



Hadronization models in GENIE

<http://www.genie-mc.org>

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Hadronization modelling

Models of exclusive single-meson production (single- π , single-K)

- Hadronic distributions calculable

But, no model of exclusive hadronic multiparticle production.

- Make predictions by stitching together modelling elements
- **Hadronization modelling** an important ingredient

Hadronization modelling

Functionally, it is a process that receives a minimal set of inputs:

- The neutrino ID ($\nu/\bar{\nu}$)
- The hit nucleon and hit quark IDs
- The interaction type (CC/NC)
- The hadronic invariant mass W

and it generates **hadronic showers!**

In the context of the GENIE empirical model, hadronic showers are produced by answering the following 3 questions:

- How many hadrons are produced?
- What are the hadron IDs?
- What are the hadron 4-momenta?

Hadronization modelling impacts:

- **Neutrino energy reconstruction**

If the detector responds differently to different hadrons, then you need to know the exact mixture of hadrons in your showers!

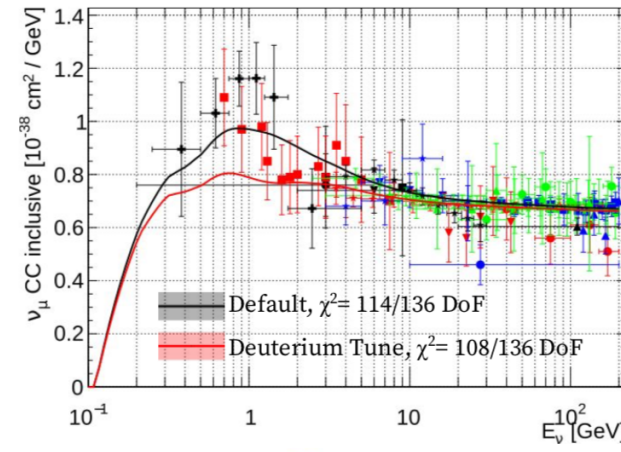
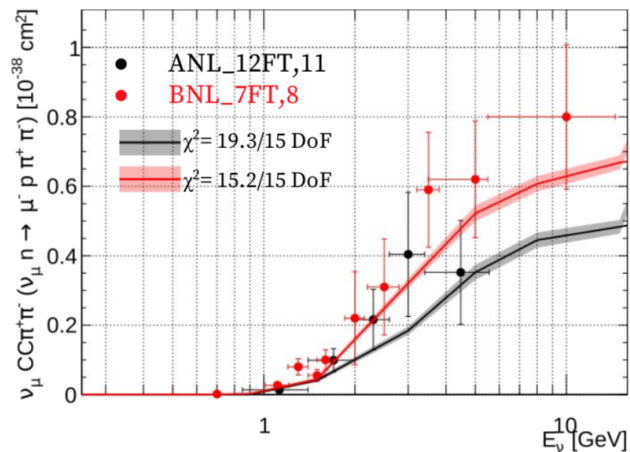
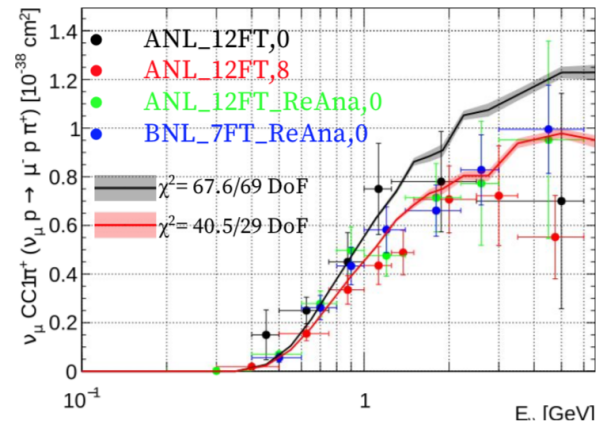
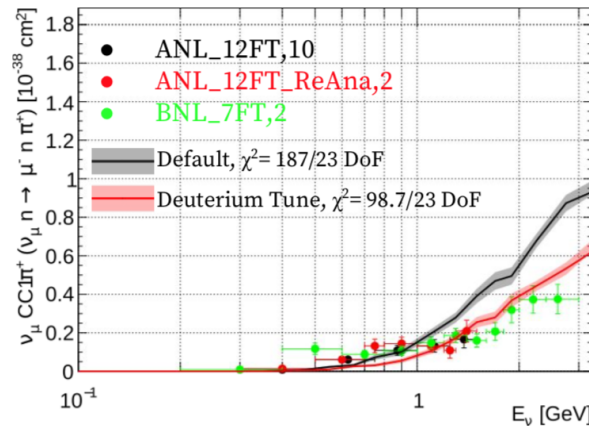
- **Efficiency calculations for event identification**

For example, hadronic showers (from NC or high- y ν_μ CC events) with large EM component could be misclassified as ν_e CC events. Shower mismodelling can impact the estimation of backgrounds.

Significance

Hadronization modelling is not just about hadron shower shapes.

String coupling with cross-section modelling (even for $1-\pi$): Decomposes the (computed) incl. cross-section into nearly all excl. cross-sections.



Hadronization modelling in GENIE

Three main elements:

- PYTHIA6, valid at higher W
- Empirical model, valid for SIS/DIS at $W < 3$ GeV
- Empirical model, specialised for DIS charm production

The above are not all the places in GENIE that produce hadrons (pre-FSI):

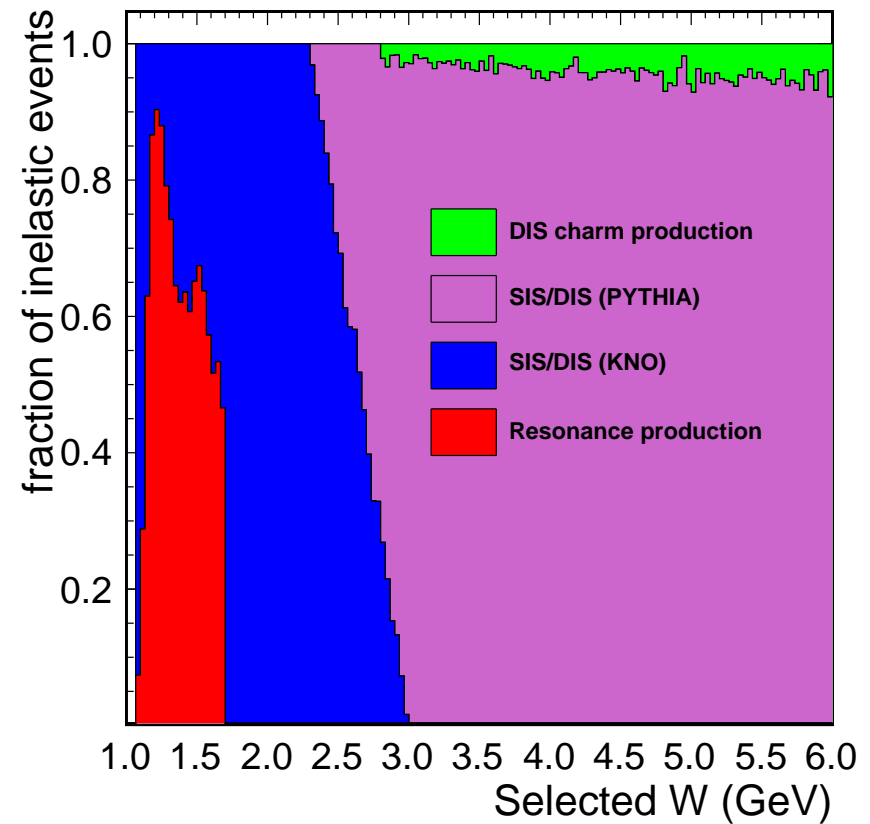
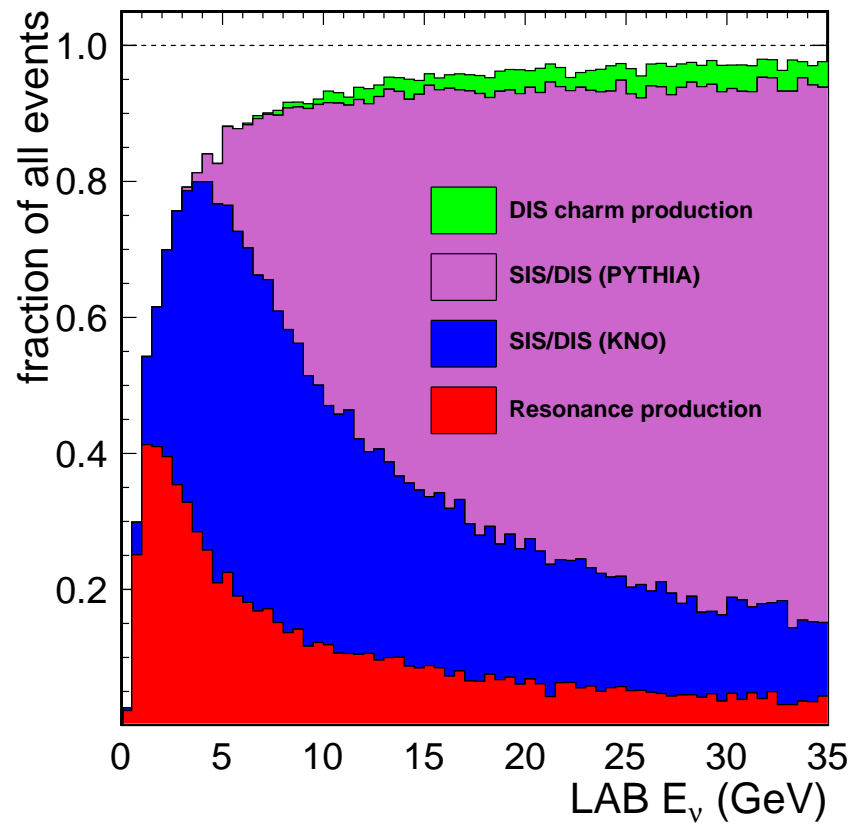
- Resonance single-pion production
- Other states from baryon resonance decays
- Single-Kaon production

These modelling elements are not thought of as parts of the family of GENIE *fragmentation models* and will not be covered.

However, it should be recognised that the distinction can be somewhat arbitrary (e.g. mixture of resonance / non-resonance contributions)

Hadronization modelling

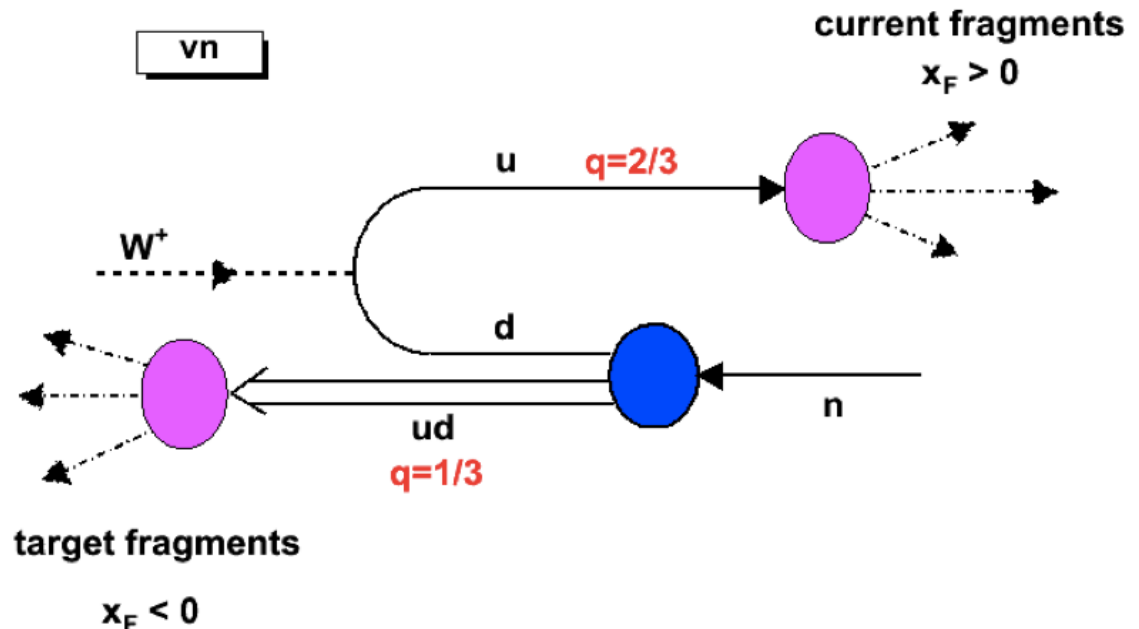
Fractions of GENIE events generated by each hadronization model:



Basic picture of hadronization

In the simple quark-parton model, the lepton interacts via W/Z exchange with one of the (anti)quarks in the nucleon:

- Struck quark hadronizes (current fragments, predominantly $x_F > 0$)
- Target remnant also hadronizes (target fragments, $x_F < 0$)
- 2 correlated hadron jets.
- Smooth transition through a central rapidity region.
- At low W , the two fragmentation regions largely overlap.



Data on neutrino-induced hadron shower characteristics

Several pieces of data exist.

- Average charged and neutral particle multiplicities
- Forward and backward hemisphere average multiplicities
- Multiplicity dispersion as function of avg multiplicity and W
- Multiplicity correlations (e.g. charged hadrons - π^0)
- Fragmentation functions (z distributions)
- x_F distributions
- p_T^2 distributions
- $x_F - p_T^2$ distributions

However, coverage of the low - W range ($W < 4\text{-}5$ GeV is poor)

- LUND string fragmentation model
- Uses the assumption of linear confinement as a starting point.
- As partons move apart, their colour flux tube gets stretched.
- Stored potential energy increases linearly with distance of colour charges.
- You can think of the "string" as the axis of the flux tube.
- The string constant is $\sim 1 \text{ GeV/fm}$.
- As the potential energy increases, the string may break producing a $q\bar{q}$ pair.
- String breaks causally disconnected; simulated in a convenient order.
- A break typically creates a meson.
- Baryons also produced; A string can break by antidiquark-diquark production, or baryons can be produced using a 'popcorn' model.
- With every break, a produced hadron takes away a fraction of the available energy/momentum.
- Continuing till some cut-off point.

Driving PYTHIA6 from GENIE

Some amount of monkey business in making **quark + diquark** assignments most certainly due to our own unfamiliarity with PYTHIA. Luckily, overall generation outcomes not sensitive to choices made.

Init state	Hit quark	Leading quark	Remnant system	PYTHIA6 assignment		Weirdness level
$\nu + p$ CC	d valence	$(d \rightarrow) u$	uu	u	uu	
$\nu + p$ CC	d sea	$(d \rightarrow) u$	$\bar{d} + uud$	u	uu	*
$\nu + p$ CC	s sea	$(s \rightarrow) u$	$\bar{s} + uud$	u	uu	**
$\nu + p$ CC	\bar{u} sea	$(\bar{u} \rightarrow) \bar{d}$	$u + uud$	u	uu	***
$\nu + n$ CC	d valence	$(d \rightarrow) u$	ud	u	ud	
$\nu + n$ CC	d sea	$(d \rightarrow) u$	$\bar{d} + udd$	u	ud	*
$\nu + n$ CC	s sea	$(s \rightarrow) u$	$\bar{s} + udd$	u	ud	**
$\nu + n$ CC	\bar{u} sea	$(\bar{u} \rightarrow) \bar{d}$	$u + uud$	u	ud	***
...
...

PYTHIA6 tuning

- NOMAD (NUX) PYTHIA6 tuning was adopted in 2007.
- Some PYTHIA6 defaults were restored in later GENIE re-tune (2010).

	PYTHIA default	NUX 2001	GENIE 2010 re-tune
$s\bar{s}$ production suppression	0.30	0.21	0.30
$\langle p_T^2 \rangle$ (GeV^2)	0.36	0.44	0.44
Non-gaussian p_T tail parameterization	0.01	0.01	0.01
Fragmentation cut-off energy (GeV)	0.80	0.20	0.20

Main issues with GENIE PYTHIA studies (circa 2007!).

- Could not find enough knobs to influence predictions.
 - How to express uncertainty?
- Could not understand change in behaviour below $W \approx 2.5 - 3.0$ GeV.
 - Limits of validity range?

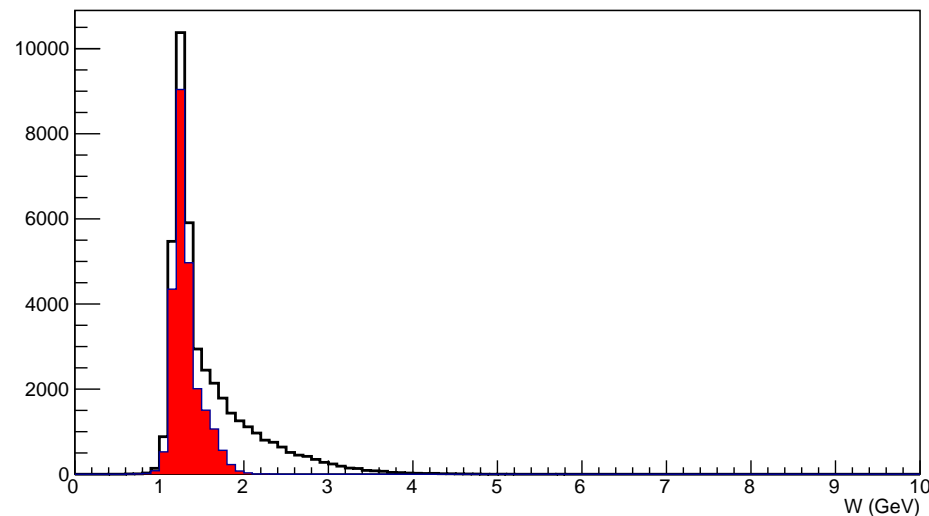
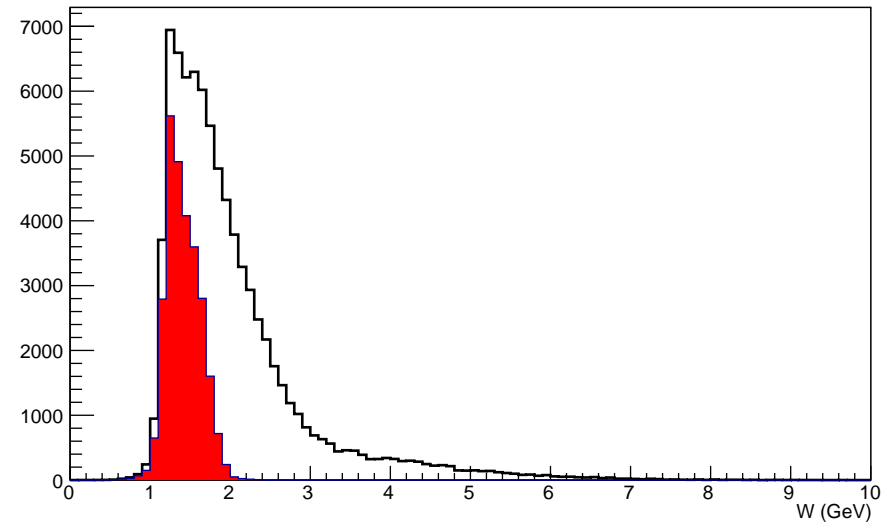
Extending the validity of GENIE model to lower W

On the right, the invariant mass distributions for inelastic events. (Distribution is smeared due to Fermi momentum.)

Up plot: DUNE, Down: HK
Red component: Resonances

Kinematic area **below 2.5 - 3.0 GeV** in invariant mass is **critically important**.

Augment PYTHIA with an **empirical GENIE model**, anchored to data and **valid in the area below 3 GeV**.
Install handles to express uncertainty.



Empirical low-W model

- An **effective KNO-based hadronization model was built** (T.Yang, H.Gallagher, P.Kehayias, C.Andreopoulos - circa 2007) for low W and was "*integrated*" with PYTHIA to cover the full kinematic space (AGKY model, Eur.Phys.J.C63:1-10,2009)
- The model was **anchored on several pieces of bubble chamber data and captures several observations on the characteristics of neutrino-induced hadron showers** (for an excellent description, see Norbert Schmitz, Adv.Ser.Direct.High Energy Phys. 2 (1988) 3-56)
- A similar, KNO-inspired model pre-existed (in neugen3). Several model improvements were installed in 2007-2008,
- Several caveats were recognized over time; few improvements were made (e.g strange baryon production by K.Hoffmann, H.Gallagher)

Empirical low-W model: How many hadrons are produced?

First order of business is to calculate the hadronic multiplicity.

First, we answer that question **on average**.

Average charged hadron multiplicities $\langle n_{ch} \rangle$ are well described by:

$$\langle n_{ch} \rangle = a + b \cdot \ln(W^2 / \text{GeV}^2)$$

The values of a,b were measured in several experiments.

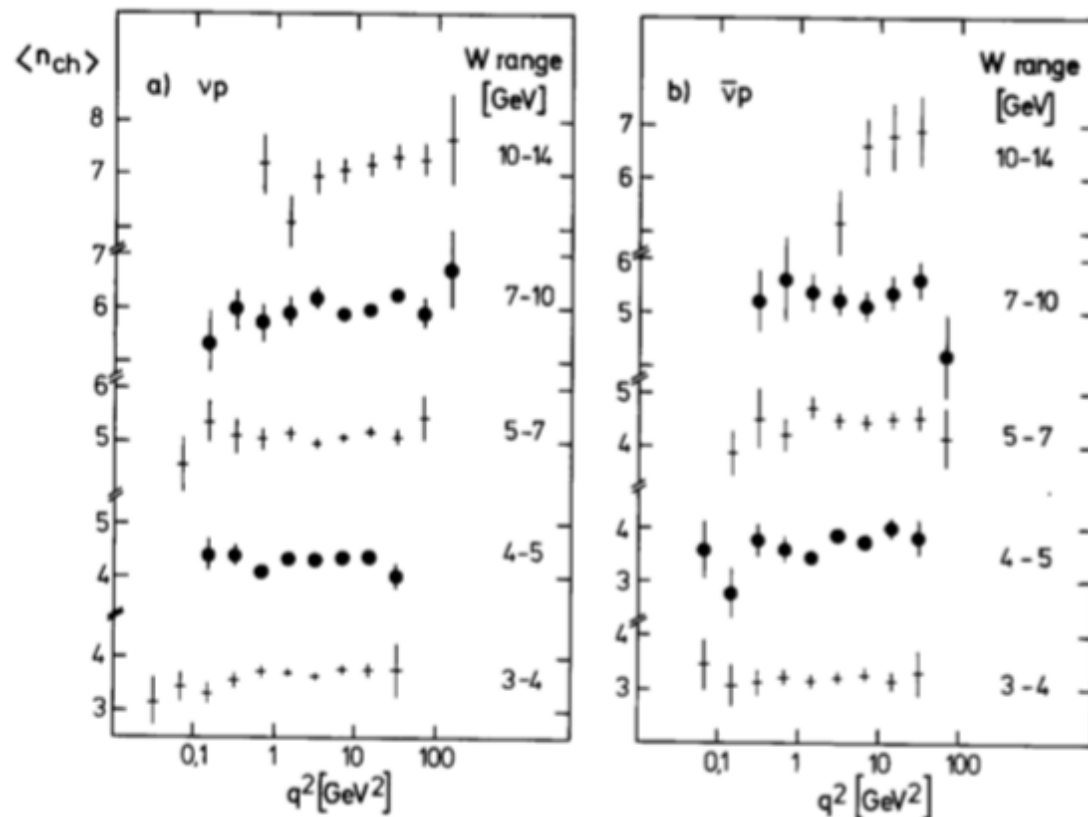
GENIE uses:

	νp	νn	$\bar{\nu} p$	$\bar{\nu} n$
a	0.40	-0.20	0.02	0.80
b	1.42	1.42	1.28	0.95

Empirical low-W model: How many hadrons are produced?

Average charged hadron multiplicities $\langle n_{ch} \rangle$ could, more generally, have an additional Q^2 dependence:

$$\langle n_{ch} \rangle = a + b \cdot \ln(W^2 / \text{GeV}^2) + b' \ln(Q^2 / \text{GeV}^2)$$

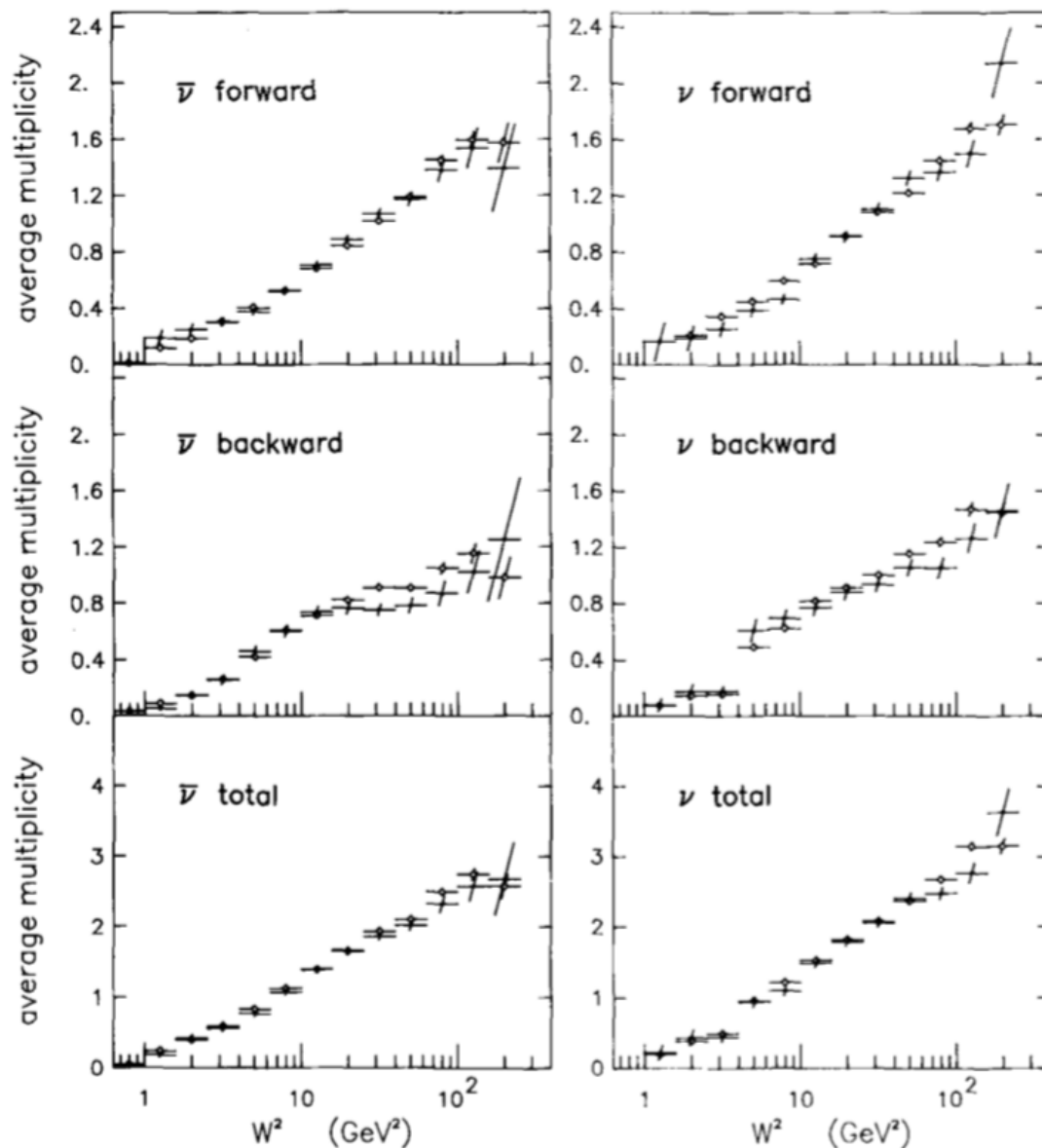


No Q^2 dependence has been observed in $\nu/\bar{\nu}$ scattering [H. Grassler et al., Nucl. Phys., **B223**, 269 (1983)].

Values of b' are 0.04 ± 0.02 for νp and 0.05 ± 0.04 for $\bar{\nu} p$

In GENIE, $b' = 0$ for all channels.

Empirical low-W model: How many hadrons are produced?



The average neutral pion multiplicity was found to be:

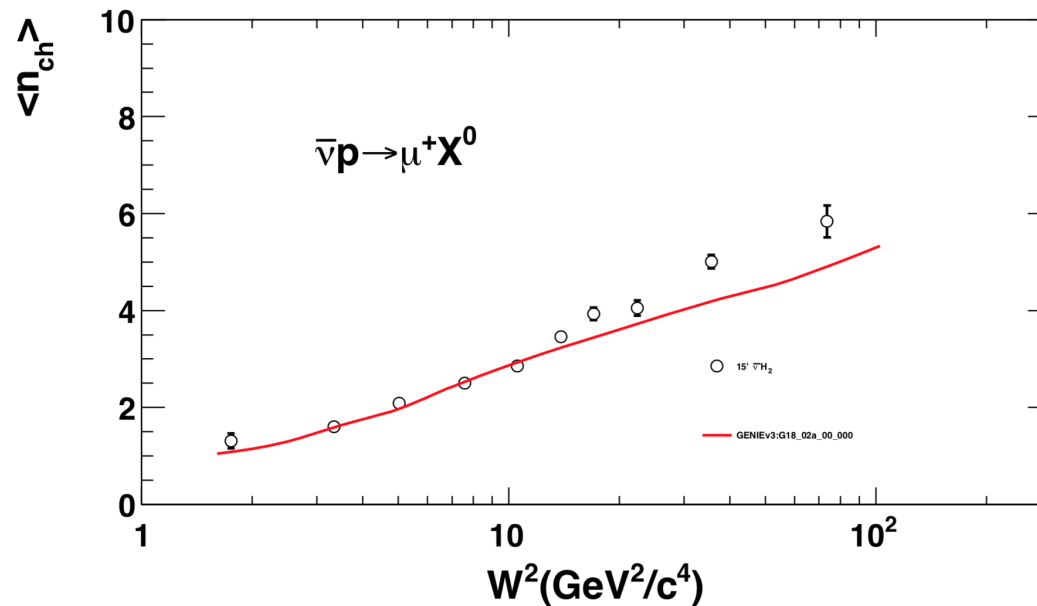
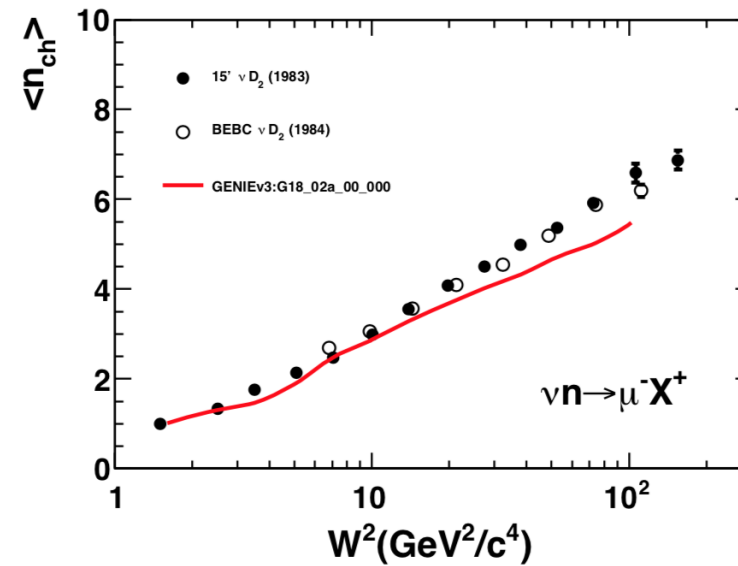
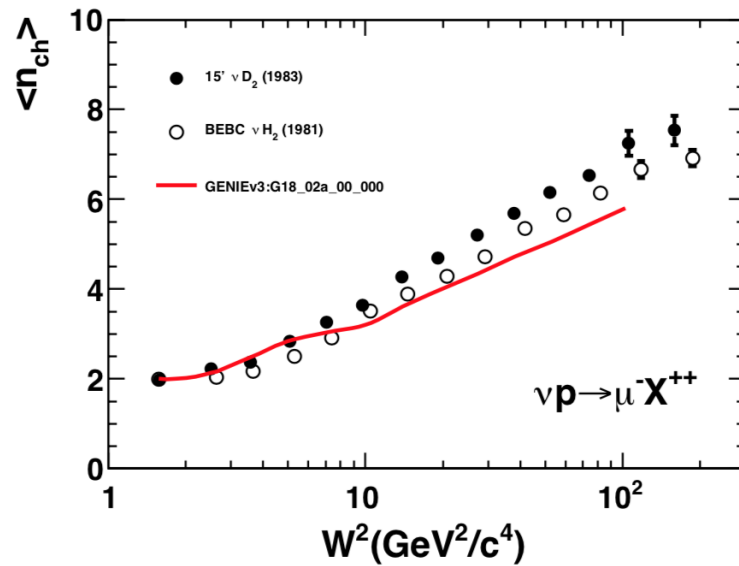
$$\frac{2 \langle n_{\pi^0} \rangle}{\langle n_{\pi^+} \rangle + \langle n_{\pi^-} \rangle} \approx 1$$

Therefore, we can write:

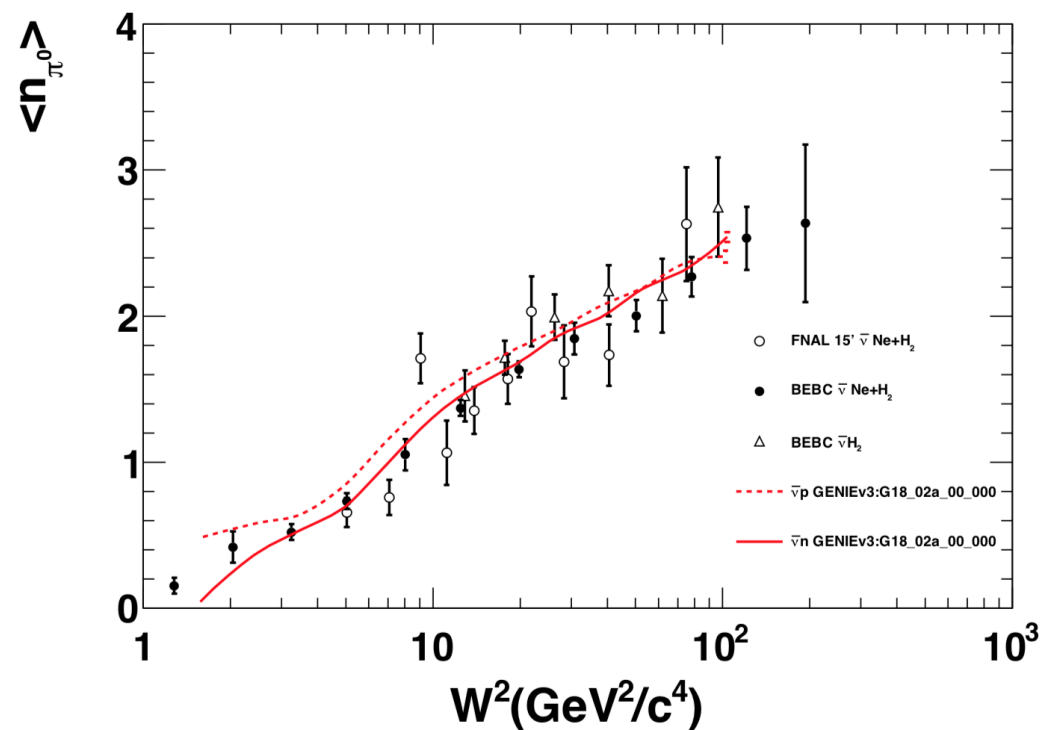
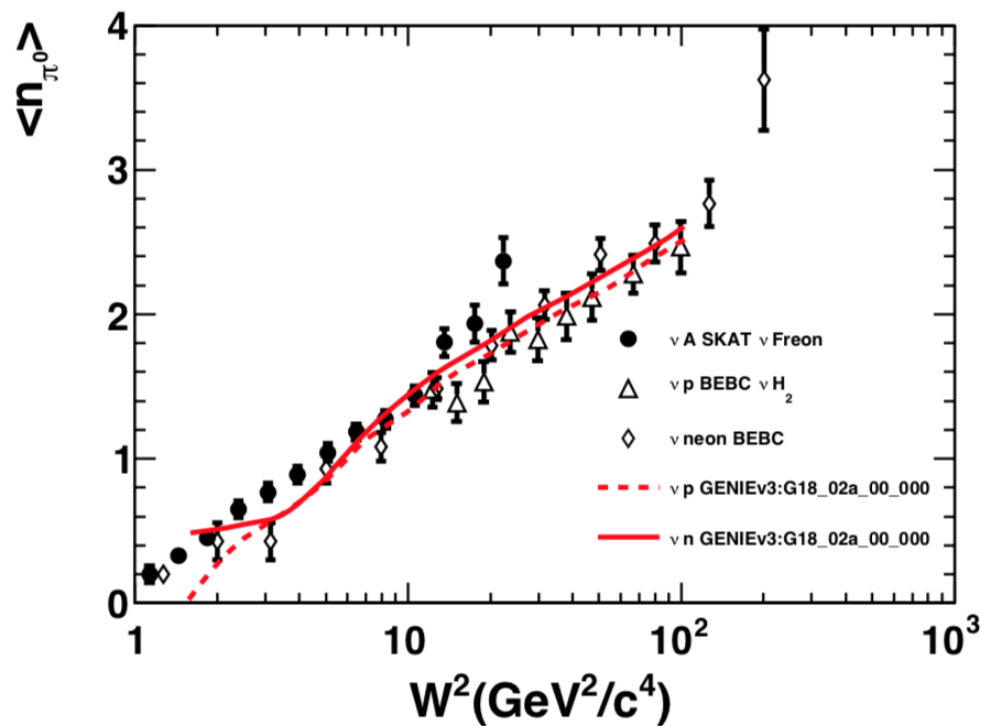
$$\langle n_{tot} \rangle \approx 1.5 \langle n_{ch} \rangle$$

On the left: Avg. π^0 multiplicity (crosses) and half the sum of the avg. charged π multiplicities (squares) [W. Wittek et al., Z. Phys. C **40**, 231 (1988)].

GENIE comparisons with average charged multiplicity data



GENIE comparisons with average neutral multiplicity data



Empirical low- W model: How many hadrons are produced?

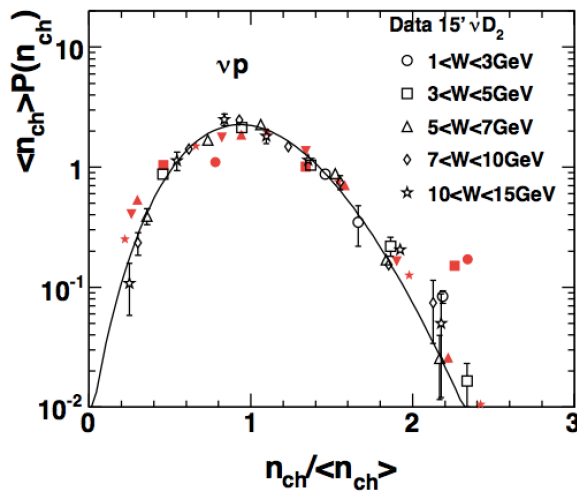
From **average** \rightarrow to **actual** multiplicities on an event-by-event basis?

- Need more information than just $\langle n_{tot} \rangle$!
- Require the probability distribution of n_{tot} , $P(n_{tot})$.

- Draw **actual** multiplicities from a Poisson distribution with given **average**?
- The particles **are not independently produced** and the actual multiplicity is not Poisson-distributed.
- **KNO scaling** to the rescue!

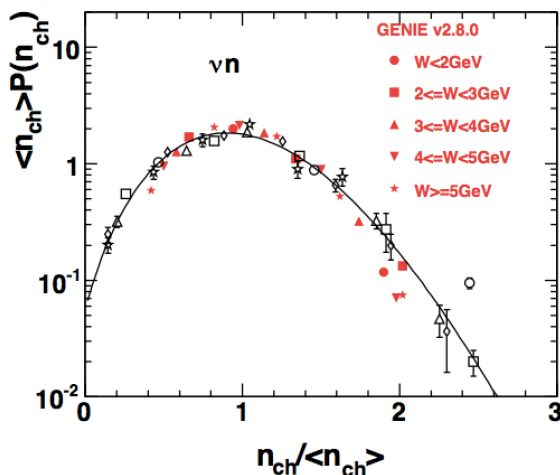
Empirical low-W model: KNO scaling

KNO scaling: $\langle n \rangle P(n) = f(n / \langle n \rangle)$ is independent of W [Z.Koba, H.B.Nielsen, P.Olesen, Nucl.Phys.B40,317(1972)]



The function $f(z = n / \langle n \rangle)$ is parameterized using the Levy function with parameter c:

$$L(z; c) = \frac{2e^{-c} c^{cz+1}}{\Gamma(cz + 1)}$$

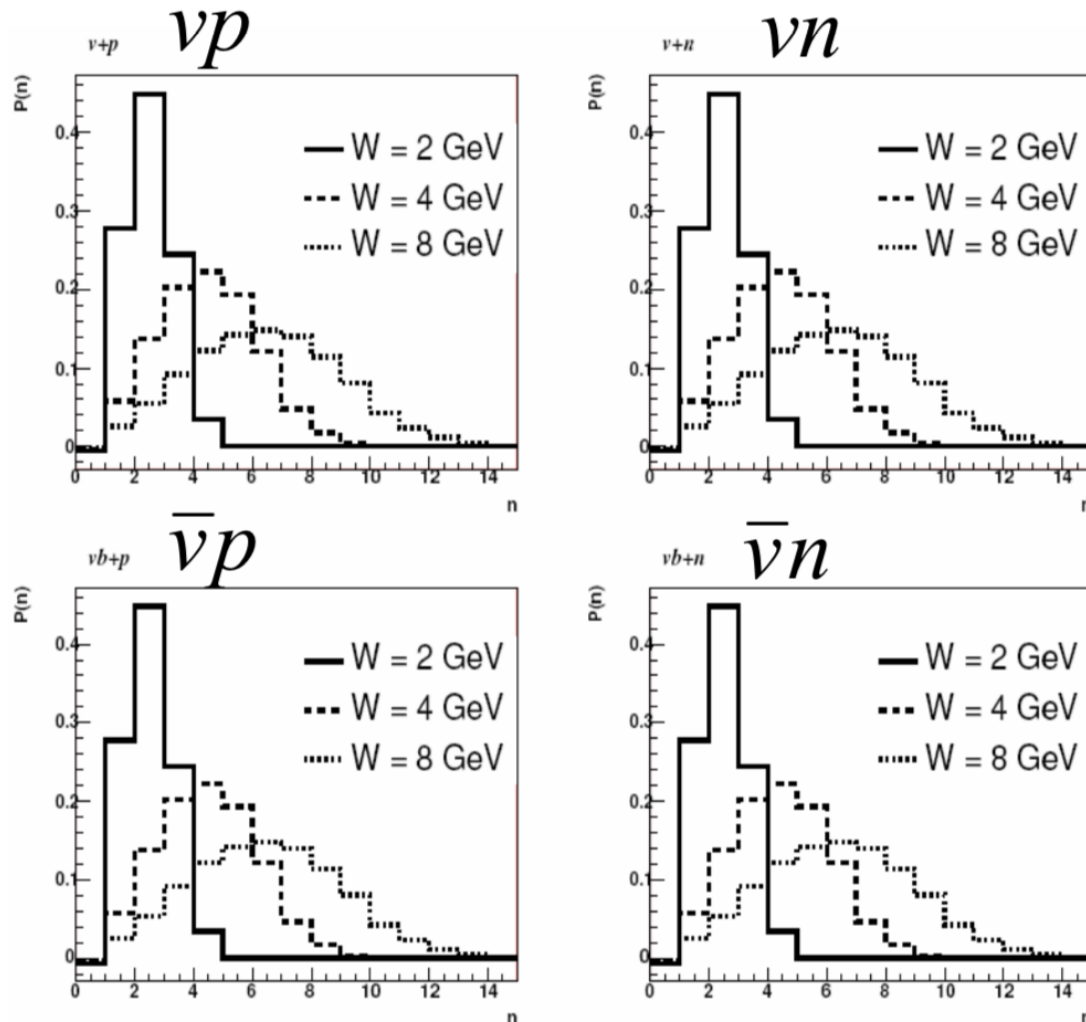


The following parameters c were determined by a GENIE fit to data:

	νp	νn	$\bar{\nu} p$	$\bar{\nu} n$
c	7.93	5.22	as in νn	as in νp

Empirical low-W model: How many hadrons are produced?

From **average** → to **actual** multiplicities on an event-by-event basis?

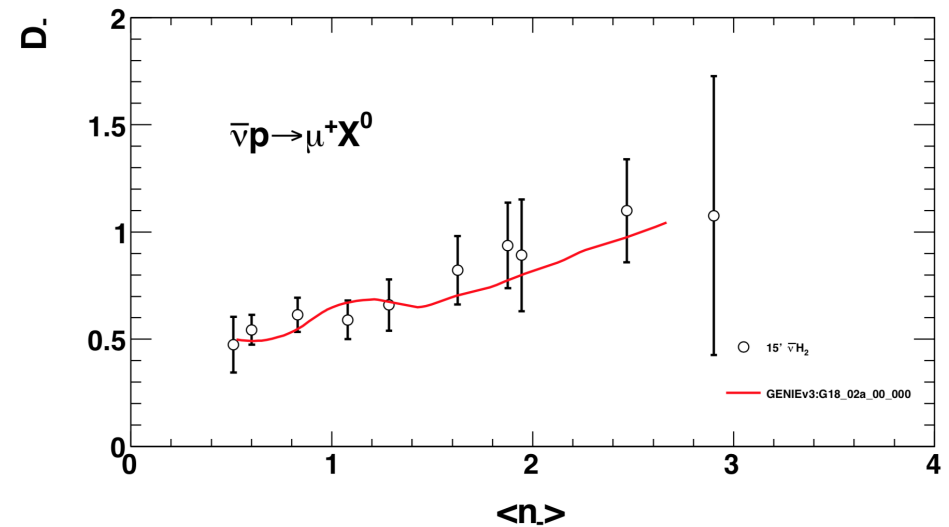
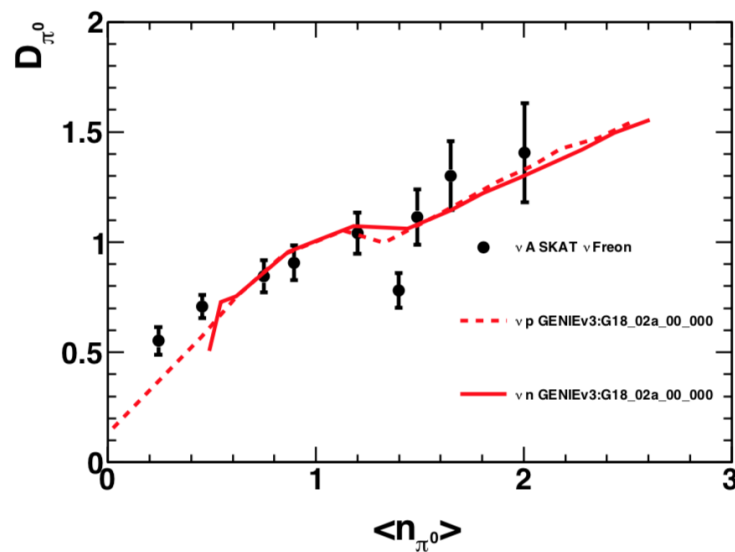
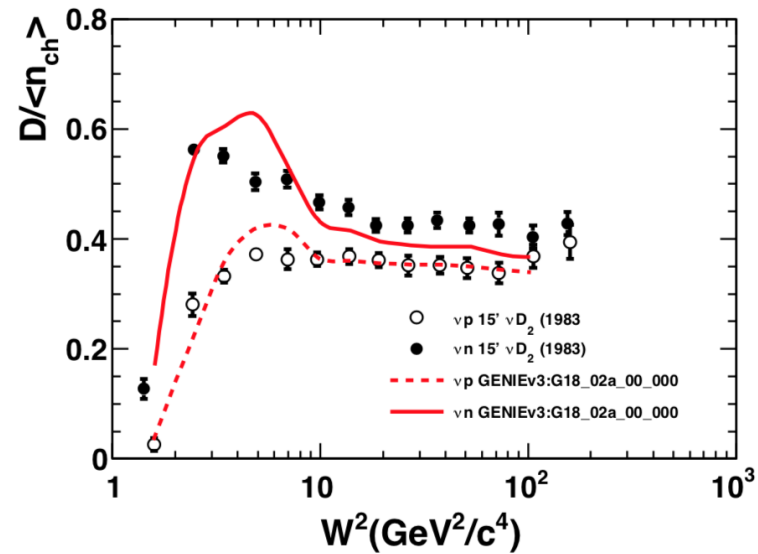
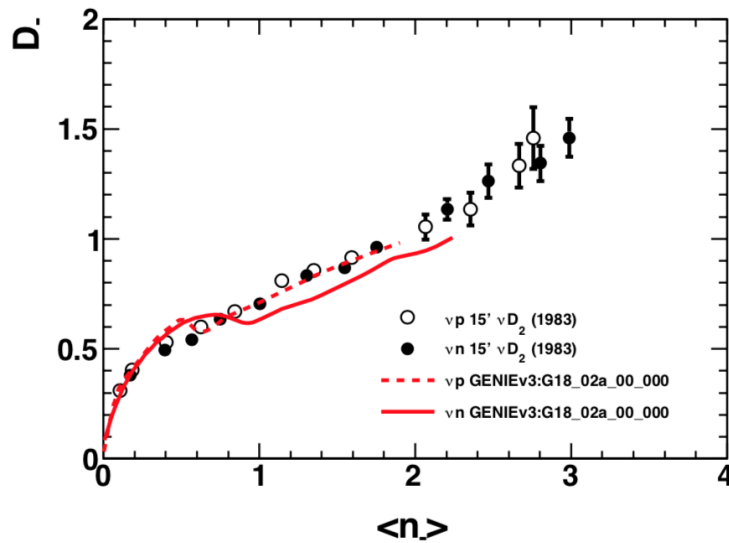


We now have

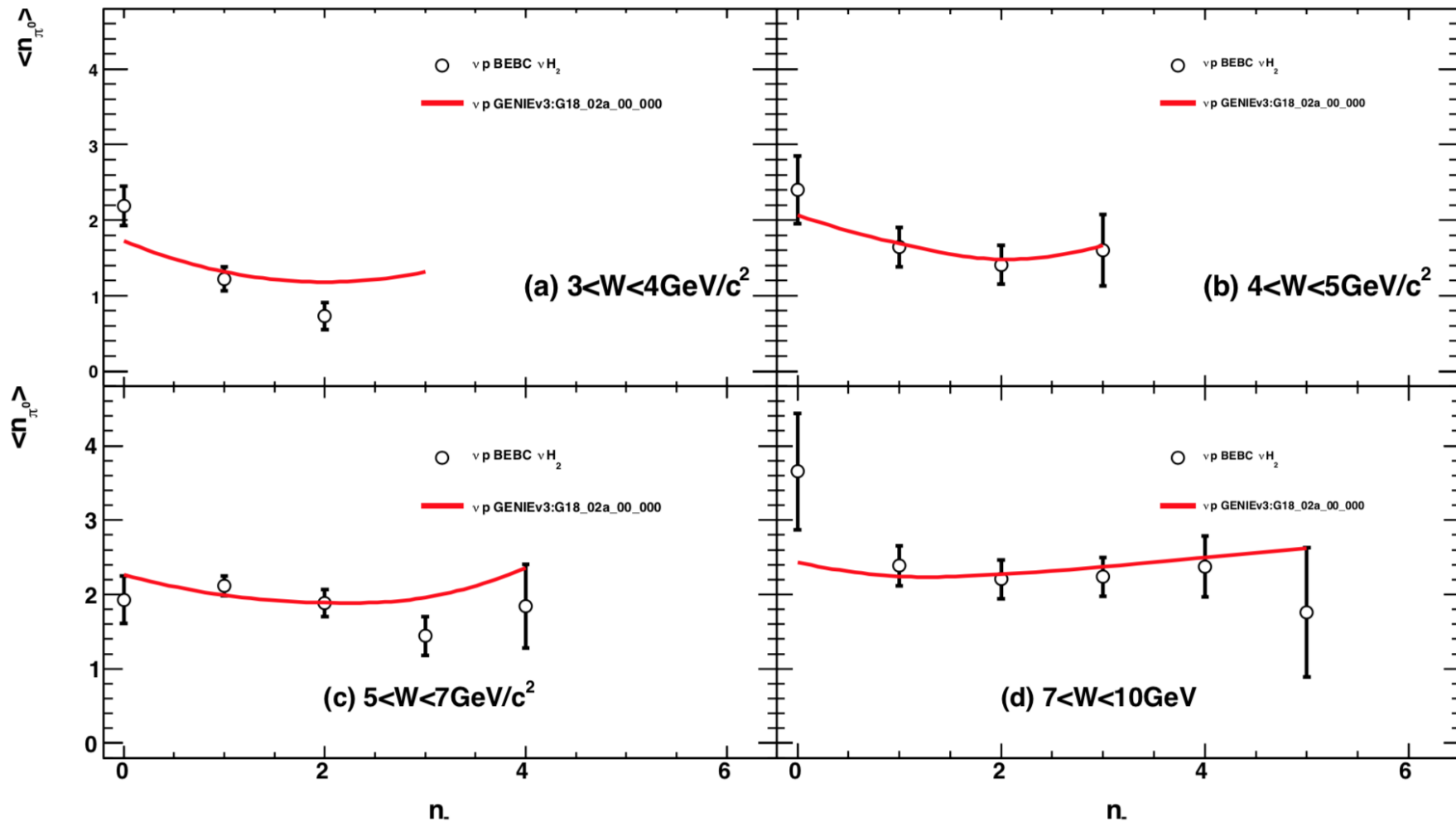
- an expression for $\langle n_{tot} \rangle P(n_{tot}) = f(n_{tot} / \langle n_{tot} \rangle)$, and
- a value for $\langle n_{tot} \rangle$

We can combine the two into an expression for the multiplicity probability distribution $P(n_{tot})$

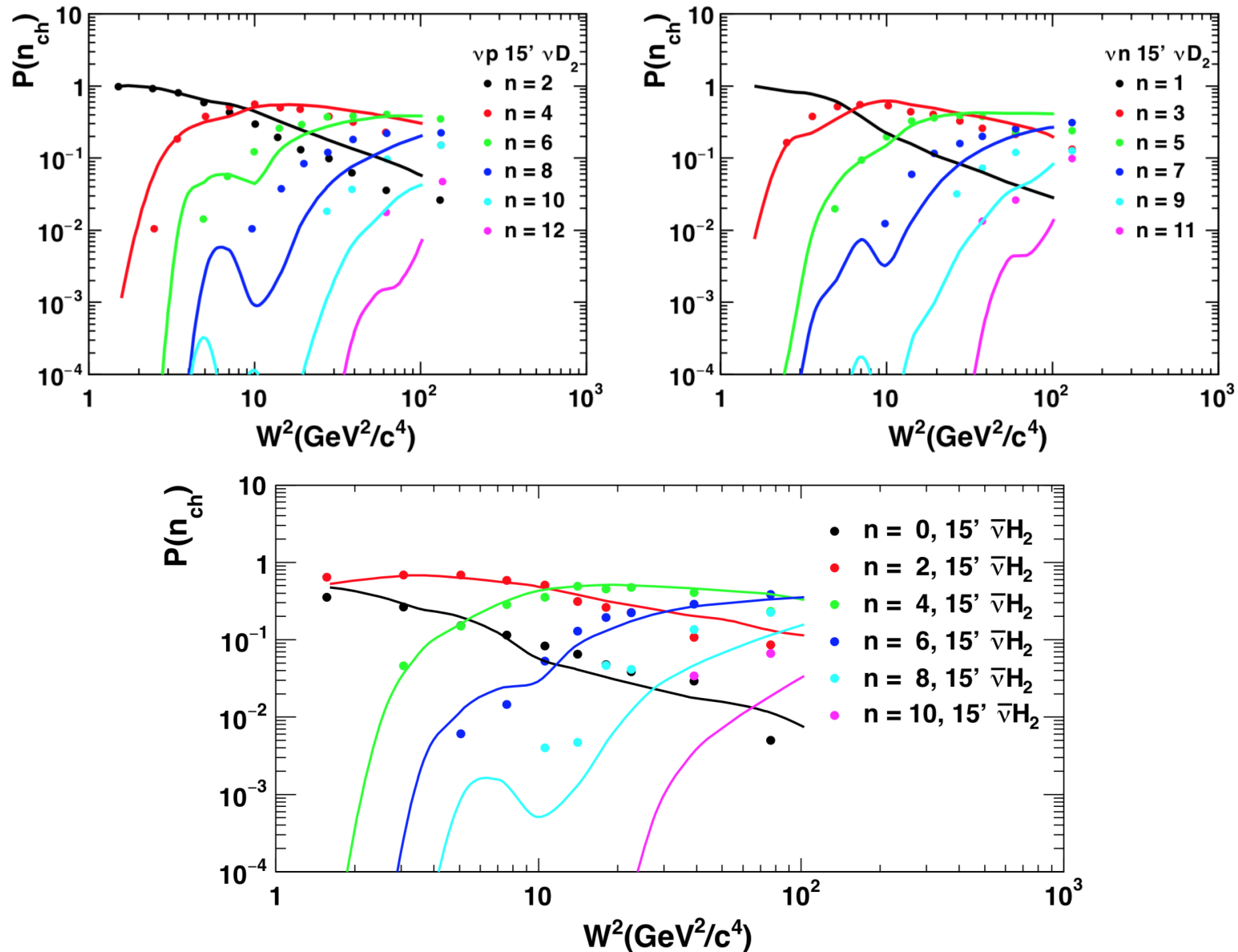
GENIE comparisons with multiplicity dispersion data



GENIE comparisons with multiplicity correlation data



GENIE comparisons with topological cross-section data



Empirical low-W model: Generating the particle spectrum

Because of kinematical constraints, it is assumed that the shower contains only 1 baryon. We decide between a p or n, with probabilities P_p and P_n ($=1-P_p$):

	n_{tot}	νp	νn	$\bar{\nu} p$	$\bar{\nu} n$
P_p	2	1.00	0.33	0.67	0.
	>2	0.67	0.50	0.50	0.33

Subsequently, one of those will be converted to a strange baryon (for ν interactions: $p \rightarrow \Sigma^+$ and $n \rightarrow \Lambda$; for $\bar{\nu}$ interactions: $p \rightarrow \Lambda$ and $n \rightarrow \Sigma^-$)

The probability for generating a strange baryon is given by:

$$\langle n_{hyperon} \rangle = a_{hyperon} + b_{hyperon} \cdot \log(W^2)$$

where

	νp	νn	$\bar{\nu} p$	$\bar{\nu} n$
$a_{hyperon}$	0.022	0.022	0.022	0.022
$b_{hyperon}$	0.042	0.042	0.042	0.042

Empirical low-W model: Generating the particle spectrum

Once a baryon (p, n or hyperon) is generated, there are $n_{tot}-1$ remaining particles. They all assigned meson IDs, with the following procedure:

Conserve strangeness:

If a hyperon was produced, add a strange meson to conserve strangeness (no $\Delta S=1$ production in hadronization; this is added separately).

Conserve charge:

Keep on adding π^+ 's or π^- 's in the hadron shower till charge is balanced.

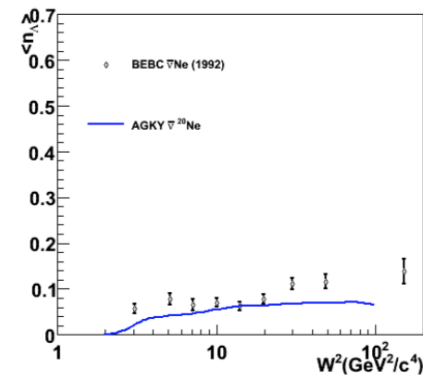
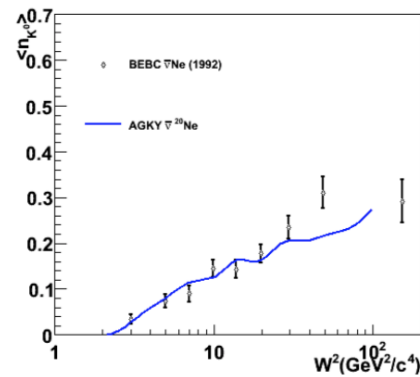
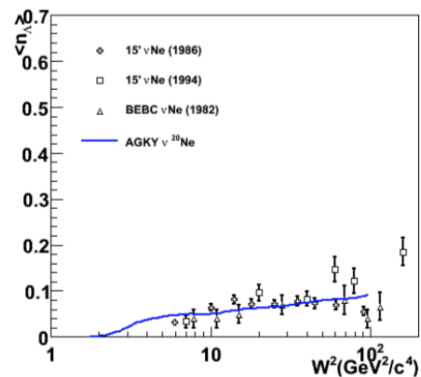
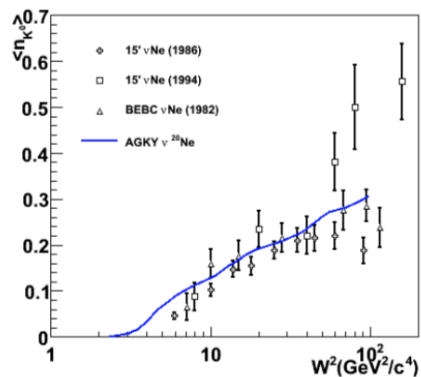
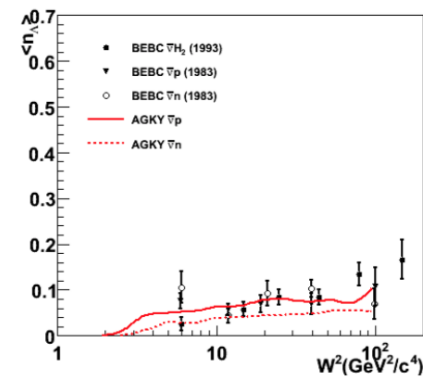
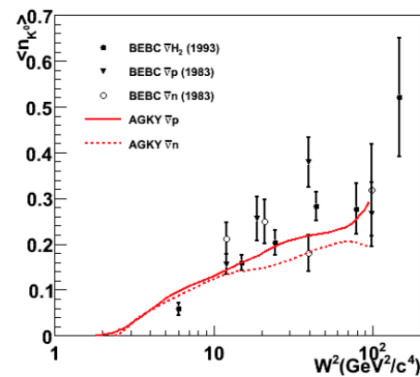
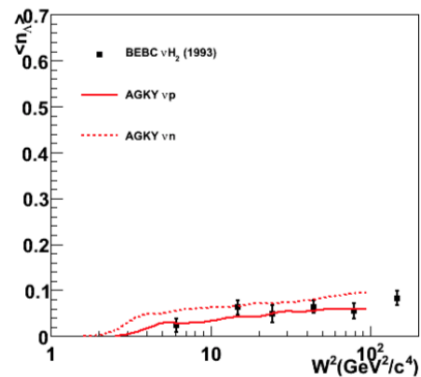
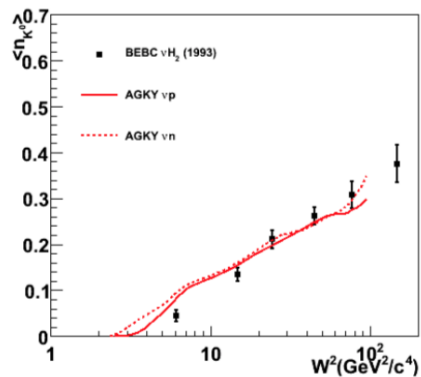
Fill-up:

Add particles in pairs with zero net strangeness and charge till all n_{tot} particle codes have been assigned.

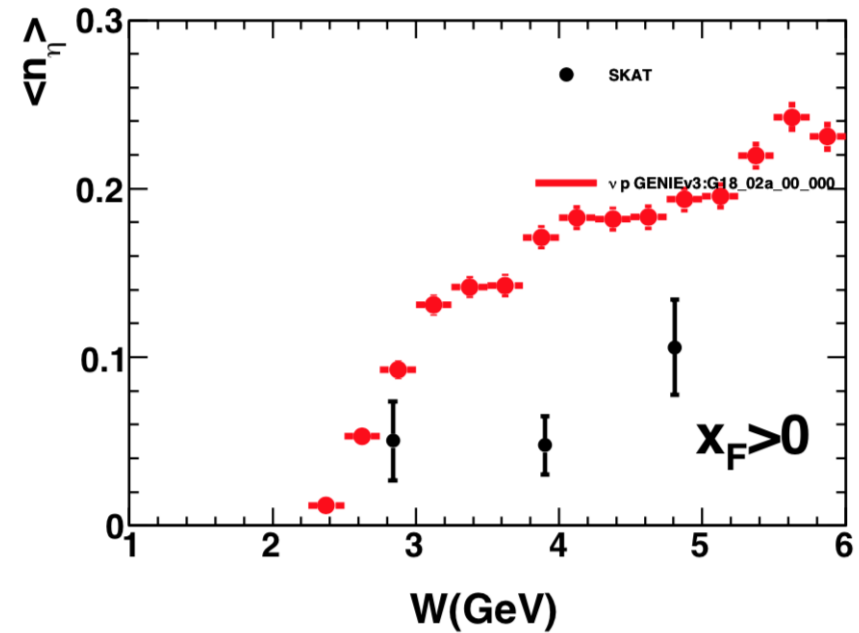
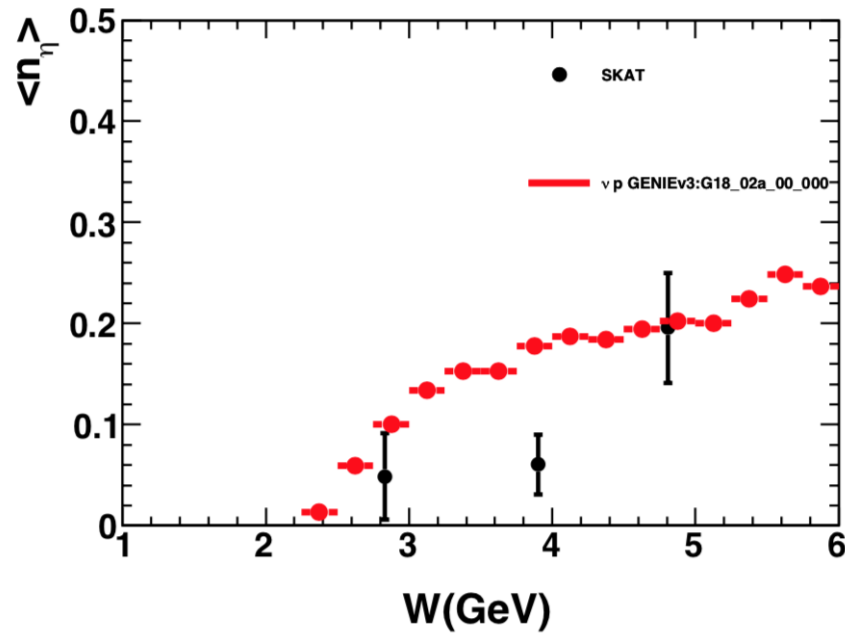
Particle pairs are added with the following probabilities:

$\pi^0\pi^0$	$\pi^+\pi^-$	$K^0\bar{K}^0$	K^+K^-	$\pi^0\eta$	$\eta\eta$
0.3133	0.6267	0.03	0.03	0.	0.

GENIE comparisons with K^0 and Λ production data



GENIE comparisons with η production data (ν)

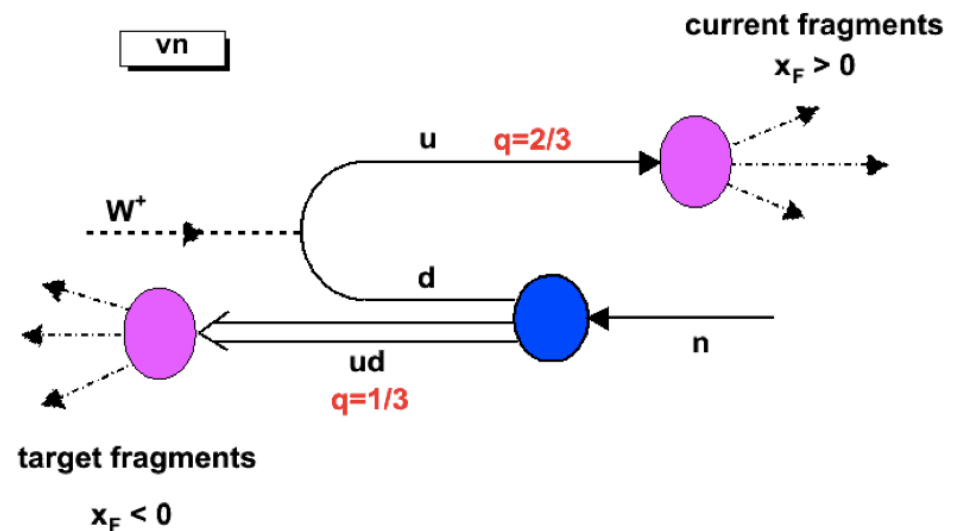


Empirical low- W model: Generating particle momenta

Early versions of the model generated momenta using phase space decays.

This fails to account for key features in the data!

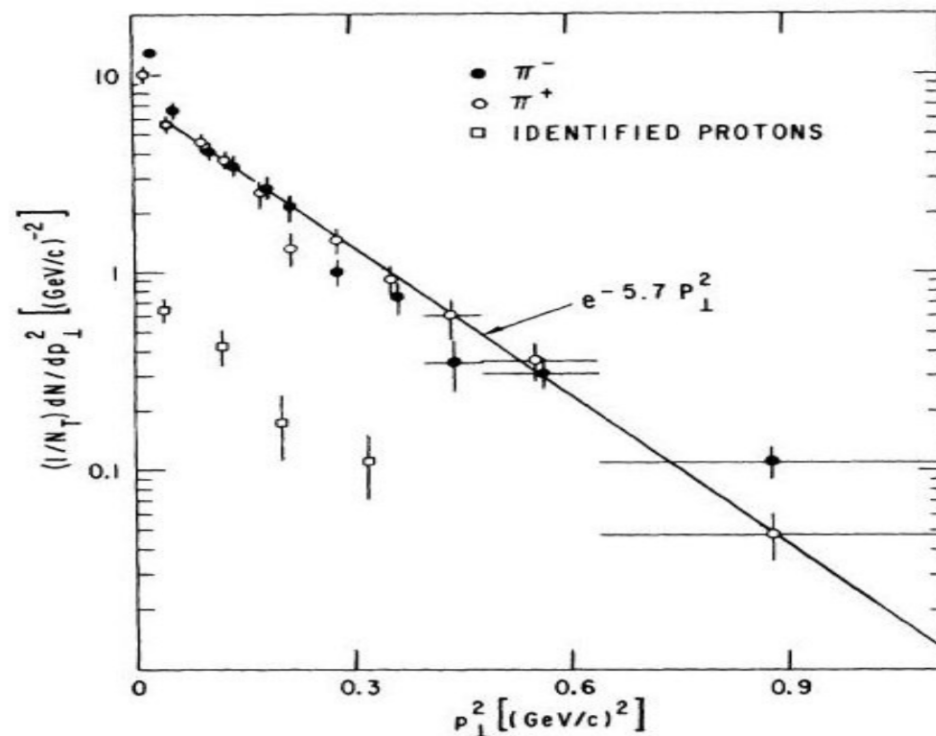
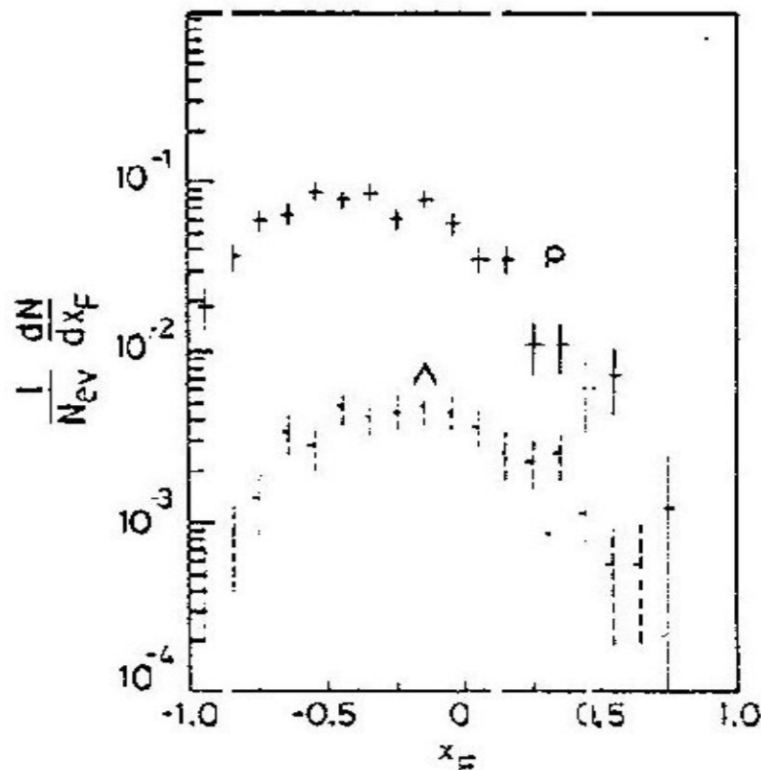
- Final-state nucleon most likely associated with target fragments.
A heavy, predominantly backwards-going particle creates a multiplicity imbalance in the two x_F hemispheres.
- The transverse momentum (p_T) and longitudinal momentum ($\propto x_F$) can not come from the same distribution.



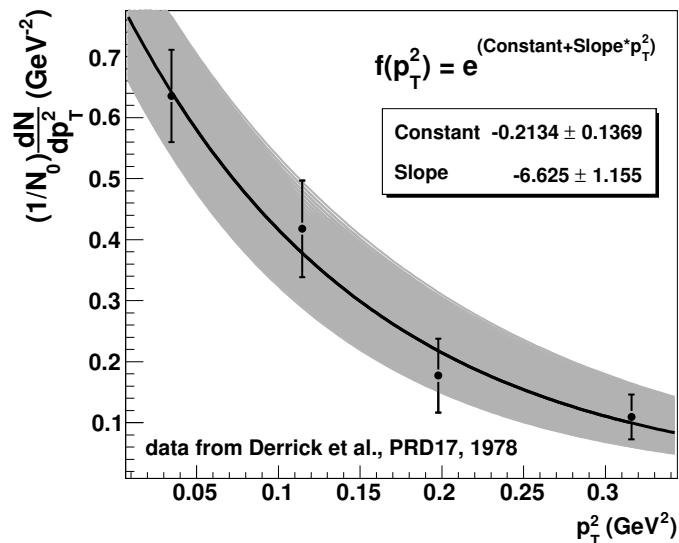
Empirical low- W model: Generating particle momenta

How to introduce an x_F asymmetry and limit p_T (less spherical hadron showers in the HCM frame)?

Looking for guidance in data we found the following in Neutrino 1982 proceedings (Cooper) - Never published (to our knowledge) but a reasonable enough starting point for our empirical model.



Empirical low-W model: Generating particle momenta



Baryon p_T and x_F distributions were parameterized from the previous data:

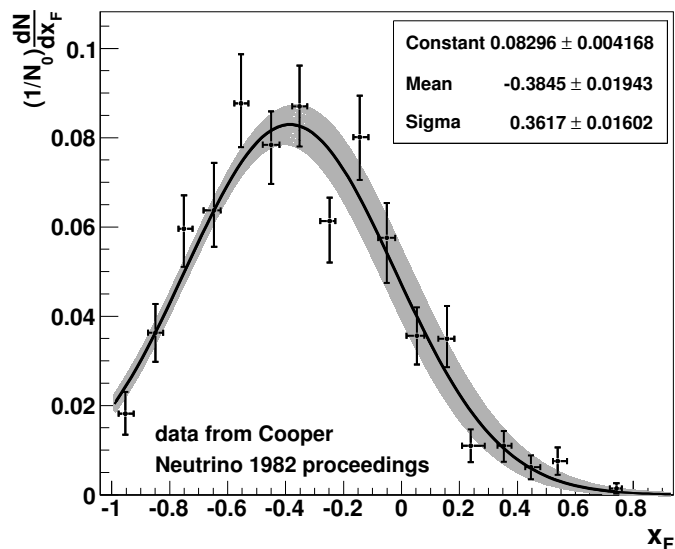
$$f(p_T^2) \propto e^{-6.625 p_T^2 / \text{GeV}^2}$$

$$f(x_F) \propto e^{-3.817(x+0.385)^2}$$

This has **important consequences** for the shower shape.

There is now some evidence that the shower shapes should have been more fwd/bkw-symmetric at lower W .

Future plan is to start shifting the x_F p.d.f. closer to 0 for lower W .



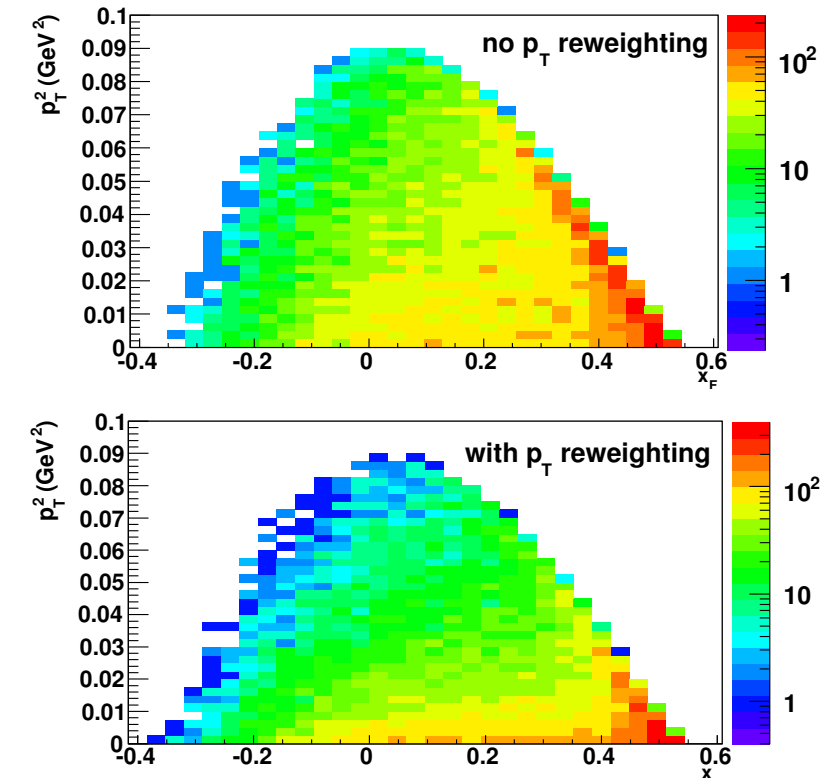
Empirical low-W model: Generating particle momenta

Momenta for the rest of the hadronic system (mass = $W - M_{baryon}$) is generated with a phase space decay.

The likelihood of higher p_T values is decreased by assigning a weight w_i to each hadron i in the decay:

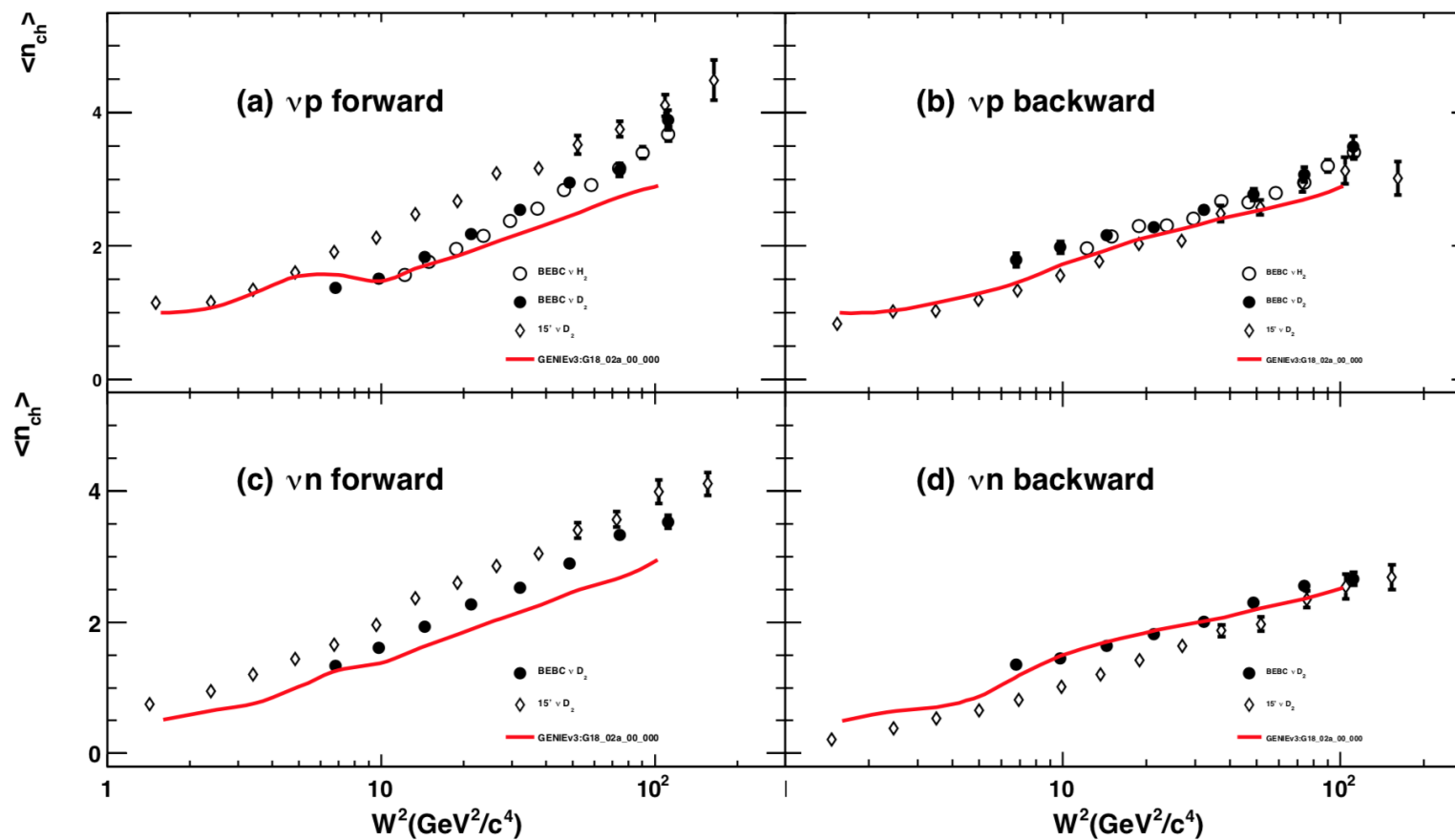
$$w_i = e^{-A \cdot p_T^i} \quad (A = 3.5 \text{ GeV}^{-1})$$

[Clegg and Donnachie, Description of Jet Structure by pt-limited Phase Space, Z. Phys. C 13: 71 (1982)]

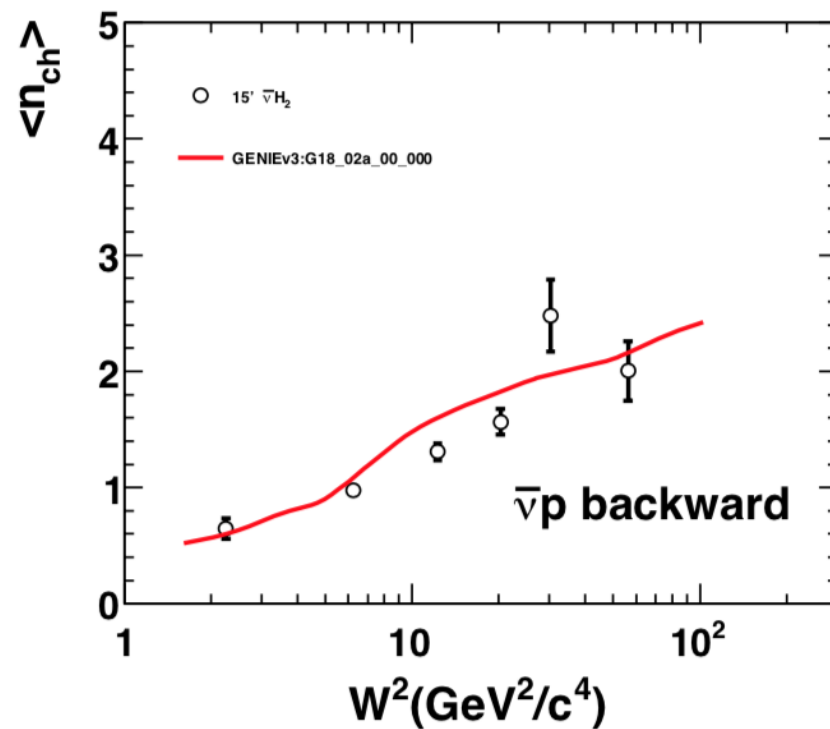
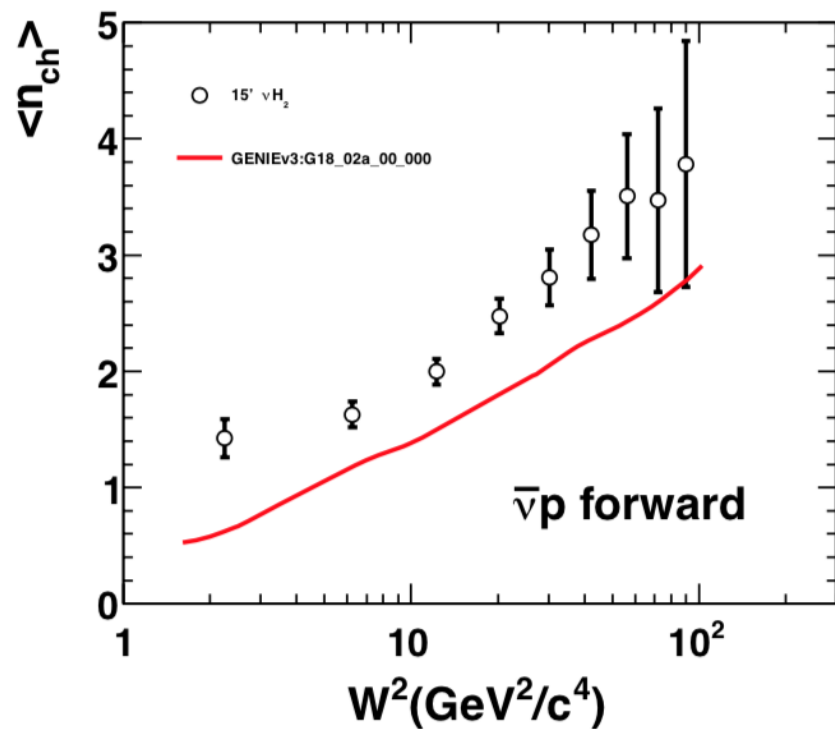


p_T is the momentum perpendicular to the direction of W^\pm/Z^0 in the HCM frame.

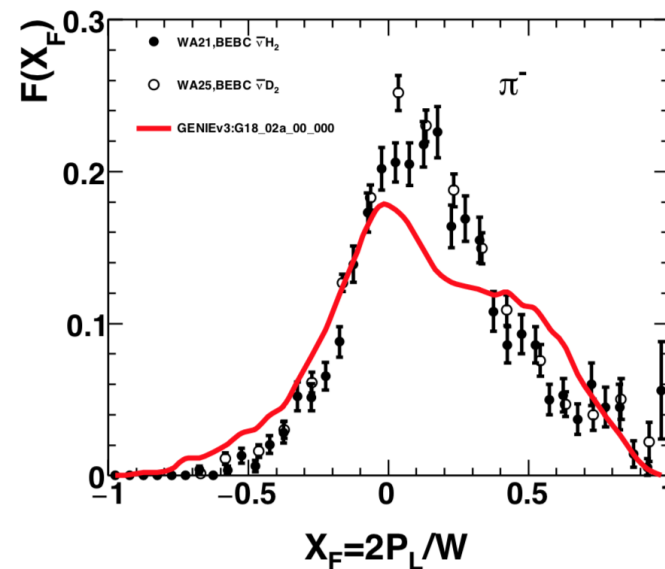
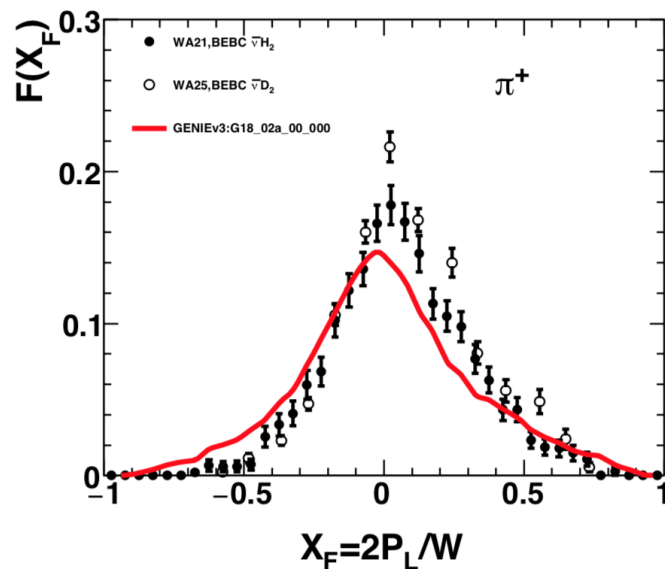
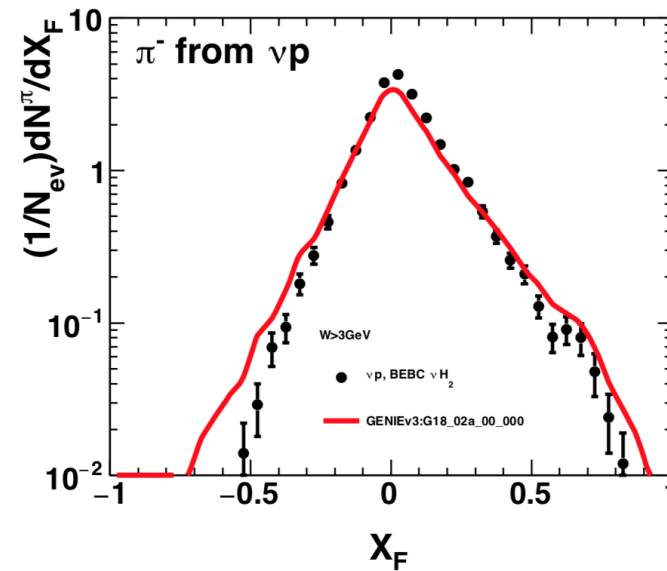
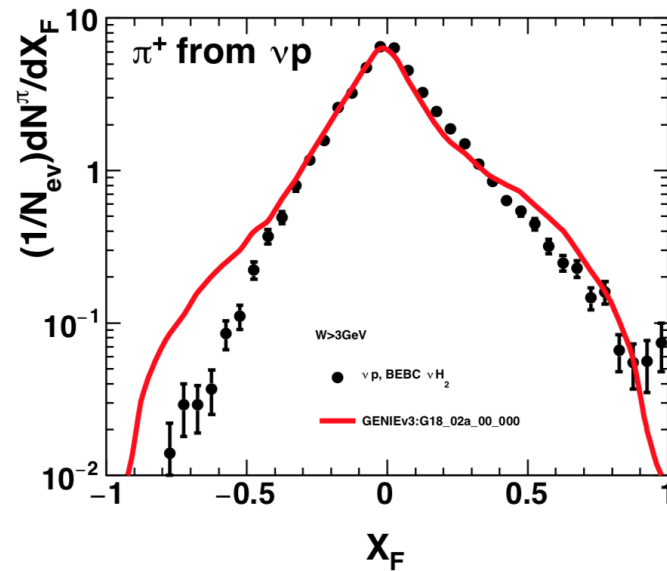
GENIE comparisons with avg. fwd/bkw multiplicity data



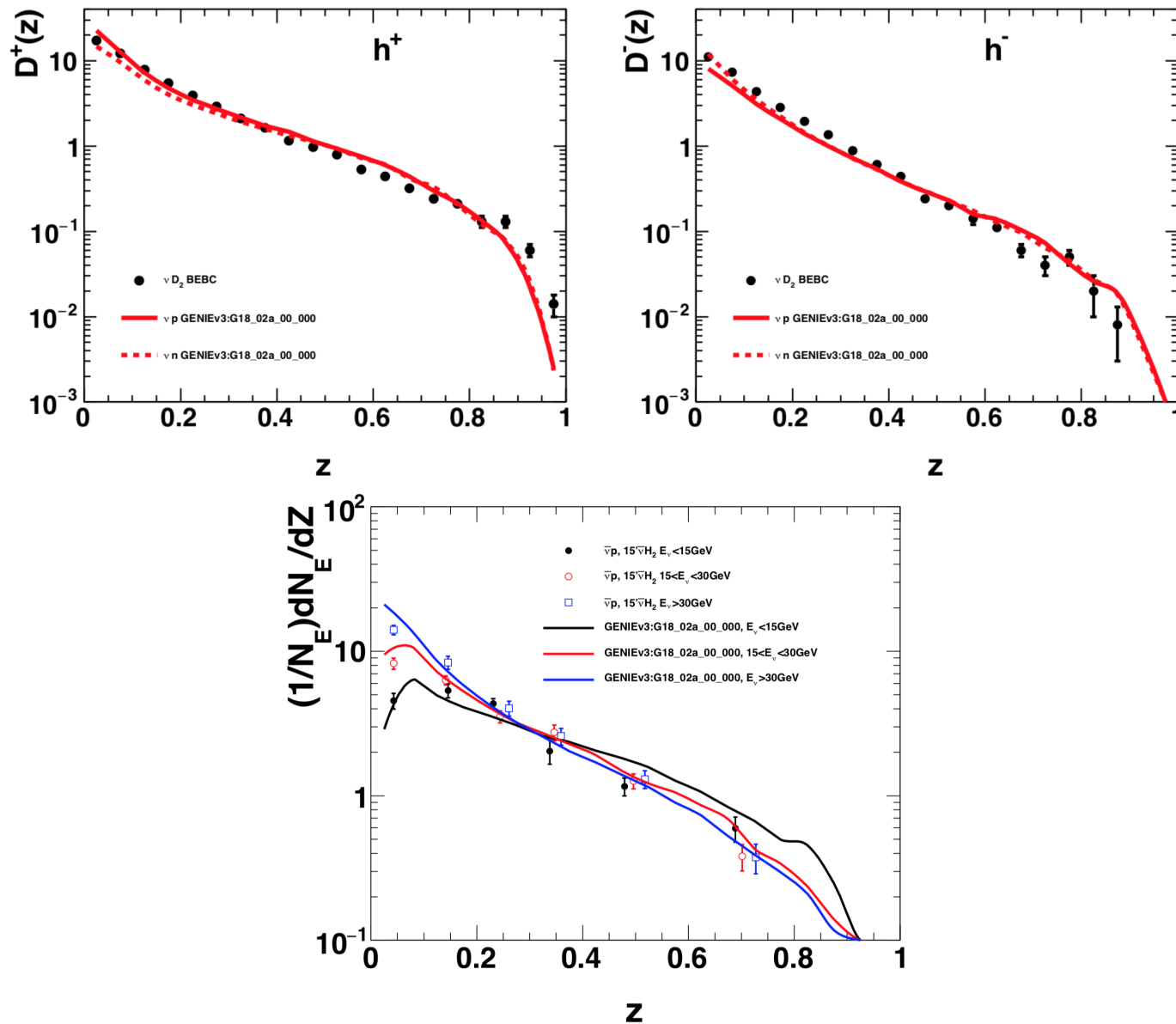
GENIE comparisons with avg. fwd/bkw multiplicity data



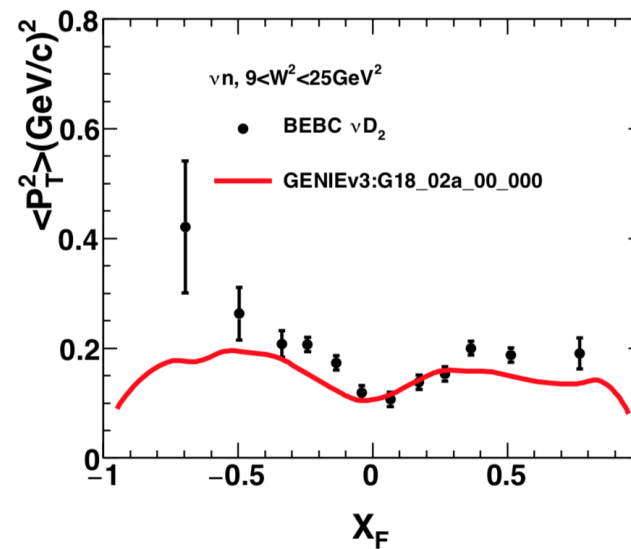
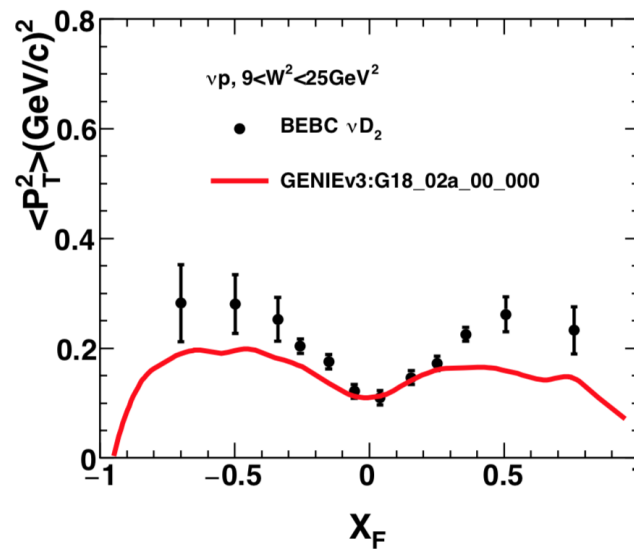
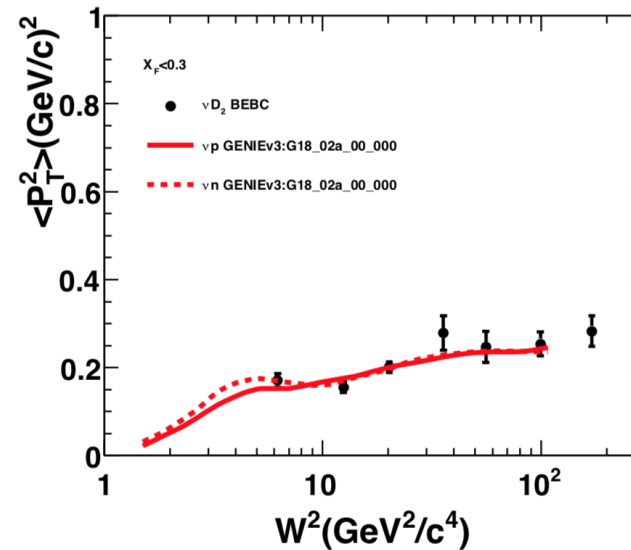
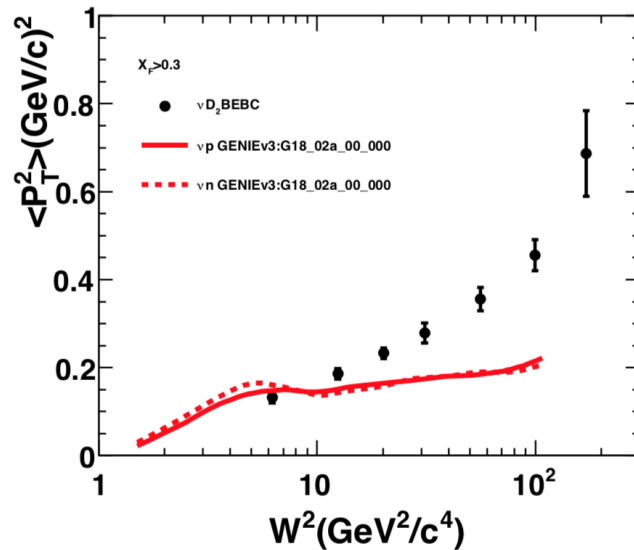
GENIE comparisons with x_F distributions



GENIE comparisons with z distributions



GENIE comparisons with p_T distributions



Empirical low-W model: Main caveats

- Experimental constraint for neutral/charged pion multiplicity

$$\frac{2 \langle n_{\pi^0} \rangle}{\langle n_{\pi^+} \rangle + \langle n_{\pi^-} \rangle} \approx 1$$

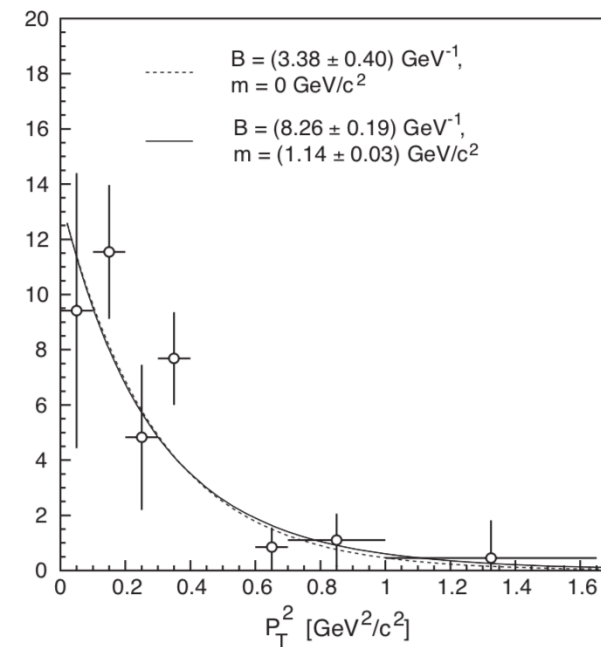
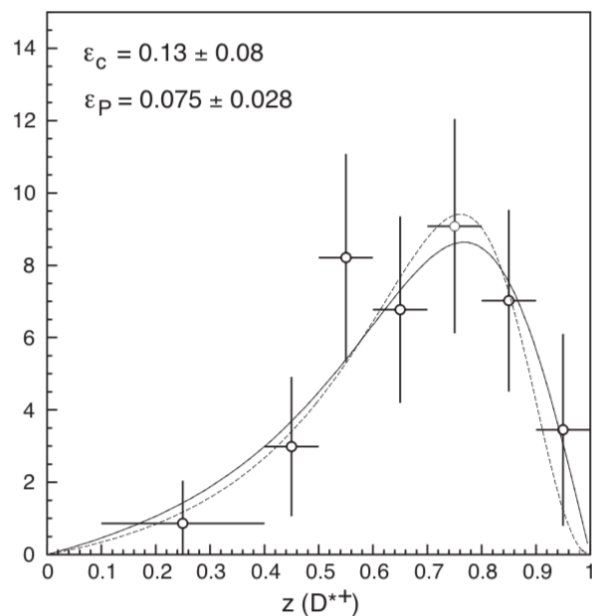
used as constraint on overall neutral/charged multiplicity.

- KNO scaling is applied for the total multiplicity distribution, rather for the charged multiplicity distribution.
- Too strong separation of target and current fragments at low W

DIS charm production

G. De Lellis et al., Phys. Reports 399 (2004)

E_ν (GeV)	f_{D^0}	f_{D^+}	$f_{D_s^+}$	$f_{A_c^+}$
5–20	0.38 ± 0.13	0.05 ± 0.06	0.17 ± 0.09	0.41 ± 0.11
20–40	0.56 ± 0.09	0.09 ± 0.07	0.19 ± 0.07	0.16 ± 0.06
40–80	0.69 ± 0.09	0.19 ± 0.08	0.08 ± 0.07	0.04 ± 0.01
> 80	0.66 ± 0.12	0.26 ± 0.10	0.00 ± 0.06	0.08 ± 0.07
> 40	0.67 ± 0.07	0.23 ± 0.03	0.03 ± 0.01	0.06 ± 0.02
> 30	0.68 ± 0.07	0.19 ± 0.06	0.07 ± 0.05	0.07 ± 0.04
> 20	0.63 ± 0.06	0.18 ± 0.04	0.10 ± 0.04	0.10 ± 0.04
> 5	0.59 ± 0.06	0.14 ± 0.04	0.11 ± 0.04	0.15 ± 0.04



DIS charm production

PYTHIA6 (di)quark assignments for the hadronization of the remnant system:

Init state	Pre-fragm quarks	Selected charm hadron	Final state charm hadron + [PYTHIA pre-fragm inputs]
νp CC	$uu(d \rightarrow)c$	D^0	$c\bar{u} (D^0) + [uu + u]$
νp CC	$uu(d \rightarrow)c$	D^+	$c\bar{d} (D^+) + [uu + d]$
νp CC	$uu(d \rightarrow)c$	D_s^+	$c\bar{s} (D_s^+) + [uu + s]$
νp CC	$uu(d \rightarrow)c$	Λ_c^+	$cud(\Lambda_c^+) + [uu + \bar{u}d]$
νn CC	$ud(d \rightarrow)c$	D^0	$c\bar{u} (D^0) + [ud + u]$
νn CC	$ud(d \rightarrow)c$	D^+	$c\bar{d} (D^+) + [ud + d]$
νn CC	$ud(d \rightarrow)c$	D_s^+	$c\bar{s} (D_s^+) + [ud + s]$
νn CC	$ud(d \rightarrow)c$	Λ_c^+	$cud(\Lambda_c^+) + [ud + \bar{u}d]$
νp CC	$uud + \bar{d}(d \rightarrow)c$	D^0	$c\bar{u} (D^0) + [uud + \bar{d} + u = uu + u]$
νp CC	$uud + \bar{d}(d \rightarrow)c$	D^+	$c\bar{d} (D^+) + [uud + \bar{d} + d = uu + d]$
...
...

In-medium effects to hadronization

- QCD version of the Landau-Pomeranchuk-Migdal effect

The QED effect [A.Migdal,Phys.Rev.103,1811(1956)] was first observed at SLAC/E-146 [P.Anthony et al, Phys.Rev.Lett.75, 1949 (1995)]

Leads to: **Smaller parton cross section within a nuclear medium**

- Cronin effect

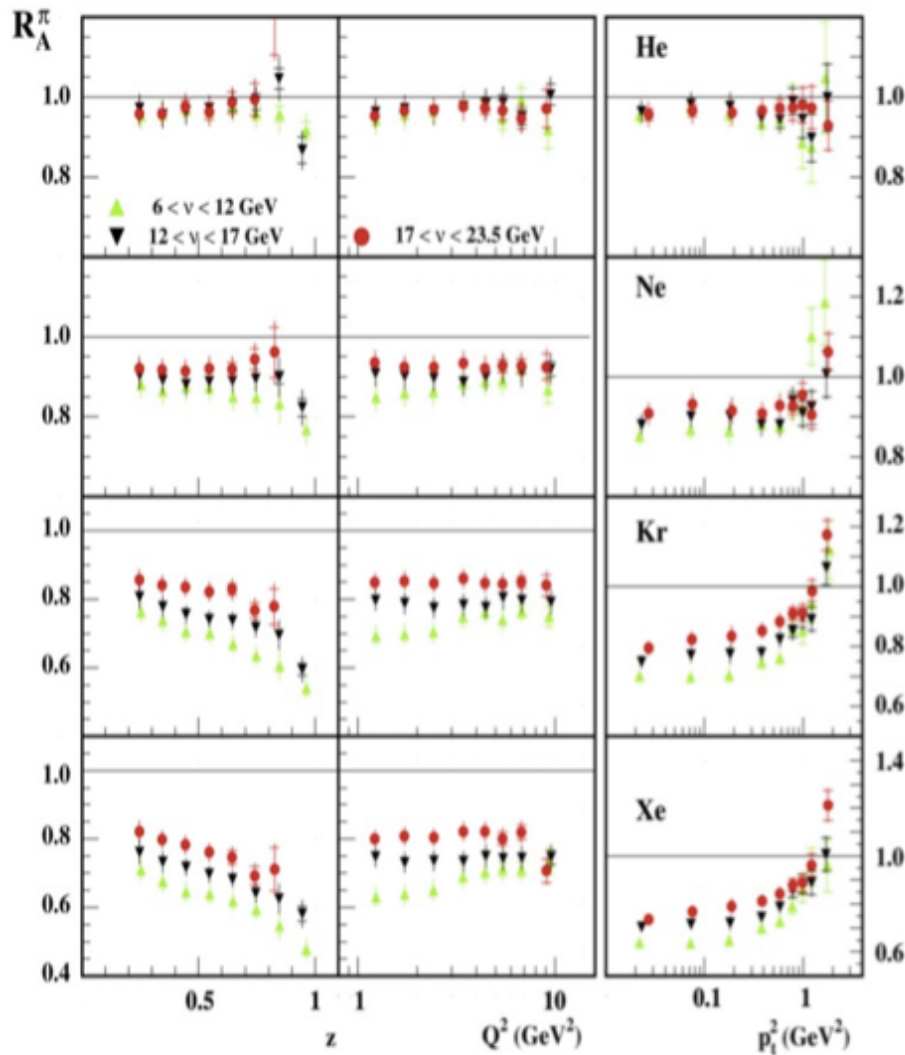
First observed at pA interactions at Fermilab [J.W.Cronin et al.,Phys.Rev.D 11, 3105 (1975)]

Leads to: **Broader pT spectrum**

In-medium effects to hadronization

- First studied using e^- DIS at SLAC
[L.S.Osborne et al., Phys.Rev.Lett. 40, 1624 (1978)]
- Similar experiments with muon beams at CERN (EMC)
[J. Ashman et al, Z.Phys.C52:361-388 (1991)]
- ... and Fermilab (E665)
[M.R.Adams et al, Phys.Rev.D50, 1836 (1994)]
- Recent data (e^- , e^+) by HERMES at DESY
[A.Airapetian et al, Euro.Phys.J.C.20,479 (2001)]
- ... and (e^-) by CLAS at JLAB
[K.Hafidi, hep-ex/0609005]

In-medium effects to hadronization

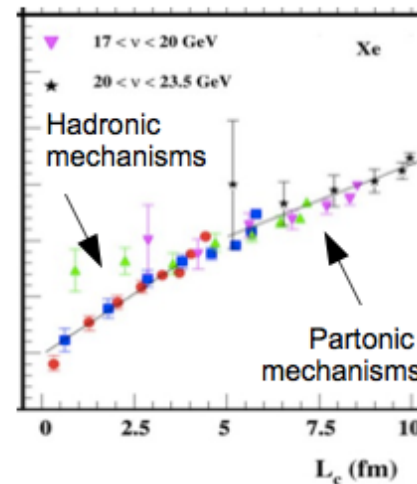


Airapetian et al., Nucl.Phys.B 780 (2007) 1-27,

A number of very interesting, quantitative results of direct interest to v MC

Airapetian et al., Nucl.Phys.B 780 (2007) 1-27, 2007
 Harlem et al., arXiv: 0704.3712

- Attenuation increases with increasing A
- Attenuation decreases with increasing v
- Attenuation increases with increasing z
- Attenuation increases slightly with Q² (pions)
- Attenuation almost independent of p_T
 - except for large p_T's (Cronin effect)

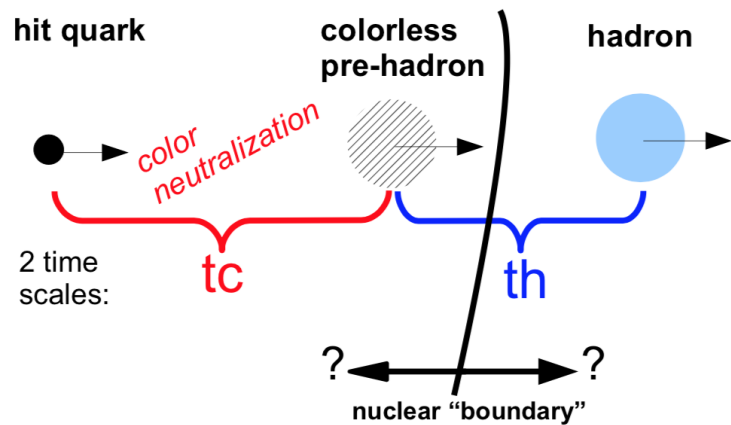


Scaling !

$$L_c = z^{0.35} (1 - \hat{z}) v / \kappa$$

In-medium effects to hadronization

Crucial point:
Hadronization time scales?



Several models:

- B.Andresson et al., Phys.Rep.97, 31 (1983)
- A.Bialas, Acta Phys Pol B 11, 475 (1980)
- M.Gyulassy et al., Nucl. Phys.B 346, 1 (1990)
- J.Cryzewski et al., Z.Phys.C56, 493 (1992)
- N.Akopov et al., Eur.Phys.J C44, 219 (2005)
- X.F.Guo et al., Phys.Rev.Lett.85, 3591 (2000)
- E.Wang et al., Phys.Rev.Lett.89, 162301 (2002)
F.Arleo, Eur.Phys.J C30, 213 (2003)
- A.Accardi et al., Nucl.Phys.A 761, 67 (2005)
- B.Z.Kopeliovich, Nucl.Phys.A740, 211 (2004)
- T.Falter et al., Phys.Rev.C70, 054609 (2004)

In GENIE, in-medium effects to hadronization are included using a single formation time of 0.342 fm/c.

We would welcome any collaborative effort to go beyond that!

Next steps for hadronization modelling in GENIE

- Showed lots of data and empirical models with lots of parameters!
- GENIE now has a very capable **global analysis of neutrino data**.
(Exercised in a retune of the free-nucleon χ -section model, and in a global fit of nuclear 0π and 1π χ -section data)
 - Relying GENIE/Comparisons product and its interfaces with Professor.
 - **Reduces computational complexity of brute-force tuning** and allows for **massive parallelisation**.
 - Can handle a **broad range of (non-reweightable) model uncertainties**.
- Understanding systematics of data shown on this talk and incorporating them in our global analysis is one of our next priorities.
- Obtain: Improved model parameters, and a **handle on model systematics** / data-driven **parameter correlations**.

Summary

- Hadronization model in GENIE
 - Simple and with several caveats, but...
 - As evidenced from data/MC comparisons, it is quite robust
 - Get's right lots of hadronic shower features
- Would love to hear inputs for the improvement of the model
- GENIE 3 was released earlier this week
 - First major new release in more than a decade!
 - An overhaul - touched all modelling aspects (except hadronization!)
 - Implemented a powerful new GENIE global analysis/fit of scattering data (and used it to inform new GENIE tunes)
- Hadronization modelling and tuning become next GENIE priorities.