

# Beyond Standard Model Phenomenology

JiJi Fan

Brown University

1st Workshop on High Energy Theory and Gender,  
CERN, 2018

BSM phenomenology covers a very broad range of topics.  
Impossible to cover in 30 mins.

Select a couple of topics related to the origin of the weak scale and Higgs physics. (the choice of speculative theories discussed in the talk has personal bias)

Not cover neutrino, flavor physics, dark matter:  
talk by Elvira Gamiz, Silvia Pascoli, Tracy Slatyer

In the era of data,

Model driven:

Construct new models with  
new experimental signals

Data driven:

Understand the theoretical  
implications of data as  
much as possible



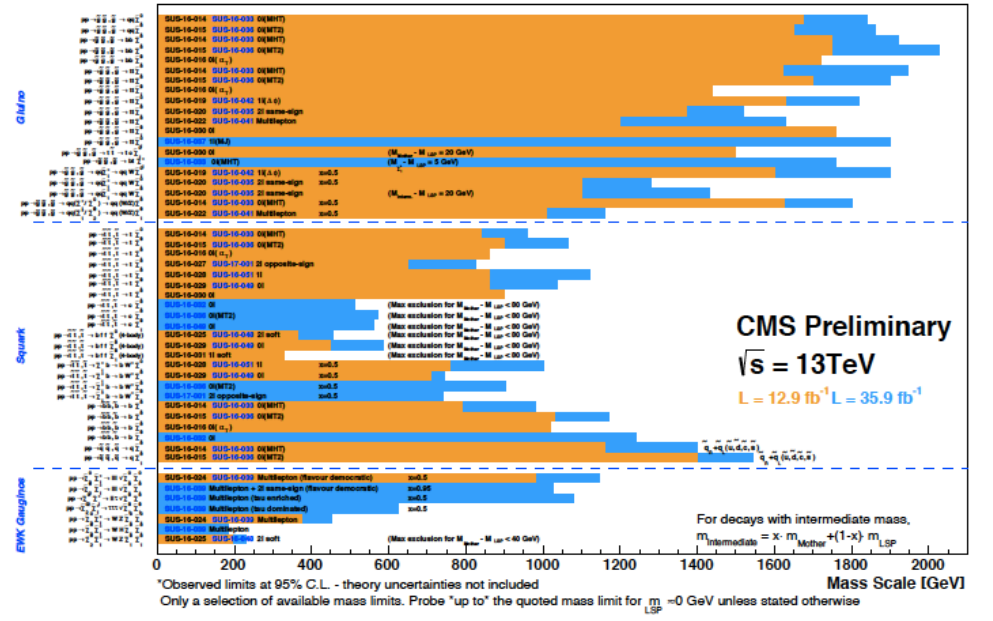
ATLAS SUSY Searches\* - 95% CL Lower Limits  
July 2018

ATLAS Preliminary  
 $\sqrt{s} = 7, 8, 13 \text{ TeV}$

Selected CMS SUSY Results\* - SMS Interpretation

ICHEP '16 - Moriond '17

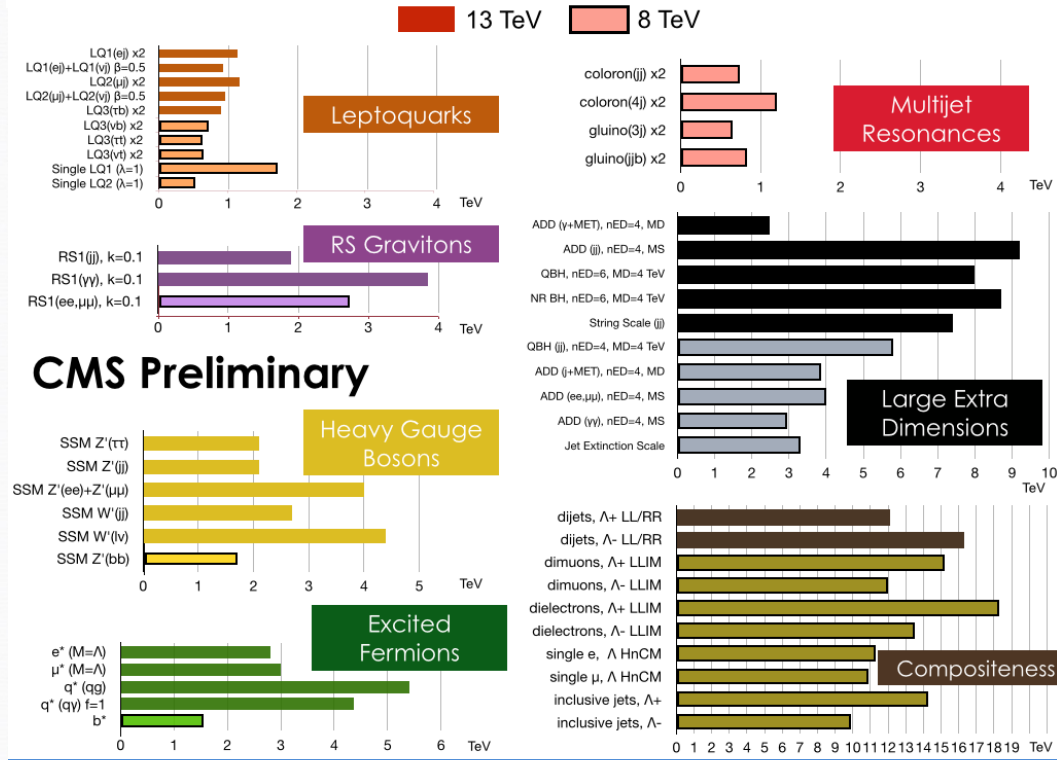
Model	$\epsilon, \mu, \tau, \gamma$	Jets	$E_{miss}^{min}$	$\int \mathcal{L} dt [fb^{-1}]$	Mass limit	Reference		
Inclusive Searches	0 mono-jet	2-6 jets	Yes	36.1	0.9	1.55	1712.0232	
		1-3 jets	Yes	36.1	0.43	0.71	1711.0301	
		0	Yes	36.1	Forbidden	0.95-1.6	2.0	1712.0232
		0	Yes	36.1	1.2	1.85	1706.0371	
3 <sup>rd</sup> gen. squarks direct production	0.2 $\epsilon, \mu$	0.2 jets	Yes	36.1	0.7	1.0	1708.0274	
		0	Yes	36.1	0.4-0.9	0.6-0.8	1709.04183, 1711.11520	
		0	Yes	36.1	0.48-0.84	0.85	1709.04183, 1711.11520	
		0	Yes	36.1	0.46	0.43	1805.1649	
EW direct	0	0	Yes	36.1	0.13-0.23	0.3	1806.0400	
		0	Yes	36.1	0.29-0.88	0.3	1804.0902	
		0	Yes	36.1	0.17	0.6	1403.594, 1806.02293	
		0	Yes	36.1	0.26	0.76	1501.0710	
Long-lived particles	0	0	Yes	36.1	0.15	0.46	1712.0819	
		0	Yes	36.1	1.6	2.4	1710.0491, 1604.04520	
		0	Yes	36.1	0.44	1.3	1607.08079	
		0	Yes	36.1	1.05	2.0	1504.05162	
RPV	0	0	Yes	36.1	0.82	1.33	1904.03568	
		0	Yes	36.1	1.05	1.9	1904.03568	
		0	Yes	36.1	0.95	1.05	1710.0371	
		0	Yes	36.1	0.92	0.61	1710.0371	



\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on

ATLAS Exotics Searches\* - 95% CL Upper Exclusion Limits  
Status: July 2018

ATLAS Preliminary  
 $\sqrt{s} = 8, 13 \text{ TeV}$



Model	$\epsilon, \mu, \tau, \gamma$	Jets	$E_{miss}^{min}$	$\int \mathcal{L} dt [fb^{-1}]$	Limit	Reference					
Extra dimensions	0	ADD $G_{KK} + g/g$	1-4	Yes	36.1	$M_{KK} = 7.7 \text{ TeV}$	1711.0301				
		ADD non-resonant $\gamma\gamma$	2 $\gamma$	Yes	36.1	$M_{KK} = 8.6 \text{ TeV}$	1703.09127				
		ADD OBH	2 $\gamma$	Yes	36.1	$M_{KK} = 6.9 \text{ TeV}$	1606.02285				
		ADD BH high $\sum p_T$	$\geq 1 \epsilon, \mu$	Yes	36.1	$M_{KK} = 8.2 \text{ TeV}$	1512.02586				
		ADD LH multijet	$\geq 3 \gamma$	Yes	36.1	$M_{KK} = 8.55 \text{ TeV}$	1707.24149				
		RS1 $G_{KK} \rightarrow \gamma\gamma$	2 $\gamma$	Yes	36.1	$M_{KK} = 4.1 \text{ TeV}$	1801.08992				
		Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel	Yes	36.1	$M_{KK} = 2.3 \text{ TeV}$	CERN-EP-2018-179				
		Bulk RS $g_{KK} \rightarrow tt$	$1 \epsilon, \mu \geq 1 b, \geq 1 J/2 \gamma$	Yes	36.1	$M_{KK} = 3.8 \text{ TeV}$	1804.10823				
		2UED / RPP	$1 \epsilon, \mu \geq 2 b, \geq 3 \gamma$	Yes	36.1	$M_{KK} = 1.8 \text{ TeV}$	1803.09678				
		Gauge bosons	0	SSM $Z' \rightarrow \ell\ell$	2 $\ell$	Yes	36.1	$M_{Z'} = 4.5 \text{ TeV}$	1707.02424		
SSM $Z' \rightarrow \tau\tau$	2 $\tau$			Yes	36.1	$M_{Z'} = 2.42 \text{ TeV}$	1709.07242				
Leptophobic $Z' \rightarrow bb$	2 $b$			Yes	36.1	$M_{Z'} = 2.1 \text{ TeV}$	1805.09299				
RS1 $G_{KK} \rightarrow \gamma\gamma$	$1 \epsilon, \mu \geq 1 b, \geq 1 J/2 \gamma$			Yes	36.1	$M_{KK} = 3.0 \text{ TeV}$	1604.10823				
SSM $W' \rightarrow \ell\nu$	$1 \epsilon, \mu$			Yes	79.8	$M_{W'} = 3.7 \text{ TeV}$	ATLAS-CONF-2018-017				
SSM $W' \rightarrow \nu\nu$	1 $\nu$			Yes	36.1	$M_{W'} = 5.6 \text{ TeV}$	1801.06992				
HVT $V' \rightarrow WW \rightarrow qqqq$ model B	$0 \epsilon, \mu, 2 J$			Yes	79.8	$M_{V'} = 4.15 \text{ TeV}$	ATLAS-CONF-2018-016				
HVT $V' \rightarrow WH/ZH$ model B	multi-channel			Yes	36.1	$M_{V'} = 2.93 \text{ TeV}$	1712.06518				
LRSM $W' \rightarrow tb$	multi-channel			Yes	36.1	$M_{W'} = 3.25 \text{ TeV}$	CERN-EP-2018-142				
CI	0			CI $qqqq$	2 $q$	Yes	36.1	$M_{CI} = 21.8 \text{ TeV}$	1703.09127		
		CI $tttt$	2 $t$	Yes	36.1	$M_{CI} = 40.0 \text{ TeV}$	1707.02424				
		CI $tttt$	$\geq 1 \epsilon, \mu \geq 1 b, \geq 1 J$	Yes	36.1	$M_{CI} = 2.57 \text{ TeV}$	CERN-EP-2018-174				
		DM	0	Axial-vector mediator (Dirac DM)	$0 \epsilon, \mu, 1-4 J$	Yes	36.1	$M_{DM} = 1.55 \text{ TeV}$	1711.0301		
				Colored scalar mediator (Dirac DM)	$0 \epsilon, \mu, 1-4 J$	Yes	36.1	$M_{DM} = 1.67 \text{ TeV}$	1711.0301		
				$W_{\chi_{1,2}}$ EFT (Dirac DM)	$0 \epsilon, \mu, 1 J, \leq 1 J$	Yes	32	$M_{DM} = 700 \text{ GeV}$	1608.02372		
				LO	0	Scalar LO 1 <sup>st</sup> gen	2 $e, \mu \geq 2 J$	Yes	32	$M_{LO} = 1.1 \text{ TeV}$	1605.06035
						Scalar LO 2 <sup>nd</sup> gen	2 $\mu, \geq 2 J$	Yes	32	$M_{LO} = 1.05 \text{ TeV}$	1605.06035
						Scalar LO 3 <sup>rd</sup> gen	$1 \epsilon, \mu, \geq 1 b, \geq 3 J$	Yes	20.3	$M_{LO} = 840 \text{ GeV}$	1508.04735
						Excited fermions/heavy quarks	0	VLO $TT \rightarrow H/Z/\gamma/Wb + X$	multi-channel	Yes	36.1
VLO $BB \rightarrow Wt/Zb + X$	multi-channel							Yes	36.1	$M_{exc} = 1.34 \text{ TeV}$	ATLAS-CONF-2018-032
VLO $T_{33} T_{33} \rightarrow Wt + X$	$2(S)_{33} \geq 3 \mu, \geq 1 b, \geq 1 J$							Yes	36.1	$M_{exc} = 1.64 \text{ TeV}$	CERN-EP-2018-171
VLO $B \rightarrow Hb + X$	$1 \epsilon, \mu, \geq 1 b, \geq 1 J$							Yes	32	$M_{exc} = 1.44 \text{ TeV}$	ATLAS-CONF-2018-072
VLO $B \rightarrow Hb + X$	$0 \epsilon, \mu, 2 \gamma, \geq 1 b, \geq 1 J$	Yes	79.8					$M_{exc} = 1.21 \text{ TeV}$	ATLAS-CONF-2018-024		
VLO $QQ \rightarrow Wq/Wq$	$1 \epsilon, \mu, \geq 4 J$	Yes	20.3					$M_{exc} = 690 \text{ GeV}$	1509.04261		
Excited fermions/leptons	0	Excited quark $q^* \rightarrow qg$	2 $g$					Yes	37.0	$M_{exc} = 8.0 \text{ TeV}$	1703.09127
		Excited quark $q^* \rightarrow q\gamma$	1 $\gamma, 1 J$	Yes	36.1			$M_{exc} = 5.3 \text{ TeV}$	1709.10440		
		Excited quark $b^* \rightarrow bg$	1 $b, 1 J$	Yes	36.1			$M_{exc} = 2.6 \text{ TeV}$	1805.09299		
		Excited lepton $l^* \rightarrow l\gamma$	3 $e, \mu, 1 J$	Yes	20.3			$M_{exc} = 3.0 \text{ TeV}$	1411.2921		
		Excited lepton $\nu^*$	3 $e, \mu, \tau$	Yes	20.3	$M_{exc} = 1.6 \text{ TeV}$	1411.2921				
		Other	0	Type III Seesaw	$1 \epsilon, \mu, \geq 2 J$	Yes	79.8	$M_{exc} = 560 \text{ GeV}$	ATLAS-CONF-2018-020		
				LRSM Majorana $\nu$	2 $e, \mu, 2 J$	Yes	36.1	$M_{exc} = 870 \text{ GeV}$	1506.06020		
				Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	2, 3 $e, \mu, \tau$ (SS)	Yes	36.1	$M_{exc} = 700 \text{ GeV}$	1710.09748		
				Higgs triplet $H^{\pm\pm} \rightarrow \tau\tau$	3 $e, \mu, \tau$	Yes	20.3	$M_{exc} = 400 \text{ GeV}$	1411.2921		
				Monopole (non-res prod)	1 $e, \mu, 1 b$	Yes	20.3	$M_{exc} = 657 \text{ GeV}$	1504.5404		
Multi-charged particles	1 $e, \mu, 1 b$			Yes	20.3	$M_{exc} = 785 \text{ GeV}$	1504.04188				
Magnetic monopoles	1 $e, \mu, 1 b$			Yes	7.0	$M_{exc} = 1.34 \text{ TeV}$	1509.08059				

\*Only a selection of the available mass limits on new states or phenomena is shown.  
†Small-radius (large-radius) jets are denoted by the letter J.

There is no confirmed signal of new physics yet. Yet before forming any strong opinion, it's worthwhile to know a bit what LHC has excluded and what are the implications of the LHC results for big physics questions such as the origin of weak scale.

A two-sentence summary of LHC results:

Strongly-interacting particles (colored particles) are strongly constrained.

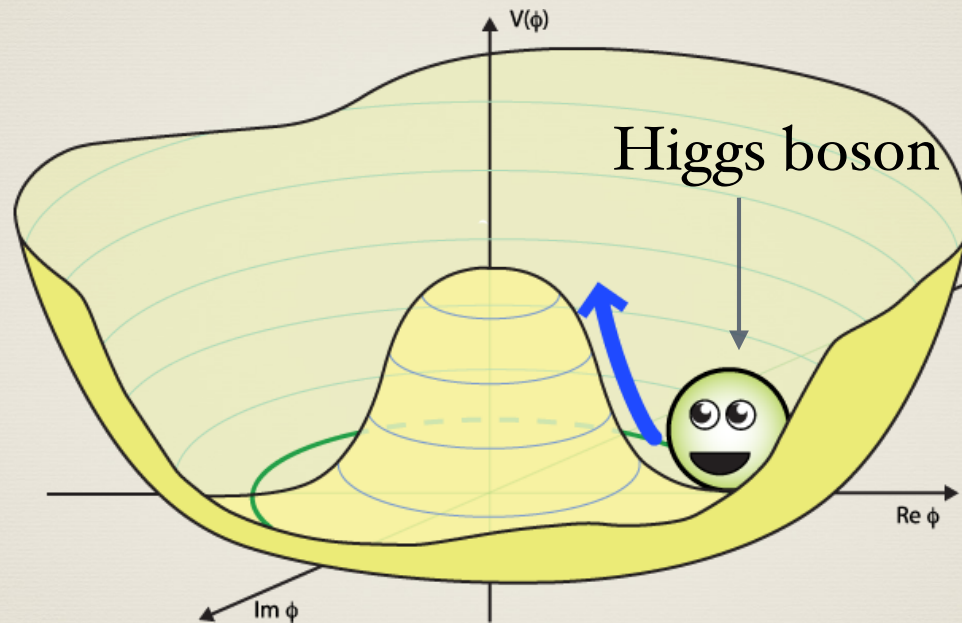
E.g., gluinos  $\sim 2$  TeV; scalar or fermionic top partners  $\sim (1 - 1.5)$  TeV;

Relatively weak constraints on weakly-interacting particles:  
depend on the final states.

(one example with essentially no constraints beyond LEP will be discussed at the end).

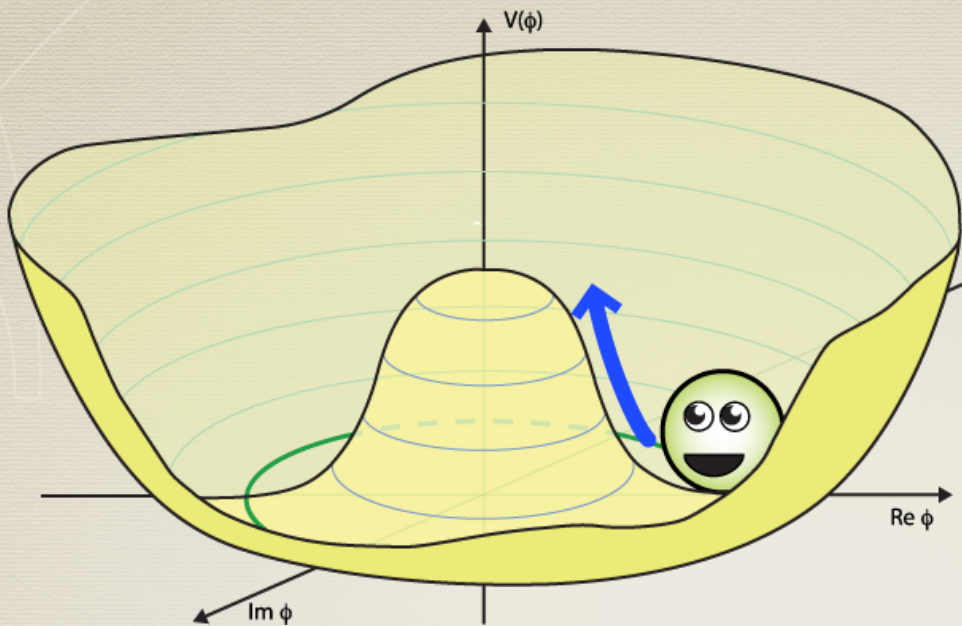
# Higgs Physics in a Nutshell

Higgs potential energy



©P. Tanedo





©P. Tanedo

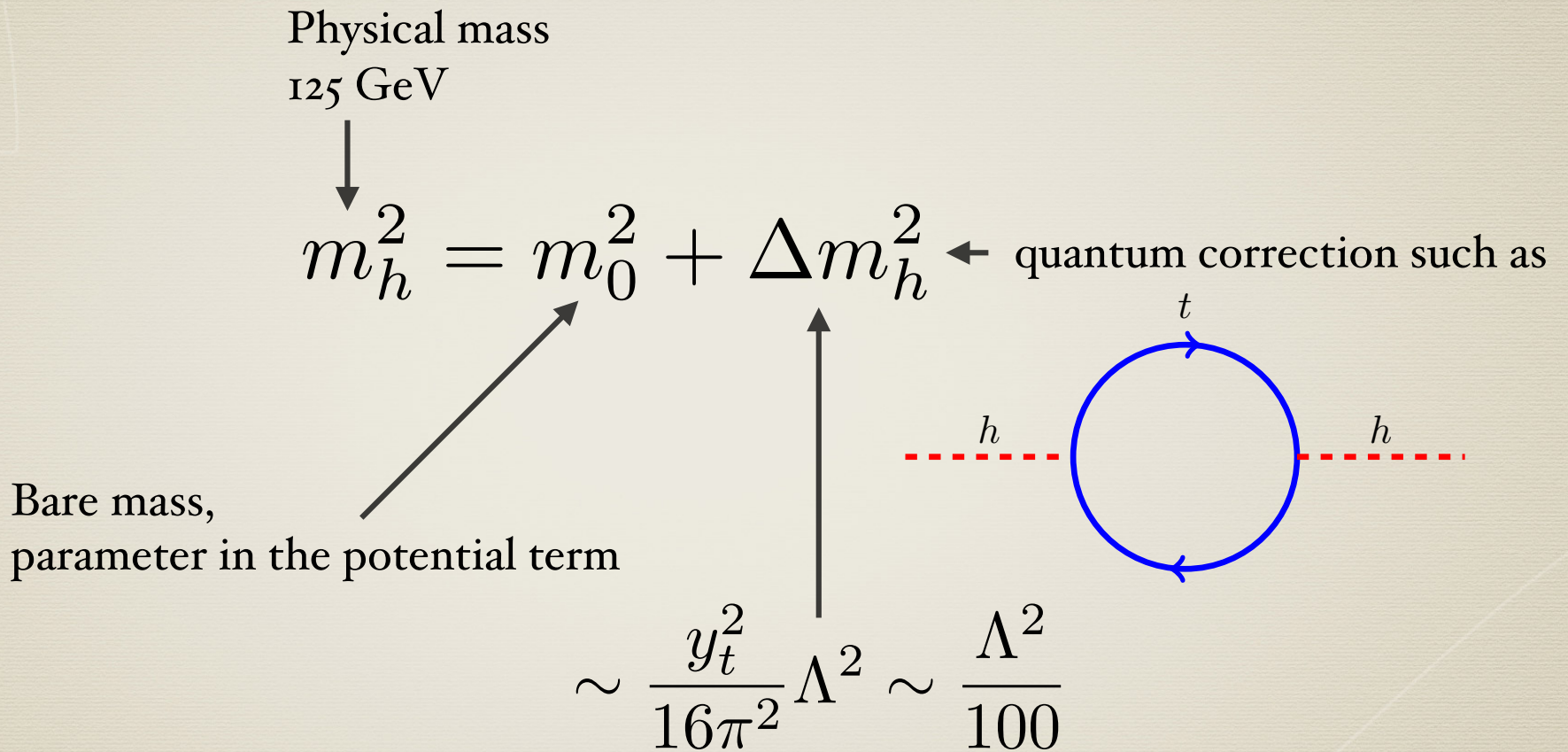
The standard model *assumes* this potential but doesn't *explain* it.

The standard model doesn't *explain* the Higgs mass itself.

It is only an **effective** description of electroweak symmetry breaking. Yet the microscopic details aren't specified.

What we really want is a **dynamical** explanation: what are the interactions driving the preference for a nonzero vacuum expectation value?

# Hierarchy Problem of an Elementary Scalar (fine-tuning problem)



$\Lambda$ : scale up to which SM is valid or **the scale of new physics that dynamically generates the Higgs potential**

# Hierarchy Problem of an Elementary Scalar (fine-tuning problem)

Physical mass  
125 GeV  $\longrightarrow$   $m_h^2 = m_0^2 + \Delta m_h^2$   $\longleftarrow$  quantum correction

Bare mass,  
parameter in the potential term

$$\sim \frac{y_t^2}{16\pi^2} \Lambda^2 \sim \frac{\Lambda^2}{100}$$

$\Lambda$ : scale up to which SM is valid or **the scale of new physics that dynamically generates the Higgs potential**

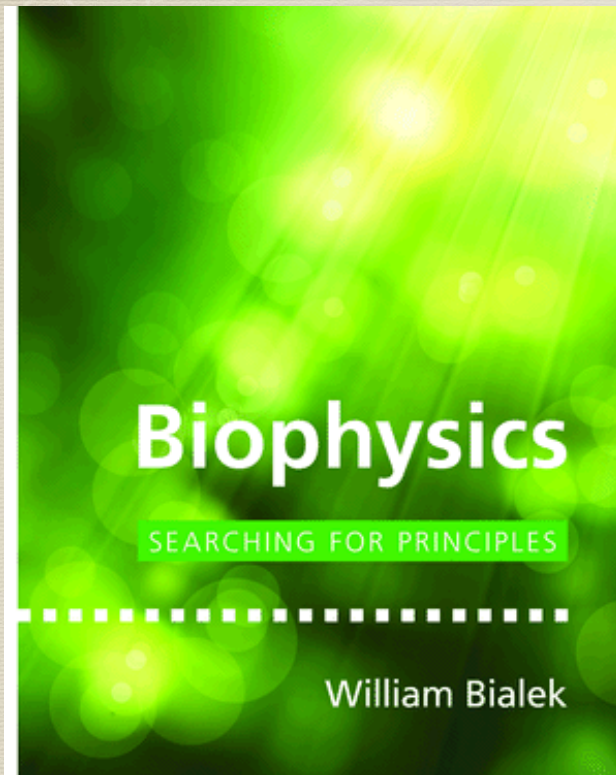
Suppose  $\Lambda = 10^{19}$  GeV, and the observed Higgs mass is 125 GeV, we need, say, a huge bare mass to cancel the quantum fluctuations

$$m_0^2 = (1,500,473, 789, 254, 211,536 \text{ GeV})^2;$$

If we miss by 10 GeV,  $m_0^2 = (1,500,473, 789, 254, 211,526 \text{ GeV})^2;$

The physical Higgs mass is  $\sim 10^9$  GeV!

No fine-tuning is a possible candidate principle not only for particle physics



## PART II CANDIDATE PRINCIPLES

### 4. Noise Is Not Negligible 127

4.1 Fluctuations and Chemical Reactions 127

4.2 Motility and Chemotaxis in Bacteria 149

4.3 Molecule Counting, More Generally 172

4.4 More about Noise in Perception 192

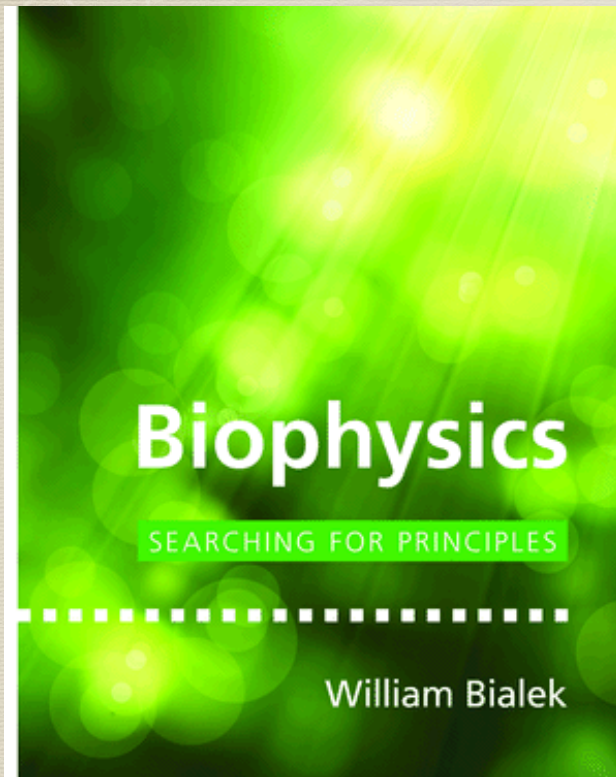
4.5 Proofreading and Active Noise Reduction 218

4.6 Perspectives 245

### 5. No Fine Tuning 247

Whether we like it or not, it could be probed **experimentally**.

No fine-tuning is a possible candidate principle not only for particle physics



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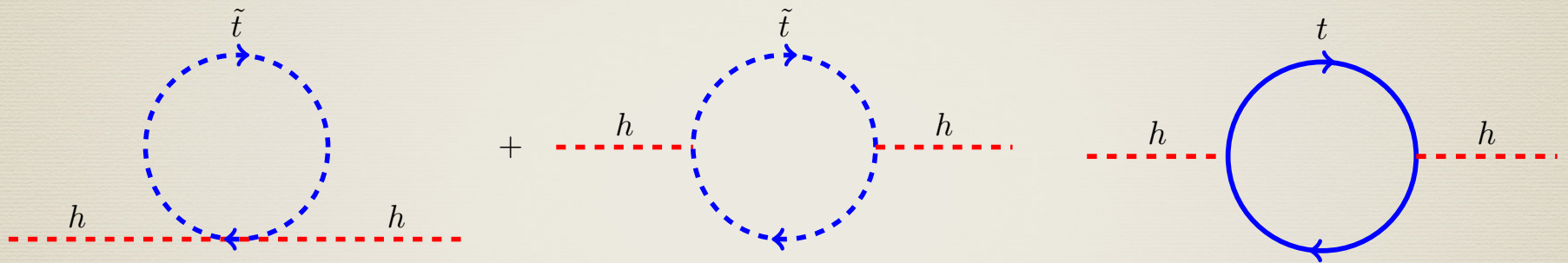
### 5. No Fine Tuning 247

Whether we like it or not, it could be probed **experimentally**. Whether the mechanism that generates the weak scale natural or unnatural, finding it will be ground-breaking!

# Electroweak Naturalness

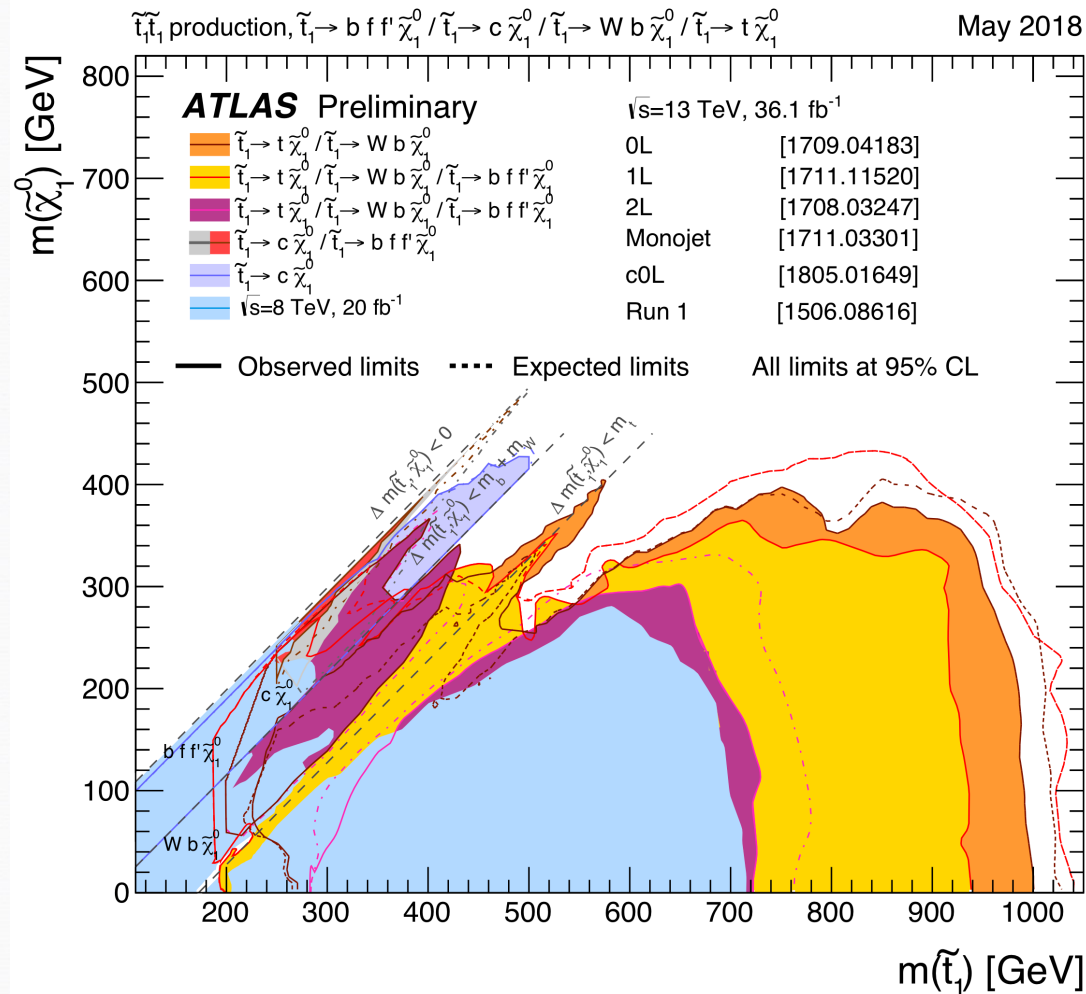
Traditional natural ways to explain the weak scale: new physics with **colored** top partners close to weak scale.

Classic examples: weak-scale SUSY and composite Higgs. In SUSY,



“**Stop**” or “**scalar top**”: cancels the biggest quantum correction from the top loop.  $\sim 10\%$  tuned if stop mass  $\sim$  TeV.

# Implications of LHC Results



Impressive reach with 13 TeV data for **simplest** stop decays (at both CMS and ATLAS): exclude stop  $\sim 1$  TeV (for neutralino below 400 GeV).

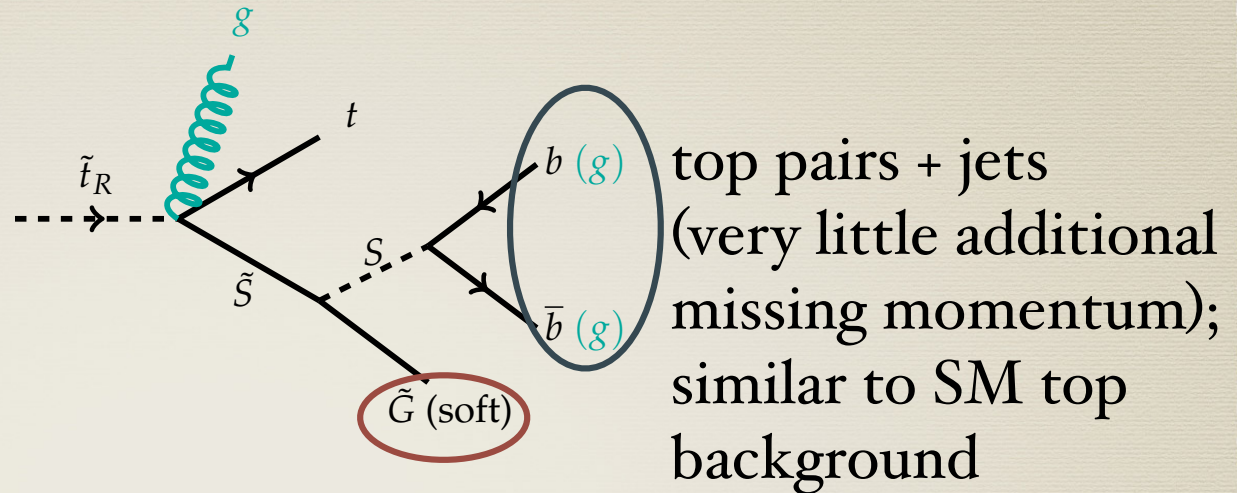
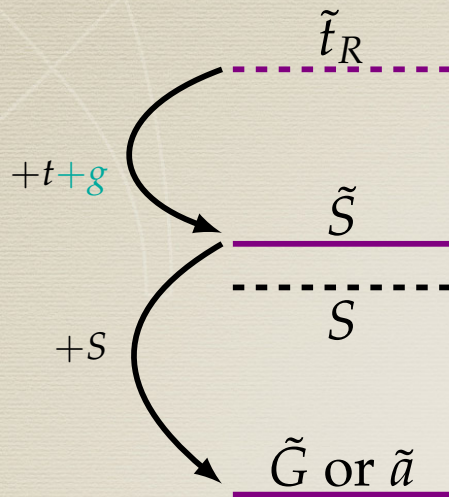
Null results teach us valuable lessons: traditional natural scenarios with electroweak fine-tuning no worse than 10% are very cornered.

There are still **loopholes** in existing searches.

The theoretical models may look more complicated and the main point is to motivate new experimental signals and searches.



# Effect of a Hidden Sector



light invisible fermion

**Stealth SUSY:** Fan, Krall, Pinner, Reece, Ruderman, 2015

Approximate SUSY in the *hidden sector* suppressing missing momentum;

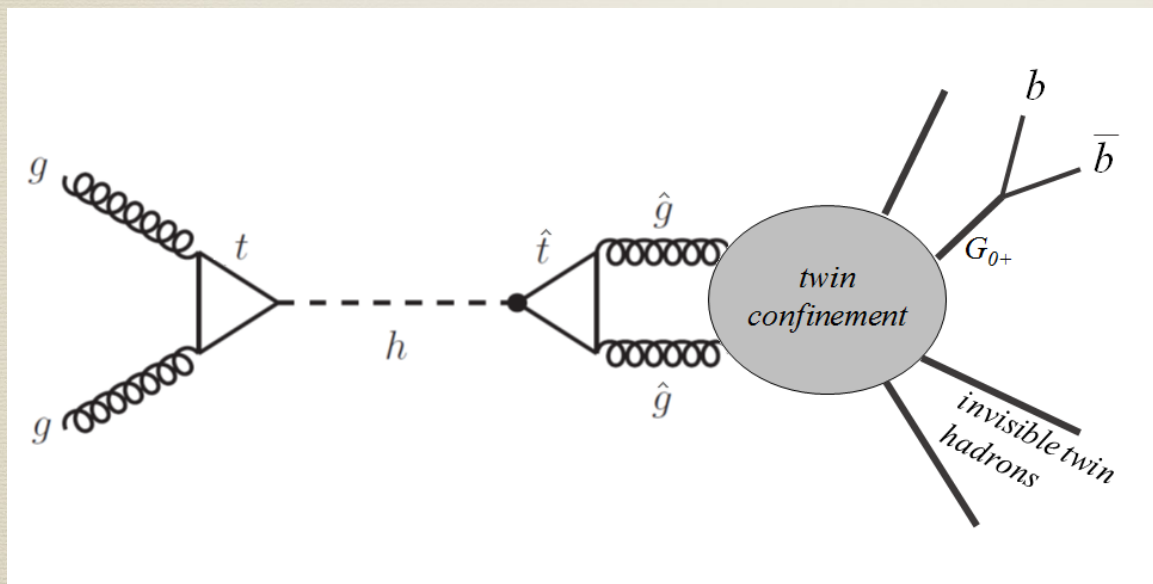
visible particles at the end of long cascades through the hidden sector have less energies.

In general, hidden valley could lead to dramatic new signatures:

Strassler, Zurek 2006

# Neutral Naturalness

Top partners do **not** feel strong dynamics.  
 Either SM gauge singlets or electroweakly charged (difficult to be found).  
 Chacko, Goh, Harnik, 2005; revived recently with many papers



Craig, Katz, Strassler,  
 Sundrum 2015

**Exotic Higgs decays:**  $gg \rightarrow h \rightarrow o^{++} + o^{++} + \dots$ ;  $o^{++} \rightarrow h^* \rightarrow b\bar{b}$   
 Long-lived, length scale  $\sim$  LHC detectors

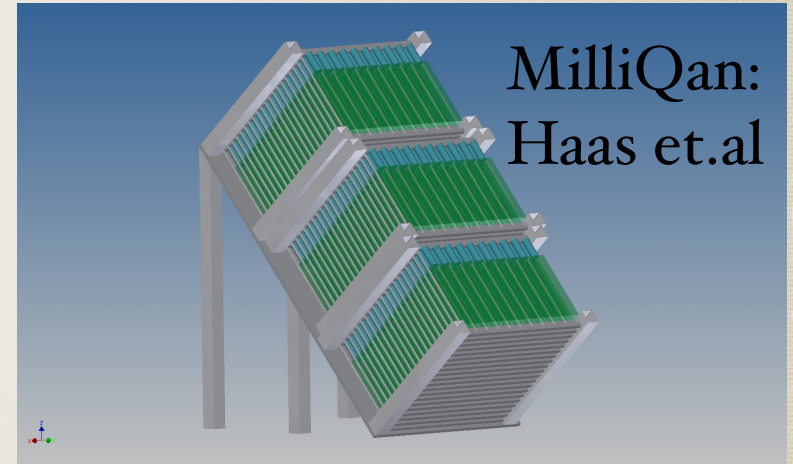
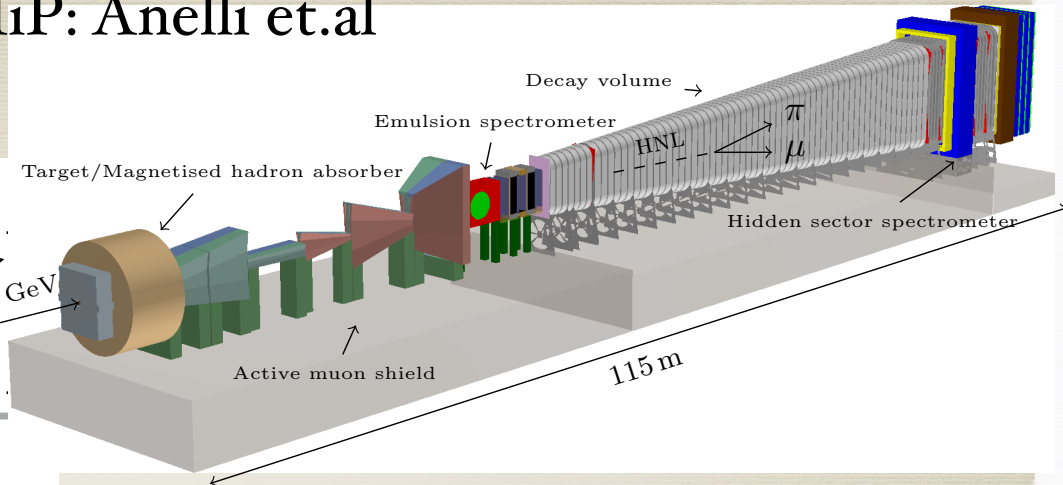
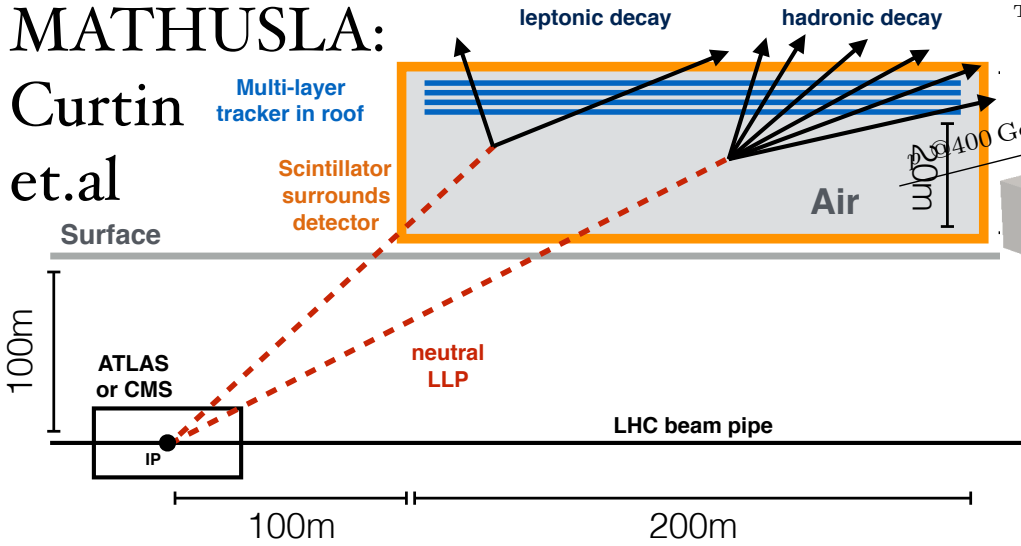
# Lifetime Frontier

SHiP: Anelli et.al

MATHUSLA:

Curtin

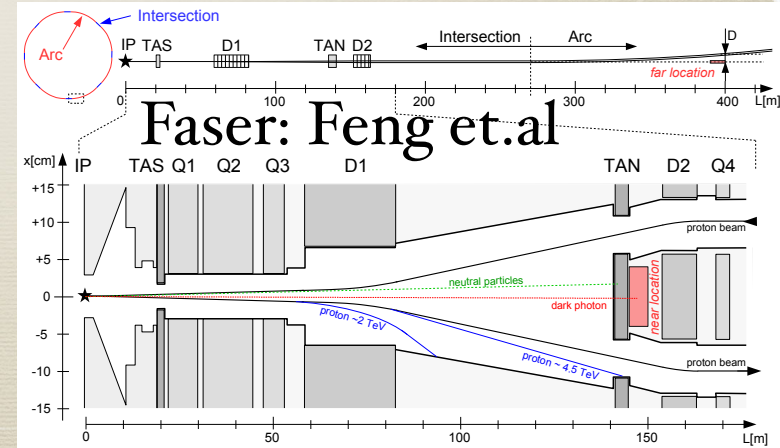
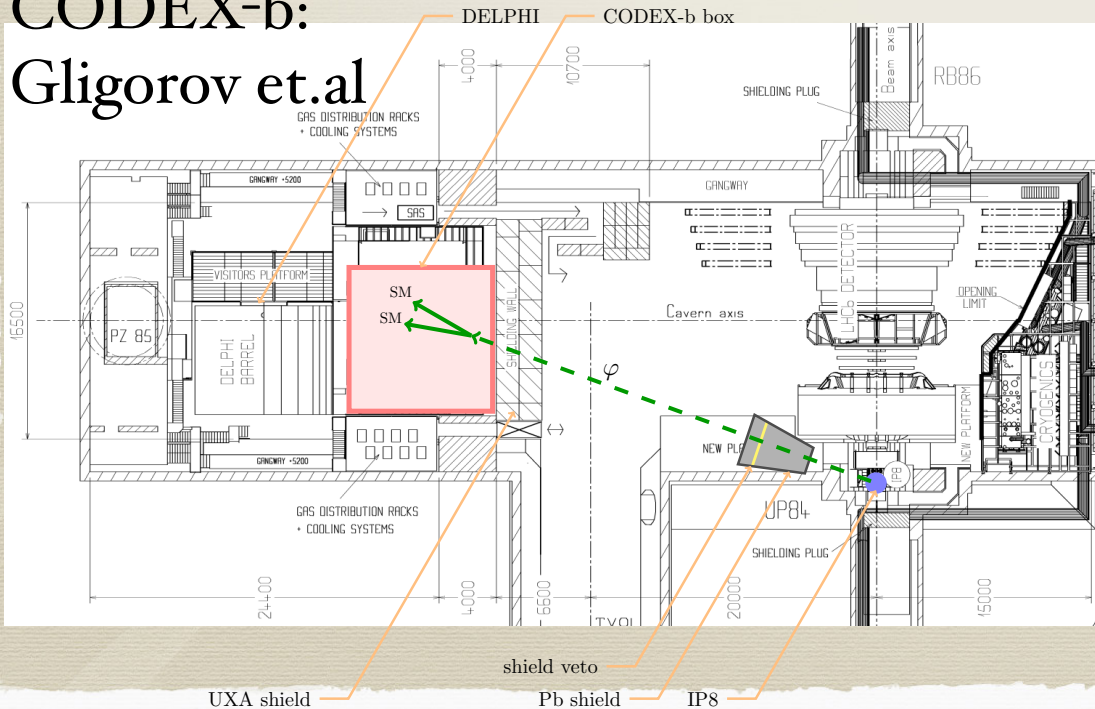
et.al



MilliQan:  
Haas et.al

CODEX-b:

Gligorov et.al



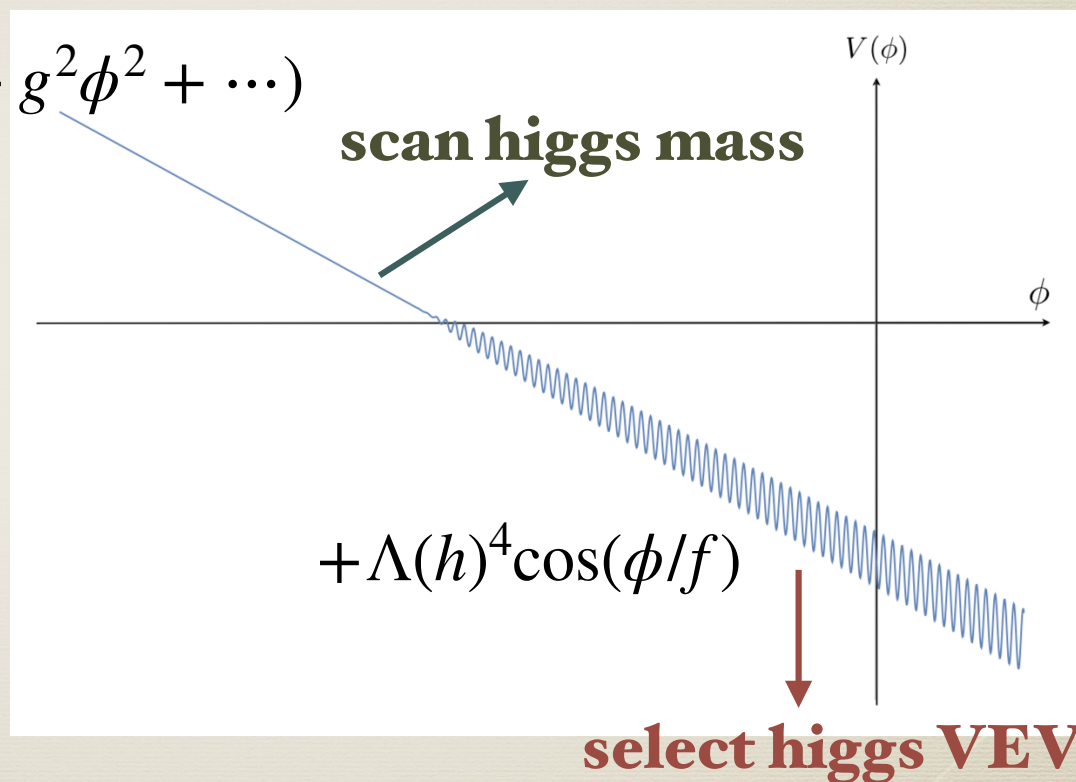
FASER: Feng et.al

# Relaxing the Little Hierarchy?

Graham, Kaplan, Rajendran 2015

Cosmological selected electroweak vacuum

$$(-M^2 + g\phi) |h|^2 + (gM^2\phi + g^2\phi^2 + \dots)$$



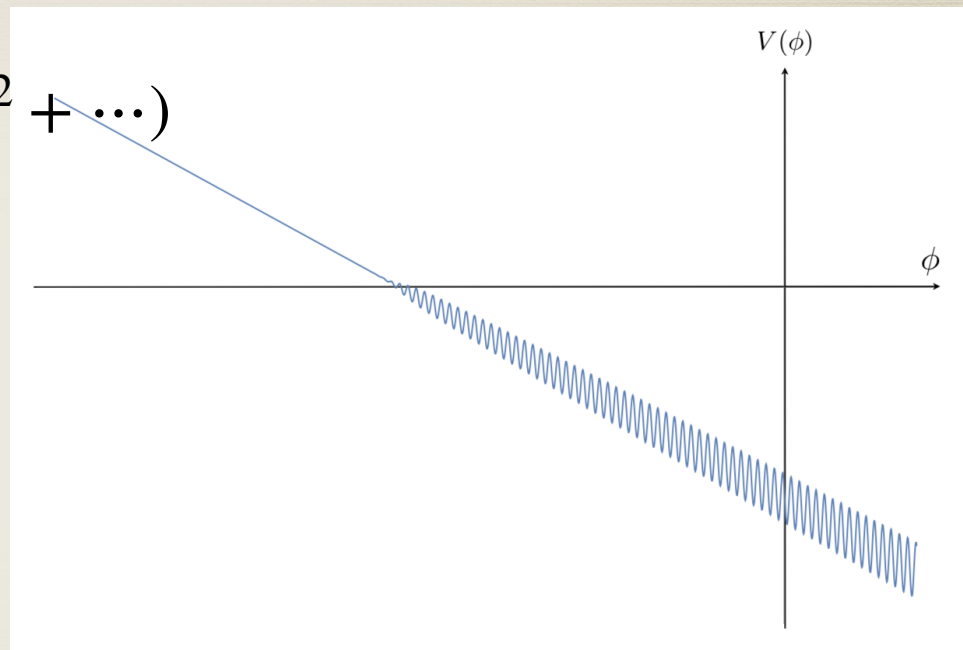
Huge energy stored in the evolving relaxation, need **dissipation**:  
inflation.

# Relaxing the Little Hierarchy?

Graham, Kaplan, Rajendran 2015

Cosmological selected electroweak vacuum

$$(-M^2 + g\phi) |h|^2 + (gM^2\phi + g^2\phi^2 + \dots) \\ + \Lambda(h)^4 \cos(\phi/f)$$



Original version requires:

*exponentially* small  $g$ ,  
*exponentially* large field range beyond the Planck scale,  
very low-scale inflation ( $H \ll \Lambda_{\text{QCD}}$ ),  
10 Giga-years of inflation...

Many further attempts based on it, to name a few,

Relaxion chiral supermultiplet with relaxino as gravitino.

Split-SUSY like spectrum with little hierarchy explained dynamically:

SUSY solves large hierarchy with relaxion solving little hierarchy

(Batell, Giudice, McCullough 2015)

Alternative *friction* during relaxation from particle production:

Smaller field range needed. Closer to plausibility?

(Hook, Marques Tavares 2016; Fonseca, Morgante, Servant, 2018)

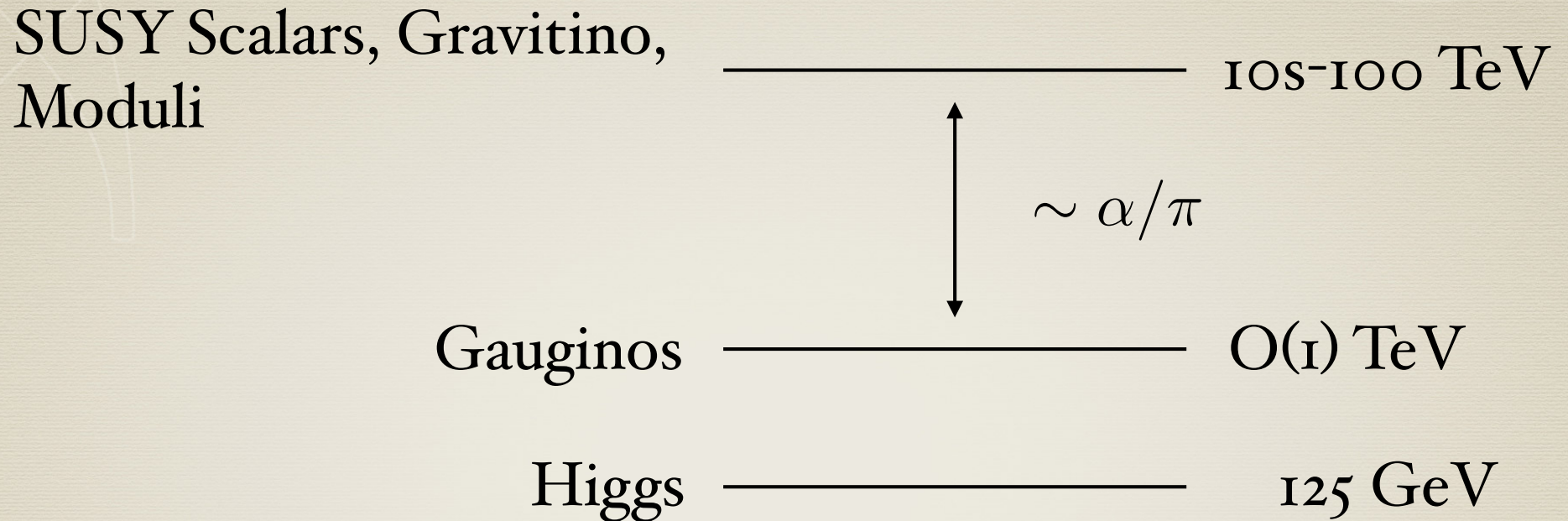
## An Intermediate View-point

Definitely weak scale  
should be explained fully  
in a natural way

Who cares about fine-tuning?  
Abandon it entirely: there  
could be other light scalars.

Higgs may be  
“meso-tuned” and no  
other random light  
scalar

# An example: mini-split



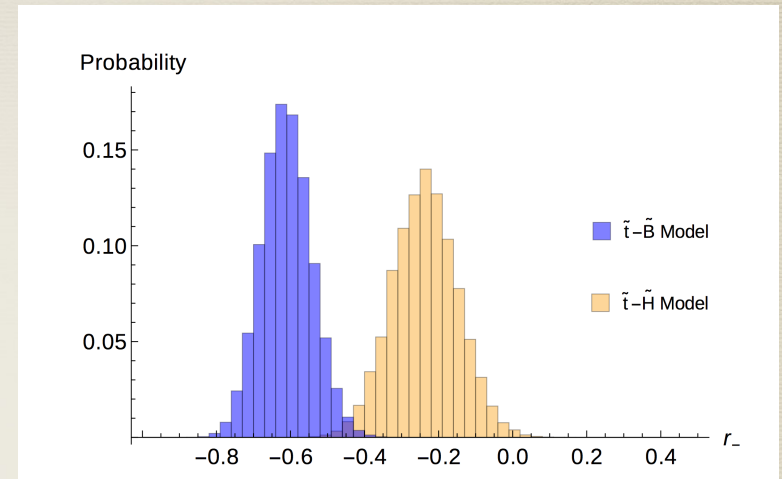
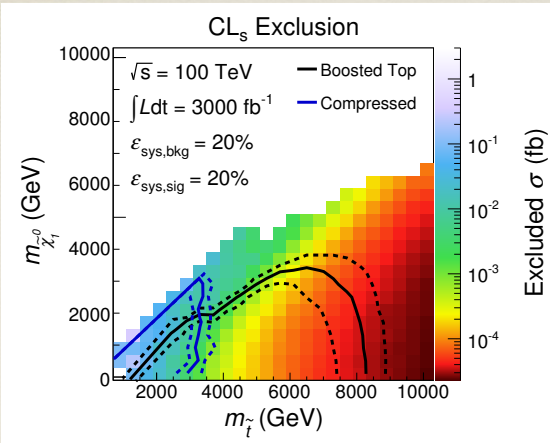
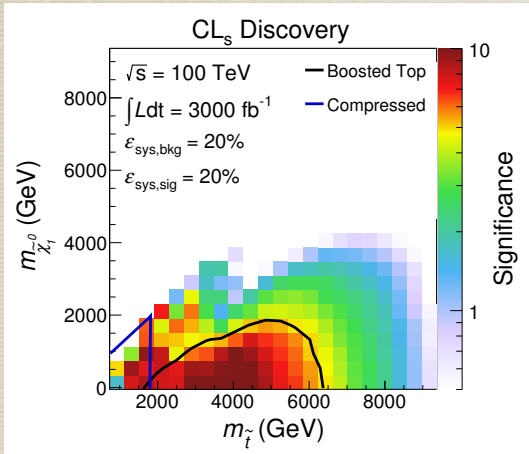
- Heavy scalars (10s of TeV) at large  $\tan \beta$ : right Higgs mass
- Loop factor: arises in AMSB (Giudice, Luty, Murayama, Rattazzi; Randall, Sundrum) and some moduli mediation
- Late-time gravitino and/or moduli decays populate nonthermal dark matter, e.g. light winos around  $O(100)$  GeV (Moroi, Randall; Kane et al.)

Many papers on “Mini-Split”: Arvanitaki et al., Hall et. al, Arkani-Hamed et al., ... 2012



# How to probe meso-tuning?

## Probe 1: future 100 TeV hadron collider



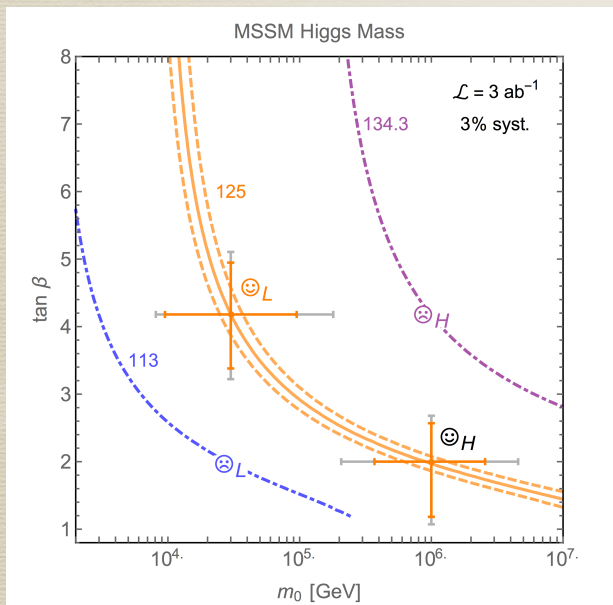
Cohen, D'Agnolo, Hance, Lou, Wacker 2014

Fan, Jaiswal, Leung 2017

Directly search for new particles

Measure their couplings precisely

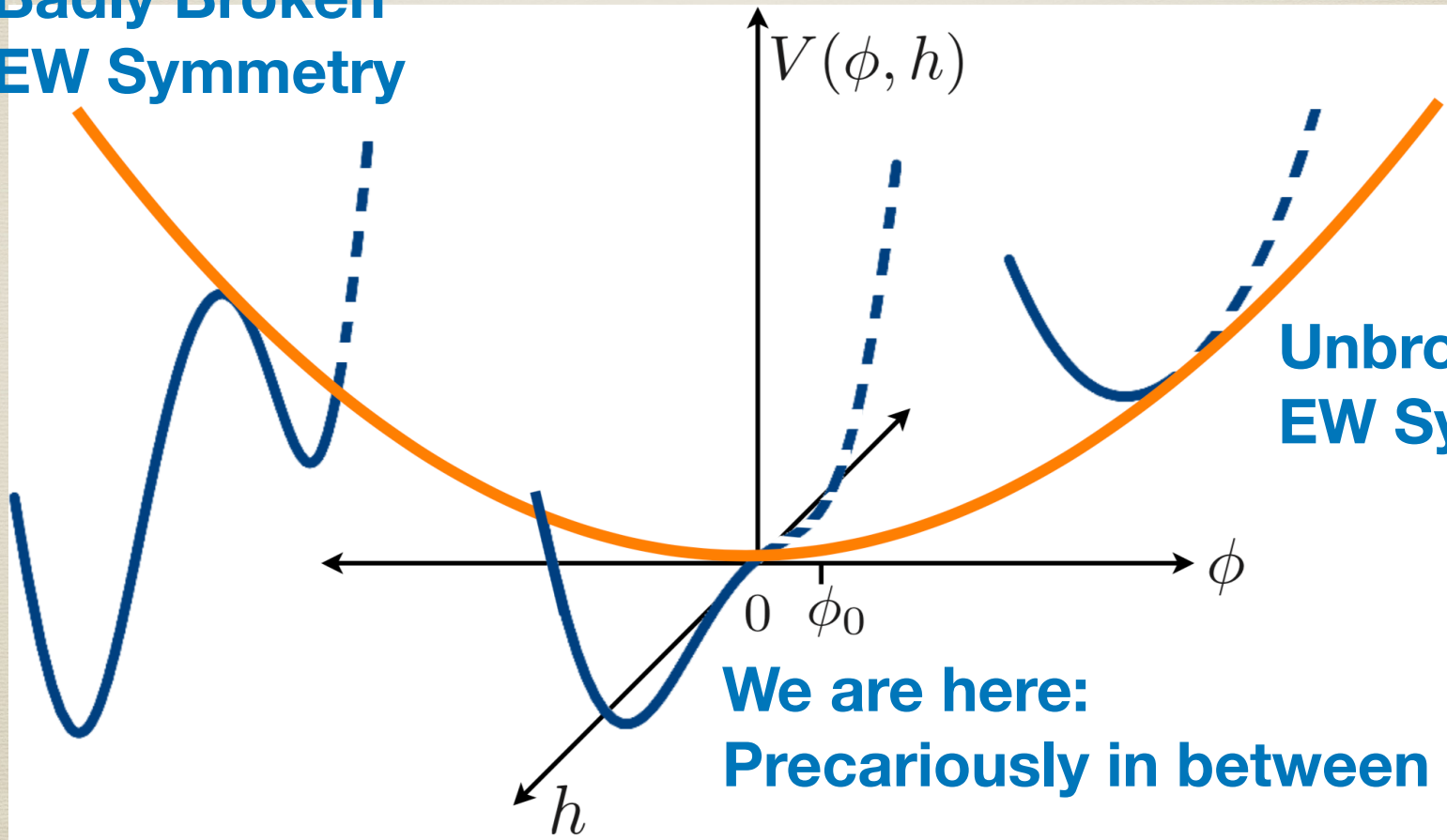
Agrawal, Fan, Reece, Xue 2017



## Probe 2: Cosmological Signal of a Fine-tuned Higgs

A time-dependent Higgs mass due to coupling to the oscillating scalar (modulus) in the early Universe.

**Badly Broken  
EW Symmetry**

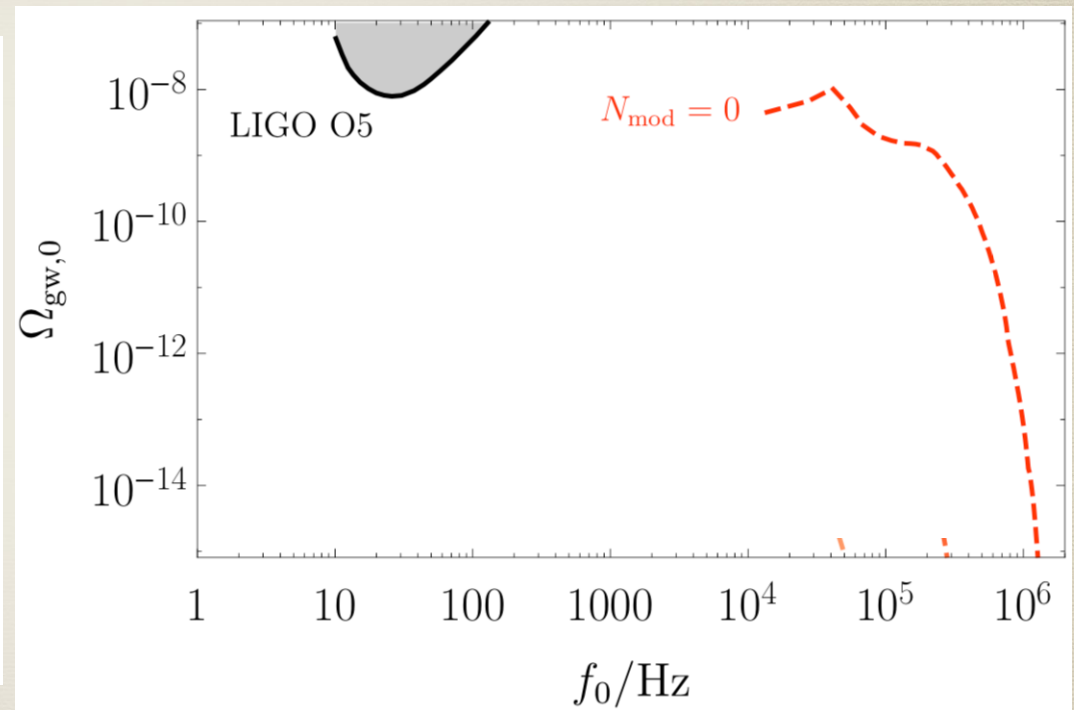
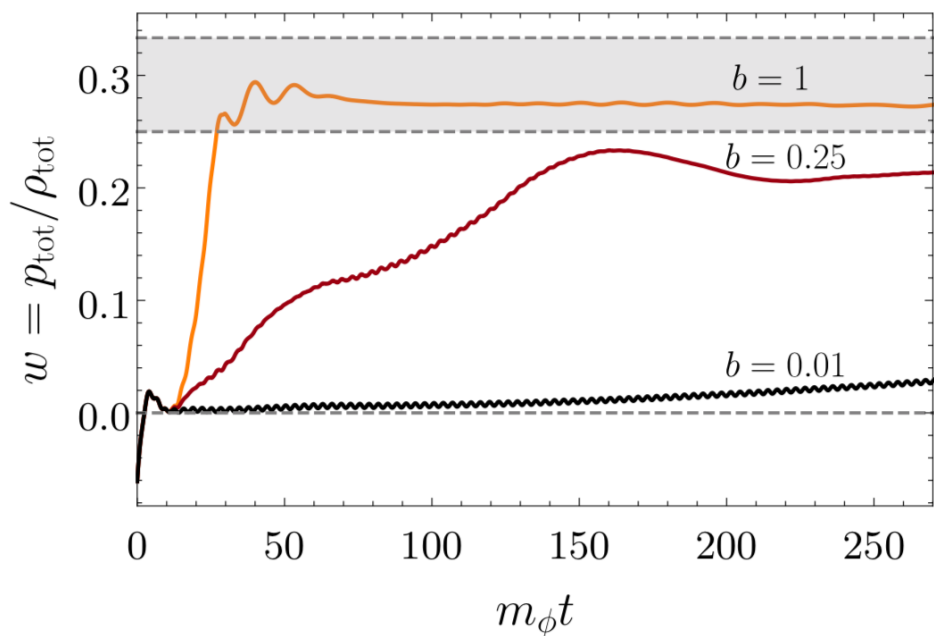


**We are here:  
Precariously in between**

If the Higgs potential is *tuned*, rapid particle production of the Higgs and fragmentation of the oscillating scalar (modulus)

**nontrivial equation of state**

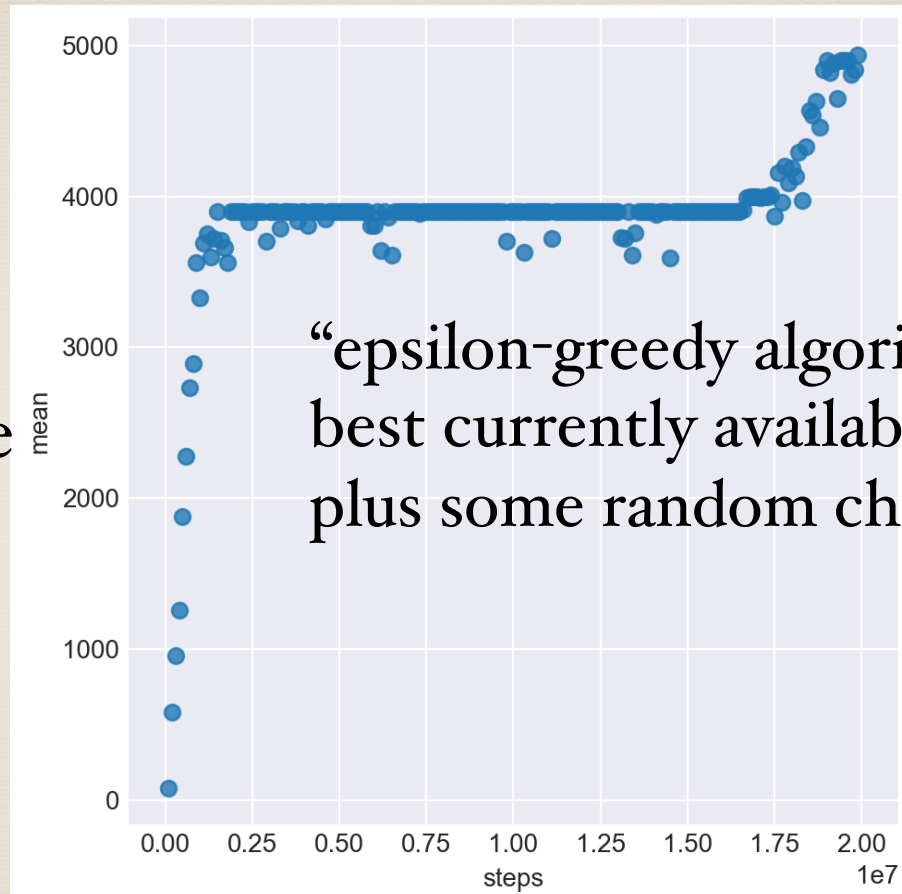
**stochastic gravitational waves**



Amin, Fan, Lozanov, Reece 1802.00444

# Concluding Remarks

Performance

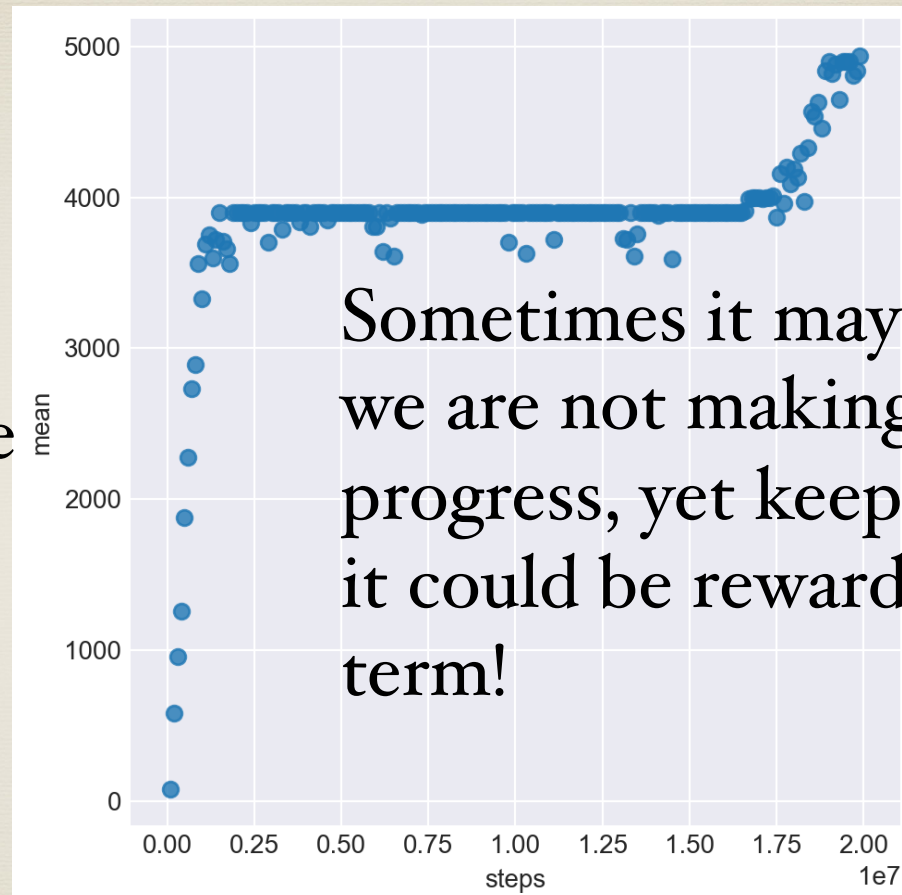


“epsilon-greedy algorithm”:  
best currently available option  
plus some random choices

Time

Reinforcement learning to scan string landscape (results thanks to Halverson)

Performance



Sometimes it may seem that we are not making that much progress, yet keep trying and it could be rewarding in the long term!

Time

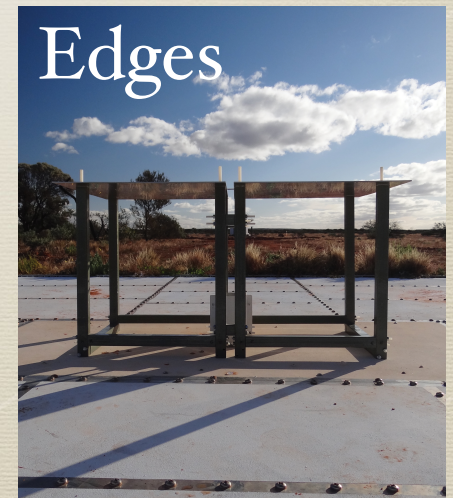
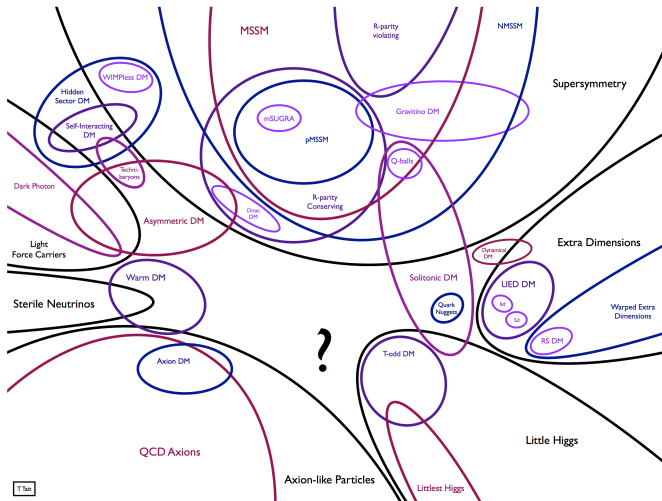
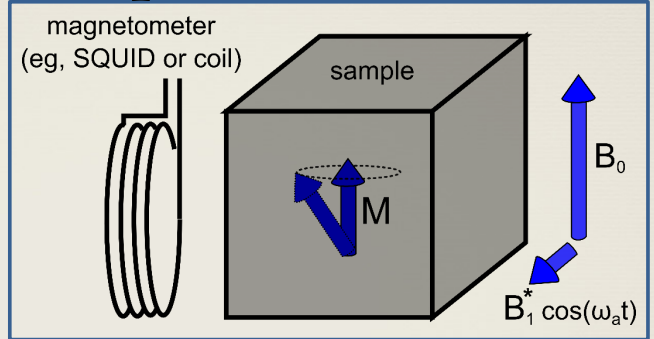
Thank you!

**Backup**

# A Little Bit on Dark Matter

Rapid expansion in both theory and experiment

## Casper



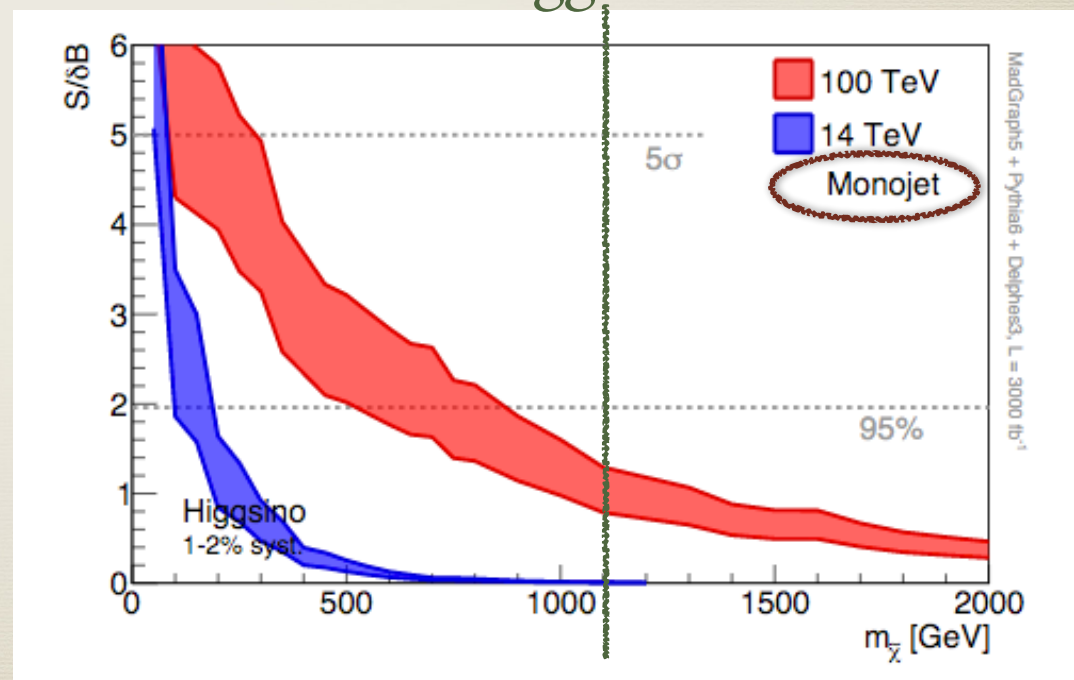
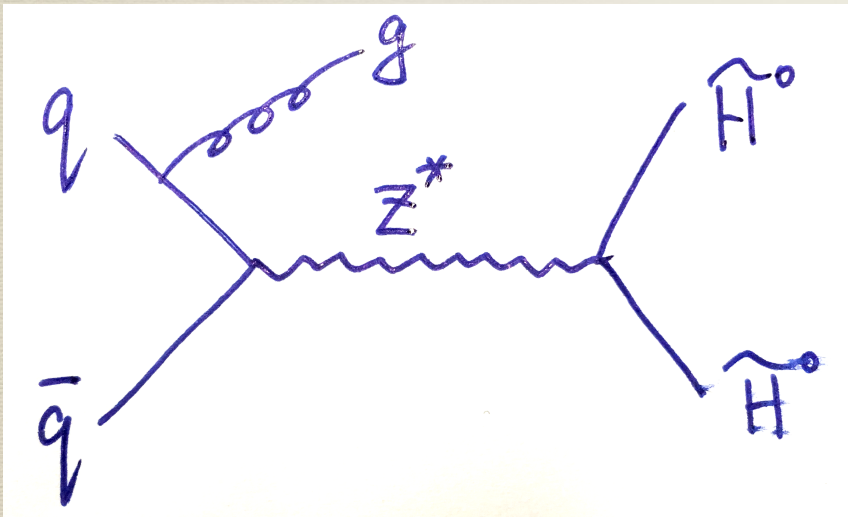
EDGES antenna in western Australia (photo credit: Judd Bowman/ASU)



# Simple WIMP at large: Higgsino DM

Simple WIMP model still **alive (elusive to all DM detections so far)**: higgsino dark matter, a fermionic electroweak doublet (fermionic copy of the Higgs doublet) with little mixing with other fermions, with the right thermal relic at 1.1 TeV.

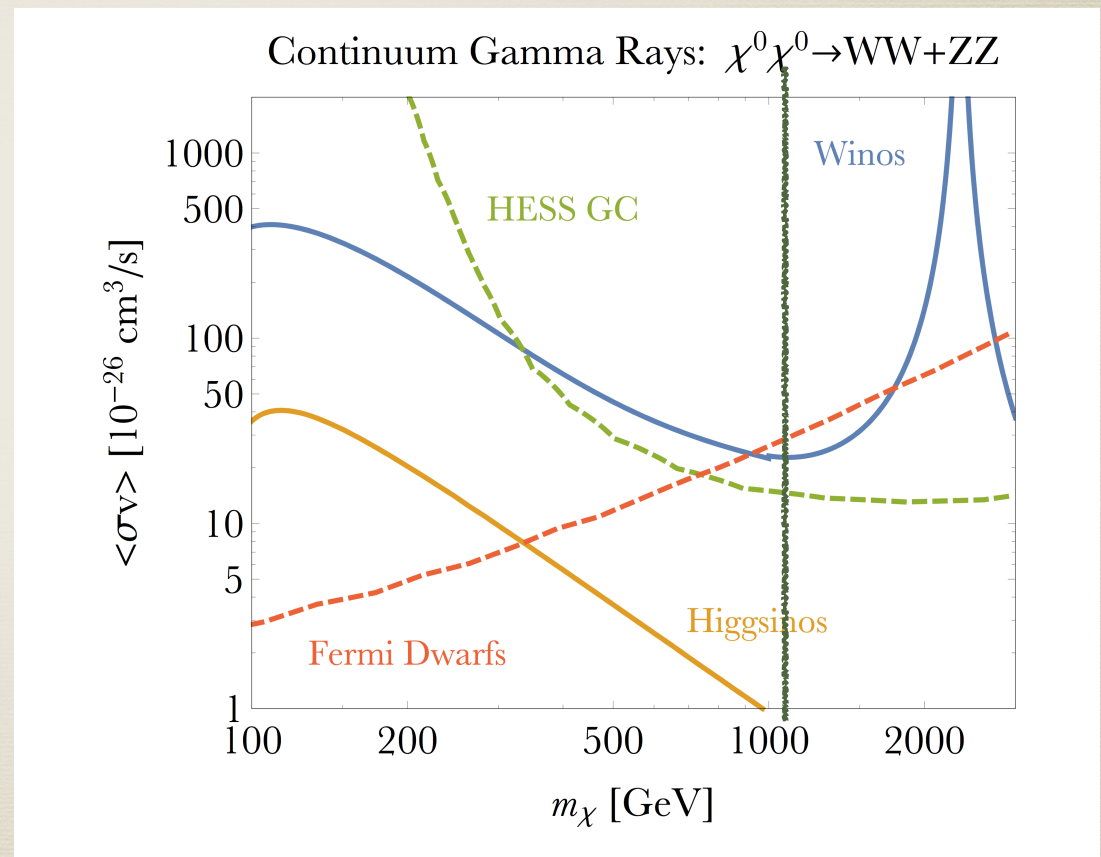
Thermal higgsino benchmark



Low, Wang:  
2014

*Direct detection:* scattering with nucleus happens at one loop level with a cross section  $\sim$  neutrino floor;

*Indirect detection:*  
about a factor of 50 below  
the current Fermi/HESS  
sensitivity.



Krall, Reece 1705.04843



Many other related studies aiming to improve the sensitivity at colliders for higgsino DM: e.g., a better tracker?

Charged and neutral higgsino nearly degenerate in mass, one-loop induced mass splitting  $\sim 360$  MeV;  
nominal decay length of charged higgsino,  $c\tau \sim \mathbf{6.6 \text{ mm}}$

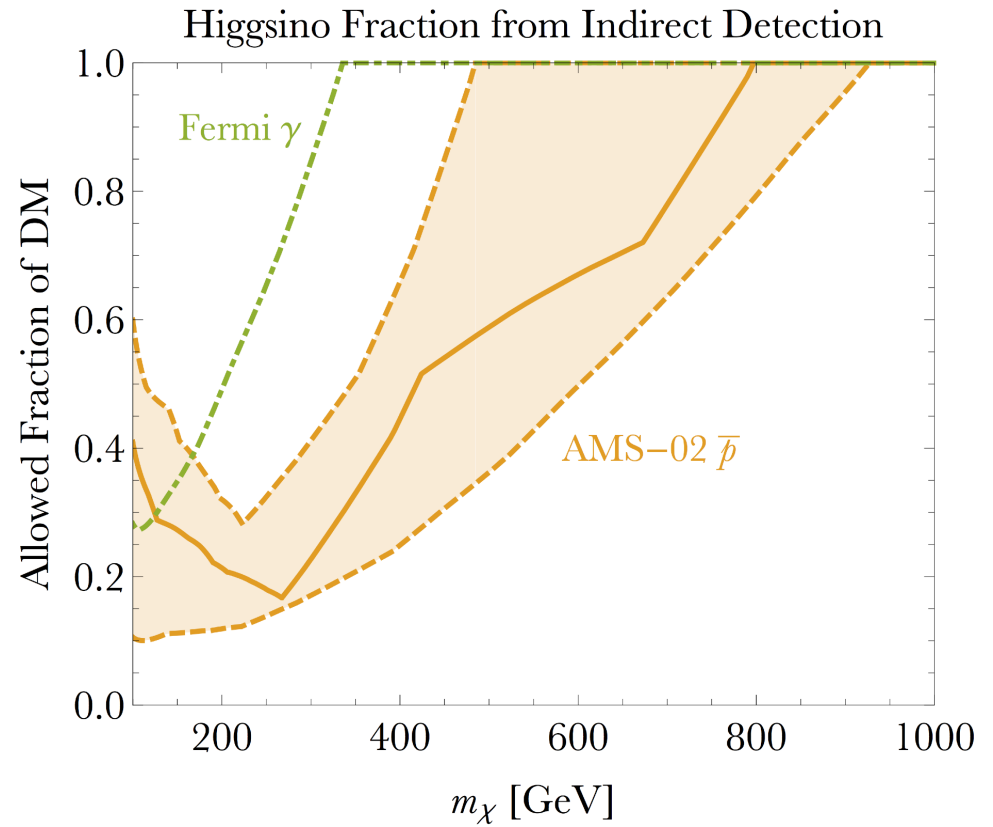
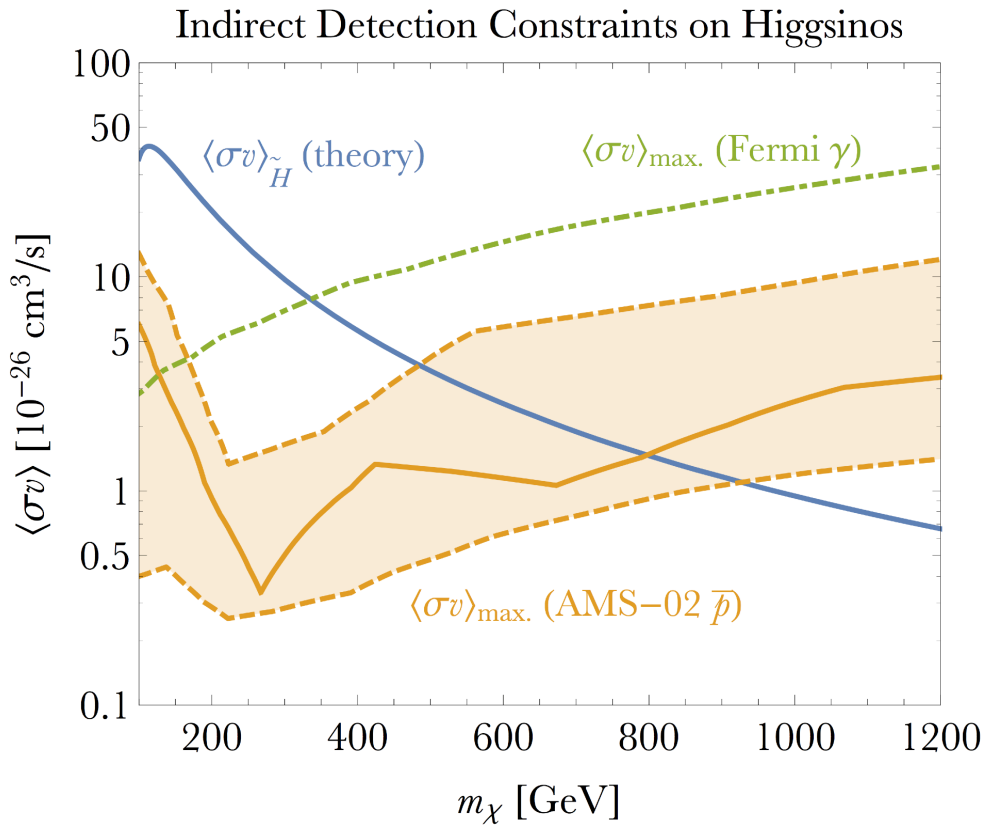
Disappearing charged track: need large boost ( $\sim 100$ ) (more easy to get large forward than transverse boost)

Increase the tracker granularity below  $r=10$  cm ( $r$ : transverse distance from the beamline): need 10 hits at  $r = 10$  cm.

In the future, may consider having a forward tracker covering  $2 \leq |\eta| \leq 4$ .

Mahbubani, Schwaller, Zurita;  
Fukuda, Nagata, Otono, Shirai, 2017

# AMS anti-proton constraint on higgsino



Krall, Reece 2017