

Work supported by the Swiss State Secretariat for Education, Research and Innovation SERI



Electron cloud

L. Mether

With contributions from: D. Astapovych, I. Bellafont, G. Iadarola, E. Wulff



FCC Week 2019

Brussels

23 - 28 June

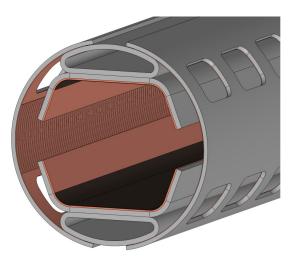


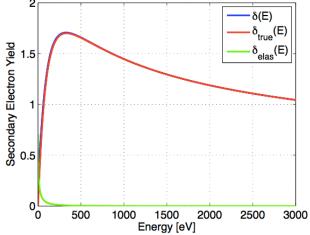
Electron cloud studies

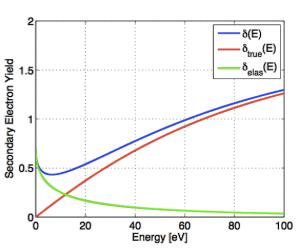


A baseline scenario for electron cloud mitigation (for CDR) has been identified based on

- o Electron cloud build-up simulations in arc dipoles, quadrupoles and drifts
 - » Full beam screen geometry
 - » Trains of 80 bunches separated by 17 empty slots for the 25 ns beam
 - » Secondary electron emission model for Cu co-laminated beam screen
- Estimations of the threshold electron density for single bunch instabilities
- Build-up studies with photoelectrons, using photoelectron distributions based on ray tracing simulations and measurements









Single bunch stability



The threshold central electron density for instability has been estimated analytically:

K. Ohmi et al

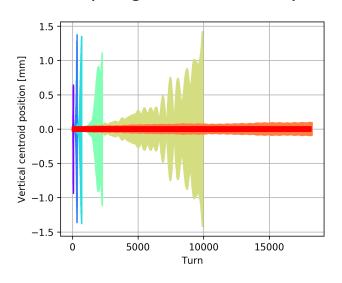
$$\rho_{e,th} = \frac{2\gamma \nu_s \omega_e \sigma_z / c}{\sqrt{3} K Q r_0 \beta L}$$

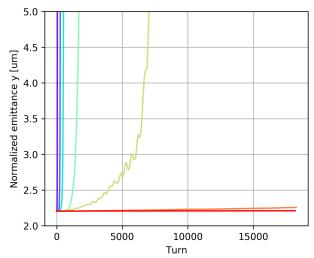
$$\rho_{e,th} = \frac{2\gamma \nu_s \omega_e \sigma_z/c}{\sqrt{3} K Q r_0 \beta L} \qquad \text{with} \quad \omega_e = \sqrt{\frac{\lambda_p r_e c^2}{\sigma_y (\sigma_x + \sigma_y)}} \; , \qquad K = \omega_e \sigma_z/c \\ Q = \min(\omega_e \sigma_z/c, 7)$$

$$K = \omega_e \sigma_z / c$$

$$Q = \min(\omega_e \sigma_z / c, 7)$$

As well as by single-bunch stability simulations in dipoles:





	e- density								
	30.0e11	-	4.0e11						
	10.0e11	-	2.5e11						
	7.0e11	-	1.0e11						
۱	5.5e11								

The estimated density thresholds are independent of the bunch spacing for fixed beam brightness

Threshold electron density

3.3 TeV	50 TeV
1 x 10 ¹¹ m ⁻³	5 x 10 ¹¹ m ⁻³

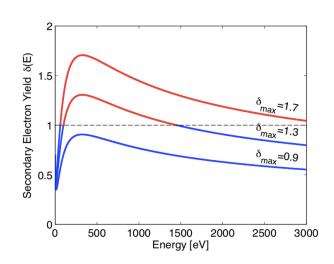


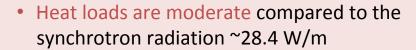
Electron cloud with 25 ns beam



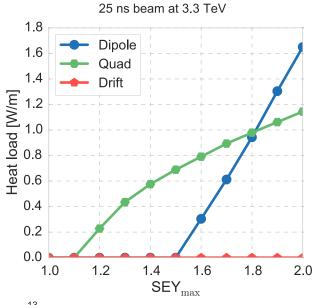
Multipacting thresholds, i.e. maximum SEY (δ_{max}) without e-cloud build-up

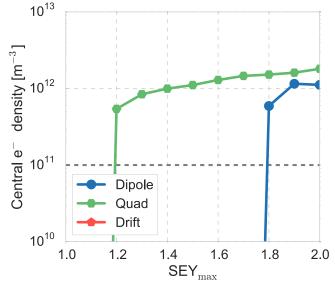
	25 ns			
E [TeV]	3.3	50		
Dipole	1.5	1.5		
Quad	1.1	1.2		
Drift	≥ 2.0	≥ 2.0		





 Stability is more critical: in quadrupoles, the central densities exceeds the instability threshold above the multipacting threshold



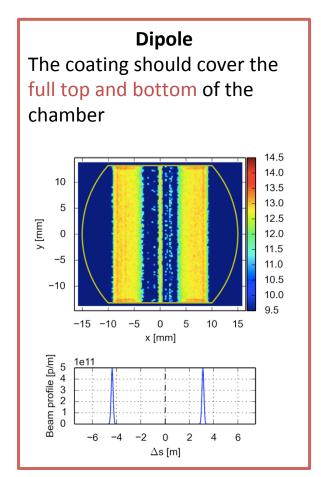


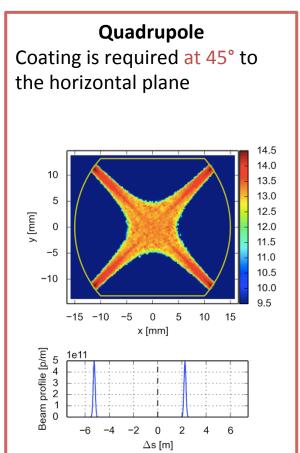


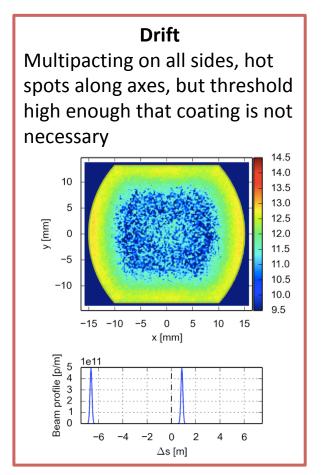
Mitigation scenario for 25 ns beam



- Ensure e-cloud suppression with amorphous carbon coating in dipoles and quadrupoles
 » SEY ≤ 1.1
- No coating needed in drifts due to high multipacting threshold









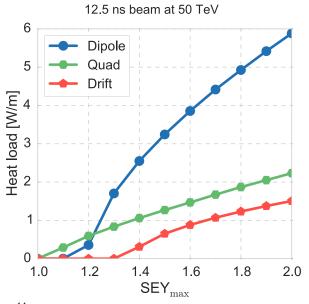
Electron cloud with alternative beams

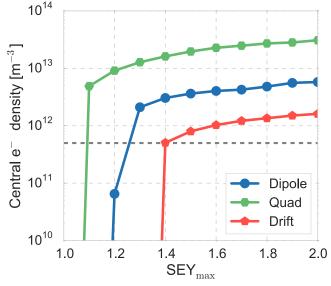


Multipacting thresholds, i.e. maximum SEY (δ_{max}) without e-cloud build-up

	25 ns		12.5 ns		5 ns	
E [TeV]	3.3	50	3.3	50	3.3	50
Dipole	1.5	1.5	1.1	1.1	1.5	1.5
Quad	1.1	1.2	1.0	1.0	1.1	1.0
Drift	≥ 2.0	≥ 2.0	1.3	1.3	1.6	1.6

- Tighter constraints on the SEY for alternative beams, in particular for 12.5 ns
- Heat loads can be significant for unconditioned surfaces
- Stability is critical: the central densities exceed the instability threshold above the multipacting threshold



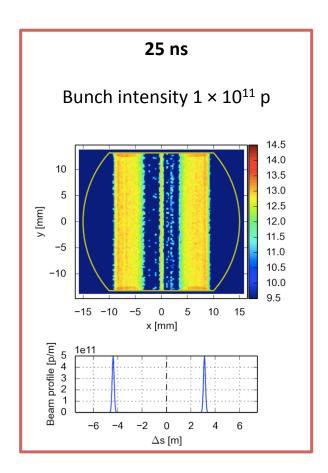


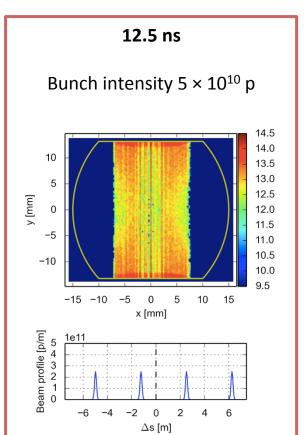


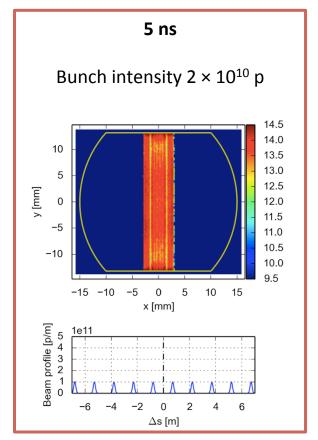
Mitigation scenario for alternative beams



- LASE surface treatments an option for decreasing the SEY beyond a-C
 - » Within constraints imposed by the impedance
- Coating or other mitigation scheme to be considered also in drifts









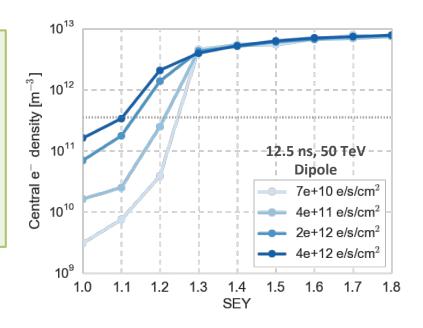
Photoelectrons



Preliminary studies showed that a high photoelectron yield could bring the central electron density above the instability threshold below the multipacting threshold

Constraints on the photoelectron flux were conveyed to the beam screen design team

The current beam screen with sawtooth was found to be compatible in dipoles and quadrupoles



Build-up studies with photoelectron flux distributions based on ray-tracing studies were implemented in the build-up simulations and results presented at the FCC Week 2018

- Confirmed the viability of the current beam screen design in dipoles and quadrupoles
- Uncertainty on amount of photoemission in drifts



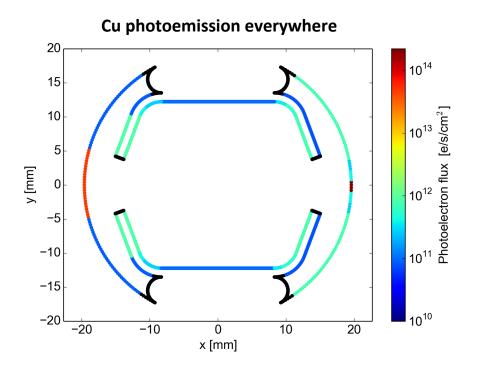
Photoelectron yield in arc magnets



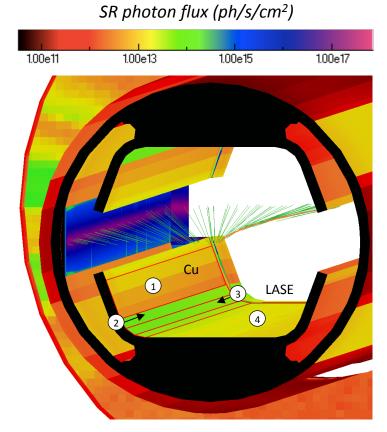
New more detailed simulations with photoelectrons based on ray-tracing simulations as well as photoelectron yield measurements on Cu and LASE surfaces

Two options for dipoles and quadrupoles:

Cu photoelectron yield everywhere



I. Bellafont



Pessimistic estimation, data for the most irradiated magnet. If using LASE on sawtooth, electron generation can be further reduced in the inner chamber



Photoelectron yield in arc magnets

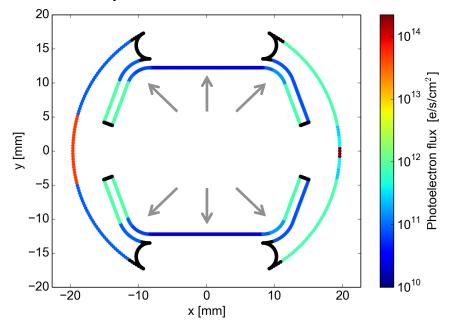


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Two options for dipoles and quadrupoles:

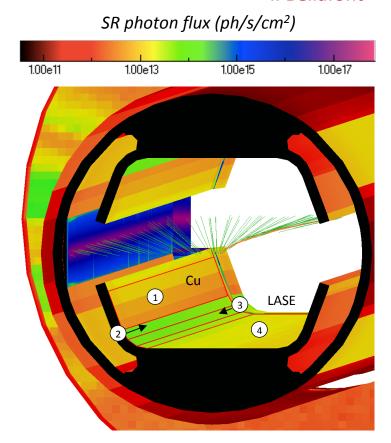
- Cu photoelectron yield everywhere
- LASE photoelectron yield in critical areas

LASE photoemission in critical areas



NB! Still using SEY model based on Cu

I. Bellafont



Pessimistic estimation, data for the most irradiated magnet. If using LASE on sawtooth, electron generation can be further reduced in the inner chamber



Effect of photoelectrons in arc magnets

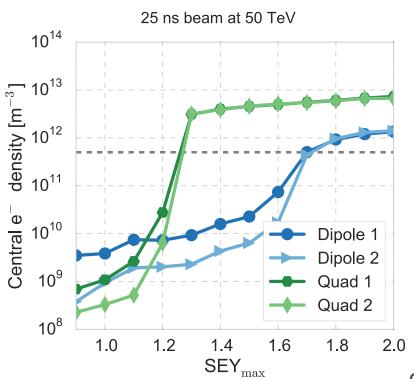


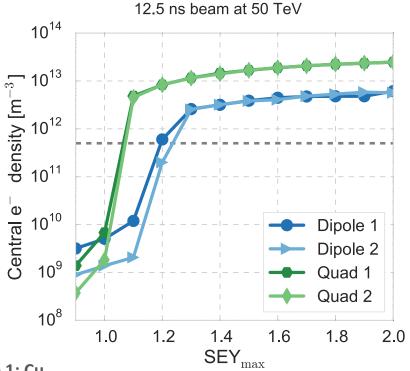
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Two options for dipoles and quadrupoles:

- Cu photoelectron yield everywhere
- LASE photoelectron yield in critical areas

No risk of beam instability below the multipacting threshold due to photoelectrons for any beam option





Case 1: Cu Case 2: LASE

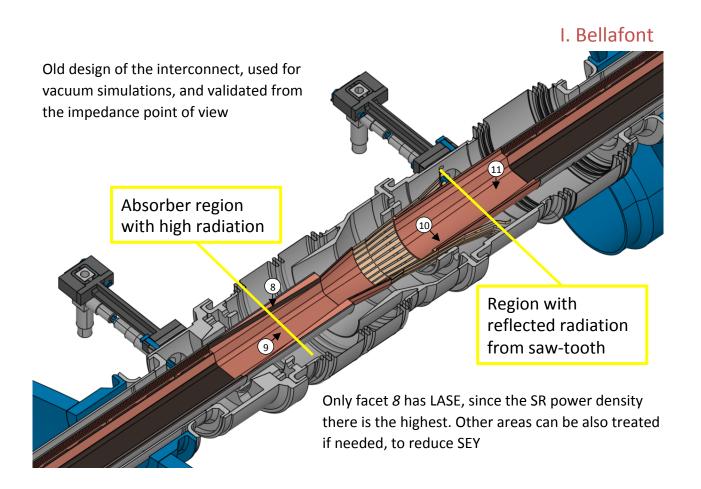


Photoelectron yield in drifts



New more detailed simulations with photoelectrons based on ray-tracing simulations as well as photoelectron yield measurements on Cu and LASE surfaces (WP4)

Two distinct areas are considered in the interconnections



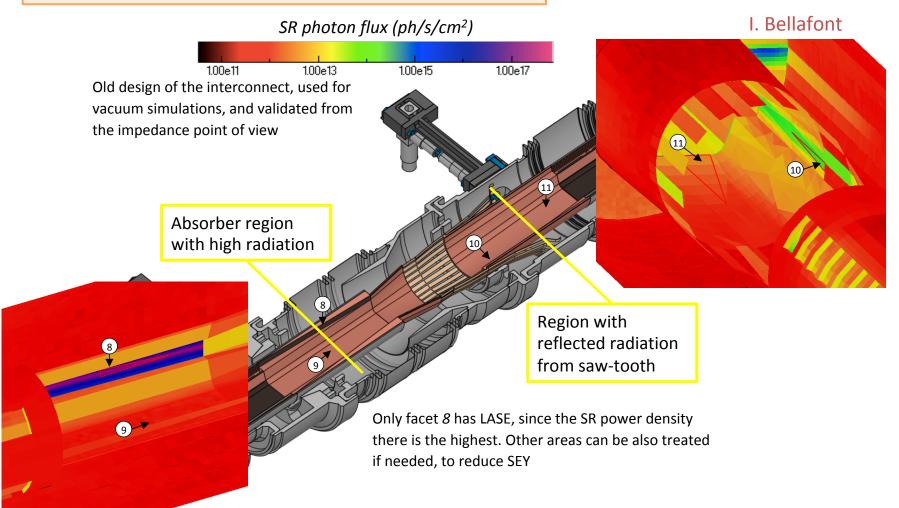


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Effect of photoelectrons in drifts

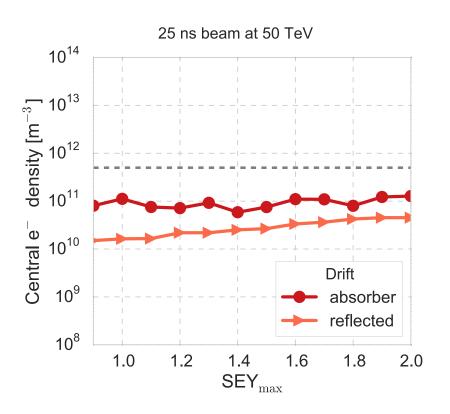


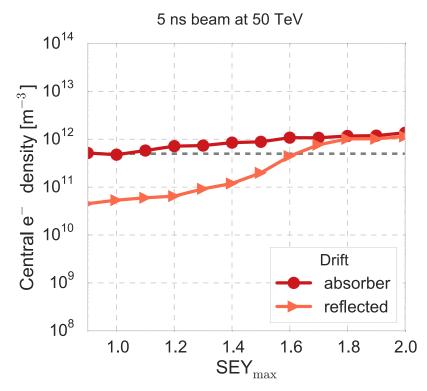
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Two distinct areas are considered in the interconnections

Each region is assumed to cover 1/3 of the interconnection length

For the 5 ns beam the electron density in the absorber region may cause instabilities





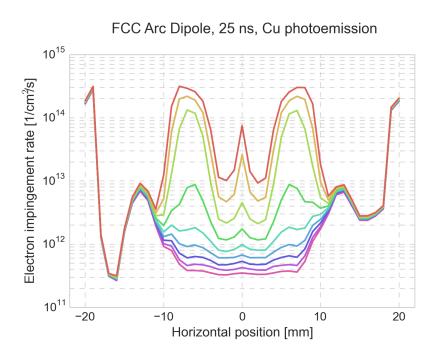


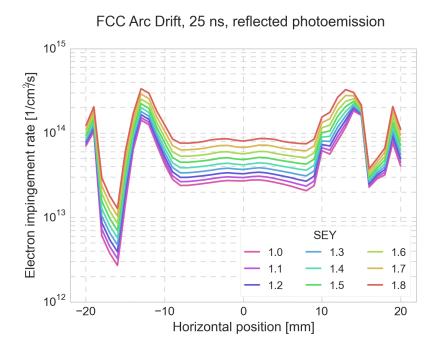
Electron impingement rates



The simulations with photoelectrons have also been used to estimate electron impingement rates and the average energies of impinging electrons

→ Used for studying the effect of electron stimulated desorption on the vacuum (WP4)





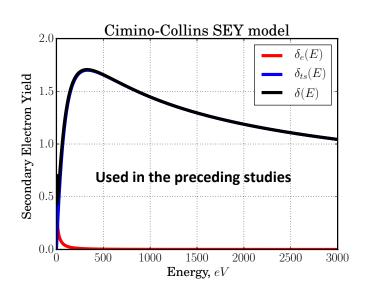
The electron impingement rate in the drifts is high, even for low SEY

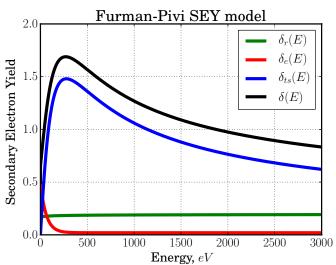
The effect of residual magnetic fields (not considered here) may alleviate the problem



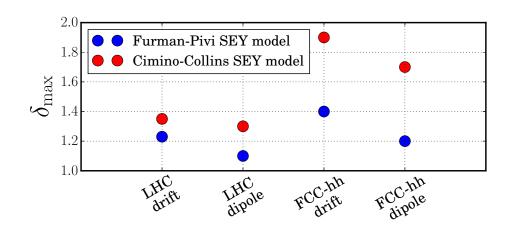


A comparison of two different SEY models for Cu surfaces was presented in Amsterdam





Work by D. Astapovych, using openECLOUD code



Significant differences in the multipacting thresholds with the two models were found

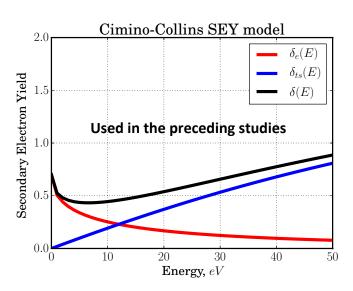
Possible causes:

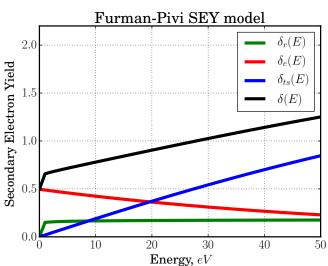
 \circ Shape of total SEY curve for given δ_{max}



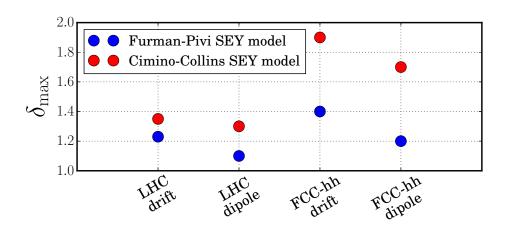


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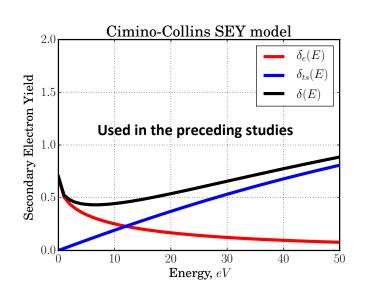
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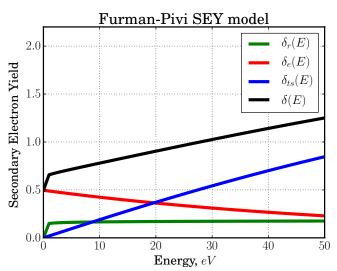
 \circ Shape of total SEY curve for given δ_{max}



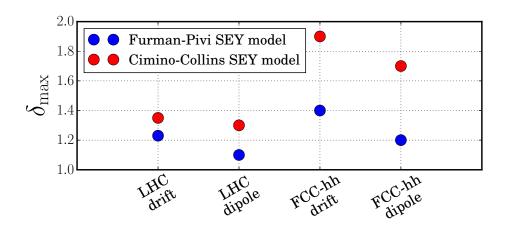


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Significant differences in the multipacting thresholds with the two models were found

Possible causes:

- \circ Shape of total SEY curve for given $\delta_{\sf max}$
- Energy spectrum of emitted electrons for rediffused component
- Numerical details e.g. representation of emitted electrons

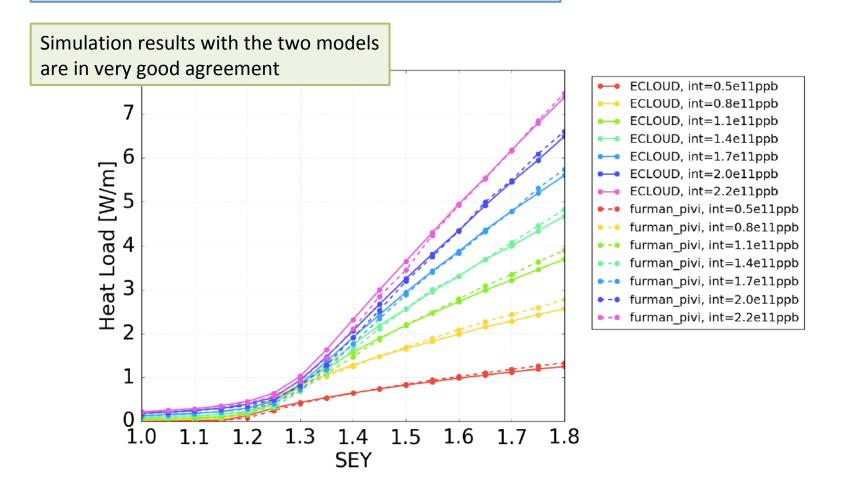




The Furman-Pivi model has recently been implemented also in the PyECLOUD code

Identified parameters for the F-P model that reproduce as well as possible both the SEY curve and the emission energy spectra of the C-C model (no rediffused electrons)

Work by E. Wulff, G. Iadarola







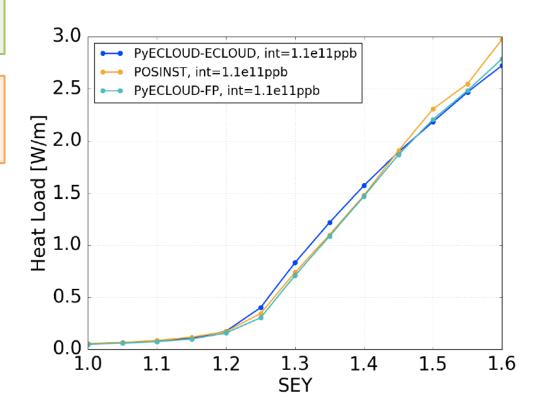
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Simulation results with the two models are in very good agreement

A very good agreements is seen also with the Furman-Pivi model using the original POSINST code







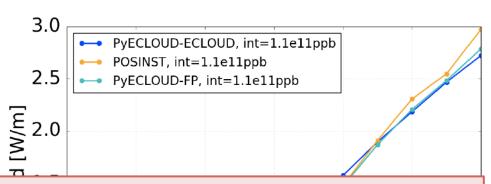
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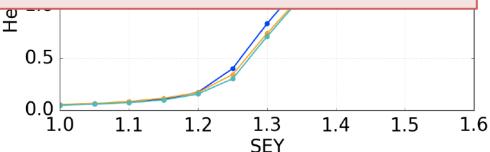
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→ We can be confident that if we use the same physical model, we get the same results regardless of SEY model parameterisation or simulation tool







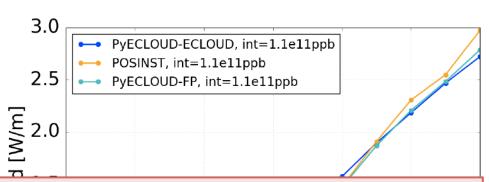
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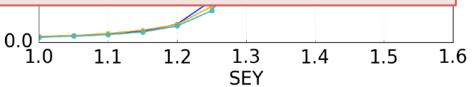
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- → We can be confident that if we use the same physical model, we get the same results regardless of SEY model parameterisation or simulation tool
- → The observed differences between the two SEY models must be due to the different physical quantities, i.e. total SEY curve and emission energy spectra, which can be measured







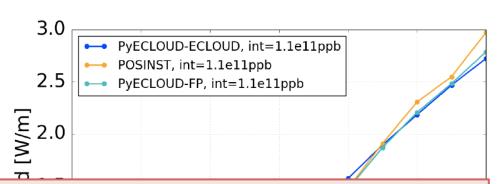
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- → The observed differences between the two SEY models must be due to the different physical quantities, i.e. total SEY curve and emission energy spectra, which can be measured
- \rightarrow The F-P parameters used in the original comparison are based on measurements of a Cu surface for a single δ_{max} , the amount of rediffused electrons is kept fixed when δ_{max} is changed not clear if this is supported by measurements

6



Conclusions



- A baseline scenario for electron cloud mitigation has been identified
 - » Amorphous carbon coating in dipoles and quadrupoles for the 25 ns beam
 - » LASE coating in quadrupoles (and dipoles) and mitigation in drifts may be needed for the alternative beams → impedance reduction of LASE still to be studied
- The effect of photoelectrons has been carefully considered
 - » Saw-tooth at synchrotron radiation impact point provides sufficient mitigation
 - » High electron densities are predicted in the region of the radiation absorbers
 - → Potentially an issue for vacuum quality and for the 5 ns beam also stability
 - → More detailed studies considering e.g. residual magnetic fields may improve prospects
- Different physical SEY models give significantly different predictions
 - » Need to take care to base our simulation models on extensive and relevant measurements
 - → Concerns also e.g. for Cu vs LASE







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