Searches for an extended Higgs sector at the LHC (theory)

Stefania Gori

University of Cincinnati

Higgs tasting workshop, 2016 Benasque, May 20th 2016

Questions for the talk

We have discovered the Higgs boson

Is it the full story for ElectroWeak Symmetry Breaking (EWSB)?

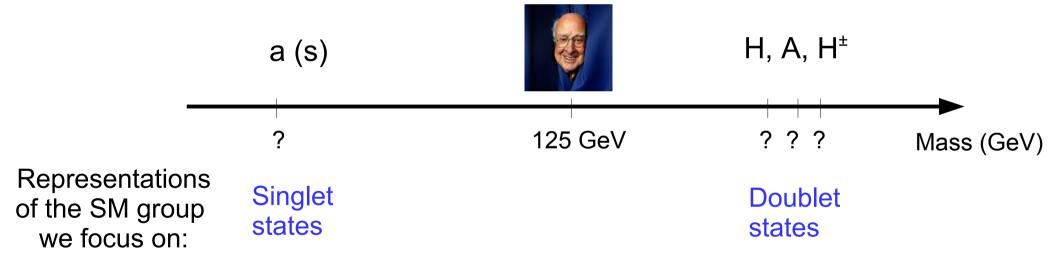
Additional Higgs bosons?(*)

^(*) Disclaimer: I will also discuss new scalars not participating to EWSB

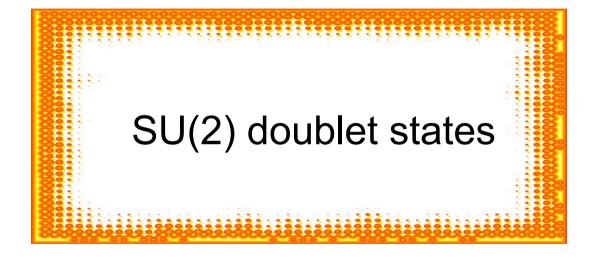
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A SM-like 125 GeV Higgs boson

$$H_1 = \begin{pmatrix} H_1^+ \\ H_1^0 \end{pmatrix}$$

$$H_1 = egin{pmatrix} H_1^+ \ H_1^0 \end{pmatrix} \qquad oldsymbol{V} \supset rac{1}{2} oldsymbol{Z}_1 (oldsymbol{H}_1^\dagger oldsymbol{H}_1)^2 + \cdots + oldsymbol{Z}_6 (oldsymbol{H}_1^\dagger oldsymbol{H}_1) oldsymbol{H}_1^\dagger oldsymbol{H}_2 + oldsymbol{h.c.} + \cdots$$

$$H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix}$$

$$H_2=egin{pmatrix} H_2^+ \ H_2^0 \end{pmatrix}$$
 The two Higgs doublets will generically mix: $\cos^2(oldsymbol{eta}-oldsymbol{lpha})=rac{oldsymbol{Z}_6^2 v^4}{(m_H^2-m_h^2)(m_H^2-oldsymbol{Z}_1 v^2)}$

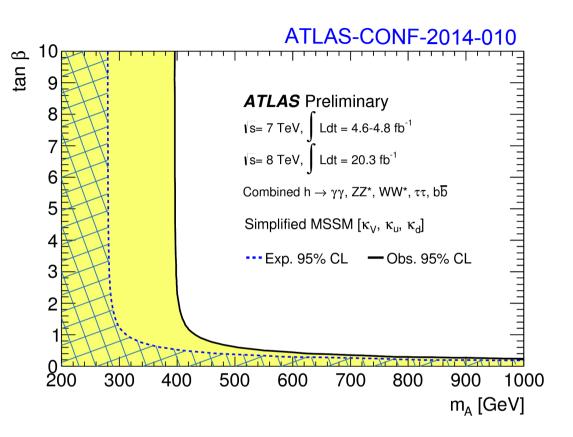


$$rac{g^2}{2}v\left[rac{Z^2}{2\cos^2 heta}+W^2
ight]\xi_V^h h,\;\; \xi_V^h=\sin(eta-lpha)$$

Mixing the two Higgs doublets induces a NP effect in the 125 GeV Higgs couplings

If <u>H is heavy</u> — Higgs with SM properties (decoupling) If Z is small Higgs with SM properties ("alignment")

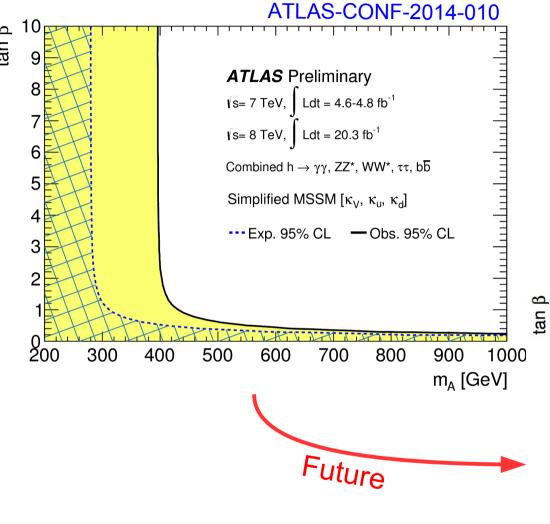
Still weak bounds



Even if we do not have alignment, we do not need much decoupling

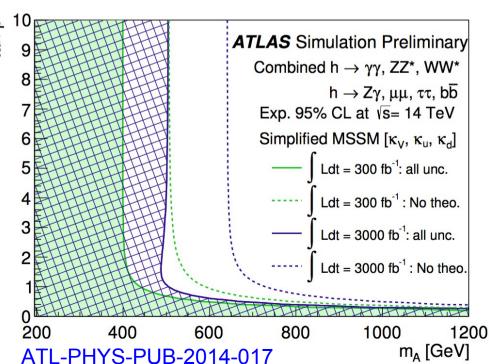
If we have some alignment, bounds are even weaker

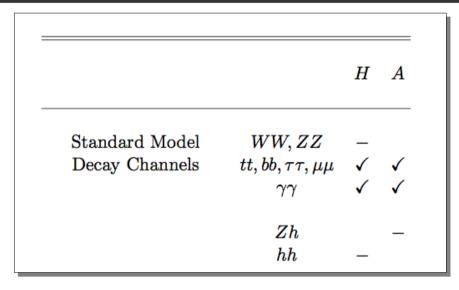
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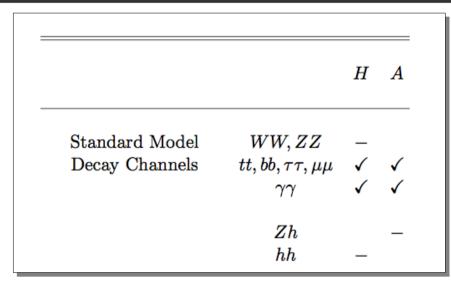
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Warning: these are not the only Interesting modes!

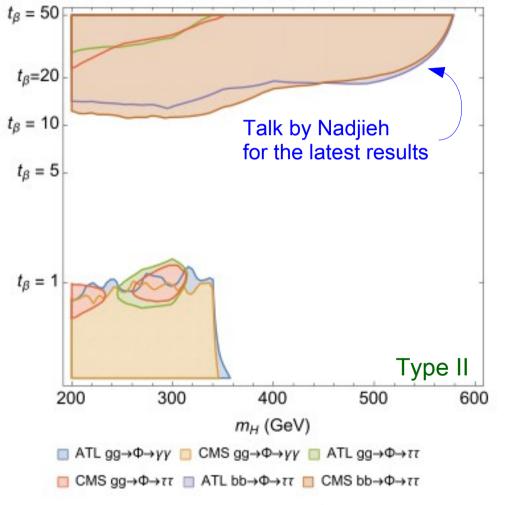


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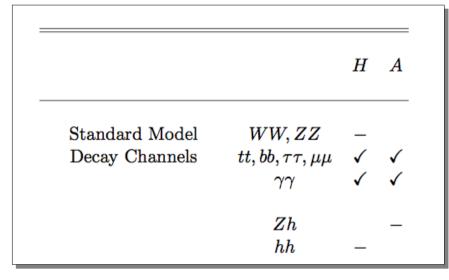
At large tanβ the heavy Higgs couplings to

- bottoms, muons and taus are enhanced
- tops are suppressed

The interpretation of the experimental searches depend on the particular type of 2HDM (e.g. Type I, II, X, Y, ...)

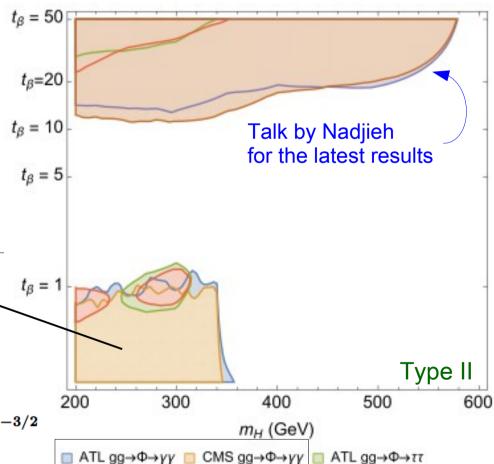


Craig et al. 1504.04630



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CMS gg→Φ→ττ ATL bb→Φ→ττ CMS bb→Φ→ττ

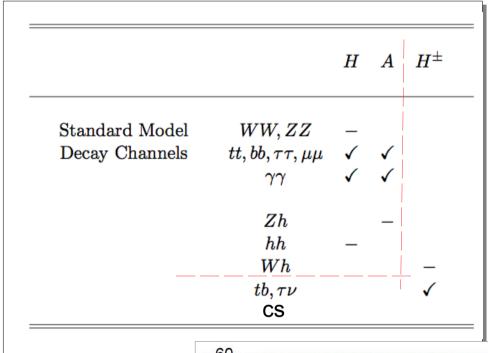
Di-photon?

Rather weak bound from H → yy

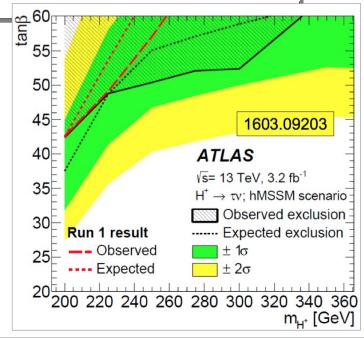
The bound switches off at ~2m, since

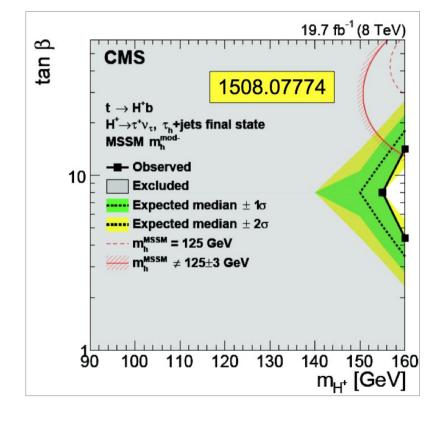
$$egin{array}{ll} {
m BR}(H o\gamma\gamma) &\simeq & rac{32lpha^2}{243\pi^2}rac{m_A^2}{4m_t^2}|A_f|^2\left(1-rac{4m_t^2}{m_H^2}
ight)^{-3/2} \ &\simeq & 10^{-5} \end{array}$$

Craig et al. 1504.04630



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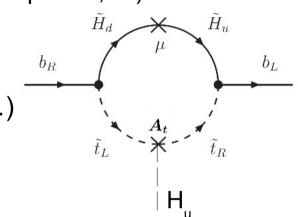


Interpretation of the experimental results

Typically the bounds in the M_A - tan β plane are determined using some specific MSSM scenarios (hMSSM, m_{hmod+} , m_{hmod-} , τ - phobic, ...).

Results depend on the specific scenario since

- Loop corrections to the Hbb (HTT) coupling
- Possible decay to EW sparticles (neutralinos, charginos, ...)

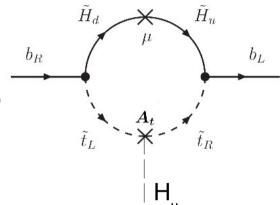


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Bounds dramatically change

if we change the flavor structure of the 2HDM:

$$\mathcal{H}_{Y} = ar{Q}_{L} X_{d1} D_{R} H_{1} + ar{L}_{L} X_{\ell 1} e_{R} H_{1} + ar{Q}_{L} X_{u1} U_{R} H_{1}^{c} + ar{Q}_{L} X_{d2} D_{R} H_{2}^{c} + ar{L}_{L} X_{\ell 2} e_{R} H_{2}^{c} + ar{Q}_{L} X_{u2} U_{R} H_{2}$$

"wrong Yukawas" (loop suppressed in the MSSM)

Generically,

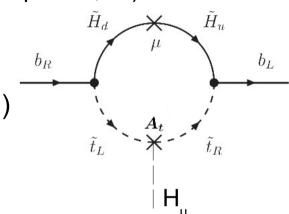
- we can have flavor changing heavy Higgs decays, weakening the previous bounds. eg. H → tc
- the diagonal heavy Higgs couplings can be very different than in a Type II 2HDM

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New interesting models:

- Flavor aligned models

- Altmannshofer, SG, Kagan, Silvestrini, Zupan, 1507.07927 See also Ghosh et al., 1508.01501
- Models with only third generation masses coming from the 125GeV Higgs boson

A flavor aligned 2HDM

$$\mathcal{H}_Y = \bar{Q}_L X_{d1} D_R H_1 + \bar{L}_L X_{\ell 1} e_R H_1 + \bar{Q}_L X_{u1} U_R H_1^c + \bar{Q}_L X_{d2} D_R H_2^c + \bar{L}_L X_{\ell 2} e_R H_2^c + \bar{Q}_L X_{u2} U_R H_2$$
Pich, Tuzon, 0908.1554

$$X_{u1} = \epsilon_t Y_u \\ X_{u2} = Y_u$$

$$X_{d2} = \epsilon_b Y_d$$

$$X_{d1} = Y_d$$

$$X_{\ell 2} = \epsilon_{\ell} Y_{\ell}$$

$$X_{\ell 1} = Y_{\ell}$$

Special case of a Minimal Flavor Violating 2HDM

	Type II	Aligned
Up	$rac{\sinlpha}{\sineta}$	$rac{\sinlpha}{\sineta}$
Down	$\frac{\cos \alpha}{\cos \beta}$	$\frac{\cos\alpha + \epsilon_b \sin\alpha}{\cos\beta + \epsilon_b \sin\beta}$
Lepton	$\frac{\cos \alpha}{\cos \beta}$	$\frac{\cos \alpha + \epsilon_{\ell} \sin \alpha}{\cos \beta + \epsilon_{\ell} \sin \beta}$

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$$\mathcal{H}_Y = ar{Q}_L X_{d1} D_R H_1 + ar{L}_L X_{\ell 1} e_R H_1 + ar{Q}_L X_{u1} U_R H_1^c + ar{Q}_L X_{d2} D_R H_2^c + ar{L}_L X_{\ell 2} e_R H_2^c + ar{Q}_L X_{u2} U_R H_2$$
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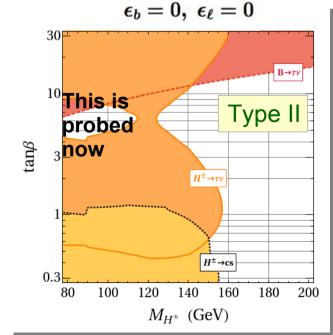
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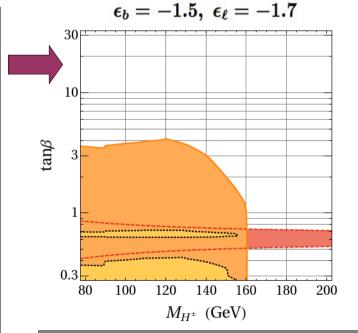
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Depending on the values of $\varepsilon_{_{\rm b}}$, $\varepsilon_{_{_{\rm T}}}$ bounds can be much weaker/stronger

Example for the charged H constraints:



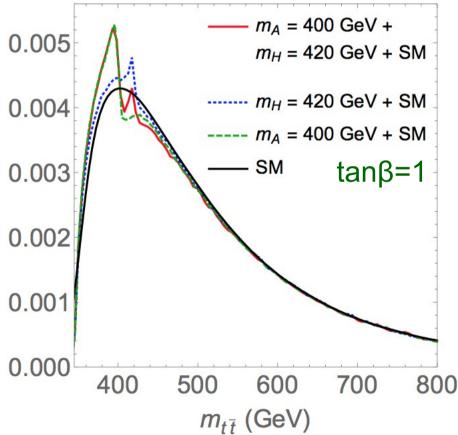


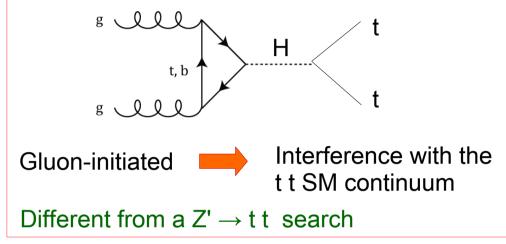
Altmannshofer, SG, Kribs, 1210.2465

The (heavy) Higgs mainly decays into tops at low/intermediate values of tanβ

1. What about t t resonance searches?

Parton level analysis





Dicus et al., 9404359

Shape analysis?
Angular correlations?

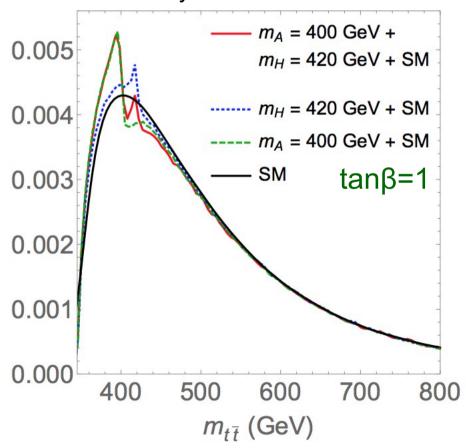
SG, IW.Kim, N.Shah, K.Zurek, 1602.02782

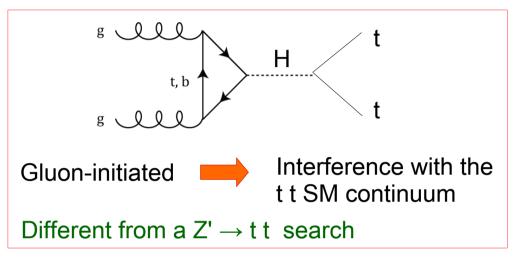
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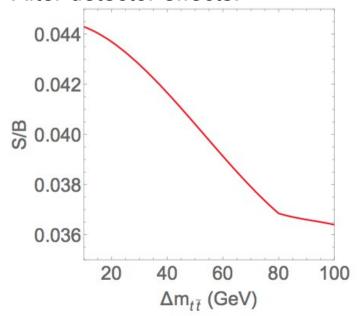
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After detector effects:

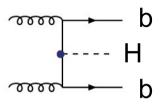


A challenging but very important search

SG, IW.Kim, N.Shah, K.Zurek, 1602.02782

The (heavy) Higgs mainly decays into tops at low/intermediate values of tanβ

2. b associated production

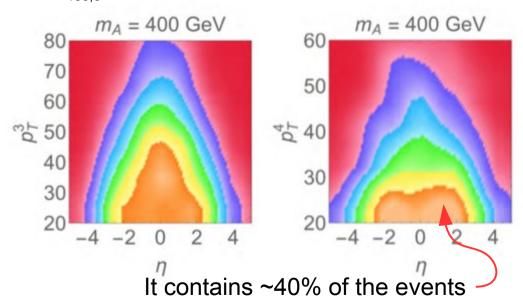


What changes?

- We do not have to worry about the interference
- the additional b quarks can be used to discriminate against the SM t t background

The challenge:

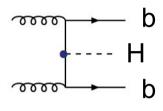
Optimize the use of this b-quarks $(S_{400.6} \sim 500 \text{ fb}; B \sim O(1 \text{nb}))$



SG, IW.Kim, N.Shah, K.Zurek, 1602.02782

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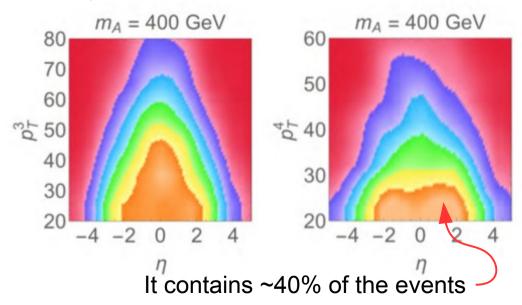


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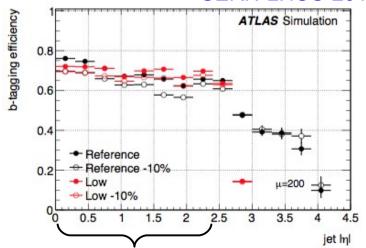
Opportunities:

- MVA or BDT analysis

See for example Hajer et al, 1504.07617

 more forward b-tagging for the HL-LHC

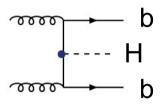
CERN-LHCC-2015-020



So far, b-tagging up to |η|

The (heavy) Higgs mainly decays into tops at low/intermediate values of tanβ

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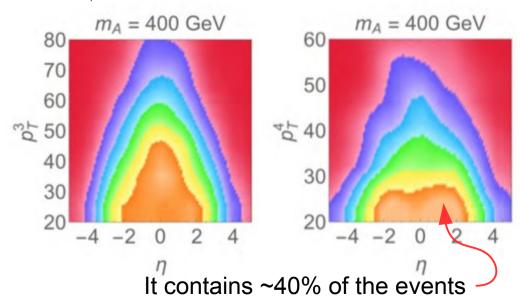


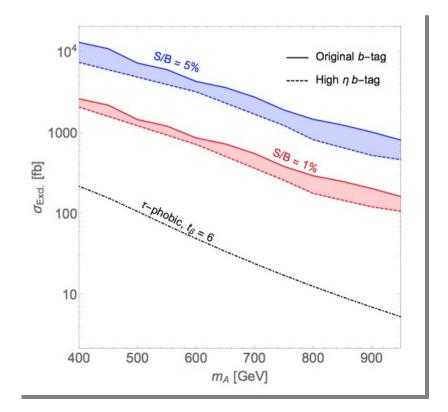
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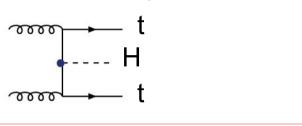




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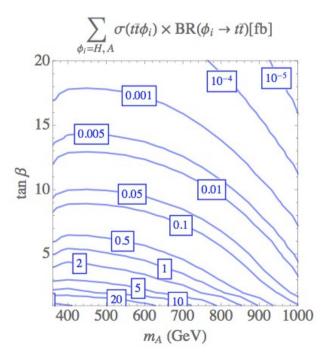


The opportunity:

We have many more objects in the signal, if compared to the t t background

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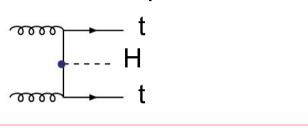
Signal rates are much smaller



SG, IW.Kim, N.Shah, K.Zurek, 1602.02782

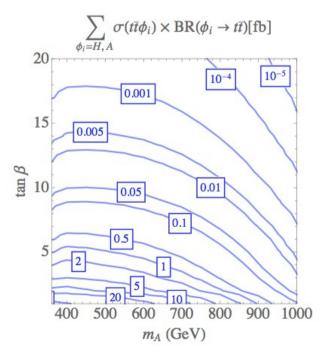
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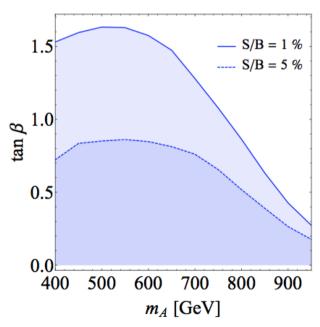
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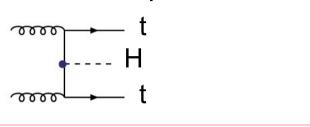
- many (b) jets + at least 1 lepton
- multi-leptons + (b) jets



1 lepton, >1b, >5jets

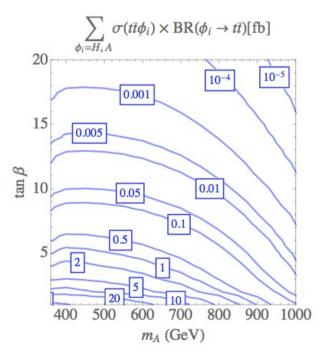
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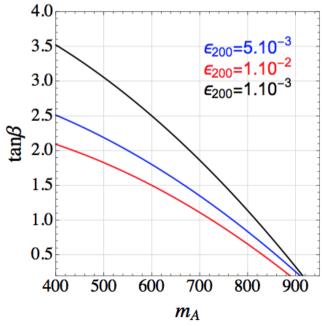
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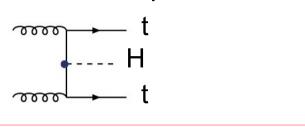
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3 leptons, >1b, >2jets

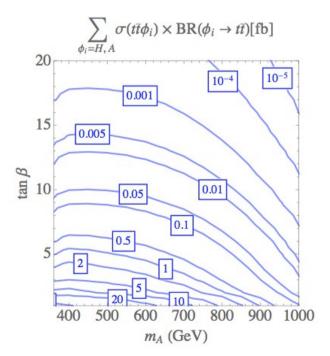
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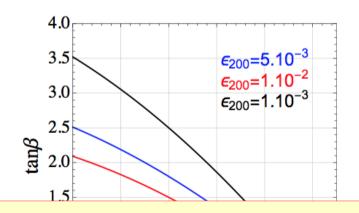
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- multi-leptons + (b) jets



Open questions:

- what about a same sign lepton analysis?
- what about t(q) H associated production?

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Theories with light scalars

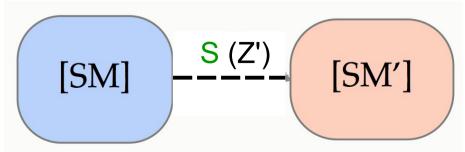
SUSY and beyond:

- NMSSM
 - (approximate) R-symmetric limit: light pseudoscalar (a).
- (approximate) PQ-symmetric limit: light pseudoscalar (a), scalar (s), singlino (χ₁).
- Folded SUSY, Twin Higgs models, fraternal twin Higgs model:
- light glueballs

Burdman et al, 0609152, Chacko et al, 0506256, Craig et al, 1501.05310

More generically, in **Simplified models**:

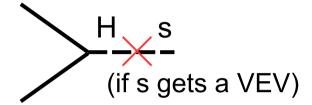
- 2HDM + scalar
- Theories for light Dark Matter with scalar mediator



How to probe these particles?

Since they are not charged under the SM symmetries...

Direct production:



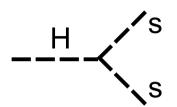
(same production mechanisms as for the SM Higgs)

Higgs portal interaction

$$rac{oldsymbol{\xi}}{2}|H|^2s^2$$

Interplay of LHC searches

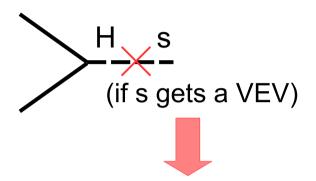
(125 GeV) Higgs decays:



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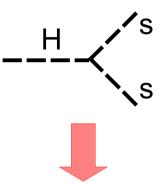
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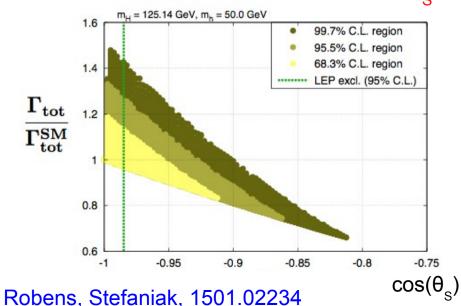


Higgs portal interaction $\frac{\xi}{2}|H|^2s^2$

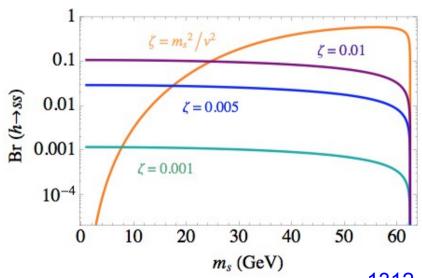
(125 GeV) Higgs decays:



The Higgs couplings to SM particles are reduced by a factor of $cos(\theta_s)$



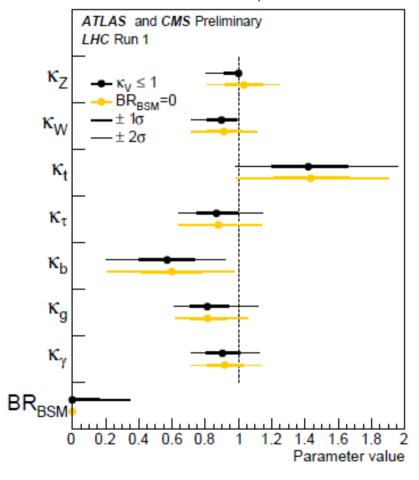
The Higgs has some "exotic" signatures



1312.4992

The exotic width of the 125 GeV Higgs

ATLAS-CONF-2015-044, CMS-PAS-HIG-15-002

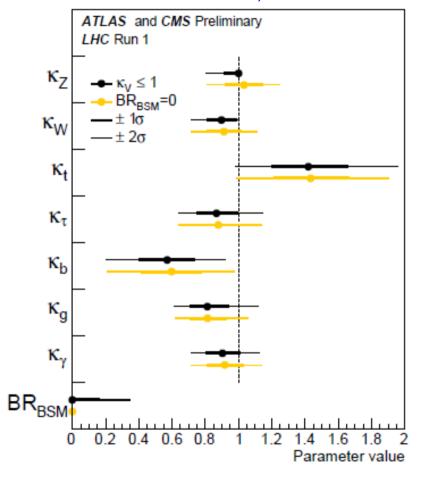


 $BR_{BSM} < 34\%$

Prospects for the HL-LHC: ~10%

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New Higgs signatures: $h \rightarrow ss \rightarrow ?$

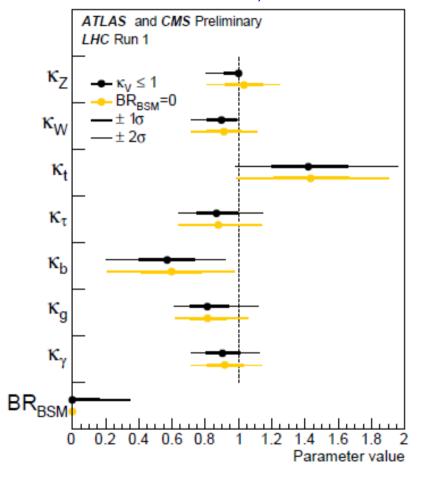
■ If $\theta_s = 0$, s is stable

Talk by Nadjieh for the latest results

Higgs invisible decay 4

The exotic width of the 125 GeV Higgs

ATLAS-CONF-2015-044, CMS-PAS-HIG-15-002



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New Higgs signatures: $h \rightarrow ss \rightarrow ?$

■ If θ_s =0, s is stable

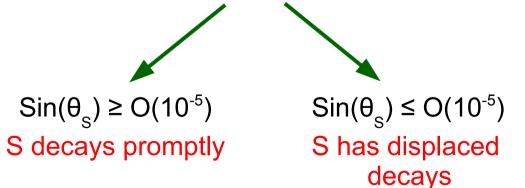
Talk by Nadjieh for the latest results

Higgs invisible decay

■ If $\theta_s \neq 0$, s will decay to SM particles

$$\Gamma(s \to f\bar{f}) = \sin^2\theta_S \frac{N_c}{8\pi} \frac{m_s m_f^2}{v^2} \beta_f^3$$

Main BRs: bb, тт, cc, ...



Prompt Decays

From Curtin, Essig, SG, Jaiswal, Katz, Liu, Liu, McKeen, Shelton, Strassler, Surujon, Tweedie, Zhong, 1312.4992

$$\mathrm{BR}(\mathcal{F}_i) \equiv \mathrm{BR}(h \to ss \to (f\bar{f})(f'\bar{f}'))$$
 BR(s)

	Projected/Current		quarks allowed	
Decay	2σ Limit	Produc-		Limit on
Mode	on $\mathrm{BR}(\mathcal{F}_{\mathrm{i}})$	tion	$\frac{\mathrm{BR}(\mathcal{F}_{\mathrm{i}})}{\mathrm{BR}(\mathrm{non-SM})}$	$\frac{\sigma}{\sigma_{\rm SM}} \cdot {\rm BR}({\rm non\text{-}SM})$
\mathcal{F}_i	7/8 [14] TeV	Mode		7/8 [14] TeV
$b ar{b} b ar{b}$	$0.7^R \ [0.2^L]$	W	0.8	0.9 [0.2]
$bar{b} au au$	$> 1 \ [0.15^L]$	V	0.1	> 1 [1]
$bar{b}\mu\mu$	$(2-7)\cdot 10^{-4}$	G	3×10^{-4}	0.6 - 1
	$[(0.6-2)\cdot 10^{-4}]$			[0.2 - 0.7]
$\tau\tau\tau\tau$	$0.2 - 0.4^R$ [U]	G	0.005	40 - 80 [U]
$ au au\mu\mu$	$(3-7) \cdot 10^{-4} ^{T} [U]$	G	3×10^{-5}	$10 - 20 [\mathrm{U}]$
$\mu\mu\mu\mu$	$1 \cdot 10^{-4} R [\text{U}]$	G	$1 \cdot 10^{-7}$	1000 [U]

Extracted bound on BR($h \rightarrow ss$)

Prompt Decays

From Curtin, Essig, SG, Jaiswal, Katz, Liu, Liu, McKeen, Shelton, Strassler, Surujon, Tweedie, Zhong, 1312.4992

$$\mathrm{BR}(\mathcal{F}_i) \equiv \mathrm{BR}(h \to ss \to (f\bar{f})(f'\bar{f}'))$$
 BR(s)

				1	
		Projected/Current		quarks allowed	
{	Decay	2σ Limit	Produc-		Limit on
	\mathbf{Mode}	on $\mathrm{BR}(\mathcal{F}_{\mathrm{i}})$	tion	$\frac{\mathrm{BR}(\mathcal{F}_{\mathbf{i}})}{\mathrm{BR}(\mathrm{non-SM})}$	$\frac{\sigma}{\sigma_{\rm SM}} \cdot {\rm BR(non\text{-}SM)}$
	\mathcal{F}_i	7/8 [14] TeV	Mode	,	7/8 [14] TeV
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	$bar{b}\mu\mu$	$(2-7) \cdot 10^{-4}$	G	3×10^{-4}	0.6 - 1
		- 3 · 10 ⁻⁴			[0.2 - 0.7]
	/ττττ	0.2 0.25	G	0.005	40 - 80 [U]
	$ au au\mu\mu$	$/(3-7)\cdot 10^{-4}$ 4 : 10	$\overline{\mathcal{G}}$	3×10^{-5}	10 - 20 [U]
	$\mu\mu\mu\mu\mu$	$1 \cdot 10^{-4}$ 4 · 10 ⁻⁵	G	$1 \cdot 10^{-7}$	1000 [U]
·14-0	41		C) HIG-15-	-011	Extracted bound

C) HIG-1

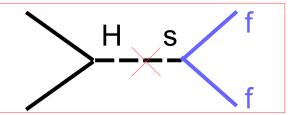
C) 1510.06534, C) HIG-13-010, C) HIG-14-022

A) 1505.01609 A) 1505.07645

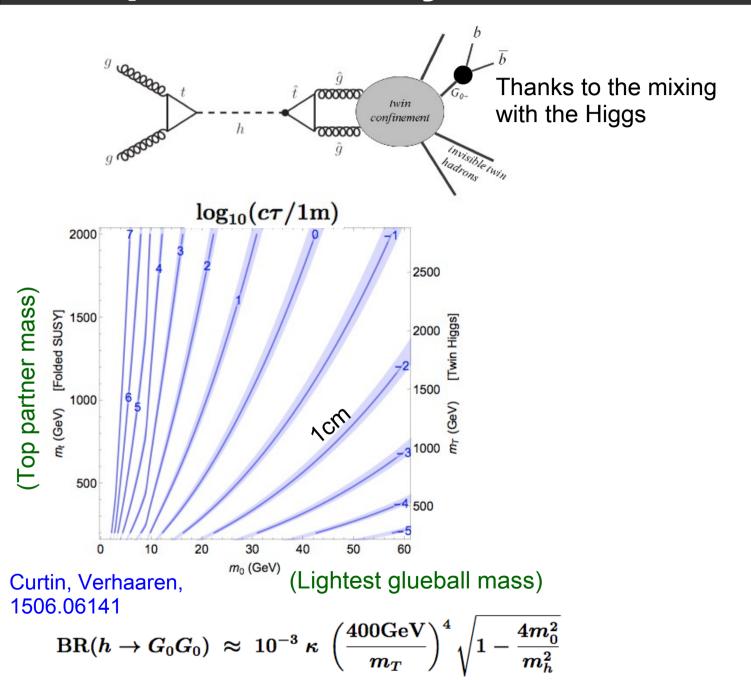
on BR(h \rightarrow ss)

Wish list:

it would be great to have a similar table for the direct production of s

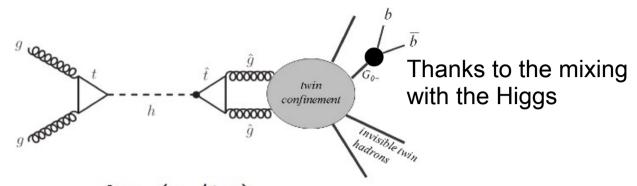


Displaced Decays



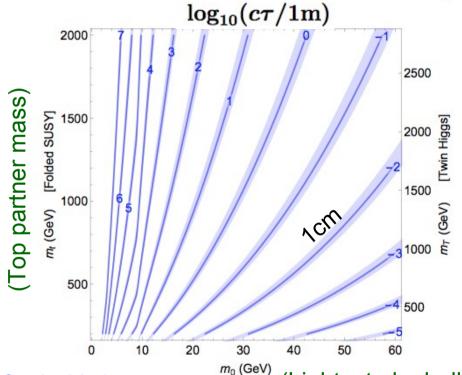
Specific example: Glueballs in Folded SUSY or Twin Higgs models

Displaced Decays



(Lightest glueball mass)

Specific example: Glueballs in Folded SUSY or Twin Higgs models



Existing displaced searches methods focus on

- 2 displaced objects or 1 displaced object
- & high-threshold associated object
- displacement longer than ~ 1mm 1cm

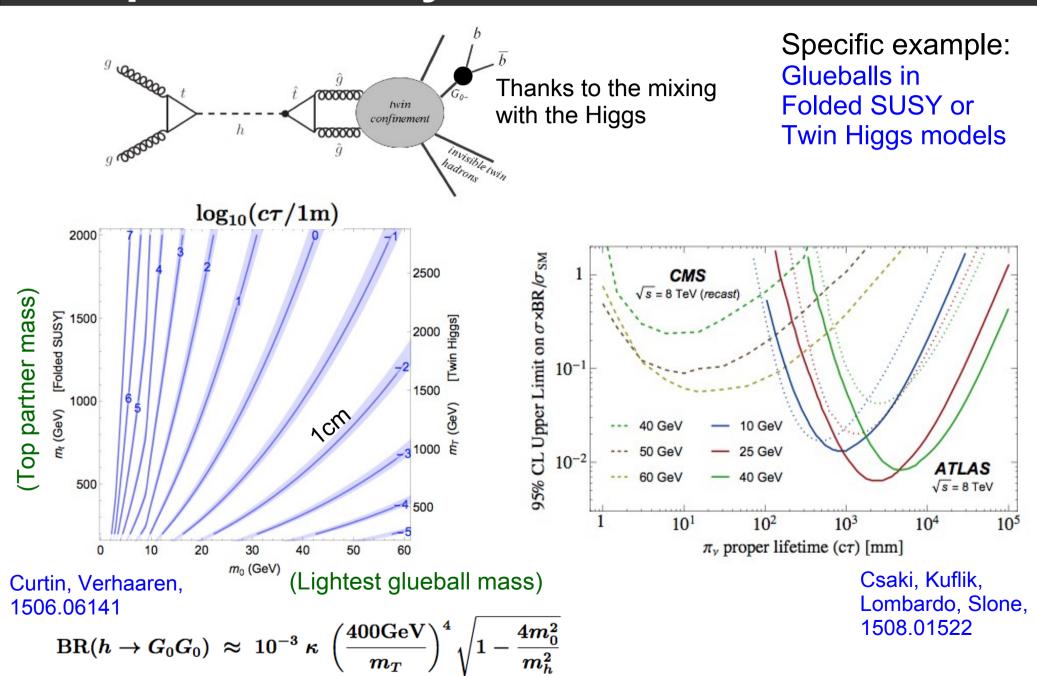
More focused Higgs exotic displaced searches?

- only 1 displaced object + thresholds suitable for the Higgs
- very short displacements
- new dedicated displaced triggers?

Curtin, Verhaaren, 1506.06141

$${
m BR}(h o G_0G_0) \;pprox \; 10^{-3} \; \kappa \; \left(rac{400{
m GeV}}{m_T}
ight)^4 \sqrt{1-rac{4m_0^2}{m_h^2}}$$

Displaced Decays



Conclusions/Take home messages

After the discovery of the 125 GeV Higgs boson, the search for additional Higgs bosons/scalars is even more motivated

Very different search strategies for SU(2) doublet or singlet states.

Many searches are on-going

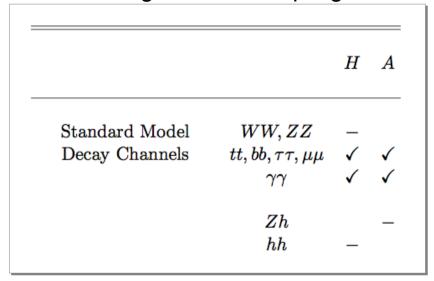
Many open questions:

- How to probe the low tanβ regime?
- How to design additional searches for light singlet scalars? (direct production vs. 125 GeV Higgs decays (4b, 2b2τ final states)
- Higgs decays to displaced scalars. How to test this scenario?

Complementarity with 125GeV Higgs coupling measurements

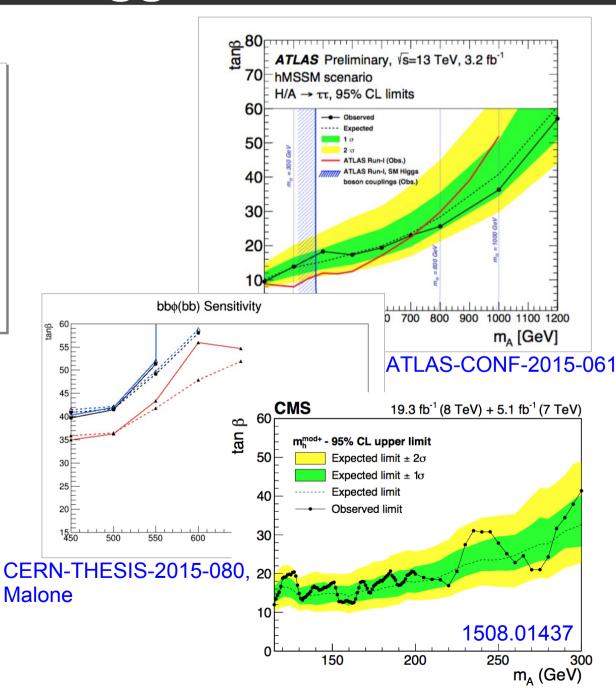
Searches for new Higgs bosons to date

In the alignment/decoupling limit:



At large tanβ the heavy Higgs couplings to

- bottoms, muons and taus are enhanced
- tops are suppressed



Interpretation of the limits

Carena et al, 1302.7033

Typically the cross section limits are interpreted in the

```
m_h^{	ext{mod}+} \ m_h^{	ext{mod}-} \ m_h^{	ext{max}} \ 	au - 	ext{phobic}
```

$$M_1=rac{5}{3}rac{s_{
m w}^2}{c_{
m w}^2}M_2$$

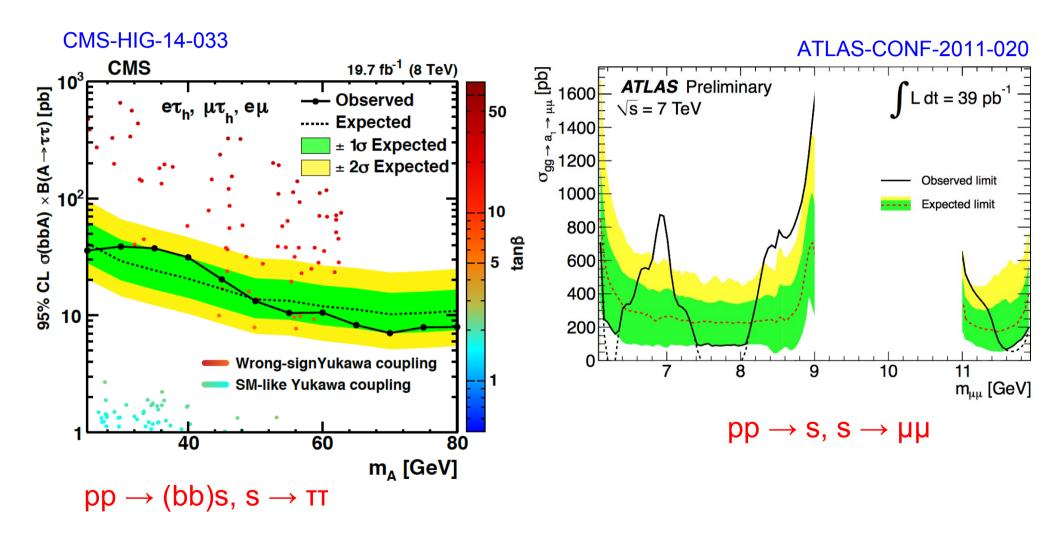
```
egin{aligned} m_t &= 173.2 \; 	ext{GeV}, \ M_{	ext{SUSY}} &= 1000 \; 	ext{GeV}, \ \mu &= 200 \; 	ext{GeV}, \ M_2 &= 200 \; 	ext{GeV}, \ X_t^{	ext{OS}} &= 1.5 \, M_{	ext{SUSY}} \; 	ext{(FD calculation)}, \ X_t^{\overline{	ext{MS}}} &= 1.6 \, M_{	ext{SUSY}} \; 	ext{(RG calculation)}, \ A_b &= A_{	au} &= A_t, \ m_{	ilde{g}} &= 1500 \; 	ext{GeV}, \ M_{	ilde{l}_3} &= 1000 \; 	ext{GeV} \; . \end{aligned}
```

```
egin{aligned} m_t &= 173.2 \; 	ext{GeV}, \ M_{	ext{SUSY}} &= 1000 \; 	ext{GeV}, \ \mu &= 200 \; 	ext{GeV}, \ M_2 &= 200 \; 	ext{GeV}, \ X_t^{	ext{OS}} &= -1.9 \, M_{	ext{SUSY}} \; 	ext{(FD calculation)}, \ X_t^{\overline{	ext{MS}}} &= -2.2 \, M_{	ext{SUSY}} \; 	ext{(RG calculation)}, \ A_b &= A_{	au} = A_t, \ m_{	ilde{g}} &= 1500 \; 	ext{GeV}, \ M_{	ilde{l}_3} &= 1000 \; 	ext{GeV} \; . \end{aligned}
```

```
egin{aligned} m_t &= 173.2 \; 	ext{GeV}, \ M_{	ext{SUSY}} &= 1000 \; 	ext{GeV}, \ \mu &= 200 \; 	ext{GeV}, \ M_2 &= 200 \; 	ext{GeV}, \ X_t^{	ext{OS}} &= 2 \, M_{	ext{SUSY}} \; 	ext{(FD calculation)}, \ X_t^{	ext{MS}} &= \sqrt{6} \, M_{	ext{SUSY}} \; 	ext{(RG calculation)}, \ A_b &= A_{	au} &= A_t, \ m_{	ilde{g}} &= 1500 \; 	ext{GeV}, \ M_{	ilde{l}_3} &= 1000 \; 	ext{GeV} \; . \end{aligned}
```

```
egin{aligned} m_t &= 173.2 \; 	ext{GeV}, \ M_{	ext{SUSY}} &= 1500 \; 	ext{GeV}, \ \mu &= 2000 \; 	ext{GeV}, \ M_2 &= 200 \; 	ext{GeV}, \ X_t^{	ext{OS}} &= 2.45 \, M_{	ext{SUSY}} \; 	ext{(FD calculation)}, \ X_t^{\overline{	ext{MS}}} &= 2.9 \, M_{	ext{SUSY}} \; 	ext{(RG calculation)}, \ A_b &= A_{	au} &= A_t \; , \ m_{	ilde{g}} &= 1500 \; 	ext{GeV}, \ \hline M_{	ilde{l}_3} &= 500 \; 	ext{GeV} \; . \end{aligned}
```

Direct searches for singlet states



Backup S.Gori