

Top/Higgs/EW processes in the SMEFT at NLO in QCD

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JHEP 1608 (2016) 039, EPJC 77 (2017) 4, 262 & JHEP 1810 (2018) 005

Higgs Couplings 2018, Tokyo

29th November 2018

Introduction

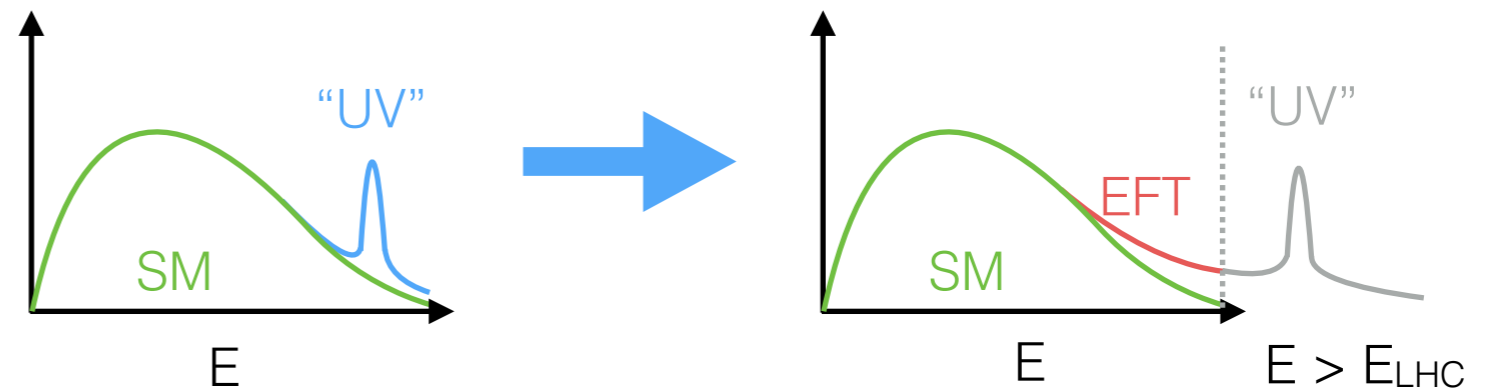
The LHC is entering a “precision era”

- No clear evidence for new physics from **direct searches**
- We are approaching the **limits of the ‘energy frontier’**
- Higgs boson discovery has completed the picture of the Standard Model (SM) Electroweak (EW) sector
- Properties consistent with SM expectations (so far)
- Complementary: **Standard Model Effective Field Theory**

Many channels are becoming **systematics dominated**

- Requires **high precision** theory input: higher order predictions
- Fixed order & interfaced with parton shower (PS)
- NLOQCD+PS standard for **SM**, also useful for **BSM** effects

SMEFT



Parametrise new physics effects at experimental energy E

- BSM states are ‘decoupled’ *i.e.* live at an energy $\Lambda \gg E$
- Generalised, gauge/Lorentz-invariant interactions between SM d.o.f

Operator expansion:

$$\mathcal{L}_{\text{eff}} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}} \quad \text{more: fields derivatives}$$

- Higher-derivative/contact operators: sensitive via large momentum flows through vertices (tails of energy distributions)

Dimension 6: 59 (76 real) - 2499 operators, depends on CP, flavor assumptions

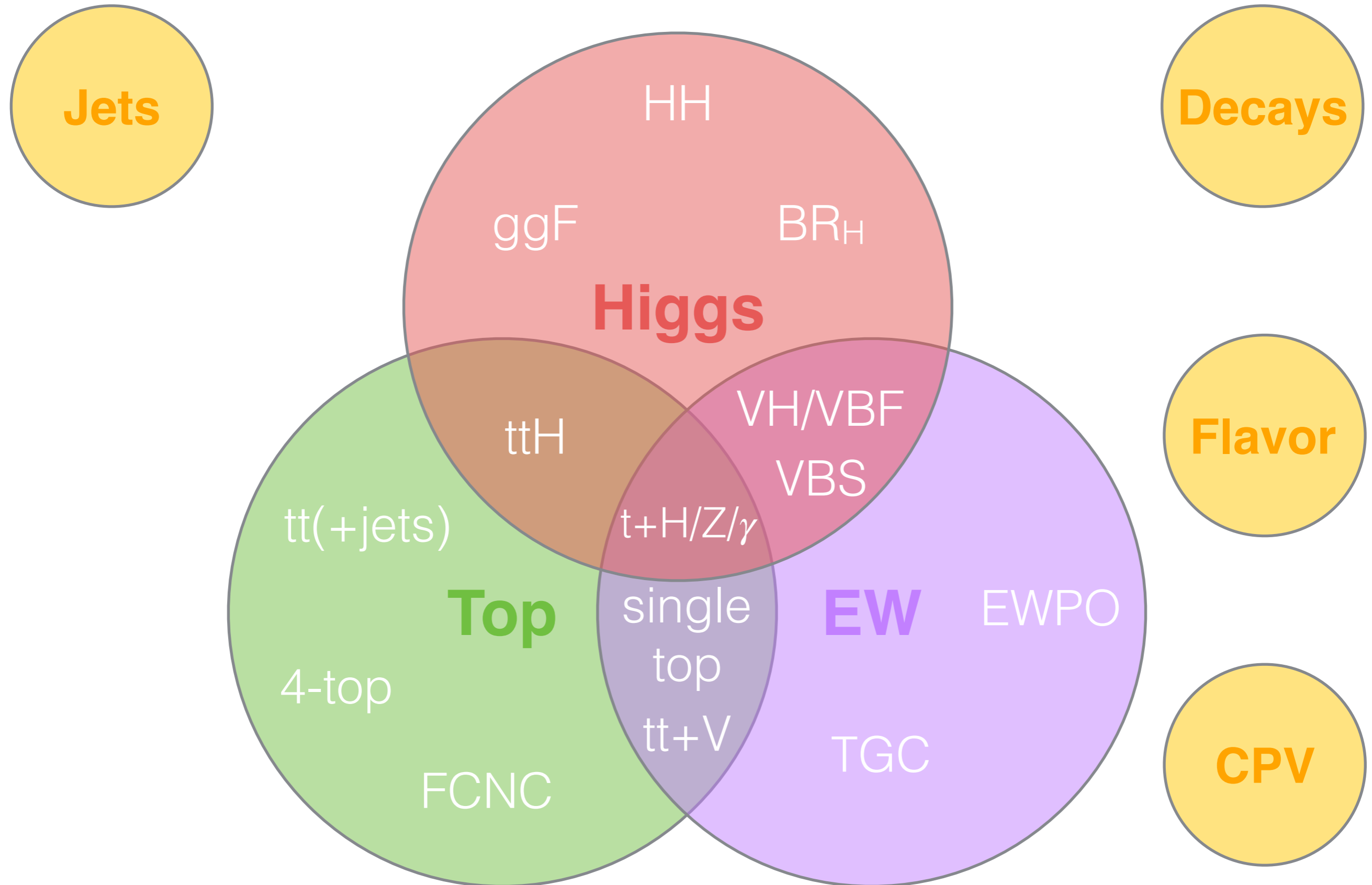
[Buchmuller & Wyler; Nucl.Phys. B268 (1986) 621]

[Grzadkowski et al.; JHEP 1010 (2010) 085]

Order-by-order, self-consistent, renormalisable QFT

- Unlike ‘Anomalous Couplings’ approach
- Consistent matching to new physics models \rightarrow patterns/correlations

SMEFT at the LHC: testing the EWSB sector



Status: LO

Increasing number of EFT interpretations in experimental analyses

Several MC implementations of SMEFT at LO

- HEL, SMEFTsim, dim6top,...
[Alloul et al.; JHEP 1404 (2014) 110]
[Brivio et al.; JHEP 1712 (2017) 070]
- Higgs Characterisation, BSMC, HiggsPO,...
[Aguilar Saavedra et al.; arXiv:1802.07237]
[Artoisenet et al.; JHEP 1311 (2013) 043]
[Falkowski et al.; EPJC 75 (2015) 12, 583]
[Greljo et al.; EPJC 77 (2017) 12, 838]

Many phenomenological analyses and global fits performed for SMEFT at LHC+...

- Separate Higgs/Gauge and top sector
[Falkowski & Riva; JHEP 1502 (2015) 039]
[Berthier & Trott; JHEP 1505 (2015) 024]
[Corbett et al.; JHEP 1508 (2015) 156]
[Buckley et al.; JHEP 1604 (2016) 015]
- LO accuracy
[Englert et al.; EPJC 76 (2016) 7, 393]
[Butter et al.; JHEP 1607 (2016) 152]
[Ellis et al.; arXiv:JHEP 1806 (2018) 146]

Going NLO

Ultimate goal: a **precision global fit** of full SMEFT including LHC observables at HL-LHC

NLO QCD(+PS) predictions

- K-factors/shapes & control over PDF + scale uncertainties

NLO EW corrections

- Potentially important but much harder
- Automation on the way with SHERPA, Madgraph5_aMC@NLO

RG-improved predictions & operator mixing

- Very helpful for cross checking NLO implementations
- Compare to full NLO calculations, assess the importance of finite terms

[Alonso, Jenkins, Manohar & Trott; JHEP 1310 (2013) 087, JHEP 1401 (2014) 035 & JHEP 1404 (2014) 159*]*

SMEFT@NLO in QCD

Motivation

- Bring SMEFT predictions to the required precision level for the LHC lifetime
- Include operators describing EWSB (top/Higgs/EW) sector interactions

Today's results: part of ongoing efforts in developing MC tool for SMEFT in top/Higgs/EW sector

- General FeynRules/NLOCT UFO implementation of 'Warsaw' basis

[Alloul et al.; *Comp. Phys. Comm.* 185 (2014) 2250] & [Degrande; *Comp. Phys. Comm.* 197 (2015) 239]

- NLOQCD + PS events in general purpose MC generators
- Generate any process including SMEFT

Model validated against existing implementations

- EW Higgs production, single-top, ttH, ttZ/ γ

[Degrande et al; *EPJC* 77 (2017) 4, 262]

[Zhang; *PRL* 116 (2016) 162002]

[Maltoni, Vryonidou & Zhang; *JHEP* 1610 (2016) 123]

[Degrande et al.; *PRD* 91 (2015) 034024]

[Bylund et al.; *JHEP* 1605 (2016) 052] [Durieux, Maltoni & Zhang; *PRD* 91 (2015) 074017]

● ●	$(\bar{L}L)(\bar{L}L)$	● ●	$(\bar{R}R)(\bar{R}R)$	● ●	$(\bar{L}L)(\bar{R}R)$
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$

Complete implementation

- No CP/B/L-violation

Flavor symmetry:

$$U(2)_Q \times U(2)_u \times U(3)_d \times U(3)_L \times U(3)_e$$

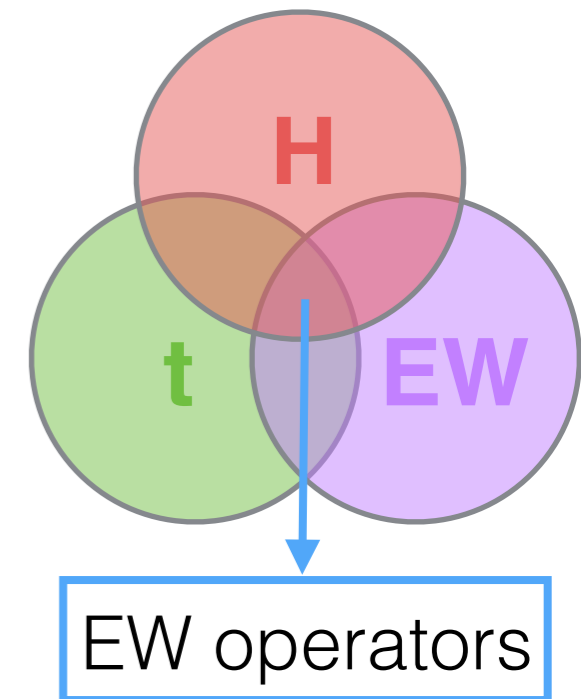
- Independent top operators
- Universal 1st & 2nd generations
- 4F: FFFF & FFff (ffff omitted)

●	$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$	X^3	φ^6 and $\varphi^4 D^2$	●	$\psi^2 \varphi^3$		
Q_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s^k q_t^j)$	Q_G	$f^{ABC} G_\mu^A G_\nu^B G_\rho^C$	Q_φ	$(\varphi^\dagger \varphi)^3$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^A G_\nu^B G_\rho^C$	$Q_{\varphi \square}$	$(\varphi^\dagger \varphi) \square (\varphi^\dagger \varphi)$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	Q_W	$\varepsilon^{IJK} W_\mu^I W_\nu^J W_\rho^K$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^I W_\nu^J W_\rho^K$				
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$						
		$X^2 \varphi^2$	●	$\psi^2 X \varphi$	● ●	$\psi^2 \varphi^2 D$	
		$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{l}_p \gamma^\mu l_r)$
		$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi) (\bar{l}_p \tau^I \gamma^\mu l_r)$
		$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{e}_p \gamma^\mu e_r)$
		$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{q}_p \gamma^\mu q_r)$
		$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi) (\bar{q}_p \tau^I \gamma^\mu q_r)$
		$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{u}_p \gamma^\mu u_r)$
		$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{d}_p \gamma^\mu d_r)$
		$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi) (\bar{u}_p \gamma^\mu d_r)$

+ bosonic ops.
 ≈ 75 parameters

Case study: tZj/tHj

[Degrande, Maltoni, KM, Vryonidou, Zhang; JHEP 1810 (2018) 005]



EW top interactions: **single top** production

- No QCD contribution
- Single top rate at 13 TeV LHC ~ 200 pb (1/4 of QCD $t\bar{t}$)
- Sensitive to **2 four-fermion** and **3 top/EW** operators that modify tbW vertex
- Limited energy dependence

Require the presence of an additional **Z** or **Higgs**

- Unique possibility of probing large set of top/Higgs/EW operators at once
- New energy growth from unitarity violating/EFT effects
- **Higher thresholds** may enhance EFT effects

Recent LHC measurement of tZj cross section at 4.2σ

[ATLAS; PLB 780 (2018) 557-577], [CMS-PAS-TOP-16-020 & PLB 779 (2018) 358-384]

- Timely moment to perform EFT sensitivity study in this pair of challenging processes & showcase model implementation

Operators

tHj
tZj
both
NLO

• \mathcal{O}_W	$\varepsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_{\rho}^{K,\mu}$	• $\mathcal{O}_{\varphi Q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi) (\bar{Q} \gamma^\mu \tau^I Q) + \text{h.c.}$
• $\mathcal{O}_{\varphi W}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) W_I^{\mu\nu} W_{\mu\nu}^I$	• $\mathcal{O}_{\varphi Q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{Q} \gamma^\mu Q) + \text{h.c.}$
• $\mathcal{O}_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$	• $\mathcal{O}_{\varphi t}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{t} \gamma^\mu t) + \text{h.c.}$
• $\mathcal{O}_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$	• $\mathcal{O}_{\varphi tb}$	$i(\tilde{\varphi} D_\mu \varphi) (\bar{t} \gamma^\mu b) + \text{h.c.}$
• $\mathcal{O}_{\varphi \square}$	$(\varphi^\dagger \varphi) \square (\varphi^\dagger \varphi)$	• $\mathcal{O}_{\varphi q}^{(1)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{q}_i \gamma^\mu q_i) + \text{h.c.}$
• $\mathcal{O}_{t\varphi}$	$(\varphi^\dagger \varphi - \frac{v^2}{2}) \bar{Q} t \tilde{\varphi} + \text{h.c.}$	• $\mathcal{O}_{\varphi q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi) (\bar{q}_i \gamma^\mu \tau^I q_i) + \text{h.c.}$
• \mathcal{O}_{tW}	$i(\bar{Q} \sigma^{\mu\nu} \tau_I t) \tilde{\varphi} W_{\mu\nu}^I + \text{h.c.}$	• $\mathcal{O}_{\varphi u}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{u}_i \gamma^\mu u_i) + \text{h.c.}$
• \mathcal{O}_{tB}	$i(\bar{Q} \sigma^{\mu\nu} t) \tilde{\varphi} B_{\mu\nu} + \text{h.c.}$	• $\mathcal{O}_{Qq}^{(3,1)}$	$(\bar{q}_i \gamma_\mu \tau_I q_i) (\bar{Q} \gamma^\mu \tau^I Q)$
• \mathcal{O}_{tG}^*	$i(\bar{Q} \sigma^{\mu\nu} T_A t) \tilde{\varphi} G_{\mu\nu}^A + \text{h.c.}$	• $\mathcal{O}_{Qq}^{(3,8)}$	$(\bar{q}_i \gamma_\mu \tau_I T_A q_i) (\bar{Q} \gamma^\mu \tau^I T^A Q)$

Constrained by electroweak precision tests (LEP)

Two blind directions in Warsaw basis:

$$\mathcal{O}_{HW} = (D^\mu \varphi)^\dagger \tau_I (D^\nu \varphi) W_{\mu\nu}^I$$

$$\mathcal{O}_{HB} = (D^\mu \varphi)^\dagger (D^\nu \varphi) B_{\mu\nu}.$$

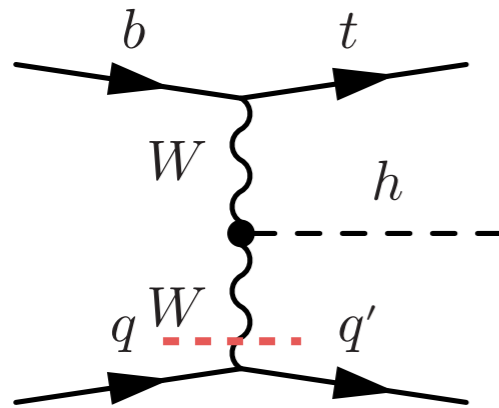
$$\frac{dC_i(\mu)}{d \log \mu} = \frac{\alpha_s}{\pi} \gamma_{ij} C_j(\mu), \quad \gamma = \begin{pmatrix} -2 & 0 & 0 \\ 0 & 2/3 & 0 \\ 0 & 0 & 2/3 \end{pmatrix}$$

Consider these two instead to assess orthogonal sensitivity of tZj/tHj

RGE

SMEFT in tHj/tZj

tHj (tZj = h → Z)

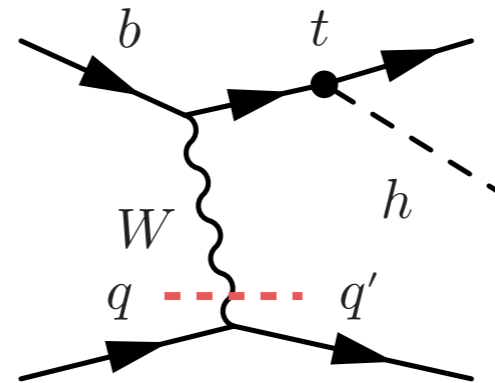


$$\mathcal{O}_{\varphi W} : \varphi^\dagger \varphi W_i^{\mu\nu} W_{\mu\nu}^i$$

HWW

TGC

$$\mathcal{O}_W : \epsilon^{ijk} W_{i,\mu\nu} W_j^{\nu\rho} W_{k,\rho}^\mu$$

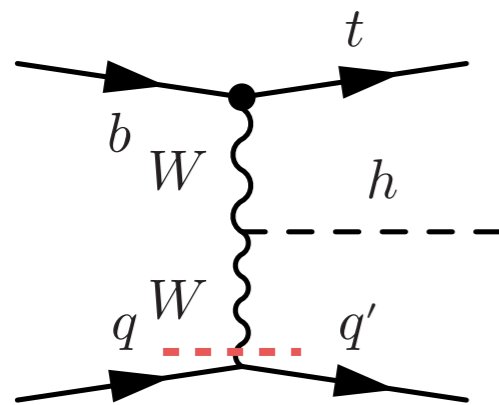


$$\mathcal{O}_{t\varphi} : (\varphi^\dagger \varphi) (\bar{Q} t) \tilde{\varphi}$$

top Yukawa

ttZ coupling

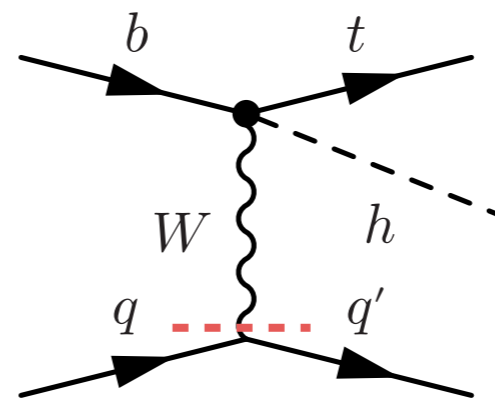
$$\mathcal{O}_{\varphi t} : i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{t} \gamma^\mu t)$$



$$\mathcal{O}_{\varphi Q}^{(3)} : i(\varphi^\dagger \overleftrightarrow{D}_\mu^i \varphi) (\bar{Q} \gamma^\mu \sigma_i Q)$$

Wtb vertex

$$\mathcal{O}_{\varphi tb} : i(\tilde{\varphi} D_\mu \varphi) (\bar{b} \gamma^\mu t)$$



$$\mathcal{O}_{\varphi Q}^{(3)} : i(\varphi^\dagger \overleftrightarrow{D}_\mu^i \varphi) (\bar{Q} \gamma^\mu \sigma_i Q)$$

Contact terms

$$\mathcal{O}_{tB} : (\bar{Q} \sigma_{\mu\nu} t) \tilde{\varphi} B^{\mu\nu}$$

Access the $bW \rightarrow tH$ & $bW \rightarrow tZ$ sub-amplitudes

- Rich interplay between EFT operators from different sectors
- Different energy growth and interference with the SM
- Potentially bring new information at high p_T

$\Lambda = 1 \text{ TeV}$ everywhere

Results

- Fixed order using Madgraph5_aMC@NLO

- NNPDF3.0 LO/NLO PDF sets
- 5-flavor scheme

$$m_t = 172.5 \text{ GeV}, \quad m_H = 125 \text{ GeV}, \quad m_Z = 91.1876 \text{ GeV}, \\ \alpha_{EW}^{-1} = 127.9, \quad G_F = 1.16637 \times 10^{-5} \text{ GeV}^{-2}.$$

- Scale choice

- t H j: $\mu_0 = (m_H + m_t)/4$
- t Z j: $\mu_0 = (m_Z + m_t)/4$

[Demartin, Maltoni & Mawatari; EPJC 75 (2015) 267]

- Uncertainties:

$$\sigma_{\pm \delta \mu_0 [\delta \mu_{EFT}]} \pm \delta_{PDF}$$

- 9 point variation of factorisation and renormalisation scale ($\mu_0/2$, μ_0 , $2\mu_0$)
- PDF uncertainties
- EFT scale variation from QCD running of operators (where relevant)

$$\sigma(\mu_0) = \sigma_{SM} + \sum_i \frac{1 \text{ TeV}^2}{\Lambda^2} C_i(\mu_0) \sigma_i(\mu_0) + \sum_{i,j} \frac{1 \text{ TeV}^4}{\Lambda^4} C_i(\mu_0) C_j(\mu_0) \sigma_{ij}(\mu_0)$$

Inclusive results: tHj

$pp \rightarrow t(\bar{t}) H j$ LHC@13 TeV

$c/\Lambda = 1$ [TeV⁻²]

σ [fb]	LO	NLO	K-factor
σ_{SM}	$57.56(4)^{+11.2\%}_{-7.4\%} \pm 10.2\%$	$75.87(4)^{+2.2\%}_{-6.4\%} \pm 1.2\%$	1.32
$\sigma_{\varphi W}$	$8.12(2)^{+13.1\%}_{-9.3\%} \pm 9.3\%$	$7.76(2)^{+7.0\%}_{-6.3\%} \pm 1.0\%$	0.96
$\sigma_{\varphi W, \varphi W}$	$5.212(7)^{+10.6\%}_{-6.8\%} \pm 10.2\%$	$6.263(7)^{+2.6\%}_{-7.8\%} \pm 1.3\%$	1.20
$\sigma_{t\varphi}$	$-1.203(6)^{+12.0\%}_{-15.6\%} \pm 8.9\%$	$-0.246(6)^{+144.5[31.4]\%}_{-157.8[19.0]\%} \pm 2.1\%$	0.20
$\sigma_{t\varphi, t\varphi}$	$0.6682(9)^{+12.7\%}_{-8.9\%} \pm 9.6\%$	$0.7306(8)^{+4.6[0.6]\%}_{-7.3[0.2]\%} \pm 1.0\%$	1.09
σ_{tW}	$19.38(6)^{+13.0\%}_{-9.3\%} \pm 9.4\%$	$22.18(6)^{+3.8[0.4]\%}_{-6.8[0.9]\%} \pm 1.0\%$	1.14
$\sigma_{tW, tW}$	$46.40(8)^{+9.3\%}_{-5.5\%} \pm 11.1\%$	$71.24(8)^{+7.4[1.5]\%}_{-14.0[6.9]\%} \pm 1.9\%$	1.54
$\sigma_{\varphi Q^{(3)}}$	$-3.03(3)^{+0.0\%}_{-2.2\%} \pm 15.4\%$	$-10.04(4)^{+11.1\%}_{-8.9\%} \pm 1.8\%$	3.31
$\sigma_{\varphi Q^{(3)}, \varphi Q^{(3)}}$	$11.23(2)^{+9.4\%}_{-5.6\%} \pm 11.2\%$	$15.28(2)^{+5.0\%}_{-10.9\%} \pm 1.8\%$	1.36
$\sigma_{\varphi tb}$	0	0	—
$\sigma_{\varphi tb, \varphi tb}$	$2.752(4)^{+9.4\%}_{-5.5\%} \pm 11.3\%$	$3.768(4)^{+5.0\%}_{-10.9\%} \pm 1.8\%$	1.54
σ_{HW}	$-3.526(4)^{+5.6\%}_{-9.5\%} \pm 10.9\%$	$-5.27(1)^{+6.5\%}_{-2.9\%} \pm 1.5\%$	1.50
$\sigma_{HW, HW}$	$0.9356(4)^{+7.9\%}_{-4.0\%} \pm 12.3\%$	$1.058(1)^{+4.8\%}_{-11.9\%} \pm 2.3\%$	1.13
σ_{tG}		$-0.418(5)^{+12.3\%}_{-9.8\%} \pm 1.1\%$	—
$\sigma_{tG, tG}$		$1.413(1)^{+21.3\%}_{-30.6\%} \pm 2.5\%$	—
$\sigma_{Qq^{(3,1)}}$	$-22.50(5)^{+8.0\%}_{-11.8\%} \pm 9.7\%$	$-20.10(5)^{+13.8\%}_{-13.3\%} \pm 1.1\%$	0.89
$\sigma_{Qq^{(3,1)}, Qq^{(3,1)}}$	$69.78(3)^{+8.0\%}_{-4.1\%} \pm 12.1\%$	$62.20(3)^{+11.5\%}_{-15.9\%} \pm 2.3\%$	0.89
$\sigma_{Qq^{(3,8)}}$	—	$0.25(3)^{+25.4\%}_{-27.1\%} \pm 4.7\%$	—
$\sigma_{Qq^{(3,8)}, Qq^{(3,8)}}$	$15.53(2)^{+8.0\%}_{-4.1\%} \pm 12.1\%$	$14.07(2)^{+11.0\%}_{-15.7\%} \pm 2.1\%$	0.91

K-factors not universal

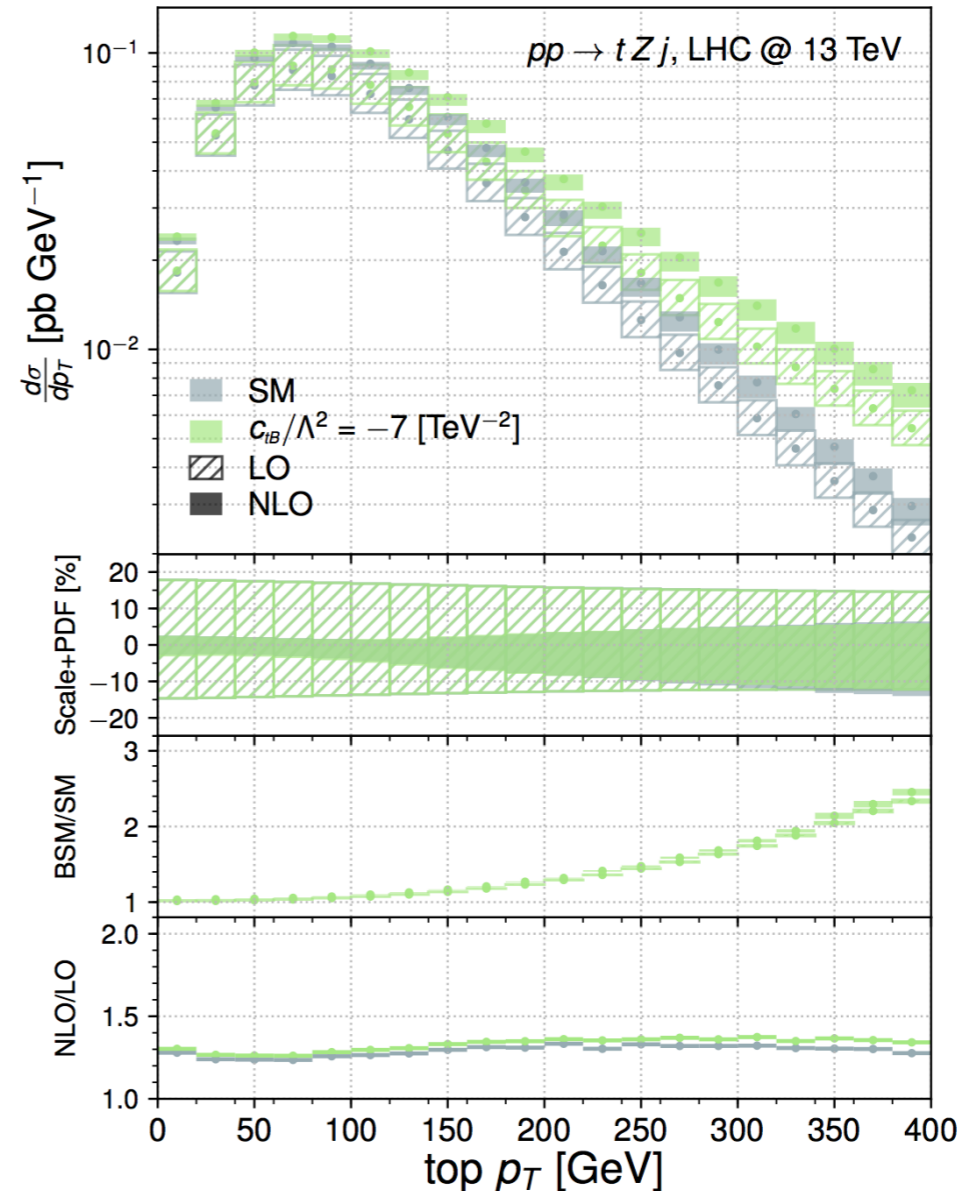
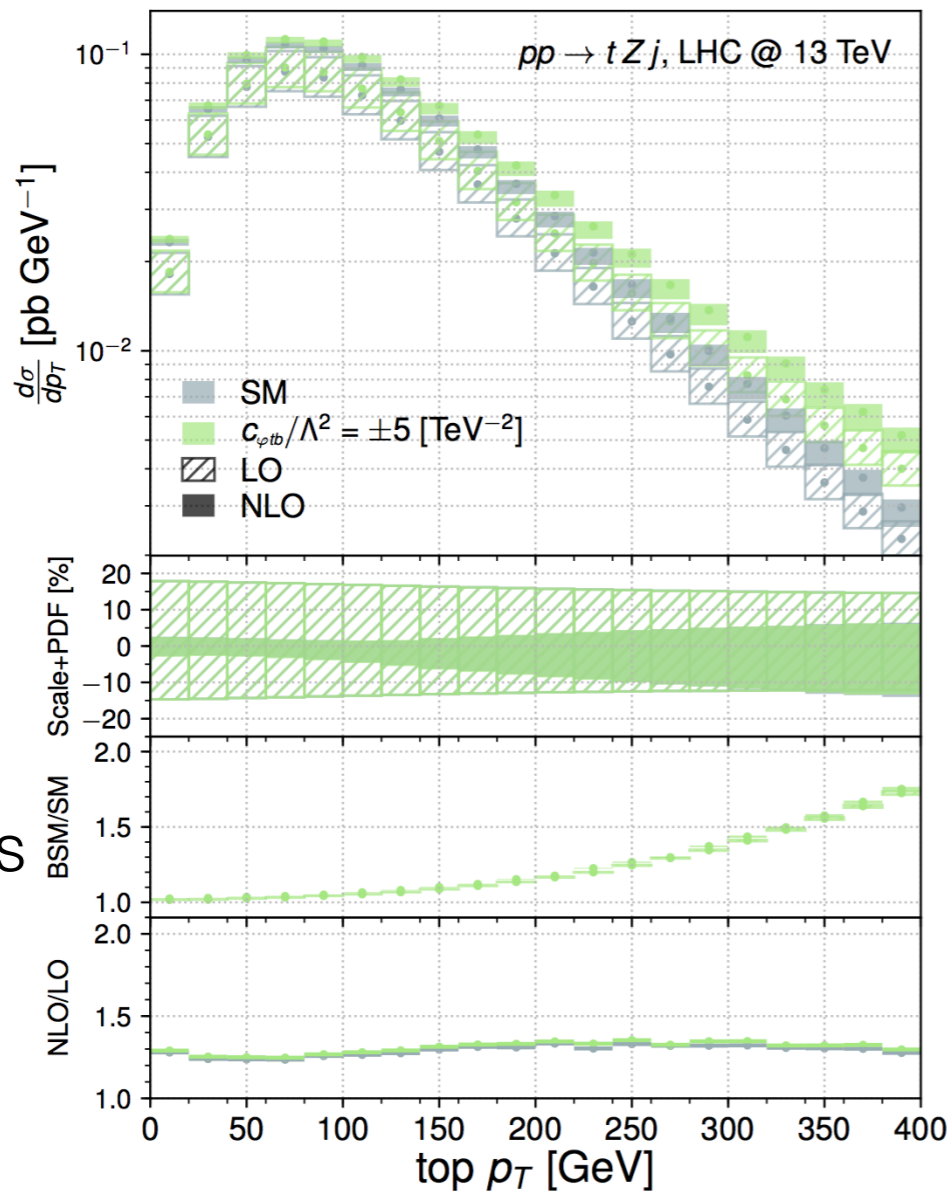
Reduction of
QCD scale/PDF
uncertainties

EFT scale uncertainty
subdominant

Some very strong
dependence on EFT
operators

O(>1) deviations within
current bounds

Differential results in tZj



Reduced
uncert.

Large effects

Potentially large deviations in the tails (saturating current limits)
 tHj process is very rare, differential results not likely at LHC

LHC sensitivity

Usual EFT story: looking at **high energy tails** increases sensitivity
 Compare to single top which has a much larger rate

$r = \sigma_i / \sigma_{SM}$	tj	tj	tZj	tZj	tHj
		$(p_T^t > 350 \text{ GeV})$		$(p_T^t > 250 \text{ GeV})$	
σ_{SM}	224 pb	880 fb	839 fb	69 fb	75.9 fb
r_{tW}	0.0275	0.024	0.016	0.010	0.292
$r_{tW,tW}$	0.0162	0.35	0.095	0.67	0.940
$r_{\varphi Q^{(3)}}$	0.121	0.121	0.192	0.172	-0.132
$r_{\varphi Q^{(3)},\varphi Q^{(3)}}$	0.0037	0.0037	0.029	0.114	0.21
$r_{\varphi tb,\varphi tb}$	0.00090	0.0008	0.0050	0.027	0.050
r_{tG}	0.0003	-0.01	0.00053	-0.0048	-0.0055
$r_{tG,tG}$	0.00062	0.045	0.0027	0.022	0.025
$r_{Qq^{(3,1)}}$	-0.353	-4.4	-0.59	-2.22	-0.39
$r_{Qq^{(3,1)},Qq^{(3,1)}}$	0.126	11.5	0.65	5.1	1.21
$r_{Qq^{(3,8)},Qq^{(3,8)}}$	0.0308	2.73	0.133	1.01	1.08

Increased sensitivity for **weak dipoles**

Consistent with $2 \rightarrow 2$ subamplitude analysis

New energy growths w.r.t single top

Single top should eventually outperform tHj/tZj for **four fermion operators**

Current sensitivity

Gauge sensitivity of these processes at LHC

Recent measurements of tZj at CMS & ATLAS

- Signal strengths 0.75 ± 0.27 & 1.31 ± 0.47 *[CMS; PLB 779 (2018) 358-384]*
[ATLAS; CERN-EP-2017-188]
- Cast into naive ‘confidence intervals’
- Ignore acceptance effects & contribution of operators to bkg processes
- Which include tW , ttV , ttH , tWZ , tHW

CMS analysis of $tHj + tHW + ttH$

[CMS-PAS-HIG-17-005]

- Combined signal strength 1.8 ± 0.67
- Take into account only modifications to tHj
- Except top Yukawa operator contribution to ttH

σ [fb]	LO	NLO
$t H j$	57.5	75.9
$t \bar{t} H$	464	507
$t H W$	14.5	15.9

Current sensitivity

tZj

TGC

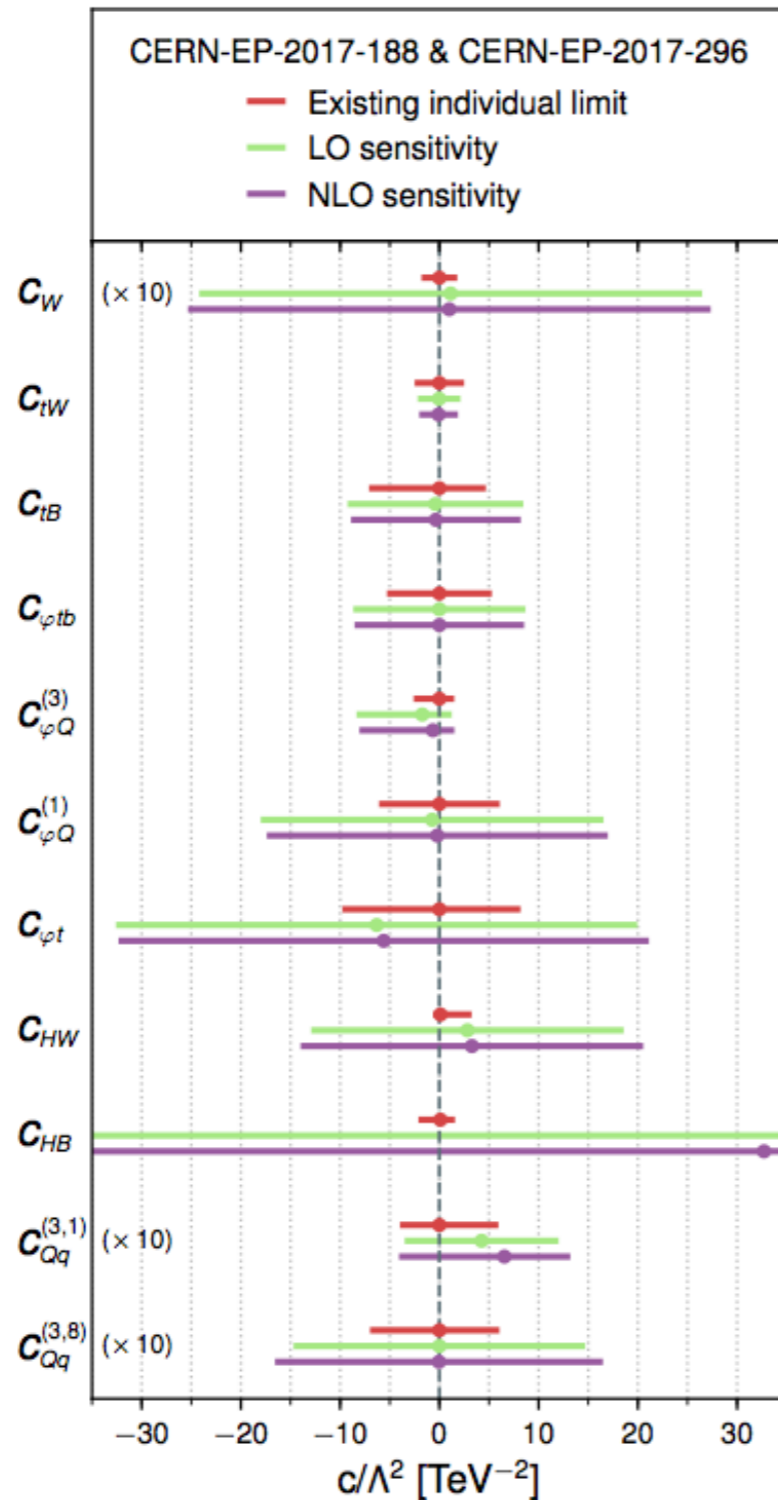
Dipoles

RHCC

Currents

LEP
orthogonal

4-fermion



tHj

Gauge-Higgs

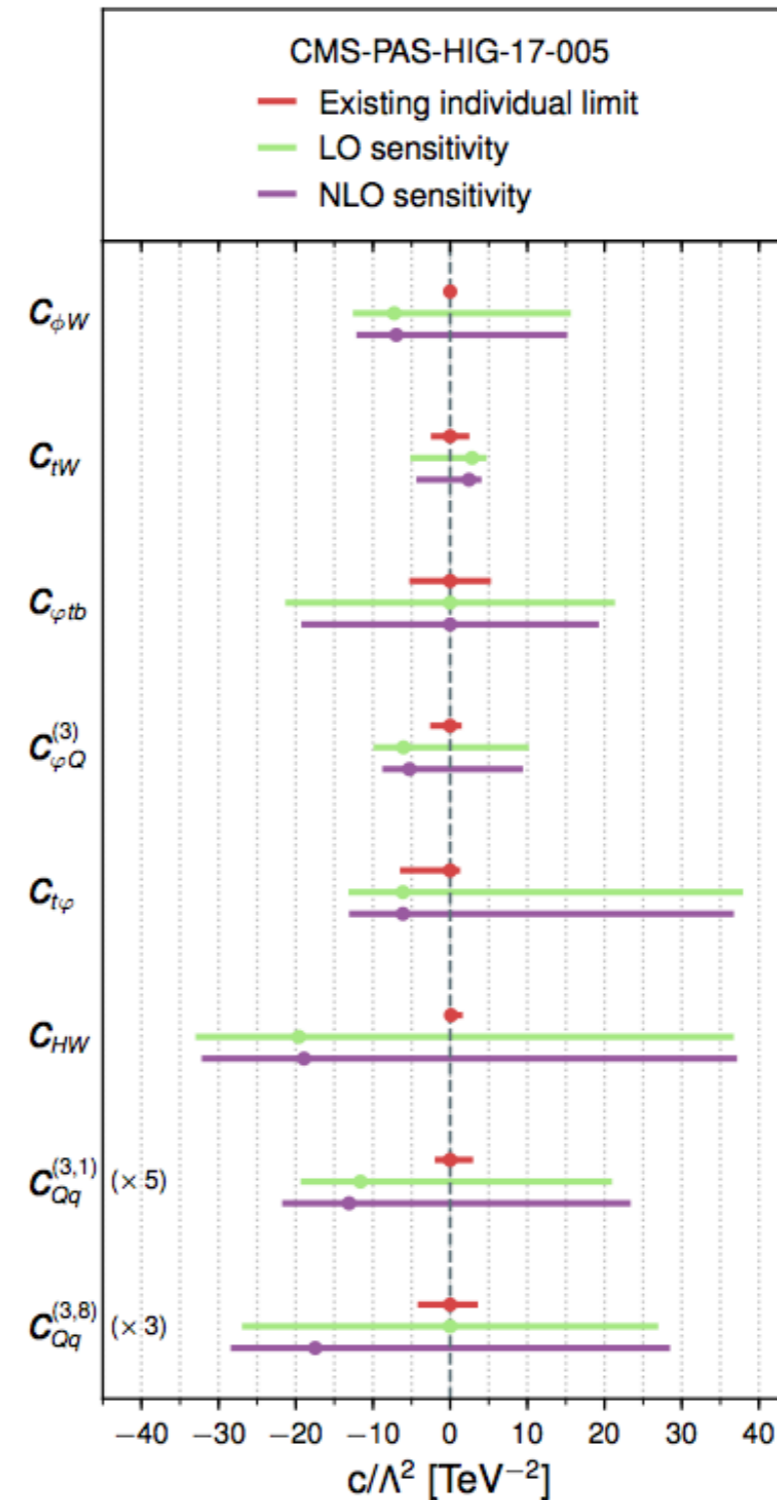
Dipole

RHCC

Currents

LEP
orthogonal

4-fermion



Future sensitivity

tZj: take high top p_T region > 250 GeV

- 10x smaller cross section \rightarrow end of Run II or HL-LHC

tHj: assume it is measured with the current tZj precision

- tHj inclusive cross section is similar to the high p_T tZj cross section
- target for HL-LHC

Start to improve on existing limits for certain operators

- Dipoles, RHCC, (top Yukawa)

NLO predictions increase sensitivity

- Bring theory uncertainties down below experimental stat. and syst.

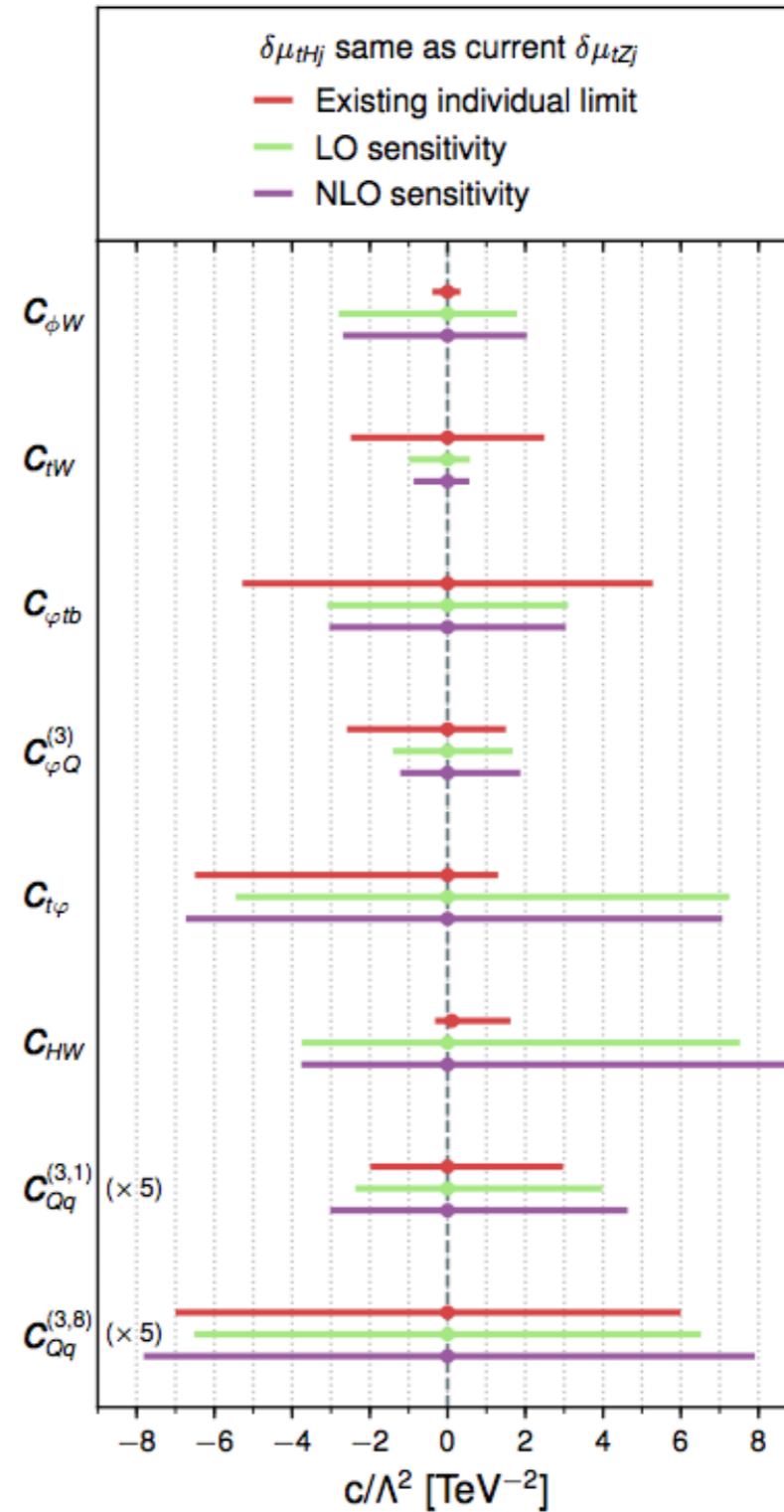
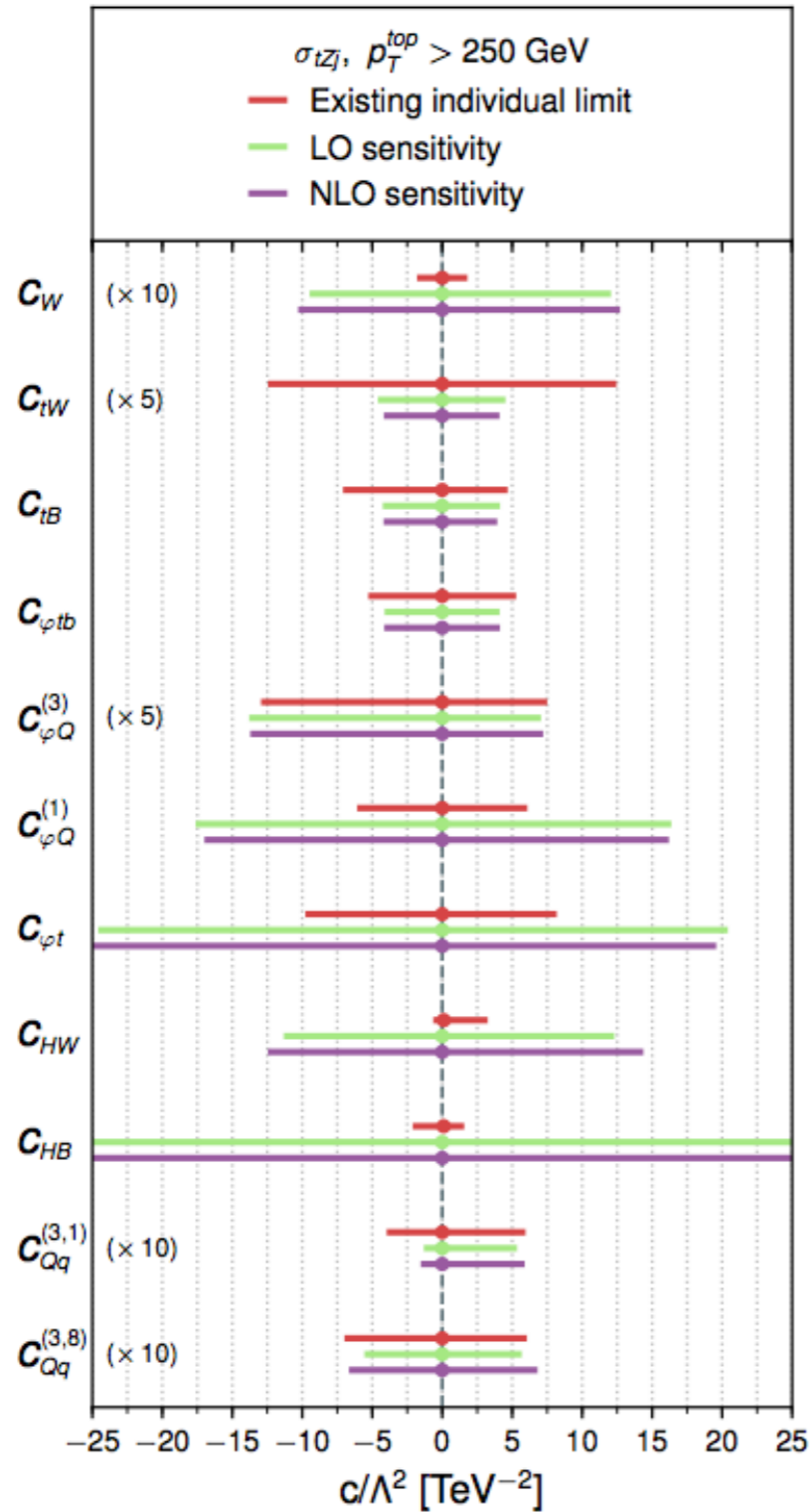
Future sensitivity

tZj

tHj

TGC
Dipoles
RHCC
Currents
LEP
orthogonal
4-fermion

Gauge-Higgs
Dipole
RHCC
Currents
LEP
orthogonal
4-fermion



Conclusion

Presented a FeynRules/NLOCT UFO implementation of **top/EW/Higgs** sector in SMEFT

- $U(3)^3 \times U(2)^2$ flavor symmetry to select **top quark** operators
- Allows for NLOQCD+PS predictions for **any relevant process** in, *e.g.*, MG5

Fixed order NLO predictions for **tZj** & **tHj** in SMEFT at LHC

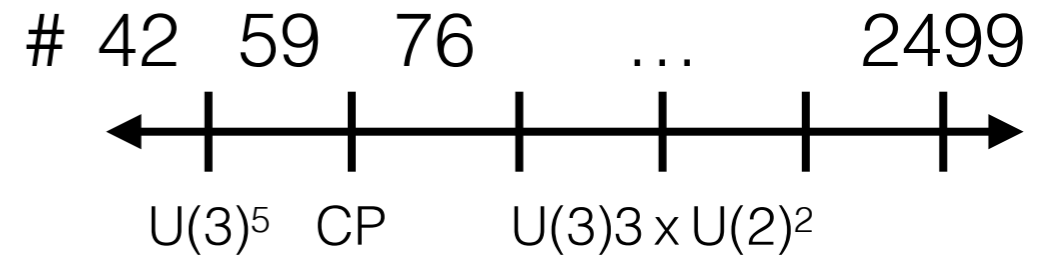
- Challenging process that showcases implementation
- Related to **mass generation/unitarity cancellations** in SM
- Complete predictions for large operator set → can be used in fits

Current & future LHC sensitivity study

- New energy growth with respect to single top for SU(2) current and RHCC
- Interesting future sensitivity at high energy to dipoles, 4F ops.

Thank you

Flavor symmetry



[Aguilar-Saavedra et al.; arXiv:1802.07237]

SM fermion sector: q^i, u^i, d^i, l^i, e^i

- 5 SU(3) x SU(2) x U(1) representations → U(3)⁵ flavor symmetry
- Only broken by Yukawa interactions

Some SMEFT operators also break it

- Chirality flipping F_{LR} structures (Yukawa-like)
- Flavor-violating (off diagonal/non-universal) entries

Starting point: flavor symmetric

- No chirality flipping & diagonal, universal structure

Controlled departures

- Minimal for top physics: U(3)₃ x U(2)₂, single out q_3, u_3
- Similar to MFV: expansion in Yukawa couplings with only y_t non-zero

[D'Ambrosio et al.; Nucl. Phys. B645 (2002) 155]

Anatomy of tHj

bW → tH Helicity amplitudes: $s \sim -t \gg v^2$

$\mathcal{O}_{\varphi tb}, \lambda_b = +$

	λ_W	0	+	-
λ_t				
+		$\sqrt{s(s+t)}$	$m_W \sqrt{-t}$	$\frac{1}{\sqrt{s}}$
-		$m_t \sqrt{-t}$	s^0	s^0

$\lambda_b, \lambda_W, \lambda_t$	SM	$\mathcal{O}_{t\varphi}$	$\mathcal{O}_{\varphi Q}^{(3)}$	$\mathcal{O}_{\varphi W}$	\mathcal{O}_{tW}	\mathcal{O}_{HW}
- , 0 , -	s^0	s^0	$\sqrt{s(s+t)}$	s^0	s^0	$\sqrt{s(s+t)}$
- , 0 , +	$\frac{1}{\sqrt{s}}$	$m_t \sqrt{-t}$	$m_t \sqrt{-t}$	$\frac{1}{\sqrt{s}}$	$\frac{m_W s}{\sqrt{-t}}$	$\frac{1}{\sqrt{s}}$
- , - , -	$\frac{1}{\sqrt{s}}$	$\frac{1}{\sqrt{s}}$	$m_W \sqrt{-t}$	$\frac{m_W s}{\sqrt{-t}}$	$m_t \sqrt{-t}$	$\frac{m_W (s+t)}{\sqrt{-t}}$
- , - , +	$\frac{1}{s}$	s^0	s^0	-	$\sqrt{s(s+t)}$	$\frac{1}{s}$
- , + , -	$\frac{1}{\sqrt{s}}$	-	$\frac{1}{\sqrt{s}}$	$\frac{m_W (s+t)}{\sqrt{-t}}$	$\frac{1}{\sqrt{s}}$	$\frac{m_W (s+t)}{\sqrt{-t}}$
- , + , +	s^0	-	s^0	s^0	s^0	$\frac{1}{s}$

Consistent with non-interference theorem in $2 \rightarrow 2$

[Cheung & Shen; PRL 115 (2015) 071601]
 [Azatov, Contino & Riva; PRD 95 (2017) 065014]

Maximum energy growth: E^2

- SU(2) current in (-,0,-) configuration → Interference growth!
- RH Charged Current (bR), Weak dipole (transverse) → no interference growth
- Subleading growth $\propto m_t E$, SM decreases with energy → no interference growth

FeynRules/NLOCT/UFO

FeynRules

[Christensen & Duhr; Comp. Phys. Comm. 180 (2009) 1614]

[Alloul et al.; Comp. Phys. Comm. 185 (2014) 2250]

- Framework: Lagrangian → Feynman rules → UFO model → MC events

Universal FeynRules Output (UFO)

[Degrande et al.; Comp. Phys. Comm. 183 (2012) 1201]

- Model file with particle content, internal/external parameters, Feynman rules, Lorentz structures, counter-terms,...
- Compatible with many MC event generators (MG5, Sherpa, Whizard,...)

NLOCT

[Degrande; Comp. Phys. Comm. 197 (2015) 239]

[Hahn; Comp. Phys. Comm. 140 (2001) 415]

- Automatic calculation of UV and R_2 counter-terms from FeynRules model
- Implemented as additional Feynman rules in the UFO format
- UV: on-shell renormalisation procedure for masses/wavefunction, $\overline{\text{MS}}$ for higher point functions
- R_2 : numerical artefacts of dimensional regularisation

ttH in SMEFT

 $(O_{t\phi}, O_{\phi G}, O_{tG})$

$$O_{t\phi} = (\varphi^\dagger \varphi) (\bar{Q}_L \tilde{\varphi} t_R)$$

$$O_{\phi G} = (\varphi^\dagger \varphi) G_{\mu\nu}^A G_A^{\mu\nu}$$

$$O_{tG} = (\bar{Q}_L \sigma_{\mu\nu} T^A t_R) \tilde{\varphi} G_A^{\mu\nu}$$

$$\frac{dC_i(\mu)}{d \log \mu} = \frac{\alpha_s}{\pi} \gamma_{ij} C_j(\mu) \quad \gamma = \begin{pmatrix} -2 & 16 & 8 \\ 0 & -7/2 & 1/2 \\ 0 & 0 & 1/3 \end{pmatrix}$$

Operators involving the **top/Higgs/gluon**

- ttH is the only direct probe of the Top-Higgs interaction
- $gg \rightarrow H$ & tt production partly constrain the Wilson coefficient space
- 3-gluon O_G and 4 fermion operators also contribute but are better constrained by tt and multi-jet measurements

Different **K-factors** among SM/dim-6 operators

Large Λ^{-4} effects in both shape & normalisation

- Scenarios where “EFT-squared” terms are large but energy is below cutoff

$$\frac{E^2}{\Lambda^2} < 1 < c_i^{(6)} \frac{E^2}{\Lambda^2} < c_i^{(6)} c_j^{(6)} \frac{E^4}{\Lambda^4}$$

EFT scale dependence

Set of running/mixing EFT couplings

- Additional source of theoretical uncertainty
- Like with α_s , can be estimated by scale variation

$$C_i(\mu) = \Gamma_{ij}(\mu, \mu_0) C_j(\mu_0) \quad \Gamma_{ij}(\mu, \mu_0) = \exp \left(\frac{-2}{\beta_0} \log \frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \gamma_{ij} \right)$$

$$\beta_0 = 11 - 2/3 n_f ,$$

$$\sigma(\mu_0) = \sigma_{SM} + \sum_i \frac{1\text{TeV}^2}{\Lambda^2} C_i(\mu_0) \sigma_i(\mu_0) + \sum_{i,j} \frac{1\text{TeV}^4}{\Lambda^4} C_i(\mu_0) C_j(\mu_0) \sigma_{ij}(\mu_0)$$

$\mu_0 \rightarrow \mu$

$$\sigma(\mu) = \sigma_{SM} + \sum_i \frac{1\text{TeV}^2}{\Lambda^2} C_i(\mu) \sigma_i(\mu) + \sum_{i,j} \frac{1\text{TeV}^4}{\Lambda^4} C_i(\mu) C_j(\mu) \sigma_{ij}(\mu)$$

$$= \sigma_{SM} + \sum_i \frac{1\text{TeV}^2}{\Lambda^2} C_i(\mu_0) \sigma_i(\mu_0; \mu) + \sum_{i,j} \frac{1\text{TeV}^4}{\Lambda^4} C_i(\mu_0) C_j(\mu_0) \sigma_{ij}(\mu_0; \mu)$$

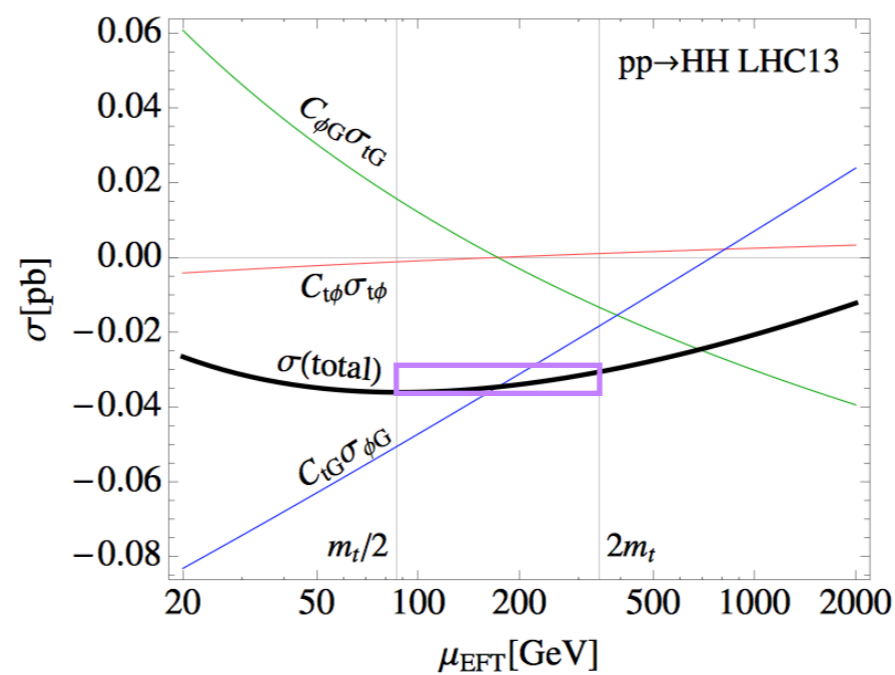
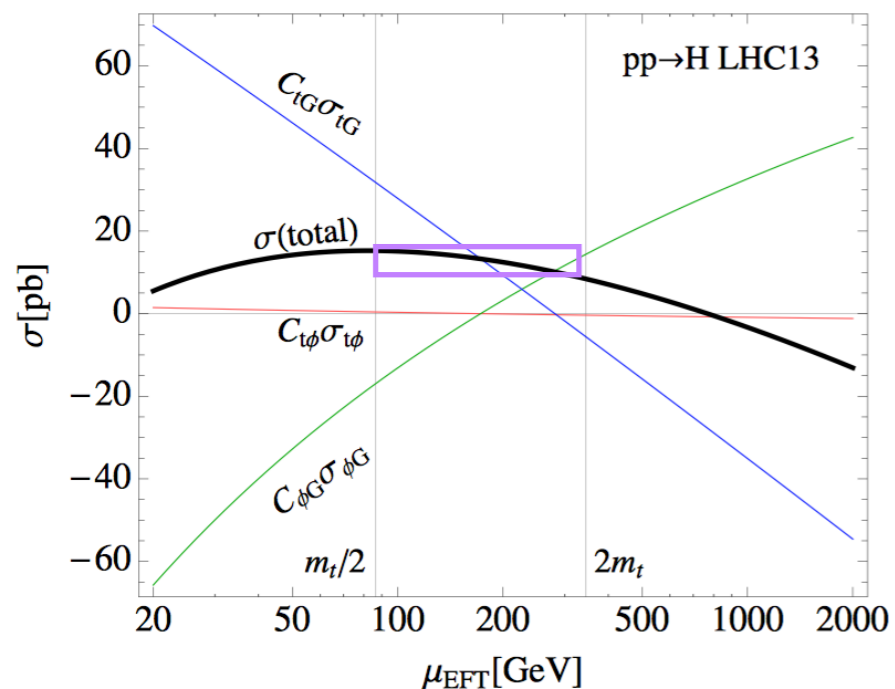
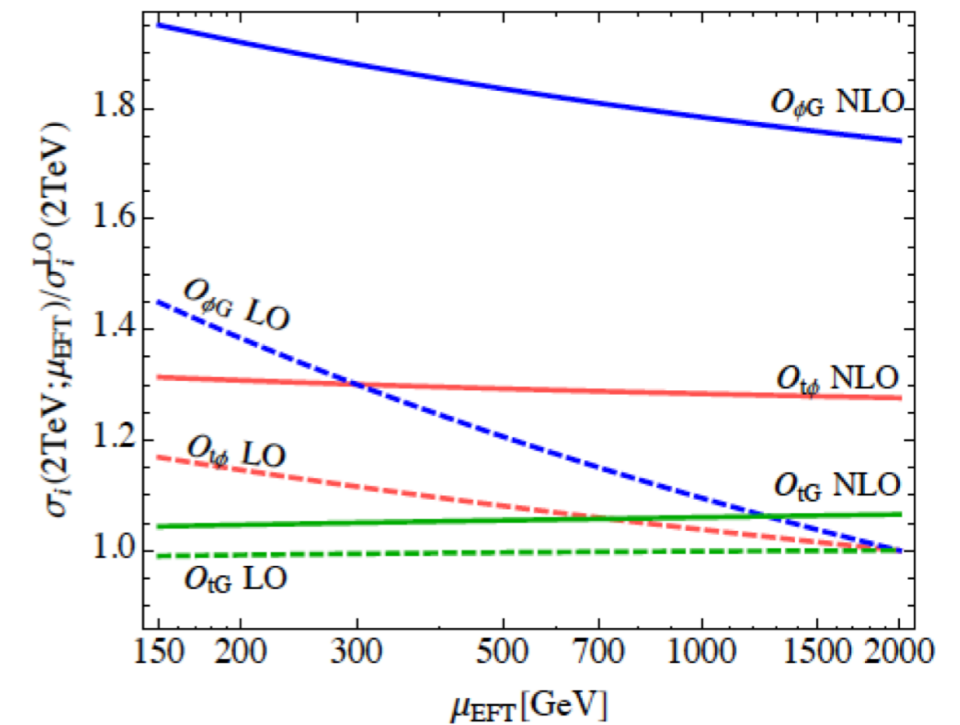
$$\sigma_i(\mu_0; \mu) = \Gamma_{ji}(\mu, \mu_0) \sigma_j(\mu) ,$$

$$\sigma_{ij}(\mu_0; \mu) = \Gamma_{ki}(\mu, \mu_0) \Gamma_{lj}(\mu, \mu_0) \sigma_{kl}(\mu)$$

EFT scale dependence

Full NLO stable under scale variation

- Large finite effects: RG improved underestimates NLO
- Scale uncertainty estimate
- Take c_i defined at scales $2\mu_0$ & $\mu_0/2$ and run back to the central scale



$\delta\mu_{\text{EFT}}$:
Does not cancel in
e.g. cross section
ratios

Existing limits

(I) : Individual
 (M) : Marginalised

Op.	TF (I)	TF (M)	RHCC (I) tree/loop	SFitter (I)	PEWM ²
\mathcal{O}_W				[-0.18,0.18]	
\mathcal{O}_{HW}				[-0.32,1.62]	
\mathcal{O}_{HB}				[-2.11,1.57]	
$\mathcal{O}_{\varphi W}$				[-0.39,0.33]	
$\mathcal{O}_{\varphi tb}$			[-5.28,5.28]/[-0.046,0.040]		
$\mathcal{O}_{\varphi Q}^{(3)}$	[-2.59,1.50]	[-4.19,2.00]			-1.0 ± 2.7 ³
$\mathcal{O}_{\varphi Q}^{(1)}$	[-3.10,3.10]				1.0 ± 2.7
$\mathcal{O}_{\varphi t}$	[-9.78,8.18]				1.8 ± 3.8
\mathcal{O}_{tW}	[-2.49,2.49]	[-3.99,3.40]			-0.4 ± 2.4
\mathcal{O}_{tB}	[-7.09,4.68]				4.8 ± 10.6
\mathcal{O}_{tG}	[-0.24,0.53]	[-1.07,0.99]			
$\mathcal{O}_{t\varphi}$				[-18.2,6.30]	
$\mathcal{O}_{Qq}^{(3,1)}$	[-0.40,0.60]	[0.66,1.24]			
$\mathcal{O}_{Qq}^{(3,8)}$	[-4.90,3.70]	[6.06,6.73]			

[TeV⁻²]

$$c_{t\varphi} \subset [-6.5, 1.3]$$

Combination of ttH @ 13 TeV

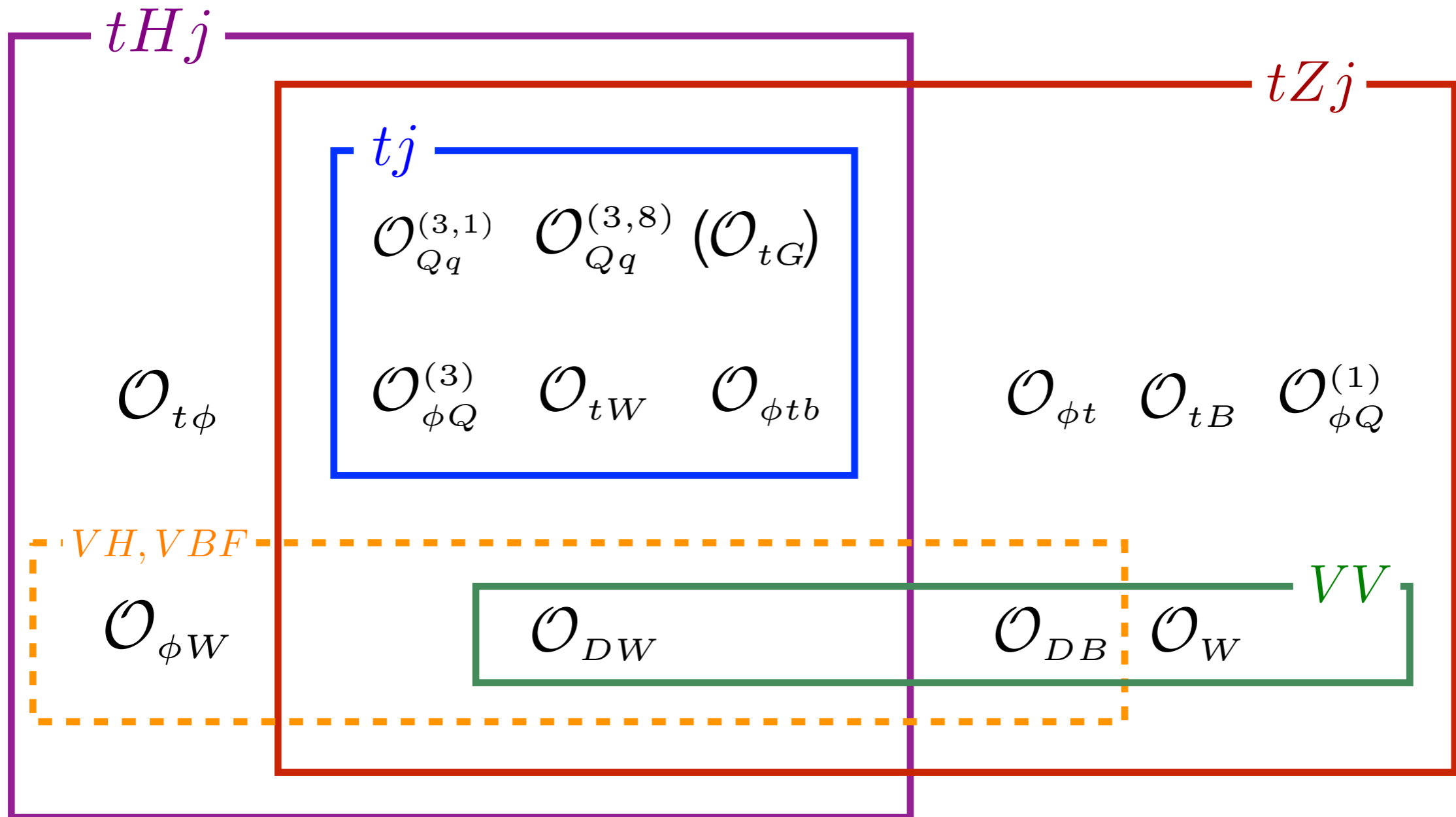
[CMS; CMS-PAS-HIG-17-003]
 [CMS; CMS-PAS-HIG-17-004]
 [ATLAS; CERN-EP-2017-281]

$$c_{Qq}^{(3,8)} \subset [-1.40, 1.20]$$

Combination of LHC single-top

[CMS; JHEP 12 (2012) 035]
 [ATLAS; PRD 90 (2014) 11, 112006]
 [CMS; JHEP 09 (2016) 027]
 [ATLAS; JHEP 04 (2017) 086]
 [ATLAS; EPJC 77 (2017) 8, 531]
 [ATLAS; PLB 756 (2016) 228-246]

Interplay



Anatomy of tZj

bW → tZ

$\lambda_b, \lambda_W, \lambda_t, \lambda_Z$	SM	$\mathcal{O}_{\varphi Q}^{(3)}$	$\mathcal{O}_{\varphi Q}^{(1)}$	$\mathcal{O}_{\varphi t}$	\mathcal{O}_{tB}	\mathcal{O}_{tW}	\mathcal{O}_W	\mathcal{O}_{HW}	\mathcal{O}_{HB}
- , 0 , - , 0	s^0	$\sqrt{s(s+t)}$	-	-	-	s^0	s^0	$\sqrt{s(s+t)}$	s^0
- , 0 , + , 0	$\frac{1}{\sqrt{s}}$	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$	$m_Z\sqrt{-t}$	$\frac{m_W(2s+3t)}{\sqrt{-t}}$	-	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$
- , - , - , 0	$\frac{1}{\sqrt{s}}$	$m_W\sqrt{-t}$	-	-	-	-	$\frac{m_W(s+2t)}{\sqrt{-t}}$	$m_W\sqrt{-t}$	$\frac{1}{\sqrt{s}}$
- , - , + , 0	$\frac{1}{s}$	s^0	s^0	s^0	s^0	$\sqrt{s(s+t)}$	s^0	s^0	$\frac{1}{\sqrt{s}}$
- , 0 , - , -	$\frac{1}{\sqrt{s}}$	$m_W\sqrt{-t}$	-	-	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$	$\frac{m_W(s+2t)}{\sqrt{-t}}$	$\frac{m_W(ss_W^2+2t)}{\sqrt{-t}}$	$\frac{m_W s}{\sqrt{-t}}$
- , 0 , - , +	$\frac{1}{\sqrt{s}}$	-	-	-	-	-	$\frac{m_W(s+t)}{\sqrt{-t}}$	$\frac{m_W(s+t)}{\sqrt{-t}}$	$\frac{m_W(s+t)}{\sqrt{-t}}$
- , 0 , + , -	s^0	s^0	s^0	-	-	s^0	s^0	s^0	s^0
- , 0 , + , +	$\frac{1}{s}$	s^0	s^0	s^0	$\sqrt{s(s+t)}$	$\sqrt{s(s+t)}$	-	s^0	s^0
- , + , - , 0	$\frac{1}{\sqrt{s}}$	-	-	-	-	-	$\frac{m_W(s+t)}{\sqrt{-t}}$	$\frac{1}{\sqrt{s}}$	$\frac{1}{\sqrt{s}}$
- , + , + , 0	s^0	s^0	-	-	-	s^0	-	s^0	$\frac{1}{s}$
- , - , - , -	s^0	s^0	s^0	-	s^0	s^0	s^0	s^0	s^0
- , - , - , +	$\frac{1}{s}$	-	-	-	-	-	$\sqrt{s(s+t)}$	s^0	s^0
- , - , + , -	$\frac{1}{\sqrt{s}}$	-	-	-	-	$\frac{m_Z(s_W^2 t - 3c_W^2(2s+t))}{\sqrt{-t}}$	-	$\frac{1}{\sqrt{s}}$	$\frac{1}{\sqrt{s}}$
- , - , + , +	-	-	-	-	$m_W\sqrt{-t}$	$m_Z\sqrt{-t}$	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$
- , + , - , -	$\frac{1}{s}$	-	-	-	-	-	$\sqrt{s(s+t)}$	s^0	s^0
- , + , - , +	s^0	s^0	s^0	-	-	-	-	s^0	s^0
- , + , + , -	$\frac{1}{\sqrt{s}}$	-	-	-	-	-	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$	$m_t\sqrt{-t}$
- , + , + , +	$\frac{1}{\sqrt{s}}$	-	-	-	-	$\frac{m_W(s+t)}{\sqrt{-t}}$	-	$\frac{1}{\sqrt{s}}$	$\frac{1}{\sqrt{s}}$

$\mathcal{O}_{\varphi tb}, \lambda_b, \lambda_t = +, +$

$\lambda_Z \backslash \lambda_W$	0	+	-
0	$\sqrt{s(s+t)}$	$m_W\sqrt{-t}$	-
+	$m_Z\sqrt{-t}$	s^0	-
-	-	-	s^0

$\mathcal{O}_{\varphi tb}, \lambda_b, \lambda_t = +, -$

$\lambda_Z \backslash \lambda_W$	0	+	-
0	-	-	s^0
+	s^0	-	-
-	s^0	-	-

Consistent with non-interference theorem in $2 \rightarrow 2$

[Cheung & Shen; PRL 115 (2015) 071601]
 [Azatov, Contino & Riva; PRD 95 (2017) 065014]

Non-interference

Sometimes, accidentally have $\sigma^{(6)}_{\text{INT}} < \sigma^{(6)}_{\text{SQ}}$

- **Non-interference** by e.g. helicity selection rules in the high energy limit

[Cheung & Shen; PRL 115 (2015) 071601]

[Azatov, Contino & Riva; PRD 95 (2017) 065014]

High energy theorem

- Many $2 \rightarrow 2$ amplitudes involving **at least one transverse gauge boson** mediated by D=6 operators do not interfere with the SM

Total Helicity

Interference?	A_4	$ h(A_4^{\text{SM}}) $	$ h(A_4^{\text{BSM}}) $
✗	VVVV	0	4,2
	VV $\phi\phi$	0	2
	VV $\psi\psi$	0	2
	V $\psi\psi\phi$	0	2
✓	$\psi\psi\psi\psi$	2,0	2,0
	$\psi\psi\phi\phi$	0	0
	$\phi\phi\phi\phi$	0	0

V = Transverse vector

ϕ = Longitudinal vector or Higgs

ψ = Fermion

$p p \rightarrow ZH, WH, WW, WZ$

Interference can be recovered:

finite mass effects

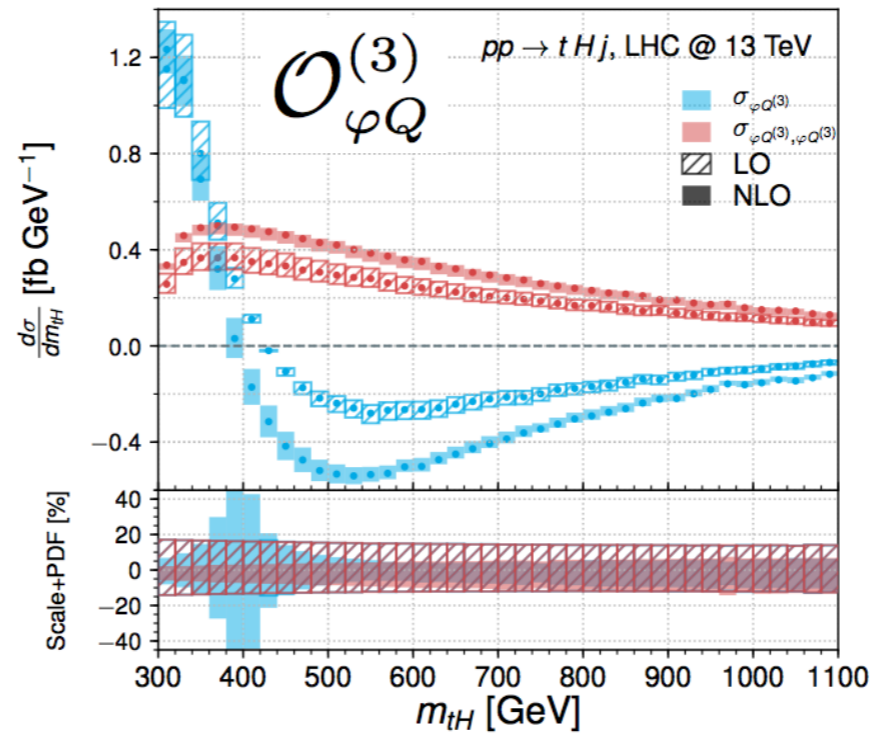
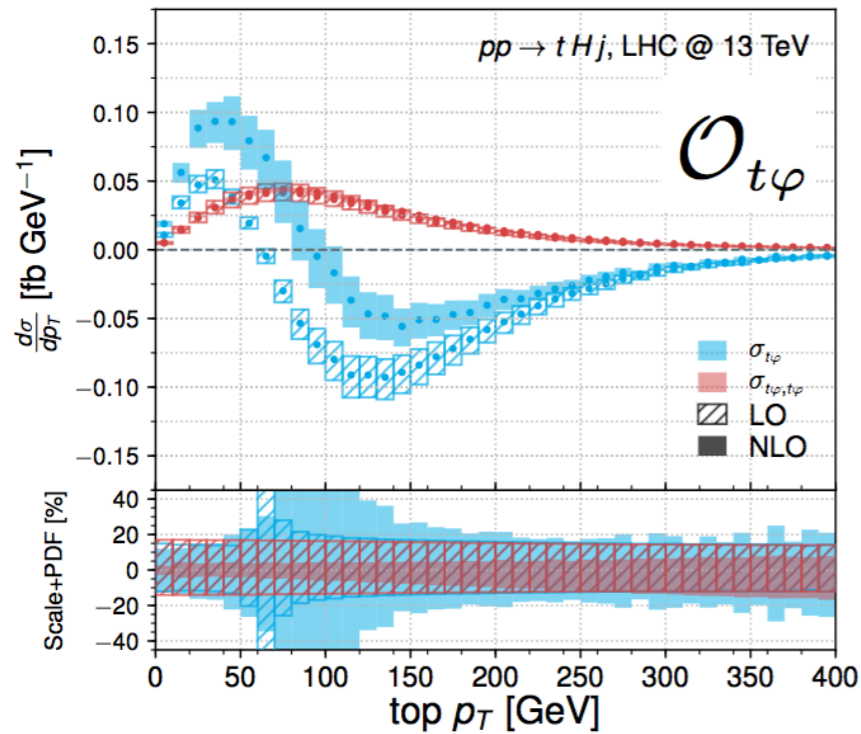
higher order corrections

higher multiplicity ($2 \rightarrow 3, 4, \dots$)

[Panico, Riva & Wulzer; CERN-TH-2017-85]

[Azatov, et al. LHEP 1710 (2017) 027]

Inclusive results: tHj



Cancellations over the PS appear/disappear for the interference contributions.

- Between top and antitop
- Strange K -factors
- Large scale uncert.

σ_{ij}	$C_{\varphi W}$	$C_{t\varphi}$	C_{tW}	$C_{\varphi Q}^{(3)}$	C_{HW}	C_{u31}	C_{u38}
$C_{\varphi W}$	—	2.752 (1.29)	12.88 (0.61)	6.384 (0.65)	-0.43 (-0.17)	—	—
$C_{t\varphi}$	2.514 (1.35)	—	-1.912 (-0.27)	-4.168 (-1.25)	-0.699 (-0.80)	—	—
C_{tW}	10.54 (0.68)	-1.772 (-0.32)	—	-26.24 (-0.79)	3.988 (0.46)	—	—
$C_{\varphi Q}^{(3)}$	5.12 (0.67)	-3.584 (-1.31)	-11.2 (-0.49)	—	4.864 (1.21)	—	—
C_{HW}	-0.402 (-0.18)	-0.6138 (-0.78)	3.124 (0.47)	3.5784 (1.10)	—	—	—
C_{u31}	-13.475 (-0.71)	5.16 (0.76)	-19.1 (-0.34)	-15.44 (-0.55)	-6.96 (-0.86)	—	4.525 (0.15)
C_{u38}	—	—	—	—	—	—	—

Inclusive results: tZj

$pp \rightarrow t(\bar{t}) Z j$ LHC@13 TeV

σ [fb]	LO	NLO	K-factor
σ_{SM}	$660.8(4)^{+13.7\%}_{-9.6\%} \pm 9.7\%$	$839.1(5)^{+1.1\%}_{-5.1\%} \pm 1.0\%$	1.27
σ_W	$-7.87(7)^{+8.4\%}_{-12.6\%} \pm 9.7\%$	$-8.77(8)^{+8.5\%}_{-4.3\%} \pm 1.1\%$	1.12
$\sigma_{W,W}$	$34.58(3)^{+8.2\%}_{-3.9\%} \pm 13.0\%$	$43.80(4)^{+6.6\%}_{-15.1\%} \pm 2.8\%$	1.27
σ_{tB}	$2.23(2)^{+14.7[0.9]\%}_{-10.7[1.0]\%} \pm 9.4\%$	$2.94(2)^{+2.3[0.4]\%}_{-3.0[0.7]\%} \pm 1.1\%$	1.32
$\sigma_{tB,tB}$	$2.833(2)^{+10.5[1.7]\%}_{-6.3[1.9]\%} \pm 11.1\%$	$4.155(3)^{+4.7[0.9]\%}_{-10.1[1.4]\%} \pm 1.7\%$	1.47
σ_{tW}	$2.66(4)^{+18.8[0.9]\%}_{-15.3[1.0]\%} \pm 11.4\%$	$13.0(1)^{+15.8[2.1]\%}_{-22.8[0.0]\%} \pm 1.2\%$	4.90
$\sigma_{tW,tW}$	$48.16(4)^{+10.0[1.7]\%}_{-5.8[1.9]\%} \pm 11.3\%$	$80.00(4)^{+7.9[1.3]\%}_{-14.7[1.6]\%} \pm 1.9\%$	1.66
$\sigma_{\varphi dtR}$	$4.20(1)^{+14.9\%}_{-10.9\%} \pm 9.3\%$	$4.94(2)^{+3.4\%}_{-6.7\%} \pm 1.0\%$	1.18
$\sigma_{\varphi dtR,\varphi dtR}$	$0.3326(3)^{+13.6\%}_{-9.5\%} \pm 9.6\%$	$0.4402(5)^{+3.7\%}_{-9.3\%} \pm 1.0\%$	1.32
$\sigma_{\varphi Q}$	$14.98(2)^{+14.5\%}_{-10.5\%} \pm 9.4\%$	$18.07(3)^{+2.3\%}_{-1.6\%} \pm 1.0\%$	1.21
$\sigma_{\varphi Q,\varphi Q}$	$0.7442(7)^{+14.1\%}_{-10.0\%} \pm 9.5\%$	$1.028(1)^{+2.8\%}_{-7.3\%} \pm 1.0\%$	1.38
$\sigma_{\varphi Q(3)}$	$130.04(8)^{+13.8\%}_{-9.8\%} \pm 9.5\%$	$161.4(1)^{+0.9\%}_{-4.8\%} \pm 1.0\%$	1.24
$\sigma_{\varphi Q(3),\varphi Q(3)}$	$17.82(2)^{+11.7\%}_{-7.5\%} \pm 10.5\%$	$23.98(2)^{+3.7\%}_{-9.3\%} \pm 1.4\%$	1.35
$\sigma_{\varphi tb}$	0	0	—
$\sigma_{\varphi tb,\varphi tb}$	$2.949(2)^{+10.5\%}_{-6.2\%} \pm 11.1\%$	$4.154(4)^{+5.1\%}_{-11.2\%} \pm 1.8\%$	1.41
σ_{HW}	$-5.16(6)^{+7.8\%}_{-12.0\%} \pm 10.5\%$	$-6.88(8)^{+6.4\%}_{-2.0\%} \pm 1.4\%$	1.33
$\sigma_{HW,HW}$	$0.912(2)^{+9.4\%}_{-5.2\%} \pm 12.0\%$	$1.048(2)^{+5.2\%}_{-12.8\%} \pm 2.1\%$	1.15
σ_{HB}	$-3.015(9)^{+9.9\%}_{-13.9\%} \pm 9.5\%$	$-3.76(1)^{+5.2\%}_{-1.0\%} \pm 1.0\%$	1.25
$\sigma_{HB,HB}$	$0.02324(6)^{+12.7\%}_{-8.5\%} \pm 9.9\%$	$0.02893(6)^{+2.3\%}_{-7.5\%} \pm 1.1\%$	1.24
σ_{tG}		$0.45(2)^{+93.0\%}_{-148.8\%} \pm 4.9\%$	—
$\sigma_{tG,tG}$		$2.251(4)^{+20.9\%}_{-30.0\%} \pm 2.5\%$	—
$\sigma_{Qq(3,1)}$	$-393.5(5)^{+8.1\%}_{-12.3\%} \pm 10.0\%$	$-498(1)^{+8.9\%}_{-3.2\%} \pm 1.2\%$	1.26
$\sigma_{Qq(3,1),Qq(3,1)}$	$462.25(3)^{+8.4\%}_{-4.1\%} \pm 12.7\%$	$545.50(5)^{+7.4\%}_{-17.4\%} \pm 2.9\%$	1.18
$\sigma_{Qq(3,8)}$	0	$-0.9(3)^{+23.3\%}_{-26.3\%} \pm 19.2\%$	—
$\sigma_{Qq(3,8),Qq(3,8)}$	$102.73(5)^{+8.4\%}_{-4.1\%} \pm 12.7\%$	$111.18(5)^{+9.3\%}_{-18.4\%} \pm 2.8\%$	1.08

$tZj \sim 10$ times bigger than tHj

NLO corrections: similar features to tHj

EFT contributions smaller relative to SM

Higgs always radiated from top/EW gauge boson
Z boson can also come from light quark leg