

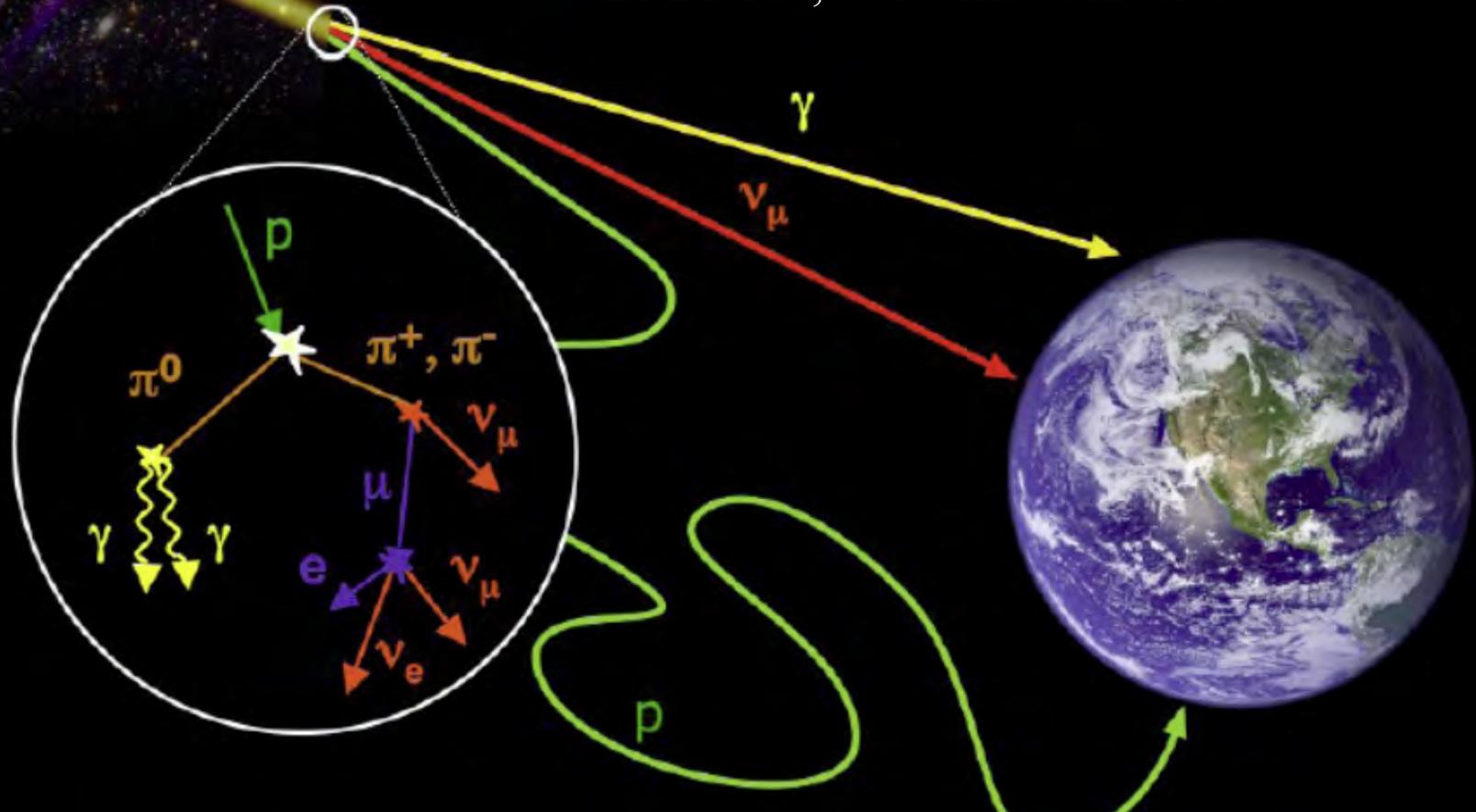
# Cosmogenic photon and neutrino fluxes in Auger era

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Astrophysical  
beam dump

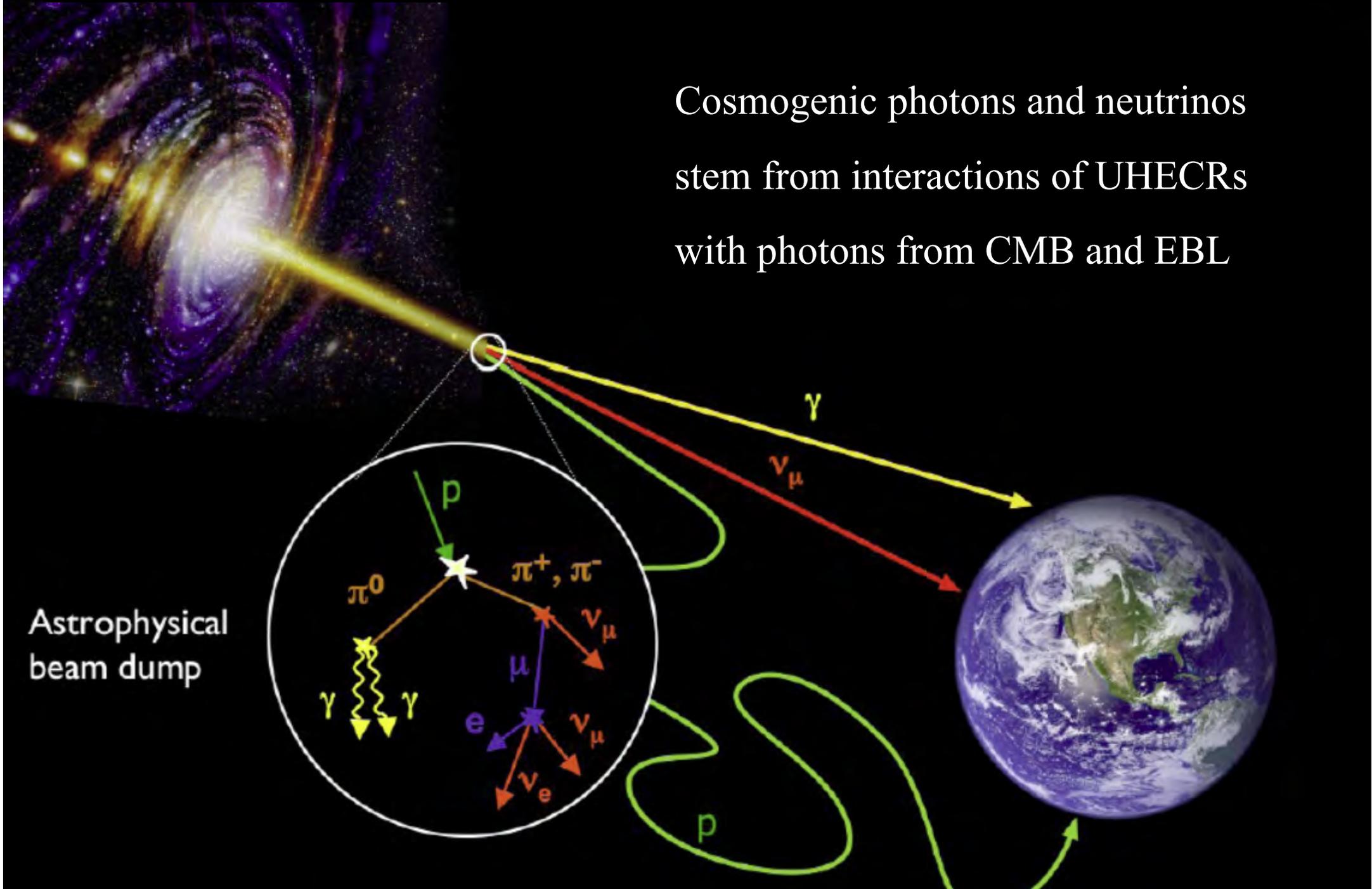


# Outline

- Propagation of UHECRs, gamma rays and neutrinos
- Idea of this work
- Pierre Auger spectrum and composition measurements
- Propagation models and simulations
- Results of the fit
- Cosmogenic photons and neutrinos
- Conclusions

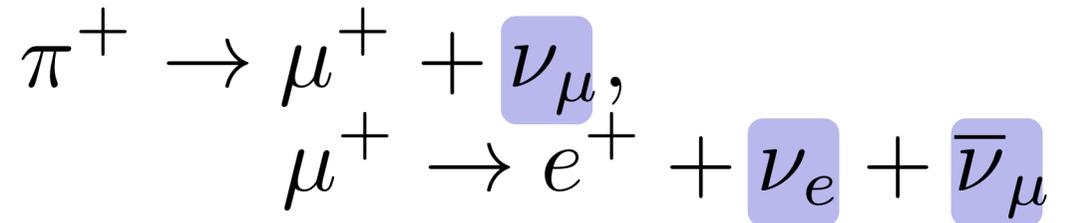
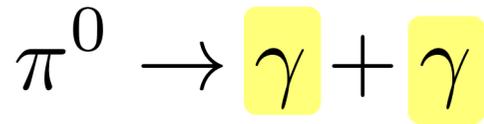
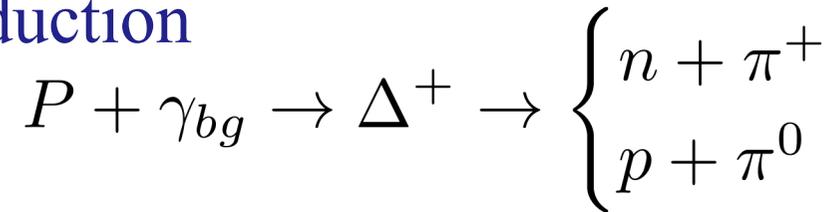
# Propagation of UHECRs, gamma rays and neutrinos

Cosmogenic photons and neutrinos stem from interactions of UHECRs with photons from CMB and EBL



# Propagation of UHECRs, gamma rays and neutrinos and

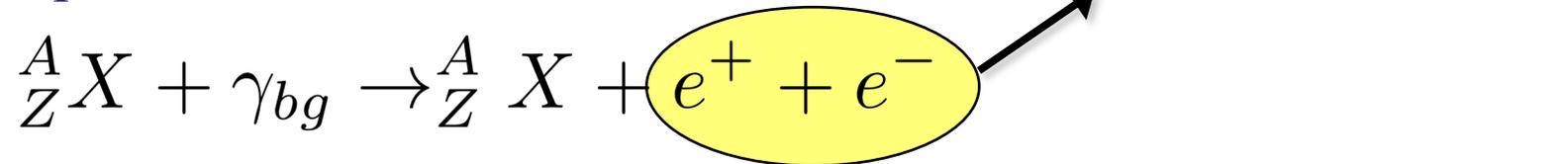
## Photopion production



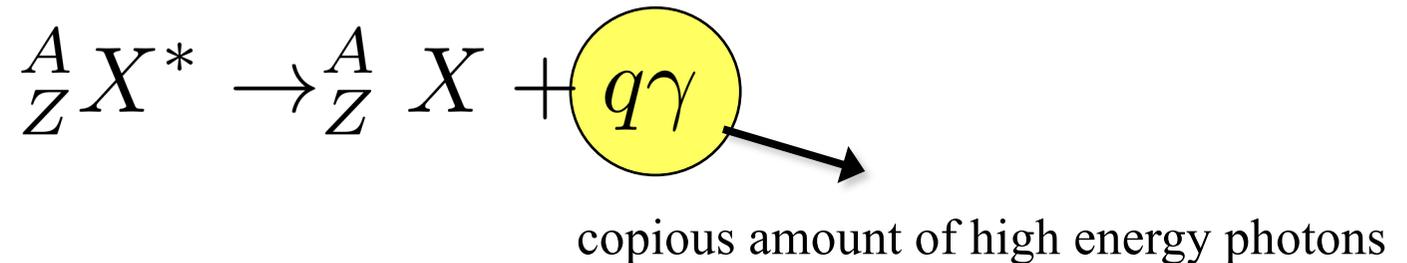
$Z > 1$ :

- Photopion : superposition of the corresponding rates for nucleons

- Bethe-Heitler pair production



- Photodisintegration: radiative decays of excited states produced in the photodisin. chain

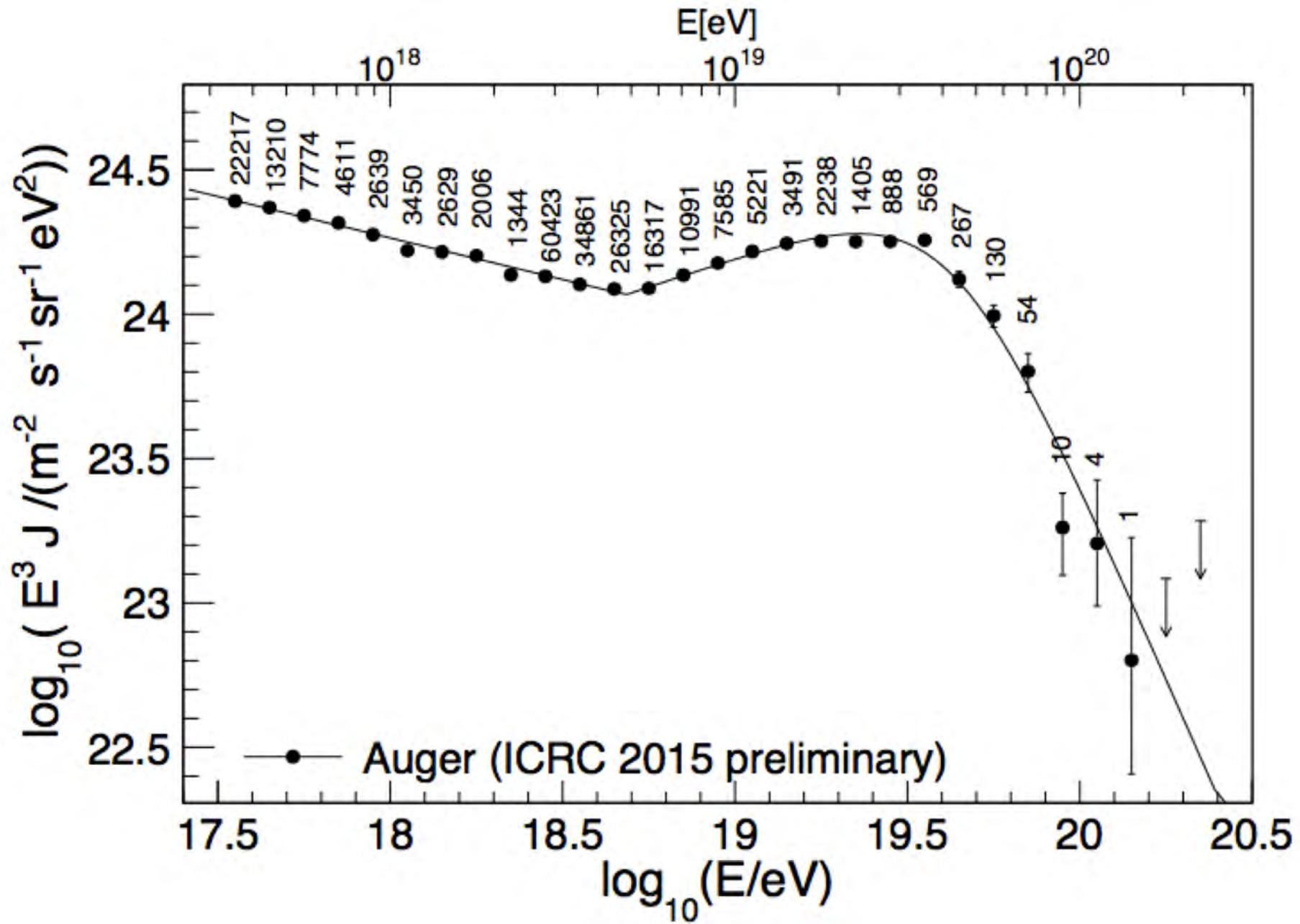


## Idea of this work

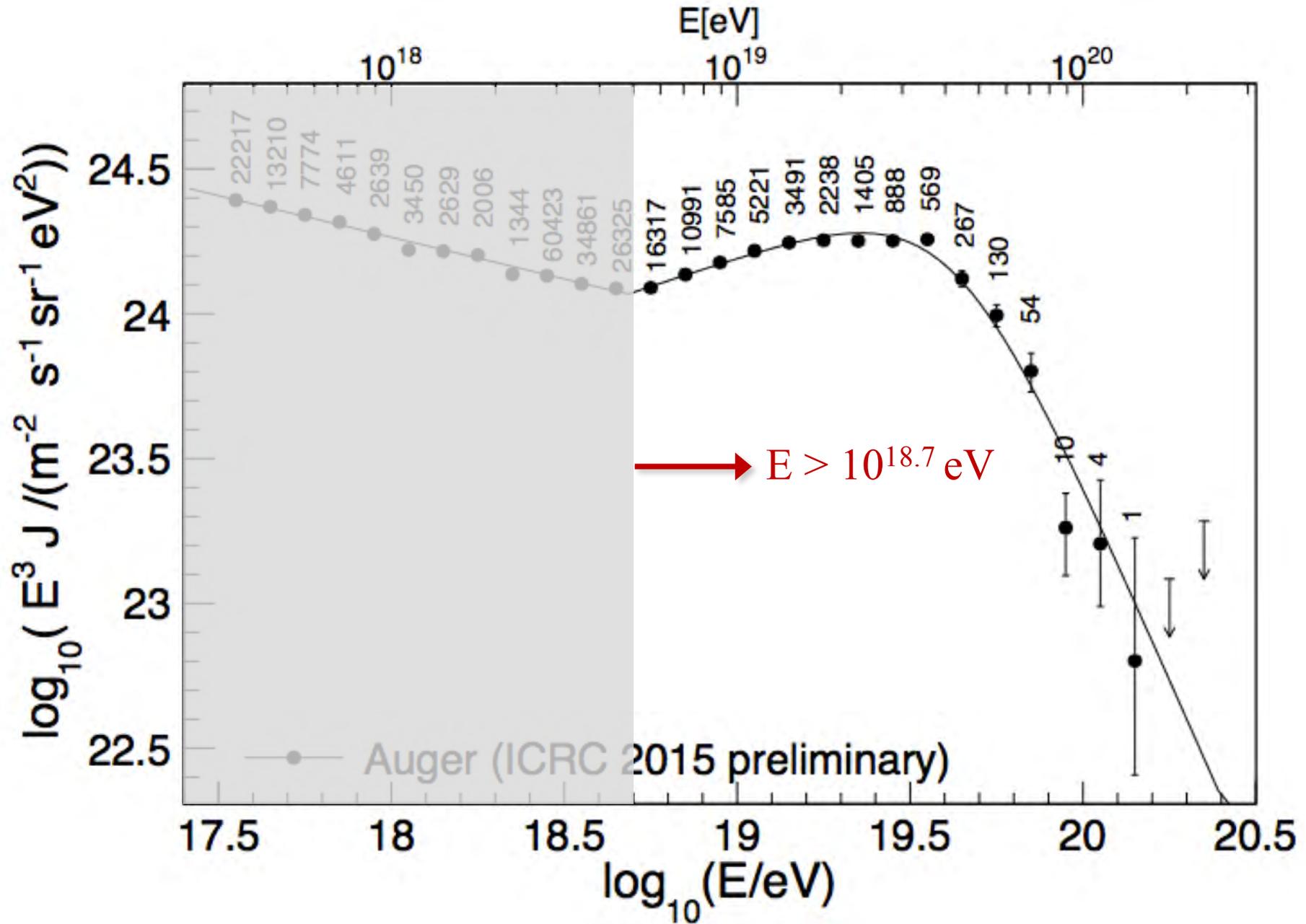
Published in JCAP 01 (2019) 002

- Perform a fit of the spectrum and composition measured by the Pierre Auger Observatory for four source emissivity scenarios
- Calculate the associated fluxes of neutrino and photons
- Prospects for the future generation of high-energy neutrino and gamma-ray observatories to constrain the sources of UHECRs

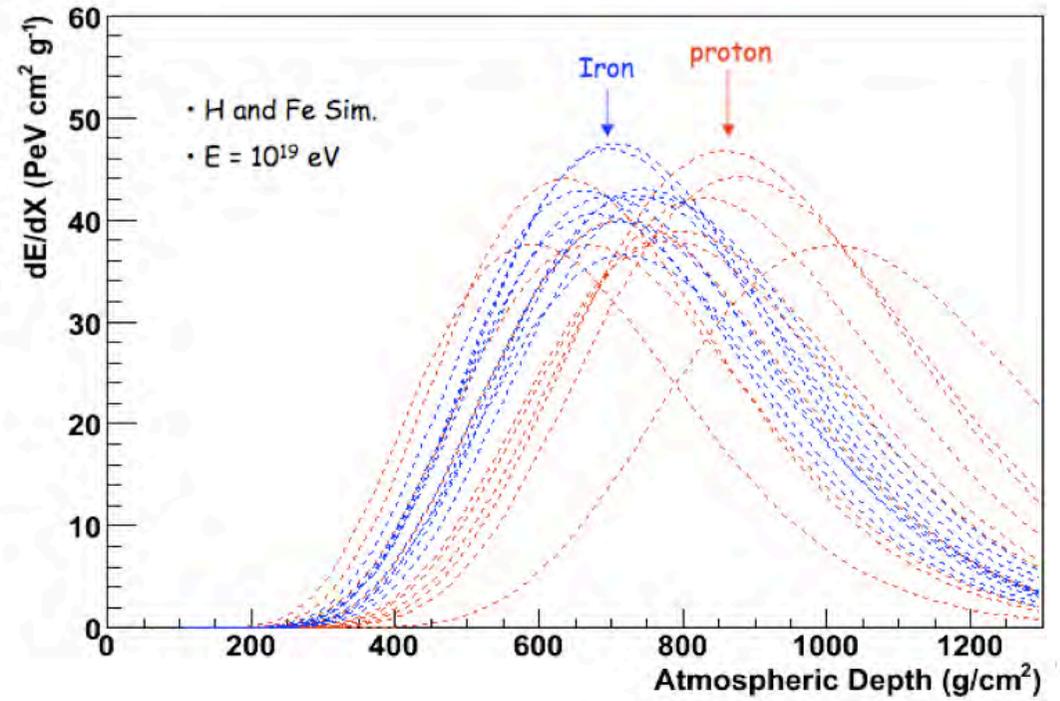
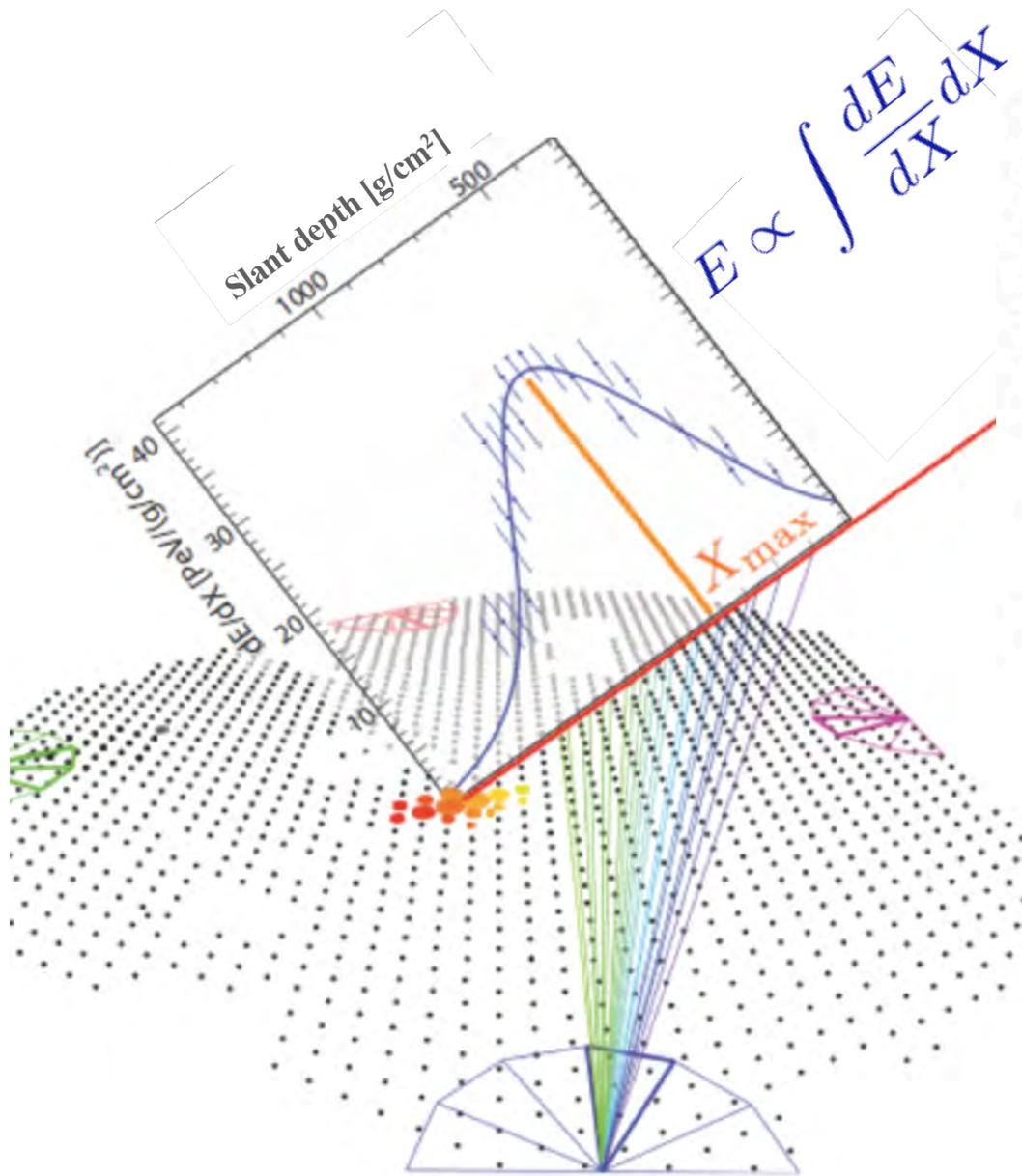
# Auger Combined Energy Spectrum



# Auger Combined Energy Spectrum



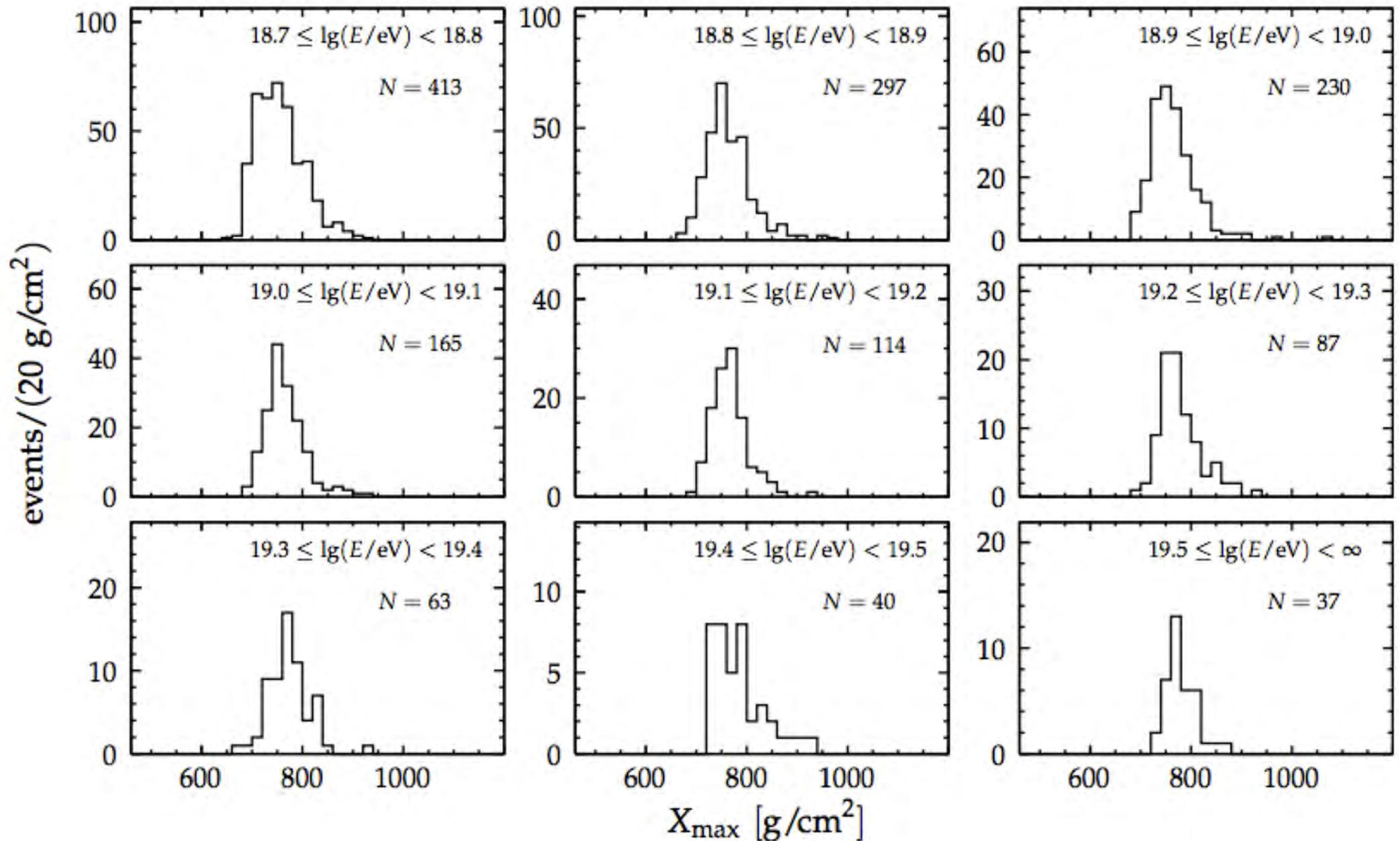
# Auger $X_{\max}$ distributions



For a given hadronic interaction model,  $X_{\max}$  is correlated with the mass of the incident cosmic ray particle

# Auger $X_{\max}$ distributions

Phys.Rev. D90 (2014) 122005



# Propagation models and simulations

We fit the Pierre Auger data on UHECR spectrum and composition to simple astrophysical scenarios:

- Distribution of sources whose comoving emissivity evolves with redshift

$$\dot{\epsilon}(z) = \dot{\epsilon}_0(1+z)^m$$

- Three other particular cases for source evolution: AGN, GRB and SFR
- Injection consisting only of H, He, N, Si, Fe nuclei
- Power-law spectrum with a rigidity-dependent broken exponential cutoff

$$J(E) = J_0 \sum_j f_j E^{-\alpha} \begin{cases} 1 & \text{if } Z_j R_{\max} > E \\ \exp\left(1 - \frac{E}{Z R_{\max}}\right) & \text{if } Z_j R_{\max} \leq E \end{cases}$$

**Free parameters ( $J_0$ ,  $\alpha$ ,  $R_{\max}$ ,  $f_H$ ,  $f_{He}$ ,  $f_N$ ,  $f_{Si}$ ) besides  $m$  in the first scenario**

# Propagation models and simulations

R. A. Batista et al. JCAP 05 (2016) 038

Simulations were performed by using the CRPropa 3 code

- Propagation strongly sensitive to Photodisintegration cross-sections and Extragalactic background light spectrum
- We adopt EBL model by Gilmore et al. and TALYS 1.8 photonuclear cross section
- The mass composition of the simulated arriving particles are inferred using the Gumbel parametrization

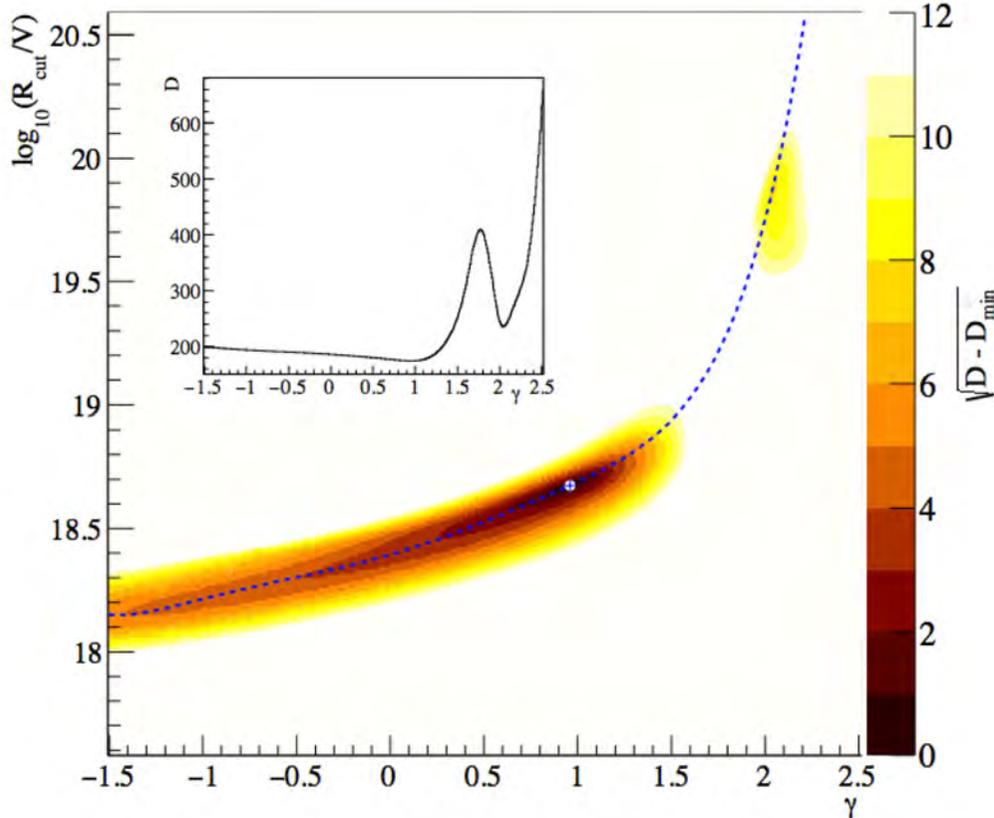
$$p(X_{\max}|\mu, \sigma, \lambda) = \frac{\lambda^\lambda \exp(-\lambda z - \lambda \exp(-z))}{\sigma \Gamma(\lambda)}$$

$$z = \frac{X_{\max} - \mu}{\sigma}$$

# Auger Combined fit (no source evolution $m = 0$ )

Auger Collaboration, JCAP04(2017)038

SPG

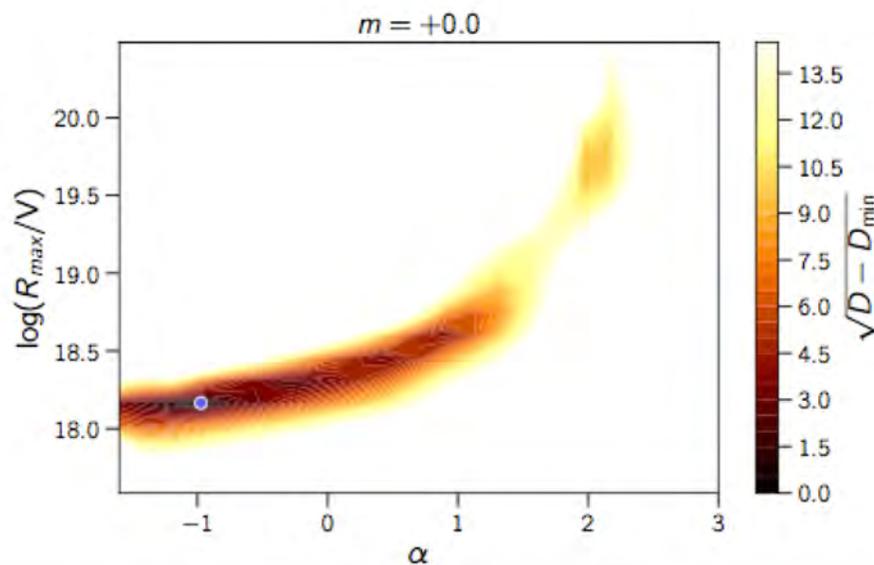
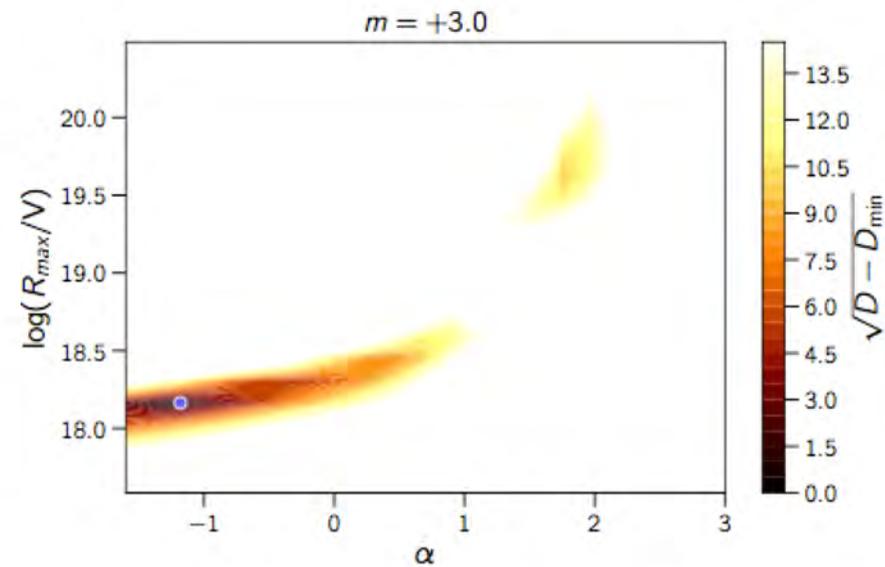
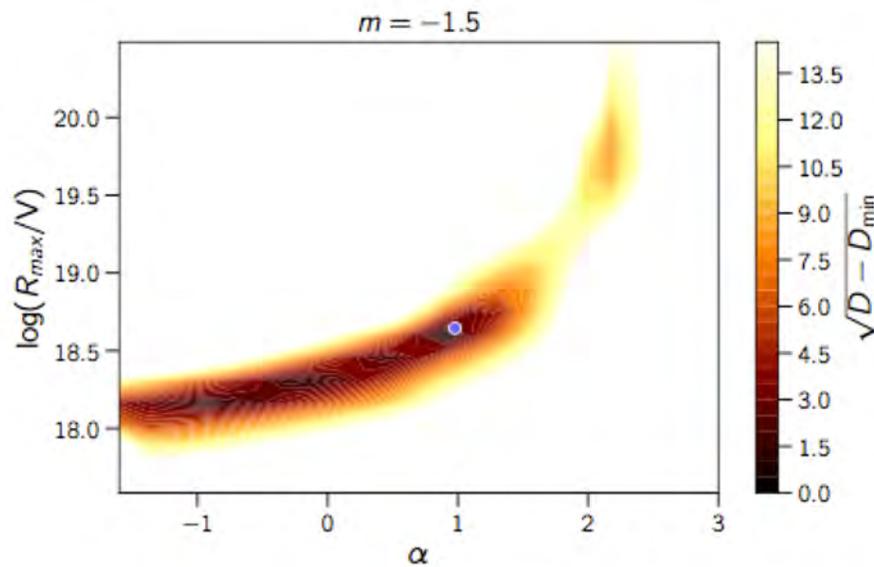


	MC code	$\sigma_{\text{photo}}$	
SPG	SimProp	PSB	2
STG	SimProp	TALYS	Gilmore 2012
SPD	SimProp	PSB	Domínguez 2011
CTG	CRPropa	TALYS	Gilmore 2012
CTD	CRPropa	TALYS	Domínguez 2011
CGD	CRPropa	Geant4	Domínguez 2011

model	$\gamma$	$\log_{10}(R_{\text{cut}}/V)$	$D$	$D(J)$	$D(X_{\text{max}})$
SPG	$+0.96^{+0.08}_{-0.13}$	$18.68^{+0.02}_{-0.04}$	174.3	13.2	161.1
STG	$+0.77^{+0.07}_{-0.13}$	$18.62^{+0.02}_{-0.04}$	175.9	18.8	157.1
SPD	$-1.02^{+0.31}_{-0.26}$	$18.19^{+0.04}_{-0.03}$	187.0	8.4	178.6
CTG	$-1.03^{+0.35}_{-0.30}$	$18.21^{+0.05}_{-0.04}$	189.7	8.3	181.4
	$+0.87^{+0.08}_{-0.06}$	$18.62 \pm 0.02$	191.9	29.2	162.7
CTD	$-1.47^{+0.28}_{*}$	$18.15^{+0.03}_{-0.01}$	187.3	8.8	178.5
CGD	$-1.01^{+0.26}_{-0.28}$	$18.21 \pm 0.03$	179.5	7.9	171.6

# Results of the fit

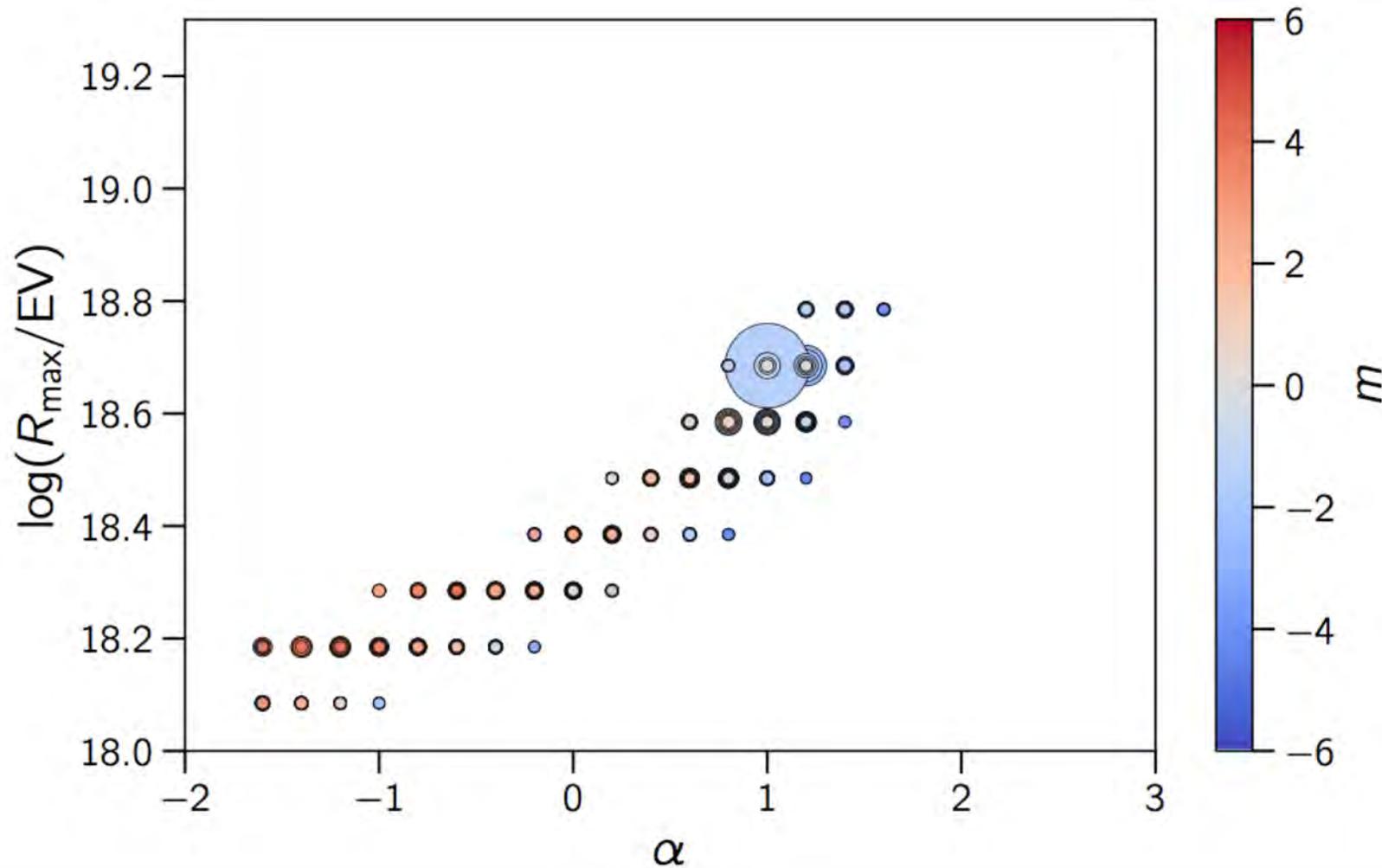
We fix the source evolution in order to obtain the best-fit spectral index ( $\alpha$ ) and maximal rigidity ( $R_{\max}$ )



→ It suggests a trend that if source evolution is accounted in the fit,  $\alpha$  becomes increasingly larger for negative values of  $m$

# Results of the fit

- Best-fit values at 99% confidence level
- Marker sizes with radii inversely proportional to the deviance



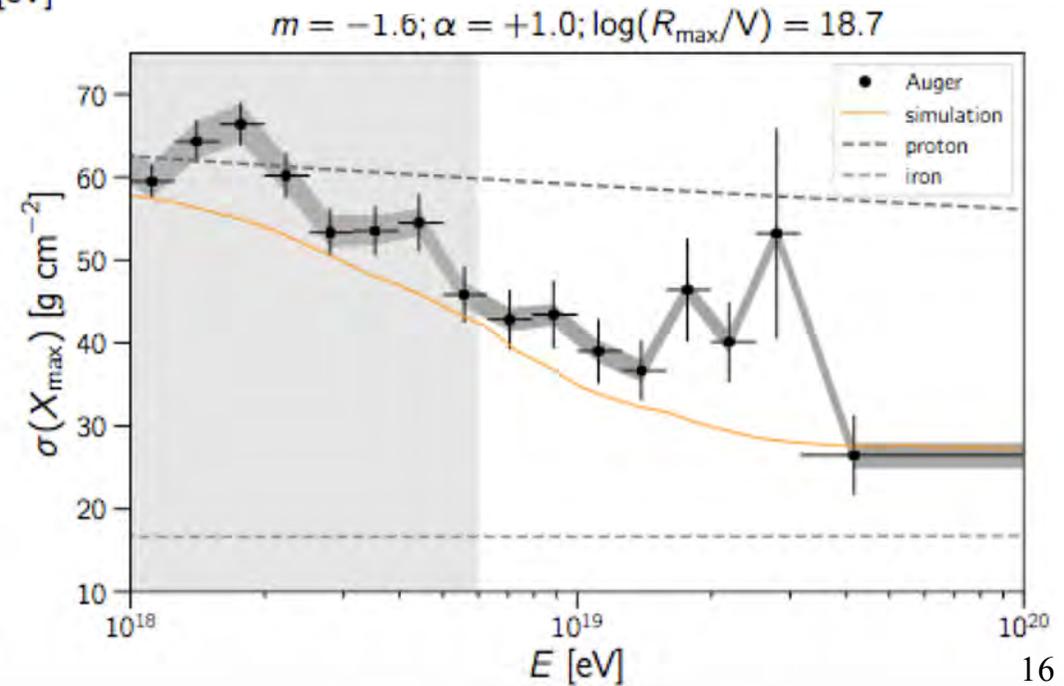
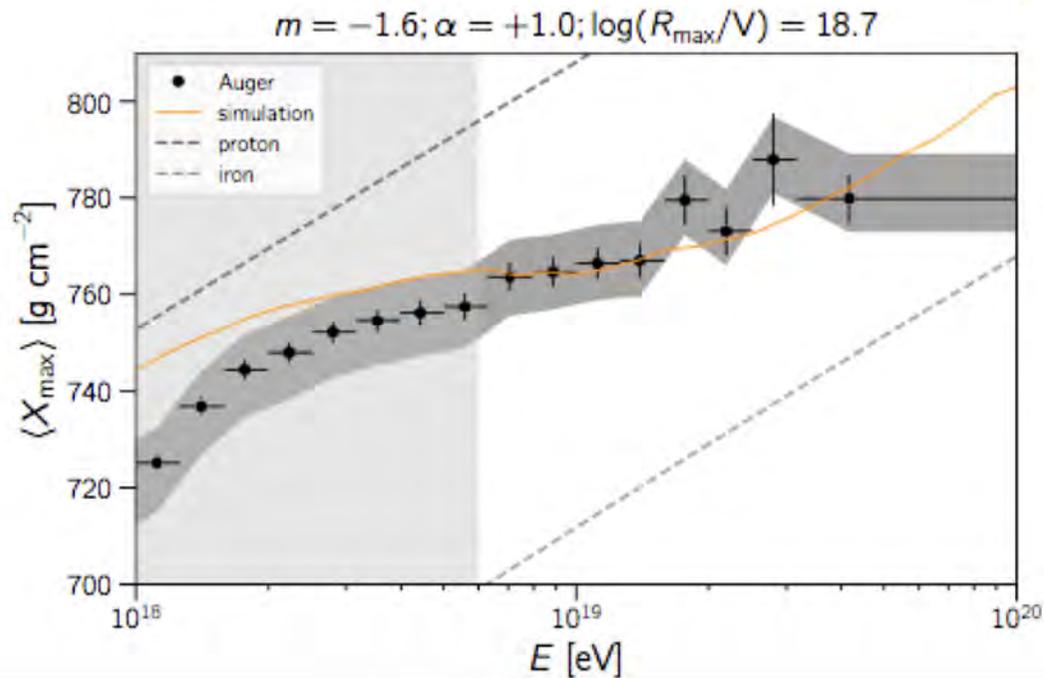
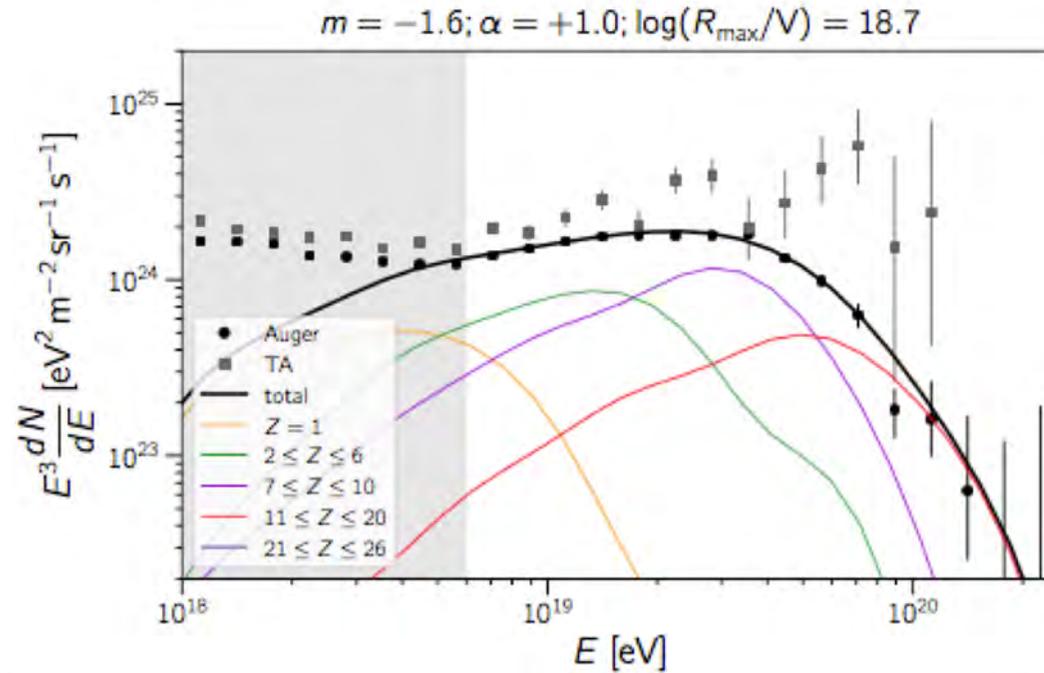
For the most common spectral indices found in the literature ( $1 \lesssim \alpha \lesssim 2$ ) scenarios with positive source evolution ( $m > 0$ ) are disfavoured

# Best fit and Confidence Level bands

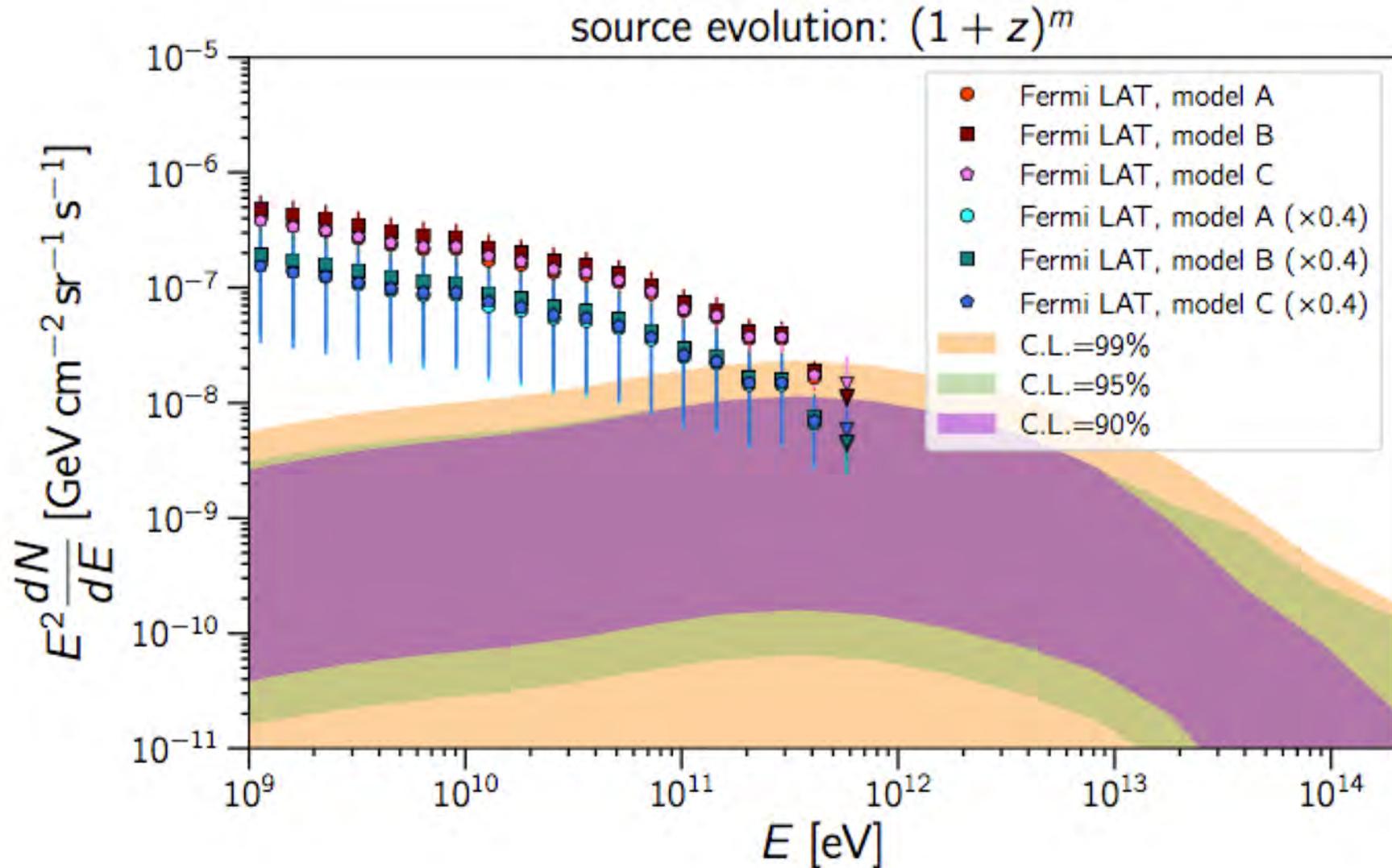
$m$	$\alpha$	$\log(\frac{R_{\max}}{V})$	$f_H$	$f_{He}$	$f_N$	$f_{Si}$	$f_{Fe}$	$\chi^2_{\text{comp}}$	$\chi^2_{\text{spec}}$	$\chi^2$
-1.6	+1.0	18.7	0.0003	0.0101	0.8906	0.0990	0.0	161.1	17.6	178.7
SFR	-1.3	18.2	0.1628	0.8046	0.0309	0.0018	0.0	184.4	19.9	204.3
AGN	-1.0	18.2	0.8716	0.0778	0.0469	0.0038	0.0	224.9	33.1	258.0
GRB	-1.5	18.2	0.5876	0.3973	0.0147	0.0004	0.0	170.8	20.1	190.9

C.L.	$D$	parameter range		
90%	< 2.71	$-4.3 \leq m \leq -0.7$	$+0.8 \leq \alpha \leq +1.2$	$18.6 \leq \log(R_{\max}/V) \leq 18.7$
95%	< 3.84	$-5.5 \leq m \leq +3.5$	$-1.6 \leq \alpha \leq +1.4$	$18.2 \leq \log(R_{\max}/V) \leq 18.7$
99%	< 6.63	$-6.0 \leq m \leq +4.2$	$-1.6 \leq \alpha \leq +1.6$	$18.1 \leq \log(R_{\max}/V) \leq 18.8$
90%	< 2.71	SFR	$-1.3 \leq \alpha \leq -1.2$	$18.2 \leq \log(R_{\max}/V) \leq 18.2$
95%	< 3.84	SFR	$-1.4 \leq \alpha \leq -0.6$	$18.2 \leq \log(R_{\max}/V) \leq 18.3$
99%	< 6.63	SFR	$-1.4 \leq \alpha \leq +0.1$	$18.2 \leq \log(R_{\max}/V) \leq 18.4$
90%	< 2.71	AGN	$-1.0 \leq \alpha \leq -0.9$	$18.2 \leq \log(R_{\max}/V) \leq 18.2$
95%	< 3.84	AGN	$-1.0 \leq \alpha \leq -0.9$	$18.2 \leq \log(R_{\max}/V) \leq 18.2$
99%	< 6.63	AGN	$-1.0 \leq \alpha \leq -0.4$	$18.2 \leq \log(R_{\max}/V) \leq 18.3$
90%	< 2.71	GRB	$-1.5 \leq \alpha \leq -1.4$	$18.2 \leq \log(R_{\max}/V) \leq 18.2$
95%	< 3.84	GRB	$-1.5 \leq \alpha \leq -0.7$	$18.2 \leq \log(R_{\max}/V) \leq 18.3$
99%	< 6.63	GRB	$-1.5 \leq \alpha \leq +0.4$	$18.1 \leq \log(R_{\max}/V) \leq 18.5$

# Best fit for $(1+z)^m$ evolution scenario

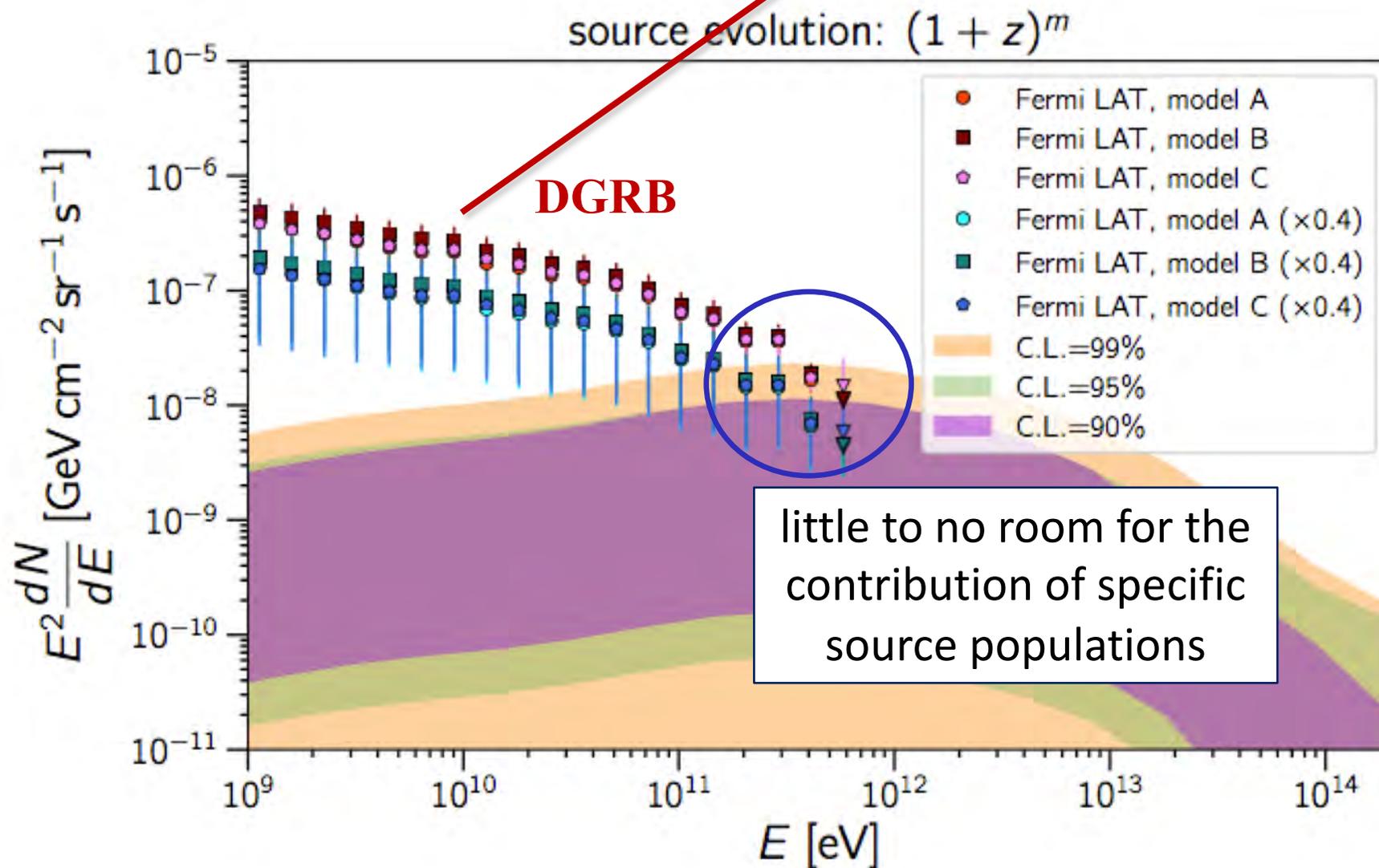


# Cosmogenic photons (coming from $z < 1$ )

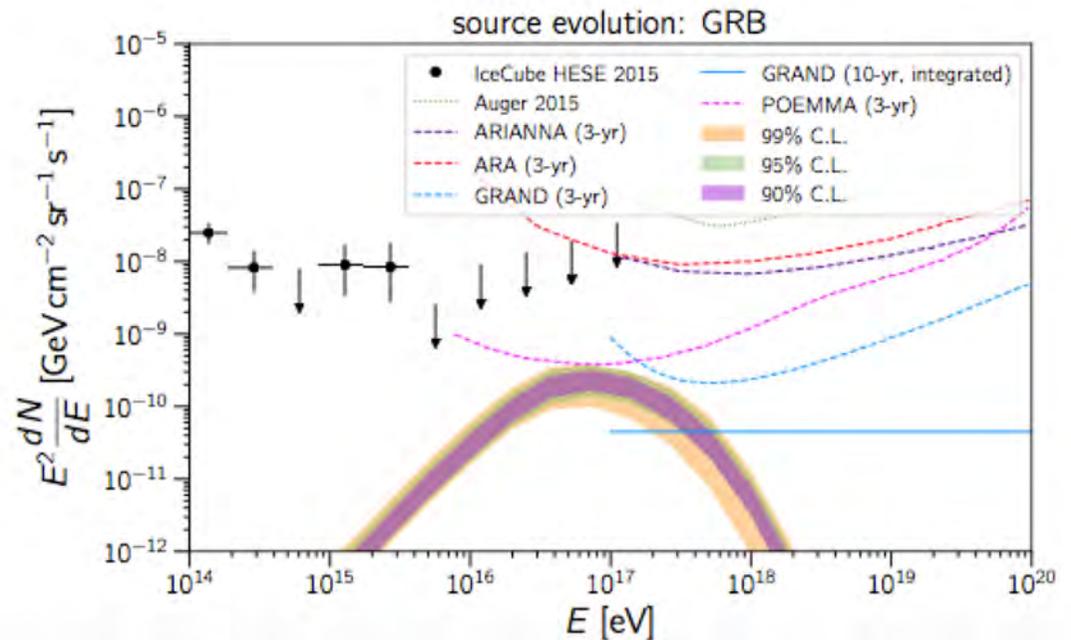
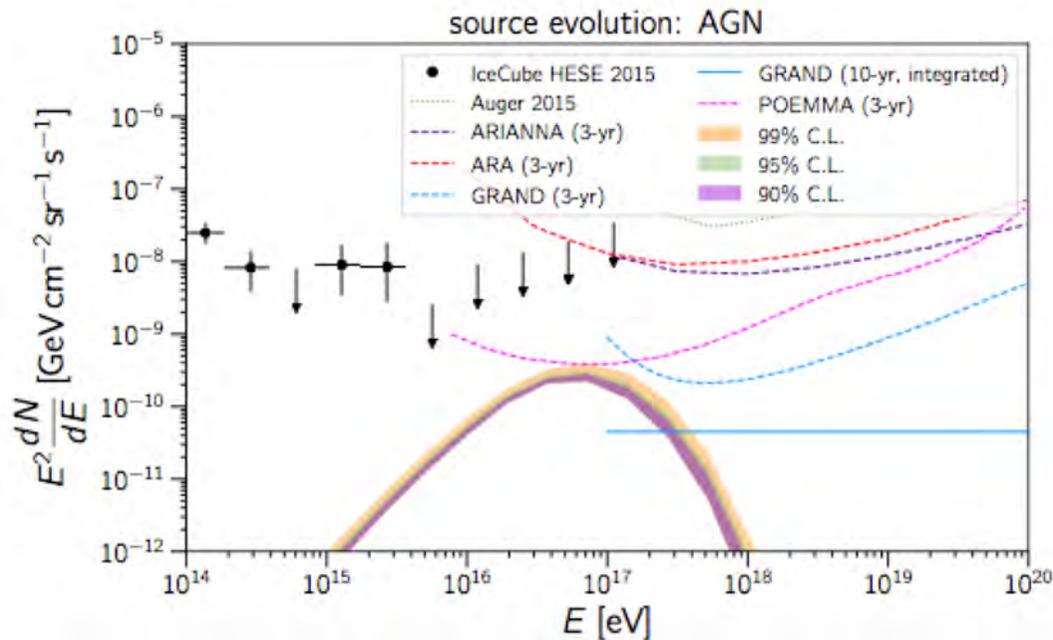
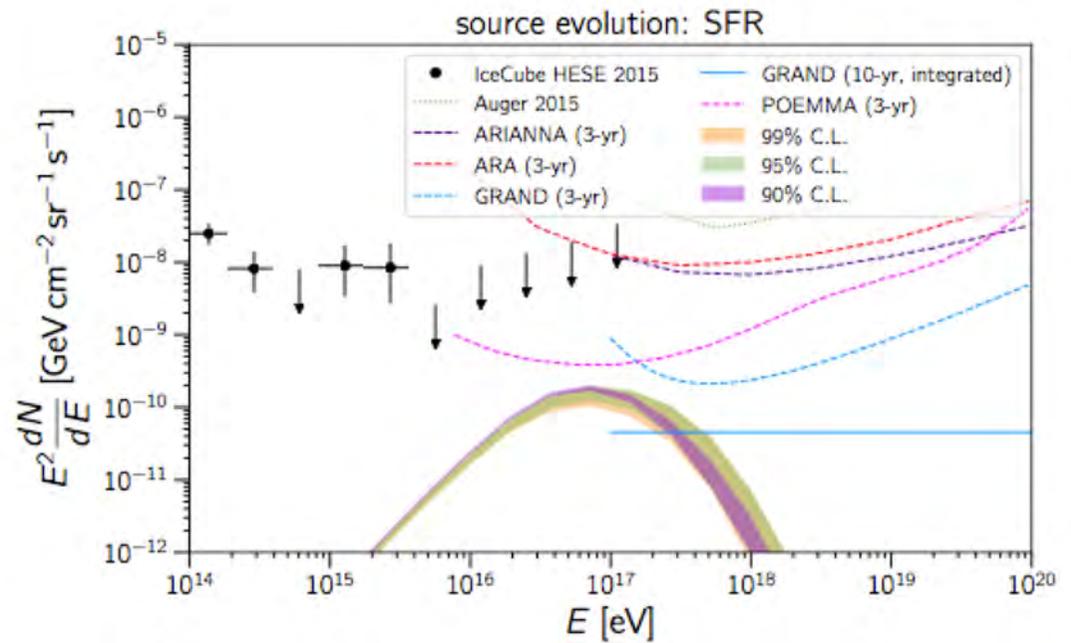
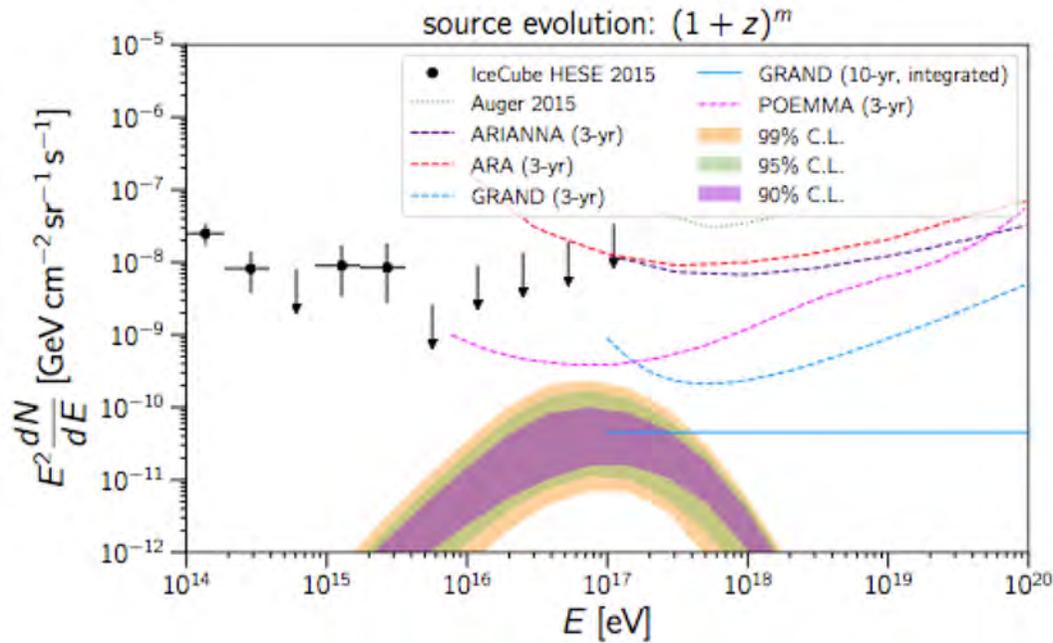


# Cosmogenic photons

Contribution from cosmogenic neutrinos + blazars, misaligned AGNs, quasar outflows, star-forming regions, decay of very and super heavy dark matter ...



# Cosmogenic neutrinos (coming from $z < 1$ )



# Conclusions

- Best fit for  $(1 + z)^m$  scenario is obtained for  $\alpha \simeq 1$ , for compositions at injection dominated by intermediate-mass nuclei (nitrogen and silicon groups). The value of  $m$  that minimises the deviance is  $m = -1.6$
- Evolutions with  $m < 0.7$  are favoured at a 90% C.L. (sources or local distribution?)
- Our best fits for  $(1 + z)^m$  with the hardest spectral indices and  $z < 1$  condition provide the lowest possible cosmogenic fluxes compatible with the observed data → **pessimistic scenario**
- Photon flux levels nearly exceed DGRB Fermi measurements → little room for contribution of specific source population; disfavour models that predict a significant contribution of dark matter decay and annihilation in this energy range.

# Conclusions

- Our cosmogenic neutrino fluxes are rather low, due to the choice of  $z < 1$ . Nevertheless, even in this pessimistic case, GRAND would be able to detect cosmogenic neutrinos for sources evolving as SFR, GRB, and AGN, at a 99% C.L., within the first 5 years of operation
- If sources evolve as  $(1 + z)^m$ , the detection of cosmogenic neutrinos would be possible but not certain
- From a different perspective, low cosmogenic fluxes could be profitable for EeV neutrino astronomy. Such a scenario would imply that the neutrinos first detected by future experiments would likely be those produced directly at the sources, via interactions of UHECRs with photon and baryon fields in the source environment.