

WW Production and Anomalous Gauge Couplings

Ian Lewis

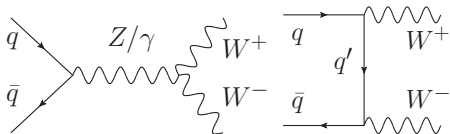
University of Kansas

Based on J. Baglio, S. Dawson, *IL Phys. Rev. D*96 (2017) 073003

May 14, 2018

HL/HE-LHC Electroweak Physics Meeting
CERN

W^+W^- production



- Informative to focus on one process.
 - Of particular interest is the electroweak sector.
 - Focus on W^+W^- production at the LHC.
 - Sensitive to anomalous trilinear gauge boson couplings (ATGCs)

W^+W^- production

- Anomalous couplings language Hagiwara, Peccei, Zeppenfeld, Hikasa NPB482 (1987):

$$\delta\mathcal{L} = -ig_{WWV} \left(g_1^V (W_{\mu\nu}^+ W^{-\mu} V^\nu - W_{\mu\nu}^- W^{+\mu} V^\nu) + \kappa^V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda^V}{M_W^2} W_{\rho\mu}^+ W^{-\mu}{}_\nu V^{\nu\rho} \right)$$

- $V = Z, \gamma$
- $g_{WWZ} = g \cos\theta_w, \quad g_{WW\gamma} = e$
- Parameterize deviations from SM:

$$g_1^Z = 1 + \delta g_1^Z \quad g_1^\gamma = 1 + \delta g_1^\gamma \quad \kappa^Z = 1 + \delta\kappa^Z \quad \kappa^\gamma = 1 + \delta\kappa^\gamma$$

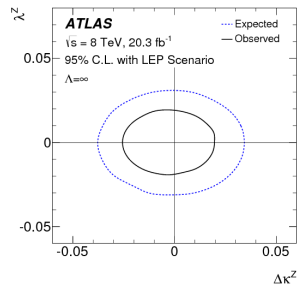
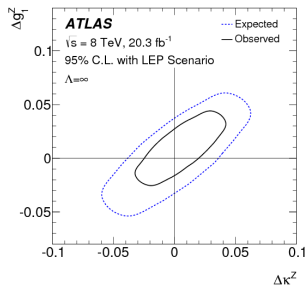
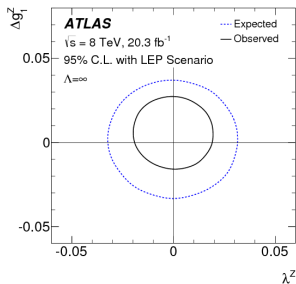
- $\lambda^Z = 0$ and $\lambda^\gamma = 0$ in SM.
- $SU(2)_L$ implies:

$$\delta g_1^\gamma = 0 \quad \lambda^\gamma = \lambda^Z \quad \delta\kappa^\gamma = \frac{\cos^2\theta_w}{\sin^2\theta_w} (\delta g_1^Z - \delta\kappa^Z)$$

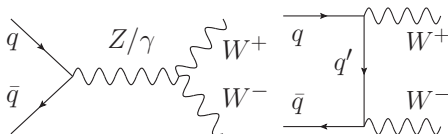
- Three independent parameters: $\lambda^Z, \delta g_1^Z, \delta\kappa^Z$

Experimental results

- ATGCs actively being searched for in W^+W^- production by both ATLAS [JHEP 1609 \(2016\) 029](#) and CMS [Phys.Lett. B772 \(2017\) 21](#)



Missing Terms



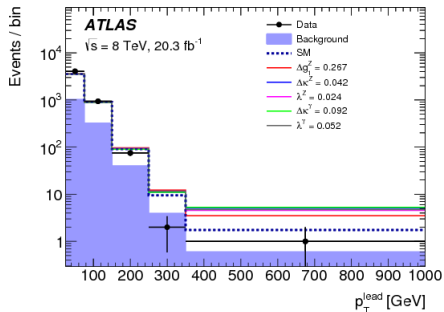
- Have not included anomalous quark gauge boson couplings.
 - Highly constrained by LEP.
 - But SM contains cancellations to unitarize amplitudes: growth with energy cancels.
 - Anomalous quark couplings can spoil cancellation and have growth with energy.
 - Garnered more attention recently [Zhang PRL118 \(2017\) 011803](#); [J. Baglio, S. Dawson I.M. Lewis, Phys. Rev. D96 \(2017\) 073003](#); [Eboli Phenomenology Symposium 2018](#)
- Parameterize via anomalous couplings:

$$\begin{aligned} \mathcal{L} = & g_Z Z_\mu \bar{q} \gamma^\mu \left\{ \left[T_3 - \sin_W^2 Q_q + \delta g_L^{Zq} \right] P_L + \left[-\sin_W^2 Q_q + \delta g_R^{Zq} \right] P_R \right\} q \\ & + \frac{g}{\sqrt{2}} \left\{ W_\mu^+ (1 + \delta g_L^W) \bar{u} \gamma^\mu P_L d + \text{hc.} \right\} \end{aligned}$$

- $SU(2)$ invariance implies $\delta g_L^W = \delta g_L^{Zu} - \delta g_L^{Zd}$.

Refit Experimental results

- ATGCs limits from ATLAS [JHEP 1609](#).
- In practice want to take differential distributions from experimental collaborations, extract constraints on anomalous couplings.
- We do not decay the W^+ .



Refit Experimental Results

- Assume strongest constraint comes from last bin.
- Scan over allowed ATGCs and determine allowed

$$\sigma(p_T^{W^+} > 500 \text{ GeV}) = \int_{500 \text{ GeV}}^{\infty} dp_T^{W^+} \frac{d\sigma}{dp_T^{W^+}}$$

- Now scan over all parameters and determine allowed regions taking into consideration LEP constraints on anomalous quark couplings [Falkowski, Riva JHEP 1502](#):

$$\delta g_L^{Zd} = (2.3 \pm 1) \times 10^{-3}$$

$$\delta g_L^{Zu} = (-2.6 \pm 1.6) \times 10^{-3}$$

$$\delta g_R^{Zd} = (16.0 \pm 5.2) \times 10^{-3}$$

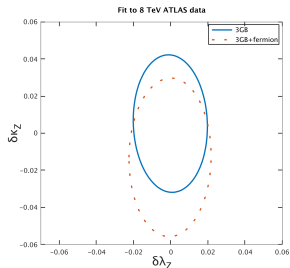
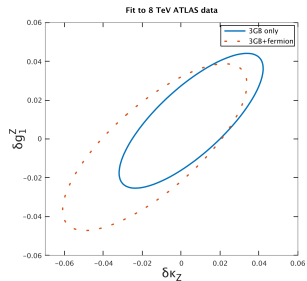
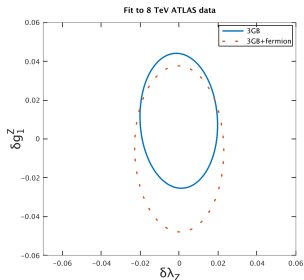
$$\delta g_R^{Zu} = (-3.6 \pm 3.5) \times 10^{-3}$$

- Accept points that fall within allowed region of $\sigma(p_T^{W^+} > 500 \text{ GeV})$.
- Have verified we can reconstruct ATLAS 2-D and 1-D fits limits on ATGCs.

Refit

- Blue: Including only ATGCs.
- Red dots: adding in anomalous quark couplings
- Inner regions allowed

Baglio, Dawson, *IL PRD96* (2017) 073003



Comment on Calculating Cross Sections

- Previous bounds found using full amplitude squared.
- Includes terms that go as Λ^{-4} .

$$|\mathcal{A}|^2 \sim \left| g_{SM} + \frac{c_{dim-6}}{\Lambda^2} \right|^2 \sim g_{SM}^2 + g_{SM} \times \frac{c_{dim-6}}{\Lambda^2} + \frac{c_{dim-6}^2}{\Lambda^4}$$

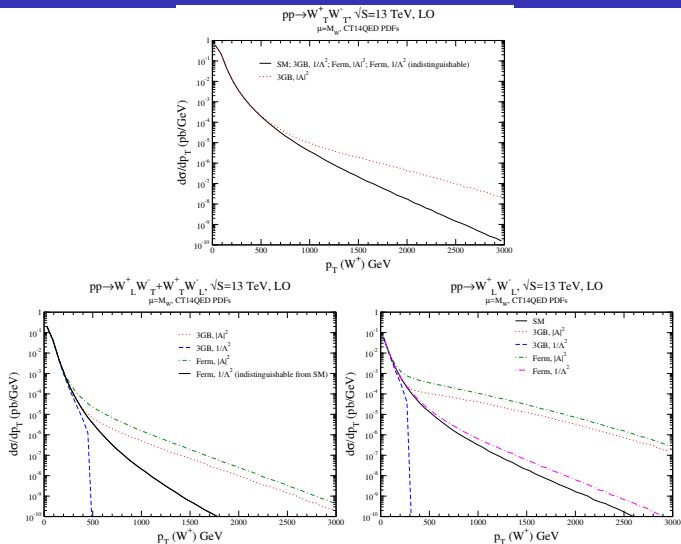
- Same order as dimension-8 contributions:

$$\begin{aligned} |\mathcal{A}|^2 &\sim \left| g_{SM} + \frac{c_{dim-6}}{\Lambda^2} + \frac{c_{dim-8}}{\Lambda^4} \right|^2 \\ &\sim g_{SM}^2 + g_{SM} \times \frac{c_{dim-6}}{\Lambda^2} + \frac{c_{dim-6}^2}{\Lambda^4} + g_{SM} \times \frac{c_{dim-8}}{\Lambda^4} + O(\Lambda^{-6}) \end{aligned}$$

Future Directions

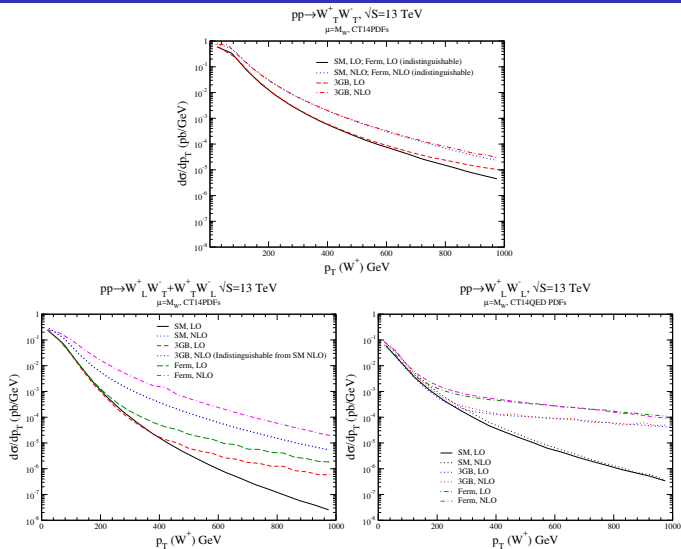
- Fully incorporate W leptonic decays at LO and NLO in QCD.
- Perform realistic collider study at NLO and fit to relevant distributions for HL and HE LHC studies.
- Perform study of importance of different $1/\Lambda^{2n}$ terms in the total cross section.
 - Relevant for transversely polarized gauge bosons. Different polarizations may be more sensitive [Panico, Riva, Wulzer Phys.Lett. B776 \(2018\) 473](#); [Azatov, Elias-Miro, Reyimuaji, Venturini JHEP 1710 \(2017\) 027](#)
 - Relevant for on-shell gauge bosons. Off-shell effects can be relevant for interference between SM and EFT [Helset, Trott JHEP 1804 \(2018\) 038](#).

Differential Distributions by Helicity



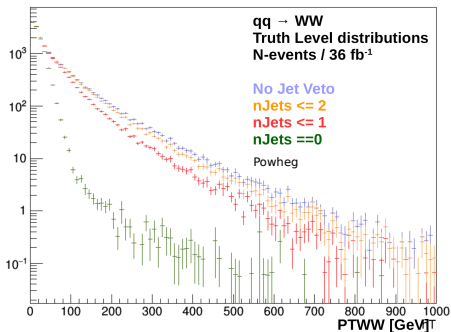
3GB: ATGCs only, **Ferm: Anomalous fermion couplings only** Baglio, Dawson, *IL PRD96* (2017) 073003

NLO QCD by Helicity including all EFT terms



Baglio, Dawson, *IL PRD96* (2017) 073003

Projections for pT-WW



Truth level distribution can be used to extract expected “unfolded” cross section for 3000 fb⁻¹

- Allow large statistical uncertainties in highest pT bin:
→ >10 events (on reco level)
- Other bins: balance out statistics, detector systematics and backgrounds
- NJets==0 with very low acceptance for high pTWW
→ detector uncertainties ~10%

- Crucial question: how will knowledge of detector performance evolve at the HL-LHC with regards to background?
- Dominant uncertainties expected to be: MET and Jet scales as well as Pileup

<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/UpgradePhysicsStudies>

<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/JetEtmisPublicResults#PubPlotsHLLHC>

Contributors: Valerie Lang, Elias Ruettinger, Kristin Lohwasser

Conclusions

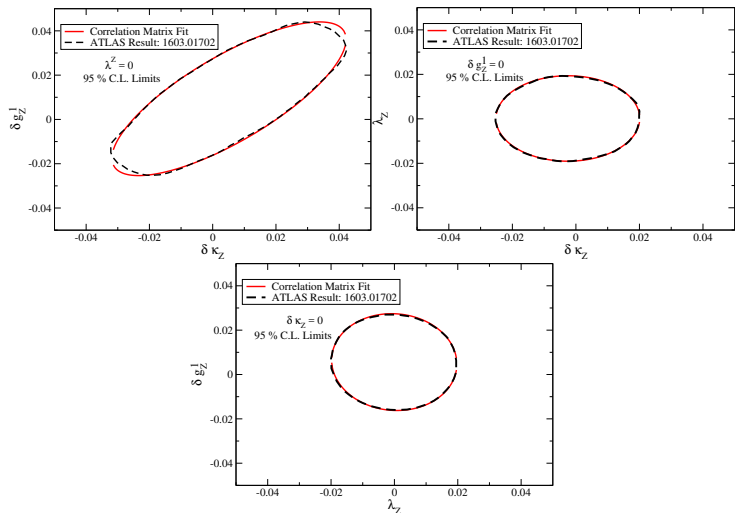
- Investigated the effects of anomalous couplings on W^+W^- production.
 - Although strongly constrained at LEP, anomalous quark-gauge boson couplings significantly change fits to anomalous couplings.
 - LHC is at higher energy, new effects arise and assumptions have to be revisited.
 - Non-interference between SM and EFT is still in effect at NLO.
 - However, interference very dependent on polarizations of Ws.
 - Public code available: WWEFT@NLO

https:

`//quark.phy.bnl.gov/Digital_Data_Archive/dawson/ww_2017/WWEFT_NLO.tar.gz`

Thank You

Verification

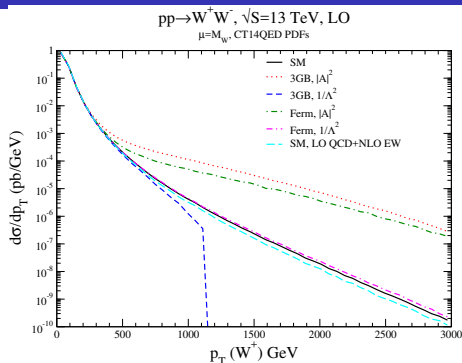


Refit Experimental Results

- Check by comparing to 1D results: set two of the ATGCs to zero:

| | 95% C.L. limit | ATLAS 95% C.L. limit JHEP 1609 |
|-------------------|------------------|--|
| δg_1^Z | [-0.0162,0.0274] | [-0.016,0.027] |
| $\delta \kappa^Z$ | [-0.0252,0.0201] | [-0.025,0.020] |
| λ^Z | [-0.0189,0.0192] | [-0.019,-0.019] |

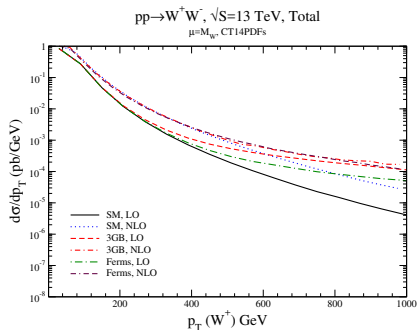
Differential Distributions



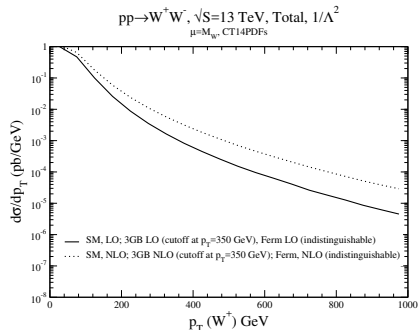
Baglio, Dawson, *IL PRD96* (2017) 073003

- $1/\Lambda^4$ terms dominate in tails and the bounds on anomalous couplings. Falkowski, Gonzalez-Alonso, Greife, Marzocca, *Sci JHEP* 1702 (2017) 115
- Ferm: ATGCs set to zero.
- 3GB: Anomalous fermion couplings set to zero.
- Assuming $C_i \lesssim 1$, anomalous couplings correspond to $\Lambda \gtrsim 2.8$ TeV.

NLO QCD Corrections



EFT Squared



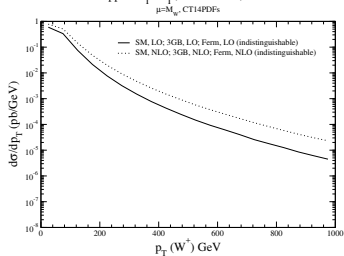
SM+1/Λ²

- “Ferm”: Anomalous trilinear gauge boson couplings set to zero.
- “3GB”: Anomalous quark couplings set to zero.
- 1/Λ⁴ contributions from EFT still dominate in tails.

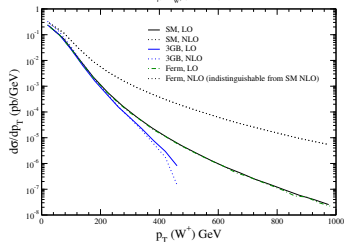
Baglio, Dawson, *IL PRD96* (2017) 073003

NLO QCD by Helicity truncating at $1/\Lambda^2$

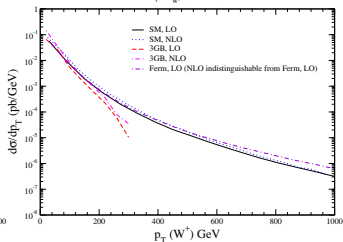
$$pp \rightarrow W_T^+ W_T^-, \sqrt{S}=13 \text{ TeV}, 1/\Lambda^2$$



$$pp \rightarrow W_L^+ W_T^- + W_T^+ W_L^-, \sqrt{S}=13 \text{ TeV}, 1/\Lambda^2$$

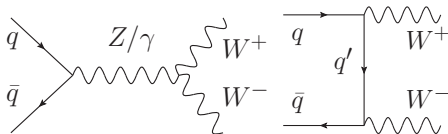


$$pp \rightarrow W_L^+ W_L^-, \sqrt{S}=13 \text{ TeV}, 1/\Lambda^2$$



Baglio, Dawson, IL PRD96 (2017) 073003

W^+W^- production



- Informative to focus on one process.
 - Of particular interest is the electroweak sector.
 - Focus on W^+W^- production at the LHC.
 - Sensitive to anomalous trilinear gauge boson couplings (ATGCs)
- Operators affecting ATGCs:

$$\begin{aligned}
 O_{3W} &= \epsilon^{abc} W_\mu^{a\nu} W_\nu^{b\rho} W_\rho^{c\mu} & O_{HD} &= |\Phi^\dagger D_\mu \Phi|^2 & O_{HWB} &= \Phi^\dagger \sigma^a \Phi W_{\mu\nu}^a B^{\mu\nu} \\
 O_{H\ell}^{(3)} &= i \left(\Phi^\dagger \overleftrightarrow{D}_\mu \sigma^a \Phi \right) \bar{\ell}_L \gamma^\mu \sigma^a \ell_L & O_{ll} &= (\bar{\ell}_L \gamma^\mu \ell_L) (\bar{\ell}_L \gamma_\mu \ell_L)
 \end{aligned}$$

Matching ATGCs in two prescriptions

- Had 5 dimension-6 operators, only three independent combinations.
- In Warsaw basis:

$$\delta g_1^Z = \frac{v^2}{\Lambda^2} \frac{1}{\cos^2 \theta_W - \sin^2 \theta_W} \left(\frac{\sin \theta_W}{\cos \theta_W} C_{HWB} + \frac{1}{4} C_{HD} + \delta v \right)$$

$$\delta \kappa^Z = \frac{v^2}{\Lambda^2} \frac{1}{\cos^2 \theta_W - \sin^2 \theta_W} \left(2 \sin \theta_W \cos \theta_W C_{HWB} + \frac{1}{4} C_{HD} + \delta v \right)$$

$$\delta \lambda^Z = \frac{v}{\Lambda^2} 3M_W C_{3W}$$

- Anomalous coupling language generic enough that any basis can be matched onto it.

W^+W^- production

- Operators affecting ATGCs:

$$\begin{aligned}
 O_{3W} &= \epsilon^{abc} W_\mu^{a\nu} W_\nu^{b\rho} W_\rho^{c\mu} & O_{HD} &= |\Phi^\dagger D_\mu \Phi|^2 & O_{HWB} &= \Phi^\dagger \sigma^a \Phi W_{\mu\nu}^a B^{\mu\nu} \\
 O_{H\ell}^{(3)} &= i \left(\Phi^\dagger \overleftrightarrow{D}_\mu \sigma^a \Phi \right) \bar{\ell}_L \gamma^\mu \sigma^a \ell_L & O_{ll} &= (\bar{\ell}_L \gamma^\mu \ell_L) (\bar{\ell}_L \gamma_\mu \ell_L)
 \end{aligned}$$

- In the EW sector have to choose input parameters: G_F, M_W, M_Z
- EFT alters relationships between other parameters and input parameters:

$$g_Z \rightarrow g_Z + \delta g_Z \quad v \rightarrow v(1 + \delta v) \quad s_W^2 \rightarrow s_W^2 + \delta s_W^2,$$

where $s_W = \sin \theta_W$, $c_W = \cos \theta_W$ and

$$g_Z = \frac{g}{\cos \theta_W} \quad s_W^2 = 1 - \frac{M_W^2}{M_Z^2} \quad G_F = \frac{1}{\sqrt{2}v^2}$$

$$\delta v = C_{H\ell}^{(3)} - \frac{1}{2} C_{\ell\ell} \quad \delta \sin^2 \theta_W = -\frac{v^2}{\Lambda^2} \frac{s_W c_W}{c_W^2 - s_W^2} \left[2s_W c_W \left(\delta v + \frac{1}{4} C_{HD} \right) + C_{HWB} \right]$$

$$\delta g_Z = -\frac{v^2}{\Lambda^2} \left(\delta v + \frac{1}{4} C_{HD} \right)$$