

CE ν NS as a probe of Z' through kinetic and mass mixing effects

Shu Liao

Collaboration with: M. Abdullah, J. Dent, B. Dutta, G. Kane, L. Strigari
arxiv: 1803.01224

Texas A&M University

Outline

Motivation

- Coherent elastic neutrino-nucleus elastic scattering

- Kinetic and mass mixing

Bounds

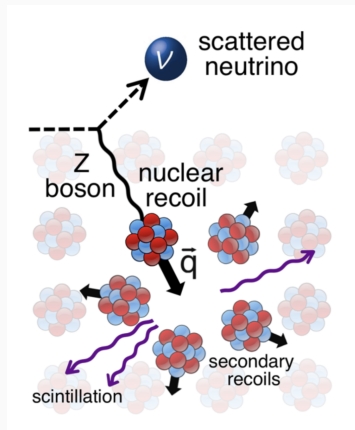
- Existing bounds

- Current bounds and future projection from CE ν NS experiments

Conclusion

Motivation

Coherent elastic neutrino-nucleus elastic scattering ($\text{CE}\nu\text{NS}$)



Coherent scattering of neutrinos from nuclei^{*}

^{*}from Akimov, D., et al. "Observation of coherent elastic neutrino-nucleus scattering." *Science* 357.6356 (2017): 1123-1126.

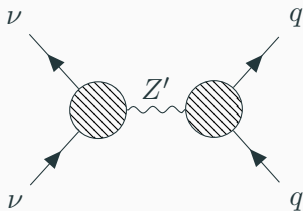
Differential cross section

Differential cross section of $\nu_\beta + N \rightarrow \nu_\alpha + N$ as a function of recoil energy E_R :

$$\frac{d\sigma}{dE_R} = \frac{G_F^2 Q_V^2}{2\pi} m_N \left(1 - \left(\frac{m_N E_R}{E_\nu^2} \right) + \left(1 - \frac{E_R}{E_\nu} \right)^2 \right) F(q^2)$$

$$Q_V^2 = \left[Z \left(\frac{1}{2} - 2 \sin^2 \theta_w \right) + N \left(-\frac{1}{2} \right) + (\text{BSMcharges}) \right]^2$$

Mixing with hypercharge gauge group



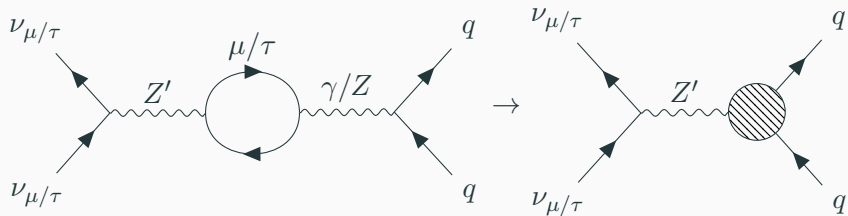
$$L_{\text{gauge}} = -\frac{1}{4}F_a^{\mu\nu}F_{a\mu\nu} - \frac{1}{4}F_b^{\mu\nu}F_{b\mu\nu} - \frac{\epsilon}{2}F_a^{\mu\nu}F_{b\mu\nu}$$

Symmetry breaking

$$\mathcal{L}_{\text{int}} = -\frac{g}{c_w} c_\alpha (t_\alpha + \eta s_w) \left(\tau_3 - \frac{t_\alpha + \eta/s_w}{t_\alpha + \eta s_w} s_w^2 Q \right) Z'_\mu \bar{f} \gamma^\mu f$$

1. Dark Z boson: $\eta \sim 0$, the interaction coupling becomes:
 $-\frac{g}{c_w} s_\alpha (\tau_3 - s_w^2 Q)$
2. Dark hypercharge boson: $s_\alpha \sim 0$, $c_\alpha \sim 1$, the interaction coupling becomes: $\frac{g}{c_w} s_w \epsilon (\tau_3 - Q) = \epsilon g' Q_Y$

$L_\mu - L_\tau$ model



Effective coupling between quark and Z' :

$$\frac{8e^2 g'}{(4\pi)^2} \frac{1}{3} \ln \frac{m_\tau}{m_\mu}$$

Bounds

1. Fixed target experiment[†]
2. Solar neutrino experiment[‡]
3. Atomic parity violation[§]
4. Neutrino trident production[¶]

[†]Phys. Rev. D38, 3375 (1988), Phys. Rev. Lett. 67, 2942 (1991)

[‡]1707.09279

[§]Phys. Rev. D85, 115019 (2012)

[¶]Phys. Rev. Lett. 66, 3117 (1991)

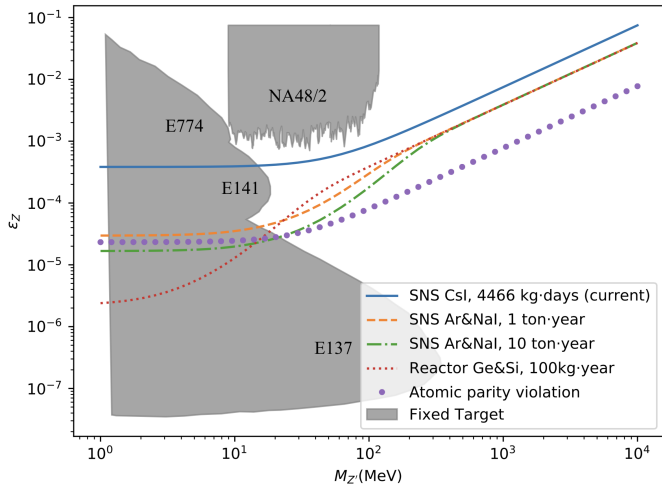
Analyzing tools

$$\chi^2 = \sum_{\text{bins, detectors}} \frac{(N_{exp} - (1 + \beta) N_{pred})^2}{N_{bg} + N_{exp}} + \left(\frac{\beta}{\sigma_{\beta}} \right)^2$$

Name	Detector	Source	Exposure	Threshold
Current (COHERENT)	CsI	SNS (20m)	4466 kg.days	4.25 keV
Future (reactor)	Ge	1GW reactor (20m)	10^5 kg.days	100 eV
	Si	1GW reactor (20m)	10^5 kg.days	100 eV
Future (accelerator)	NaI	SNS (20m)	1 or 10 ton.year	2 keV
	Ar	SNS (20m)	1 or 10 ton.year	30 keV

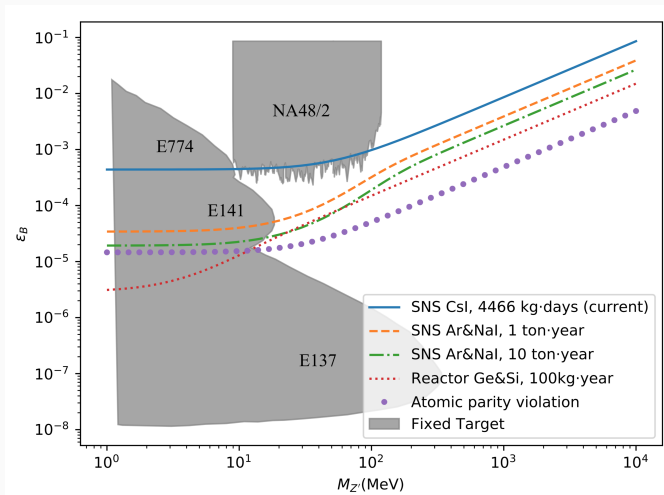
Experimental configurations used in this analysis

Current bounds and future projection for dark Z



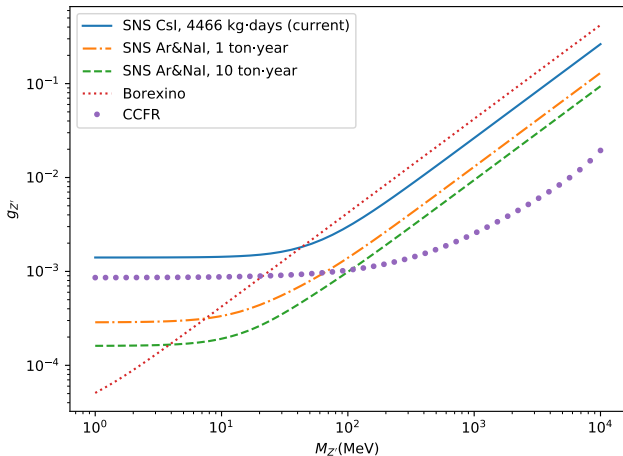
The current and future bounds on the mixing ϵ_Z in the dark Z case are plotted as a function of the Z' mass $M_{Z'}$.

Current bounds and future projection for dark hypercharge



The current and future bounds on the mixing ϵ_B in the dark hypercharge case are plotted as a function of the Z' mass $M_{Z'}$.

Current bounds and future projection for $L_\mu - L_\tau$



The current and future bounds on the coupling $g_{Z'}$ in the $L_\mu - L_\tau$ model are plotted as a function of the Z' mass $M_{Z'}$.

Conclusion

Conclusion

- The capability of $\text{CE}\nu\text{NS}$ to probe light $U(1)$ gauge boson through mixing effect is complementary to fixed target experiments and competitive to atomic parity violation experiments.
- $\text{CE}\nu\text{NS}$ dominates Brexino and CCFR in $L_\mu - L_\tau$ scenario where mixing is generated through fermion loop.