



Introduction to Particle Physics and Cosmology

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Georgian Teachers' Programme

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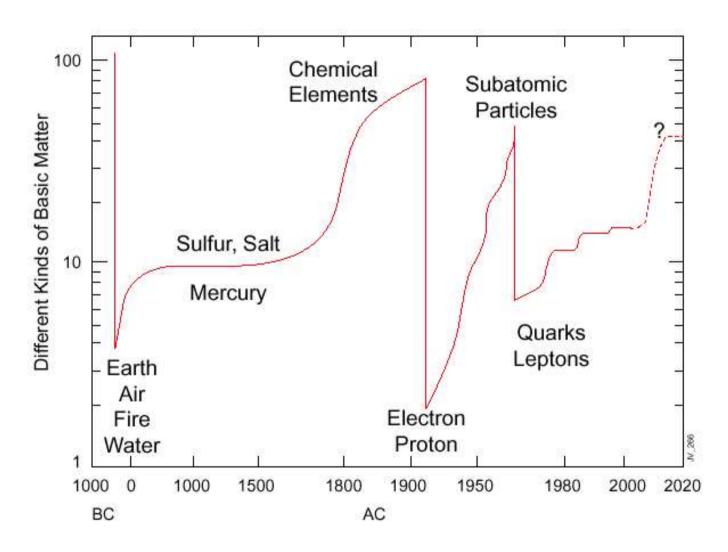


- Concepts and methods of particle physics
- ◆ The Standard Model (SM)
- ◆ CERN and its accelerator structure
- ◆ Large Hadron Collider (LHC)
- Proton-proton collisions, luminosity and triggers
- lacktriangle Lepton pair production: $J/\psi, \Upsilon, Z, \ldots$
- lacktriangle Other SM measurements: $t\bar{t}$, W^{\pm} , ZZ, ...
- ♦ Observation of the Higgs Boson
- ◆ Links with Cosmology
- Dark matter and dark energy
- ◆ Unanswered questions in particle physics and in cosmology
- ◆ Summary and outlook



Constituents of Matter



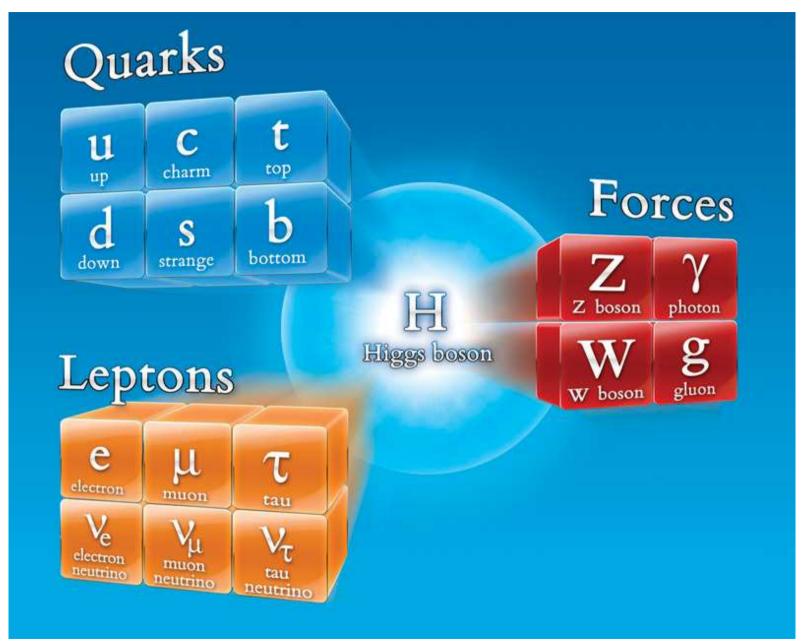


From http://teachers.web.cern.ch/teachers/archiv/HST2002/webgroup/mcclean/Introduction to Particle Physics.ppt



Fundamental constituents of the Standard Model







Particle Physics — What's This About?



'Elementary' Particles — $e, p, n, \nu, \mu, \tau, \gamma, W, Z \dots$ and their interactions.

You should already know a few things about them.

Is Particle Physics a difficult subject?

Compared to other areas of physics (nuclear, solid state, bio-...) and other sciences (botany, chemistry, zoology, medicine) PP is actually **very simple**:

- ◆ Particles have (relatively) few properties ('quantum numbers').
- ◆ These properties usually have few discrete values.
- ◆ Particles obey very simple, relatively few, well-defined laws.
- ◆ All elementary particles of the same type are absolutely identical.



Why does PP Seem So Hard Then?



◆ The world of particles is so far from our everyday experience, that all these simple properties and simple laws may look and seem unnatural and weird;

What can we do?

'Friendly' names: strangeness, charm, colour, top, bottom... Find analogies and simple rules

◆ Many mathematical methods used to describe the world of particles are quite advanced (Group Theory, Quantum Field Theory, Advanced Statistics . . .)

What can we do?
Use simplified maths, skip derivations...

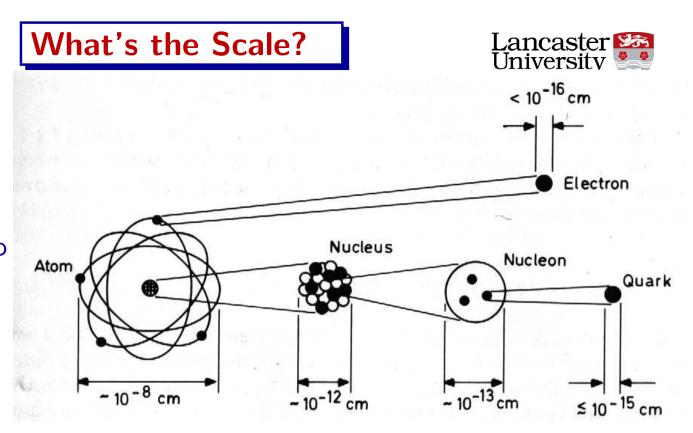
♦ Your intuition fails to work

What can we do?

Build our intuition by solving lots of various problems



'Elementary' Particles: the smallest constituents of matter (known so far): leptons and quarks, and also the interaction carriers: photons γ , gluons g, W^{\pm} and Z^0 bosons.



Well-established models and theories at present exclude gravitational interactions:

- 1. quantum theory of gravity has not been built yet;
- 2. may (should!) be tied to properties of space-time at tiny scales;
- 3. too weak to matter for particles under 'usual' circumstances.

However, weak, electromagnetic and strong interactions are understood and described reasonably well.



Is SI System Useful in Particle Physics?



Main properties of particles: mass m, charge e, spin s.

For an electron in SI system:

$$m_e = 9.109 \times 10^{-31} \text{ kg}$$
 $e = -1.602 \times 10^{-19} \text{ C}$
 $s_z = \pm \hbar/2 = \pm (1/2) \times 1.055 \times 10^{-34} \text{ J} \cdot \text{s}$

Particle physicists do not use SI system. Instead, a particle physicist would write:

$$m_e = 0.51 \text{ MeV}/c^2$$

 $e = -1 \text{ proton charge}$
 $s_z = \pm 1/2$

The last equation suggests: in particle physics

$$\hbar = 1.055 \times 10^{-34} \text{ J} \cdot \text{s} = 1$$

which, for one thing, states that in particle physics the product of units of [energy] and [time] is dimensionless.



Can we Make it Even Simpler?



So, it's natural to choose units such that $\hbar = 1$. This means that

[energy]
$$\times$$
 [time] $=$ 1 and also [momentum] \times [distance] $=$ 1

Now, remember the relativistic relation between Energy E, momentum \mathbf{p} and mass m:

$$E^2 = \mathbf{p}^2 \, c^2 + m^2 \, c^4$$

Relativistic particles move with speeds close to speed of light. Carrying all these huge factors like $(300000000 \text{ m/s})^2$ around will be avoided in a system of units where c=1, which simply means that [new unit of time] is [old unit of time]/c.

The choice $\hbar = 1$ and c = 1 would mean that

- ◆ Energy, momentum and mass are measured in the same units
- ◆ Angular momentum is dimensionless
- ◆ Time and distance are measured in the same units
- ◆ Energy is inverse of time
- ◆ One needs just one dimesional unit, which is usually chosen as the unit of energy
- ♦ In Particle Physics this is 1 GeV



Natural System of Units



The system of units with $\hbar=1$ and c=1 is called the Natural system:

```
1 \text{ unit of length} = 1 \text{ GeV}^{-1} \simeq 0.1978 \text{ fm}
1 \text{ unit of time} = 1 \text{ GeV}^{-1} \simeq 0.6588 \cdot 10^{-24} \text{ s}
1 \text{ unit of energy} = 1 \text{ GeV}
1 \text{ unit of momentum} = 1 \text{ GeV} sometimes \text{GeV}/c
1 \text{ unit of mass} = 1 \text{ GeV} sometimes \text{GeV}/c^2
```

Note: 1 GeV = 1000 MeV and $(1 \text{ GeV})^{-1} = (1000 \text{ MeV})^{-1}$, but $1000 \text{ GeV}^{-1} = 1 \text{ MeV}^{-1}$

One more unit: **barn b** for cross section: $1 \text{ b} = 10^{-24} \text{ cm}^2$.

One barn is far too big a unit for particle physics:

$$1 \text{ b} = 10^3 \text{ mb} = 10^6 \ \mu \text{b} = 10^9 \text{ nb} = 10^{12} \text{ pb} = 10^{15} \text{ fb}$$

The cross sections of most interesting processes in particle physics are usually measured in femtobarns fb.

Rare processes have smaller cross sections, and vice-versa.



Generations and masses

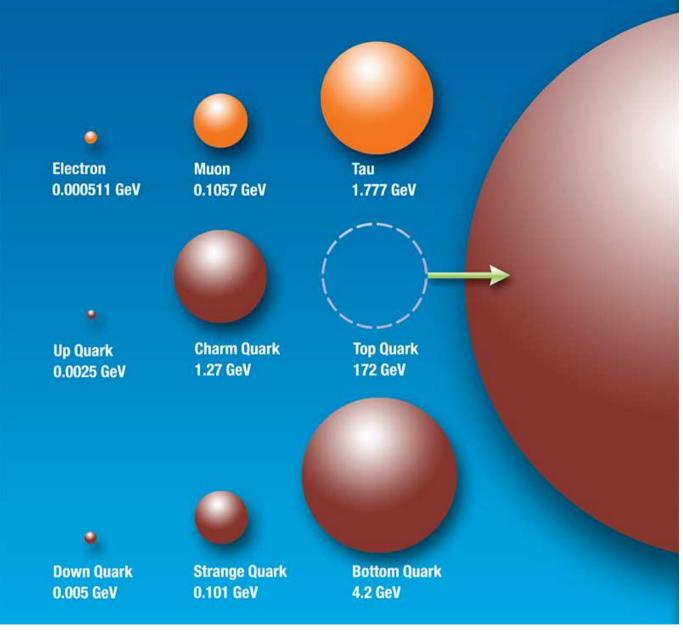


Three "generations"

Getting heavier and heavier

Top quark especially heavy

No clue why...





CERN 'overview'



Birdseye view of CERN and neighbourhood

Alps, lake Geneva, Geneva airport

LHC ring shown as the red line





CERN accelerator complex

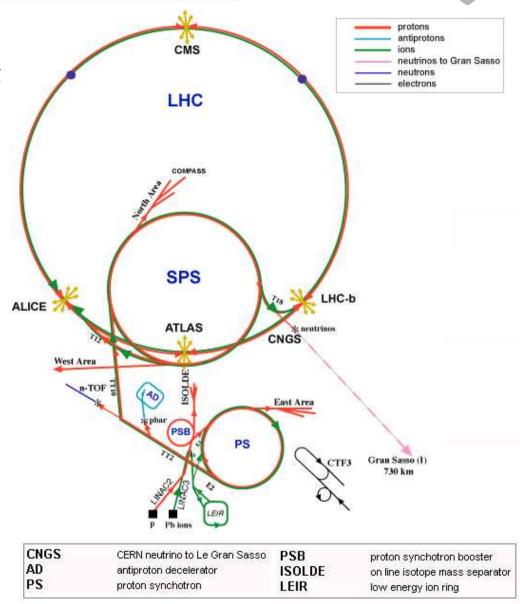


A very long chain of accelerators, culminating in the Large Hadron Collider (LHC)

Producing beams of protons, ions, antiprotons. . . even neutrinos!

Lots of experiments, all very interesting and important

I will only touch the one I know better...



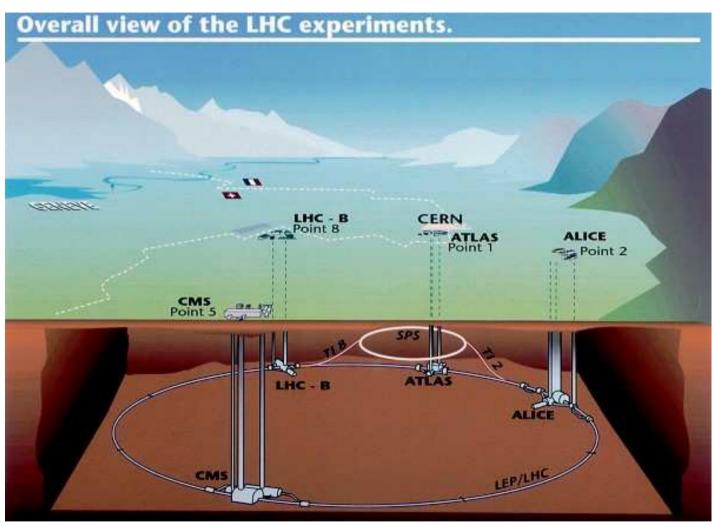


The Large Hadron Collider (LHC)



LHC is the flagship of CERN research programme, colliding two proton beams with energy of 13 TeV

One of the largest and most complicated engineering constructions in human history



Two multi-purpose experiments: ATLAS and CMS

Others – such as LHCb and ALICE – are more specialised

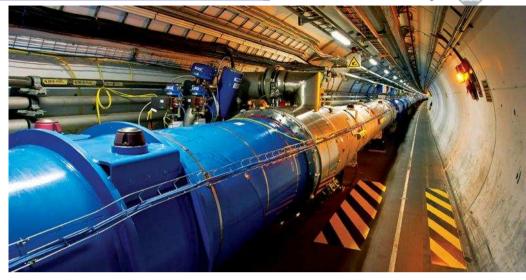


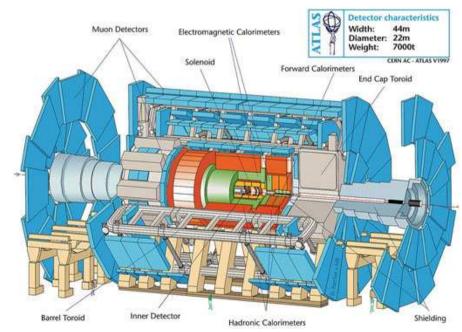
LHC tunnel, ATLAS and CMS



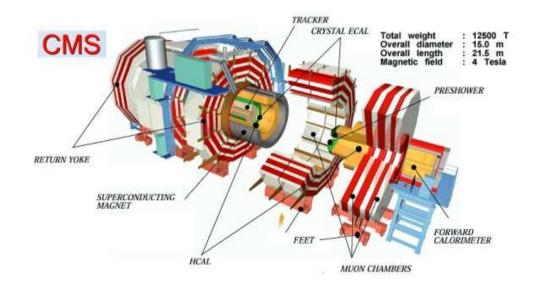
Tunnel 27 km long 100 m under the surface 2000 magnets of various types

Two huge multi-purpose experimental installations: ATLAS and CMS









(page 15) V. Kartvelishvili (Lancaster U)



Is LHC really a proton - proton collider?

tarcet

gluon

virtual pion



High energy of constituents is needed to produce something new and interesting

A proton is a bunch of quarks and gluons, each carrying a fraction of energy 13 TeV of pp collision energy barely enough to produce a 2 TeV object...



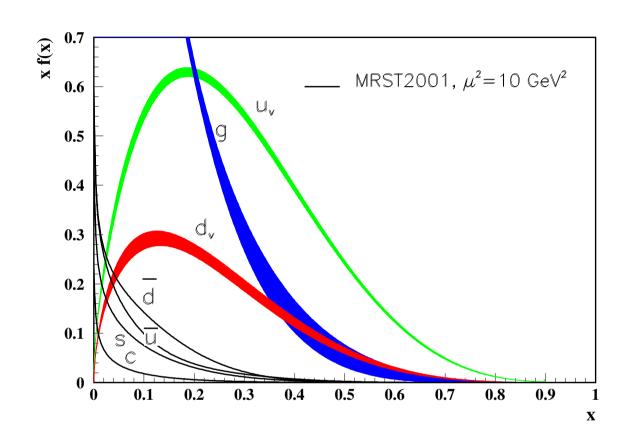
Quark and gluon distributions in a proton



Only 30% of proton energy is carried by the three constituent uud quarks

Most of proton energy is carried by gluons

The "sea" of quark-antiquark pairs is also important



$$M^2 = x_1 \times x_2 \times (13 \ TeV)^2$$

$$d\sigma \sim f_1(x_1) \times f_2(x_2) \times \hat{\sigma}(M^2)$$



Cross sections and units



- ◆ The intensity of various collisions is measured in terms of the cross section for particular reactions
- Cross section is the effective area which needs to be crossed by a test particle to get scattered
- ◆ Since early days of nuclear physics, measured in barns

$$1 \text{ barn} = 10^{-28} \text{ m}^2 = 100 \text{ fm}^2$$

is about the size of lead or uranium nucleus

- ◆ Total cross section of proton-proton collisions is about 100 millibarn at 7 TeV
- ♦ Interesting processes like Higgs production have much smaller probabilities, and hence much smaller cross sections, measured in picobarns (10^{-12} barn) or femtobarns (10^{-15} barn) or even attobarns (10^{-18} barn)
- ◆ The smaller the cross section of a process, the fewer events you get
- ♦ Integrated luminosity of 100 pb⁻¹ means that if the cross section is 1 pb, you will see 100 events



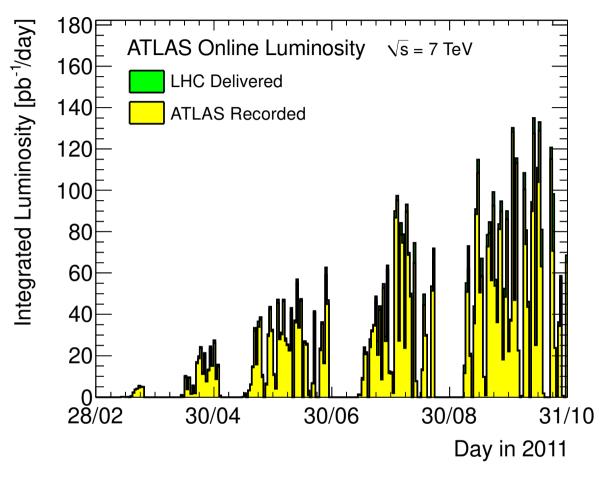
Luminosity



In early days of LHC: 100's of collisions / sec

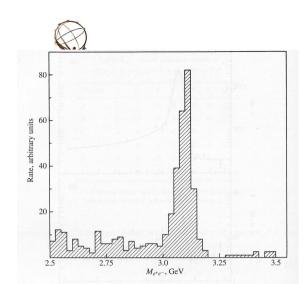
Now: many millions / sec

No time for viewing events one-by one...



Full computing power of CERN only allows to reconstruct "just" a few hundred events per second

Very careful selection ("triggering") of potentially interesting events is required!



1974: discovery of J/ψ



← Discovery 1: Ting's group

$$pN \to e^+e^-X$$

at $P_{\rm lab}=30~{\rm GeV/c}$

[Aubert et al., PRL, 6/11/1974]

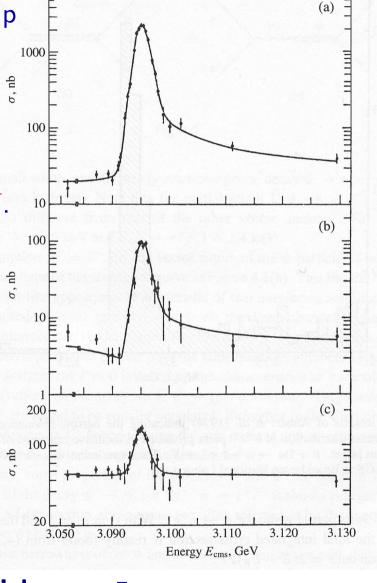
Found a peak in e^+e^- inv.mass at 3.1 GeV, called it J.

Discovery 2: Richter's group ⇒

- $(a) e^+e^- \rightarrow \text{hadrons}$
- $(b) e^+e^- \rightarrow \mu^+\mu^-$
- (c) $e^+e^- \rightarrow e^+e^-$

[Augustin et al., PRL, 7/11/1974]

Found a peak in all these three cross-sections, at the c.m.s. energy 3.1 GeV; called it ψ .



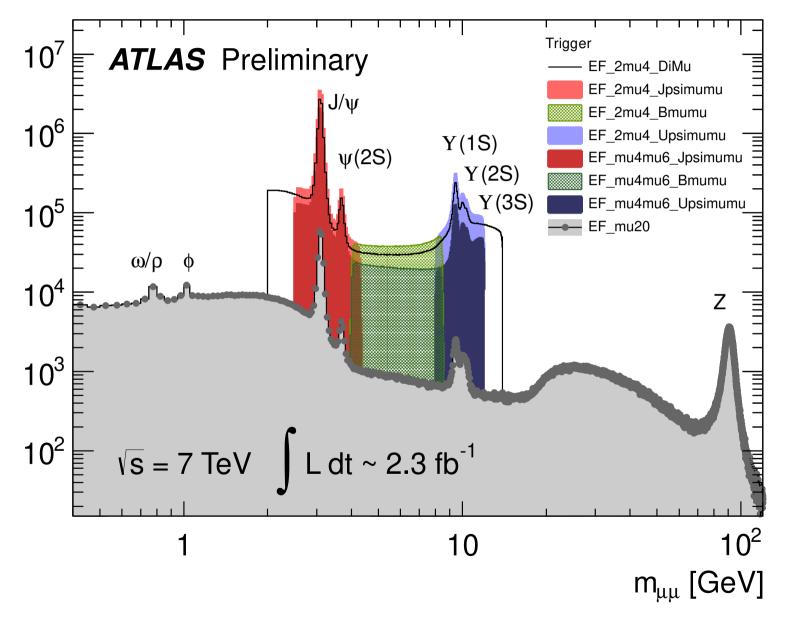
Now we know: J/ψ is a bound state of charm-anticharm, $c\bar{c}$.



History of 20th century Particle Physics in one plot



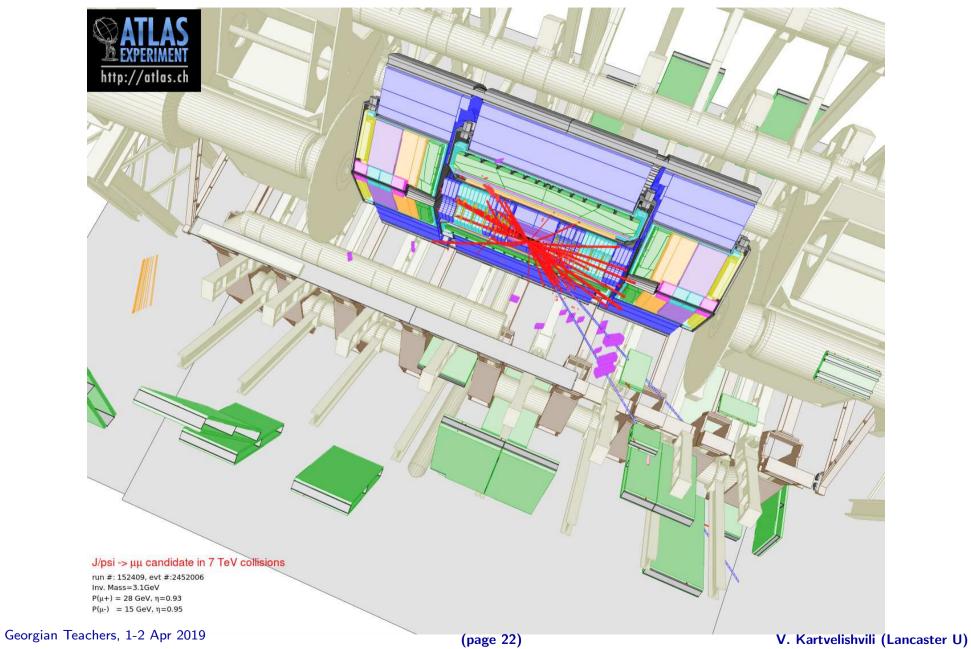






$pp \to J/\psi(\to \mu^+\mu^-) + X$

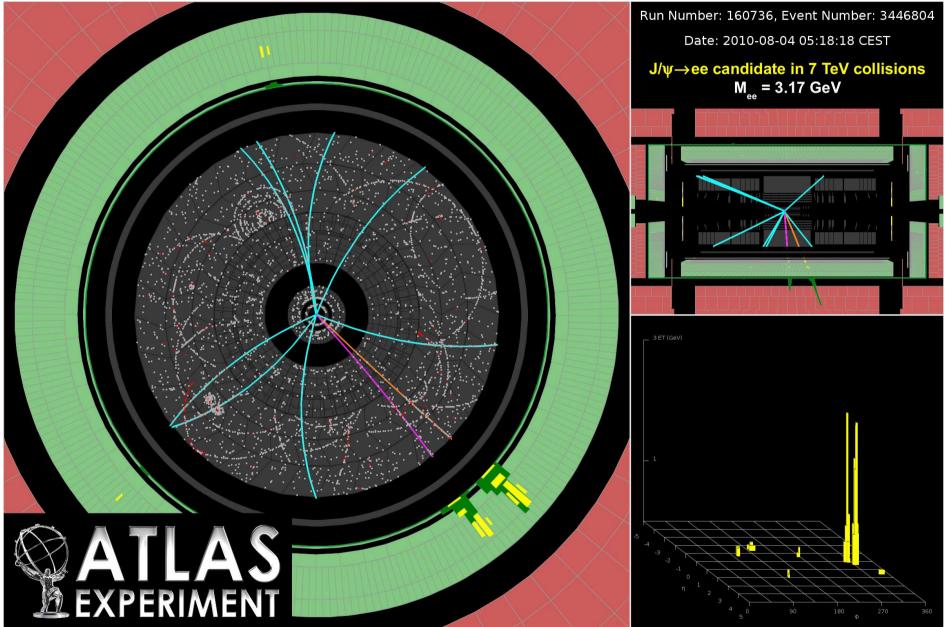






$pp \to J/\psi(\to e^+e^-) + X$

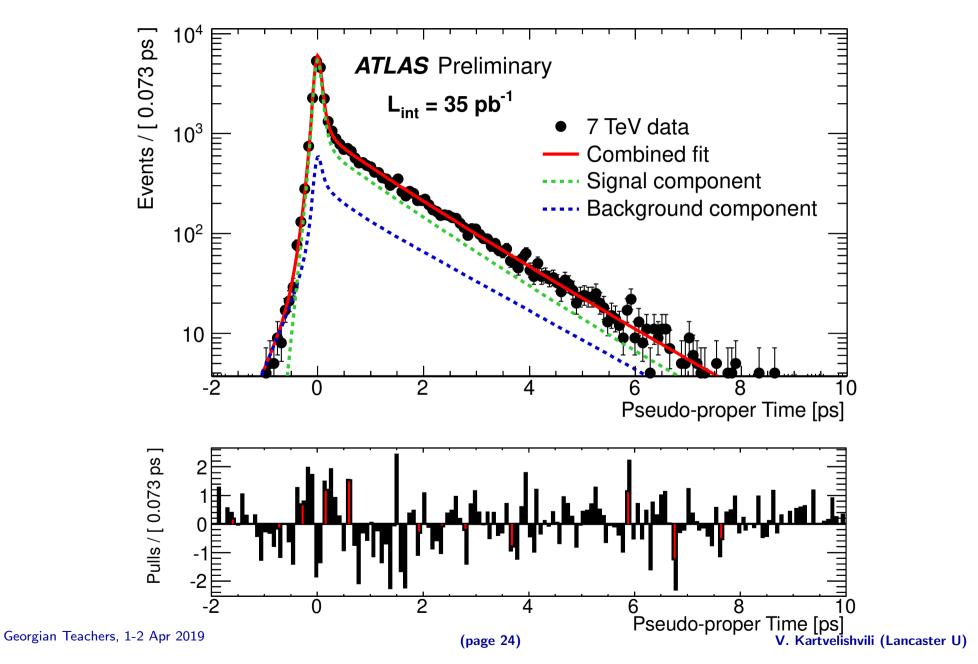






Proper Decay Time of the J/ψ vertex

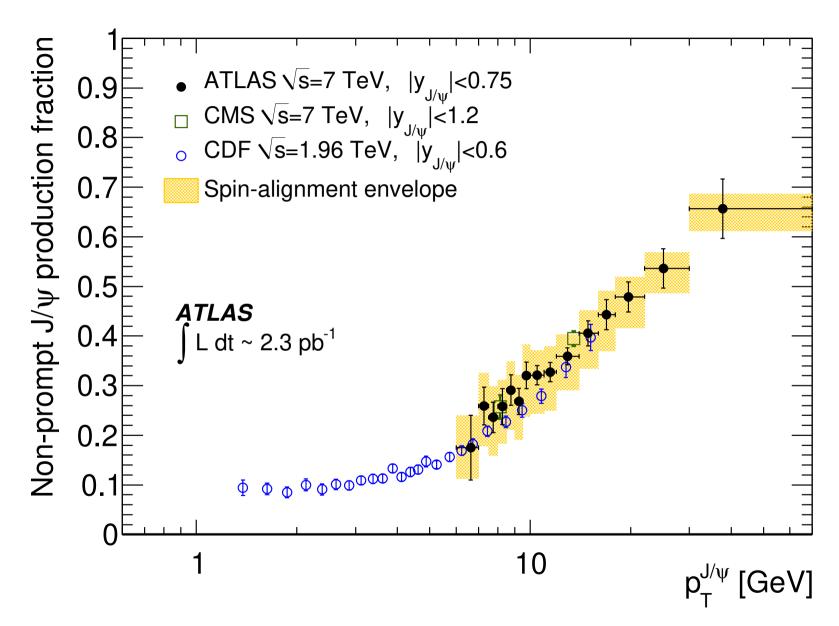






Fraction of non-promptly produced J/ψ

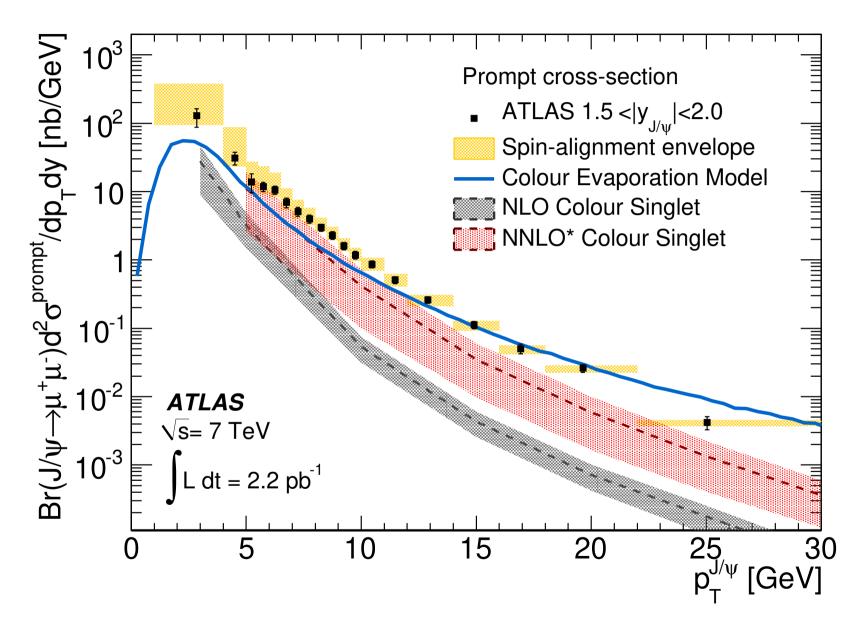






p_T dependence of prompt J/ψ

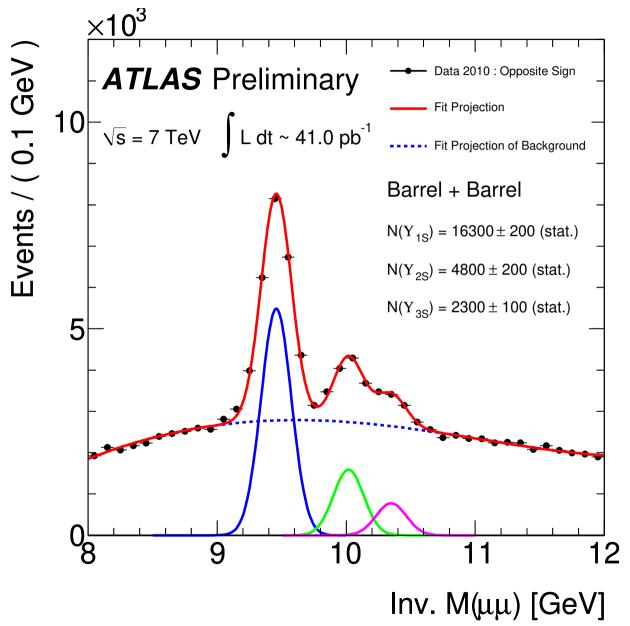






$b \bar b$ bound states: Υ system







Spectroscopy of bb **mesons**



Observed bottomonium radiative decays in ATLAS, L = 4.4 fb¹

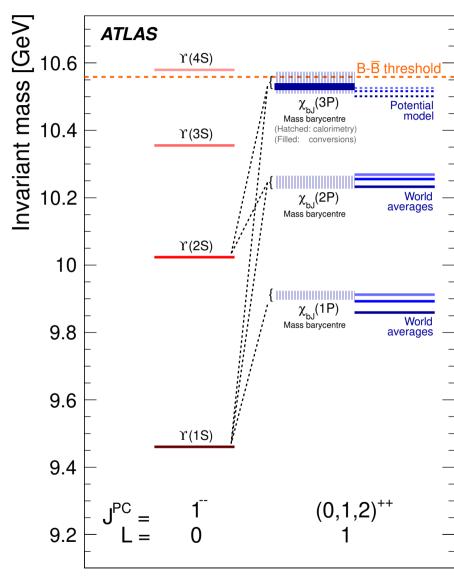
Spectroscopy similar to hydrogen atom

 $\Upsilon(1S)$: ground state

 $\Upsilon(2S,3S)$: radial excitations

Three families of χ_b : orbital excitations, L=1

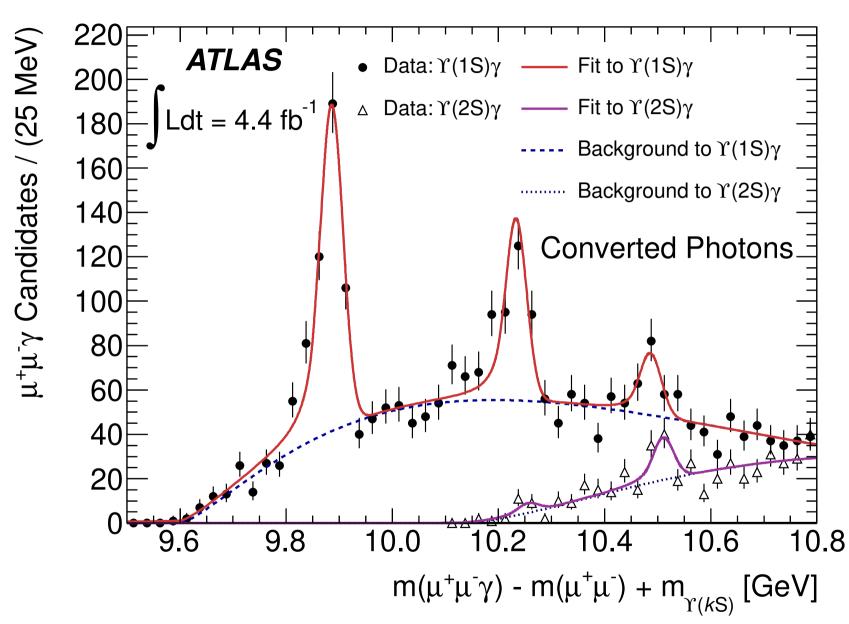
Until 22 December 2011, only $\chi_b(1P)$ and $\chi_b(2P)$ were observed





All three χ_b peaks as seen by ATLAS

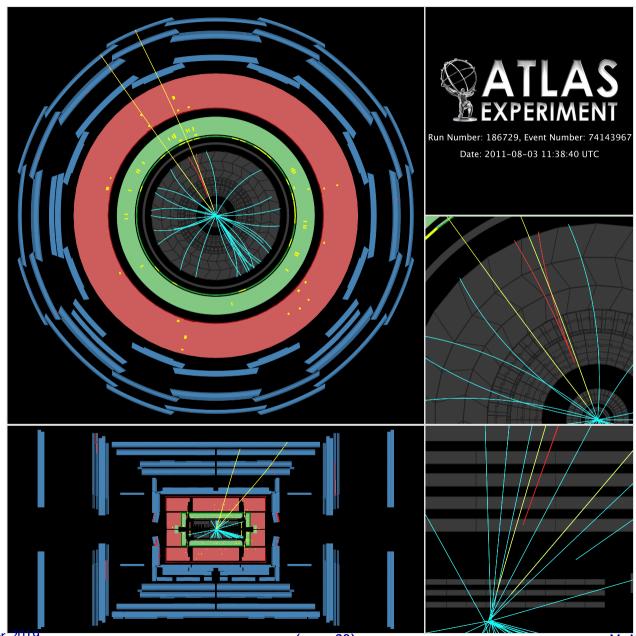






Event with $\chi_b(3P)$ candidate

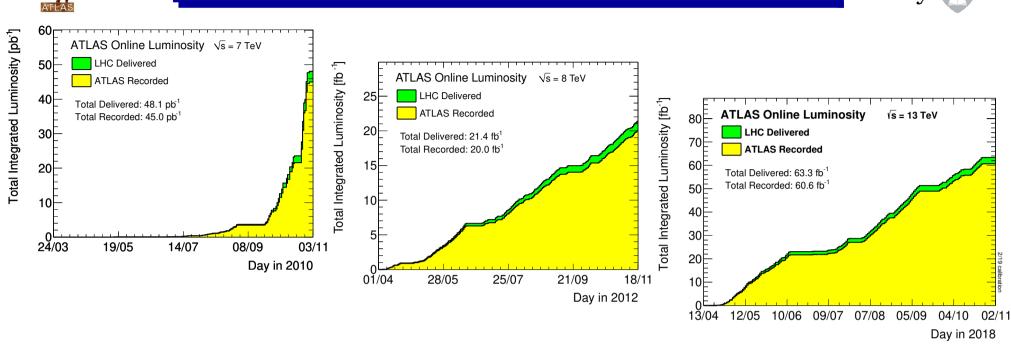






Integrated luminosity in 2010, 2012, 2018





Look at the scales on y-axes: 1 fb⁻¹ = 1000 pb⁻¹

Dramatic progress over the years, meaning that one can now access less and less frequent processes...

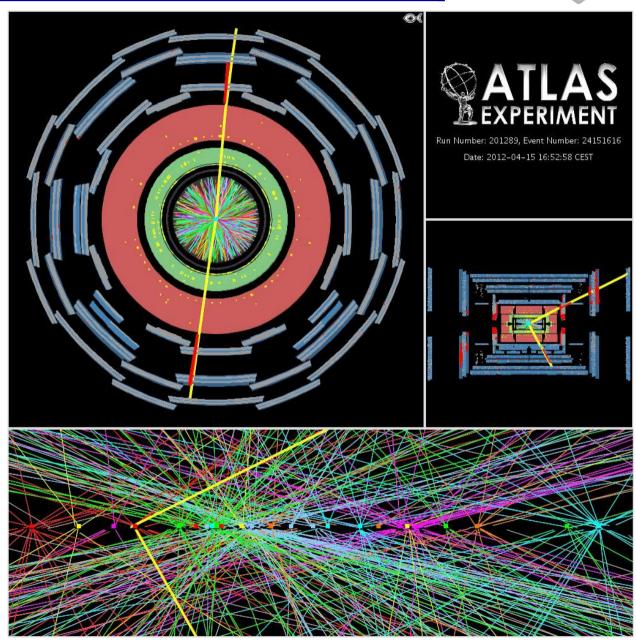
...and at a higher and higher energy!



$Z \rightarrow \mu^+ \mu^-$ candidate at high luminosity



There are 20+ collisions in one bunch crossing, with a $Z \to \mu^+\mu^-$ candidate produced in one of them.





W^+W^- pair production

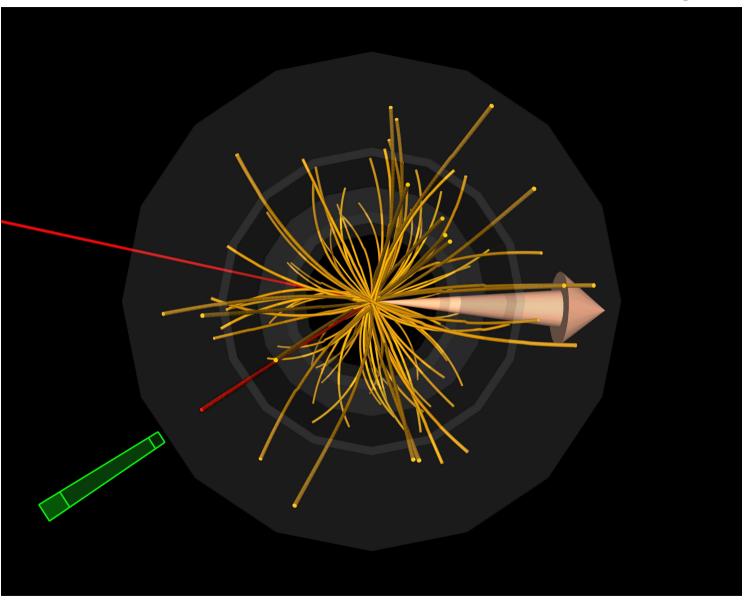


$$W^+ \to \mu^+ \nu_\mu$$

$$W^- \to e^- \bar{\nu}_e$$

Neutrinos escape detection

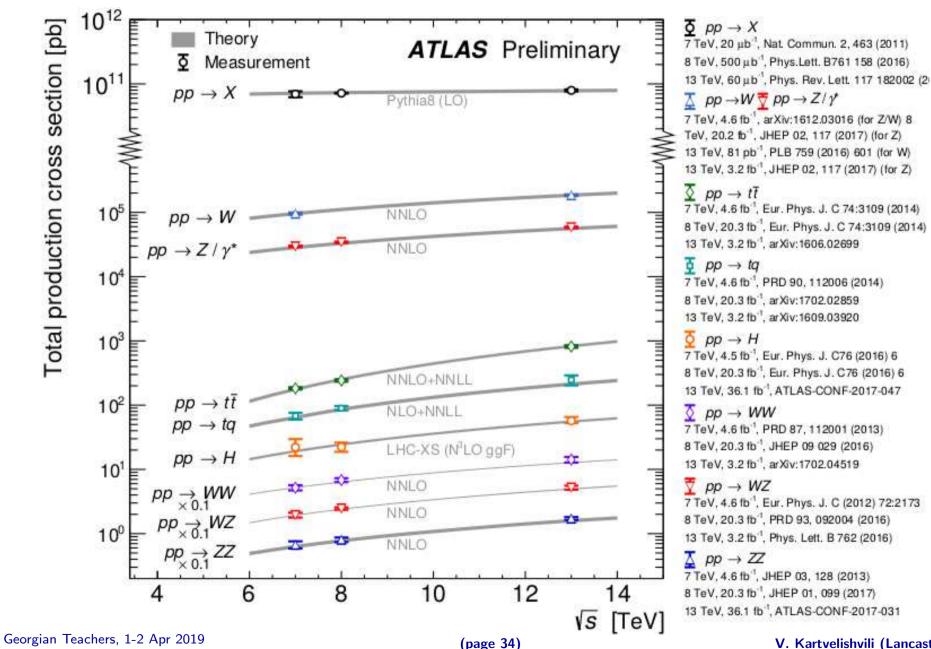
 \Rightarrow missing P_T





Standard Model cross sections vs theory

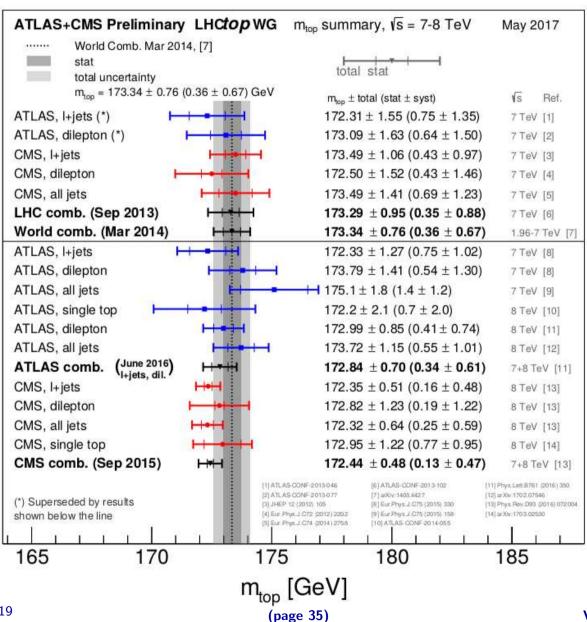






Top quark mass measurement

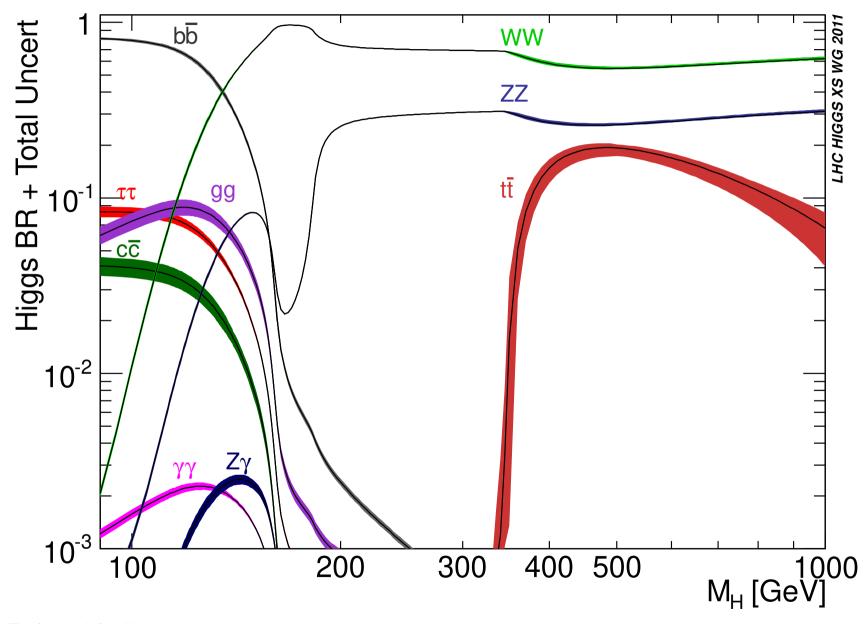






Decay modes of the Standard Model Higgs Boson

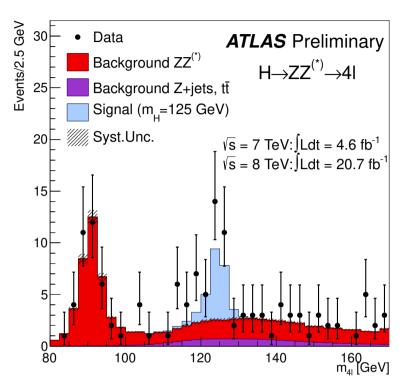


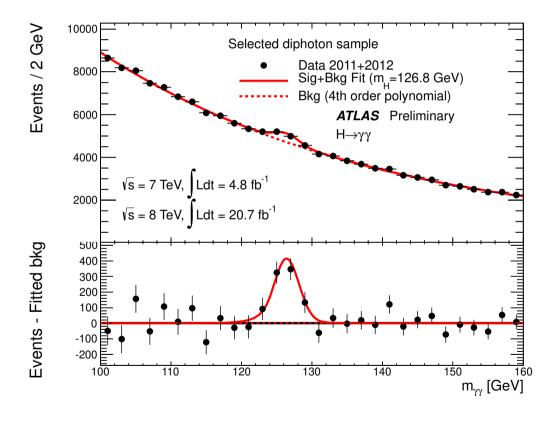




Higgs boson observation



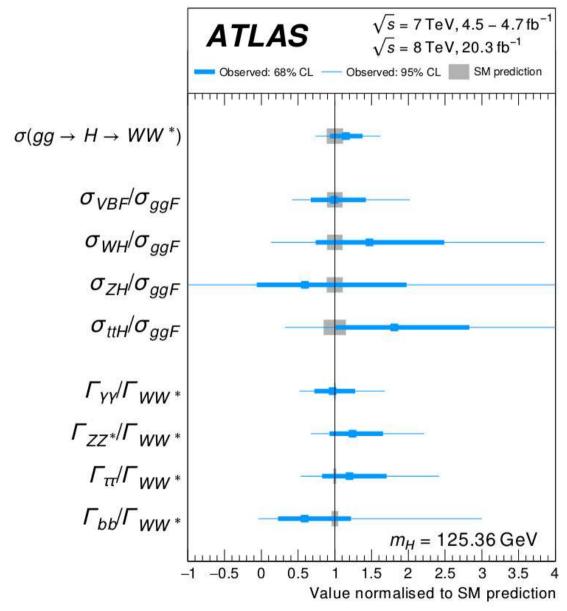






Higgs decay Branching Ratios vs SM







Questions to the Standard Model



There are three types of interactions in the Standard Model, and the variety of gauge bosons, the interaction carriers: γ for electromagnetic, W^{\pm}, Z^0 for weak, g for strong.

- ♦ Why are these three types so different and the fourth, gravity, even more so?
- ♦ Why are there three generations of quarks and leptons?
- ♦ Why fractional electric charges of quarks?
- ♦ Why are the fermion masses so different?
- ♦ What determines the mixing of various generations?

These and many more questions cannot be answered within SM.

We need a bigger theory...



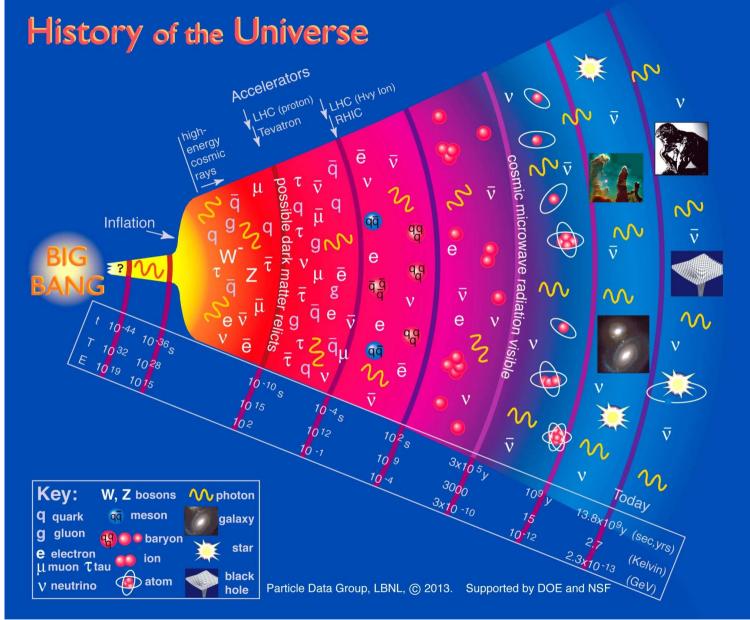
Cosmology: source of inspiration



- lacktriangle Universe is made up of $\sim 10^{11}$ galaxies; each galaxy contains $10^{10}-10^{12}$ stars
- ◆ Cosmology: science about the history of the Universe
- ◆ Assumption: laws of physics have not changed along the way
- ◆ Method 1: observe the Universe evolution NOW and try to extrapolate backward
- ♦ Method 2: assume some starting point (the Big Bang) and extrapolate forward
- ◆ The overall established picture in modern cosmology is arguably as stable and solid as the Standard Model in Particle Physics, but it also has its unanswered questions
- ◆ The hope (from both camps) is that the answers may be shared!









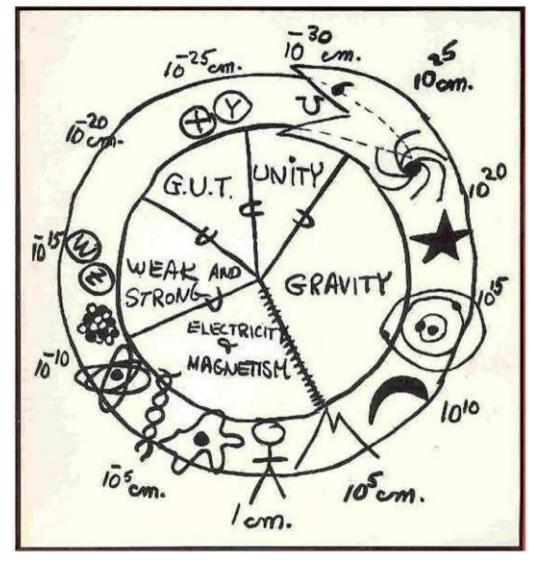
Glashow's serpent



As usual, "natural" system of units:

- \bullet $\hbar = 1, c = 1, k_B = 1$
- lacktriangle distance \sim time
- Energy $\sim 1/\text{distance}$
- lacktriangle Temperature \sim Energy
- ♦ Hence, Planck's mass

$$M_p = \sqrt{\frac{\hbar c}{G_N}} = 10^{19} \text{ GeV}$$
$$\sim 10^{-25} \text{ cm}$$





Expanding Universe



Experimental fact: Universe is expanding

Light from distant galaxies is red-shifted (Doppler effect)

The larger the distance, the more the shift (can be measured precisely)

The light wave expands with space, hence the shift towards lower frequency

Hubble constant: 70 km/s per Megaparsec

Once, the Universe was 3000 times smaller – and 3000 times hotter than today

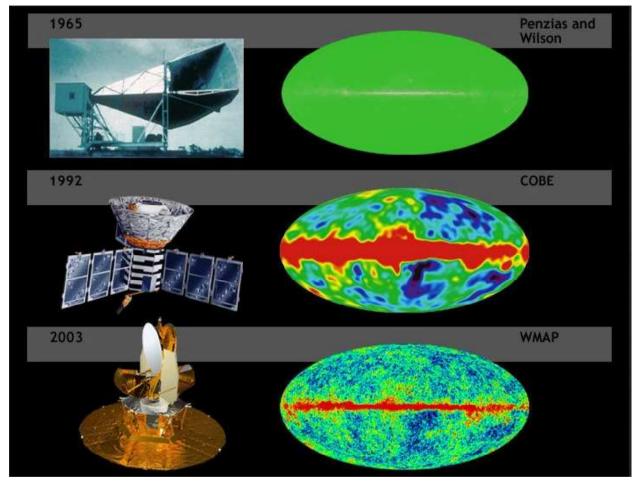
Cosmic Microwave Background $2.7~{\rm K}$ today: photons wandering in space since then

Almost isotropic (same in all directions) – but NOT EXACTLY!



CMB anisotropy





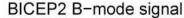
Ripples from times 300 000 years ago, at the level of 10^{-3}

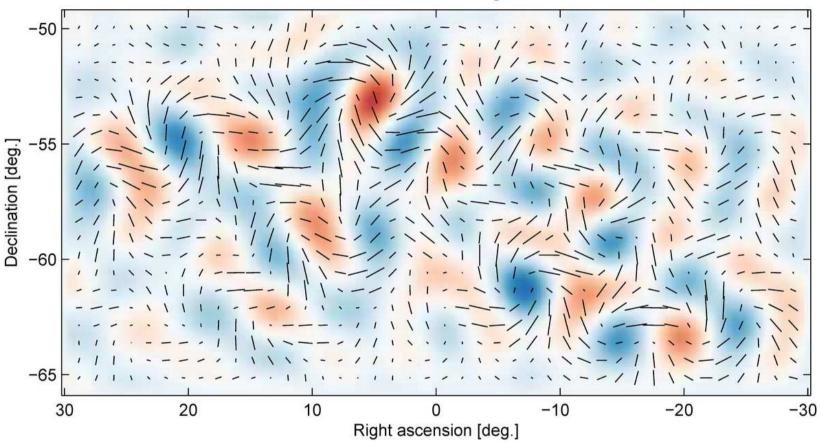
These small non-uniformities may be signals from the seeds of galaxy formation



Polarisation fluctuations







Possible signs of gravitational waves from the Big Bang?



Bariogenesis



Once the Universe was a billion times smaller and hotter than today

Light chemical elements were formed: He^4 , D, He^3 , Li, ...

Relative abundance of these elements can be predicted by theory

Depends on density of matter and number of types of particles

Does not seem to be enough to stop expansion, or even to form the galaxies like ours:



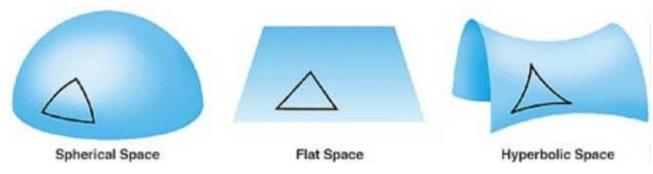


Cosmological inflation

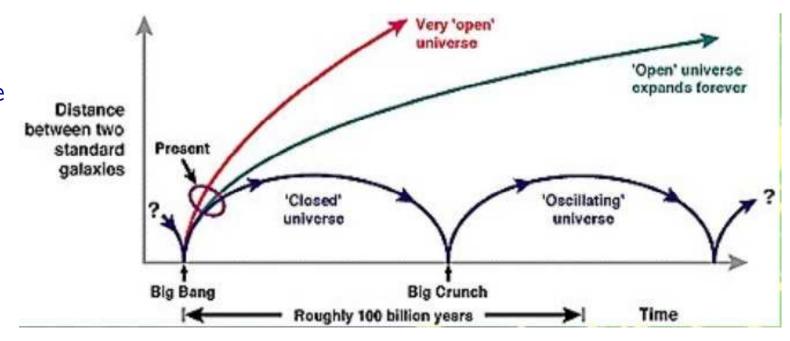


Basic idea: very early, about maybe 10^{-35} s after the Big Bang, the expansion was exponentially fast

Can explain why the universe looks almost flat now



Fate of the Universe depends on this:





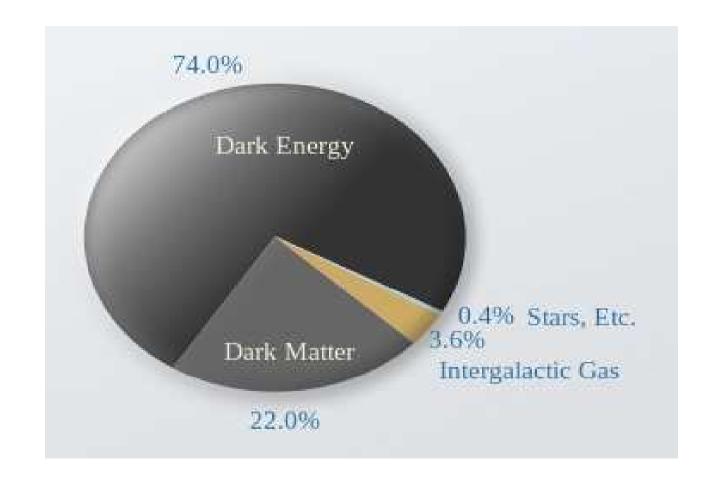
Energy density budget of the Universe



There is some critical value of the energy density which keeps the balance between expansion and contraction of the universe.

 $\Omega=1$ corresponds to a flat universe – close to what we see today

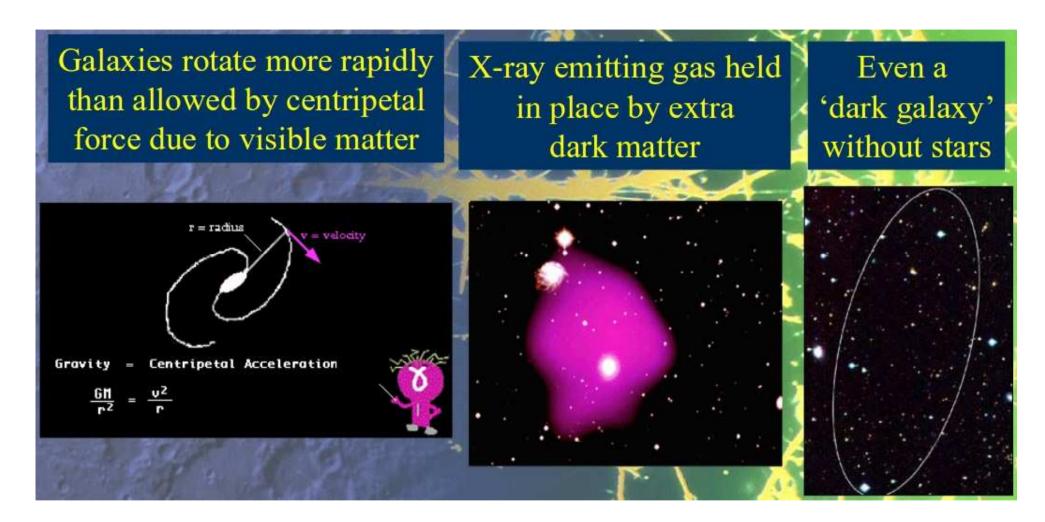
Latest measurements show that there are different components to this density:





Evidence for Dark Matter – I

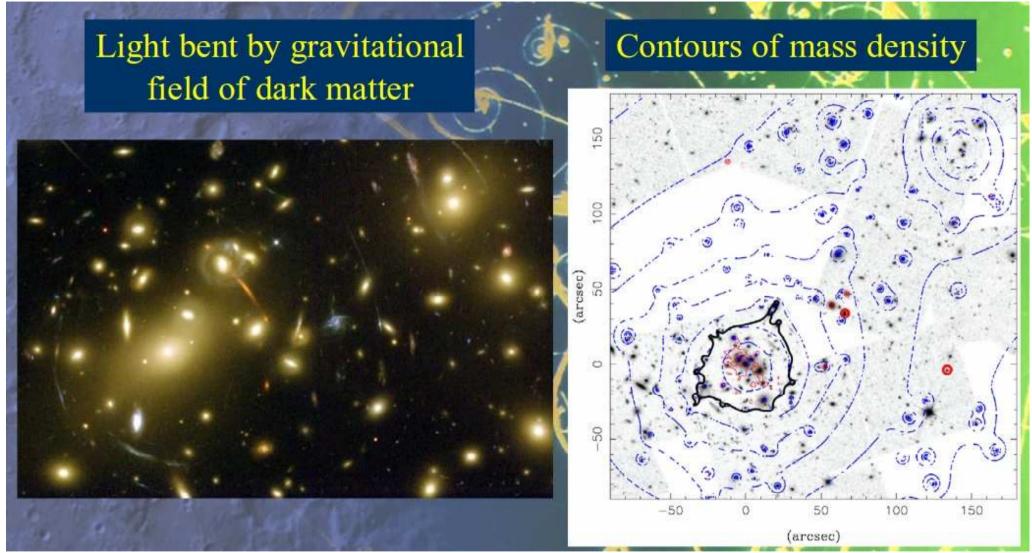






Evidence for Dark Matter – II

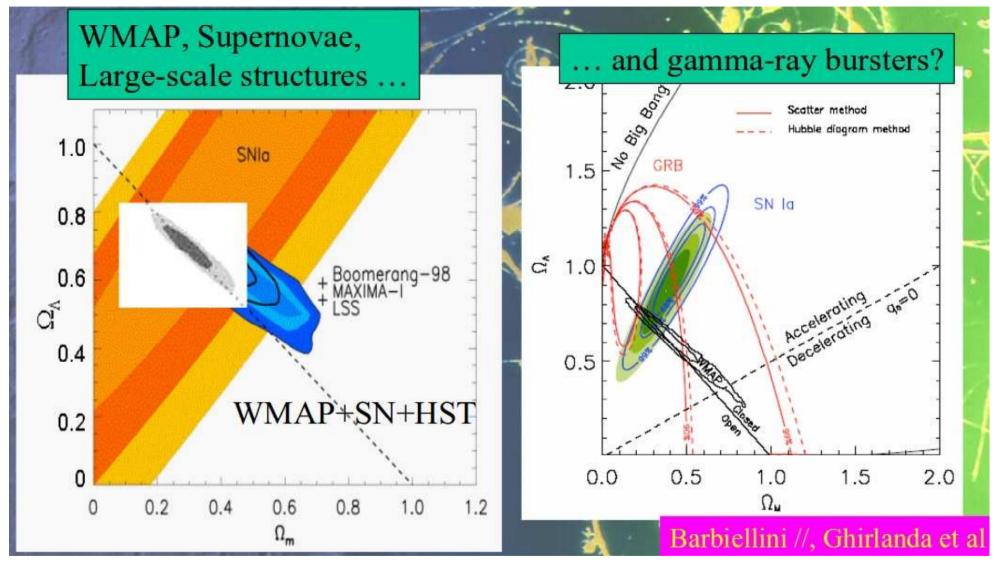






Experimental data on components of Ω







Unresolved questions in Cosmology



- Why is the Universe so big and old?
 ~ 13.000.000.000 years
- Why is its geometry nearly Euclidean?
 almost flat: density nearly critical
- Where did the matter come from?
 1 proton for every 1,000,000,000 photons
- How did structures form?
 ripples + invisible dark matter?
- What is the dark matter?
- What is the dark energy?

The hope is that Particle Physics can help answer at least some of these!



Beyond the Standard Model



- ◆ Is there a bigger symmetry group, which will become visible at higher energies?
 - ⇒ Grand Unification
- ◆ Or maybe the Poincaré-Lorentz invariance group can be extended to include anticummutation relations?
 - ⇒ Supersymmetry
- ◆ Or maybe our space-time has more than 3+1 dimensions, some of which are "compactified"?
 - ⇒ Large extra dimensions

These, and many other, theories exist — and predict some observable effects.

Physicists are searching for them, in a hope to answer some of the questions...



Supersymmetry searches: lower limits



ATLAS SUSY Searches* - 95% CL Lower Limits

ATLAS Preliminary

 \sqrt{s} = 7, 8 TeV

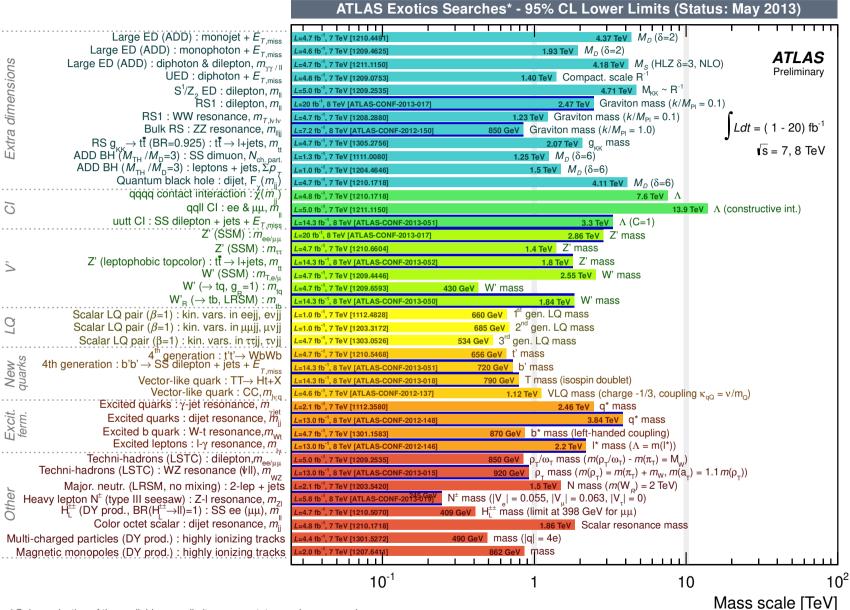
| | Status: ICHEP 2014 Model | $e,\mu,	au,\gamma$ | Jets E_1^{i} | ^{miss} ∫£ dt[fb | Mass limit | | $\sqrt{s} = 7, 8 \text{ TeV}$ Reference |
|------------------------------|---|---|---|---|---|---|---|
| Codmon Contrology | GMSB (* NLSP) GMSB (* NLSP) GGM (bino NLSP) GGM (wino NLSP) GGM (higgsino-bino NLSP) GGM (higgsino NLSP) Gravitino LSP | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ | 3-6 jets | Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 4.7 Yes 20.3 Yes 20.3 Yes 4.8 Yes 4.8 Yes 4.8 Yes 5.8 Yes 10.5 | $ \begin{array}{c ccccccccccccccccccccccccccccccccccc$ | $\begin{array}{l} n(\bar{q}) \\ (\bar{q}) \\ = 0 \text{ GeV, } m(1^{st} \text{ gen. } \bar{q}) = m(2^{nd} \text{ gen. } \bar{q}) \\ = 0 \text{ GeV} \\ < 200 \text{ GeV, } m(\bar{\chi}^{z}) = 0.5(m(\bar{\chi}^{0}_{1}) + m(\bar{g})) \\ = 0.5(m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) \\ = 0.5(m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) \\ = 0.5(m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}_{1}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}_{1}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}_{1}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) + m(\bar{\chi}^{0}) \\ = 0.5(m(\bar{\chi}^{0}) + m$ | 1405.7875 ATLAS-CONF-2013-062 1308.1841 1405.7875 1405.7875 ATLAS-CONF-2013-062 ATLAS-CONF-2013-089 1208.4688 1407.0603 ATLAS-CONF-2014-001 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-152 ATLAS-CONF-2012-147 |
| 3 rd gen. | $ \begin{array}{ccc} \bar{g} \rightarrow b \bar{b} \bar{\chi}_1^0 \\ \bar{g} \rightarrow t \bar{t} \bar{\chi}_1^0 \\ \bar{g} \rightarrow t \bar{t} \bar{\chi}_1^0 \\ \bar{g} \rightarrow t \bar{t} \bar{\chi}_1^0 \\ \bar{g} \rightarrow b \bar{t} \bar{\chi}_1^1 \end{array} $ | 0 0 7 0-1 <i>e</i> , μ 0-1 <i>e</i> , μ | 7-10 jets \ | Yes 20.1 Yes 20.3 Yes 20.1 Yes 20.1 | $ \begin{array}{cccc} \tilde{g} & & & & & & & \\ \tilde{g} & & & & & & & \\ \tilde{g} & & & & & & \\ \end{array} $ | <400 GeV <350 GeV <400 GeV <300 GeV | 1407.0600 1308.1841 1407.0600 1407.0600 |
| 3 rd gen. squarks | $\begin{array}{c} \tilde{b}_{1}\tilde{b}_{1},\tilde{b}_{1}\to b\tilde{\chi}_{1}^{0}\\ \tilde{b}_{1}\tilde{b}_{1},\tilde{b}_{1}\to b\tilde{\chi}_{1}^{0}\\ \tilde{b}_{1}\tilde{b}_{1},\tilde{b}_{1}\to b\tilde{\chi}_{1}^{1}\\ \tilde{i}_{1}\tilde{i}_{1}(\text{light}),\tilde{i}_{1}\to b\tilde{\chi}_{1}^{\pm}\\ \tilde{i}_{1}\tilde{i}_{1}(\text{light}),\tilde{i}_{1}\to b\tilde{\chi}_{1}^{\pm}\\ \tilde{i}_{1}\tilde{i}_{1}(\text{inedium}),\tilde{i}_{1}\to b\tilde{\chi}_{1}^{\pm}\\ \tilde{i}_{1}\tilde{i}_{1}(\text{medium}),\tilde{i}_{1}\to b\tilde{\chi}_{1}^{\pm}\\ \tilde{i}_{1}\tilde{i}_{1}(\text{heavy}),\tilde{i}_{1}\to t\tilde{\chi}_{1}^{0}\\ \tilde{i}_{1}\tilde{i}_{1}(\text{heavy}),\tilde{i}_{1}\to t\tilde{\chi}_{1}^{0}\\ \tilde{i}_{1}\tilde{i}_{1},\tilde{i}\to c\tilde{\chi}_{1}^{0}\\ \tilde{i}_{1}\tilde{i}_{1},\text{outual}(\text{aGMSB})\\ \tilde{i}_{2}\tilde{i}_{2},\tilde{i}_{2}\to \tilde{i}_{1}+Z \end{array}$ | 1-2 e, μ 2 e, μ 2 e, μ 0 1 e, μ 0 | 0-3 b 1 1-2 b 1 0-2 jets 2 2 jets 3 2 b 1 1 b 2 no-jet/c-tag 1 1 b 1 | Yes 20.1 Yes 20.3 Yes 4.7 Yes 20.3 Yes 20.3 Yes 20.1 Yes 20.1 Yes 20.1 Yes 20.3 Yes 20.3 Yes 20.3 | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ | <90 GeV =2 $m(\tilde{\chi}_1^0)$ =55 GeV , $m(\tilde{t}_1)$ << $m(\tilde{\chi}_1^\pm)$ =1 GeV <200 GeV , $m(\tilde{t}_1^+)$ =5 GeV =0 GeV =0 GeV $m(\tilde{k}_1^0)$ =85 GeV >500 GeV $m(\tilde{k}_1^0)$ =85 GeV >750 GeV <200 GeV | 1308.2631 1404.2500 1208.4305, 1209.2102 1403.4853 1403.4853 1308.2631 1407.0583 1406.1122 1407.0608 1403.5222 |
| EW | $\begin{array}{c} \begin{array}{c} \overline{\ell}_{LR}\overline{\ell}_{LR},\overline{\ell}\rightarrow \ell \chi_1^0 \\ \overline{\chi}_1^+\overline{\chi}_1^-,\overline{\chi}_1^+\rightarrow \overline{\ell}\nu(\ell \bar{\nu}) \\ \overline{\chi}_1^+\overline{\chi}_1^-,\overline{\chi}_1^+\rightarrow \overline{\ell}\nu(\ell \bar{\nu}) \\ \overline{\chi}_1^+\overline{\chi}_2^-,\overline{\chi}_1^+,\overline{\ell}\nu(\bar{\nu}) \\ \overline{\chi}_1^+\overline{\chi}_2^-,\overline{\chi}_1^+\overline{\ell}\nu(\bar{\nu}) \\ \overline{\chi}_1^+\overline{\chi}_2^-,\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+\overline{\chi}_2^+\overline{\chi}_2^+\overline{\chi}_2^+,\overline{\chi}_2^+$ | 2 e, µ 2 e, µ 2 τ 3 e, µ 2-3 e, µ 1 e, µ 4 e, µ | 0 \\ - \\ 0 \\ 0 \\ 2 b \\ | Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 Yes 20.3 | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ | =0 GeV =0 GeV, $m(\tilde{\ell}, \tilde{\gamma})$ =0.5 $(m(\tilde{k}_1^+)+m(\tilde{k}_1^0))$ =0 GeV, $m(\tilde{\tau}, \tilde{\gamma})$ =0.5 $m(\tilde{k}_1^+)+m(\tilde{k}_1^0))$ $m(\tilde{k}_1^0)$ =0, $m(\tilde{\ell}, \tilde{\gamma})$ =0.5 $m(\tilde{k}_1^0)+m(\tilde{k}_1^0))$ = $m(\tilde{k}_2^0)$, $m(\tilde{k}_1^0)$ =0, sleptons decoupled = $m(\tilde{k}_2^0)$, $m(\tilde{k}_1^0)$ =0, sleptons decoupled $n(\tilde{k}_1^0)$ =0, $m(\tilde{\ell}, \tilde{\gamma})$ =0.5 $m(\tilde{k}_2^0)+m(\tilde{k}_1^0))$ | 1403.5294 1403.5294 1407.0350 1402.7029 1403.5294, 1402.7029 ATLAS-CONF-2013-093 1405.5086 |
| Long-lived | Direct $\tilde{X}_1^+\tilde{X}_1^-$ prod., long-lived \tilde{X}_1^+ Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{X}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(\tilde{e}, \tilde{\mu}) + \tilde{\tau}(\tilde{e}, $ | 0 1 | 1-5 jets \ | Yes 20.3 Yes 27.9 - 15.9 Yes 4.7 - 20.3 | $ar{m{g}}$ 832 GeV $\mbox{m}(ar{m{k}}^0)$ 10 <tar> $ar{m{k}}^0$ 230 GeV 0.4<τ()</tar> |)-m($\tilde{\chi}_1^0$)=160 MeV, τ($\tilde{\chi}_1^+$)=0.2 ns =100 GeV, 10 μs<τ(\tilde{g})<1000 s n β <50 ($\tilde{\chi}_1^0$)<2 ns cτ<156 mm, BR(μ)=1, m($\tilde{\chi}_1^0$)=108 GeV | ATLAS-CONF-2013-069 1310.6584 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092 |
| //00 | $ \begin{array}{l} LFV \ pp \rightarrow \tilde{\mathbf{v}}_{\tau} + X, \tilde{\mathbf{v}}_{\tau} \rightarrow e + \mu \\ LFV \ pp \rightarrow \tilde{\mathbf{v}}_{\tau} + X, \tilde{\mathbf{v}}_{\tau} \rightarrow e(\mu) + \tau \\ Bilinear \ RPV \ CMSSM \\ \tilde{\mathbf{X}}_{1}^{+} \tilde{\mathbf{X}}_{1}^{-}, \tilde{\mathbf{X}}_{1}^{+} \rightarrow W\tilde{\mathbf{X}}_{0}^{+}, \tilde{\mathbf{X}}_{0}^{+} \rightarrow ee\tilde{\mathbf{v}}_{\mu}, e\mu\tilde{\mathbf{v}}_{e} \\ \tilde{\mathbf{X}}_{1}^{+} \tilde{\mathbf{X}}_{1}^{-}, \tilde{\mathbf{X}}_{1}^{+} \rightarrow W\tilde{\mathbf{X}}_{0}^{+}, \tilde{\mathbf{X}}_{0}^{+} \rightarrow ee\tilde{\mathbf{v}}_{\mu}, e\mu\tilde{\mathbf{v}}_{e} \\ \tilde{\mathbf{y}} \rightarrow qqq \\ \tilde{\mathbf{g}} \rightarrow \tilde{\mathbf{i}}_{1}t, \tilde{\mathbf{i}}_{1} \rightarrow bs \end{array} $ | $\begin{array}{c} 4 \ e, \mu \\ 3 \ e, \mu + \tau \\ 0 \end{array}$ | - 0-3 b \ - - \ - 6-7 jets | - 4.6 - 4.6 Yes 20.3 Yes 20.3 Yes 20.3 - 20.3 Yes 20.3 | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ | 0.10, λ_{132} =0.05 0.10, $\lambda_{1(2)33}$ =0.05 $m(\hat{g})$, c_{TLSP} <1 mm >0.2× $m(\hat{\tau}_1^2)$, λ_{121} ≠0 >0.2× $m(\hat{\tau}_1^2)$, λ_{133} =0 =BR(b)=BR(c)=0% | 1212.1272 1212.1272 1404.2500 1405.5086 1405.5086 ATLAS-CONF-2013-091 1404.250 |
| 104+C | _ | 2 e, μ (SS) 0 π | 2 b \ | - 4.6 Yes 14.3 Yes 10.5 | sgluon 350-800 GeV | imit from 1110.2693 :80 GeV, limit of<687 GeV for D8 | 1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147 |
| | $\sqrt{s} = 7 \text{ TeV}$ full data | \sqrt{s} = 8 TeV partial data | $\sqrt{s} = 8 \text{ To}$ full data | | 10 ⁻¹ 1 | Mass scale [TeV] | |

^{*}Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.



Exotics searches: lower limits





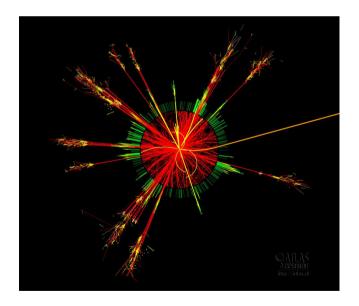
^{*}Only a selection of the available mass limits on new states or phenomena shown



Summary and outlook



- ◆ Huge amount of work has been done by CERN experiments
- Antimatter has been created and studied in some detail
- ◆ The Higgs boson discovered in 2012 so far looks like the Standard Model Higgs
- ◆ The Standard Model is standing strong no SUSY, no sign of any exotics either...
- ◆ Some data still to be analysed, and much more data is still to come
- ♦ Hoping for many fascinating discoveries in the near future!





Web Resources



- 1. Lancaster Particle Physics Package for A-level students:
 - http://www.hep.lancs.ac.uk/package/
 Some basic stuff worth a look or two (feedback welcome)
- 2. Paricle Physics in the UK website, plenty of info and links: http://hepweb.rl.ac.uk/ppUK/
- 3. FNAL (Fermi National Accelerator Laboratory), home of the Tevatron: http://www.fnal.gov/
- 4. CERN (European Centre for Nuclear Research), home of LEP and LHC: http://public.web.cern.ch/public/
- 5. The ultimate resource: Particle Data Group website http://pdg.lbl.gov The official reference for all particle data. Many useful review articles, too