

Workshop: Low- x gluon structure of nuclei
and signals of saturation at LHC
CERN, March 27th 2018

Lessons from EIC: how to determine the gluon density

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EXCELENCIA
MARÍA
DE MAEZTU



XUNTA
DE GALICIA

Contents:

I. Introduction.

2. Present status of nPDFs:

- Available sets.
- Further constrains from the LHC.

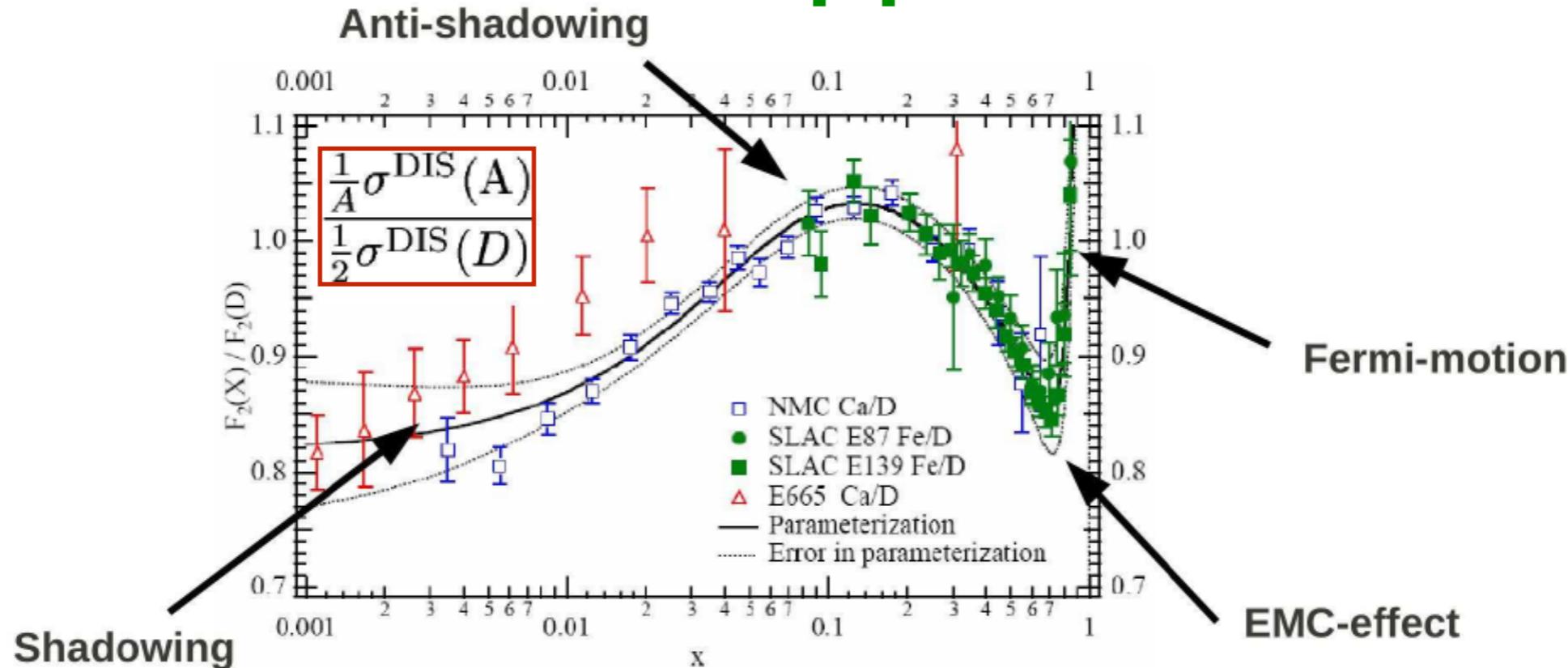
3. Nuclear PDFs from EICs:

- Kinematics.
- The method.
- Constraints on nPDFs.

4. Summary.

See the talks by Vadim Guzey, Juan Rojo and Ilkka Helenius.

Collinear approach:



- Bound nucleon \neq free nucleon: search for process independent nPDFs that realise this condition, assuming collinear factorisation.

$$\sigma_{\text{DIS}}^{\ell+A \rightarrow \ell+X} = \sum_{i=q,\bar{q},g} f_i^A(\mu^2) \otimes \hat{\sigma}_{\text{DIS}}^{\ell+i \rightarrow \ell+X}(\mu^2)$$

Nuclear PDFs, obeying
the standard DGLAP Usual perturbative
coefficient functions

$$f_i^{p,A}(x, Q^2) = R_i^A(x, Q^2) f_i^p(x, Q^2)$$

$$R = \frac{f_i/A}{Af_i/p} \approx \frac{\text{measured}}{\text{expected if no nuclear effects}}$$

Collinear approach:

Anti-shadowing

- At an ep/eA collider:
 - PDF of a single nucleus possible, no need of ratios that would be obtained a posteriori.
 - Same method of extraction in both ep and eA.
 - Physics beyond standard collinear factorisation can be studied in a single setup, with size effects disentangled from energy effects and a large lever arm in x at perturbative Q^2 .
- Bound nucleon \neq free nucleon: search for process independent nPDFs that realise this condition, assuming collinear factorisation.

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Nuclear PDFs, obeying
the standard DGLAP

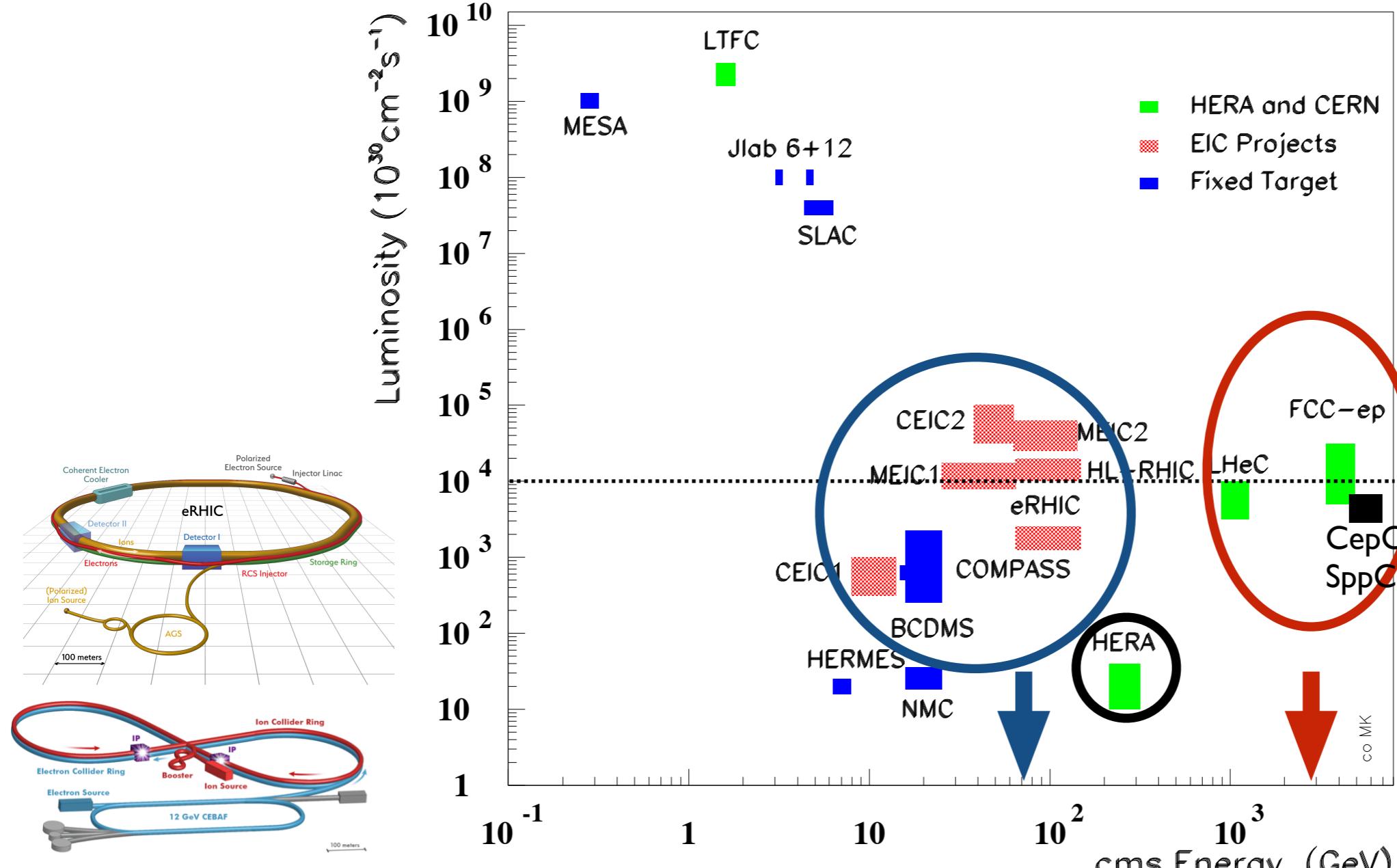
Usual perturbative
coefficient functions

$$f_i^{p,A}(x, Q^2) = R_i^A(x, Q^2) f_i^p(x, Q^2)$$

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Summary of machines:

Lepton-proton/nucleus scattering facilities



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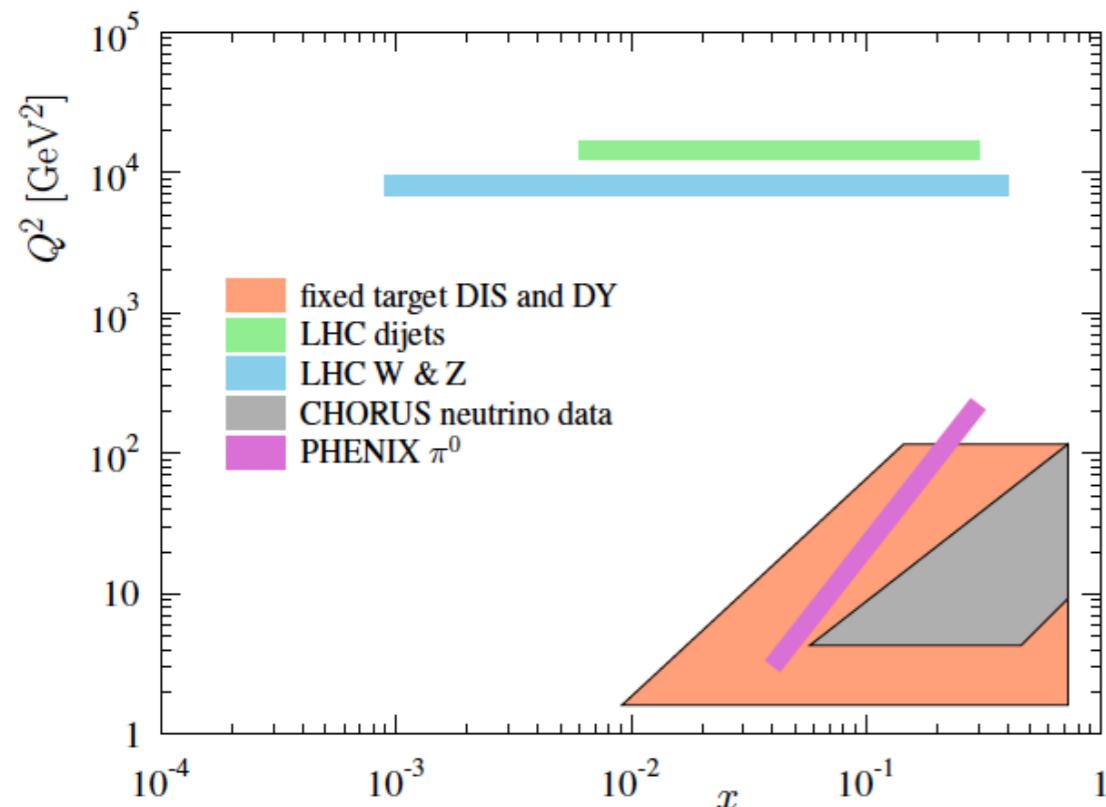
Available sets:

SET		HKN07 PRC76 (2007) 065207	EPS09 JHEP 0904 (2009) 065	DSSZ PRD85 (2012) 074028	nCTEQ15 PRD93 (2016) 085037	KA15 PRD93 (2016) 014036	EPPS16 EPJC C77 (2017)163
data	eDIS	✓	✓	✓	✓	✓	✓
	DY	✓	✓	✓	✓	✓	✓
	π^0	✗	✓	✓	✓	✗	✓
	vDIS	✗	✗	✓	✗	✗	✓
	pPb	✗	✗	✗	✗	✗	✓
# data	1241	929	1579	740	1479	1811	
order	NLO	NLO	NLO	NLO	NNLO	NLO	
proton PDF	MRST98	CTEQ6.I	MSTW2008	~CTEQ6.I	JR09	CTI4NLO	
mass scheme	ZM-VFNS	ZM-VFNS	GM-VFNS	GM-VFNS	ZM-VFNS	GM-VFNS	
comments	$\Delta\chi^2=13.7$, ratios, <u>no EMC for gluons</u>	$\Delta\chi^2=50$, ratios, <u>huge shadowing-antishadowing</u>	$\Delta\chi^2=30$, ratios, <u>medium-modified FFs for π^0</u>	$\Delta\chi^2=35$, PDFs, valence <u>flavour sep.</u> , <u>not enough sensitivity</u>	PDFs, <u>deuteron</u> <u>sep., not enough sensitivity</u>	$\Delta\chi^2=52$ flavour sep., ratios, LHC pPb data	

Available sets:

SET	HKN07 PRC76 (2007)	EPS09 JHEP 0904	DSSZ PRD85 (2012) 074028	nCTEQ15 PRD93 (2016) 085037	KA15 PRD93 (2016) 014036	EPPS16 EPJC C77 (2017)163
<ul style="list-style-type: none"> • Centrality dependence (EPS09s) • not from data but from the A-dependence of the parameters. • Several models provide it: Vogt et al., FGS, Ferreiro et al.,... 						
PDF	MRST98	CTEQ6.1	1579	740	1479	1811
mass scheme	ZM-VFNS	ZM-VFNS	NLO	NLO	NNLO	NLO
comments	$\Delta\chi^2=13.7$, ratios, <u>no EMC for gluons</u>	$\Delta\chi^2=50$, ratios, <u>huge shadowing-antishadowing</u>	$\Delta\chi^2=30$, ratios, <u>medium-modified FFs for π^0</u>	$\Delta\chi^2=35$, PDFs, valence <u>flavour sep.</u> , not enough sensitivity	PDFs, <u>deuteron</u> data included	$\Delta\chi^2=52$ flavour sep., ratios, LHC pPb data

EPPS16:

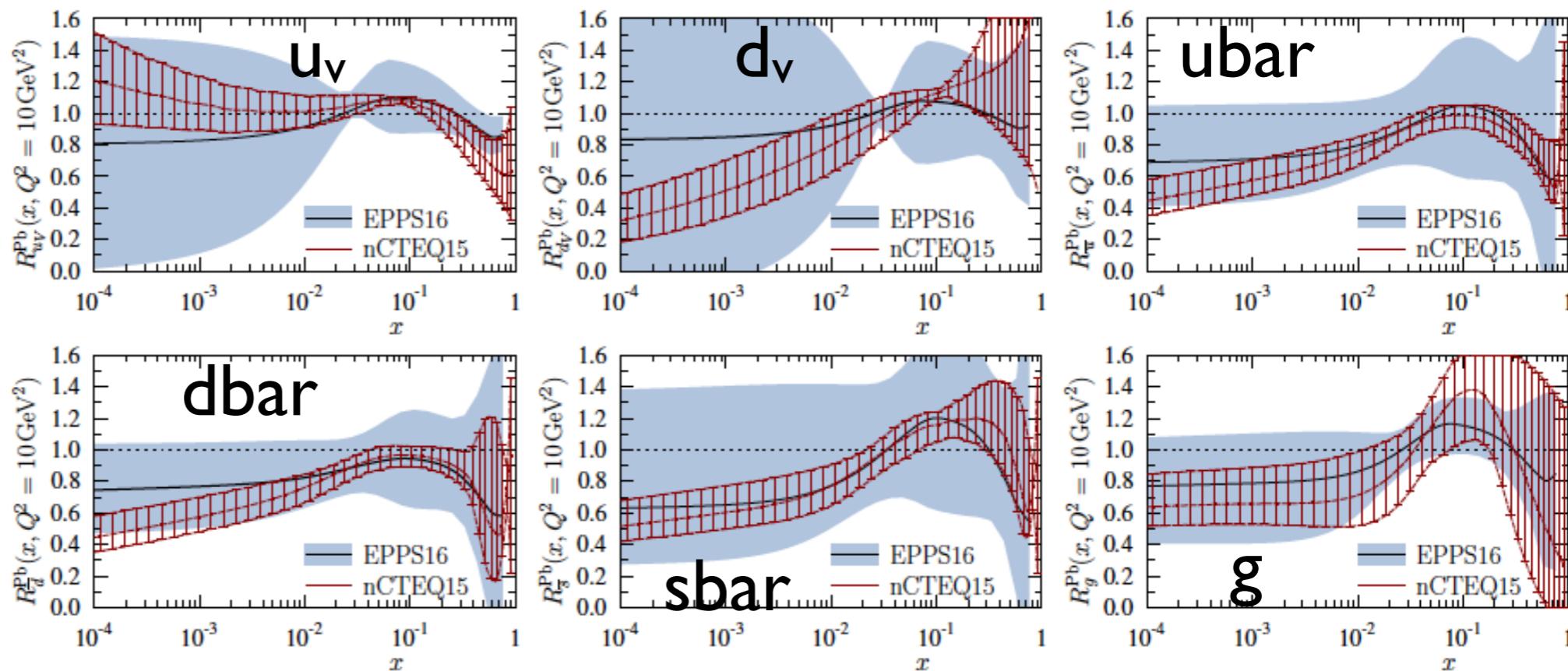


- Most Pb data from CHORUS, 30 Pb points from pPb@LHC: fit for a single nucleus not possible.

Experiment	Observable	Collisions	Data points	χ^2	Ref.
SLAC E139	DIS	e^- He(4), e^- D	21	12.2	[69]
CERN NMC 95, re.	DIS	μ^- He(4), μ^- D	16	18.0	[70]
CERN NMC 95	DIS	μ^- Li(6), μ^- D	15	18.4	[71]
CERN NMC 95, Q^2 dep.	DIS	μ^- Li(6), μ^- D	153	161.2	[71]
SLAC E139	DIS	e^- Be(9), e^- D	20	12.9	[69]
CERN NMC 96	DIS	μ^- Be(9), μ^- C	15	4.4	[72]
SLAC E139	DIS	e^- C(12), e^- D	7	6.4	[69]
CERN NMC 95	DIS	μ^- C(12), μ^- D	15	9.0	[71]
CERN NMC 95, Q^2 dep.	DIS	μ^- C(12), μ^- D	165	133.6	[71]
CERN NMC 95, re.	DIS	μ^- C(12), μ^- D	16	16.7	[70]
CERN NMC 95, re.	DIS	μ^- C(12), μ^- Li(6)	20	27.9	[70]
FNAL E772	DY	pC(12), pD	9	11.3	[73]
SLAC E139	DIS	e^- Al(27), e^- D	20	13.7	[69]
CERN NMC 96	DIS	μ^- Al(27), μ^- C(12)	15	5.6	[72]
SLAC E139	DIS	e^- Ca(40), e^- D	7	4.8	[69]
FNAL E772	DY	pCa(40), pD	9	3.33	[73]
CERN NMC 95, re.	DIS	μ^- Ca(40), μ^- D	15	27.6	[70]
CERN NMC 95, re.	DIS	μ^- Ca(40), μ^- Li(6)	20	19.5	[70]
CERN NMC 96	DIS	μ^- Ca(40), μ^- C(12)	15	6.4	[72]
SLAC E139	DIS	e^- Fe(56), e^- D	26	22.6	[69]
FNAL E772	DY	e^- Fe(56), e^- D	9	3.0	[73]
CERN NMC 96	DIS	μ^- Fe(56), μ^- C(12)	15	10.8	[72]
FNAL E866	DY	pFe(56), pBe(9)	28	20.1	[74]
CERN EMC	DIS	μ^- Cu(64), μ^- D	19	15.4	[75]
SLAC E139	DIS	e^- Ag(108), e^- D	7	8.0	[69]
CERN NMC 96	DIS	μ^- Sn(117), μ^- C(12)	15	12.5	[72]
CERN NMC 96, Q^2 dep.	DIS	μ^- Sn(117), μ^- C(12)	144	87.6	[76]
FNAL E772	DY	pW(184), pD	9	7.2	[73]
FNAL E866	DY	pW(184), pBe(9)	28	26.1	[74]
CERN NA10★	DY	π^- W(184), π^- D	10	11.6	[49]
FNAL E615★	DY	π^+ W(184), π^- W(184)	11	10.2	[50]
CERN NA3★	DY	π^- Pt(195), π^- H	7	4.6	[48]
SLAC E139	DIS	e^- Au(197), e^- D	21	8.4	[69]
RHIC PHENIX	π^0	dAu(197), pp	20	6.9	[28]
CERN NMC 96	DIS	μ^- Pb(207), μ^- C(12)	15	4.1	[72]
CERN CMS★	W±	pPb(208)	10	8.8	[43]
CERN CMS★	Z	pPb(208)	6	5.8	[45]
CERN ATLAS★	Z	pPb(208)	7	9.6	[46]
CERN CMS★	dijet	pPb(208)	7	5.5	[34]
CERN CHORUS★	DIS	ν Pb(208), $\bar{\nu}$ Pb(208)	824	998.6	[47]
Total			1811	1789	

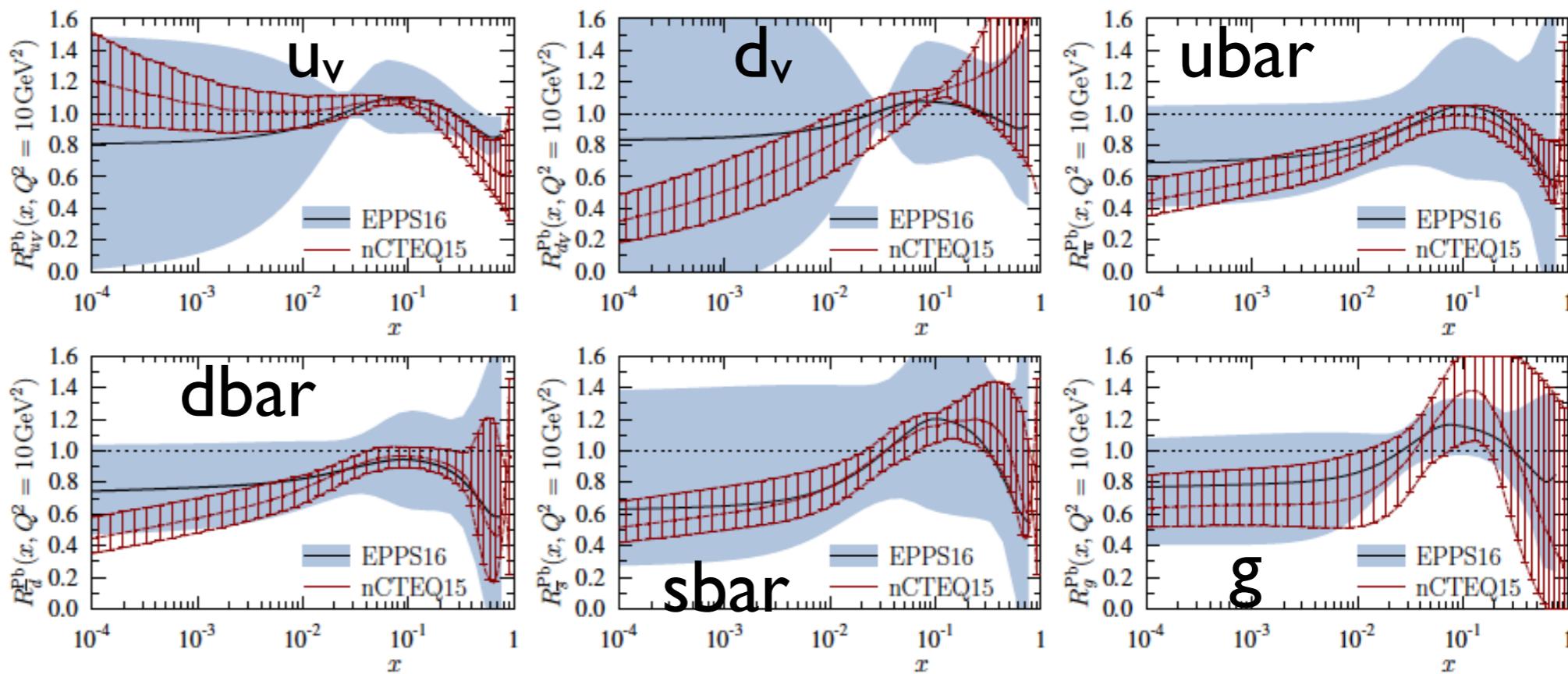
EPPS16:

- nCTEQ15 vs. EPPS16: note the parametrisation bias.

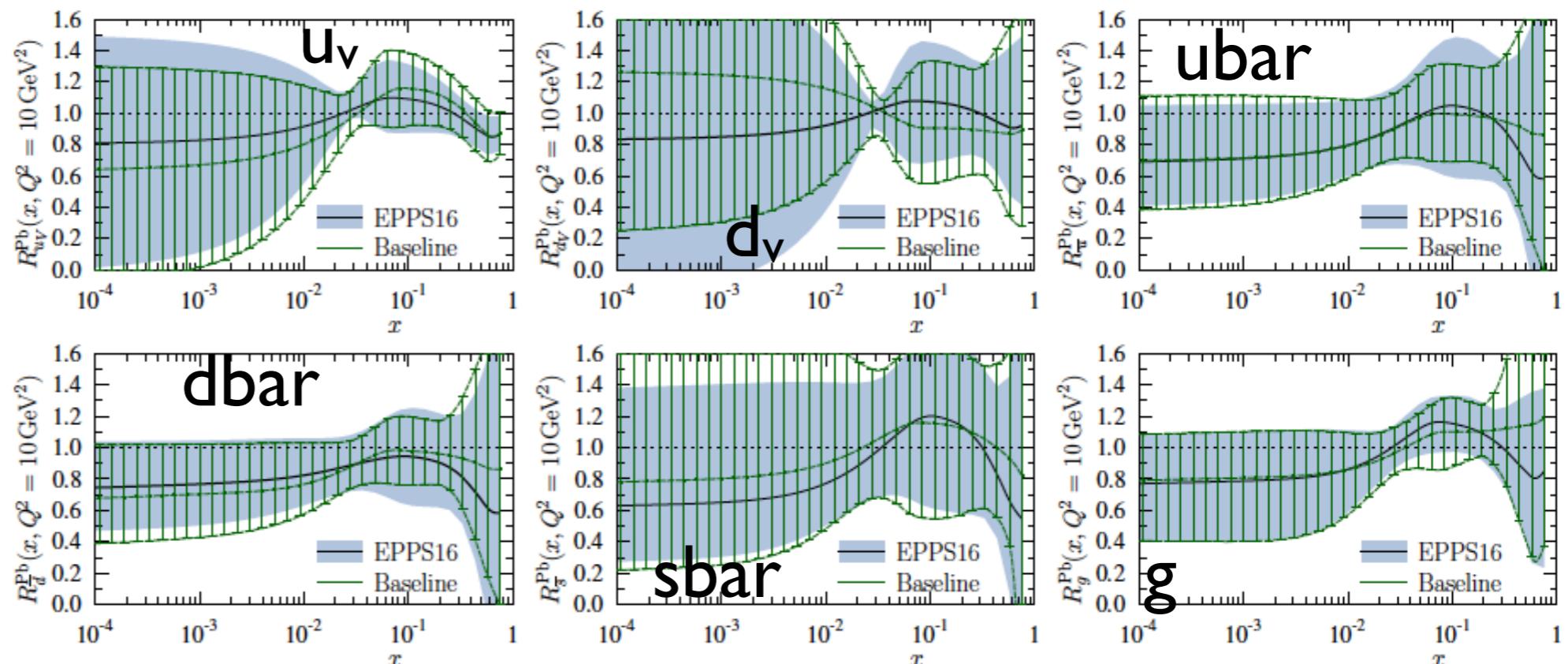


EPPS16:

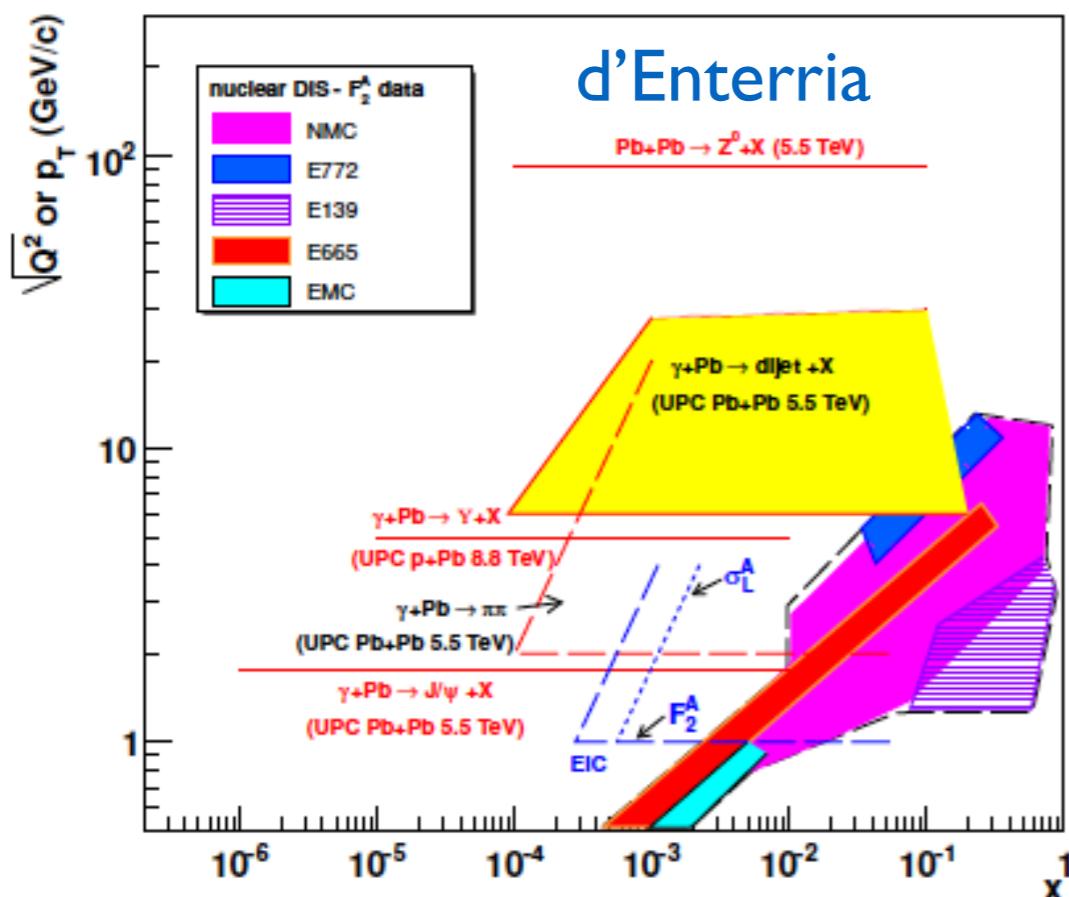
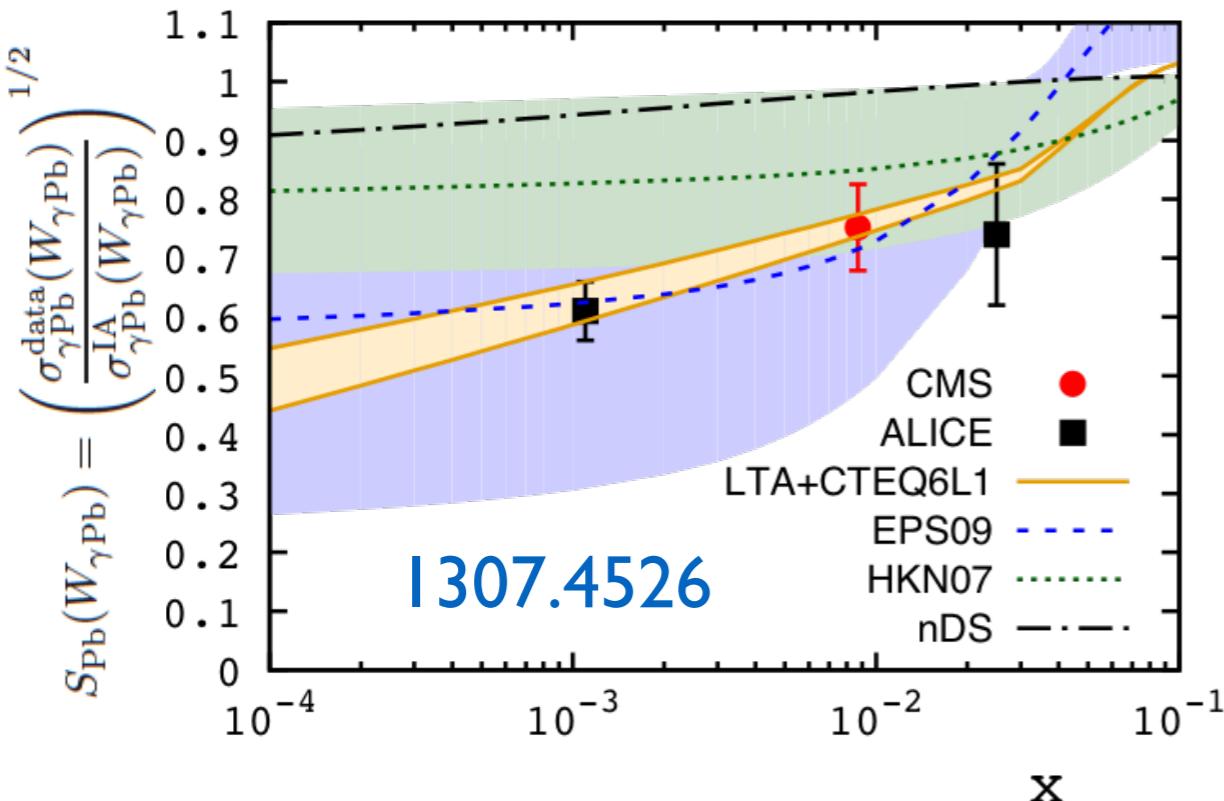
- nCTEQ15 vs. EPPS16: note the parametrisation bias.



- Presently available LHC data seem not to have a large effect: large-x glue (baseline=no v, no LHC data).



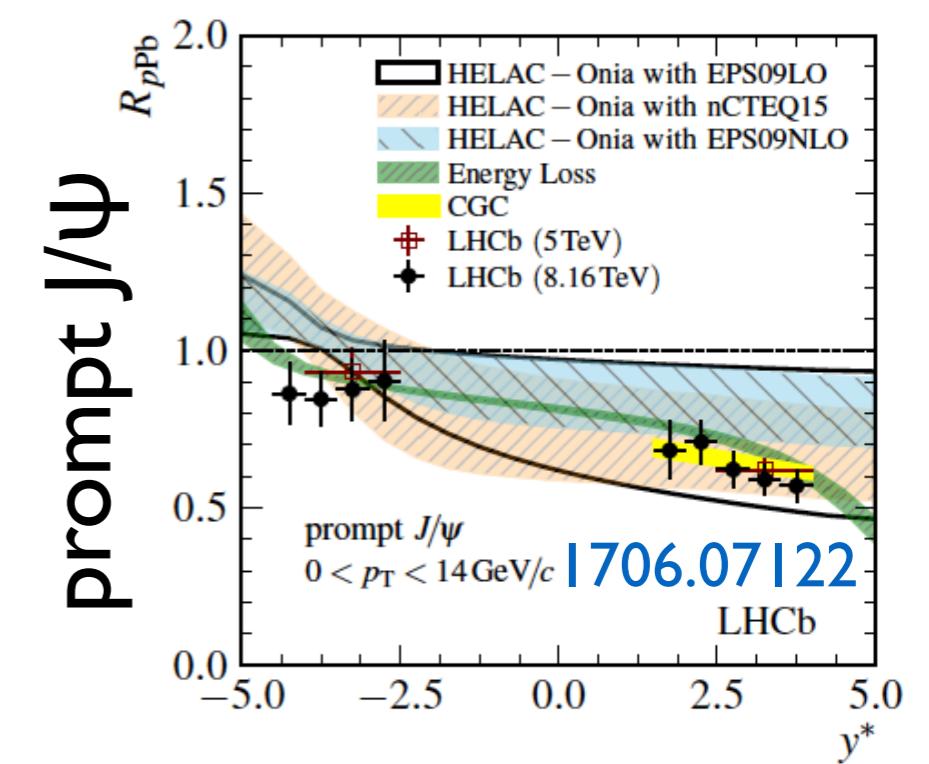
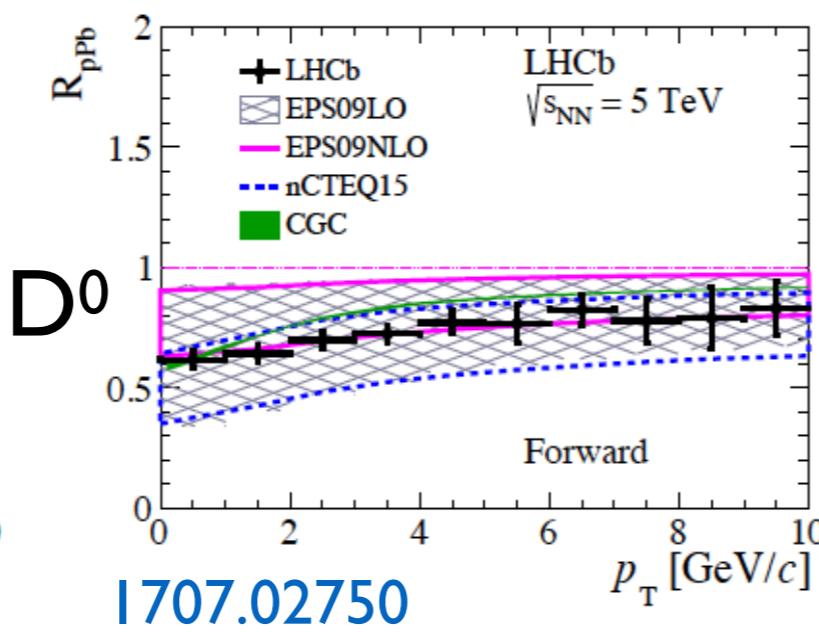
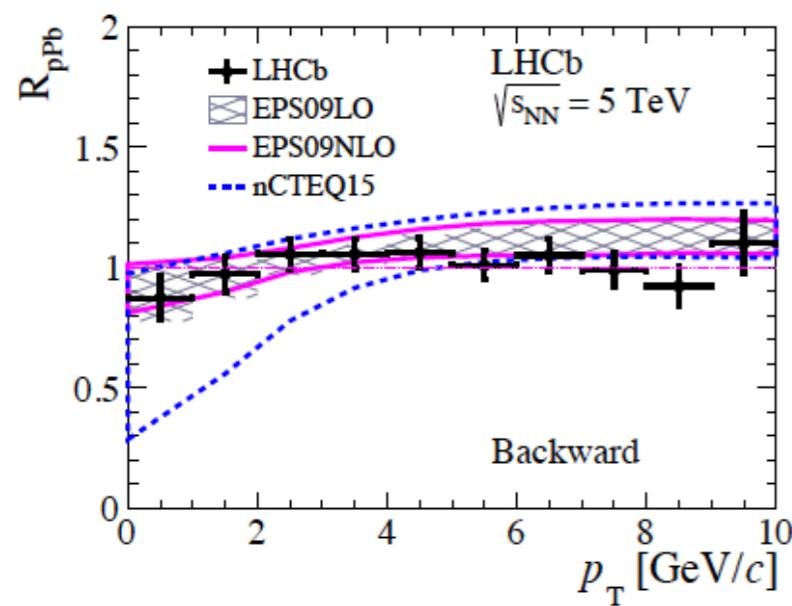
UPCs:



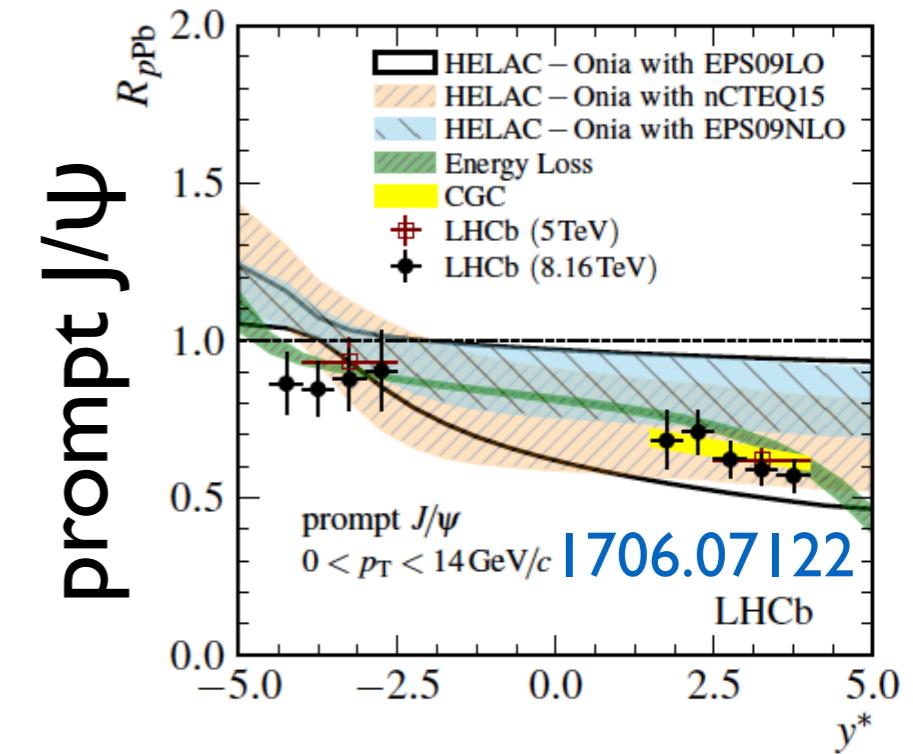
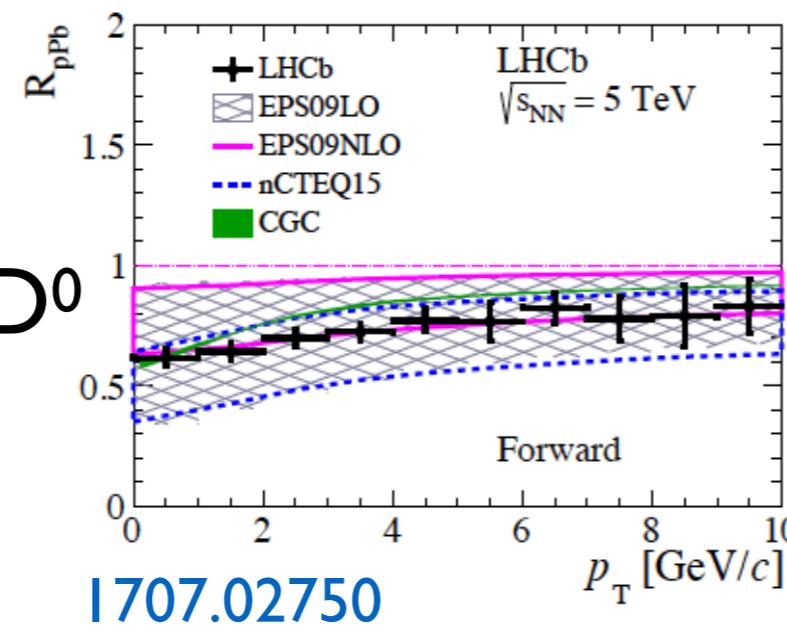
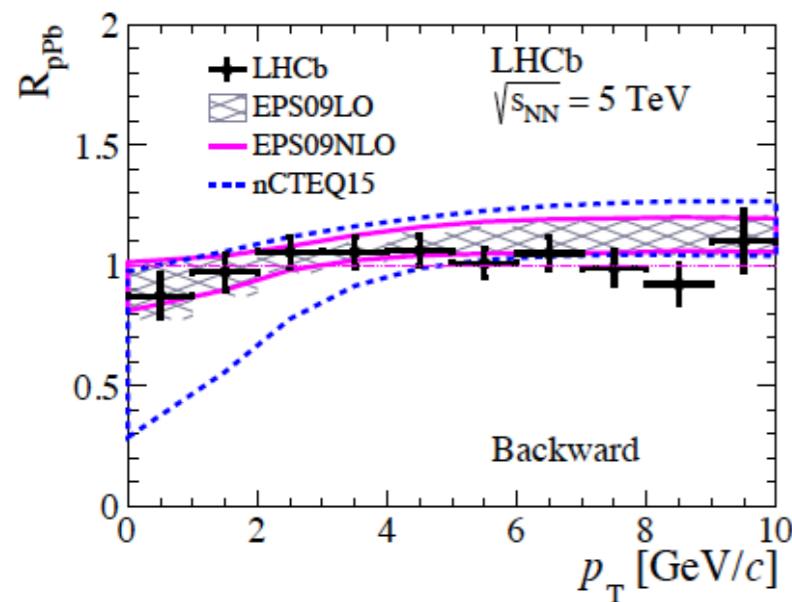
- UPCs offer possibilities for constraining both nPDFs: they were the first indication of nuclear shadowing.

- Uncertainties on the precision and applicability of standard collinear factorisation exist for many of the processes currently studied e.g. exclusive VM production: work still required for this data to be included in global fits.

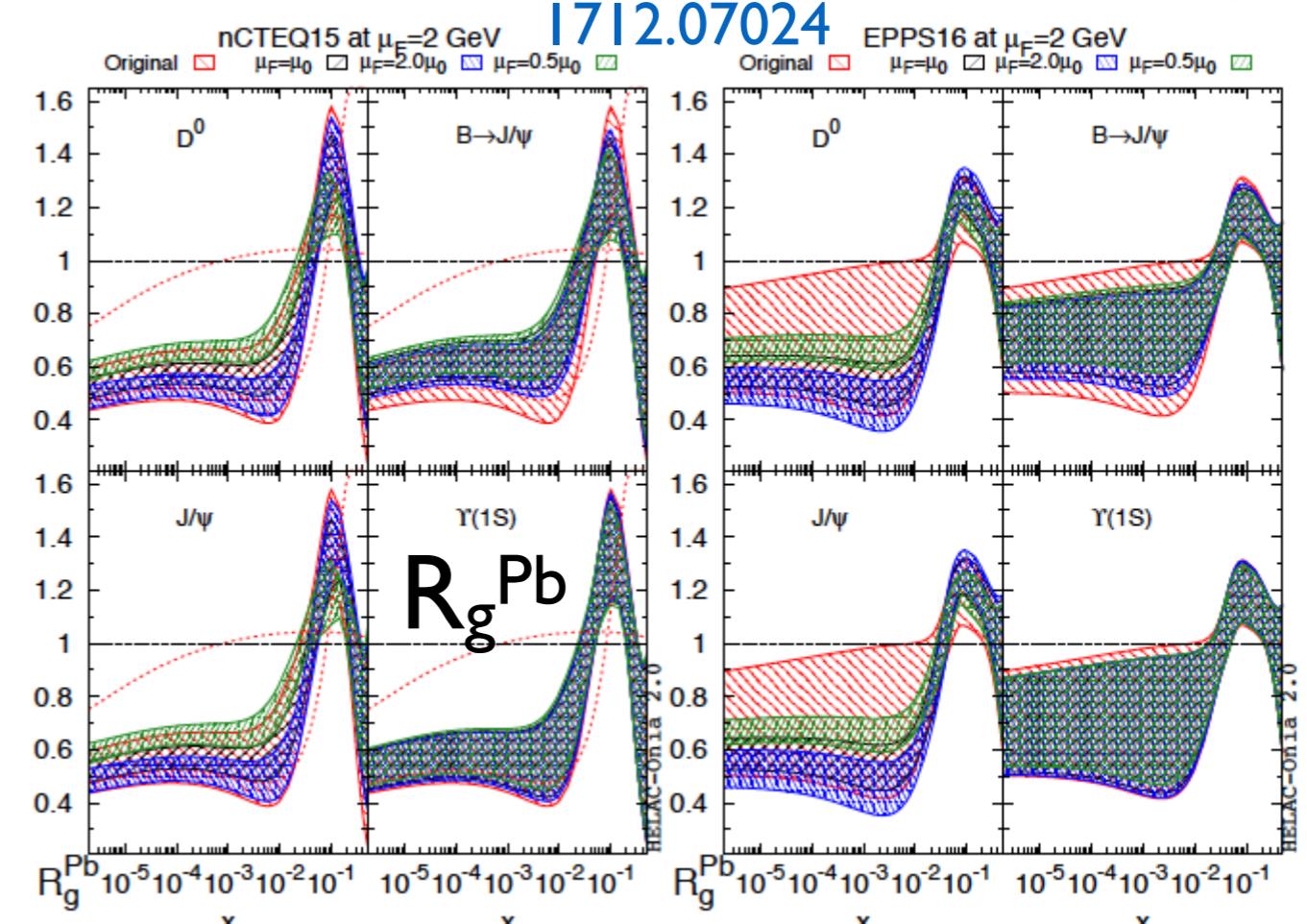
Open charm and J/ ψ :



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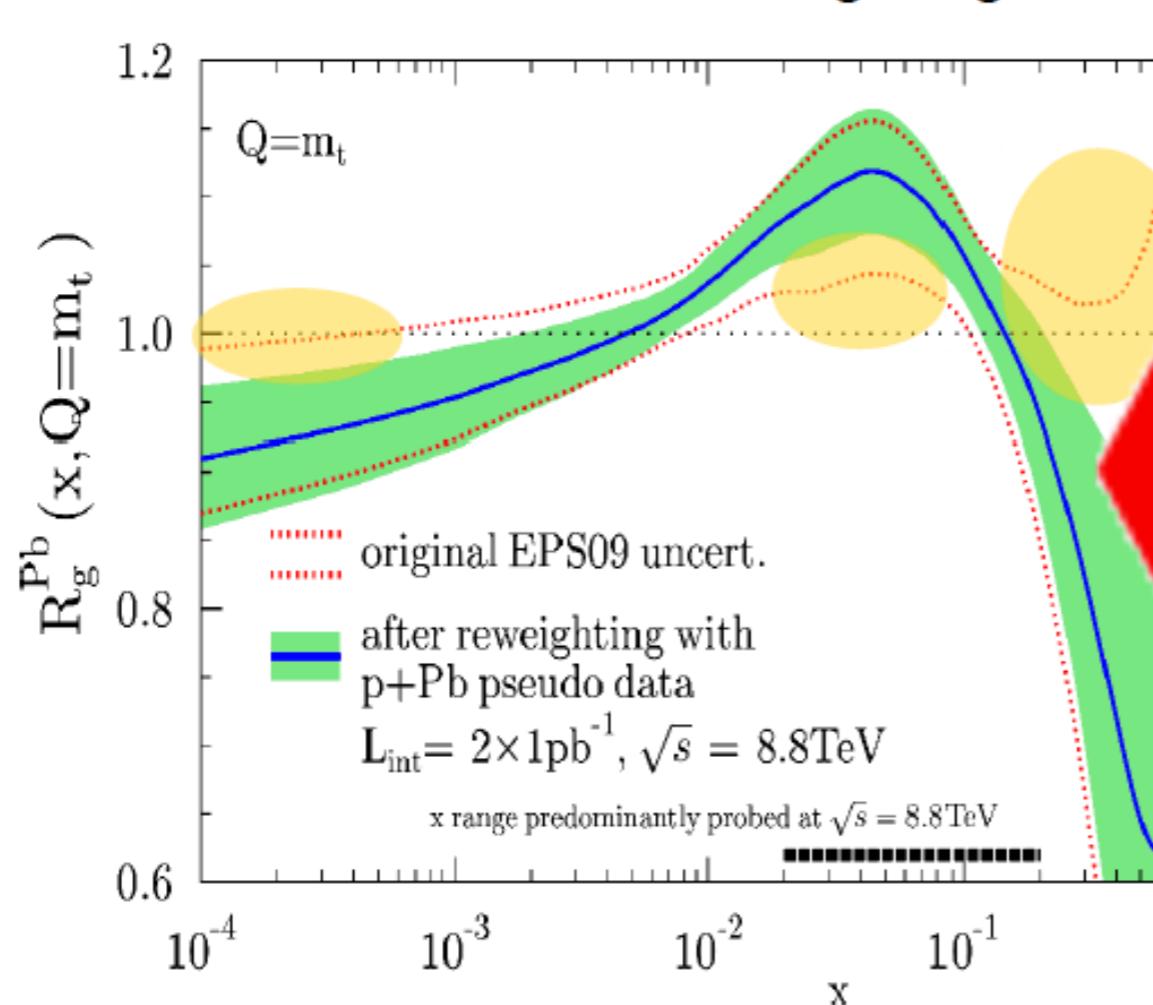
- Theoretical control in PT over forward D or J/ ψ under debate, even in pp: **scales**, DPS, non-linear dynamics, ..., see 1610.09373 vs. 1710.05935.
- E.g. quarkonium: superposition of nPDFs + eloss/absorption + comovers for ψ' + ...
- Collectivity (flow for D in pPb as for charged hadrons in pPb and PbPb?) would limit the use of low p_T data for extraction of nPDFs.



Tops at HL-LHC:

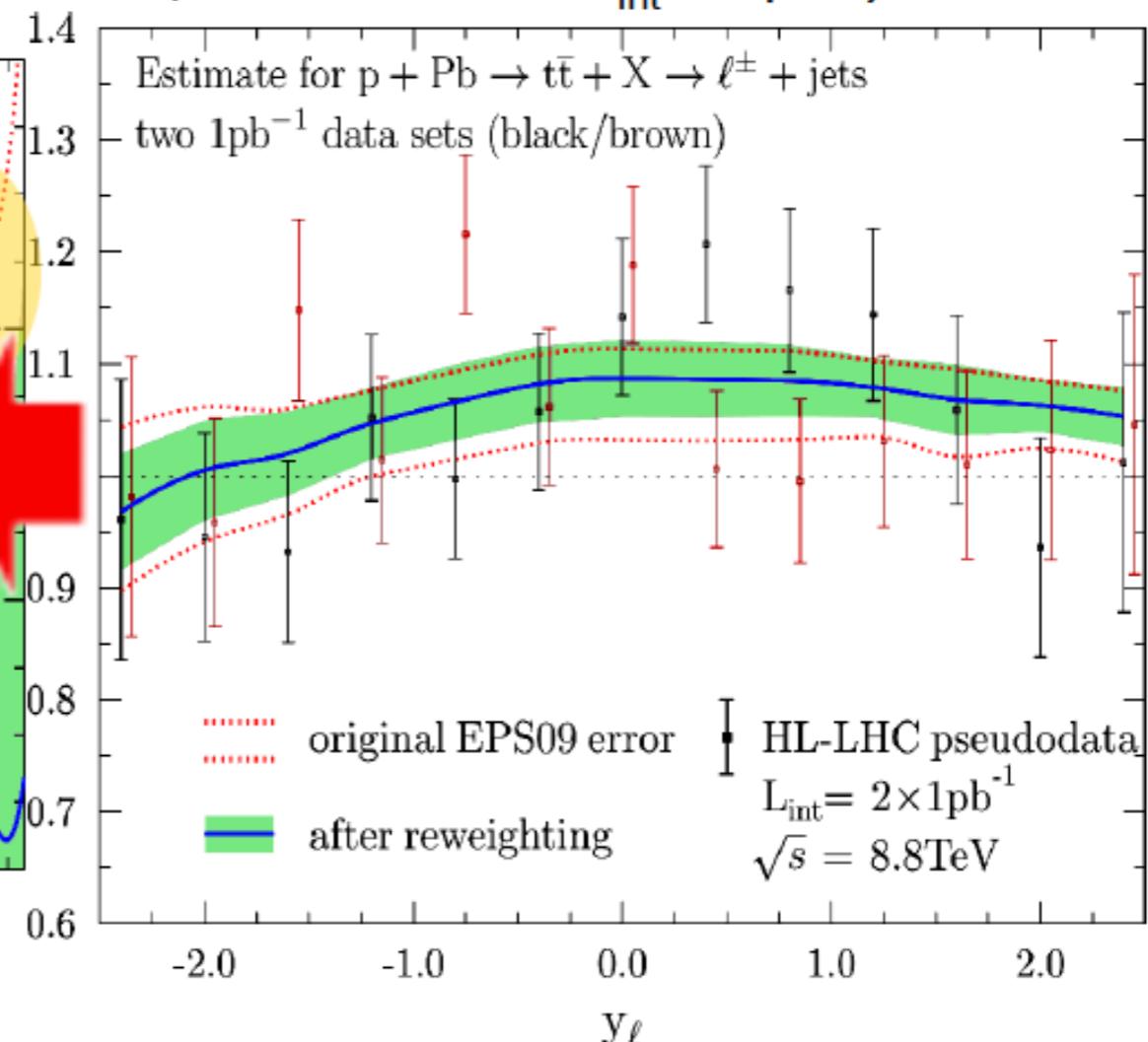
pPb \rightarrow ttbar+X (8.8 TeV): Gluon density constraints

- Improved gluon density via Hessian PDF reweighting



- ~50% reduction in uncertainties at antishadowing ($x \sim 0.05$) and EMC ($x \sim 0.4$) regions.

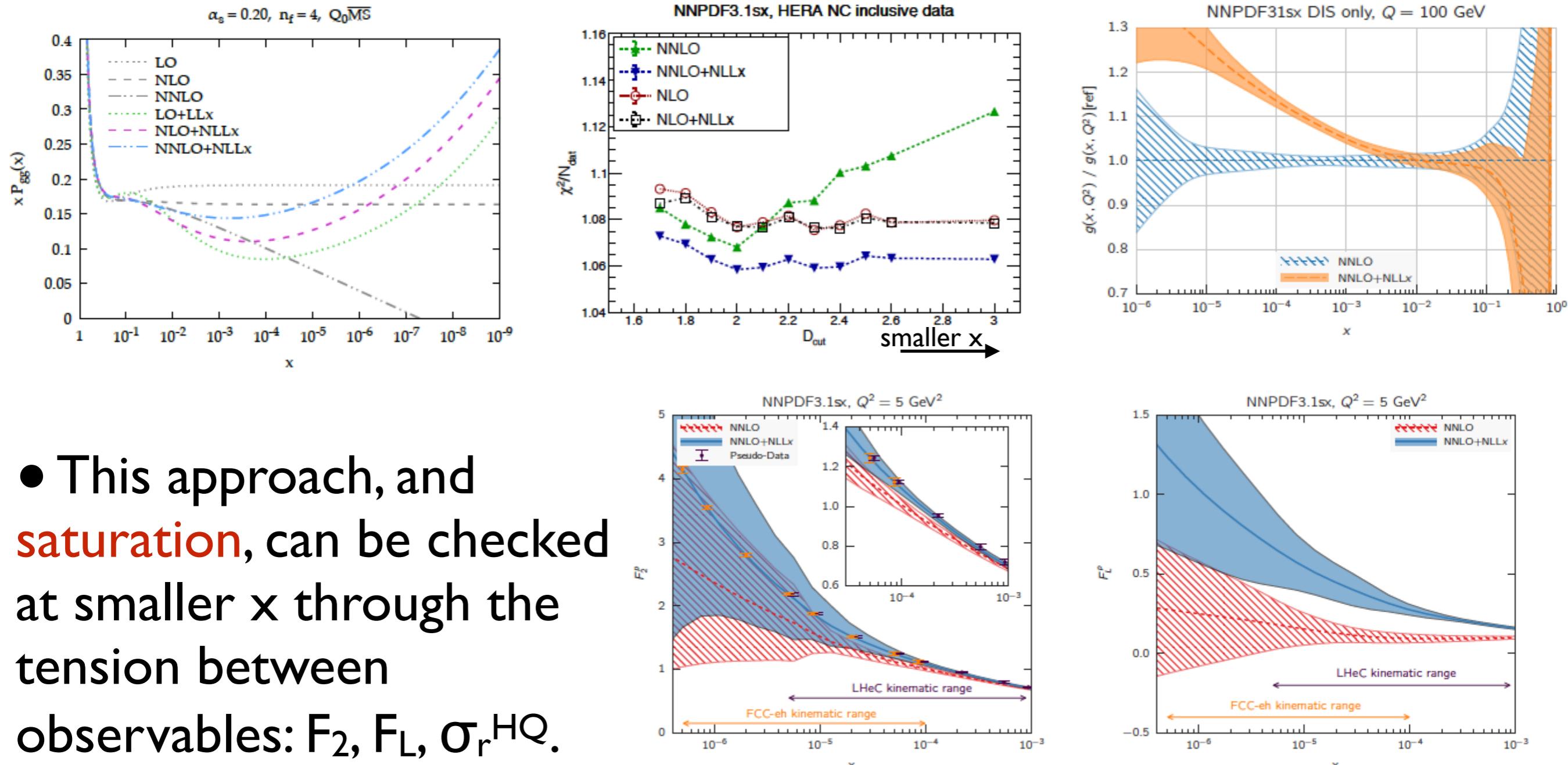
- Isolated lepton y -distrib. after cuts:
(Pseudodata for $L_{\text{int}} = 1 \text{ pb}^{-1}$)



- (5% syst. uncertainties dominate now)
- nPDF R_{pPb} effects (lepton): $\pm 10\%$
- $L_{\text{int}} = 2 \times 1 \text{ pb}^{-1}$: ~35% constraining power

Resummation:

- Resummation has been suggested ([I710.05935](#)) to cure the problem seen in HERA data of a worsening of the PDF fit quality with decreasing x and Q^2 : the problem lies in F_L .



- This approach, and **saturation**, can be checked at smaller x through the tension between observables: F_2 , F_L , σ_r^{HQ} .

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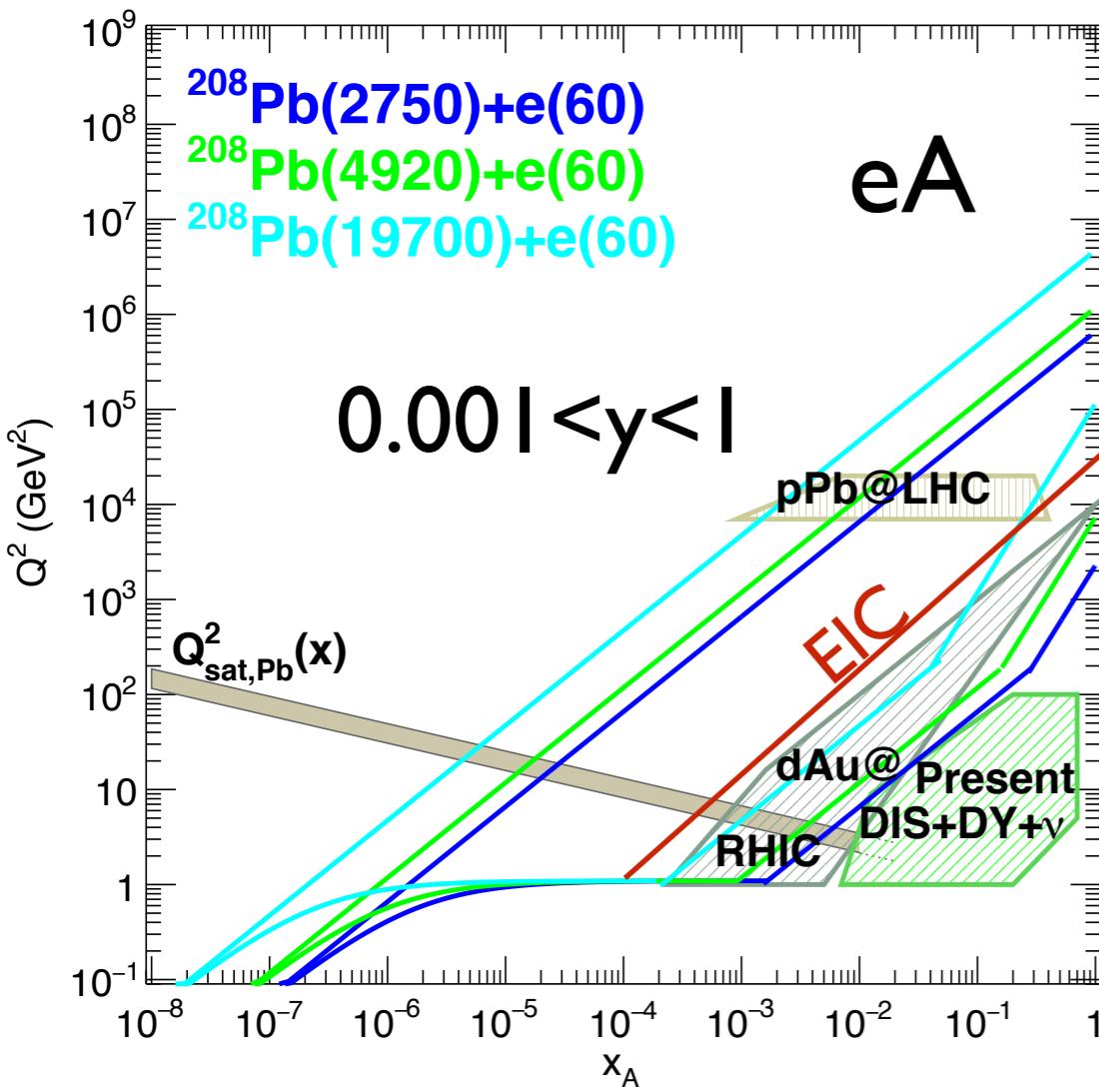
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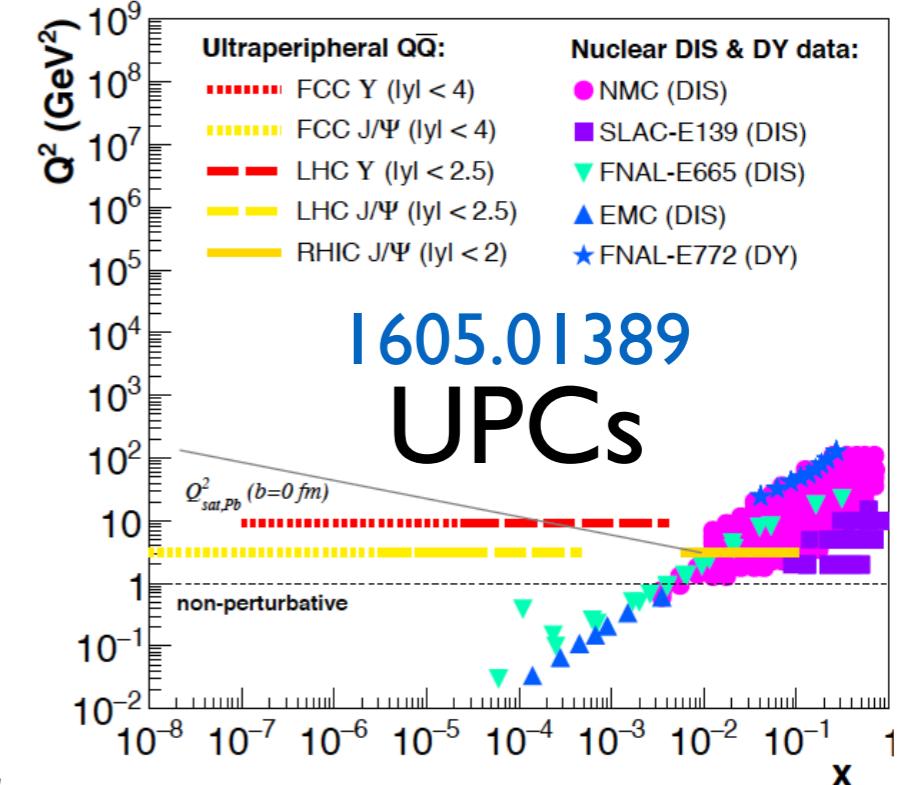
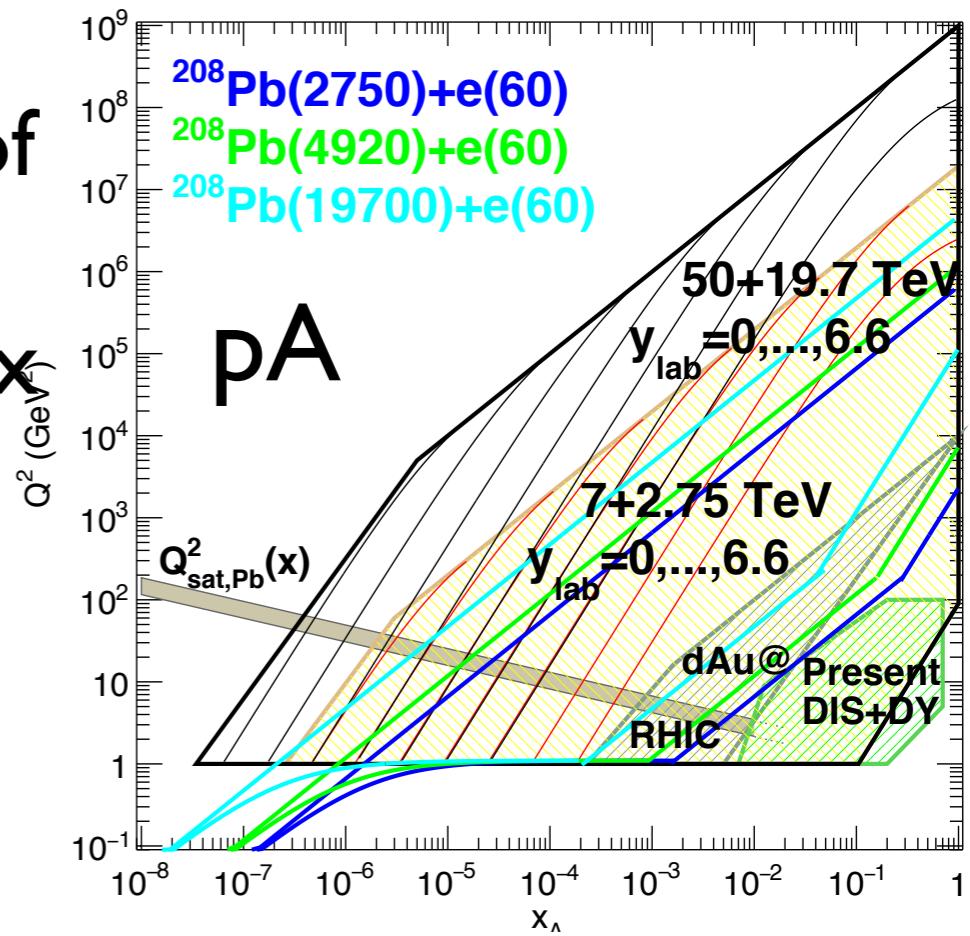
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EICs versus pA:



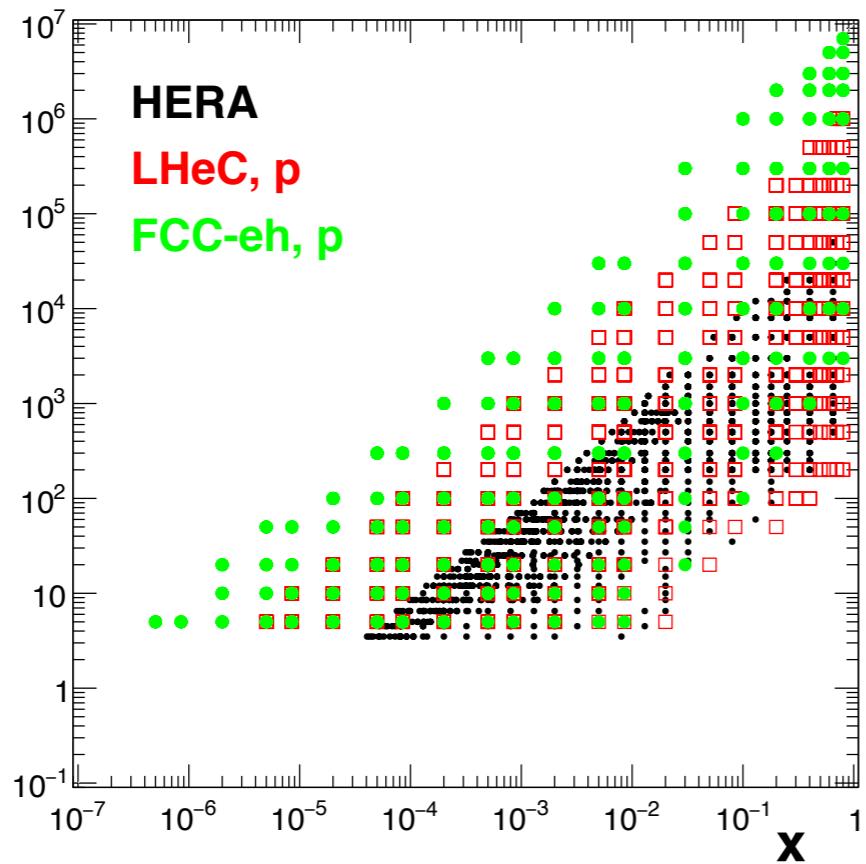
- Extension of 4-5 orders of magnitude in x and Q^2 wrt existing DIS data.



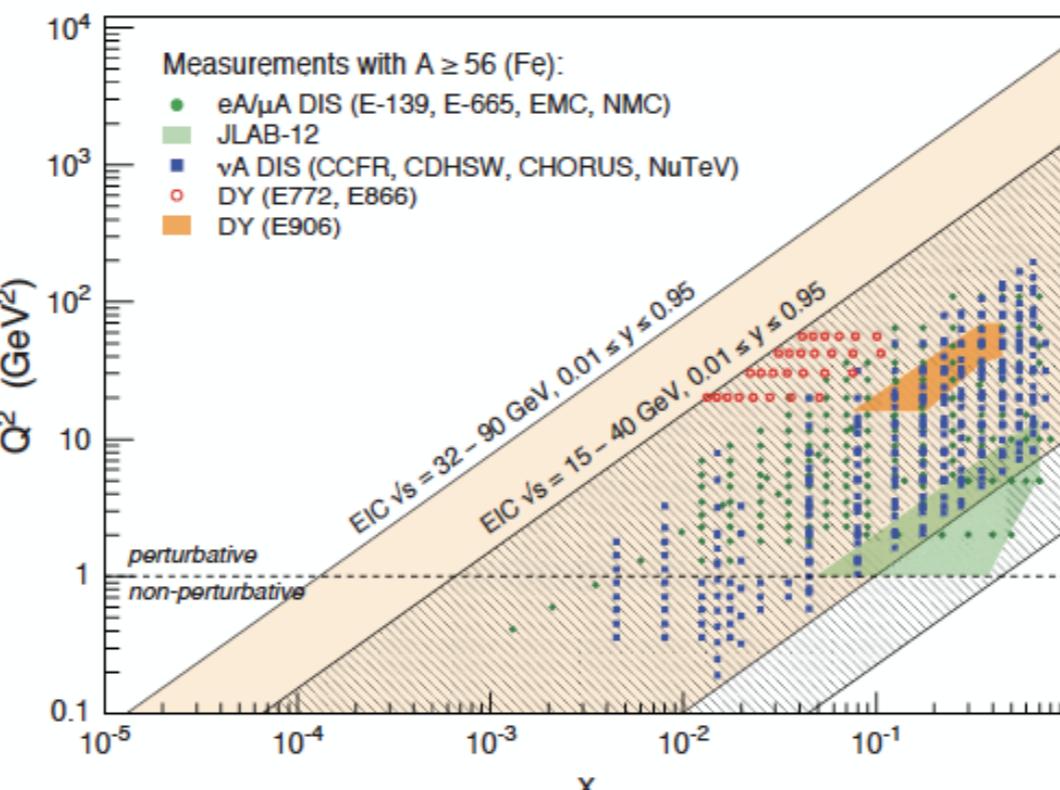
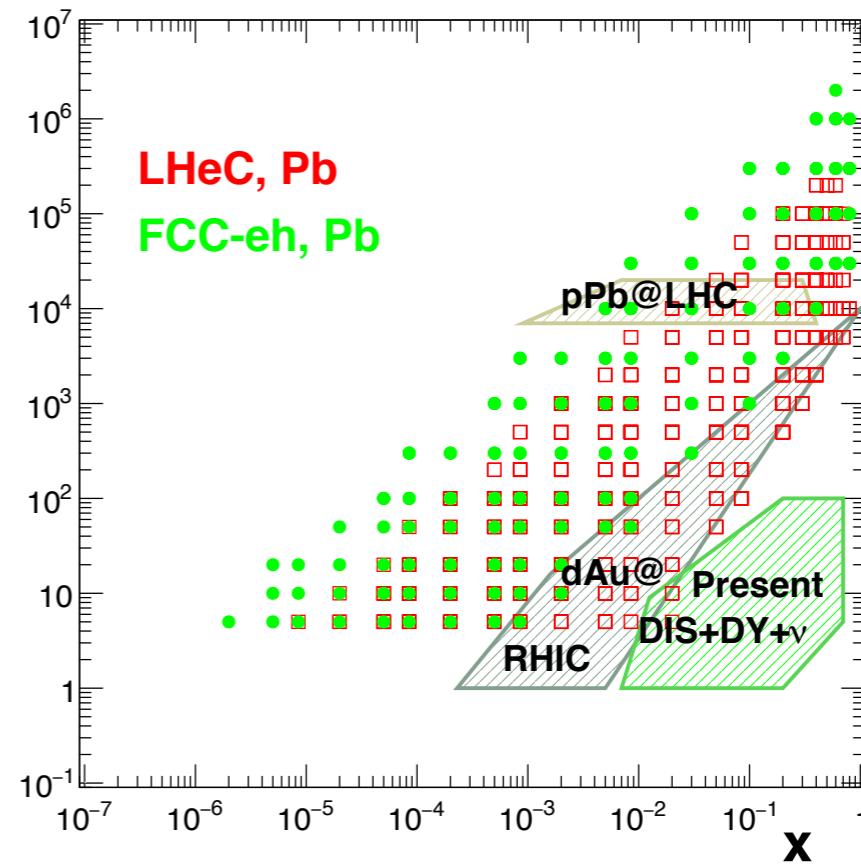
- DIS offers:
 - A clean experimental environment: low multiplicity, no pileup, fully constrained kinematics x, Q^2 by reconstructing the outgoing lepton;
 - A more controlled theoretical setup: many 1st-principles calculations.

LHeC/FCC-eh vs. EIC:

Q^2 (GeV 2)

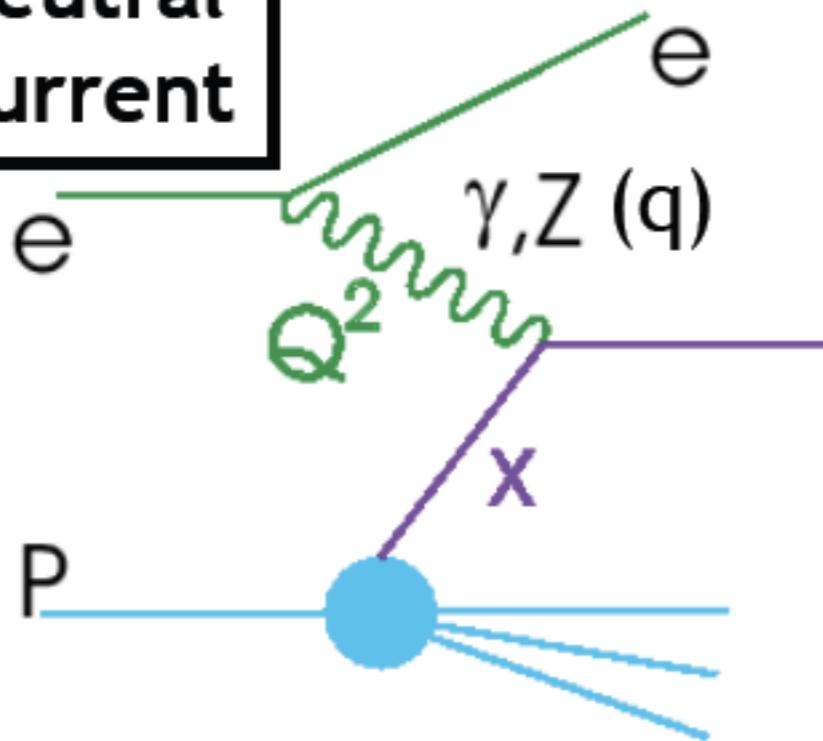


Q^2 (GeV 2)



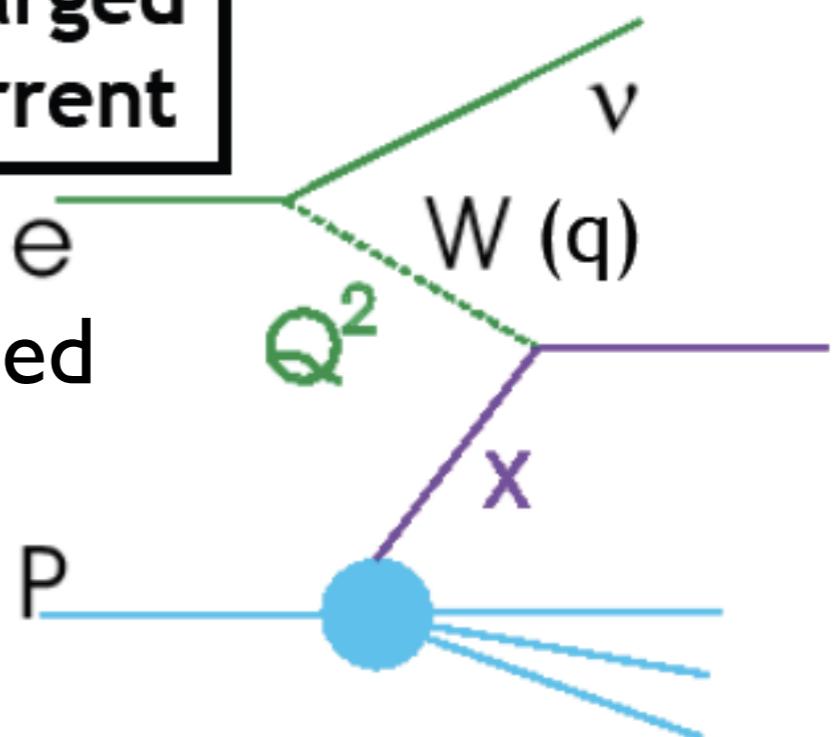
Some relations:

Neutral Current



Charged Current

Fully constrained kinematics!



$$\frac{d^2\sigma_{NC}}{dx dQ^2} = \frac{2\pi\alpha^2 Y_+}{Q^4 x} \cdot \sigma_{r,NC}$$

$$\frac{d^2\sigma_{CC}^\pm}{dx dQ^2} = \frac{1 \pm P}{2} \cdot \frac{G_F^2}{2\pi x} \cdot \left[\frac{M_W^2}{M_W^2 + Q^2} \right]^2 Y_+ \cdot \sigma_{r,CC}$$

$$\sigma_{r,NC} = F_2 + \frac{Y_-}{Y_+} x F_3 - \frac{y^2}{Y_+} F_L,$$

$$Y_\pm = 1 \pm (1 - y)^2$$

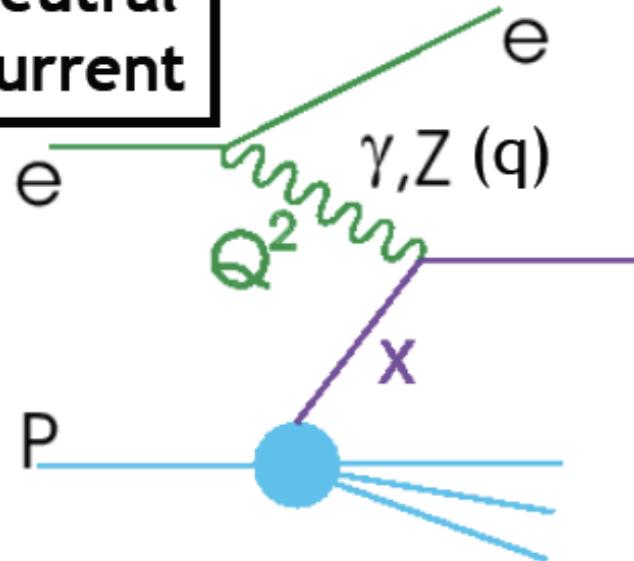
$$\sigma_{r,CC}^\pm = W_2^\pm \mp \frac{Y_-}{Y_+} x W_3^\pm - \frac{y^2}{Y_+} W_L^\pm$$

$$F_2^\pm = F_2 + \kappa_Z (-v_e \mp P a_e) \cdot F_2^{\gamma Z} + \kappa_Z^2 (v_e^2 + a_e^2 \pm 2 P v_e a_e) \cdot F_2^Z$$

$$x F_3^\pm = \kappa_Z (\pm a_e + P v_e) \cdot x F_3^{\gamma Z} + \kappa_Z^2 (\mp 2 v_e a_e - P (v_e^2 + a_e^2)) \cdot x F_3^Z$$

Extraction of PDFs:

Neutral Current



$$\frac{d^2\sigma_{NC}}{dx dQ^2} = \frac{2\pi\alpha^2 Y_+}{Q^4 x} \cdot \sigma_{r,NC}$$

$$\sigma_{r,NC} = F_2 + \frac{Y_-}{Y_+} x F_3 - \frac{y^2}{Y_+} F_L,$$

- Ignoring EW contributions:

$$F_2(x, Q^2) \propto \sum x q(x, Q^2)$$

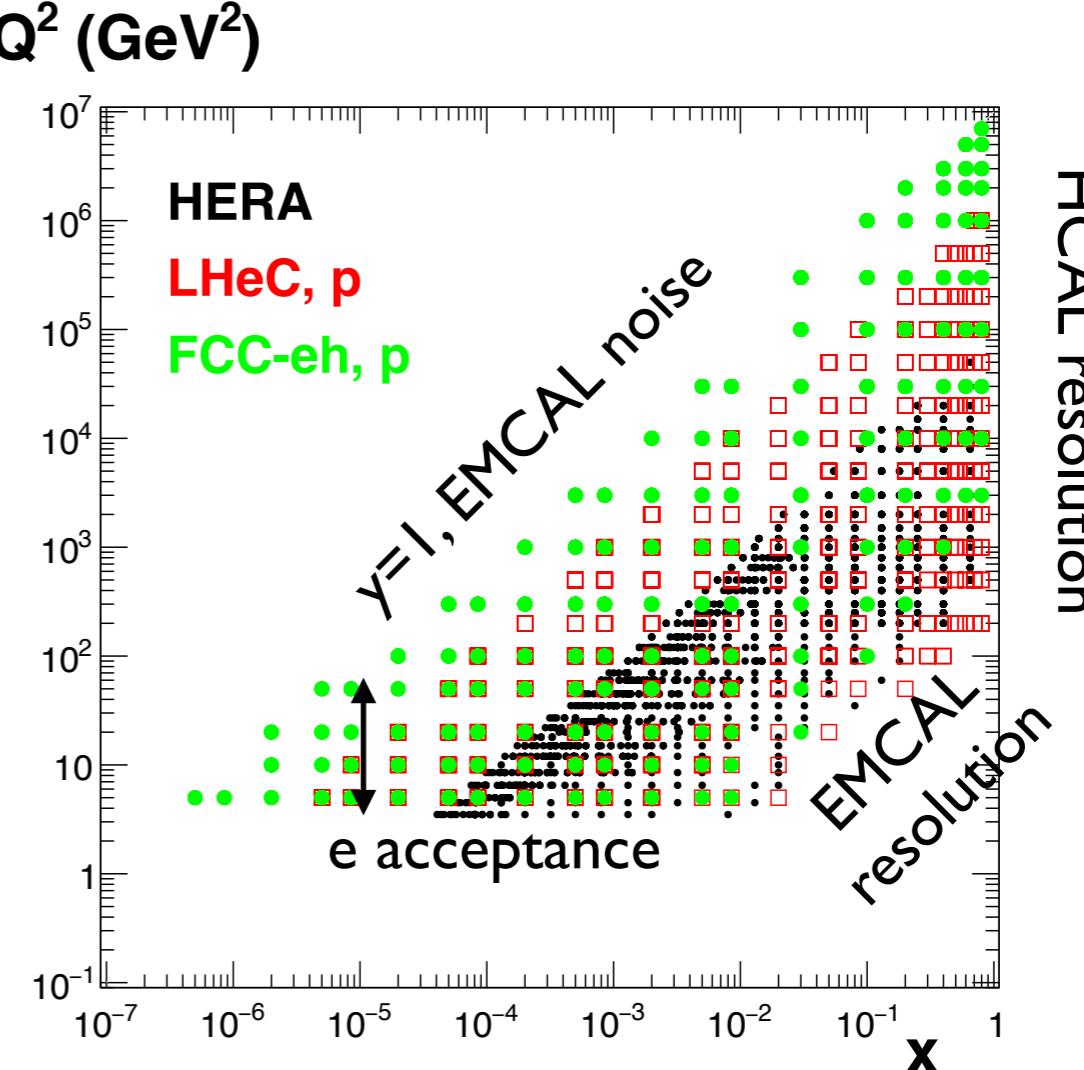
$$\frac{\partial F_2(x, Q^2)}{\partial \log Q^2} \propto x g(x, Q^2)$$

$$F_L(x, Q^2) \propto x g(x, Q^2) - F_2(x, Q^2)$$

Sensitive to the sea.

Via DGLAP: requires lever arm in Q^2 .

Via DGLAP: requires lever arm in s
(different y at fixed x, Q^2): better σ_{red} .



HCAL resolution

EPPS16@LHeC (I):

The LHeC pseudodata

- Assume $\mathcal{L}_{\text{ep}} = 10 \text{ fb}$, $\mathcal{L}_{\text{ePb}} = 1 \text{ fb}$ (per nucleon)
- The assumed energy configs: $\sqrt{s_p} = 7 \text{ TeV}$, $\sqrt{s_{\text{Pb}}} = 2.75 \text{ TeV}$ (per nucleon) on $E_e = 60 \text{ GeV}$ electrons.
- The pseudodata are here obtained from ratios of reduced cross sections σ^i and relative point-to-point ($\delta_{\text{uncor.}}^i$) and normalization ($\delta_{\text{norm.}}^i$) uncertainties as

$$R_i = R_i(\text{EPS09}) \times [1 + \delta_{\text{uncor.}}^i r^i + \delta_{\text{norm.}}^i r^{\text{norm.}}]$$

where

$$R_i(\text{EPS09}) = \frac{\sigma_{\text{ePb}}^i(\text{CTEQ6.6} + \text{EPS09})}{\sigma_{\text{ep}}^i(\text{CTEQ6.6})},$$

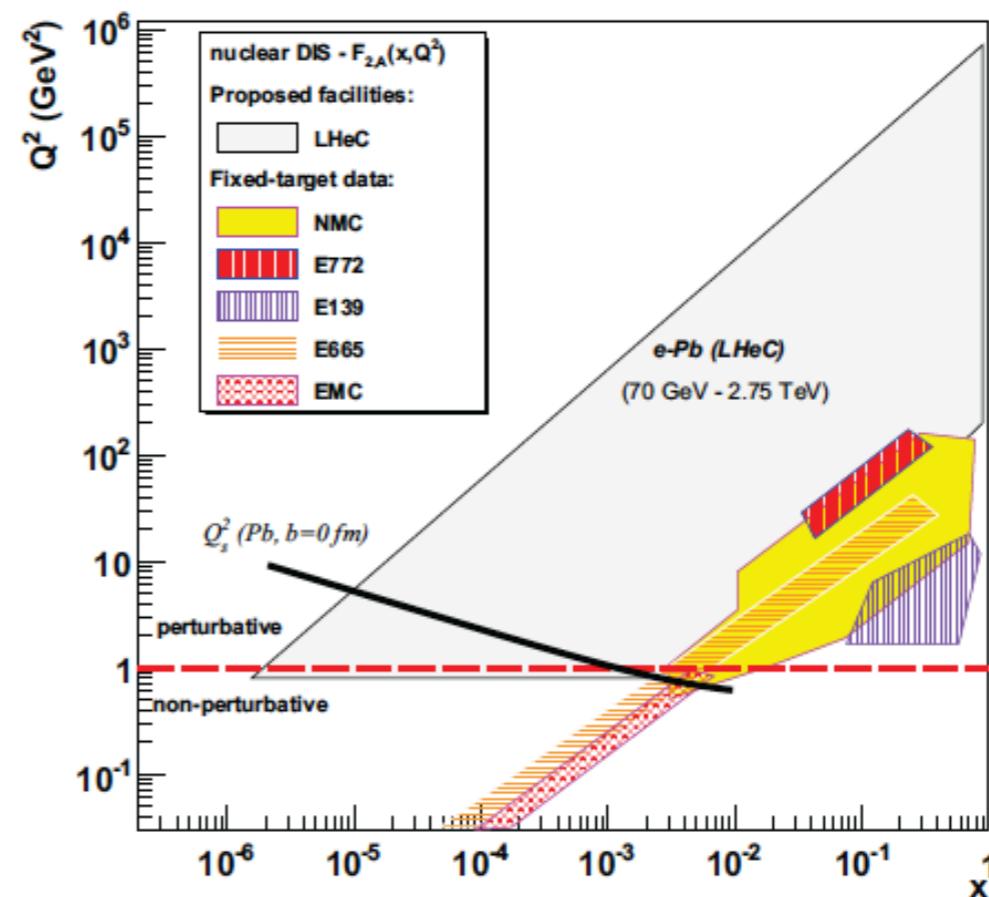
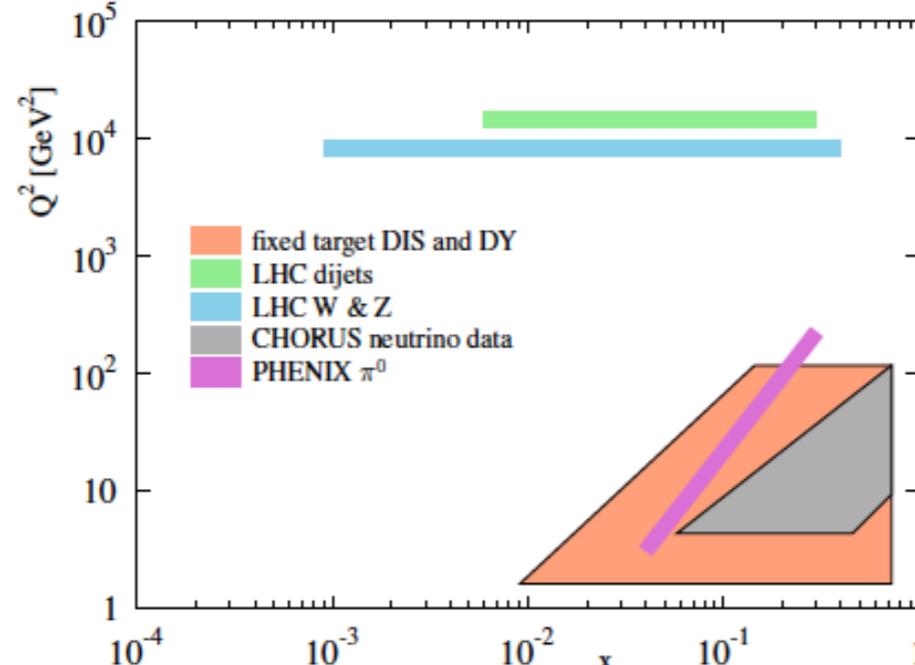
and r^i and $r^{\text{norm.}}$ are Gaussian random numbers.

- In EPS09 $R_{u_V} \approx R_{d_V}$, $R_{\bar{u}} \approx R_{\bar{d}} \approx R_{\bar{s}}$ (free in EPPS16, but would not expect large deviations from this)

EPPS16@LHeC (I):

The analysis framework

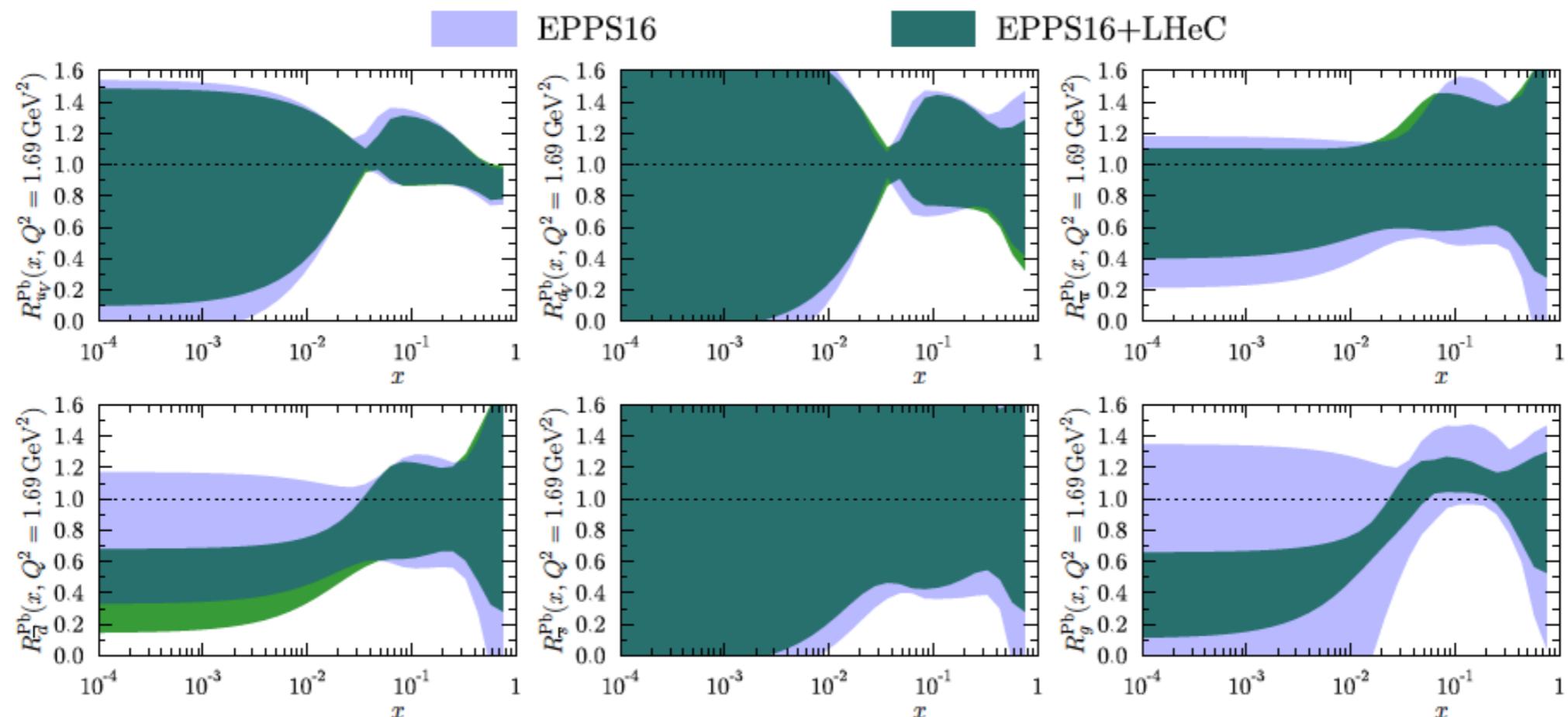
- The fit framework same as in the EPPS16 analysis [EPJ C77, 163]
- Include the same data as in EPPS16 plus LHeC (NC and CC) pseudo data.
- Hessian uncertainty analysis with $\Delta\chi^2 = 52$ (as in EPPS16)



EPPS16@LHeC (I):

The effect of LHeC pseudodata

- The improvement after adding the LHeC data ($Q^2 = 1.69 \text{ GeV}^2$)



nPDFs@LHeC (II):

The effect of LHeC pseudodata

- Why it's so hard to pin down the flavor dependence?
 - Take the valence up-quark distribution u_V^A as an example:

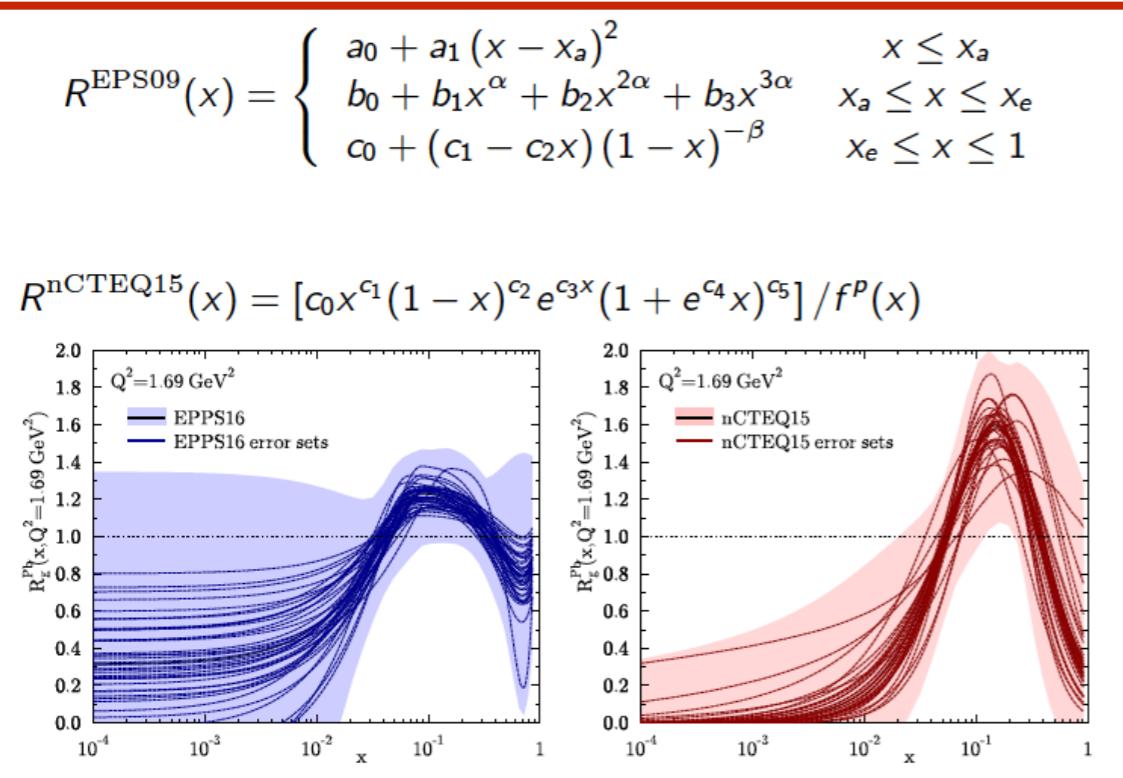
$$u_V^A = \frac{Z}{A} R_{u_V} u_V^{\text{proton}} + \frac{A-Z}{A} R_{d_V} d_V^{\text{proton}}$$

- Write this in terms of average modification R_V and the difference δR_V

$$R_V \equiv \frac{R_{uV} u_V^{\text{proton}} + R_{dV} d_V^{\text{proton}}}{u_V^{\text{proton}} + d_V^{\text{proton}}}, \quad \delta R_V \equiv R_{uV} - R_{dV}$$

- The effects of flavour separation (i.e. δR_V here) are suppressed in cross sections — but also so in most of the nPDF applications.

nPDFs@LHeC (II):

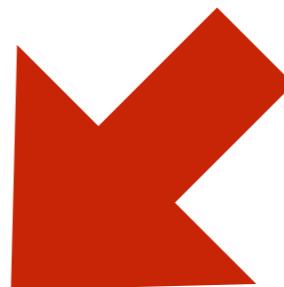
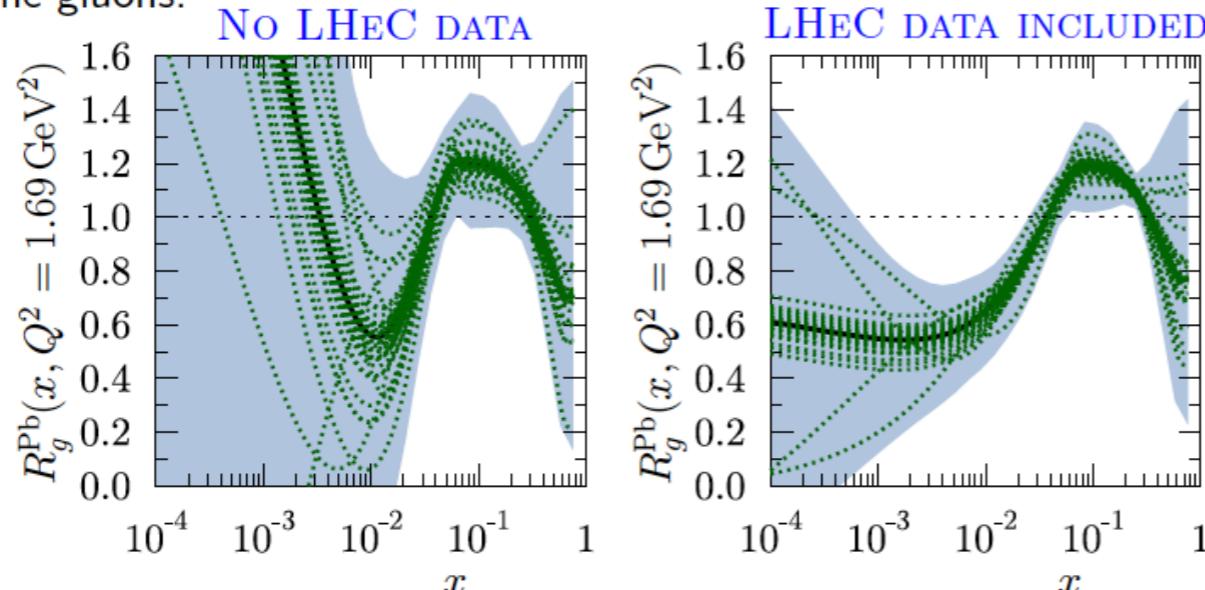


$$R(x \leq x_a) = a_0 + a_1(x - x_a)^2 + \sqrt{x}(x_a - x) \left[a_2 \log \left(\frac{x}{x_a} \right) + a_3 \log^2 \left(\frac{x}{x_a} \right) + a_4 \log^3 \left(\frac{x}{x_a} \right) + \dots \right]$$

or

$$R(x \leq x_a) = a_0 + (x - x_a)^2 [a_1 + a_2 x^\alpha + a_3 x^{2\alpha} + a_4 x^{3\alpha} + \dots], \quad \alpha \ll 1$$

- Would need Monte-Carlo methods to more reliably map the uncertainties
➡ Further work needed
- Despite all the shortcomings, a typical result using a more flexible form for the gluons:



xFitter:

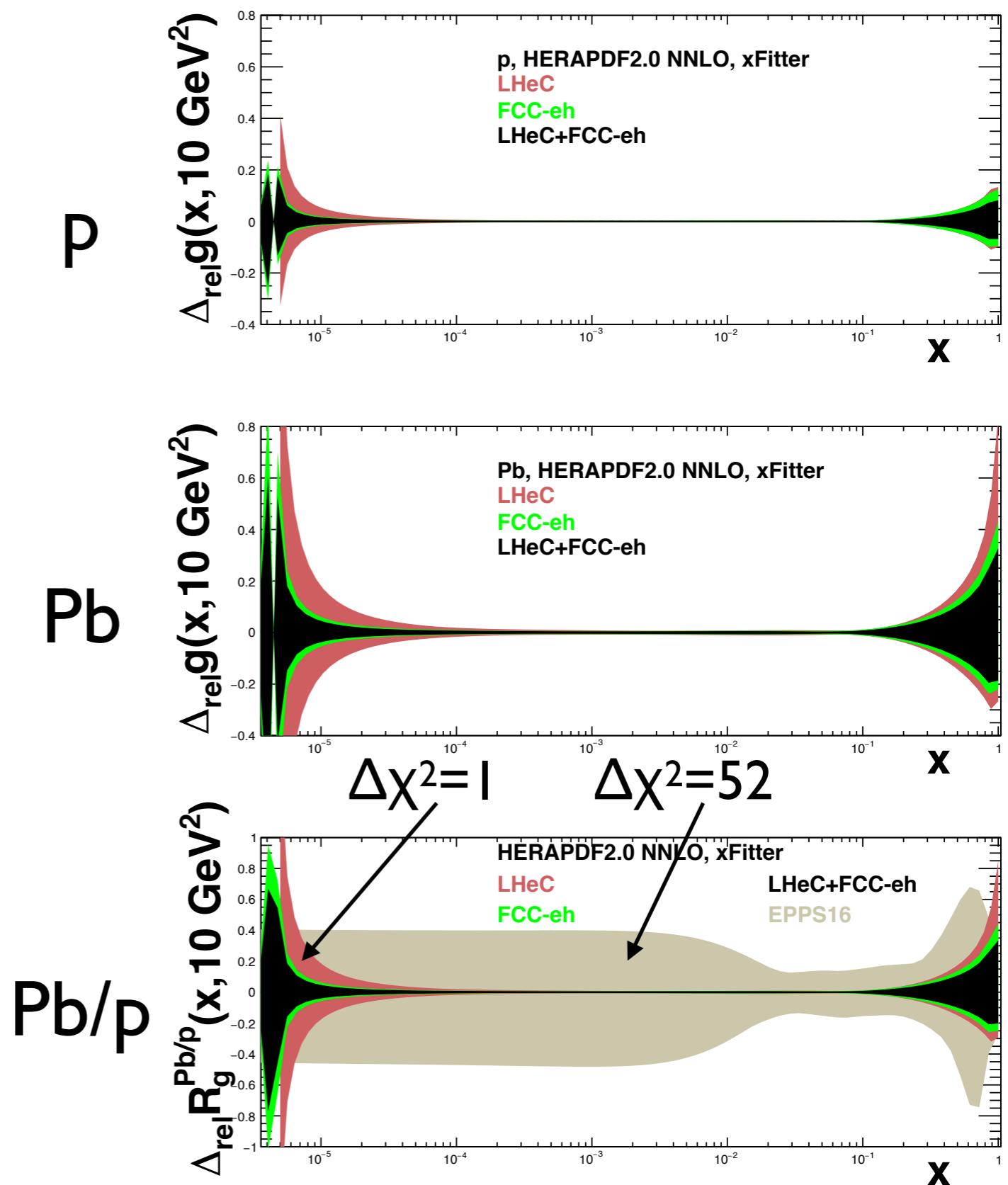
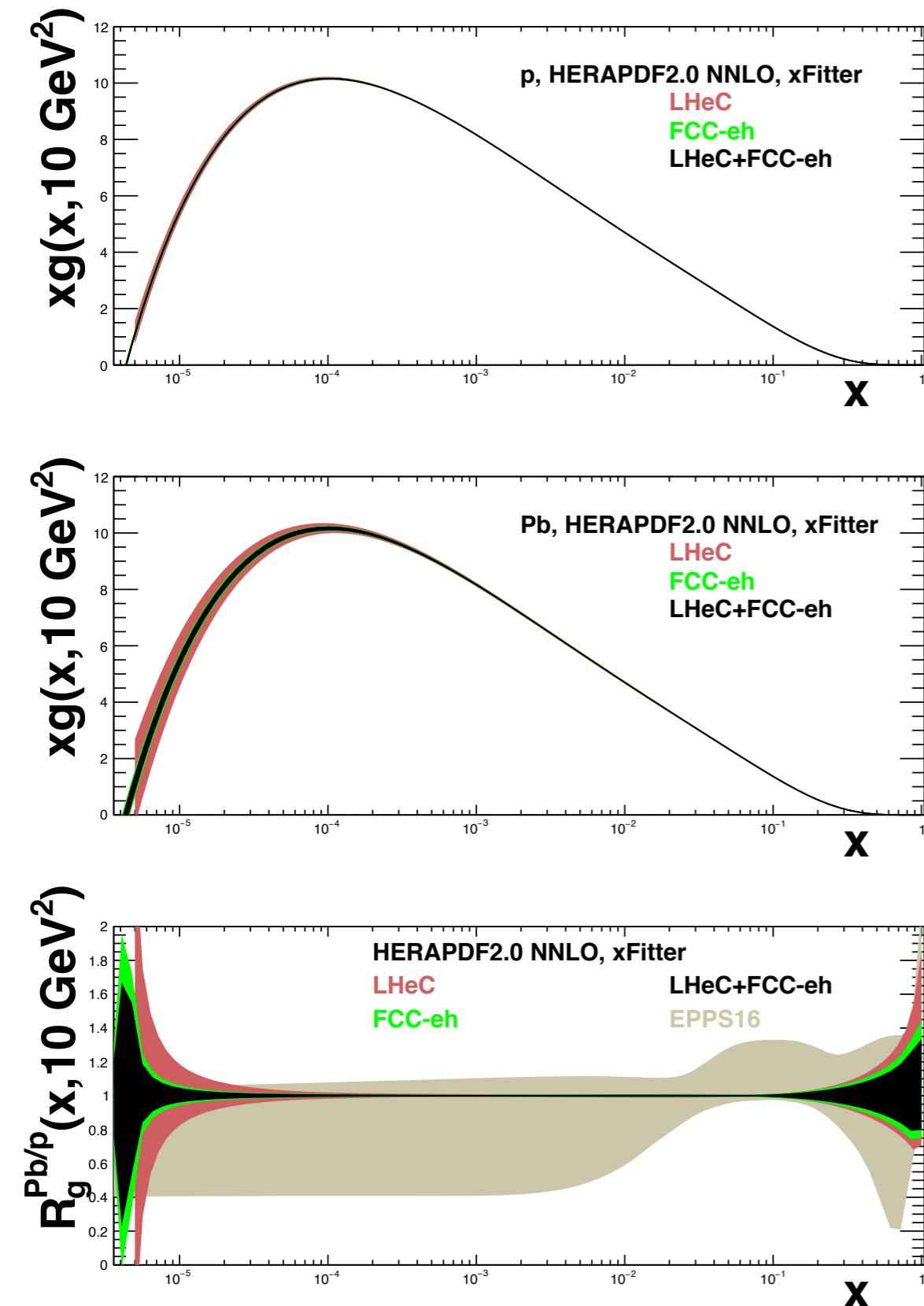
- Extraction of Pb-only PDFs by fitting pseudodata, using xFitter ([I410.4412](#)) I.2.2 to estimate the ‘ultimate’ precision that could be achieved ([P. Agostini, NA](#)):
→ HERAPDF2.0-type parametrisation ([I506.06042](#), 14 parameters), NNLO evolution, RTOPT mass scheme, $\alpha_s=0.118$.

$$\begin{aligned}xg(x) &= A_g x^{B_g} (1-x)^{C_g} - A'_g x^{B'_g} (1-x)^{C'_g}, \\xu_v(x) &= A_{u_v} x^{B_{u_v}} (1-x)^{C_{u_v}} \left(1 + E_{u_v} x^2\right), \\xd_v(x) &= A_{d_v} x^{B_{d_v}} (1-x)^{C_{d_v}}, \\x\bar{U}(x) &= A_{\bar{U}} x^{B_{\bar{U}}} (1-x)^{C_{\bar{U}}} (1 + D_{\bar{U}} x), \\x\bar{D}(x) &= A_{\bar{D}} x^{B_{\bar{D}}} (1-x)^{C_{\bar{D}}}.\end{aligned}$$

$$xU = xu + xc, \quad x\bar{U} = x\bar{u} + x\bar{c}, \quad xD = xd + xs, \quad x\bar{D} = x\bar{d} + x\bar{s}$$

- Central pseudodata values from HERAPDF2.0: no parametrisation bias.
- Standard xFitter/HERAPDF treatment of correlated/uncorrelated systematics.
- Only data with $Q^2 \geq 3.5 \text{ GeV}^2$, initial evolution scale 1.9 GeV^2 .
- Proton PDFs extracted in the same setup for consistency.

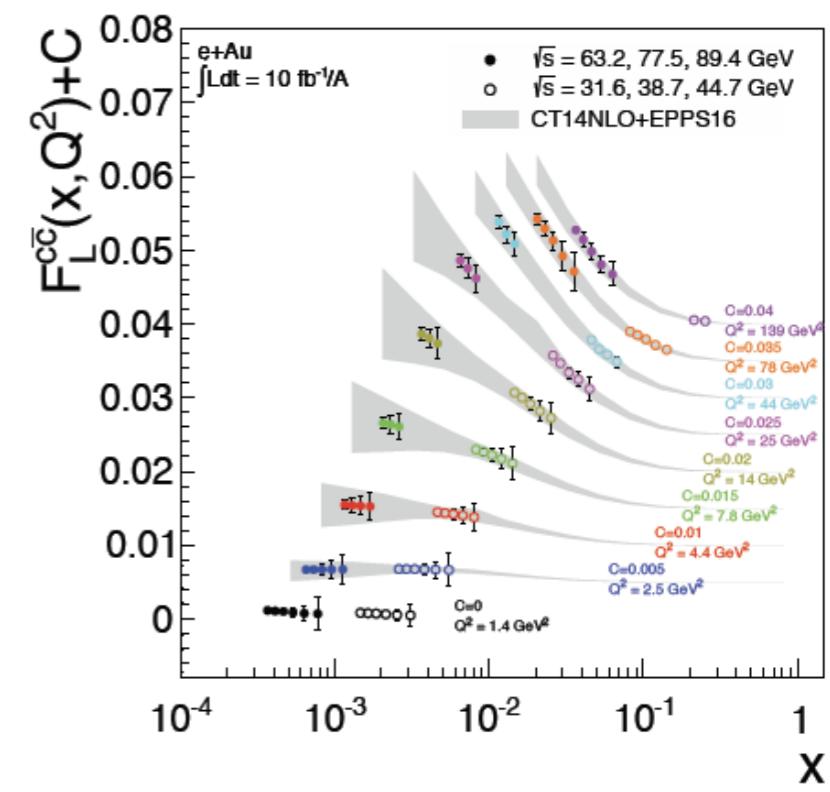
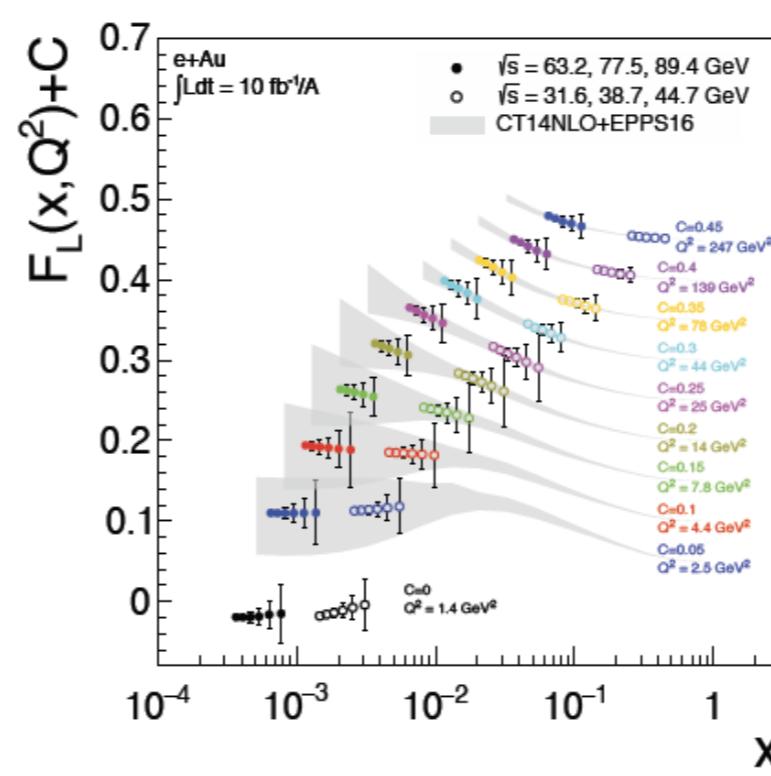
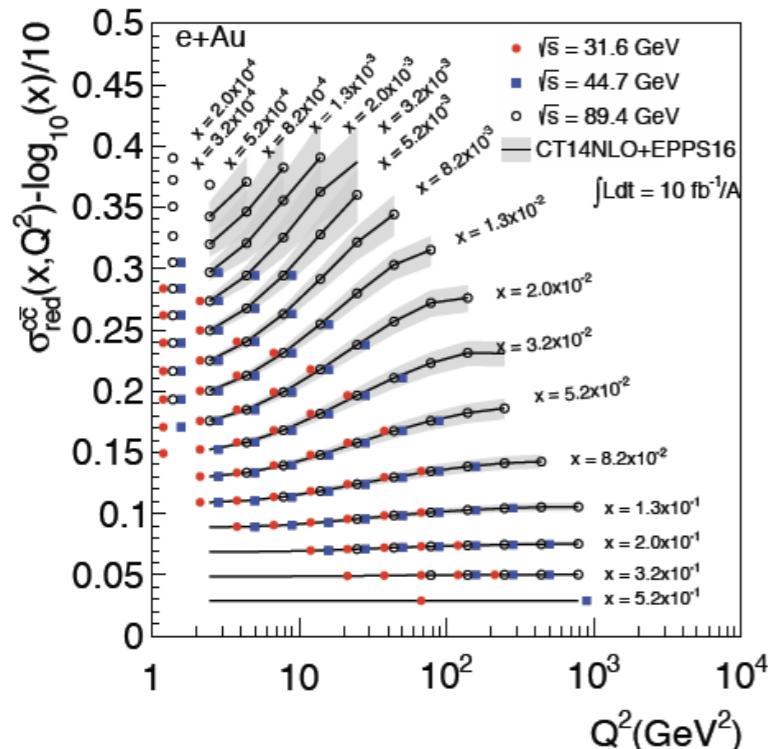
Results: gluon



EPPS16@EIC:

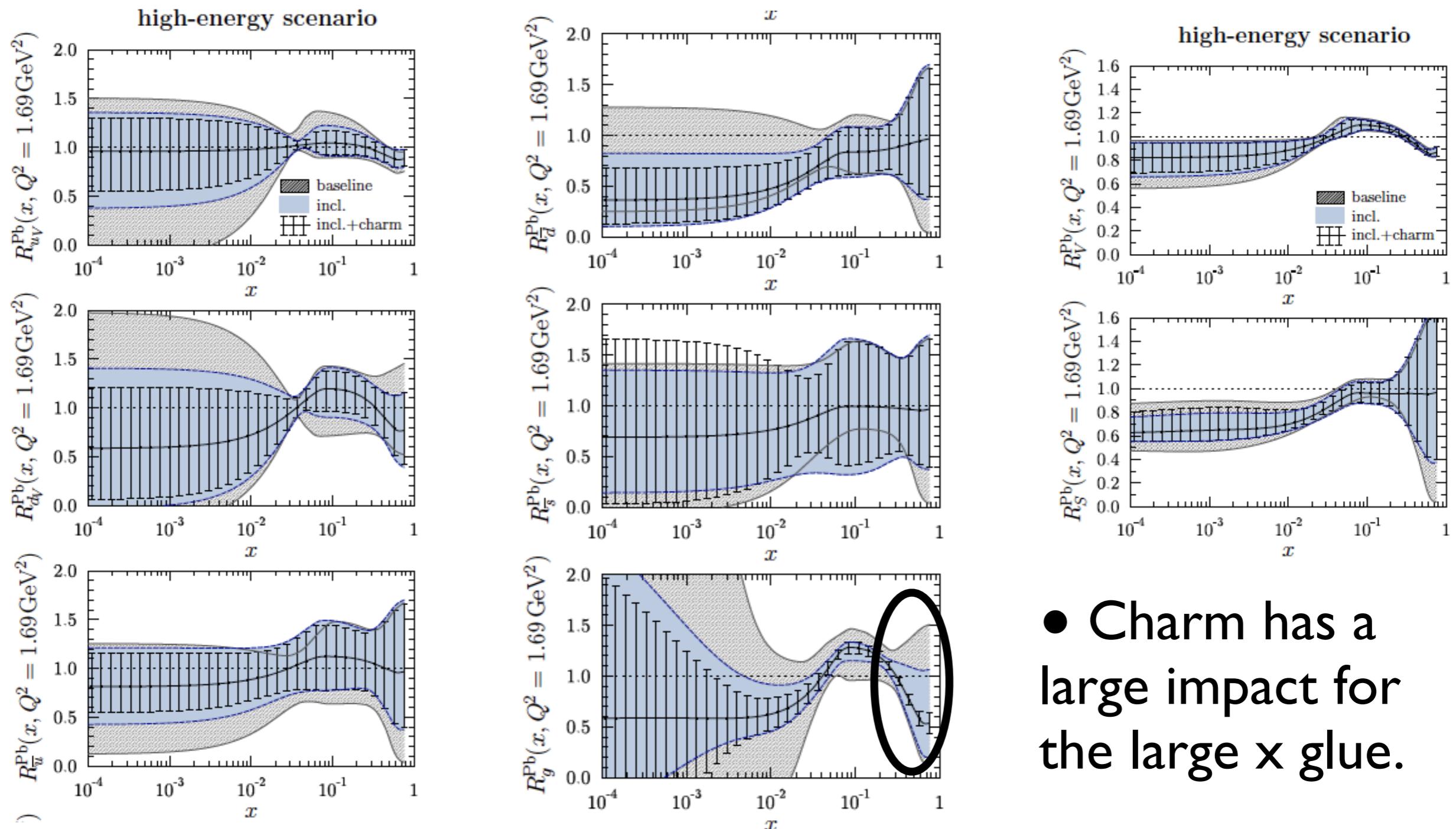
- Following a very similar approach to that shown for the LHeC (1708.05654):

→ Pseudodata generated with EPS09, uncertainties as achieved at HERA.
 → Impact of low (5 GeV) and high (20 GeV) electron energies, and of charm.



EPPS16@EIC:

- Very similar approach to that shown for the LHeC (1708.05654):
- Pseudodata generated with EPS09, uncertainties as at HERA.
- Impact of low (5 GeV) and high (20 GeV) E_e , and of charm.



- Charm has a large impact for the large x glue.

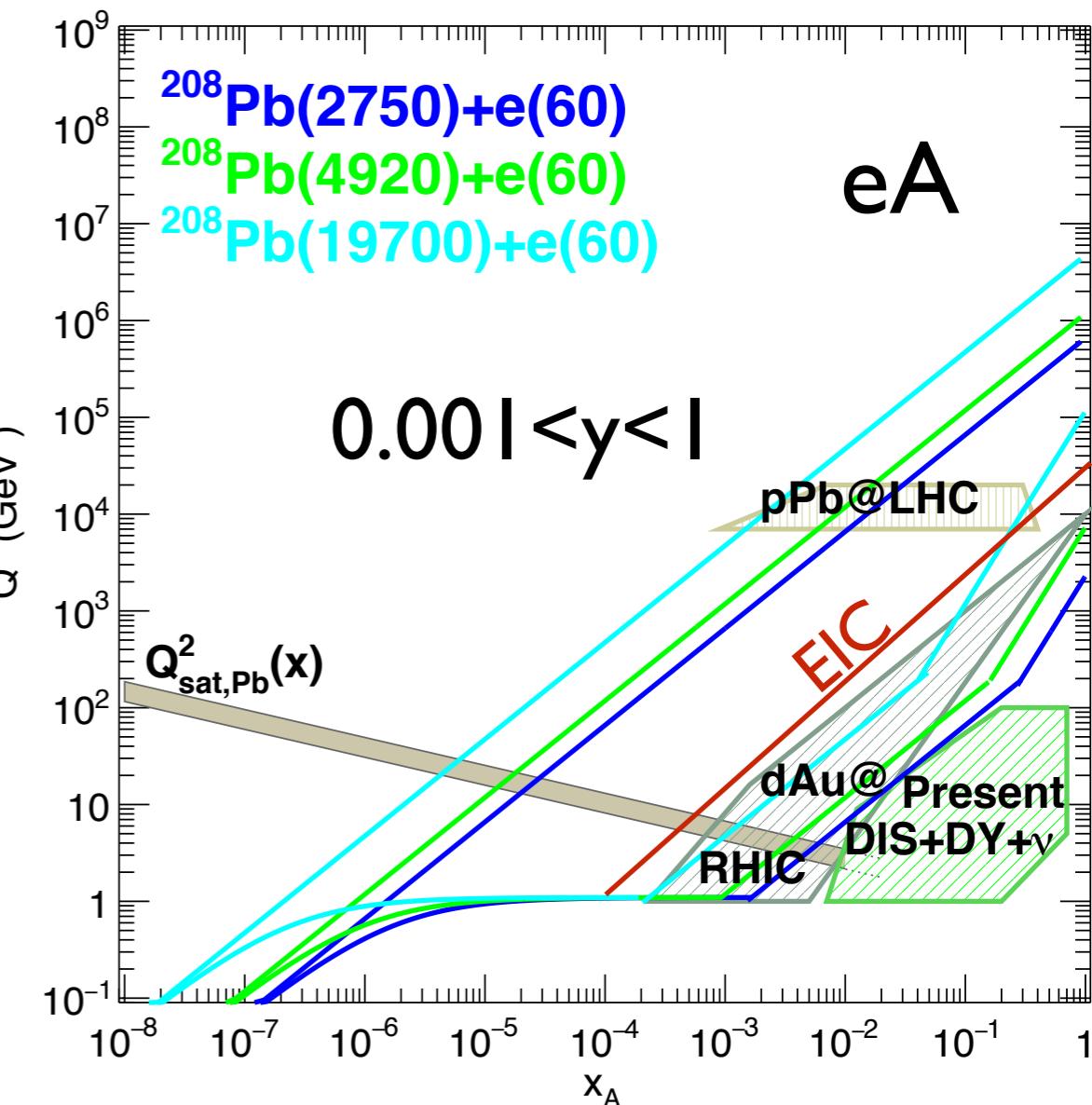
Summary:

- EICs are the ideal places to determine the nPDFs: fully constrained kinematics, well controlled th. & exp. setup.

- Limitation: do not cover as much as hadronic colliders, and luminosity may be important for quantitative studies e.g. impact of high x on low x .

- The EIC will not cover the kinematic region for the LHC or for future pA/AA machines.

- pA cannot be challenged in terms of kinematic reach: tests of collinear factorisation and its eventual breaking.
- Establishing the existence of a new regime of QCD will most probably be a quantitative issue demanding both ep/eA and pp/pA.



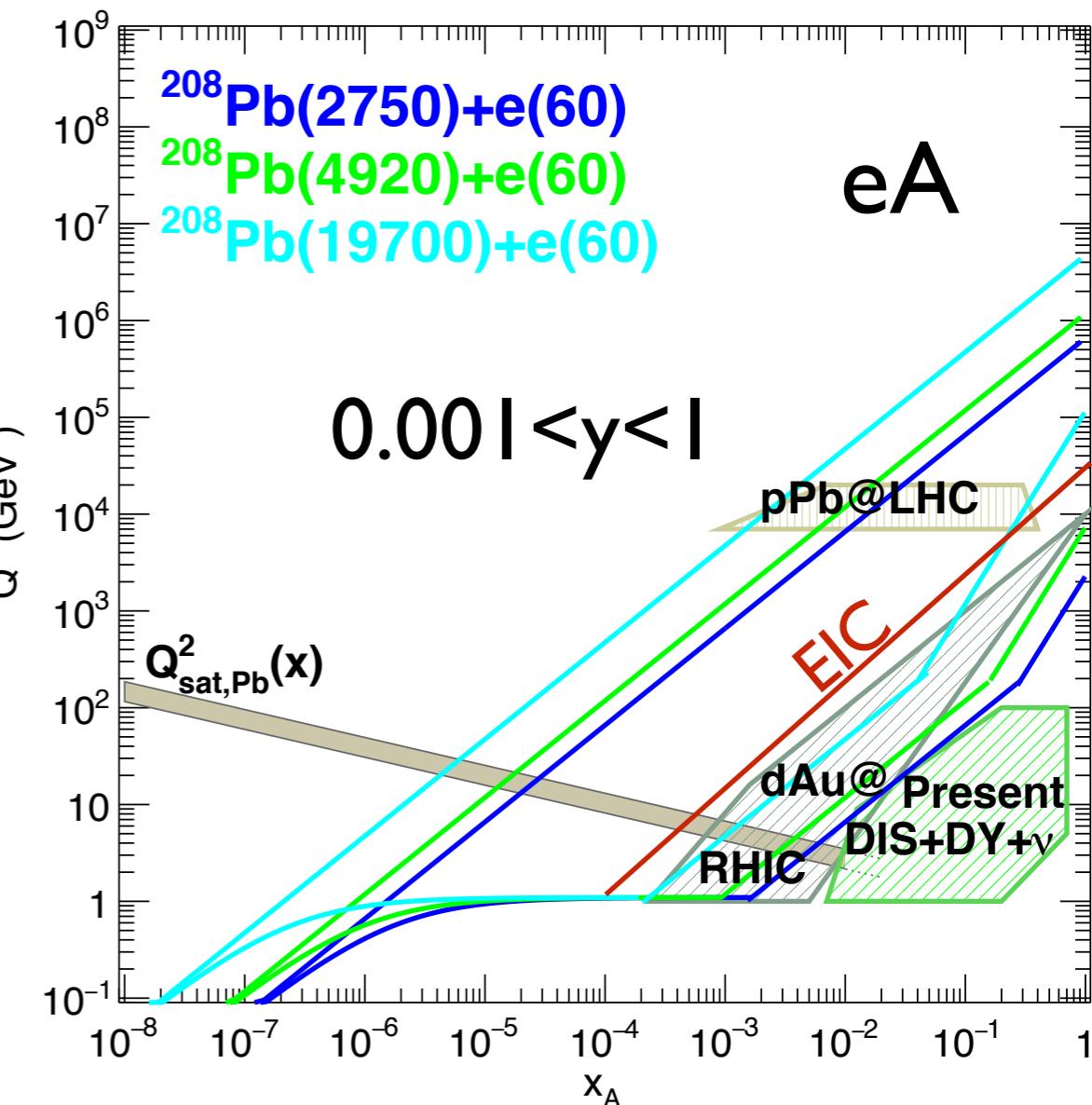
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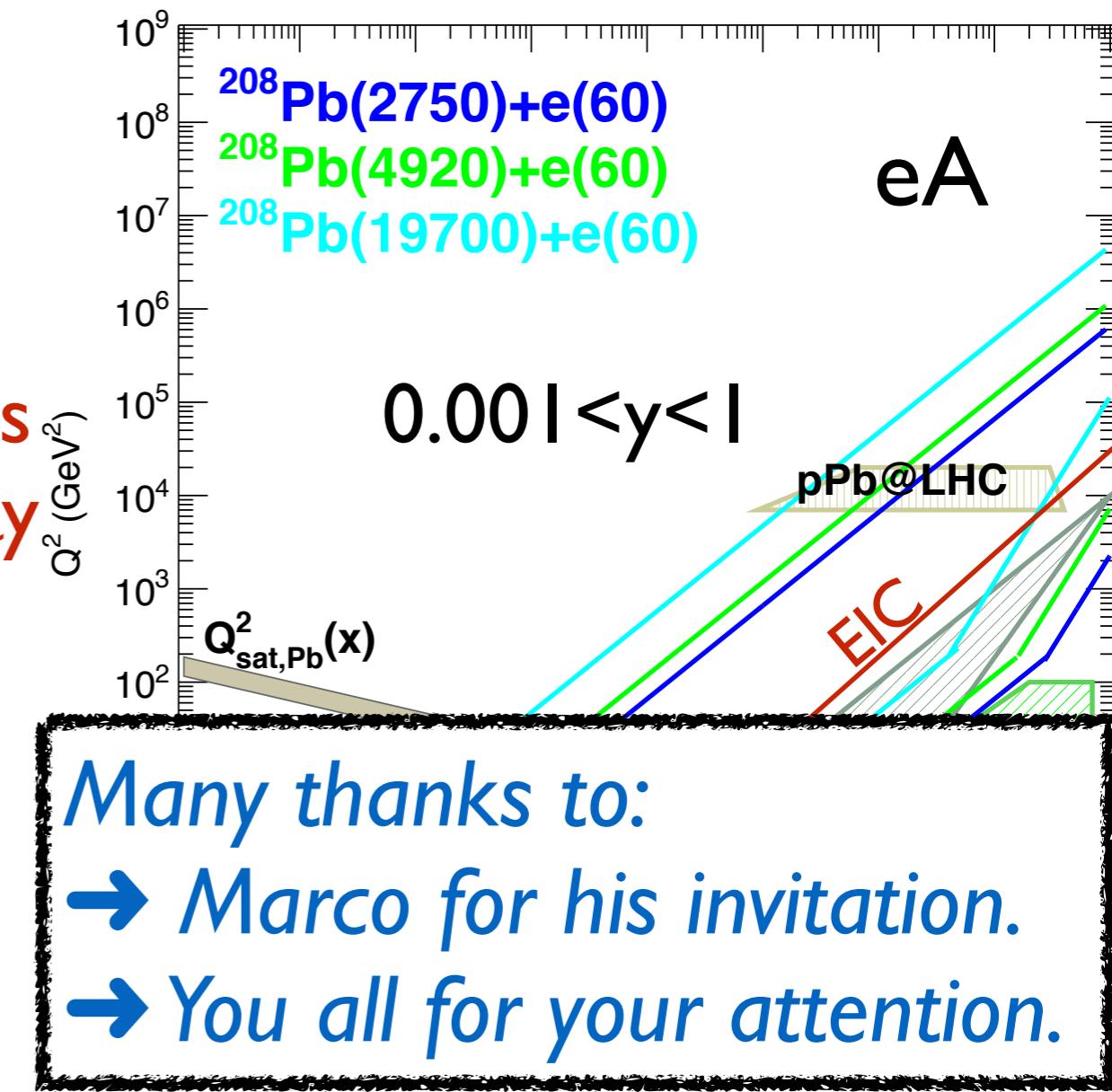


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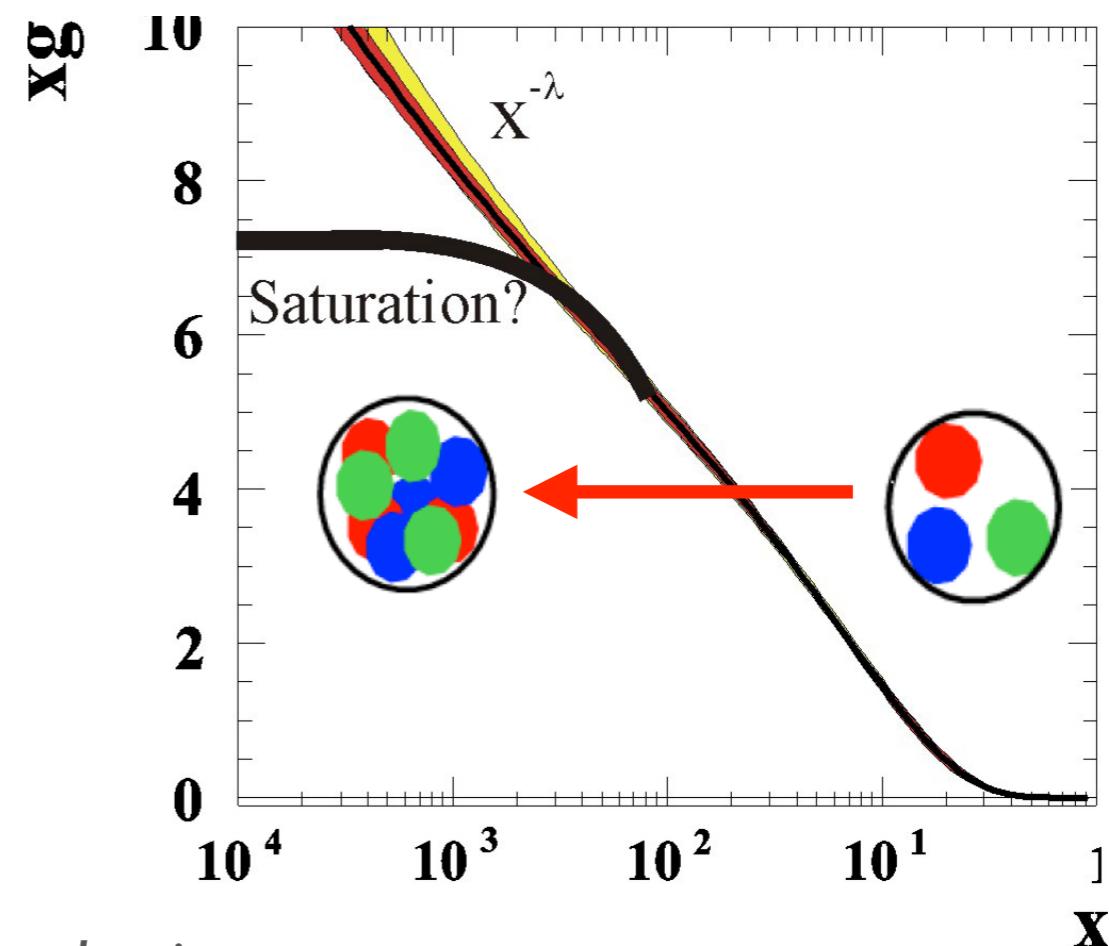
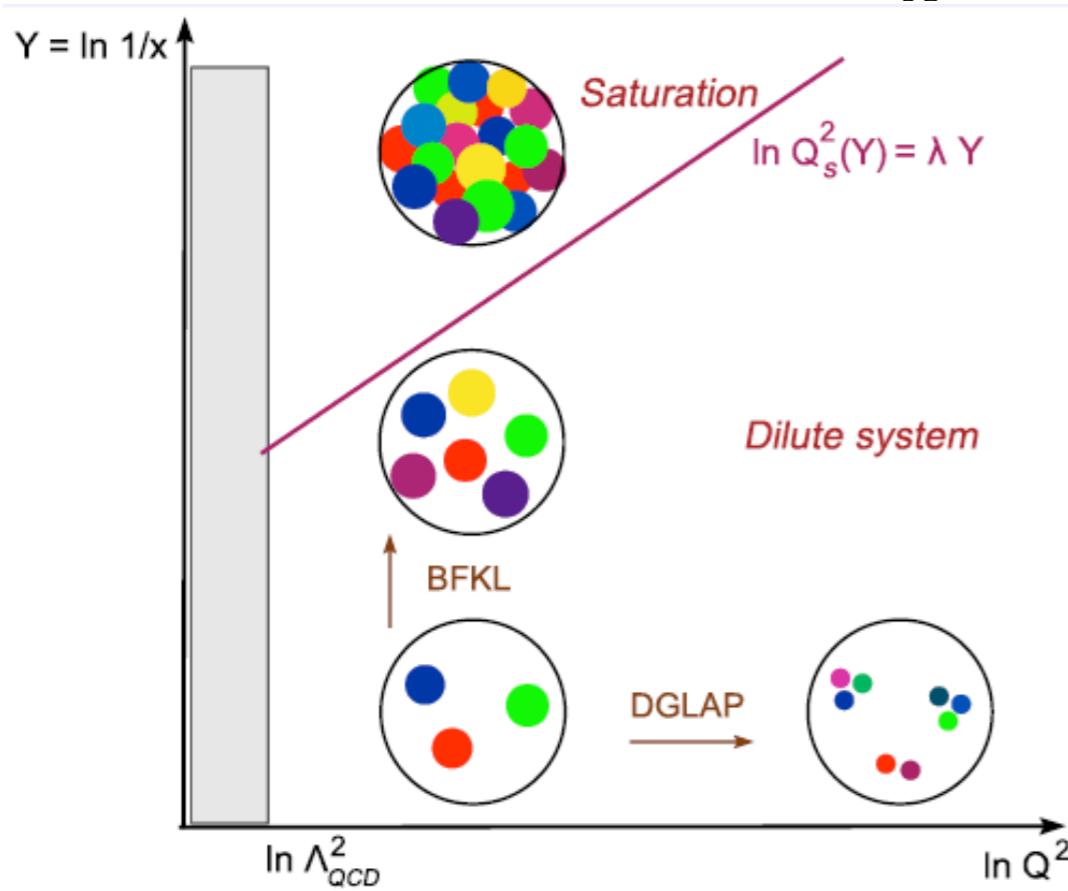
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Backup:

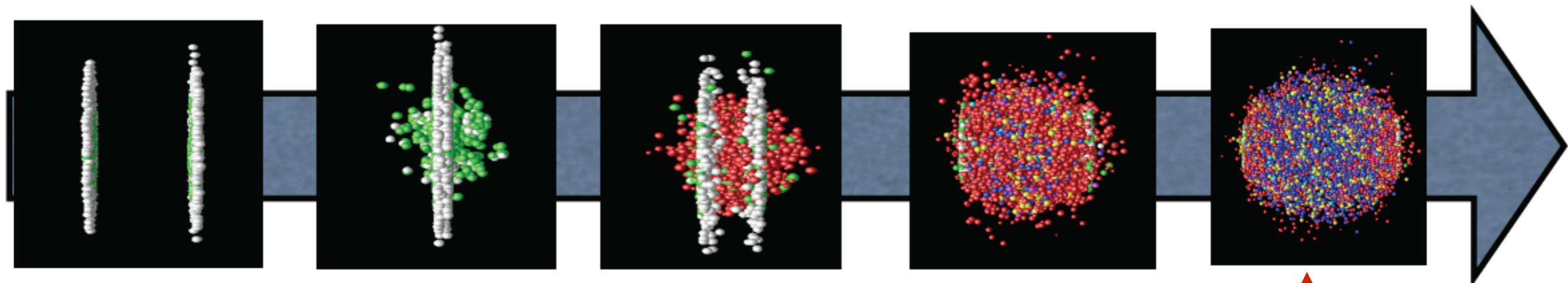
h/A wave function:

- Standard fixed-order perturbation theory (DGLAP, linear evolution) must eventually fail:
 - Large logs e.g. $\alpha_s \ln(1/x) \sim 1$: resummation (BFKL,CCFM,ABF,CCSS).
 - High-density: $x \downarrow, A \uparrow \Rightarrow$ non-linear regime, recombination balancing splitting: saturation, perturbative (CGC) or non.

$$\frac{xG_A(x, Q_s^2)}{\pi R_A^2 Q_s^2} \sim 1 \implies Q_s^2 \propto A^{1/3} x^{-0.3}$$



Relevance for HIC:



Gluons from saturated nuclei



Glasma?



QGP



Reconfinement

- Nuclear wave function at small x : **nuclear structure functions.**

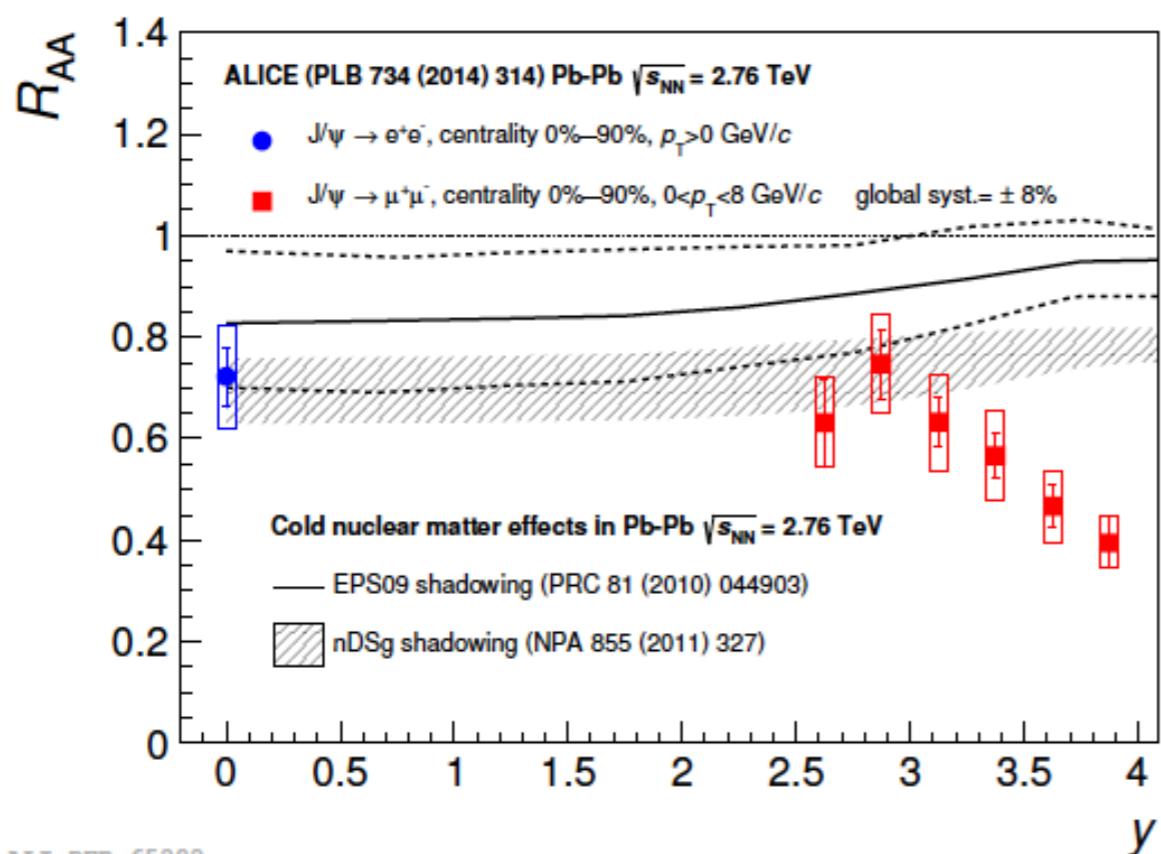
- Particle production at the very beginning: which factorisation in eA ?
- How does the system behave as \sim isotropised so fast?: initial conditions for plasma formation to be studied in eA .

- Probing the medium through energetic particles (jet quenching etc.): modification of QCD radiation and hadronization in the nuclear medium.

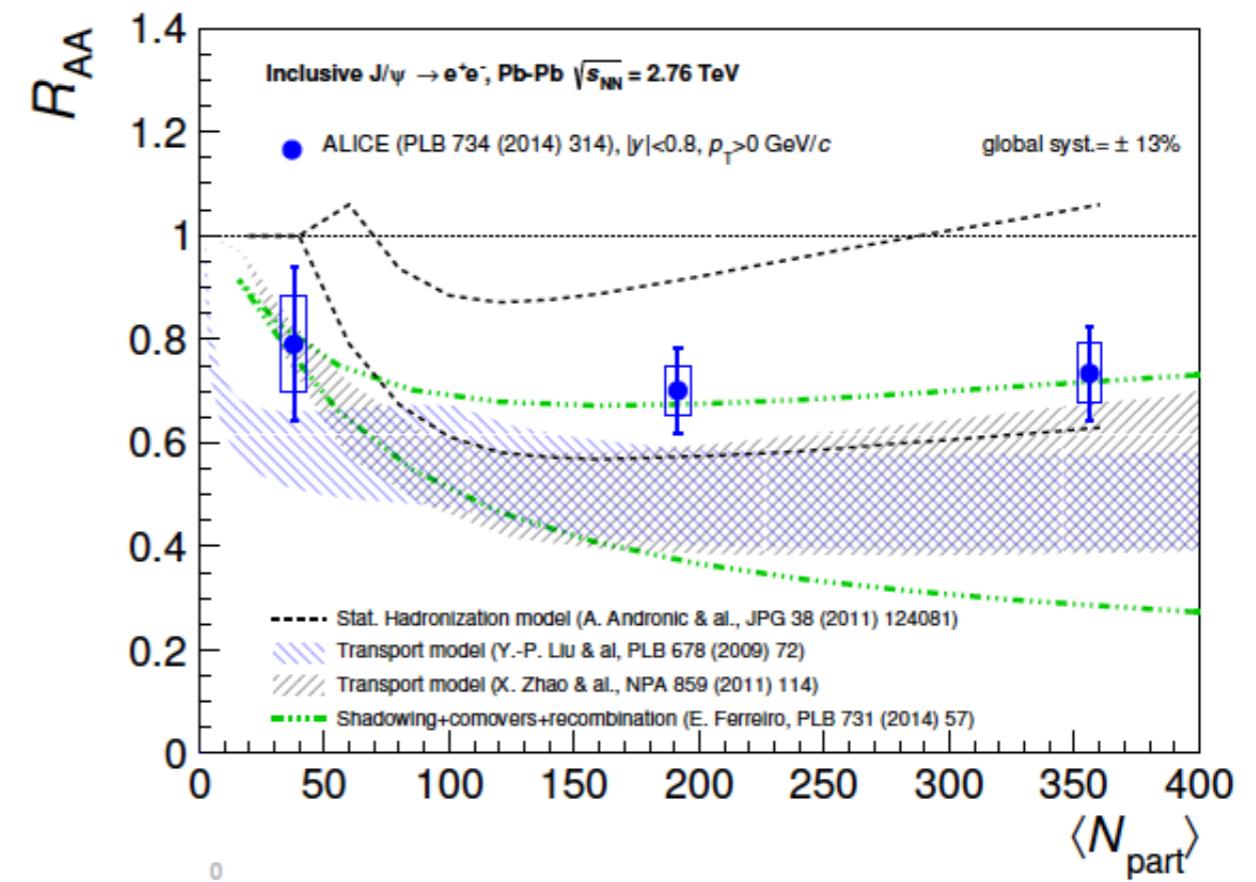
nPDFs for HIC:

- Lack of data \Rightarrow large uncertainties for the nuclear glue at small scales and x : problem for benchmarking in HIC in order to extract ‘medium’ parameters.

$$R = \frac{f_i/A}{Af_i/p} \approx \frac{\text{measured}}{\text{expected if no nuclear effects}}$$

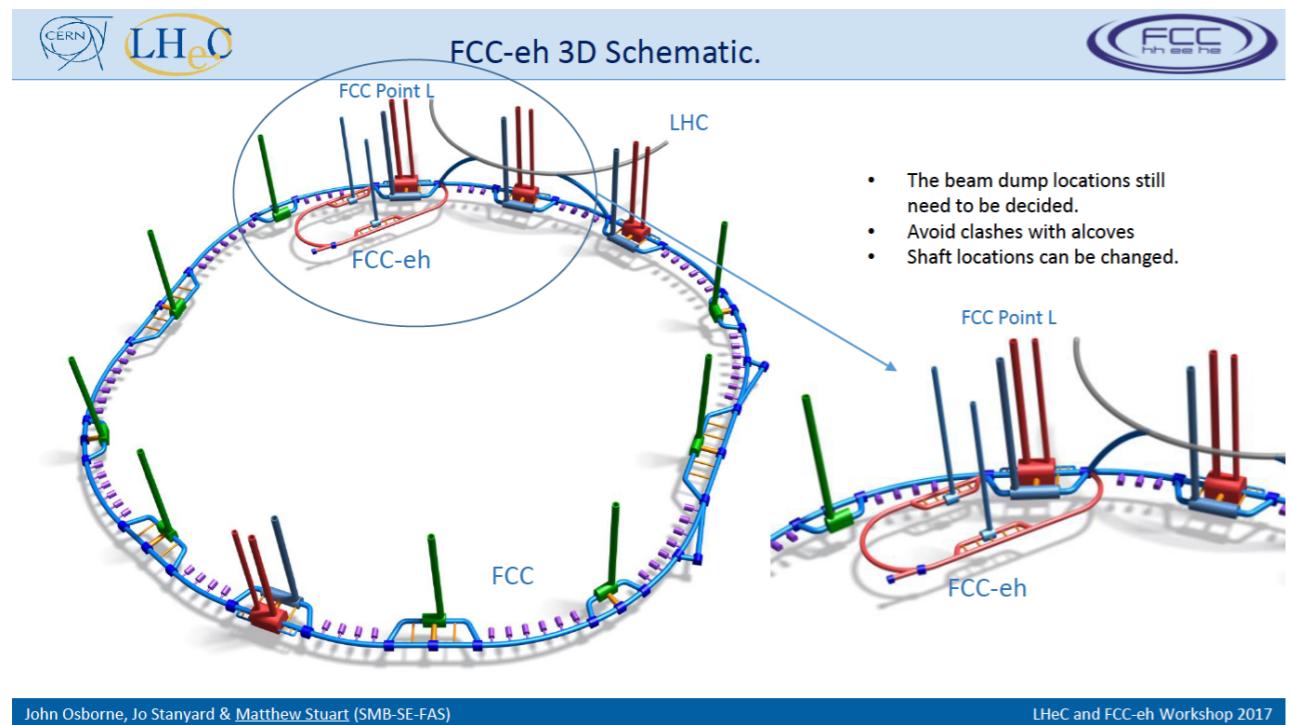


ALI-DER-65282



1506.03981

FCC-eh (I):



John Osborne, Jo Stanyard & Matthew Stuart (SMB-SE-FAS)

LHeC and FCC-eh Workshop 2017

- eA could run either concurrently with pA/AA or in dedicated mode.

parameter [unit]	LHeC (HL-LHC)	eA at HE-LHC	FCC-he
E_{Pb} [PeV] CERN-ACC-2017-0019	0.574	1.03	4.1
E_e [GeV]	60	60	60
$\sqrt{s_{eN}}$ electron-nucleon [TeV]	0.8	1.1	2.2
bunch spacing [ns]	50	50	100
no. of bunches	1200	1200	2072
ions per bunch [10^8]	1.8	1.8	1.8
$\gamma \epsilon_A$ [μm]	1.5	1.0	0.9
electrons per bunch [10^9]	4.67	6.2	12.5
electron current [mA]	15	20	20
IP beta function β_A^* [cm]	7	10	15
hourglass factor H_{geom}	0.9	0.9	0.9
pinch factor H_{b-b}	1.3	1.3	1.3
bunch filling H_{coll}	0.8	0.8	0.8
luminosity [$10^{32} \text{cm}^{-2}\text{s}^{-1}$]	7	18	54
Integrated lumi. in 10 y. (fb^{-1}) ~~	6	15	45

eD at LHeC:

$$L_{eN} = A L_{eA} > \sim 3 \times 10^{31} \text{ cm}^{-2}\text{s}^{-1}$$

(old CDR number)

- 100 times larger luminosity than HERA, full HERA integrated luminosity in less than a month.

Extracting PDFs:

PDFs, or nuclear effects on them, parametrised at initial scale $Q_0 \gg \Lambda_{\text{QCD}}$ employing sum rules (parametrisation biases)

DGLAP evolution, available up to NNLO

PDFs at all required scales

NO \Rightarrow vary parameters

Minimum?

Final PDFs with uncertainties

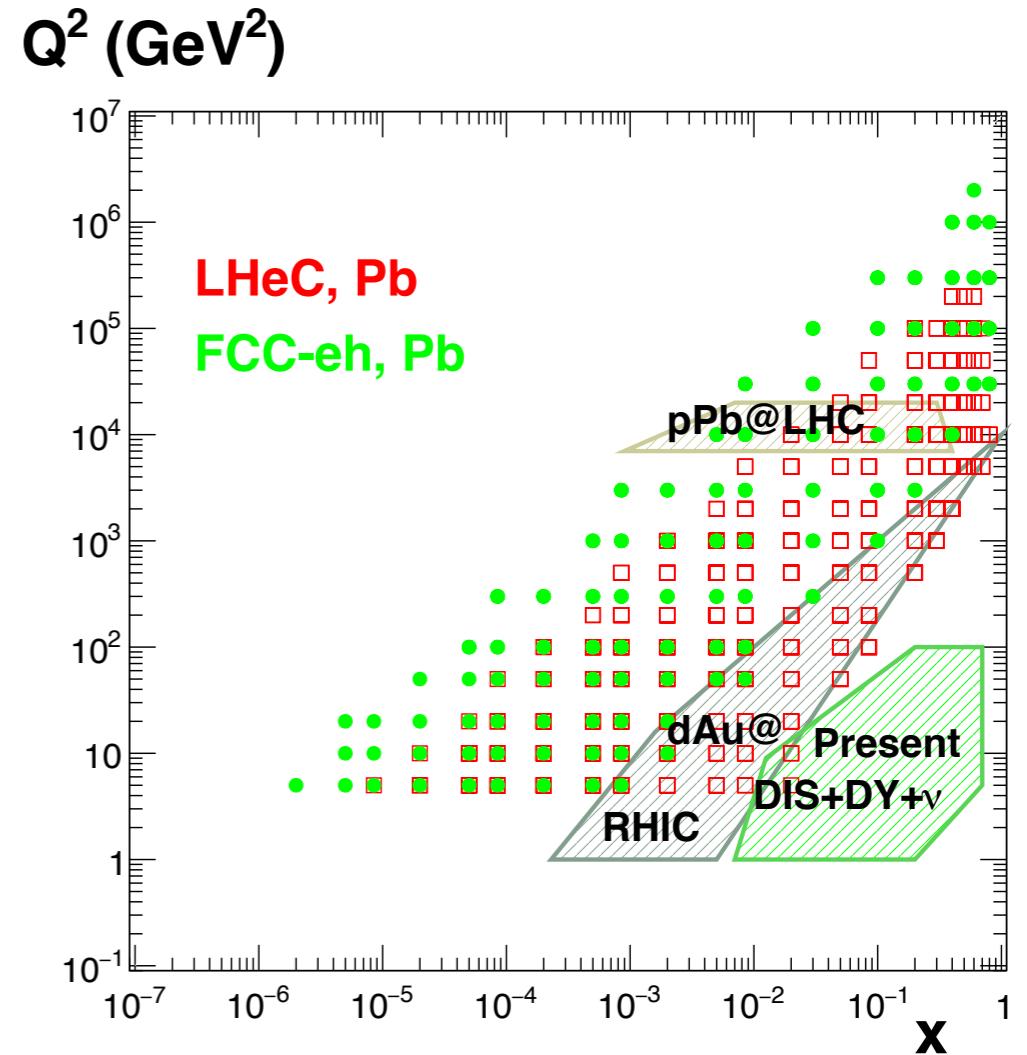
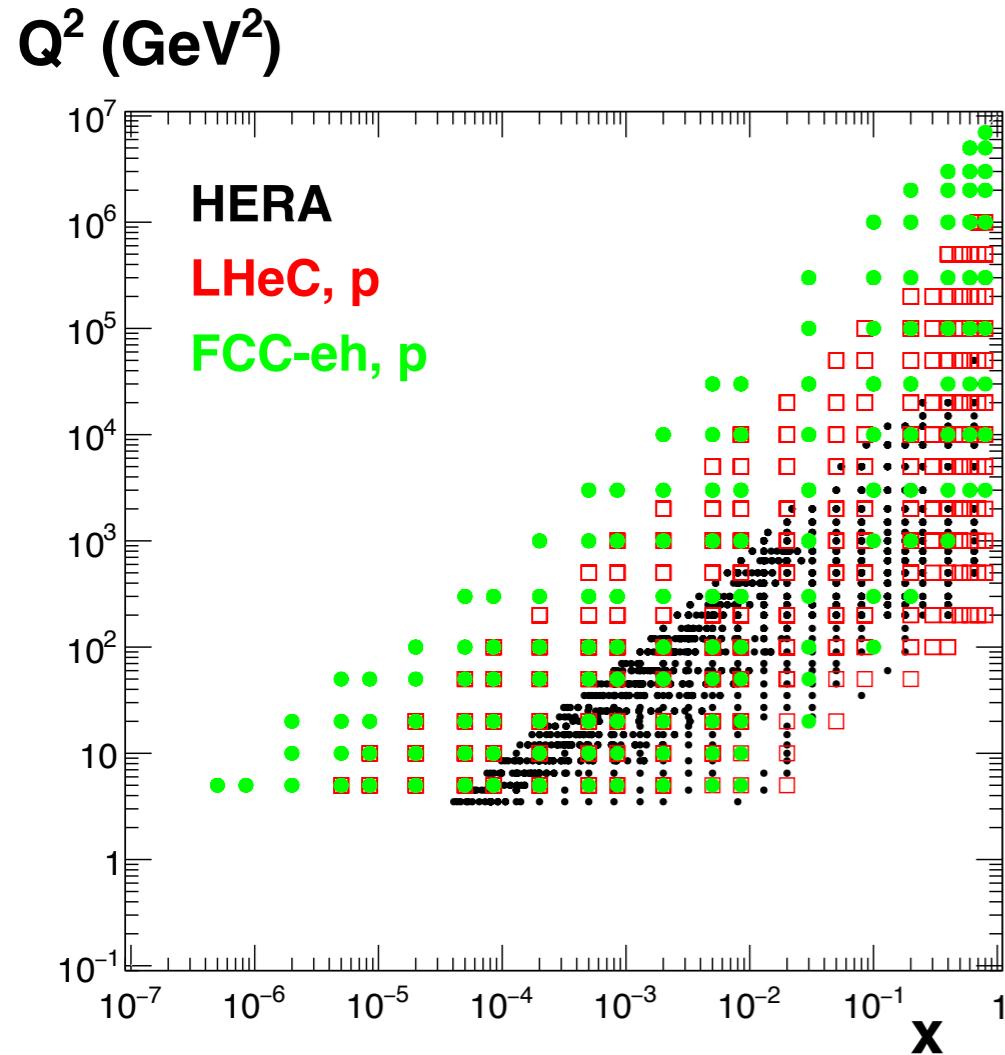
YES

Evaluation of the criterium for comparison data/theory, (treatment of errors, tolerance criteria for different data sets)

Calculation of observables in collinear factorisation, compatible with evolution

Comparison with data that are available and for which pQCD can be considered reliable (e.g. scale dependency)

Pseudodata:



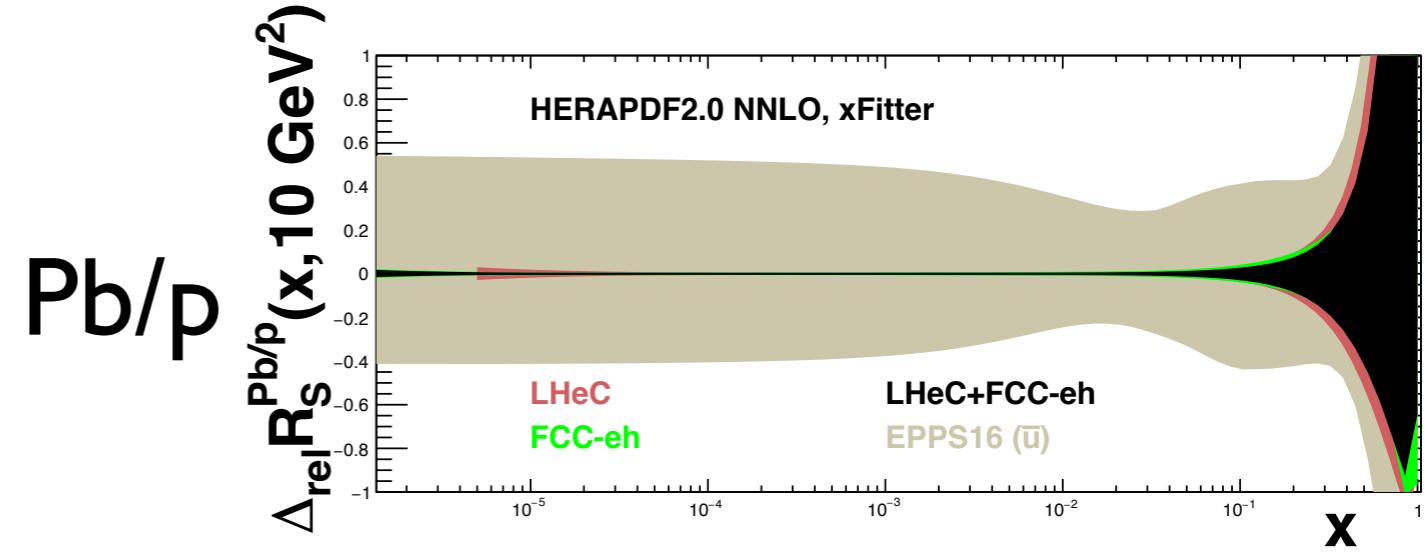
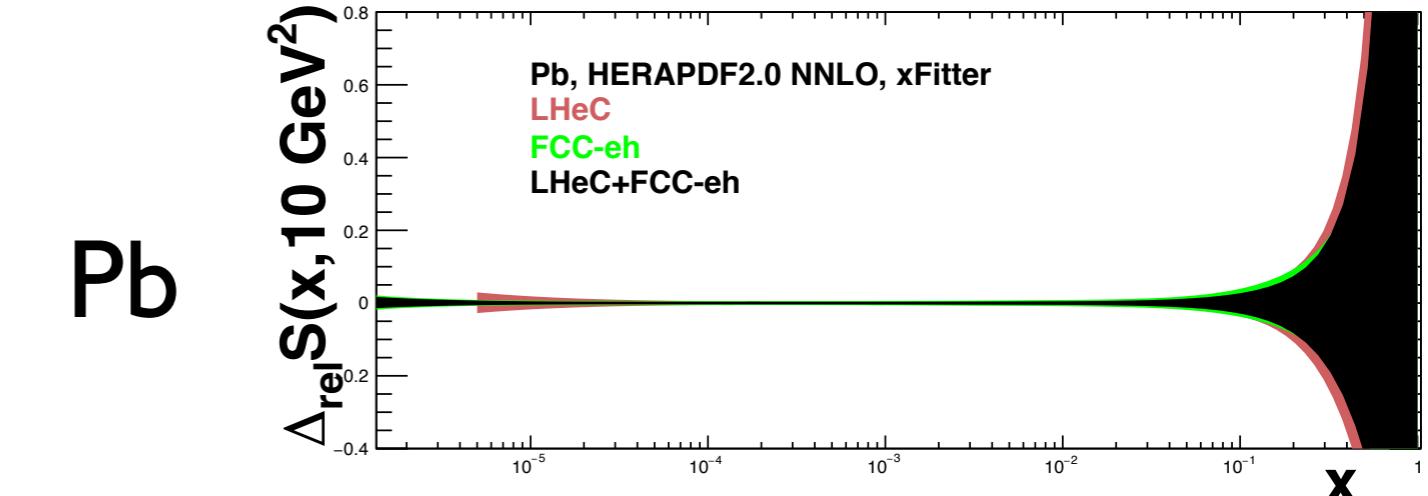
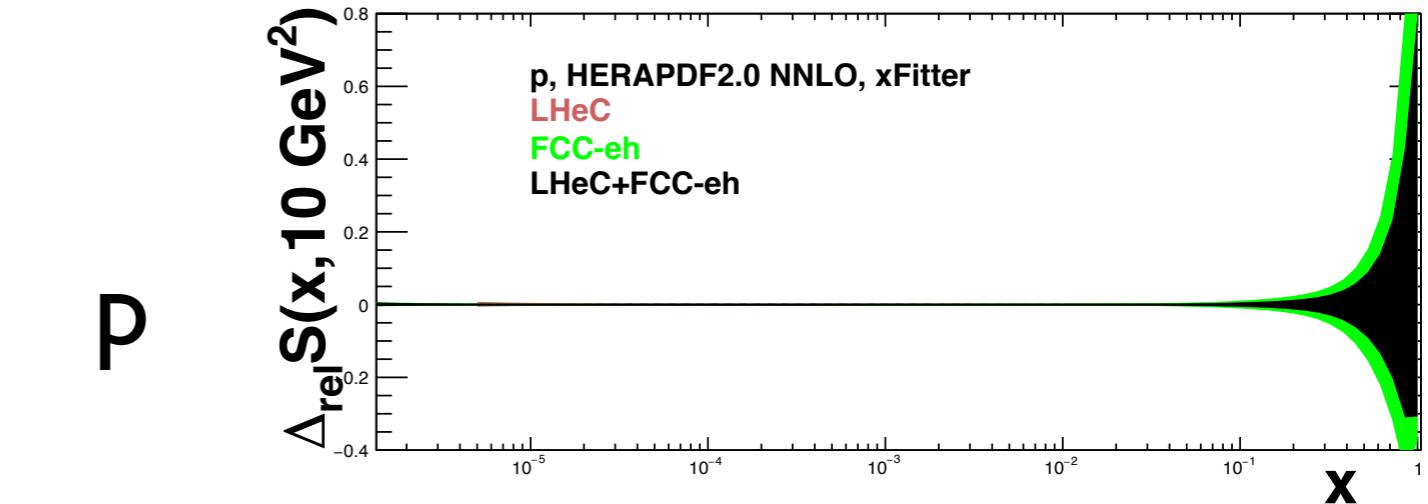
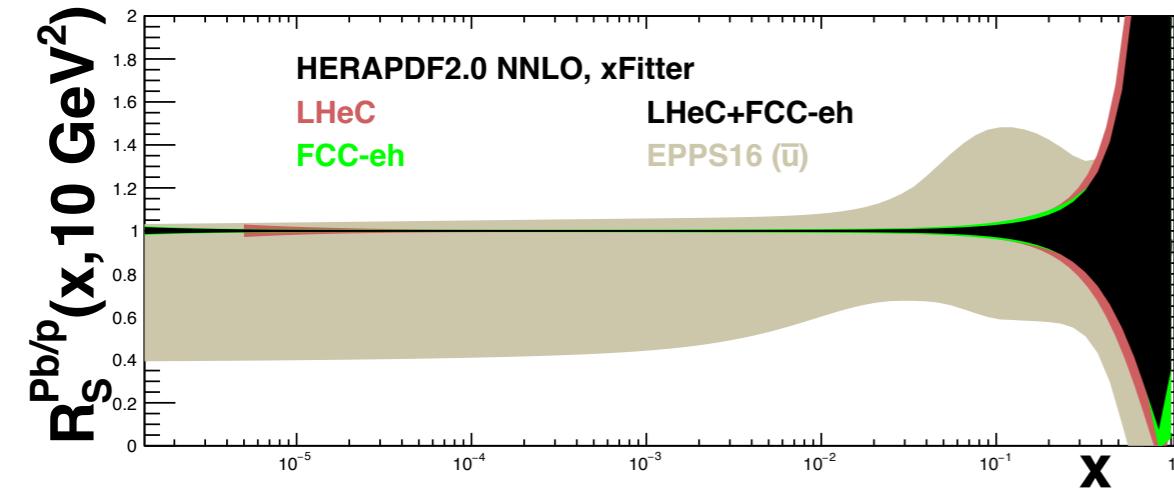
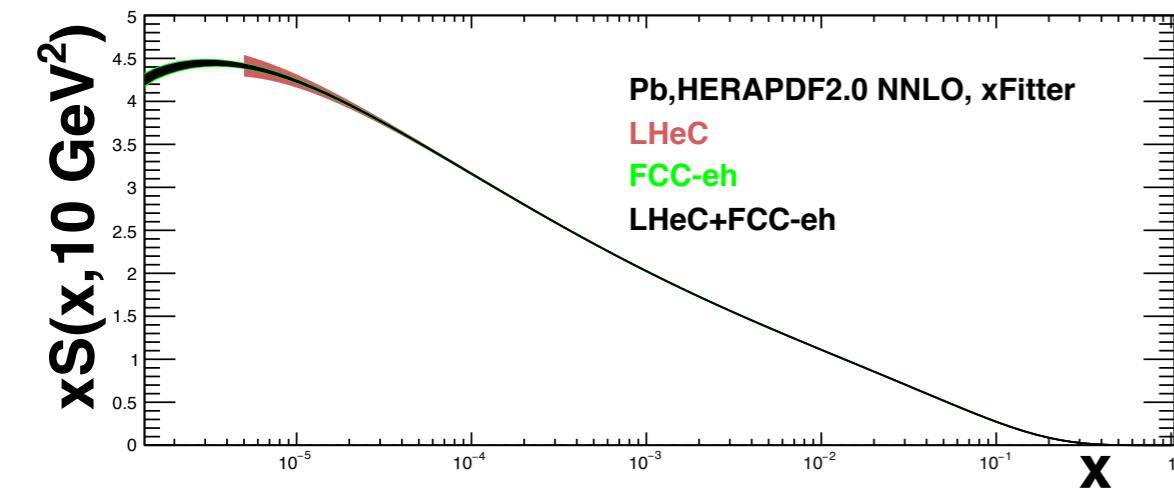
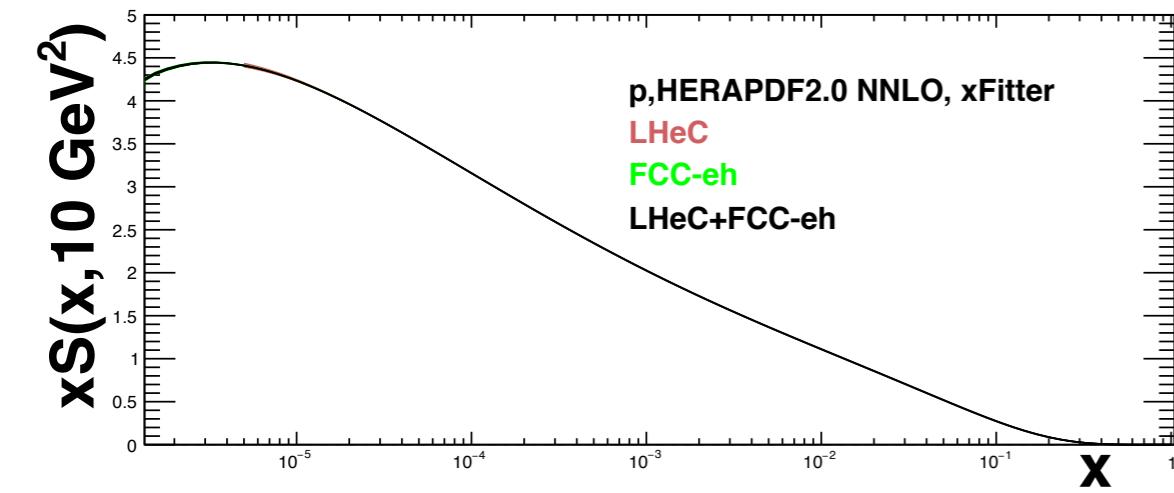
- Pseudodata generated using a code (**Max Klein**) validated with the HI MC.
- Cuts: $|\eta_{\max}|=5$, $0.95 < y < 0.001$.
- Error assumptions \sim factor 2 better than at HERA (luminosity uncertainty kept aside).
- Stat./syst. errors (ePb@FCC-eh) from 0.1/1.2% (small x , NC) to 37/6% (large x & Q^2 , CC).

Source of uncertainty	Error on the source or cross section
scattered electron energy scale	0.1 %
scattered electron polar angle	0.1 mrad
hadronic energy scale	0.5 %
calorimeter noise ($y < 0.01$)	1-3 %
radiative corrections	1-2 %
photoproduction background	1 %
global efficiency error	0.7 %

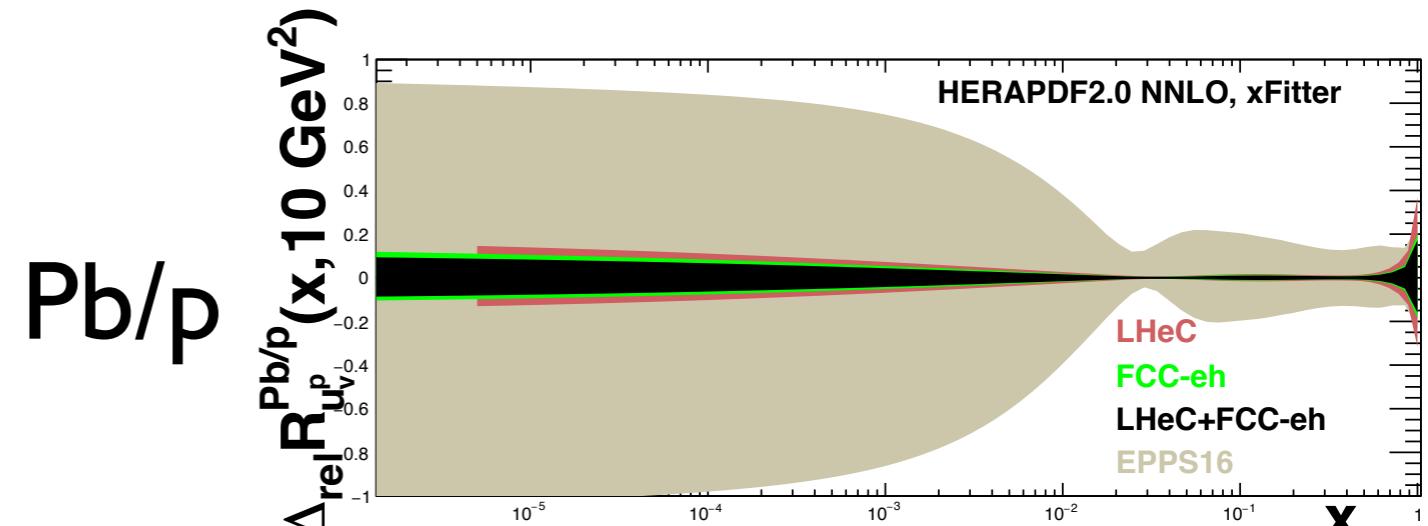
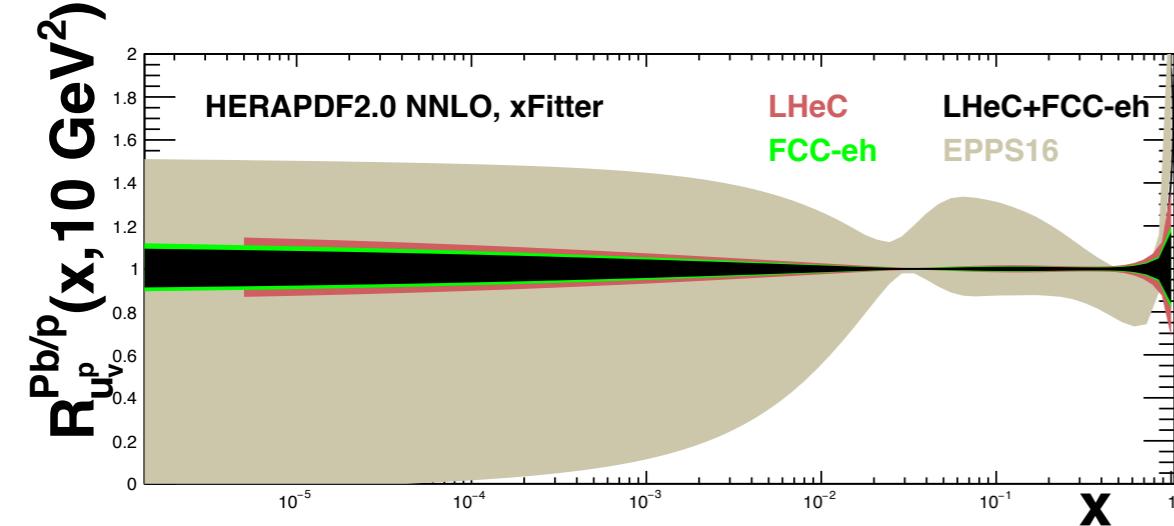
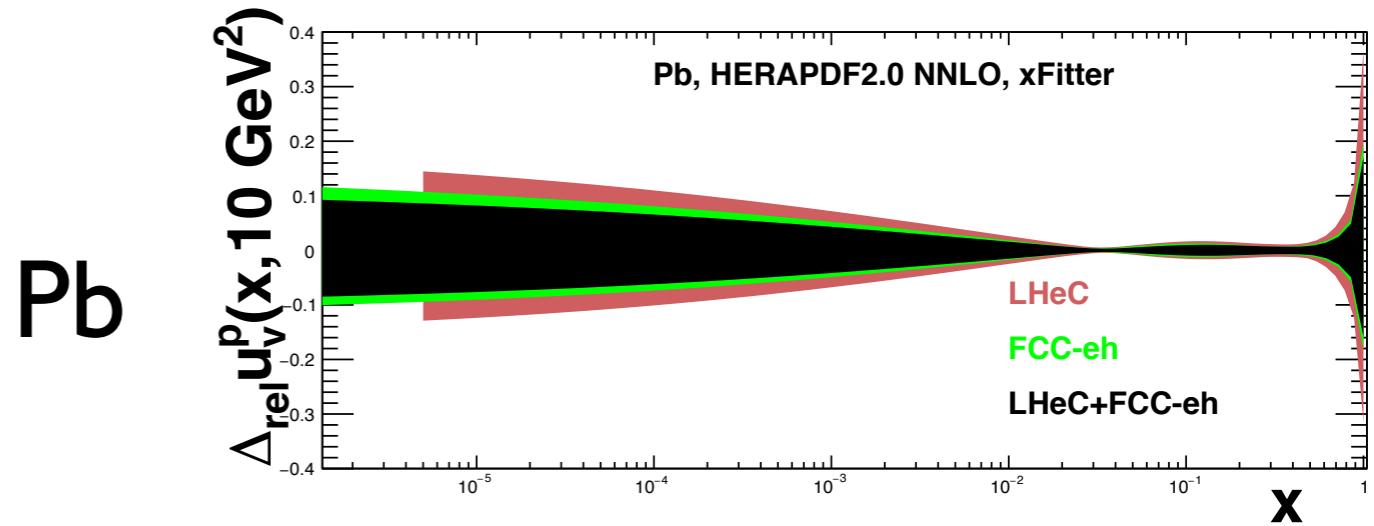
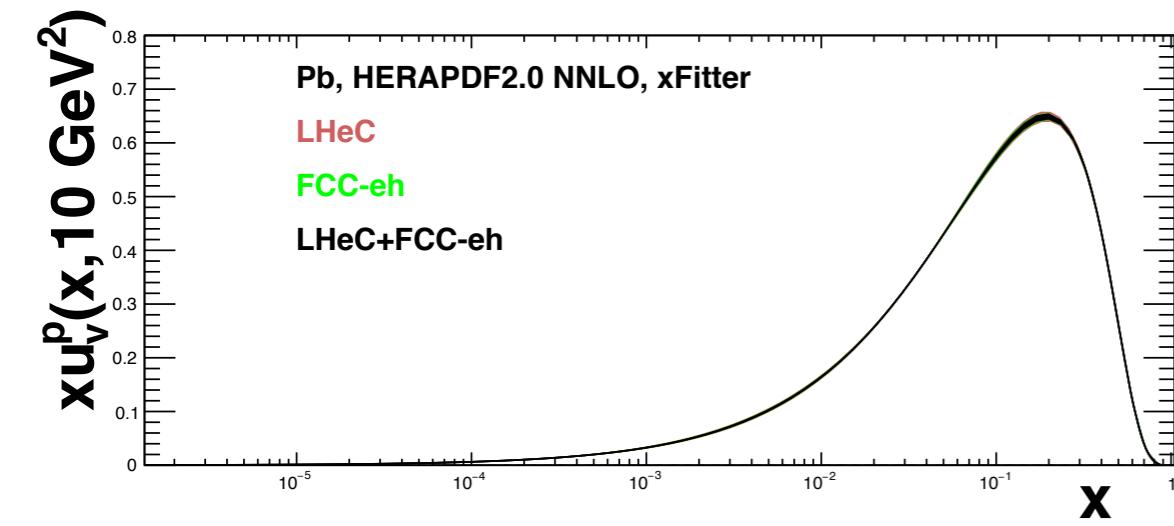
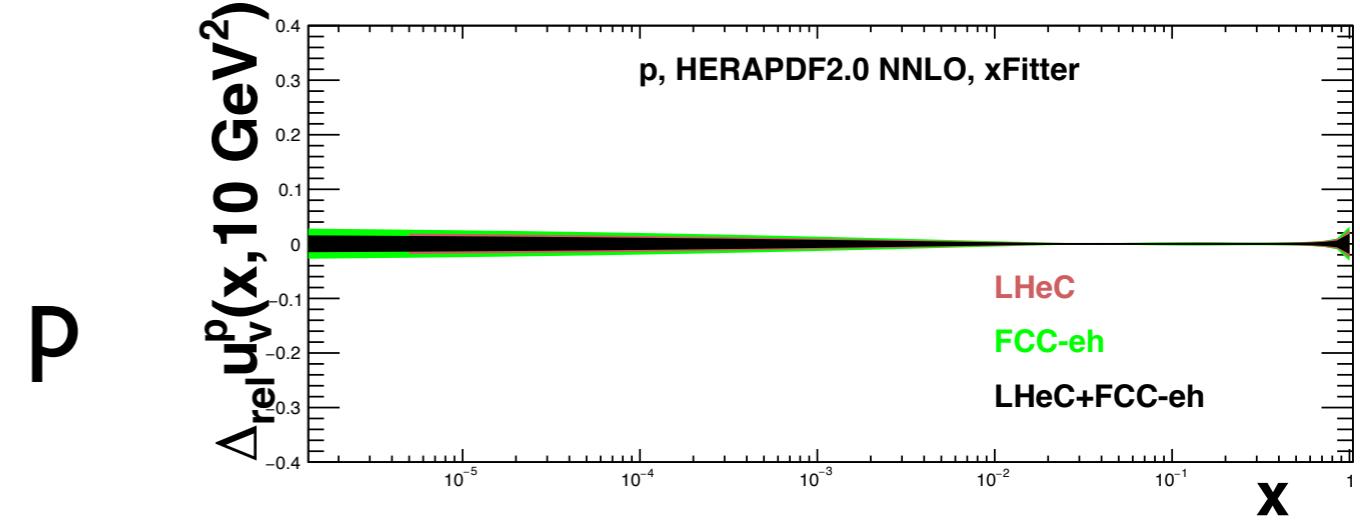
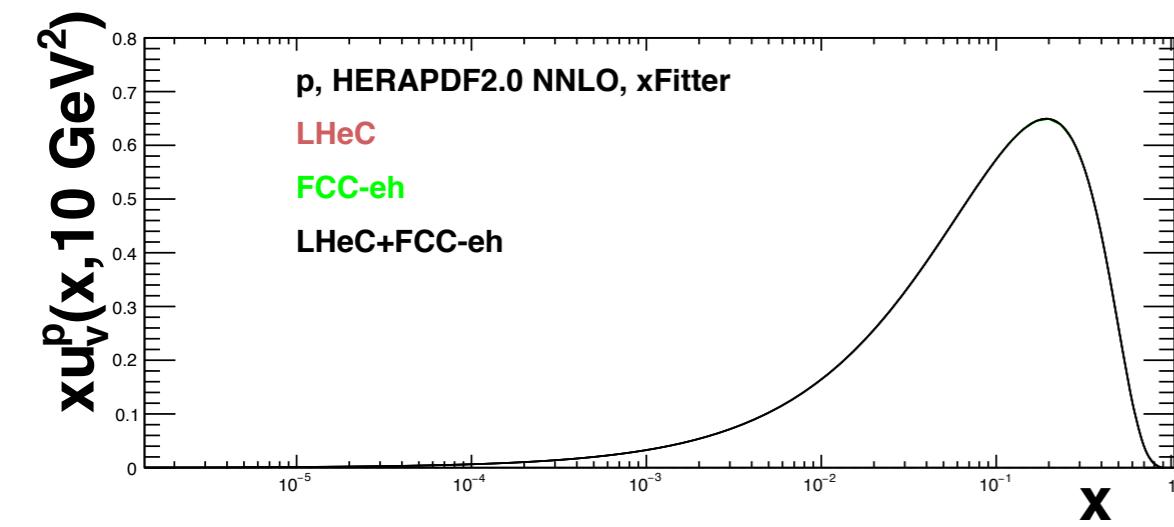
Pseudodata:

	E_e (GeV)	E_h (TeV/nucleon)	Polarisation	Luminosity (fb⁻¹)	NC/CC	# data
ep@LHeC, 1005 data points for Q²≥3.5 GeV²	60 (e ⁻)	1 (p)	0	100	CC	93
	60 (e ⁻)	1 (p)	0	100	NC	136
	60 (e ⁻)	7 (p)	-0.8	1000	CC	114
	60 (e ⁻)	7 (p)	0.8	300	CC	113
	60 (e ⁺)	7 (p)	0	100	CC	109
	60 (e ⁻)	7 (p)	-0.8	1000	NC	159
	60 (e ⁻)	7 (p)	0.8	300	NC	159
	60 (e ⁺)	7 (p)	0	100	NC	157
ePb@LHeC, 484 data points for Q²≥3.5 GeV²	20 (e ⁻)	2.75 (Pb)	-0.8	0.03	CC	51
	20 (e ⁻)	2.75 (Pb)	-0.8	0.03	NC	93
	26.9 (e ⁻)	2.75 (Pb)	-0.8	0.02	CC	55
	26.9 (e ⁻)	2.75 (Pb)	-0.8	0.02	NC	98
	60 (e ⁻)	2.75 (Pb)	-0.8	1	CC	85
	60 (e ⁻)	2.75 (Pb)	-0.8	1	NC	129
ep@FCC-eh, 619 data points for Q²≥3.5 GeV²	20 (e ⁻)	7 (p)	0	100	CC	46
	20 (e ⁻)	7 (p)	0	100	NC	89
	60 (e ⁻)	50 (p)	-0.8	1000	CC	67
	60 (e ⁻)	50 (p)	0.8	300	CC	65
	60 (e ⁺)	50 (p)	0	100	CC	60
	60 (e ⁻)	50 (p)	-0.8	1000	NC	111
	60 (e ⁻)	50 (p)	0.8	300	NC	110
	60 (e ⁺)	50 (p)	0	100	NC	107
ePb@FCC-eh, 150 data points for Q²≥3.5 GeV²	60 (e ⁻)	20 (Pb)	-0.8	10	CC	58
	60 (e ⁻)	20 (Pb)	-0.8	10	NC	101

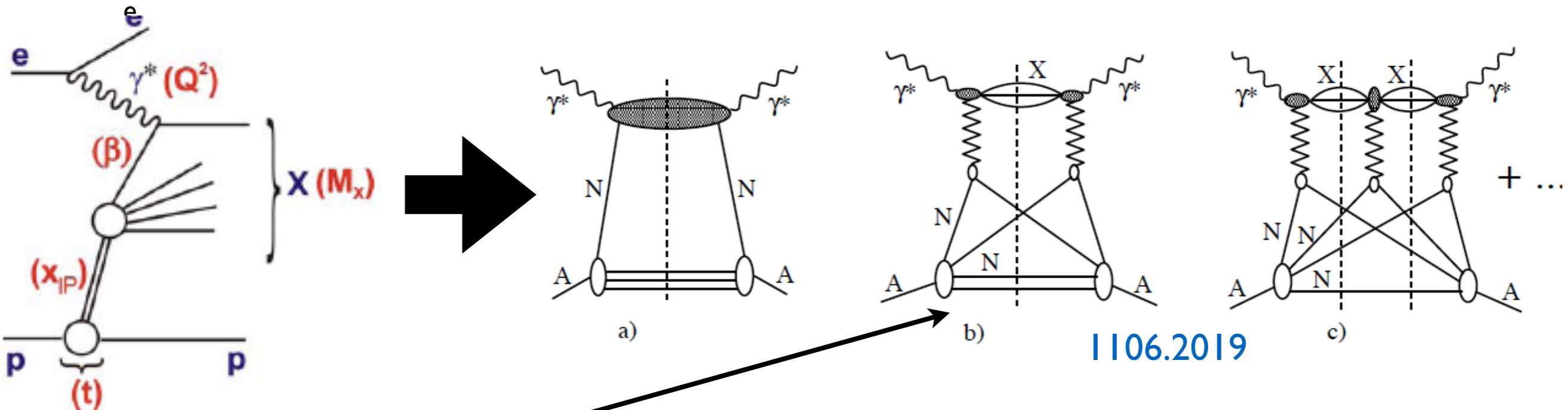
Results: sea



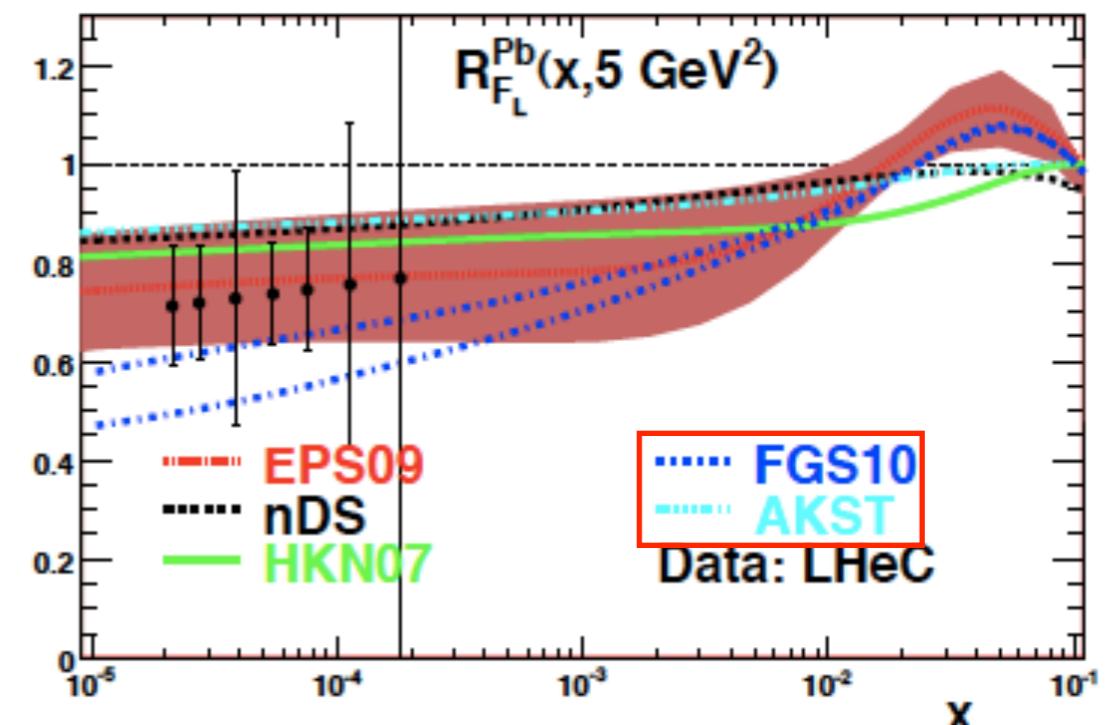
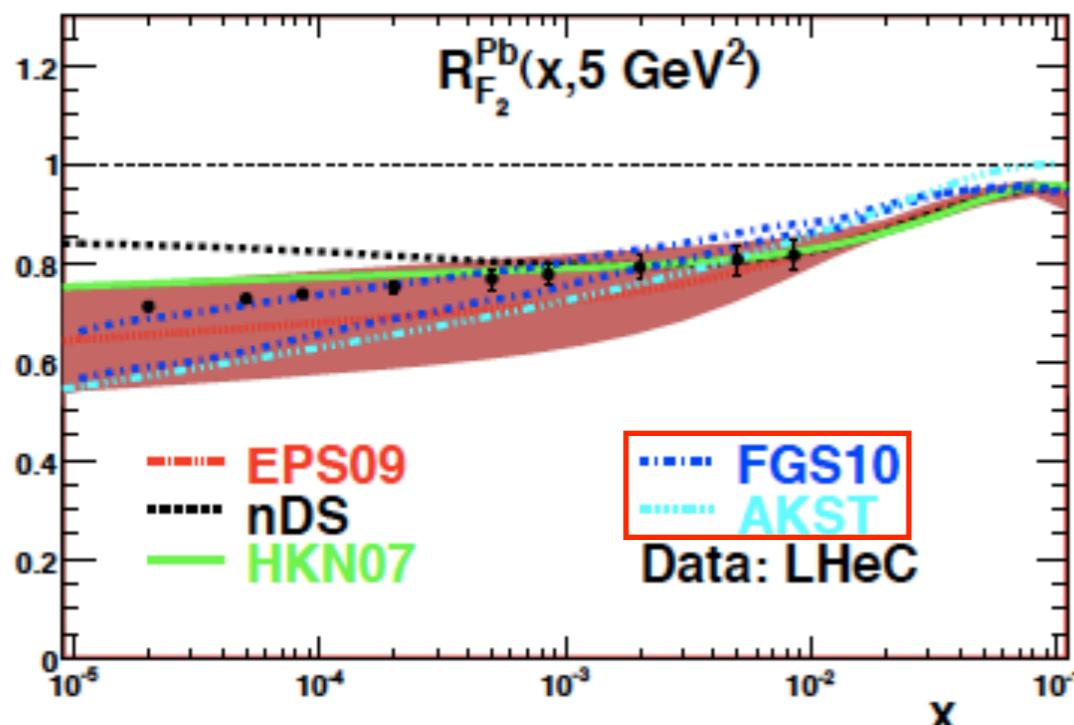
Results: valence



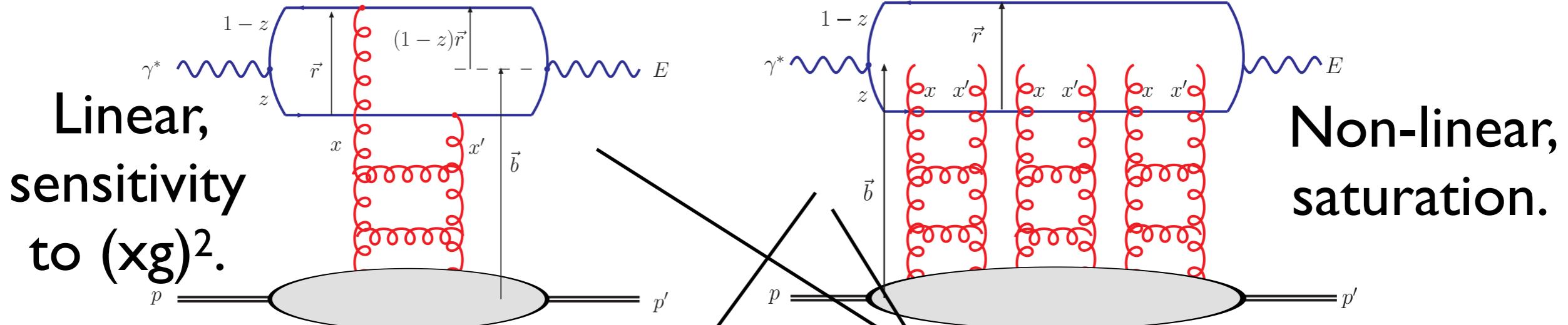
Diffraction in ep and shadowing:



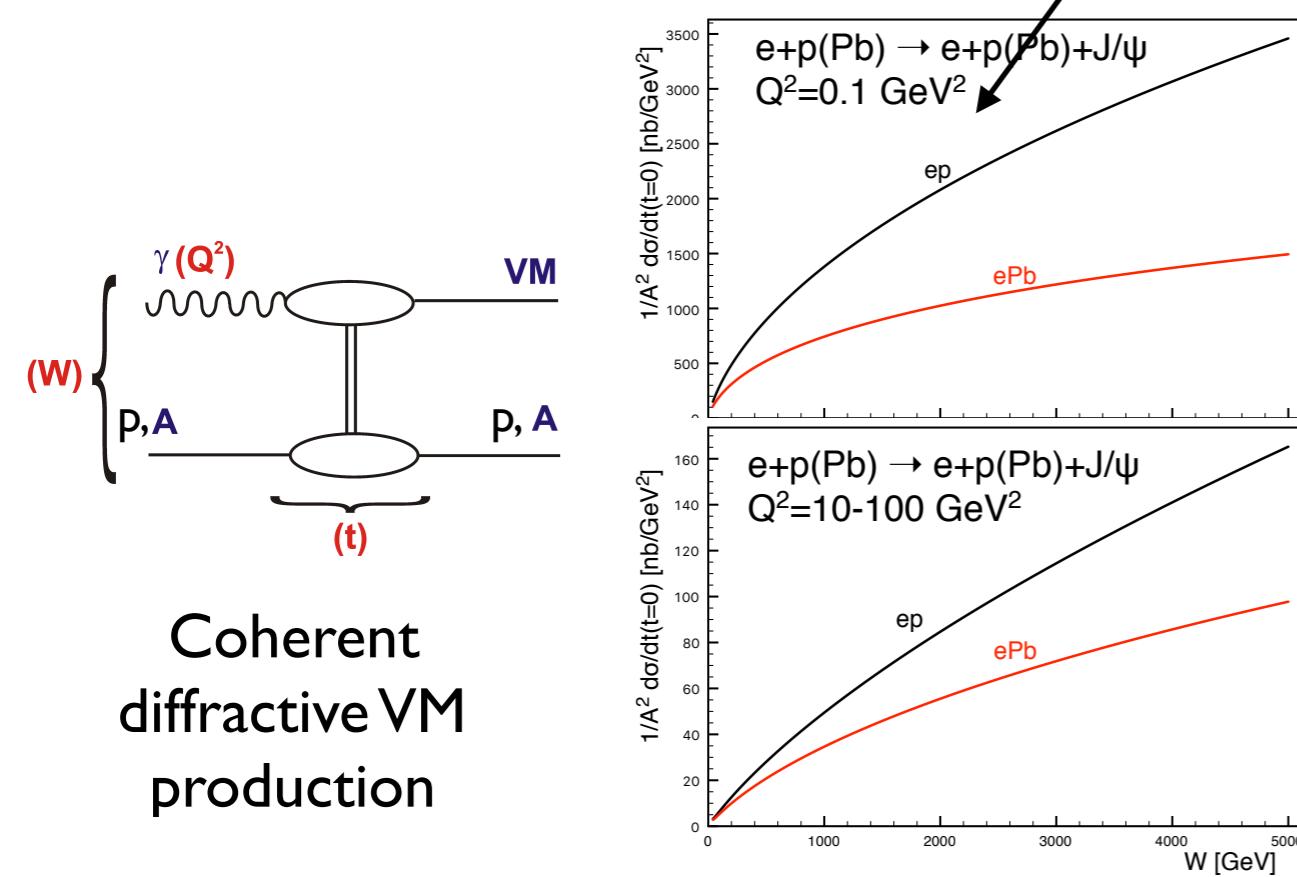
- Diffraction in ep is linked to nuclear shadowing through basic QFT (Gribov): eD to test and set the ‘benchmark’ for new effects.



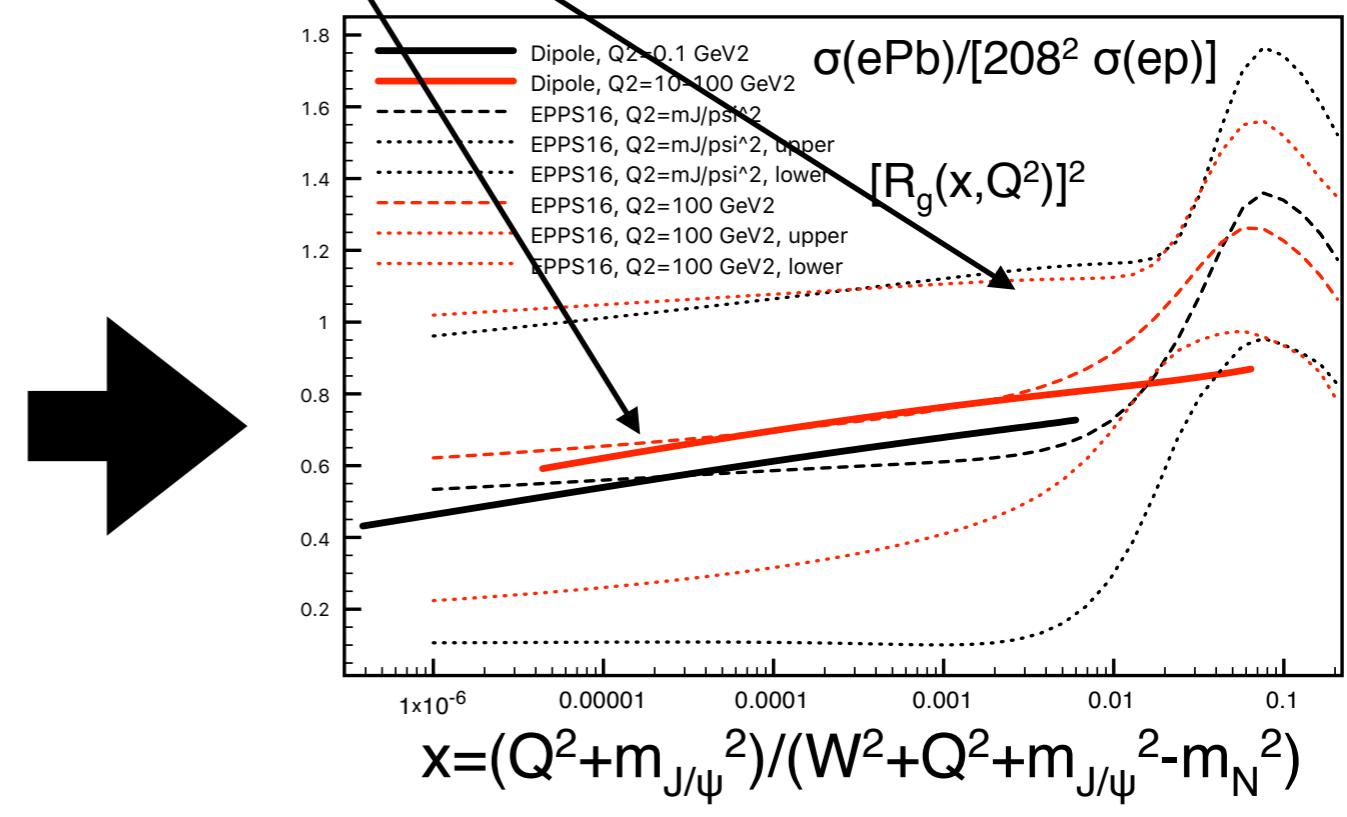
Exclusive VMs:



- The magnitude of nuclear shadowing needs not be different in collinear and non-collinear approaches.



Coherent
diffractive VM
production



Mantysaari, Paukkunen