$$d\sigma_{\text{LO}} = B(\Phi_B) \left[\Delta^{\mathcal{K}}(t_0, t_{\min}) + \underbrace{\int_{t_{\min}}^{t_0} d\Phi_{R|B} \mathcal{K}(\Phi_{R|B}) \Delta^{\mathcal{K}}(t_0, t) \Theta(Q_{\text{cut}} - Q)}_{\text{PS regime}} \right. \\ \left. + \underbrace{\int_{t_{\min}}^{t_0} d\Phi_{R|B} \frac{R(\Phi_R)}{B(\Phi_B)} \Delta^{\mathcal{K}}(t_0, t) \Theta(Q - Q_{\text{cut}})}_{\text{ME regime}} \right]$$

→ [...] ≈ 1 to (N)LL accuracy only

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