

Global SMEFT Fit

Tevong You



Based on J. Ellis, C. Murphy, V. Sanz and TY, JHEP (1803.03252)

Introduction

- Motivation
- SM EFT framework
- EWPT and TGCs
- Updated Run 2 SM EFT fit
- Conclusion


Why SM EFT?

Assuming a SM Higgs and **decoupled new physics** at higher energies, the SM EFT is the next phenomenological framework

The TeV Scale

What effective theory captures everything we know experimentally about weak interactions?

1933–1982 4-fermion interactions



$$\sim G_F E^2 \quad \Rightarrow \Lambda \sim \text{TeV}$$

1982–2011 SM without Higgs



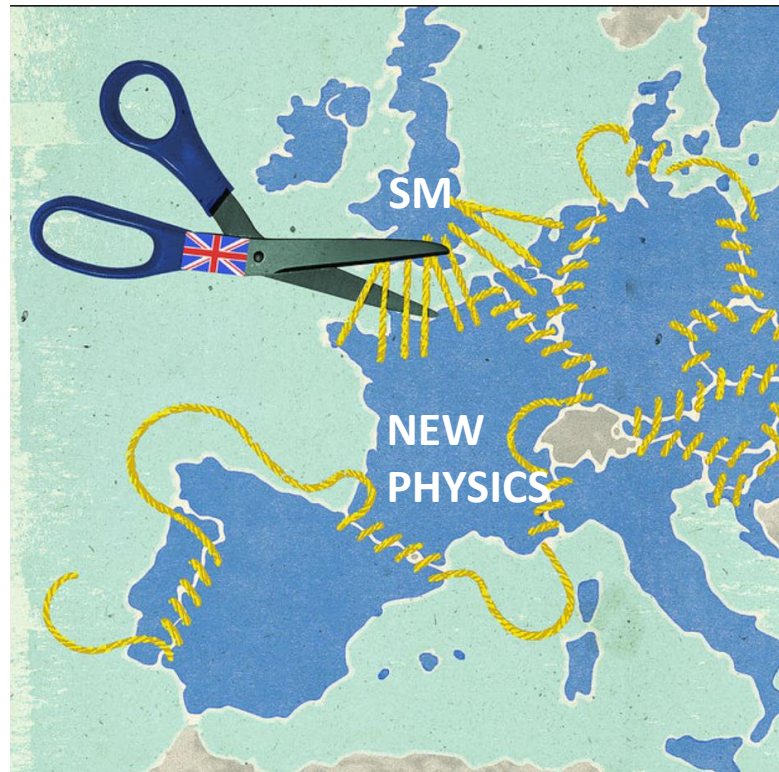
$$\sim \frac{g^2 E^2}{m_W^2} \quad \Rightarrow \Lambda \sim \text{TeV}$$

2012–now SM + higher-dimension operators?

$$\Rightarrow \Lambda \lesssim M_P?$$

SMEXIT:

decoupling new physics



SMEFT framework

- New physics appear to be **decoupled at higher energies**
- Given particle content, write down *all* terms **allowed by symmetries...**

	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$
Q_L	3	2	$\frac{1}{6}$
q_R^u	3	1	$\frac{2}{3}$
q_R^d	3	1	$-\frac{1}{3}$
L_L	1	2	$-\frac{1}{2}$
l_R	1	1	-1
ϕ	1	2	$\frac{1}{2}$



$$\mathcal{L}_{SM} = \mathcal{L}_m + \mathcal{L}_g + \mathcal{L}_h + \mathcal{L}_y \quad ,$$

$$\mathcal{L}_m = \bar{Q}_L i \gamma^\mu D_\mu^L Q_L + \bar{q}_R i \gamma^\mu D_\mu^R q_R + \bar{L}_L i \gamma^\mu D_\mu^L L_L + \bar{l}_R i \gamma^\mu D_\mu^R l_R$$

$$\mathcal{L}_G = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu}$$

$$\mathcal{L}_H = (D_\mu^L \phi)^\dagger (D^{L\mu} \phi) - V(\phi)$$

$$\mathcal{L}_Y = y_d \bar{Q}_L \phi q_R^d + y_u \bar{Q}_L \phi^c q_R^u + y_L \bar{L}_L \phi l_R + \text{h.c.} \quad ,$$

- ...Including **higher-dimensional** operators!

$$+ \quad \boxed{\mathcal{L}_{SM}^{\text{dim-6}} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i}$$

- Generated by **new physics at scale $\Lambda \gg v$**

SMEFT framework

- Unique dim-5 Weinberg operator, violates lepton number, predicts neutrino masses
- Modulo flavour structure, there are **59** dim-6 (CP-even) operators in a **non-redundant** basis

Buchmuller and Wyler (Nucl. Phys. B 268 (1986) 621)
 Gradkowski et al [arXiv:1008.4884]

- **~19** operators relevant for EWPT, TGC, and Higgs physics

EWPTs	Higgs Physics	TGCs
$\mathcal{O}_W = \frac{ig}{2} \left(H^\dagger \overset{\leftrightarrow}{D}^\mu H \right) D^\nu W_{\mu\nu}^a$		
$\mathcal{O}_B = \frac{ig'}{2} \left(H^\dagger \overset{\leftrightarrow}{D}^\mu H \right) \partial^\nu B_{\mu\nu}$		$\mathcal{O}_{3W} = g \frac{\epsilon_{abc}}{3!} W_\mu^{a\nu} W_{\nu\rho}^b W^{c\rho\mu}$
$\mathcal{O}_T = \frac{1}{2} \left(H^\dagger \overset{\leftrightarrow}{D}_\mu H \right)^2$	$\mathcal{O}_{HW} = ig(D^\mu H)^\dagger \sigma^a (D^\nu H) W_{\mu\nu}^a$	
$\mathcal{O}_{LL}^{(3)l} = (\bar{L}_L \sigma^a \gamma^\mu L_L) (\bar{L}_L \sigma^a \gamma_\mu L_L)$	$\mathcal{O}_{HB} = ig'(D^\mu H)^\dagger (D^\nu H) B_{\mu\nu}$	
$\mathcal{O}_R^e = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{e}_R \gamma^\mu e_R)$	$\mathcal{O}_g = g_s^2 H ^2 G_{\mu\nu}^A G^{A\mu\nu}$	
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$\mathcal{O}_L^{(3)q} = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \sigma^a \gamma^\mu Q_L)$	$\mathcal{O}_f = y_f H ^2 \bar{F}_L H^{(c)} f_R + \text{h.c.}$	
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In SILH basis (Giudice et al. hep-ph/0703164), adopted from Pomarol and Riva (1308.1426)

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Combinations of operators constrained in EWPT more easily set to zero in Higgs and TGCs

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Operators
 constrained by
 measurements at
 per cent level or
 worse

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Operators benefit from per mille precision at LEP

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LEP EWPT Example

- (Pseudo-)Observables

$$\Gamma_Z^L = \Gamma_{had}^L + 3\Gamma_e^L + 3\Gamma_\nu^L \quad R_l = \frac{\Gamma_{had}^L}{\Gamma_Z^L} \quad \sigma_{had} = 12\pi \frac{\Gamma_e^L \Gamma_{had}^L}{\hat{m}_Z^2 \Gamma_Z^2} \quad A_{FB}^f = \frac{3}{4} A_e A_f \quad M_W = c_W M_Z$$

$$R_q = \frac{\Gamma_q}{\Gamma_{had}^L}$$

- Depends on

$$\Gamma_f = \frac{\sqrt{2} G_F M_Z^2 \hat{M}_Z}{G\pi} \left[(g_L^f)^2 + (g_R^f)^2 \right] \quad A_f = \frac{(g_L^f)^2 - (g_R^f)^2}{(g_L^f)^2 + (g_R^f)^2}$$

$$g^f = T_f^3 - Q_f s_W^2$$

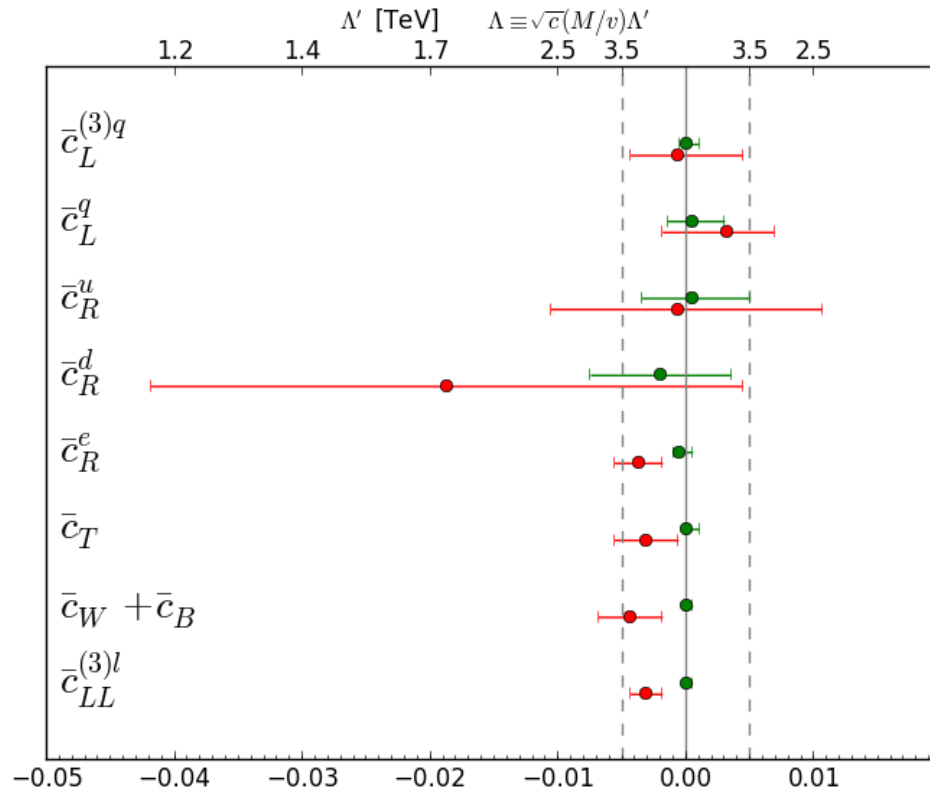
$$s_W^2 \equiv \frac{1}{2} - \frac{1}{2} \sqrt{1 - \frac{4\pi\alpha}{\sqrt{2}G_F M_Z^2}}$$

- Dim-6 operators can modify observables directly through Zff couplings contributions or indirectly through redefinitions of input observables

$$m_Z^2 = (m_Z^2)^0 (1 + \pi_{ZZ}^0) \quad G_F = G_F^0 (1 - \pi_{uw}^0) \quad \alpha(m_Z) = \alpha^0(m_Z) (1 + \pi_{\gamma\gamma}^0)$$

LEP EWPT Example

- Individual (green) and marginalised (red) 95% CL limits



Ellis, Sanz and T.Y. 1410.7703

- 8 (combinations of) operators probed by EWPT

Triple-Gauge-Couplings in Diboson

- Assume SM Z and W couplings to fermions in diboson measurements
- Interpret in anomalous TGC framework:

$$\mathcal{L}_{\text{TGC}} =$$

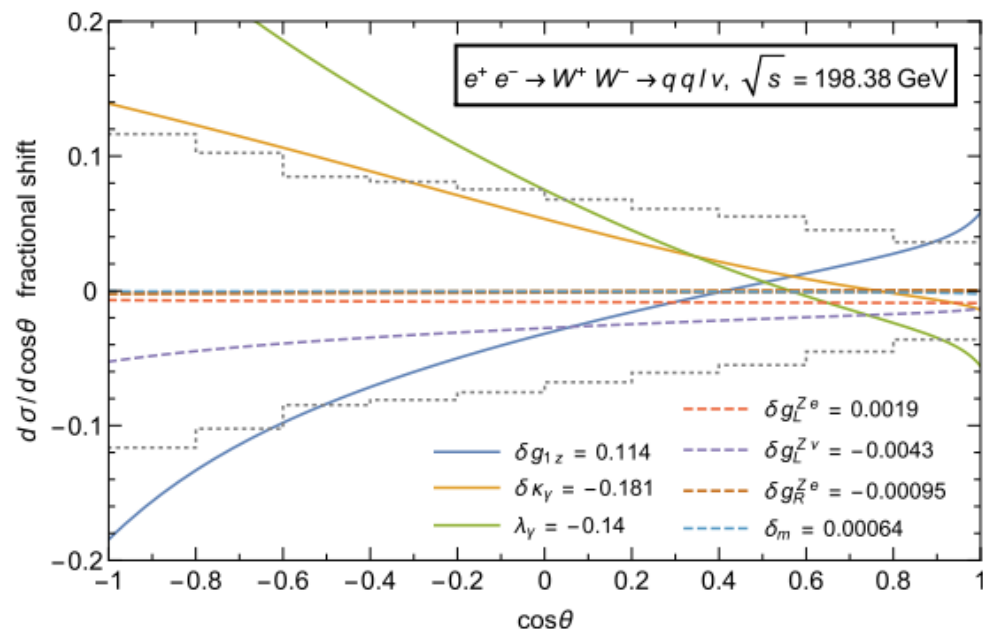
$$ig \left\{ (W_{\mu\nu}^+ W^{-\mu} - W_{\mu\nu}^- W^{+\mu}) [(1 + \delta g_{1z}) c_\theta Z^\nu + s_\theta A^\nu] \right.$$

$$+ \frac{1}{2} W_{[\mu}^+ W_{\nu]}^- [(1 + \delta \kappa_z) c_\theta Z^{\mu\nu} + (1 + \delta \kappa_\gamma) s_\theta A^{\mu\nu}]$$

$$\left. + \frac{1}{m_W^2} W_\mu^{+\nu} W_\nu^{-\rho} (\lambda_z c_\theta Z_\rho^\mu + \lambda_\gamma s_\theta A_\rho^\mu) \right\},$$

Z. Zhang, 1610.01618

- Justified at LEP:



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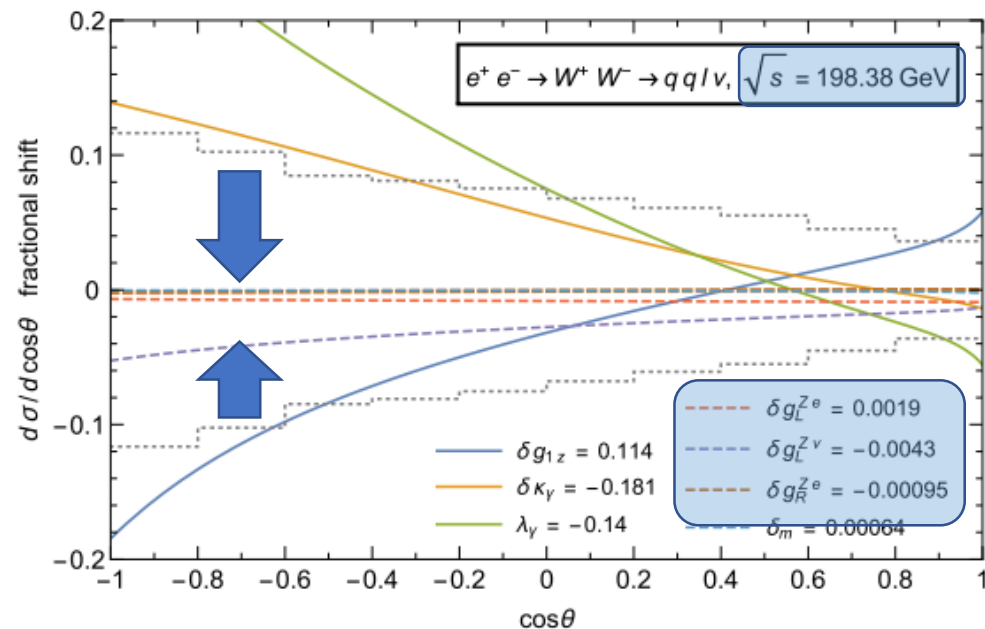
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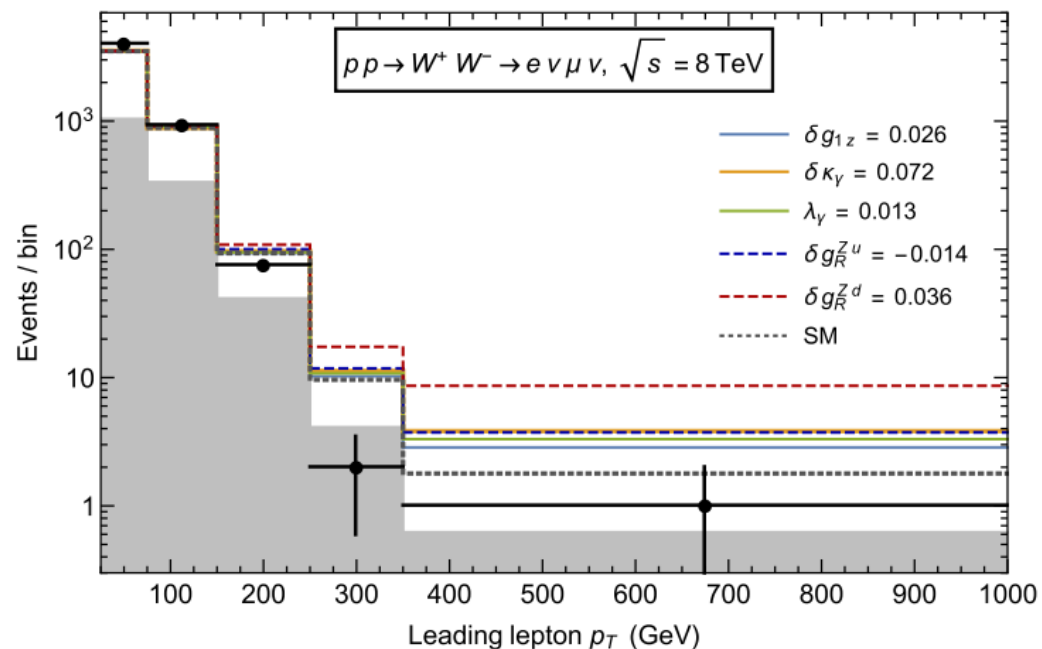
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- But not at high p_T :



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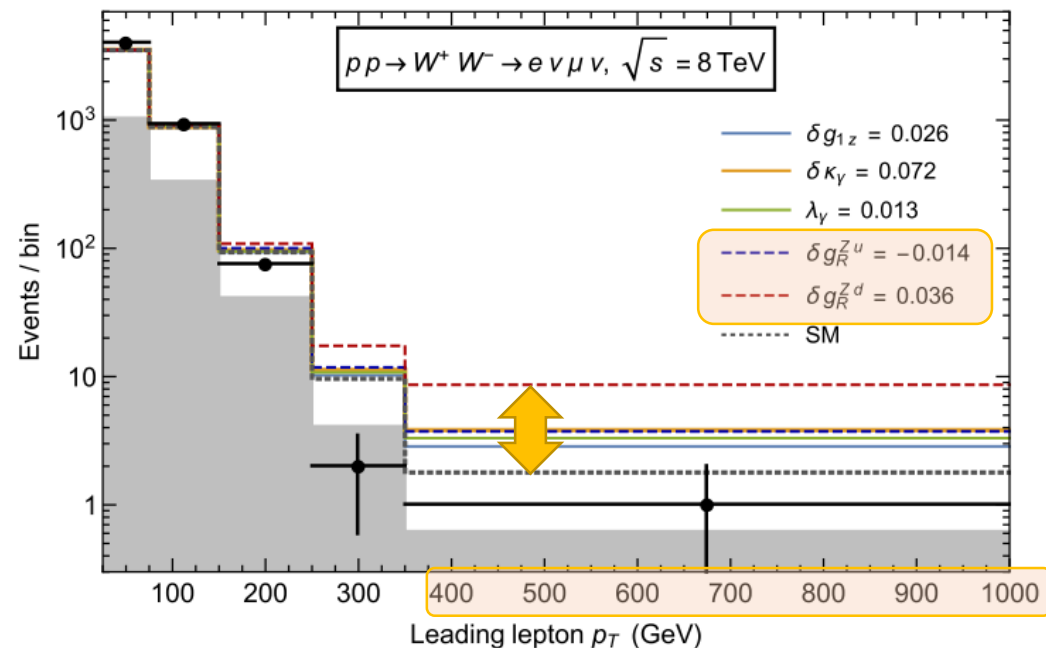
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Note: quadratic dim-6 effect
unless diboson SM-BSM
interference recovered

Azatov et al 1607.05236, 1707.08060, Panico, Riva,
Wulzer 1708.07823, Bellazzini, Riva 1806.09640



Updated Global SMEFT Fit

J. Ellis, C. Murphy, V. Sanz and TY, JHEP 1803.03252

- Combine **EWPT, diboson, Higgs** data
- Fit to 20 dim-6 CP-even operators **simultaneously**
- Present results in **Warsaw** and **SILH** basis
- Match to **simplified models**

Updated Global SMEFT Fit

- SILH basis

$$\begin{aligned}
\mathcal{L}_{\text{SMEFT}}^{\text{SILH}} \supset & \frac{\bar{c}_W}{m_W^2} \frac{ig}{2} \left(H^\dagger \sigma^a \overleftrightarrow{D}^\mu H \right) D^\nu W_{\mu\nu}^a + \frac{\bar{c}_B}{m_W^2} \frac{ig'}{2} \left(H^\dagger \overleftrightarrow{D}^\mu H \right) \partial^\nu B_{\mu\nu} + \frac{\bar{c}_T}{v^2} \frac{1}{2} \left(H^\dagger \overleftrightarrow{D}_\mu H \right)^2 \\
& + \frac{\bar{c}_{ll}}{v^2} 2(\bar{L}\gamma_\mu L)(\bar{L}\gamma^\mu L) + \frac{\bar{c}_{He}}{v^2} (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{e}_R\gamma^\mu e_R) + \frac{\bar{c}_{Hu}}{v^2} (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{u}_R\gamma^\mu u_R) \\
& + \frac{\bar{c}_{Hd}}{v^2} (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{d}_R\gamma^\mu d_R) + \frac{\bar{c}'_{Hq}}{v^2} (iH^\dagger \sigma^a \overleftrightarrow{D}_\mu H)(\bar{Q}_L\sigma^a\gamma^\mu Q_L) \\
& + \frac{\bar{c}_{Hq}}{v^2} (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{Q}_L\gamma^\mu Q_L) + \frac{\bar{c}_{HW}}{m_W^2} ig(D^\mu H)^\dagger \sigma^a (D^\nu H)W_{\mu\nu}^a + \frac{\bar{c}_{HB}}{m_W^2} ig'(D^\mu H)^\dagger (D^\nu H)B_{\mu\nu} \\
& + \frac{\bar{c}_{3W}}{m_W^2} g^3 \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^b W^{c\rho\mu} + \frac{\bar{c}_g}{m_W^2} g_s^2 |H|^2 G_{\mu\nu}^A G^{A\mu\nu} + \frac{\bar{c}_\gamma}{m_W^2} g'^2 |H|^2 B_{\mu\nu} B^{\mu\nu} \\
& + \frac{\bar{c}_H}{v^2} \frac{1}{2} (\partial^\mu |H|^2)^2 - \sum_{f=e,u,d} \frac{\bar{c}_f}{v^2} y_f |H|^2 \bar{F}_L H^{(c)} f_R \\
& + \frac{\bar{c}_{3G}}{m_W^2} g_s^3 f_{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu} - \frac{\bar{c}_{uG}}{m_W^2} 4g_s y_u H^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} T_a u_R G_{\mu\nu}^A.
\end{aligned} \tag{6}$$

Updated Global SMEFT Fit

- Warsaw basis

$$\begin{aligned}
 \mathcal{L}_{\text{SMEFT}}^{\text{Warsaw}} \supset & \frac{\bar{C}_{Hl}^{(3)}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu^I H) (\bar{l} \tau^I \gamma^\mu l) + \frac{\bar{C}_{Hl}^{(1)}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{l} \gamma^\mu l) + \frac{\bar{C}_{ll}}{v^2} (\bar{l} \gamma_\mu l) (\bar{l} \gamma^\mu l) \\
 & + \frac{\bar{C}_{HD}}{v^2} |H^\dagger D_\mu H|^2 + \frac{\bar{C}_{HWB}}{v^2} H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu} \\
 & + \frac{\bar{C}_{He}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{e} \gamma^\mu e) + \frac{\bar{C}_{Hu}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{u} \gamma^\mu u) + \frac{\bar{C}_{Hd}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{d} \gamma^\mu d) \\
 & + \frac{\bar{C}_{Hq}^{(3)}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu^I H) (\bar{q} \tau^I \gamma^\mu q) + \frac{\bar{C}_{Hq}^{(1)}}{v^2} (H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{q} \gamma^\mu q) + \frac{\bar{C}_W}{v^2} \epsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_{\text{SMEFT}}^{\text{Warsaw}} \supset & \frac{\bar{C}_{eH}}{v^2} (H^\dagger H) (\bar{l} e H) + \frac{\bar{C}_{dH}}{v^2} (H^\dagger H) (\bar{q} d H) + \frac{\bar{C}_{uH}}{v^2} (H^\dagger H) (\bar{q} u \tilde{H}) \\
 & + \frac{\bar{C}_G}{v^2} f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu} + \frac{\bar{C}_{H\Box}}{v^2} (H^\dagger H) \Box (H^\dagger H) + \frac{\bar{C}_{uG}}{v^2} (\bar{q} \sigma^{\mu\nu} T^A u) \tilde{H} G_{\mu\nu}^A \\
 & + \frac{\bar{C}_{HW}}{v^2} H^\dagger H W_{\mu\nu}^I W^{I\mu\nu} + \frac{\bar{C}_{HB}}{v^2} H^\dagger H B_{\mu\nu} B^{\mu\nu} + \frac{\bar{C}_{HG}}{v^2} H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}.
 \end{aligned}$$

Observables

- LEP and SLC EWPTs, M_W from ATLAS, Tevatron

Observable	Measurement	Ref.	SM Prediction	Ref.
Γ_Z [GeV]	2.4952 ± 0.0023	[39]	2.4943 ± 0.0005	[38]
σ_{had}^0 [nb]	41.540 ± 0.037	[39]	41.488 ± 0.006	[38]
R_ℓ^0	20.767 ± 0.025	[39]	20.752 ± 0.005	[38]
$A_{\text{FB}}^{0,\ell}$	0.0171 ± 0.0010	[39]	0.01622 ± 0.00009	[114]
$\mathcal{A}_\ell(P_\tau)$	0.1465 ± 0.0033	[39]	0.1470 ± 0.0004	[114]
$\mathcal{A}_\ell(\text{SLD})$	0.1513 ± 0.0021	[39]	0.1470 ± 0.0004	[114]
R_b^0	0.021629 ± 0.00066	[39]	0.2158 ± 0.00015	[38]
R_c^0	0.1721 ± 0.0030	[39]	0.17223 ± 0.00005	[38]
$A_{\text{FB}}^{0,b}$	0.0992 ± 0.0016	[39]	0.1031 ± 0.0003	[114]
$A_{\text{FB}}^{0,c}$	0.0707 ± 0.0035	[39]	0.0736 ± 0.0002	[114]
\mathcal{A}_b	0.923 ± 0.020	[39]	0.9347	[114]
\mathcal{A}_c	0.670 ± 0.027	[39]	0.6678 ± 0.0002	[114]
M_W [GeV]	80.387 ± 0.016	[40]	80.361 ± 0.006	[114]
M_W [GeV]	80.370 ± 0.019	[94]	80.361 ± 0.006	[114]

- LEP WW measurements
- ATLAS WW high p_T overflow bin

Observables

- ATLAS+CMS Higgs Run 1

Production	Decay	Signal Strength	Production	Decay	Signal Strength
ggF	$\gamma\gamma$	$1.10^{+0.23}_{-0.22}$	Wh	$\tau\tau$	-1.4 ± 1.4
ggF	ZZ	$1.13^{+0.34}_{-0.31}$	Wh	bb	1.0 ± 0.5
ggF	WW	0.84 ± 0.17	Zh	$\gamma\gamma$	$0.5^{+3.0}_{-2.5}$
ggF	$\tau\tau$	1.0 ± 0.6	Zh	WW	$5.9^{+2.6}_{-2.2}$
VBF	$\gamma\gamma$	1.3 ± 0.5	Zh	$\tau\tau$	$2.2^{+2.2}_{-1.8}$
VBF	ZZ	$0.1^{+1.1}_{-0.6}$	Zh	bb	0.4 ± 0.4
VBF	WW	1.2 ± 0.4	tth	$\gamma\gamma$	$2.2^{+1.6}_{-1.3}$
VBF	$\tau\tau$	1.3 ± 0.4	tth	WW	$5.0^{+1.8}_{-1.7}$
Wh	$\gamma\gamma$	$0.5^{+1.3}_{-1.2}$	tth	$\tau\tau$	$-1.9^{+3.7}_{-3.3}$
Wh	WW	$1.6^{+1.2}_{-1.0}$	tth	bb	1.1 ± 1.0
pp	$Z\gamma$	$2.7^{+4.6}_{-4.5}$	pp	$\mu\mu$	0.1 ± 2.5

Observables

- ATLAS+CMS Higgs Run 2

	Production	Decay	Sig. Stren.		Production	Decay	Sig. Stren.
[96]	1-jet, $p_T > 450$	$b\bar{b}$	$2.3^{+1.8}_{-1.6}$	[105]	pp	$\mu\mu$	-0.1 ± 1.4
[97]	Zh	$b\bar{b}$	0.9 ± 0.5	[106]	Zh	$b\bar{b}$	$0.69^{+0.35}_{-0.31}$
[97]	Wh	$b\bar{b}$	1.7 ± 0.7	[106]	Wh	$b\bar{b}$	$1.21^{+0.45}_{-0.42}$
[98]	$t\bar{t}h$	$b\bar{b}$	$-0.19^{+0.80}_{-0.81}$	[107]	$t\bar{t}h$	$b\bar{b}$	$0.84^{+0.64}_{-0.61}$
[99]	$t\bar{t}h$	$1\ell + 2\tau_h$	$-1.20^{+1.50}_{-1.47}$	[108]	$t\bar{t}h$	$2\ell os + 1\tau_h$	$1.7^{+2.1}_{-1.9}$
[99]	$t\bar{t}h$	$2\ell ss + 1\tau_h$	$0.86^{+0.79}_{-0.66}$	[108]	$t\bar{t}h$	$1\ell + 2\tau_h$	$-0.6^{+1.6}_{-1.5}$
[99]	$t\bar{t}h$	$3\ell + 1\tau_h$	$1.22^{+1.34}_{-1.00}$	[108]	$t\bar{t}h$	$3\ell + 1\tau_h$	$1.6^{+1.8}_{-1.3}$
[100]	$t\bar{t}h$	$2\ell ss$	$1.7^{+0.6}_{-0.5}$	[108]	$t\bar{t}h$	$2\ell ss + 1\tau_h$	$3.5^{+1.7}_{-1.3}$
[100]	$t\bar{t}h$	3ℓ	$1.0^{+0.8}_{-0.7}$	[108]	$t\bar{t}h$	3ℓ	$1.8^{+0.9}_{-0.7}$
[100]	$t\bar{t}h$	4ℓ	$0.9^{+2.3}_{-1.6}$	[108]	$t\bar{t}h$	$2\ell ss$	$1.5^{+0.7}_{-0.6}$
[101]	0-jet	WW	$0.9^{+0.4}_{-0.3}$	[109]	VBF	WW	$1.7^{+1.1}_{-0.9}$
[101]	1-jet	WW	1.1 ± 0.4	[109]	Wh	WW	$3.2^{+4.4}_{-4.2}$
[101]	2-jet	WW	1.3 ± 1.0	[110]	$B(h \rightarrow \gamma\gamma) / B(h \rightarrow 4\ell)$		$0.69^{+0.15}_{-0.13}$
[101]	VBF 2-jet	WW	1.4 ± 0.8	[110]	0-jet	4ℓ	$1.07^{+0.27}_{-0.25}$
[101]	Vh 2-jet	WW	$2.1^{+2.3}_{-2.2}$	[110]	1-jet, $p_T < 60$	4ℓ	$0.67^{+0.72}_{-0.68}$
[101]	Wh 3-lep	WW	-1.4 ± 1.5	[110]	1-jet, $p_T \in (60, 120)$	4ℓ	$1.00^{+0.63}_{-0.55}$
[102]	ggF	$\gamma\gamma$	$1.11^{+0.19}_{-0.18}$	[110]	1-jet, $p_T \in (120, 200)$	4ℓ	$2.1^{+1.5}_{-1.3}$
[102]	VBF	$\gamma\gamma$	$0.5^{+0.6}_{-0.5}$	[110]	2-jet	4ℓ	$2.2^{+1.1}_{-1.0}$
[102]	$t\bar{t}h$	$\gamma\gamma$	2.2 ± 0.9	[110]	“BSM-like”	4ℓ	$2.3^{+1.2}_{-1.0}$
[102]	Vh	$\gamma\gamma$	$2.3^{+1.1}_{-1.0}$	[110]	VBF, $p_T < 200$	4ℓ	$2.14^{+0.94}_{-0.77}$
[103]	ggF	4ℓ	$1.20^{+0.22}_{-0.21}$	[110]	Vh lep	4ℓ	$0.3^{+1.3}_{-1.2}$
[104]	0-jet	$\tau\tau$	0.84 ± 0.89	[110]	$t\bar{t}h$	4ℓ	$0.51^{+0.86}_{-0.70}$
[104]	boosted	$\tau\tau$	$1.17^{+0.47}_{-0.40}$				
[104]	VBF	$\tau\tau$	$1.11^{+0.34}_{-0.35}$				

Observables

Including kinematical
information facilitated
by **STXS**

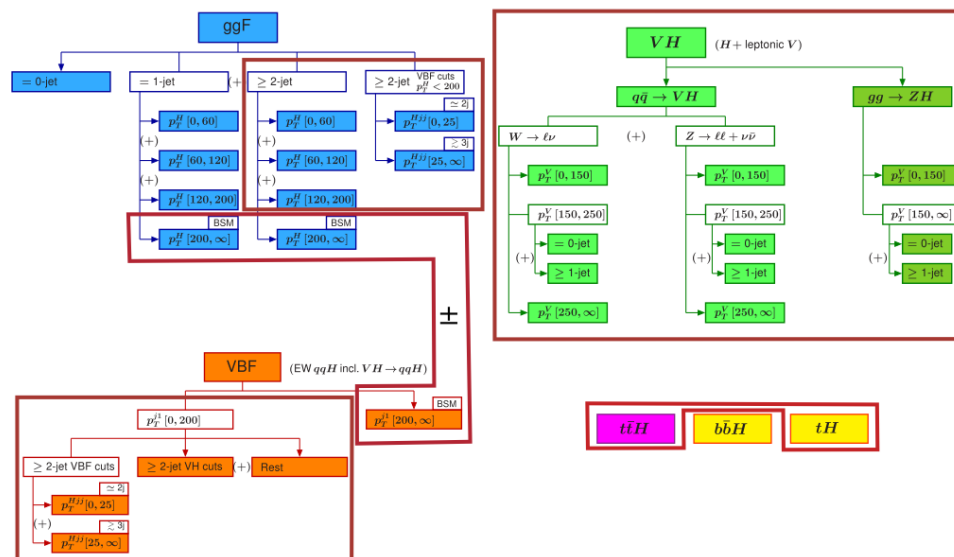
- ATLAS+CMS Higgs Run 2

	Production	Decay	Sig. Stren.		Production	Decay	Sig. Stren.
[96]	1-jet, $p_T > 450$	$b\bar{b}$	$2.3^{+1.8}_{-1.6}$	[105]	pp	$\mu\mu$	-0.1 ± 1.4
[97]	Zh	$b\bar{b}$	0.9 ± 0.5	[106]	Zh	$b\bar{b}$	$0.69^{+0.35}_{-0.31}$
[97]	Wh	$b\bar{b}$	1.7 ± 0.7	[106]	Wh	$b\bar{b}$	$1.21^{+0.45}_{-0.42}$
[98]	$t\bar{t}h$	$b\bar{b}$	$-0.19^{+0.80}_{-0.81}$	[107]	$t\bar{t}h$	$b\bar{b}$	$0.84^{+0.64}_{-0.61}$
[99]	$t\bar{t}h$	$1\ell + 2\tau_h$	$-1.20^{+1.50}_{-1.47}$	[108]	$t\bar{t}h$	$2\ell os + 1\tau_h$	$1.7^{+2.1}_{-1.9}$
[99]	$t\bar{t}h$	$2\ell ss + 1\tau_h$	$0.86^{+0.79}_{-0.66}$	[108]	$t\bar{t}h$	$1\ell + 2\tau_h$	$-0.6^{+1.6}_{-1.5}$
[99]	$t\bar{t}h$	$3\ell + 1\tau_h$	$1.22^{+1.34}_{-1.00}$	[108]	$t\bar{t}h$	$3\ell + 1\tau_h$	$1.6^{+1.8}_{-1.3}$
[100]	$t\bar{t}h$	$2\ell ss$	$1.7^{+0.6}_{-0.5}$	[108]	$t\bar{t}h$	$2\ell ss + 1\tau_h$	$3.5^{+1.7}_{-1.3}$
[100]	$t\bar{t}h$	3ℓ	$1.0^{+0.8}_{-0.7}$	[108]	$t\bar{t}h$	3ℓ	$1.8^{+0.9}_{-0.7}$
[100]	$t\bar{t}h$	4ℓ	$0.9^{+2.3}_{-1.6}$	[108]	$t\bar{t}h$	$2\ell ss$	$1.5^{+0.7}_{-0.6}$
[101]	0-jet	WW	$0.9^{+0.4}_{-0.3}$	[109]	VBF	WW	$1.7^{+1.1}_{-0.9}$
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[101]	VBF 2-jet	WW	1.4 ± 0.8	[110]	0-jet	4ℓ	$1.07^{+0.27}_{-0.25}$
[101]	Vh 2-jet	WW	$2.1^{+2.3}_{-2.2}$	[110]	1-jet, $p_T < 60$	4ℓ	$0.67^{+0.72}_{-0.68}$
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[102]	ggF	$\gamma\gamma$	$1.11^{+0.19}_{-0.18}$	[110]	1-jet, $p_T \in (120, 200)$	4ℓ	$2.1^{+1.5}_{-1.3}$
[102]	VBF	$\gamma\gamma$	$0.5^{+0.6}_{-0.5}$	[110]	2-jet	4ℓ	$2.2^{+1.1}_{-1.0}$
[102]	$t\bar{t}h$	$\gamma\gamma$	2.2 ± 0.9	[110]	“BSM-like”	4ℓ	$2.3^{+1.2}_{-1.0}$
[102]	Vh	$\gamma\gamma$	$2.3^{+1.1}_{-1.0}$	[110]	VBF, $p_T < 200$	4ℓ	$2.14^{+0.94}_{-0.77}$
[103]	ggF	4ℓ	$1.20^{+0.22}_{-0.21}$	[110]	Vh lep	4ℓ	$0.3^{+1.3}_{-1.2}$
[104]	0-jet	$\tau\tau$	0.84 ± 0.89	[110]	$t\bar{t}h$	4ℓ	$0.51^{+0.86}_{-0.70}$
[104]	boosted	$\tau\tau$	$1.17^{+0.47}_{-0.40}$				
[104]	VBF	$\tau\tau$	$1.11^{+0.34}_{-0.35}$				

STXS

- Simplified Template Cross-Sections
- Sub-division into kinematic regions for production processes

ATLAS preliminary

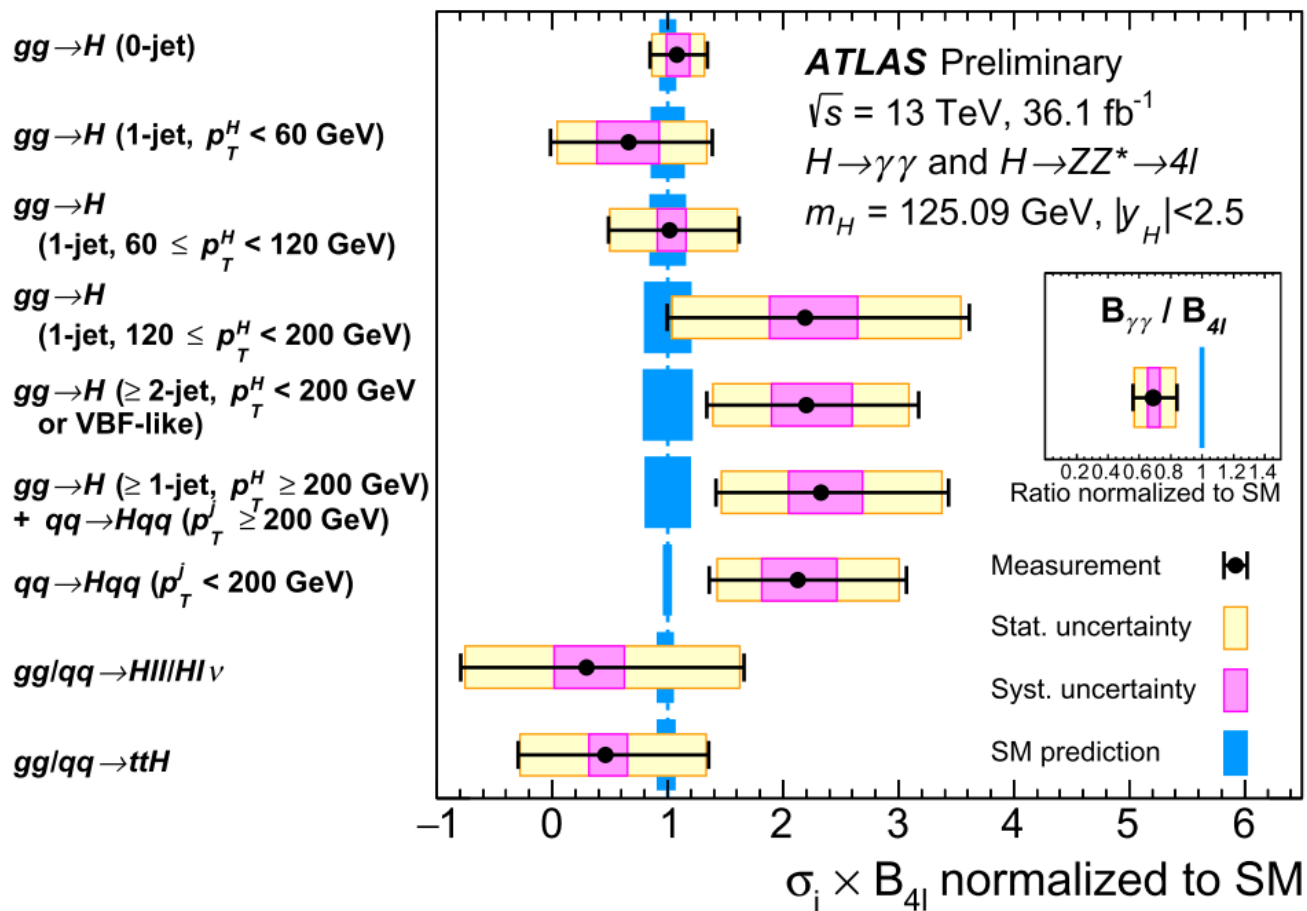


ATLAS-CONF-2017-047

- Facilitates combination and interpretation

STXS

• STXS measurements



ATLAS-CONF-2017-047

STXS

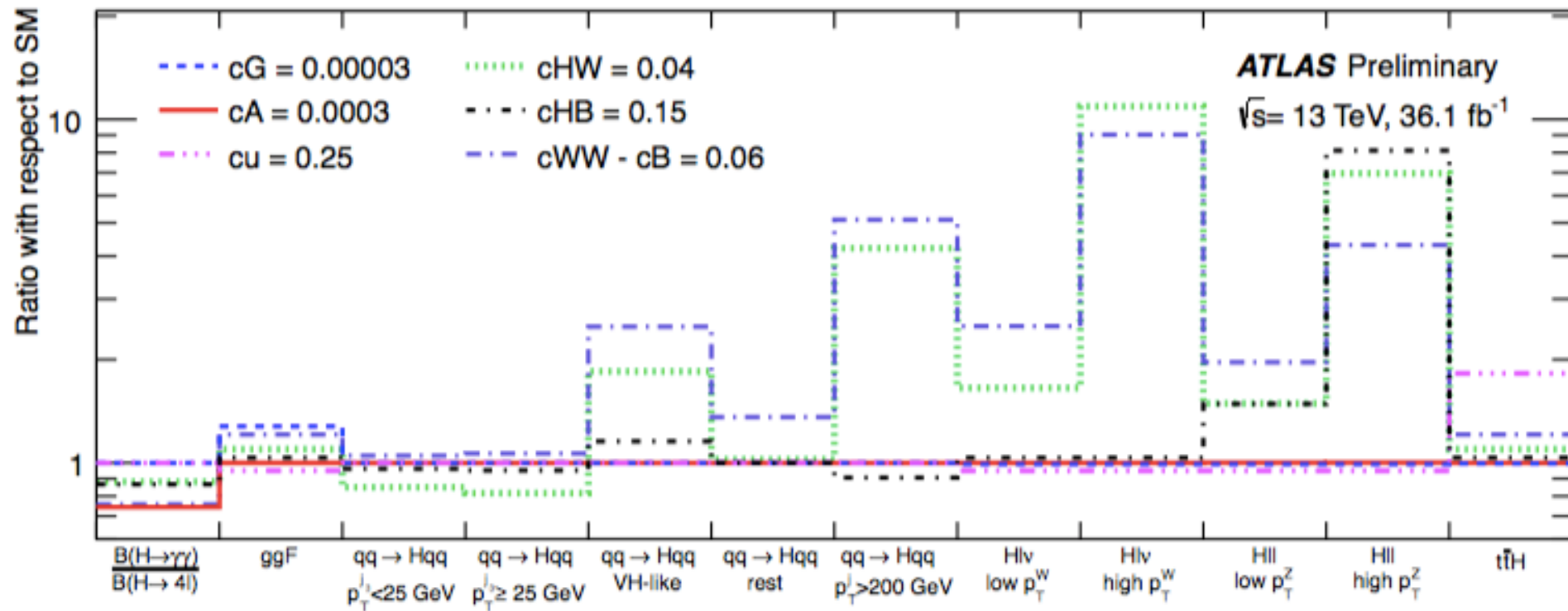
- STXS dim-6 predictions

Cross-section region	$\sum_i A_i c_i$
$gg \rightarrow H$ (0-jet)	
$gg \rightarrow H$ (1-jet, $p_T^H < 60$ GeV)	$56c'_g$
$gg \rightarrow H$ (1-jet, $60 \leq p_T^H < 120$ GeV)	
$gg \rightarrow H$ (1-jet, $120 \leq p_T^H < 200$ GeV)	$56c'_g + 18c3G + 11c2G$
$gg \rightarrow H$ (1-jet, $p_T^H \geq 200$ GeV)	$56c'_g + 52c3G + 34c2G$
$gg \rightarrow H$ (\geq 2-jet, $p_T^H < 60$ GeV)	$56c'_g$
$gg \rightarrow H$ (\geq 2-jet, $60 \leq p_T^H < 120$ GeV)	$56c'_g + 8c3G + 7c2G$
$gg \rightarrow H$ (\geq 2-jet, $120 \leq p_T^H < 200$ GeV)	$56c'_g + 23c3G + 18c2G$
$gg \rightarrow H$ (\geq 2-jet, $p_T^H \geq 200$ GeV)	$56c'_g + 90c3G + 68c2G$
$gg \rightarrow H$ (\geq 2-jet VBF-like, $p_T^{j3} < 25$ GeV)	$56c'_g$
$gg \rightarrow H$ (\geq 2-jet VBF-like, $p_T^{j3} \geq 25$ GeV)	$56c'_g + 9c3G + 8c2G$
$qq \rightarrow Hqq$ (VBF-like, $p_T^{j3} < 25$ GeV)	$-1.0cH - 1.0cT + 1.3cWW - 0.023cB - 4.3cHW$ $-0.29cHB + 0.092cHQ - 5.3cPHQ - 0.33cHu + 0.12cHd$
$qq \rightarrow Hqq$ (VBF-like, $p_T^{j3} \geq 25$ GeV)	$-1.0cH - 1.1cT + 1.2cWW - 0.027cB - 5.8cHW$ $-0.41cHB + 0.13cHQ - 6.9cPHQ - 0.45cHu + 0.15cHd$
$qq \rightarrow Hqq$ ($p_T^j \geq 200$ GeV)	$-1.0cH - 0.95cT + 1.5cWW - 0.025cB - 3.6cHW$ $-0.24cHB + 0.084cHQ - 4.5cPHQ - 0.25cHu + 0.1cHd$
$qq \rightarrow Hqq$ ($60 \leq m_{jj} < 120$ GeV)	$-0.99cH - 1.2cT + 7.8cWW - 0.19cB - 31cHW$ $-2.4cHB + 0.9cHQ - 38cPHQ - 2.8cHu + 0.9cHd$
$qq \rightarrow Hqq$ (rest)	$-1.0cH - 1.0cT + 1.4cWW - 0.028cB - 6.2cHW$ $-0.42cHB + 0.14cHQ - 6.9cPHQ - 0.42cHu + 0.16cHd$
$gg/q\bar{q} \rightarrow t\bar{t}H$	$-0.98cH + 2.9cu + 0.93cG + 310cuG$ $+27c3G - 13c2G$

Hays, Sanz, Zemaityte
[LHCHXSWG-INT-2017-01]

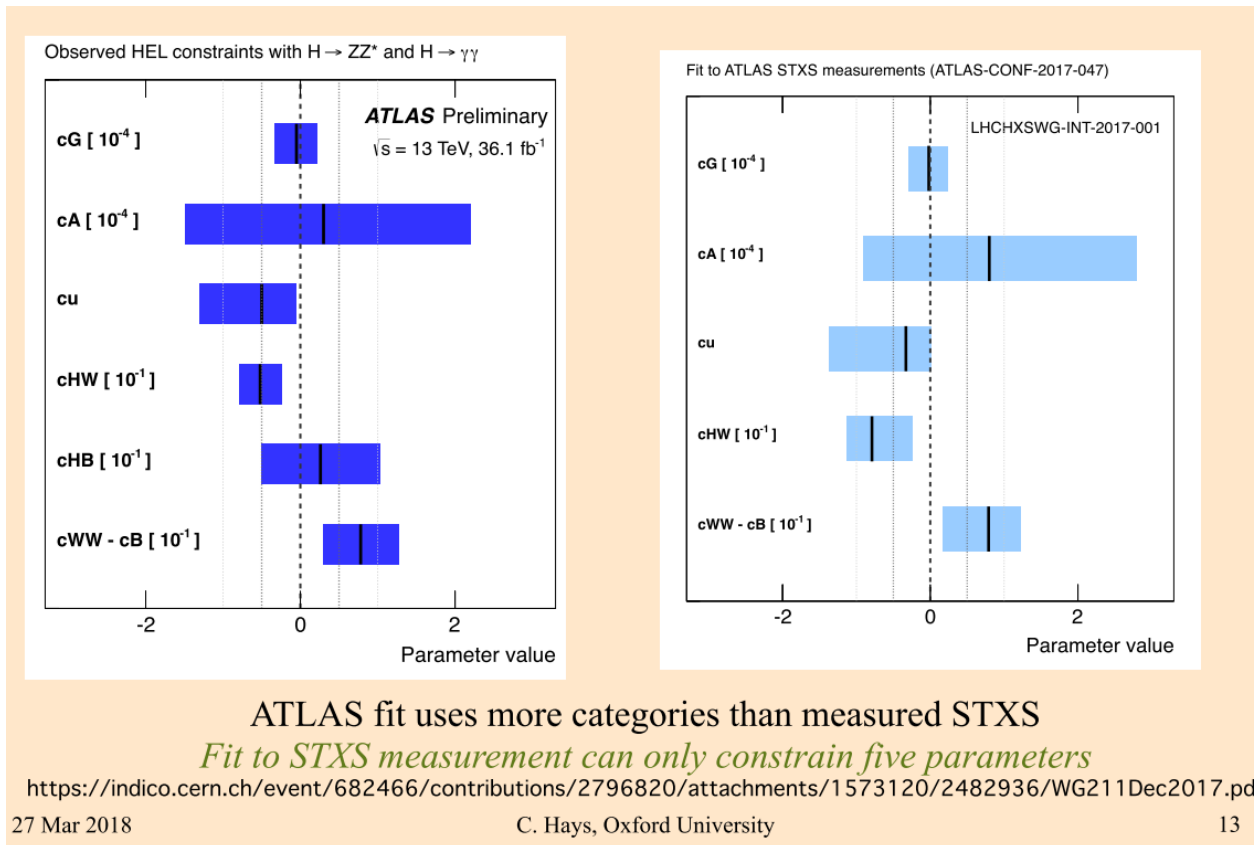
STXS

- STXS dim-6 predictions



STXS

- Good agreement with optimised non-STXS fit



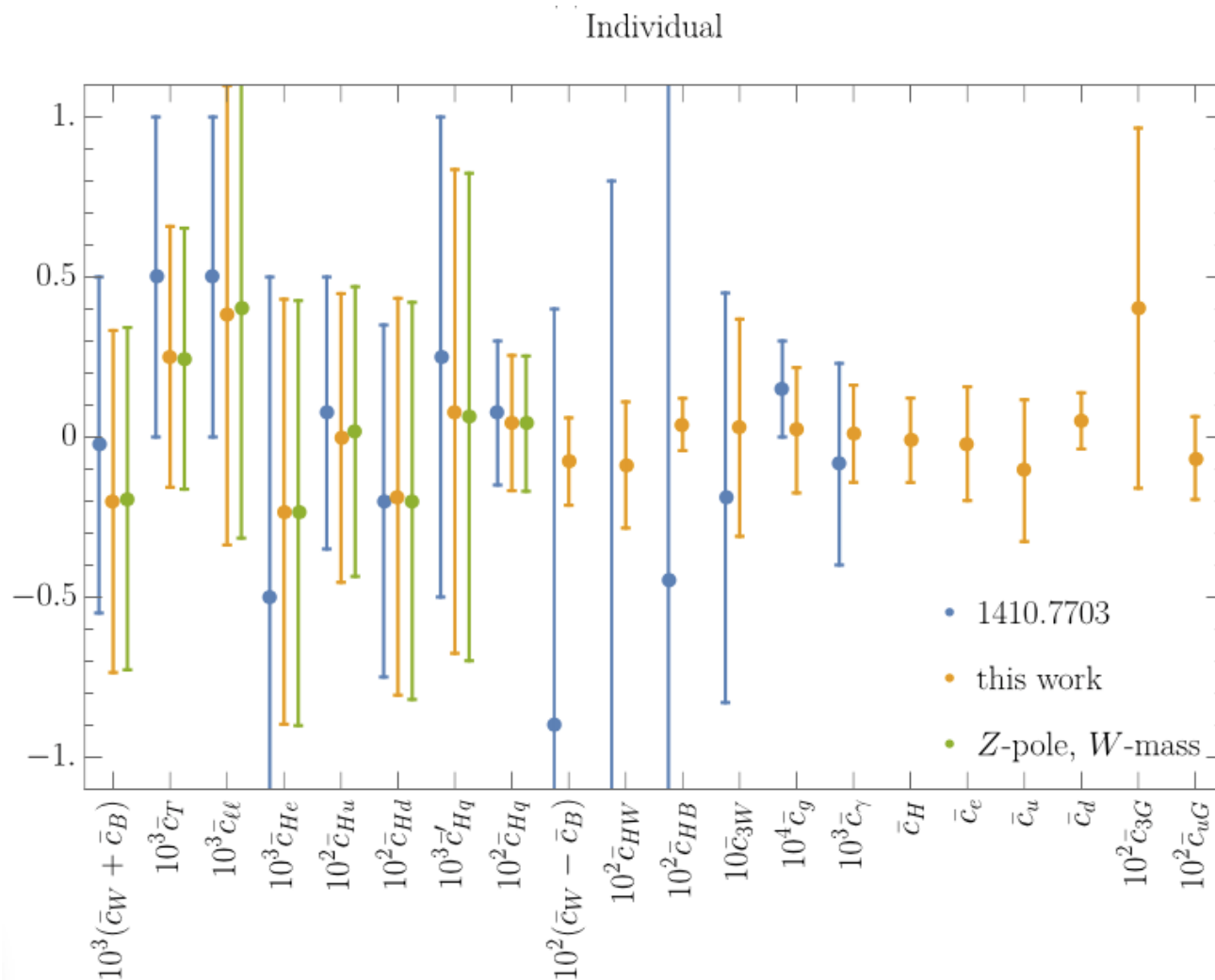
- See also **VHbb** case (ATLAS-CONF-2018-053)

De Blas, Lohwasser, Musella, Mimasu (1803.10379)

Talks by Douglas Schaefer and Chris Hays

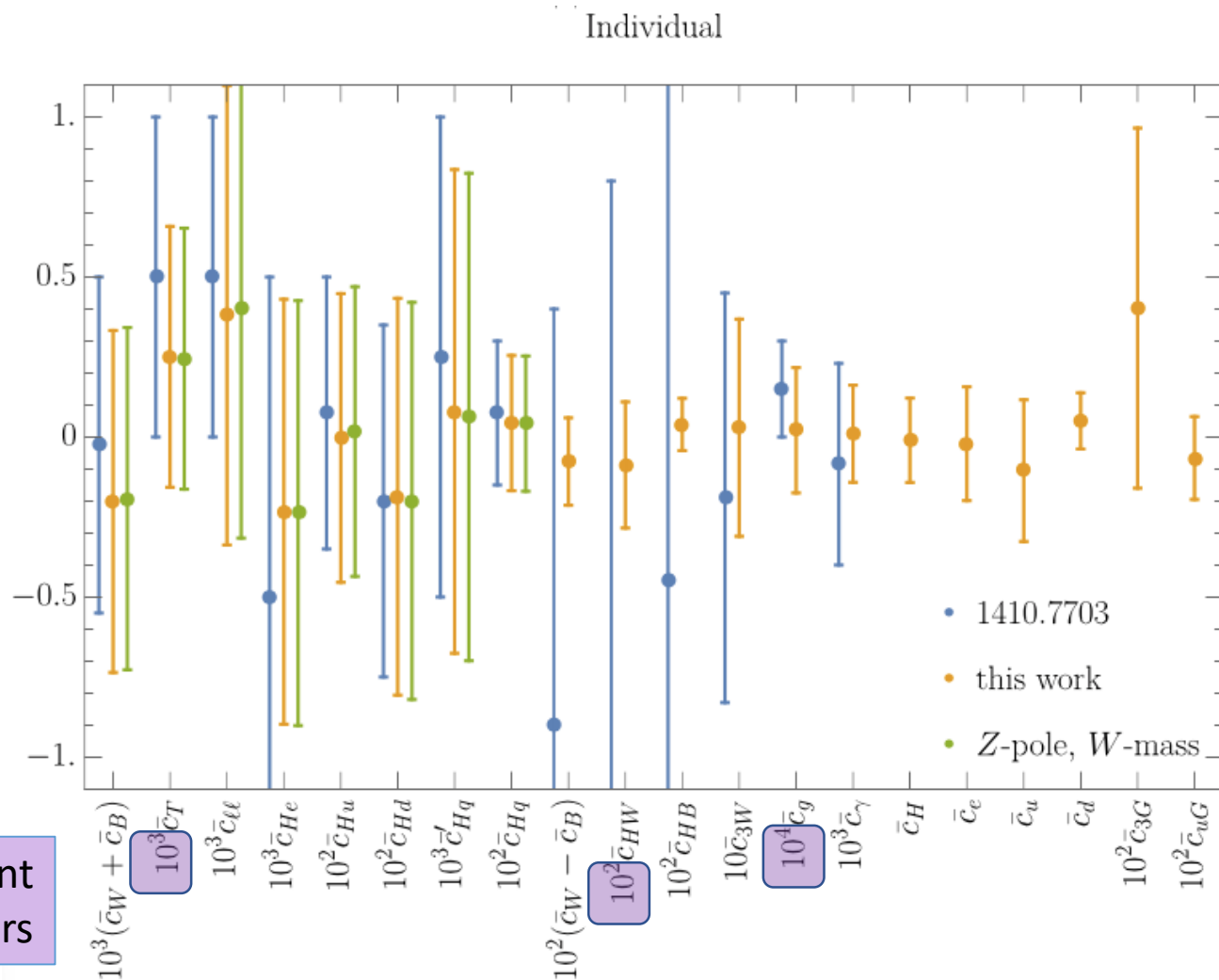
Results

- **SILH** basis, fit each operator **individually**



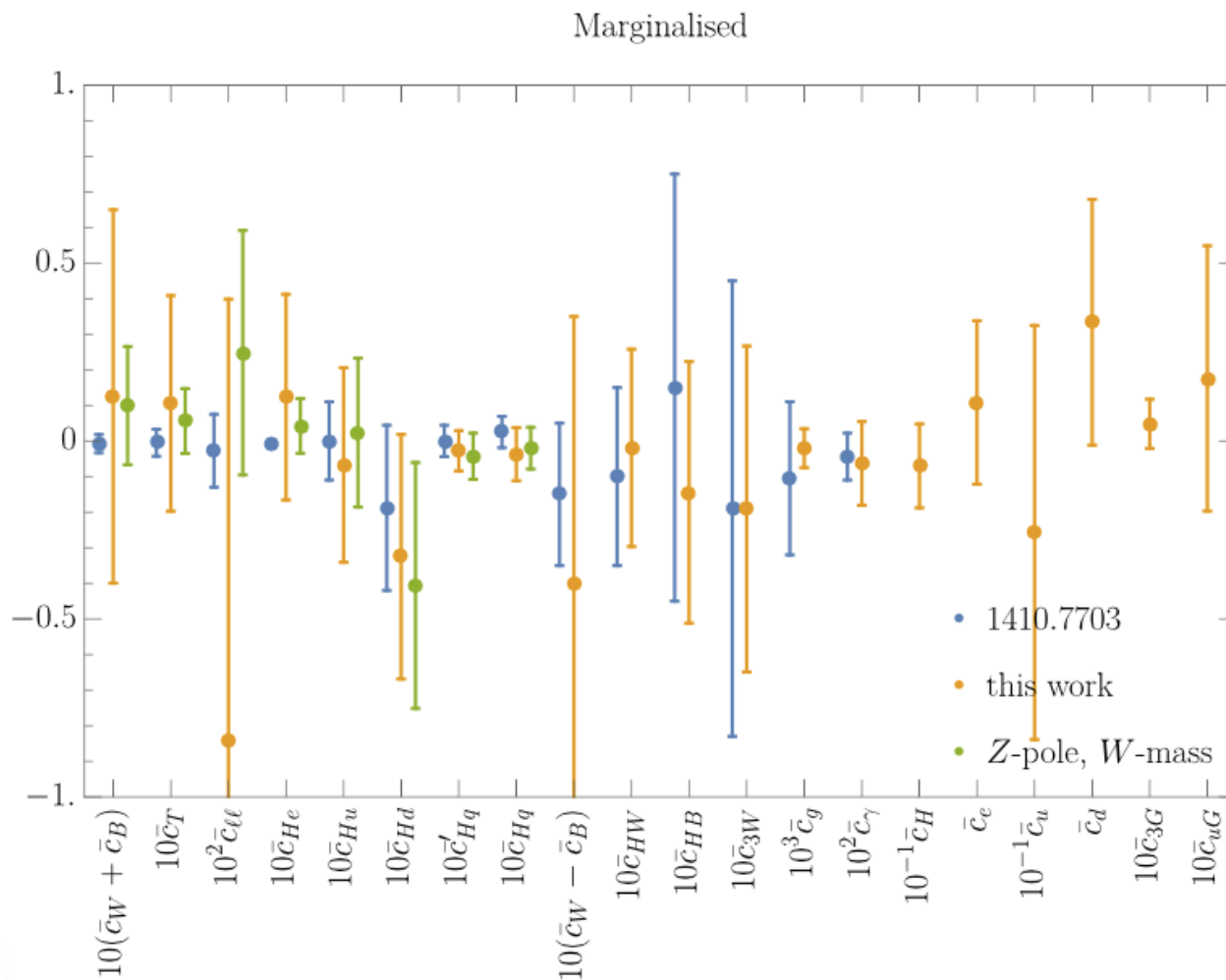
Results

- **SILH** basis, fit each operator **individually**



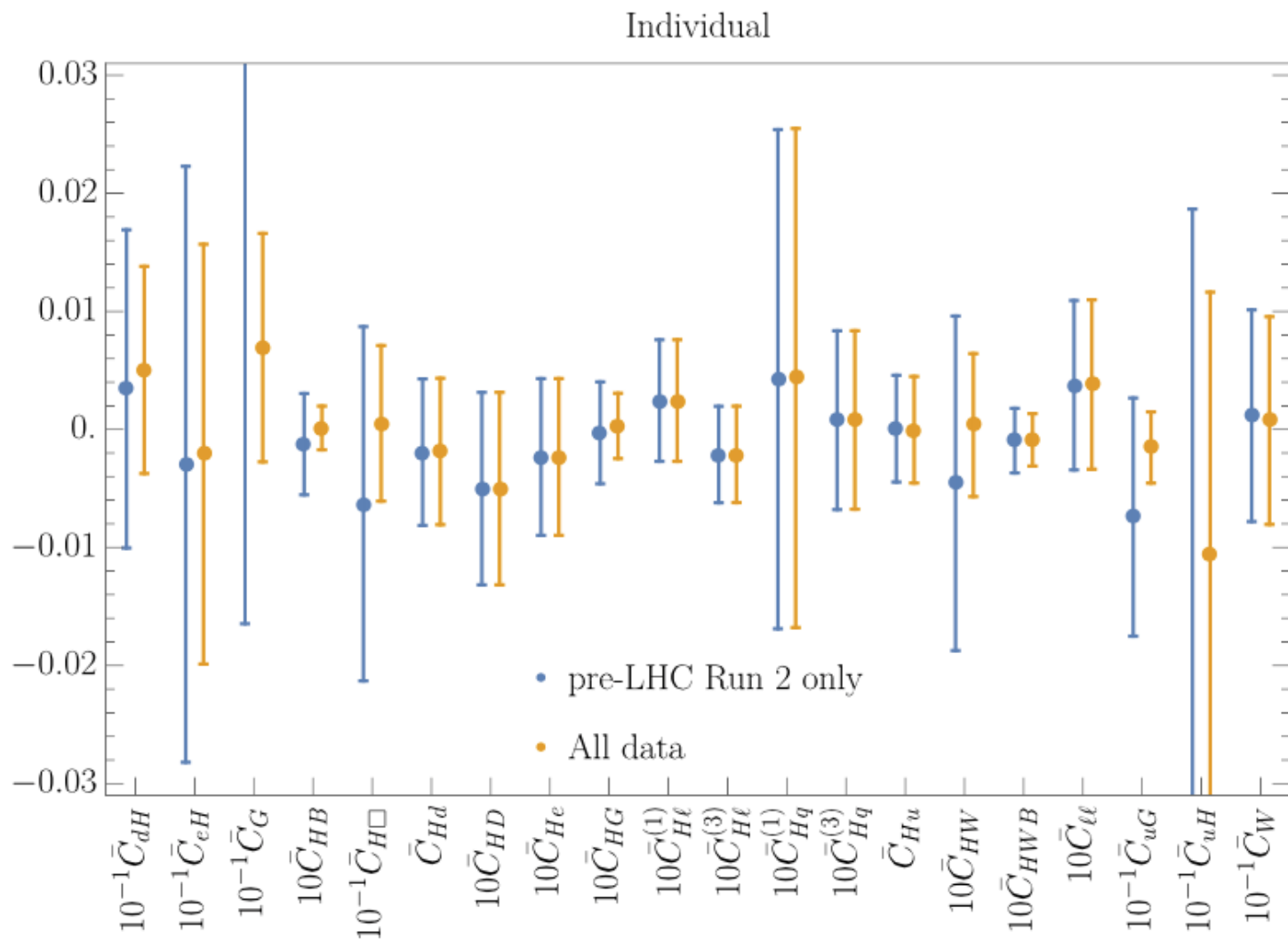
Results

- **SILH** basis, fit *all* operators **simultaneously**



Results

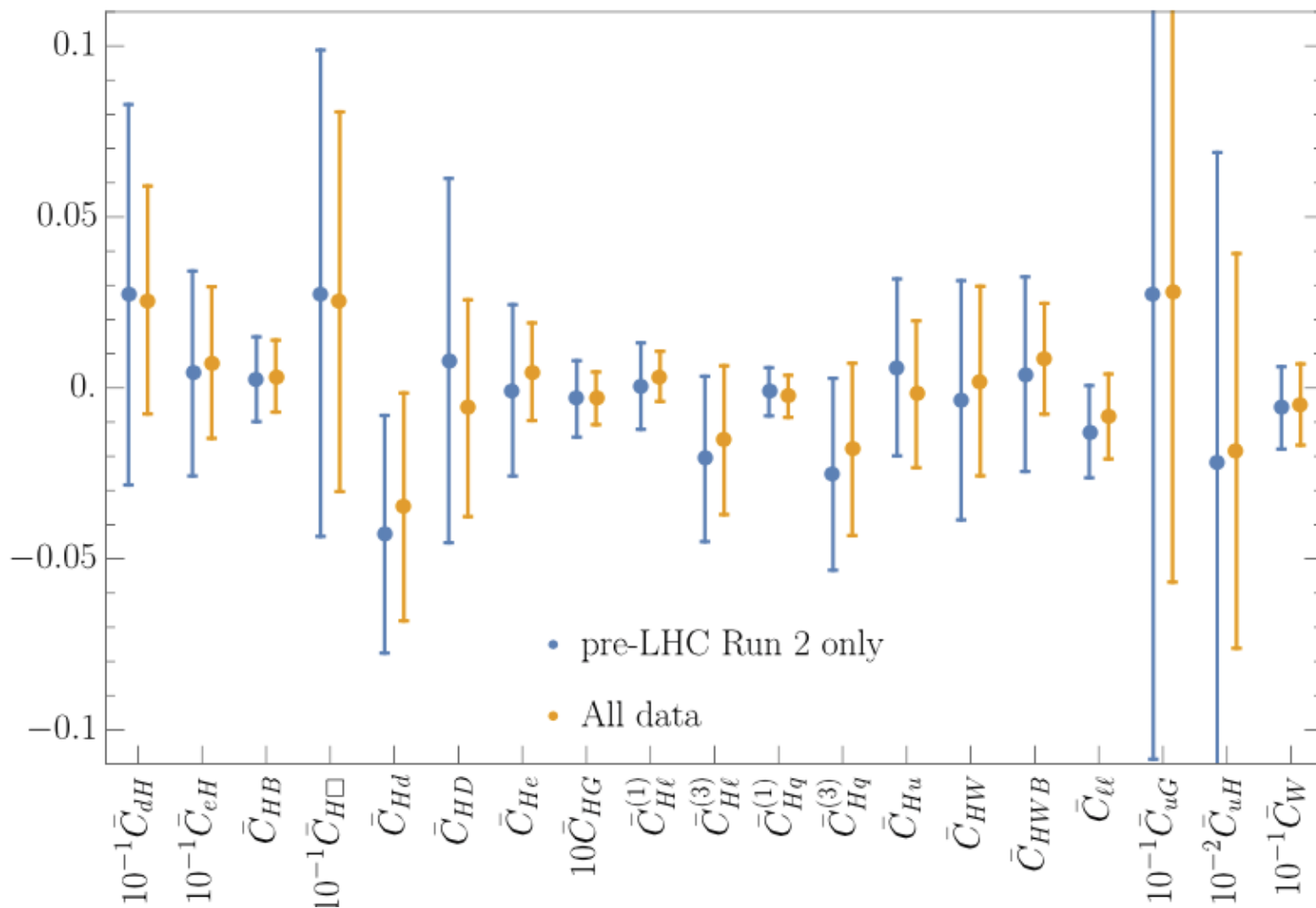
- **Warsaw** basis, fit each operator **individually**



Results

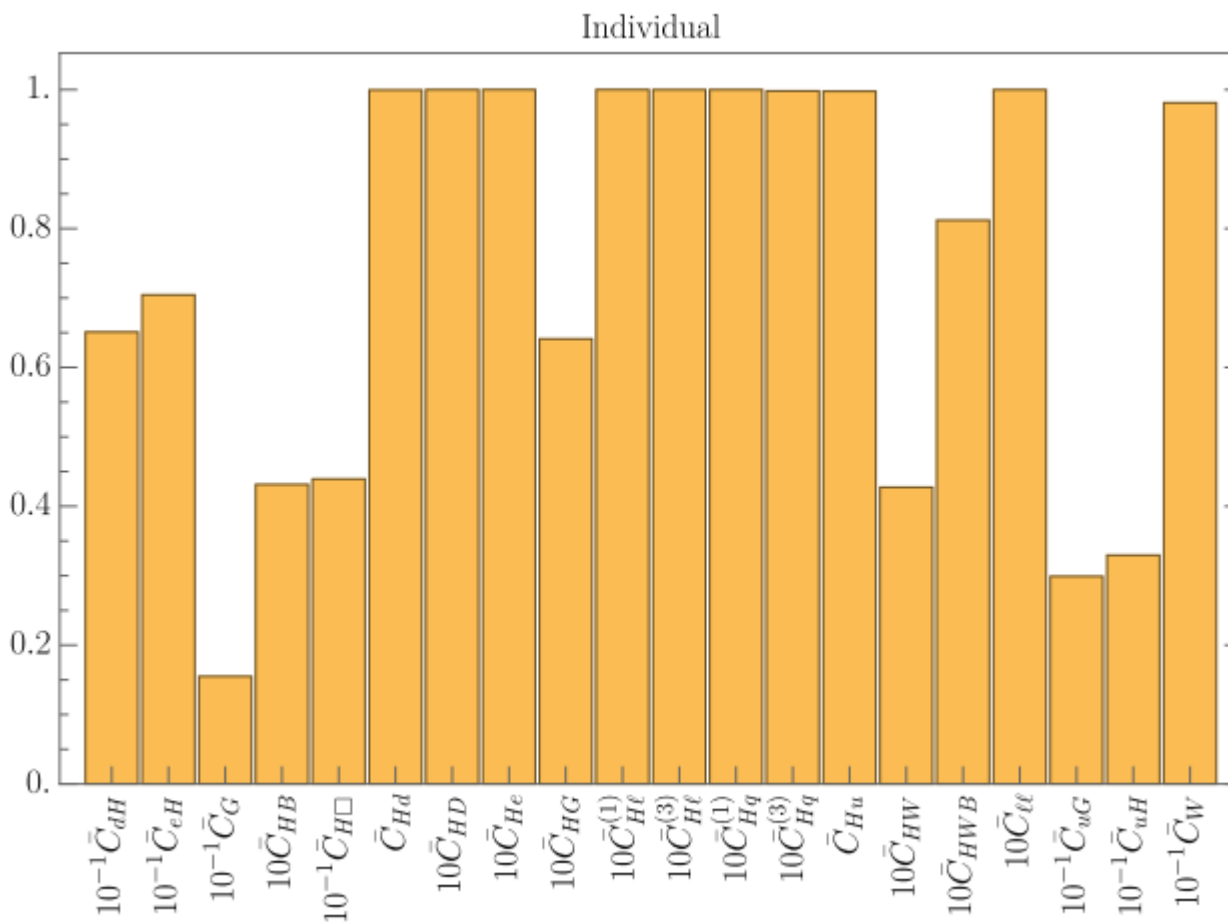
- **Warsaw basis, fit *all* operators simultaneously**

Marginalised



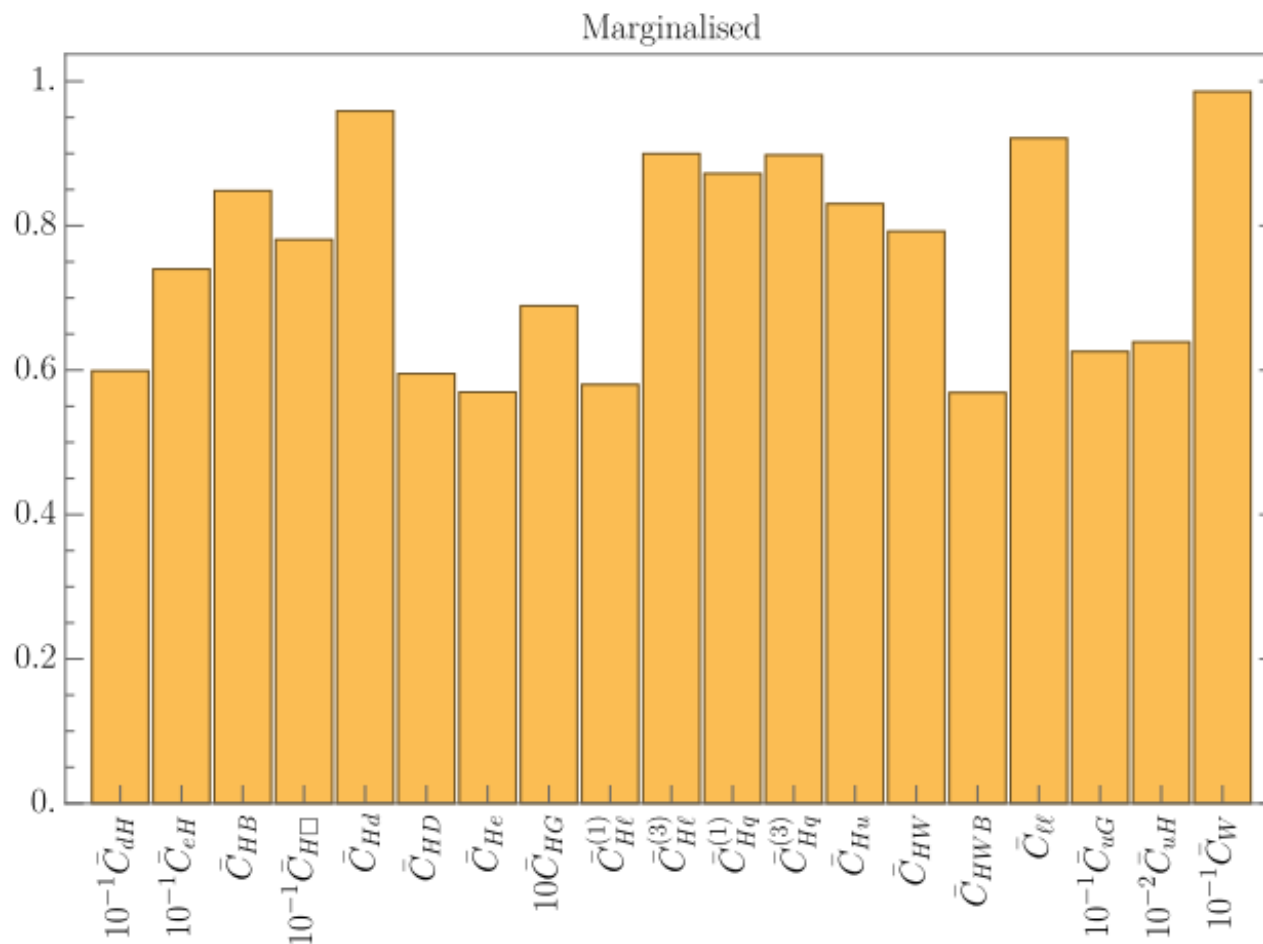
Results

- **Warsaw** basis, improvement from Run 1 to 2 (lower is better) for **individual** fit



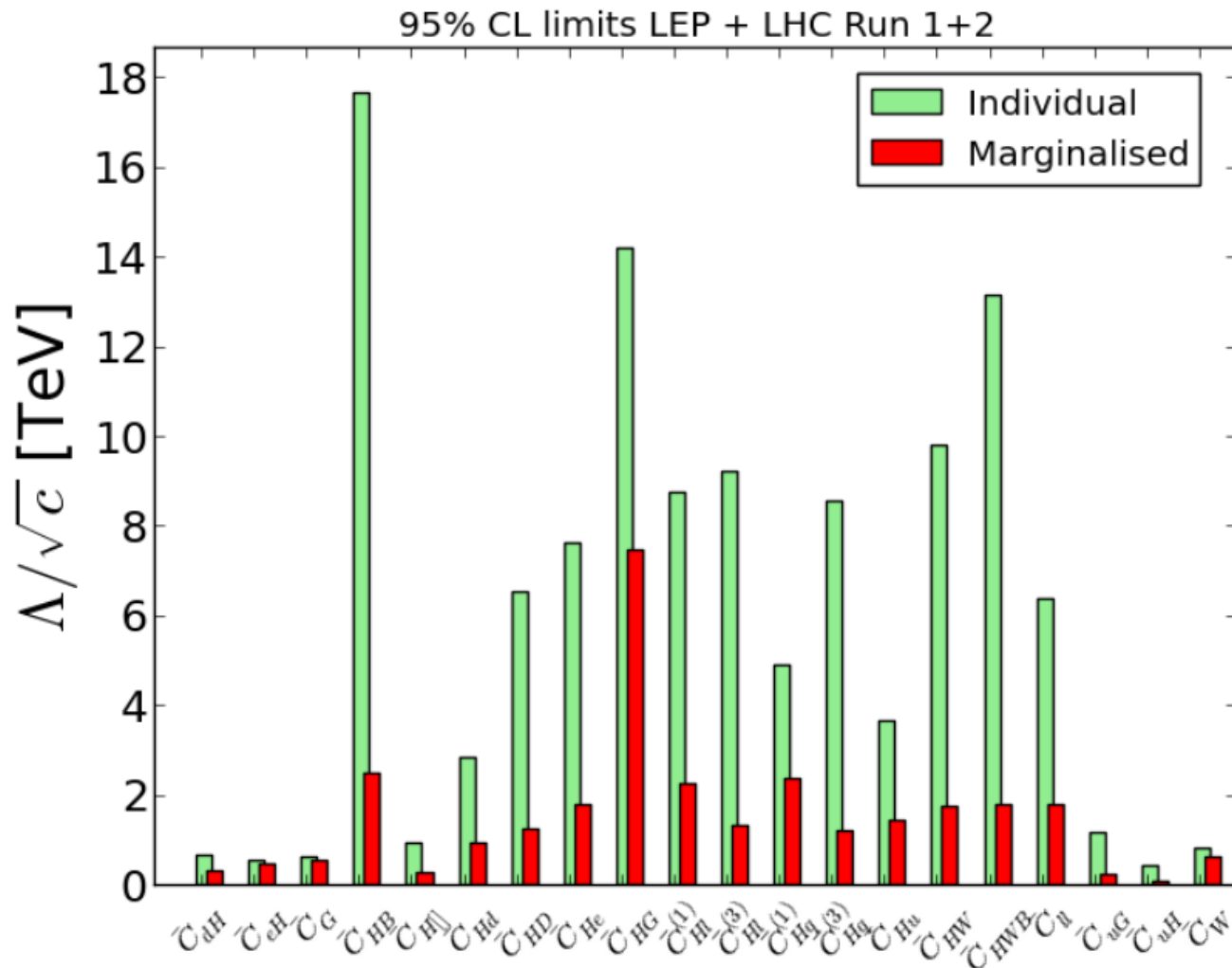
Results

- **Warsaw** basis, improvement from Run 1 to 2 (lower is better) for **marginalised** fit



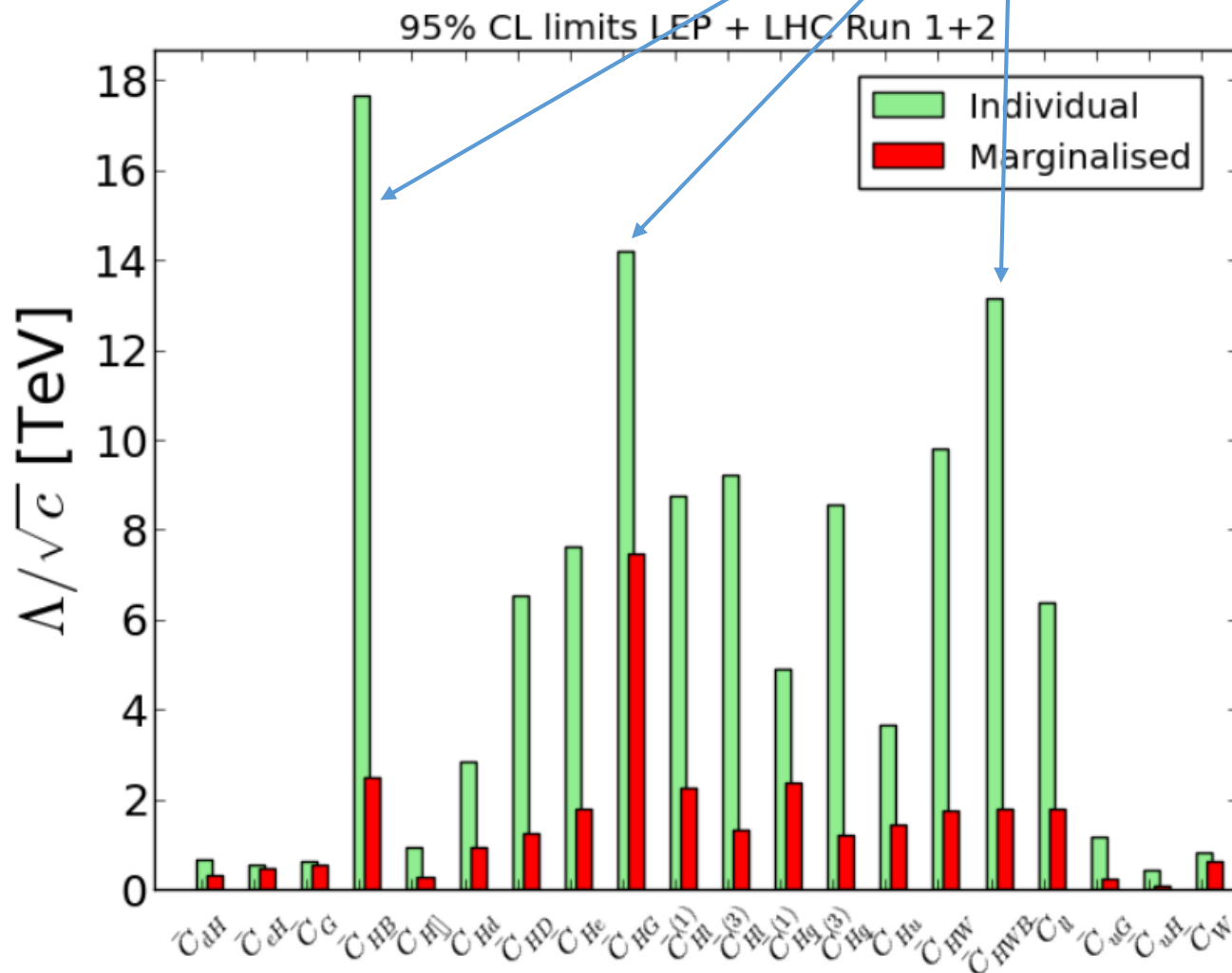
Results

- **Warsaw basis, summary**



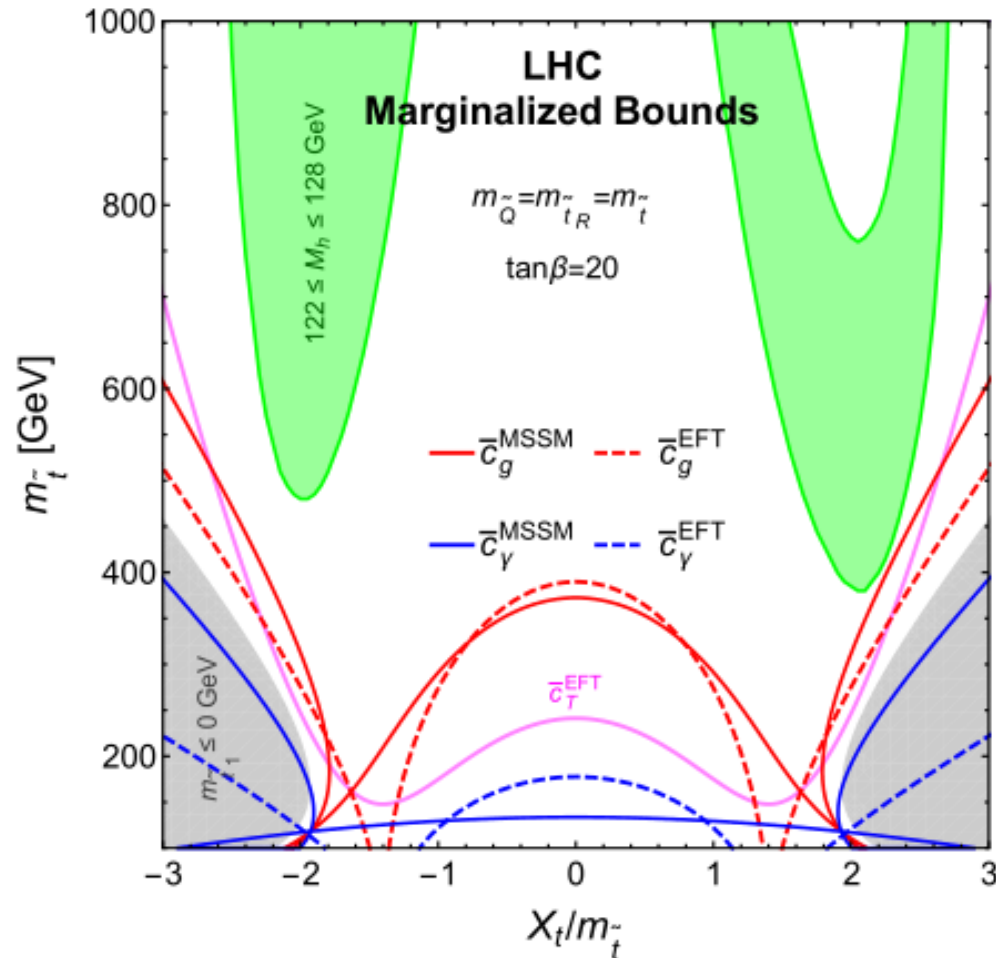
Results

- **Warsaw basis, summary**



Results

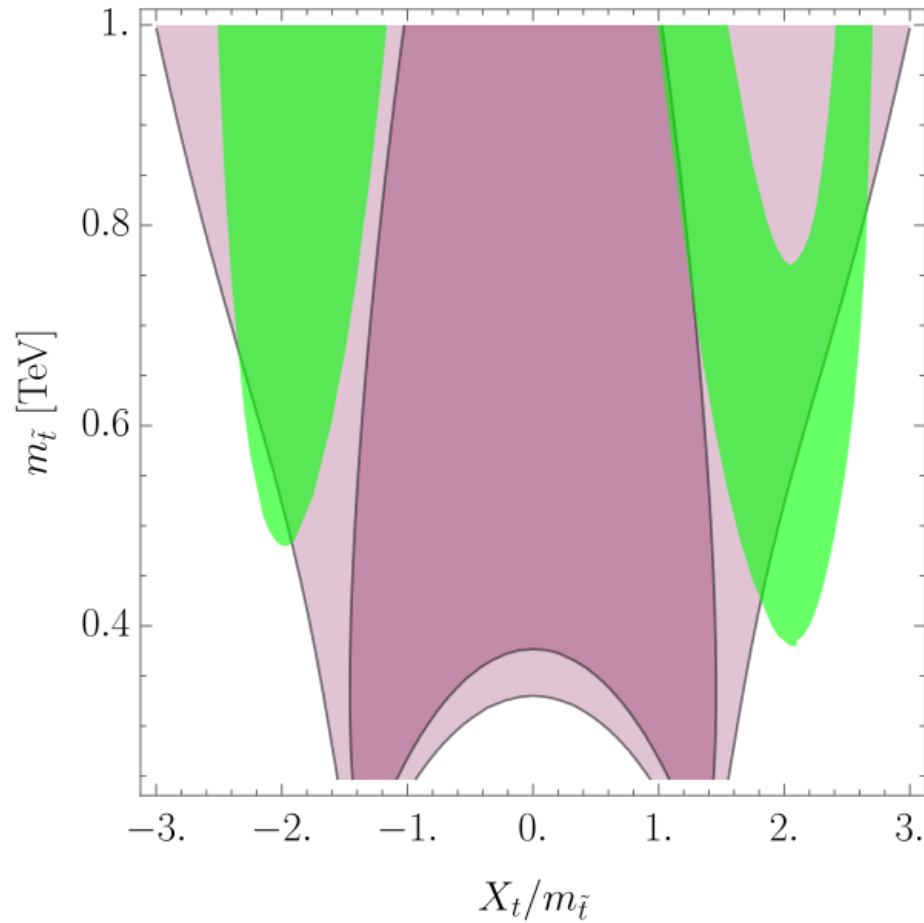
- Simplified models: **stops** (Run 1)



Results

- Simplified models: **stops** (Run 2)

$$\tan \beta = 20$$



Results

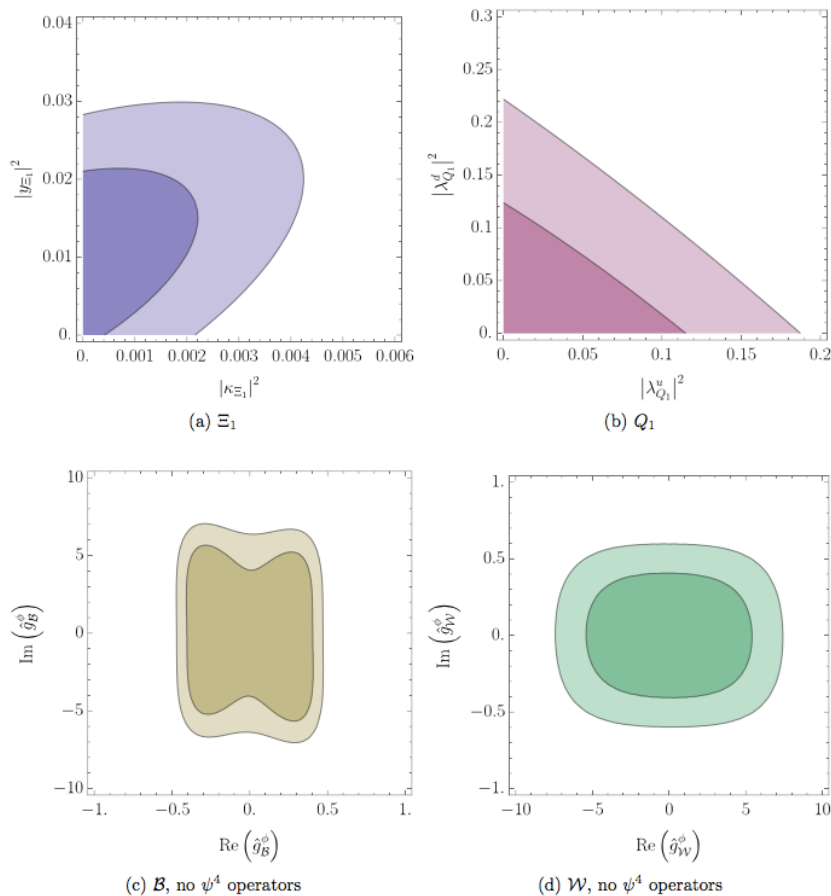
- Simplified models: **renormalisable SM extensions**

Name	Spin	$SU(3)$	$SU(2)$	$U(1)$	Name	Spin	$SU(3)$	$SU(2)$	$U(1)$
\mathcal{S}	0	1	1	0	Δ_1	$\frac{1}{2}$	1	2	$-\frac{1}{2}$
\mathcal{S}_1	0	1	1	1	Δ_3	$\frac{1}{2}$	1	2	$-\frac{1}{2}$
φ	0	1	2	$\frac{1}{2}$	Σ	$\frac{1}{2}$	1	3	0
Ξ	0	1	3	0	Σ_1	$\frac{1}{2}$	1	3	-1
Ξ_1	0	1	3	1	U	$\frac{1}{2}$	3	1	$\frac{2}{3}$
\mathcal{B}	1	1	1	0	D	$\frac{1}{2}$	3	1	$-\frac{1}{3}$
\mathcal{B}_1	1	1	1	1	Q_1	$\frac{1}{2}$	3	2	$\frac{1}{6}$
\mathcal{W}	1	1	3	0	Q_5	$\frac{1}{2}$	3	2	$-\frac{5}{6}$
\mathcal{W}_1	1	1	3	1	Q_7	$\frac{1}{2}$	3	2	$\frac{7}{6}$
N	$\frac{1}{2}$	1	1	0	T_1	$\frac{1}{2}$	3	3	$-\frac{1}{3}$
E	$\frac{1}{2}$	1	1	-1	T_2	$\frac{1}{2}$	3	3	$\frac{2}{3}$

- Classification and tree-level matching dictionary

Results

- Simplified models: renormalisable SM extensions



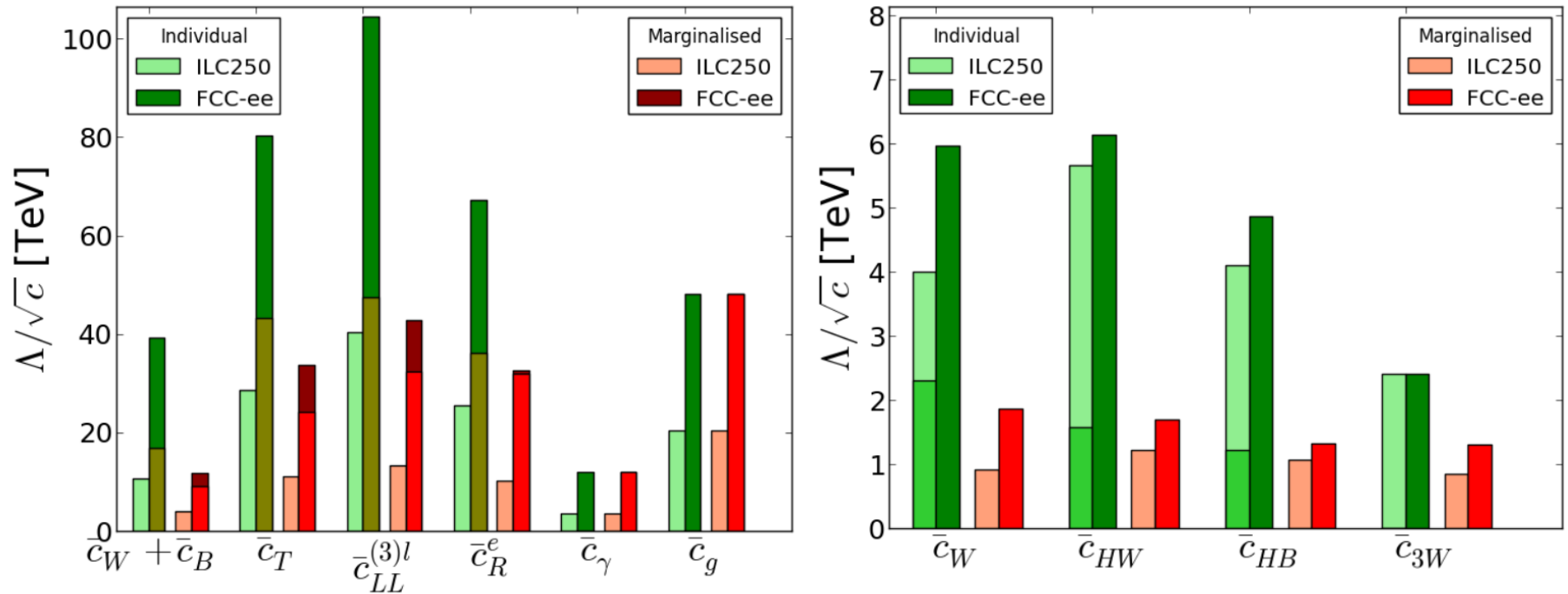
Model	χ^2	χ^2/n_d	Coupling	Mass / TeV
SM	157	0.987	-	-
\mathcal{S}_1	156	0.986	$ y_{\mathcal{S}_1} ^2 = (6.3 \pm 5.9) \cdot 10^{-3}$	$M_{\mathcal{S}_1} = (9.0, 49)$
φ , Type I	156	0.986	$Z_6 \cdot \cos \beta = -0.64 \pm 0.59$	$M_{\varphi} = (0.9, 4.3)$
Ξ	155	0.984	$ \kappa_{\Xi} ^2 = (4.2 \pm 3.4) \cdot 10^{-3}$	$M_{\Xi} = (12, 35)$
N	155	0.978	$ \lambda_N ^2 = (1.8 \pm 1.2) \cdot 10^{-2}$	$M_N = (5.8, 13)$
\mathcal{W}_1	155	0.984	$ \hat{g}_{\mathcal{W}_1}^{\phi} ^2 = (3.3 \pm 2.7) \cdot 10^{-3}$	$M_{\mathcal{W}_1} = (4.1, 13)$
E	157	0.993	$ \lambda_E ^2 < 1.2 \cdot 10^{-2}$	$M_E > 9.2$
Δ_3	156	0.990	$ \lambda_{\Delta_3} ^2 < 1.9 \cdot 10^{-2}$	$M_{\Delta_3} > 7.3$
Σ	157	0.992	$ \lambda_{\Sigma} ^2 < 2.9 \cdot 10^{-2}$	$M_{\Sigma} > 5.9$
Q_5	156	0.990	$ \lambda_{Q_5} ^2 < 0.18$	$M_{Q_5} > 2.4$
T_2	157	0.992	$ \lambda_{T_2} ^2 < 7.1 \cdot 10^{-2}$	$M_{T_2} > 3.8$
\mathcal{S}	157	0.993	$ y_{\mathcal{S}} ^2 < 0.32$	$M_{\mathcal{S}} > 1.8$
Δ_1	157	0.993	$ \lambda_{\Delta_1} ^2 < 5.7 \cdot 10^{-3}$	$M_{\Delta_1} > 13$
Σ_1	157	0.993	$ \lambda_{\Sigma_1} ^2 < 7.3 \cdot 10^{-3}$	$M_{\Sigma_1} > 12$
U	157	0.993	$ \lambda_U ^2 < 2.8 \cdot 10^{-2}$	$M_U > 6.0$
D	157	0.993	$ \lambda_D ^2 < 1.4 \cdot 10^{-2}$	$M_D > 8.4$
Q_7	157	0.993	$ \lambda_{Q_7} ^2 < 7.7 \cdot 10^{-2}$	$M_{Q_7} > 3.6$
T_1	157	0.993	$ \lambda_{T_1} ^2 < 0.13$	$M_{T_1} > 3.0$
\mathcal{B}_1	157	0.993	$ \hat{g}_{\mathcal{B}_1}^{\phi} ^2 < 2.4 \cdot 10^{-3}$	$M_{\mathcal{B}_1} > 21$

- Streamlines process of interpreting limits on BSM parameter space

Future e+e- Constraints

ILC and FCC-ee

J. Ellis and T.Y. [arXiv:1510:04561]



- Future precision sensitive to TeV scale, even for loop-induced operators
- One-loop matching simplified by a **Universal One-Loop Effective Action**

Henning, Lu, Murayama, 1412.1837;

Drozd, J. Ellis, Quevillon, TY, 1512.03003;

S.A.R. Ellis, Quevillon, TY, Z. Zhang, 1604.02445, 1706.07765.

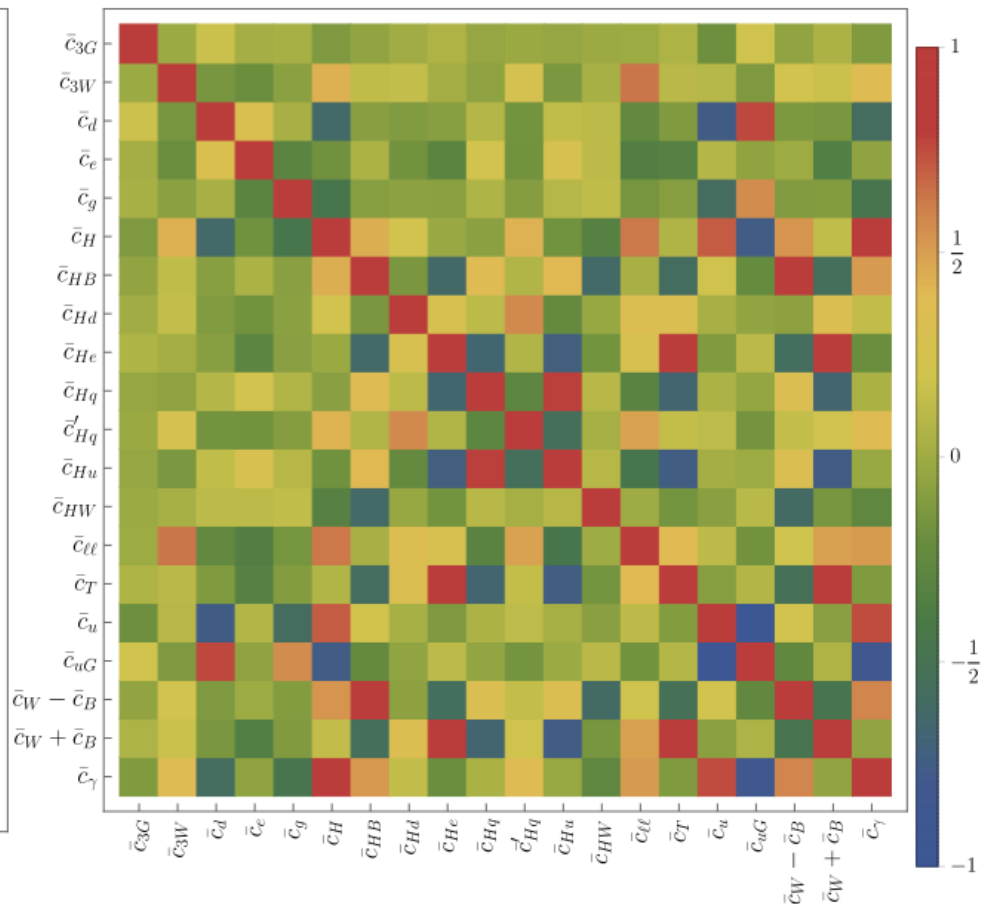
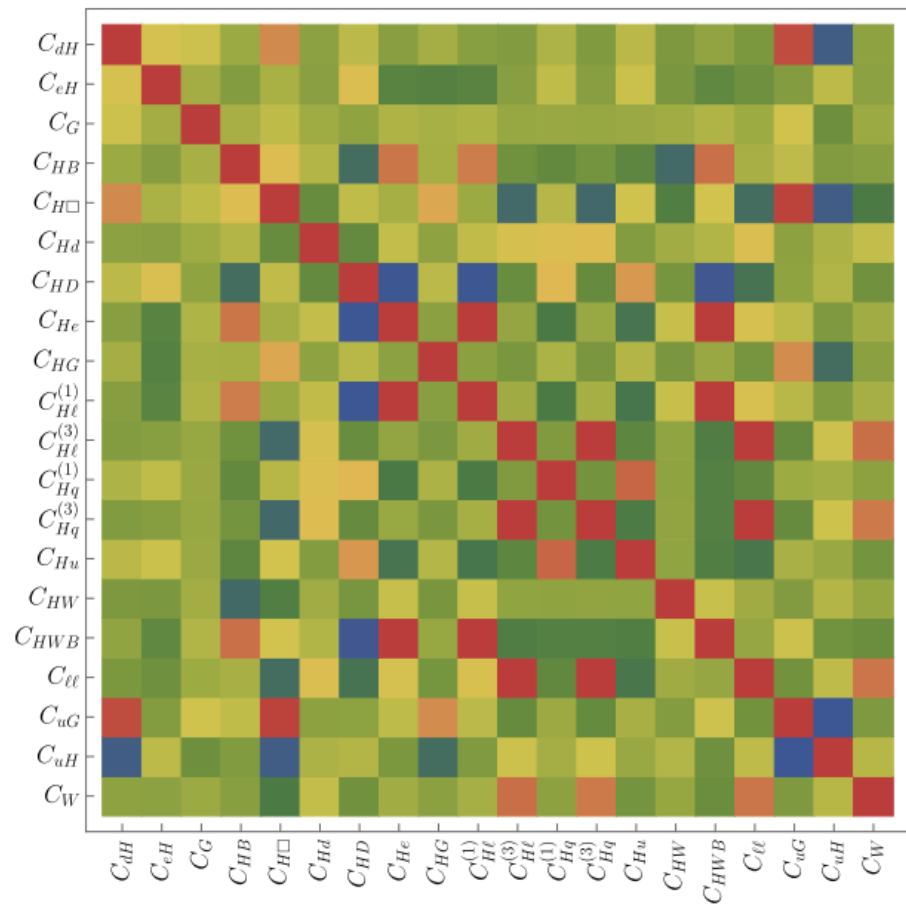
Conclusion

- SM EFT framework is the **Fermi theory of the 21st century**
- **Systematic classification** of decoupled new physics
- **Correlates measurements and eases interpretation**
- Finding **patterns of deviations** will give clues to a deeper underlying theory at higher energies

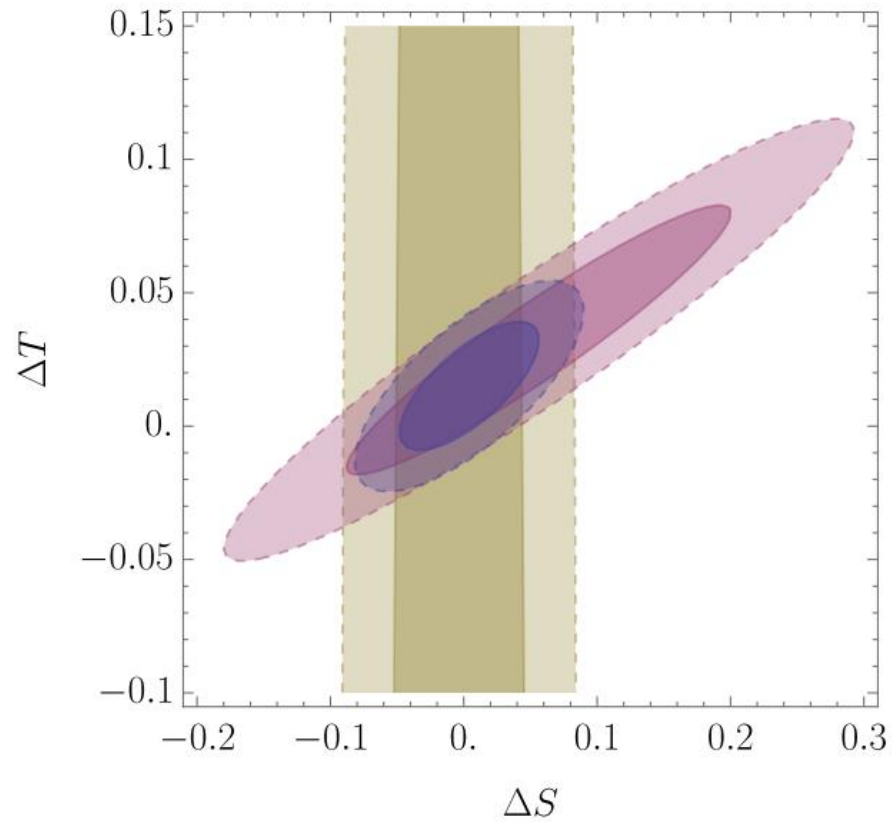
Backup

Coefficient	Z-pole + m_W	WW at LEP2	Higgs Run1	Higgs Run2	LHC WW high- p_T
\bar{C}_{dH}	×	×	42.4	57.6	×
\bar{C}_{eH}	×	×	49.6	50.4	×
\bar{C}_G	×	×	2.4	97.6	×
\bar{C}_{HB}	×	×	18.6	81.4	×
$\bar{C}_{H\Box}$	×	×	19.3	80.7	0.01
\bar{C}_{Hd}	99.85	×	0.04	0.1	×
\bar{C}_{HD}	99.92	0.06	×	×	×
\bar{C}_{He}	99.99	0.01	×	×	×
\bar{C}_{HG}	×	×	41.1	58.9	0.03
$\bar{C}_{H\ell}^{(1)}$	99.97	0.03	×	×	×
$\bar{C}_{H\ell}^{(3)}$	99.56	0.41	×	×	0.01
$\bar{C}_{Hq}^{(1)}$	99.98	×	×	×	×
$\bar{C}_{Hq}^{(3)}$	98.5	0.96	0.19	0.31	0.07
\bar{C}_{Hu}	99.3	×	0.2	0.42	0.04
\bar{C}_{HW}	×	×	18.3	81.7	×
\bar{C}_{HWB}	57.7	0.02	8.2	34.1	×
$\bar{C}_{\ell\ell}$	99.66	0.3	×	0.01	×
\bar{C}_{uG}	×	×	8.9	91.1	×
\bar{C}_{uH}	×	×	10.9	89.1	×
\bar{C}_W	×	96.2	×	×	3.8

Backup



Backup



$$\frac{v^2}{\Lambda^2} C_{HWB} = \frac{g_1 g_2}{16\pi} \Delta S, \quad \frac{v^2}{\Lambda^2} C_{HD} = -\frac{g_1 g_2}{2\pi (g_1 + g_2)} \Delta T,$$

Higgs constraints on dim-6 operators

- Operators affect Higgs signal strength measurements, differential distributions

