

Perturbativity constraints in $U(1)_{B-L}$ and Left-Right Models

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In collaboration with
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- Introduction & Motivation
- Theoretical Constraints
- Bounds in $U(1)_{B-L}$ model
- Bounds in Minimal LRSM
- Conclusions

- The Standard Model(SM) has been highly successful but needs extension to include new physics such as tiny neutrino masses, DM and baryon asymmetry.

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- Many TeV scale extensions introduce extended gauge groups like extra $U(1)$'s or $SU(2) \times U(1)$.
- Our results apply to a subclass of these gauge extensions of SM, where the generators of the extra gauge groups contribute to the electric charge.

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- The motivation is to embed the TeV-scale gauge extension into a larger gauge symmetry at GUT scale.
- We'll specifically focus on $U(1)_{B-L}$ & minimal LRSM, and discuss the implications for gauge boson searches.

Theoretical Constraint on Gauge Couplings

- Consider a SM extension: $SU(2)_L \times U(1)_X \times U(1)_Z$ such that:

$$Q = I_{3L} + I_X + \frac{Q_Z}{2}$$

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- Then requiring that coupling g_Z is perturbative at breaking scale,

$$\Rightarrow \boxed{r_g \equiv \frac{g_X}{g_L} > \tan \theta_W \left(1 - \frac{4\pi}{g_Z^2} \frac{\alpha_{EM}}{\cos^2 \theta_W} \right)^{-1/2}}$$

$U(1)_{B-L}$ model

- Particle content of the $SU(2)_L \times U(1)_{I_{3R}} \times U(1)_{B-L}$ model:

	$SU(2)_L$	$U(1)_{I_{3R}}$	$U(1)_{B-L}$
Q	2	0	$\frac{1}{3}$
u_R	1	$+\frac{1}{2}$	$\frac{1}{3}$
d_R	1	$-\frac{1}{2}$	$\frac{1}{3}$
L	2	0	-1
N	1	$+\frac{1}{2}$	-1
e_R	1	$-\frac{1}{2}$	-1
H	2	$-\frac{1}{2}$	0
Δ_R	1	-1	2

$U(1)_{B-L}$ model

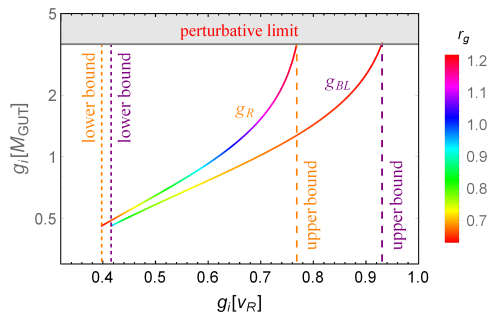
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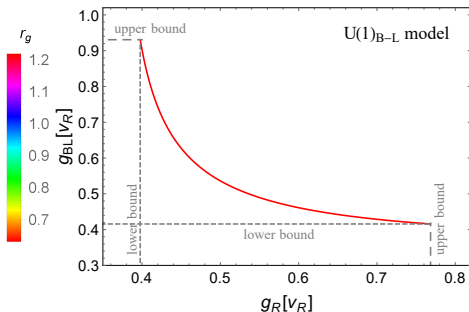
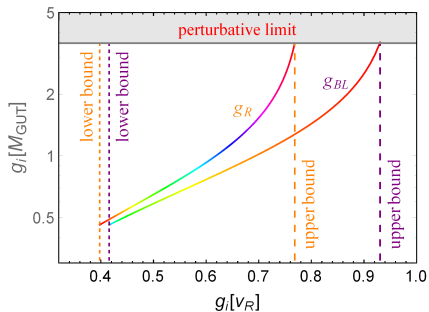
- The RGEs for the gauge couplings of the two $U(1)$'s are respectively

$$16\pi^2\beta(g_{I_{3R}}) = \frac{9}{2}g_{I_{3R}}^3, \quad 16\pi^2\beta(g_{BL}) = 3g_{BL}^3.$$

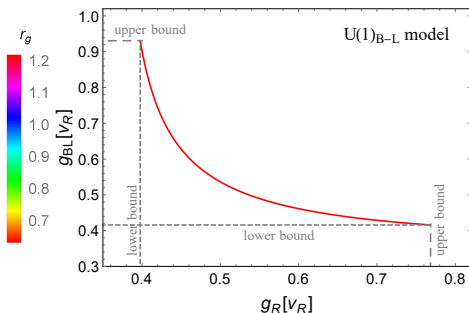
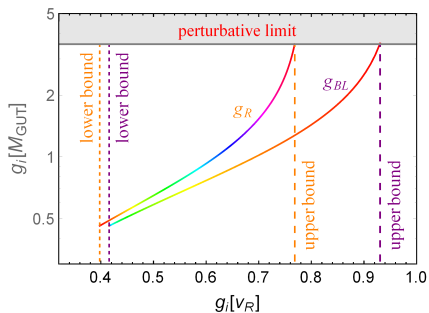
$SU(2)_L \times U(1)_{I3R} \times U(1)_{B-L}$ (Gauge Couplings)



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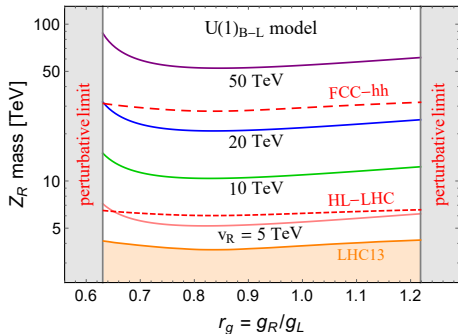
$SU(2)_L \times U(1)_{I3R} \times U(1)_{B-L}$ (Gauge Couplings)



$$0.398 < g_R < 0.768; \quad 0.416 < g_{BL} < 0.931, \quad \text{with } 0.631 < r_g < 1.218$$

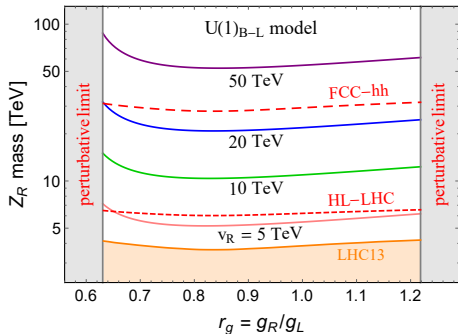
at $v_R = 5$ TeV

$SU(2)_L \times U(1)_{I3R} \times U(1)_{B-L}$ (Z_R searches)

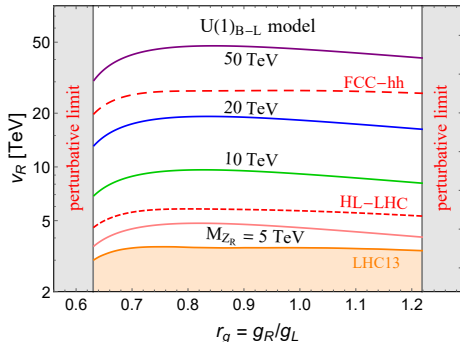


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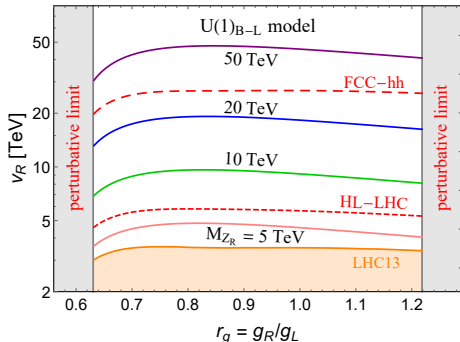
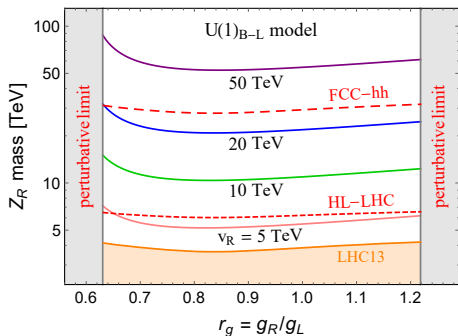
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$SU(2)_L \times U(1)_{I3R} \times U(1)_{B-L}$ (Z_R searches)



collider	M_{Z_R} [TeV]	v_R [TeV]
LHC13	[3.6, 4.2]	[3.02, 3.57]
HL-LHC	[6.0, 6.6]	[4.60, 5.82]
FCC-hh	[27.9, 31.8]	[19.9, 26.8]

Minimal LRSM

- Particle content of the minimal LRSM based on the gauge group $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$:

	$SU(2)_L$	$SU(2)_R$	$U(1)_{B-L}$
$Q_L \equiv \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	2	1	$\frac{1}{3}$
$Q_R \equiv \begin{pmatrix} u_R \\ d_R \end{pmatrix}$	1	2	$\frac{1}{3}$
$\psi_L \equiv \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	2	1	-1
$\psi_R \equiv \begin{pmatrix} N \\ e_R \end{pmatrix}$	1	2	-1
$\Phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}$	2	2	0
$\Delta_R = \begin{pmatrix} \frac{1}{\sqrt{2}}\Delta_R^+ & \Delta_R^{++} \\ \Delta_R^0 & -\frac{1}{\sqrt{2}}\Delta_R^+ \end{pmatrix}$	1	3	2

- The RGEs for the gauge couplings in the minimal LRSM are ¹

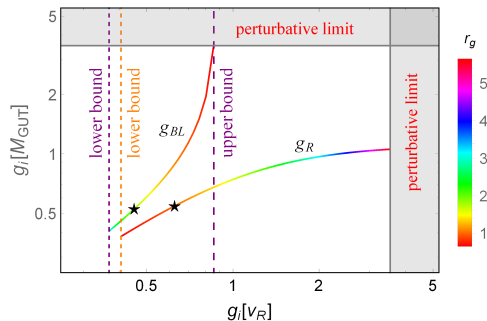
$$16\pi^2\beta(g_L) = -3g_L^3,$$

$$16\pi^2\beta(g_R) = -\frac{7}{3}g_R^3,$$

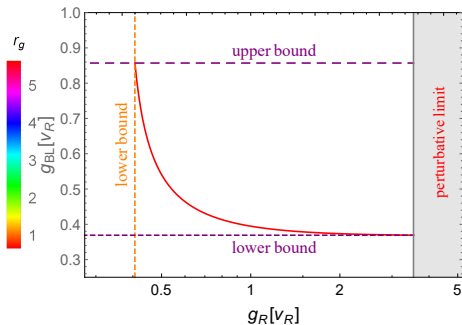
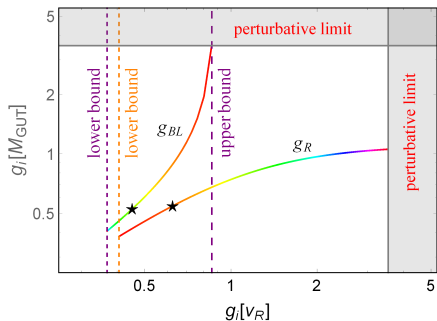
$$16\pi^2\beta(g_{BL}) = \frac{11}{3}g_{BL}^3$$

¹I. Z. Rothstein, Nucl. Phys. B358, 181 (1991)

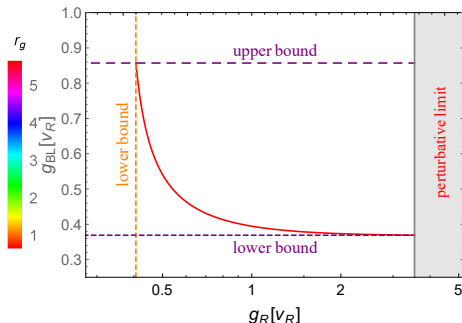
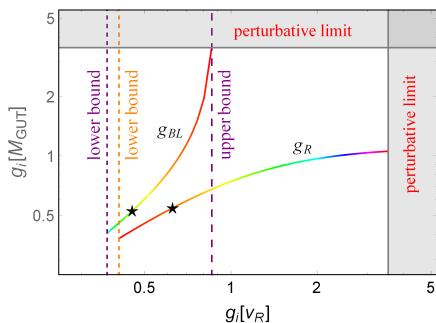
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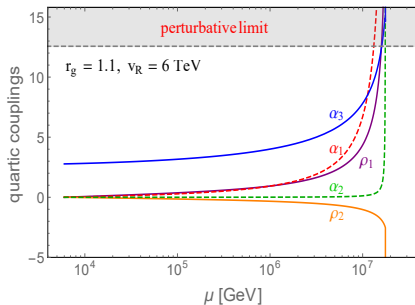
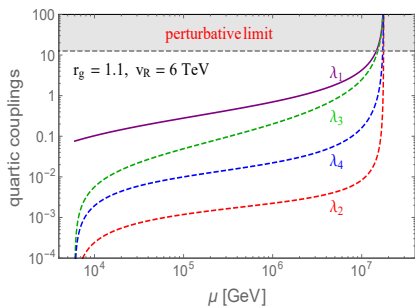
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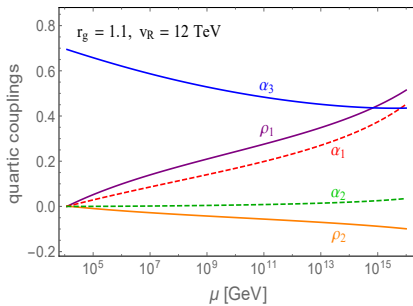
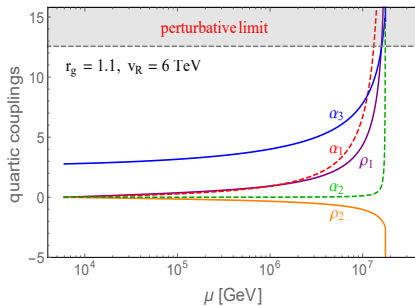
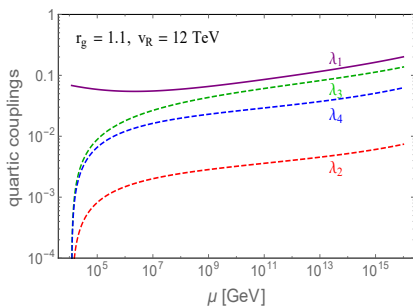
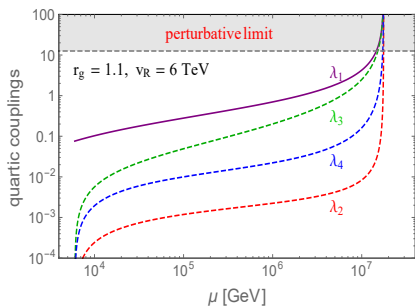
$$0.406 < g_R < \sqrt{4\pi}; \quad 0.369 < g_{BL} < 0.857, \quad \text{with } 0.648 < r_g < 5.65$$

at $v_R = 10 \text{ TeV}$

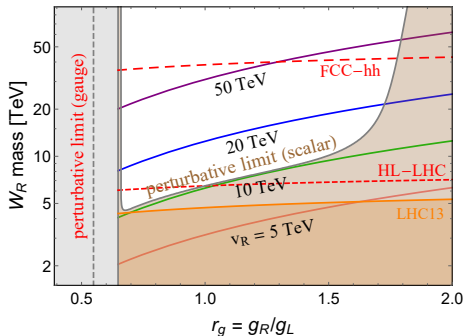
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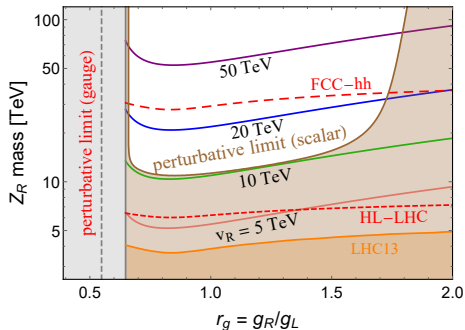
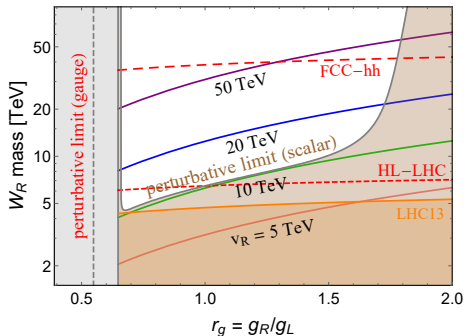


$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$ (Z_R and W_R searches)



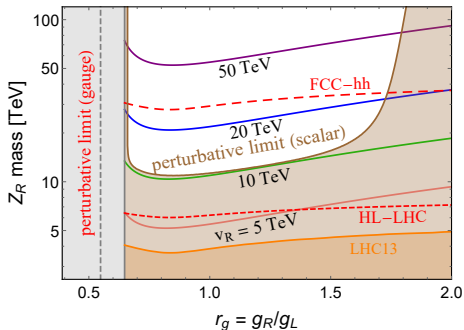
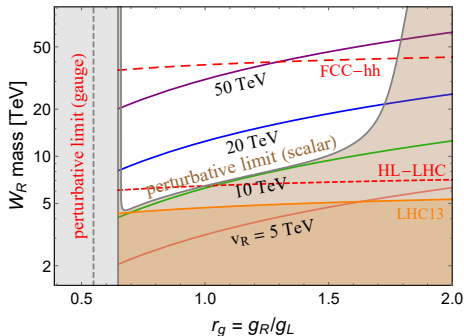
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(arXiv: 1809.11105)
(arXiv: 1803.11116)

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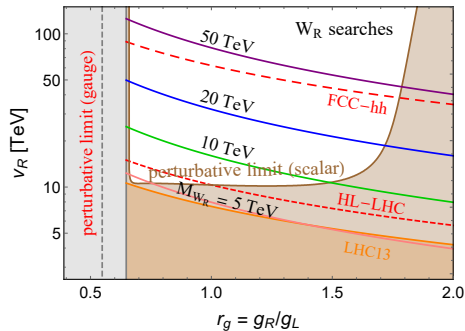
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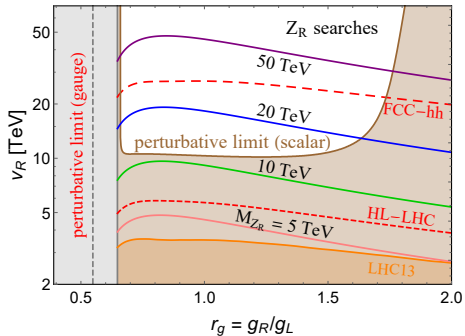
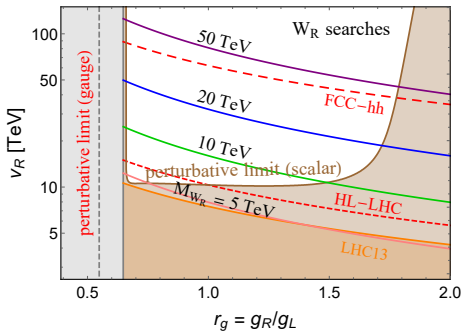


collider	W_R searches		Z_R searches	
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- In case, Z_R is found in HL-LHC run then couldn't be from minimal LRSM.
- The results can be generalized to other gauge group extensions.