# STATUS OF STANDARD MODEL PREDCTTIONS FOR RD(") 

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## LEPTON UNIVERSALITY VIOLATION?

> Deviations in $\mathrm{B} \rightarrow \mathrm{D}^{(*)}$ tv decays found in multiple measurements over the last years, almost $4 \sigma$ disagreem with SM prediction

- Other hints of lepton universality violations in other decay modes


$$
\left.\left.\begin{array}{ll}
\left.\mathrm{R}(J / \psi)\right|_{\text {exp }}=\frac{\mathrm{BR}\left(B_{c} \rightarrow J / \psi \tau \nu\right)}{\mathrm{BR}\left(B_{c} \rightarrow J / \psi \ell \nu\right)}=0.71 \pm 0.17 \pm 0.18 & \text { vs }
\end{array} \quad \mathrm{R}(J / \psi)\right|_{t h}=0.25-0.28\right)
$$

Is it New Physics? Interesting BSM interpretations $\rightarrow$ see talks in later sessions
> To assess discrepancy one need up-to-date predictions for the SM, with careful assessment of theoretical uncertainties
> Uncertainties come from form factors

- FFs determined by combination of
> data
> lattice QCD
> theoretical modeling
> subset of FFs affect also $\mathrm{V}_{\mathrm{cb}}$ exclusive determination and long standing discrepancy there between exclusive and inclusive determinations


## $B \rightarrow D, D^{*}$ DECAYS: NOTATION

$$
\begin{aligned}
\frac{\mathrm{d} \Gamma(\bar{B} \rightarrow D l \nu)}{\mathrm{d} w}= & \frac{G_{F}^{2}\left|V_{c b}\right|^{2} \eta_{\mathrm{EW}}^{2} m_{B}^{5}}{48 \pi^{3}}\left(w^{2}-1\right)^{3 / 2} r_{D}^{3}\left(1+r_{D}\right)^{2} \mathcal{G}(w)^{2} \\
\frac{\mathrm{~d} \Gamma\left(\bar{B} \rightarrow D^{*} l \nu\right)}{\mathrm{d} w}= & \frac{G_{F}^{2}\left|V_{c b}\right|^{2} \eta_{\mathrm{EW}}^{2} m_{B}^{5}}{48 \pi^{3}}\left(w^{2}-1\right)^{1 / 2}(w+1)^{2} r_{D^{*}}^{3}\left(1-r_{D^{*}}\right)^{2} \\
& \times\left[1+\frac{4 w}{w+1} \frac{1-2 w r_{D^{*}}+r_{D^{*}}^{2}}{\left(1-r_{D^{*}}\right)^{2}}\right] \mathcal{F}(w)^{2}
\end{aligned}
$$

with $r_{D^{(*)}}=m_{D^{(*)}} / m_{B}$ and $w=v \cdot v^{\prime}=\frac{m_{B}^{2}+m_{D^{(*)}}^{2}-q^{2}}{2 m_{B} m_{D^{(*)}}}$

$$
\begin{aligned}
\mathcal{G}(w)=h_{+} & -\frac{1-r_{D}}{1+r_{D}} h_{-}, \\
\mathcal{F}(w)^{2}=h_{A_{1}}^{2} & \left\{2\left(1-2 w r_{D^{*}}+r_{D^{*}}^{2}\right)\left(1+R_{1} \frac{w-1}{w+1}\right)+\left[\left(1-r_{D^{*}}\right)+(w-1)\left(1-R_{2}\right)\right]^{2}\right\} \\
& \times\left[\left(1-r_{D^{*}}\right)^{2}+\frac{4 w}{w+1}\left(1-2 w r_{D^{*}}+r_{D^{*}}^{2}\right)\right]^{-1},
\end{aligned}
$$

$$
R_{1}(w)=\frac{h_{V}}{h_{A_{1}}}, \quad R_{2}(w)=\frac{h_{A_{3}}+r_{D^{*}} h_{A_{2}}}{h_{A_{1}}} . \quad \begin{aligned}
& R_{i} \text { are angular distributions } \rightarrow \\
& \text { can be accessed experimentally }
\end{aligned}
$$

## $B \rightarrow D, D *$ DECAYS: NOTATION

in case of $\tau$ decays one extra form factor in SM (more with NP)
$\rightarrow$ define other ratios:

$$
R_{3}(w)=\frac{h_{A_{3}}-r_{D^{*}} h_{A_{2}}}{h_{A_{1}}}, \quad R_{0}(w)=\frac{h_{A_{1}}(w+1)-h_{A_{3}}\left(w-r_{D^{*}}\right)-h_{A_{2}}\left(1-w r_{D^{*}}\right)}{\left(1+r_{D^{*}}\right) h_{A_{1}}}
$$

enter in rate suppressed by factors of $m_{\tau}{ }^{2} / m_{B}{ }^{2}$

## Determination of Form Factors?

## THEORY INPUTS

## BGL: UNITARITY CONSTRAINTS

> Boyd, Grinstein, Lebed ('95) (BGL): relate FFs to two point functions via dispersion relations, crossing symmetry, quarkhadron duality:

(Blanschke factors)
$>$ unitarity $\rightarrow$ constraints on $a_{n}$, e.g. for single channel: $\quad \sum_{n=0}^{\infty}\left|a_{n}\right|^{2} \leq 1$.
Can be used directly to fit spectra

## HQET (+ UNITARITY)

- FFs are related by heavy quark symmetry (HQS)
$>$ HQET $\rightarrow$ organize expansion in powers of $a_{s}, \Lambda / m_{b}, \Lambda / m_{c}$
> relations among form factors
> can be used to relate form factors measurements in e, $\mu$ to additional ffs in $\tau$
> At LO: everything proportional to Isgur-Wise function $\xi$ or 0
$\rightarrow$ At $\mathrm{O}\left(\Lambda / \mathrm{m}_{\mathrm{b}}, \Lambda / \mathrm{m}_{\mathrm{c}}\right)$ : subleading IW functions: $\chi_{2}, \chi_{3}, \eta$
$>$ HQET + z-parameterization from unitarity:
> Compute FF at $\mathrm{O}\left(\mathrm{a}_{\mathrm{s}}, \Lambda / \mathrm{m}_{\mathrm{b}}, \Lambda / \mathrm{m}_{\mathrm{c}}\right)$
> Taylor expand IW functions

$$
\frac{\mathcal{G}(w)}{\mathcal{G}\left(w_{0}\right)} \simeq 1-8 a^{2} \rho_{*}^{2} z_{*}+\left(V_{21} \rho_{*}^{2}-V_{20}\right) z_{*}^{2} .
$$

$$
\hat{\chi}_{2}(w) \simeq \hat{\chi}_{2}(1)+\hat{\chi}_{2}^{\prime}(1)(w-1), \quad \hat{\chi}_{3}(w) \simeq \hat{\chi}_{3}^{\prime}(1)(w-1), \quad \eta(w) \simeq \eta(1)+\eta^{\prime}(1)(w-1),
$$

Fit input shapes to 6 parameters (3 slopes, 2 intercepts $+\rho *$ )

## CLN: UNITARITY + HQET + QCD SUM RULES

> Caprini Lellouch Neubert (CLN ‘98): Use NLO HQET and further constraints from QCD sum rules:
> subleading IW functions determined
> Only two parameters: normalization and $\rho^{*}$
> Uncertainties small: $<2 \%$ (?) $\rightarrow$ mostly neglected in experimental analyses (e.g. fix slopes to CLN prediction and float intercepts, ...)

## EXPERIMENTAL \& LATTICE INPUTS

## EXPERIMENTAL \& LATIICE ACCESS TO SPECTRA

> B $\rightarrow$ Dlv: Belle, Lattice:
1510.03657


Lattice: good at small recoil, Exp: good at large recoil

## EXPERIMENTAL ACCESS TO SPECTRA

> $\mathrm{B} \rightarrow \mathrm{D}^{*} \mathrm{lv}$ (2017, Belle):


No lattice results yet for spectra (only preliminary info for finite lattice spacing)

## NORMALIZATION AT ZERO RECOIL

> Lattice measurements at zero recoil:

$$
\mathcal{G}(1)_{\mathrm{LQCD}}=1.054(8), \quad \mathcal{F}(1)_{\mathrm{LQCD}}=0.906(13),
$$

FIT RESULTS

## BGL FIT IN B $\rightarrow$ D* $^{*}$ USING BELLE SPECTRA




FNAL/MILC D $=$ Lattice $B \rightarrow D+H Q S$
> HQET predict $\mathrm{R}_{1,2}=1+\mathrm{O}\left(\Lambda / \mathrm{m}_{\mathrm{b}, \mathrm{c}}, \mathrm{a}_{\mathrm{s}}\right)$, slopes small

- BGL seem to suggest large HQS violations, not seen from lattice
> Using lattice to extract $\mathrm{V}_{\mathrm{cb}}$ :
(if $\sim 100 \%$ correl $\rightarrow$ more than $5 \sigma$ discrepancy)

$$
\begin{array}{ll}
\left|V_{c b}\right|_{\mathrm{CLN}}=(38.2 \pm 1.5) \times 10^{-3}, & {[1]} \\
\left|V_{c b}\right|_{\mathrm{BGL}}=\left(41.7_{-2.1}^{+2.0}\right) \times 10^{-3}, & {[3]} \\
\left|V_{c b}\right|_{\mathrm{BGL}}=\left(41.9_{-1.9}^{+2.0}\right) \times 10^{-3}, & {[4]}
\end{array}
$$

> Some tension between data + lattice + HQS

## OTHER FIT COMBINATIONS

Use NLO HQET and:

| fix norm. to | Fit | QCDSR | Lattice QCD |  |  | Belle Data |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | $\mathcal{F}(1)$ | $f_{+, 0}(1)$ | $f_{+, 0}(w>1)$ |  |
| lattice zero- | $\mathrm{L}_{w=1}$ | - | $\checkmark$ | $\checkmark$ | - | $\checkmark$ |
| recoil results | $\mathrm{L}_{w=1}+\mathrm{SR}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | - | $\checkmark$ |
| float norm. independently | NoL | - | - | - | - | $\checkmark$ |
|  | NoL+SR | $\checkmark$ | - | - | - | $\checkmark$ |
| fit $\xi$ to lattice $B \rightarrow D$, use lattice for $B \rightarrow D^{*}$ norm | $\mathrm{L}_{w \geq 1}$ | - | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
|  | $\mathrm{L}_{w \geq 1}+\mathrm{SR}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
|  | th: $\mathrm{L}_{w \geq 1}+\mathrm{SR}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | - |

" + SR": use QCD sum rules for priors on subleading IW functions

$$
\begin{gathered}
\hat{\chi}_{2}^{\text {ren }}(1)=-0.06 \pm 0.02, \quad \hat{\chi}_{2}^{\prime \text { ren }}(1)=0 \pm 0.02, \quad \hat{\chi}_{3}^{\prime \text { ren }}(1)=0.04 \pm 0.02, \\
\eta(1)=0.62 \pm 0.2, \quad \eta^{\prime}(1)=0 \pm 0.2 .
\end{gathered}
$$

## FIT RESULTS

|  | $\mathrm{L}_{w=1}$ | $\mathrm{~L}_{w=1}+\mathrm{SR}$ | NoL | $\mathrm{NoL}+\mathrm{SR}$ | $\mathrm{L}_{w \geq 1}$ | $\mathrm{~L}_{w \geq 1}+\mathrm{SR}$ | th: $\mathrm{L}_{w \geq 1}+\mathrm{SR}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\chi^{2}$ | 40.2 | 44.0 | 38.7 | 43.1 | 49.0 | 53.8 | 7.4 |
| dof | 44 | 48 | 43 | 47 | 48 | 52 | 4 |
| $\left\|V_{c b}\right\| \times 10^{3}$ | $38.8 \pm 1.2$ | $38.5 \pm 1.1$ | - | - | $39.1 \pm 1.1$ | $39.3 \pm 1.0$ | - |
| $\mathcal{G}(1)$ | $1.055 \pm 0.008$ | $1.056 \pm 0.008$ | - | - | $1.060 \pm 0.008$ | $1.061 \pm 0.007$ | $1.052 \pm 0.008$ |
| $\mathcal{F}(1)$ | $0.904 \pm 0.012$ | $0.901 \pm 0.011$ | - | - | $0.898 \pm 0.012$ | $0.895 \pm 0.011$ | $0.906 \pm 0.013$ |
| $\bar{\rho}_{*}^{2}$ | $1.17 \pm 0.12$ | $1.19 \pm 0.07$ | $1.06 \pm 0.15$ | $1.19 \pm 0.08$ | $1.33 \pm 0.11$ | $1.24 \pm 0.06$ | $1.24 \pm 0.08$ |
| $\hat{\chi}_{2}(1)$ | $-0.26 \pm 0.26$ | $-0.07 \pm 0.02$ | $0.36 \pm 0.62$ | $-0.06 \pm 0.02$ | $0.13 \pm 0.22$ | $-0.06 \pm 0.02$ | $-0.06 \pm 0.02$ |
| $\hat{\chi}_{2}^{\prime}(1)$ | $0.21 \pm 0.38$ | $-0.00 \pm 0.02$ | $0.14 \pm 0.39$ | $-0.00 \pm 0.02$ | $-0.36 \pm 0.28$ | $-0.00 \pm 0.02$ | $-0.00 \pm 0.02$ |
| $\hat{\chi}_{3}^{\prime}(1)$ | $0.02 \pm 0.07$ | $0.05 \pm 0.02$ | $0.18 \pm 0.19$ | $0.04 \pm 0.02$ | $0.09 \pm 0.07$ | $0.05 \pm 0.02$ | $0.04 \pm 0.02$ |
| $\eta(1)$ | $0.30 \pm 0.04$ | $0.30 \pm 0.03$ | $-0.56 \pm 0.80$ | $0.35 \pm 0.14$ | $0.30 \pm 0.04$ | $0.30 \pm 0.03$ | $0.31 \pm 0.04$ |
| $\eta^{\prime}(1)$ | 0 (fixed) | $-0.12 \pm 0.16$ | 0 (fixed) | $-0.11 \pm 0.18$ | $0($ fixed $)$ | $-0.05 \pm 0.09$ | $0.05 \pm 0.10$ |
| $m_{b}^{1 S}[\mathrm{GeV}]$ | $4.70 \pm 0.05$ | $4.70 \pm 0.05$ | $4.71 \pm 0.05$ | $4.70 \pm 0.05$ | $4.71 \pm 0.05$ | $4.71 \pm 0.05$ | $4.71 \pm 0.05$ |
| $\delta m_{b c}[\mathrm{GeV}]$ | $3.40 \pm 0.02$ | $3.40 \pm 0.02$ | $3.40 \pm 0.02$ | $3.40 \pm 0.02$ | $3.40 \pm 0.02$ | $3.40 \pm 0.02$ | $3.40 \pm 0.02$ |

no signs of strong tensions, $V_{c b}$ is still "low", data prefers a lower $\eta$ than QCDSR input

$$
L_{w>=1}+S R
$$

spectra:






## FITS \& R ${ }_{1,2}$




## R(D) AND R(D*)




## R(D) AND R(D*)

| Scenario | $R(D)$ | $R\left(D^{*}\right)$ | Correlation |
| :--- | :---: | :---: | :---: |
| $\mathrm{L}_{w=1}$ | $0.292 \pm 0.005$ | $0.255 \pm 0.005$ | $41 \%$ |
| $\mathrm{~L}_{w=1}+\mathrm{SR}$ | $0.291 \pm 0.005$ | $0.255 \pm 0.003$ | $57 \%$ |
| NoL | $0.273 \pm 0.016$ | $0.250 \pm 0.006$ | $49 \%$ |
| $\mathrm{NoL}+\mathrm{SR}$ | $0.295 \pm 0.007$ | $0.255 \pm 0.004$ | $43 \%$ |
| $\mathrm{~L}_{w \geq 1}$ | $0.298 \pm 0.003$ | $0.261 \pm 0.004$ | $19 \%$ |
| $\mathrm{~L}_{w \geq 1}+\mathrm{SR}$ | $\mathbf{0 . 2 9 9} \pm \mathbf{0 . 0 0 3}$ | $\mathbf{0 . 2 5 7} \pm \mathbf{0 . 0 0 3}$ | $\mathbf{4 4 \%}$ |
| th:L $\mathrm{L}_{w \geq 1}+\mathrm{SR}$ | $0.306 \pm 0.005$ | $0.256 \pm 0.004$ | $33 \%$ |
| Data [9] | $0.403 \pm 0.047$ | $0.310 \pm 0.017$ | $-23 \%$ |
| Refs. [48, 52, 54] | $0.300 \pm 0.008$ | - | - |
| Ref. [53] | $0.299 \pm 0.003$ | - | - |
| Ref. [34] | - | $0.252 \pm 0.003$ | - |

- Reduced uncert on SM predictions
- Consistency between different fits
- Discrepancy with data still present and sizable


## NLO HQET FOR SM+NP

> NLO HQET calculation also for form factors entering BSM contributions




## CONCLUSIONS

> Experimental data in e, $\mu$ \& Lattice results are improving determination of $\mathrm{B} \rightarrow \mathrm{D}\left(^{*}\right)$ form factors
> Apparent "tension" in current inputs between HQS, lattice and Belle B $\rightarrow \mathrm{D}^{*}$ distributions (can't self-consistently use lattice+BGL to extract $\mathrm{V}_{\mathrm{cb}}$ )
> Future lattice $\mathrm{B} \rightarrow \mathrm{D}^{*}$ spectra and Belle II data (and non-unfolded BGL Belle fit?) will have something to say on this
> Updated $\mathrm{R}(\mathrm{D}), \mathrm{R}\left(\mathrm{D}^{*}\right)$ predictions still show large discrepancy with measurements
> Updated BSM predictions for $\mathrm{R}(\mathrm{D}), \mathrm{R}\left(\mathrm{D}^{*}\right)$
> Results included in Hammer package

BACKUP

## FORM FACTOR DEFINITIONS

> $\mathrm{B} \rightarrow \mathrm{D}$ :

$$
\begin{aligned}
\langle D| \bar{c} b|\bar{B}\rangle & =\sqrt{m_{B} m_{D}} h_{S}(w+1), \\
\langle D| \bar{c} \gamma^{5} b|\bar{B}\rangle & =\langle D| \bar{c} \gamma^{\mu} \gamma^{5} b|\bar{B}\rangle=0, \\
\langle D| \bar{c} \gamma^{\mu} b|\bar{B}\rangle & =\sqrt{m_{B} m_{D}}\left[h_{+}\left(v+v^{\prime}\right)^{\mu}+h_{-}\left(v-v^{\prime}\right)^{\mu}\right], \\
\langle D| \bar{c} \sigma^{\mu \nu} b|\bar{B}\rangle & =i \sqrt{m_{B} m_{D}}\left[h_{T}\left(v^{\prime \mu} v^{\nu}-v^{\prime \nu} v^{\mu}\right)\right],
\end{aligned}
$$

> $\mathrm{B} \rightarrow \mathrm{D}^{*}$ :

$$
\begin{aligned}
\left\langle D^{*}\right| \bar{c} b|\bar{B}\rangle & =0, \\
\left\langle D^{*}\right| \bar{c} \gamma^{5} b|\bar{B}\rangle & =-\sqrt{m_{B} m_{D^{*}}} h_{P}\left(\epsilon^{*} \cdot v\right), \\
\left\langle D^{*}\right| \bar{c} \gamma^{\mu} b|\bar{B}\rangle & =i \sqrt{m_{B} m_{D^{*}}} h_{V} \varepsilon^{\mu \nu \alpha \beta} \epsilon_{\nu}^{*} v_{\alpha}^{\prime} v_{\beta}, \\
\left\langle D^{*}\right| \bar{c} \gamma^{\mu} \gamma^{5} b|\bar{B}\rangle & =\sqrt{m_{B} m_{D^{*}}}\left[h_{A_{1}}(w+1) \epsilon^{* \mu}-h_{A_{2}}\left(\epsilon^{*} \cdot v\right) v^{\mu}-h_{A_{3}}\left(\epsilon^{*} \cdot v\right) v^{\prime \mu}\right], \\
\left\langle D^{*}\right| \bar{c} \sigma^{\mu \nu} b|\bar{B}\rangle & =-\sqrt{m_{B} m_{D^{*}}} \varepsilon^{\mu \nu \alpha \beta}\left[h_{T_{1}} \epsilon_{\alpha}^{*}\left(v+v^{\prime}\right)_{\beta}+h_{T_{2}} \epsilon_{\alpha}^{*}\left(v-v^{\prime}\right)_{\beta}+h_{T_{3}}\left(\epsilon^{*} \cdot v\right) v_{\alpha} v_{\beta}^{\prime}\right] .
\end{aligned}
$$

## NLO HQET FF EXPRESSIONS

$$
\begin{aligned}
& \hat{h}_{+}=1+\hat{\alpha}_{s}\left[C_{V_{1}}+\frac{w+1}{2}\left(C_{V_{2}}+C_{V_{3}}\right)\right]+\left(\varepsilon_{c}+\varepsilon_{b}\right) \hat{L}_{1}, \\
& \hat{h}_{-}=\hat{\alpha}_{s} \frac{w+1}{2}\left(C_{V_{2}}-C_{V_{3}}\right)+\left(\varepsilon_{c}-\varepsilon_{b}\right) \hat{L}_{4}, \\
& \hat{h}_{S}=1+\hat{\alpha}_{s} C_{S}+\left(\varepsilon_{c}+\varepsilon_{b}\right)\left(\hat{L}_{1}-\hat{L}_{4} \frac{w-1}{w+1}\right), \\
& \hat{h}_{T}=1+\hat{\alpha}_{s}\left(C_{T_{1}}-C_{T_{2}}+C_{T_{3}}\right)+\left(\varepsilon_{c}+\varepsilon_{b}\right)\left(\hat{L}_{1}-\hat{L}_{4}\right) . \\
& \hat{h}_{V}
\end{aligned}=1+\hat{\alpha}_{s} C_{V_{1}}+\varepsilon_{c}\left(\hat{L}_{2}-\hat{L}_{5}\right)+\varepsilon_{b}\left(\hat{L}_{1}-\hat{L}_{4}\right), \quad \text {, } \begin{aligned}
\hat{h}_{A_{1}} & =1+\hat{\alpha}_{s} C_{A_{1}}+\varepsilon_{c}\left(\hat{L}_{2}-\hat{L}_{5} \frac{w-1}{w+1}\right)+\varepsilon_{b}\left(\hat{L}_{1}-\hat{L}_{4} \frac{w-1}{w+1}\right), \\
\hat{h}_{A_{2}} & =\hat{\alpha}_{s} C_{A_{2}}+\varepsilon_{c}\left(\hat{L}_{3}+\hat{L}_{6}\right), \\
\hat{h}_{A_{3}} & =1+\hat{\alpha}_{s}\left(C_{A_{1}}+C_{A_{3}}\right)+\varepsilon_{c}\left(\hat{L}_{2}-\hat{L}_{3}+\hat{L}_{6}-\hat{L}_{5}\right)+\varepsilon_{b}\left(\hat{L}_{1}-\hat{L}_{4}\right), \\
\hat{h}_{P} & =1+\hat{\alpha}_{s} C_{P}+\varepsilon_{c}\left[\hat{L}_{2}+\hat{L}_{3}(w-1)+\hat{L}_{5}-\hat{L}_{6}(w+1)\right]+\varepsilon_{b}\left(\hat{L}_{1}-\hat{L}_{4}\right) \\
\hat{h}_{T_{1}} & =1+\hat{\alpha}_{s}\left[C_{T_{1}}+\frac{w-1}{2}\left(C_{T_{2}}-C_{T_{3}}\right)\right]+\varepsilon_{c} \hat{L}_{2}+\varepsilon_{b} \hat{L}_{1}, \\
\hat{h}_{T_{2}} & =\hat{\alpha}_{s} \frac{w+1}{2}\left(C_{T_{2}}+C_{T_{3}}\right)+\varepsilon_{c} \hat{L}_{5}-\varepsilon_{b} \hat{L}_{4}, \\
\hat{h}_{T_{3}} & =\hat{\alpha}_{s} C_{T_{2}}+\varepsilon_{c}\left(\hat{L}_{6}-\hat{L}_{3}\right) .
\end{aligned}
$$

## HOET RI EXPRESSIONS

$$
\begin{aligned}
& R_{1}(1) \simeq 1.34-0.12 \eta(1) \\
& R_{2}(1) \simeq 0.98-0.42 \eta(1)-0.54 \hat{\chi}_{2}(1) \\
& R_{1}^{\prime}(1) \simeq-0.15+0.06 \eta(1)-0.12 \eta^{\prime}(1) \\
& R_{2}^{\prime}(1) \simeq 0.01-0.54 \hat{\chi}_{2}^{\prime}(1)+0.21 \eta(1)-0.42 \eta^{\prime}(1) \\
& R_{3}(1) \simeq 1.19-0.26 \eta(1)-1.20 \hat{\chi}_{2}(1) \\
& R_{0}(1) \simeq 1.09+0.25 \eta(1) \\
& R_{3}^{\prime}(1) \simeq-0.08-1.20 \hat{\chi}_{2}^{\prime}(1)+0.13 \eta(1)-0.26 \eta^{\prime}(1) \\
& R_{0}^{\prime}(1) \simeq-0.18+0.87 \hat{\chi}_{2}(1)+0.06 \eta(1)+0.25 \eta^{\prime}(1)
\end{aligned}
$$

