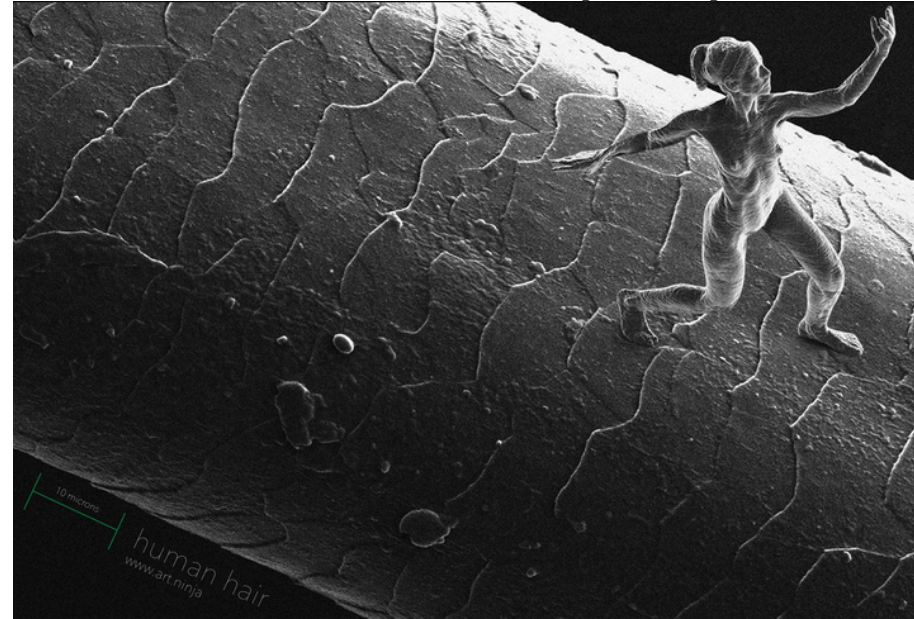
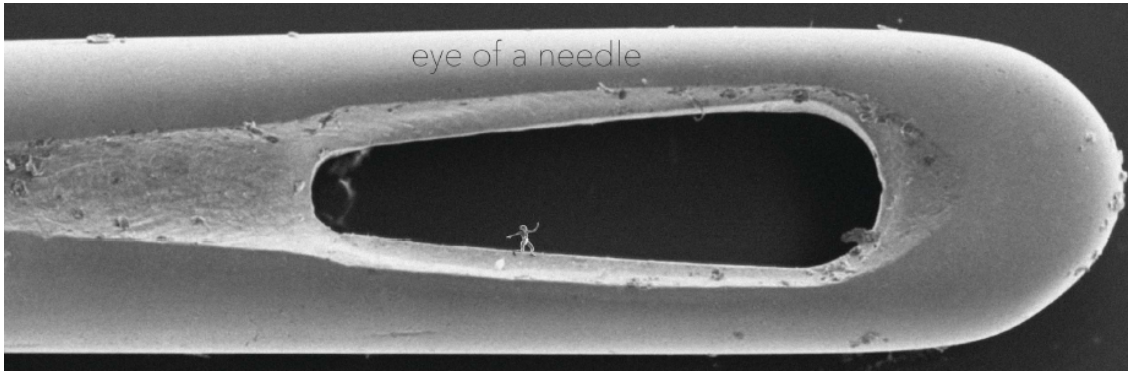


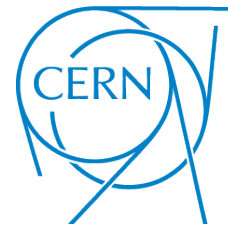
Art by Jonty Hurwitz



## High resolution 3D characterization of silicon detectors using a **Two Photon Absorption** Transient Current Technique



Marcos Fernández<sup>(1)</sup>, Iván Vila



Michael Moll  
Moritz Wiehe

eman ta zabal zazu



UPV EHU

Raúl Montero



Rogelio Palomo

(1) Also visiting scientist at CERN-SSD

# Outline

What is TCT and TPA?

The setup(s)

Examples of application in Si:

Diodes

HVCMOS



Irradiated and non-irradiated

Summary

**Two Photon Absorption-Transient Current Technique (TPA-TCT)** is a new **technique** to characterize semiconductor detectors using a **point-like laser probe**, so called “**voxel**”. A voxel can be **scanned in the three coordinates**, thus obtaining true **3D spatial resolution**.

### Transient Current Technique

R&D on radiation  
hardness of Si sensors



### Two Photon Absorption

Test of Single Event  
Effects on electronics



Ultrafast molecular  
spectroscopy



+

**TPA-TCT**



This mixture has been possible via **RD50** collaboration

# What is TCT?



Technique to characterize a material via the transport of **excess carriers** generated typically using a **laser** beam.

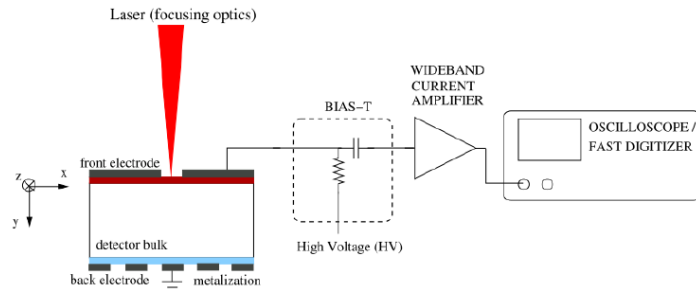
**E-field** profile, **space charge**, charge collection efficiency, trapping.. can be reached.

Induced current  
pulse is **time**  
**resolved**, measured,  
and analyzed.

# What is TCT?

Technique to characterize a material via the transport of **excess carriers** generated typically using a **laser beam**.

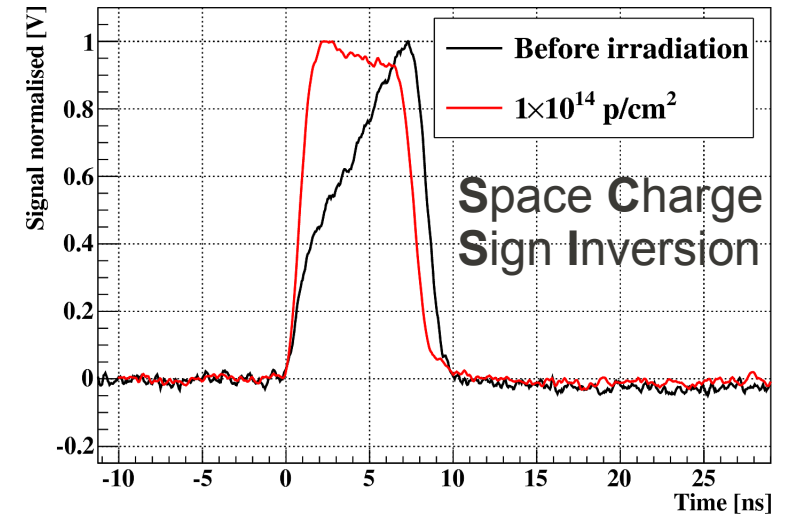
**E-field profile, space charge, charge collection efficiency, trapping..** can be reached.



Induced current pulse is **time resolved**, measured, and analyzed.

From **early 1990s** widely used for measurements of radiation effects in semiconductors.

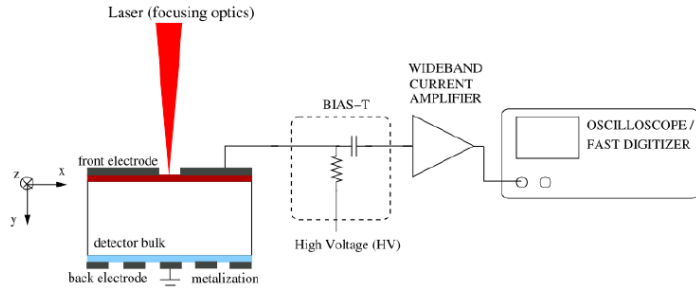
Workhorse for **ROSE** and **RD50** collaborations.



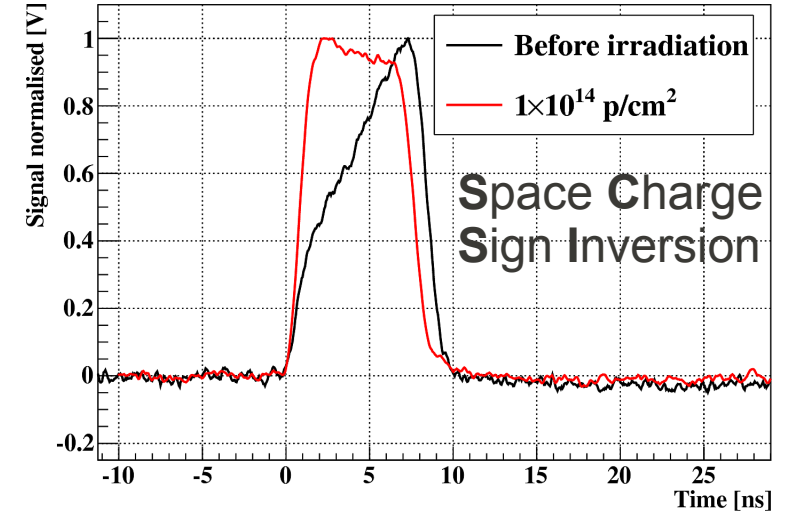
# What is TCT?

Technique to characterize a material via the transport of **excess carriers** generated typically using a **laser beam**.

**E-field profile, space charge, charge collection efficiency, trapping..** can be reached.



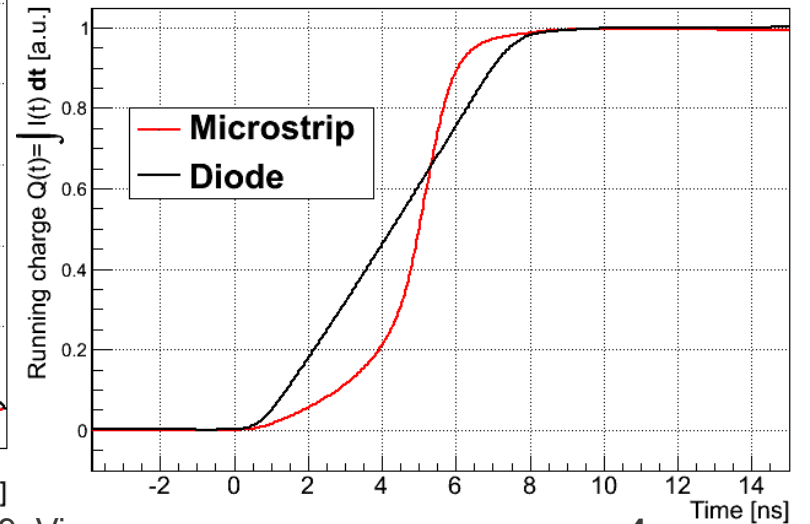
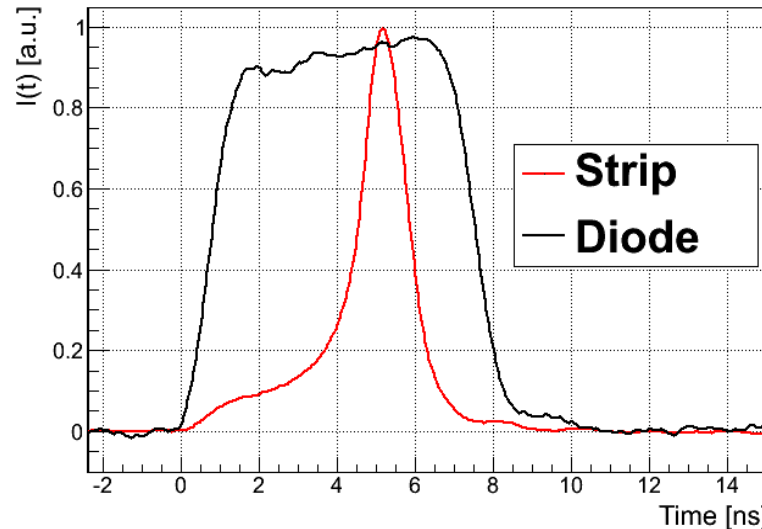
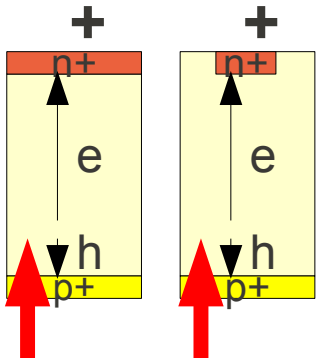
Induced current pulse is **time resolved**, measured, and analyzed.



From **early 1990s** widely used for measurements of radiation effects in semiconductors.

Workhorse for **ROSE** and **RD50** collaborations.

Sensitive to product of **drift** and **weighting field**

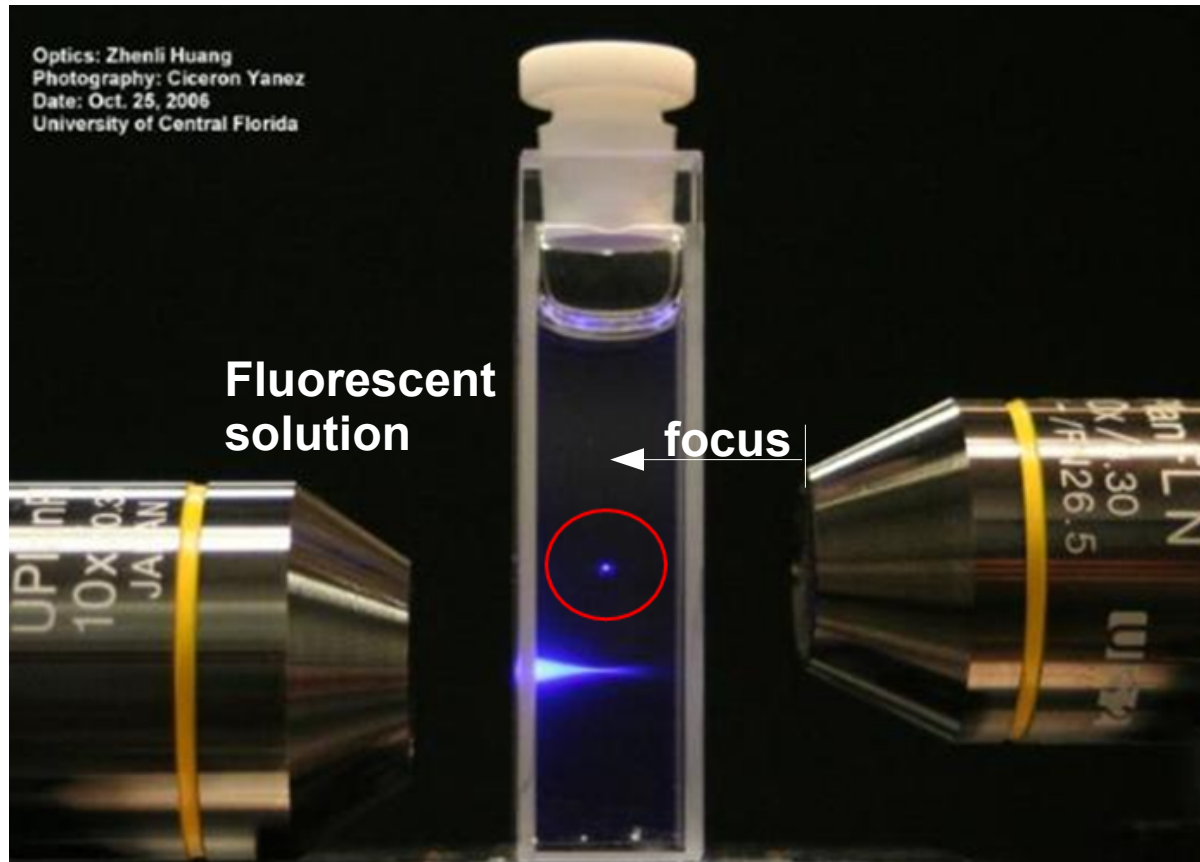


# What is TPA?

- TPA is a **non-linear effect** shown by any material when illuminated with a **high intensity** source (for instance, a laser). For certain wavelengths, light absorption (=signal) only happens **at the focus** of the beam. No photons are absorbed “out of focus”. The more light is focused, the better “point-like” signal generation volume.
- The physical phenomena exploited is the **simultaneous absorption of 2 photons** in the material

## OLD Single Photon Absorption

Continuous energy deposition (no spatial resolution along beam prop. dir.)



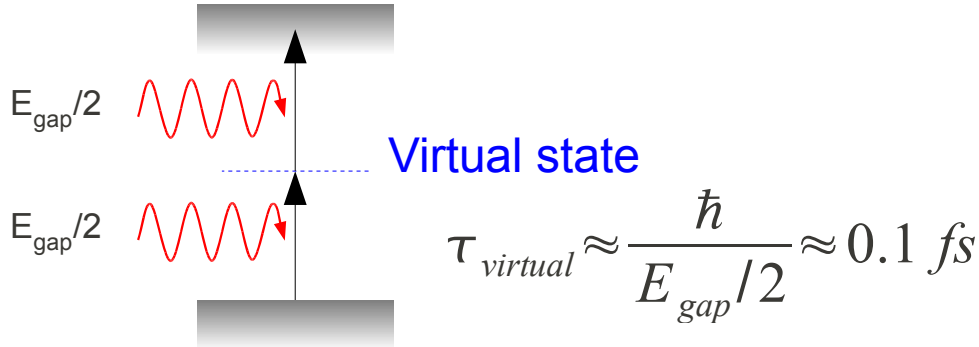
## NEW Two Photon Absorption

Energy confinement

Two photons from one laser !!

## Two Photon Absorption

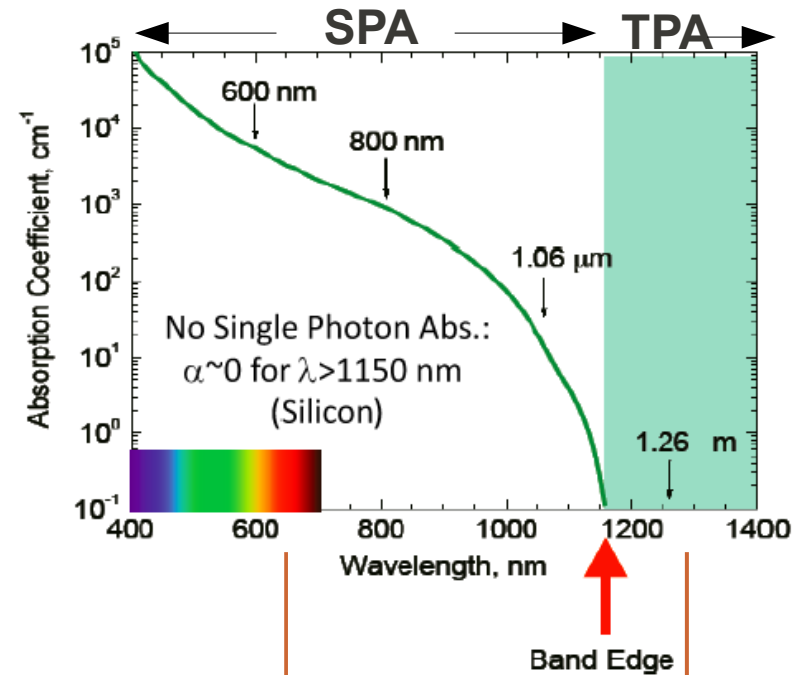
**TPA:** Conventionally, no excitation if  $E_{\text{photon}} < E_{\text{gap}} \sim 1 \text{ eV}$ .  
But, if **TWO** photons arrive in  $\sim 100$  attoseconds:



Two Photons ( $E > E_{\text{gap}}/2$ ) must be:

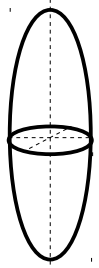
- 1) coincident in time (pulsed mode-locked **fs-lasers**)
- 2) and in space (microfocusing)

## Absorption in Si



## Spot size

$\sigma = 10 \mu\text{m}$



$\sigma = 1 \mu\text{m}$

Measured using knife edge technique.

This is the **3D probe** we scan across the volume of the detector

**x100 objective**

$$\frac{dN(r, z)}{dt} = \alpha \frac{I(r, z)}{\hbar\omega} + \beta_2 \frac{I^2(r, z)}{2\hbar\omega}$$

### SPA

Single Photon Absorption

Negligible  $\lambda \geq 1100 \text{ nm}$

### TPA

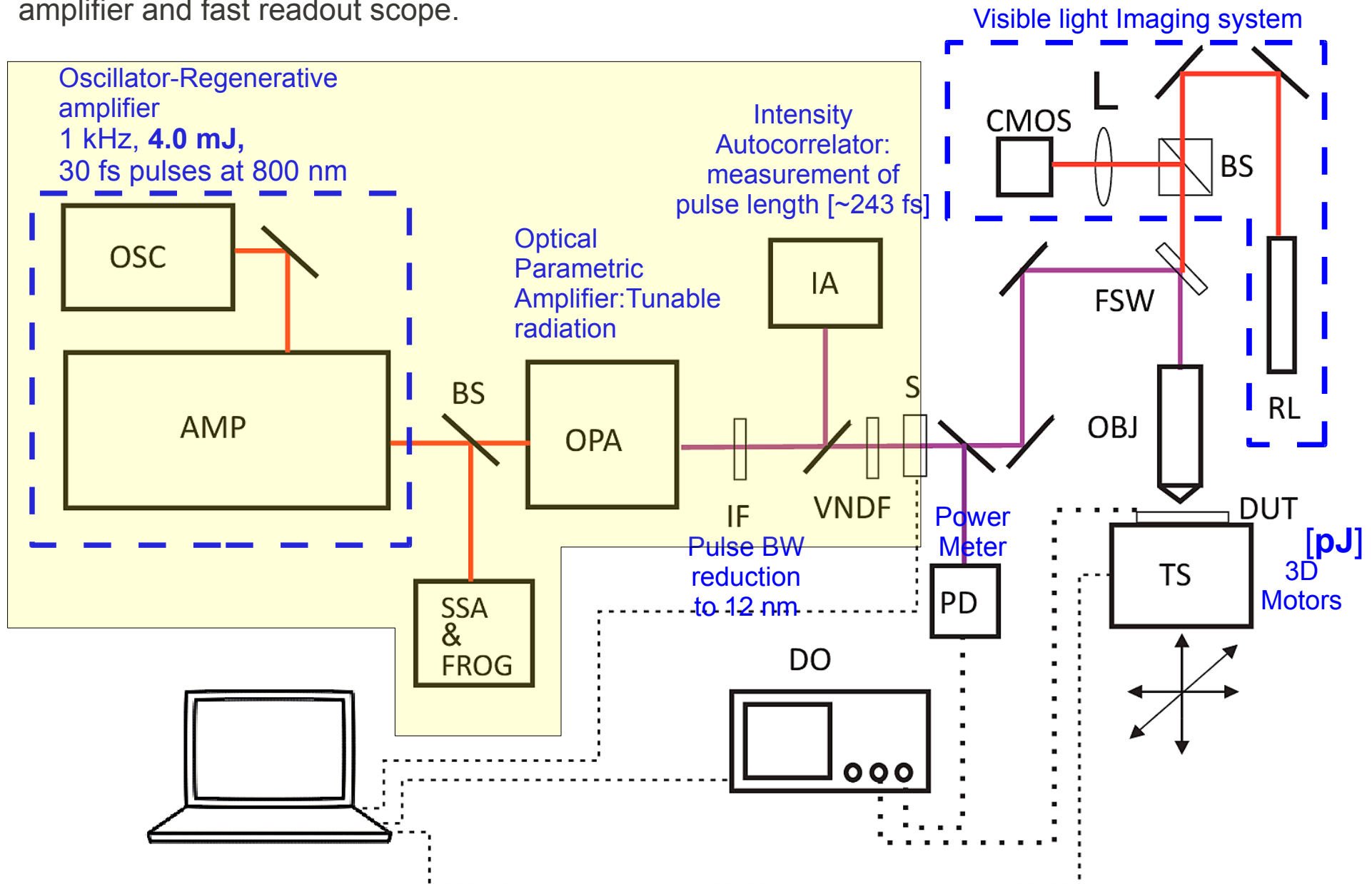
Two Photon Absorption  
Point-like resolution

Boosted by using short pulsed fs lasers, where  $I^2 \gg I$  for the same average power laser.



# TPA-TCT in UPV laser facility

- **Tunable** wavelength, energy for the sensor (after attenuation) [ $\sim$ pJ]. Then typical **TCT readout**: top/side injection, 3D system, current amplifier and fast readout scope.



# New CERN TPA-TCT



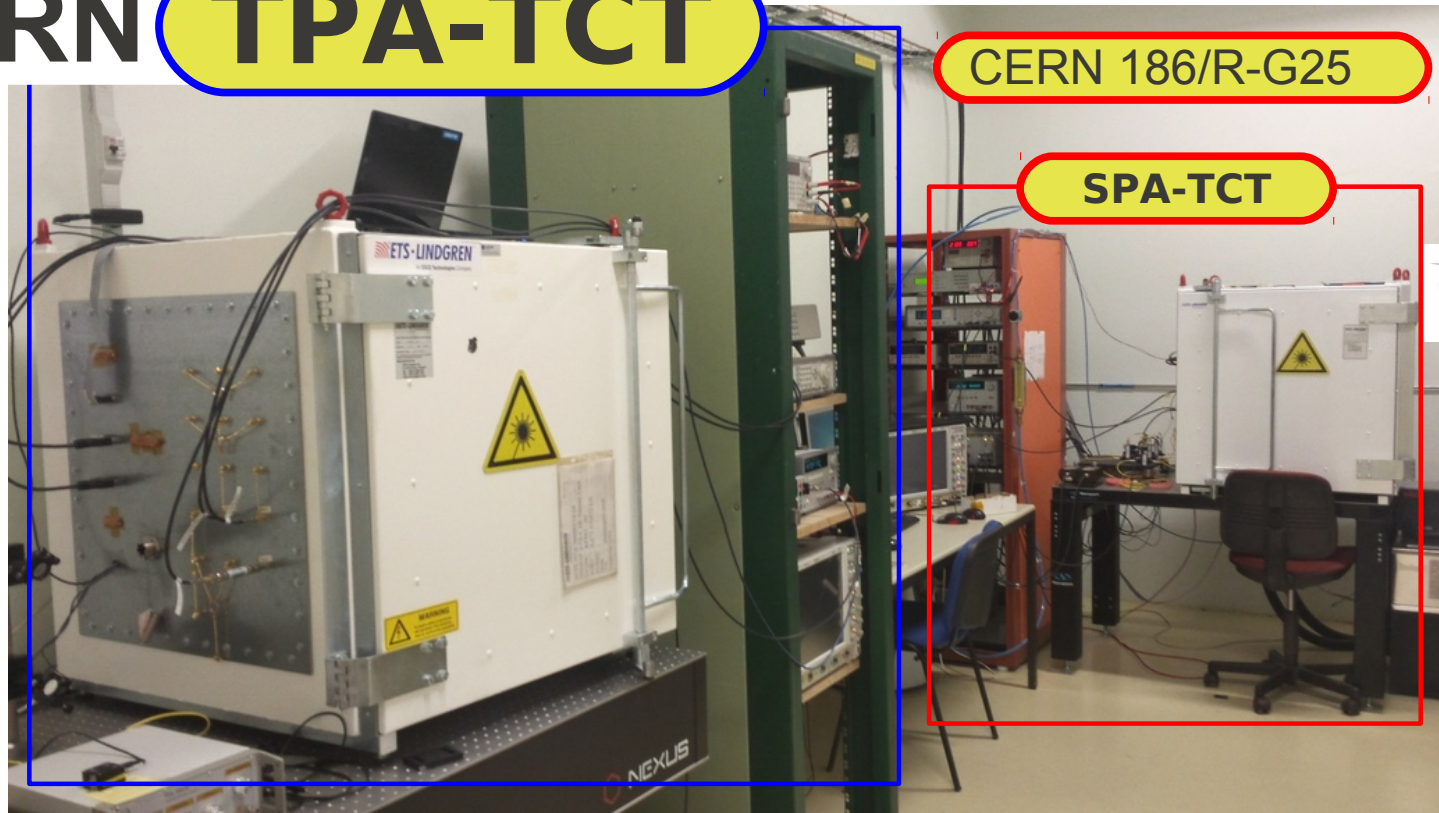
TPA-TCT presented at 2016 CERN EP-**Knowledge Transfer** innovation day  
→ It was **selected for funding**: Proposed Project on non-destructive QC of semiconductors

## Status:

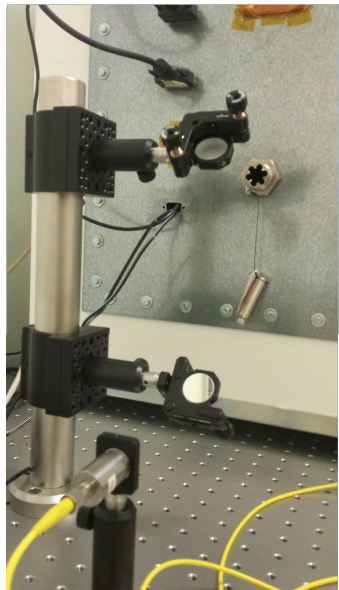
- The core of this system is a **custom fibered laser** working at  $1.5 \mu\text{m}$ . It is being delivered this March 2019 after  $\sim 1$  year development by Laser company.
- Rest of the system is based in our already existing SPA-TCT at CERN-SSD lab (known as TCT+).
- Some key differences / improvements:
  - **Improved positioning system** (6 degrees of freedom hexapod system)
  - **Optimal cooling and sample support**. It can be rotated for edge-side injection
- German Doctoral Student (M.W., Gentner Program) responsible for construction of this demonstrator.
- Upon completion of the project access to this laser via RD50 collaboration.

# New CERN TPA-TCT

Status as of  
Feb 15th, 2019



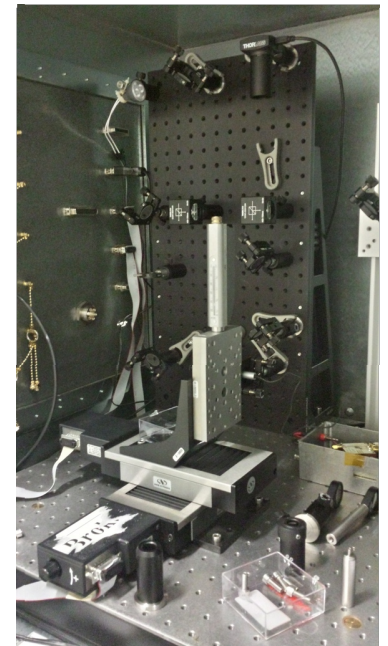
Light injection  
system



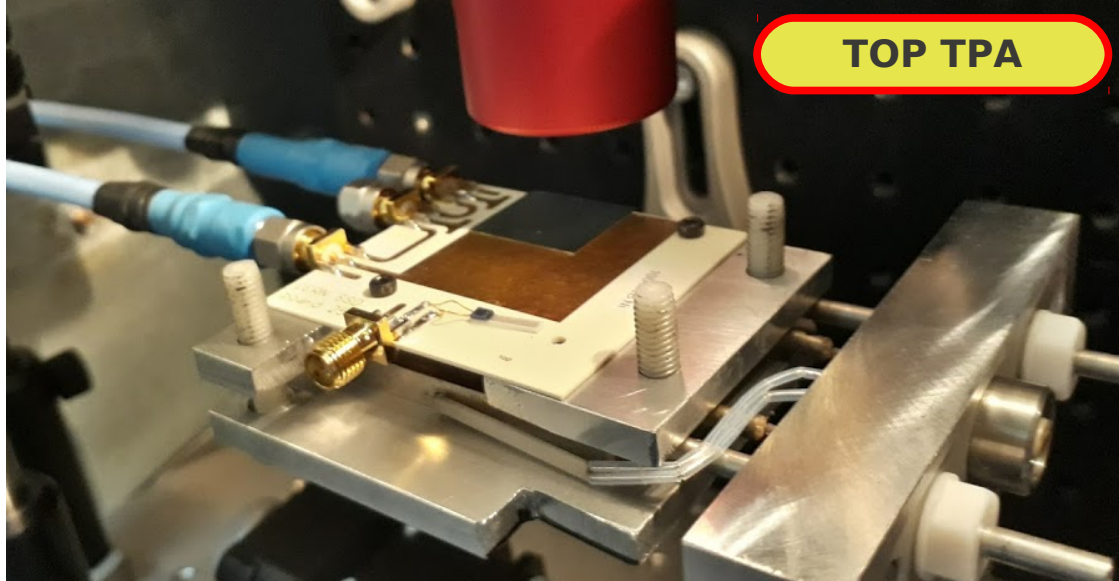
6 degrees of  
freedom, high  
load capacity  
motion system



Laser routing,  
microfocusing,  
visible laser  
imaging  
system,  
cooling, stage  
system inside  
Faraday Cage.

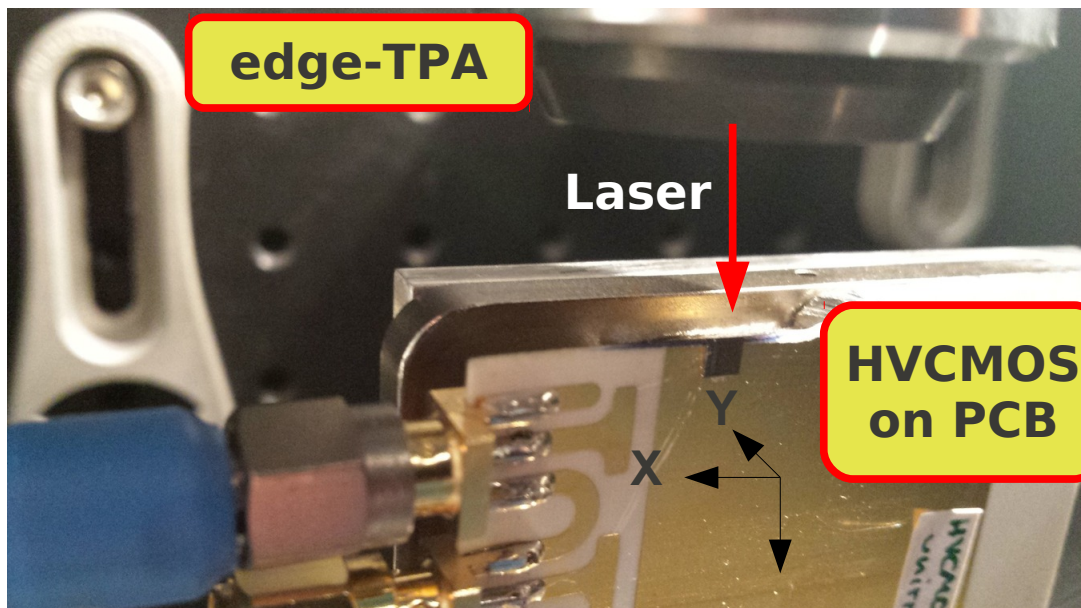


# TPA-TCT light injection



In both illumination configurations, the focus is fixed and the detector is moved

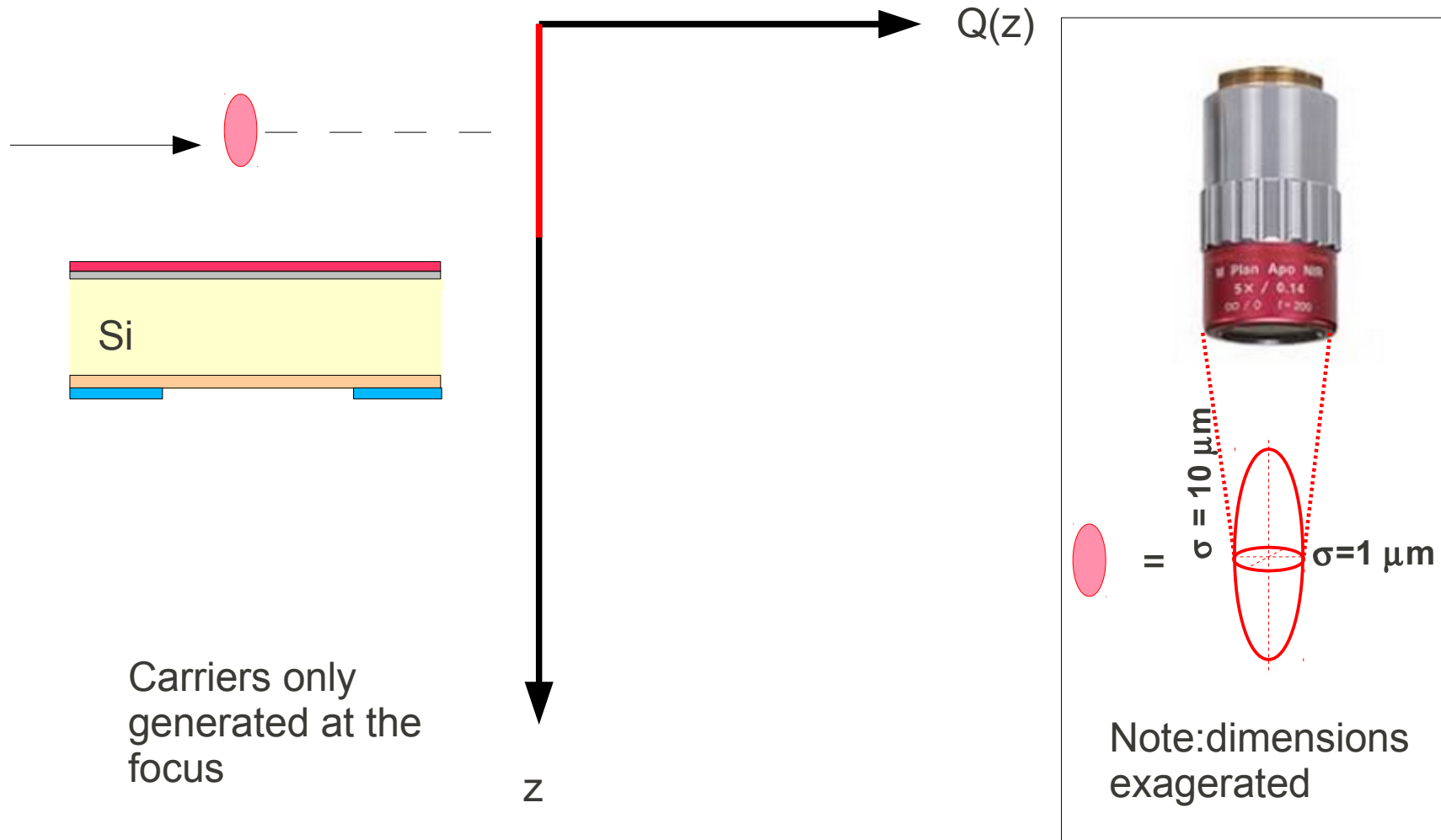
**Top-TPA:** detector is scanned in vertical  $\rightarrow$  depth scan (10  $\mu\text{m}$  resol.)



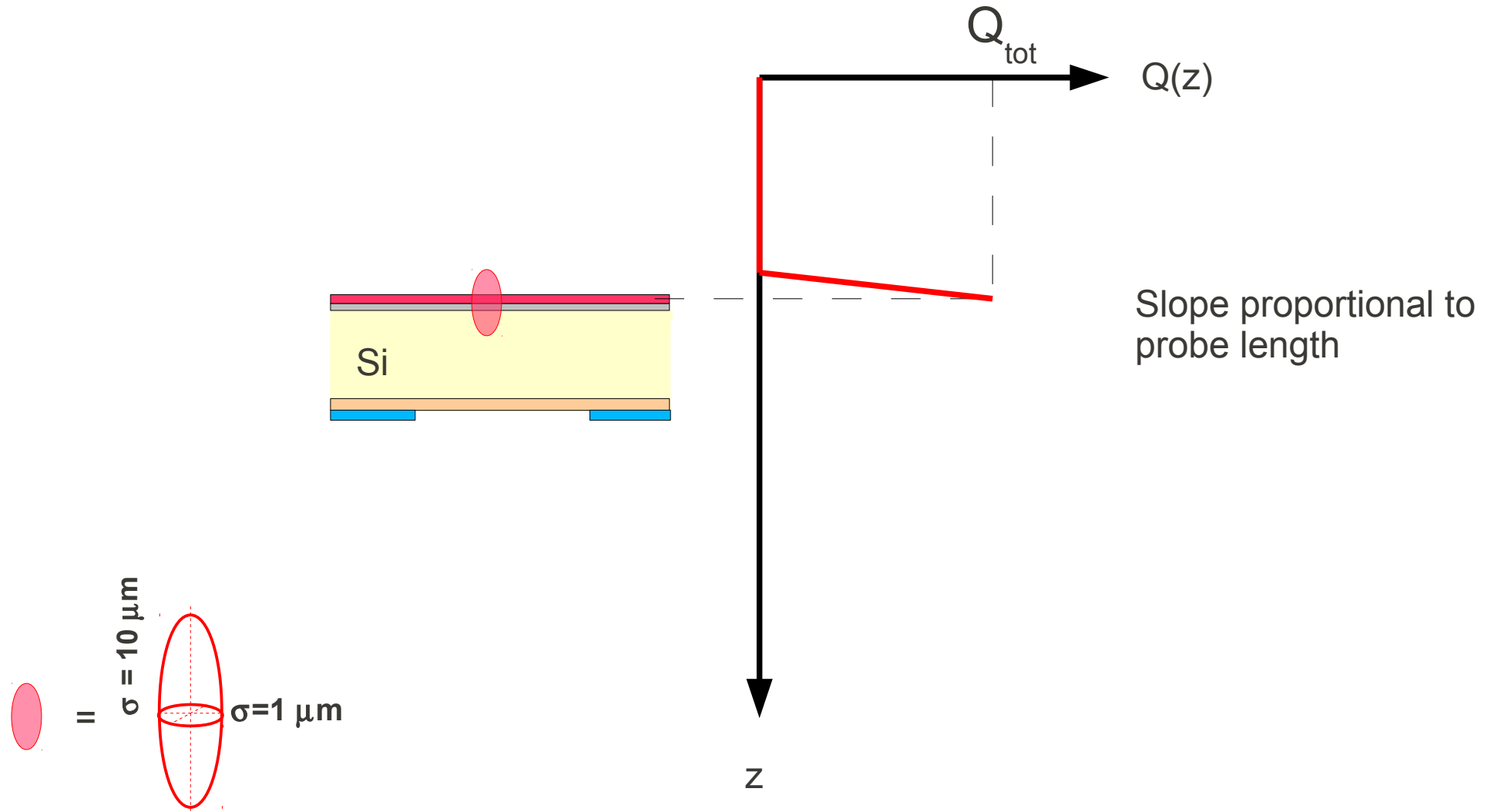
Rotating the detector by 90 degrees, we can, for instance, map the depletion region with 1  $\mu\text{m}$  resolution: **edge-TPA**

# TPA top injection

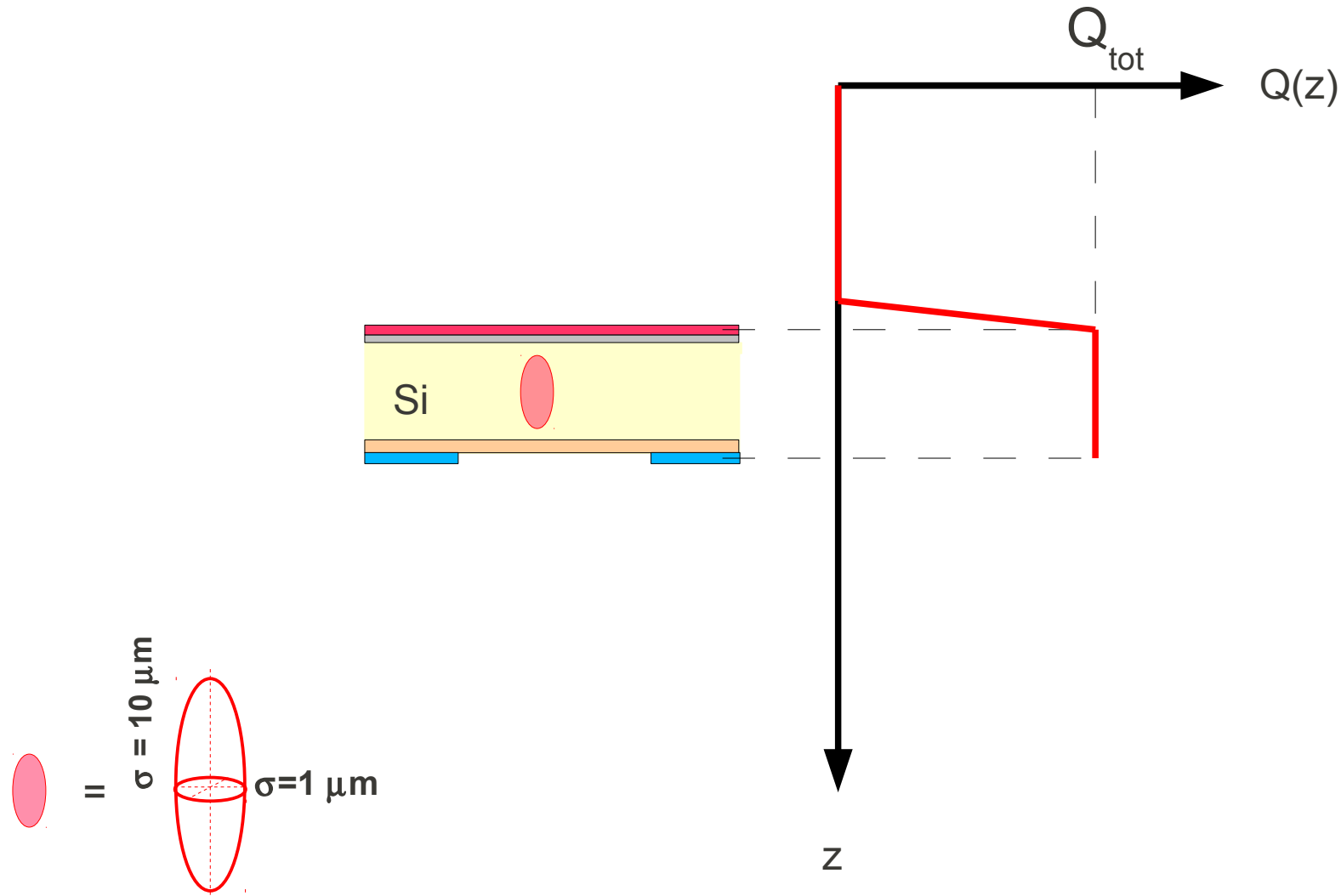
- **TPA:** For simplicity, here moving the focus wrt the detector. In real measurements, the detector moves.
- Assuming fully depleted detector
- Laser probe  $\rightarrow$  Excess charge carriers  $\rightarrow$  Induced Current  $\rightarrow$  Showing integrated current vs position



# TPA top injection



# TPA top injection

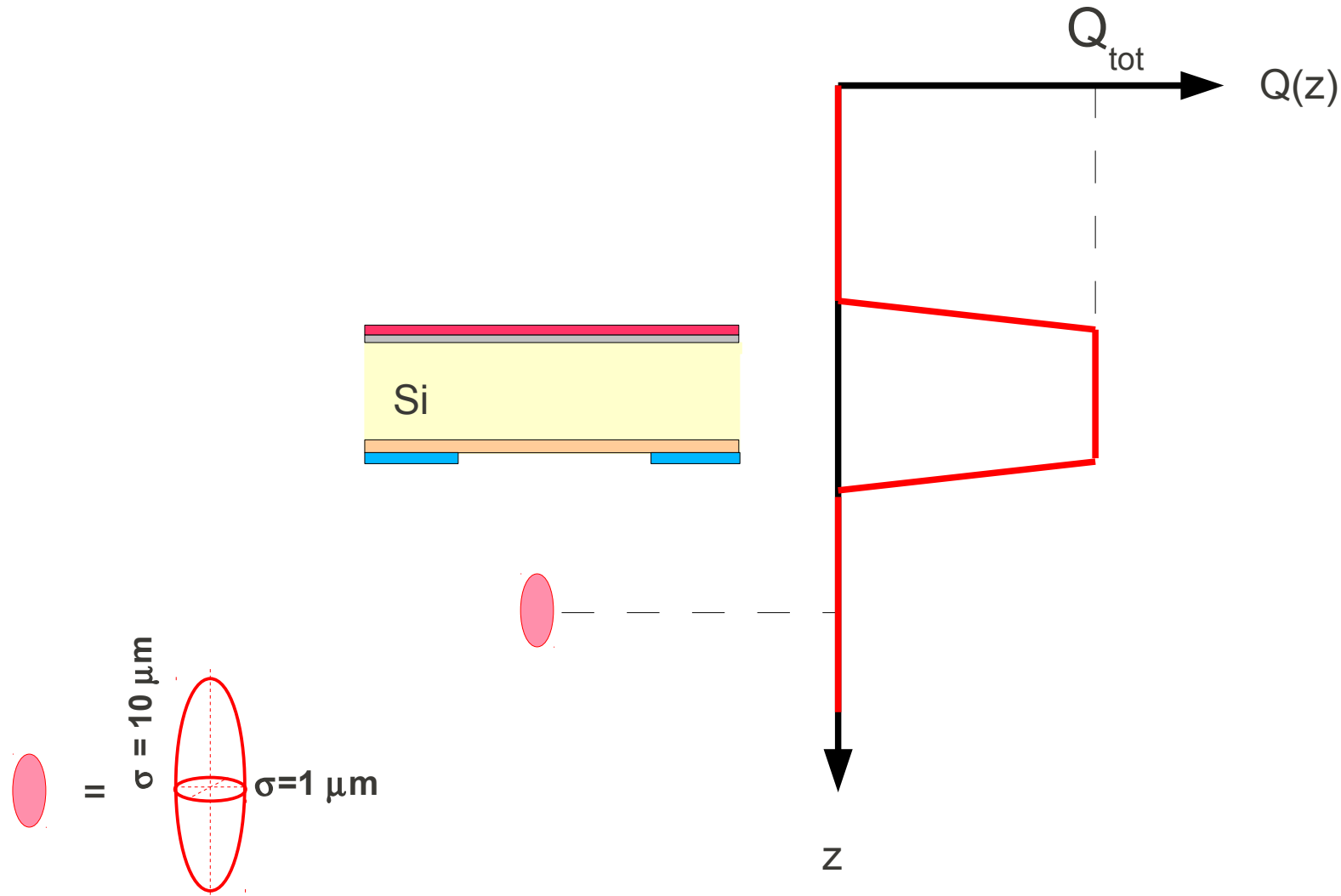


# TPA top injection

**No light absorbed before/after the focus**

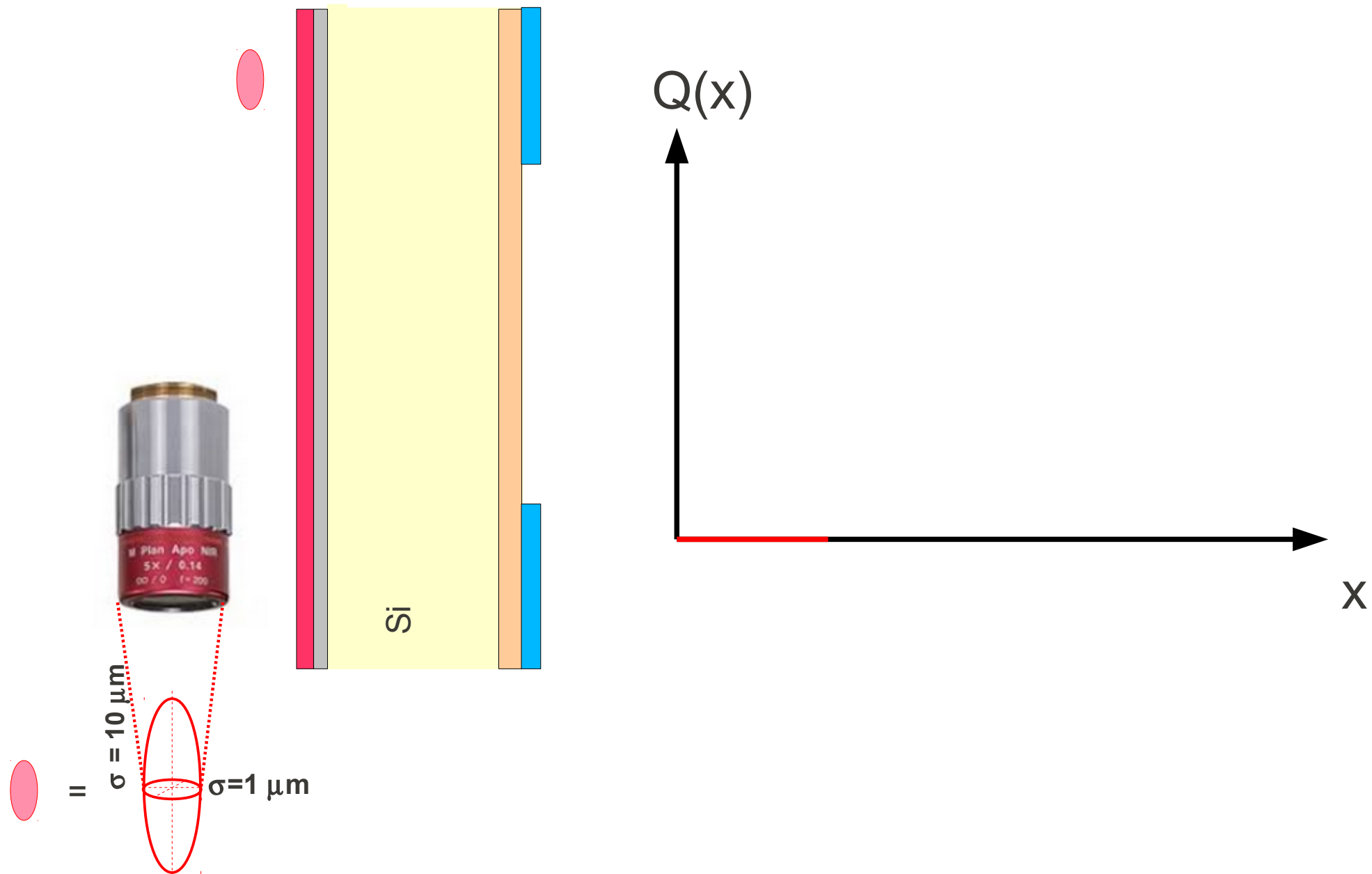
A depth scan of the detector is done from the top!

Measurements not affected by side effects, guard rings...

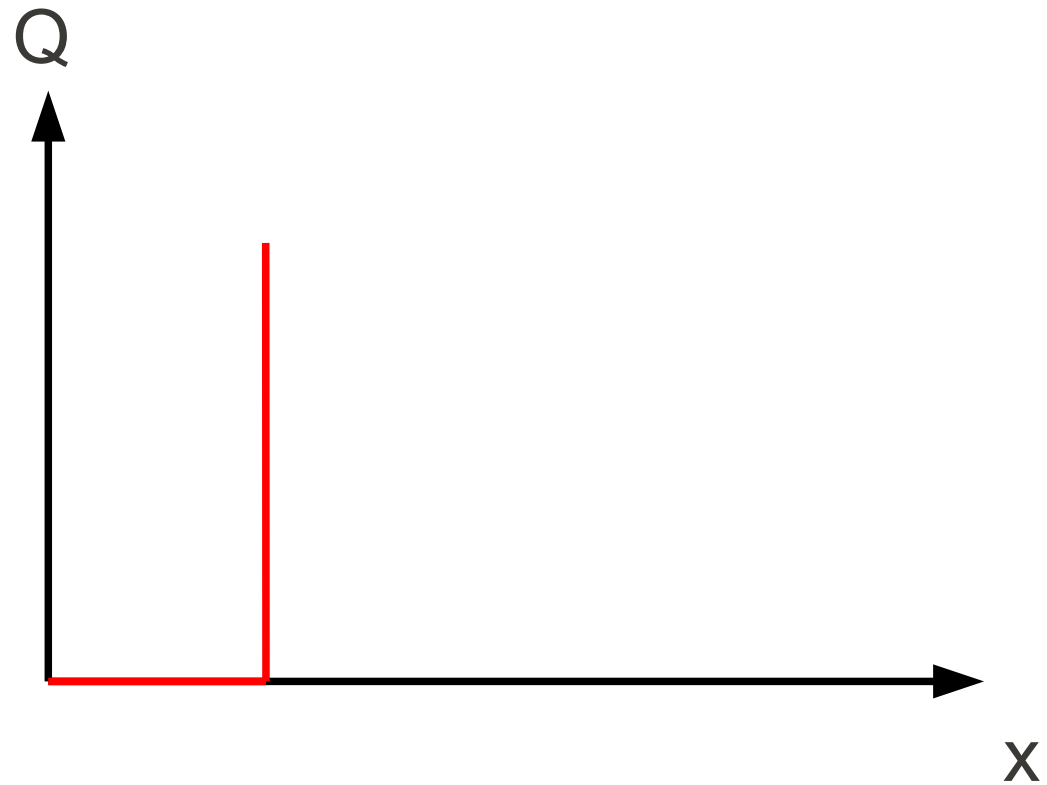
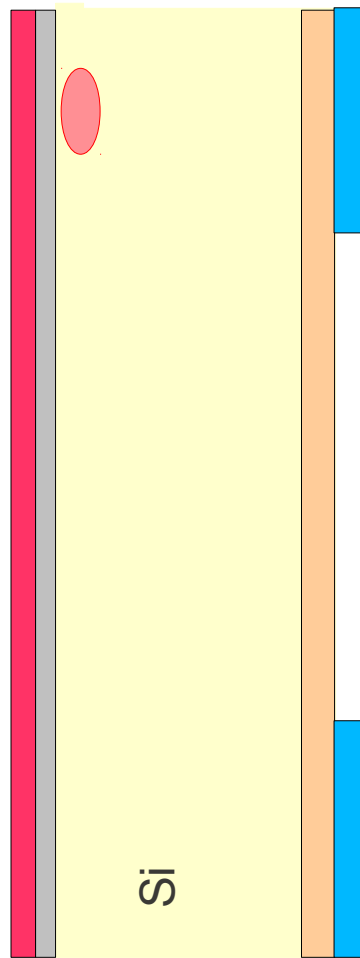




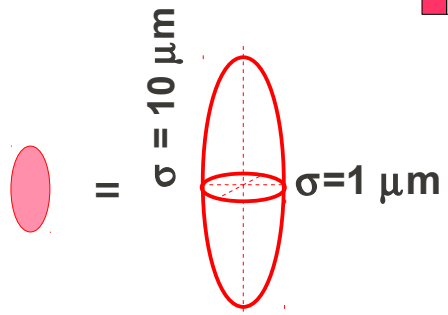
# TPA edge injection



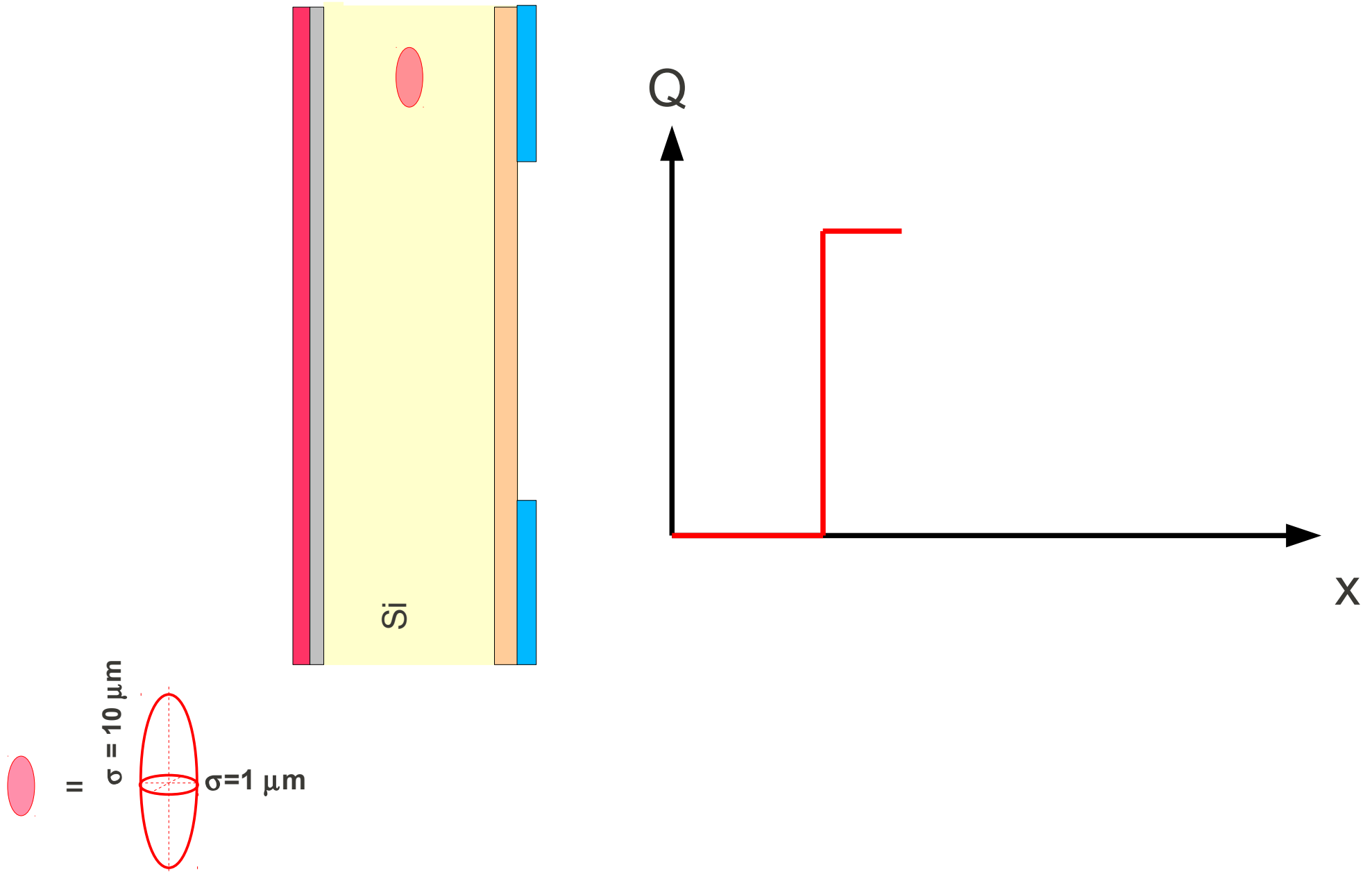
# TPA edge injection



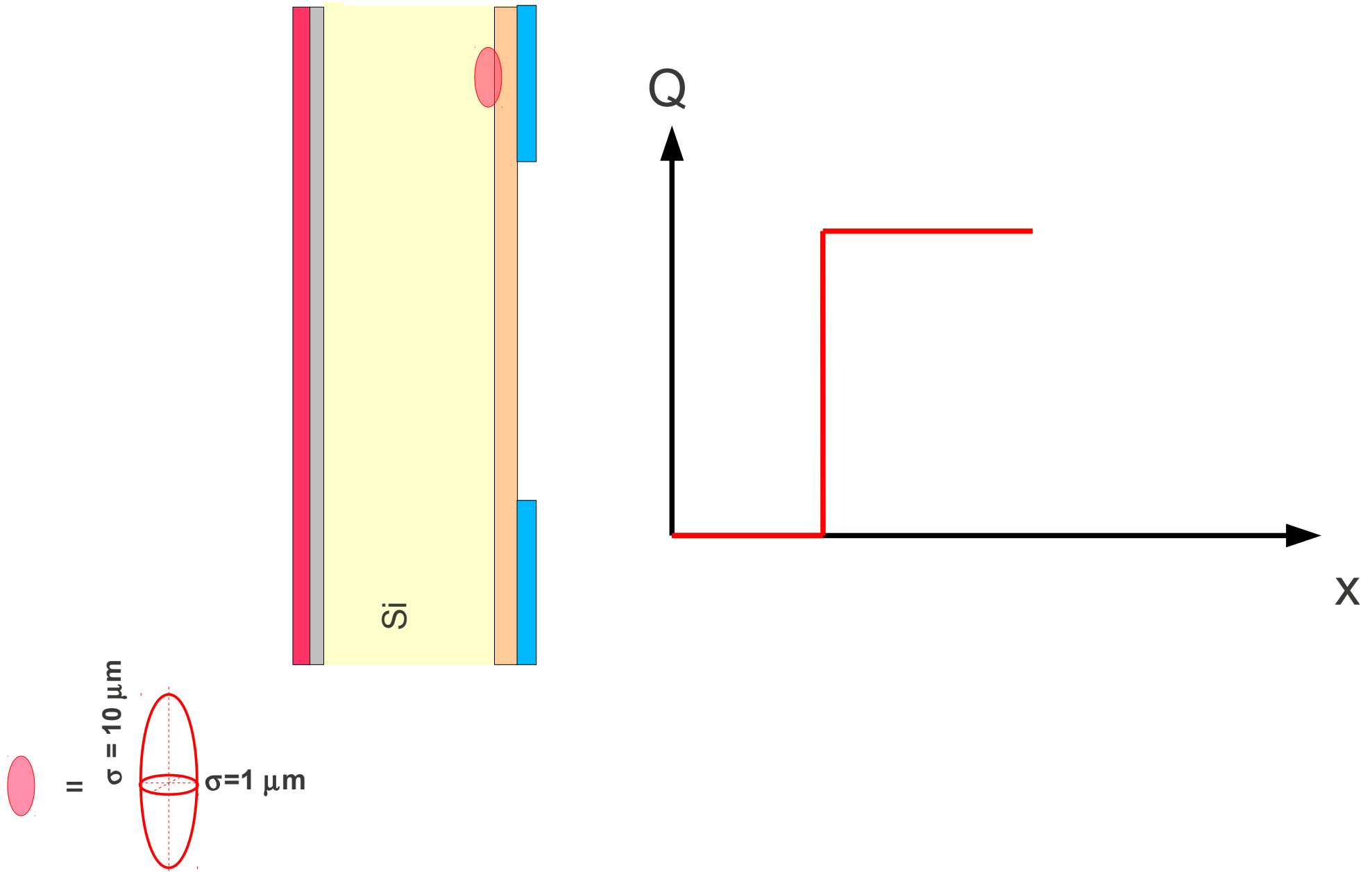
Very sharp transition because of 1  $\mu\text{m}$  beam waist



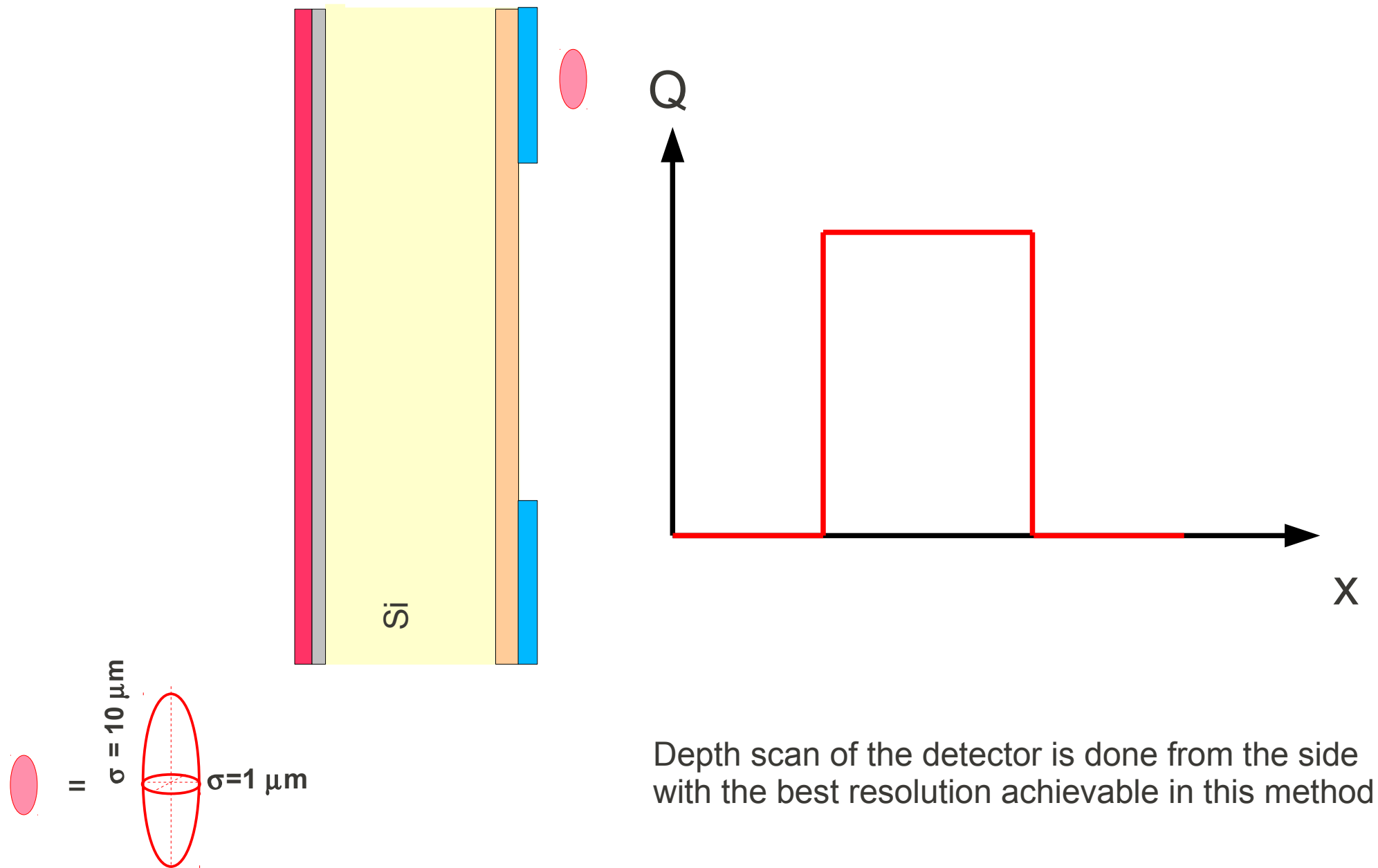
# TPA edge injection



# TPA edge injection

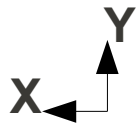


# TPA-TCT edge injection

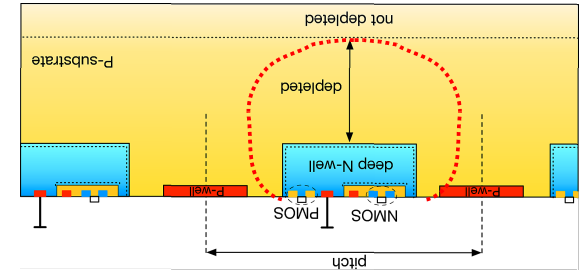


Depth scan of the detector is done from the side with the best resolution achievable in this method

# Edge-TPA: unirradiated HVCMOS

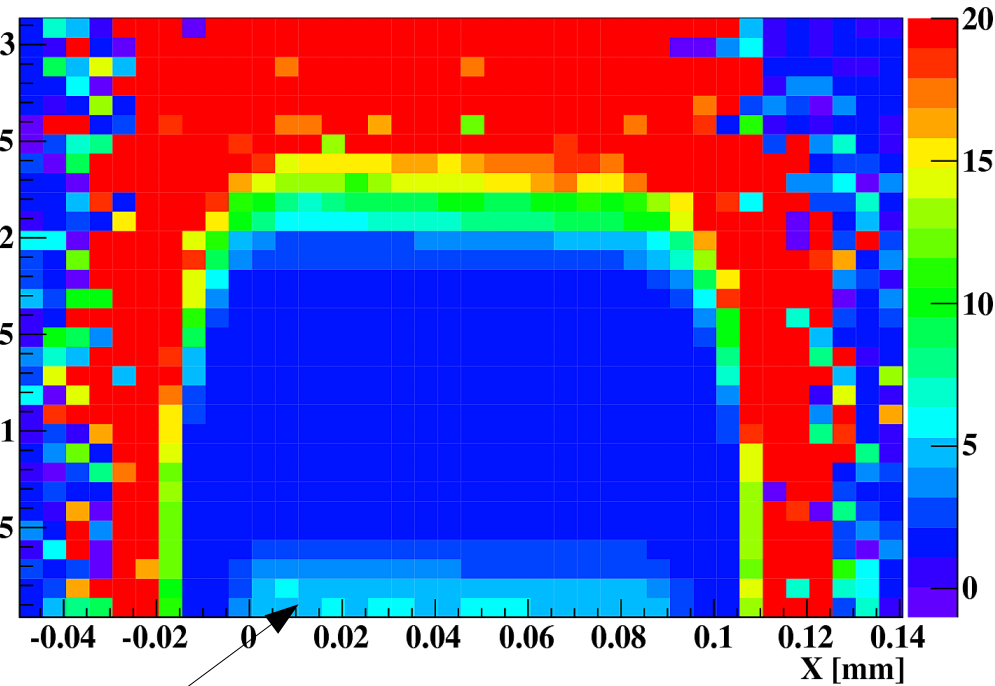
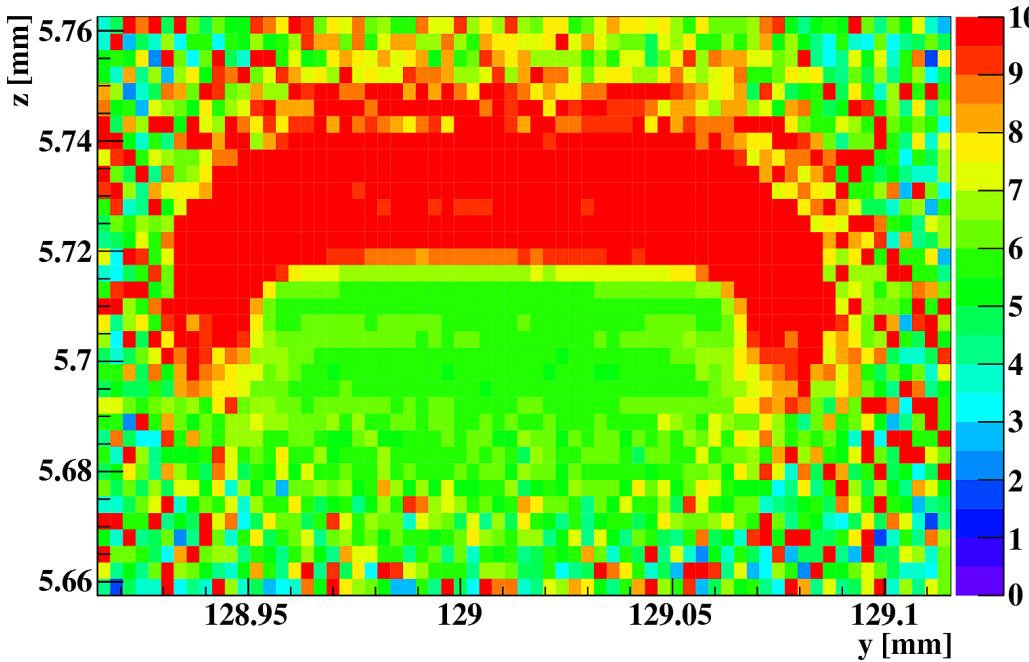


- HVCMOS is a partially depleted sensor, built on commercial CMOS technology. Substrate of **very low resistivity** → very narrow depletion width (~10 μm at 100 V). Challenging detector for SPA-TCT



## Standard SPA-TCT

## TPA-TCT



Spatially continuous laser source, ps pulses

$\lambda=1064$  nm,  $\sigma=10$  μm, 200 ps-pulsed laser

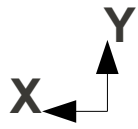
Point-like laser source, fs pulses

$\lambda=1300$  nm,  $\sigma=1$  μm, 200 fs-pulsed laser

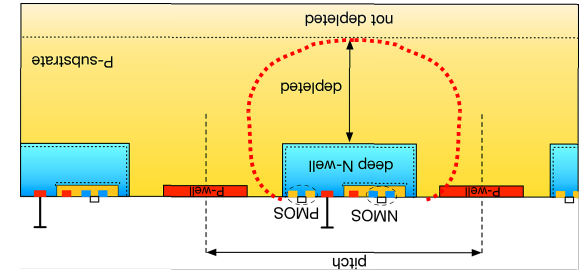
- Showing 2D map of the **collection time** of charge carriers (time lapse till 98% charge is collected). Very sensitive to creation point of carriers.

## Substructures resolved with TPA-TCT

# Edge-TPA: unirradiated HVCMOS

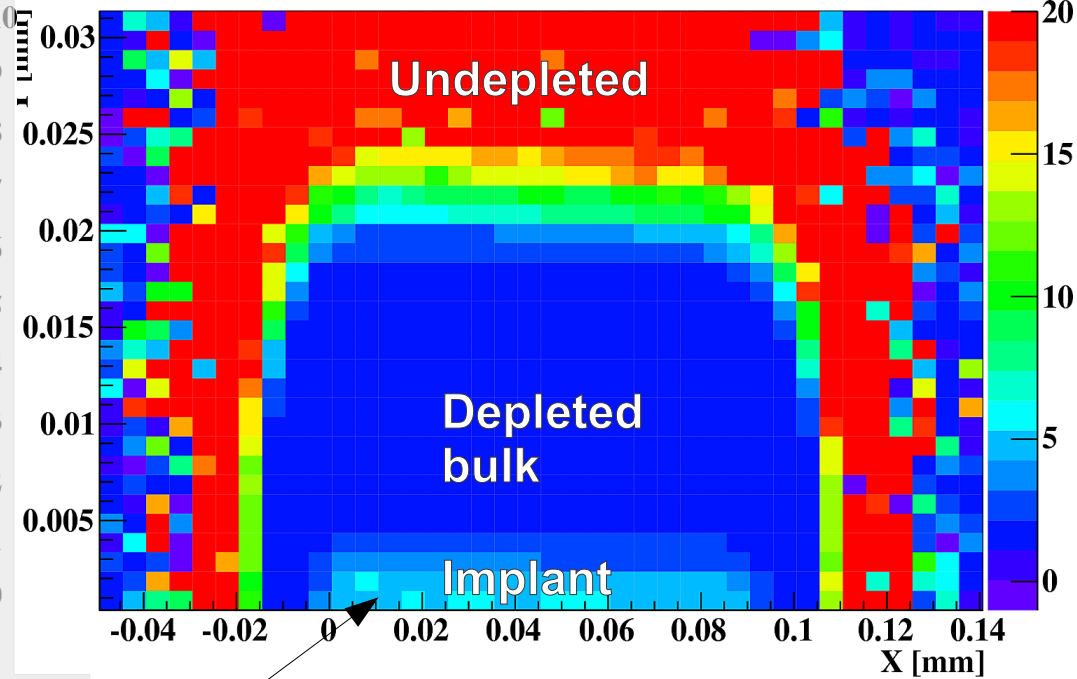
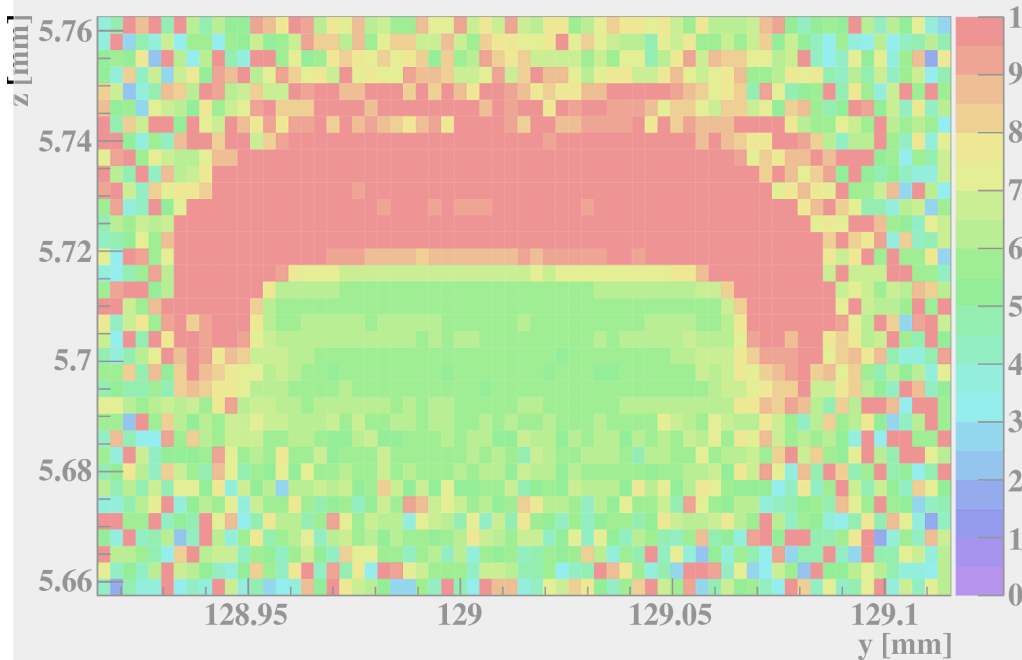


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## Standard SPA-TCT

## TPA-TCT



Spatially continuous laser source, ps pulses

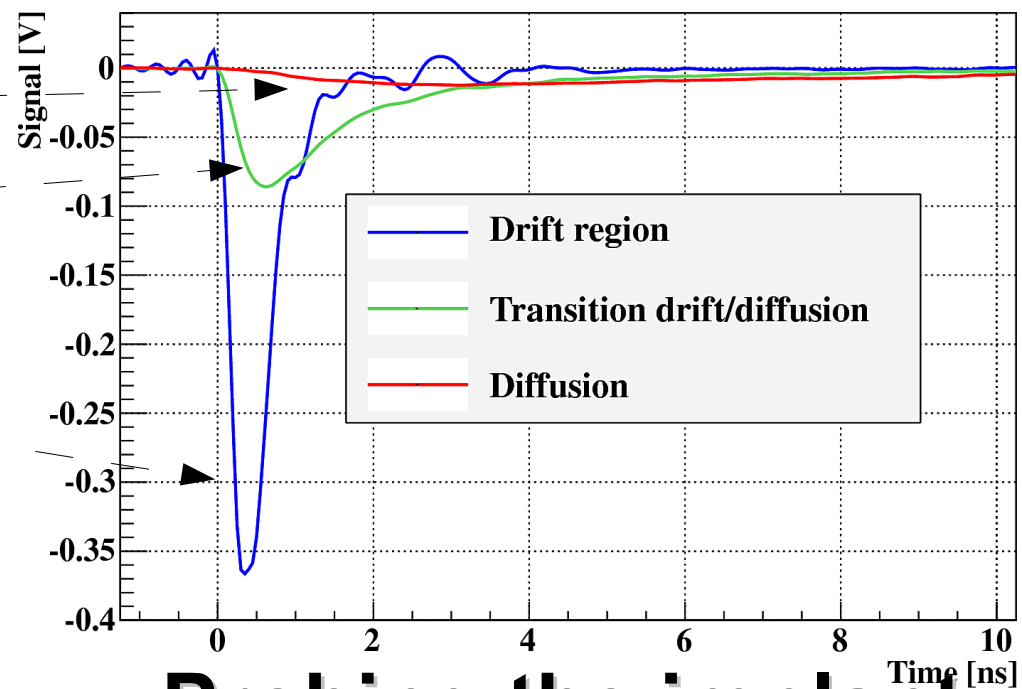
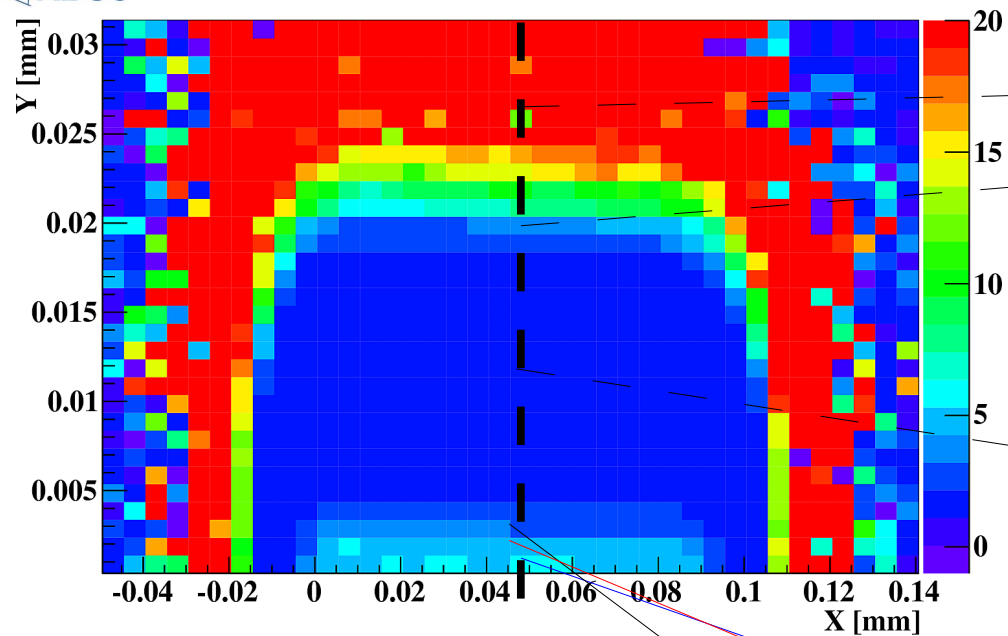
$\lambda=1064 \text{ nm}$ ,  $\sigma=10 \mu\text{m}$ , 200 ps-pulsed laser

Point-like laser source, fs pulses

$\lambda=1300 \text{ nm}$ ,  $\sigma=1 \mu\text{m}$ , 200 fs-pulsed laser

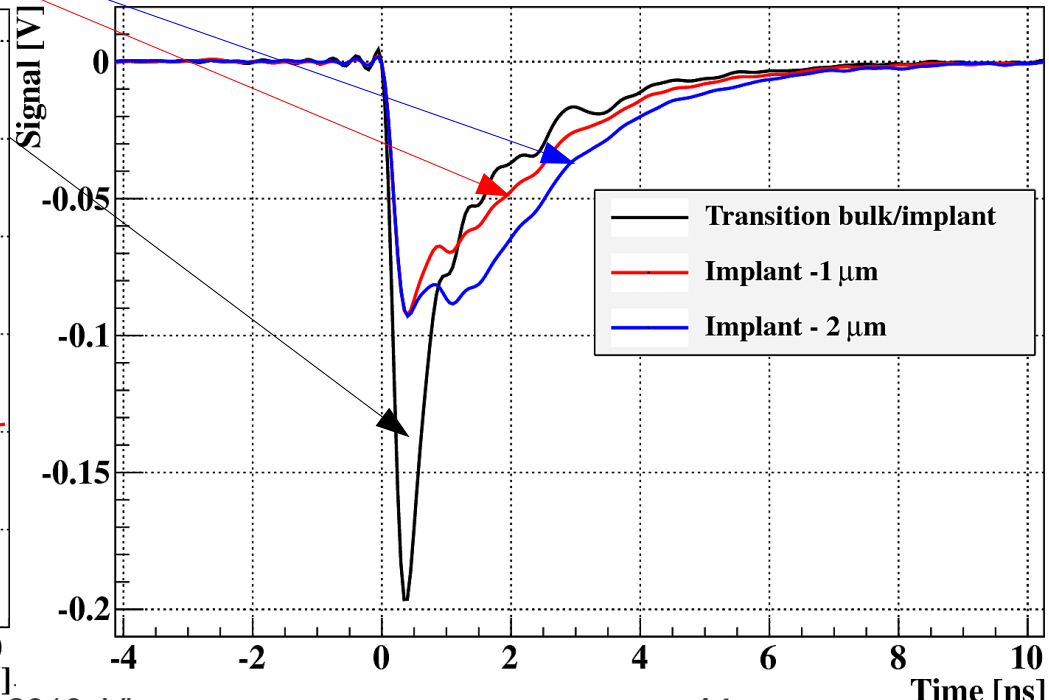
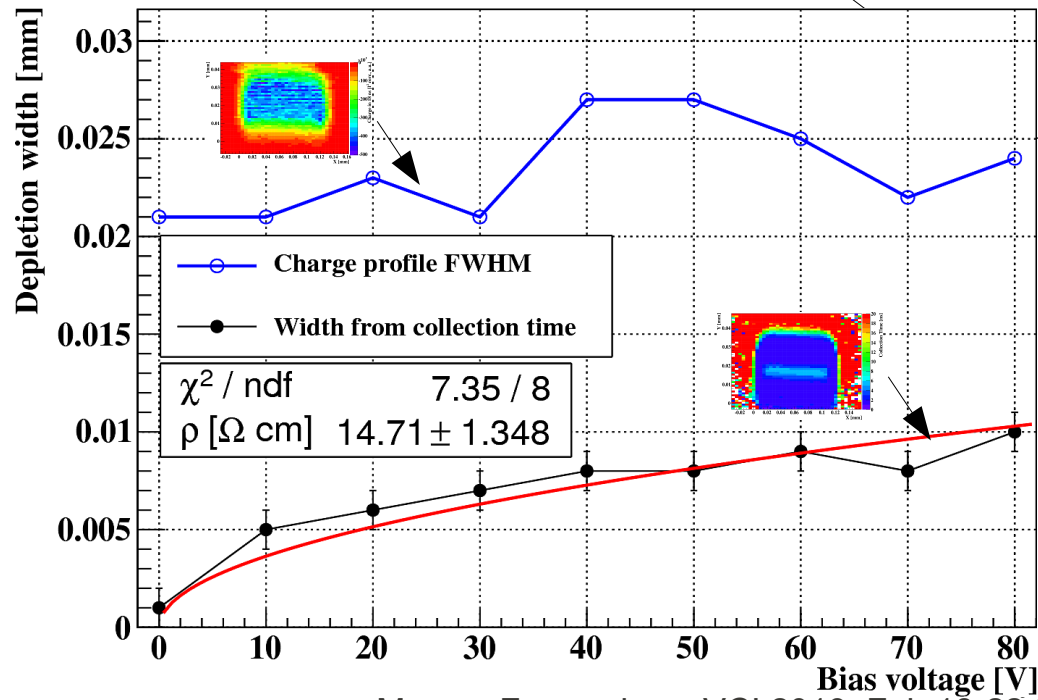
- Showing 2D map of the **collection time** of charge carriers (time lapse till 98% charge is collected)

## Substructures resolved with TPA-TCT



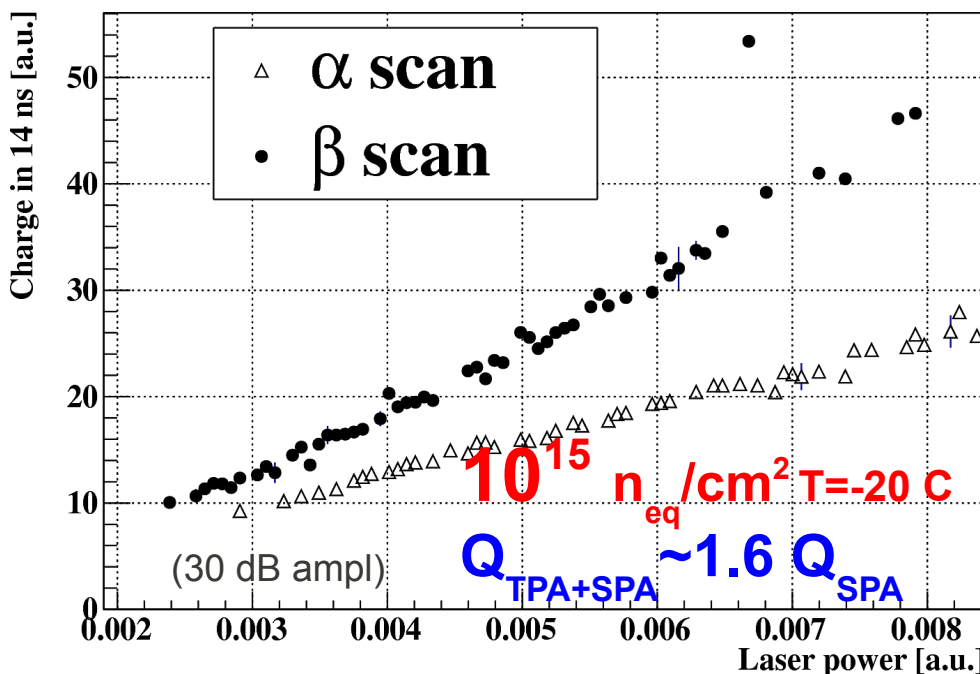
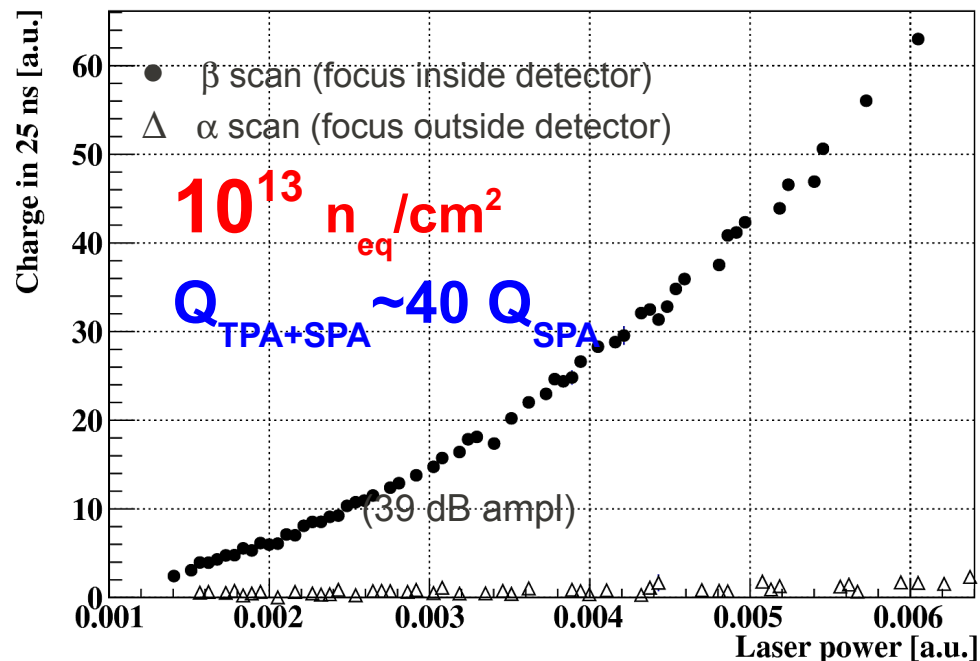
## Doping concentration

## Probing the implant





-Sum\$((volt-BlineMean)\*( time>=23 && time<48 )):LPower



- In Silicon, ionizing radiation creates **Deep Energy Levels** within the bandgap, that **increase Single Photon Linear Absorption**. More free charge carriers but less photons for TPA

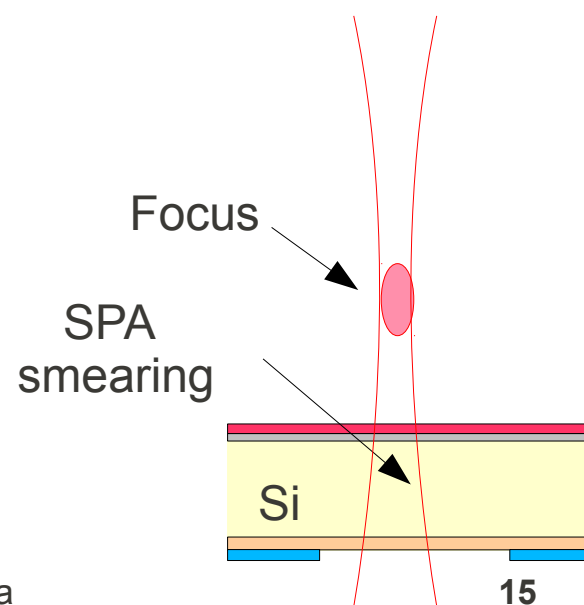
- Linear absorption **smears** spatial resolution along the beam propagation direction.

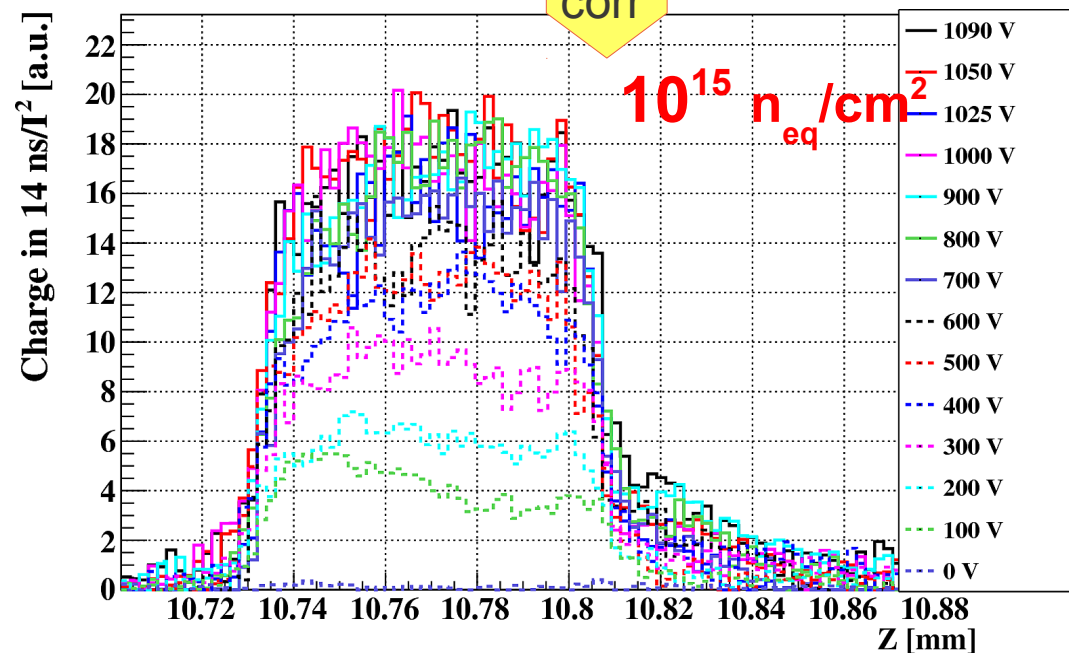
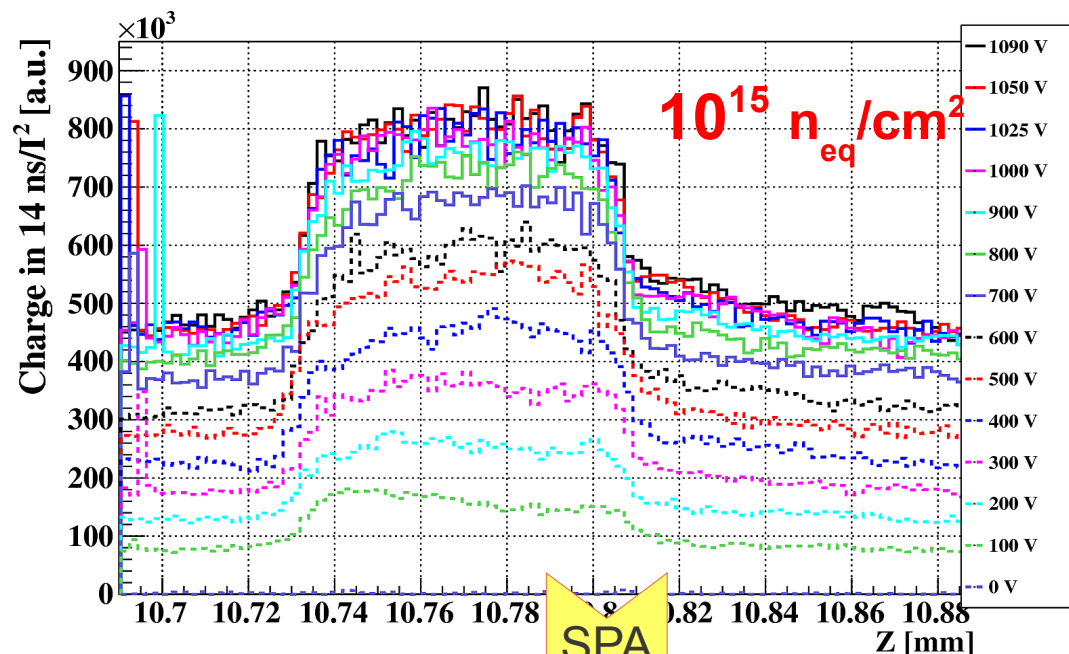
- Clear signature:

- 1) carriers collected even when focus is outside the detector.

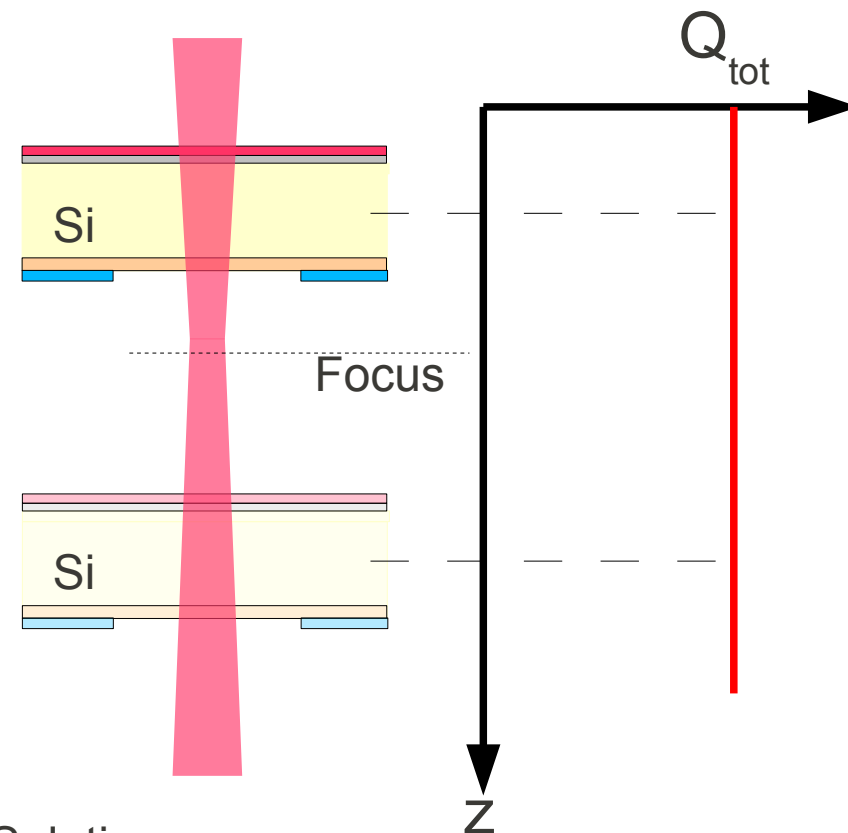
- 2) Linear increase of signal as power increases

$$\frac{dN(r, z)}{dt} = \alpha \frac{I(r, z)}{\hbar\omega} + \beta_2 \frac{I^2(r, z)}{2\hbar\omega}$$





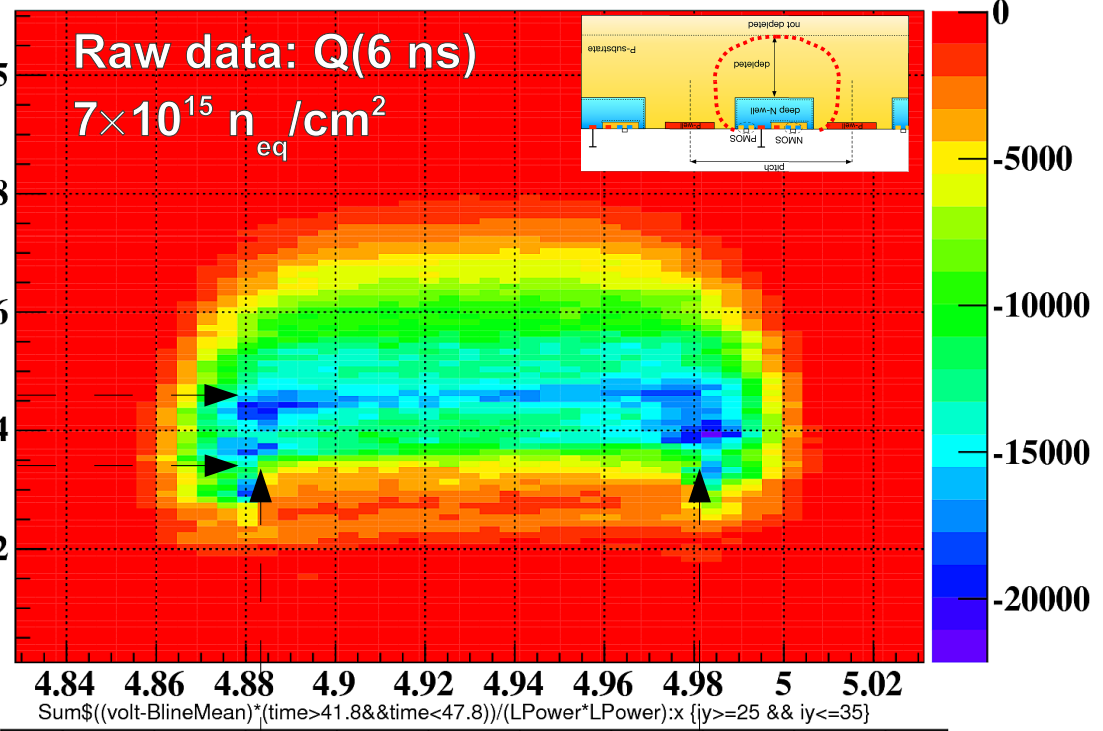
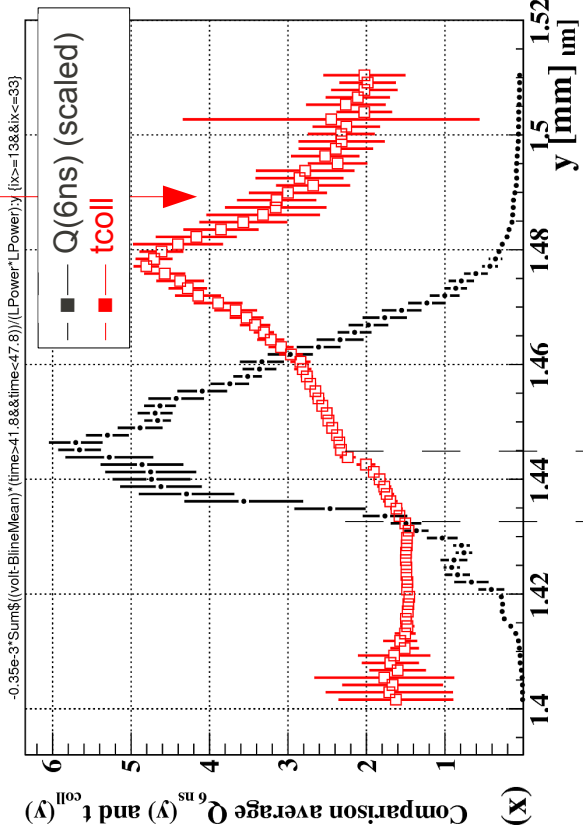
In TPA, the increase of linear absorption can be subtracted exploiting the fact that it is Z-invariant



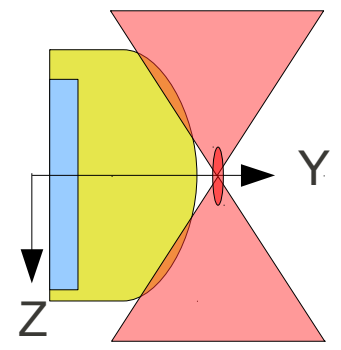
Solution:

- 1) place focus outside the detector  $\rightarrow$  measure  $\alpha$ -induced signal:  $I(t, \alpha)$
- 2) Subtract  $I(t, \alpha)$  (waveform-by-waveform) to measurements with focus inside the detector

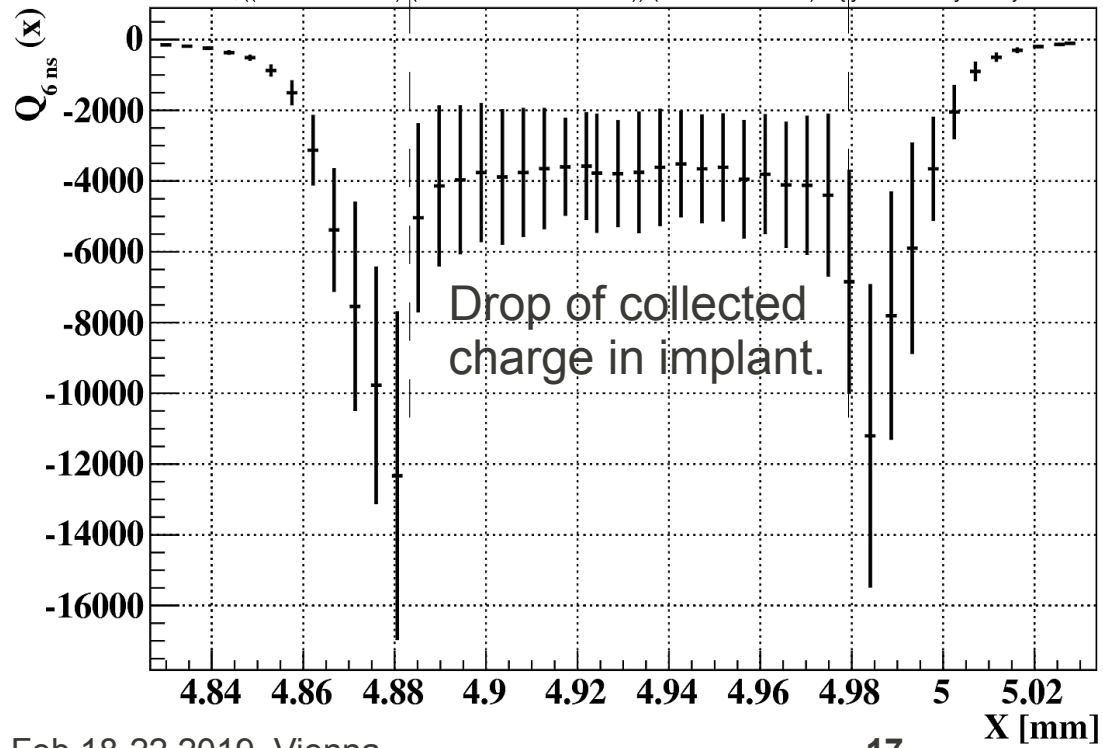
# TPA in irradiated HVCMOS



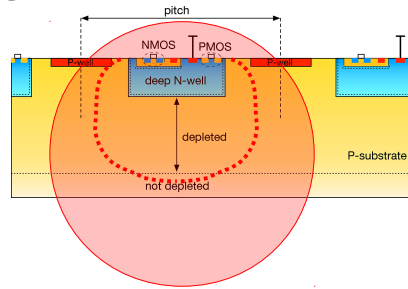
- No sharp decrease of  $Q(y)$  or collection time, due to  $\alpha$ -contamination.



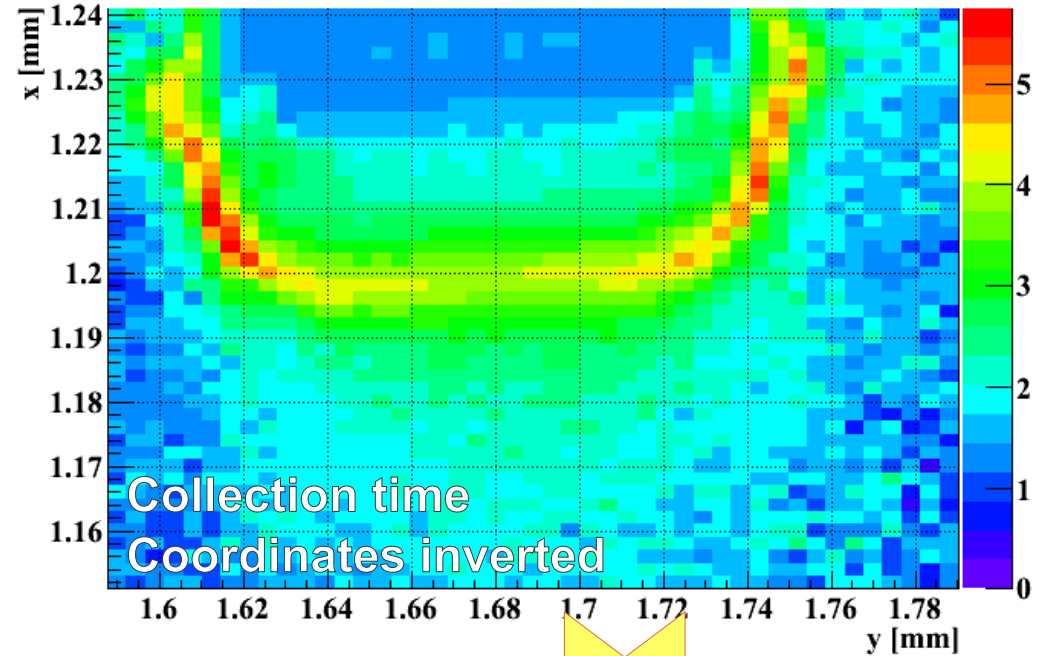
- edge-TPA can overestimate depletion width in irradiated devices if this effect is not accounted for.



Due to strong beam focusing/divergence, the  $\alpha$ -contribution can overflow the depleted region:

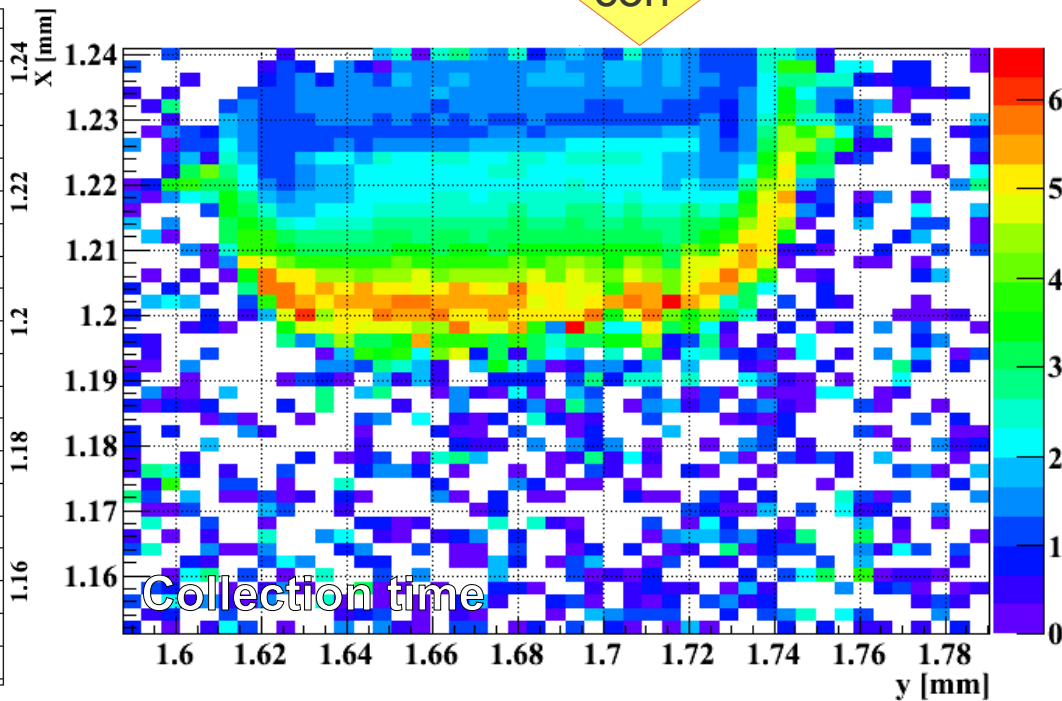
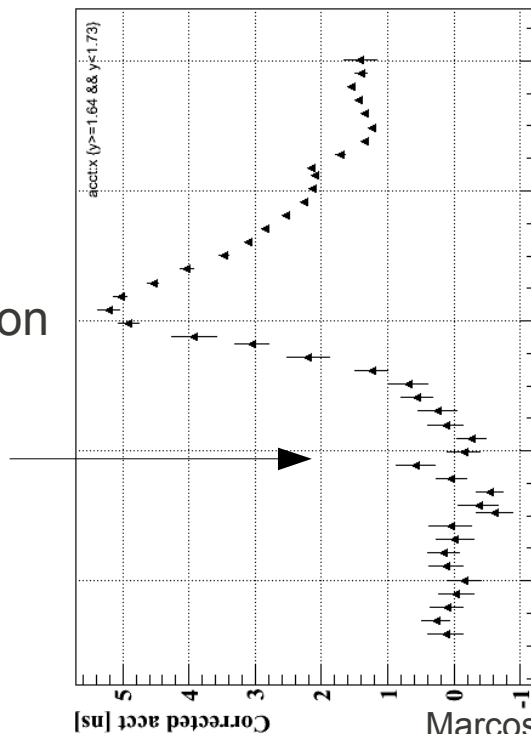


Diode-like correction method not applicable  $\Rightarrow$  Developed a different one (see backup) based on 2 consecutive scans of the same region at different powers



SPA  
corr

Because of  $\alpha$  correction, collection time drops to zero behind depleted region



# Summary

- TPA-TCT is a new material characterization technique providing 1  $\mu\text{m}$  resolution in 2D and up to 10  $\mu\text{m}$  along the beam direction → Optimal configuration for edge injection
- Physical phenomena exploited is the simultaneous absorption of two photons at the focus of a fs-laser beam.
- Demonstrated performance on a low resistivity HVCMOS with 15  $\mu\text{m}$  depletion depth
- Increase of linear absorption with irradiation in Si (intrinsic to Si) reduces TPA contrast. This smearing can be corrected either by moving the focus outside of the sample (diodes) or by dual scan at different power (devices with very small depletion width).
- From the experience gained on the demonstrator system in Bilbao, we are building an optimized setup for Silicon measurements at CERN. Laser delivery this next month. Project funded by 2016 CERN Knowledge Transfer Program.
- All these activities have been carried out in the framework of CERN-RD50 collaboration

# TPA -TCT extra information

1) NIMA Vol 845, 11 February 2017, Pages 69-71

<https://doi.org/10.1016/j.nima.2016.05.070>

2) Journal of Instrumentation, Volume 12, January 2017

<http://iopscience.iop.org/article/10.1088/1748-0221/12/01/C01038>

3) I. Vila, CERN Detector Seminar, 26<sup>th</sup> January 2018

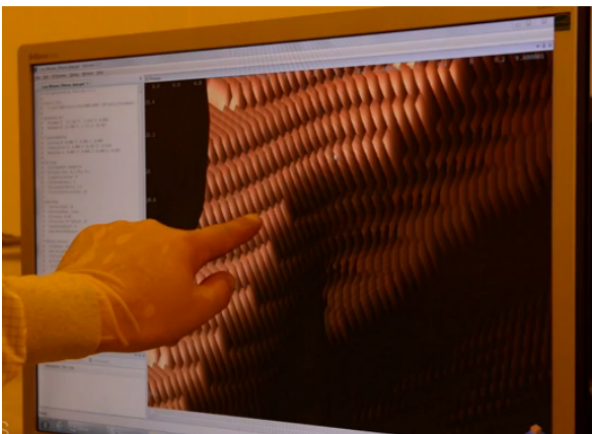
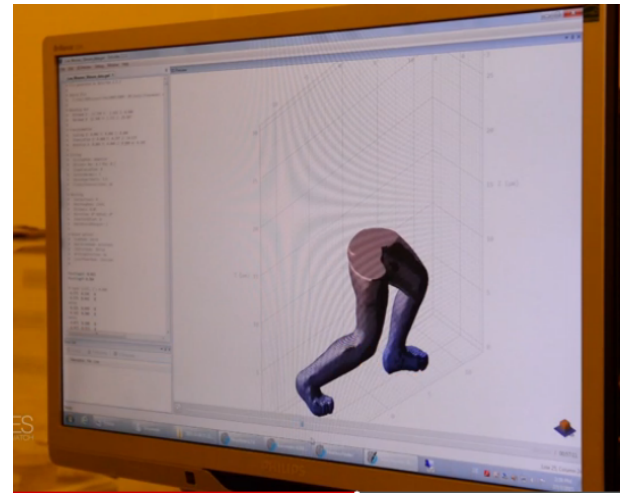
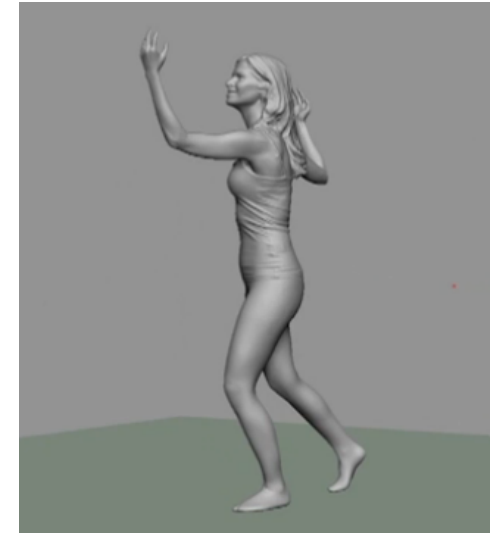
<https://indico.cern.ch/event/697958/>

4) M. Fernandez, Seminaire de physique corpusculaire, University of Geneva,  
30<sup>th</sup> Nov 2016

<http://dpnc.unige.ch/seminaire/talks/fernandez.pdf>

**BACKUPS**

**3D microfabrication:** using **focused ultrashort laser pulses** on the volume of a **photoresist**, the pulses initiate **polymerization**. After illumination of the structure and **development** (washing out the non-illuminated regions) the **polymerized** material remains in the prescribed **3D** form.



Screenshots from "Is this the world's smallest sculpture?" CNN "Ones to watch"

Further reading: **Two Photon Polymerization: a New Approach to Micromachining**, Photonics spectra, October 2006

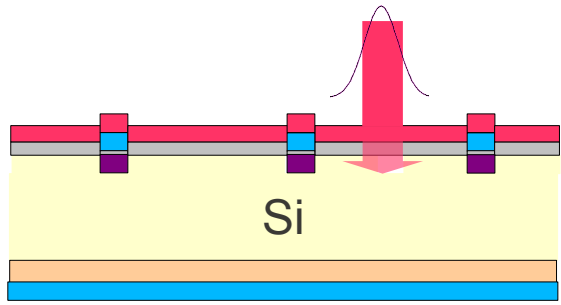


# Common characterization techniques of semiconductor radiation detectors

Technique	Type	Spatial resolution	Efficiency	Reach	Requirements
IV/CV	Electrical	Not applicable	No	$I_{leak}$ , $V_{dep} (\propto N_{eff})$ , $V_{br}$ , $C_{end}$	Needle system Cold chuck pA-meter, LCR meter...
TCT (Transient Current Technique)	Optical	2D	Yes	$N_{eff}$ , $\tau_{e,h}$ (low $\Phi_{eq}$ ) uniformity	ps-pulsed laser stepper motors Amplifier, scope
Radioactive source	Source	None	Yes	$V_{dep}$ , CCE	Radioactive source, trigger system
IBIC	Source	2D	Yes	2D CCE	Accelerator
Test beam	Source	2D	Yes	Setup/sensor dependent: 2D CCE,...	Accelerator, beam time, telescope, trigger,

MIP capabilities: Test beam, infrared TCT

Laser characterization techniques

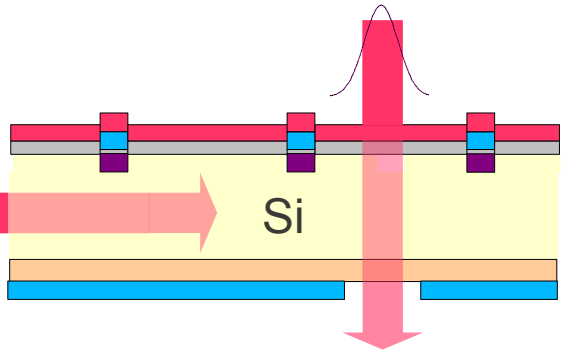


### SPA-TCT Red

Employing short absorption length laser (red for Si), all carriers deposited in few  $\mu\text{m}$  from surface. Allows to study drift of one kind of carriers. No **spatial resolution along beam direction**.

**1 photon=1 e-h pair**

SPA: Single Photon Absorption  
TCT: Transient Current Technique



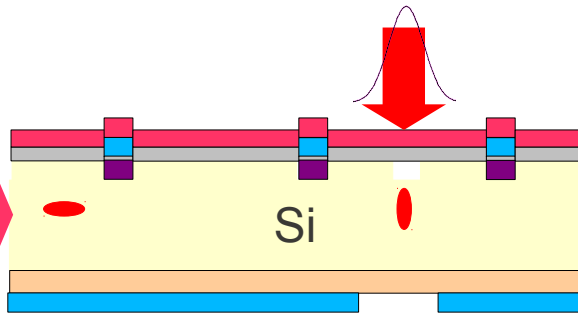
### SPA-TCT Infrared

Using long absorption length laser (infrared for Si). Homogeneous distribution along "Rayleigh length". Similar to MIPs, though different  $dE/dx$ .

Incidence can be from **top, bottom or edge**.

**Edge: lateral spatial resolution.**

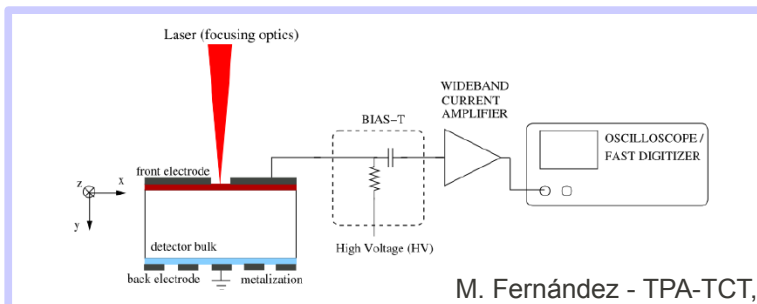
**1 photon=1 e-h pair**



### Two Photon Absorption (TPA-TCT)

Point-like energy deposition  $\rightarrow$  3D spatial resolution  
**Novel technique** developed by **IFCA, CERN, US, UPV**

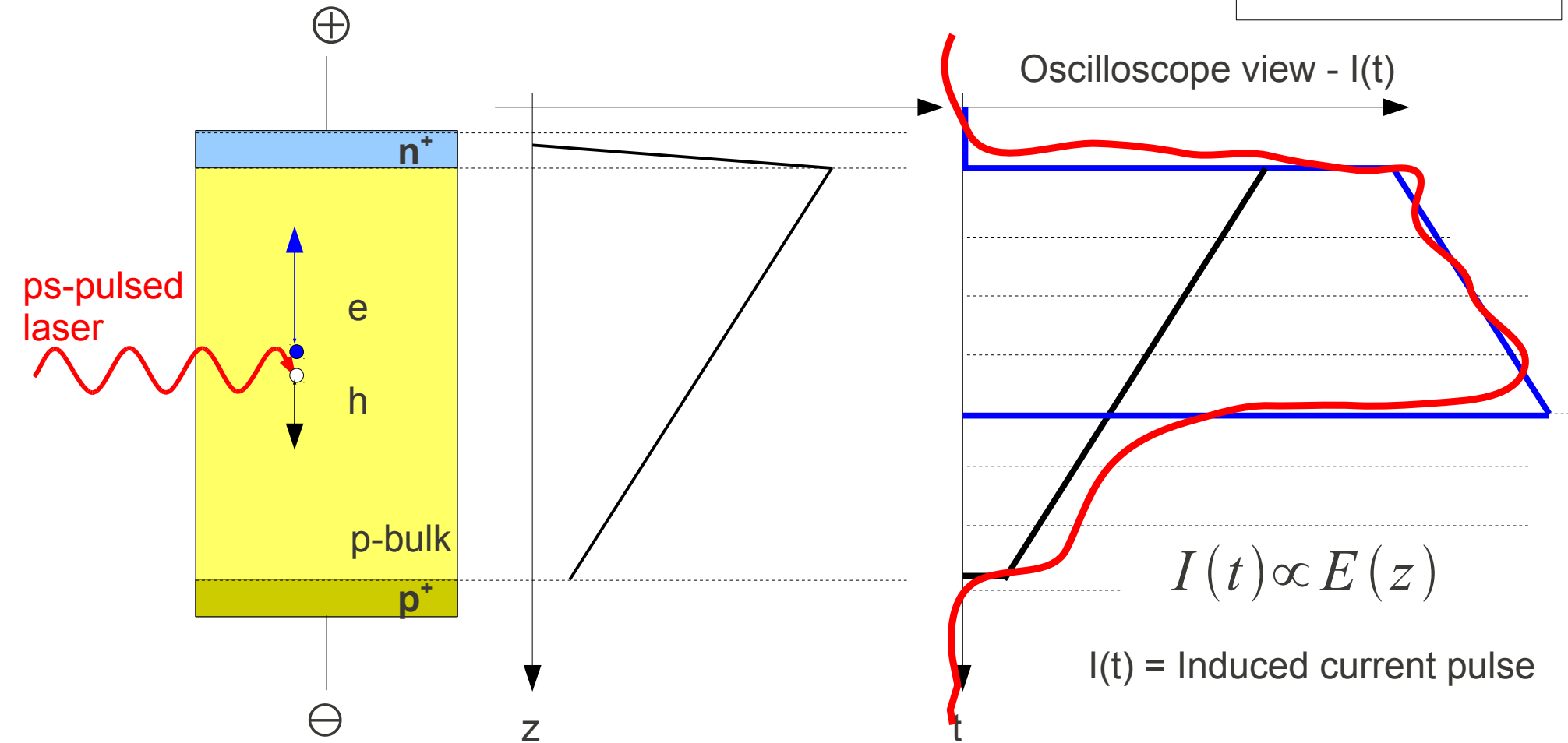
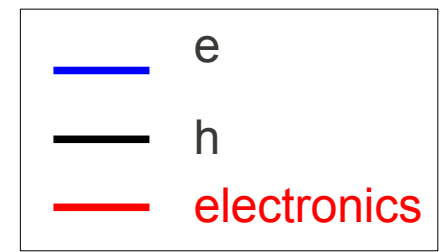
**2 photons=1 e-h pair**



### Transient Current Techniques

Applicable to both pad/segmented detectors  
Simple readout  
DAQ directly on digital scope

# TCT: Transient Current Technique



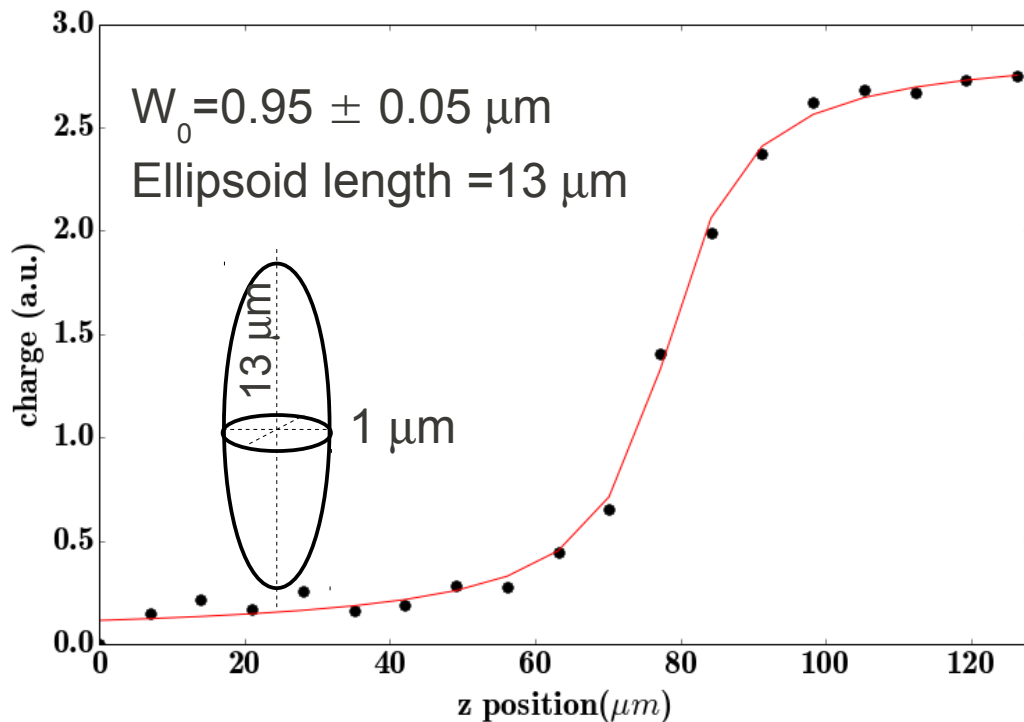
$$I(t) = N_{eh} A q_e v_{drift} E_W \propto v_{drift} = \mu(E) E \Rightarrow I(t) \propto E(z)$$

**Assumption:** 1D, overdepleted, non-irradiated diode

# Evidences for TPA process

- 1) Collected charge varies quadratically with power
- 2) Z-scan is not Z-invariant.

Then characterize the excitation volume:



Ellipsoid is completely described by waist ( $w_0$ ),  $\lambda$  and  $\beta$ .

$$w(z) = w_0 \sqrt{\frac{\lambda z}{\pi w_0^2 n}}$$

$$I(z) = \frac{2P}{\pi w^2(z)} e^{\frac{-2r^2}{w^2(z)}}$$

$$\frac{dN(r, z)}{dt} = \frac{\beta_2 I^2(r, z)}{2\hbar\omega}$$

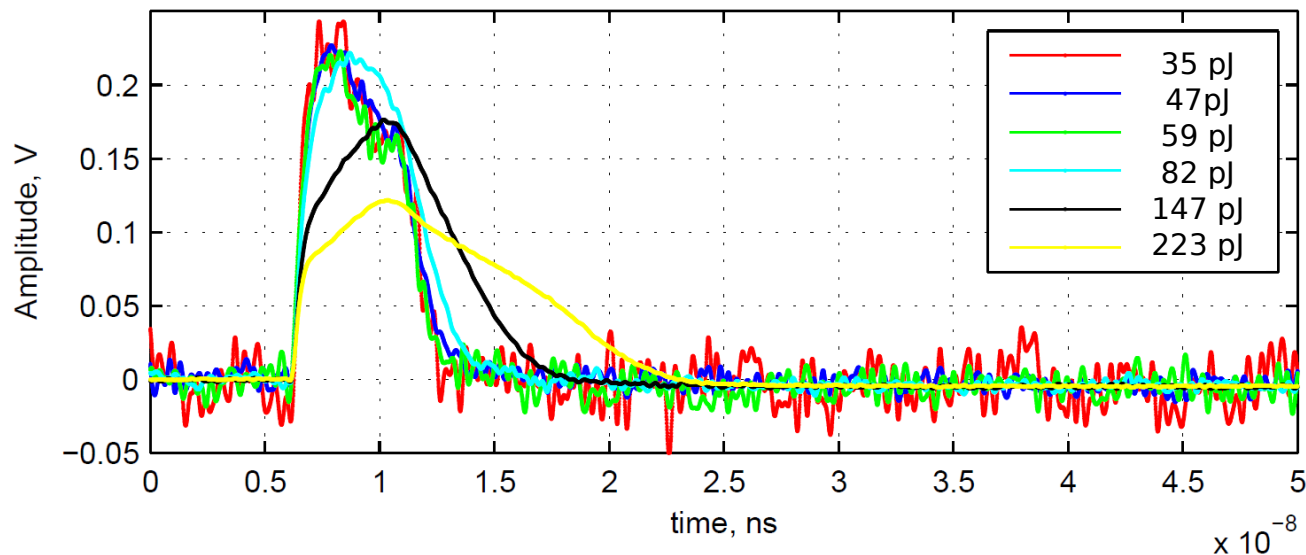
$$t \sim t_p \rightarrow N(z) = \int_{-\infty}^{\infty} 2\pi r \cdot t_p \cdot N(r, z) dr$$

An **edge-TPA** scan is optimum, because spatial resolution is  $\sim 1 \mu\text{m}$

Try to scan **pads from the edge** → Get active area very close to the border

⇒ Work in a power regime where  $\beta \gg \alpha$  but **without** producing **plasma**.

Focalización dentro 15um-FZ285223B-Diferente potencia-Vbias=200v



■ For this detector.  
Laser power < 80 pJ  
⇒ no plasma effect

The primary equations governing pulse propagation and carrier generation in a semiconductor material are [15], [16]

$$\frac{dI(r, z)}{dz} = -\alpha I(r, z) - \beta_2 I^2(r, z) - \sigma_{ex} N I(r, z) \quad (1)$$

$$\frac{d\Phi(r, z)}{dz} = \beta_1 I(r, z) - \gamma_1 N(r, z) \quad (2)$$

$$\frac{dN(r, z)}{dt} = \frac{\alpha I(r, z)}{\hbar\omega} + \frac{\beta_2 I^2(r, z)}{2\hbar\omega} \quad (3)$$

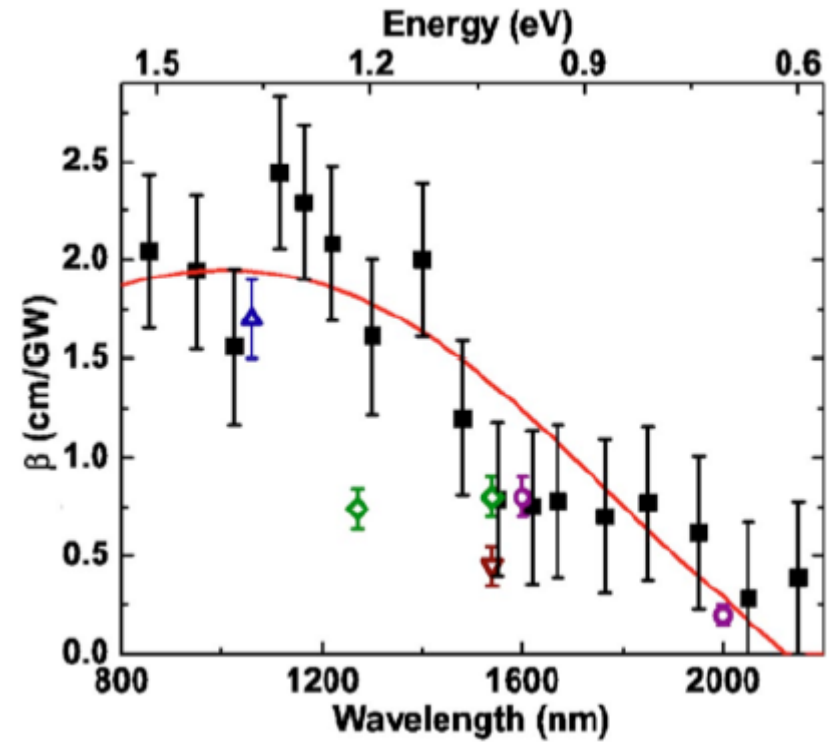
where  $I$  is pulse irradiance,  $N$  is the density of free carriers, and  $\Phi$  is the phase.  $\alpha$  is the linear absorption coefficient,  $\beta_2$  is the two-photon absorption coefficient that is proportional to the imaginary part of  $\chi^{(3)}$  (the third-order nonlinear-optical susceptibility),  $\sigma_{ex}$  is the absorptivity of laser-generated free carriers,  $\beta_1$  is proportional to the real part of  $\chi^{(3)}$ ,  $\gamma_1$  describes the refraction due to free carriers, and  $z$  is the depth in the material.

$$N_{1P}(z) = \frac{\alpha}{\hbar\omega} \exp(-\alpha z) \int_{-\infty}^{\infty} I_o(z, t) dt$$

$$N_{2P}(z) = \frac{\beta_2}{2\hbar\omega} \int_{-\infty}^{\infty} I^2(z, t) dt$$

ered carefully. When nonlinear absorption is the only loss mechanism in a material, the irradiance as a function of depth is given by

$$I(z) = \frac{I_o}{1 + \beta_2 I_o z}. \quad (9)$$



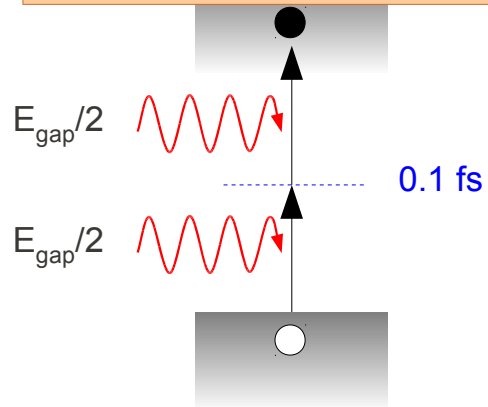
Applied Physics Letters, vol. 90, no. 19, p. 191104, 2007.

the longitudinal dependence of the beam radius  $w(z)$  is

$$w(z) = w_o \left[ 1 + \left( \frac{\lambda z}{\pi w_o^2 n} \right)^2 \right]^{1/2}. \quad z_o = \pm \frac{\pi n w_o^2}{\lambda}.$$

The parameter  $2z_o$  defines the propagation distance over which the beam is reasonably well collimated in the vicinity of  $w_o$ . In silicon ( $n \approx 3.51$ ), for  $1.26 \mu\text{m}$  light and  $0.8 \mu\text{m}$  beam radius (that of the present study),  $2z_o$  is  $\sim 11.2 \mu\text{m}$ .

### Non irradiated



$$\frac{dN(r, z)}{dt} \propto \frac{\beta_2 I^2(r, z)}{2 \hbar \omega}$$

**Two Photon Absorption**

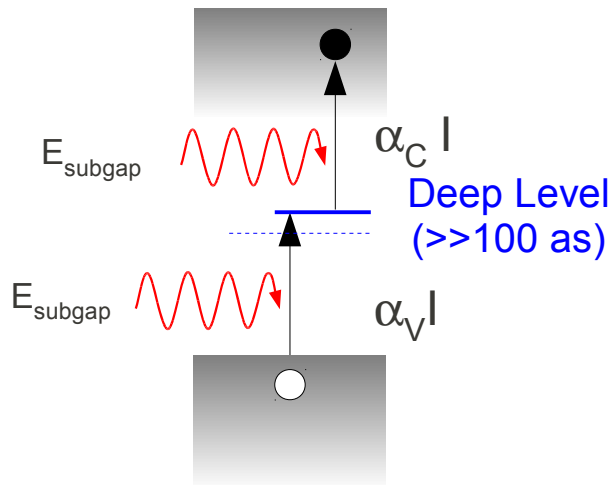
**TPA:** Needed two (simultaneous) photons to produce 1 e-h pair. Mediated by virtual state

**Resonant TPA:** Two (sequential) photons to produce 1 e-h pair. Mediated by Deep Level (DL)

**Two SPA:** Unpaired holes (in BV) or electrons (BC). Generation is proportional to number of DLs (it grows with fluence) and laser intensity (more photons available for transition).

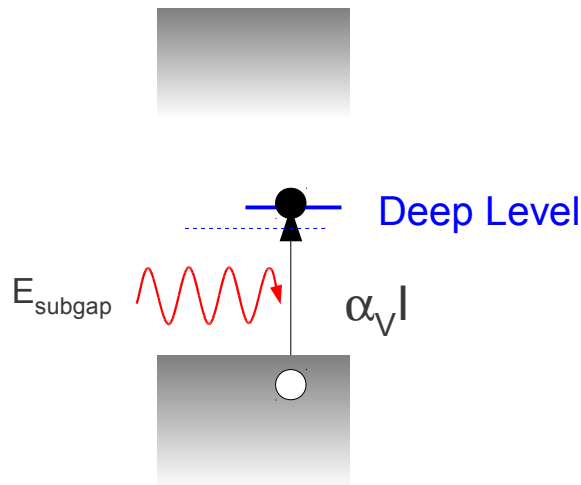
### Irradiated

#### Resonant (sequential) TPA

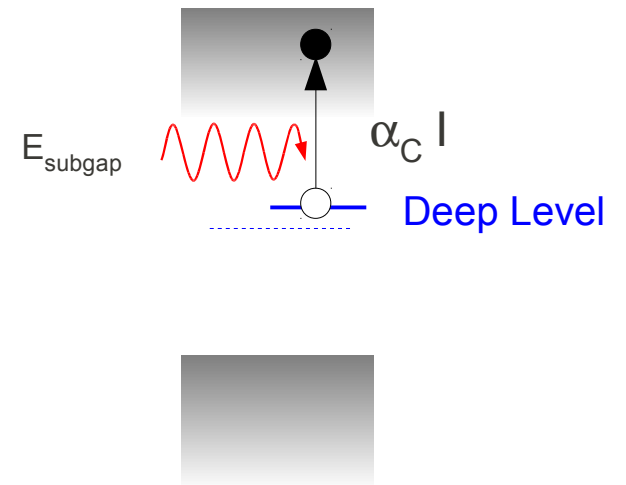


$$\frac{dN(r, z)}{dt} \propto \frac{\alpha_V \cdot \alpha_C I^2}{2 \hbar \omega}$$

#### Two Single Photon Absorption



$$\frac{dN(r, z)}{dt} \propto \frac{\alpha_V(\Phi) I}{\hbar \omega}$$



$$\frac{dN(r, z)}{dt} \propto \frac{\alpha_C(\Phi) I}{\hbar \omega}$$

# HVCMOS: correction of radiation induced SPA signal

Generation rate of e-h pairs per unit volume:

$$\frac{dN(r, z; i)}{dt} = \alpha \frac{I(r, z; i)}{\hbar \omega} + \frac{\beta_2 I^2(r, z; i)}{2 \hbar \omega}$$

with  $I(r, z) = \frac{2P}{\pi w(z)^2} \exp \frac{-2r^2}{w(z)^2}$

A Ge photodiode (linear response at 1300 nm) is used to measure laser power, which is proportional to the integral of the irradiance:  $I(r, t; i)$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(r, z) dr dz \propto P \quad \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} I(r, z)^2 dr dz \propto P^2$$

Then:

$$Q_{DUT} = q_e \int_r \int_z N(r, z; i) dr dz = q_e N(i)$$

This is the charge measured in the detector (easy to calculate)



# HVCMOS: correction of radiation induced SPA signal

Assuming a “rectangular” laser pulse of duration  $t_p$  ( $\sim 30$  fs)

$$\int \alpha \frac{I(r, z; i)}{\hbar \omega} dt + \int \frac{\beta_2 I^2(r, z; i)}{2 \hbar \omega} dt = \alpha \frac{P t_p}{\hbar \omega} + \frac{\beta_2 t_p P^2}{2 \hbar \omega}$$

So finally:

$$Q_{DUT} = \frac{q_e \alpha t_p}{\hbar \omega} P + \frac{q_e \beta_2 t_p}{2 \hbar \omega} P^2 \quad \left\{ \begin{array}{l} Q_\alpha = \frac{q_e \alpha t_p}{\hbar \omega} \\ Q_\beta = \frac{q_e \beta_2 t_p}{2 \hbar \omega} \end{array} \right.$$

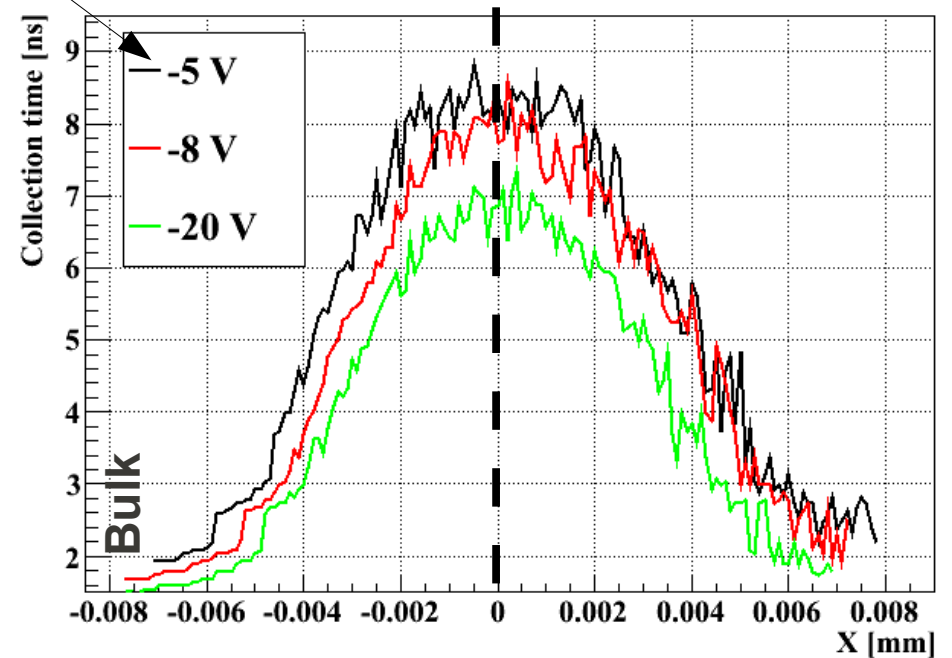
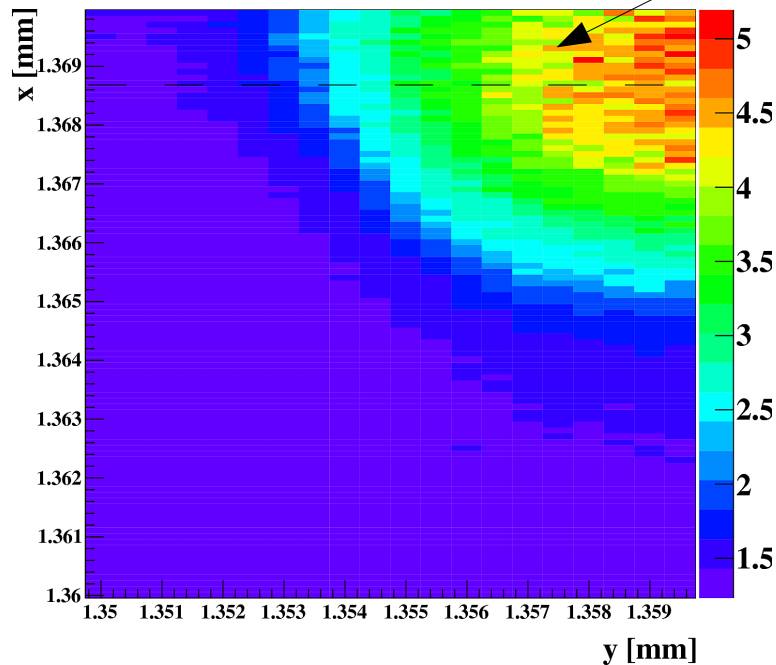
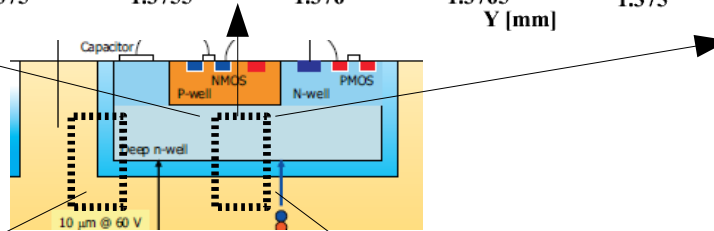
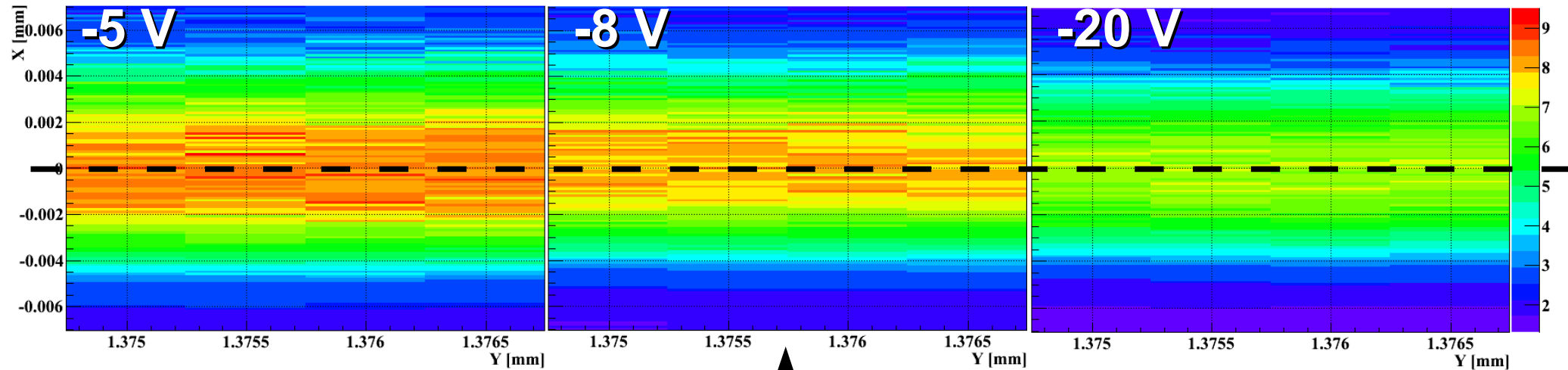
With  $Q_\alpha$  and  $Q_\beta$  the SPA and TPA contributions to the signal (they should be constants)

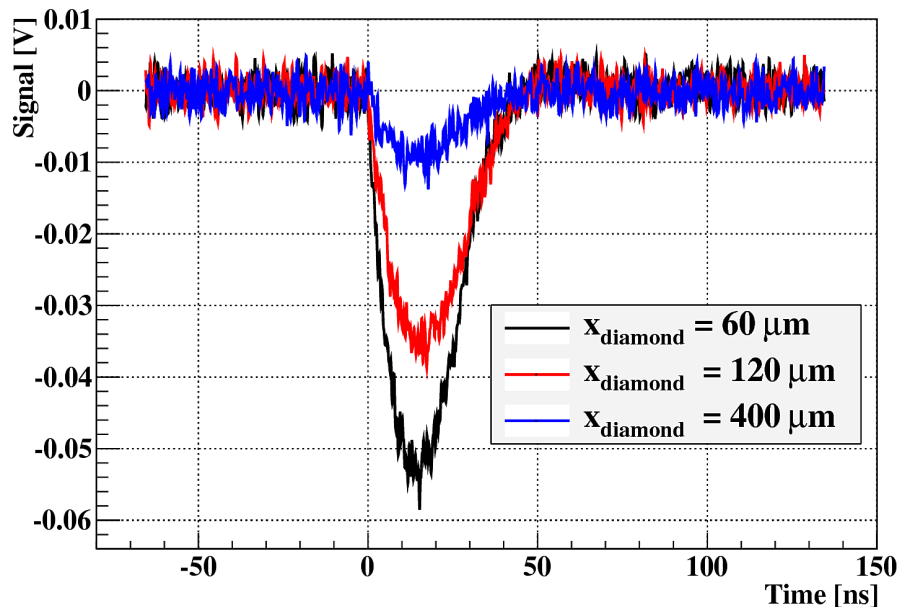
Now, we can take 2 identical measurements at 2 different intensities and calculate  $Q_\beta$

$$\left. \begin{array}{l} Q_1 = Q_\beta P_1^2 + Q_\alpha P_1 \\ Q_2 = Q_\beta P_2^2 + Q_\alpha P_2 \end{array} \right\} \quad Q_\beta = \frac{\frac{Q_1}{P_1} - \frac{Q_2}{P_2}}{P_1 - P_2} \quad Q_\alpha = \frac{\frac{Q_1}{P_1} - \frac{Q_2}{P_2}}{\frac{1}{P_1} - \frac{1}{P_2}}$$

The  $\alpha$  corrected signal is:  $Q_\beta P_1^2 = Q_1 - Q_\alpha P_1$

# Fun: High Resolution imaging inside implant

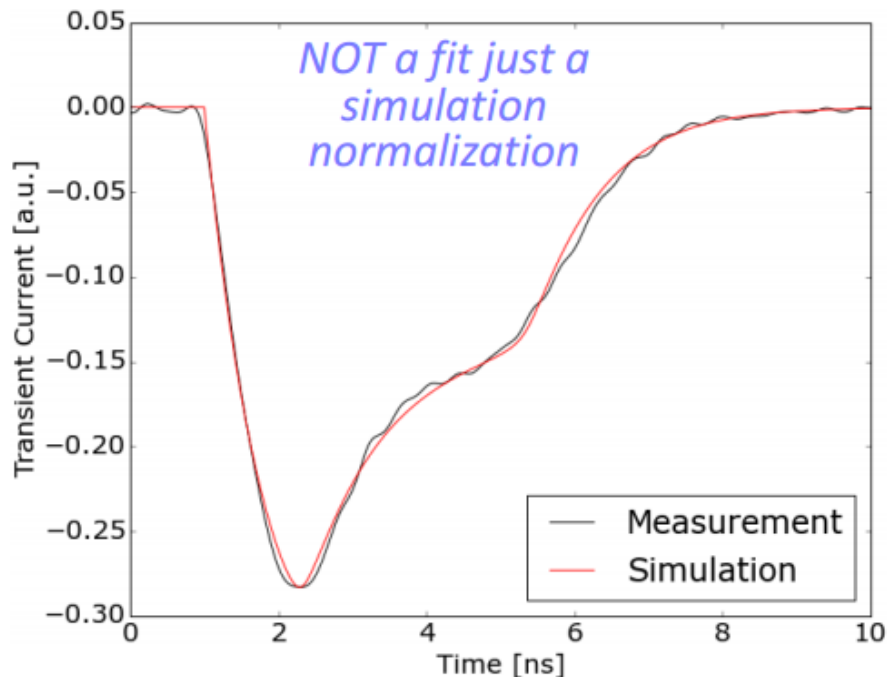




- Applying standard TCT to diamond because would require a **UV laser**.
- Using  $\lambda_{\text{TPA}} = 400 \text{ nm}$  we did TPA-TCT in diamond
- Picture to the left is just a proof of principle: we got a signal in a diamond device.

That's a first laser TCT in diamond ever !

## DAQ, simulation and analysis packages



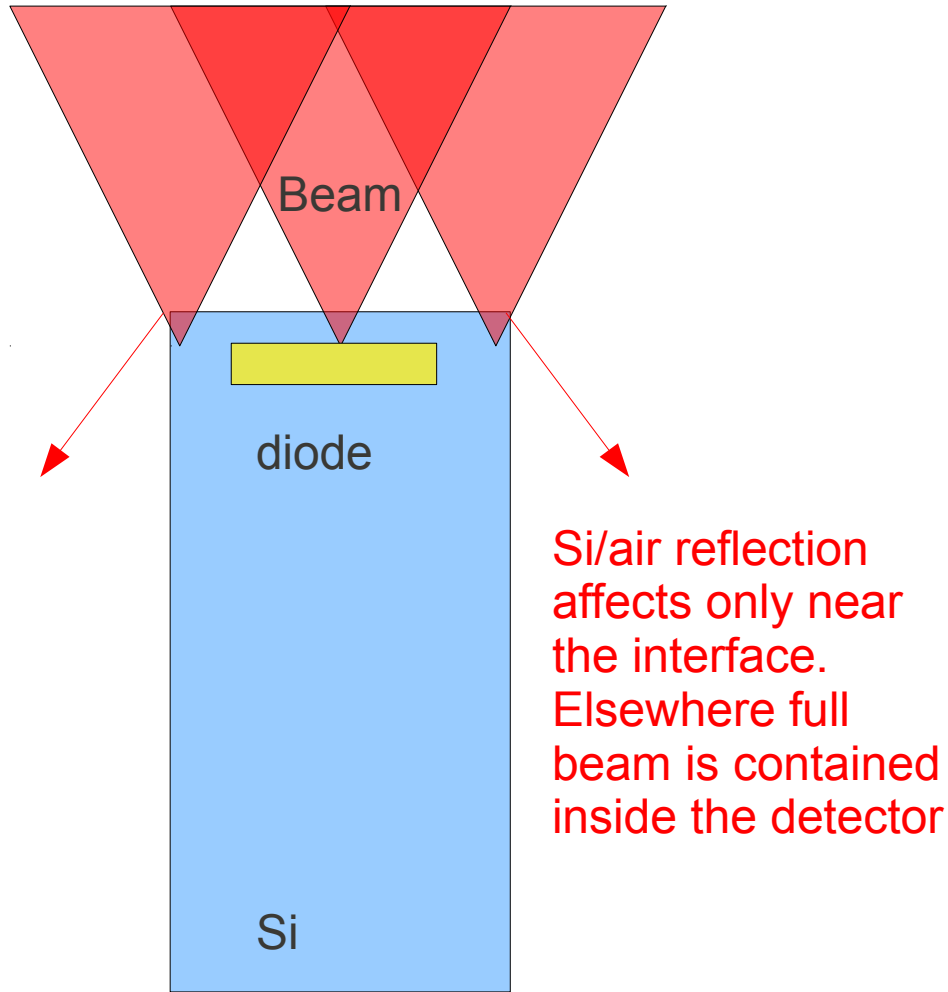
- We have **developed** a simulation package called **TRACS**:

<https://github.com/JulesDoc/Tracs>

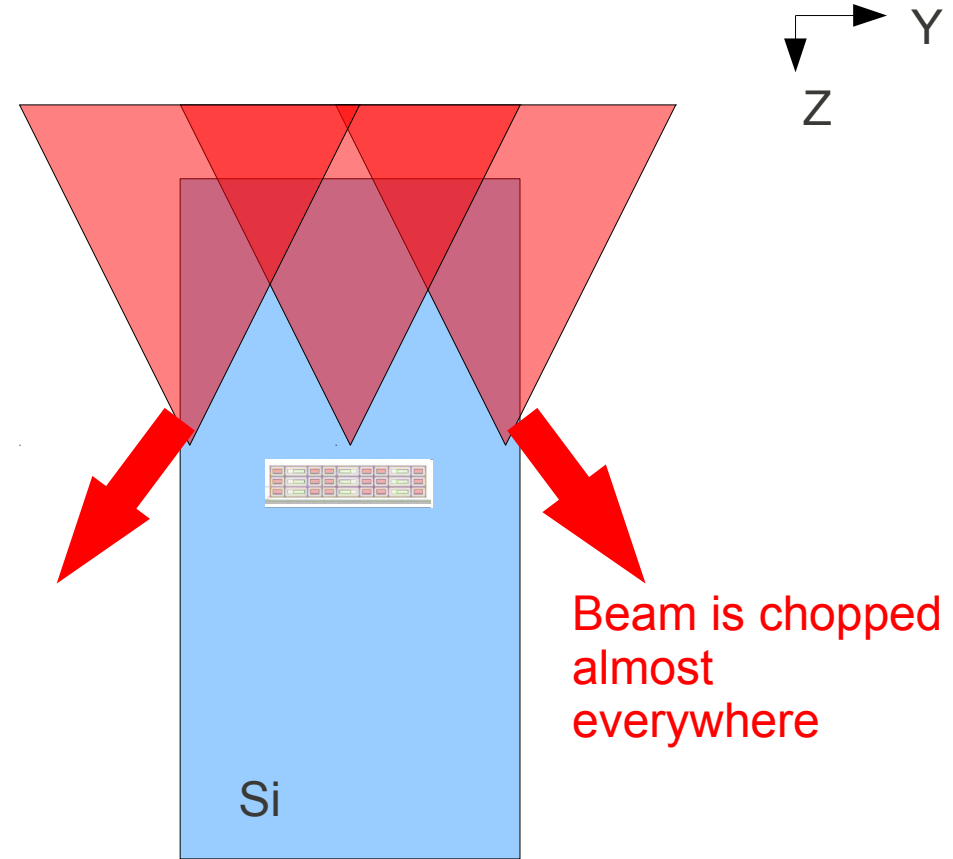
that simulates any TCT, including TPA-TCT.

- We also have a complete **ROOT-based** software package to do **3D analysis** of TPA data.

# Beam coupling to deep motifs

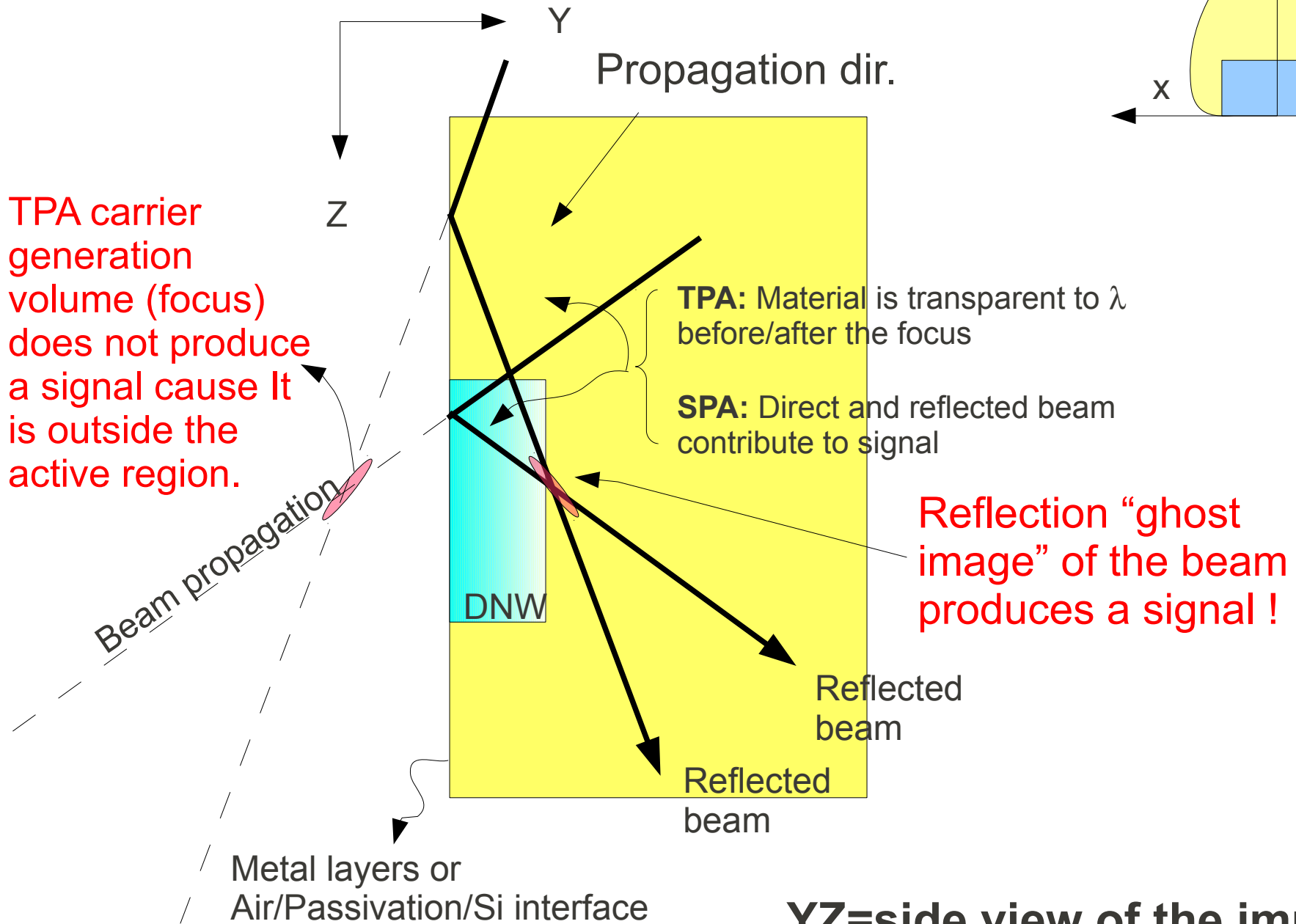


Shallow motif to scan  
 Small loss due to reflection at the borders  
 Approximately constant coupling of energy to focus.  
**Good layout** for TPA measurement

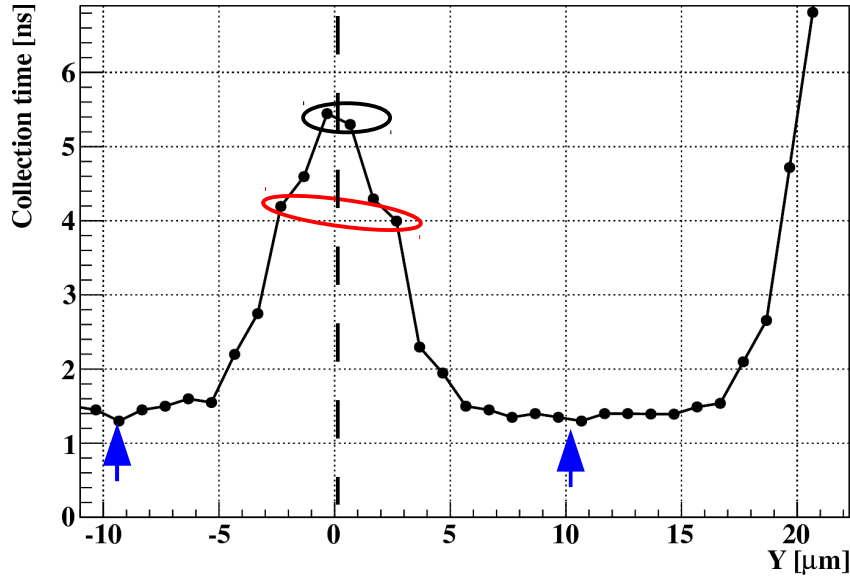


Deep motif to scan  
 Very asymmetric coupling to the focus  
**Bad layout** for TPA measurement

# Beam reflection at interfaces



# Comparison of direct and ghost signals



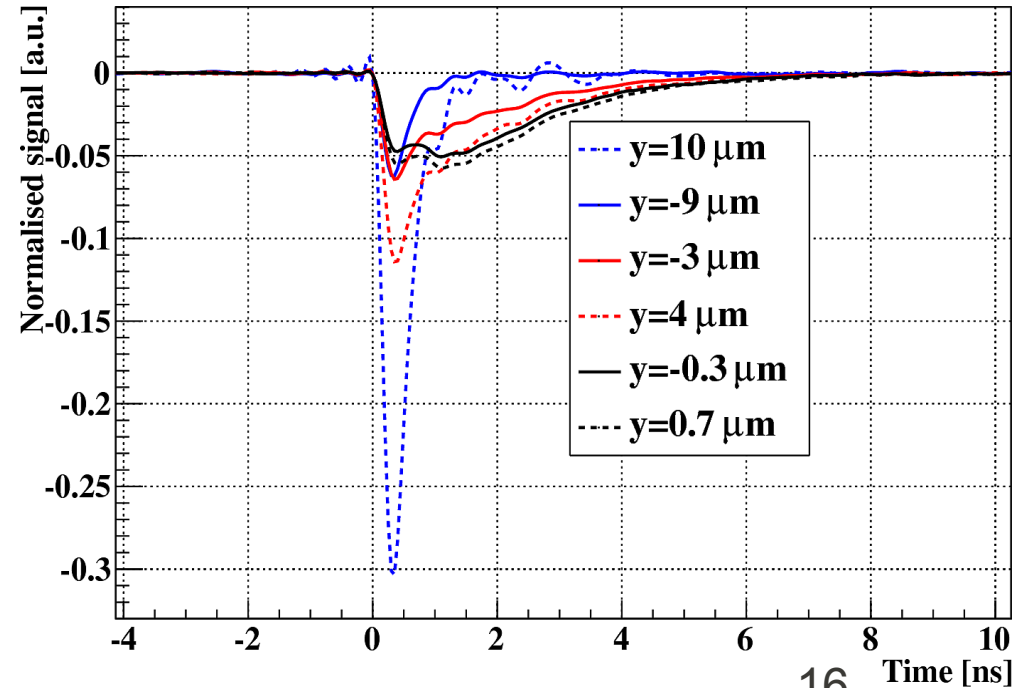
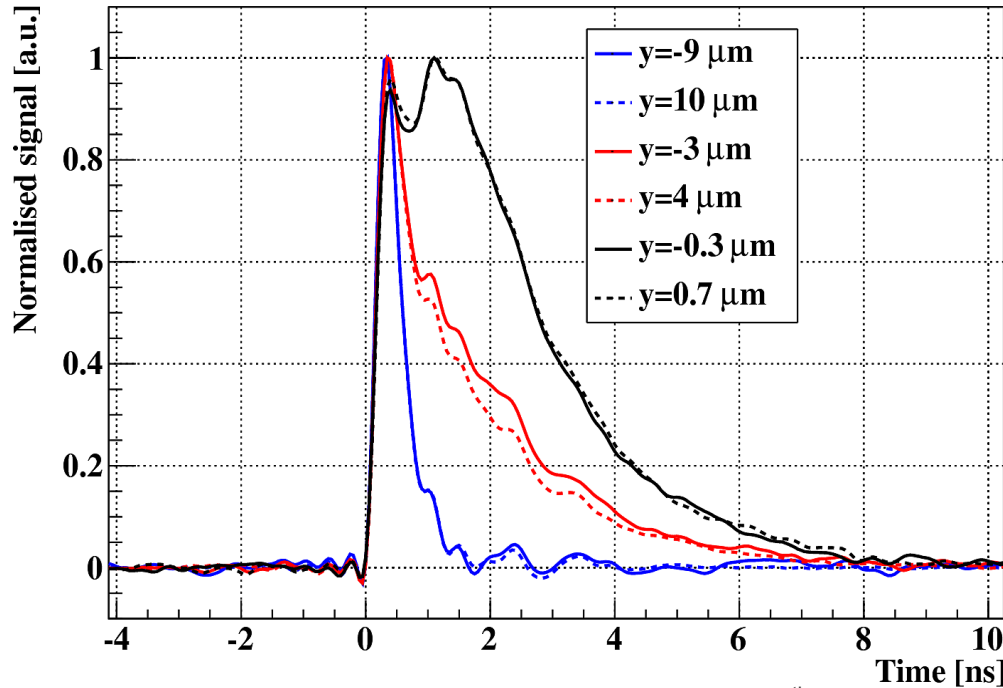
Choosing pairs of waveforms:

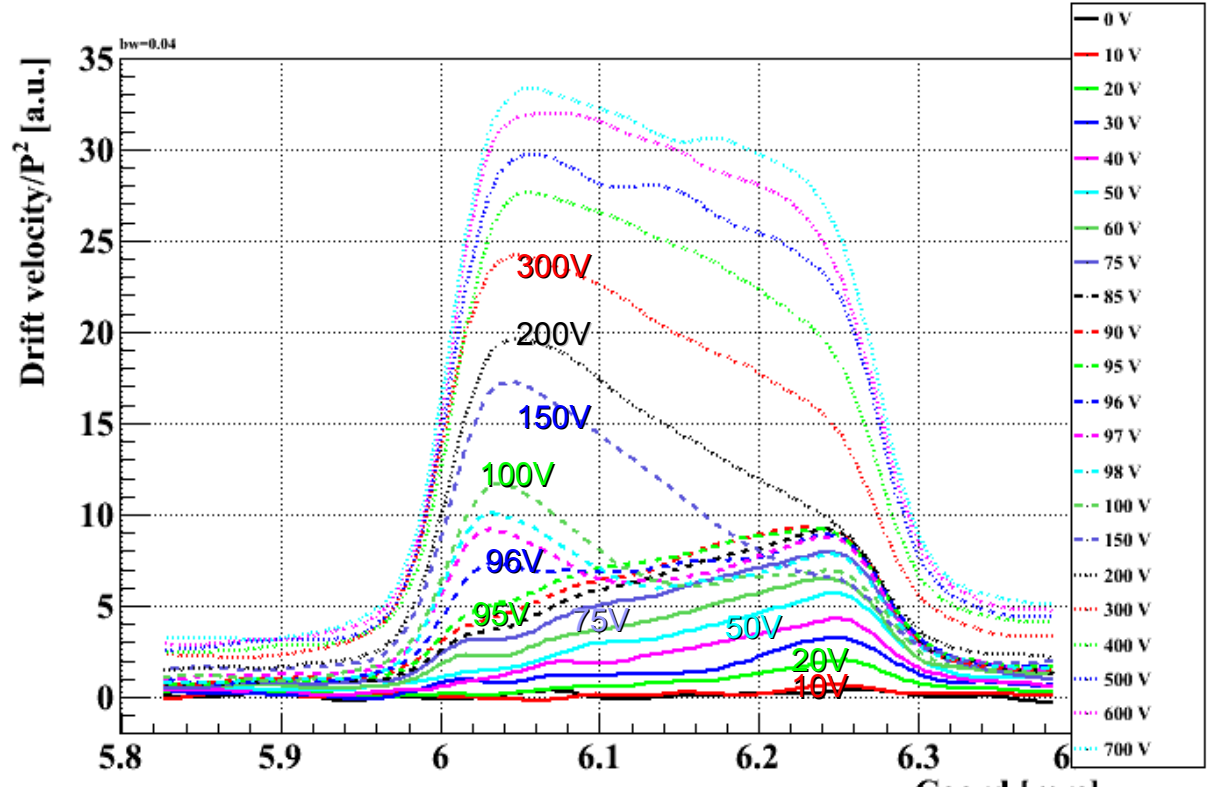
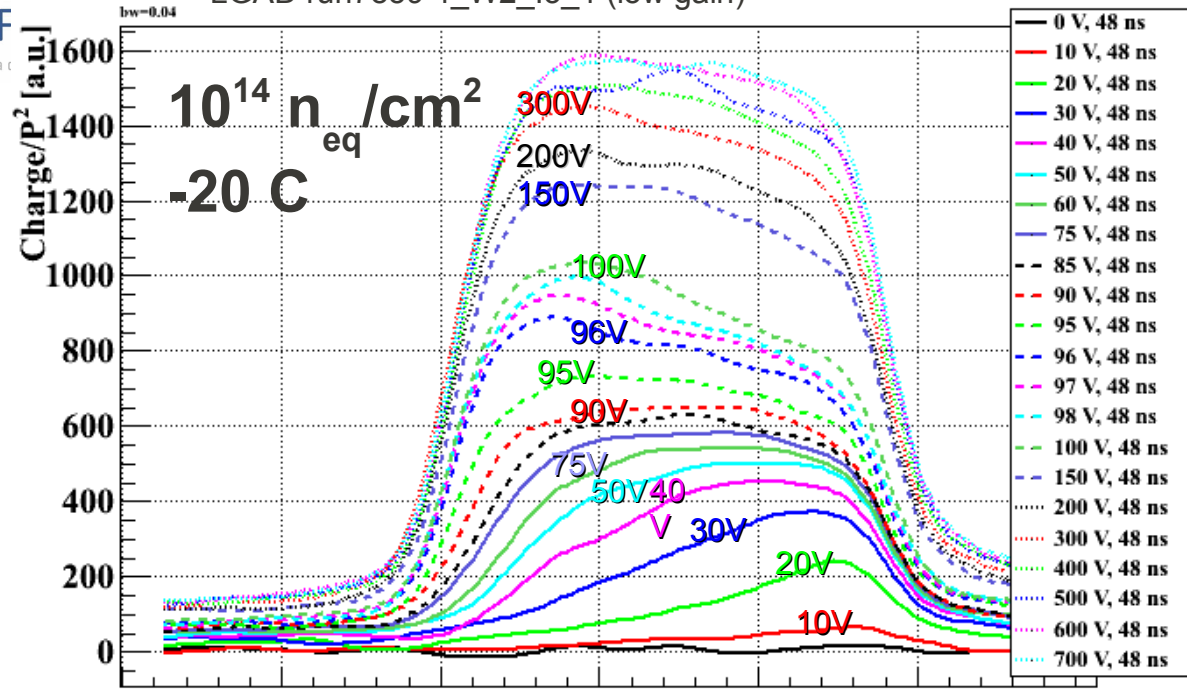
- 1) at the interface
- 2)  $\pm 2 \mu\text{m}$  away
- 3)  $\pm 10 \mu\text{m}$  away

Seen reflection of pulses well inside the bulk.

In SPA-edge-TCT this reflection can not be easily resolved because the beam is continuous.

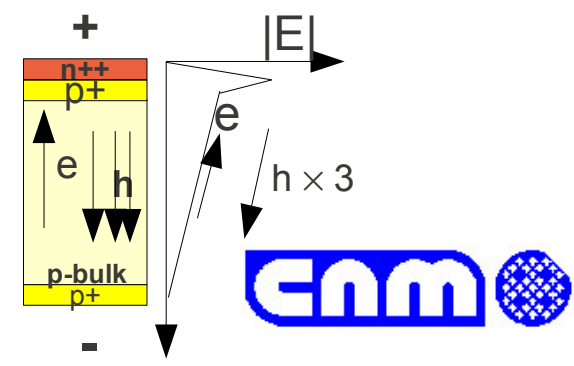
I. Vila, proceedings Pixel 2016, Sept. 2016





# Top-TPA on LGADs

- Top TPA on LGAD
- PS protons:  $10^{14} n_{eq}/cm^2$



**0-95 V:** Charge collection starts from the back.

SCSI:  $p \rightarrow n$

Inverted p-type device

**96-200V:** Front junction develops and overcomes back junction

**>200V:** Overdepletion

- Proves existence of double junction mechanism in LGADs. Impact on Acceptor Removal interpretation

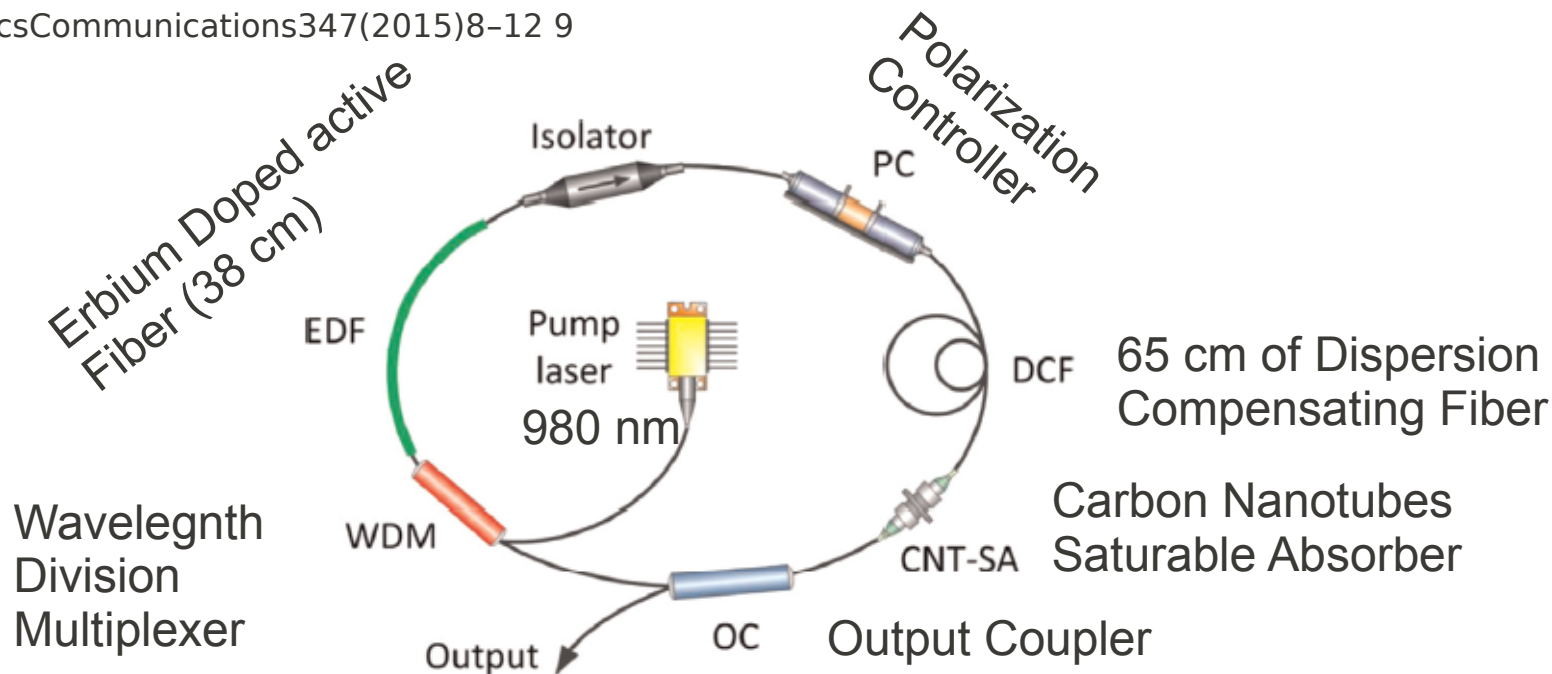


Fig. 1. Setup of the dispersion-managed, Er-doped seed laser.

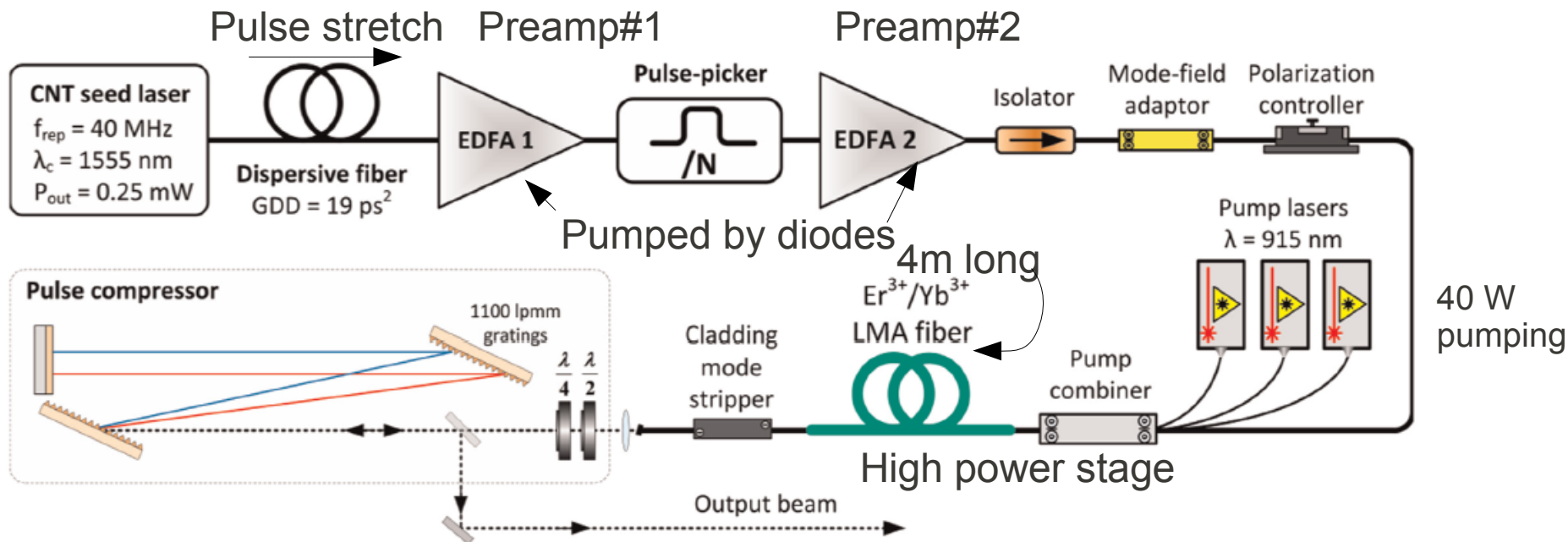


Fig. 3. Schematic of the complete CPA setup.



**Z scan Technique  
Funny and Useful!**

**SubSurface Laser Engraving (SSLE)**

Typically in BK7 Glass (Borosilicate doped with potassium)

Also with pure quartz (SiO<sub>2</sub>)

Pico or FemtoSecond Laser, 1064 nm (SiO<sub>2</sub>), 532 nm (BK7)

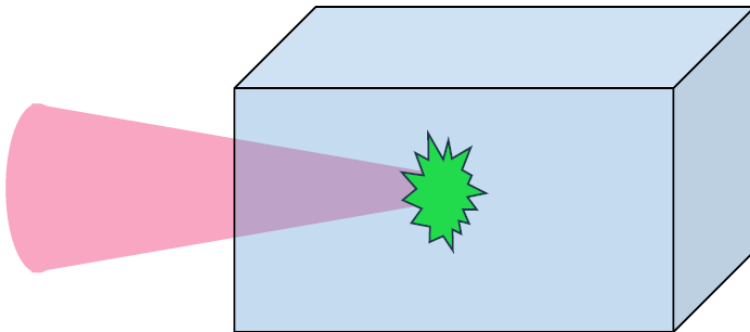
Multi-Photon Absorption

Free electron creation in the focus point

FotoChemistry in Solids:

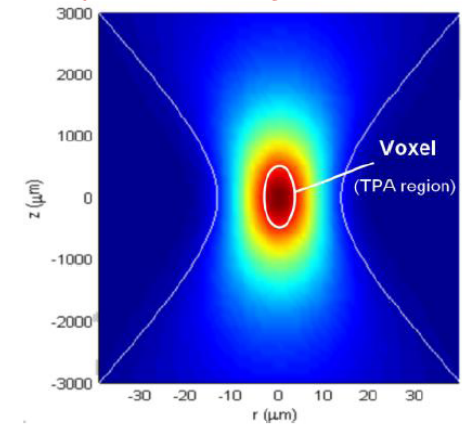
Index of refraction changes,

Color centers

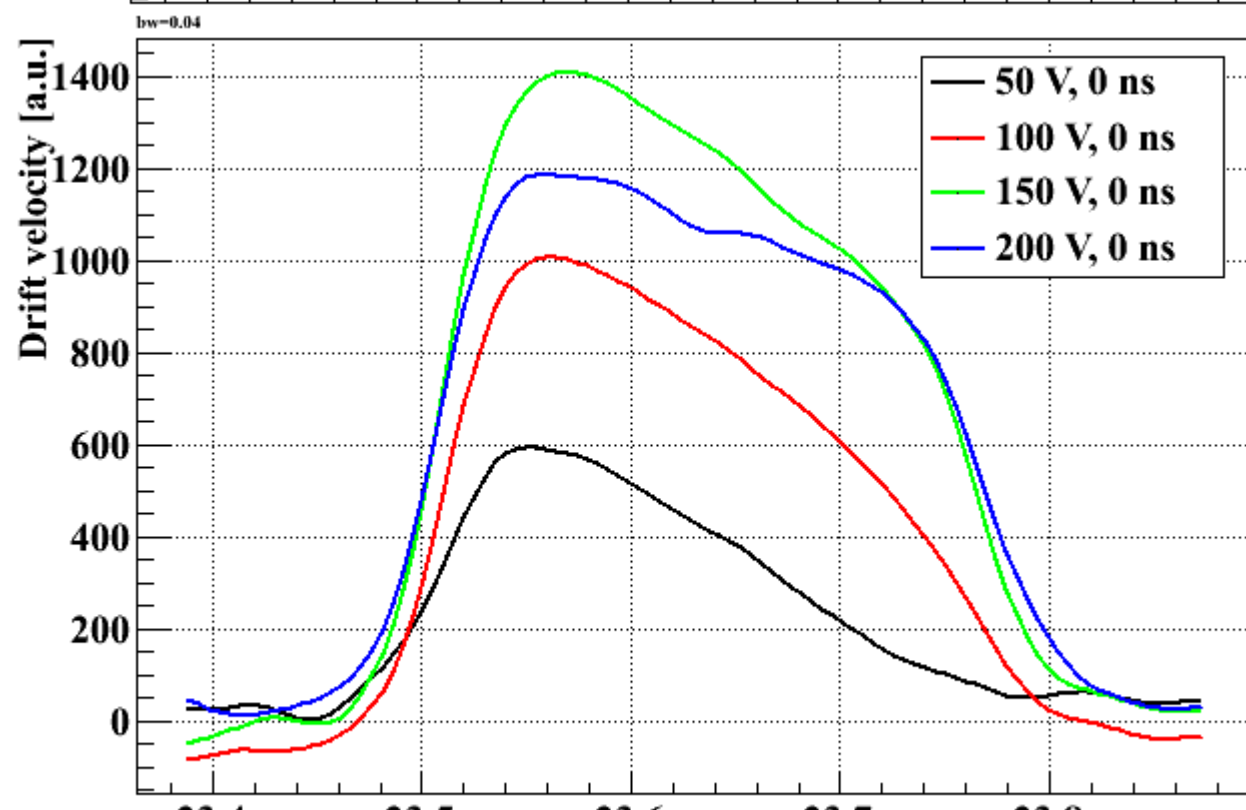
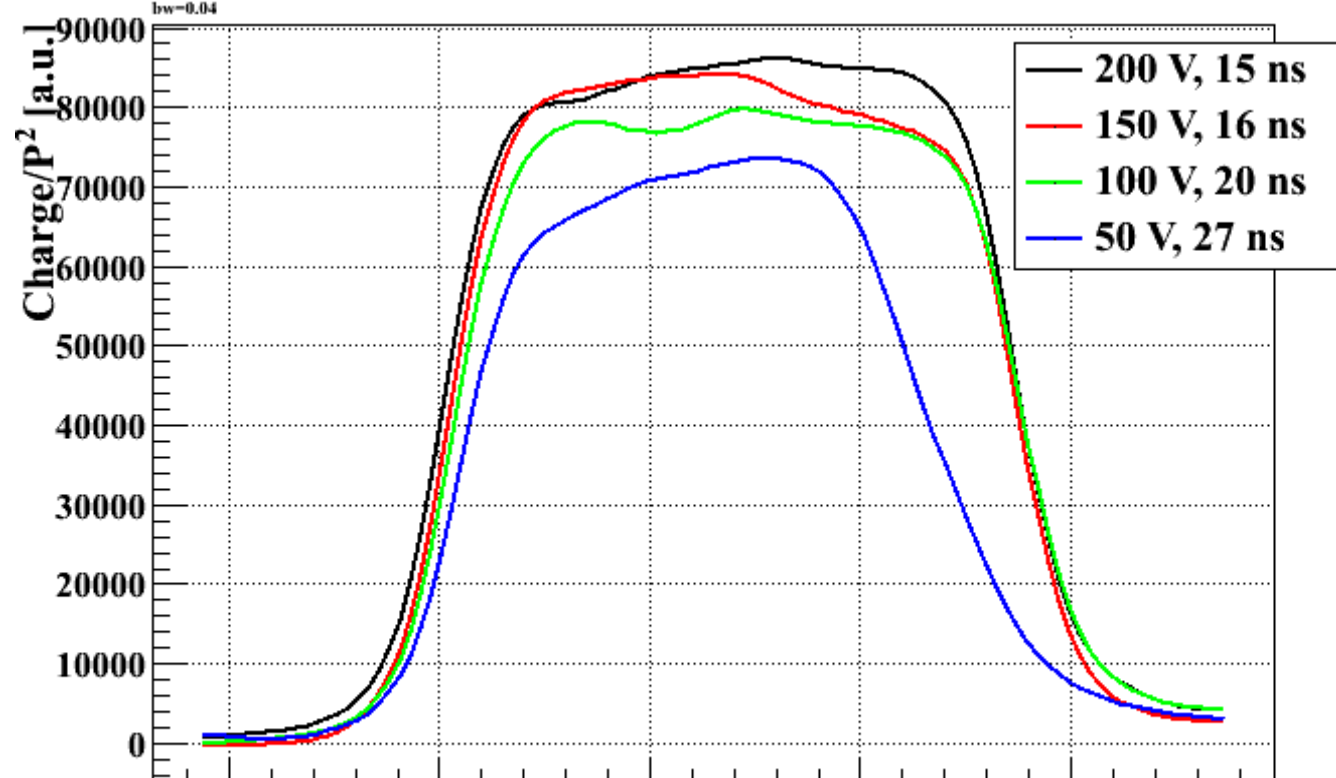


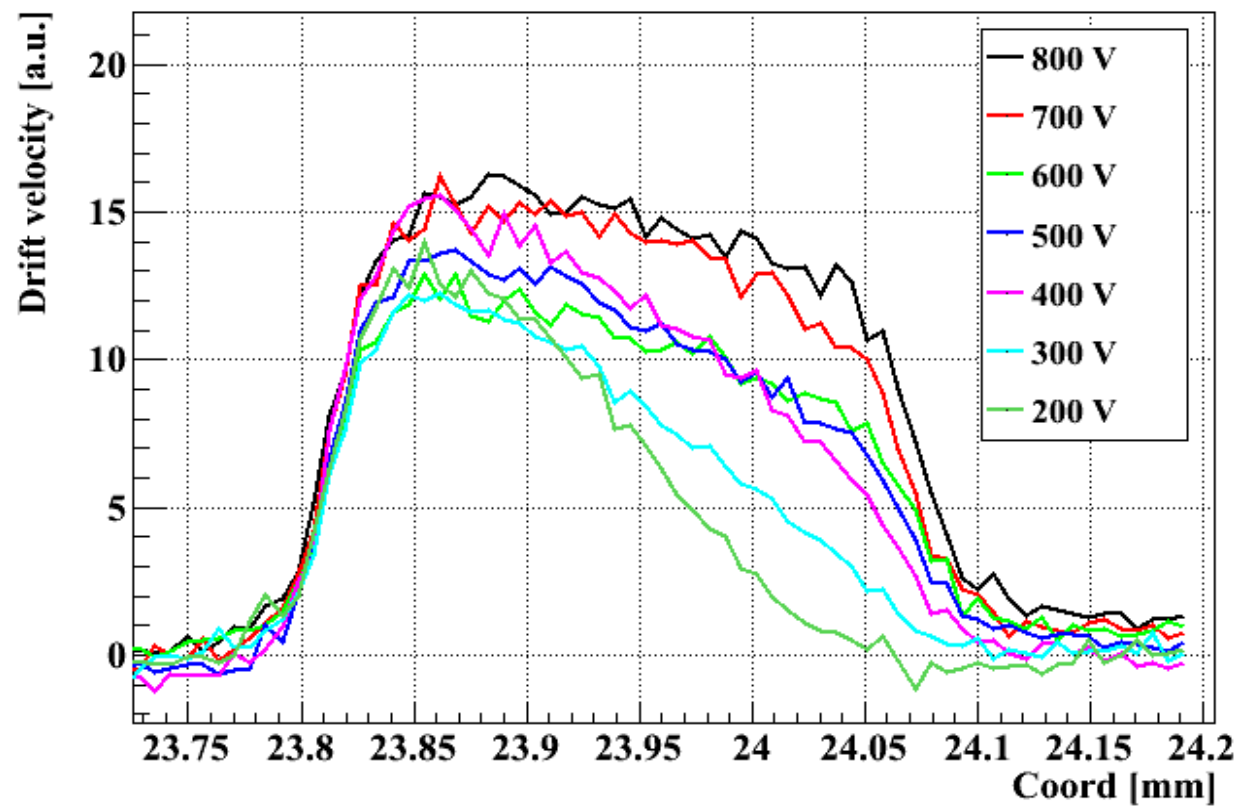
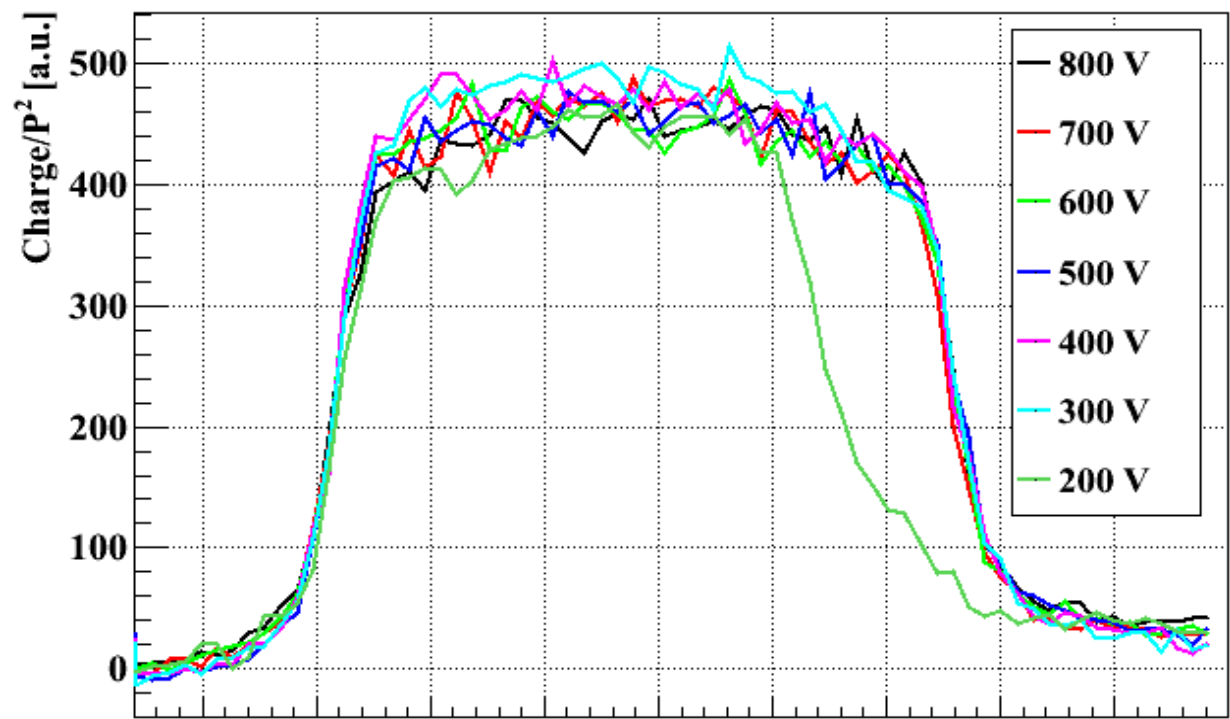
TJDP-532K Machine (532 nm, BK7 crown glass)

<http://www.tianjunlaser.com/>

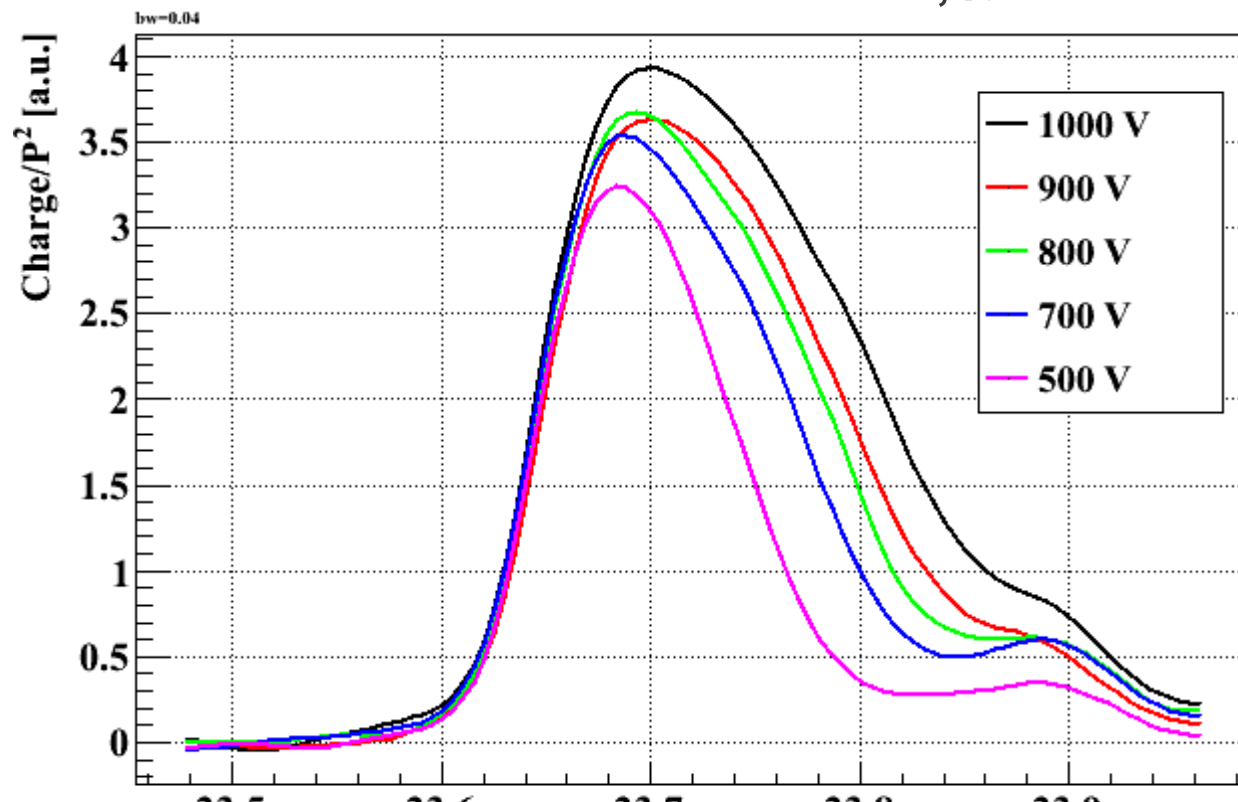
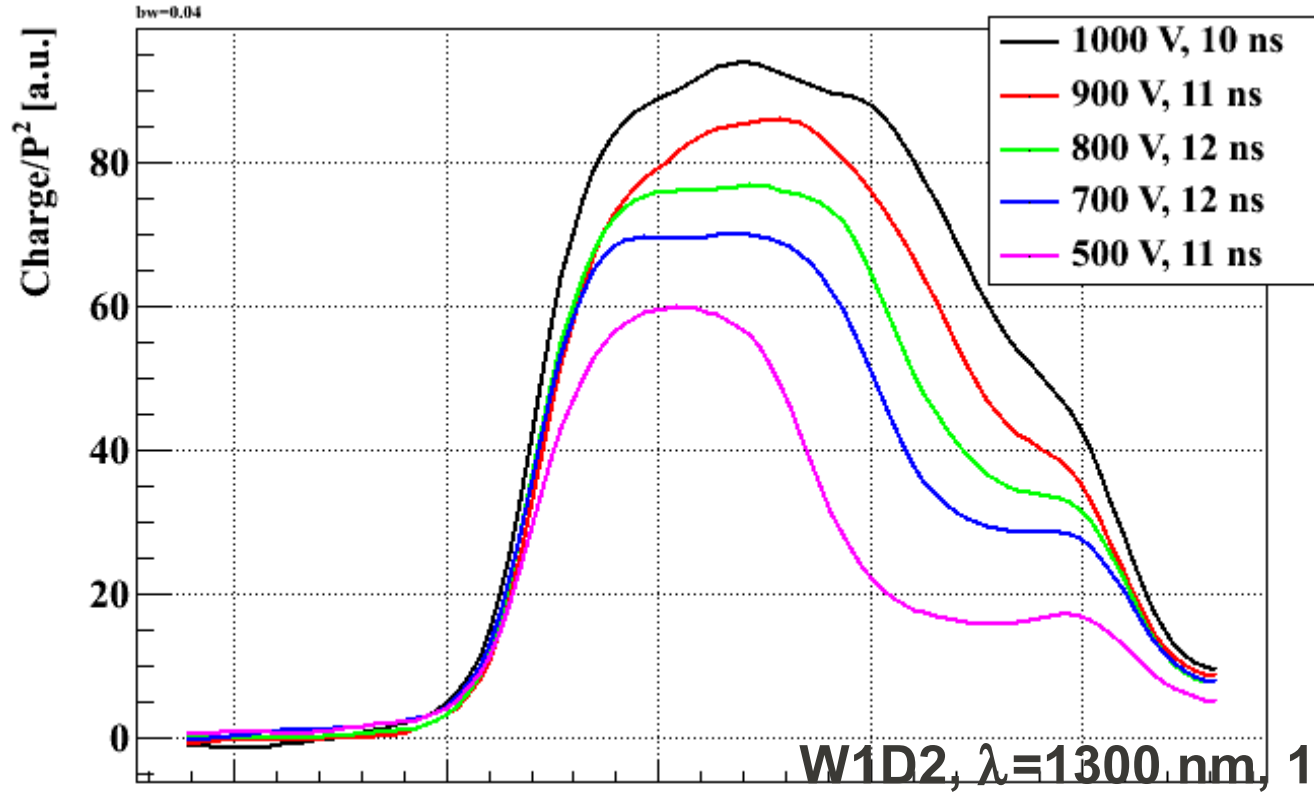


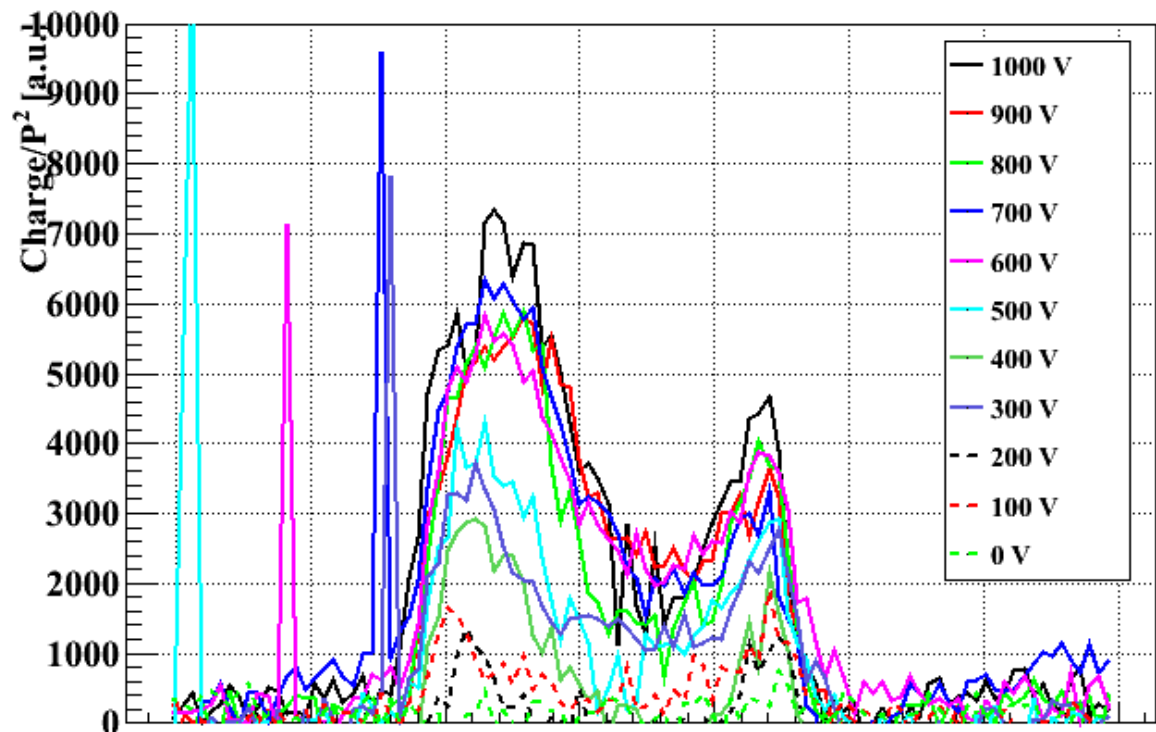
*Two-Photon Photopolymerization and 3D Litographic Microfabrication.* H.B.Sun and S.Kawata. APS (2004) 170 pp 169-273, Springer-Verlag.  
*Femtosecond Laser Litography in Organic and Non-Organic Materials,* F.Jipa et al., Chap.3, Nanotechnology and Nanomaterials, "Updates in Advanced Litography", ed. by S.Hosaka, INTECH, 2013.



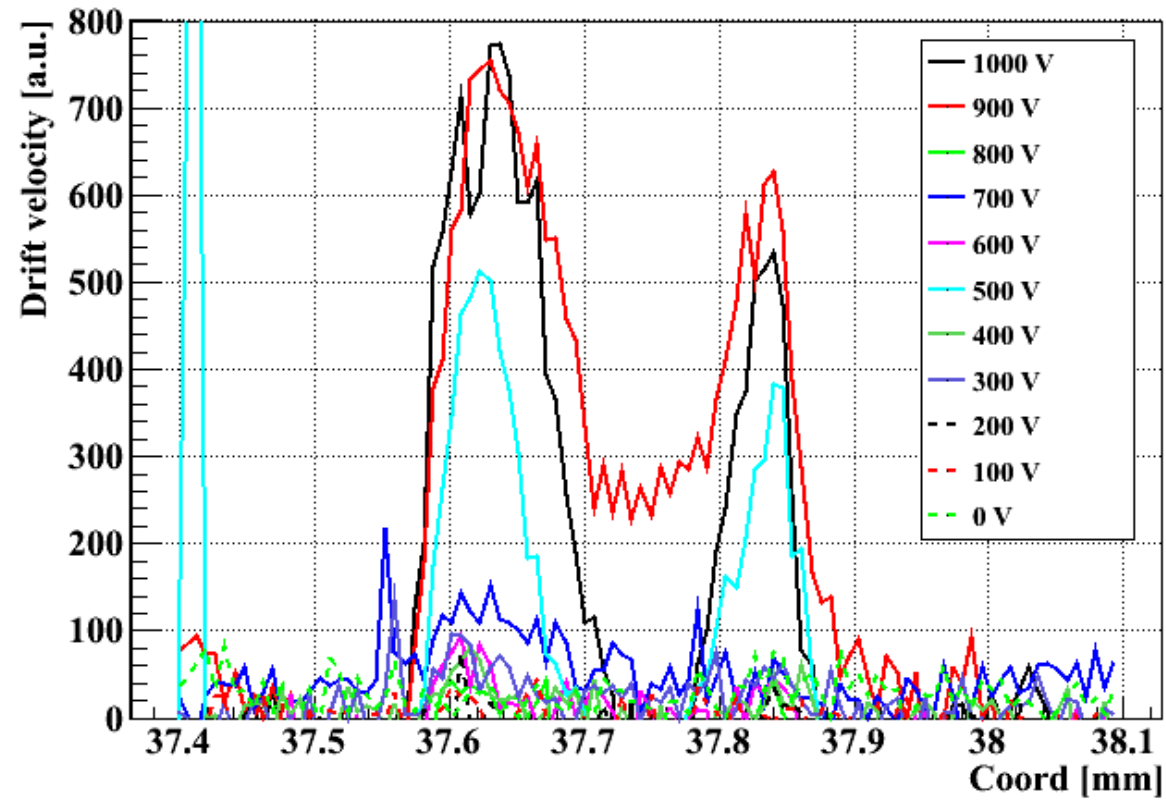


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