Logarithmic accuracy of parton showers

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in collaboration with
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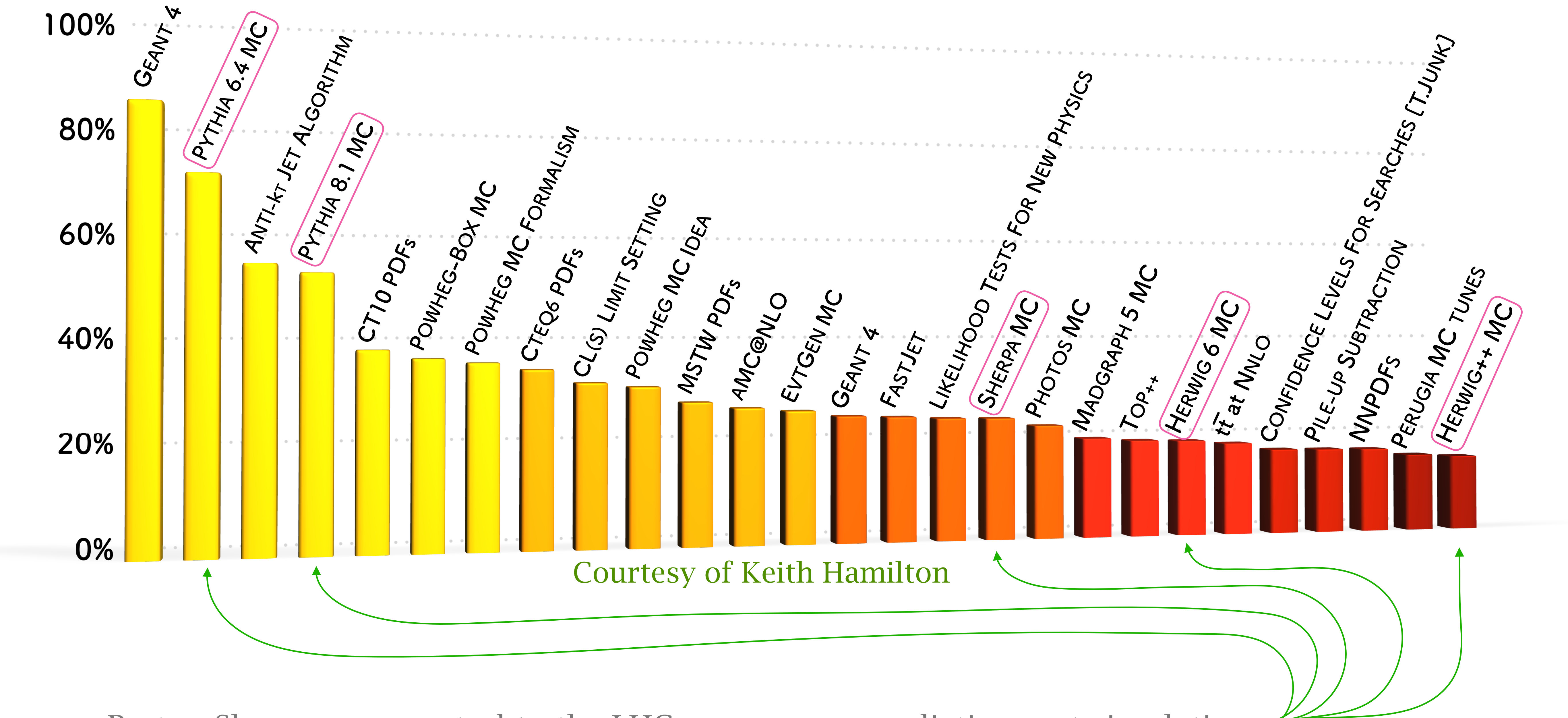
- Percentage of ATLAS+CMS+LHCb papers citing a given article since Jan '14 (w/o self citations)



- Parton Showers are central to the LHC programme: realistic event simulations

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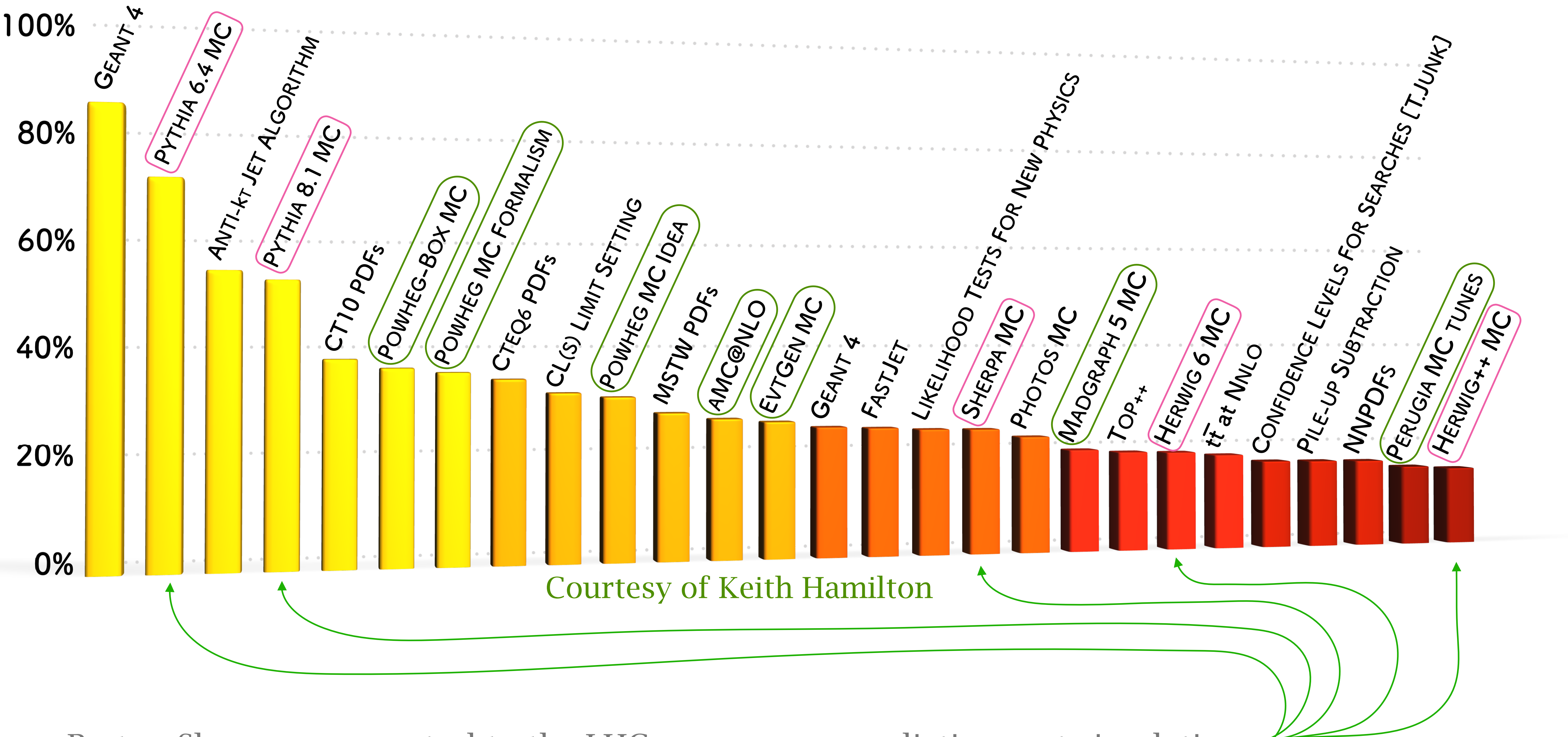

- Parton Showers are central to the LHC programme: realistic event simulations
- Used in essentially all event generators

Resummations vs. Parton Showers

- Both frameworks provide an all-order calculation for collider observables
- Several differences in the way this is formulated
- The higher logarithmic accuracy of current resummations comes with a lower versatility

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Treatment of radiation	• Several simplifications: amplitudes, phase space, observable • All calculations derived in the on-shell/singular limit (only logarithms)	• Radiation is described fully exclusively. Provide full set of final-state momenta • Full momentum conservation necessary (e.g. initial condition for hadronisation)

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Observable dependence	• Tailored to the observable, e.g. global vs. non-global, specific approximations in each case	• A simple shower should be accurate for a broad family of observables at once

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Observable dependence	• Tailored to the observable, e.g. global vs. non-global, specific approximations in each case	• A simple shower should be accurate for a broad family of observables at once
Logarithmic Accuracy	• Higher logarithmic orders achieved thanks to the above simplifications in the formulation	• Currently unknown. The goal of this talk is to initiate a formal study of this point

NLL resummation

- To understand (and ultimately improve) the logarithmic accuracy of PS, crucial to build a systematic connection to resummation

- Use the technology of numerical resummations to approach the problem

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e.g. e+e- -> q qbar + X at NLL
```

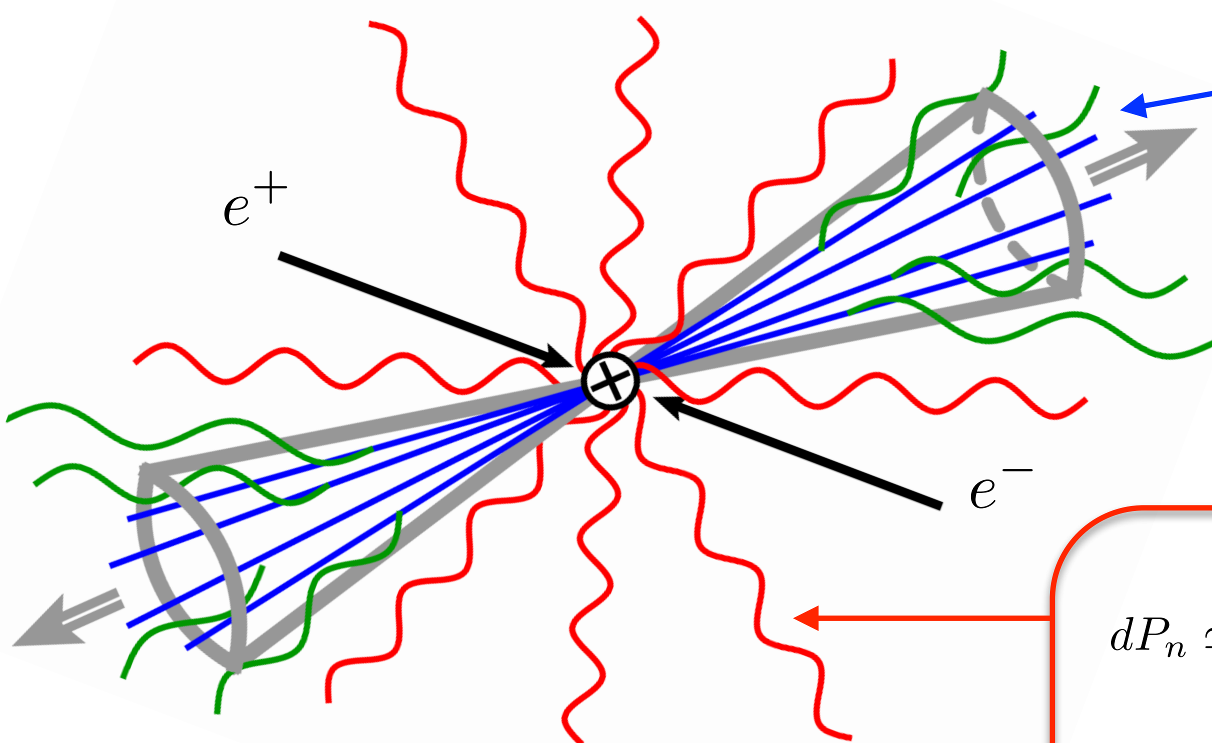

Image by T. Becher et al.

$$dP_n \simeq \frac{C_F^n}{n!}\prod_{i=1}^{n}\left(\frac{\alpha_s^{\rm CMW}(p_{\perp,i}^2)}{\pi}\frac{dp_{\perp,i}}{p_{\perp,i}}dz_i P_{q\to qg}(z_i)\frac{d\phi_i}{2\pi}\right)$$

- collinear limit described by independent emissions strongly ordered in angle

[Catani et al. '91-'93; Banfi, Salam, Zanderighi '01-'04]

$$dP_n \simeq \frac{1}{n!}\prod_{i=1}^{n}\frac{\alpha_s}{\pi}\frac{d\omega_i}{\omega_i}\frac{d^2\Omega}{4\pi}N_c\sum_{\pi_n}\frac{p_1\cdot p_2}{(p_1\cdot k_{i_1})(k_{i_1}\cdot k_{i_2})\dots(k_{i_n}\cdot p_2)}$$

- soft wide angle limit described by a shower of *soft* colour dipoles strongly ordered in energy

[Dasgupta, Salam '01; Banfi, Marchesini, Smye '02]

Parton Showers

- Main defining features (at least for LO showers)

1. Ordering variable: generate emissions in sequence according to a kinematic variable v (e.g. k_t, angle, virtuality).

2. Branching probability: state S_n with n partons at a given v found with a probability $P(S_n, v)$

   ➡ This probability evolves with the ordering variable as

$$\frac{dP(S_n, v)}{d\ln 1/v} = -f(S_n, v)P(S_n, v)$$

This evolution equation accounts for real and virtual corrections (unitarity)

3. Kinematic mapping: state S_{n+1} obtained from a state S_n via a mapping $\mathcal{M}(S_n \to S_{n+1}; v)$

   ➡ Is a function of all partons involved in the branching. It defines how the recoil is absorbed by other partons in the event. E.g. for a *local* recoil scheme

$$S_{n+1} = \mathcal{M}(S_n, v; \underbrace{i, j}_{\text{emitters}}, \underbrace{z, \phi}_{\text{emission}})$$

   ➡ The map is accompanied by the relative probabilities of all possible new states, i.e.

$$f(S_n, v) = \sum_{i,j} \int dv' dz d\phi \, \frac{d\mathcal{P}(S_n, v'; i, j, z, \phi)}{dv' dz d\phi} \delta(\ln v'/v) \qquad \sum_{i,j} d\mathcal{P}(S_n, v; i, j, z, \phi) \simeq \frac{d\Phi_{n+1}}{d\Phi_n} \frac{|M^2(S_{n+1})|}{|M^2(S_n)|}$$

A case study: dipole showers

- Several designs available…

global recoil

local recoil

angular ordering

kt ordering

virtuality ordering

antenna shower

dipole shower

A case study: dipole showers

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- We focus on k_t-ordered dipole showers with local recoil
  - Most common design today
  - Ability to reproduce non-global logarithms at LL, for which different solutions might fail

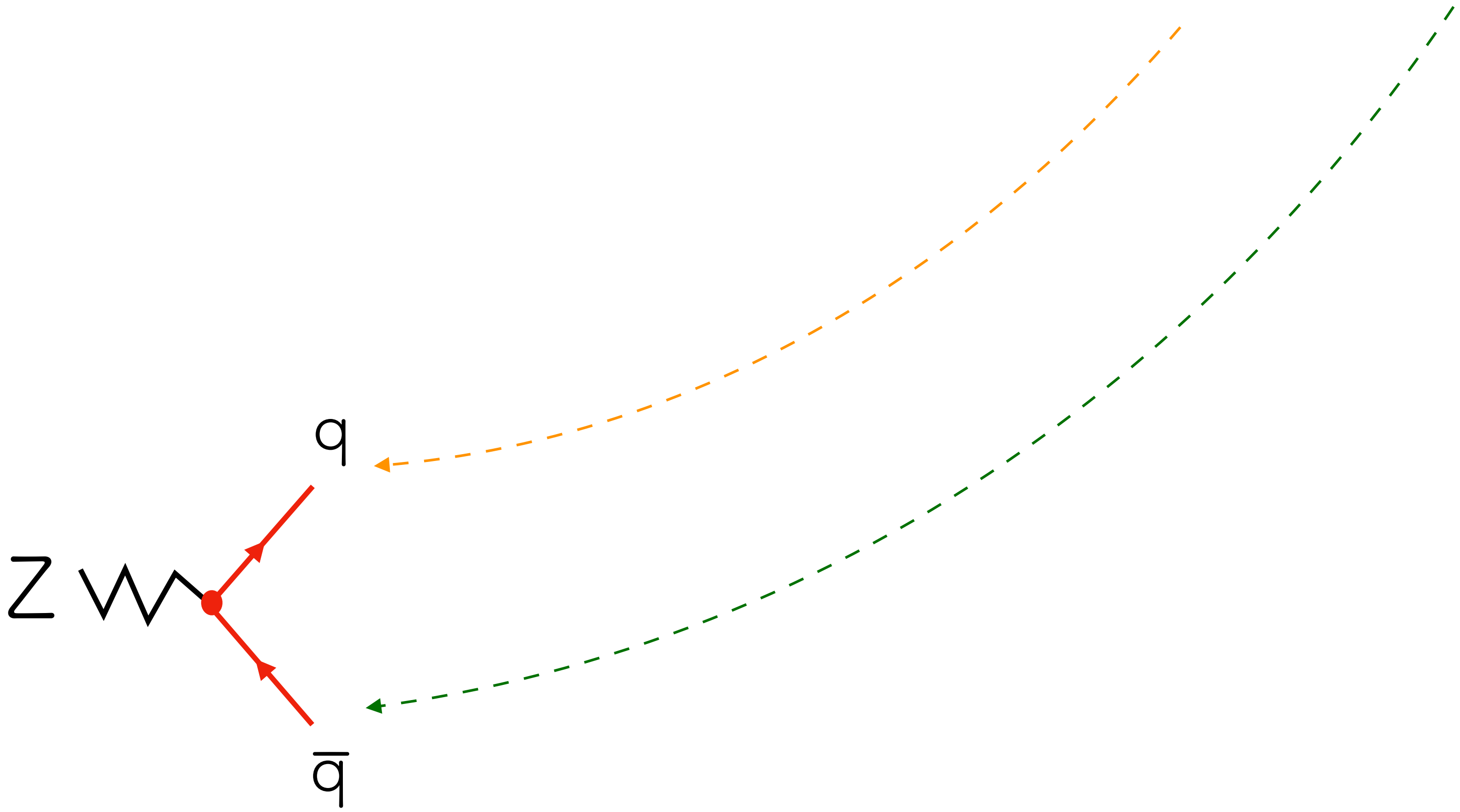
see e.g. [Banfi, Corcella, Dasgupta ’06]

- Consider the designs of `Pythia8`’s shower and `Dire` as a case study

[Sjostrand, Skands ’04]
[Hoeche, Prestel ‘15]

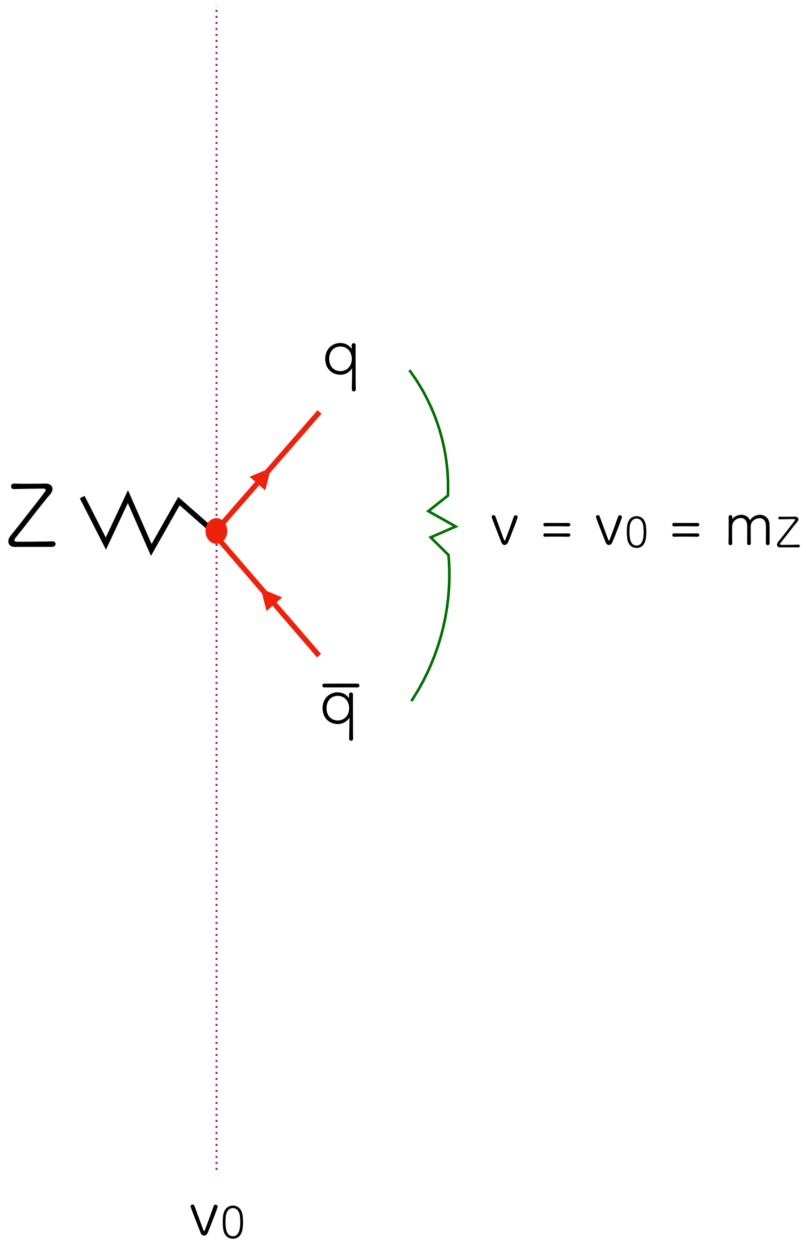
Dipole showers

- Events are viewed throughout as a collection of colour-anticolour dipole ends



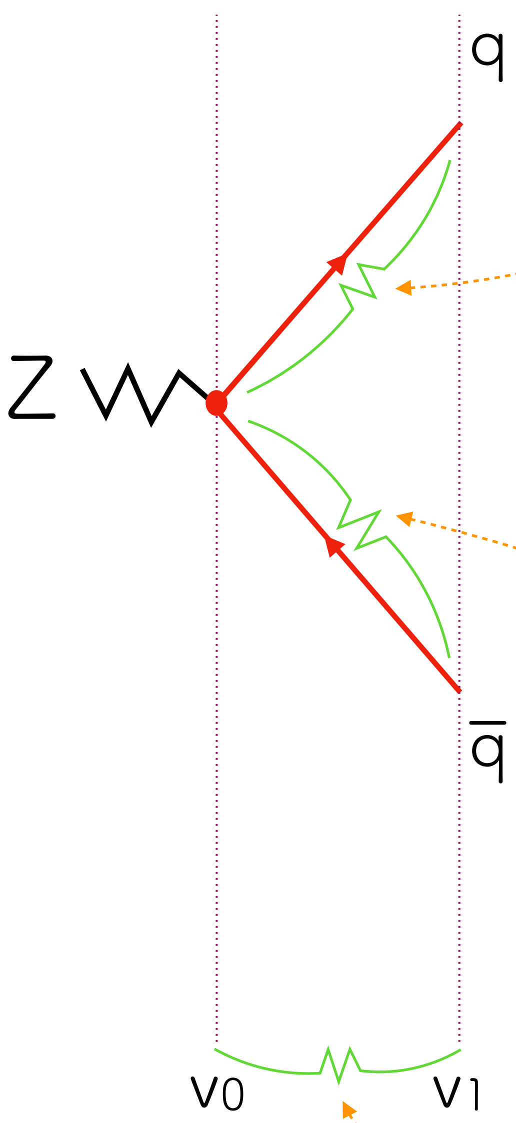
Dipole showers: evolution variable

- Ordering variable v : smallest $p_{\perp}$ separation (resolution) between any pair of partons
- Zooming out to smaller v values more partons get resolved



Dipole showers: branching

- Branching probability: evolution equation solved in terms of a Sudakov form factor



P[q dipole end doesn't emit in $v_0 \rightarrow v_1$]

$$\Delta_q(v_0, v_1) = \exp\left[-\int_{v_1}^{v_0} d\mathcal{P}_{q[\bar{q}]\to qg[\bar{q}]}\right]$$

P[q dipole end doesn't emit in $v_0 \rightarrow v_1$]

$$\Delta_{\bar{q}}(v_0, v_1) = \exp\left[-\int_{v_1}^{v_0} d\mathcal{P}_{\bar{q}[q]\to \bar{q}g[q]}\right]$$

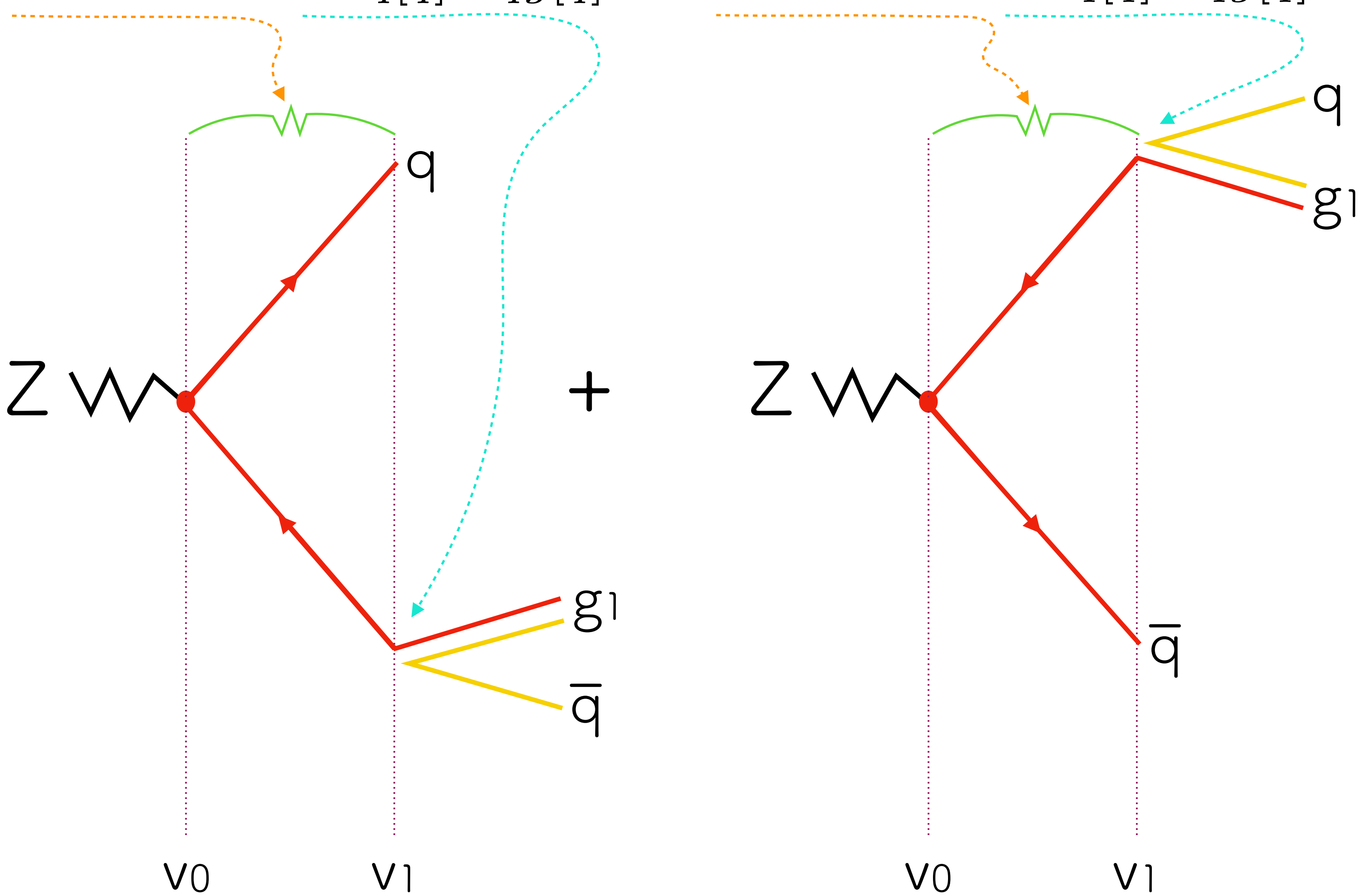
P[nothing in event emits in $v_0 \rightarrow v_1$]

$$\Delta(v_0, v_1) = \prod_{\substack{\text{dipole} \\ \text{ends } \tilde{i}}} \Delta_{\tilde{i}}(v_0, v_1)$$

Dipole showers: branching

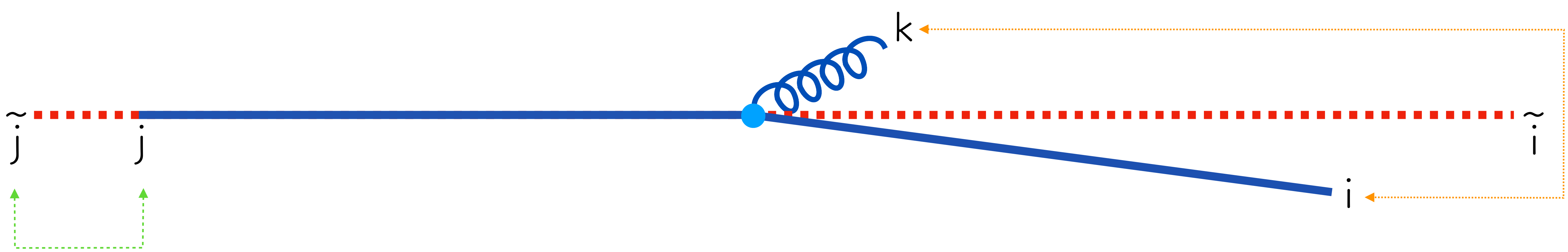
- Branching probability: evolution equation solved in terms of a Sudakov form factor

$$\Delta\left(v_0, v_1\right)\, d\mathcal{P}_{\bar{q}[q]\to\bar{q}g[q]} + \Delta\left(v_0, v_1\right)\, d\mathcal{P}_{q[\bar{q}]\to qg[\bar{q}]}$$



Dipole showers: local recoil

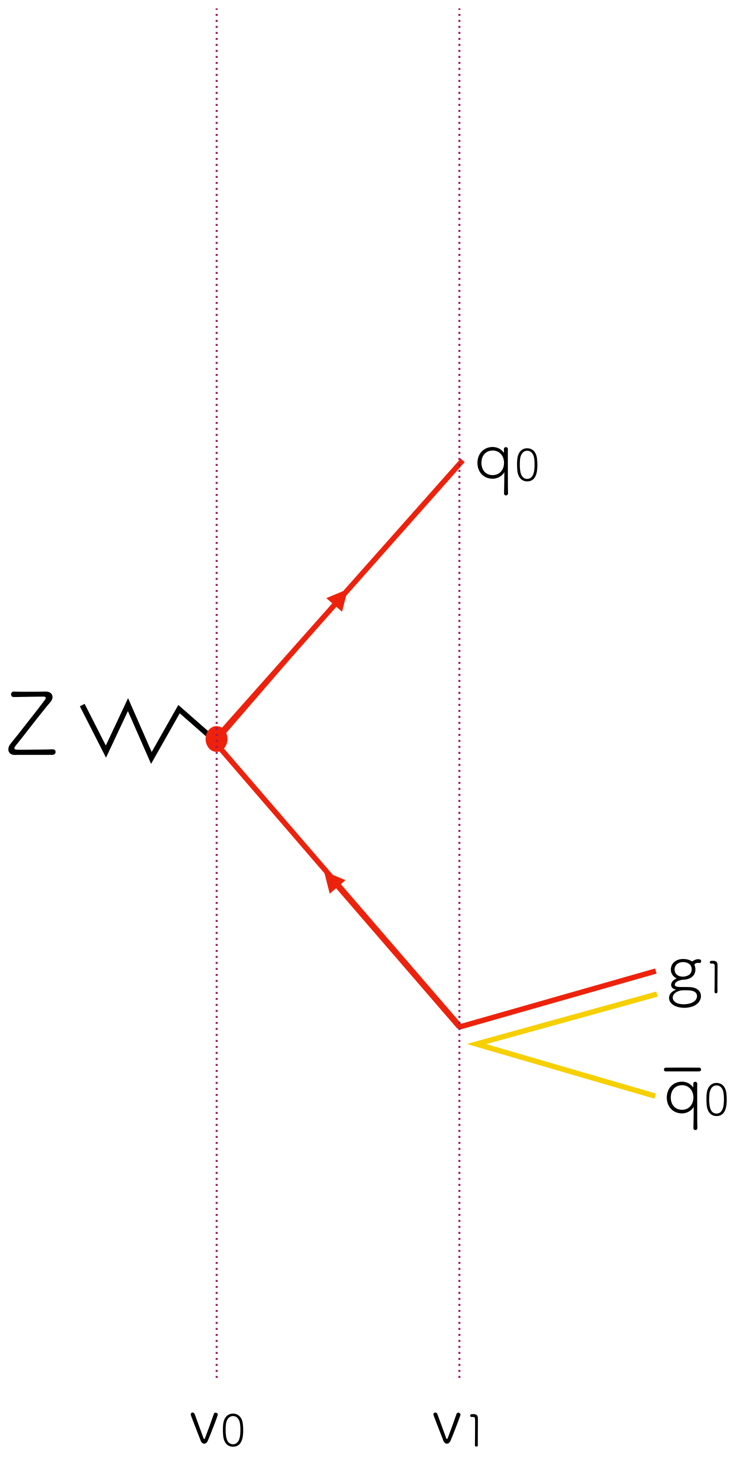
- Kinematic mapping: to ensure momentum conservation, the recoil is assigned locally (within the dipole)
- the *emitter* i takes the recoil of k in the $\tilde{i}\,\tilde{j}$ C.O.M. frame
- residual longitudinal recoil absorbed by the *spectator* j



$$\tilde{p}_i + \tilde{p}_j \xrightarrow{\tilde{p}_i \to p_i + p_k} p_i + p_j + p_k$$

$$\begin{aligned} p_i^\mu &= \tilde{z}\,\tilde{p}_i^\mu + y\,(1-\tilde{z})\,\tilde{p}_j^\mu + k_\perp \\ p_k^\mu &= (1-\tilde{z})\,\tilde{p}_i^\mu + y\,\tilde{z}\,\tilde{p}_j^\mu - k_\perp^\mu \\ p_j^\mu &= (1-y)\,\tilde{p}_j^\mu \end{aligned}$$

Dipole showers: iterate



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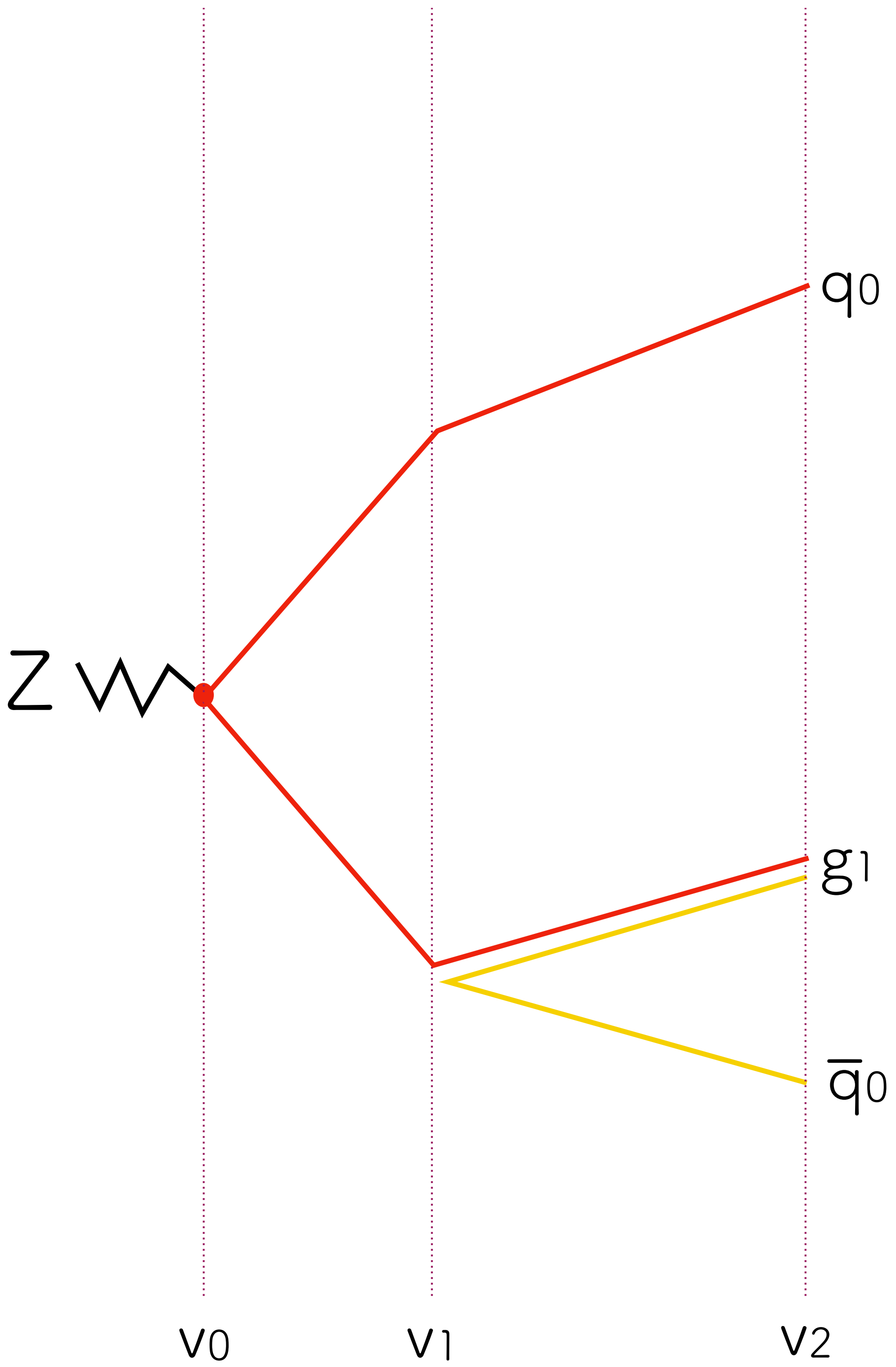

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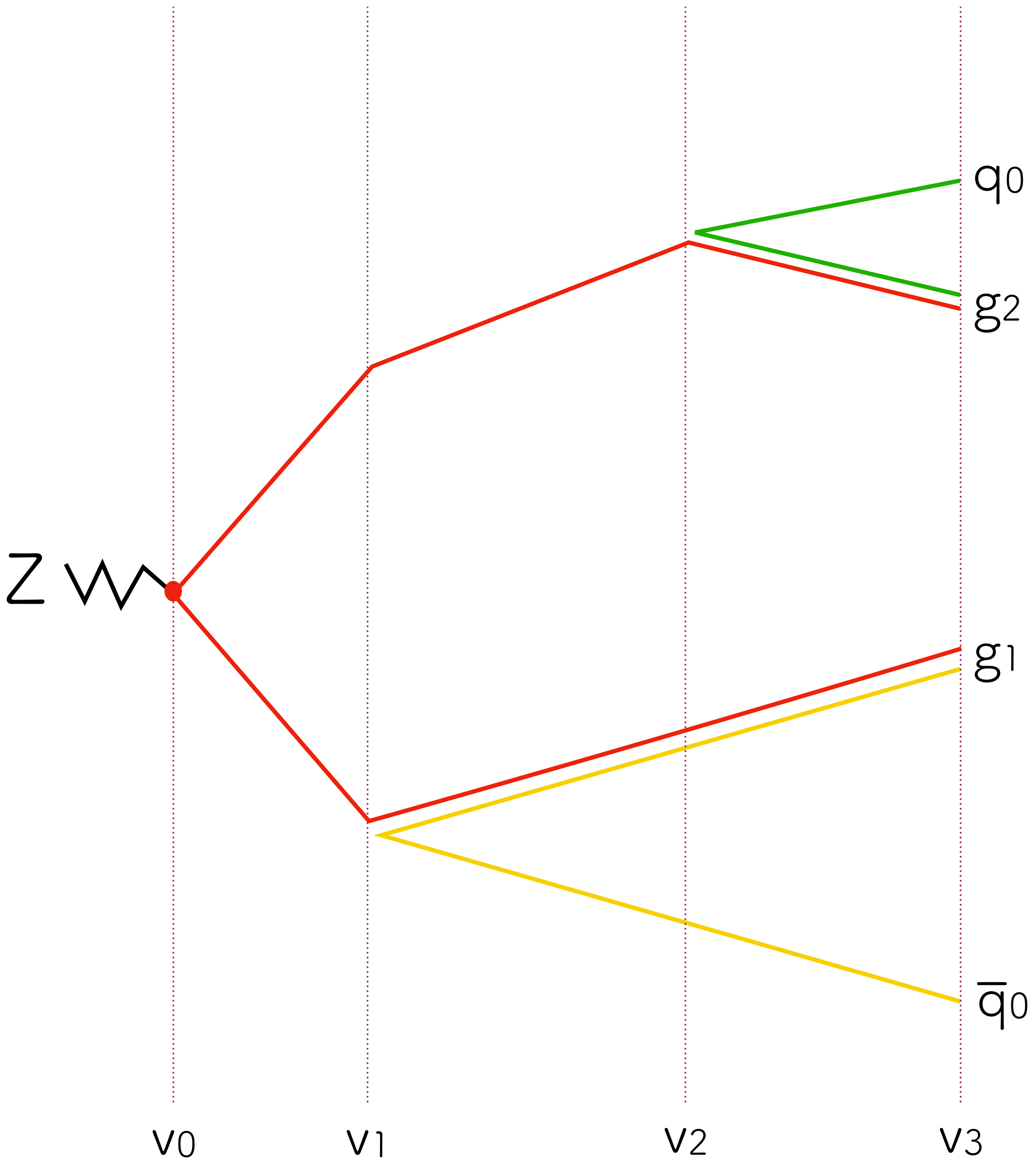

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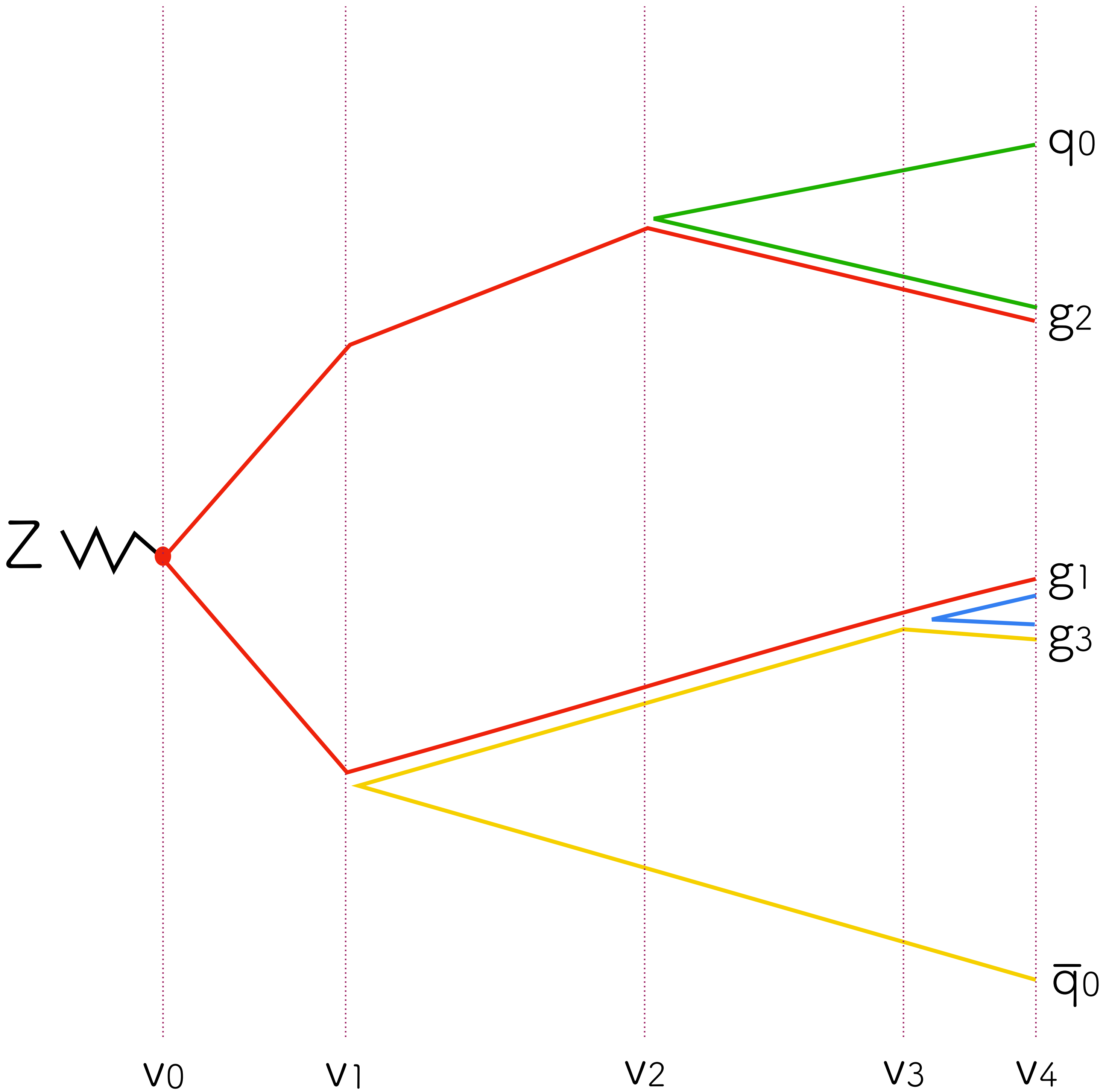

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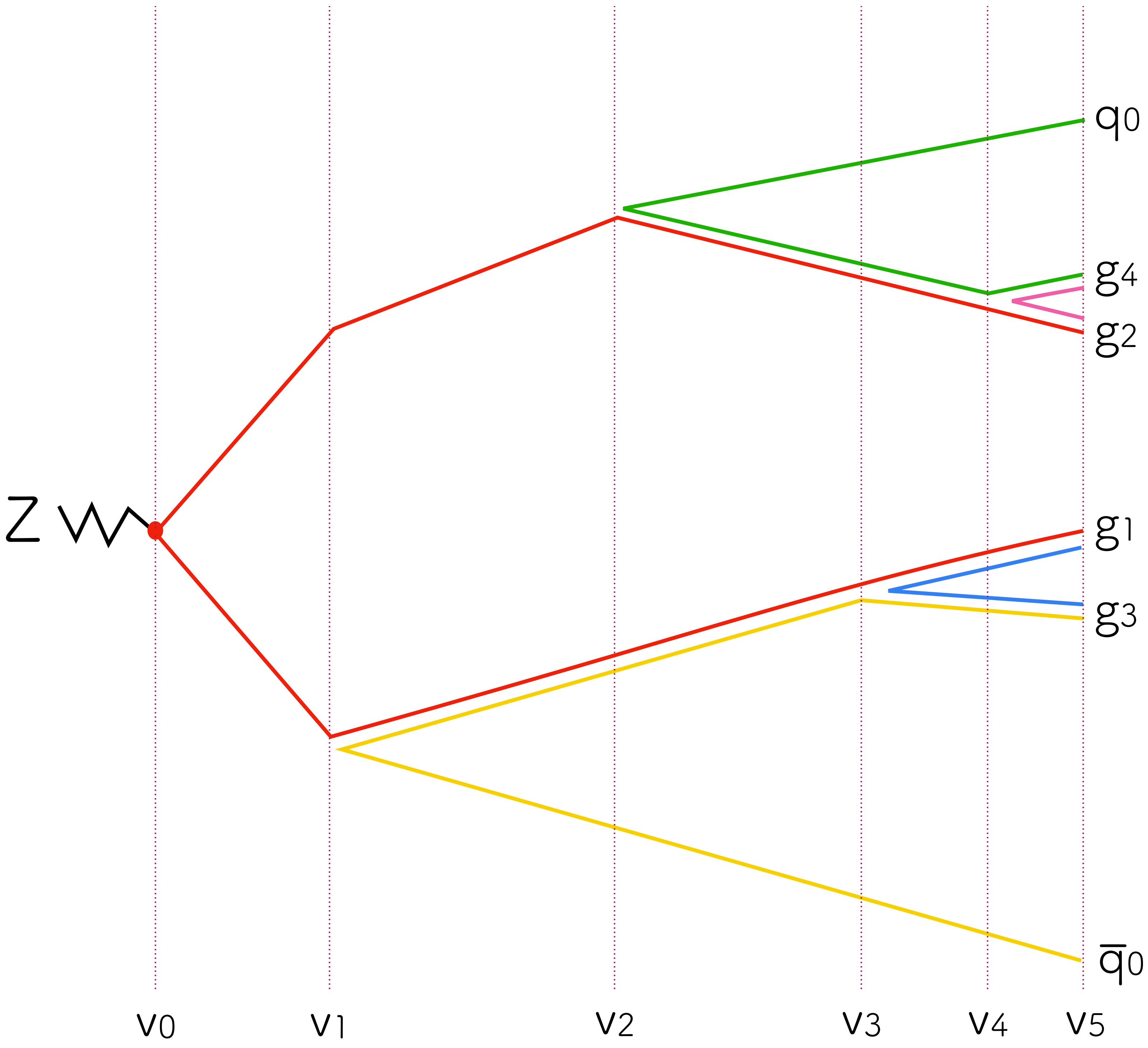

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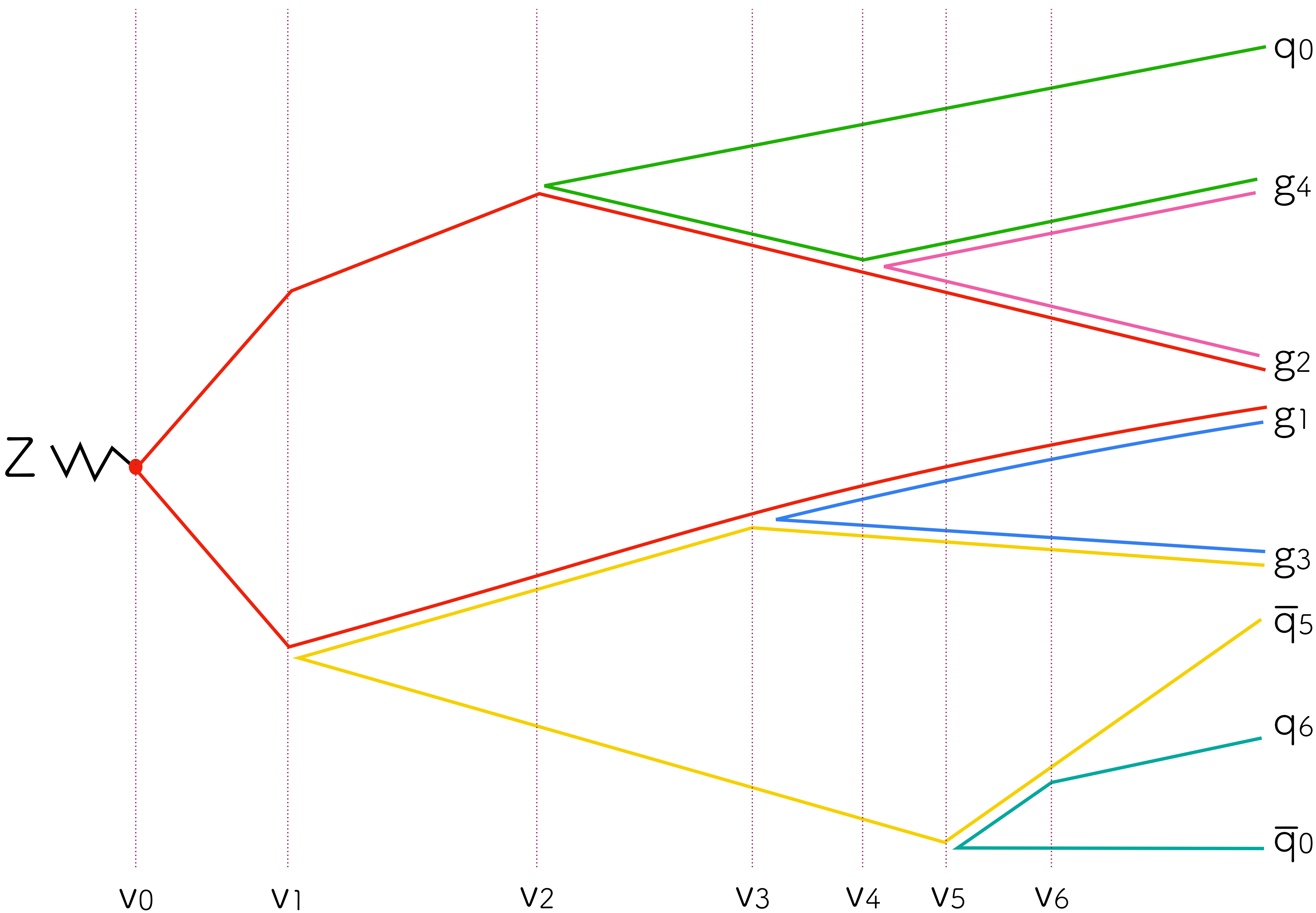

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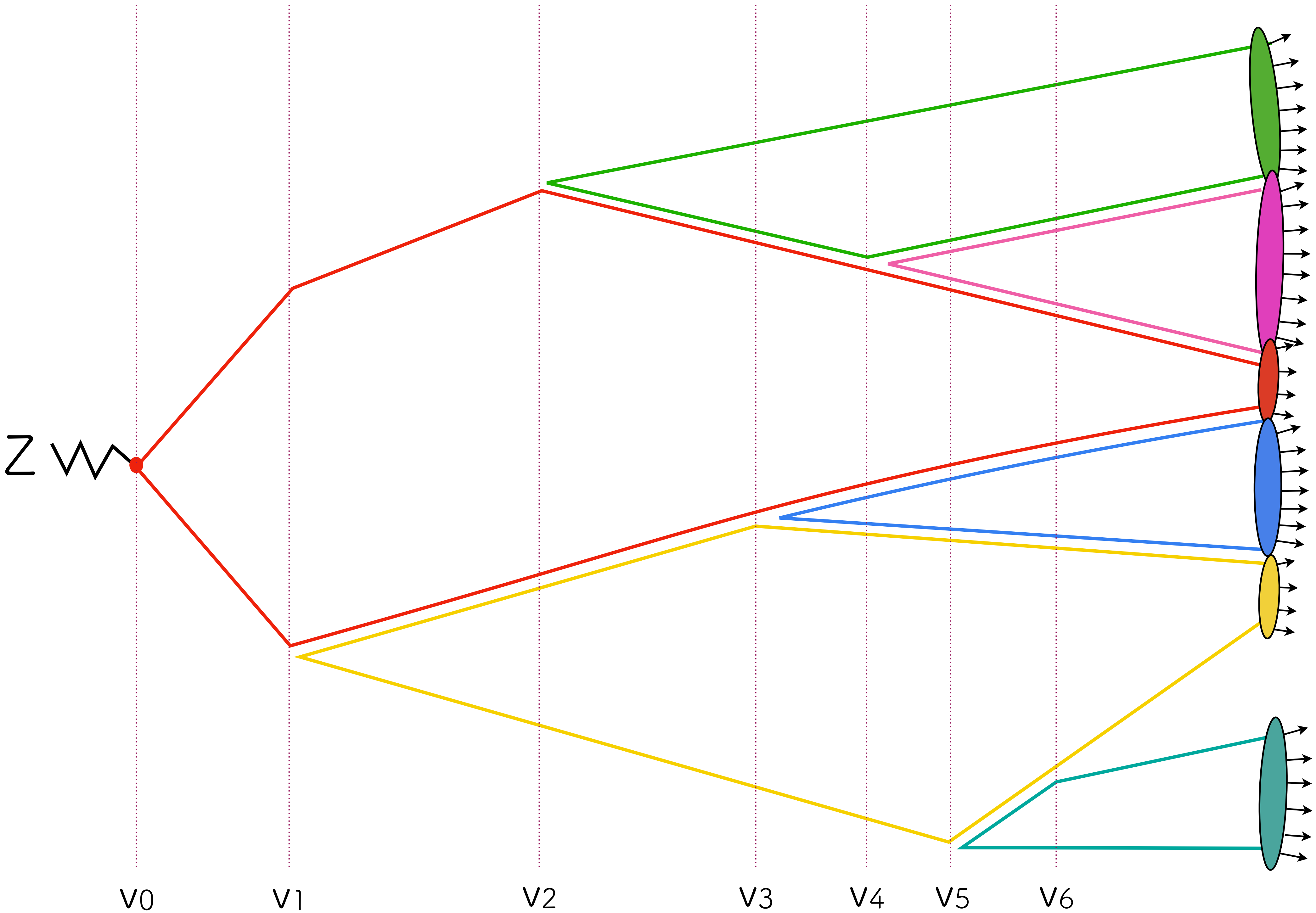
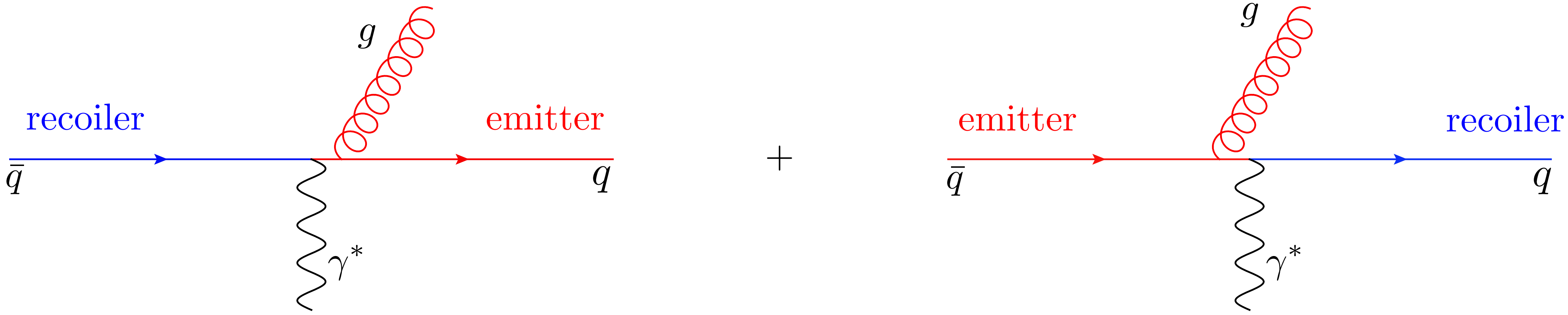


Single soft emission

Single emission: soft limit

- Both showers divide the dipole into two parts, at zero rapidity in the dipole's rest frame



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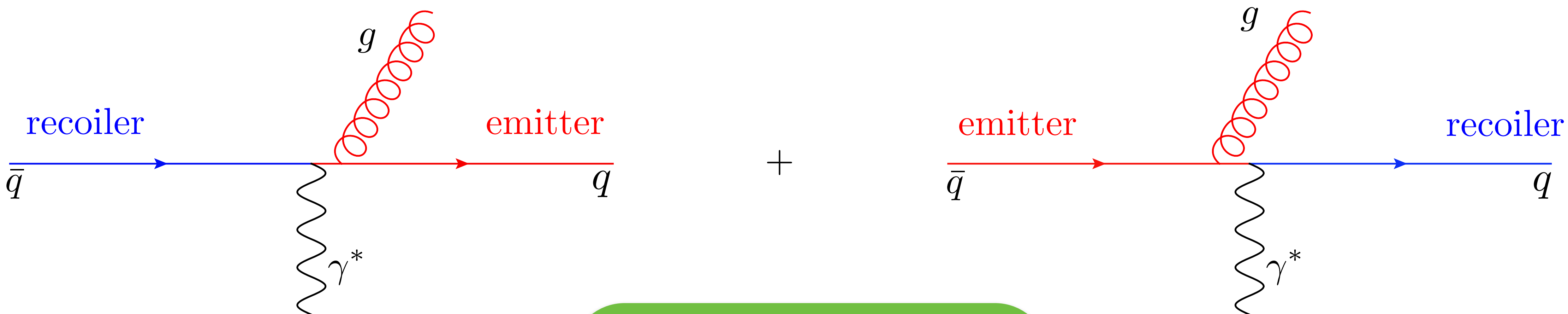


e.g. emission off the quark

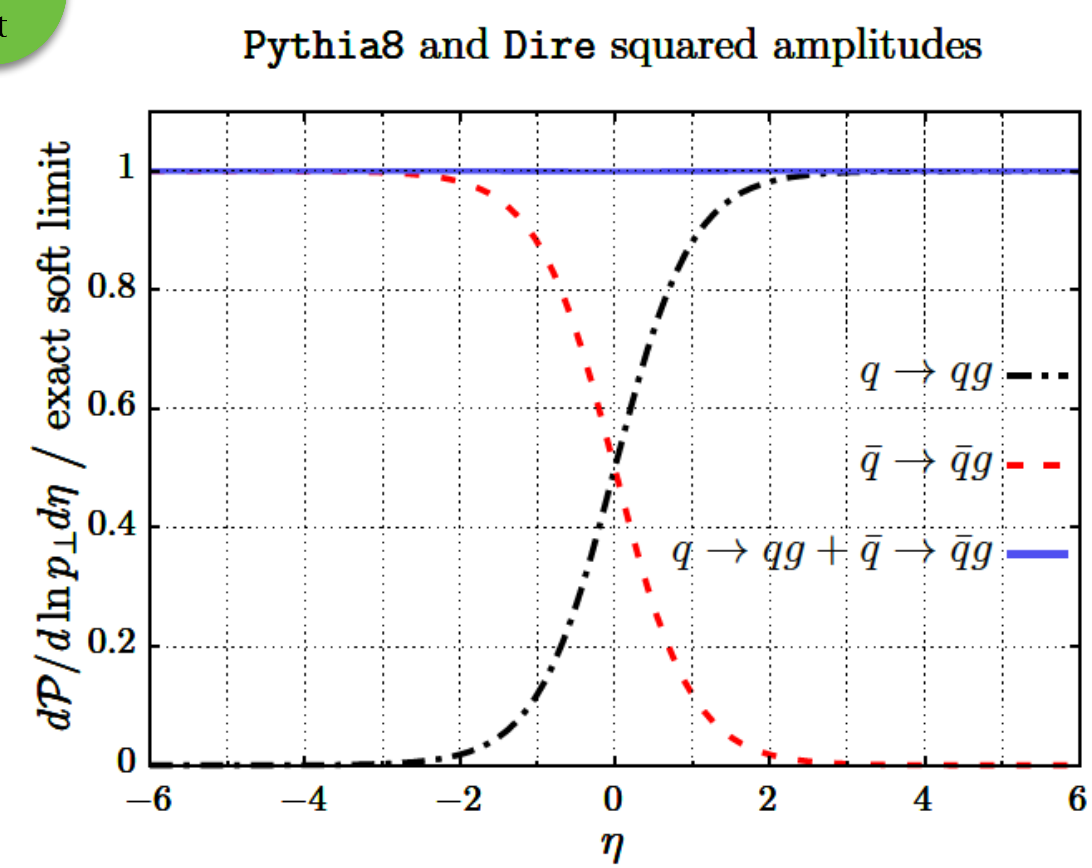
Scale of the coupling should be the physical dipole k_t

Pythia

$$d\mathcal{P}_{q\to qg} = \frac{2\alpha_s(p^2_{\perp,\mathrm{evol}})C_F}{\pi}\frac{dp_\perp}{p_\perp}d\eta\left(\frac{e^{2\eta}}{1+e^{2\eta}}\right)$$

Dire

$$d\mathcal{P}_{q\to qg} = \frac{2\alpha_s(t)C_F}{\pi}\frac{dp_\perp}{p_\perp}d\eta\left(\frac{e^{2\eta}}{1+e^{2\eta}}\right)$$



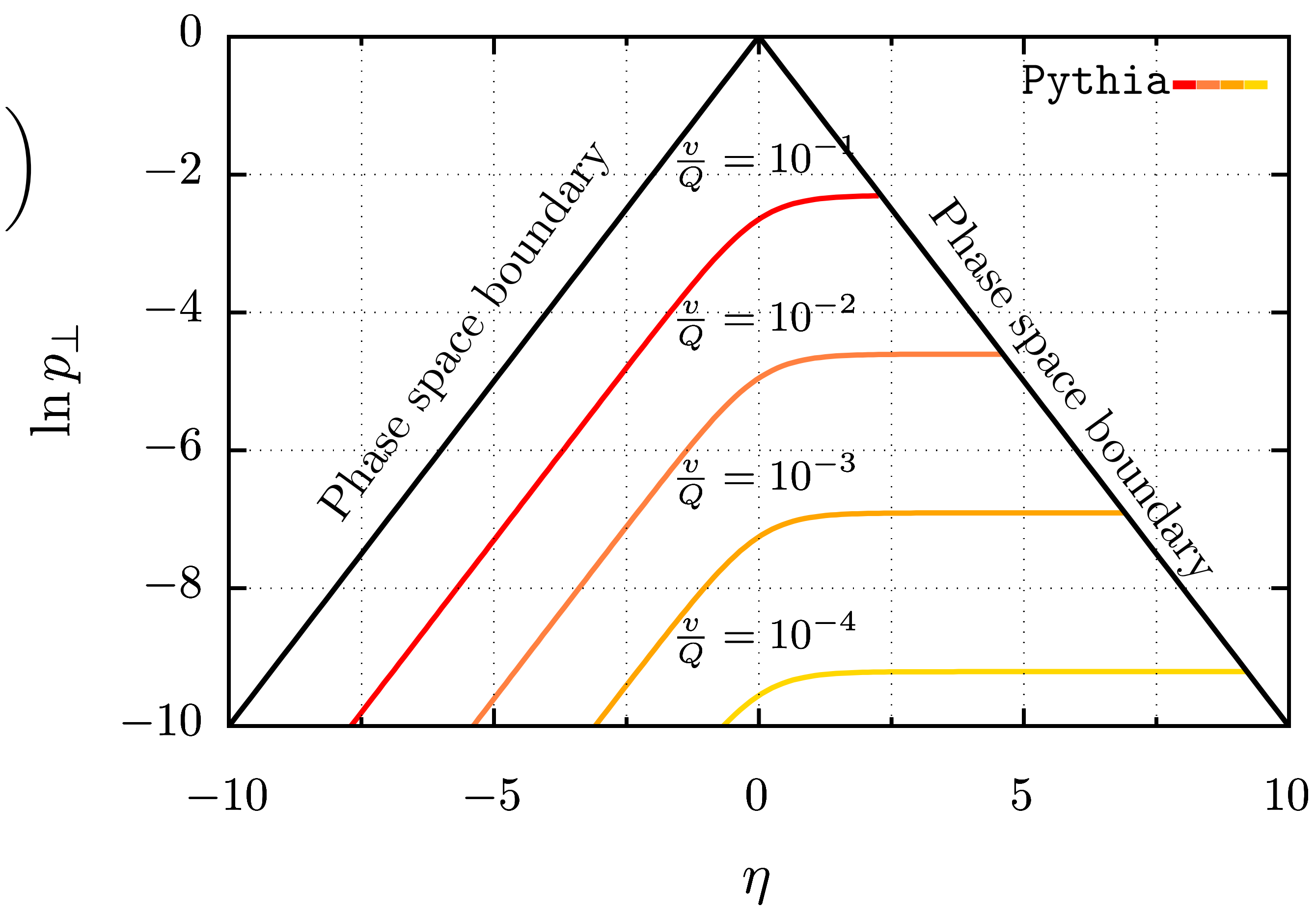
Single emission: soft limit

- `Pythia` case:

$$\eta = \frac{1}{2}\ln\left[\frac{(1-z)^2}{\rho^2_{\perp,\mathrm{evol}}} - 1\right], \qquad |p_\perp^2| = p^2_{\perp,\mathrm{evol}}\left(\frac{e^{2\eta}}{1+e^{2\eta}}\right)$$

$$d\mathcal{P}_{q\to qg} = \frac{2\alpha_s(p^2_{\perp,\mathrm{evol}})C_F}{\pi}\frac{dp_\perp}{p_\perp}d\eta\left(\frac{e^{2\eta}}{1+e^{2\eta}}\right)$$

Constant evolution variable contours in the Lund plane



- Correct matrix element for a single emission is reproduced up to running coupling effects

$$d\mathcal{P}_{q\to qg} + d\mathcal{P}_{\bar{q}\to \bar{q}g} = \frac{2\alpha_s C_F}{\pi}\frac{dp_\perp}{p_\perp}d\eta$$

Not true anymore with running coupling in the soft-wide-angle region (NNLL effect)

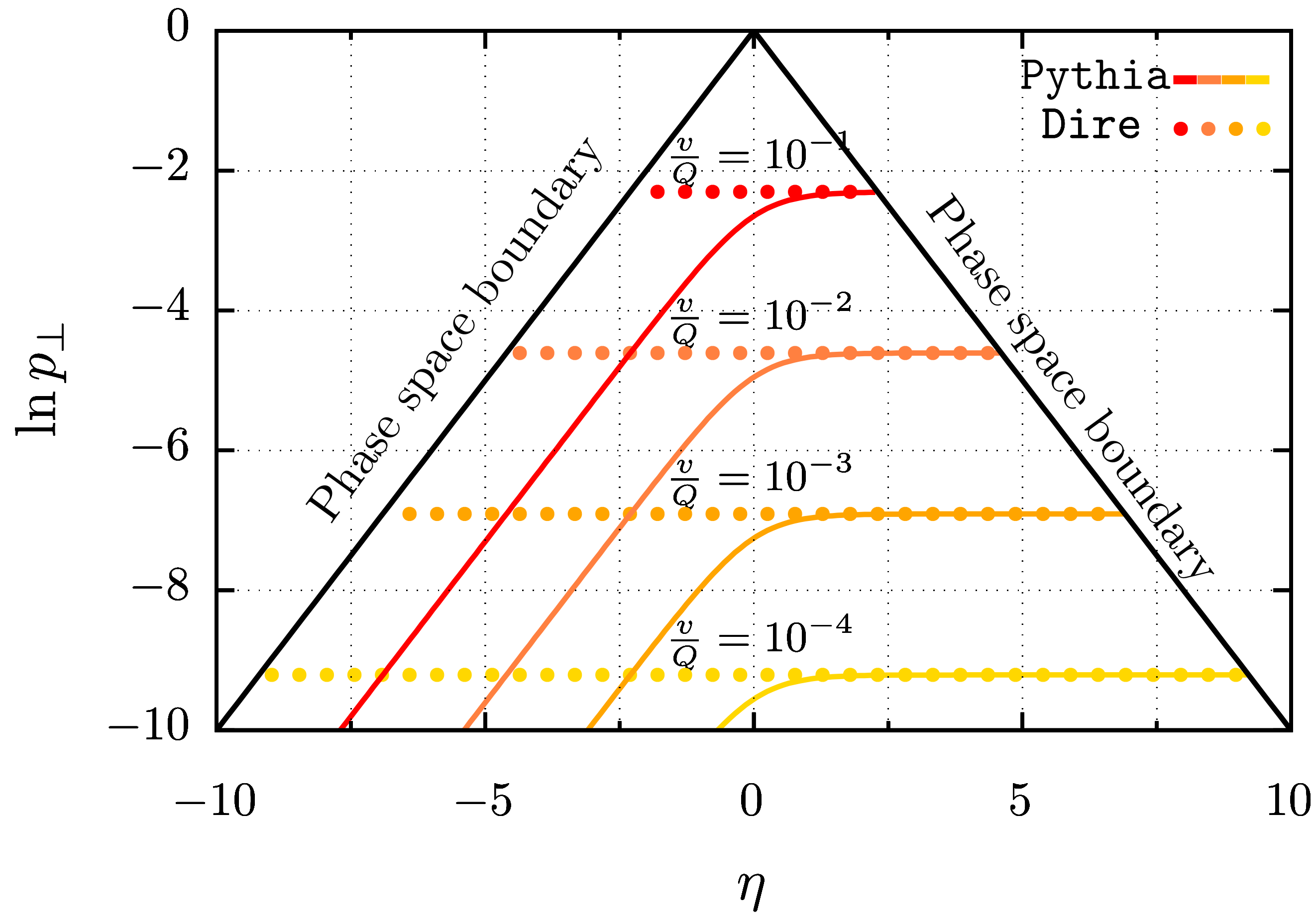
- Non-zero (although suppressed) probability to have an emission with zero transverse momentum even if $p_{\perp,\mathrm{evol}} \neq 0$. This creates a new suppression mechanism in competition with the usual Sudakov suppression. In practice, unlikely to be of phenomenological interest

Single emission: soft limit

Constant evolution variable contours in the Lund plane

- `Dire` case:

$$\eta = \frac{1}{2}\ln\left[\frac{(1-z)^2}{\kappa^2}\right], \qquad |p_\perp^2| = t$$



$$d\mathcal{P}_{q\to qg} = \frac{2\alpha_s(t)C_F}{\pi}\frac{dp_\perp}{p_\perp}d\eta\left(\frac{e^{2\eta}}{1+e^{2\eta}}\right)$$

- Correct matrix element for a single emission is reproduced including running coupling effects*

$$d\mathcal{P}_{q\to qg} + d\mathcal{P}_{\bar{q}\to\bar{q}g} = \frac{2\alpha_s(|p_\perp^2|)C_F}{\pi}\frac{dp_\perp}{p_\perp}d\eta$$

* CMW scheme available both in Pythia and Dire

Multiple soft emissions

Multiple emissions: soft limit

- We now consider two soft-collinear emissions (g_1 and g_2 with $v_1 > v_2$) in the limit where they are strongly ordered in angle. This approximation is relevant at NLL for all global, rIRC safe observables.
- From the resummation one expects that both gluons are emitted off the initial $q\bar{q}$ dipole with

$$dP_2 = \frac{C_F^2}{2!} \prod_{i=1,2} \left(\frac{2\alpha_s(p_{\perp,i}^2)}{\pi} \frac{dp_{\perp,i}}{p_{\perp,i}} d\eta_i \frac{d\phi_i}{2\pi} \right)$$

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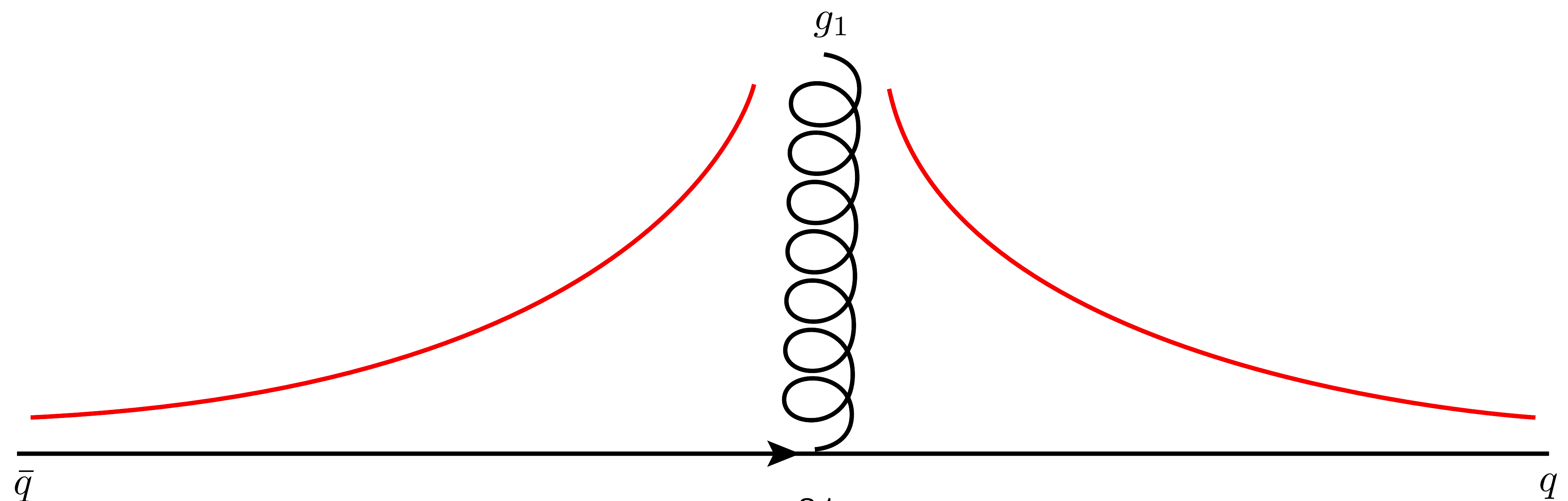
- Instead, the dipole-shower algorithm assigns the second emission to the first gluon in a portion of phase space in which it's collinear to the quarks: implications on logarithmic accuracy

e.g.

- start with an emission g_1



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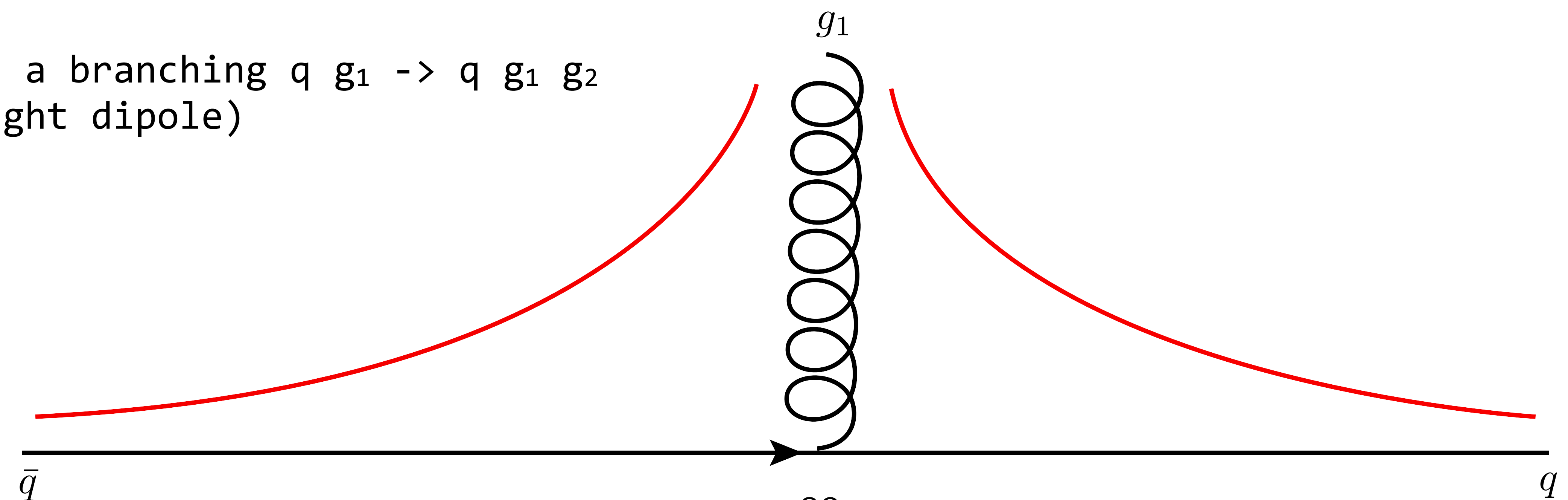

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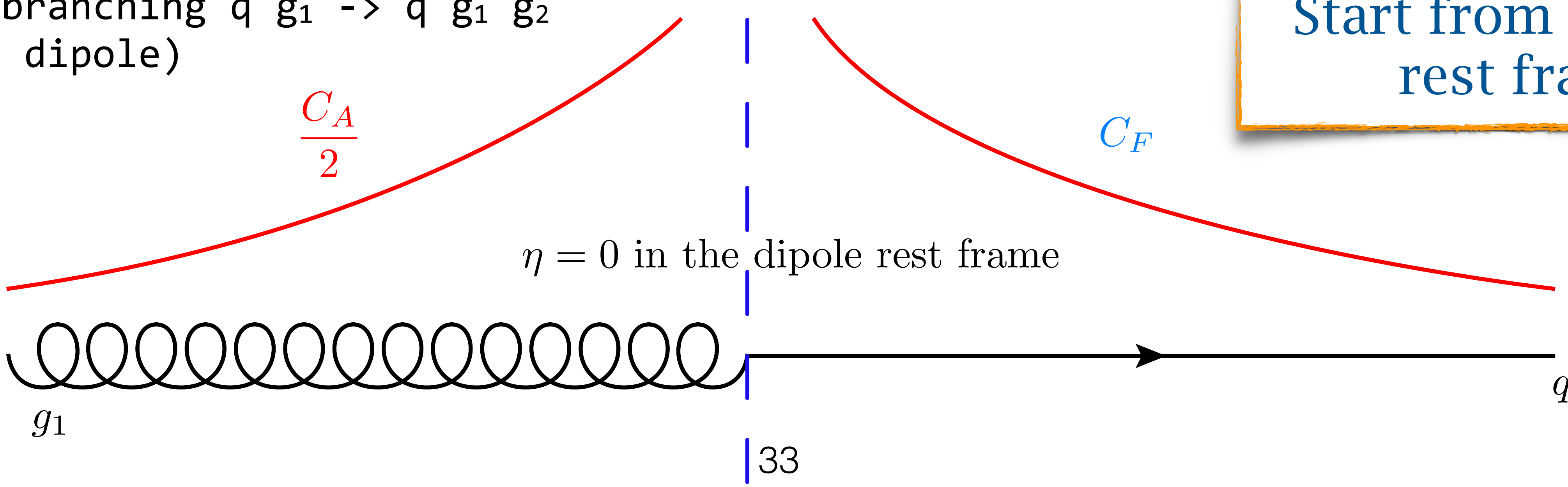
Start from dipole's rest frame



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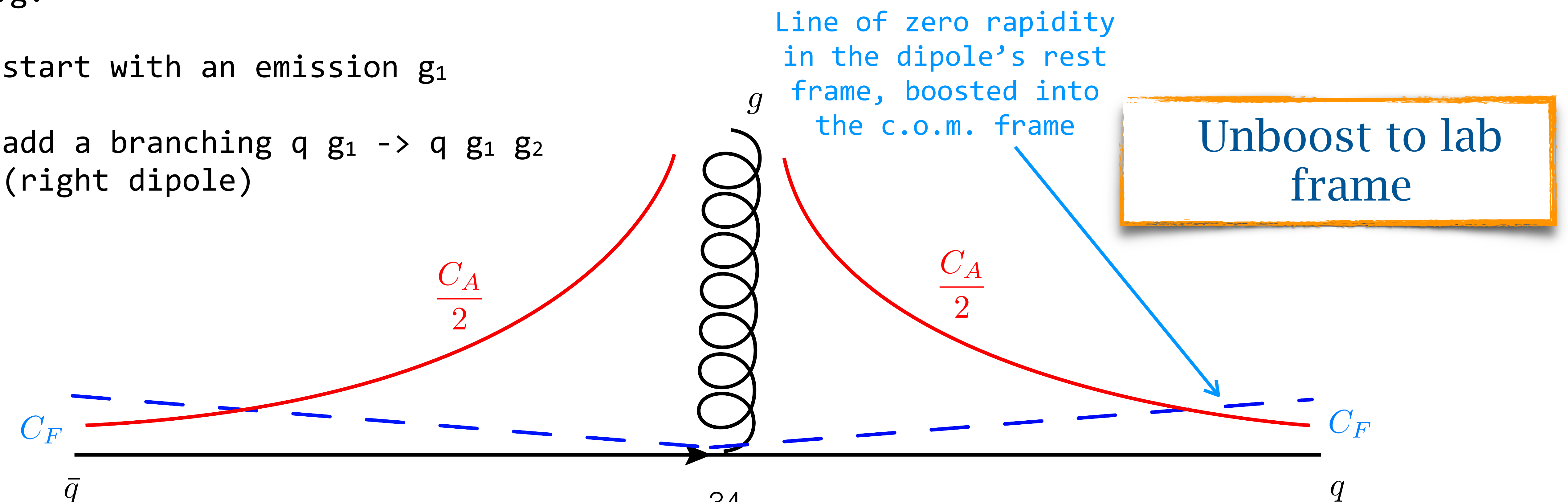


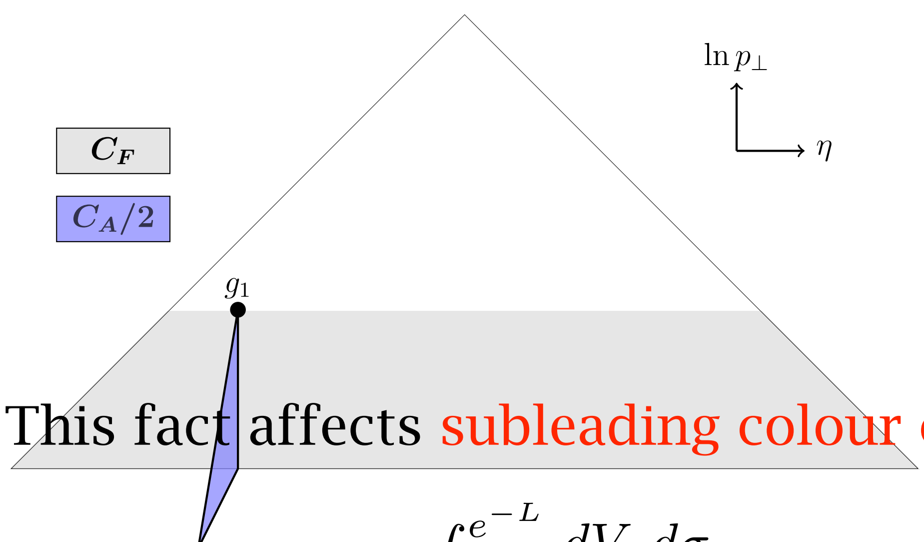
Double strong ordering

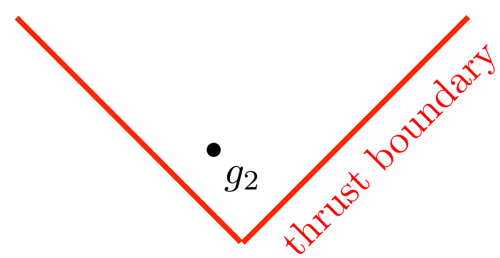
- Start by considering the limit where (in addition to angles) the ordering variable is strongly ordered, i.e. the kinematic of g_1 is not affected by the much softer g_2

$$v_1 \gg v_2$$

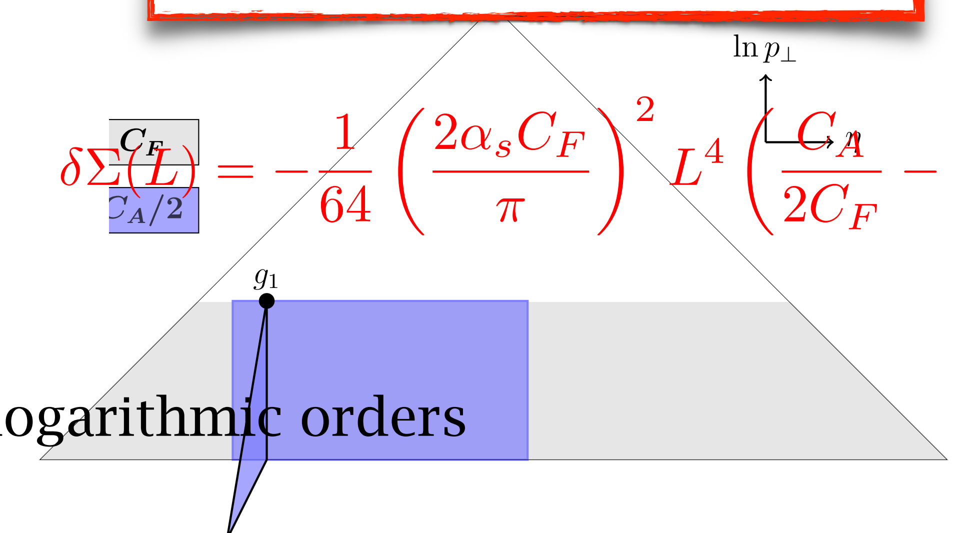
- However, the colour charge for the second emission depends on the above partitioning

Correct radiation pattern





Difference between dipole and correct pattern



$$\delta\Sigma(L) = -\frac{1}{64}\left(\frac{2\alpha_s C_F}{\pi}\right)^2 L^4 \left(\frac{C_A}{2C_F} - 1\right)$$

- This fact affects subleading colour corrections at different logarithmic orders

$$\Sigma(L) \equiv \int_0^{e^{-L}} \frac{dV}{\sigma_B}\frac{d\sigma}{dV} = \exp\left[L g_1(\alpha_s L) + g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \cdots\right] + \mathcal{O}\left(\alpha_s e^{-L}\right)$$

for an observable $V(p, \{\text{Born momenta}\}) \propto p_\perp^a e^{-|\eta_p| b}$

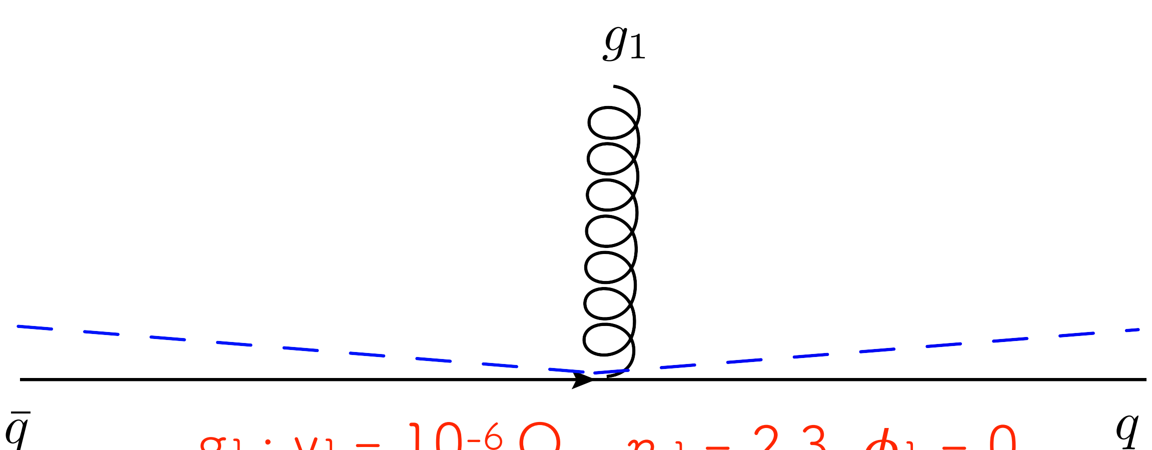
- Observables with $b = 0$ (e.g. p_t, k_t jet rates,…) are affected at NLL
- Observables with $b \neq 0$ (e.g. thrust, jet mass, …) are affected at LL

Single strong ordering: kinematics

- When the ordering variables are of the same order ($v_1 \gtrsim v_2$) the first emission g_1 is affected by the second (g_2) when this is far from g_1 in the lab frame
- The kinematics of the first emission is thus affected also by these *recoil* effects (transverse recoil + conservation of dipole’s invariant mass)
- Eventually reflected in the observables

e.g.

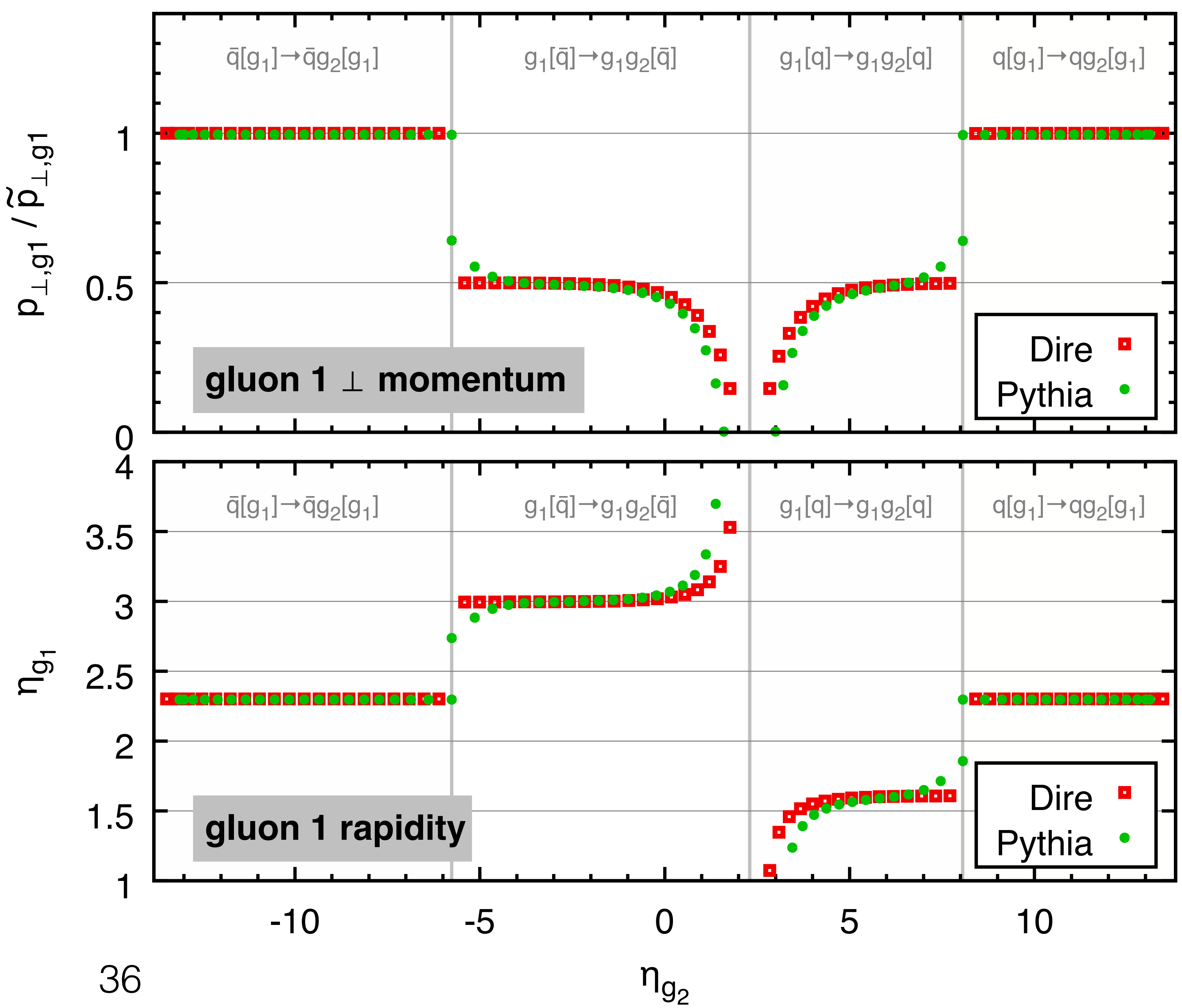
- start with an emission g_1



g_1 : $v_1 = 10^{-6}$ Q , $\eta_1 = 2.3$, $\phi_1 = 0$

g_2 : $v_2 = 0.5\ v_1$, $\phi_2 = 0$, scan in η_2

impact of gluon-2 emission on gluon-1 momentum



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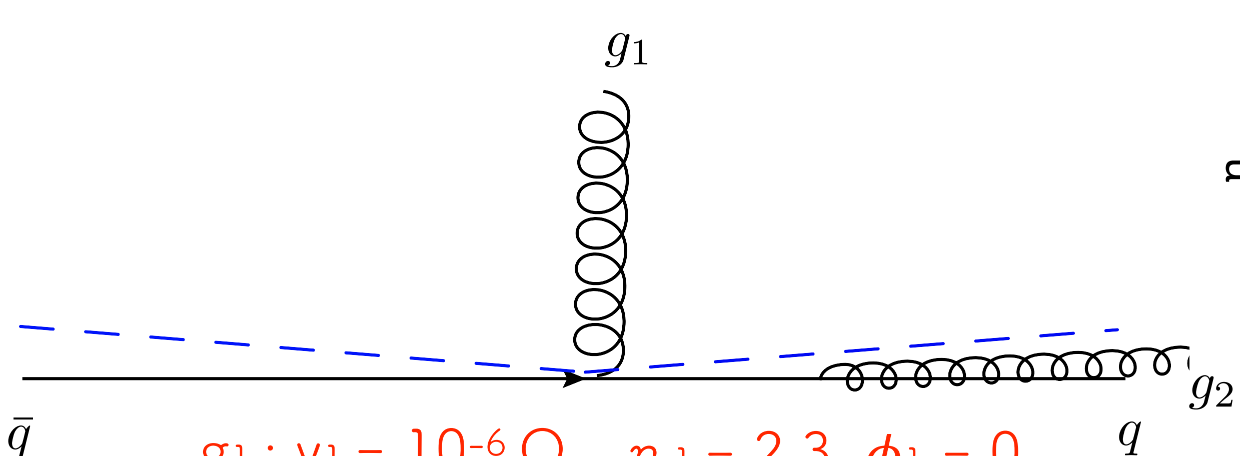
- The kinematics of the first emission is thus affected also by these *recoil* effects (transverse recoil + conservation of dipole's invariant mass)
- Eventually reflected in the observables

e.g.

- start with an emission g_1
- add a second emission g_2

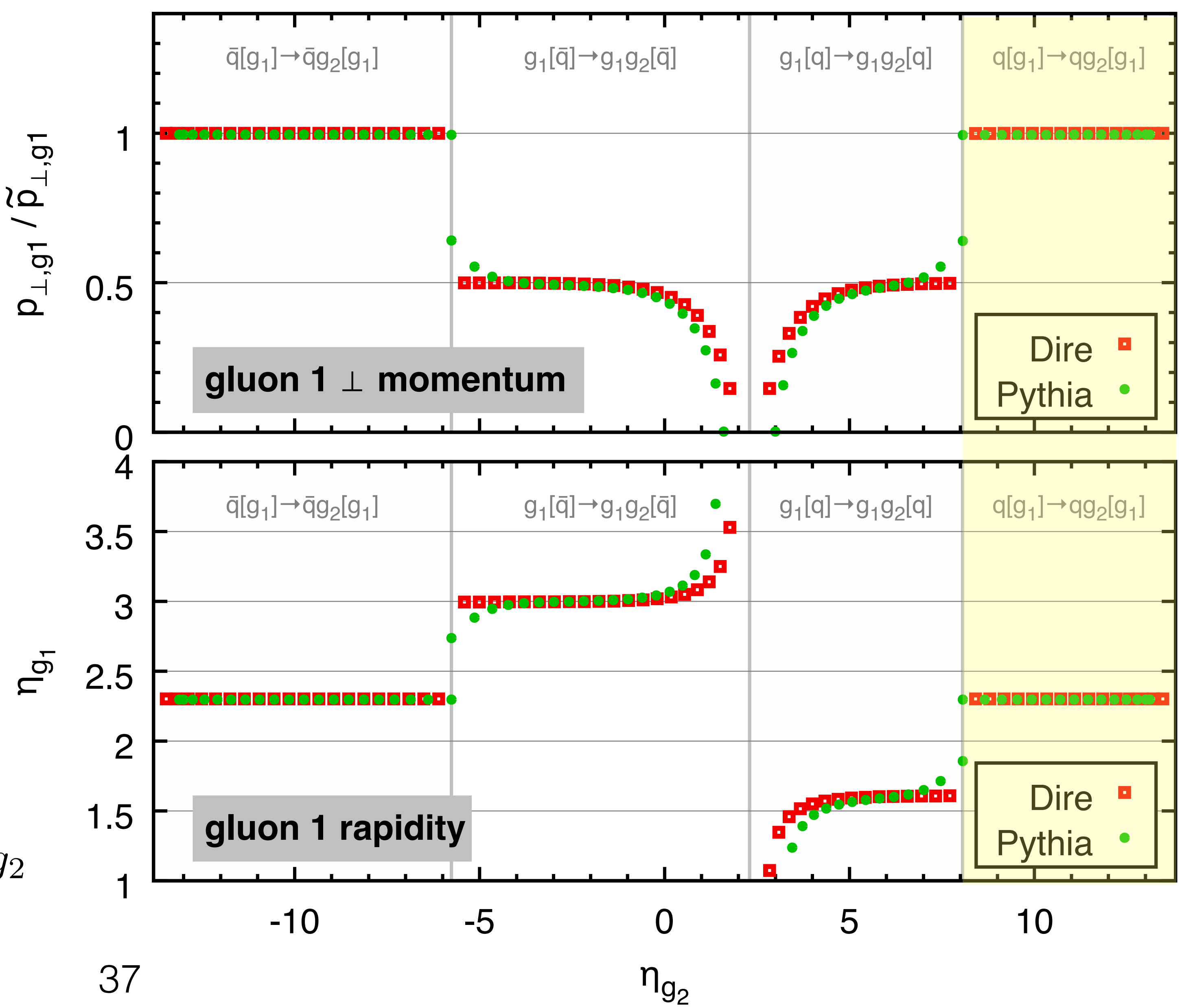
$q[g_1] \to q g_2[g_1]: \; \boldsymbol{p}_{\perp,g_1} = \tilde{\boldsymbol{p}}_{\perp,g_1}, \eta_{g_1} = \tilde{\eta}_{g_1}$



g_1 : $v_1 = 10^{-6}$ Q , $\eta_1 = 2.3$, $\phi_1 = 0$

g_2 : $v_2 = 0.5\, v_1$, $\phi_2 = 0$, scan in η_2

impact of gluon-2 emission on gluon-1 momentum



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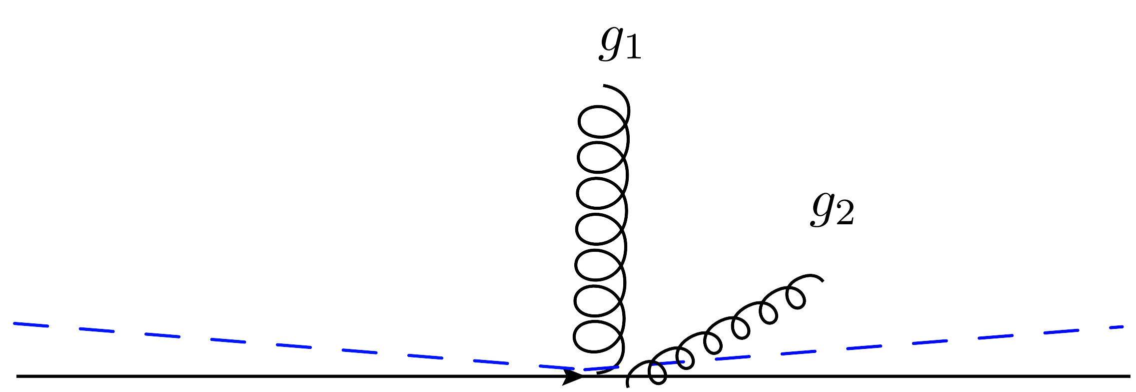

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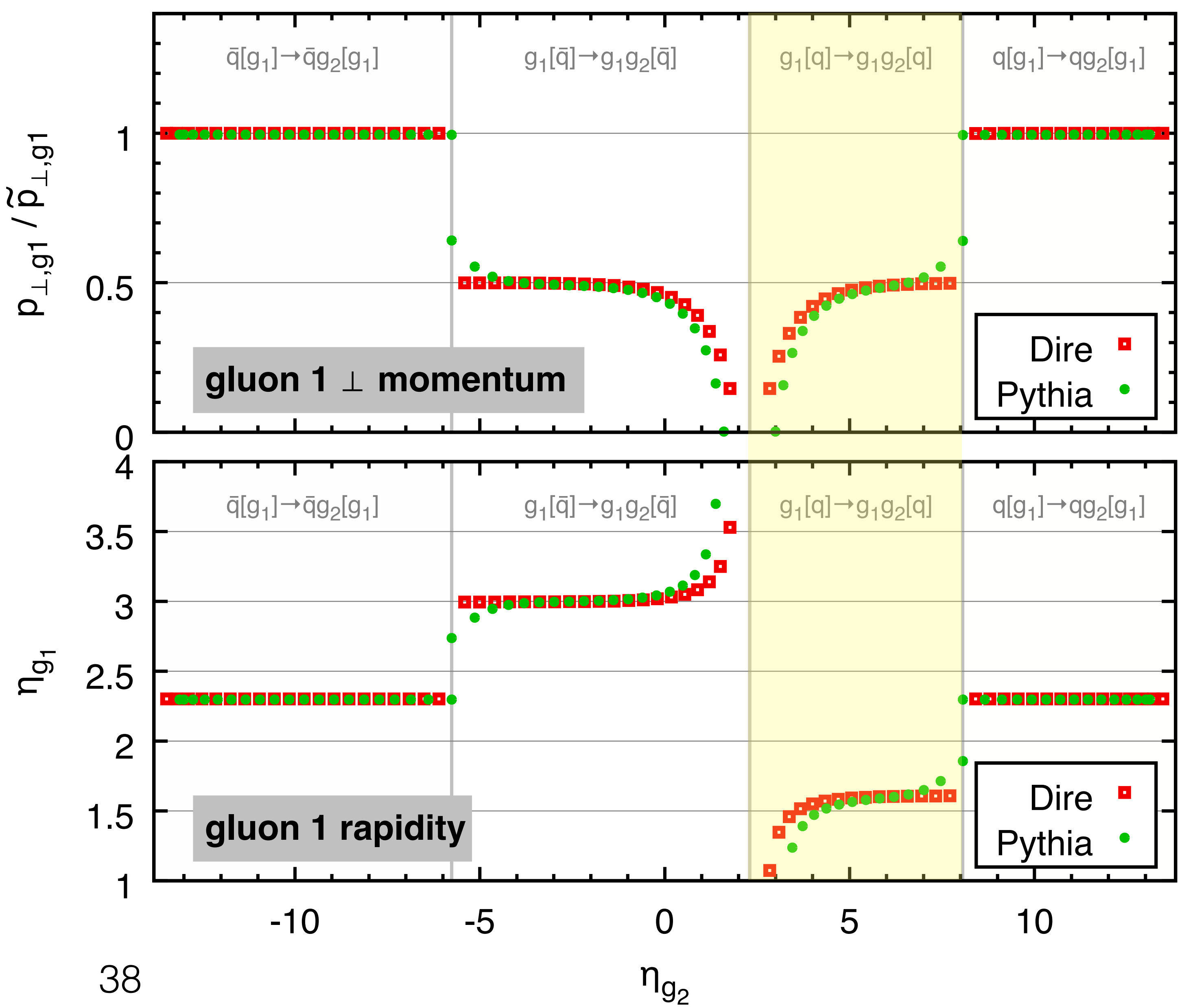

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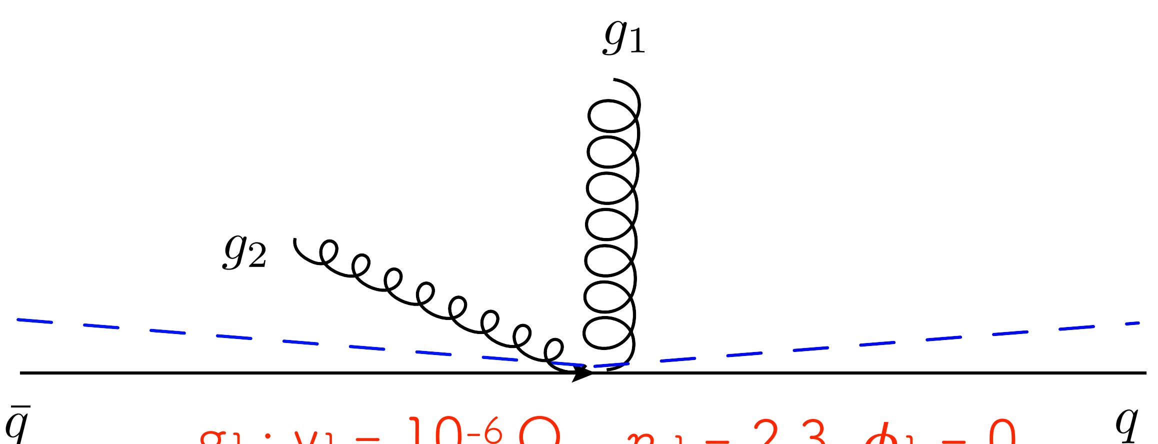

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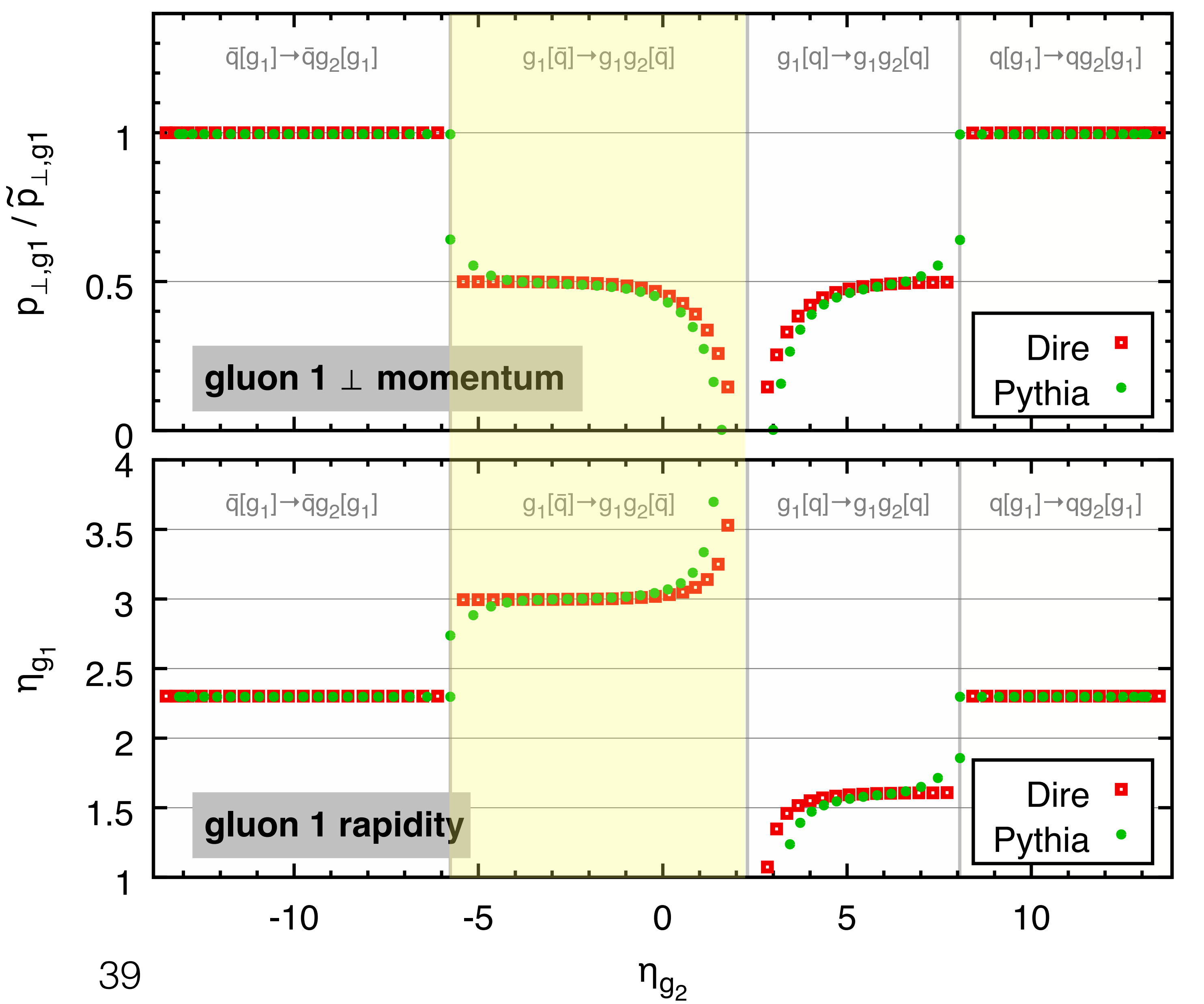

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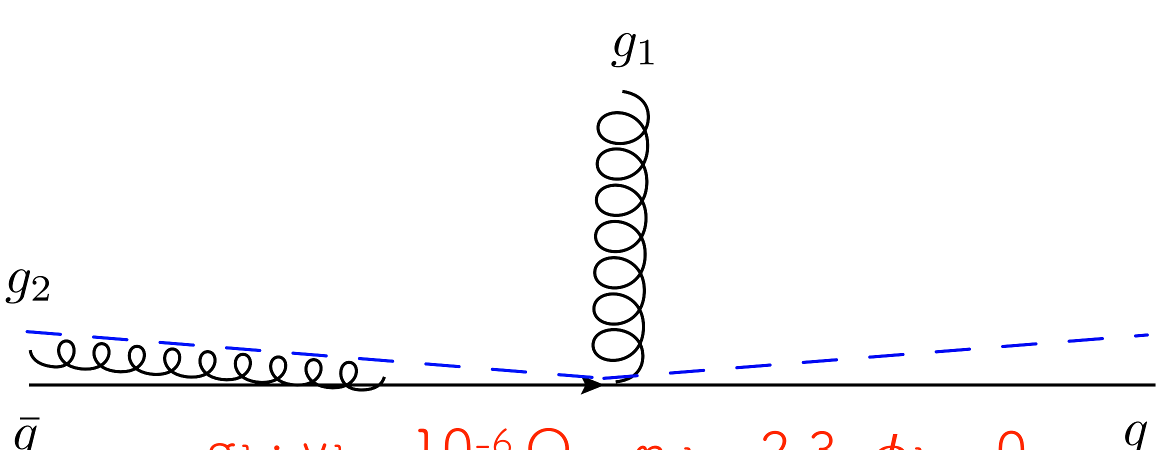

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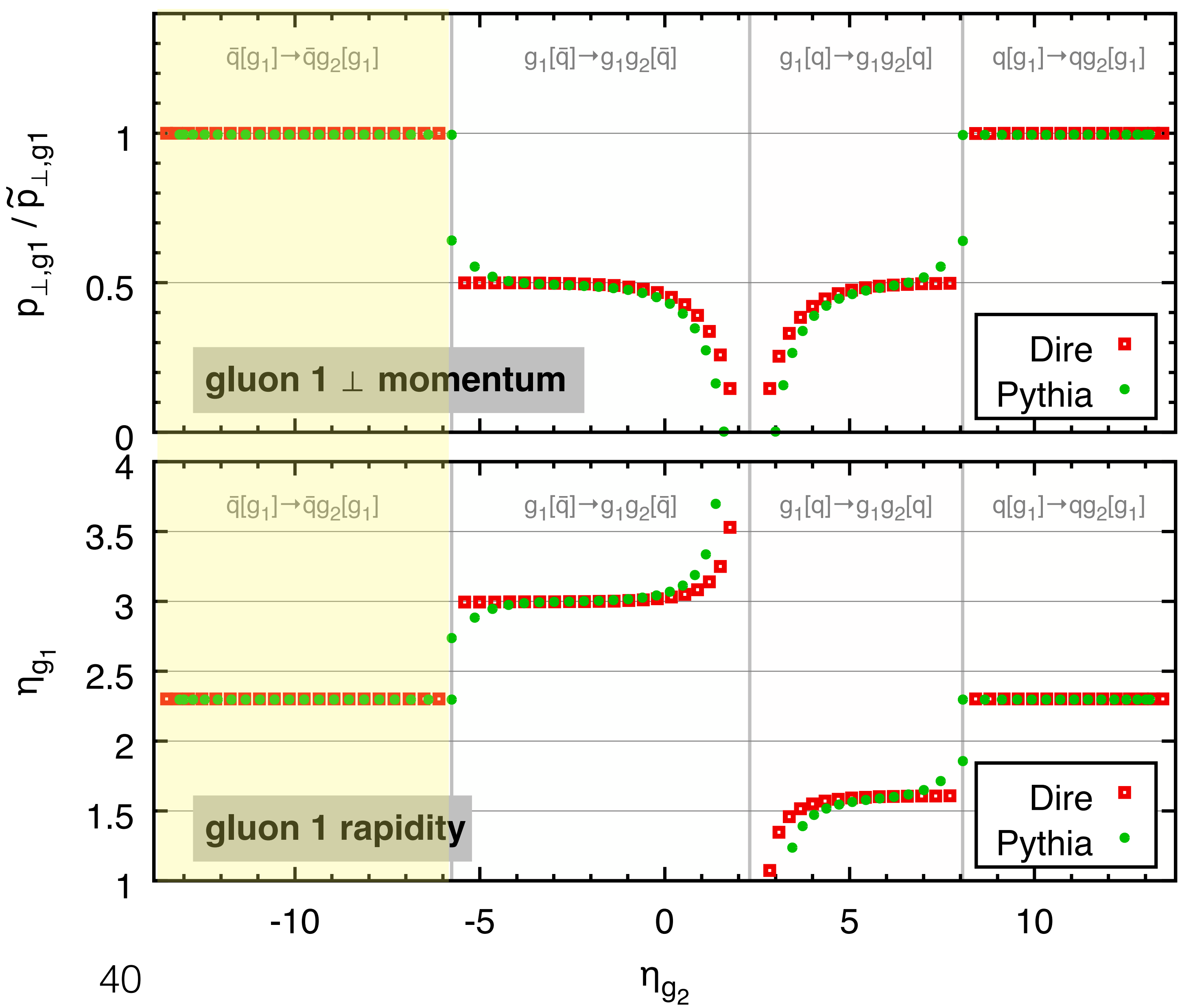
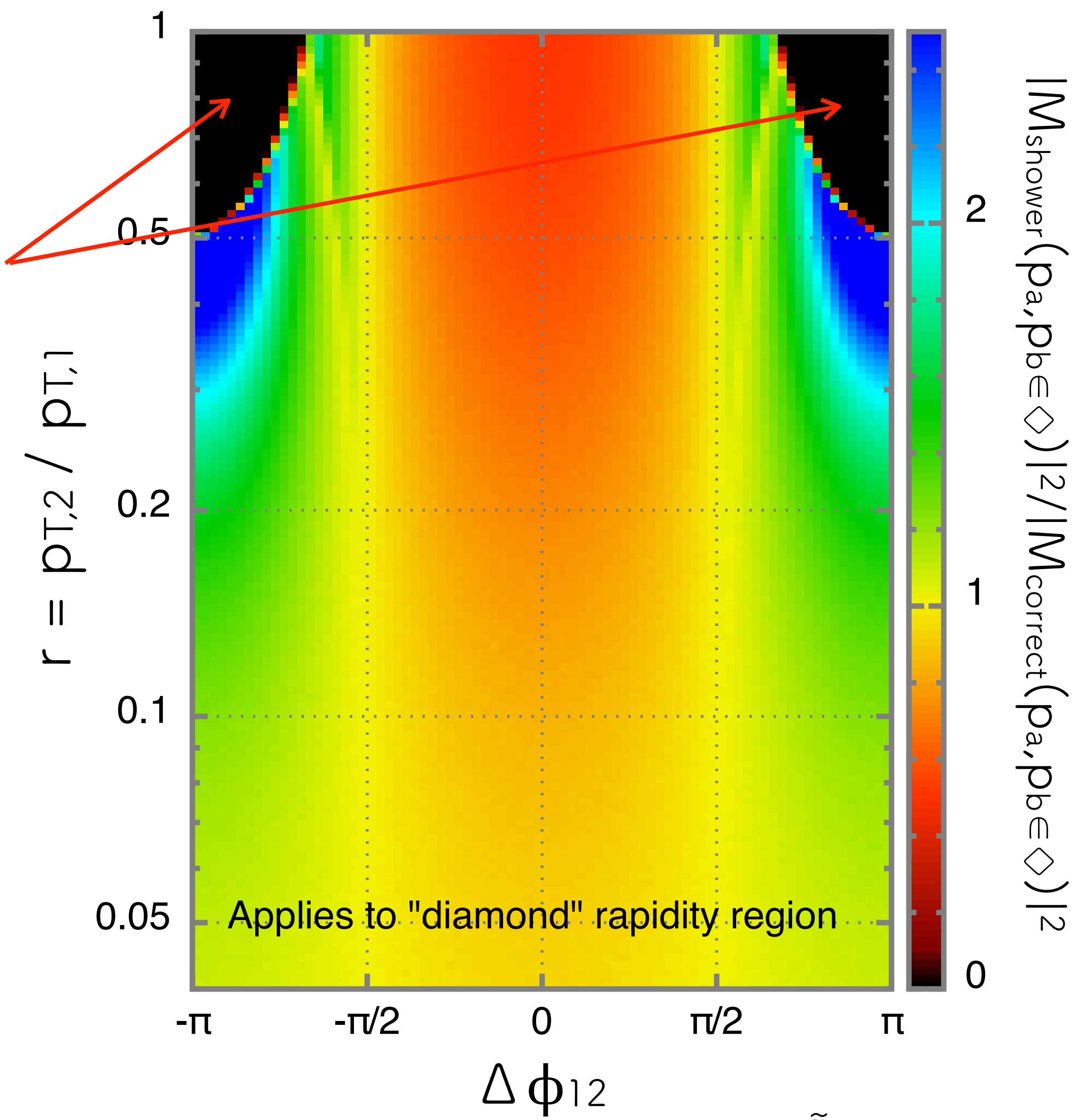


Single strong ordering: matrix element

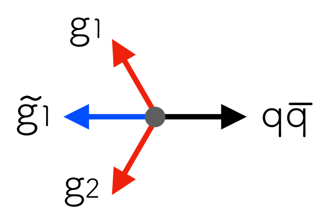
- As a consequence, starting from second order, the effective matrix element differs from the NLL prediction
- Effects can be large for observables sensitive to exclusive regions of phase space
- This mechanism affects the pattern of subsequent real radiation, and virtual corrections, at all higher orders

e.g. at α_s^2

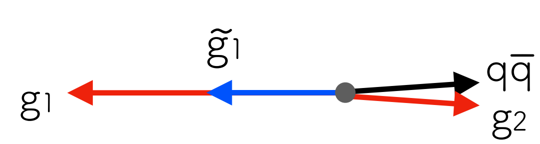
dipole-shower double-soft ME / correct result



➡E.g. $r = 1$, $|\Delta\phi| > \pm 2\pi/3$:



➡E.g. $|\Delta\phi| = \pm\pi$, $r > 0.5$:



Single strong ordering

- Occurs in a region relevant to NLL (leading colour) for all rIRC safe, global observables

e.g. 3-jet resolution in Cambridge algorithm

(angular ordered clustering of soft and/or collinear radiation)

$$\delta\Sigma^{(2\,\text{emissions})}(L) = \left(C_F \frac{2\alpha_s}{\pi}\right)^2 \int_0^1 \frac{dv_1}{v_1} \int_{\ln v_1}^{\ln 1/v_1} d\eta_1 \int_0^{v_1} \frac{dv_2}{v_2} \int_{\ln v_2}^{\ln 1/v_2} d\eta_2 \int_0^{2\pi} \frac{d\phi_1}{2\pi} \int_0^{2\pi} \frac{d\phi_2}{2\pi} \times \left[\Theta\left(e^{-L} - V(p_1^{\text{shower}}, p_2)\right) - \Theta\left(e^{-L} - V(p_1^{\text{correct}}, p_2)\right)\right]$$

$$V(p_1^{\text{correct}}, p_2) = v_1 \qquad V(p_1^{\text{shower}}, p_2) = \max\left(v_2,\ \sqrt{v_1^2 + v_2^2 - 2v_1 v_2 \cos\phi_{12}}\right)$$

$$\delta\Sigma^{\text{cam}}(L) = -0.18277\,\bar{\alpha}^2 L^2 + \mathcal{O}\left(\bar{\alpha}^2 L\right)$$

Single strong ordering

- Occurs in a region relevant to NLL (leading colour) for all rIRC safe, global observables

e.g. for a sample of observables

Observable	$\mathrm{NLL}_{\ln\Sigma}$ discrepancy
$1-T$	$0.116^{+0.004}_{-0.004}\,\bar{\alpha}^3L^3$
vector p_t sum	$-0.349^{+0.003}_{-0.003}\,\bar{\alpha}^3L^3$
B_T	$-0.0167335\,\bar{\alpha}^2L^2$
y_3^{cam}	$-0.18277\,\bar{\alpha}^2L^2$
FC_1	$-0.066934\,\bar{\alpha}^2L^2$

Occasionally the effect is postponed to NNLL at second order after azimuthal integration, and it shows up at NLL at third order

Conclusions

- A single shower must be accurate for different observables
  - necessary to develop a correspondence ingredients of the shower (branching probability, mapping, ordering), all-order amplitudes, and the logarithmic order
- We initiated such a study considering the family of dipole showers with local recoil
  - Asymptotic limits of the shower equations to establishing a connection to resummation
  - Differences in regions of phase space relevant for LL (subleading N_c) and NLL (leading N_c) in global, rIRC safe observables
- Ideally future developments should come with statements about how a given choice affect the all-order logarithmic structure
  - Further developments necessary to test the accuracy of a shower at all orders
  - Establish a solid basis for the development of algorithms with higher accuracy
- Impact of tuning and pre-asymptotic effects important (perhaps dominant for some designs in phenomenological applications). Still a lot to understand

Thank you for listening

CAESAR: ordering variable

- The study of the logarithmic accuracy of parton showers requires a careful comparison with resummed calculations. The starting point is to build a resummation framework that is suitable for a MC formulation

- global and recursively IRC safe observables at NLL: CAESAR [Banfi, Salam, Zanderighi ’01-’04]
  - resummation given by a shower of independent emissions off the Born legs strongly ordered in angle

e.g. $e^+e^- \to p_1\, p_2 + X$

Use observable $v_i = V(k_i)$ as evolution variable
(not strictly necessary, it leads to a simpler structure)

$$dP_n \simeq \frac{C_F^n}{n!} \prod_{i=1}^{n} \left(\frac{\alpha_s^{\rm CMW}(p_{\perp,i}^2)}{\pi} \frac{dp_{\perp,i}}{p_{\perp,i}} dz_i P_{q\to qg}(z_i) \frac{d\phi_i}{2\pi} \right)$$

$$\Sigma(L) = \int_0^{e^{-L}} dv' \frac{d\sigma}{dv'} \sim e^{-R(L)} \mathcal{F}_{\rm NLL}(\alpha_s L)$$

Do

- Sudakov radiator *R(v)* computed at NLL
- Effect of multiple emissions evaluated with LL (soft-collinear) matrix elements and observable

$$dP_n = \frac{C_F^n}{n!} \prod_{i=1}^{n} \left(\frac{2\alpha_s(p_{\perp,i}^2)}{\pi} \frac{dp_{\perp,i}}{p_{\perp,i}} \frac{dz_i}{1-z_i} \frac{d\phi_i}{2\pi} \right)$$

$$\mathcal{F}_{\rm NLL}(v) = \langle \Theta(1 - \lim_{v\to 0} \frac{V_{\rm sc}(\{\tilde{p}\},\{k_i\})}{v}) \rangle$$

Dipole showers: mapping

- The map is defined by (local recoil)

$$\tilde{p}_i + \tilde{p}_j \xrightarrow{\tilde{p}_i \to p_i + p_k} p_i + p_j + p_k$$

$$\begin{aligned} p_i^\mu &= \tilde{z}\, \tilde{p}_i^\mu + y\,(1-\tilde{z})\, \tilde{p}_j^\mu + k_\perp \\ p_k^\mu &= (1-\tilde{z})\, \tilde{p}_i^\mu + y\, \tilde{z}\, \tilde{p}_j^\mu - k_\perp^\mu \\ p_j^\mu &= (1-y)\, \tilde{p}_j^\mu \end{aligned}$$

see backup for branching probabilities

Pythia

- Evolution variable and branching:

$$v \equiv p_{\perp,\mathrm{evol}}$$

$$\rho_{\perp,\mathrm{evol}}^2 = \frac{p_{\perp,\mathrm{evol}}^2}{(\tilde{p}_i + \tilde{p}_j)^2}\,, \qquad y = \frac{\rho_{\perp,\mathrm{evol}}^2}{z(1-z)}\,, \qquad \tilde{z} = \frac{(1-z)\left(z^2 - \rho_{\perp\mathrm{evol}}^2\right)}{z\,(1-z) - \rho_{\perp\mathrm{evol}}^2}$$

$$\rho_{\perp,\mathrm{evol}} \le z \le 1 - \rho_{\perp,\mathrm{evol}}$$

- k_t and rapidity of emission w.r.t. the emitter

$$\eta = \ln \frac{(1-\tilde{z})Q}{|k_\perp|}\,, \qquad |k_\perp^2| = \frac{\left(z^2 - \rho_{\perp\mathrm{evol}}^2\right)\left((1-z)^2 - \rho_{\perp\mathrm{evol}}^2\right)}{\left(z\,(1-z) - \rho_{\perp\mathrm{evol}}^2\right)^2}$$

Dire

- Evolution variable and branching:

$$v \equiv \sqrt{t}$$

$$\kappa^2 = \frac{t}{(\tilde{p}_i + \tilde{p}_j)^2}\,, \qquad y = \frac{\kappa^2}{1-z}\,, \qquad \tilde{z} = \frac{z-y}{1-y}$$

$$\frac{1}{2} - \sqrt{\frac{1}{4} - \kappa^2} \le z \le \frac{1}{2} + \sqrt{\frac{1}{4} - \kappa^2}$$

- k_t and rapidity of emission w.r.t. the emitter

$$\eta = \ln \frac{(1-\tilde{z})Q}{|k_\perp|}\,, \qquad |k_\perp^2| = (1-z)\,\frac{z(1-z) - \kappa^2}{(1-z-\kappa^2)^2}\, t$$

Dipole showers: branchings

- We focus on k_t-ordered dipole showers with local recoil (most common design today)
  - Consider the designs of `Pythia8`'s shower and `Dire`. The map is defined by

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`Pythia`

$$d\mathcal{P}_{q\to qg} = \frac{\alpha_s(p^2_{\perp,\mathrm{evol}})}{2\pi} \frac{dp^2_{\perp\mathrm{evol}}}{p^2_{\perp\mathrm{evol}}}\, dz\, \frac{d\phi}{2\pi}\, C_F \left(\frac{1+z^2}{1-z}\right)$$

$$d\mathcal{P}_{g\to gg} = \frac{\alpha_s(p^2_{\perp,\mathrm{evol}})}{2\pi} \frac{dp^2_{\perp\mathrm{evol}}}{p^2_{\perp\mathrm{evol}}}\, dz\, \frac{d\phi}{2\pi}\, \frac{C_A}{2} \left[\frac{1+z^3}{1-z}\right]$$

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$$\mathcal{D} = (1-x)^2(1+x)\,, \qquad x \equiv \frac{(p_i+p_k)^2}{(\tilde{p}_i+\tilde{p}_j)^2}$$

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Dire

$$\begin{aligned} d\mathcal{P}_{q\to qg} &= \frac{\alpha_s(t)}{2\pi}\,\frac{dt}{t}\,dz\,\frac{d\phi}{2\pi}\,C_F\,\left[2\frac{1-z}{(1-z)^2+\kappa^2} - (1+z)\right] \\ d\mathcal{P}_{g\to gg} &= \frac{\alpha_s(t)}{2\pi}\,\frac{dt}{t}\,dz\,\frac{d\phi}{2\pi}\,\frac{C_A}{2}\,\left[2\frac{1-z}{(1-z)^2+\kappa^2} - 2 + z(1-z)\right] \\ d\mathcal{P}_{g\to q\bar{q}} &= \frac{\alpha_s(t)}{2\pi}\,\frac{dt}{t}\,dz\,\frac{d\phi}{2\pi}\,\frac{n_f T_R}{2}\,\left[1-2z(1-z)\right] \end{aligned}$$

Difference between shower and NLL

$$\delta\Sigma^{(2\,\text{emissions})}(L) = \left(C_F \frac{2\alpha_s}{\pi}\right)^2 \int_0^1 \frac{dv_1}{v_1} \int_{\ln v_1}^{\ln 1/v_1} d\eta_1 \int_0^{v_1} \frac{dv_2}{v_2} \int_{\ln v_2}^{\ln 1/v_2} d\eta_2 \int_0^{2\pi} \frac{d\phi_1}{2\pi} \int_0^{2\pi} \frac{d\phi_2}{2\pi} \times$$
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$$\delta\Sigma^{(3\,\text{emissions})}(L) = \left(C_F \frac{2\alpha_s}{\pi}\right)^3 \int_0^1 \frac{dv_1}{v_1} \int_0^{v_1} \frac{dv_2}{v_2} \int_0^{v_2} \frac{dv_3}{v_3} \int_{\ln v_1}^{\ln 1/v_1} d\eta_1 \int_{\ln v_2}^{\ln 1/v_2} d\eta_2 \int_{\ln v_3}^{\ln 1/v_3} d\eta_3 \times$$
$$\times \int_0^{2\pi} \frac{d\phi_1}{2\pi} \int_0^{2\pi} \frac{d\phi_2}{2\pi} \int_0^{2\pi} \frac{d\phi_3}{2\pi} \times$$
$$\times \Big[\Theta(e^{-L} - V(p_1^{\text{shower}}, p_2^{\text{shower}}, p_3)) - \Theta(e^{-L} - V(p_1^{\text{correct}}, p_2^{\text{correct}}, p_3))$$
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