### UNCERTAINTIES IN THE TAGS MEASUREMENTS M. Estienne<sup>1</sup>, M. Fallot<sup>1</sup>, L. Giot<sup>1</sup>, V. Guadilla<sup>1</sup>, L. Le Meur<sup>1</sup>, A. Porta<sup>1</sup> And the TAGS collaboration <sup>1</sup> SEN group, SUBATECH (CNRS/IN2P3, Ecole des Mines de Nantes, Université de Nantes),

4, rue A. Kastler, 44307 Nantes cedex 3, France fallot@subatech.in2p3.fr

### « Workshop NACRE » Cadarache June 2018





### Outline

- Introduction: the recent experiments & their setup
- The TAGS technique
- Data Analysis
- Contaminants
- Detector Response
- Study of impact of sources of systematic uncertainties in the case of <sup>92</sup>Rb (Z. Issoufou 's PhD), and <sup>87,88</sup>Br, <sup>94</sup>Rb (Rocinante detector)
- Study of impact of sources of systematic uncertainties with the DTAS detector (V. Guadilla's (2017) and L. Le Meur's (2018) PhDs)



# TAGS Solution to Pandemonium Effect

#### Pandemonium effect\*\* :

Due to the use of Ge detectors to measure the decay schemes: lower efficiency at higher energy  $\rightarrow$  underestimate of  $\beta$  branches towards high energy excited states: overestimate of the high energy part of the FP  $\beta$  spectra

 $\Rightarrow$  Solution is Total Absorption γ-ray Spectroscopy (TAGS) Big cristal, 4π => A TAGS is a calorimeter !



\*\* J.C.Hardy et al., Phys. Lett. B, 71, 307 (1977)



2 TAGS arrays developed by the Valencia team (Spain, B. Rubio, J.L. Tain, A. Algora et al.): Rocinante (12 BaF2) & DTAS (18 NaI)

# **TAGS** Experimental Campaigns

- Decay Total Absorption Spectrometer (DTAS) for FAIR (IFIC Valencia): used in Jyväskylä in Feb.
   2014 for our reactor antineutrino proposal: 18 modules 15x15x25 cm<sup>3</sup> Nal(TI) + 5" PMT
  - 12 nuclei for antineutrinos measured & 11 for decay heat
  - See V. Guadilla's talk at JEFF/CHANDA meeting in Nov. 2017
- BAF<sub>2</sub> TAGS (Surrey-Valencia): used for the 2009 measurement at IGISOL-JYFLTRAP: <sup>86</sup>Br, <sup>87</sup>Br, <sup>88</sup>Br, <sup>91</sup>Rb, <sup>92</sup>Rb, <sup>93</sup>Rb, <sup>94</sup>Rb
  - <sup>92,93</sup>Rb results already shown at last meetings, see
     A. Zakari-Issoufou et al., PRL 115, 102503 (2015)
  - □ <sup>87</sup>Br, <sup>88</sup>Br, <sup>94</sup>Rb E. Valencia et al. PRC 95 024320 (2017)
  - □ <sup>86</sup>Br, <sup>91</sup>Rb S. Rice et al. PRC 96 014320 (2017)

# Antineutrino Proposal in Jyväskylä: Subatech-IFIC collaboration

V.Guadilla et al.,, Nucl. Inst. and Meth. B, in press. Online (2015) : http:// www.sciencedirect.com/science/article/pii/S0168583X15012628

Pure beams required: Use of the double Penning trap from JYFL



2 TAGS arrays developed by the Valencia team (Spain, B. Rubio, J.L. Tain, A. Algora et al.):

## **TAGS** Technique

⇒ Solution is Total Absorption γ-ray Spectroscopy (TAGS) Big cristal,  $4\pi$  => A TAS is a calorimeter !

**Observable:**  $\beta$ **-intensity** =>  $\beta$ **-strength:** 



- Spectrum must be clean
- Response must be accurately known
- Solution of inverse problem must be stable



2 TAGS arrays developed by the Valencia team (Spain, B. Rubio, J.L. Tain, A. Algora et al.): Rocinante (12 BaF2) & DTAS (18 NaI)

NIM A430 (1999) 333 NIM A571 (2007) 719 NIM A430 (1999) 488 NIM A571 (2007) 728

# Total Absorption Spectroscopy (TAS)

Big cristal,  $4\pi => A$  TAS is a calorimeter !



 $d_i =$ 





### **Observable:**

β-intensity => β-strength: An ideal TAS would give directly the β-intensity I<sub>β</sub> which is linked with the β-strength S<sub>β</sub>:

$$S_{i} = \frac{I_{i}}{f(Q_{\beta} - E_{i})T_{1/2}} \left[s^{-1}\right]$$

Statement of the problem:

Relation between TAS data and the  $\beta$ -intensity distribution:

$$\sum_{j} R_{ij} f_{j} \implies R_{j} =$$

Deconvolution (Inverse problem) algorithms

 $\mathbf{R}_{\mathbf{j}} = \sum_{k=0}^{j-1} b_{jk} \mathbf{g}_{\mathbf{j}\mathbf{k}} \otimes \mathbf{R}_{\mathbf{k}}$ 

Monte Carlo simulations + Nuclear statistical model

- Spectrum must be clean
- Response must be accurately known
- Solution of inverse problem must be stable

NIM A430 (1999) 333NIM A571 (2007) 719NIM A430 (1999) 488NIM A571 (2007) 728

# Experimental setup at Jyväskylä (142Cs, 99Y, 138I, 96,96mY)



V. Guadilla et al., Nucl. Instrum. Methods B 376 (2016), p. 334



- DTAS = 18 crystals of NaI(TI)
  - ~90% efficiency for a 1 MeV gamma
  - ΔE/E ~ 5% at 1.3 MeV
- β detector = plastic detector
  - In coincidence with γ → suppression of the background
  - 30% detection efficiency
  - HPGe detector
    - Allow identification of possible contaminants coming from the decay chain

Why Jyväskylä IGISOL-4 facility ?

- → Because of the JYFLTRAP, a double Penning Trap
- $\rightarrow$  Mass resolution of  $\delta m/m \sim 10^{\text{-6}}$
- → A very pure beam is needed

L. Le Meur's PhD



→ Solve the Inverse Problem Solved by an iterative procedure based on the **f**<sub>i</sub>: β feeding Bayes Theorem J.L. Tain, D. Cano-Ott, Nucl. Inst. and Meth. in Phys. Res. A 571 (2007) 728 Energy E, depends on the knowledge of the  $r_j = \sum_{k=0}^{\infty} b_{jk} g_{jk} * r_k$ daughter convolution : Models dependent  $R_i^{\beta^-} = b^- * r_i$ (Strengths E1/E2/M1, levels density) r<sub>k</sub> : response for k level  $\mathbf{b}_{ik}$ : branching ratio for  $j \rightarrow k$  transition  $\mathbf{g}_{ik}$ :  $\gamma$ -response for  $j \rightarrow k$  transition  $\gamma$ -response  $\mathbf{g}_{ik}$  and  $\beta$ -response **b** are simulated (Geant4)

L. Le Meur's PhD

## Data analysis

Aim of TAS analysis =  $\beta$  feeding

### → Requires clean spectrum

### → Response matrix calculation (R<sub>ii</sub>)

probability that feeding at a level j gives counts in data channel *i* of the data spectrum

Recursive

- **b**<sup>•</sup> : response to the e<sup>•</sup> emission in
- $\beta$ -decay to level j

- **b**<sub>ik</sub> calculation separates in two part :
  - Known (discrete) / Unknown (continue)



 $d_i = \sum R_{ij} f_j$ 

d, : Experimental data

**R**<sub>..</sub>: Detector response matrix

### Background

- Spectra are dominated by  $\alpha$  peaks
- Spectra recorded in between sources
- Normalisation by acquisition time





### Pile-Up

- False events stored in data
- constrained random coincidences
- 1<sup>st</sup> order normalisation

Zakari-Issoufou's PhD

### Contamination from daughters should be well estimated and removed

$${}^{92}\text{Rb}(4.48s) \rightarrow {}^{92}\text{Sr}(2.6h) \rightarrow {}^{92}\text{Y}$$
  
$${}^{93}\text{Rb}(5.84s) \rightarrow {}^{93}\text{Sr}(7.4mn) \rightarrow {}^{93}\text{Y}$$
  
$${}^{93}\text{Rb}(1.39\%) \rightarrow {}^{92}\text{Sr}(2.6h) \rightarrow {}^{92}\text{Y}$$



- TAS spectra from daughter is calculated  $\Rightarrow \vec{d} = R\vec{f}$
- Integral dictated by Bateman

neutron emission is simulated

• 
$$N_{\beta n\gamma} = P_n N_d \varepsilon_{\beta n\gamma}$$

Zakari-Issoufou's PhD

→ Solve the Inverse Problem Solved by an iterative procedure based on the **f**<sub>i</sub>: β feeding Bayes Theorem J.L. Tain, D. Cano-Ott, Nucl. Inst. and Meth. in Phys. Res. A 571 (2007) 728 Energy E, depends on the knowledge of the  $r_j = \sum_{k=0}^{\infty} b_{jk} g_{jk} * r_k$ daughter convolution : Models dependent  $R_i^{\beta^-} = b^- * r_i$ (Strengths E1/E2/M1, levels density) r<sub>k</sub> : response for k level  $\mathbf{b}_{ik}$ : branching ratio for  $j \rightarrow k$  transition  $\mathbf{g}_{ik}$ :  $\gamma$ -response for  $j \rightarrow k$  transition  $\gamma$ -response  $\mathbf{g}_{ik}$  and  $\beta$ -response **b** are simulated (Geant4)

L. Le Meur's PhD

## Data analysis

Aim of TAS analysis =  $\beta$  feeding

### → Requires clean spectrum

### → Response matrix calculation (R<sub>ii</sub>)

probability that feeding at a level j gives counts in data channel *i* of the data spectrum

Recursive

- **b**<sup>•</sup> : response to the e<sup>•</sup> emission in
- $\beta$ -decay to level j

- **b**<sub>ik</sub> calculation separates in two part :
  - Known (discrete) / Unknown (continue)



 $d_i = \sum R_{ij} f_j$ 

d, : Experimental data

**R**<sub>..</sub>: Detector response matrix



$$\begin{split} T_{XL}(E_{\gamma}) &= 2\pi E_{\gamma}^{2L+1} f_{XL}(E_{\gamma}) \\ f_{XL}(E_{\gamma}) &= \frac{26\times 10^{-8}}{2L+1} \sigma_0 \Gamma E_{\gamma}^{3-2L} \frac{\Gamma_0}{(E_{\gamma}^2 - E_0^2)^2 + E_{\gamma}^2 \Gamma_0^2} \\ T_{XL} \ : \ \text{Transmission coefficient} \\ f_{XL} \ : \ \text{Strength function}, \ XL &= \text{E1,M1,E2} \\ E_{\gamma} \ : \ \text{Transition energy} \end{split}$$

Resonance parameters  $\sigma_0, \Gamma_0, E_0$  are obtained from RIPL https://www-nds.iaea.org/RIPL-3

### Zakari-Issoufou's PhD

#### Level densities



Egidy et al. PRC 72 (2005), Dilg et al. NPA 212 (1973), Gilbert et al. Canadian Journal of Physics 43 (1965)

- Semi-empiric models for levels densities:
  - Back Shifted Fermi Gas
  - Constant Temperature
  - Gilbert Cameron
- Parameters obtained from fits
  - experimental data for known
  - theoretical calculation for unknown (HFB)
- Chose the one that is more in agreement with experimental data in the lower energy part

A.-A. Zakari-Issoufou's PhD

E. Valencia, S. Rice Phd work







- Precise description of the TAS geometry
- $\bullet~{\rm GEANT4}$  simulation of  $\beta$  and  $\gamma$

e c

 Compare simulated spectra to experiment for validation

Zakari-Issoufou's PhD

### The case of <sup>92</sup>Rb







Jobotech

# Impact of sources of systematic uncertainties



### A.-A. Zakari-Issoufou's PhD

FIGURE 3.33 – Courbe des valeurs de  $\chi^2$  obtenues lors de la résolution du problème inverse pour différents paramètres de normalisation des empilements et différents taux de contamination du <sup>92</sup>Sr. On trouve le minimum pour un facteur de 0.90 de la valeur que nous avions estimée par le calcul pour les empilements.

#### Impact de la normalisation des empilements et de la contamination du noyau fils



**FIGURE 3.34** – Comparaison de la distribution d'alimentation bêta obtenue (rouge) sous les conditions de normalisation des contaminants qui minimisent la valeur de  $\chi^2$  entre le spectre reconstruit et les données expérimentales et le résultat que nous avions adopté (bleu). En bas l'écart relatif entre les deux distributions, avec des variations plus prononcées en fin de spectre où la contribution des empilements est relativement plus importante.

## Impact des paramètres d'étalonnage



FIGURE 3.35 – Comparaison des alimentations bêtas obtenues (rouge) en modifiant la pente de la relation d'étalonnage en énergie de 1% avec le résultat adopté (bleu). Les écart relatifs sont compatibles avec les incertitudes statistiques.

#### Impact de l'étalonnage



FIGURE 3.36 – Comparaison des alimentations bêtas obtenues (rouge) en prenant les bornes supérieures des paramètres de la relation d'étalonnage en résolution avec le résultat adopté (bleu). L'impact est plus prononcé que celui de l'étalonnage en énergie.

## Impact de l'épaisseur du Si



FIGURE 3.37 – En (a) superposition entre les données expérimentales et la reconstruction à partir des alimentations bêtas obtenues lorsque l'épaisseur du silicium est augmenté de 10%. En (b) cette alimentation (rouge) est comparée avec l'alimentation bêta de référence (bleu).

## Impact des densités de niveau



FIGURE 3.38 – Comparaison des alimentations bêta obtenue pour le <sup>92</sup>Rb, en utilisant des densités de niveaux en moyenne 2 fois plus élevées que celle utilisées dans l'analyse de référence dont les alimentations sont représentées en bleu.

# Tableau récapitulatif

# Pour l'instant : Somme quadratique des incertitudes, malgré les corrélations entre paramètres

$E_{level}$	$I_{\beta}$	Incertitudes[%]						Total[0/1
		Statistique	Soustraction	$E_{Cal}$	$R_{Cal}$	$\Delta_{Si}$	Densité	10(a1[70]
0	87.5	0.23	0.16	0.02	0.05	2.77	0.57	2.84
814.97998	1.05836	0.43	0.80	0.21	1.87	7.40	13.70	15.71
1384.79004	0.713029	0.40	3.47	0.24	1.44	10.81	8.04	14.00
1778.32996	0.69934	0.37	0.71	0.67	0.15	23.25	-0.98	23.29
1860	0.0956618	0.41	0.76	0.88	3.25	0.04	-2.33	4.19
1900	0.0765636	0.41	1.22	0.03	1.37	16.95	-19.52	25.92
1940	0.082986	0.41	1.35	0.38	0.50	25.45	-23.74	34.83
1980	0.110346	0.40	1.41	0.63	0.23	32.55	-20.89	38.71
2020	0.133679	0.41	1.48	0.71	0.57	29.90	-19.18	35.57
2060	0.14103	0.41	1.65	0.68	0.64	23.69	-18.09	29.87
2100	0.126544	0.41	1.89	0.50	0.25	18.44	-18.32	26.07
2140	0.0972783	0.41	2.21	0.19	0.67	15.82	-20.34	25.88
2180	0.082471	0.41	2.35	0.02	1.23	13.80	-22.30	26.36
2220	0.0768429	0.41	2.37	0.05	1.08	13.80	-24.53	28.27
2260	0.0793313	0.41	2.31	0.37	0.13	16.83	-25.22	30.41
2300	0.0831362	0.41	2.25	0.60	0.54	20.13	-26.21	33.14
2340	0.0859162	0.41	2.15	0.69	0.74	26.55	-26.73	37.75
2380	0.0801968	0.41	2.11	0.55	0.19	33.03	-28.13	43.44
2420	0.0729589	0.41	2.06	0.46	0.06	39.35	-29.77	49.39
2460	0.0670209	0.41	1.97	0.60	0.28	45.28	-30.58	54.68
2500	0.060217	0.41	1.85	0.67	0.56	45.85	-30.35	55.03
2540	0.0556243	0.41	1.72	0.48	0.05	42.96	-29.86	52.35

Impact of several sources of systematic uncertainty on the shape of the intensity distribution.

In total the blue areas are 14 solutions for <sup>87</sup>Br, 13 for <sup>88</sup>Br, and 15 for <sup>94</sup>Rb.



J. Agramunt et al., Nucl. Instrum. Methods Phys.  $\operatorname{Res}^{10^{\circ}}_{...}$ Sect. A 807, 69 (2016). E. Valencia et al. PRC 95 024320 (2017) FIG. 4.

FIG. 4. (Color online) Beta intensity distributions from TAGS. The thin black line is the adopted solution, the light blue filled region indicates the spread of solutions due to the systematic effects investigated. See text for details.

## Experiment @ Jyväskylä in 2014



#### **Reactor antineutrino proposal:**

- 12 nuclei for antineutrinos measured & 11 for decay heat
- First use of new DTAS (17+1 Nal) developed by IFIC (Valencia)
- Successful use of the new IGISOL-4 facility
- First use of precision trap with IGISOL-4

V.Guadilla et al.,, Nucl. Inst. and Meth. B, in press. Online (2015) : http://www.sciencedirect.com/ science/article/pii/S0168583X15012628 Analysis on-going

### How?

- Geant4 simulation code
- Event generator for the sources
- Geometry of the set-up
- Non-proportionality NaI(TI)







イロト イヨト イヨト イヨ

-

25

### How?

- Geant4 simulation code
- Event generator for the sources
- Geometry of the set-up
- Non-proportionality NaI(TI)







< ロ ト < 同 ト < 三 ト

JEFF-CHANDA ND Workshop

21/11/2017 9 / 43

26

### How?

- Geant4 simulation code
- Event generator for the sources
- Geometry of the set-up
- Non-proportionality NaI(TI)



< ロ > < 同 > < 三 >

### Segmentation can be useful to study different multiplicities:



21/11/2017 9 / 43

27

### How?

- Geant4 simulation code
- Event generator for the sources
- Geometry of the set-up
- Non-proportionality NaI(TI)

LaBr<sub>3</sub>:Ce J.L. Tain et al., NIMA 774 (2015) 17

**BaF**<sub>2</sub> E. Valencia et al., PRC 95 (2017) 024320

### **Nal(TI)** $\rightarrow$ Response of DTAS to neutrons (inelastic, capture...):



Víctor Guadilla

JEFF-CHANDA ND Workshop

21/11/2017 9 / 43

## <sup>137</sup>I: errors

Following E. Valencia et al., PRC 95 (2017) 024320: envelope of solutions compatible with data

- Daughter normalization
- $\beta$ -n normalization
- Summing-pileup normalization
- Known level scheme
- Deconvolution algorithm
- $\beta$  detector efficiency
- β-gated vs. singles vs.
   neutron veto



• Branching ratio variations: parameters of level density,  $\gamma$ -strength function, spin-parity values of known level scheme, better reproduction of known  $\gamma$ -intensities etc.

21/11/2017 17 / 43

- - E - E

200

## <sup>137</sup>I: errors

Following E. Valencia et al., PRC 95 (2017) 024320: envelope of solutions compatible with data

- Daughter normalization
- $\beta$ -n normalization
- Summing-pileup normalization
- Known level scheme
- Deconvolution algorithm
- $\beta$  detector efficiency
- β-gated vs. singles vs.
   neutron veto





200

17 / 43

21/11/2017

## L. Le Meur's PhD thesis

Input parameters for branching ratio matrix :

- γ-strength (E1, M1, E2)...
- ... and parametrization  $(E_{R}, \Gamma_{R}, \sigma_{R})$
- Level densities and parameter a
- Nucleus deformation β
- Limit between Known and Unknown part (here = 1880 keV)
  - Depends on the knowledge of the daughter nucleus

Impact of input models for the γ-strength (generalized Lorentzian, QRPA-Gogny v0, 1 and 2 from S. Péru-Désenfants and M. Martini)



# Outlooks

- Study the propagation of uncertainties with correlations
- Extract correlation/covariance matrices
- Include the TAGS data and their uncertainties in the evaluated databases...
- We regularly include the new TAGS data in summation calculations to compare with integral data and extract new priority lists (decay heat, antineutrinos, Beff)

=> cf. Our proposal to the next European Nuclear Data project + next NACRE 2019-2021

### TAS COLLABORATION

IFIC Valencia: A. Algora, B. Rubio, J.A. Ros, V. Guadilla, J.L. Tain, E. Valencia, A.M. Piza, S. Orrigo, M.D. Jordan, J. Agramunt

SUBATECH Nantes: J.A. Briz, M. Fallot, A. Porta, A.-A. Zakari-Issoufou, M. Estienne, T. Shiba, A.S. Cucoanes

U. Surrey: W. Gelletly

IGISOL Jyvaskyla: H. Penttilä, Äystö, T. Eronen, A. Kankainen, V. Eloma, J. Hakala, A. Jokinen, I. Moore, J. Rissanen, C. Weber

CIEMAT Madrid: T. Martinez, L.M. Fraile, V. Vedia, E. Nacher

IPN Orsay: M. Lebois, J. Wilson

BNL New-York: A. Sonzogni Istanbul Univ.: E. Ganioglu Special thanks to the young researchers working in the project:
L. Le Meur, V. Guadilla, A. -A. Zakari-Issoufou, E. Valencia, S. Rice, J.A. Briz
Discussions with from: A. Algora, J. L. Tain, B. Rubio, S. Cormon, A. Cucoanes, M. Estienne, M. Fallot, L. Giot, A. Porta, T. Shiba, ...are acknowledged

### TOTAL ABSORPTION SPECTROSCOPY FOR NUCLEAR STRUCTURE, NUCLEAR ASTROPHYSICS, NEUTRINO AND REACTOR PHYSICS

Spokespersons: A. Algora, M. Fallot, A. Porta, B. Rubio, J.-L. Taín

And collaborators:

A.S. Cucoanes, M. Estienne, A. Onillon, A.-A. Zakari-Issoufou

1 Subatech, CNRS/IN2P3, University of Nantes, EMN, Nantes, France

J. Agramunt, C. Domingo-Pardo, V. Guadilla, A. Montaner, S. Orrigo

2 IFIC,CSIC-Univ. Valencia, Valencia, Spain

#### M. Martini

Department of Physics and Astronomy, Ghent University, Proeftuinstraat 86, B-9000 Gent, Belgium

CEA, DAM, DIF, F-91680 Arpajon, France

Institut d'Astronomie et d'Astrophysique, CP-226, Université Libre de Bruxelles, B-1050, Brussels, Belgium **S. Péru** 

CEA, DAM, DIF, F-91680 Arpajon, France

#### A. Fantina, S. Goriely

Institut d'Astronomie et d'Astrophysique, CP-226, Université Libre de Bruxelles, B-1050, Brussels, Belgium

#### J. Margueron

Institut de Physique Nucléaire de Lyon, Villeurbanne, France

#### S. Franchoo, A. Gottardo, P. Halipré, E. Khan, M. Lebois, D. Verney, J. Wilson

Institut de Physique Nucléaire d'Orsay, Orsay, France

Zs. Dombrádi, D. Sohler, M. Csatlos, A. Krasznahorkay, Zs. Fülöp

Inst. of Nuclear Research of the Hungarian Academy of Sciences, Debrecen, Hungary

#### P. Regan, Zs. Podolyák, W. Gelletly

Univ. of Surrey, Guildford, United Kingdon

#### S. Judge, C. Larijani

The National Physics Laboratory, Teddington, Surrey, UK

#### Ela Ganioglu

Department of Physics, University of Istanbul, Turkey

#### PAC ALTO Jan. 24 2014





## **Bayes theorem**

Where :

$$P(f_j|d_i) = \frac{P(d_i|f_j)P(f_j)}{\sum_{j=1}^m P(d_i|f_j)P(f_j)}$$

 $P(f_j) \ (\equiv f_j / \sum_{j=1}^m f_j)$  is the a priori probability on the feedings  $f_j$  $P(d_i|f_j)$  is the condition probability that  $d_i$  is due to the parameter  $f_j$  and is equivalent to  $R_{ij}$ 

 $P(f_j|d_i)$  is the a posteriori conditional probability that parameter  $f_i$  was the cause of  $d_i$ 

$$\hat{f}_j = \frac{1}{\sum_{i=1}^n R_{ij}} \sum_{i=1}^n P(f_j | d_i) \hat{d}_i, \quad j = 1, \dots, m \quad \rightarrow \text{Relation between expected feedings and data}$$

Include the latter formula into the first gives an iterative algorithm

$$f_j^{(s+1)} = \frac{1}{\sum_{i=1}^n R_{ij}} \sum_{i=1}^n \frac{R_{ij}f_j^{(s)}d_i}{\sum_{k=1}^m R_{ik}f_k^{(s)}}, \quad j = 1, \dots, m.$$