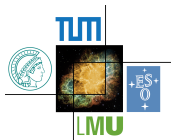


Theoretical prospects on V_{ub} and V_{cb}

Martin Jung



DFG Deutsche
Forschungsgemeinschaft

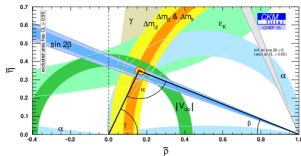
Talk at the XIII Meeting on B Physics: “Synergy between LHC
and SUPERKEKB in the Quest for New Physics”
Marseille, France, 1st of October 2018

Importance of (semi-)leptonic hadron decays

In the Standard Model:

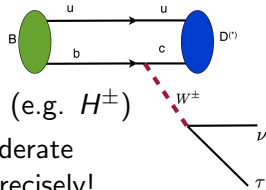
- Tree-level, $\sim |V_{ij}|^2 G_F^2 FF^2$
- Determination of $|V_{ij}|$ (7/9)

➡ This talk



Beyond the Standard Model:

- Leptonic decays $\sim m_l^2$
 - ➡ large relative NP influence possible (e.g. H^\pm)
- NP in semi-leptonic decays small/moderate
 - ➡ Need to understand the SM very precisely!



For instance isospin breaking in $\Upsilon(4S) \rightarrow B\bar{B}$ [MJ'15]

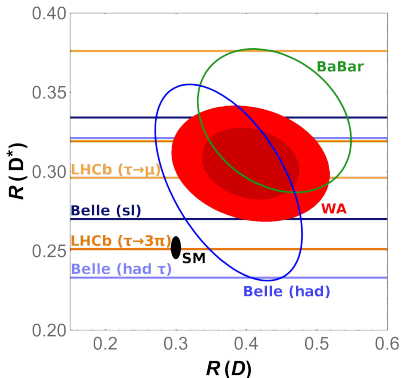
Key advantages:

- Large rates
- Minimal hadronic input \Rightarrow systematically improvable
- Differential distributions \Rightarrow large set of observables

Lepton-non-Universality in $b \rightarrow c\tau\nu$ 2018

[Talks tomorrow by D. Buttazzo + A. Morris]

$$R(X) \equiv \frac{\text{Br}(B \rightarrow X\tau\nu)}{\text{Br}(B \rightarrow X\ell\nu)}$$



contours: 68% CL
filled: 95(68)% CL

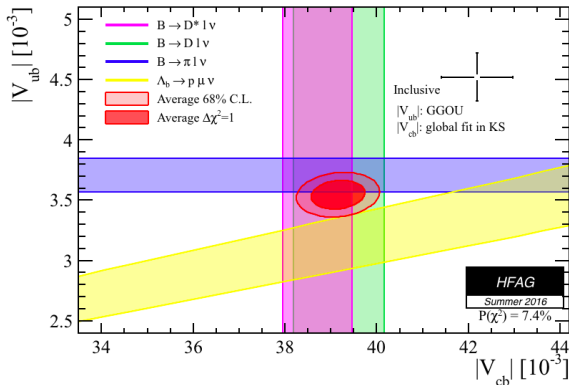
- $R(D^{(*)})$:
2× LHCb, 4× Belle recently
➡ average $\sim 4\sigma$ from SM
- τ -polarization ($\tau \rightarrow \text{had}$) [1608.06391]
- $B_c \rightarrow J/\psi\tau\nu$ [1711.05623] : huge
- Differential rates from Belle, BaBar
- Total width of B_c
- $b \rightarrow X_c\tau\nu$ by LEP
- **New:** $F_L(D^*)$ [Belle@CKM'18]

Present e, μ results **separately!**

$|V_{xb}|$: inclusive versus exclusive

Example of complementarity between LHC and SuperKEKB!

Long-standing problem, motivation for NP [e.g. Voloshin'97]:



- Very hard to explain by NP [Crivellin/Pokorski'15] (but see [Colangelo/de Fazio'15])
- Suspicion: experimental/theoretical systematics?

Comments regarding systematics and fitting [MJ/Straub'18]

Present (and future!) precision renders small effects important:

- d'Agostini effect:
assuming systematic uncertainties \sim (exp. cv) introduces bias
 - ➡ e.g. $1-2\sigma$ shift in $|V_{cb}|$ in Belle 2010 binned data
- Rounding in a fit with strong correlations and many bins:
 - ➡ 1σ between fit to Belle 2017 data from paper vs. HEPdata
- Problem: how to provide unfolded data independent from the (precise) signal hypothesis?
 - ➡ Independent of form factor parametrization (later)
 - ➡ Independent of potential NP contributions
(more severe for $b \rightarrow c\tau\nu$)

BR measurements and isospin violation [MJ 1510.03423]

Detail due to high precision and small NP

➡ Relevant for $\sigma_{\text{BR}}/\text{BR} \sim \mathcal{O}(\%)$

Branching ratio measurements require normalization. . .

- B factories: depends on $\Upsilon \rightarrow B^+B^-$ vs. $B^0\bar{B}^0$
- LHCb: normalization mode, usually obtained from B factories

Assumptions entering this normalization:


- PDG: assumes $r_{+0} \equiv \Gamma(\Upsilon \rightarrow B^+B^-)/\Gamma(\Upsilon \rightarrow B^0\bar{B}^0) \equiv 1$
- LHCb: assumes $f_u \equiv f_d$, uses $r_{+0}^{\text{HFAG}} = 1.058 \pm 0.024$
(also usually used for sl analyses by B factories)

Both approaches problematic:

- Potential large isospin violation in $\Upsilon \rightarrow BB$ [Atwood/Marciano'90]
- Measurements in r_{+0}^{HFAG} assume isospin in exclusive decays
 - ➡ This is one thing we want to test!
 - ➡ Avoiding this assumption yields $r_{+0} = 1.035 \pm 0.038$

Inclusive V_{cb} determination

Inclusive $B \rightarrow X_c \ell \nu$ calculated in operator product expansion:

- Systematic expansion in $1/m_{b,c}$ and α_s
- State of the art: [Alberti/Becher/Bigi/Biswas/Boos/Czarnecki/Ewerth/Gambino/Lunghi/Mannel/Melnikov/Nandi/Pak/Pivovarov/Rosenthal/...]
 - $\mathcal{O}(\alpha_s^0)$: parametrization up to $1/m^5$
(proliferation of hadronic parameters from $1/m^4$)
 - $\mathcal{O}(\alpha_s^1)$: up to $1/m^2$, $1/m^3$ work in progress [Gambino+]
 - $\mathcal{O}(\alpha_s^2)$: leading order
-  Consistent fit, seen e.g. in quark-mass determination

$$|V_{cb}| = (42.00 \pm 0.64) \times 10^{-3} \text{ [Gambino +' 16]}$$

Prospects: [Gambino@CKM]

- α_s/m^3 underway, α_s^3 “feasible” (total rate), necessary?
- Weak+e/m effects require attention (\rightarrow theory vs. experiment)
- Lattice determination of local B matrix elements

[Kronfeld/Simone,Gambino/Melis/Simula]

 Improvements in sight, several steps to qualitative new level

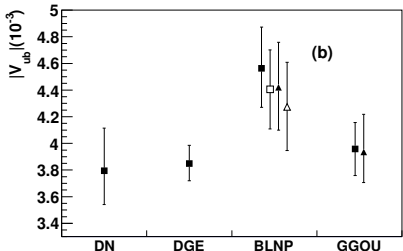
Inclusive V_{ub} determination

In a perfect world: analogous to V_{cb} inclusive.

The difference between theory and practice is that in theory, there is no difference, but in practice there is.

$$|V_{ub}|^2/|V_{cb}|^2 \sim 1\%$$

- ➡ Truly inclusive measurement flooded with $b \rightarrow c$ background
- ➡ Non-local OPE, hadronic **functions** instead of **parameters**
 - Leading shape function **universal**, extracted from $B \rightarrow X_s \gamma$
 - Subleading SF treatment: new approaches NNVub + SIMBA
 - Moments of SFs related to hadronic parameters in $B \rightarrow X_c \ell \nu$



Meanwhile. . .

- New BaBar analysis 2017
- High E_ℓ -region critical
- 3/4 methods: **lower** $|V_{ub}|$
- ➡ Effect in other studies?

Exp+Theo collaboration essential

Exclusive V_{ub} determinations

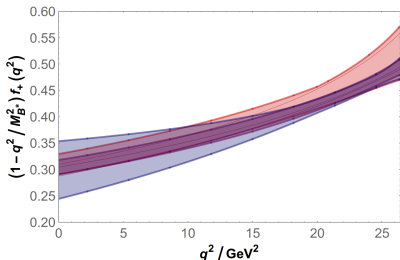
Traditionally: $B \rightarrow \pi \ell \nu$, $B \rightarrow \rho$ (weaker)

➡ Complementarity: $B \rightarrow \pi(\rho)$ @ Belle II, new modes @ LHCb

➡ $|V_{ub}/V_{cb}|$ via $\Lambda_b \rightarrow p$ vs. $\Lambda_b \rightarrow \Lambda_c$, $B_s \rightarrow K^{(*)}$ vs. $B_s \rightarrow D_s$, $B_c(?)$

Larger kinematical range accessible

➡ combine lattice and LCSR via pseudodata / BCL coefficients



Determination over full kin. range

LCSR: NLO twist 2+3, LO higher-twist,

NNLO [Khodjamirian+,Ball+,Bharucha]

Lattice: Immense recent progress!

➡ extending q^2 range, 2+1+1, ...

[Celis/MJ/Li/Pich'17]

$B \rightarrow \rho$: Recent LCSR results [Bharucha+'15], issue: ρ theo/exp

➡ $B \rightarrow \pi \pi \ell \nu$ description [Faller+,Cheng+,Feldmann+,Kim+,Böer+,Kang+,Meißner+]

$\Lambda_b \rightarrow p$: first $|V_{ub}/V_{cb}|$ result, improvement (exp+theo) ongoing

$B_s \rightarrow K$: form factor improvements [Khodjamirian+,FNAL/MILC,RBC/UKQCD]

➡ Excellent prospects, a lot to gain from LHC + SuperKEK-B!

Exclusive $b \rightarrow c$ determinations

Dominated by $B \rightarrow D^{(*)}$, $B_s \rightarrow D_s^{(*)}$ possible (competitive?)

Heavy to heavy transition \rightarrow HQET domain, lattice difficult

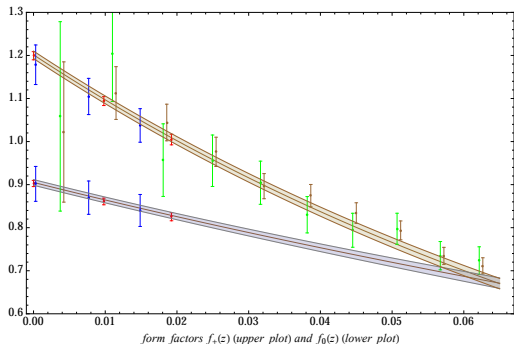
$B \rightarrow D, D^*$ form factors:

- **Unitarity + analyticity** used to obtain expansion in z
 - ↳ $|z| \ll 1$, used in CLN + BGL parametrization alike
 - ↳ BGL: formulation such that coefficients ≤ 1 as well
- CLN: uses **heavy-quark limit** to relate $B^* \rightarrow D^{(*)}$ FFs
 - ↳ Expansion in $1/m_{b,c}$ and α_s (known to NLO)
 - LO: unique function [Isgur/Wise], $1/m$ 3(4) additional functions
 - $1/m^2$ structure known [Falk+], but 1 unknown function per FF
 - ↳ extremely efficient parametrization up to $1/m$ and α_s
- Up to 2015: typically V_{cb} from CLN-parametrization fit
 - ↳ 2 problems: CLN error estimate optimistic + ignored by exp.

Recent discussion: how large are $1/m_c^2$ corrections?

V_{cb} from $B \rightarrow D$

2015: Unfolded $B \rightarrow D\ell\nu$ spectra [Belle] + finite recoil LQCD [HPQCD,MILC]



Analysis by Bigi/Gambino:

- Improved unitarity constraints
 - Lattice data “contradict” CLN (sensitivity to higher $1/m$ orders)
- ➡ $|V_{cb}| = 40.49(96) \times 10^{-3}$, compatible with V_{cb}^{incl} and $B \rightarrow D^*$

V_{cb} from $B \rightarrow D^*$

2017: Prel. unfolded spectrum (4 variables) from Belle

- ➡ However, in this case no finite-recoil FFs available from lattice
- ➡ w/ Belle results SM fit in BGL possible (including lattice (+LCSR))

Results: [Bigi+,Grinstein+]

- Both CLN and BGL yield excellent fits
 - ➡ $|V_{cb}^{\text{CLN}}| = 38.2(15) \times 10^{-3}$
 - ➡ $|V_{cb}^{\text{BGL}}| = 41.7(21)[40.4(17)] \times 10^{-3}$ w/ or w/o LCSR
 - ➡ BGL 1 – 2 σ higher, larger difference than expected!
 - ➡ Intriguing result, but requires confirmation exp. + lattice

[1809.03290]: New Belle result

➡ $|V_{cb}^{\text{CLN}}| = 38.4(2)(6)(5)10^{-3}$ $|V_{cb}^{\text{BGL}}| = 42.3(3)(7)(6)10^{-3}$

Uncertainties due to parametrization were underestimated

➡ Using BGL, there is no indication of a V_{cb} puzzle

➡ Lattice data should resolve the issue within the year

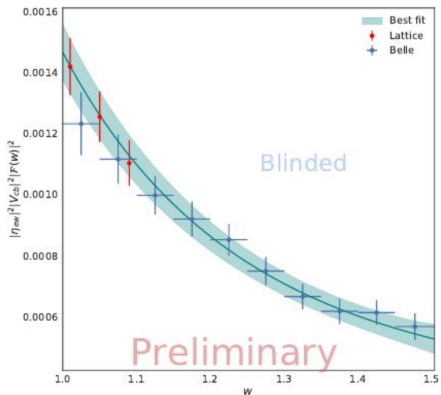
N.B.: This discussion relates to SM $R(D, D^*)$ predictions

Status lattice calculation of $B \rightarrow D^*$

Chris Monahan @ CKM:

FNAL/MILC: first (blind) 2+1 results, on MILC AsqTad ensembles

First result for $R(D^*)$ soon...



$b \rightarrow c$ Form Factors beyond the SM

Only $V_{cb} \times \text{FF}(q^2)$ extracted from data

SM: fit to data + normalization from lattice/LCSR/... $\rightarrow |V_{cb}|$

NP: can affect the q^2 -dependence, introduces additional FFs

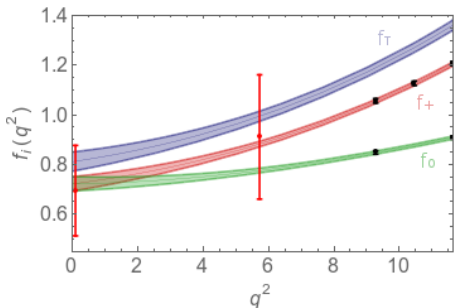
➡ To determine general NP, FF shapes needed from theory

In [MJ/Straub'18], we use all available theory input:

- Unitarity bounds (using results from [BGL,Bigi/Gambino(/Schacht)'16'17])
- LQCD for $f_{+,0}(q^2)$ ($B \rightarrow D$), $h_{A_1}(q^2_{\text{max}})$ ($B \rightarrow D^*$)
[HPQCD'15,'17,Fermilab/MILC'14,'15]
- LCSR for $R_{1,2}(0)$, $h_{A_1}(w = w_{\text{max}}, 1.3)$, $G(w = w_{\text{max}}, 1.3)$ [Faller+'08]

HQET relations up to $\mathcal{O}(\alpha_s, 1/m_{b,c})$ plus $1/m_{c,b}^2$ subset, mostly à la [Bernlocher+'17], but w/o CLN

- relation between slope and curvature



Conclusions

Absence of clear NP signals \rightarrow new challenges

- Issues like isospin breaking now center of attention
- V_{cb}^{incl} stable and still improvable (theory homework)
- V_{ub}^{incl} : BaBar result needs to be understood
 - ➡ No resolution, but “suggestive”
- V_{ub}^{excl} : Theoretical and experimental progress in parallel
 - ➡ New modes + improved existing ones
 - ➡ expect significant improvement!
- V_{cb}^{excl} : $B \rightarrow D$ lattice sensitive to $1/m^2$ corrections
 - ➡ Improved $V_{cb} + R(D)$ determination w/ BGL
- $B \rightarrow D^*$ awaits first finite-recoil LQCD calculation, BGL vs. CLN
 - ➡ V_{xb} puzzles severely reduced
- NP analyses require lattice determinations also of non-SM FFs!

THANK YOU!

Implications of the Higgs EFT for Flavour: $q \rightarrow q' l \nu$

$b \rightarrow c \tau \nu$ transitions (SM: $C_{V_L} = 1, C_{i \neq V_L} = 0$):

$$\mathcal{L}_{\text{eff}}^{b \rightarrow c \tau \nu} = -\frac{4G_F}{\sqrt{2}} V_{cb} \sum_j^5 C_j \mathcal{O}_j, \quad \text{with}$$

$$\mathcal{O}_{V_{L,R}} = (\bar{c} \gamma^\mu P_{L,R} b) \bar{\tau} \gamma_\mu \nu, \quad \mathcal{O}_{S_{L,R}} = (\bar{c} P_{L,R} b) \bar{\tau} \nu,$$

$$\mathcal{O}_T = (\bar{c} \sigma^{\mu\nu} P_L b) \bar{\tau} \sigma_{\mu\nu} \nu.$$

- All operators are independently present already in the linear EFT
- However: Relations between **different** transitions:
 C_{V_R} is **lepton-flavour universal** [see also Cirigliano+'09]
Relations between charged- and neutral-current processes, e.g.
 $\sum_{U=u,c,t} \lambda_{Us} C_{S_R}^{(U)} = -\frac{e^2}{8\pi^2} \lambda_{ts} C_S^{(d)}$ [see also Cirigliano+'12, Alonso+'15]
- These relations are again absent in the non-linear EFT

Matching for $b \rightarrow c\ell\nu$ transitions

$$C_{V_L} = -\mathcal{N}_{CC} \left[C_L + \frac{2}{v^2} c_{V5} + \frac{2V_{cb}}{v^2} c_{V7} \right],$$

$$C_{V_R} = -\mathcal{N}_{CC} \left[\hat{C}_R + \frac{2}{v^2} c_{V6} \right],$$

$$C_{S_L} = -\mathcal{N}_{CC} (c'_{S1} + \hat{c}'_{S5}),$$

$$C_{S_R} = 2\mathcal{N}_{CC} (c_{LR4} + \hat{c}_{LR8}),$$

$$C_T = -\mathcal{N}_{CC} (c'_{S2} + \hat{c}'_{S6}),$$

where $\mathcal{N}_{CC} = \frac{1}{2V_{cb}} \frac{v^2}{\Lambda^2}$, $C_L = 2c_{LL2} - \hat{c}_{LL6} + \hat{c}_{LL7}$ and $\hat{C}_R = -\frac{1}{2}\hat{c}_{Y4}$.

LO and NLO in linear and non-linear HEFT

Linear EFT

Building blocks $\psi_f, X_{\mu\nu}, D_\mu, H$

Finite powers of fields

H -interactions symmetry-restricted

LO:

- Terms of dimension 4
- ➔ SM (renormalizable)

NLO:

- 59 ops. (w/o flavour)
[Buchmüller+'86, Grzadkowski+'10]

Non-linear EFT

Building blocks $\psi_f, X_{\mu\nu}, D_\mu, U, h$
($U = \exp(2i\Phi/v)$)

Arbitrary powers of Φ, h : $U, f(h/v)$

U -interactions symmetry-restricted

LO:

- Tree-level h, U interactions
+ $SU(2)_{L+R}, g_{X-h}$ weak
- ➔ SM + $f_i(h/v)$, non-renorm.

NLO:

- ~ 100 ops. (w/o flavour)
[Buchalla+'14]

- Non-linear EFT **generalizes** linear EFT
- LO EFT predictive, justification for κ framework

$|V_{cb}|$: Recent developments

Recent Belle $B \rightarrow D, D^* \ell \nu$ analyses

Recent lattice results for $B \rightarrow D$

[FNAL/MILC, HPQCD, RBC/UKQCD (ongoing)]

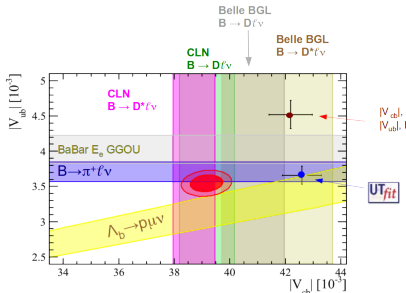
➡ $B \rightarrow D$ between incl. + $B \rightarrow D^*$

New lattice result for $B \rightarrow D^*$ [HPQCD]

➡ V_{cb}^{incl} cv, compatible with old result

$B \rightarrow D^* \ell \nu$ re-analyses with CLN,
 $|V_{cb}| = 39.3(1.0)10^{-2}$ [Bernlochner+'17]

+ BGL [Bigi+, Grinstein+'17] (Belle only),
 $|V_{cb}| = 40.4(1.7)10^{-2}$



[Plot modification by M. Rotondo]

Theoretical uncertainties previously underestimated, in two ways:

- $1/m_c^2$ contributions likely underestimated in CLN
- Uncertainty given in CLN ignored in experimental analyses
- ➡ Inclusive-exclusive tension softened

Experimental analyses used

Decay	Observable	Experiment	Comment	Year
$B \rightarrow D(\mathbf{e}, \mu)\nu$	BR	BaBar	global fit	2008
$B \rightarrow D\ell\nu$	$\frac{d\Gamma}{dw}$	BaBar	hadronic tag	2009
$B \rightarrow D(\mathbf{e}, \mu)\nu$	$\frac{d\Gamma}{dw}$	Belle	hadronic tag	2015
$B \rightarrow D^*(\mathbf{e}, \mu)\nu$	BR	BaBar	global fit	2008
$B \rightarrow D^*\ell\nu$	BR	BaBar	hadronic tag	2007
$B \rightarrow D^*\ell\nu$	BR	BaBar	untagged B^0	2007
$B \rightarrow D^*\ell\nu$	BR	BaBar	untagged B^\pm	2007
$B \rightarrow D^*(\mathbf{e}, \mu)\nu$	$\frac{d\Gamma_{L,\tau}}{dw}$	Belle	untagged	2010
$B \rightarrow D^*\ell\nu$	$\frac{d\Gamma}{d(w, \cos\theta_V, \cos\theta_I, \phi)}$	Belle	hadronic tag	2017

Different categories of data:

- Only total rates vs. differential distributions
- e, μ -averaged vs. individual measurements
- Correlation matrices given or not
- ➡ Sometimes presentation prevents use in non-universal scenarios 😞
- ➡ Recent Belle analyses (mostly) exemplary 😊

NP in semileptonic decays - Setup and tree-level scenarios

EFT for $b \rightarrow c \ell \nu_{\ell'}$ transitions (no light ν_R , SM: $C_j^{\ell\ell'} = 0$):

$$\mathcal{L}_{\text{eff}}^{b \rightarrow c \ell \nu} = -\frac{4G_F}{\sqrt{2}} V_{cb} \sum_j^5 \sum_{\ell, \ell' = e, \mu, \tau} \left[\delta_{\ell\ell'} \delta_{jV_L} + C_j^{\ell\ell'} \right] \mathcal{O}_j^{\ell\ell'}, \quad \text{with}$$

$$\mathcal{O}_{V_{L,R}}^{\ell\ell'} = (\bar{c} \gamma^\mu P_{L,R} b) \bar{\ell} \gamma_\mu \nu_{\ell'}, \quad \mathcal{O}_{S_{L,R}}^{\ell\ell'} = (\bar{c} P_{L,R} b) \bar{\ell} \nu_{\ell'}, \quad \mathcal{O}_T^{\ell\ell'} = (\bar{c} \sigma^{\mu\nu} P_L b) \bar{\ell} \sigma_{\mu\nu} \nu_{\ell'}.$$

NP models typically generate **subsets** (never C_T alone)

➡ Full classification possible for tree-level mediators [Freytsis+'15]:

Model	C_{V_L}	C_{V_R}	C_{S_R}	C_{S_L}	C_T	$C_{S_L} = 4C_T$	$C_{S_L} = -4C_T$
Vector-like singlet	×						
Vector-like doublet		×					
W'	×						
H^\pm			×	×			
S_1	×						×
R_2						×	
S_3	×						
U_1	×		×				
V_2			×				
U_3	×						

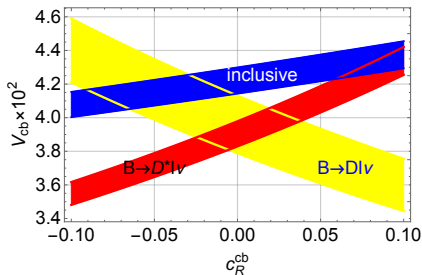
Right-handed vector currents [MJ/Straub'18]

Usual suspect for tension inclusive vs. exclusive [e.g. Voloshin'97]

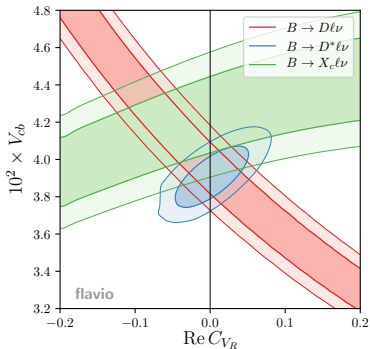
SMEFT: $C_{V_R}^{\ell\ell'}$ is **lepton-flavour-universal** [Cirigliano+'10, Catà/MJ'15]

➡ All available data can be used in SMEFT context

➡ Violation could signal non-linear realization of EWSB [Catà/MJ'15]



[Plot: updated from Crivellin/Pokorski'14]



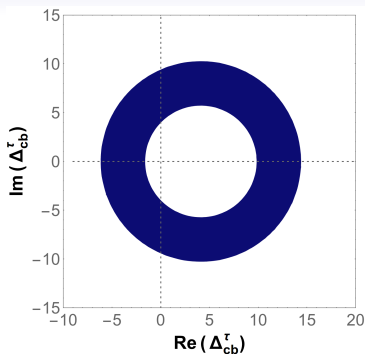
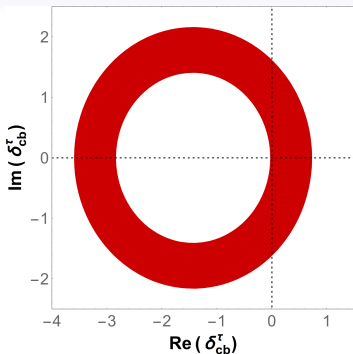
Impact of differential distributions:

V_{cb} and C_{V_R} can be determined **individually** in $B \rightarrow D^*$

➡ Tension smaller, but is **not** improved by C_{V_R}

➡ C_{V_R} in SMEFT cannot explain $b \rightarrow c \tau \nu$ data

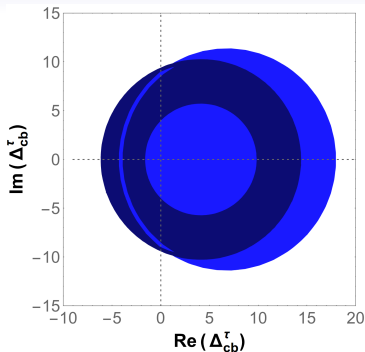
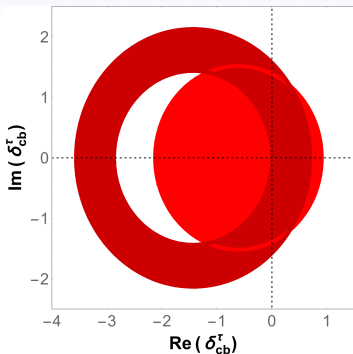
$b \rightarrow c\tau\nu$ data and scalar NP [Celis/MJ/Li/Pich'17]



$R(D), R(D^*)$: trivially explainable, but strange

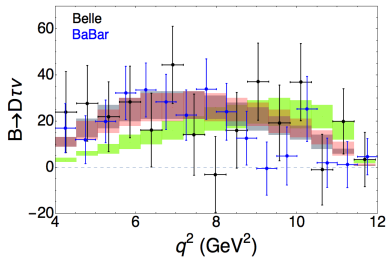
- $R(D) : \delta_{cb}^I \equiv \frac{(C_{S_L} + C_{S_R})(m_B - m_D)^2}{m_l(\bar{m}_b - \bar{m}_c)}$, $R(D^*) : \Delta_{cb}^I \equiv \frac{(C_{S_L} - C_{S_R})m_B^2}{m_l(\bar{m}_b + \bar{m}_c)}$
- $R(D)$ compatible with SM at $\sim 2\sigma$
- Preferred scalar couplings from $R(D^*)$ huge ($|C_{S_L} - C_{S_R}| \sim 1 - 5$)
- Can't go beyond circles with just $R(D, D^*)$!

$b \rightarrow cTV$ data and scalar NP [Celis/MJ/Li/Pich'17]

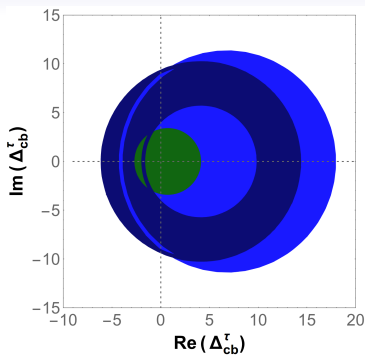
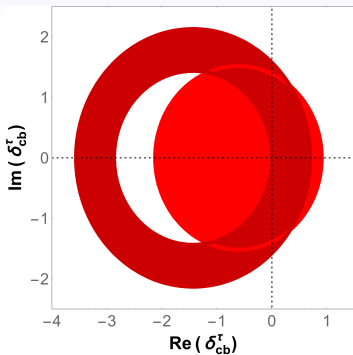


Differential rates:

- compatible with SM and NP
- already now constraining, especially in $B \rightarrow DTV$
- “theory-dependence” of data needs addressing [Bernlochner+'17]



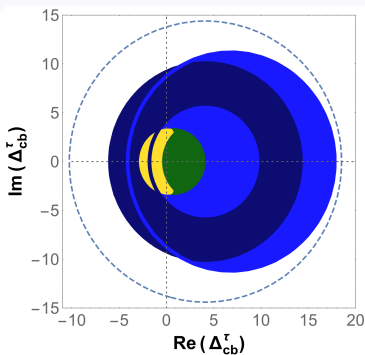
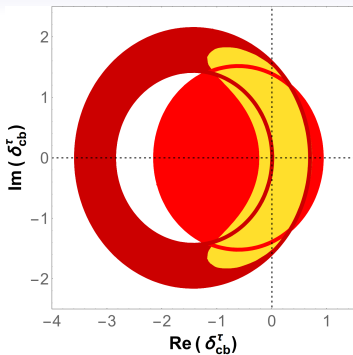
$b \rightarrow c\tau\nu$ data and scalar NP [Celis/MJ/Li/Pich'17]



Total width of B_c :

- $B_c \rightarrow \tau\nu$ is an obvious $b \rightarrow c\tau\nu$ transition
 - ➡ not measurable in foreseeable future
 - ➡ can oversaturate total width of B_c ! [X.Li+'16]
- Excludes second real solution in Δ_{cb}^τ plane (even scalar NP for $R(D^*)$? [Alonso+'16, Akeroyd+'17])

$b \rightarrow c\tau\nu$ data and scalar NP [Celis/MJ/Li/Pich'17]



τ polarization:

- So far not constraining (shown: $\Delta\chi^2 = 1$)
- Differentiate NP models: with scalar NP [Celis/MJ/Li/Pich'13]

$$X_2^{D^{(*)}}(q^2) \equiv R_{D^{(*)}}(q^2) \left[A_\lambda^{D^{(*)}}(q^2) + 1 \right] = X_{2,SM}^{D^{(*)}}(q^2)$$

Consistent explanation in 2HDMs possible, flavour structure?

Differentiating models with $b \rightarrow c\tau\nu$ observables

Large $R(D^*)$ possible with NP in V_L ($\hat{R}(X) = R(X)/R(X)_{SM}$):

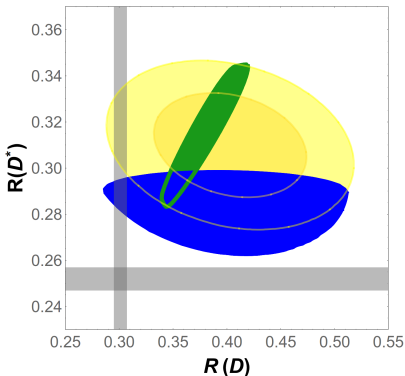
- trivial prediction: $\hat{R}(D) = \hat{R}(D^*) = \hat{R}(\Lambda_c) = \dots \stackrel{exp}{\sim} 1.25$
- can be related to anomaly in $B \rightarrow K^{(*)}\ell^+\ell^-$ modes
- $\hat{R}(X_c) = 0.99 \pm 0.10$ measured by LEP, oversaturation
- issues with $\tau \rightarrow \mu\nu\nu$ [Feruglio+'16] and $b\bar{b} \rightarrow X \rightarrow \tau^+\tau^-$ [Faroughy+'16]

Differentiating models with $b \rightarrow c\tau\nu$ observables

Large $R(D^*)$ possible with NP in V_L ($\hat{R}(X) = R(X)/R(X)_{SM}$):

- trivial prediction: $\hat{R}(D) = \hat{R}(D^*) = \hat{R}(\Lambda_c) = \dots \stackrel{exp}{\sim} 1.25$
- can be related to anomaly in $B \rightarrow K^{(*)}\ell^+\ell^-$ modes
- $\hat{R}(X_c) = 0.99 \pm 0.10$ measured by LEP, oversaturation
- issues with $\tau \rightarrow \mu\nu\nu$ [Feruglio+'16] and $b\bar{b} \rightarrow X \rightarrow \tau^+\tau^-$ [Faroughy+'16]

Fit results for the two scenarios for $B \rightarrow D^{(*)}\tau\nu$:

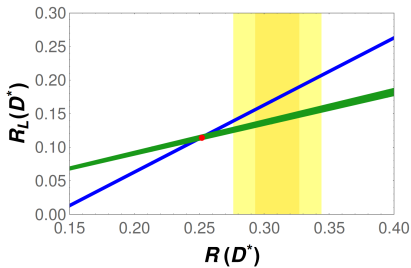
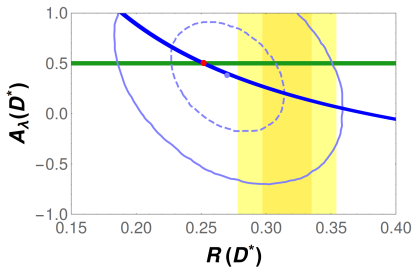


Differentiating models with $b \rightarrow c\tau\nu$ observables

Large $R(D^*)$ possible with NP in V_L ($\hat{R}(X) = R(X)/R(X)_{SM}$):

- trivial prediction: $\hat{R}(D) = \hat{R}(D^*) = \hat{R}(\Lambda_c) = \dots \stackrel{exp}{\sim} 1.25$
- can be related to anomaly in $B \rightarrow K^{(*)}\ell^+\ell^-$ modes
- $\hat{R}(X_c) = 0.99 \pm 0.10$ measured by LEP, oversaturation
- issues with $\tau \rightarrow \mu\nu\nu$ [Feruglio+'16] and $b\bar{b} \rightarrow X \rightarrow \tau^+\tau^-$ [Faroughy+'16]

Fit predictions for polarization-dependent $B \rightarrow D^*\tau\nu$ observables:

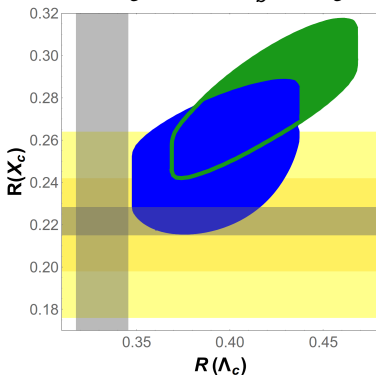


Differentiating models with $b \rightarrow c\tau\nu$ observables

Large $R(D^*)$ possible with NP in V_L ($\hat{R}(X) = R(X)/R(X)_{SM}$):

- trivial prediction: $\hat{R}(D) = \hat{R}(D^*) = \hat{R}(\Lambda_c) = \dots \stackrel{exp}{\sim} 1.25$
- can be related to anomaly in $B \rightarrow K^{(*)}\ell^+\ell^-$ modes
- $\hat{R}(X_c) = 0.99 \pm 0.10$ measured by LEP, oversaturation
- issues with $\tau \rightarrow \mu\nu\nu$ [Feruglio+'16] and $b\bar{b} \rightarrow X \rightarrow \tau^+\tau^-$ [Farouhy+'16]

Fit predictions for $B \rightarrow X_{c\tau\nu}$ and $\Lambda_b \rightarrow \Lambda_{c\tau\nu}$:



SM and left-handed vector operators

As a crosscheck, produce SM values (using data from HEPdata):

$$V_{cb}^{B \rightarrow D} = (39.6 \pm 0.9)10^{-3} \quad V_{cb}^{B \rightarrow D^*} = (39.0 \pm 0.7)10^{-3}$$

➡ low compared to BGL analyses, compatible with recent results

NP in $\mathcal{O}_{V_L}^{\ell\ell'}$: can be absorbed via $\tilde{V}_{cb}^\ell = V_{cb} \left[|1 + C_{V_L}^\ell|^2 + \sum_{\ell' \neq \ell} |C_{V_L}^{\ell\ell'}|^2 \right]^{1/2}$

Only subset of data usable

$B \rightarrow D, D^*$ in agreement

No sign of LFNU

➡ constrained to be $\lesssim \% \times V_{cb}$

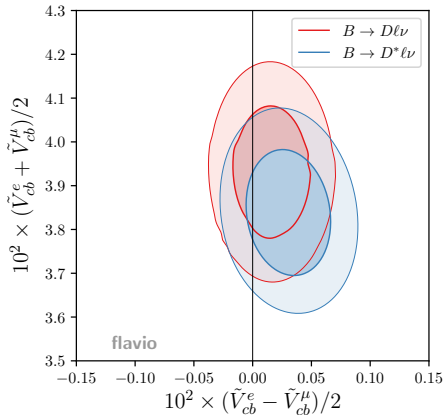
In the following:

- e and μ analyzed separately
- ➡ Usable in different contexts
- Full FF constraints used

🎨 Plots created with **flavio**

+ independently double-checked

➡ Open source, adaptable



Scalar operators

For $m_\ell \rightarrow 0$, no interference with SM

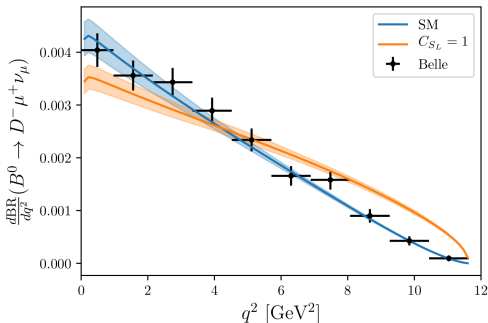
➡ For fixed V_{cb} , scalar NP **increases** rates

Close to $q^2 \rightarrow q_{\max}^2$ in the SM: $\frac{d\Gamma(B \rightarrow D\ell\nu)}{dq^2} \propto f_+^2 (q^2 - q_{\max}^2)^{3/2}$

With scalar contributions: $\frac{d\Gamma(B \rightarrow D\ell\nu)}{dq^2} \propto f_0^2 |C_{S_R} + C_{S_L}|^2 (q^2 - q_{\max}^2)^{1/2}$

➡ Endpoint very sensitive to scalar contributions! [see also Nierste+'08]

Scalar contributions ruled out by the distributions ($\Gamma_1 = \Gamma_2$):



Scalar operators

For $m_\ell \rightarrow 0$, no interference with SM

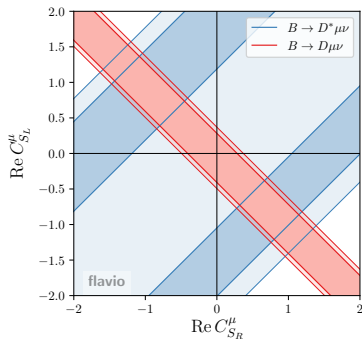
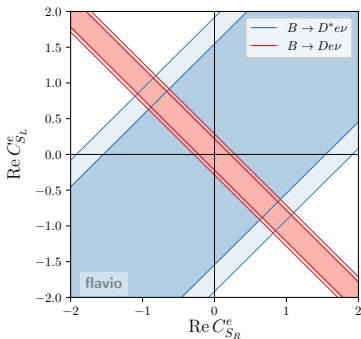
➡ For fixed V_{cb} , scalar NP **increases** rates

Close to $q^2 \rightarrow q_{\text{max}}^2$ in the SM: $\frac{d\Gamma(B \rightarrow D\ell\nu)}{dq^2} \propto f_+^2 (q^2 - q_{\text{max}}^2)^{3/2}$

With scalar contributions: $\frac{d\Gamma(B \rightarrow D\ell\nu)}{dq^2} \propto f_0^2 |C_{S_R} + C_{S_L}|^2 (q^2 - q_{\text{max}}^2)^{1/2}$

➡ Endpoint very sensitive to scalar contributions! [see also Nierste+'08]

Fit with scalar couplings (generic $C_{S_{L,R}}$):



Slightly favours large contributions in muon couplings with $C_{S_R}^\mu \approx -C_{S_L}^\mu$

Scalar operators

For $m_\ell \rightarrow 0$, no interference with SM

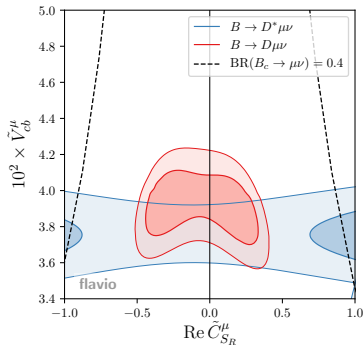
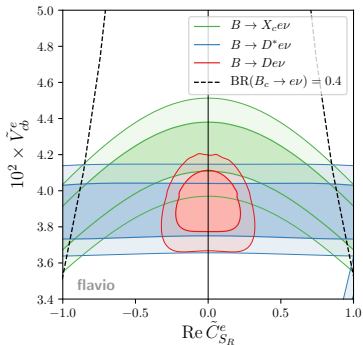
➡ For fixed V_{cb} , scalar NP **increases** rates

Close to $q^2 \rightarrow q_{\text{max}}^2$ in the SM: $\frac{d\Gamma(B \rightarrow D\ell\nu)}{dq^2} \propto f_+^2 (q^2 - q_{\text{max}}^2)^{3/2}$

With scalar contributions: $\frac{d\Gamma(B \rightarrow D\ell\nu)}{dq^2} \propto f_0^2 |C_{S_R} + C_{S_L}|^2 (q^2 - q_{\text{max}}^2)^{1/2}$

➡ Endpoint very sensitive to scalar contributions! [see also Nierste+'08]

Also for LQ U_1 (or V_2): $B \rightarrow D$ stronger than $B \rightarrow D^*$, X_C :



Possible large contribution in $C_{S_R}^\mu$ excluded by $B \rightarrow D$

Tensor operators

For $m_\ell \rightarrow 0$, no interference with SM

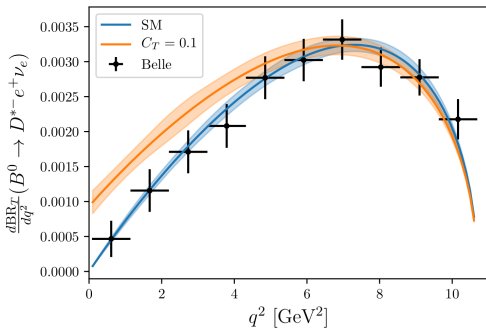
➡ For fixed V_{cb} , tensor contributions **increase** rates

Close to $q^2 \rightarrow q_{\min}^2$:

$$\frac{d\Gamma_T(B \rightarrow D^* \ell \nu)}{dq^2} \propto q^2 C_{V_L}^2 (A_1(0)^2 + V(0)^2) + 16m_B^2 C_T^2 T_1(0)^2 + O\left(\frac{m_{D^*}^2}{m_B^2}\right)$$

➡ Endpoint ($q^2 \sim 0$) very sensitive to tensor contributions!

Tensor contributions ruled out by the distributions ($\Gamma_1 = \Gamma_2$):



Tensor operators

For $m_\ell \rightarrow 0$, no interference with SM

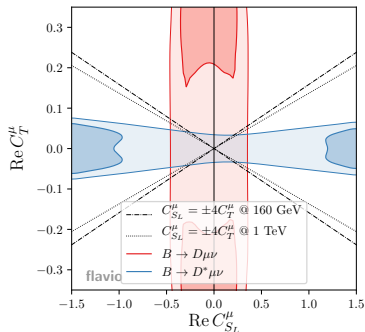
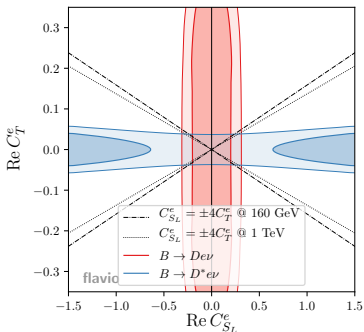
➡ For fixed V_{cb} , tensor contributions **increase** rates

Close to $q^2 \rightarrow q_{\min}^2$:

$$\frac{d\Gamma_T(B \rightarrow D^* \ell \nu)}{dq^2} \propto q^2 C_{V_L}^2 (A_1(0)^2 + V(0)^2) + 16m_B^2 C_T^2 T_1(0)^2 + O\left(\frac{m_{D^*}^2}{m_B^2}\right)$$

➡ Endpoint ($q^2 \sim 0$) very sensitive to tensor contributions!

Fit for generic C_{S_L} and C_T (including LQs S_1 and R_1):



$B \rightarrow D^*$ favours large contributions in $C_{S_L}^{e,\mu}$, ruled out by $B \rightarrow D$