

# Tile Multiple Readout and Beyond

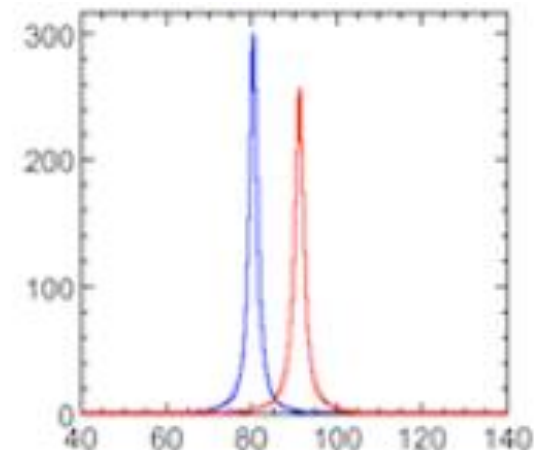
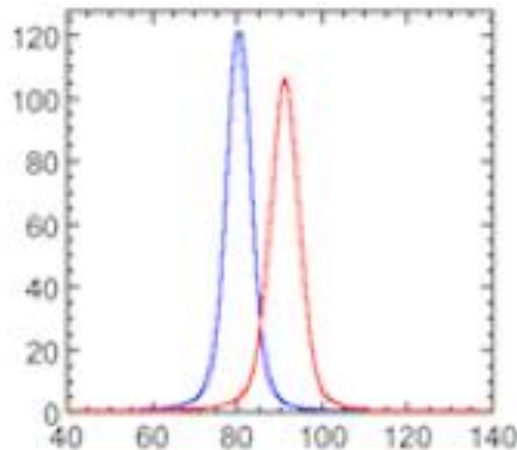
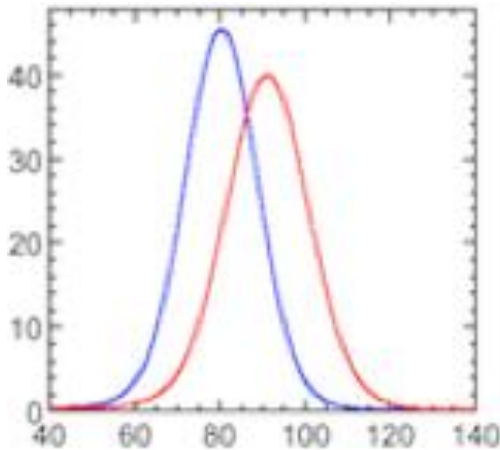
Yasar Onel - Iowa

David R Winn – Fairfield



**FCC** Week **2019**  
**BRUSSELS** 24-28 June

- - **Jet-Jet masses:** Goal for future experiments: SM  $Z \rightarrow \text{jetjet}$ ;  $W \rightarrow \text{jetjet}$
- Ratio  $W, Z \rightarrow jj$  to  $W, Z \rightarrow \text{leptons}$   $\sim 6-7$
- Reconstruct AND Separate (+SM,  $E_{\text{Tmiss}}$ , jet tags, V-V scattering, BSM,  $W'$ ,  $Z'$ ...)
- Separation of W from Z:  $\sigma_{E_{\text{jet}}}/E_{\text{jet}} \sim 3\%$  necessary at 100 GeV, with typical single particle energies  $\sim 10$  GeV [*ASIDE: during collision crossing times which may be as small as  $\sim 10$ 's of ns, pileup events  $\sim 200/\text{crossing}$  and raddam exceeding 50-100MRad.] A 3%-4% jet energy resolution from 50-500 GeV gives  $2.6-2.3\sigma$  W/Z separation. J-J mass resolution is very important In searches for heavy  $W'/Z'$ , vector boson scattering, triple VVV....*



- W/Z  $\rightarrow$  jet-jet separation: **Left** - calorimeter  $\sigma_E/E = 60\%/\sqrt{E}$ ; **Middle**  $\sigma_E/E = 22\%/\sqrt{E}$  (3% @ 50 GeV)  $\sim 2.6\sigma$  separation; **Right** - perfect resolution:  $\sim 4.5\sigma$  separation.

# Dual Readout: Cerenkov Compensation

## First form of multiple readout

- 1<sup>st</sup> Quantitative MC Study:
  - GEANT MC in 1988 [*Compensating Hadron Calorimeters with Cerenkov Light, D.R. Winn, W.Worstell, IEEE Trans. Nuc. Sci.,V1 NS-36, 334(1989)*]
    - Idea: Use differences in response to e-m fluctuations between
      - Cerenkov Medium (transparent – LAr, H<sub>2</sub>O, SiO<sub>2</sub> ....) vs
      - Ionization Medium (scintillator, LAr ion collection,..)to reduce hadron shower fluctuations and make e/h → 1
- DREAM Collaboration/Richard Wigmans et al.
  - Excellent progress in real tests!
  - Thorough Analysis of Dual Parallel Fiber Calorimeters!
- MC: 18%/√E seems possible....But DREAM: 30%/√E:  
Parallel Scintillating+Q Fibers



# Parallel Fiber Deficits - 1

- **1. Constant Term—unavoidable issue** – scintillator light attenuation in  $\sim 2+$  m fibers.
- **2. Pointing/Projective Geometry problematic** in a practical parallel fiber calorimeter over a substantial solid angle. The mechanics + fiber packing of fully projective  $(\theta, \varphi)$  very difficult for (pitch, yaw) more than  $\sim 5^\circ$ . Streaming down fiber holes lowered the resolution in DREAM, even at a  $2^\circ$  pitch. Packing extra fibers from the back or conical fibers:  $\rightarrow$ Constant term,  $\rightarrow$ Calibration Issues
- **3. Scintillator Fiber & Photodetector Raddam:** At present, there are no good examples of scintillator fibers which have proven sufficient raddam resistance or speed to be useful for hadron calorimetry at many future colliders or high flux.
- **4. Fiber Bundle & Photodetector Punchthrough:** Huge fiber bundles,  $>33\%$  of the back of the fiber dual calorimeter area, are directly behind the calorimeter. Large punchthrough backgrounds are generated by these fibers, photodetectors ( $\sim 1/800$  incident  $\pi/K$  quasi-elastic scatter through a  $10 L_{\text{int}}$  calorimeter).



# Parallel Fiber Deficits - 2

- **4. E-M and Hadronic Components of Incident Jets:** Parallel fibers: ~no ability to detect + separate *incident direct e-m component inside of a jet, since there is no longitudinal segmentation.*
- **5. High Resolution EM Front End.** The parallel fiber dual readout jet calorimeter: ~no ability to make a *compensated high-Z high sampling EM front end.*
- **6. Calibration:** Parallel fiber geometry difficult to calibrate, as radiation damage & attenuation varies w/ length. (Contrast w/ longitudinally segmented calorimeters)
- **7. Timing & Pileup:** Longitudinal fibers store the information of jet/em showers: the signal is over the time for the light to traverse the fibers. The light generated *at the back of the calorimeter arrives at the photodetector first.* Thus Fiber calorimeters measure the falling edge of the shower, a less precise measurement

- **8. Longitudinal Segmentation:** Fiber dual readout is incompatible with true longitudinal segmentation, even with waveform electronics, and cannot be easily rebuilt for front readout or implement 4,5 above.
- **9. Radiation Damage: No ability to Repair front end damage.**
- **10. Cerenkov Fiber Index of Refraction:** High Radiation Resistant Cerenkov fibers are limited to quartz, with  $n=1.46$  yield an  $h/e_c \sim 0.25-0.20$  – limiting resolution. Lower index  $n < 1.4$  fibers yielding a lower  $h/e_c$  ratio are not conveniently available (Ex: silica aerogels, Teflon AF, Siloxanes, fluoride glass)
- **11. Cost:** the cost of tiles is significantly less per mass or volume of sensitive material than that of fibers, and the cost of a fabricated tile absorber matrix is considerably less than the parallel fiber Swiss cheese.

## **12. No Particle Flow/Energy Flow Calorimetry:**

**Parallel Fibers are incompatible with high granularity** - improving jet  $\delta\theta/\theta$ , core ID of jets, isolation/ID of leptons/photons in jets and pileup, and neutral particle ( $K^0$ ,  $n$ ) ID, especially under pileup. **Tile readout:** fully compatible with highly granular calorimetry, easily added to particle flow calorimeters

## **13. No Other Sensors for Dual Readout, Triple or Multiple Readouts:**

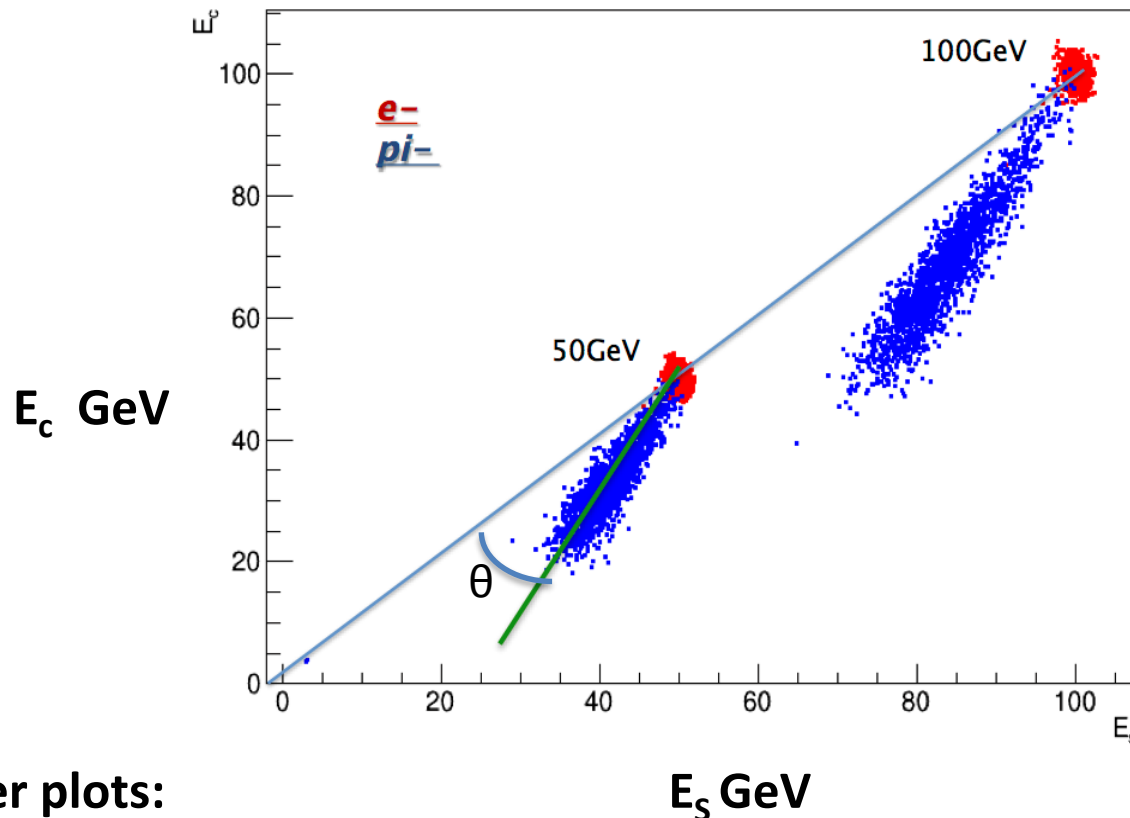
*Parallel fibers cannot use other sensors which could further separate e-m and hadronic components*

Ionization detectors

- Solids – Si, Diamond, GaAs,..;
- Liquids- LArgon, Liq. Scintillators
- Gasses – micromegas, TRD
- $\beta \rightarrow 1$  sensitive detectors such TRD, or ultra-low-index materials(aerogels  $n \sim 1.1$ ,  $MgF_2$ , water  $n \sim 1.33$ , perfluoro-, silicones,..);
- Secondary Emission sensors with higher response to slow particles  $\beta \rightarrow 0$  and minimal response to minimum ionizing energy (new large MCP);
- Inorganic non-hydrogenous scintillators (LYSO,  $PbWO_4$ , ZnO:Ga et al.),
- Neutron-Enhanced:  $^6Li$ ,  $^{10}B$ ,  $^3He$ , fissionable... containing materials.

- GEANT4 MC on a simple tile calorimeter: 0.5 cm thick each of quartz, plastic scintillator, and Cu absorber tiles.
- Two energies (50, 100 GeV) each of 1000 electrons (red dots) and of ~800-1000 pions (blue dots) were sent into the 50x50 cm calorimeter, 12.2 Lint deep (8 Lint of Cu, 1.5 Lint polystyrene, 2.7 Lint Quartz = 12.2 Lint Length ~ 3.6m)
- $N_{\text{photons}}$  325-650nm generated in the Cerenkov and in the scintillator tiles were counted. 0.5% at random were assigned as converted to p.e.
- Scintillator photons ~120x Cerenkov photons; photostatistics not limiting factors.
- Means of histograms of the electron shower p.e. in quartz and in scintillator were used to convert/normalize the number of collected p.e. in Cerenkov light and in Scintillator light to normalized energies  $E_{\text{Cerenkov}}$  and  $E_{\text{Scintillator}}$ , and then plotted as a scatter plot of  $E_C$  vs  $E_S$  for each electron.
- Pions of 50, 100 GeV were then simulated, converted to  $E_C$  vs  $E_S$





- **Scatter plots:**

*Electrons:*  $E_c$  vs  $E_s$  (red) lie along line shown schematically as  $E_c = E_s$ .

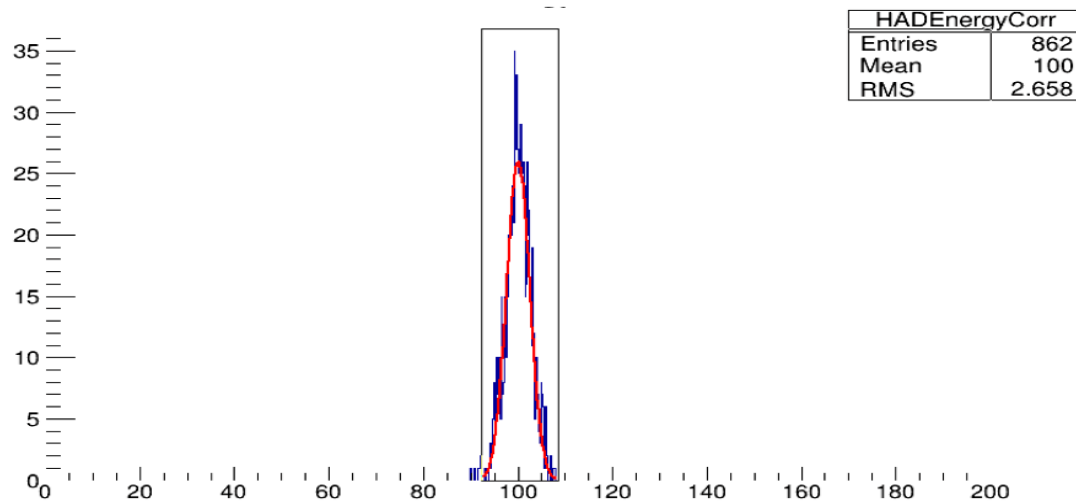
*Pions:*  $E_c$  vs  $E_s$  scatter-plotted (blue) lie mainly below the  $E_c = E_s$  electron line with correlation between  $E_c$  vs  $E_s$  fitted as a line (green, 50 GeV points at an angle  $\theta$ ).

- ***As the shower fluctuates more to hadrons,  $E_c$  falls faster than  $E_s$ .***

- ***A Simple analysis:*** Linear fit to hadron scatter points (Green line), with slope  $R$ , corrects the energy: Project the scatter points as a histogram perpendicular to the linear correlation, the energy distribution becomes Gaussian & narrower.

- Pion Energy  $E$  (*first order*):  $E = E_s + [a \text{ correction term proportional to the difference } (E_s - E_c)]$ :  

$$E = E_s + \alpha(E_s - E_c)$$
with  $\alpha$  given by slope  $R$  as  $R = (1 + \alpha)/\alpha$  or  $\alpha = 1/(1 - R)$ .
  - The angle between the line  $E_c = E_s$  and fitted  $\pi$  scatter plot line:  $\theta = \arctan(R) - \pi/4$ .
- $(E_s - E_c)$  grows as shower fluctuates into nuclear/hadronic energies.
- As slope  $R$  gets steeper, the correction term  $\alpha(E_s - E_c)$  becomes more important. When Cerenkov  $E_c$  is the same as scintillation  $E_s$  (e's or  $\pi$ 's exchange to  $\pi^0$ 's), then  $(E_s - E_c) \sim 0$ ,  $E = E_s = E_c$



***(mean, rms) = (100, 2.66) GeV***

***→  $\sigma E/E$  that enables  $W \rightarrow \text{jet-jet}$  separated from  $Z \rightarrow \text{jet-jet}$***

- **Higher order terms-**  $\alpha_2(E_s - E_c)^2 + \alpha_3(E_s - E_c)^3 + \dots$  and energy dependent  $\alpha_n$  – there is a continuous mapping (vector field) of the points in  $E_c$  vs  $E_s$  space to the line  $E_c = E_s = E$ .

## – Rules of Thumb:

- (0) An intrinsic limit of normal hadron calorimetry:  $\sigma_E/E > 11-13\%/ \sqrt{E}$ , given by the ratio of detectable neutron energy to the fluctuations in lost nuclear binding energy.
- (1) Contrast between  $h_i/e_i$  ( $i$ =ionization) and  $h_c/e_c$  ( $C$ =Cerenkov) for hadrons  $h$  and e-m energy: the ratio of ratios  $[h_i/e_i]/[h_c/e_c] \geq 4$  in order to reach incident hadron energy resolutions below  $30\%/ \sqrt{E}$ , with  $18\%/ \sqrt{E}$  being a reasonable target to achieve using plastic scintillator and low index materials;
- (2)  $h_i/e_i$ : as large as possible  $\rightarrow$  hydrogenous or n-sensitized ionization detection media.
- (3) e-m energy resolution in Cerenkov light  $< 70\%/ \sqrt{E}$  to achieve  $< 20\%/ \sqrt{E}$ ;
- (4) Resolution scales  $\sim \sqrt{f_{\text{sample}}/f_{\text{frequency}}}$ .
- (5) Compensation can be achieved by enhancing neutron (hydrogenous or n-absorbing) or ion fragment sensitivity and/or by suppression electromagnetic component by tuning the absorber thickness relative to sampling media ( $f_{\text{sample}}$  typically  $\sim 1/10$  but at a loss of potential ultimate resolution).

- **Adding sensor tiles** relatively insensitive to MIPs, OR more sensitive to  $\gamma\beta \rightarrow 0$  increases the contrast between e-m and hadronic energy (enhancing the low energy hadronic signal) – one such sensor is Secondary Emission; its signal scales as  $dE/dx$ , with a MIP SE signal  $\sim 100x$  less than that of the energy of the peak signal (peak signal for protons occurs at  $\sim 200\text{KeV}$  -  $n+p \rightarrow p+n$  knock-on protons).
- **Homogeneous non-hydrogenous dense inorganic scintillators** (LYSO,  $\text{PbWO}_4$ ,  $\text{CeF}_3$ )
  - $h_i/e_i \sim 0.4$  and  $h_c/e_c \sim 0.25$ , or  $[h_i/e_i]/[h_c/e_c] \sim 1.6$ :
    - > *Homogeneous calorimeters cannot achieve dual readout compensation better than  $\sim 50\text{-}60\%/ \sqrt{E}$  on hadrons, even with perfect separation between scintillator & Cerenkov light in the homogeneous detector. [Note: LAr/Ch4]*

**Theoretical  $\sim 15\text{-}18\%/ \sqrt{E}$  on jets:** scintillator sensors with  $h_i/e_i \sim 0.6\text{-}0.8$  (likely hydrogenous &  $n$ -sensitive), and Cerenkov sensors with  $h_c/e_c \leq 0.2$  are needed. To achieve  $h_c/e_c < 0.2$ , lower  $n$  (index of refraction) Cerenkov radiators are required (i.e.  $\beta_{\text{thresh}} \rightarrow 1$ ), but require enough photons to achieve an e-m resolution  $< 70\%/ \sqrt{E}(\text{GeV})$  or  $N_{pe} > 2 \text{ pe/GeV}$ .

**Extend E-M Response by higher sensitivity to  $\beta \rightarrow 1$**   
**Results in high contrast ratio with  $h_c/e_c > 0.15$  (i.e.  $e_c/h_c > 6-7$ )**  
**( Lessens low energy Hadron, n, and nuclear fragment Sensitivity)**

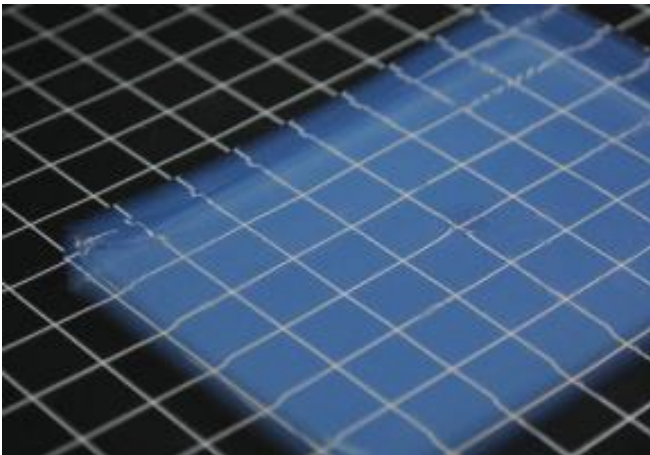
## (A) TRD

- Straw tubes,..... Low mass issue for calorimetry)

## (B) Low index Tiles

**(1.1 < n < 1.35) tiles :**

- silica aerogels (n=1.05-1.3)
- TeflonAF (n=1.29, 12 Mrad) (amorphous form; water-clear)
- polysiloxanes (n=1.35, 100 Mrad)
- $MgF_2$  (1.37);



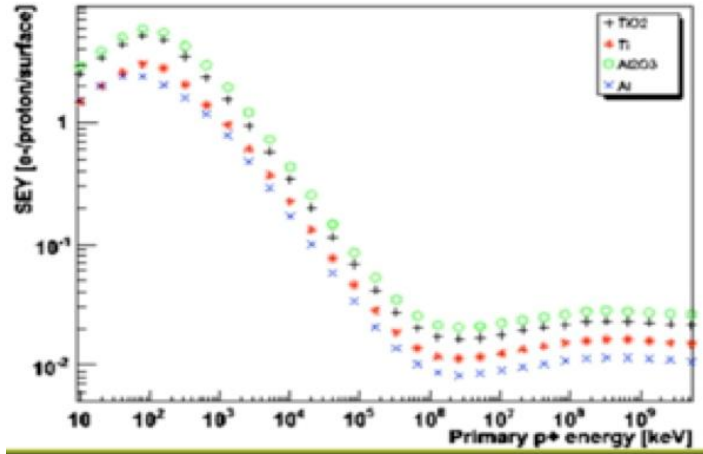


# Beyond: Multiple Particle Flow Readout -2

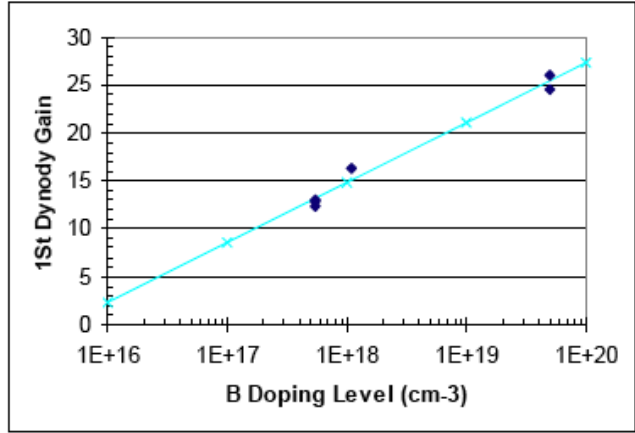
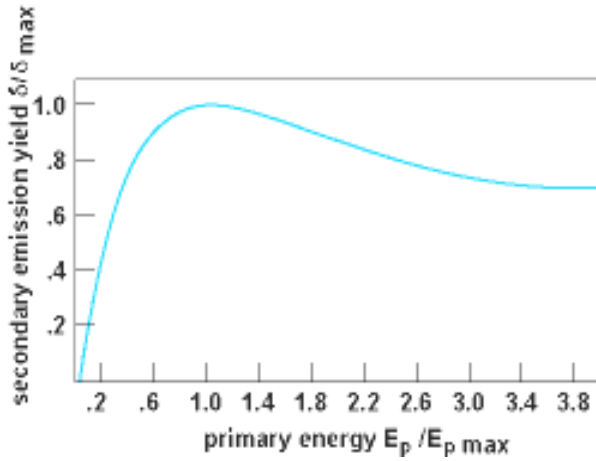
- **Secondary Emission (SE).** Secondary Emission(SE) tiles are more sensitive to  $\gamma\beta \rightarrow 0$  particles than to MIPs - scales as  $dE/dx$ . MIP SE signal  $\sim 100-200x$  less than at peak gb SE signal – **the opposite of Cerenkov light**. SE tiles for correction for heavy fragments, lost neutron energies, slow hadrons. (triple/quadruple readout.) 1-2 MeV alpha particle: secondary yield  $\sim$  at max.

LHC SE Monitors:  
SE Yield vs  
 $E_p$  (KeV)

$10^{20}$  p/cm<sup>2</sup> !



12x12 MCP  
FNAL Test Beam



B-doped  
Nanoxtal  
Diamond  
1 $\mu$ m thick  
On dynode



- **Triple Readout and beyond:** 3 tiles to improve dual readout: non-hydrogenous scintillator, hydrogenous/neutron-sensitive scintillator, 2 indices of Cerenkov tile(s), SE tiles.... Compare less-sensitive neutron scintillators [non-hydrogenous scintillators; inorganic and perfluorocompounds] to more neutron-sensitive H or n-absorbing/converting scintillator tiles.
- **Combined Dual/Triple Readout with Particle Flow:** Add Cerenkov tiles, TRD tiles, SE tiles, or others to existing Particle/Energy flow calorimeter prototypes.
- **Neutron-enhanced detecting scintillator tiles** thin film coatings -  $^{10}\text{B}$ ,  $^6\text{Li}$ , hydrogenous materials [ $^6\text{LiH}$ ] – thin clear film, buffered w/ alumina films; interesting:  $\text{Li}^6\text{B}^{10}\text{H}_4$  which would be transparent if deposited as thin films between clear buffers.  $^{10}\text{B}$  SE yield dynodes.
- **Liquids:** very large homogeneous detectors: LB, cosmic neutrinos or proton decay
  - 1) water “tiles”(n=1.29-1.31 TeflonAF light pipe) + LS tiles – no absorber
  - 2) LArgon drifted ions + Cerenkov light detection. The index n good e/h contrast; scintillation light at 128nm will not penetrate PMT windows.

# Multiple Readout

- Multiple Readout can be tuned for best  $\sigma E/E$ , timing, rate, and radiation resistance
- Multiple Tile Readout can enhance Energy, Intensity, and Cosmic Experiments
- Multiple Tile Readout enables Radiation-Resistance
- Multiple Tile Readout compatible with Energy Flow high granularity calorimetry.
- Multiple Readout can be added to Calice or the CMS endcap HGCal and other existing calorimeters