

# Stable Sexaquark as Dark Matter



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# Stable Sexaquark as Dark Matter



How could we have missed a  
stable particle made of quarks?

[Hints from Astrophysics]

Primordial Nucleosynthesis

Dark-Matter to Ordinary-Matter ratio

Detecting S dark matter

# Unique among multi-quark states:

Fermi statistics is compatible with a  
totally symmetric spatial wave function AND

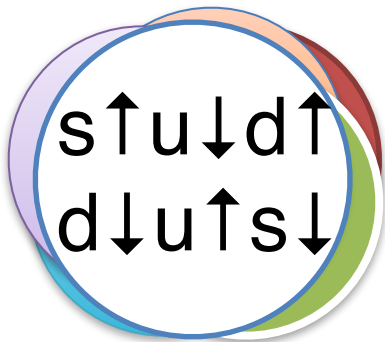
antisymmetric (singlet) in:

*color*

*flavor*

*spin*

totally symmetric in space



S

**(Most-Attractive Channel)<sup>3</sup> :**

6-quark,  $Q=0$ ,  $B=2$

*Spin-0, scalar*

*Flavor singlet*

$m_S < 2 \text{ GeV}???$

Same quark content as H-dibaryon (Jaffe 1977), but different physics: **not a loosely bound di- $\Lambda$ !**

*mass  $\sim 2150 \text{ MeV}$  in bag model*

# Why consider $m_S \sim 2 m_p$ ?

- Light quarks almost massless, i.e. relativistic
    - $m_{u,d} \approx 0$ ,  $m_s = 91 \text{ MeV}$
  - S has same QNs as ground state glueball
    - why not  $m_S \approx m_{\text{glueball}} + 180 \text{ MeV} = (1.5-1.7) + 0.18 \text{ GeV} \approx 2 m_p$
  - 3x di-quark mass = 1.2 - 2-ish GeV
  - $m_S < 2 (m_p + m_e)$  : S is absolutely stable
  - $m_S > 2 (m_p - 8 \text{ MeV})$  : nuclei are stable
  - triple-singlet (color, flavor, spin): MAC, lattice, almost all models  $\Rightarrow m_S < 2 m_\Lambda$
  - *extensive experimental searches **exclude** weak-lifetime &  $m > 2 \text{ GeV}$*
- bound state exists and mass  $< 2 \text{ GeV}$  ( $\tau > \tau_{\text{Univ}}$  or stable)



# Stable Sexaquark Hypothesis

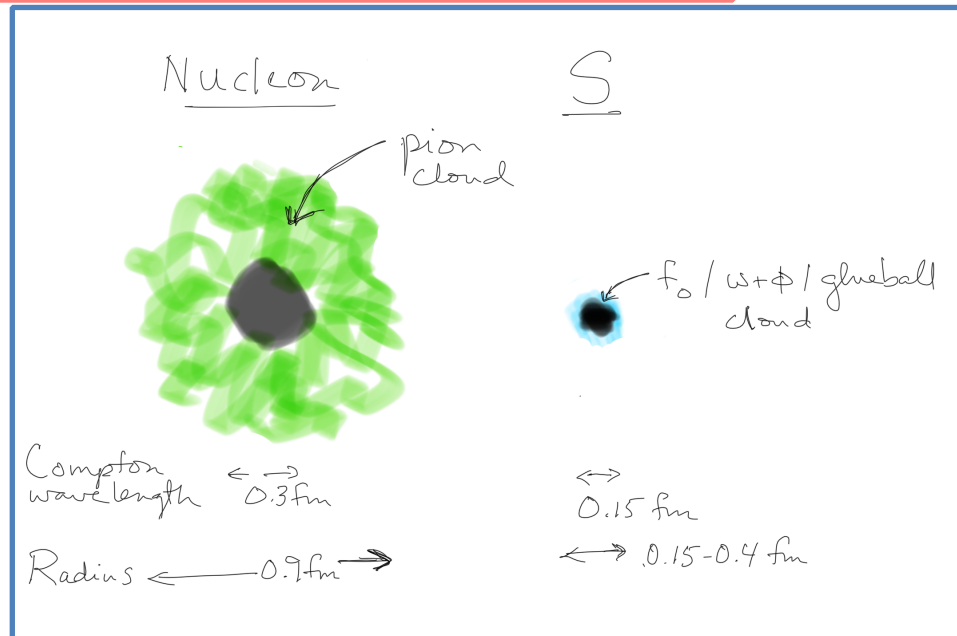
6	sexa- <sup>[19]</sup>	–	sen- <sup>[20]</sup>	sext- <sup>[21]</sup>	hex- <sup>[22]</sup>	hexakis- hexaplo- hexad- e.g. <a href="#">hexahedron</a>	hect- <sup>[23]</sup> hectai-	shat-
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<sup>a b</sup> Sometimes Greek *hexa-* is used in Latin compounds, such as [hexadecimal](#), due to [taboo avoidance](#) with the English word [sex](#).  
[https://en.wikipedia.org/wiki/Numeral\\_prefix](https://en.wikipedia.org/wiki/Numeral_prefix)

## **Crucial fact:**

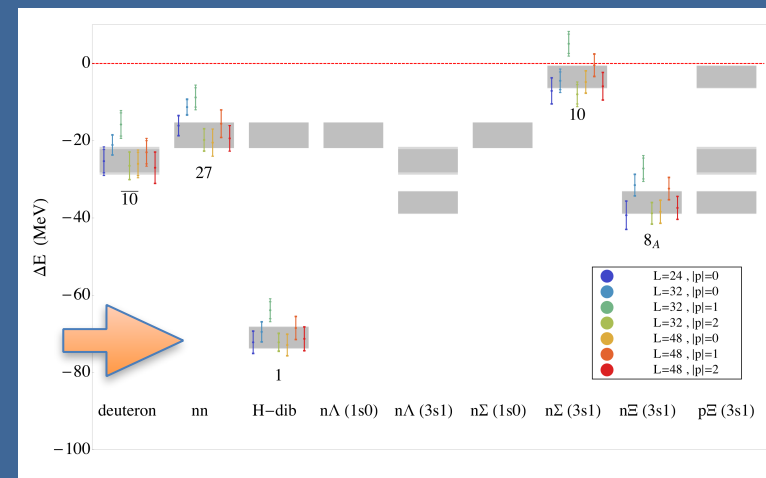
***S does not couple to pions. ∴ much smaller than usual hadrons*** (proton, pion,...)

6-quark,  $Q=0$ ,  $B=2$   
*Spin-0, scalar*  
*Flavor singlet*  
 $m \sim 1.7\text{-}2 \text{ GeV}$



# Stable S?

- $\tau > \tau_{\text{Univ}}$  if:
  - $M_S < 2 m_p + 2 m_e = 1877.6 \text{ MeV} \rightarrow \text{absolutely stable}$
  - if  $M_S < m_p + m_e + m_\Lambda$ ,  $\Gamma \sim G_F^4 \times (\text{wave function overlap})^2$
- **Lattice predicts binding (Beane+13)**
  - ( $m_q = 850 \text{ MeV}$  so not realistic)
  - 80 MeV binding
- **Experiments exclude decaying S**  
**=> it must be STABLE ! ; -)**



# Conditions on QCD Dark Matter

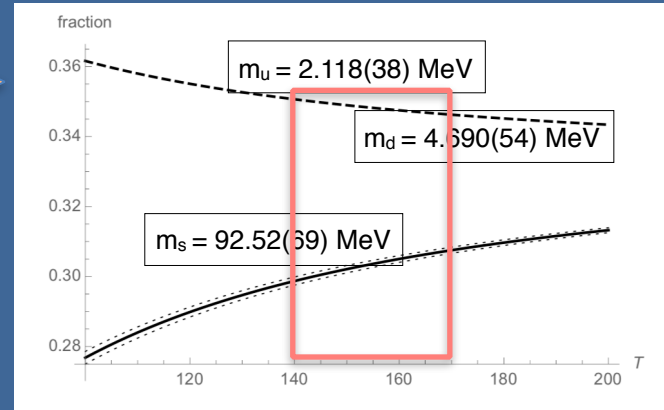
- ✓  $\tau_{\text{DM}} > \tau_{\text{Univ}}$  , cold, neutral
- ✓ primordial nucleosynthesis
- ✓ Particle must not be already excluded
  - accelerator searches
  - exotic isotopes
  - **DM searches**
  - indirect impacts (heating planets, helioseismology,...)
  - stability of nuclei
  - equation of state of neutron stars (and their stability)
- ✓ Correct relic density (for natural  $m_{\text{DM}}$  &  $\sigma_{\text{f.o.}}$ )

# $\Omega_{DM} / \Omega_b$

follows from **stat mech** , **quark masses** & temperature of **QGP-hadronization** transition

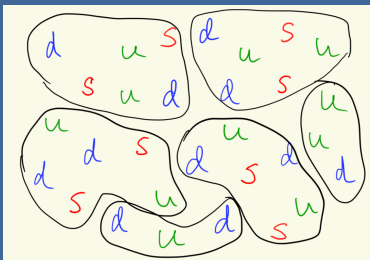
$$\frac{\Omega_{DM}}{\Omega_b} = \frac{y_b \kappa_s 3f_s}{1 - \kappa_s 3f_s}$$

$m_s / (2m_p)$

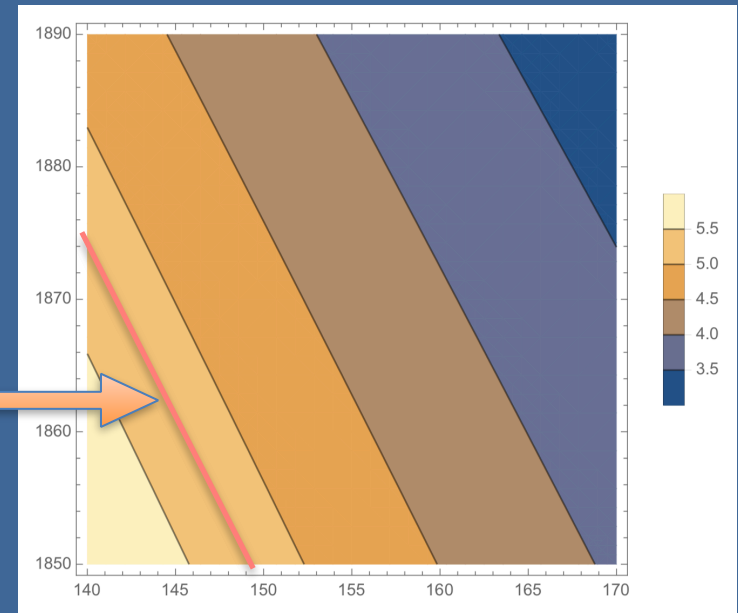


$$\kappa_s(m_s, T) = \frac{1}{1 + (r_{\Lambda, \Lambda} + r_{\Lambda, \Sigma} + 2r_{\Sigma, \Sigma} + 2r_{N, \Xi})}$$

$$r_{1,2} \equiv \exp[-(m_1 + m_2 - m_s)/T]$$



$$\Omega_{DM} / \Omega_b = 5.3 \pm 0.1$$

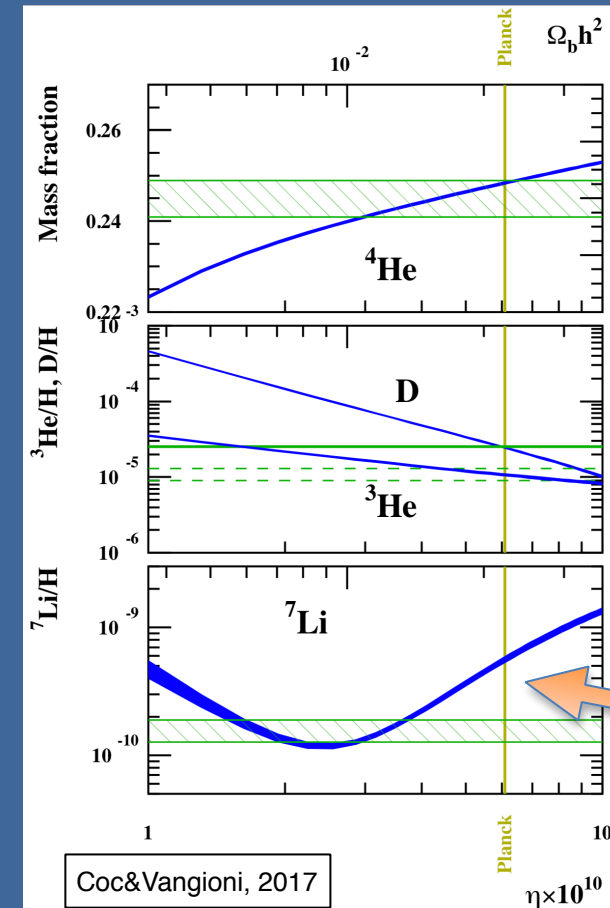


Prediction also applies to strange quark nuggets...

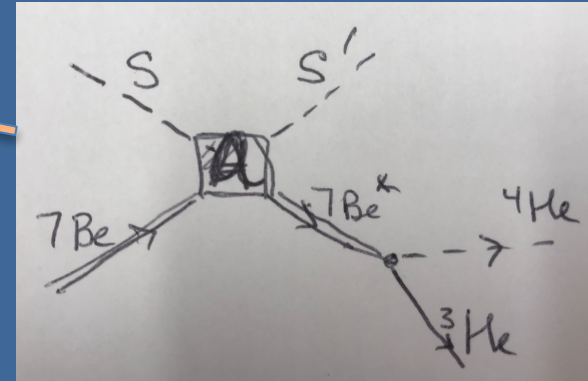
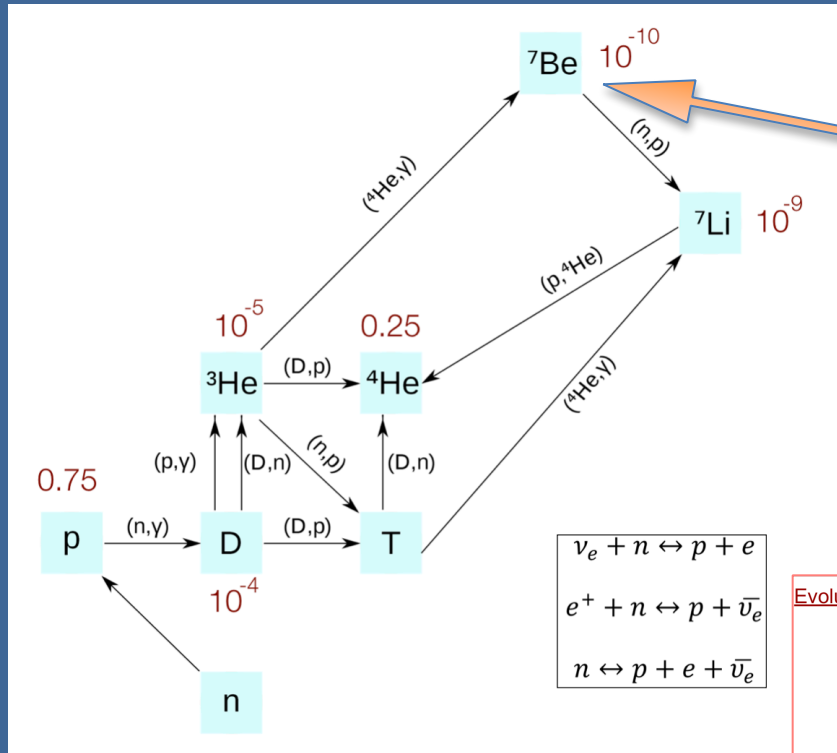
# Problem with primordial ${}^7\text{Li}$

(GRF + Richard Galvez, in preparation)

- Big Bang Nucleosynthesis works brilliantly *except*
  - Predicted abundance of  ${}^7\text{Li} = (5.61 \pm 0.26) 10^{-10}$
  - Observed abundance of  ${}^7\text{Li} = (1.58 \pm 0.31) 10^{-10}$
- Discrepancy is now very serious:
  - Nuclear rates all well-measured
  - $\eta = n_b/n_r = (6.58 \pm 0.02) 10^{-10}$  from CMB
  - Astrophysics now secure (Spite plateau):
    - small scatter
    - flat > 3 decades of low metallicity
- **S solves the puzzle**
  - No other (reasonable) solution known

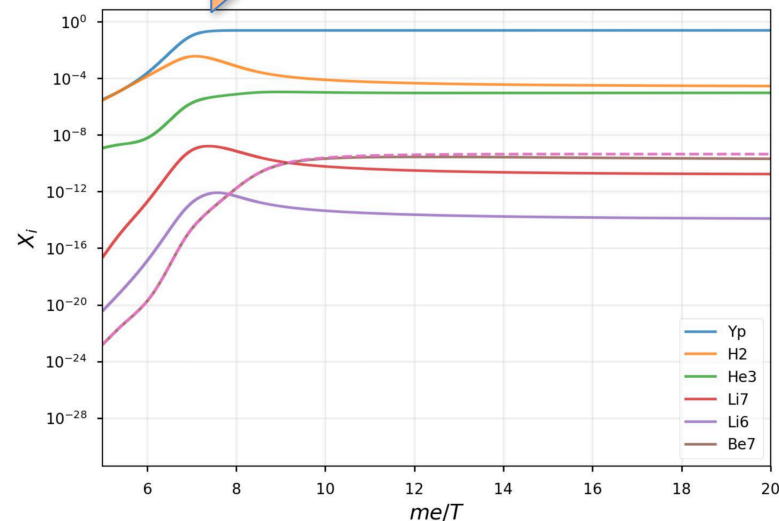


# S dark matter breaks up ${}^7\text{Li}$ & ${}^7\text{Be}$



- **Solves the puzzle**
  - **Doesn't affect He or d**
  - E threshold for breakup =  
1.58, 2.46, 4.47, 5.75, 19.3 [2.2] MeV
- ${}^7\text{Be}$   ${}^7\text{Li}$   ${}^3\text{He}$  T  ${}^4\text{He}$  [d]

Evolution of abundances.



Standard  ${}^7\text{Be}$   
case is dashed  
line

# Dark Matter with Hadronic Interactions

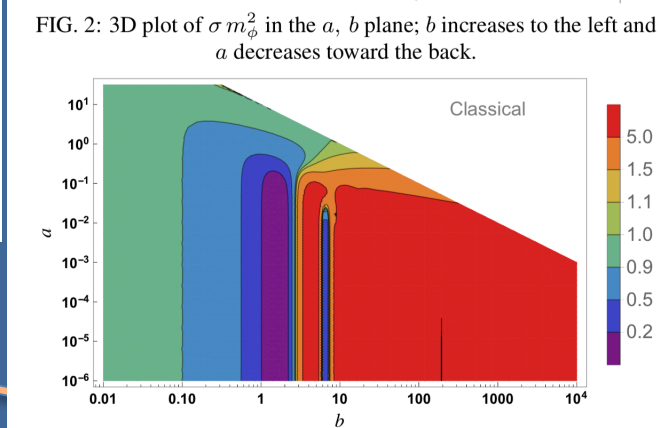
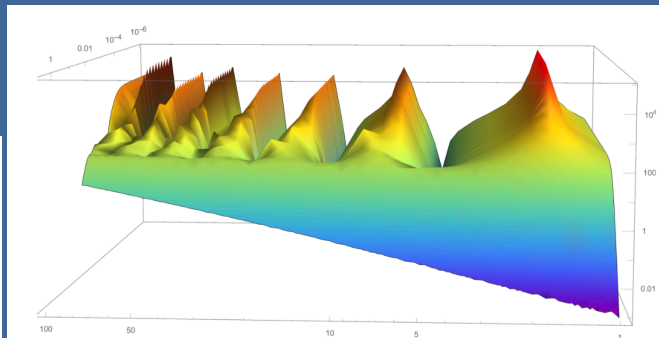
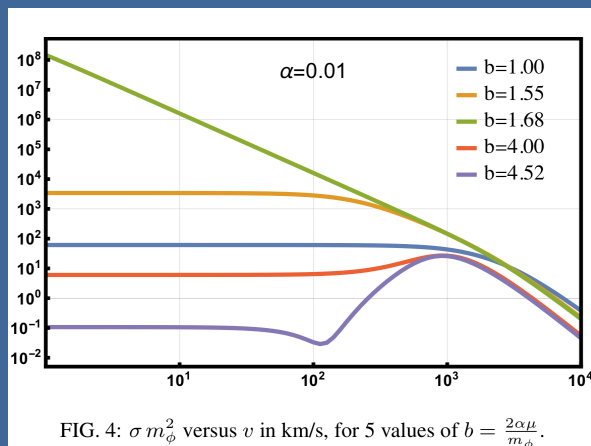
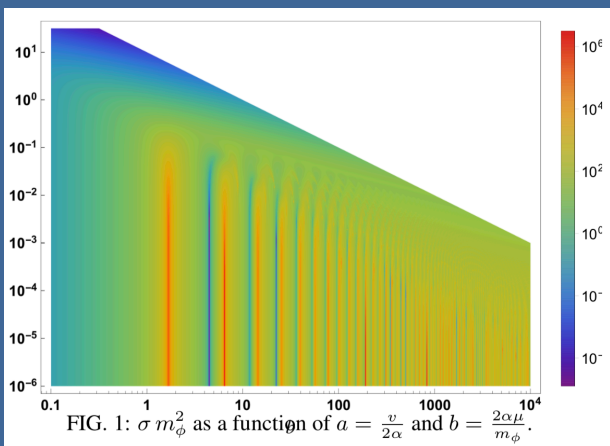
(GRF + Xingchen Xu, to appear shortly)

$$V(r) = \frac{\alpha}{r} e^{-r m_\phi}$$

$m_\phi = 1$  GeV (flavor-singlet  $\omega$ - $\phi$  combo), sourced by p or A

- $v/c$  (DM)  $\sim 10^{-3}$   $10^3$  km/s (galaxy clusters) down to 1 km/s (atm &  $z = 17$ )
  - must solve Schroedinger Eqn. **Born approximation may fail badly**
  - cross section depends only on combos  $a = \frac{v}{2\alpha}$  and  $b = \frac{2\alpha\mu}{m_\phi}$

$$a = \frac{v}{2\alpha} \text{ and } b = \frac{2\alpha\mu}{m_\phi}$$

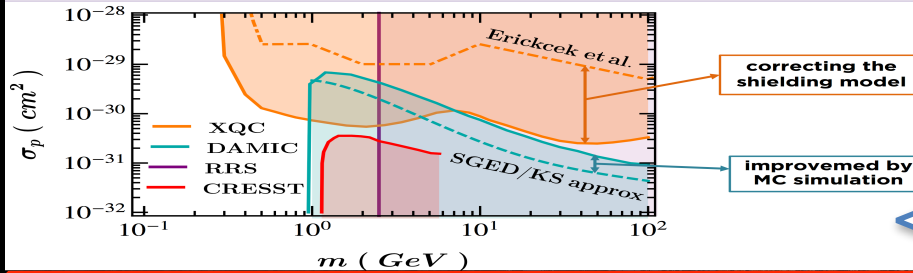


$10^{-26}$  -  $10^{-25}$   $\text{cm}^2$



# Stable S as Dark Matter

Traditional Window is Now Closed !



If Born Approximation good & XQC efficiency is perfect

Closing the window on  $\sim \text{GeV}$  Dark Matter with moderate ( $\sim \mu\text{b}$ ) interaction with nucleons

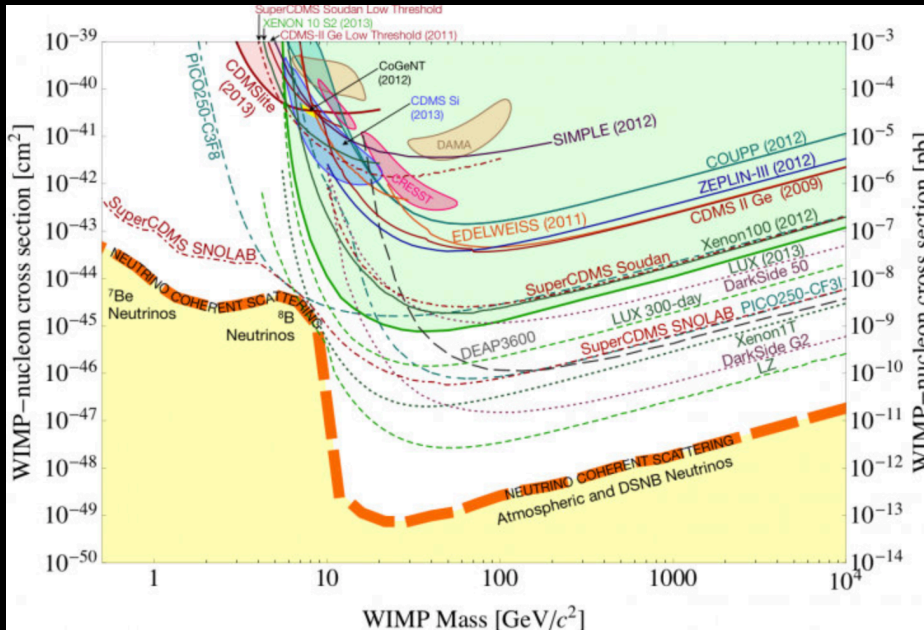
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**Abstract.** We improve limits on the spin-independent scattering cross section of Dark Matter on nucleons, for DM in the 300 MeV – 100 GeV mass range, based on the DAMIC and XQC experiments. Our results close the window which previously existed in this mass range, for a DM-nucleon cross section of order  $\sim \mu\text{b}$ , assuming the standard velocity distribution.

Shielded (e.g. underground) detectors are not sensitive

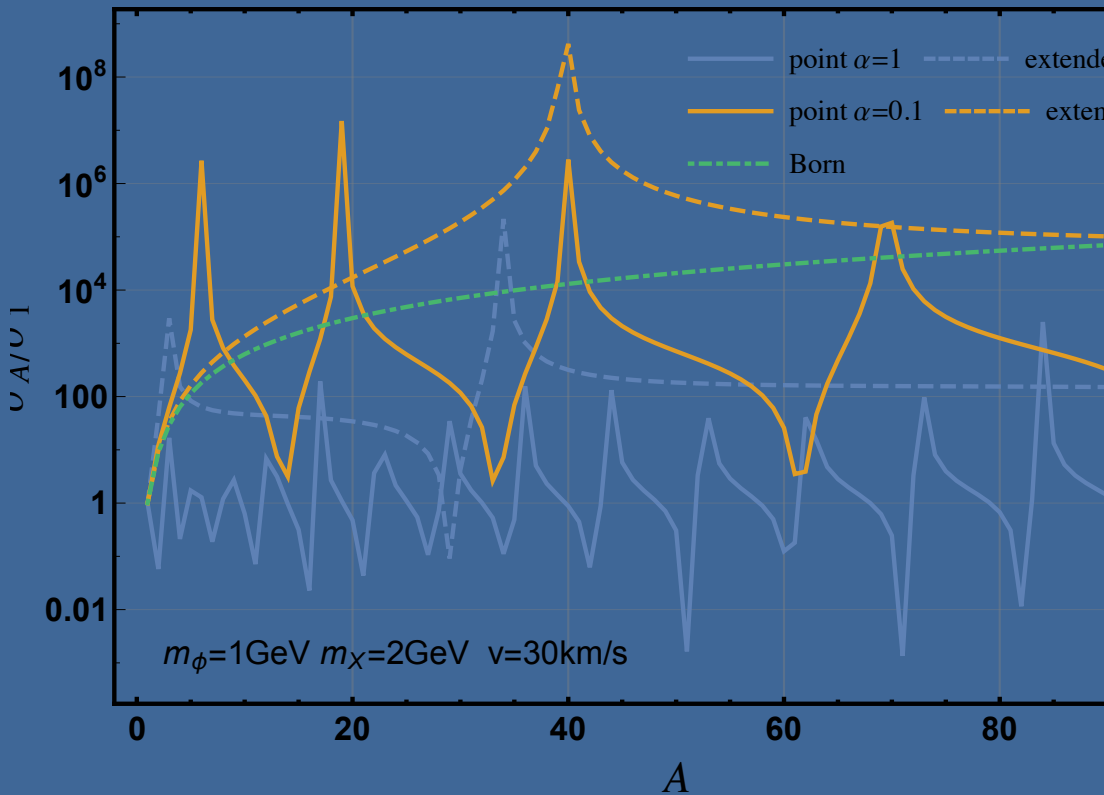




# Plenty of Room for HIDM

(GRF + Xingchen Xu, to appear shortly)

Caution:  $A$ -dependence very sensitive to nuclear form factor. Born approximation often misleading, by orders of magnitude.



Allowed regions of coupling from best Direct Detection

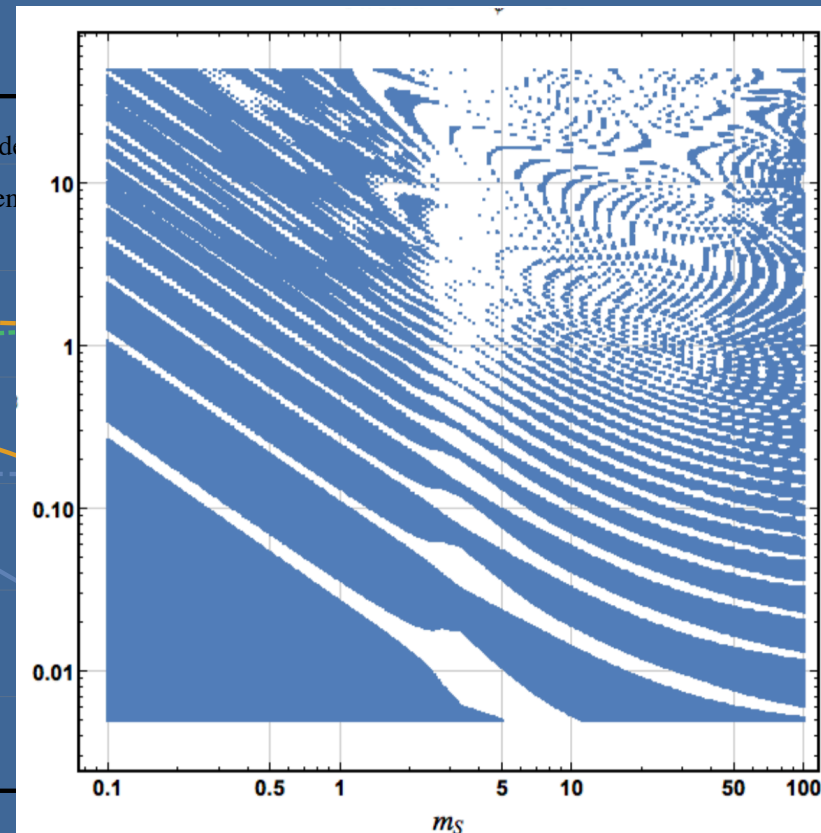


FIG. 7: Allowed regions (blue) in the coupling-DM mass plane  $\alpha$  (vertical axis) and  $m_{DM}$  in GeV (horizontal axis) from XQC using

# S has not been discovered because it is elusive

- **Many negative searches, but all are inapplicable.** They either:
  - looked for H-dibaryon through decays (but S is stable)
  - restricted to mass  $> 2$  GeV (but  $m_S < 2$  GeV)
  - required fast production in  $S=-2$  hypernuclei (but small overlap with baryons)
- **Wavefunction overlap with baryons is very small.** Extremely rare fluctuation required for  $S \leftrightarrow \Lambda\Lambda$ ;  $S \leftrightarrow NN$  is  $G_F^4$  smaller  $\Rightarrow$ 
  - nuclei can be stable ( $\tau > 10^{29}$  yr) even for  $m_S > 2 m_p$
  - hard to produce in fixed target experiments
- **S is similar to** (much more copious) **neutrons**
- **Promising accelerator detection strategies**
  - **Apparent lack of baryon number and strangeness conservation:**
    - $\Delta B = \pm 2$  with  $\Delta S = \mp 2$
  - **Reconstruct missing mass, e.g.:**
    - $\Upsilon \rightarrow \Lambda \Lambda \bar{S}$  (+ pions)  $M_{S^2} = (p_Y - p_{\Lambda 1} - p_{\Lambda 2} - \sum p_{\pi_i})^2$

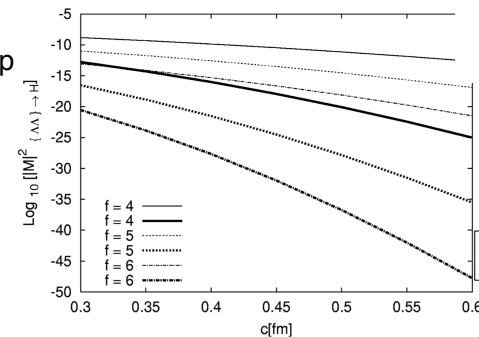
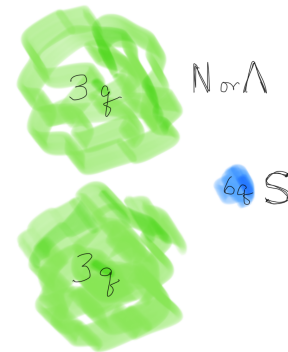


FIG. 1.  $\text{Log}_{10}$  of  $|M|^2_{\Lambda\Lambda \rightarrow H}$  versus hard core radius in femtometers, for ratio  $f = R_H / R_H$  and two values of the Isgur-Karl oscillator parameter:  $\alpha_H = 0.406$  GeV (thick lines) and  $\alpha_H = 0.221$  GeV (thin lines).

GRF+G.Zaharijas 2004

# Experimental searches so far

Looking for Jaffe's H-dibaryon (same QN but assumed to be unstable and  $r \sim 1$  fm)

- Require  $M > 2$  GeV:
  - Gufstafson+ FNAL1976 : Beam-dump + tof *Limit on production of neutral stable strongly interacting particle with mass  $> 2$  GeV.*
  - Carroll+ BNL 1978: No narrow missing mass peak above 2 GeV in  $pp \rightarrow K K X$
- Require H-dibaryon decay:
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  - Bernstein+ FNAL 1988: Limit on production of neutral with  $10^{-8} < \tau < 2 \times 10^{-6}$  s
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  - Kim+ Belle 2013: no narrow resonance in  $\Upsilon \rightarrow \Lambda p K$
- Limits from production in doubly-strange hypernuclei:
  - Ahn+ BNL 2001
  - Takahashi+ KEK 2001

## Search for Six-Quark States

A. S. Carroll, I.-H. Chiang, R. A. Johnson, T. F. Kycia, K. K. Ki,  
L. S. Littenberg, and M. D. Marx  
Brookhaven National Laboratory, Upton, New York 11973

and

R. Cester, R. C. Webb, and M. S. Witherell  
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(Received 26 July 1978)

We have searched the missing-mass spectrum of the reaction  $pp \rightarrow K^+ K^+ X$  for a narrow six-quark resonance in the mass range 2.0–2.5 GeV/ $c^2$ . No narrow structure was observed. Upper limits for the production cross section of such a state depend upon mass and vary from 30 to 130 nb.

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PHYSICAL REVIEW LETTERS

29 APRIL 1996

## Search for the Weak Decay of an H Dibaryon

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<sup>8</sup>College of William and Mary, Williamsburg, Virginia 23187

(Received 8 December 1995)

We have searched for a neutral H dibaryon decaying via  $H \rightarrow \Lambda n$  and  $H \rightarrow \Sigma^0 n$ . Our search has yielded two candidate events from which we set an upper limit on the H production cross section. Normalizing to the inclusive  $\Lambda$  production cross section, we find  $(d\sigma_H/d\Omega)/(d\sigma_\Lambda/d\Omega) < 6.3 \times 10^{-6}$  at 90% C.L., for an H of mass  $\approx 2.15$  GeV/ $c^2$ . [S0031-9007(96)00050-6]

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PHYSICAL REVIEW LETTERS

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## Production of $\Lambda\Lambda^4\text{H}$ Hypernuclei

J. K. Ahn,<sup>1,2</sup> S. Ajimura,<sup>10</sup> H. Aikawa,<sup>9</sup> B. Bassalleck,<sup>9</sup> A. Berdoz,<sup>2</sup> D. Carman,<sup>2</sup> R. E. Chrien,<sup>1</sup> C. A. Davis,<sup>8,14</sup> P. Eugenio,<sup>2</sup> H. Fischer,<sup>3</sup> G. B. Franklin,<sup>2</sup> J. Franz,<sup>2</sup> T. Fukuda,<sup>12</sup> L. Gan,<sup>4</sup> H. Hotchi,<sup>12</sup> A. Ichikawa,<sup>7</sup> K. Imai,<sup>7</sup> S. H. Kahana,<sup>1</sup> P. Khaustov,<sup>2</sup> T. Kishimoto,<sup>10</sup> P. Koran,<sup>2</sup> H. Kohri,<sup>10</sup> A. Kourepin,<sup>10</sup> K. Kubota,<sup>12</sup> M. Landry,<sup>8</sup> M. May,<sup>1</sup> C. Meyer,<sup>2</sup> Z. Meziani,<sup>10</sup> S. Minami,<sup>10</sup> T. Miyachi,<sup>12</sup> T. Nagae,<sup>2</sup> J. Nakano,<sup>12</sup> H. Ota,<sup>5</sup> K. Paschke,<sup>2</sup> P. Pile,<sup>1</sup> M. Prokhabatlov,<sup>6</sup> B. P. Quinn,<sup>2</sup> V. Rasin,<sup>6</sup> A. Rusek,<sup>1</sup> H. Schmitt,<sup>3</sup> R. A. Schumacher,<sup>2</sup> M. Sekimoto,<sup>3</sup> K. Shileev,<sup>6</sup> Y. Shimizu,<sup>10</sup> R. Sutter,<sup>1</sup> T. Tamagawa,<sup>12</sup> L. Tang,<sup>4</sup> K. Tanida,<sup>12</sup> K. Yamamoto,<sup>2</sup> M. L. Yuan<sup>4</sup>

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<sup>10</sup>Department of Physics, Osaka University, Toyonaka, Osaka 560-0043, Japan

<sup>11</sup>Department of Physics, Temple University, Philadelphia, Pennsylvania 19122

<sup>12</sup>Department of Physics, University of Tokyo, Tokyo 113-0033, Japan

<sup>13</sup>Department of Physics, Pusan National University, Pusan 609-735, Korea

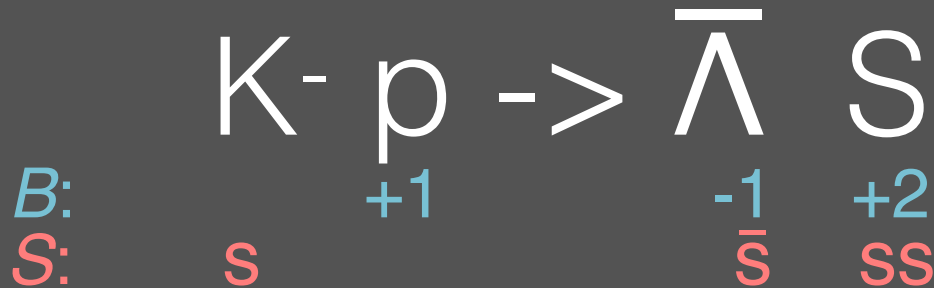
<sup>14</sup>TRIUMF, 4004 Wesbrook Mall, Vancouver, BC, Canada V6T 2A3

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(Received 14 May 2001; published 5 September 2001)

An experiment demonstrating the production of double- $\Lambda$  hypernuclei in  $(K^-, K^+)$  reactions on  $^9\text{Be}$  was carried out at the D6 line in the BNL alternating-gradient synchrotron. The technique was the observation of pions produced in sequential mesonic weak decay, each pion associated with one unit of strangeness change. The results indicate the production of a significant number of the double hypernucleus  $\Lambda\Lambda^4\text{H}$  and the twin hypernuclei  $^4\text{H}$  and  $^3\text{H}$ . The relevant decay chains are discussed and a simple model of the production mechanism is presented. An implication of this experiment is that the existence of an  $S = -2$  dibaryon more than a few MeV below the  $\Lambda\Lambda$  mass is unlikely.

Classic Approach: would be great, but *very* low rate



- $\bar{\Lambda}$  is a gold-plated signature :  $\bar{\Lambda} \rightarrow \pi^+ \bar{p}$ 
  - Easy to ID & reconstruct 4-momentum
  - $c\tau = 8 \text{ cm}$     all  $\bar{\Lambda}$  are ID'd
- S: undetected, but 4 momentum determined
  - $p_S = p_K + p_p - p_{\bar{\Lambda}}$
  - NA61: est.  $\sim 20 \text{ MeV}$  accuracy on “missing-mass” of S
  - For  $p_{\text{beam}} < 5.35 \text{ GeV}/c$  , no conventional source of  $\bar{\Lambda}$  's
- NA61:  $9 \text{ GeV}/c$   $K^-$  beam

# Experimental Searches

- Require  $M > 2 \text{ GeV}$ :
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VOLUME 76, NUMBER 18 PHYSICAL REVIEW LETTERS 29 APRIL 1996

**Search for the Weak Decay of an  $H$  Dibaryon**

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 (Received 8 December 1995)

We have searched for a neutral  $H$  dibaryon decaying via  $H \rightarrow \Lambda n$  and  $H \rightarrow \Sigma^0 n$ . Our search has yielded two candidate events from which we set an upper limit on the  $H$  production cross section. Normalizing to the inclusive  $\Lambda$  production cross section, we find  $(d\sigma_H/d\Omega)/(d\sigma_\Lambda/d\Omega) < 6.3 \times 10^{-6}$  at 90% C.L., for an  $H$  of mass  $\approx 2.15 \text{ GeV}/c^2$ . [S0031-9007(96)00050-6]

VOLUME 87, NUMBER 13 PHYSICAL REVIEW LETTERS 24 SEPTEMBER 2001

**Production of  $_{\Lambda\Lambda}^4\text{H}$  Hypernuclei**

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An experiment demonstrating the production of double- $\Lambda$  hypernuclei in  $(K^-, K^+)$  reactions on  $^9\text{Be}$  was carried out at the D6 line in the BNL alternating-gradient synchrotron. The technique was the observation of pions produced in sequential mesonic weak decay, each pion associated with one unit of strangeness change. The results indicate the production of a significant number of the double hypernucleus  $_{\Lambda\Lambda}^4\text{H}$  and the twin hypernuclei  $_{\Lambda\Lambda}^4\text{H}$  and  $_{\Lambda\Lambda}^4\text{He}$ . The relevant decay chains are discussed and a simple model of the production mechanism is presented. An implication of this experiment is that the existence of an  $S = -2$  dibaryon more than a few MeV below the  $\Lambda\Lambda$  mass is unlikely.

# Discovery Strategy

- Apparent lack of B and S conservation:

- $\Delta B = \pm 2 \quad + \quad \Delta S = \mp 2$

- Reconstruct missing mass, e.g.:

- $\Upsilon \rightarrow \Lambda \Lambda \bar{S}$  (+ pions)  $M_S^2 = (p_\Upsilon - p_{\Lambda 1} - p_{\Lambda 2} - \sum p_{\pi i})^2$

- LHC:  $\bar{S} + N \rightarrow \bar{\Lambda} K^+$   $M_S^2 = (p_{\bar{\Lambda}} + p_K - p_N)^2$

- Snolab nuclei:  $pn \rightarrow S e^+ \nu$   $G_F^4 \quad \tau > 10^{29} \text{ yr}$

$$\Upsilon \rightarrow \Lambda \Lambda \bar{S} \text{ \& } \bar{\Lambda} \bar{\Lambda} S$$

(+ pions)

- $\Upsilon$  is **localized** source of  $ggg$

$\Rightarrow$  production of  $S$  is (relatively) enhanced

- Many  $\times 10^8$  events collected (CLEO, Babar, Belle)

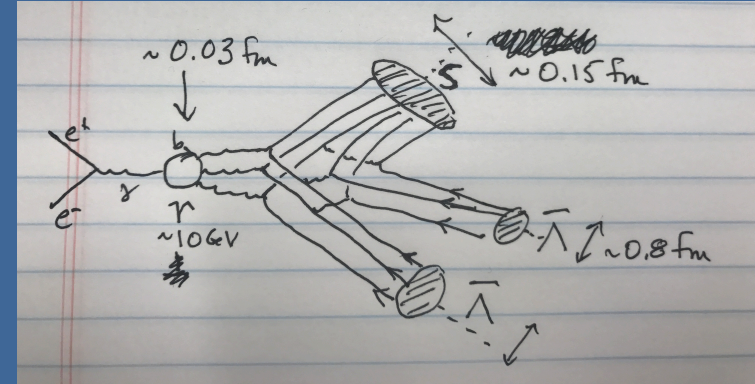
- detectors pretty hermetic, have good mass resolution,  $O(10 \text{ MeV})$
- $\Lambda$  decays quickly to  $p\pi$  so easy to ID.  $c\tau = 8 \text{ cm}$

- Can MEASURE  $m_S$  via missing mass

- **Very clean**

- Main bkg is  $K_S K_S K_L K_L$  (+ pions)
  - $K_S$ 's mis-ID'd as  $\Lambda$ 's and  $K_L$ 's escaping before decay : *negligible for Belle*
  - rare and can model accurately
  - $K_S K_S K_L K_L$  (+ pions) *is measurable*, from  $K^+ K^+ K^- K^-$  (+ pions)
- “Conspiracy” of missed particles producing  $\Delta B = \pm 2$ ,  $\Delta S = \mp 2$  very hard

**Background does not have narrow peak in missing mass!**



# LHC

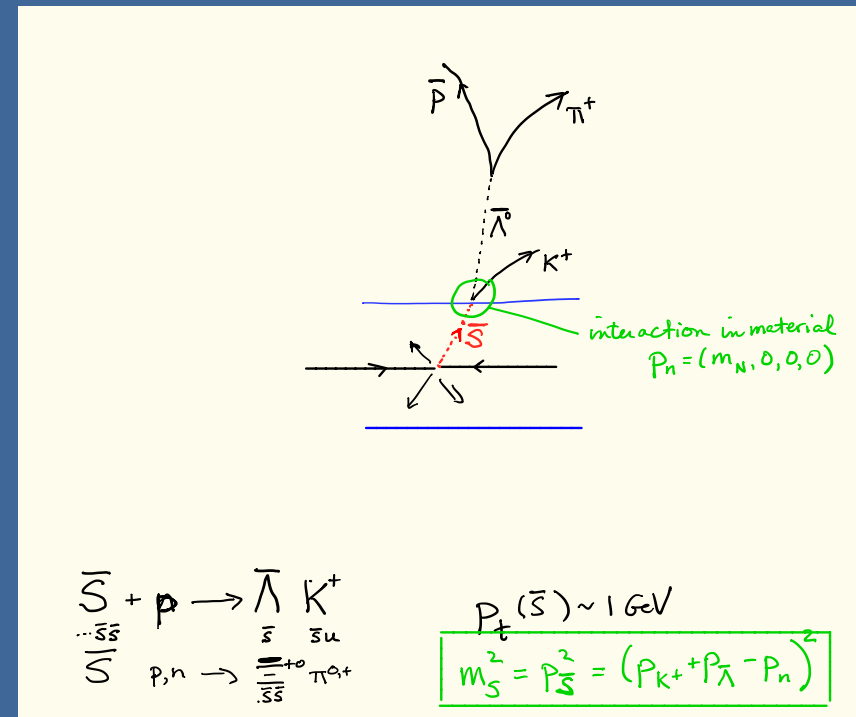
- Hadronic collisions: low production rate due to small wfn overlap
- Find a needle in a haystack ( $10^{11}$  recorded events; potential for trigger >>> more)
- Statistical examination of correlation between

$$\Delta B = \pm 2, \Delta S = \mp 2$$

- **$\bar{S}$  annihilation in tracker, tag by**

$$\bar{E}^{+,0} \rightarrow \bar{\Lambda} \pi^{+,0} \quad (c\tau = 57 \text{ cm}) \quad \bar{\Lambda} \rightarrow \bar{p} \pi^+$$

$$\text{or } \bar{\Lambda} K^+$$



- 2nd exponential in scattering-length distribution of n-like interactions, due to S



# Cosmology & structure formation

- **DM-baryon interaction: momentum transfer => *slight drag on DM during structure formation***
  - Dvorkin, Blum, Kamionkowski (2014):
    - **Ly-alpha forest:  $\sigma < \sim 10$  mb if v-indept — no problem for S**
  - Buen-Abad, Marques-Tavares, Schmalz (2015):
    - **momentum transfer helps reconcile  $H_0$  &  $\sigma_8$**
  - Boring or an opportunity? To be determined...
- **S-S self interactions + S-baryon interactions:**
  - may have similar benefits as Self Interacting DM
    - core-cusp, “too-big-to-fail” & missing sub-halos problems.

# Galaxies & Clusters

**DM-gas scattering provides a source of heating, needed for:**

- Milky Way's extended hot gas halo —  $2 \times 10^6$  K
- Quenching star formation
- Avoiding “cooling flow catastrophe” in X-ray clusters

# Key points to take home

- ***There may a tightly bound 6-quark state  $S = uuddss$*** 
  - Unique, symmetric structure  $\Rightarrow$  other hadrons don't provide guidance
    - mass is not driven by chiral symmetry breaking (unlike baryons)
    - constituent quark model probably completely misleading
  - ***If  $M_S < 2 m_p + 2 m_e$  ,  $S$  is absolutely stable***
- ***If  $S$  is stable, its an excellent Dark Matter candidate***
  - Relic abundance is natural. **EXPLAINS 7Li Discrepancy in BBN and Dark Matter to baryon ratio**
  - Usual WIMP detection strategy isn't applicable.
  - Baryon asymmetric Universe is expected
  - May reconcile tension in  $H_0$  &  $\sigma_8$  and explain astrophysics puzzles ("quenching", core-cusp, DM rotation curves...)
- ***$S$  may be waiting to be discovered in existing  $\Upsilon$ -decays or LHC experiments... mass can be accurately measured in  $\Upsilon$ -decay exclusive final states.***
- ***SDM will be challenging to detect, but not impossible. Astrophysical and cosmological effects may allow it to be constrained, excluded or confirmed.***

# Backup Slides

# Strategies to detect DM

## *if DM is comprised of S's*

- With  $v_{\text{rel}} = 30 \text{ km/s}$  ,  $KE_{\text{DM}} \sim 10 \text{ eV}$

- $\langle E_{\text{deposit}} \rangle / KE_{\text{DM}} = 0.12$  for Si target
  - $= 0.02$  for Hg target
  - $= 0.44$  for H or He target

- Energy-loss length of S in Earth crust:

$$\lambda = 2 \text{ cm} / \sigma_{10mb}$$

- Present detectors shielded or too high threshold (new CRESST expt has  $E_{\text{th}} = 19 \text{ eV}$ , but 30 cm shielding  $\Rightarrow$  not sensitive)

- Heating rate liquid He:  $\sim \text{nW/mol} \sim \text{CR muon energy deposit rate}$ .
- can't *shield* muons & other CRs: veto? but what about neutrons?

### Results on MeV-scale dark matter from a gram-scale cryogenic calorimeter operated above ground

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July 24, 2017

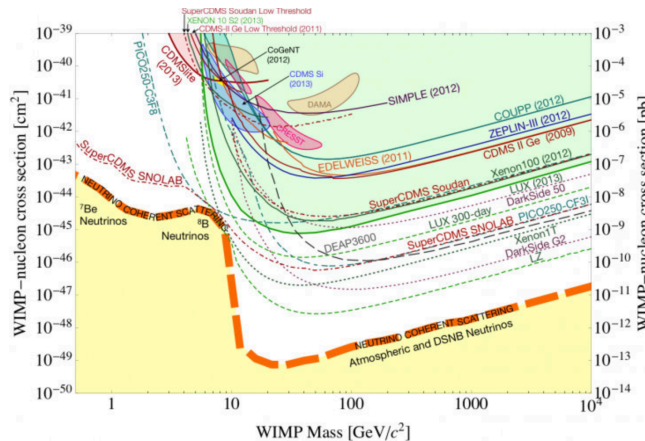
**Abstract.** Models for light dark matter particles with masses below  $1 \text{ GeV}/c^2$  are a natural and well-motivated alternative to so-far unobserved weakly interacting massive particles. Gram-scale cryogenic calorimeters provide the required detector performance to detect these particles and extend the direct dark matter search program of CRESST. A prototype  $0.5 \text{ g}$  sapphire detector developed for the  $\nu$ -cleus experiment has achieved an energy threshold of  $E_{\text{th}} = (19.7 \pm 0.9) \text{ eV}$ . This is one order of magnitude lower than for previous devices and independent of the type of particle interaction. The result presented here is obtained in a setup above ground without significant shielding against ambient and cosmogenic radiation. Although operated in a high-background environment, the detector probes a new range of light-mass dark matter particles previously not accessible by direct searches. We report the first limit on the spin-independent dark matter particle-nucleon cross section for masses between  $140 \text{ MeV}/c^2$  and  $500 \text{ MeV}/c^2$ .

1707.06749v1 [astro-ph.CO] 21 Jul 2017

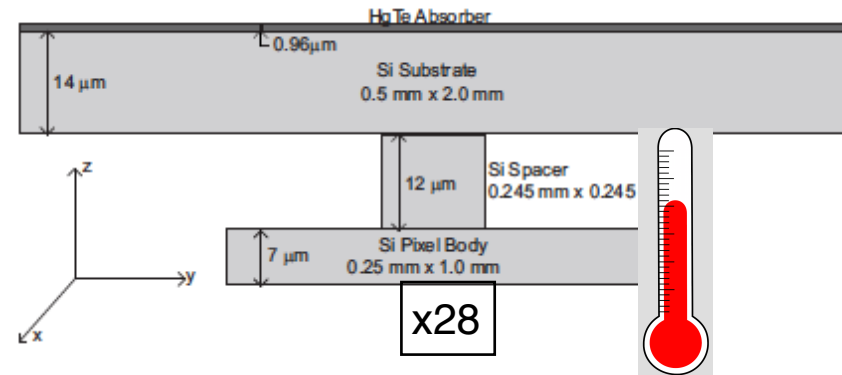
# Directly Detecting the S

- **Light, slow-moving, strongly interacting DM is not visible in current detectors:**

- $KE_{DM} = 500 \text{ eV} (m_{DM} / 2 m_p) v^2 / (220 \text{ km/s})^2$
- $\langle E_{deposit} \rangle / KE_{DM} = 0.12$  for Si target  
 $= 0.02$  for Hg target  
 $= 0.44$  for H or He target
- Energy-loss length in Earth crust:  $\lambda = 2 \text{ cm} / \sigma_{10mb}$
- **CRESST, Xenon1T, LUX, DAMIC are not sensitive**



# X-Ray Quantum Calorimeter (XQC)



- On sounding rocket, 200 km above earth
- Best limit for high x-secn (McCammon+02, Wandelt+02, GF+Zaharijas05, Erickcek+07, Mahdawi & GF 17)
  - sensitive to X-rays with  $E \gtrsim 29$  eV
  - 100 sec flight,  $\sim 100$  events
  - nuclear recoil  $\Rightarrow$  X-rays, which thermalize (*assumption*)

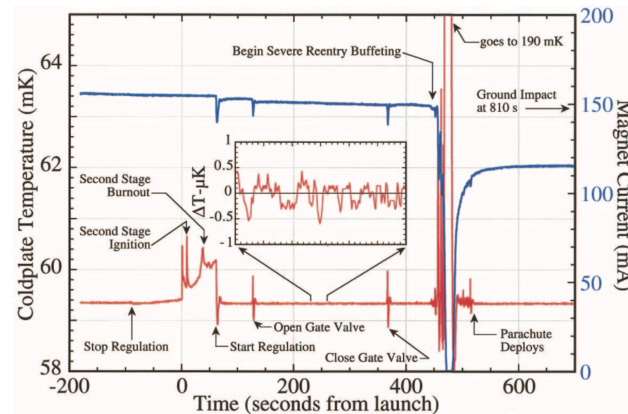
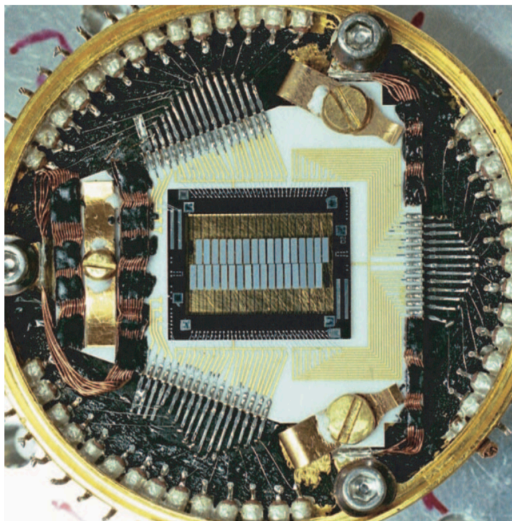
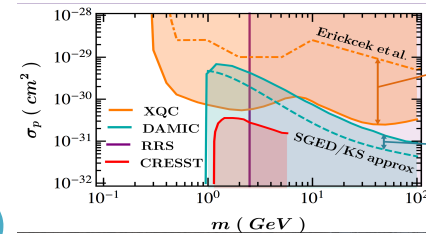


FIG. 7.—In-flight performance of the temperature control system, showing the coldplate temperature and magnet current. Temperature fluctuations during data taking are about 210 nK rms. The gate-valve motor is located on the vacuum jacket and caused the most serious thermal disturbance up to reentry. Accelerations during reentry exceeded 20 g with tumbling at  $\sim 1$  Hz, introducing heat to the cold stage faster than it could be removed. Temperature regulation is recovered once tumbling stops, allowing calibration data to be obtained.

## Calibrate with X-rays

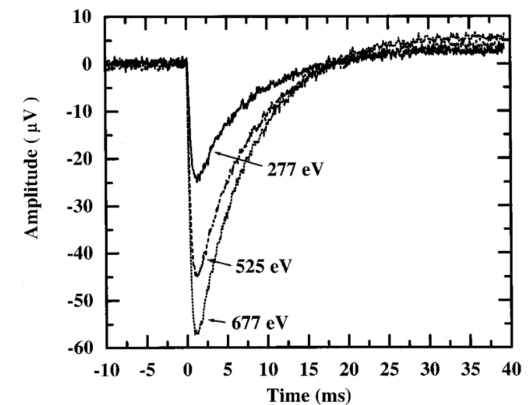


FIG. 8.—Unfiltered X-ray pulses from the gate-valve calibration source.

# Challenge to Direct Detection of S

- Taking XQC sensitivity at face value:

**XQC above atmosphere is best.  $E_{\text{thresh}} = 30 \text{ eV}$**

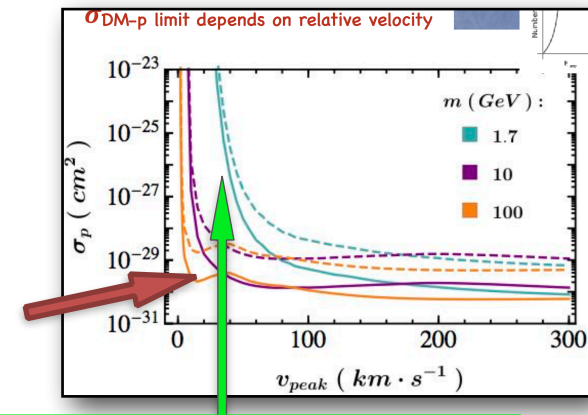
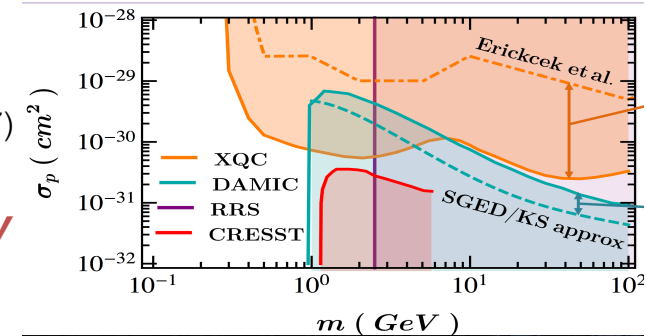
(McCammon+02, Wandelt+02, GF+Zaharijas05, Erickcek+07, Mahdawi+GRF 2017)

- Mahdawi+GRF 2017:  $\sigma_{\text{DM}} < 10^{-29} \text{ cm}^2$  *for standard velocity dispersion* — **SDM has  $\sigma_{\text{DM}} \sim 10\text{-}100 \text{ mb}$**

- $v_{\text{min}, \text{XQC}} = 100 \text{ km/s } (2 m_p / m_{\text{DM}})^{1/2}$

- **BUT:**

- **velocity distribution?**
  - (gas co-rotates, red vel  $\sim 10 \text{ km/s}$ )
- **does  $E_{\text{rec}}$  really thermalize?**



**Co-rotation reduces SDM signal**



# Closer look at XQC sensitivity

Silicon nucleus recoil:  $KE_{\max} \sim 500 \text{ eV} \Rightarrow v_{\max} \sim 20 \text{ km/s} \ll v_e \Rightarrow$

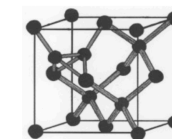
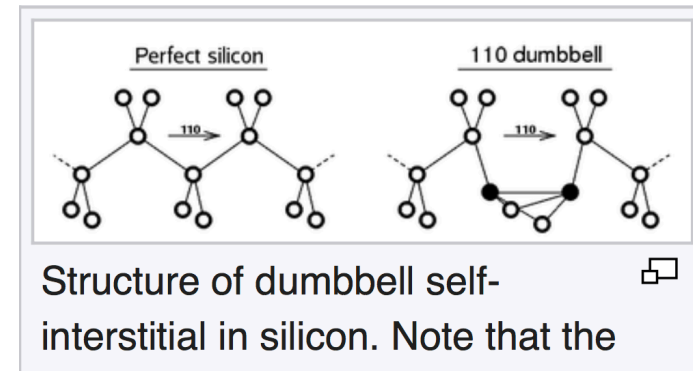
- atomic interaction is adiabatic  $\Rightarrow$
- negligible ionization.

Si atom moving in semiconductor crystal:

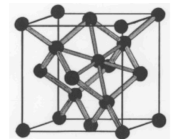
- rearranges covalent bonds
- produces interstitial defects
- 500eV atom produces Frenkel pairs (V+I)
  - $E_{\text{Fp\_min}} = 5 \text{ eV}$
  - $E_{\text{migration}} \sim 0.1 \text{ eV}$
- Cascade energy loss producing
  - $N \sim (KE_{\text{rec}} / 5 \text{ eV})$  Frenkel pairs,

**$\sim 2\%$  of  $KE_{\text{rec}}$  goes to thermalization**

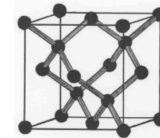
**$\Rightarrow KE_{\text{rec,min}} > 1.5 \text{ keV} \Rightarrow KE_{\text{DM,min}} > 6 \text{ keV} \Rightarrow$   
 $v_{\text{DM,min}} > 300 \text{ km/s}$**



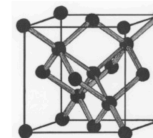
(a) Split-(110)



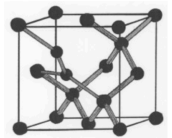
(b) Hexagonal



(c) Perfect crystal



(d) Tetrahedral



(e) Concerted exchange

**Figure 1.** Illustrations of the split-(110), hexagonal, and tetrahedral self-interstitial defects, together with the perfect crystal and the saddle point of Pandey's concerted exchange.