

LHCb prospects for measurements of UT angles

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Synergy between LHC and SuperKEKB in the Quest for New Physics
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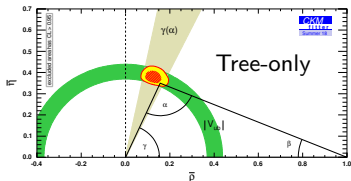
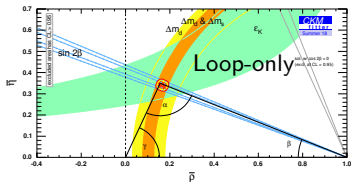
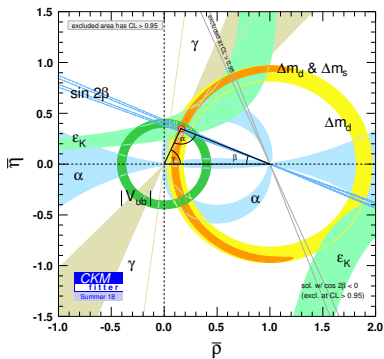
On behalf of the LHCb collaboration

Unitarity Triangle measurements

Cabibbo-Kobayashi-Maskawa matrix

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \simeq \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

Sensitivity to NP comes from the global consistency of various measurements

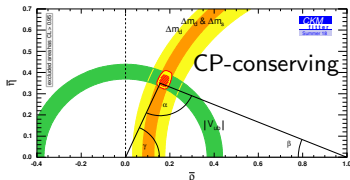
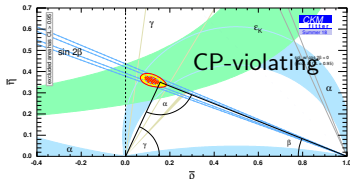
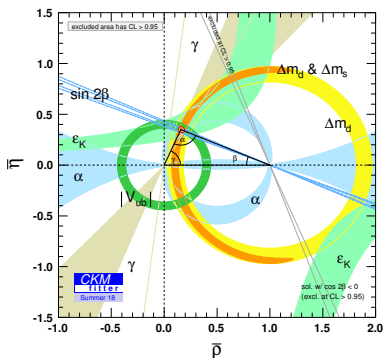


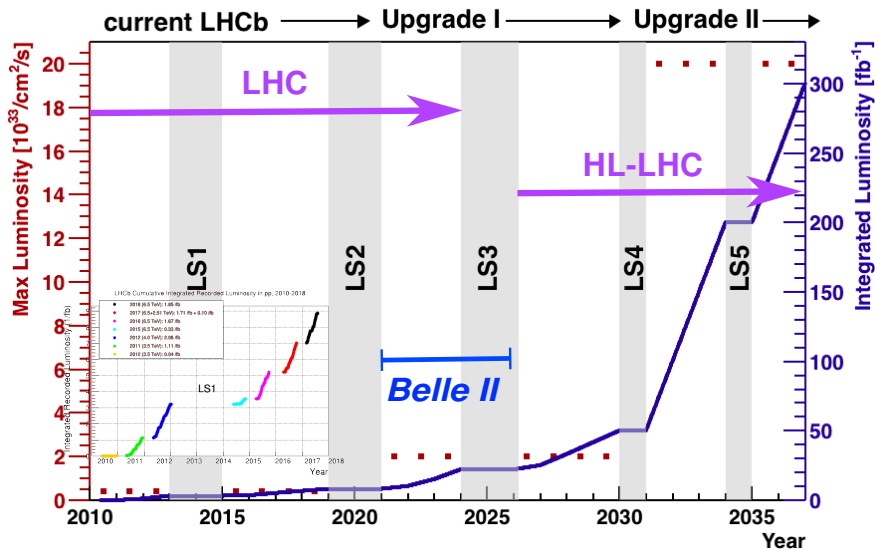
Unitarity Triangle measurements

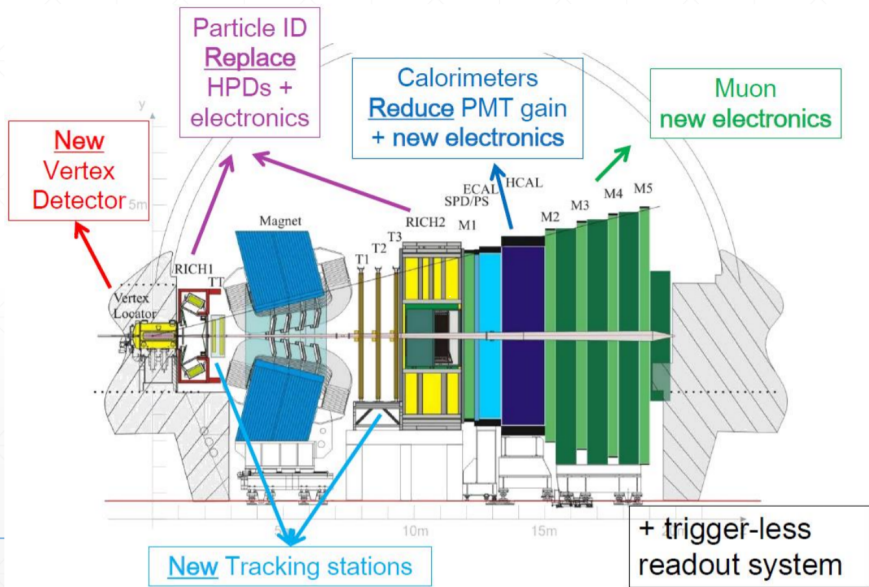
Cabibbo-Kobayashi-Maskawa matrix

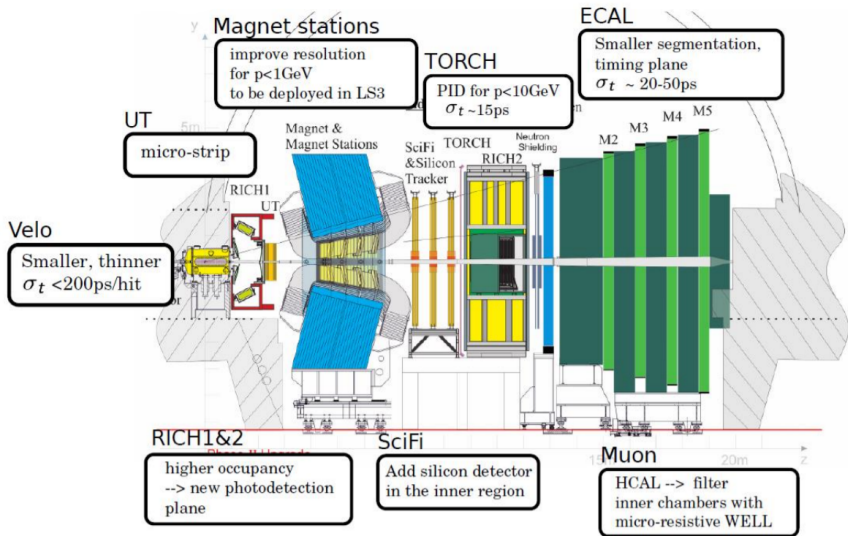
$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \simeq \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

Sensitivity to NP comes from the global consistency of various measurements



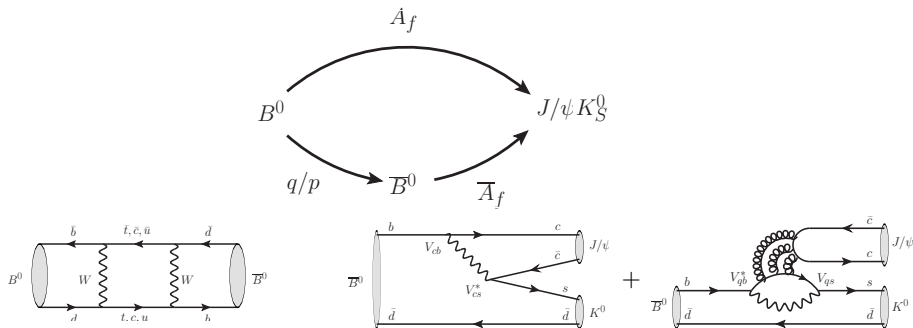






β

Interference between the decays with and without mixing



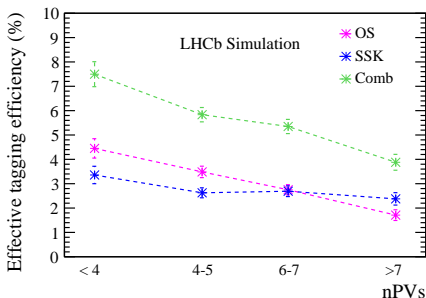
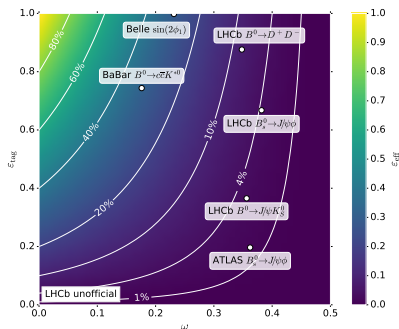
Time-dependent CP asymmetry

$$A_{CP}(t) = \frac{\Gamma(\bar{B}^0(t) \rightarrow f) - \Gamma(B^0(t) \rightarrow f)}{\Gamma(\bar{B}^0(t) \rightarrow f) + \Gamma(B^0(t) \rightarrow f)} = \frac{S \sin(\Delta m t) + C \cos(\Delta m t)}{\cosh(\frac{1}{2} \Delta \Gamma t) + A_{\Delta \Gamma} \sinh(\frac{1}{2} \Delta \Gamma t)}$$

$\overset{= \sin 2\beta}{S}$ $\overset{\approx 0}{C}$
 $\underset{1}{\cosh(\frac{1}{2} \Delta \Gamma t)}$ $\underset{0}{A_{\Delta \Gamma} \sinh(\frac{1}{2} \Delta \Gamma t)}$

Performance critically depends on flavour tagging performance.

$$\sigma \propto \frac{1}{\sqrt{\epsilon_{\text{eff}} N}}; \quad \epsilon_{\text{eff}} = \epsilon_{\text{tag}}(1 - 2\omega); \quad \epsilon_{\text{tag}} = \frac{N_{\text{tagged}}}{N_{\text{tagged}} + N_{\text{untagged}}}; \quad \omega = \frac{N_{\text{right}}}{N_{\text{right}} + N_{\text{wrong}}}$$

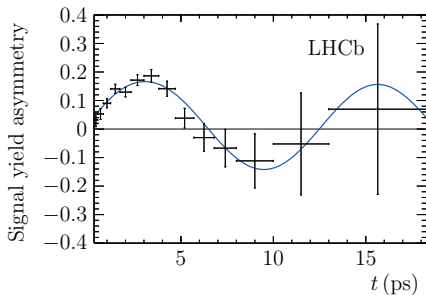


With Upgrade II: expect $\langle N_{PV} \rangle \sim 50$.

Timing in VELO/tracker is essential to recover Run I performance.

~ 50 ps will be needed to reach reasonable track-PV mis-association rate.

Measured in $B^0 \rightarrow J/\psi K_S^0$ with
 $J/\psi \rightarrow \mu^+\mu^-$ and e^+e^- , and
 $B^0 \rightarrow \psi(2S)K_S^0$, $\psi(2S) \rightarrow \mu^+\mu^-$.



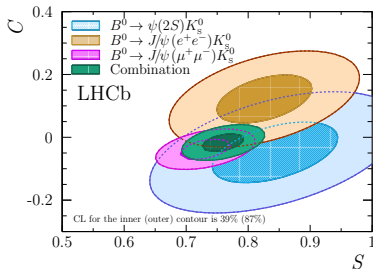
Combined result:

$$S_{c\bar{c}K_S^0} = 0.760 \pm 0.034$$

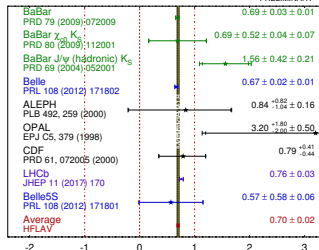
$$C_{c\bar{c}K_S^0} = -0.017 \pm 0.029$$

Precision similar to B factories. Dominated by statistics.

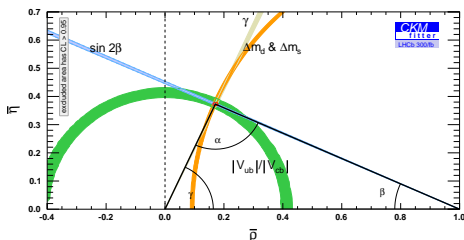
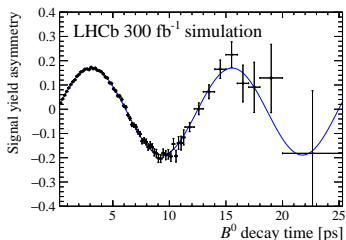
[PRL 115, 031601 (2015)] [JHEP 11 (2017) 170]



$\sin(2\beta) \equiv \sin(2\phi_1)$ **HFLAV**
 March 2018
 PRELIMINARY



[LHCb-PUB-2018-009: "Physics case for an LHCb Upgrade II"]

Assuming FT performance can be kept at Run I level: $\sigma(\sin 2\beta) = 0.003$ for 300 fb^{-1} .

Penguin contribution becomes a major systematic uncertainty.

- $SU(3)$ symmetry to constrain penguin contribution, e.g. $B_s \rightarrow J/\psi K_S^0$.
- $B^0 \rightarrow D\pi^+\pi^-$ mode: tree-level process; no penguin pollution. In addition, measures $\sin 2\beta$ and $\cos 2\beta$ independently.

- $D \rightarrow KK, \pi\pi$: TD Dalitz plot analysis (model uncertainty)

$$\sigma(\sin 2\beta) \sim 0.007, \sigma(\cos 2\beta) \sim 0.017 \quad \text{for } 300 \text{ fb}^{-1}$$

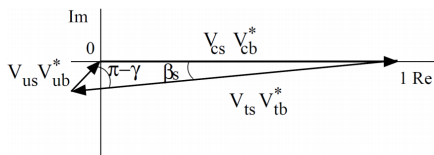
[JPG 36 (2009) 025006]

- $D \rightarrow K_S^0 hh$: MI approach, TD analysis of two correlated Dalitz plots

$$\text{Estimated } \sigma(\beta) \sim (2-3)^\circ \quad \text{with } 50 \text{ fb}^{-1}$$

[JHEP 03(2018) 195]

ϕ S



β_s : similar to β , but in B_s^0 system.

Measurable phase between mixing and decay:

$$\phi_s = \phi_M - 2\phi_D = -2\beta_s + \Delta\phi_s^{\text{peng}} + \delta_s^{\text{NP}}$$

- Sensitivity to New Physics
- Input measurement for other CKM studies (e.g. $B_s^0 \rightarrow D_s K$, see later)

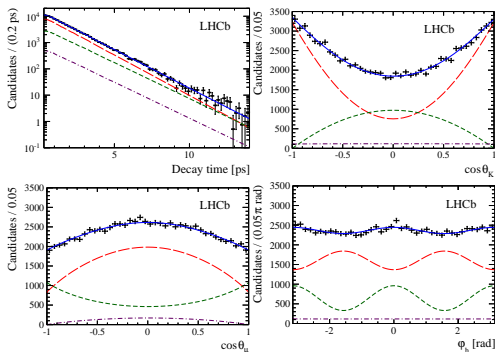
Similar TD formalism, but more complicated because of finite $\Delta\Gamma_s$:

$$\frac{d\Gamma_{B_s^0(\bar{B}_s^0) \rightarrow f}(t)}{dt} \propto e^{-\Gamma_s t} \left[\cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) \right. \\ \left. \pm C_f \cos(\Delta m_s t) \mp S_f \sin(\Delta m_s t) \right]$$

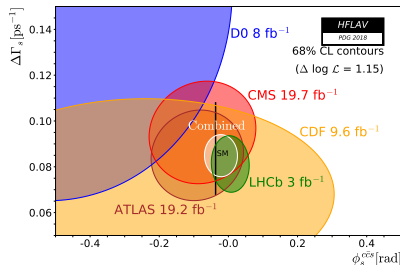
$A_f^{\Delta\Gamma} = -\frac{2|\lambda_f| \cos \phi_s}{1+|\lambda_f|^2}$
 $C_f = \frac{1-|\lambda_f|^2}{1+|\lambda_f|^2}$
 $S_f = \frac{2|\lambda_f| \sin \phi_s}{1+|\lambda_f|^2}$

- $K^+ K^-$ can be in P wave (ϕ) or S wave
- 3 P waves (\mathcal{CP} -odd or \mathcal{CP} -even), angular analysis to distinguish them
- Ambiguity $\varphi_s \leftrightarrow \pi - \varphi_s$ is resolved by measuring the P wave strong phase as a function of m_{KK} .
- Combined with $B_s^0 \rightarrow J/\psi \pi^+ \pi^-$, $B_s^0 \rightarrow \psi(2S) K^+ K^-$

Decay time and helicity distributions:



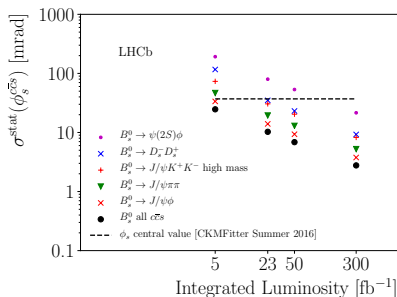
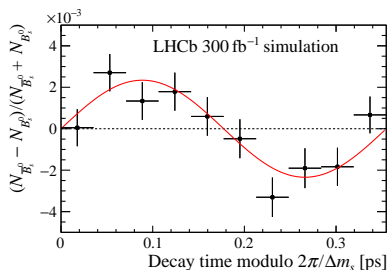
$$\begin{aligned} \phi_s &= -58 \pm 49 \pm 6 \text{ mrad} \\ \Gamma_s &= 0.6603 \pm 0.0027 \pm 0.0015 \\ \Delta\Gamma_s &= 0.0805 \pm 0.0091 \pm 0.0032 \\ \Delta m_s &= 17.711_{-0.057}^{+0.055} \pm 0.0032 \\ |\lambda| &= 0.964 \pm 0.019 \pm 0.007 \end{aligned}$$



[LHCb-PUB-2018-009: "Physics case for an LHCb Upgrade II"]

Important to combine several $b \rightarrow c\bar{c}s$ modes to control systematic uncertainties and penguin contribution.

Several modes are expected to measure non-zero $-\beta_s^{SM} = -36.4 \pm 1.2$ mrad.

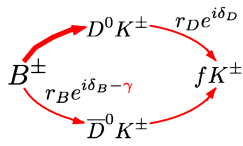
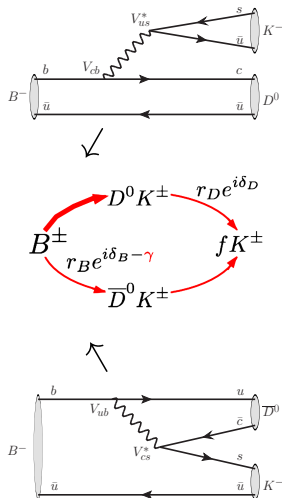


Expected sensitivity with 300 fb⁻¹: **4 mrad** for $B_s^0 \rightarrow J/\psi \phi$ and **3 mrad** for all modes combined

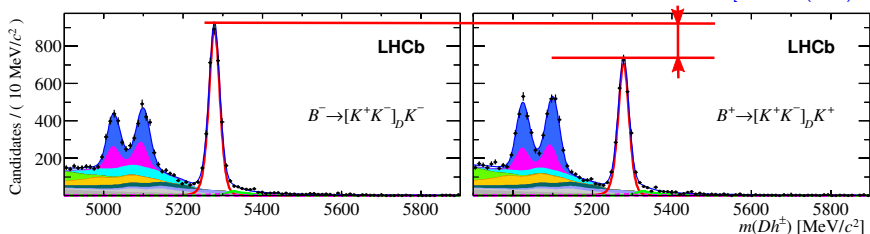
γ

- Measured entirely from tree decays.
- All hadronic parameters can be constrained from experiment \Rightarrow theoretically very clean (uncertainty $< 10^{-7}$ [Brod, Zupan, JHEP 1401 (2014) 051])
- Combination of many different modes:
 - Time-integrated asymmetries in $B \rightarrow DK$, $B \rightarrow DK^*$, $B \rightarrow DK\pi$ with $D \rightarrow hh, hhhh$ (“ADS”, “GLW”)
 - Dalitz plot analyses of $D^0 \rightarrow K_S^0 h^+ h^-$ from $B \rightarrow DK$, $B \rightarrow DK^*$ (“GGSZ”)
 - Time-dependent analyses of $B_s^0 \rightarrow D_s K$, $B^0 \rightarrow D\pi$
- Experimentally, just entering precision measurement regime ($< 10\%$)

[LHCb-CONF-2016-001]



[PLB 777 (2018) 16]



CP -violating rate for $B^\pm \rightarrow D(\rightarrow f)K^\pm$ decays:

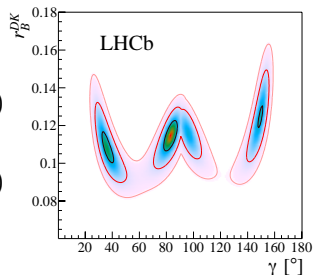
$$\Gamma(B^\pm \rightarrow D(\rightarrow f)K^\pm) = r_D^2 + r_B^2 + 2\kappa r_D r_B \cos(\delta_B - \delta_D \pm \gamma)$$

r_B : ratio of $b \rightarrow u$ and $b \rightarrow c$ amplitudes

r_D : ratio of $D^0 \rightarrow f$ and $\bar{D}^0 \rightarrow f$ amplitudes ($\equiv 1$ for D_{CP})

δ_B and δ_D : corresponding strong phase differences

κ : coherence factor ($\equiv 1$ for 2-body decays)



Combination of several D and B decay modes to constrain γ

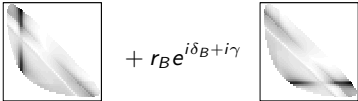
Information on γ from Dalitz plot analysis of $D \rightarrow K_S^0 \pi^+ \pi^-$ from $B^\pm \rightarrow DK^\pm$.

Dalitz plot density: $d\sigma(m_+^2, m_-^2) \sim |A|^2 dm_+^2 dm_-^2$, where $m_\pm^2 = m_{K_S^0 \pi^\pm}^2$

Flavour D amplitude: $A_D(m_+^2, m_-^2)$

Amplitude of $D \rightarrow K_S^0 \pi^+ \pi^-$ from $B^+ \rightarrow DK^+$:

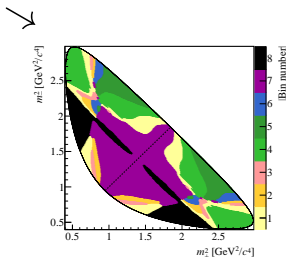
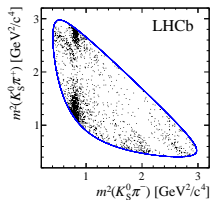
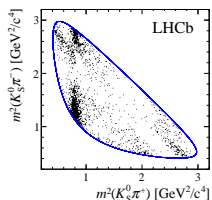
$$A_B(m_+^2, m_-^2) = A_D(m_+^2, m_-^2) + r_B e^{i\delta_B + i\gamma} A_D(m_-^2, m_+^2)$$

$$= \text{[Dalitz Plot 1]} + r_B e^{i\delta_B + i\gamma} \text{[Dalitz Plot 2]}$$


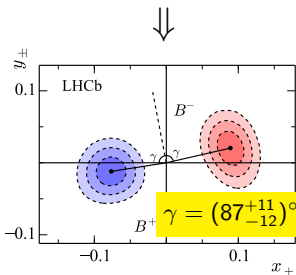
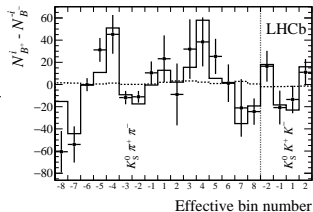
Need to know $A_D(m_+^2, m_-^2)$, both amplitude and phase (or, more precisely, phase difference between (m_+^2, m_-^2) and (m_-^2, m_+^2)).

Model-dependent: obtain A_D from $D \rightarrow K_S^0 \pi^+ \pi^-$ fit to the isobar model \Rightarrow model uncertainty

Model-independent: obtain phase difference info from $e^+ e^- \rightarrow D^0 \bar{D}^0$ decays (CLEO, BES-III)



Run II (2015, 2016): [\[JHEP 08 \(2018\) 176\]](#)

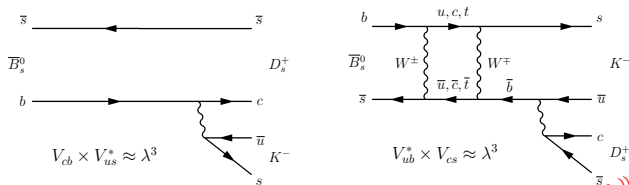


Combination with Run I: $\gamma = (80_{-9}^{+10})^\circ$

Main syst uncertainties (total $\sim 30\%$ of stat):

- Strong phase measurement (CLEO)
- Dalitz plot efficiency correction
- Mass fit shapes

Interference between $b \rightarrow u$ and $b \rightarrow c$ amplitude from B_s^0 mixing.
 Comparable magnitudes $r = \left| \frac{p}{q} \frac{A_f}{A_{\bar{f}}} \right| \simeq 0.4$.



Time-dependent decay rates for $B_s^0(\bar{B}_s^0) \rightarrow f$ (similarly for \bar{f})

$$\frac{d\Gamma_{B_s^0(\bar{B}_s^0) \rightarrow f}(t)}{dt} \propto e^{-\Gamma_s t} \left[\cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + A_f \Delta\Gamma \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) \right. \\ \left. \pm C_f \cos(\Delta m_s t) \mp S_f \sin(\Delta m_s t) \right]$$

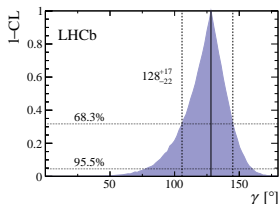
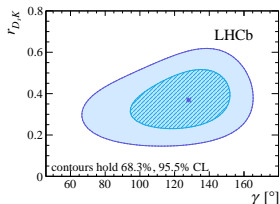
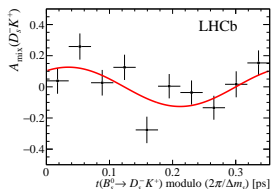
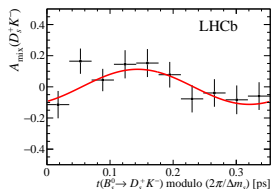
Measure $\gamma - 2\beta_s, \delta, r$

$$= \frac{1-r^2}{1+r^2}$$

$$= \frac{2r \sin(\delta - (\gamma - 2\beta_s))}{1+r^2}$$

Similar technique with $B^0 \rightarrow D\pi$ (but negligible $\Delta\Gamma_d$, small $r \simeq 0.02 \Rightarrow$ only two observables S_f, \bar{S}_f).

Measure $2\beta + \gamma$ with the external input for r (from $SU(3) B^0 \rightarrow D_s \pi$)



Relies on input $-2\beta_s = -0.030 \pm 0.033 \Rightarrow \gamma = (128^{+17}_{-22})^\circ$ (stat-limited).

Systematic uncertainties: background, Δm_s , time acceptance, resolution, flavour tagging. All data-driven.

Safe to assume $1/\sqrt{N}$ scaling $\Rightarrow \sim 1^\circ$ with 300 fb^{-1} .

Many modes are combined to constrain γ :[\[LHCb-CONF-2018-002\]](#)

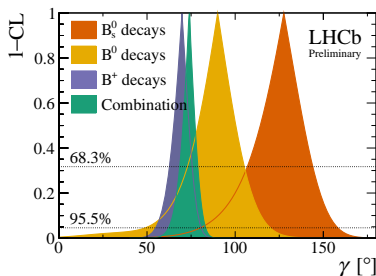
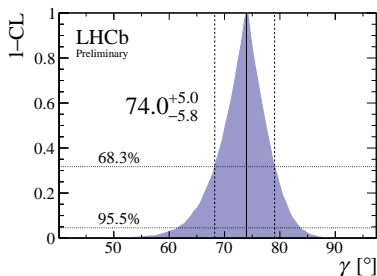
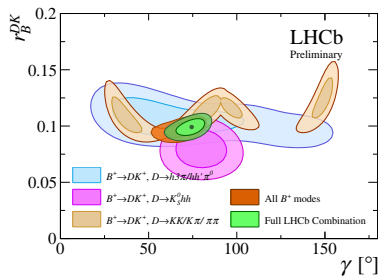
B decay	D decay	Method	Ref.	Dataset [†]	Status since last combination [3]
$B^+ \rightarrow DK^+$	$D \rightarrow h^+h^-$	GLW	[14]	Run 1 & 2	Minor update
$B^+ \rightarrow DK^+$	$D \rightarrow h^+h^-$	ADS	[15]	Run 1	As before
$B^+ \rightarrow DK^+$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	GLW/ADS	[15]	Run 1	As before
$B^+ \rightarrow DK^+$	$D \rightarrow h^+h^-\pi^0$	GLW/ADS	[16]	Run 1	As before
$B^+ \rightarrow DK^+$	$D \rightarrow K_s^0 h^+h^-$	GGSZ	[17]	Run 1	As before
$B^+ \rightarrow DK^+$	$D \rightarrow K_s^0 h^+h^-$	GGSZ	[18]	Run 2	New
$B^+ \rightarrow DK^+$	$D \rightarrow K_s^0 K^+\pi^-$	GLS	[19]	Run 1	As before
$B^+ \rightarrow D^*K^+$	$D \rightarrow h^+h^-$	GLW	[14]	Run 1 & 2	Minor update
$B^+ \rightarrow DK^{*+}$	$D \rightarrow h^+h^-$	GLW/ADS	[20]	Run 1 & 2	Updated results
$B^+ \rightarrow DK^{*+}$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	GLW/ADS	[20]	Run 1 & 2	New
$B^+ \rightarrow DK^+\pi^+\pi^-$	$D \rightarrow h^+h^-$	GLW/ADS	[21]	Run 1	As before
$B^0 \rightarrow DK^{*0}$	$D \rightarrow K^+\pi^-$	ADS	[22]	Run 1	As before
$B^0 \rightarrow DK^+\pi^-$	$D \rightarrow h^+h^-$	GLW-Dalitz	[23]	Run 1	As before
$B^0 \rightarrow DK^{*0}$	$D \rightarrow K_s^0 \pi^+\pi^-$	GGSZ	[24]	Run 1	As before
$B_s^0 \rightarrow D_s^\mp K^\pm$	$D_s^+ \rightarrow h^+h^-\pi^+$	TD	[25]	Run 1	Updated results
$B_s^0 \rightarrow D_s^\mp \pi^\pm$	$D^+ \rightarrow K^+\pi^-\pi^+$	TD	[26]	Run 1	New

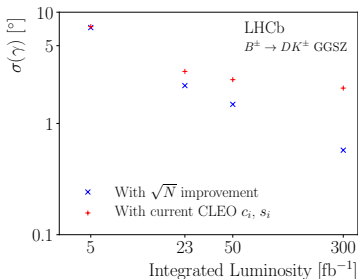
[†] Run 1 corresponds to an integrated luminosity of 3 fb^{-1} taken at centre-of-mass energies of 7 and 8 TeV
. Run 2 corresponds to an integrated luminosity of 2 fb^{-1} taken at a centre-of-mass energy of 13 TeV .

State of the art: γ combination

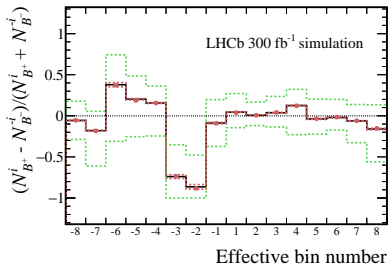
- $D \rightarrow hh$ (ADS/GLW) provide strong constraints in r_B, δ_B, γ space, but ambiguities and non-gaussian uncertainties.
- $D \rightarrow K_S^0 hh$ modes break ambiguities
- Different correlation patterns result in combined γ uncertainty better than just plain average.
- Different analysis approaches (rates, Dalitz, time-dep) allow better control of systematic uncertainties.

[LHCb-CONF-2018-002]





[LHCb-PUB-2018-009: "Physics case for an LHCb Upgrade II"]



- Critical uncertainty: CLEO measurement of strong phase difference in bins. Currently: $\simeq 2^\circ$.
- Further reduction is possible:
 - Expect BES-III to contribute with around $\times 4$ larger dataset.
 - Technique to obtain $D^0 - \bar{D}^0$ phase difference from charm mixing fits at LHCb [JHEP 10 (2012) 185]
 - Use other $B \rightarrow DX$ decays to overconstrain phase difference, such as $B \rightarrow DK\pi$, $D \rightarrow K_S^0 \pi \pi$ [PRD 97, 056002 (2018)]
 - $B \rightarrow DK$ decays themselves constrain phase difference for sufficiently large dataset [preliminary toy MC studies]
- Other uncertainties depend on control or MC samples.

Assume that $1/\sqrt{N}$ scaling is valid: $\sigma(\gamma) < 1^\circ$ with 300 fb⁻¹

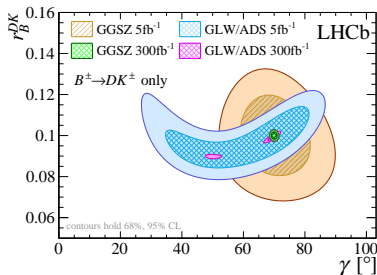
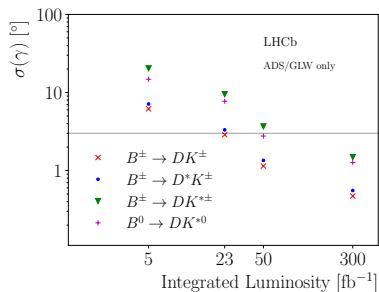
Main systematic uncertainties with rate and asymmetry measurements:

- Production and instrumentation asymmetries
- Backgrounds and their asymmetries.

All data-driven, so assumed to scale with data sample.

Additional subtle point to be taken into account:

- Charm mixing and CP violation in charm
- Matter effects for K_S^0 final states

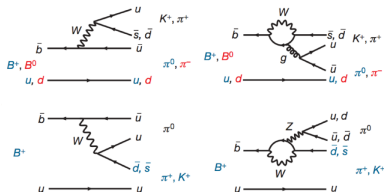


Estimated $\sigma(\gamma) \simeq 0.35^\circ$ with 300 fb^{-1} (!)

- New modes not used in the combination yet:
 - Multibody B and D decays
 - Modes with neutrals, both fully and partially-reconstructed.
 - $B_c \rightarrow DD_s$: Large CPV with $r_B \simeq 1$, but very low yields (~ 1 event in Run 1)
 - b baryons, e.g. $\Lambda_b^0 \rightarrow D\Lambda$, DpK : complication due to S and P -wave amplitudes with possibly different strong parameter. Precision depends on Λ_b^0 polarisation.
- New strategies with existing modes:
 - “Unified” approach for GLW/ADS/GGSZ with common treatment of systematic uncertainties [\[arXiv:1804.05597\]](#)
 - Fourier analysis instead of binning for model-independent GGSZ approach: squeeze last bits of statistical power [\[EPJC \(2018\) 78: 121\]](#)

$\alpha, \beta(s), \gamma$ from charmless decays

- Extraction of γ and β_s from time-dependent asymmetries of $B^0 \rightarrow \pi^+ \pi^-$ and $B_s^0 \rightarrow K^+ K^-$
- Use U -spin symmetry to disentangle contributions from different topologies



Data sample	$C_{\pi^+\pi^-}$	$S_{\pi^+\pi^-}$	$C_{K^+K^-}$	$S_{K^+K^-}$	$A_{K^+K^-}^{\Delta\Gamma}$
Run 1 (3fb ⁻¹ 112)	$-0.34 \pm 0.06 \pm 0.01$	$-0.63 \pm 0.05 \pm 0.01$	$0.20 \pm 0.06 \pm 0.02$	$0.18 \pm 0.06 \pm 0.02$	$-0.79 \pm 0.07 \pm 0.10$
	σ (stat.)				
Run 1-3 (23 fb ⁻¹)	0.015	0.013	0.015	0.015	0.018
Run 1-6 (300fb ⁻¹)	0.004	0.004	0.004	0.004	0.005

Are an input to global analysis which allows to measure γ and $-2\beta_s$.

Other charmless measurements:

- $\alpha = \pi - \beta - \gamma$ from $B \rightarrow \pi\pi, B \rightarrow \rho\rho$: main sensitivity will come from Belle II (neutral final states)
- $B_{60} \rightarrow \pi^+ \pi^- \pi^0$ TD Dalitz plot analysis, measure α . Expect large yields at LHCb.
- $B_{(s)}^0 \rightarrow K_S^0 hh$: determine CP violating parameters for resonant contributions, related to γ and β_s

- UT angle measurements provide essential information about CP violation mechanism, together with other CKM measurements, a sensitive probe for physics beyond SM.
- No major show-stopper for LHCb Upgrade I and Upgrade II.
- But **a lot** of efforts will be needed to make use of huge statistics of 300 fb^{-1} .
 - Computing, storage, bookkeeping
 - Analysis infrastructure
 - Personpower for many additional modes
 - Subtle systematic effects at sub-% level
 - Combination with cross-feed between tens (or hundreds) of modes

All of that is doable, but again, **a lot** of efforts will be needed.

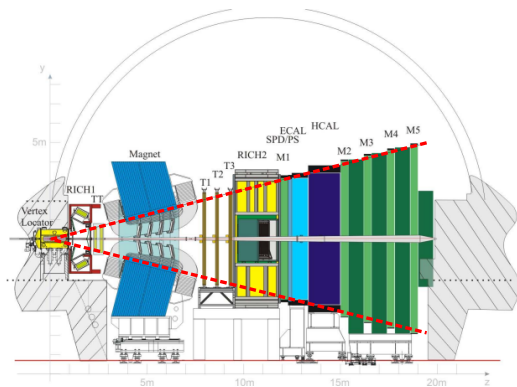
[LHCb-PUB-2018-009: "Physics case for an LHCb Upgrade II"]

Observable	Current LHCb	LHCb 2025	Belle II	Upgrade II	ATLAS & CMS
EW Penguins					
R_K ($1 < q^2 < 6 \text{ GeV}^2 c^4$)	0.1 274	0.025	0.036	0.007	–
R_{K^*} ($1 < q^2 < 6 \text{ GeV}^2 c^4$)	0.1 275	0.031	0.032	0.008	–
R_ϕ, R_{pK}, R_π	–	0.08, 0.06, 0.18	–	0.02, 0.02, 0.05	–
CKM tests					
γ , with $B_s^0 \rightarrow D_s^+ K^-$	$(^{+17}_{-22})^\circ$ 136	4°	–	1°	–
γ , all modes	$(^{+5.0}_{-5.8})^\circ$ 167	1.5°	1.5°	0.35°	–
$\sin 2\beta$, with $B^0 \rightarrow J/\psi K_s^0$	0.04 609	0.011	0.005	0.003	–
ϕ_s , with $B_s^0 \rightarrow J/\psi \phi$	49 mrad 44	14 mrad	–	4 mrad	22 mrad 610
ϕ_s , with $B_s^0 \rightarrow D_s^+ D_s^-$	170 mrad 49	35 mrad	–	9 mrad	–
ϕ_s^{ss} , with $B_s^0 \rightarrow \phi \phi$	154 mrad 94	39 mrad	–	11 mrad	Under study 611
a_{sl}^s	33×10^{-4} 211	10×10^{-4}	–	3×10^{-4}	–
$ V_{ub} / V_{cb} $	6% 201	3%	1%	1%	–
$B_s^0, B^0 \rightarrow \mu^+ \mu^-$					
$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)/\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$	90% 264	34%	–	10%	21% 612
$\tau_{B_s^0 \rightarrow \mu^+ \mu^-}$	22% 264	8%	–	2%	–
$S_{\mu\mu}$	–	–	–	0.2	–
$b \rightarrow c\ell\bar{\nu}_\ell$ LUV studies					
$R(D^*)$	0.026 215 217	0.0072	0.005	0.002	–
$R(J/\psi)$	0.24 220	0.071	–	0.02	–
Charm					
$\Delta A_{CP}(KK - \pi\pi)$	8.5×10^{-4} 613	1.7×10^{-4}	5.4×10^{-4}	3.0×10^{-5}	–
A_Γ ($\approx x \sin \phi$)	2.8×10^{-4} 240	4.3×10^{-5}	3.5×10^{-4}	1.0×10^{-5}	–
$x \sin \phi$ from $D^0 \rightarrow K^+ \pi^-$	13×10^{-4} 228	3.2×10^{-4}	4.6×10^{-4}	8.0×10^{-5}	–
$x \sin \phi$ from multibody decays	–	$(K3\pi) 4.0 \times 10^{-5}$	$(K_s^0 \pi\pi) 1.2 \times 10^{-4}$	$(K3\pi) 8.0 \times 10^{-6}$	–

Backup

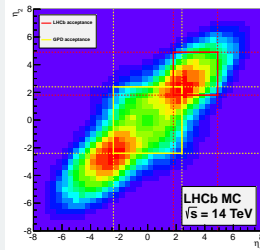


One-arm spectrometer optimised for studies of beauty and charm decays at LHC



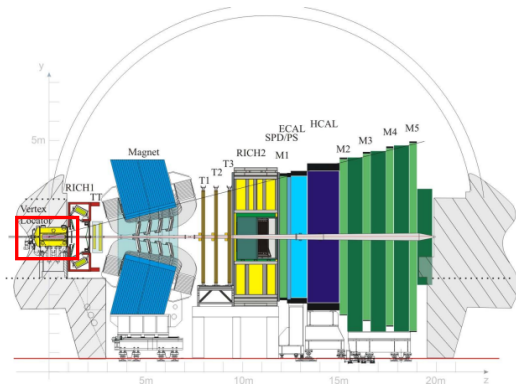
Rapidity coverage

$$2 < \eta < 5$$

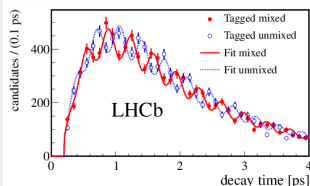


- Covers forward region (maximum of c and b production)

One-arm spectrometer optimised for studies of beauty and charm decays at LHC



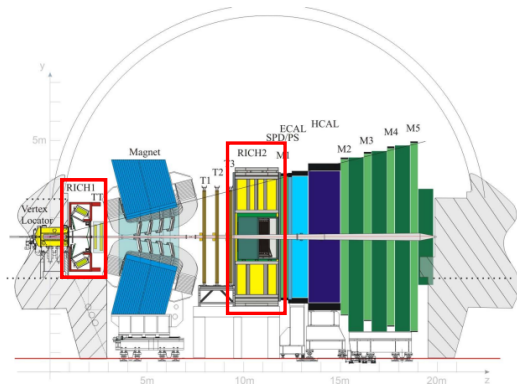
Vertexing

 B_s^0 oscillations with $B_s^0 \rightarrow D_s \pi$


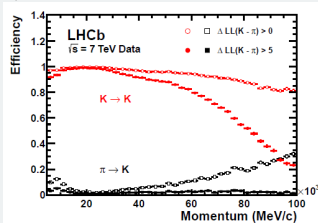
[New J. Phys. 15 (2013) 053021]

- Covers forward region (maximum of c and b production)
- Good vertexing: measure B^0 and B_s^0 oscillations, reject prompt background

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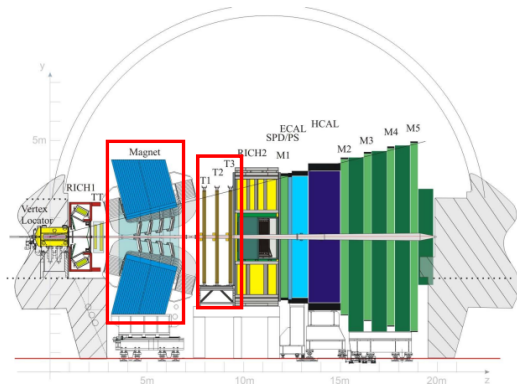
PID

 K/π ID efficiency and misID rate

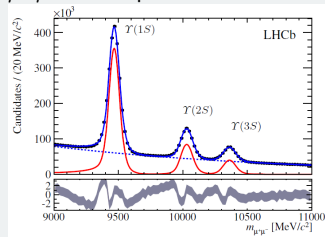
[EPJ C73 (2013) 2431]

- Covers forward region (maximum of c and b production)
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- Particle identification: flavour tagging, misID background

One-arm spectrometer optimised for studies of beauty and charm decays at LHC



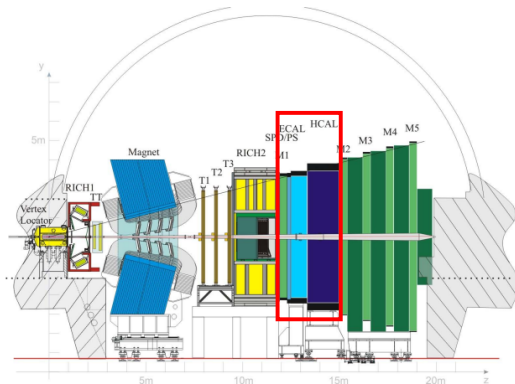
Tracking

 $\mu^+\mu^-$ mass spectrum

[PRL 111 (2013) 101805]

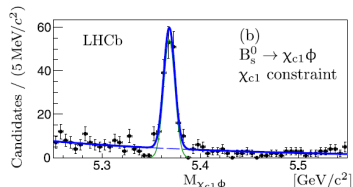
- Covers forward region (maximum of c and b production)
- Good vertexing: measure B^0 and B_s^0 oscillations, reject prompt background
- Particle identification: flavour tagging, misID background
- High-resolution tracking

One-arm spectrometer optimised for studies of beauty and charm decays at LHC



Calorimetry

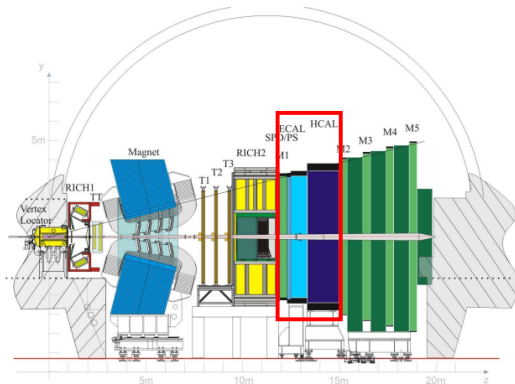
$$B_s^0 \rightarrow \chi_{c1} \phi, \chi_{c1} \rightarrow J/\psi \gamma$$



[Nucl. Phys. B874 (2013) 663]

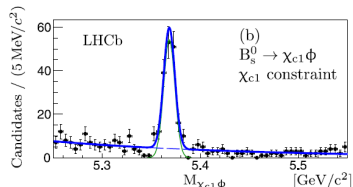
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- Calorimetry: reconstruct neutrals (π^0, γ) in the final state

One-arm spectrometer optimised for studies of beauty and charm decays at LHC



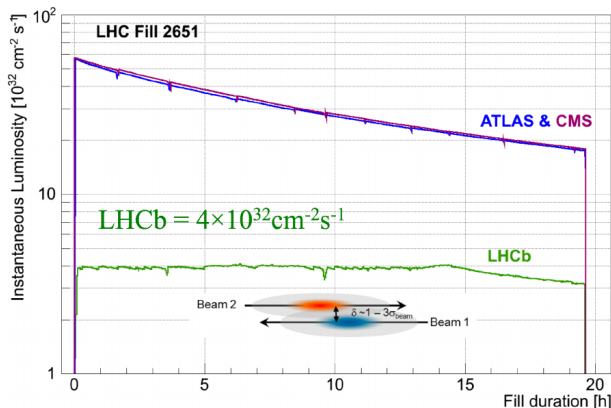
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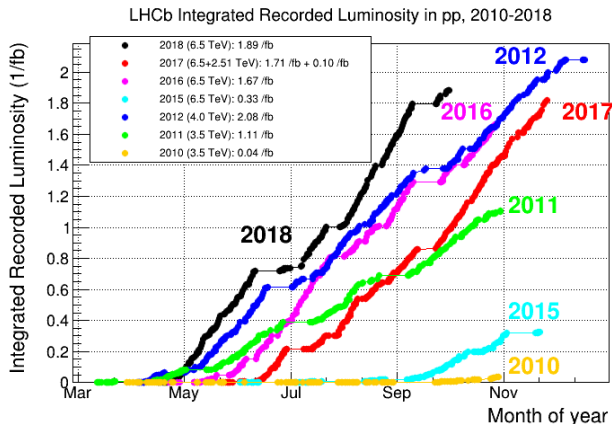


[Nucl. Phys. B874 (2013) 663]

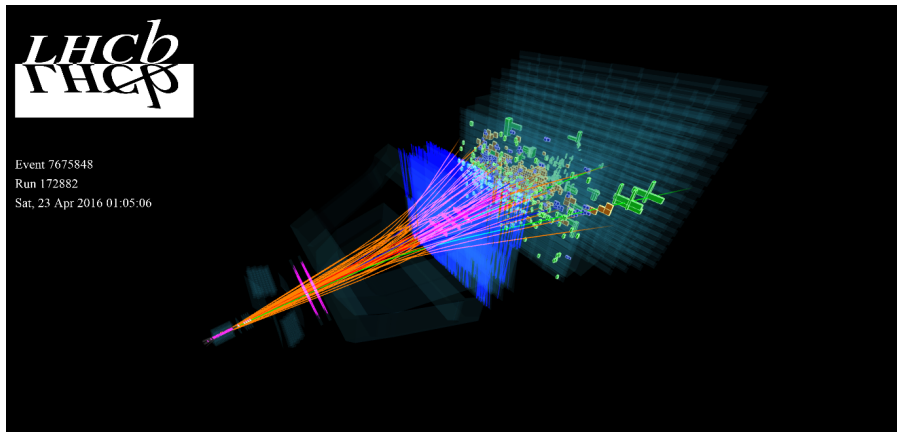
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- High-resolution tracking
- Calorimetry: reconstruct neutrals (π^0, γ) in the final state
- Efficient trigger, including fully hadronic modes



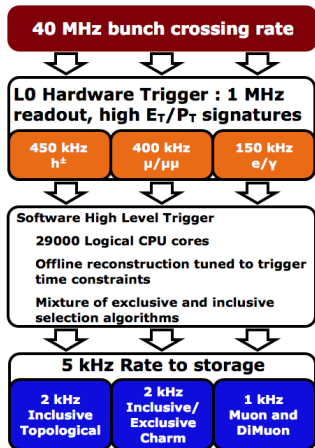
3 fb^{-1} in 2011 and 2012 (Run 1, $\sqrt{s} = 7, 8 \text{ TeV}$): **Most of results in this talk**
 2 fb^{-1} in 2015 and 2016 (Run 2, $\sqrt{s} = 13 \text{ TeV}$, higher b CS): **Analyses ongoing**
 1.71 fb^{-1} (2017) and 1.89 fb^{-1} (2018 so far) at 13 TeV



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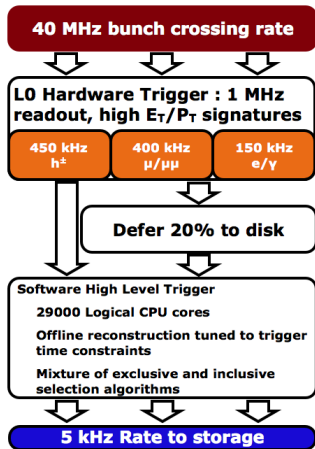


Trigger is a crucial elements in experiments at hadron machines. Need to work in a very difficult environment with hundreds of tracks in each beam crossing.



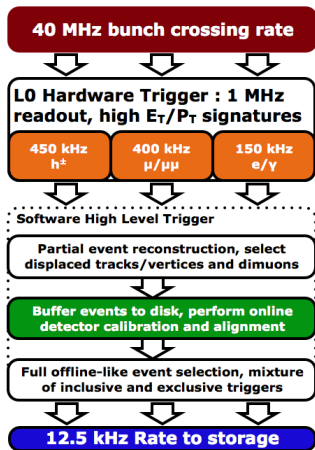
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- 2011 and early 2012: increased trigger bandwidth (compared to design 2 kHz) to accommodate charm
- 2012: *deferred trigger* configuration: keep the trigger farm busy between fills
- Since 2015: *split trigger*
 - All 1st stage (HLT1) output stored on disk
 - Used for real-time calibration and alignment
 - 2nd stage (HLT2) uses offline-quality calibration
 - 5 kHz of 12 kHz to Turbo stream:
 - Candidates produced by trigger are stored
 - No raw event \Rightarrow smaller event size
 - Used for high-yield channels (charm, J/ψ , ...)

Time-dependent measurements

Measure lifetime based on vertex displacement from the primary vertex of pp interaction.

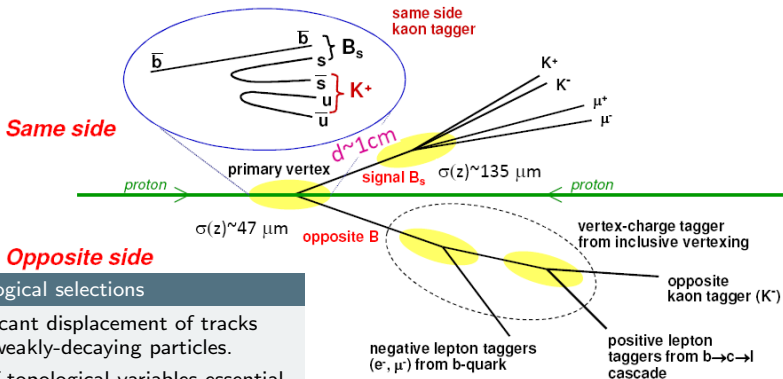
Large boost provides excellent time resolution ($\sigma_t \simeq 45$ fs)

Flavor tagging

Need to identify B flavour *at production time* (different from flavour at decay time due to oscillations).

Use decay products of the opposite-side B (OS) and π , K associated with same-side B (SS).

Effective tagging power $\epsilon_{\text{tag}} D^2 = 3.7\%$.



Topological selections

Significant displacement of tracks from weakly-decaying particles.

Use of topological variables essential to reduce combinatorial background.