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Astro-connections to lepton flavor and universality violation

XIII Meeting on B Physics: Synergy between LHC and SUPERKEKB in the Quest for New Physics

Tuesday, October 2, 2018

- **Lepton flavor physics with cosmic rays: (1) the positron excess**
- **Lepton flavor physics with cosmic rays: (2) DAMPE**
- **Dark matter and lepton flavor universality violation**

➤ **Lepton flavor physics with cosmic rays: (1) the positron excess**

➤ Lepton flavor physics with cosmic rays: (2) DAMPE

➤ Dark matter and lepton flavor universality violation

Most “famous” charged leptons
in **Astrophysics**:
Anomalous **HE Positrons**

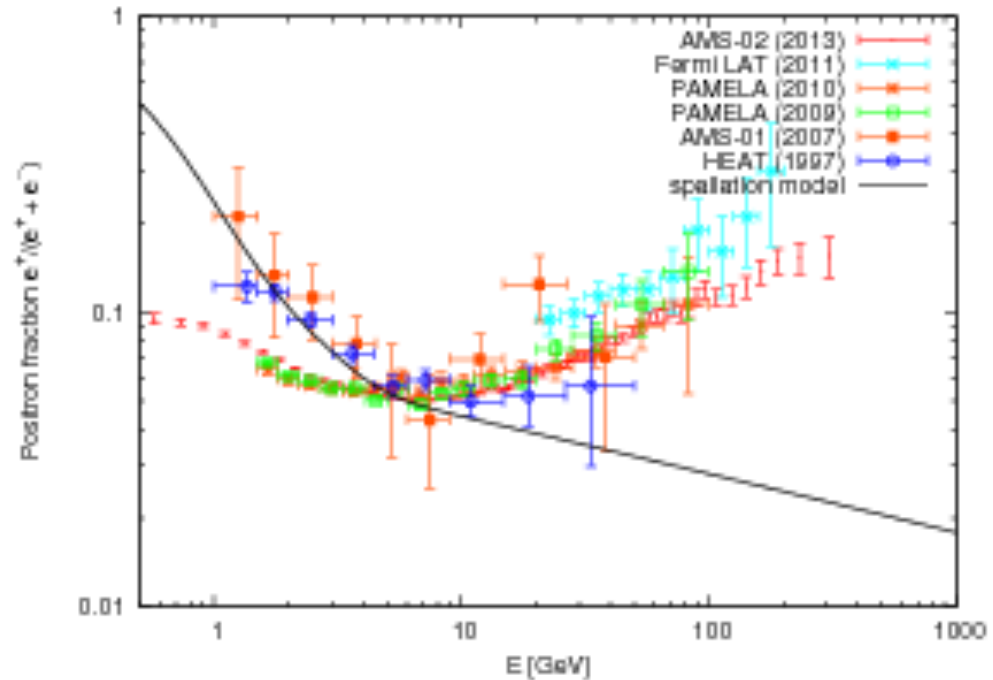
...if indeed from New Physics,
require LF universality violation

Rising Positron Fraction with energy cut-off at Dark Matter particle mass, envisioned ~30 years ago, as smoking gun for Dark Matter searches

[Tylka 1989, Turner and Wilczek, 1990]

First hint of a rising positron fraction >20 year old!

- ✓ **HEAT 1997**
- ✓ **Pamela 2009**
- ✓ **Fermi 2010**
- ✓ **AMS 2013, 2015**



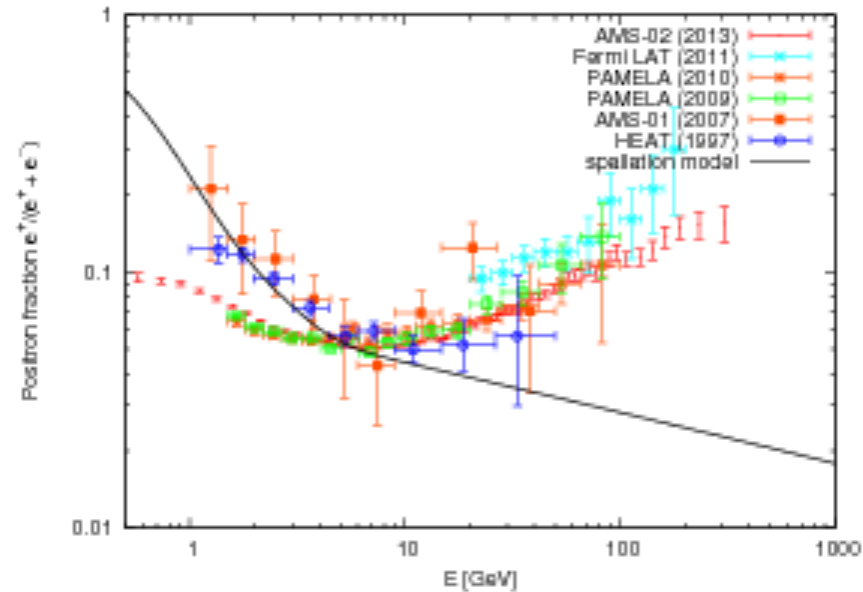
Decreasing positron fraction assumes
exclusive secondary origin

[Physics: $D(E) \sim E^\delta$]

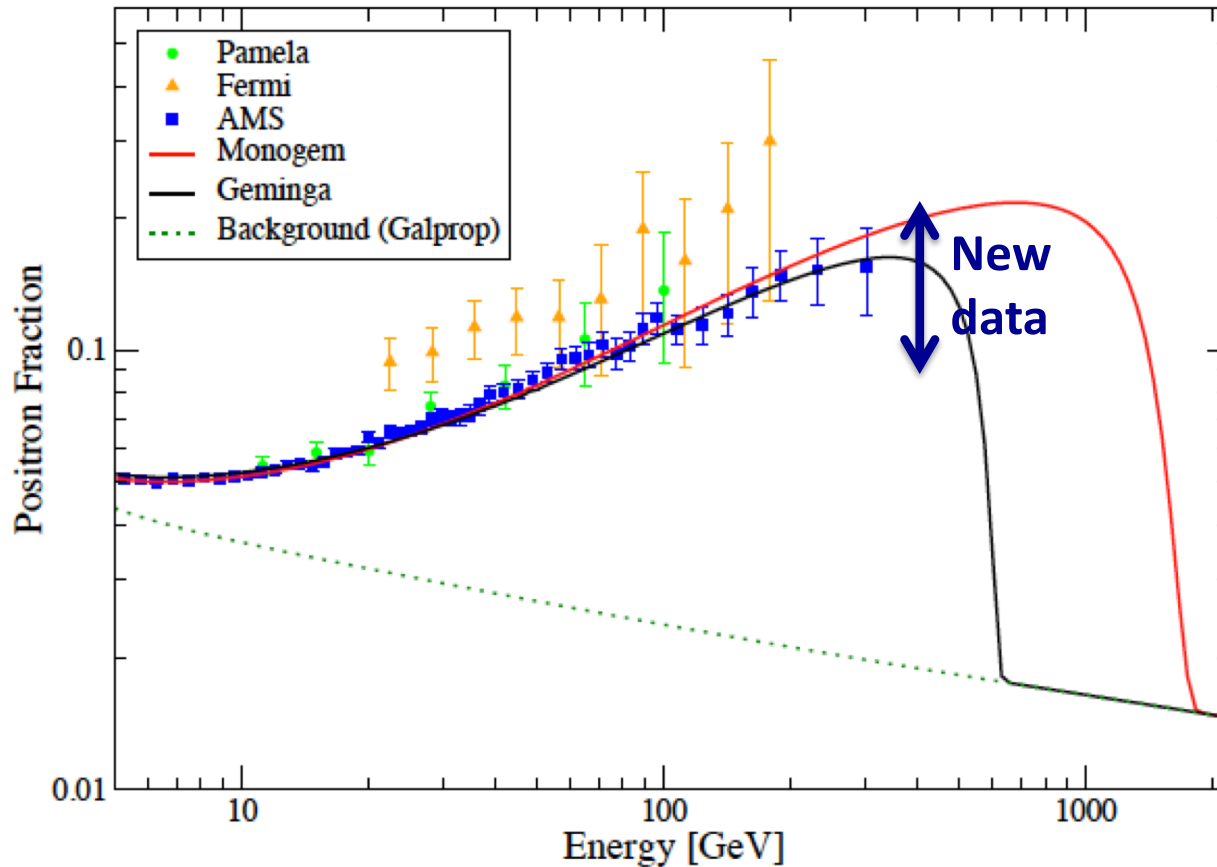
Hence **rising** fraction = **excess**

Caveats:

- in-source secondary **reacceleration**
(**ruled out** by B/C)
- **primary** production (e.g. PSR)



PSRs work perfectly well



hardly any free parameter!

Cutoff (unlike what Sam Ting says)
is **not** a smoking gun for DM!

$$\frac{dE}{dt} = -bE^2 \quad \int_{\infty}^{E_{\max}} \frac{dE}{E^2} = -bT_{PSR} \quad E_{\max} = \frac{1}{bT_{PSR}}$$

Observing a **cutoff** will likely
help pinpointing **relevant PSR(s)**

Known **PSR** OK

While **AMS** continues to increase **statistics***, a **new** high-impact observational **result** appeared!

* and to **claim** that a cutoff is a smoking gun for dark matter – **FAKE NEWS!**



Moon (To Scale)

Geminga



PSR B0656+14

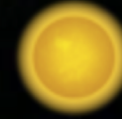




Extended gamma-ray sources around pulsars constrain the origin of the positron flux at Earth

A. U. Abeysekara,¹ A. Albert,² R. Alfaro,³ C. Alvarez,⁴ J. D. Álvarez,⁵ R. Arceo,⁴
J. C. Arteaga-Velázquez,⁵ D. Avila Rojas,³ H. A. Ayala Solares,⁶ A. S. Barber,¹
N. Bautista-Figueroa,⁷ A. Benford,³ E. Belmont-Moreno,³ S. V. Berzini,⁸ D. Berlyan,⁹ A. Bernal,¹⁰

demonstrate that these pulsars are indeed local sources of accelerated leptons, but the measured tera-electron volt emission profile constrains the diffusion of particles away from these sources to be much slower than previously assumed. **We demonstrate that the leptons emitted by these objects are therefore unlikely to be the origin of the excess positrons, which may have a more exotic origin.**



Home / News / **New pulsar result supports particle dark matter**



New pulsar result supports particle dark matter

The nature of dark matter remains elusive, but astronomers are now one step closer to the answer.

By Robert Naeye | Published: Thursday, November 16, 2017

dark matter and energy

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NEWS SPACE 17 NOVEMBER 2017

Pulsars out for positrons leaving dark matter

New modelling discounts a leading theory on anti-matter
Andrew Masterson reports.

Sky-high observatory sheds light on origin of excess anti- matter

New study excludes nearby pulsars, points to dark matter as possible culprit

Chris Cesare, Miguel Mostafá, Gail McCormick
November 16, 2017

UNIVERSITY PARK, Pa. — The High-Altitude Water Cherenkov (HAWC) Observatory in Mexico, built and operated by an international team that

from dark matter after all

My key **problem**: (while writing numerous papers on the dark matter interpretation) I have a **decade-old emotional attachment** to the **pulsar** interpretation, that **named names...**

Dissecting Pamela (and ATIC) with Occam's Razor: existing, well-known Pulsars naturally account for the "anomalous" Cosmic-Ray Electron and Positron Data

Stefano Profumo^{1,2}

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²*Santa Cruz Institute for Particle Physics, University of California, Santa Cruz, CA 95064, USA*

(Dated: April 14, 2018)

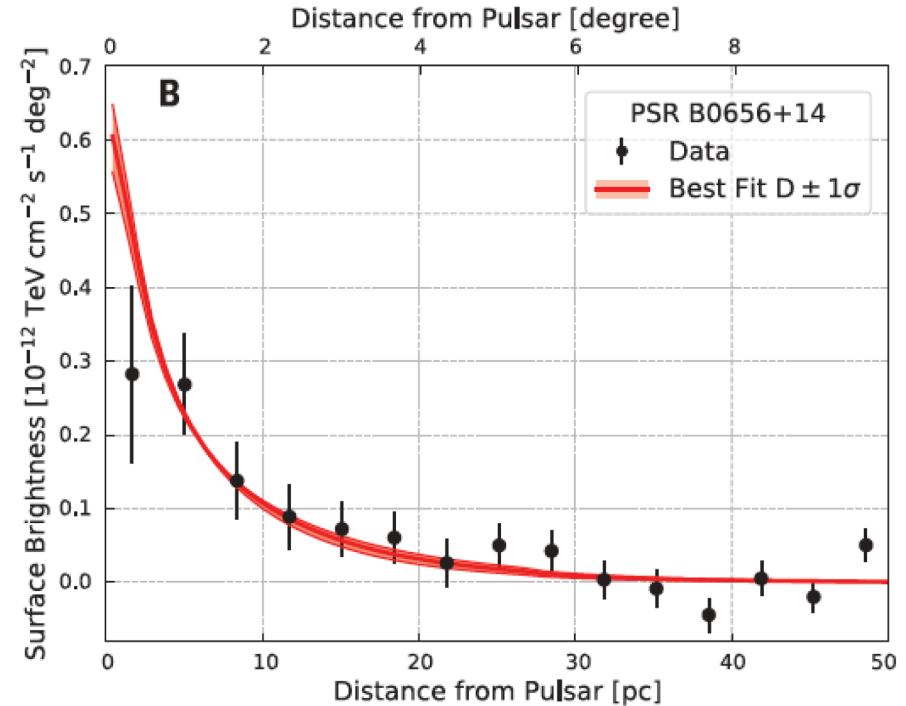
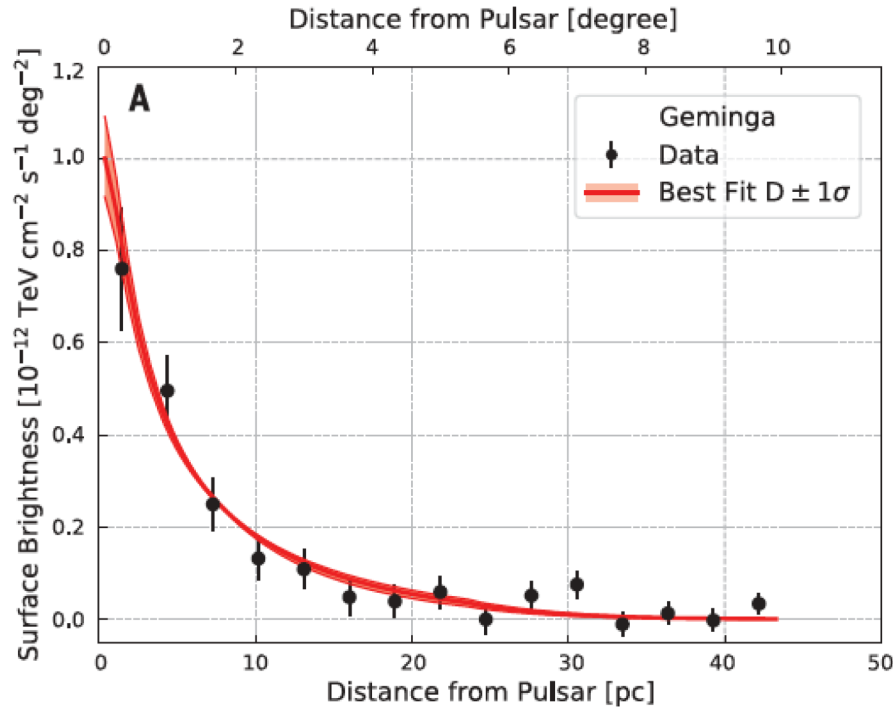
We argue that both the positron fraction measured by PAMELA and the peculiar spectral features reported in the total differential electron-positron flux measured by ATIC have a very natural explanation in electron-positron pairs produced by nearby pulsars. While this possibility was pointed

Name	Distance [kpc]	Age [yr]	\dot{E} [ergs/s]	E_{out} [ST]	E_{out} [CCY]	E_{out} [HR]	E_{out} [ZC]	$f_{e\pm}$	g
Geminga [J0633+1746]	0.16	3.42×10^5	3.2×10^{34}	0.360	0.344	0.013	0.053	0.005	0.70
Monogem [B0656+14]	0.29	1.11×10^5	3.8×10^{34}	0.084	0.456	0.004	0.372	0.015	0.14

simple theoretical models for estimating the energy output, the diffusion setup and the injection spectral index of electron-positron pairs, and by (2) considering all known pulsars (as given in the ATNF catalogue). It appears unlikely that a single pulsar be responsible for both the PAMELA result and for the ATIC excess, although two sources are enough to naturally explain both of the experimental results. The PAMELA data favor mature pulsars (age $\sim 2 \times 10^6$ yr), with a distance of 0.8-1 kpc, or a younger and closer source like Geminga or the SNR Loop I. The ATIC data require a larger (and marginally unlikely) energy output, and favor an origin associated to powerful, more distant (1-2 kpc) and younger (age $\sim 5 \times 10^5$ yr) pulsars. We list several candidate pulsars that can

23 Dec 2008

Key observational result: angular surface brightness



Gamma-ray energies as large as 20 TeV \rightarrow e^+e^- as energetic as **100 TeV**

100 TeV is deep in **KN regime** for starlight

\rightarrow only relevant photons: **CMB**

\rightarrow Direct measurement of **e^+e^- diffusion!**

Inferred **diffusion** coefficients:

	Geminga	PSR B0656+14
D_{100} (diffusion coefficient of 100-TeV electrons from joint fit of two PWNe) ($\times 10^{27}$ square centimeters per second)	4.5 ± 1.2	4.5 ± 1.2
D_{100} (diffusion coefficient of 100-TeV electrons from individual fit of PWN) ($\times 10^{27}$ square centimeters per second)	$3.2_{-1.0}^{+1.4}$	15_{-9}^{+49}

...versus **ISM** diffusion coefficient (GALPROP, AMS-02...)

$$D_{100}^{\text{ISM}} \simeq 3.86 \times 10^{28} \left(\frac{E_e}{\text{GeV}} \right)^{0.33} \text{ cm}^2/\text{s} \rightarrow 1,720 \times 10^{27} \text{ cm}^2/\text{s}$$

...thus the inferred **PWN** diffusion coefficient is **100-500 times smaller** than the ISM effective value!

(notice also the injected **power** in electrons is consistent with being a fraction of O(0.1) of spin-down power)

What does this mean?

Go back to electron **transport**

Transport model: diffusive (**Brownian**) motion + energy losses

$$\frac{\partial \psi}{\partial t} = \vec{\nabla} \cdot (D(\vec{x}, E) \vec{\nabla} \psi) + \frac{\partial}{\partial E} (P(E) \psi) + Q$$

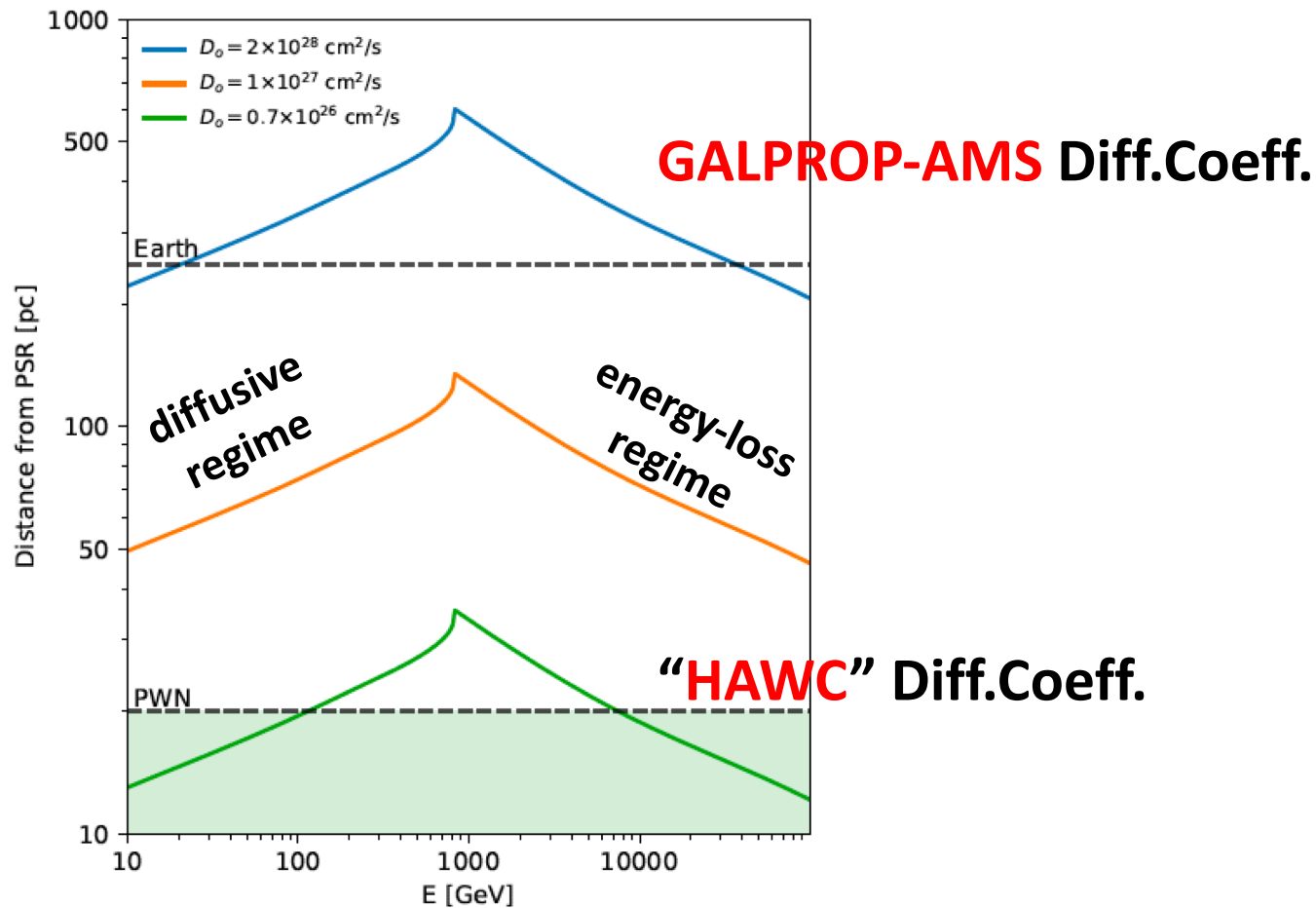
This PDE has a known **Green** function

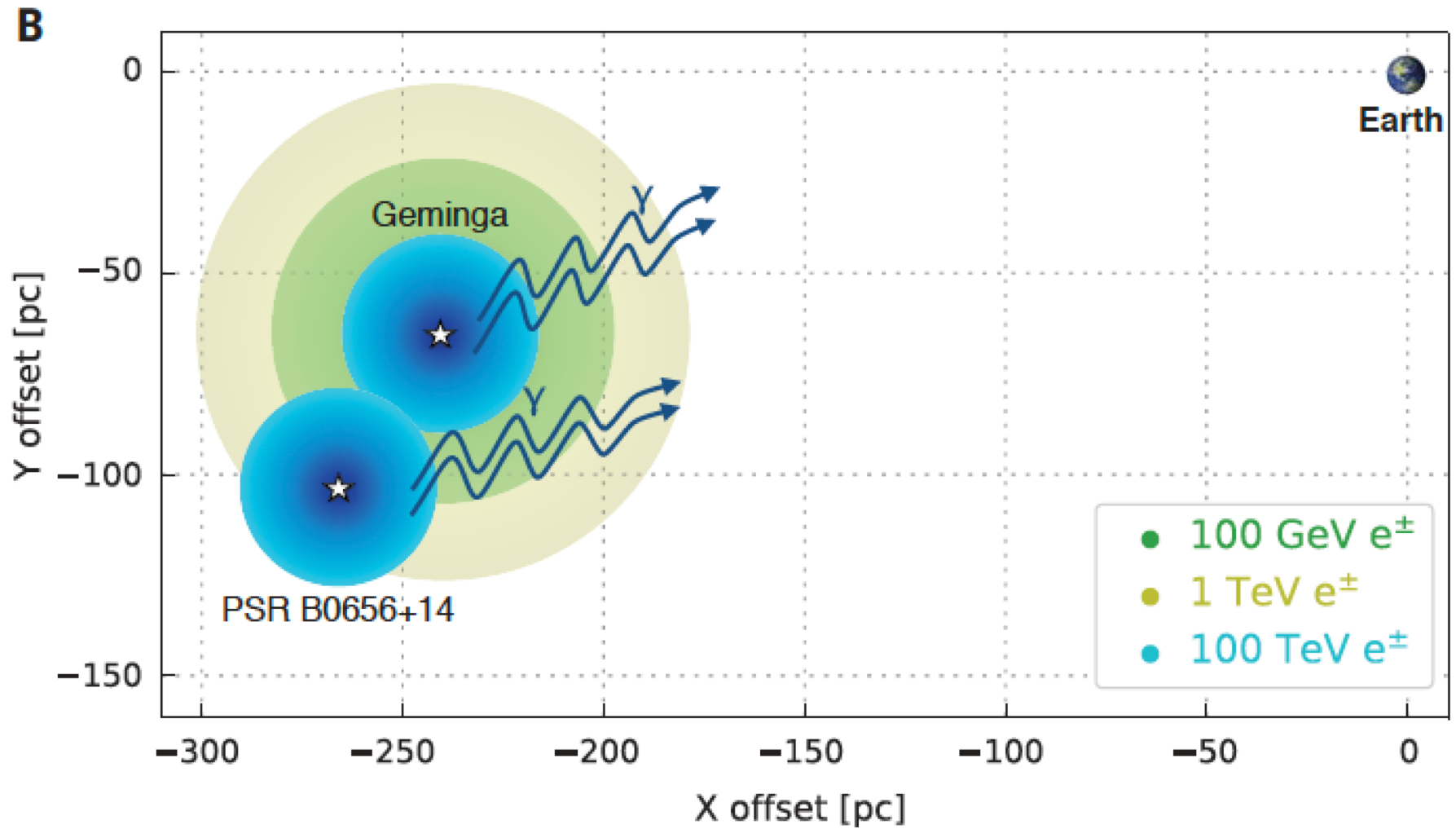
$$\psi(t, r, E) = \frac{N_o(E_o) P(E_o)}{\pi^{3/2} P(E) r_{\text{diff}}^3} e^{-r^2/r_{\text{diff}}^2}, \quad r_{\text{diff}}^2 = \int_E^{E_o} D(x)/P(x) dx.$$

$$\psi(t, r, E) = \frac{N_o(E_o)P(E_o)}{\pi^{3/2}P(E)r_{\text{diff}}^3} e^{-r^2/r_{\text{diff}}^2},$$

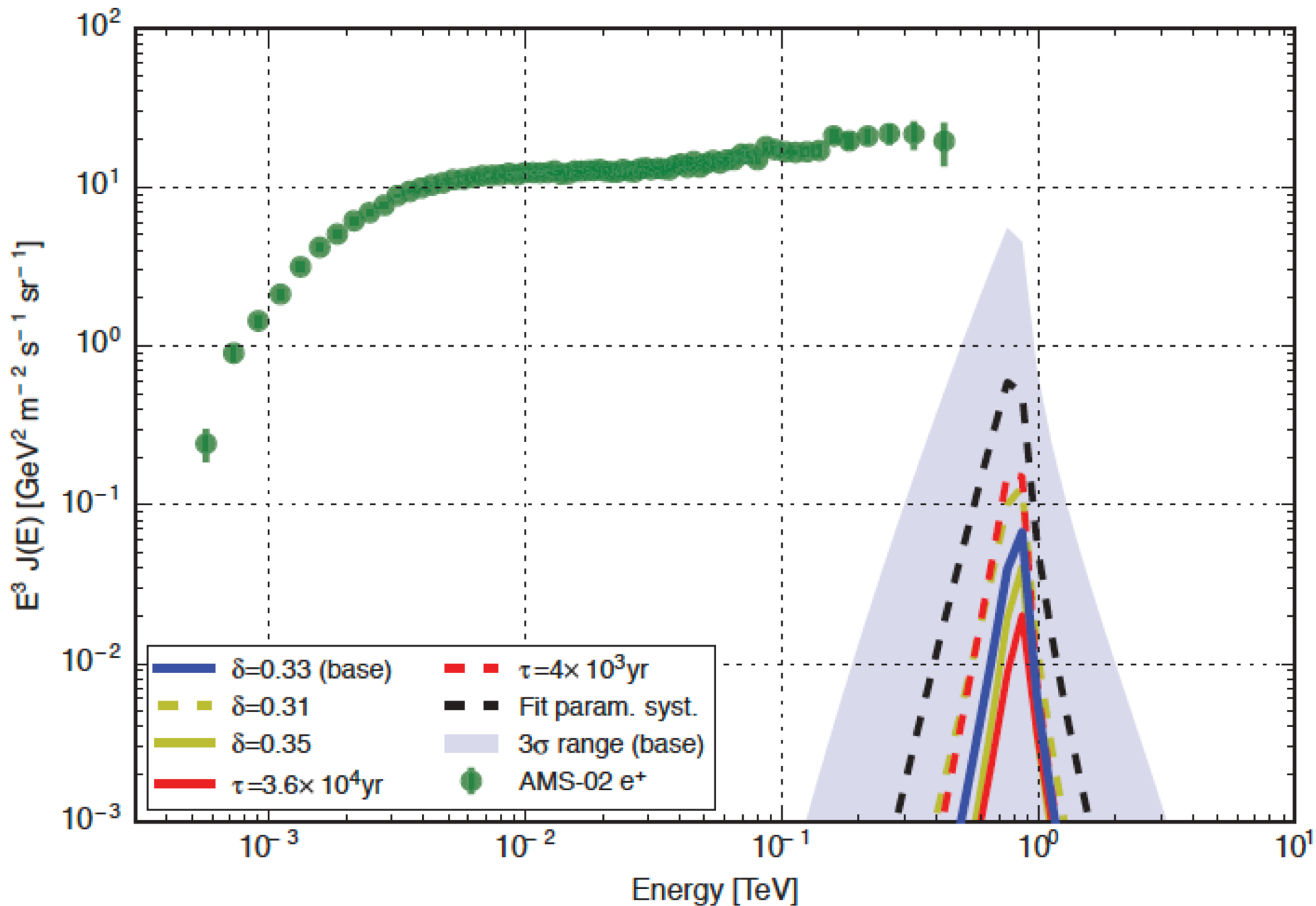
$$r_{\text{diff}}^2 = \int_E^{E_o} D(x)/P(x)dx.$$

What is the **diffusion radius** for relevant electron energies?





* Abaysekara et al (HAWC Coll.) 2017



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Extended gamma-ray sources around pulsars constrain the origin of the positron flux at Earth

A. U. Abeysekara,¹ A. Albert,² R. Alfaro,³ C. Alvarez,⁴ J. D. Álvarez,⁵ R. Arceo,⁴ J. C. Arteaga-Velázquez,⁵ D. Avila Rojas,³ H. A. Ayala Solares,⁶ A. S. Barber,¹

Geminga and PSR B0656+14 are the oldest pulsars for which a tera-electron volt nebula has so far been detected. Under our assumption of isotropic and homogeneous diffusion, the dominant source of the positron flux above 10 GeV cannot be either Geminga or PSR B0656+14. Under the unlikely situation that the field is nearly

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NEWS SPACE 17 NOVEMBER 2017

Pulsars out for positron count, leaving dark matter in the frame

New modelling discounts a leading theory on antimatter particles hitting Earth. Andrew Masterson reports.

Is this conclusion **plausible?**

Very probably **NO**.

Two **key** arguments:

1. **Lifetime** of TeV electrons is **short**: $\tau_e \sim 3 \times 10^5 \text{ yr} \times (1 \text{ TeV}/E_e)$.

We observe directly CR electrons with energies **>20 TeV**

$$d \lesssim \sqrt{D\tau_e}$$

for HAWC's Diffusion vCoeff., this means a source within 10-20 pc.

Such a source however **doesn't exist!**

Is this conclusion **plausible**?

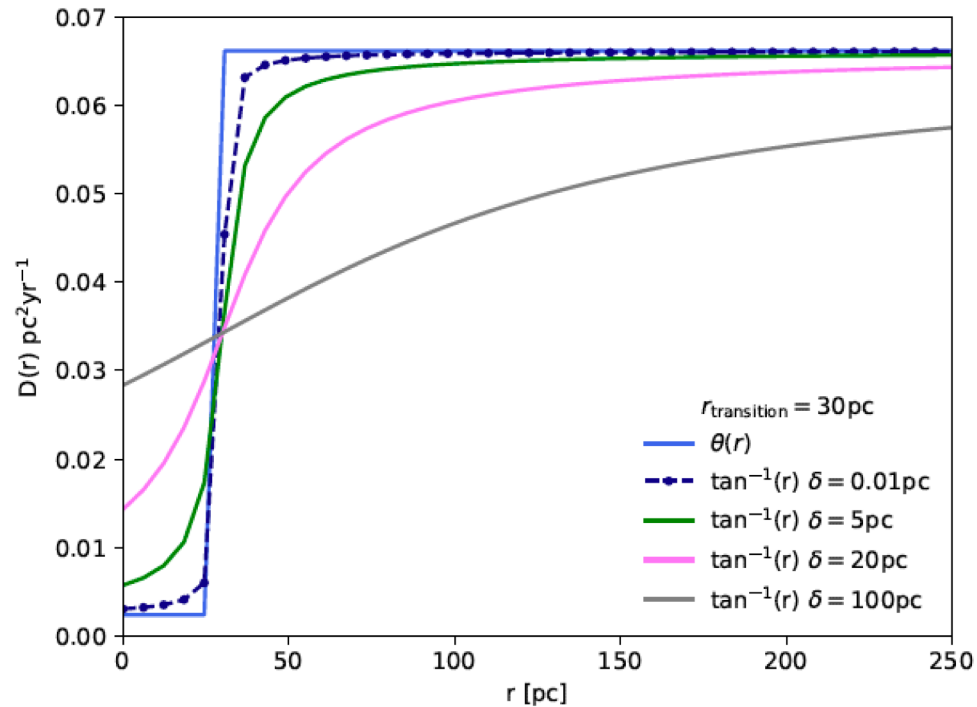
Very probably **NO**.

Two **key** arguments:

2. Models of CR emission **predict** inefficient diffusion near sources

Alfven waves generated by cosmic rays induce a net force that **suppresses diffusion** near the sites of cosmic-ray acceleration and, more generally, where cosmic-ray fluxes are larger

What happens to the local electron flux if indeed **diffusion** is **not homogeneous**?



$$D_{\theta}(r) = D_1 \theta(r_T - r) + D_2 \theta(r - r_T)$$

$$D_T(r) = D_1 + \frac{(D_2 - D_1)}{\pi} \left(\tan^{-1} \left(\frac{r - r_T}{\delta} \right) + \frac{\pi}{2} \right)$$

What happens to the local electron flux
if indeed **diffusion** is **not homogeneous**?

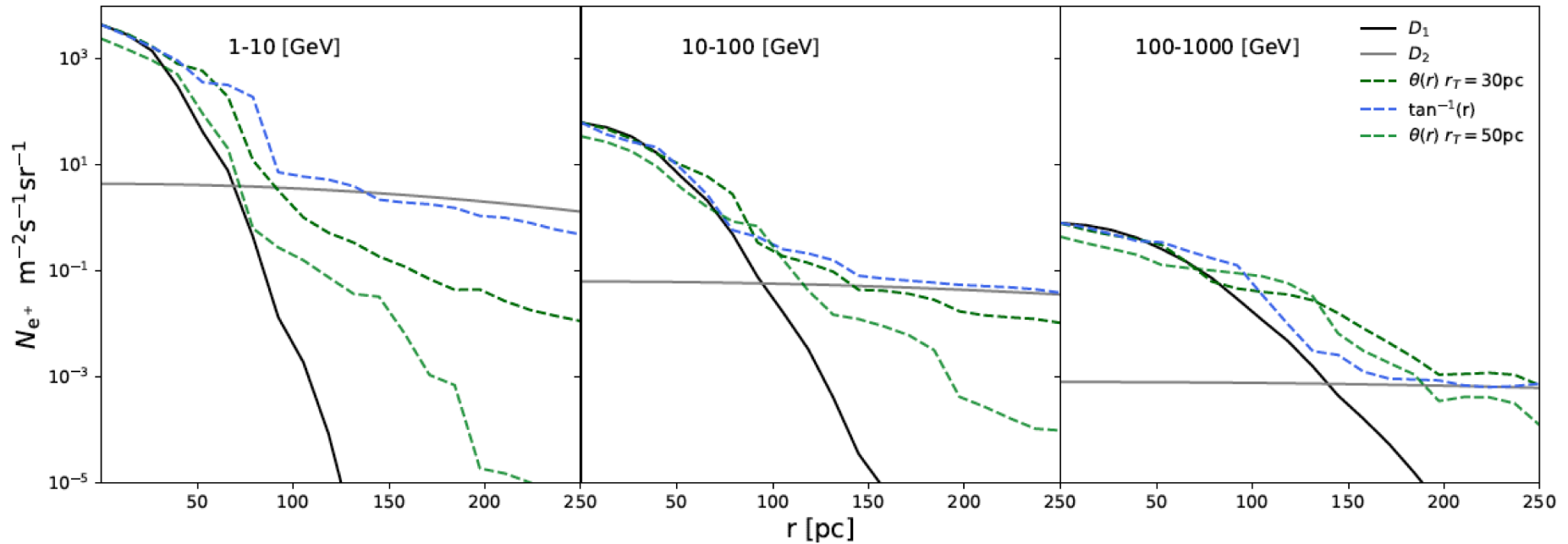
$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$$

$$f(t, x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dx' f_o(x') \left\{ \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{(x-x')^2}{4Dt}} \right\}$$

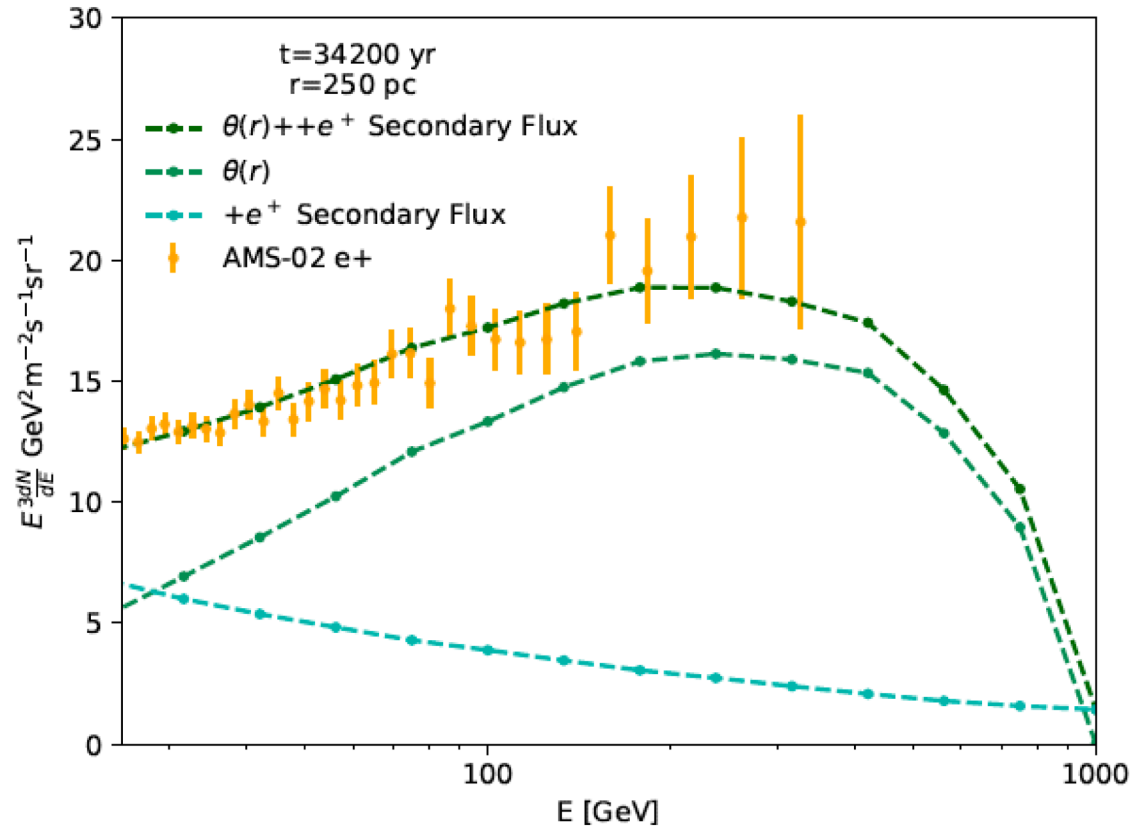
$$x = x_o + \eta \sqrt{2D\Delta t},$$

$$\Delta t \rightarrow \Delta u(E) \quad \Delta u(E_i) = \int_{E_i + \Delta E}^{E_i} \frac{D(x)}{P(x)} dx,$$

What happens to the local electron flux if indeed **diffusion** is **not homogeneous**?



What happens to the local electron flux if indeed **diffusion is not homogeneous**?



How can we **test** inhomogeneous diffusion? Does it **matter**, globally on Galactic scales?

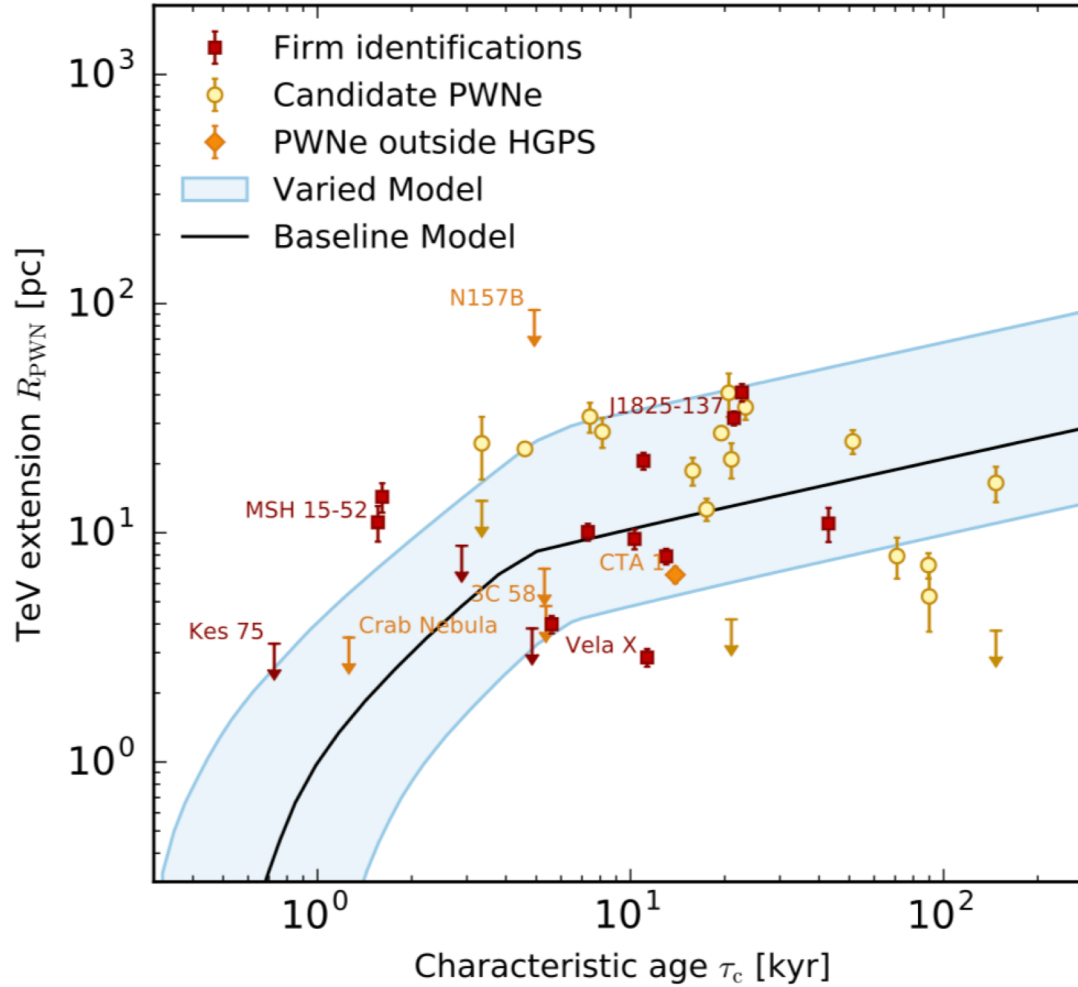
Estimate the **volume** of regions of inefficient diffusion

1. How **big** is a **PWN** as a function of time?

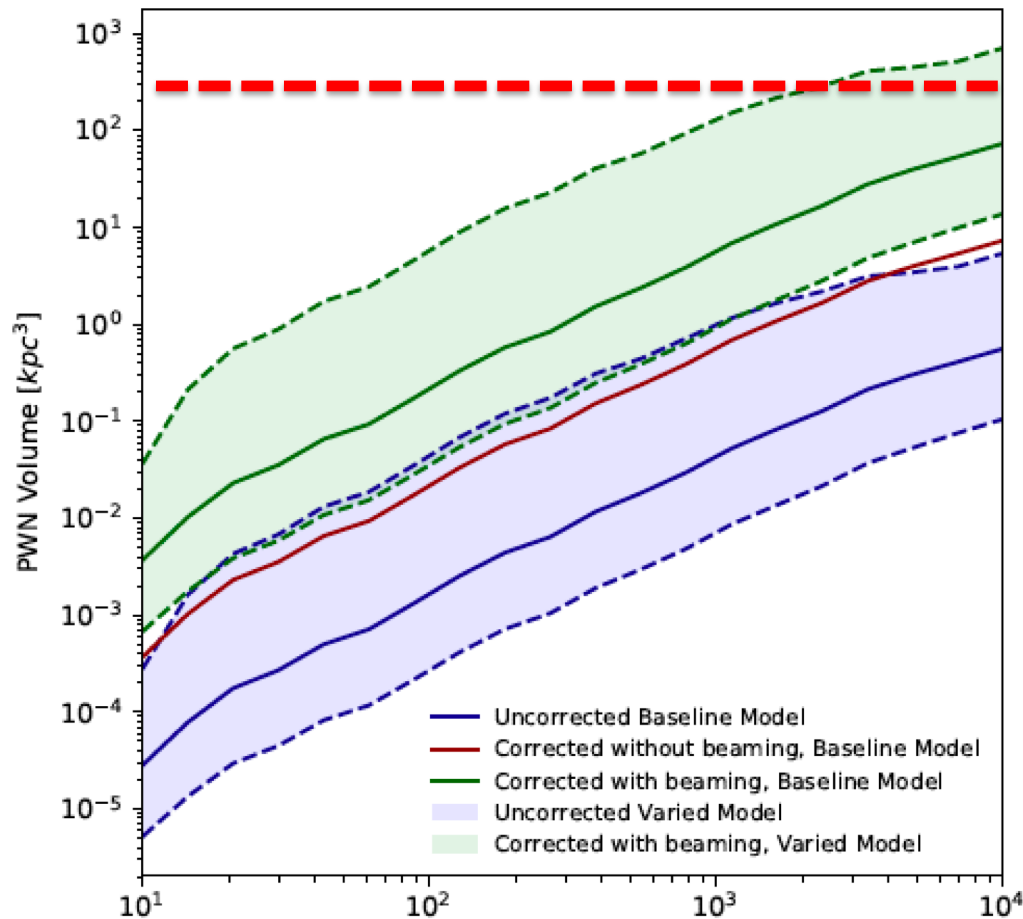
The population of TeV pulsar wind nebulae in the H.E.S.S. Galactic Plane Survey

H.E.S.S. Collaboration, H. Abdalla¹, A. Abramowski², F. Aharonian^{3,4,5}, F. Ait Benkhali³, A.G. Akhperjanian^{†6,5}, T. Andersson¹⁰, E.O. Angüiner⁷, M. Arrieta¹⁵, P. Aubert²⁴, M. Backes⁸, A. Balzer⁹, M. Barnard¹, Y. Becherini¹⁰, J. Becker Tjus¹¹, D. Berge¹², S. Bernhard¹³, K. Bernlöhr³, R. Blackwell¹⁴, M. Böttcher¹, C. Boisson¹⁵, J. Bolmont¹⁶, P. Bordas³, J. Bregeon¹⁷, F. Brun²⁶, P. Brun¹⁸, M. Bryan⁹, T. Bulik¹⁹, M. Capasso²⁹, J. Carr²⁰, S. Carrigan^{‡,3}, S. Casanova^{21,3}, M. Cerruti¹⁶, N. Chakraborty³, R. Chalme-Calvet¹⁶, R.C.C. Chaves^{17,22}, A. Chen²³, I. Chevalier²⁴, M. Chrétien¹⁶, S. Colafrancesco²³, G. Colonna²⁵, B. Condon²⁶, J. Conrad^{27,28}

1. How **big** is a **PWNe** as a function of time?



...but of course the sample is **incomplete**... (beaming+detectability)
 ...and we don't know when PWN run **out of steam**...



$$\langle V \rangle_{\text{ISM}} \simeq 500 \text{ kpc}^3 \left(\frac{R_h}{20 \text{ kpc}} \right)^2 \left(\frac{z_h}{0.2 \text{ kpc}} \right)$$

so, does this **matter**?

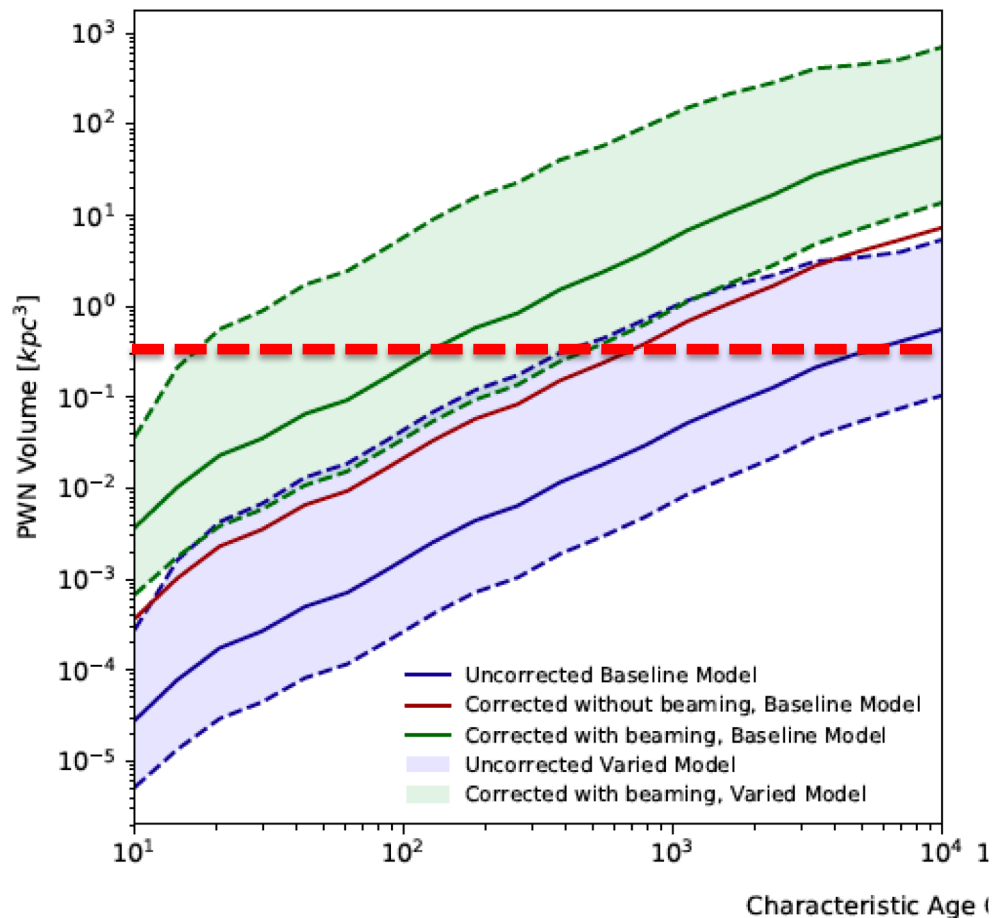
well, the time spent in inefficient diffusion pockets is potentially **much larger** than volume ratios!

$$\langle L \rangle \sim \sqrt{D \cdot t}, \quad t \sim \frac{\langle L \rangle^2}{D} \sim \frac{V^{2/3}}{D}$$

$$\frac{t_{\text{PWN}}}{t_{\text{ISM}}} \sim \left(\frac{\langle V \rangle_{\text{PWN}}}{\langle V \rangle_{\text{ISM}}} \right)^{2/3} \quad \frac{D_{\text{ISM}}}{D_{\text{PWN}}} \sim 10^{-2} \left(\frac{\langle V \rangle_{\text{PWN}}}{\langle V \rangle_{\text{ISM}}} \right)^{2/3}$$

$$\langle V \rangle_{\text{PWN}} \gtrsim 0.5 \text{ kpc}^3$$

$$\langle V \rangle_{\text{PWN}} \gtrsim 0.5 \text{ kpc}^3$$



...CR's are **guaranteed** to spend most of their time in
pockets of **inefficient diffusion!**

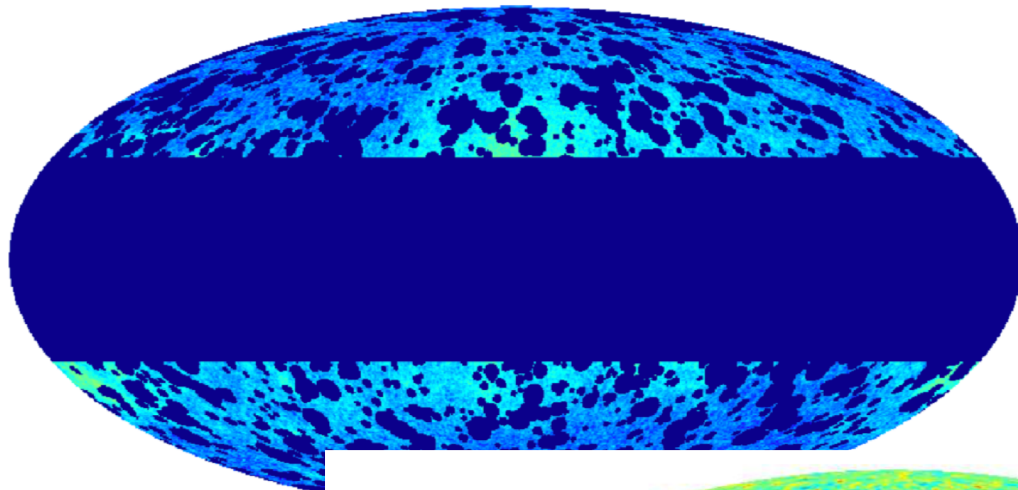
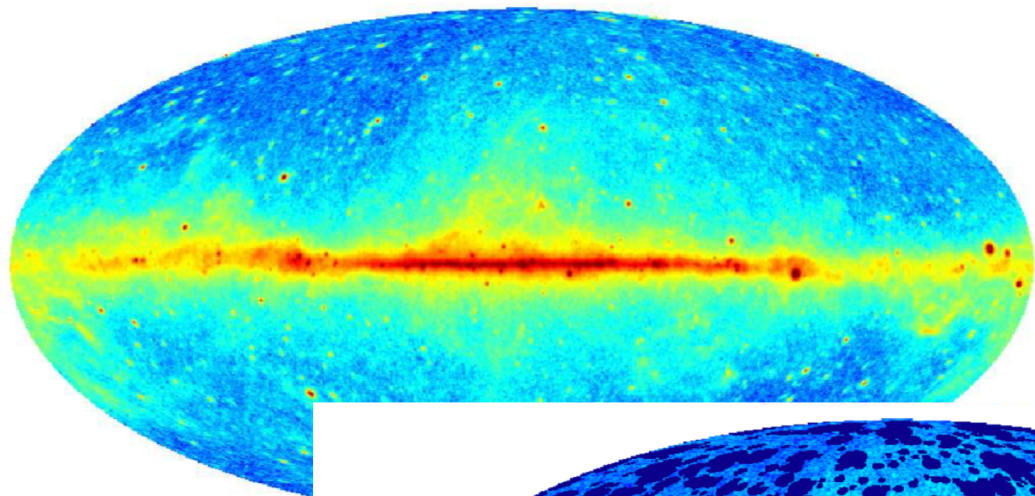
...OK, but how can we **test** this?

if a large fraction of CR electrons are **trapped** in inefficient diffusion pockets, those pockets will be **illuminated** by energy-loss **radiative processes** (radio, IC, brems)

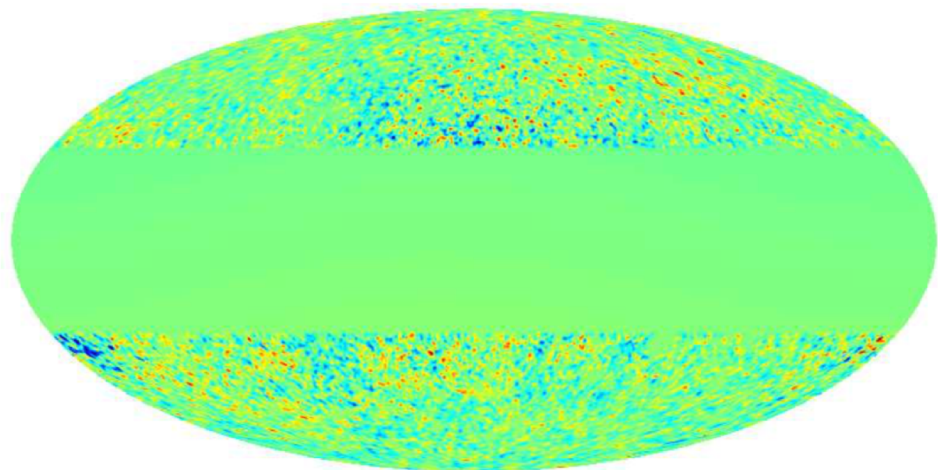
$$\theta \sim \frac{R_{\text{PWN}}}{d_{\text{PWN}}}, \quad \text{theta ranges from } \mathbf{\text{few degrees}} \text{ to } \mathbf{0.1} \text{ degrees}$$

can use any frequency from **radio** (with additional magnetic field uncertainties) to **X-ray** to **gamma rays**

Can use simple **angular power spectrum**, or **wavelet** transforms, **Poissonian** noise analysis

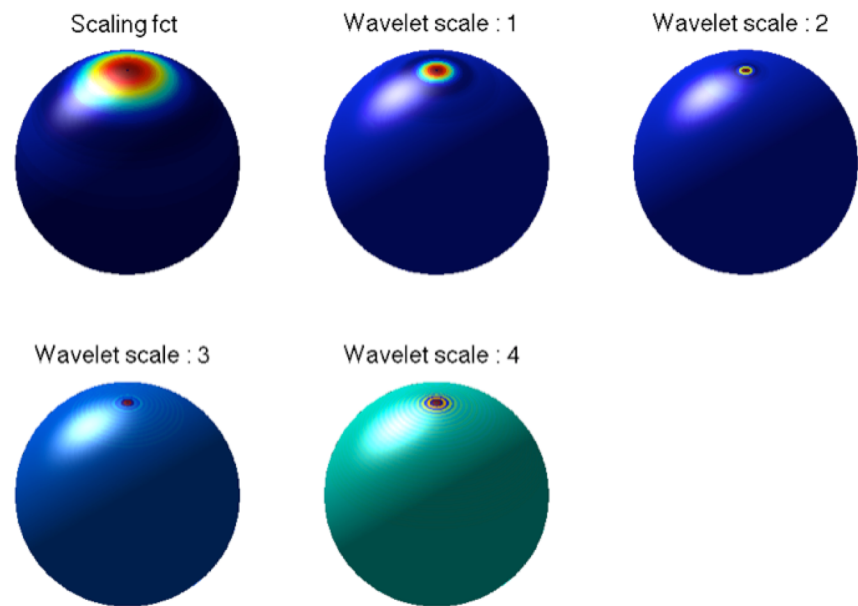
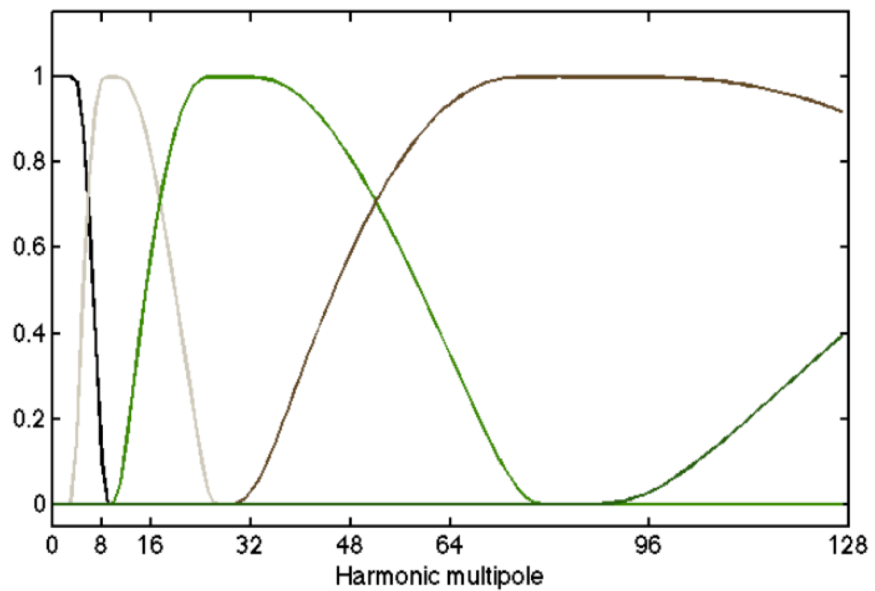


-7.4

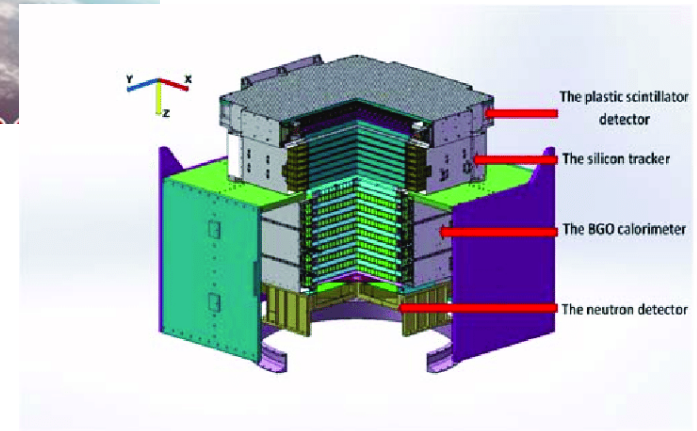
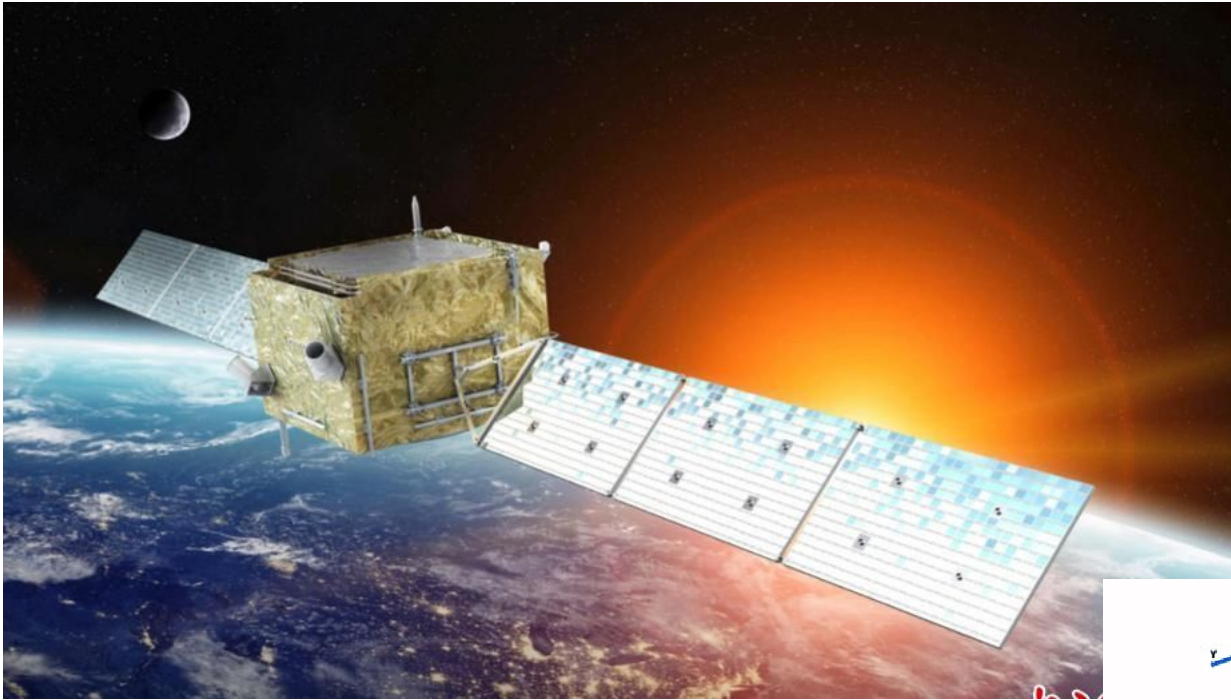


-5.0e-07  5.0e-07 Intensity [$\text{cm}^{-2}\text{s}^{-1}\text{sr}^{-1}$]

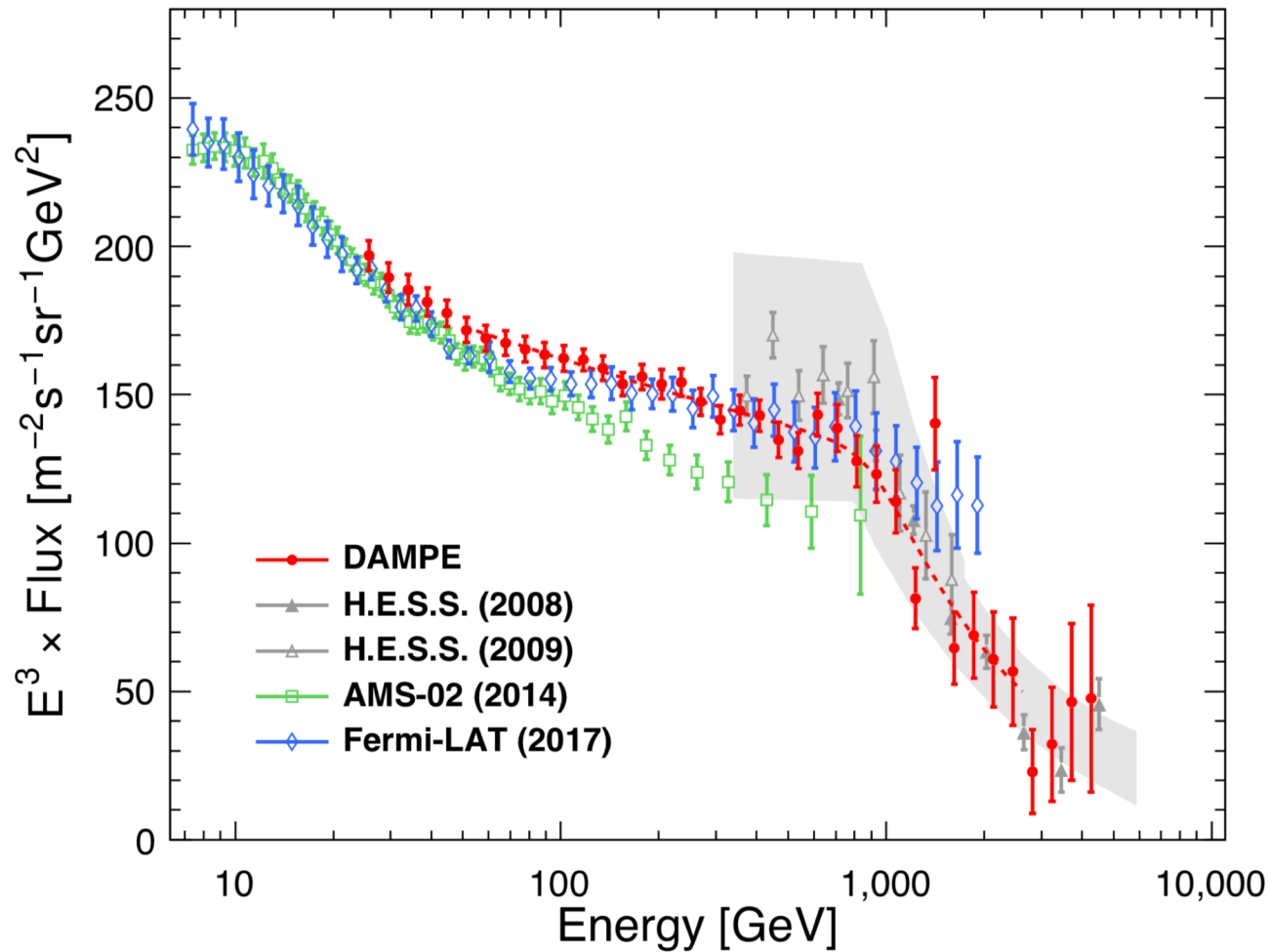
* Fornasa et al



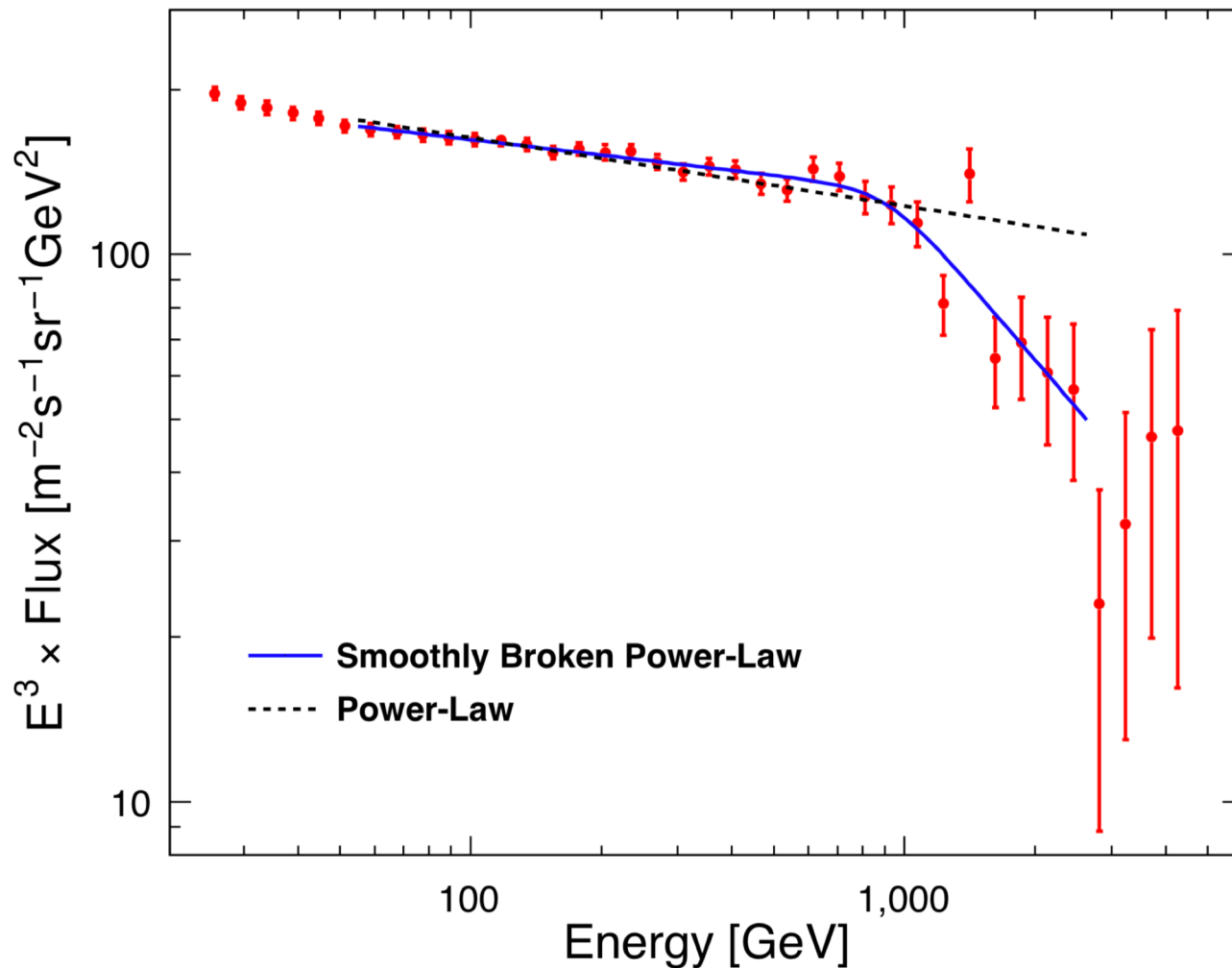
- Lepton flavor physics with cosmic rays: (1) the positron excess
- **Lepton flavor physics with cosmic rays: (2) DAMPE**
- Dark matter and lepton flavor universality violation



The Dark Matter Particle Explorer, or **DAMPE**, is a Chinese Academy of Sciences satellite launched on **17 December 2015**



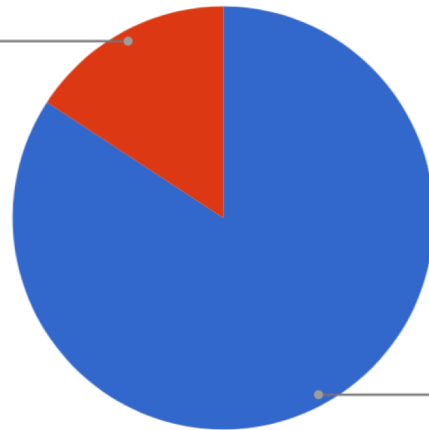
Key result: high-statistics measurement of HE CR electron+positron flux



Key result: high-statistics measurement of HE CR electron+positron flux

...and prompted an interesting **socio-nationalistic** phenomenon

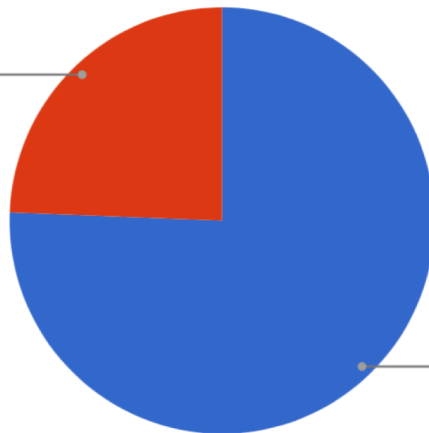
Non-Chinese
15.8%



Chinese Institutions
84.2%

**Nov/2017 cites
(total=19)**

Non-Chinese
24.4%



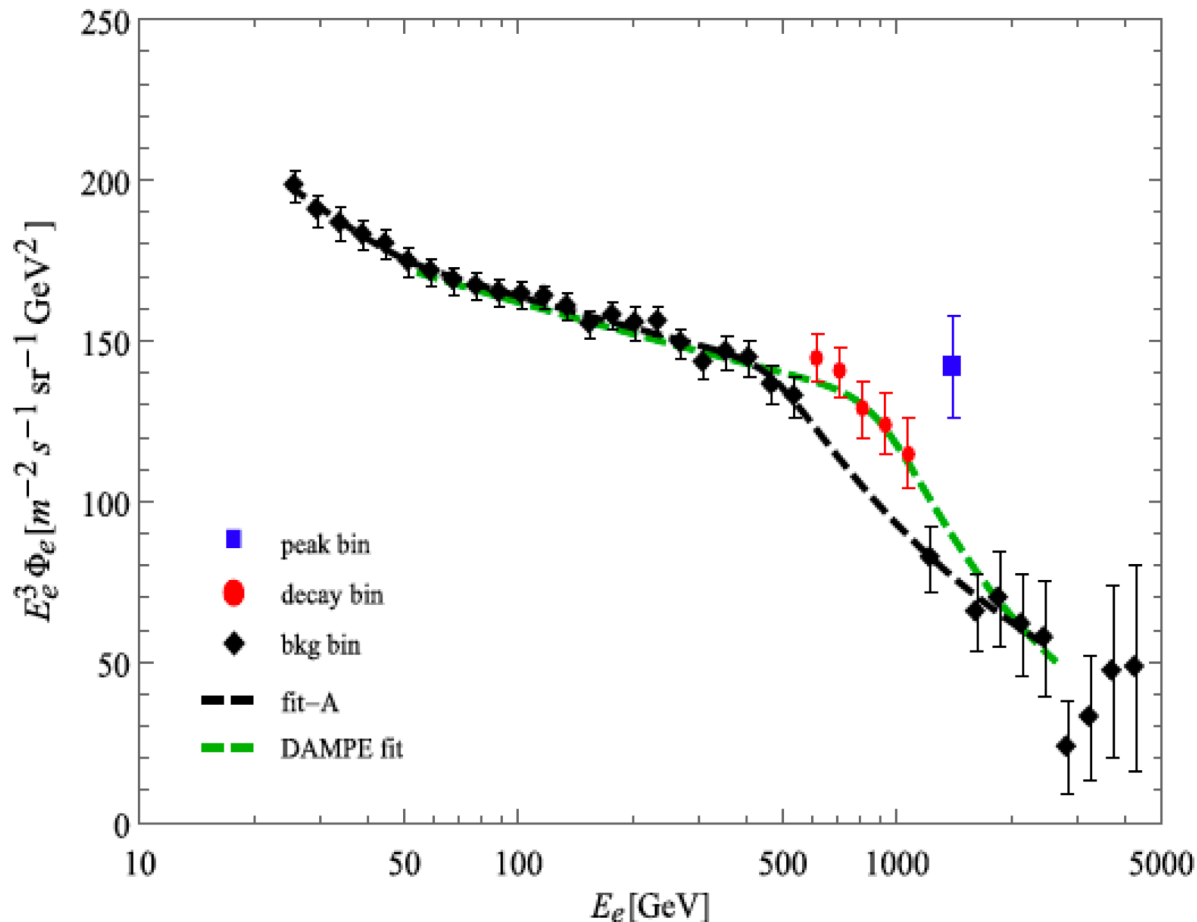
Chinese Institutions
75.6%

**Nov+Dec/2017 cites
(total=41)**

Assume indeed there is a **bump + line** – so what?

Ge et al (2018) “pretend” there’s a line+bump

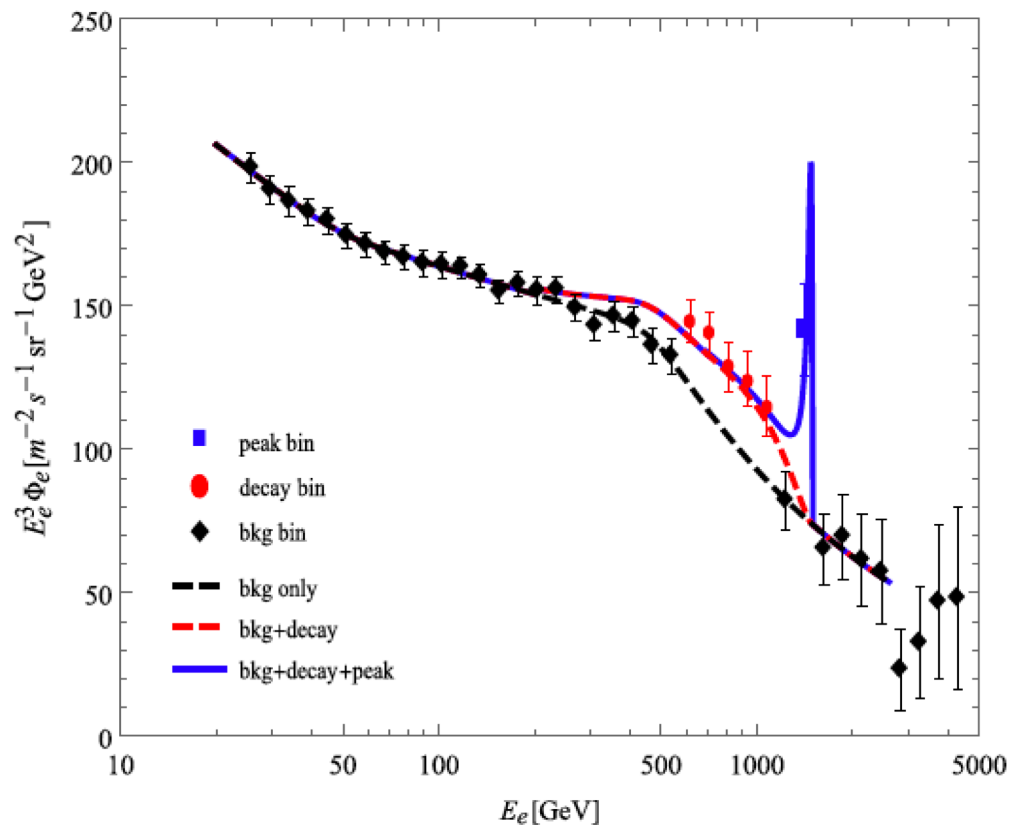
background: double-broken power law



Assume indeed there is a **bump + line** – so what?

Ge et al (2018) “pretend” there’s a line+bump

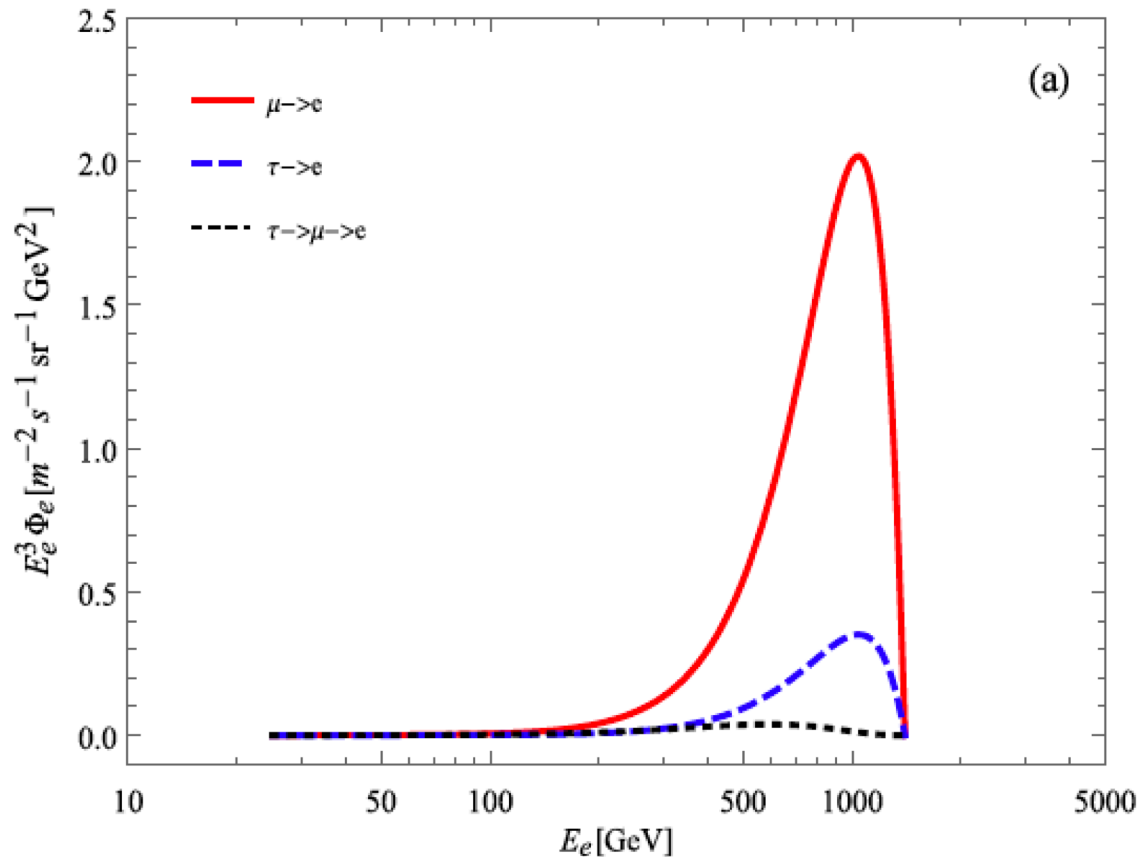
signal: **lepto-philic 1.4 TeV DM** annihilation in a nearby clump, with **some non-flavor universal BR** into **e:μ:τ**



Assume indeed there is a **bump + line** – so what?

τ decay produces much fewer and softer electrons than μ decay – hard to constrain/negligible:

use **$N_e:N_\mu:N_\tau=1:y:0$** , fit for **y**

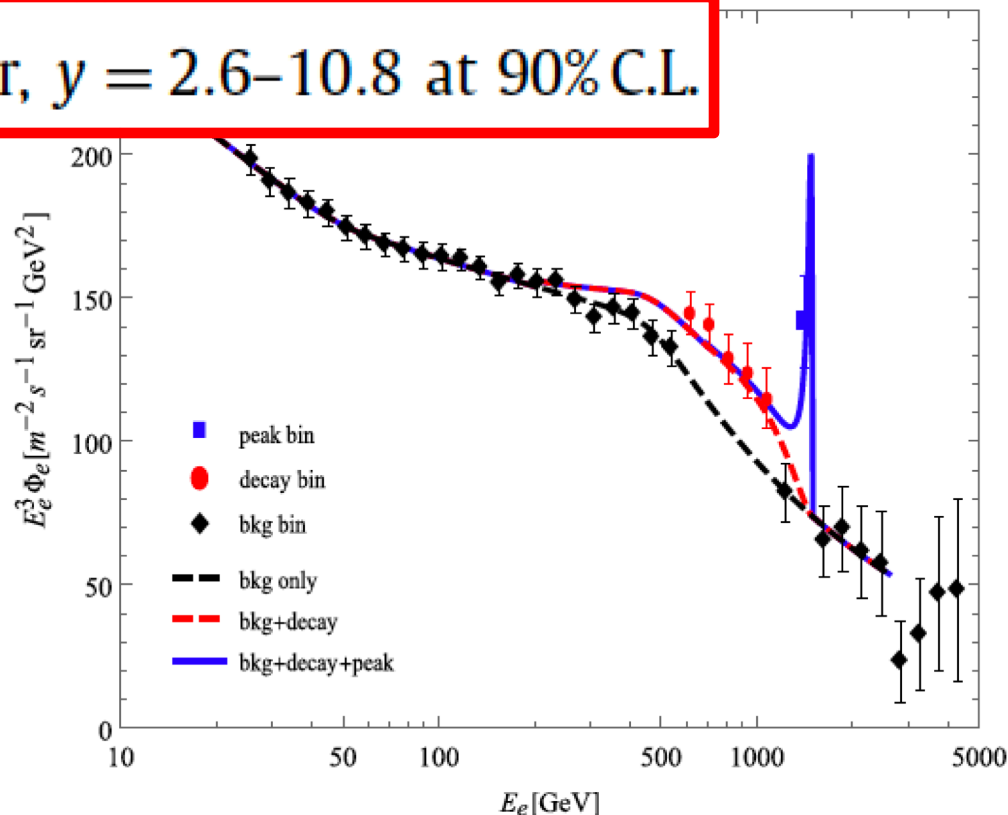


Assume indeed there is a **bump** + **line** – so what?

τ decay produces much fewer and softer electrons than μ decay – hard to constrain/negligible:

use $N_e:N_\mu:N_\tau=1:y:0$, fit for y

$$y = 8.6^{+1.4}_{-2.5}, \text{ or, } y = 2.6-10.8 \text{ at } 90\% \text{ C.L.}$$



Assume indeed there is a **bump + line** – so what?

What do we **learn**?

1. **Flavor structure**: e.g. implemented by Dirac fermion DM with scalar mediator,

$$\mathcal{L}_\chi \supset \lambda_j S_j \overline{\chi}_L \ell_{Rj} + \text{h.c.},$$

2. **Distance** to the **clump** fixed by **width** of the “line”, density distribution fixed by height of line

$d \sim 200$ pc, $r_s \sim 100$ pc (we are partly within clump!)
[therefore the local DM density should be high(er)!]

- Lepton flavor physics with cosmic rays: (1) the positron excess
- Lepton flavor physics with cosmic rays: (2) DAMPE
- **Dark matter and lepton flavor universality violation**

Augment L_μ - L_τ Z' model for B anomalies with a DM particle

Dressing $L_\mu - L_\tau$ in Color

Wolfgang Altmannshofer,^{1,*} Stefania Gori,^{1,†} Maxim Pospelov,^{1,2,‡} and Itay Yavin^{1,3,§}

(another case of a cool title killed by boring journal editors:
“Quark flavor transitions in L_μ - L_τ models”)

$$\mathcal{L}_{Z'} = -\frac{1}{4} (Z')_{\alpha\beta} (Z')^{\alpha\beta} + |D_\alpha \Phi|^2 + V(\Phi) + g' Z'_\alpha J_{Z'}^\alpha . \quad J_{Z'}^\alpha (\text{lep}) = Q_\ell (\bar{\ell}_2 \gamma^\alpha \ell_2 - \bar{\ell}_3 \gamma^\alpha \ell_3 + \bar{\mu}_R \gamma^\alpha \mu_R - \bar{\tau}_R \gamma^\alpha \tau_R)$$

$$\mathcal{L}_m = m_Q \bar{Q}_L \tilde{Q}_R + m_D \bar{\tilde{D}}_L D_R + m_U \bar{\tilde{U}}_L U_R + \text{h.c.}$$

$$\begin{aligned} \mathcal{L}_{\text{mix}} = & \Phi \bar{\tilde{D}}_R (Y_{Qb} b_L + Y_{Qs} s_L + Y_{Qd} d_L) \\ & + \Phi \bar{\tilde{U}}_R (Y_{Qt} t_L + Y_{Qc} c_L + Y_{Qu} u_L) \\ & + \Phi^\dagger \tilde{U}_L (Y_{Ut} t_R + Y_{Uc} c_R + Y_{Uu} u_R) \\ & + \Phi^\dagger \tilde{D}_L (Y_{Db} b_R + Y_{Ds} s_R + Y_{Dd} d_R) + \text{h.c.} \end{aligned}$$

Augment $L_\mu - L_\tau$ Z' model for B anomalies with a DM particle

Dressing $L_\mu - L_\tau$ in Color

Wolfgang Altmannshofer,^{1,*} Stefania Gori,^{1,†} Maxim Pospelov,^{1,2,‡} and Itay Yavin^{1,3,§}

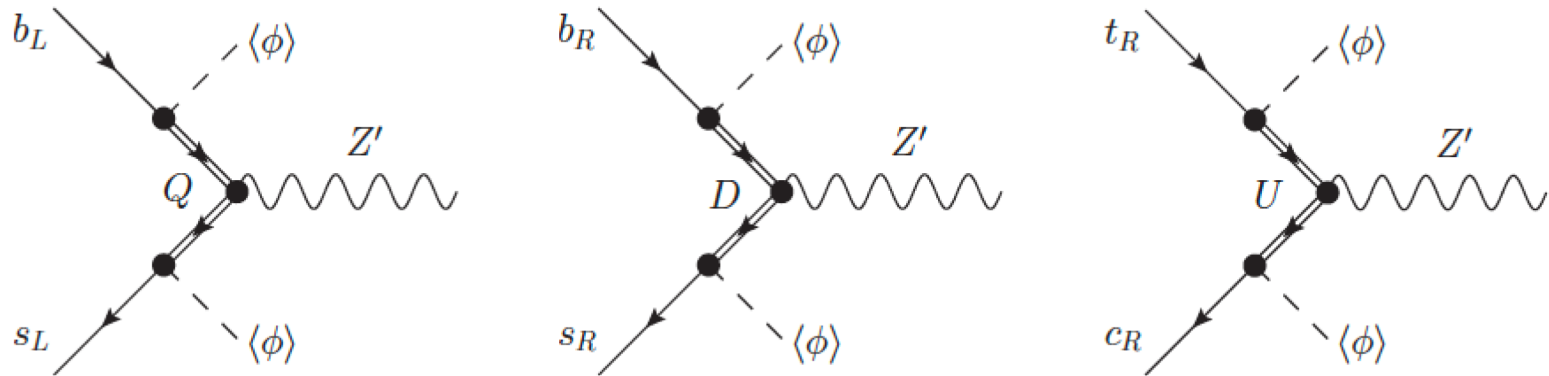


FIG. 1. Example diagrams in the high energy theory that lead to flavor-changing effective couplings of the Z' to SM quarks.

Augment L_μ - L_τ Z' model for B anomalies with a DM particle

Explaining Dark Matter and B Decay Anomalies with an $L_\mu - L_\tau$ Model

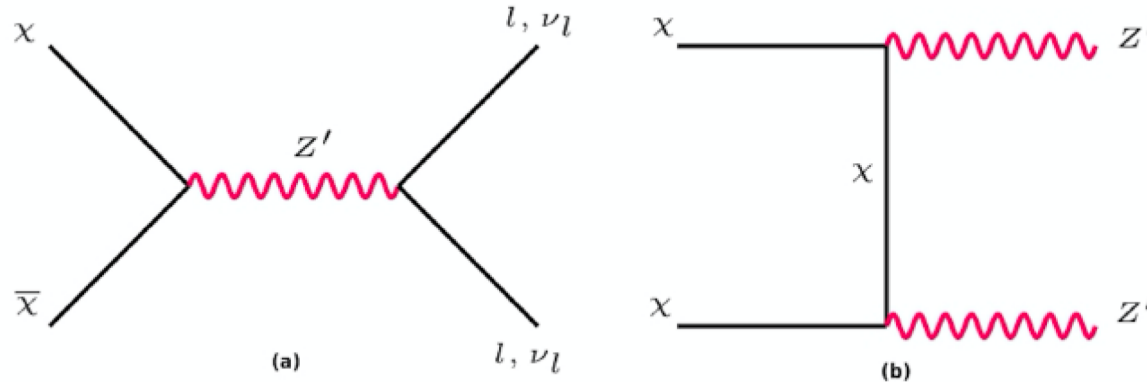
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$$\mathcal{L}_{\text{fermions}} \supset q_l g' (\bar{\mu} \gamma_\alpha \mu - \bar{\tau} \gamma_\alpha \tau + \bar{\nu}_\mu \gamma_\alpha P_L \nu_\mu - \bar{\nu}_\tau \gamma_\alpha P_L \nu_\tau) Z'^\alpha$$

$$\mathcal{L}_{\text{dark}} \supset q_\chi g' \bar{\chi} \gamma_\alpha \chi Z'^\alpha$$

Consider $q_l = q_\chi$ and $q_l \gg q_\chi$ limits; without loss of generality, set $q_l = 1$

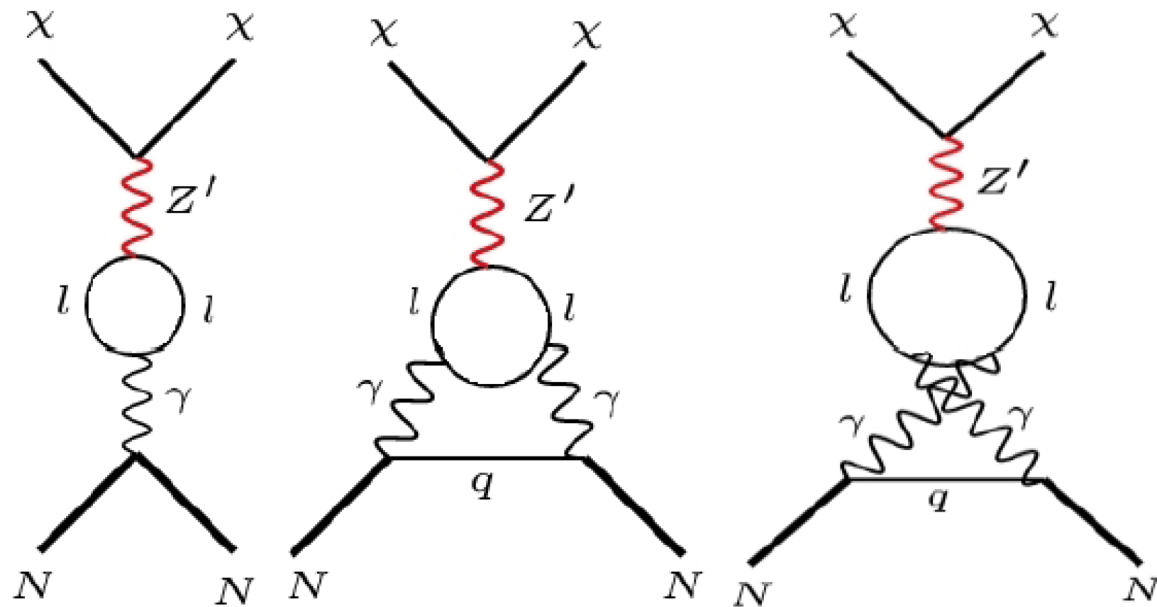
DM Phenomenology: indirect/relic density and direct detection



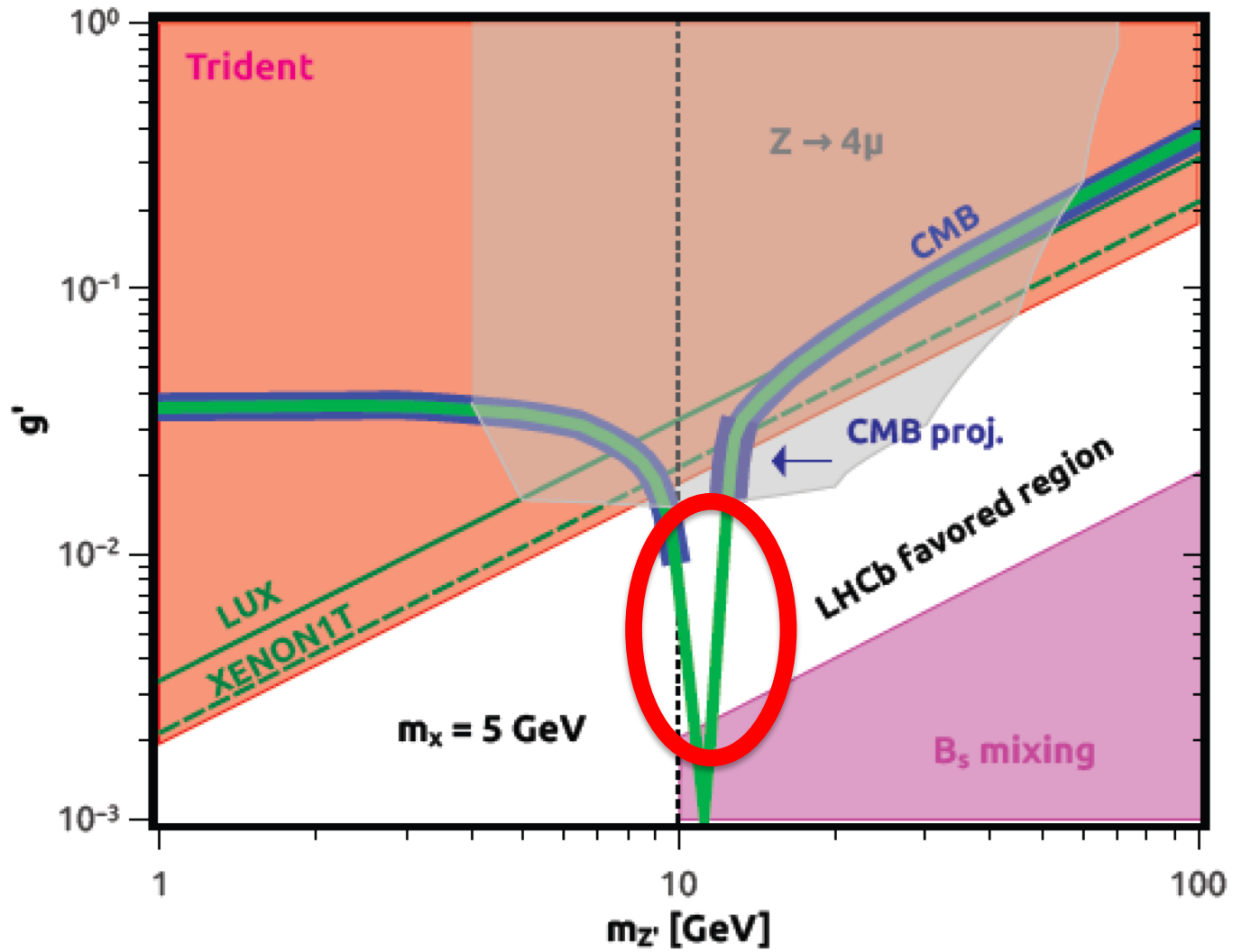
$$\sigma v (\chi \bar{\chi} \rightarrow l^+ l^-) \approx \frac{q_\chi^2 q_l^2 g'^4}{2\pi} \frac{m_\chi^2}{(4m_\chi^2 - m_{Z'}^2)^2}$$

...expect **3 regimes** depending on $m_\chi, m_{Z'}$ hierarchy

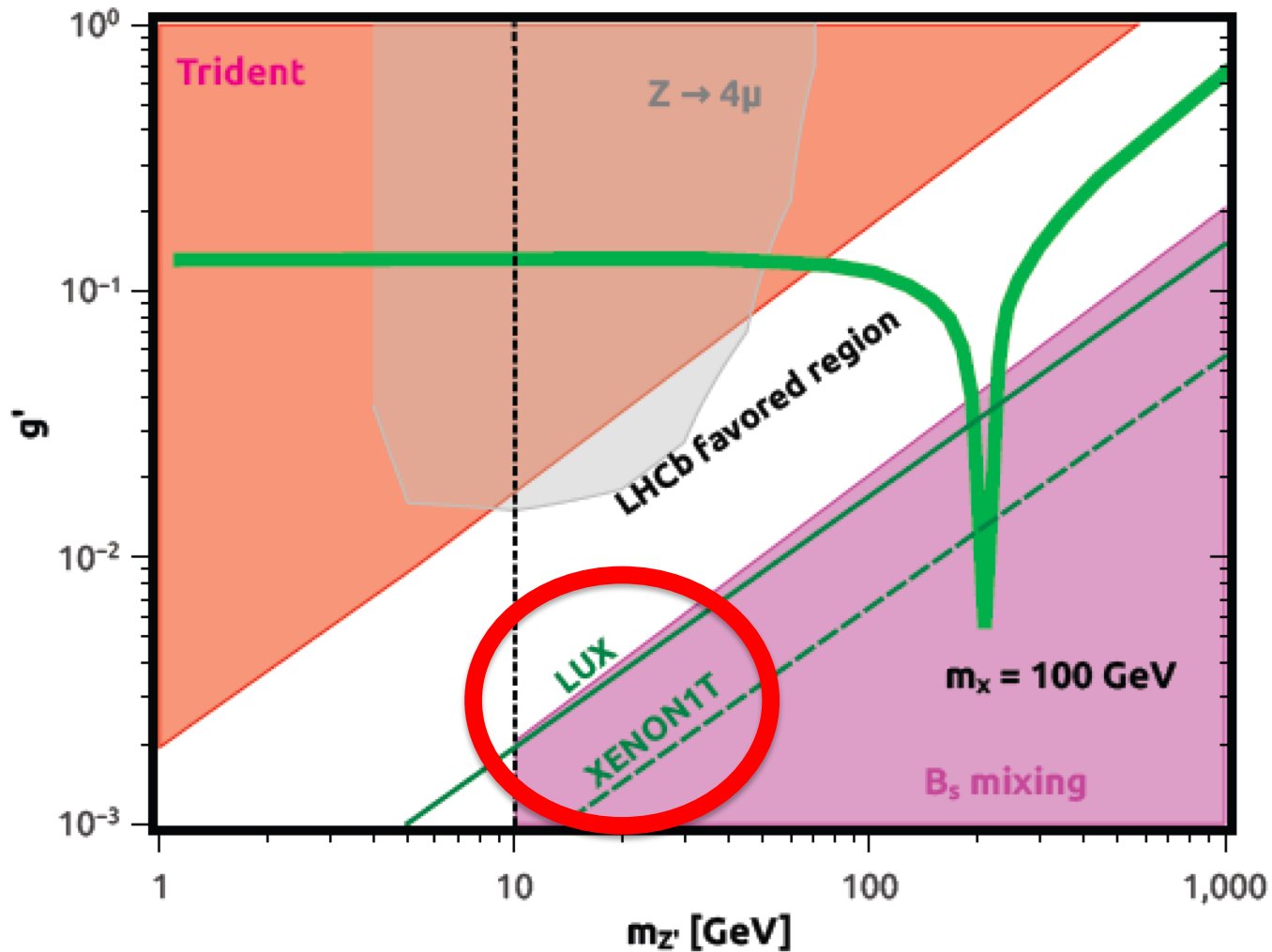
DM Phenomenology: indirect/relic density and direct detection



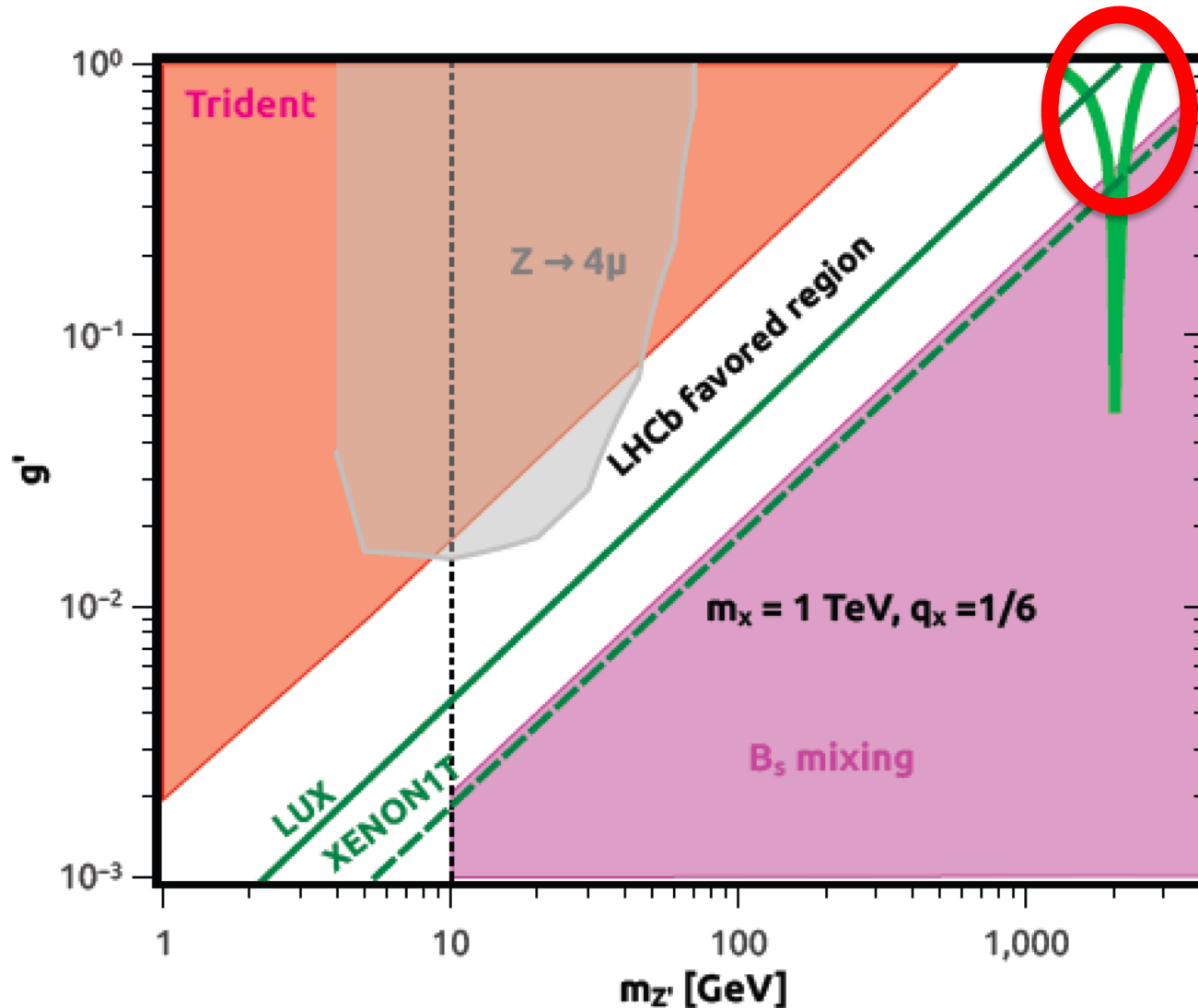
Light DM, large DM charge, $q_\chi = 1$



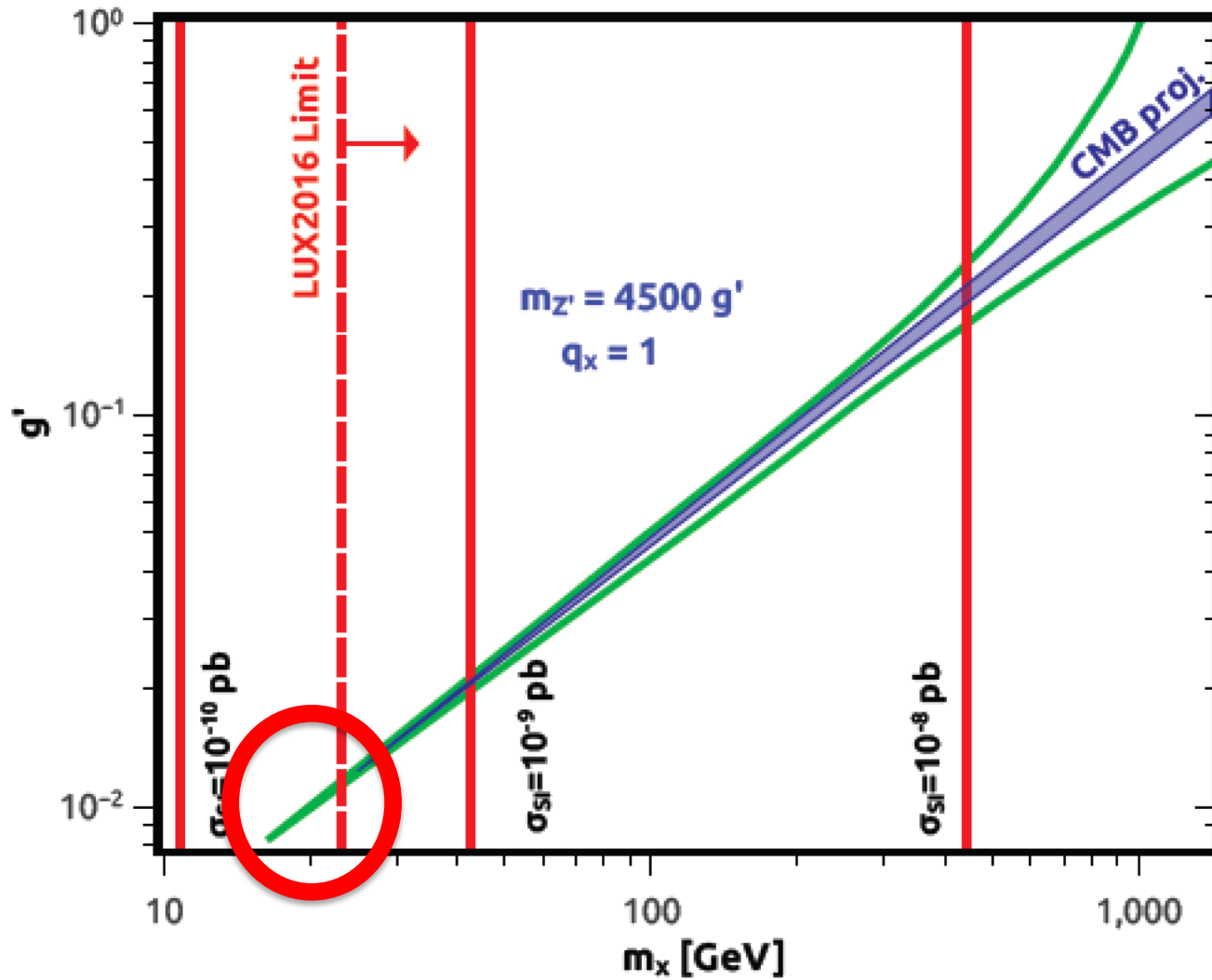
Heavy DM, large DM charge, $q_\chi = 1$: tough luck!



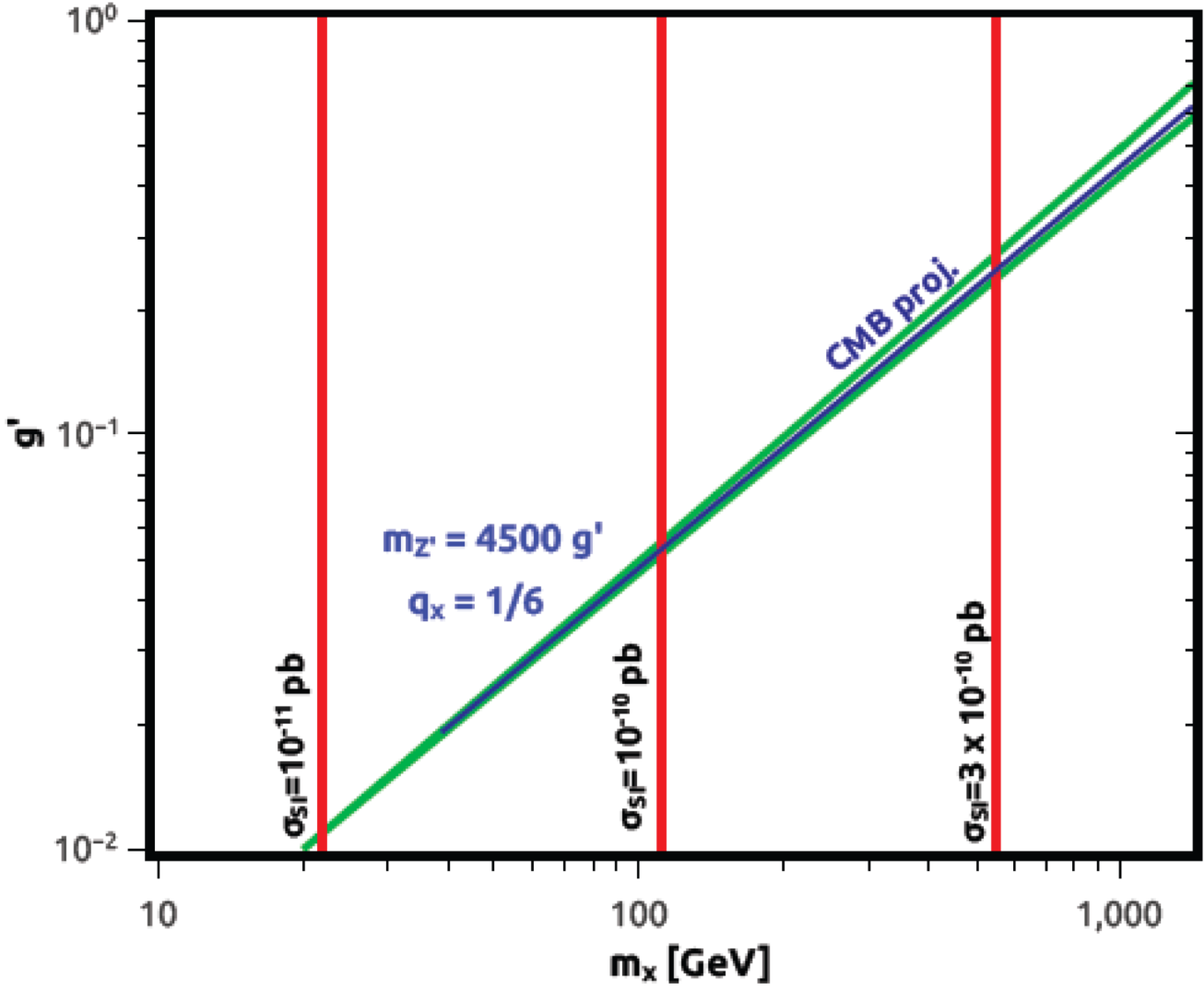
Even **Heavier** DM, **smaller** DM charge, $q_\chi = 1/6$: barely OK (for now)



Orthogonal view, large DM charge $q_\chi = 1$: ouch!



Orthogonal view, small DM charge $q_\chi = 1/6$



- **Lepton flavor physics with cosmic rays: (1) the positron excess**
 - ✓ *nearby **pulsars** are the **likely source** of anomalous CR positrons*
 - ✓ *if so, **diffusion is not homogeneous**; we know how to test it!*

- **Lepton flavor physics with cosmic rays: (2) DAMPE**
 - ✓ *premature, but proof of principle of **flavor physics model building** from **astro data!***

- **Dark matter and lepton flavor universality violation**
 - ✓ *if Z' from L_μ - L_τ is the right explanation to B anomalies, and if Z' is the portal for DM, **very predictive scenario, testable** with direct detection*