Exploring doubly charged Higgs bosons collider signals in light of low energy constrains

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Plan:

- Extensions of the Standard Model
- 2 Higgs Triplet Model
- (3) $H^{\pm\pm} I I'$ coupling (LFV)
- Experimental constraints
- **5** $H^{\pm\pm}$ decay (HTM and LRSM)
- 6 Pair production and collider signals

Summary

Extensions of the Standard Model

• Higgs Triplet Model (HTM) One additional triplet:

$$\Delta = rac{1}{\sqrt{2}} \left(egin{array}{cc} w_\Delta^+ & \sqrt{2} \delta^{++} \ v_\Delta + h_\Delta + i z_\Delta & -w_\Delta^+ \end{array}
ight)$$

 Doubly charged scalar particles exist in other Standard Model's extention, for example Left Right Symmetric Model (LRSM)

Extensions of the Standard Model

HTM	MLRSM		
Type II See-Saw	Three heavy neutrinos		
SU(2) imes U(1)	$SU(2)_R imes SU(2)_L imes U(1)_{B-L} \ W_1,\ W_2,\ Z_1,\ Z_2,\ \gamma$		
$h, H, A, H^{\pm}, H^{\pm\pm}$	$\begin{array}{c} \textbf{h}, \textbf{H}_1, \textbf{H}_2, \textbf{H}_3, \textbf{A}_1, \textbf{A}_2 \\ \textbf{H}_1^{\pm}, \textbf{H}_2^{\pm}, \ \textbf{H}_1^{\pm\pm}, \textbf{H}_2^{\pm\pm} \end{array}$		

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$V = -m_{\Phi}^{2} \left(\Phi^{\dagger} \Phi \right) + M^{2} \operatorname{Tr} \left(\Delta^{\dagger} \Delta \right) + \left\{ \mu \left(\Phi^{T} i \sigma_{2} \Delta^{\dagger} \Phi \right) + \text{h.c.} \right\} \\ + \frac{\lambda}{4} \left(\Phi^{\dagger} \Phi \right)^{2} + \lambda_{1} \left(\Phi^{\dagger} \Phi \right) \operatorname{Tr} \left(\Delta^{\dagger} \Delta \right) + \lambda_{2} \left\{ \operatorname{Tr} \left(\Delta^{\dagger} \Delta \right) \right\}^{2} \\ + \lambda_{3} \operatorname{Tr} \left[\left(\Delta^{\dagger} \Delta \right)^{2} \right] + \lambda_{4} \left(\Phi^{\dagger} \Delta \Delta^{\dagger} \Phi \right)$

• Potential stability

Unitarity

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HTM - Lagrangian

$$\begin{split} M_A^2 &= \frac{\mu}{\sqrt{2}v_\Delta} (v_{\Phi}^2 + 4v_{\Delta}^2) \\ M_h^2 &= \lambda v_{\Phi}^2 \cos^2 \alpha + \left(\frac{\mu v_{\Phi}^2}{\sqrt{2}v_\Delta} + 2v_{\Delta}^2 (\lambda_2 + \lambda_3)\right) \sin^2 \alpha \\ &+ 2 \left(v_{\Phi} v_{\Delta} (\lambda_1 + \lambda_4) - \sqrt{2}\mu v_{\Phi}\right) \cos \alpha \sin \alpha \\ M_H^2 &= \lambda v_{\Phi}^2 \sin^2 \alpha + \left(\frac{\mu v_{\Phi}^2}{\sqrt{2}v_\Delta} + 2v_{\Delta}^2 (\lambda_2 + \lambda_3)\right) \cos^2 \alpha \\ &- 2 \left(v_{\Phi} v_{\Delta} (\lambda_1 + \lambda_4) - \sqrt{2}\mu v_{\Phi}\right) \cos \alpha \sin \alpha \\ M_{H^{\pm}}^2 &= \frac{(2\sqrt{2}\mu - \lambda_4 v_\Delta)}{4v_\Delta} (v_{\Phi}^2 + 2v_{\Delta}^2) \\ M_{H^{\pm\pm}}^2 &= \frac{\mu v_{\Phi}^2}{\sqrt{2}v_\Delta} - \frac{\lambda_4}{2} v_{\Phi}^2 - \lambda_3 v_{\Delta}^2 \end{split}$$

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Unitarity, stability



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$H^{\pm\pm} - I - I'$ coupling

$$\mathcal{L}_{Y} = \frac{1}{2} f_{\ell \ell'} L_{\ell}^{T} C^{-1} i \sigma_{2} \Delta L_{\ell'} + \text{h.c.}$$

$$\mathcal{L}_{
u} = rac{1}{2}\,ar{
u}_\ell\,rac{
u_\Delta}{\sqrt{2}}\,f_{\ell\ell'}\,
u_{\ell'}$$



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$H^{\pm\pm} - I - I'$ coupling



$$f = \frac{1}{\sqrt{2}v_{\Delta}} V_{PMNS}^* D_{\nu} V_{PMNS}^{\dagger} \qquad D_{\nu} = \frac{1}{2} \operatorname{diag}\{m_1, m_2, m_3\}$$

$$V_{\Delta} \iff f_{\parallel \prime} \iff \theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP}$$

$$m_1, m_2, m_3$$

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$f_{ll'}$ coupling and parameter's constraints

- $10^{-7} \leq f_{ll'} \leq \sqrt{4\pi}$ [hep-ex/0309076]
- $v_{\Delta} \sim 1 \ {
 m GeV}$ [Phys.Rev. D21 (Mar, 1980) 1404-1409], [arXiv:0712.4053]

$$\rho = \frac{M_W^2}{M_Z^2 \cos \theta_W} = \frac{1 + 2\frac{v_A^2}{v_\phi^2}}{1 + 4\frac{v_A^2}{v_\phi^2}} = 1.00040 \pm 0.00024$$

• $\sum_{i=1}^{n} m_i \le 2 \text{ eV}$ - Iritium decay [PDG] • $\sum_{i=1}^{n} m_i \le 0.23 \text{ eV}$ - astrophysics [arXiv:1303.5076]

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Neutrino oscillation data

	Normal hierarchy			Inverted hierarchy		
	Best fit:	σ	bfp $\pm 2\sigma$	Best fit:	σ	bfp $\pm 2\sigma$
$\sin^2 \theta_{12}$	0.306	$^{+0.012}_{-0.012}$	$0.282 \div 0.330$	0.306	+0.012 -0.012	0.282 ÷ 0.330
$\sin^2 \theta_{23}$	0.441	+0.027 -0.021	0.399 ÷ 0.495	0.587	+0.020 -0.024	0.539 ÷ 0.627
$\sin^2 \theta_{13}$	0.02166	+0.00075 -0.00075	0.02016 ÷ 0.02316	0.02179	+0.00076 -0.00076	0.02027 ÷ 0.02331
δ _{CP} [°]	261	+51 -59	143 ÷ 363	277	+40 -46	185 ÷ 357
$\frac{\Delta m_{21}^2}{10^{-5} \text{eV}^2}$	7.50	$^{+0.19}_{-0.17}$	$7.16 \div 7.88$	7.50	$^{+0.19}_{-0.17}$	7.16 ÷ 7.88
$\frac{\Delta m_{3l}^2}{10^{-3} \mathrm{eV}^2}$	+2.524	+0.039 -0.040	2.445 ÷ 2.602	-2.514	+0.038 -0.041	-2.596 ÷ -2.438

[www.nu-fit.org , arXiv:1611.01514]

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Experimental constraints

- High energy:
 - Bhabha scattering: $f_{ee}^2 \leq 6.0 imes 10^{-6} M_{H^{\pm\pm}}$

[Phys.Rev. D40 (1989) 1521], [hep-ph/0304069]



Experimental constraints

• Low energy:



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Experimental constraints

• $(g-2)_{\mu}$ $\Delta a_{(muong-2)}$ = (29.3 ± 9.0) × 10⁻¹⁰ • μ to e (for Au) $BR(\mu N \rightarrow eN^*) < 7.0 \times 10^{-13}$ Radiative LFV decays 10^{-13} $\mathsf{BR}(\mu \to e\gamma) < 5.7 \times$ $\mathsf{BR}(au o e\gamma)$ < 3.3 imes 10^{-8} $\mathsf{BR}(\tau \rightarrow \mu \gamma) < 4.4 \times 10^{-8}$ LFV three body decays $\mathsf{BR}(\tau \to \mu \gamma)$ < 4.4 × 10⁻⁸ $\mathsf{BR}(au o 3\mu)$ < 2.1 imes 10⁻⁸ $\mathsf{BR}(au o e \mu^+ \mu^-)$ < 2.7 imes 10^{-8} [arXiv:1512.03581] . [Nucl. Phys. B299 (1988) 1-6] $\mathsf{BR}(\mu \to 3e) < 1.0 \times 10^{-12}$

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 v_{Δ} vs $M_{H^{\pm\pm}}$

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$H^{\pm\pm}$ decay

•
$$H^{\pm\pm} \rightarrow I_i I_j$$

• $H^{\pm\pm} \rightarrow W^{\pm} W^{\pm}$
• $H^{\pm\pm} \rightarrow H^{\pm} W^{\pm}$

 $\bullet \hspace{0.2cm} H^{\pm\pm} \to H^{\pm}H^{\pm}$



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$H^{\pm\pm}$ decay



$H^{\pm\pm}$ pair production in lepton colliders



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$H^{\pm\pm}$ pair production in hadron colliders



For MLRSM see arXiv:1311.4144

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$pp \rightarrow H^{++}H^{--} \rightarrow$ 4-leptons signal

Luminosity 25 fb⁻¹, $\sqrt{s} = 14$ TeV

			НТМ				
$M_{H^{\pm\pm}}$	SM	MLRSM	NH		I	IH	
			$\alpha_1 = 0, \ \alpha_2 = 0$	$\alpha_1 = 0, \ \alpha_2 = \frac{\pi}{2}$	$\alpha_1=$ 0, $\alpha_2=$ 0	$\alpha_1 = 0, \ \alpha_2 = \frac{\pi}{2}$	
400	2.9	30	7.3	2.6	33	20	
600		4.4	1.0	0.4	4.9	2.9	

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- $H^{\pm\pm}$ pair production in colliders gives possibility for clean BSM 4I signals
- For $\theta_{13} \neq 0$ the strongest limit on v_t comes from the $\mu \rightarrow 3e$ LFV process
- T channel contribution to the $H^{\pm\pm}$ pair production in lepton colliders is negligible due to the low energy constraints
- Heavy gauge bosons (RH currents) do not influence the total number of event for the 4-lepton signal in hadron colliders

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Thank you for your attention

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U_{PMNS}

$$V_{PMMS} = \begin{bmatrix} c_{12}c_{13}e^{i\alpha_1} & s_{12}c_{13}e^{i\alpha_2} & s_{13}e^{-i\delta}CP \\ (-s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta}CP)e^{i\alpha_1} & (c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta}CP)e^{i\alpha_2} & s_{23}c_{13} \\ (s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta}CP)e^{i\alpha_1} & (-c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta}CP)e^{i\alpha_2} & c_{23}c_{13} \end{bmatrix}$$
(1)

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ρ and T parameters

$$\Delta \rho = \frac{\prod_{WW}(0)}{M_W^2} - \frac{\prod_{ZZ}(0)}{M_Z^2}$$
$$T = \frac{1}{\alpha} \left(\frac{\prod_{WW}(0)}{M_W^2} - \frac{\prod_{ZZ}(0)}{M_Z^2} \right) = \frac{\rho - 1}{\alpha} = 0.05 \pm 0.12$$
$$\Delta T = \frac{1}{4\pi \sin \theta_W^2 M_W^2} \left(F(M_{H^{\pm}}^2, M_A^2) + F(M_{H^{\pm\pm}}^2, M_{H^{\pm}}^2) \right)$$
$$\Delta T < 0.2$$
$$F(x, y) = \frac{x + y}{2} - \frac{xy}{x - y} \ln\left(\frac{x}{y}\right)$$

 $(g - 2)_{\mu}$



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$H^{\pm\pm}$ decay



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Potential stability

$$\begin{array}{rcl} \lambda &> 0\\ \lambda_2 + \lambda_3 &\geq 0\\ \lambda_2 + \frac{\lambda_3}{2} &\geq 0\\ \lambda_1 + \sqrt{\lambda(\lambda_2 + \lambda_3)} &\geq 0\\ \lambda_1 + \lambda_4 + \sqrt{\lambda(\lambda_2 + \lambda_3)} &\geq 0\\ \left[& |\lambda_4|\sqrt{\lambda_2 + \lambda_3} - \lambda_3\sqrt{\lambda} &\geq 0\\ & \left[& |\lambda_4|\sqrt{\lambda_2 + \lambda_3} - \lambda_3\sqrt{\lambda} &\geq 0 \\ & \left[& |\lambda_4|\sqrt{\lambda_2 + \lambda_3} - \lambda_3\sqrt{\lambda} &\geq 0 \\ & \end{array} \right] \end{array}$$

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