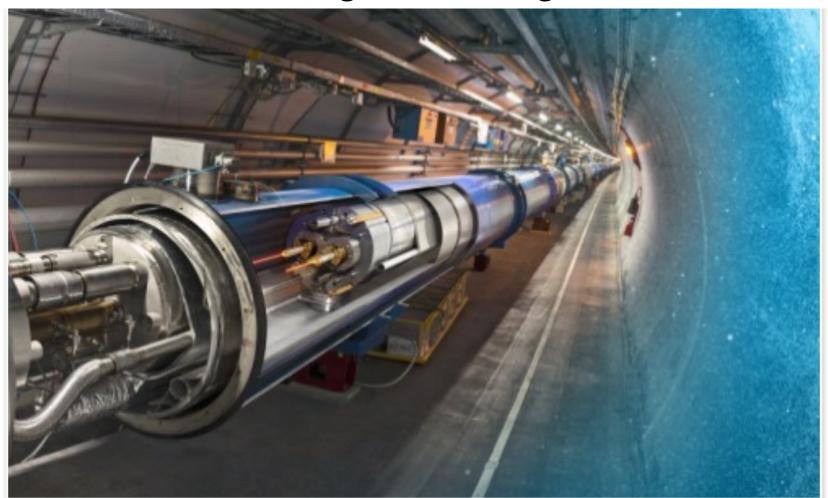


Annual SPS meeting 28-31 August 2018, EPFL





LANDAU DAMPING LIMITATIONS IN HIGH ENERGY HADRON COLLIDERS

Dr. Claudia Tambasco

(EPFL Thesis n. 7867)

Supervisors: Prof. Rivkin, T. Pieloni

Particle Accelerator Physics Laboratory (LPAP)

Institute of Physics

EPF Lausanne, Switzerland



THE LARGE HADRON COLLIDER



LHC is the world's largest and most powerful particle accelerator. Its main application is highenergy physics research thanks to 4 experiments where the beams collide.

Event Rate = $L \cdot \sigma_{ev}$

Luminosity:

$$L = \frac{N^2 f_{\text{rev}} n_b}{2\pi \sigma^2}$$

Geneva Lake

CMS

LHC

LHC

LHC

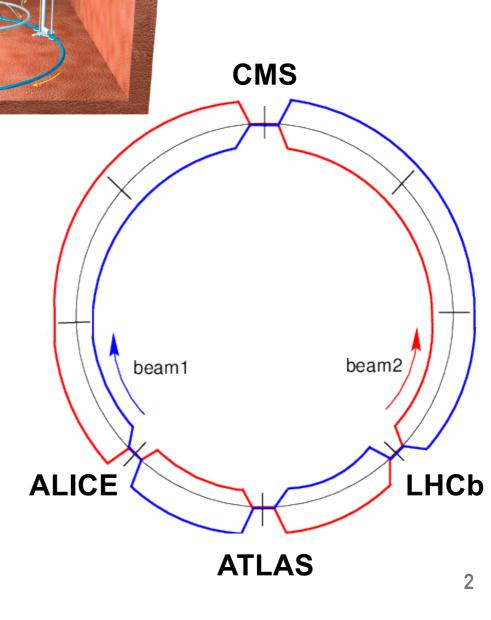
LHC

- Large number of bunches
- High Intensity
- Small r.m.s. beam size

High brightness (N/σ)

Reached Luminosity in the LHC (2017):

$$L = 2.2 \times 10^{34} \, \text{cm}^{-2} \, \text{s}^{-1}$$





BASIC PRINCIPLES OF BEAM DYNAMICS IN A SYNCHROTRON

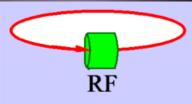


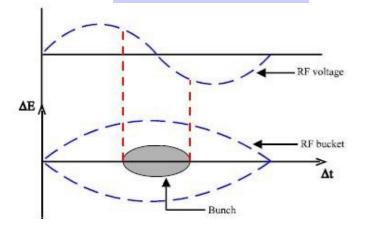
Basic principle: LORENTZ FORCE $\vec{F} = q\vec{E} + q\vec{v}\vec{x}\vec{B}$ | Magnetic force | Magnetic

Charged particles

RF Cavities accelerate and bunch the particles

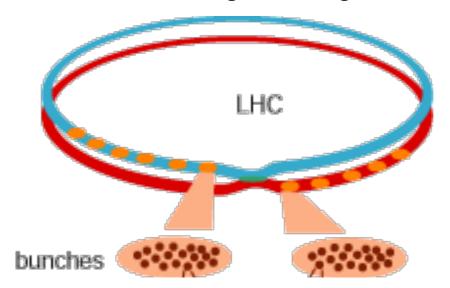






Longitudinal Plane

The **RF cavities** accelerate the particles and bunch the beams along the longitudinal direction



~2800 bunches per beam

The particles oscillate in each bunch with a frequency proportional to

Synchrotron tune: Qs

LHC synchrotron tunes: Q_s=0.005 (injection energy) Q_s=0.002 (top energy)



BASIC PRINCIPLES OF BEAM DYNAMICS IN A SYNCHROTRON



Magnetic fields

- Dipoles: bend and guide the particles along the accelerator
- Quadrupoles: focus the particles (avoid beam divergency)
- Sextupoles: correct for chromatic effects
- Octupoles: provide beam stability (see next)
- and higher order magnetic fields

dipole Dipole Focusing quadrupole Defocusing quadrupole Closed orbit RF cavities **Schematic Top View**

Transverse Plane

Particle motion ≈ oscillations in the transverse plane under the action of periodic external fields produced by the machine lattice (arrangement of magnets):

Betatron frequency ω_β: oscillation frequency in the transverse plane



Betatron tune: Q_β

H=64.31 V=59.32 (collisions)



COLLECTIVE EFFECTS



Additional (unwanted) electromagnetic fields perturb the beam dynamics and induce coherent motion of particles inside the beam → coherent beam instabilities (loss of stable motion)

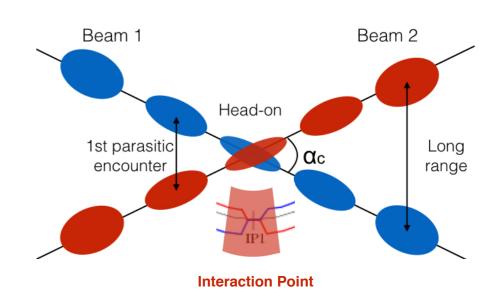
Stronger for higher brightness beams

$$L = \frac{N^2 f_{\text{rev}} n_b}{2\pi \sigma^2}$$

EM fields from other charge distribution

Beam – beam effects in colliders, EM interaction of two beams

when sharing common pipe

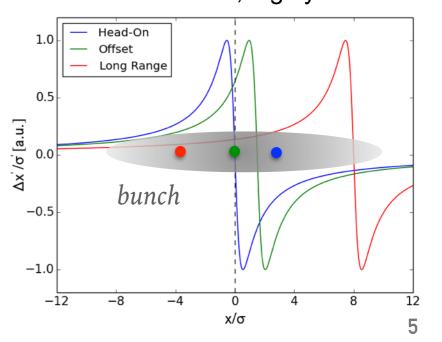


Self-induced EM fields

■ EM interaction of the beam with the surrounding environment

Wake fields can remain trapped and decay

Beam-Beam force, highly non-linear





COLLECTIVE EFFECTS



Additional (unwanted) electromagnetic fields perturb the beam dynamics and induce coherent motion of particles inside the beam → coherent beam instabilities (loss of stable motion)

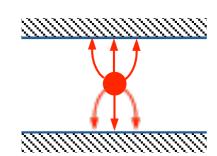
Stronger for higher brightness beams

$$L = \frac{N^2 f_{\text{rev}} n_b}{2\pi \sigma^2}$$

EM fields from other charge distribution

■ Beam – beam effects in colliders, EM interaction of two beams

when sharing common pipe



Self-induced EM fields



Wake fields can remain trapped and decay



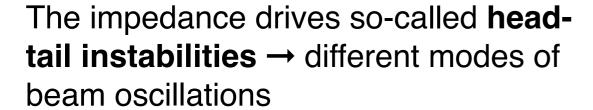




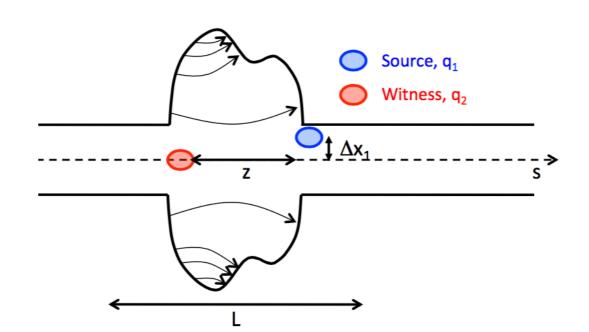
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)



- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift







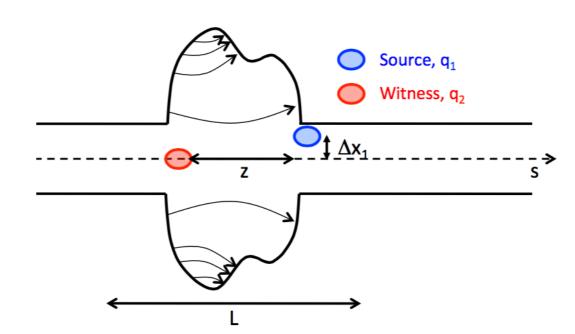
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

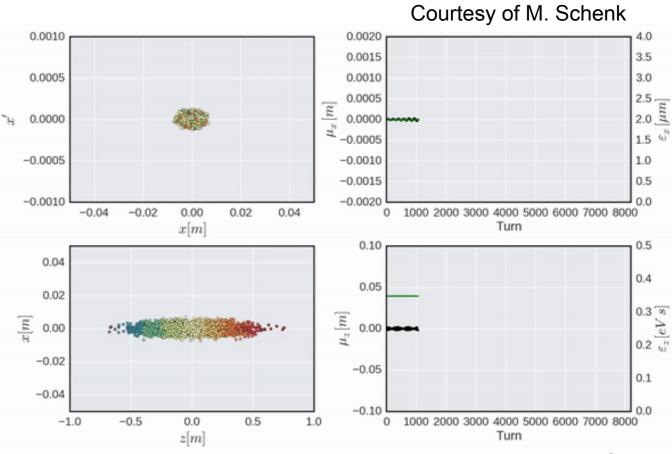
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- $Im(\Delta Q)$: instability rise time
- $Re(\Delta Q)$: real coherent tune shift









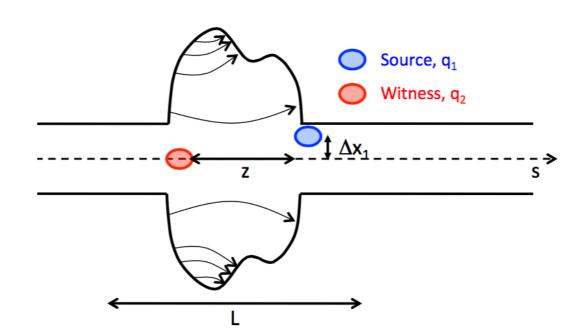
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

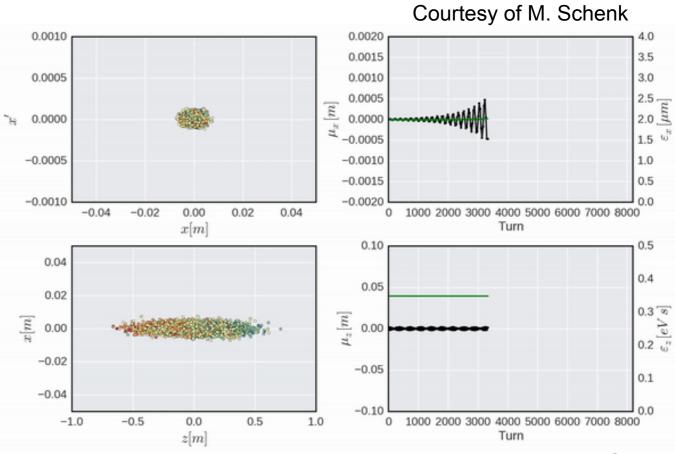
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









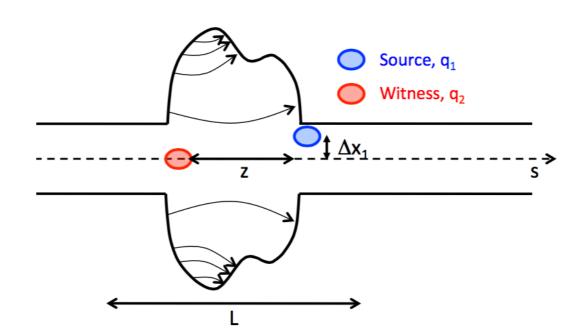
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

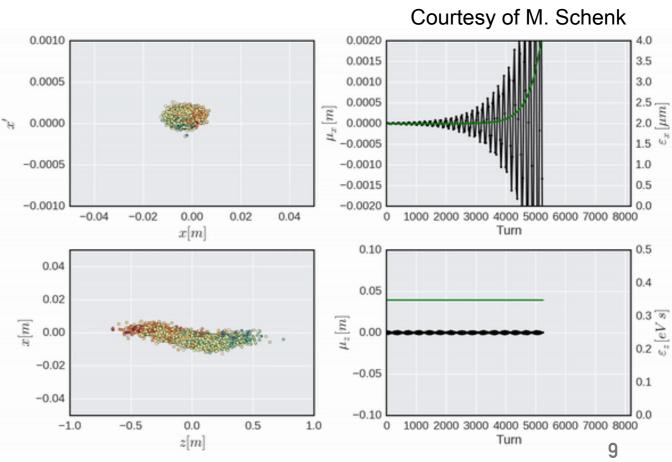
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









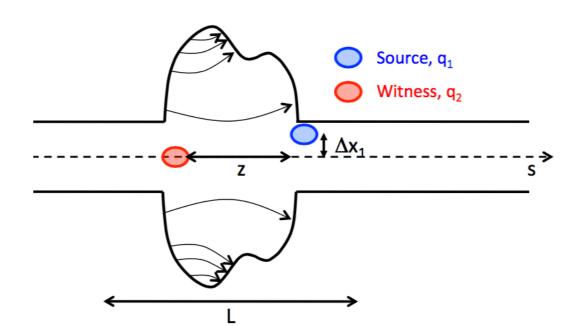
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

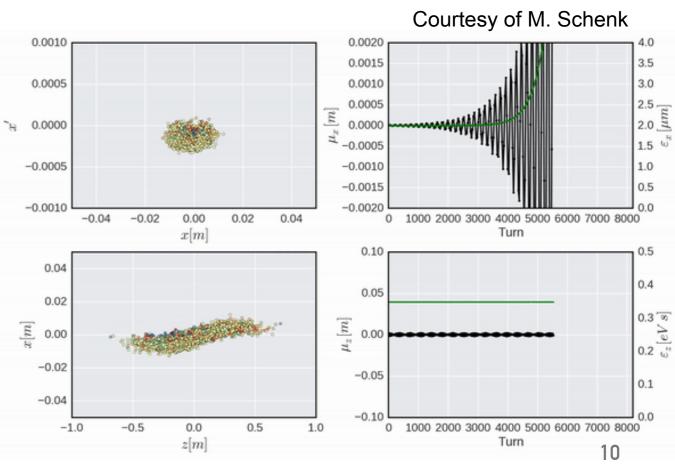
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- Re(ΔQ): real coherent tune shift









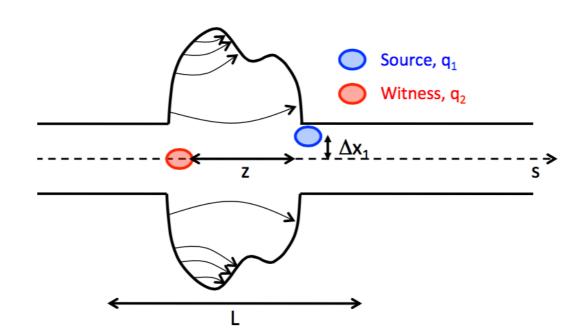
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

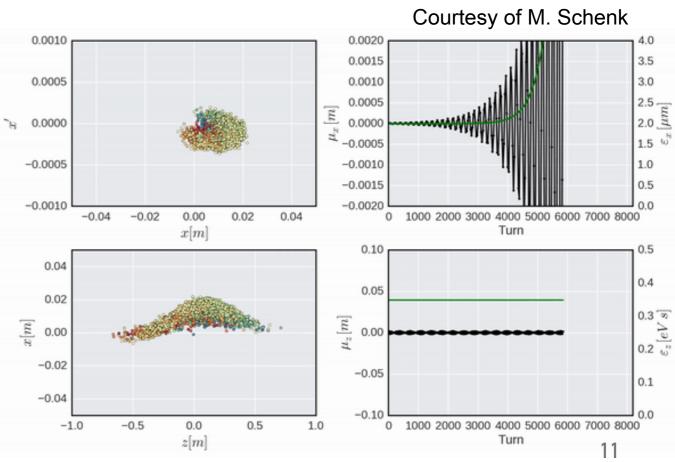
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









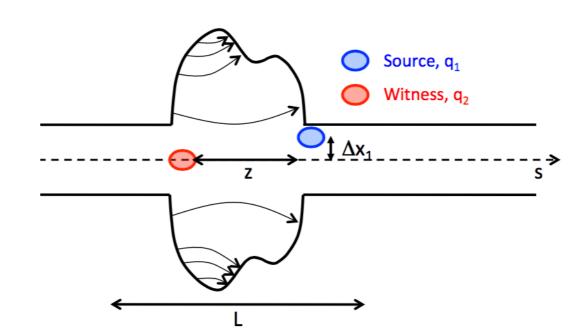
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

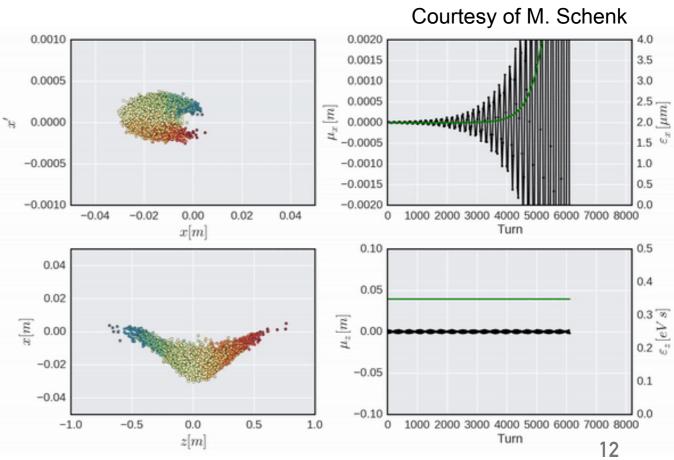
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









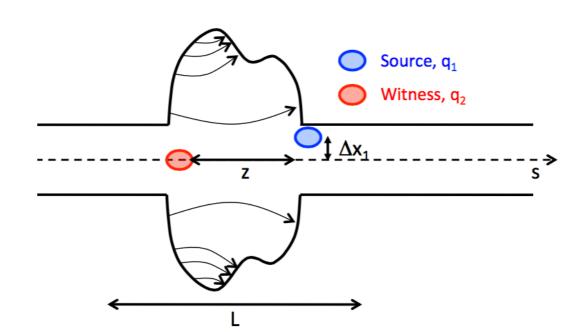
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

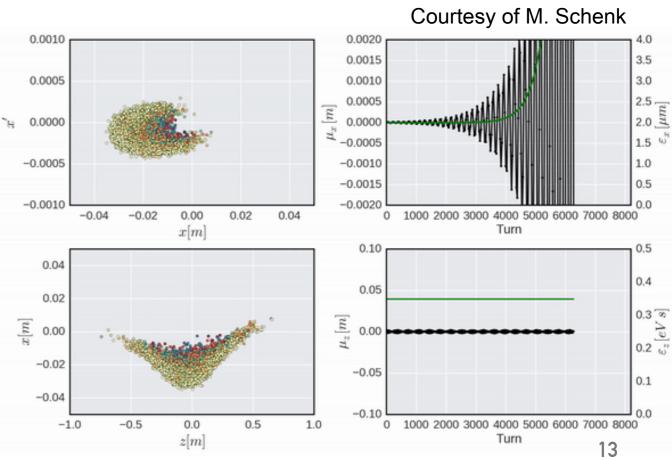
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- Re(ΔQ): real coherent tune shift









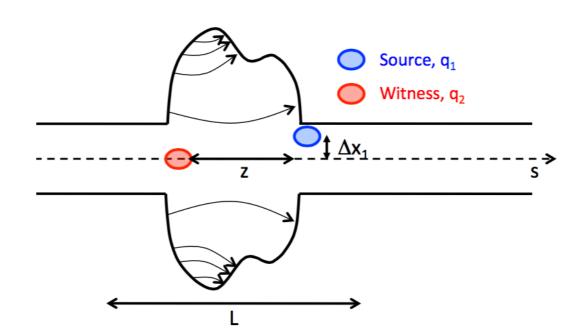
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

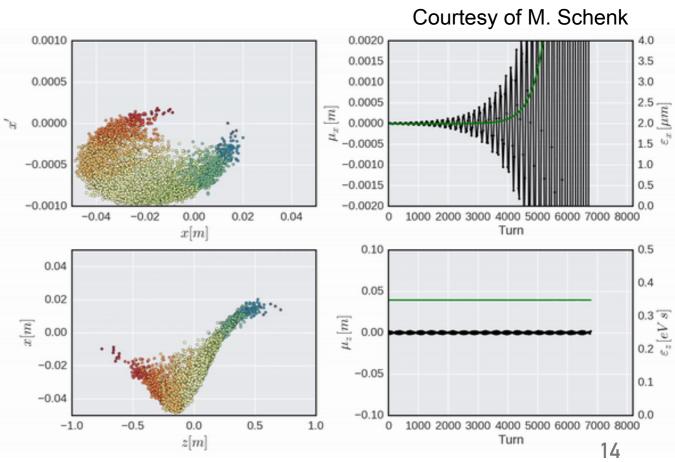
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









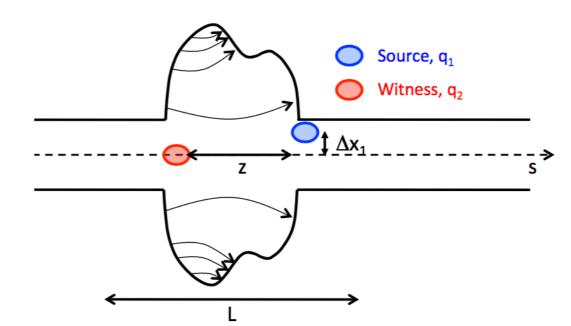
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

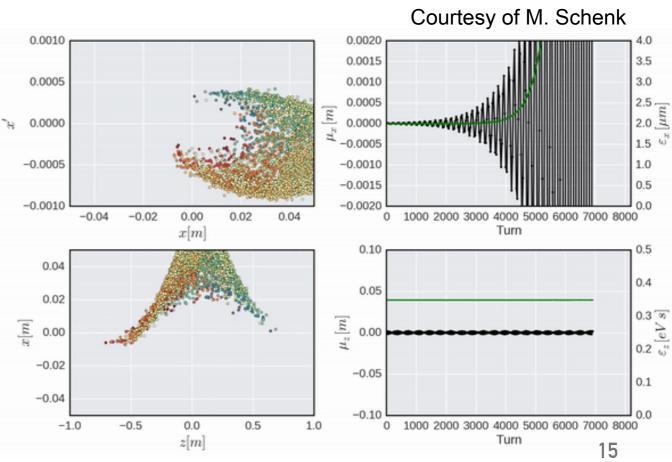
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









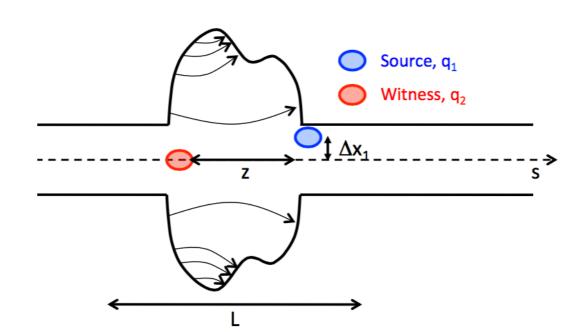
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

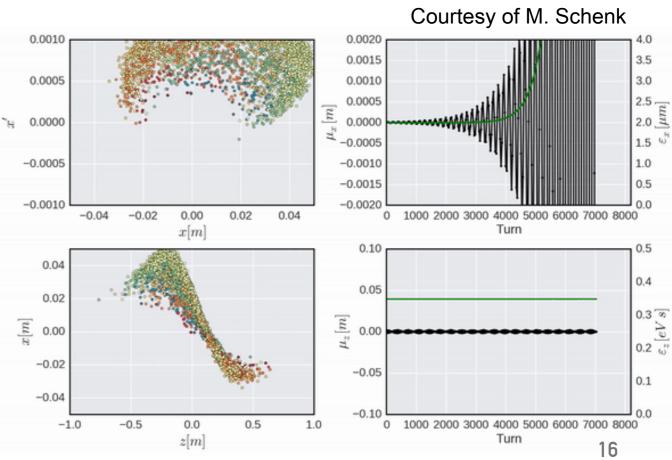
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- $Re(\Delta Q)$: real coherent tune shift









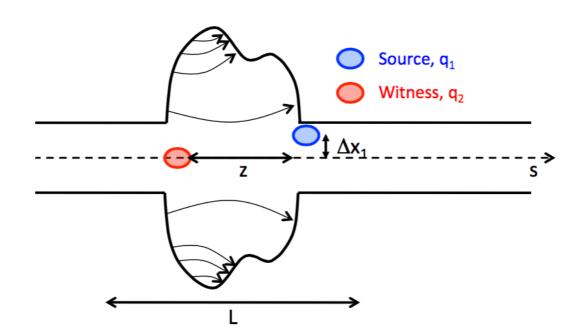
The source particle induces electromagnetic wake fields (**impedance**) that act back on the trailing (witness) particles

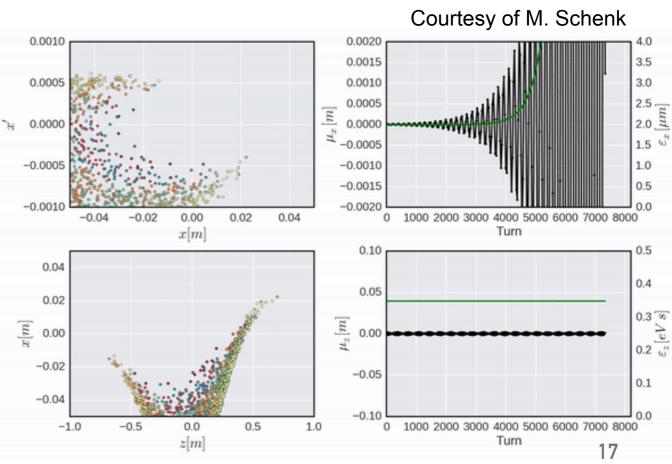
Strong for:

- High brightness beams (N/σ)
- Small aperture elements (collimator 5.5σ)

The impedance drives so-called **head- tail instabilities** → different modes of beam oscillations

- Im(ΔQ): instability rise time
- Re(ΔQ): real coherent tune shift







COHERENT INSTABILITIES AT LHC



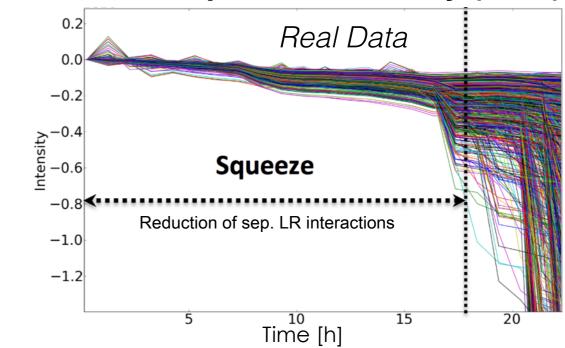
Several coherent instabilities since the first run:

- Emittance blow-up (increase of the rms beam size)
- Loss of intensity
- Coherent oscillations of single bunches

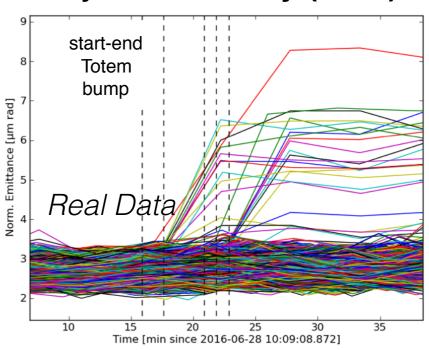
→ REDUCTION OF LUMINOSITY REACH



End of squeeze instability (2012)



Adjust instability (2016)





COHERENT INSTABILITIES AT LHC



Several coherent instabilities since the first run:

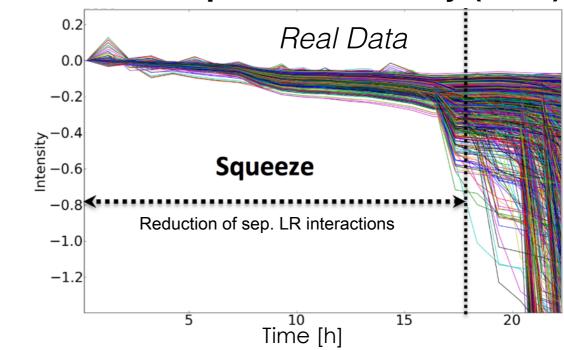
- Emittance blow-up (increase of the rms beam size)
- Loss of intensity
- Coherent oscillations of single bunches

→ REDUCTION OF LUMINOSITY REACH

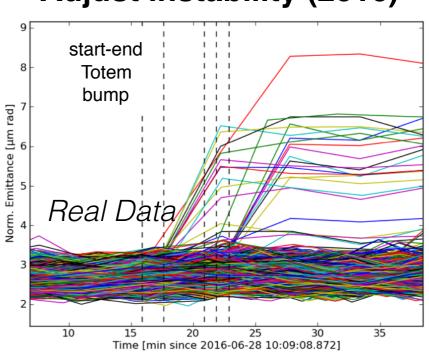


Mitigation techniques: High chromaticity Change of Transverse feedback impedance modes Landau damping Passive mitigation

End of squeeze instability (2012)



Adjust instability (2016)

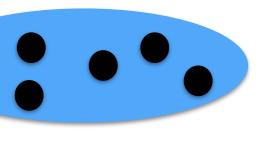






This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

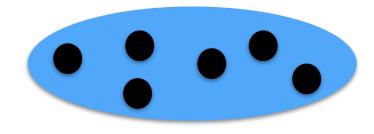






This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

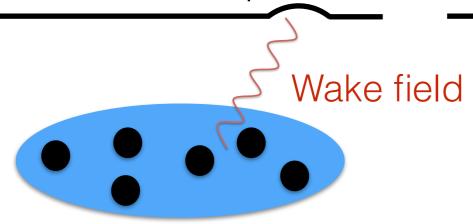






This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

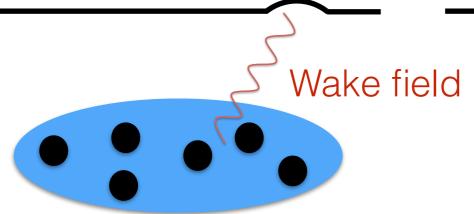






This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

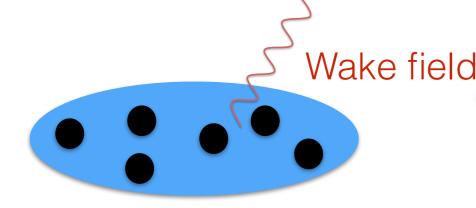


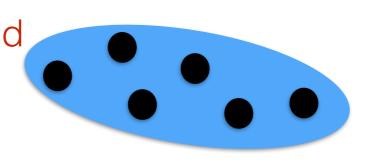




This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies



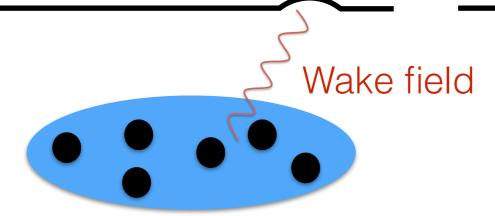


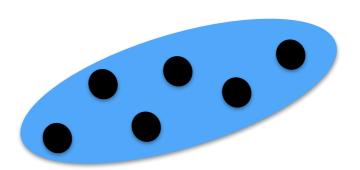




This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies



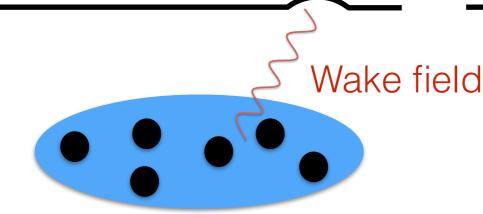


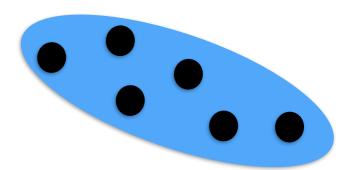




This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies





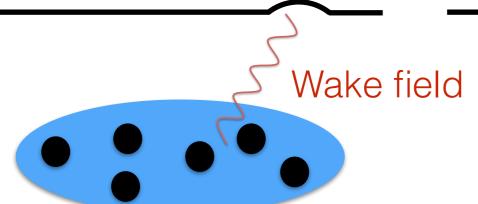




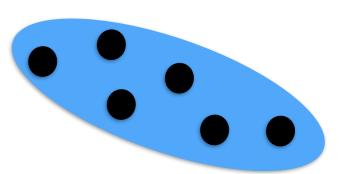
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

Energy of the wake transferred to the coherent motion of the beam



Head tail oscillation: coherent mode (organized motion of the particles)



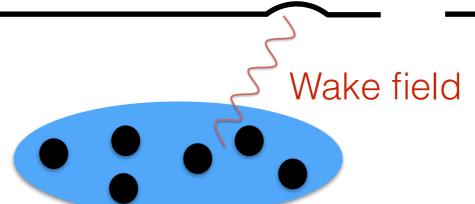




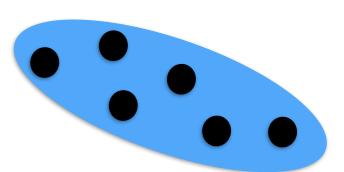
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

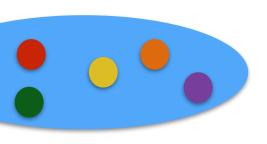
Energy of the wake transferred to the coherent motion of the beam



Head tail oscillation: coherent mode (organized motion of the particles)



Different oscillation frequencies



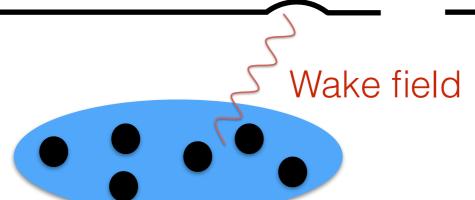




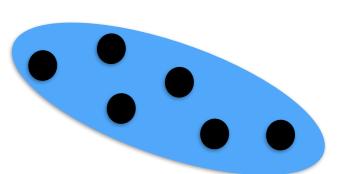
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

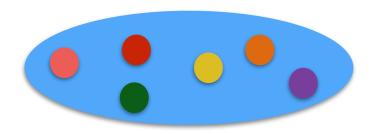
Energy of the wake transferred to the coherent motion of the beam



Head tail oscillation: coherent mode (organized motion of the particles)



Different oscillation frequencies



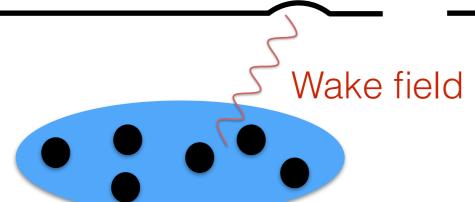




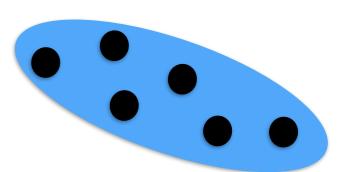
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

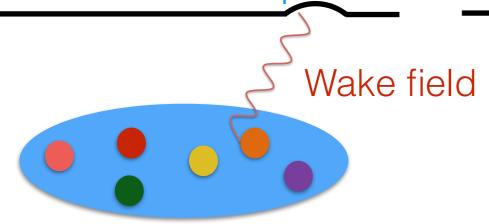
Energy of the wake transferred to the coherent motion of the beam



Head tail oscillation: coherent mode (organized motion of the particles)



Different oscillation frequencies



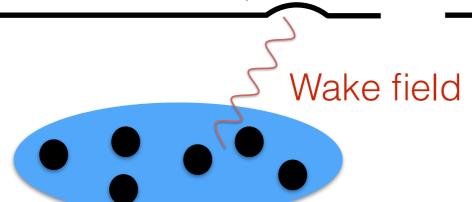




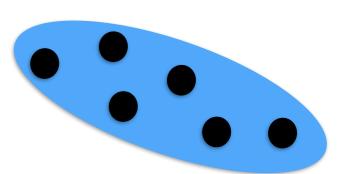
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

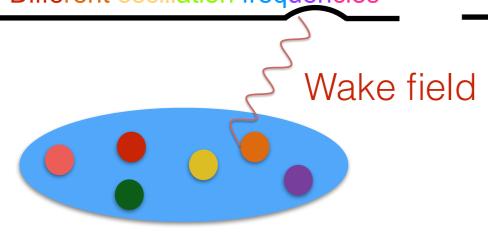
Energy of the wake transferred to the coherent motion of the beam



Head tail oscillation: coherent mode (organized motion of the particles)



Different oscillation frequencies



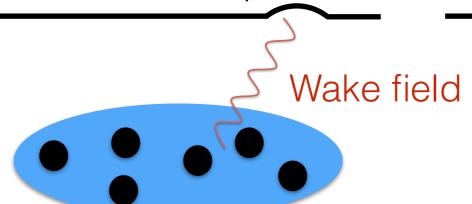




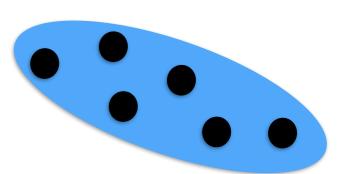
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

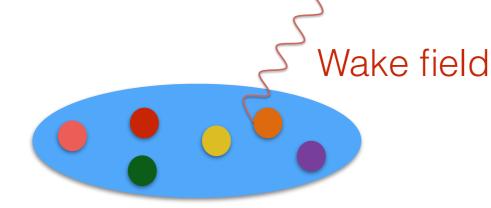
Energy of the wake transferred to the coherent motion of the beam

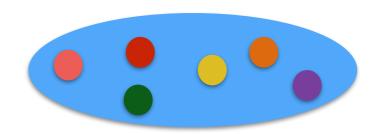


Head tail oscillation: coherent mode (organized motion of the particles)



Different oscillation frequencies





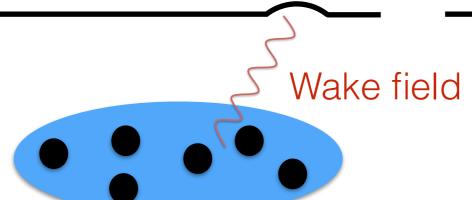




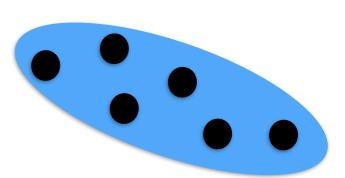
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

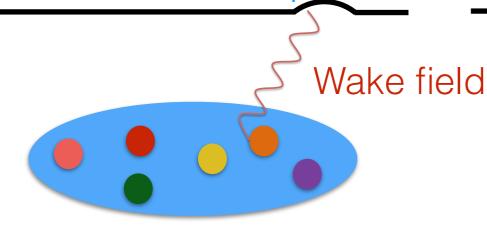
Energy of the wake transferred to the coherent motion of the beam

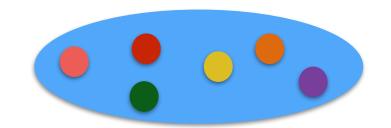


Head tail oscillation: coherent mode (organized motion of the particles)



Different oscillation frequencies





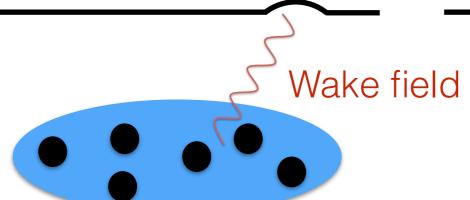




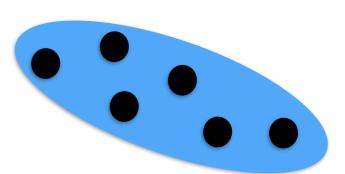
This mechanism **prevents an instability from developing** (region of stability) In accelerators it is provided by the **diversification of oscillation frequencies** of the particles in the beams (**tune spread**)

Same oscillation frequencies

Energy of the wake transferred to the coherent motion of the beam

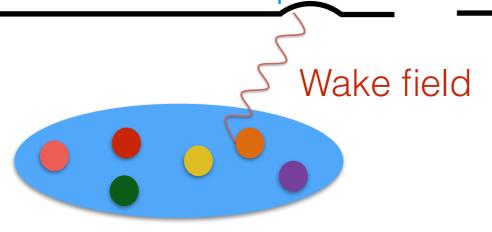


Head tail oscillation: coherent mode (organized motion of the particles)



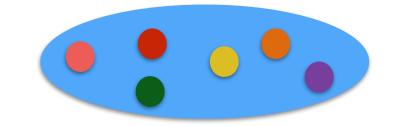
Different oscillation frequencies

The energy of the wake **is not transferred** to the coherent motion of the beam



Landau damping: the coherent oscillation does not develop

The coherent motion is Landau damped!





LANDAU DAMPING AND STABILITY DIAGRAMS



The Landau damping of head-tail instabilities is quantified by solving the dispersion integral:

Particle distribution

$$SD^{-1} = \int_0^\infty \int_0^\infty \frac{J_{x,y}}{Q_0 - (q_{x,y}(J_x,J_y)) - i\epsilon} dJ_x dJ_y$$
 Stable region

Detuning with amplitude (tune spread)
Octupoles + beam-beam (any non-linearities)

 J_x , $J_y \rightarrow$ particle amplitudes (in units of beam size)

 $q_x,q_y \rightarrow$ particle transverse tunes







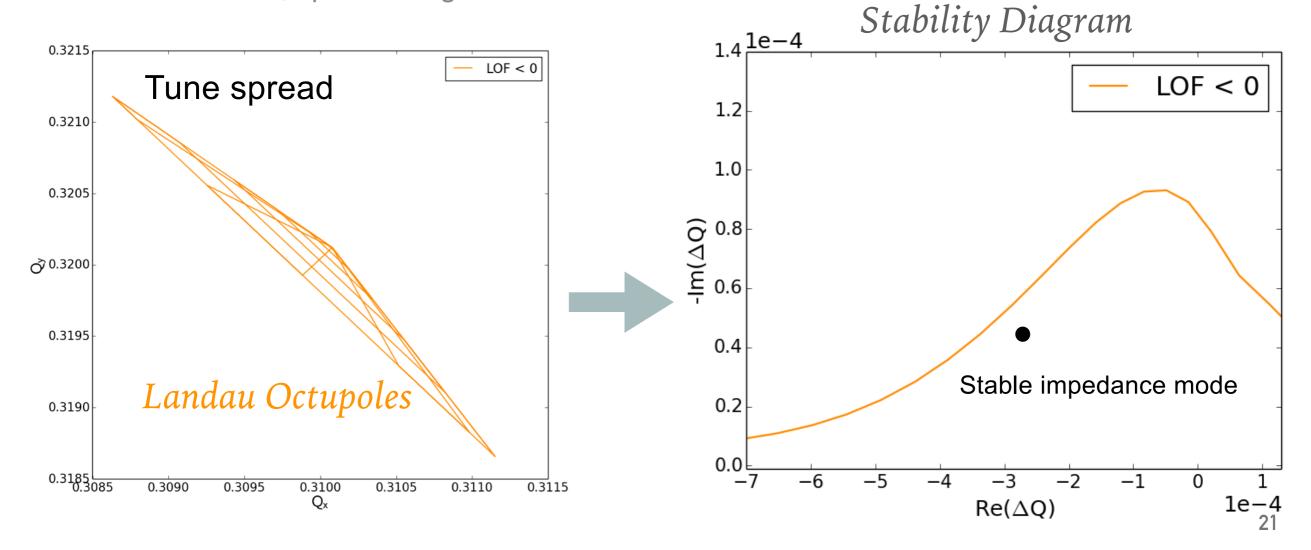


- → Any non-linearities that cause detuning with amplitude (tune spread)
- Main source of tune spread: Landau octupole magnets (linear detuning with amplitude) Non-linear effects:
- beam-beam interactions
- non-linear errors from magnets
- electron cloud, space charge...





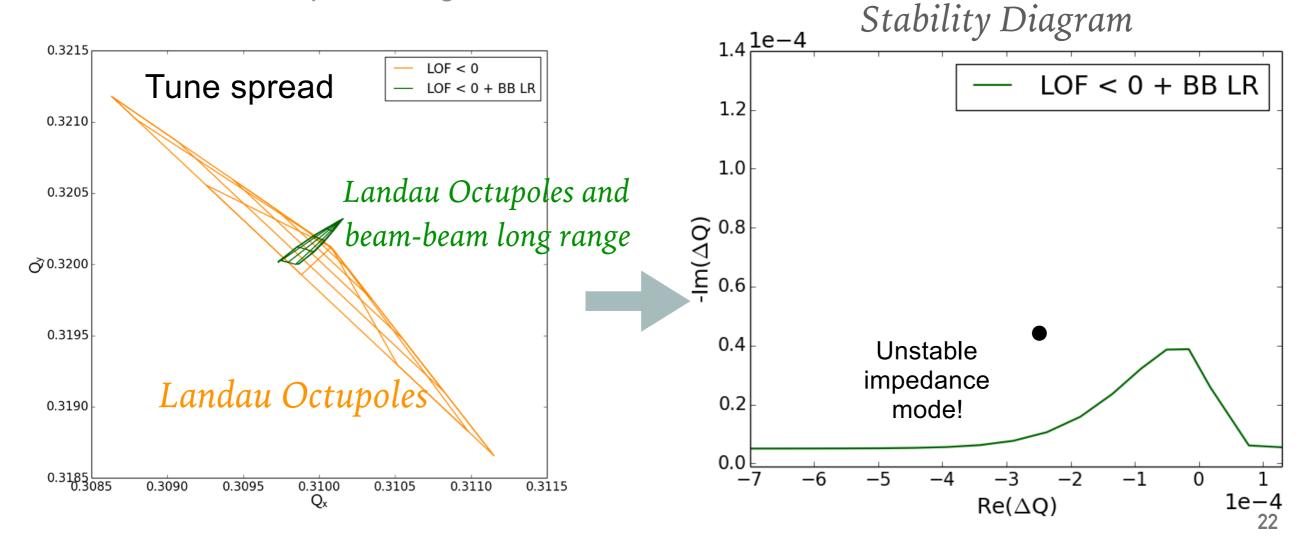
- → Any non-linearities that cause detuning with amplitude (tune spread)
- Main source of tune spread: Landau octupole magnets (linear detuning with amplitude) Non-linear effects:
- beam-beam interactions
- non-linear errors from magnets
- electron cloud, space charge...







- → Any non-linearities that cause detuning with amplitude (tune spread)
- Main source of tune spread: Landau octupole magnets (linear detuning with amplitude) Non-linear effects:
- beam-beam interactions
- non-linear errors from magnets
- electron cloud, space charge...

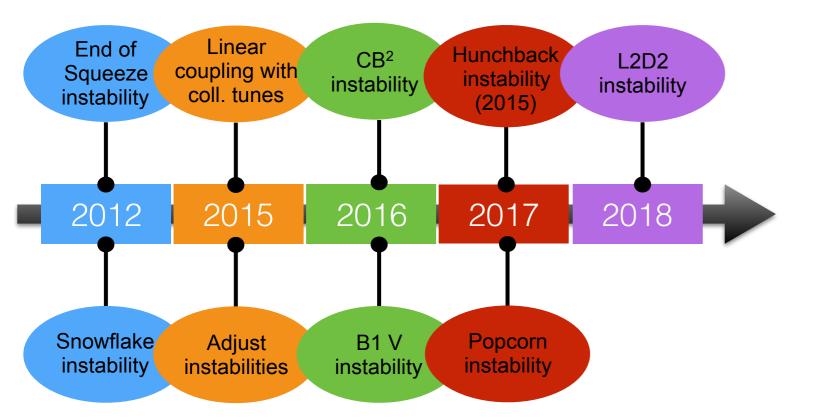




LANDAU DAMPING AND INSTABILITIES IN THE LHC



Models predict stability margins but coherent instabilities present in the LHC during Physics runs

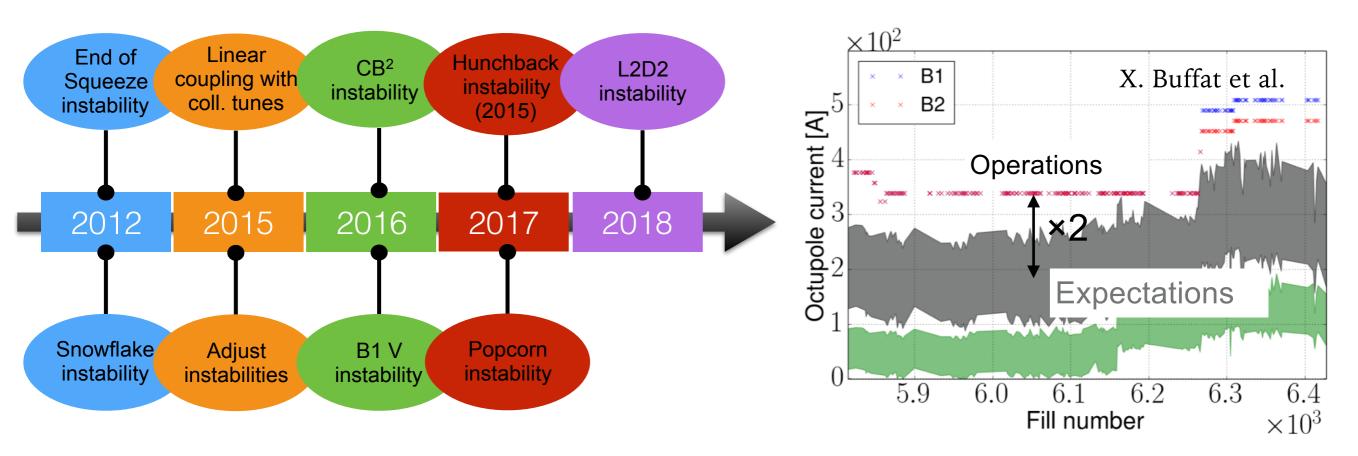




LANDAU DAMPING AND INSTABILITIES IN THE LHC



Models predict stability margins but coherent instabilities present in the LHC during Physics runs



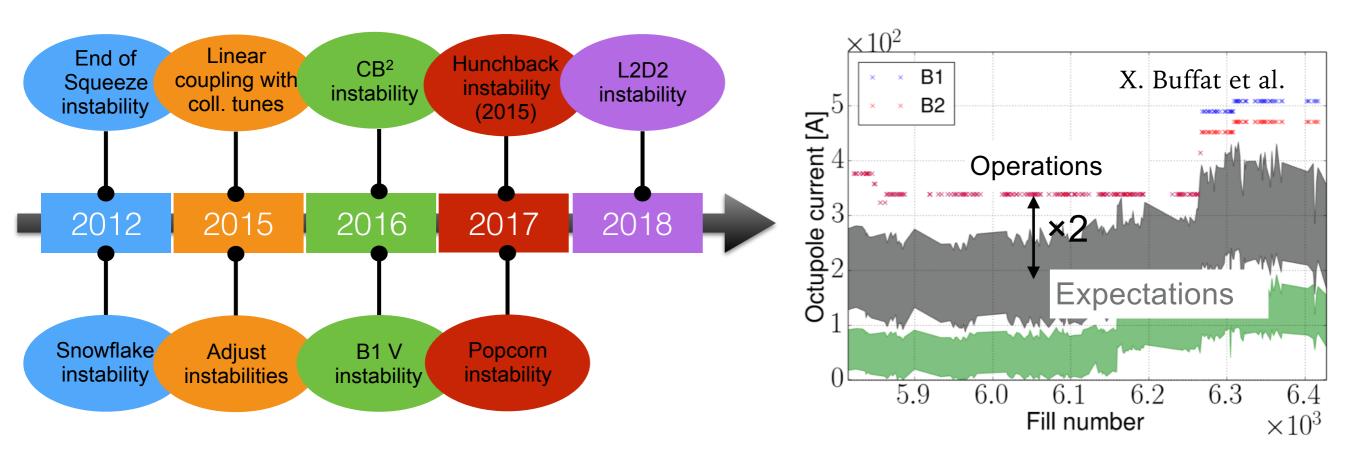
- A factor 2 in the tune spread provided by Landau octupoles is required to stabilize the beams during operations → good understanding of the mechanisms is still needed
- The tune spread provided by the octupole is limited by the beam brightness



LANDAU DAMPING AND INSTABILITIES IN THE LHC



Models predict stability margins but coherent instabilities present in the LHC during Physics runs



- A factor 2 in the tune spread provided by Landau octupoles is required to stabilize the beams during operations → good understanding of the mechanisms is still needed
- The tune spread provided by the octupole is limited by the beam brightness

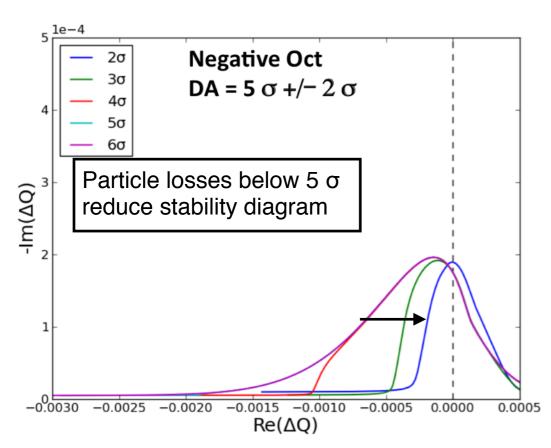
The understanding of the limitations of the models is fundamental in the perspective of future projects that aims to double LHC beam intensities (HL-LHC)

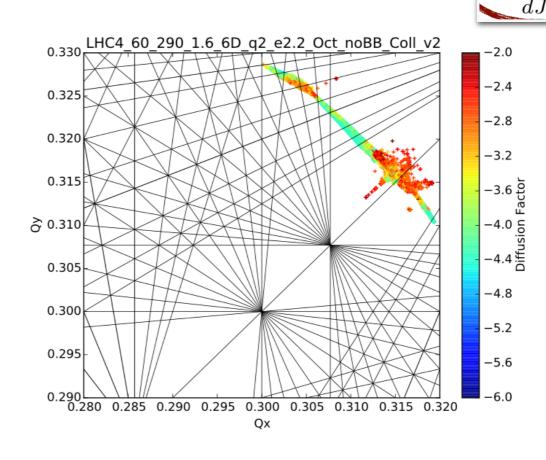


IMPACT OF PARTICLE DISTRIBUTION ON BEAM STABILITY



In presence of diffusive mechanisms the particle distribution changes





- Beam beam long range interactions excite resonances
- Particles are trapped on the resonances and the frequency particle distribution is changing

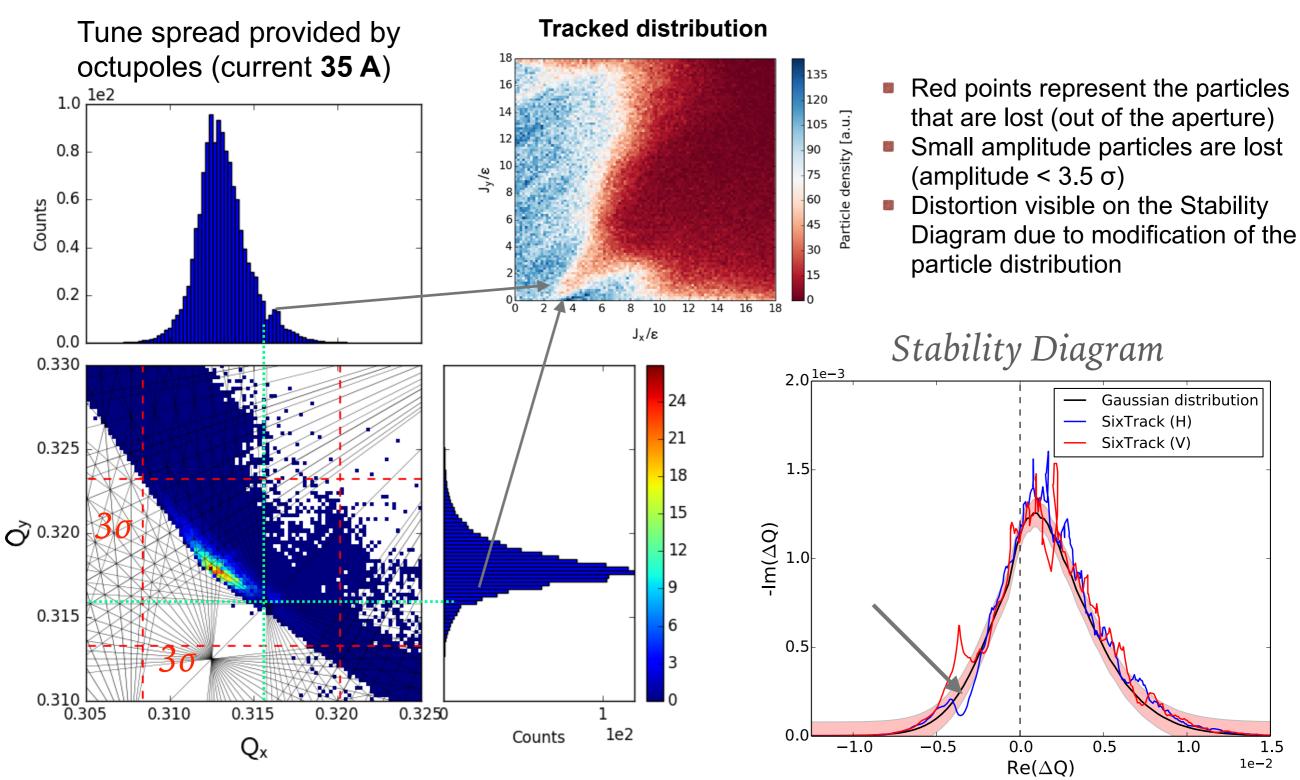
EXTENSION OF MODELS:

Characterize the impact of realistic lattice on particle distribution for the computation of Landau damping → done for the first time!



IMPACT OF PARTICLE DISTRIBUTION ON BEAM STABILITY







BEAM TRANSFER FUNCTION TO MEASURE BEAM STABILITY

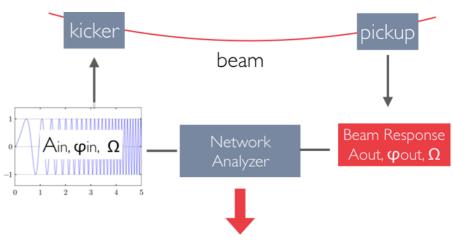


BENCHMARK WITH EXPERIMENTAL DATA:

Installation of Beam Transfer Function system in the LHC to quantify and measure beam stability

Beam Transfer Function measurements are direct measurements of the dispersion integral:

BTF
$$\propto \int_0^\infty \int_0^\infty \frac{J_{x,y} \frac{d\Psi_{x,y}(J_x,J_y)}{dJ_{x,y}}}{Q_0 - q_{x,y}(J_x,J_y) - i\epsilon} dJ_x dJ_y$$

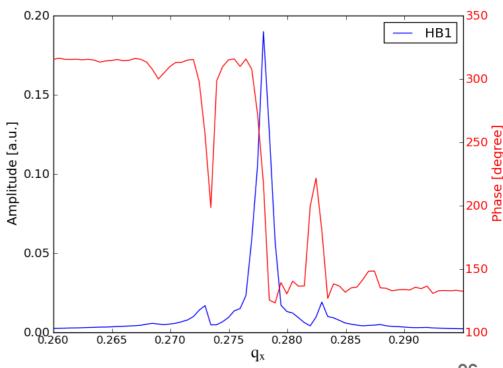


BTF: Amplitude (Ω), phase (Ω)

BTF can experimentally verify the stability

→ direct measurements of SD!

- Tune and chromaticity measurements
- Coherent mode observations
- Sensitive to particle distribution changes
- Tune spread of the beams





FITTING METHOD TO RECONSTRUCT STABILITY DIAGRAM FROM MEASUREMENTS



Uncalibrated system and difficulty to reconstruct stability diagrams from BTF measurements

→ Solution: develop a fitting method that allows quantitative comparisons of measurements with expectations

BTF (complex response) Amplitude (Q)

Phase (Q)

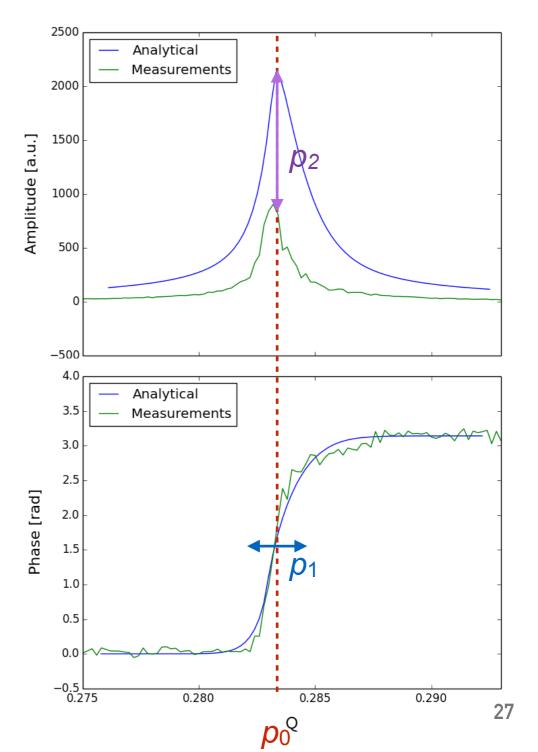
SD
$$_{\times}$$
 1/BTF = A⁻¹ e^{-i ϕ}

Fitting method allows to compare measurements w.r.t. models (reference case, i.e. octupoles)

$$Q_{fit} = p_0 + p_1 \cdot (Q_{analyt} - Q_0)$$

$$A_{fit} = p_2 / p_1 \cdot A_{analyt}$$

 p_0 = Tune p_1 = Tune spread factor w.r.t. a reference case independent from calibration factor, (phase slope) p_2 = Amplitude factor: calibration, proportionality constant





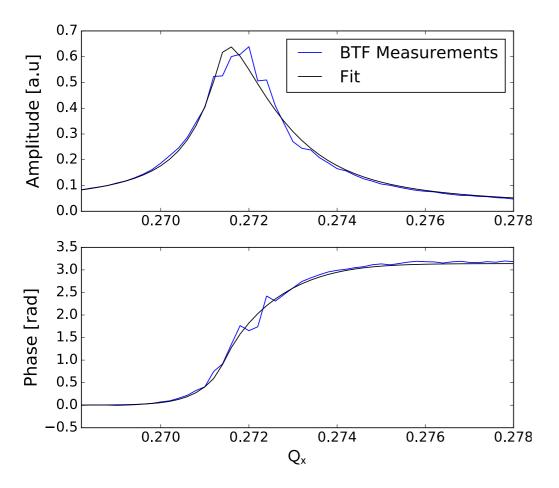
MEASUREMENTS OF STABILITY DIAGRAMS

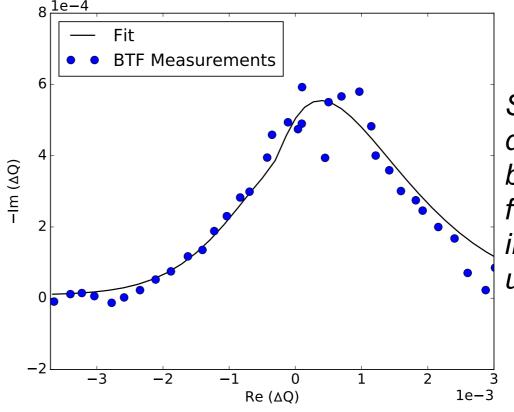


RECONSTRUCTION OF STABILITY DIAGRAM FROM BTF MEASUREMENTS

$$Q_{fit} = p_0 + p_1 \cdot (Q_{analyt} - Q_0)$$

$$A_{fit} = p_2 / p_1 \cdot A_{analyt}$$





Stability
diagrams have
been measured
for the first time
in the LHC by
using BTFs

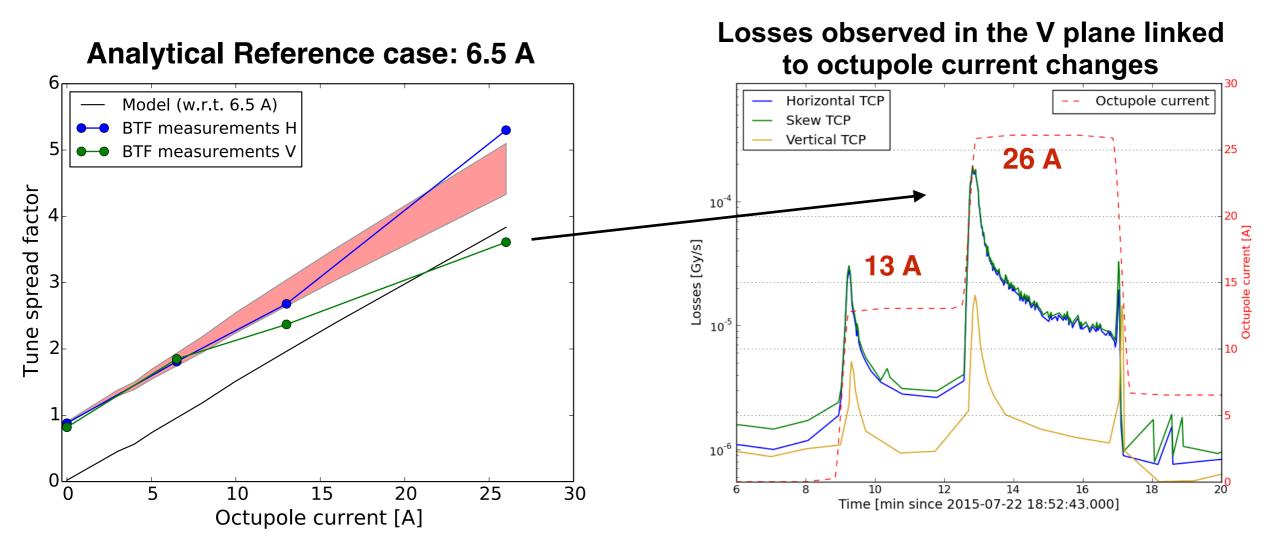
Tune spread p₁=1.71

(10 A octupole current@injection energy)



MEASURED TUNE SPREAD AND EFFECT OF PARTICLE LOSSES





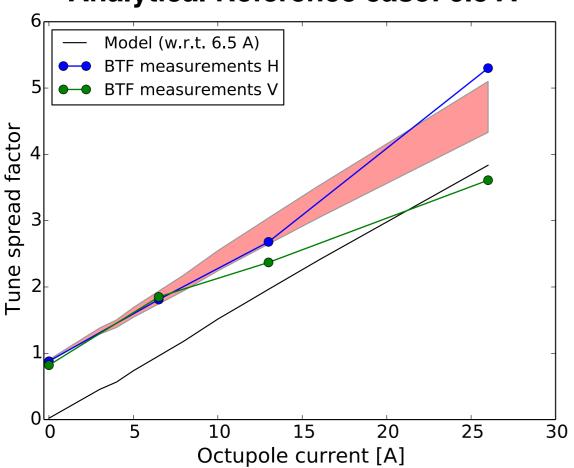
- Fitting method to compare measurements and expectations (tune spread factor)
- Case with no octupoles: consistent with optics measurements in 2015 of spread from magnet non-linearities (equivalent to 5 A octupole spread)
- Linear trend reproduced

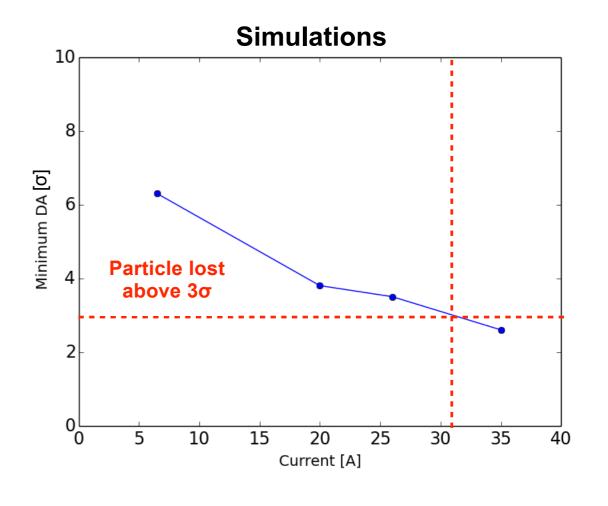


MEASURED TUNE SPREAD AND EFFECT OF PARTICLE LOSSES



Analytical Reference case: 6.5 A





- According to simulations, particles above 3 σ are lost considering the measured tune spread factor that corresponds to ~31 A (= p₁ × ref. case)
- \blacksquare Particle losses at 3 σ reduce Landau damping in the beams

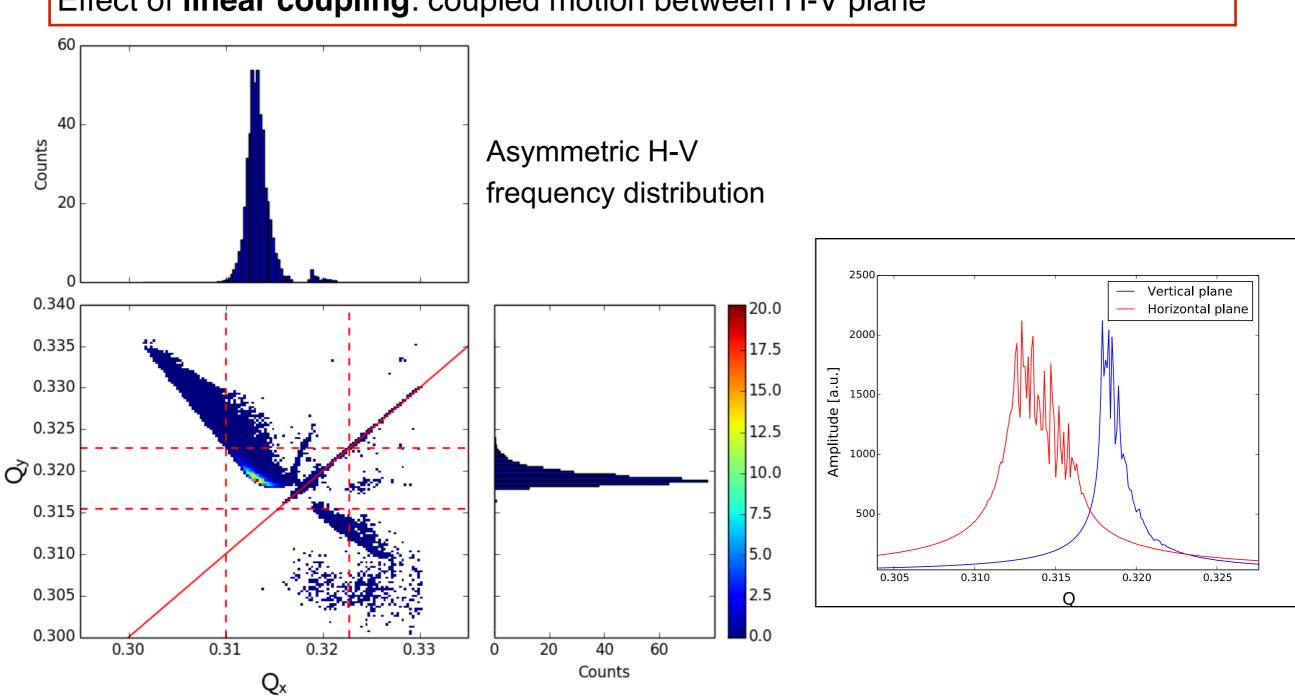
INCREASING THE TUNE SPREAD IS BENEFICIAL FOR LANDAU DAMPING AS LONG AS NO PARTICLE LOSSES ARE PRESENT (FIRST EXPERIMENTAL OBSERVATION)



EFFECT OF LINEAR COUPLING



Effect of linear coupling: coupled motion between H-V plane

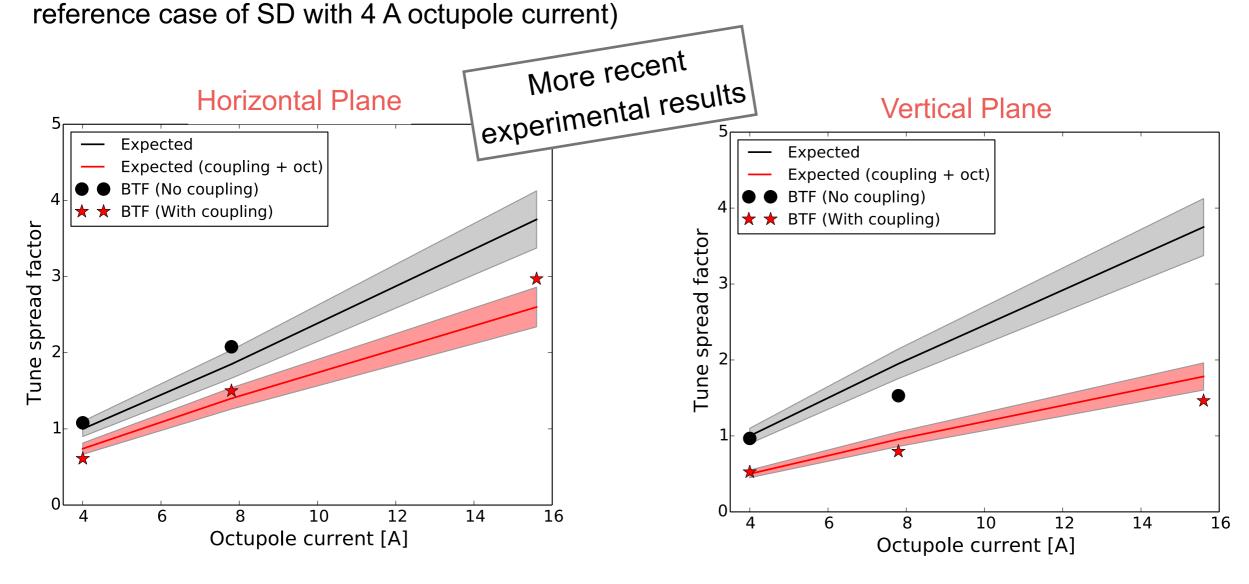




MEASURED EFFECTS OF LINEAR COUPLING ON STABILITY



Fitting function method is applied to measure tune spread from BTFs (w.r.t to an analytical



Quantitative comparison w.r.t expectations (with and without linear coupling)

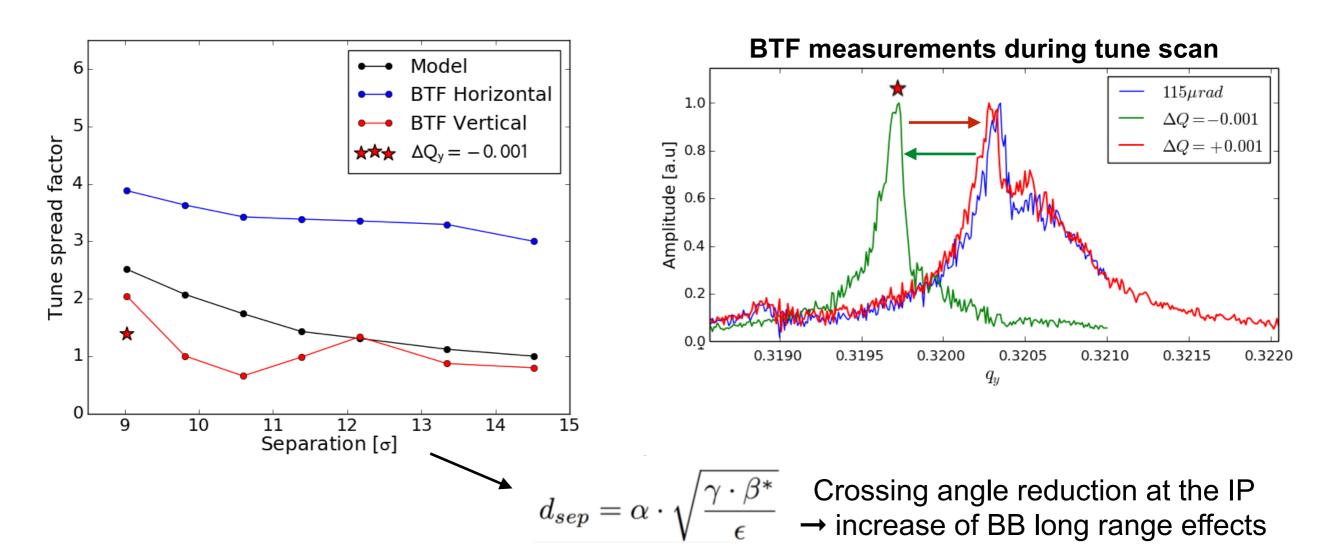
→ BTF measurements agree well with expectations!



BEAM-BEAM INTERACTIONS AND LINEAR COUPLING



Measured beam-beam Long Range contribution on beam stability as a function of BB LR separation

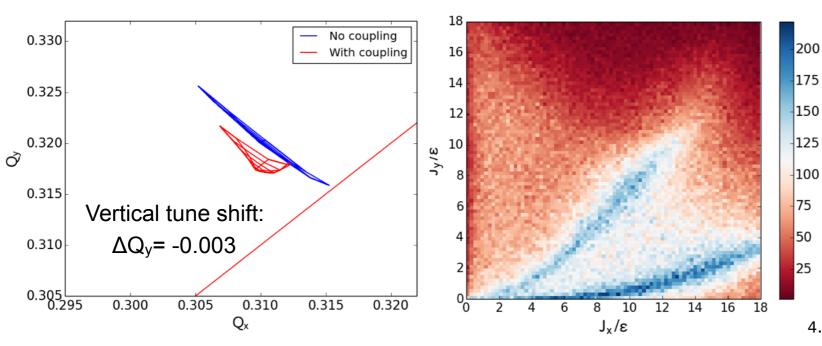


Change of tune has a strong impact on Landau damping indicating that other mechanisms may play a role → **linear coupling**



BTF IN THE PRESENCE OF BEAM-BEAM INTERACTIONS AND LINEAR COUPLING



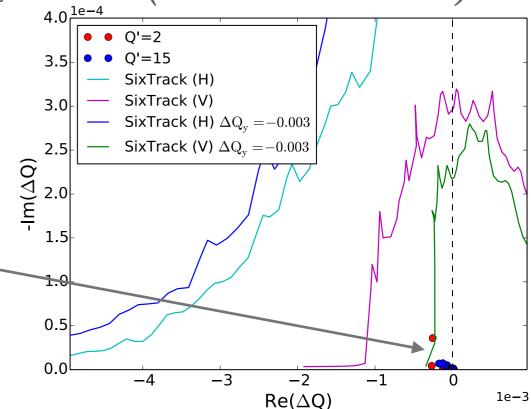


Stability Diagram (tracked distribution)

Particle density [a.u.]

- Strong dependency on tune value
- Sharp cut visible in the vertical SD (0 3 σ particles approach the diagonal)
- Modes can become unstable in the vertical plane

ASYMMETRIC H-V STABILITY DUE TO PARTICLE
DISTRIBUTION MODIFICATIONS IN THE PRESENCE OF
LINEAR COUPLING AND BEAM-BEAM INTERACTIONS THAT
CAN EXPLAIN THE 2012 INSTABILITIES

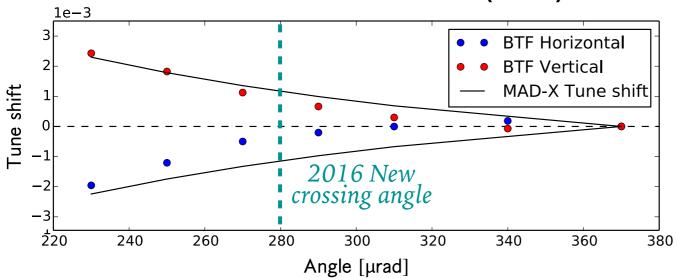




BTF MEASUREMENTS WITH BEAM-BEAM INTERACTIONS



BTF measured tune shifts (2016)



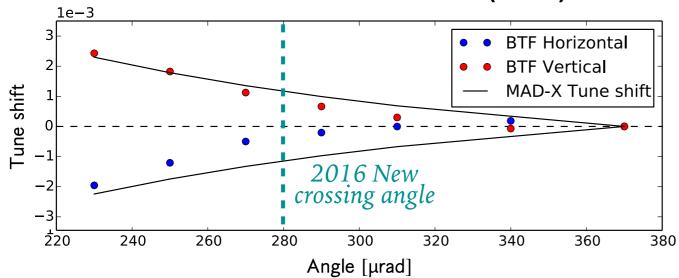
- Unexpected tune shifts were observed in the BTF response due to LR beam-beam effect
- The tunes at the LHC are only measured on the non-colliding beams: not observed during regular operations



BTF MEASUREMENTS WITH BEAM-BEAM INTERACTIONS

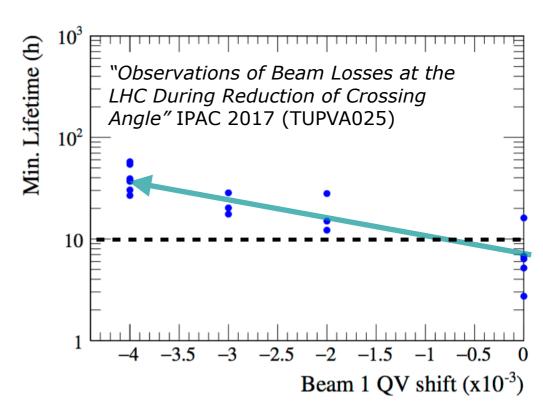


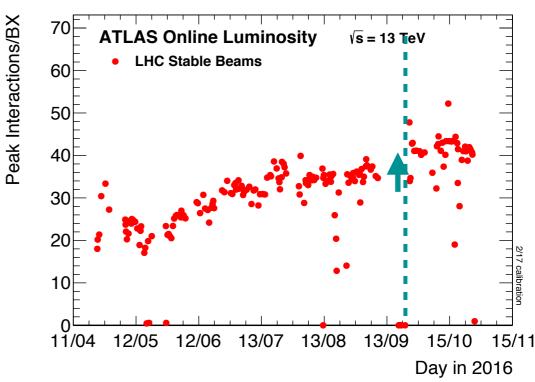
BTF measured tune shifts (2016)



- Unexpected tune shifts were observed in the BTF response due to LR beam-beam effect
- The tunes at the LHC are only measured on the non-colliding beams: not observed during regular operations

CORRECTION OF LONG-RANGE INDUCED TUNE SHIFT
DURING OPERATIONS WITH DIRECT INCREASE OF BEAM
LIFETIMES → INCREASE OF 10% INTEGRATED LUMINOSITY

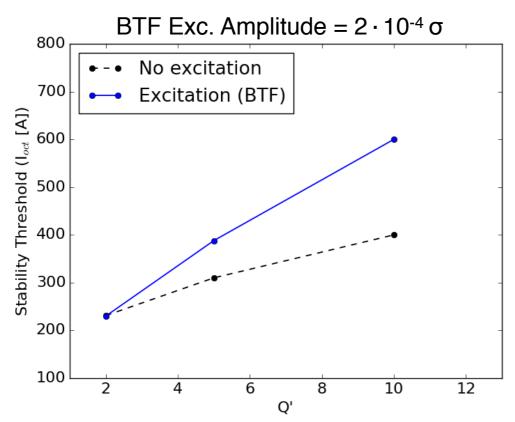






WHAT'S NEXT?

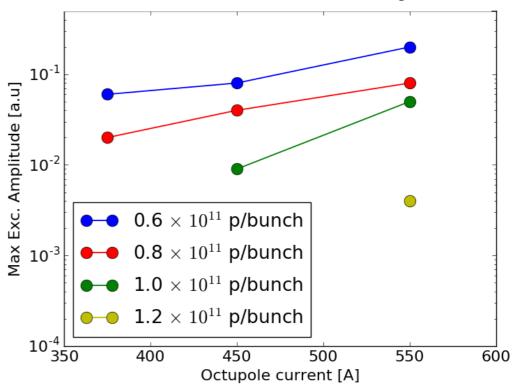




Higher octupole current required in the presence of external excitations (noise, small amplitude external excitation) → could explain the need for higher octupole strength during LHC operation

- A coherent instability linked to BTF excitation has been observed
- The instability was fully reproduced and it is linked to the increase of the impedance in 2017

Excitations at the stability limit



Higher intensity bunches are more sensitive to external excitations (confirmed by recent measurements 2018)

RECENT MEASUREMENTS CONFIRMED THAT THE BEAMS ARE MUCH MORE SENSITIVE TO ANY EXTERNAL SOURCE OF NOISE OR EXCITATION IF IMPEDANCE EFFECTS ARE STRONGER



CONCLUSIONS



- Extended models: the particle distribution affects Landau stability → visible effects due to
 particle distributions after tracking in a realistic lattice configuration
- First reconstruction of measured Stability Diagram in the LHC by using BTF system
- First tune spread measurements in the LHC and comparison to expectations by using the fitting method → measured effects of linear coupling in good agreement with expectations
- It was measured the impact of particle losses on beam stability
- Increase of beam lifetime and luminosity (10% increase in integrated luminosity) by means of correction of Long-Range induced tune shifts measured by BTFs
- Asymmetric Horizontal Vertical beam stability (observed by BTF measurements) reproduced in simulations in the presence of linear coupling and beam-beam interactions
- Experimental studies are planned in the LHC to measure the Landau damping of the beams
 including beam-beam interactions, Landau octupoles, transverse linear coupling and in the
 presence of noise.





Thank you for your attention!





Thank you for your attention!

Particle Accelerator Physics Laboratory (LPAP)

BE/Beam Instrumentation (BI) Group (CERN)

BE/ABP/Hadron Synchrotron Coherent (HSC) effects (CERN)

BE/Operation (OP) Group (CERN)



BACK UP

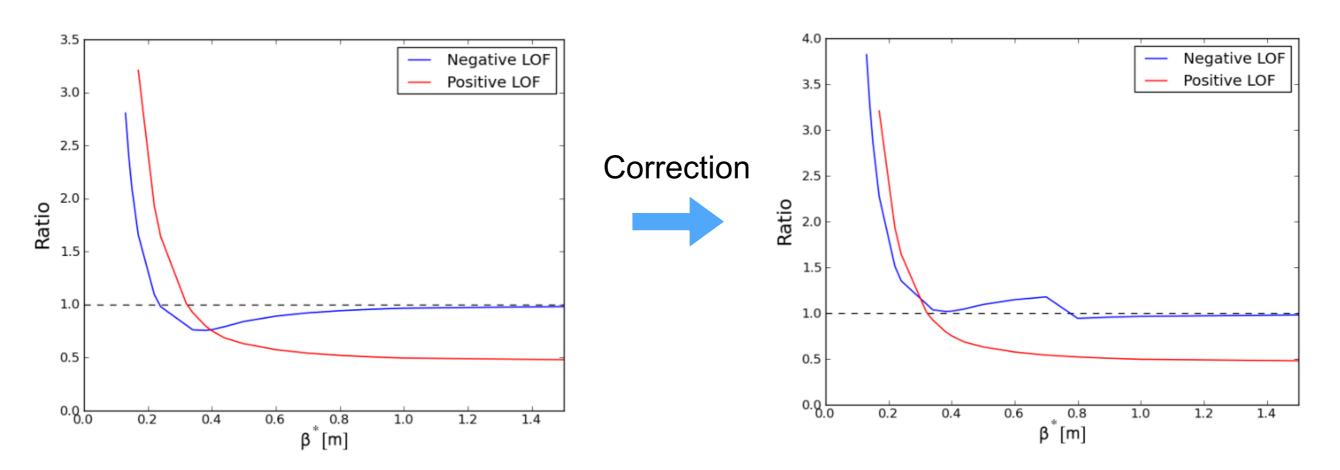




STABILITY STUDIES FOR HL-LHC



The beam stability during the full operational cycle has been studied and maximized



By applying a correction of the β - function (8%) in the arcs (octupole magnets) from β^* =70 cm the stability reduction is compensated and stability is maximized

RESULT: PROPOSED OPERATIONAL SCENARIO WITH MAXIMUM LANDAU STABILITY DURING OPERATIONS

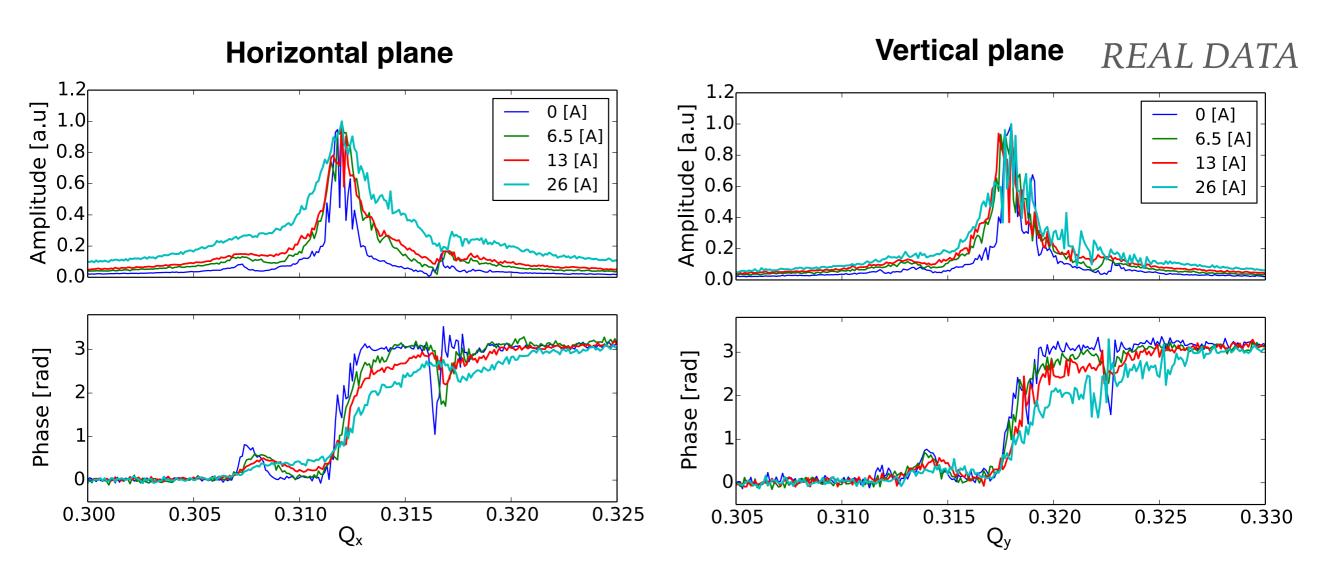
ACCEPTED AS A BASELINE SCENARIO OF HL-LHC



BTF MEASUREMENTS WITH OCTUPOLE MAGNETS



Tune spread given by Landau octupoles and lattice non-linearities



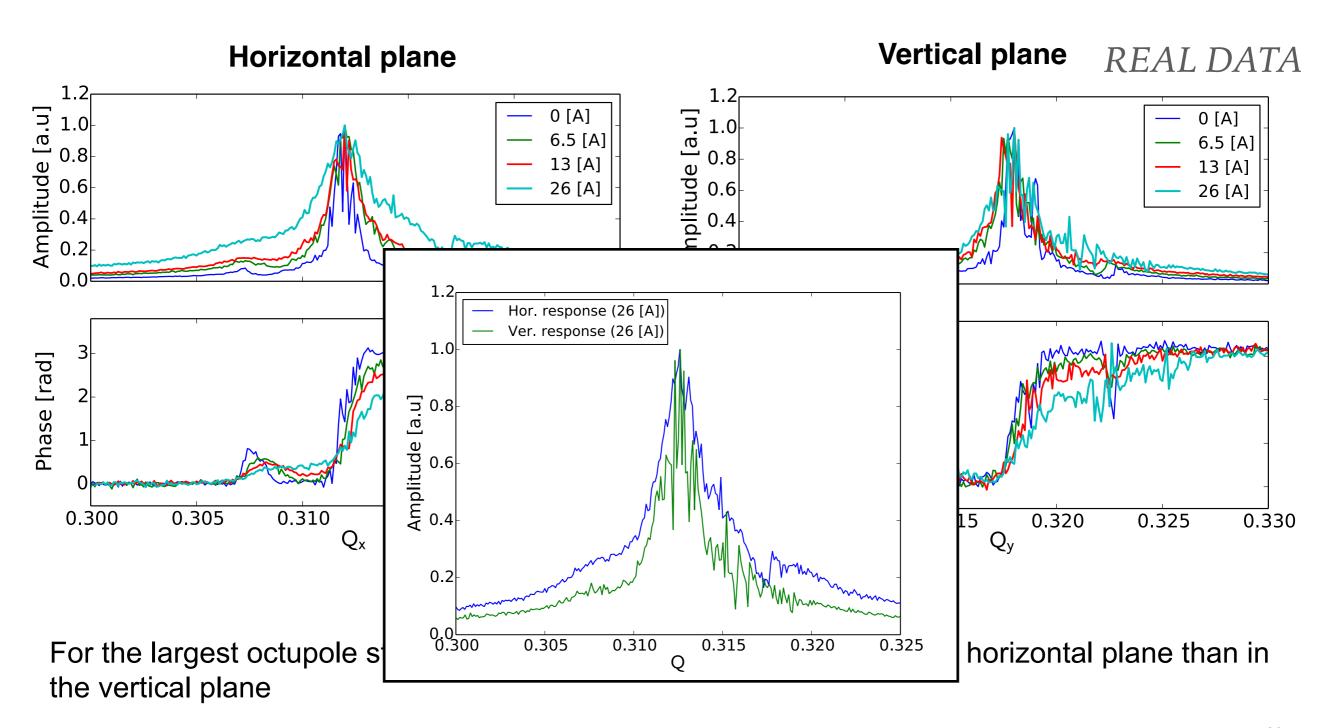
For the largest octupole strength (26 A) larger spread measured in the horizontal plane than in the vertical plane



BTF MEASUREMENTS WITH OCTUPOLE MAGNETS



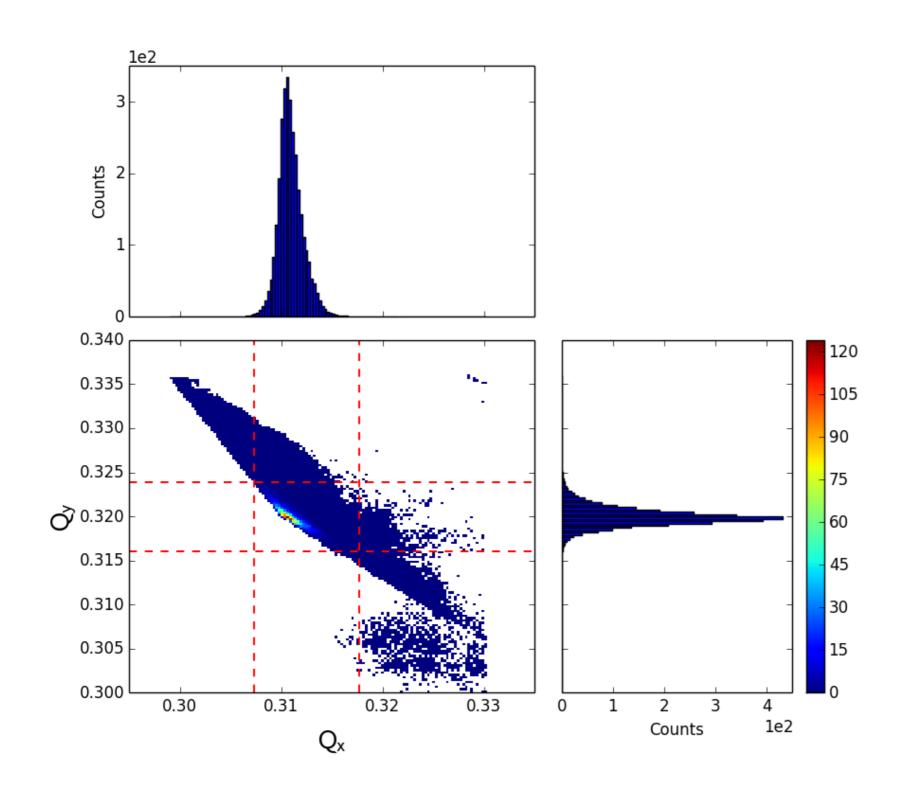
Tune spread given by Landau octupoles and lattice non-linearities





BTF MEASUREMENTS WITH OCTUPOLE MAGNETS





No drastic change in the frequency distribution and it can not explain H-V asymmetry in the BTF amplitude