Hadron+jet correlations in ALICE

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Hard scattering in heavy-ion collisions



- Hard scattered partons produce collimated sprays of particles (back-to-back, pT balanced)
- Jet is a phenomenological object defined via algorithm
- Well understood theoretically in pQCD in elementary reactions
- Hard scattering occurs in early stages of heavy-ion collision
- Jet quenching



Jets in ALICE





- Charged jets: tracks $|\eta| < 0.9$, $0^{\circ} < \varphi < 360^{\circ}$, $p_{T}^{const} > 150 \text{ MeV}/c$
- Full jets: tracks + EMCAL/DCAL clusters, $|\eta| < 0.7$,

EMCAL: 80° < φ < 180°, DCAL: 260° < φ < 327°

Jet reconstruction: anti-k_T algorithm (FastJet package [1])

Given jet R, charged jet acceptance is $|\eta_{\rm jet}| < 0.9 - R$

[1] Cacciari et al., Eur. Phys. J. C 72 (2012) 1896.

Mean background density correction





 $\label{eq:product} \begin{array}{l} \bullet \quad \mbox{Background energy density ρ estimated by area-based method $$_{[1]}$} \\ \rho = \mbox{median}_{k_{\rm T}\,{\rm jets}} \{ p_{{\rm T},{\rm jet}} / A_{\rm jet} \} \end{array}$

event by event

$$p_{\mathsf{T},\mathsf{jet}}^{\mathsf{corr}} = p_{\mathsf{T},\mathsf{jet}} - \rho imes A_{\mathsf{jet}}$$

[1] Cacciari et al., Phys. Lett. B 659 (2008) 119.

Corrections of raw jet spectra





- Detector response: based on GEANT + PYTHIA
- Response matrix:

two effects are assumed to factorize $R_{\text{full}} \left(p_{\text{T,jet}}^{\text{rec}}, p_{\text{T,jet}}^{\text{part}} \right) = \delta p_{\text{t}} \left(p_{\text{T,jet}}^{\text{rec}}, p_{\text{T,jet}}^{\text{det}} \right) \otimes R_{\text{instr}} \left(p_{\text{T,jet}}^{\text{det}}, p_{\text{T,jet}}^{\text{part}} \right)$

- R⁻¹_{full} obtained with Bayesian [2] and SVD [3] unfolding with RooUnfold [4]
- [1] ALICE collab., JHEP 1203 (2012) 053
- [2] D'Agostini, Nucl.Instrum.Meth.A362 (1995) 487
- [3] Höcker and Kartvelishvili, Nucl.Instrum.Meth.A372 (1996) 469
- [4] http://hepunx.rl.ac.uk/~adye/software/unfold/RooUnfold.html







- Hard scattering, rare process embedded in large background
- Spectrum of reconstructed jets at low p_T dominated by combinatorial jets
- Suppression of combinatorial jets by high-p_T jet constituent requirement results in fragmentation bias on jets

Hadron-jet coincidence measurement





 $\mathsf{TT} = \mathsf{trigger} \; \mathsf{track}$

 $TT{X,Y}$ means $X < p_{T,trig} < Y \text{ GeV}/c$

- h-jet correlation allows to suppress combinatorial jets including multi parton interaction without imposing fragmentation bias
- \diamond Data driven approach allows to measure jets with large R and low $p_{\rm T}$
- \diamond In events with a high- p_{T} trigger hadron analyze recoiling away side jets ${}_{\rm [1]}$

 $|arphi_{\mathsf{trig}} - arphi_{\mathsf{jet}} - \pi| < \mathsf{0.6} \; \mathsf{rad}$

 \diamond Assuming combinatorial jets are independent of trigger $p_{\rm T}$



$$\Delta_{\text{recoil}} = \frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{jet}}}{d\rho_{\text{T,jet}}^{\text{ch}} d\eta} \Big|_{\rho_{\text{T,trig}} \in \text{TT}\{20,50\}} - \frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{jet}}}{d\rho_{\text{T,jet}}^{\text{ch}} d\eta} \Big|_{\rho_{\text{T,trig}} \in \text{TT}\{8,9\}}$$

$$\diamond \text{ Link to theory } \frac{1}{N_{\text{trig}}^{\text{ch}}} \frac{d^2 N_{\text{jet}}^{\text{AA}}}{d\rho_{\text{T,jet}}^{\text{ch}} d\eta_{\text{jet}}} \Big|_{\rho_{\text{T,trig}} \in \text{TT}} = \left(\frac{1}{\sigma^{\text{AA} \rightarrow h+X}} \cdot \frac{d^2 \sigma^{\text{AA} \rightarrow h+jet+X}}{d\rho_{\text{T,jet}}^{\text{ch}} d\eta_{\text{jet}}}\right) \Big|_{\rho_{\text{T,h}} \in \text{TT}}$$

$$\uparrow 0^{10^2} \frac{10^2 (10^2 - 0.00\% \text{ Pb-Pb} \sqrt{\text{s}} = 2.76 \text{ TeV})}{A \text{LICE} - 0.10\% \text{ Pb-Pb} \sqrt{\text{s}} = 2.76 \text{ TeV}} \frac{1}{A + 0.6} \frac{1}{20,50} - 17(8,9)} \frac{10^2 (10^2 - 0.00\% \text{ Pb-Pb} \sqrt{\text{s}} = 2.76 \text{ TeV})}{10^4 + 0.50} \frac{\pi - \Delta \varphi < 0.6}{17(20,50) - 17(8,9)} \frac{1}{10^4 + 0.50} \frac{1}{20,50} - 17(8,9)} \frac{1}{10^4 + 0.50} \frac{1}{10^4 + 0.50$$





- Reference $\Delta_{\text{recoil}}^{\text{PYTHIA}}$ from PYTHIA Perugia 10
- Suppression of the recoil jet yield
- Magnitude of the suppression similar for different R

More details in ALICE collab., JHEP 09 (2015), 170

Ratios of recoil jet yields obtained with different R





- Red band: variation in observable calculated using PYTHIA tunes
- No evidence for significant energy redistribution w.r.t. PYTHIA up to jets with R = 0.5

QGP signatures in small systems



- Indication of collective effects in p-Pb
- Is there jet quenching in p-Pb?
- Considerations
 - $\diamond \Delta E \propto \hat{q} L^2$

BDMPS, Nucl. Phys. B483 (1997) 291

 $\diamond ~ \hat{q} |_{ extsf{pPb}} = rac{1}{7} \hat{q} |_{ extsf{PbPb}}$

K.Tywoniuk, Nucl.Phys. A 926 (2014) 85-91

 $\diamond ~\hat{q}|_{ ext{PbPb}} = (1.9 \pm 0.7) \, ext{GeV}^2/ ext{fm}$

JET Collaboration, Phys.Rev. C 90, 014909 (2014)

 $\diamond ~ \hat{q} |_{\mathsf{Cold\,Nuclear\,Matter}} \approx 0.02 \, \mathsf{GeV}^2 / \mathsf{fm}$

W.T.Deng, X.N.Wang, Phys.Rev. C 81, 024902 (2010)

 $\Delta E = (8 \pm 2_{\text{stat}}) \text{ GeV}/c \text{ medium-induced } E \text{ transport to } R > 0.5 \text{ in Pb-Pb}$ ALICE, JHEP 09 (2015) 170



ALICE, Phys.Lett. B 719 (2013) 29-41

Jet R_{dAu} in d+Au at $\sqrt{s_{NN}} = 200 \text{ GeV}$

$$R_{
m dAu} = rac{{
m d}N_{
m jets}^{
m cent}/{
m d}p_{
m T}}{T_{
m dAu}\cdot{
m d}\sigma_{
m pp}/{
m d}p_{
m T}}$$

R_{dAu} for MB compatible with unity

Event Activity strongly affects R_{dAu}

EA from BBC in Au-going direction 3 $< \left| \eta \right| <$ 3.9

EA = Event Activity

PHENIX, Phys. Rev. C94, 064901 (2016)



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Event Activity biased jet measurements in p-Pb at LHC



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Jet $\mathit{R}_{\rm pPb}$ in p–Pb at $\sqrt{\mathit{s}_{\rm NN}}=5.02\,{\rm TeV}$

EA from $E_{\rm T}$ in Pb-going direction $-4.9 < \eta < -3.2$

Caveats:

- *T*_{pPb}, *T*_{dAu} assume EA correlated with geometry (Glauber modeling)
- Conservation laws and fluctuations

Kordell, Majumder, arXiv:1601.02595v1

Alternative:

h-jet correlations conditional yields





Calculable at NLO pQCD [1] $\underbrace{\frac{1}{N_{\text{trig}}^{\text{AA}}} \frac{d^2 N_{\text{jet}}^{\text{AA}}}{dp_{\text{T,jet}}^{\text{ch}} d\eta_{\text{jet}}} \Big|_{p_{\text{T,trig}} \in \text{TT}}}_{\text{measured}} = \underbrace{\left(\frac{1}{\sigma^{\text{AA} \rightarrow \text{h} + X}} \cdot \frac{d^2 \sigma^{\text{AA} \rightarrow \text{h} + \text{jet} + X}}{dp_{\text{T,jet}}^{\text{ch}} d\eta_{\text{jet}}}\right) \Big|_{p_{\text{T,h}} \in \text{TT}}}_{\text{from theory}}$ In case of no nuclear effects $\frac{1}{N_{\text{trig}}^{\text{AA}}} \frac{d^2 N_{\text{jet}}^{\text{AA}}}{dp_{\text{trig}}^{\text{ch}} d\eta_{\text{jet}}} \Big|_{\sigma_{\text{T,iet}} \in \text{TT}} = \left(\frac{1}{\sigma^{\text{pp} \rightarrow \text{h} + X}} \cdot \frac{d^2 \sigma^{\text{pp} \rightarrow \text{h} + \text{jet} + X}}{dp_{\text{trig}}^{\text{ch}} d\eta_{\text{jet}}}\right) \Big|_{\sigma_{\text{T,iet}} \in \text{TT}} \times \frac{T_{\text{AA}}}{T_{\text{AA}}}$

- > This coincidence observable is self-normalized, no requirement of $T_{\rm AA}$ scaling
- No requirement to assume correlation between Event Activity and collision geometry, no Glauber modeling

Event Activity in p–Pb at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$









High-p_T track requirement (TT) biases event to large EA
 Similar EA bias for TT 6–7 GeV/c and 12–50 GeV/c







- Correction via unfolding for local bkgd. fluct. and instrumental effects
- \blacktriangleright Systematic uncertainties on Δ_{recoil} : tracking efficiency $4\text{--}10\,\%$ other sources $<4\,\%$

Ratios of Event Activity biased Δ_{recoil} distributions



ZNA





Ratio



compatible with unity

Systematic	uncertainties:
unfolding	3–8 %
other sourc	es < 4 %

Correlated syst. uncert. in numerator and denominator cancel

0.61-0

Syst. uncert.

20 25 30

0.4 GeV/c spectrum jet shift

35

45

50

40

 $p_{\rm T,jet}^{\rm ch}$ (GeV/c)





• Express the suppression in terms of energy shift \overline{s}



Parameterize

$$\Delta_{\text{recoil}}|_{50-100\%} = a \exp\left(-\frac{p_{\text{T,jet}}^{\text{ch}}}{b}\right)$$

 Assume parton energy loss causes average shift of Δ_{recoil} by s̄ independent of p^{ch}_{T,iet}

$$\Delta_{
m recoil}|_{0-20\%} = a \exp\left(-rac{
ho_{
m T,jet}^{
m ch}+ar{s}}{b}
ight)$$

the same a and b as for $\Delta_{
m recoil}|_{50-100\,\%}$

$$\frac{\Delta_{\rm recoil}|_{0-20\%}}{\Delta_{\rm recoil}|_{50-100\%}} = \exp\left(-\frac{\bar{s}}{b}\right)$$



Limits on energy transport out of R = 0.4 cone in p-Pb





- Shift for high EA (0–20%) relative to low EA (50–100%) p–Pb
 - $$\begin{split} \bar{s} &= (-0.06 \pm 0.34_{\rm stat} \pm 0.02_{\rm syst}) \; \text{GeV}/c \; \text{for VOA} \\ \bar{s} &= (-0.12 \pm 0.35_{\rm stat} \pm 0.03_{\rm syst}) \; \text{GeV}/c \; \text{for ZNA} \\ \bar{s} &= (8 \pm 2_{\rm stat}) \; \text{GeV}/c \; \text{in Pb-Pb} \qquad \text{Alice, Jhep 09 (2015) 170} \end{split}$$
- Medium-induced charged energy transport out of R = 0.4 cone is less than 0.4 GeV/c (one sided 90% CL)



 h+jet technique alows to measure jet quenching in heavy-ion collisions and small systems

Summary

- does not require the assumption that Event Activity is correlated with collision geometry
- provides systematically well-controlled comparison of jet quenching as a function of Event Activity
- ▶ Pb–Pb at $\sqrt{s_{NN}} = 2.76$ TeV: suppression of recoil jet yield, but no evidence of intra-jet broadening of energy profile out to R = 0.5
- ▶ p-Pb at $\sqrt{s_{\rm NN}} = 5.02 \,{\rm TeV}$: no significant quenching effects are observed when comparing recoil jet yield for low and high Event Activity for both EA metrics. At 90% CL, medium-induced charged energy transport out of R = 0.4 cone is less than $0.4 \,{\rm GeV}/c$