the QCD phase boundary and the production of rare, loosely bound objects

- the hadron resonance gas and (u,d,s) hadron production
- Dashen-Ma-Bernstein taken at face value
- experimental determination of the QCD phase boundary
- loosely bound objects
- coalescence vs thermal production
- outlook

pbm
international workshop on
heavy flavor production in high energy nuclear collisions
CCNU, Wuhan
Oct. 8, 2018

phenomenology results obtained in collaboration with
Anton Andronic, Krzysztof Redlich, and Johanna Stachel

arXiv:1710.09425,

Nature 561(2018) no.7723, 321-330

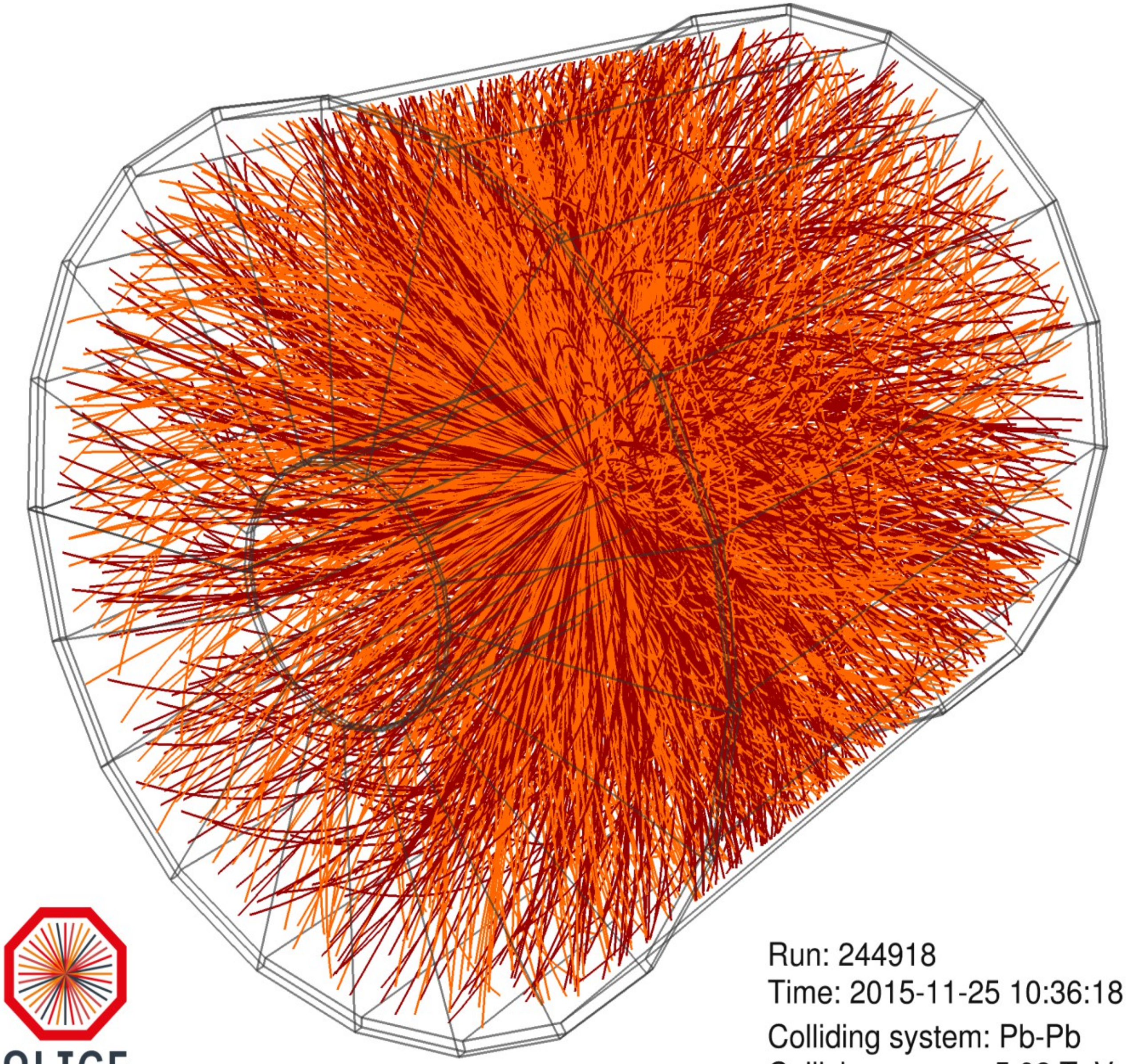
first PbPb collisions at LHC at √s = 5.02 A TeV

Run1: 3 data taking campaigns
pp, pPb, Pb—Pb
> 150 publications

Run2 with 13 TeV pp
Pb—Pb run 5 TeV/u
p-Pb Run at 5 and 8 TeV

Now running with 13 TeV pp

Nov. 2018: PbPb 5 TeV/u

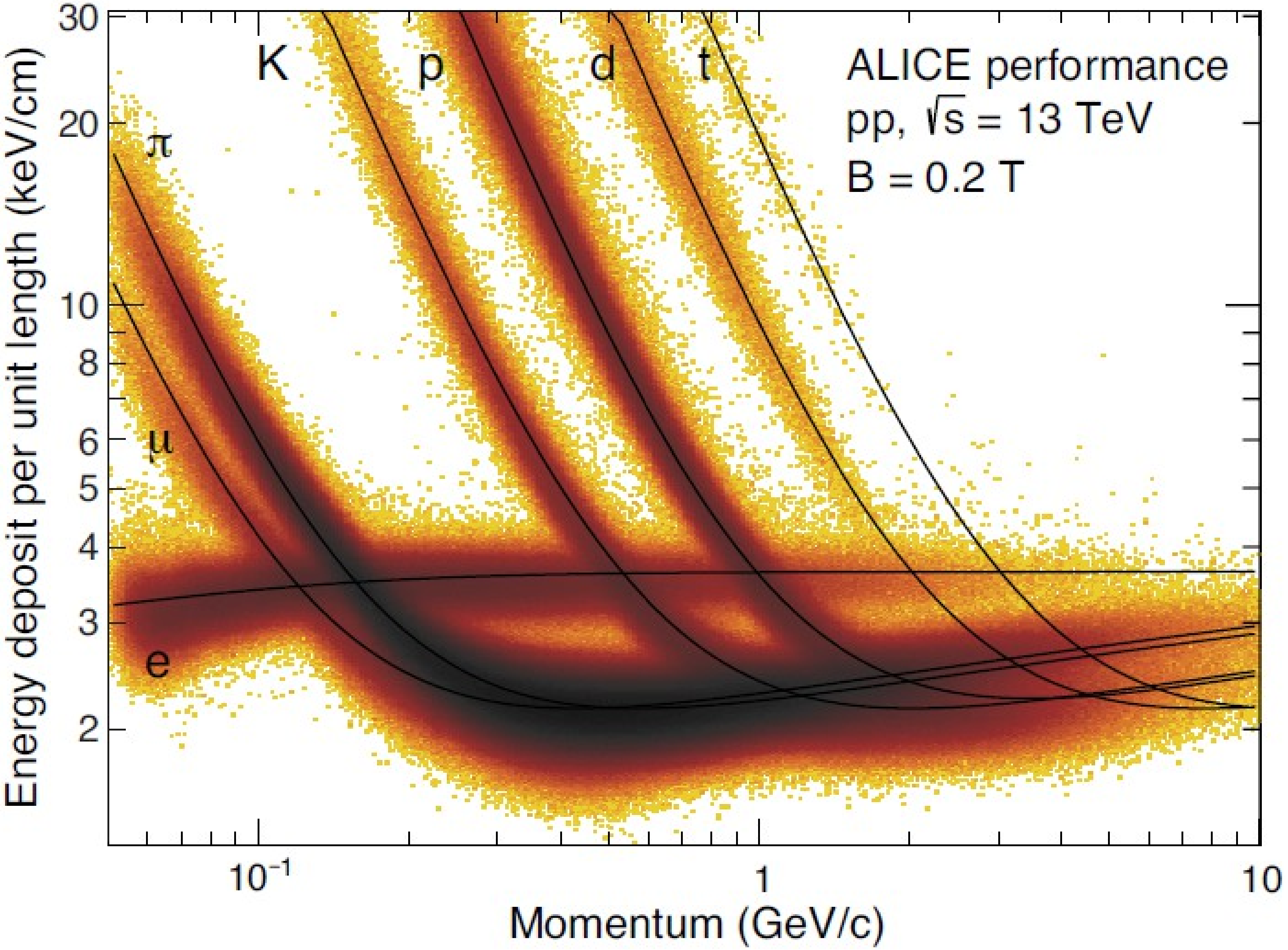


Run: 244918
Time: 2015-11-25 10:36:18
Colliding system: Pb-Pb
Collision energy: 5.02 TeV

ALICE

particle identification with the ALICE TPC

from 50 MeV to 50 GeV



now PDG standard

hadron production and the QCD phase boundary

part 1: the hadron resonance gas

duality between hadrons and quarks/gluons (I)

Z: full QCD partition function

all thermodynamic quantities derive from QCD partition functions

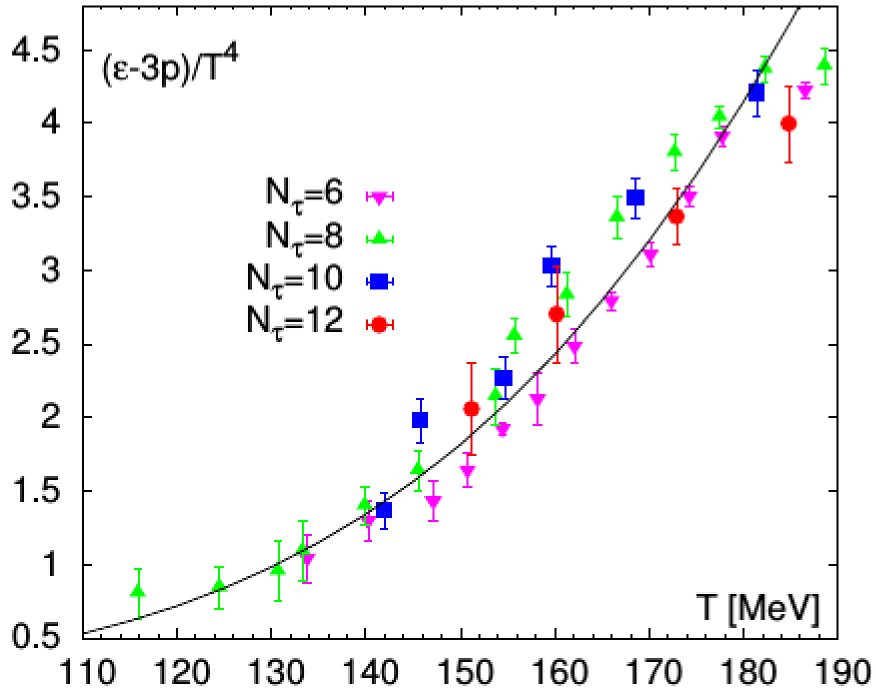
for the pressure we get:

$$\frac{p}{T^4} \equiv \frac{1}{VT^3} \ln Z(V,T,\mu)$$

comparison of trace anomaly from LQCD
Phys.Rev. D90 (2014) 094503
HOTQCD coll.

with hadron resonance gas prediction
(solid line)

LQCD: full dynamical quarks with realistic
pion mass

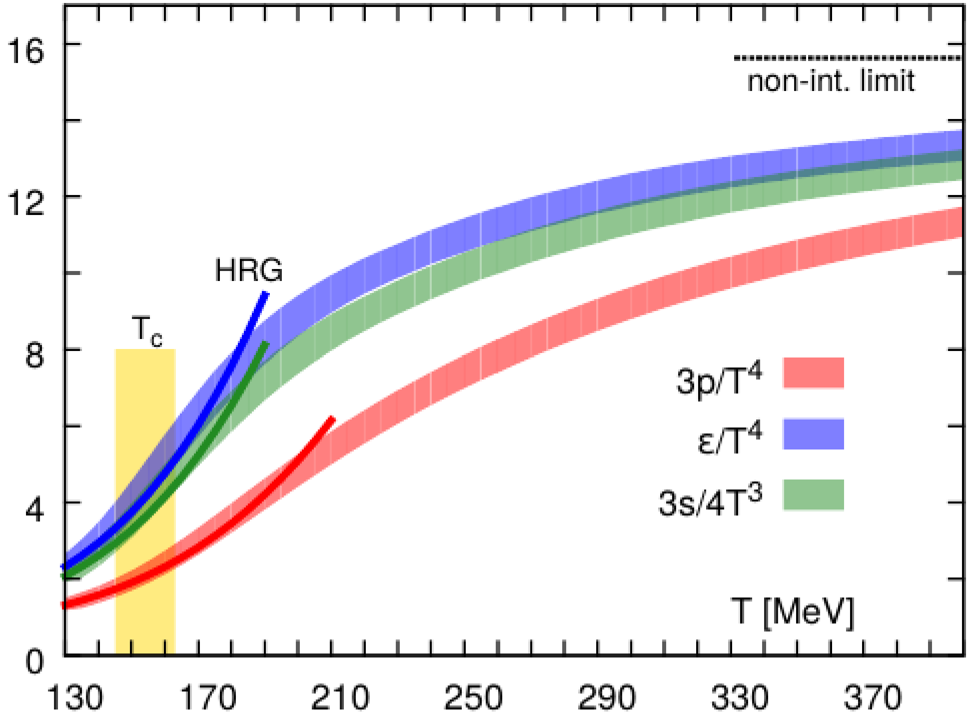


duality between hadrons and quarks/gluons (II)

comparison of equation of state from LQCD
Phys.Rev. D90 (2014) 094503
HOTQCD coll.

with hadron resonance gas predictions (colored lines)

essentially the same results also from Wuppertal-Budapest coll.
Phys.Lett. B730 (2014) 99-104



pseudo-critical temperature

T_c = 156 ± 1.5 MeV, arXiv:1807.05607

very new and improved

$\varepsilon_{crit} = 340\ \mathrm{MeV/fm^3}$

$\varepsilon_{nucl} = 450\ \mathrm{MeV/fm^3}$

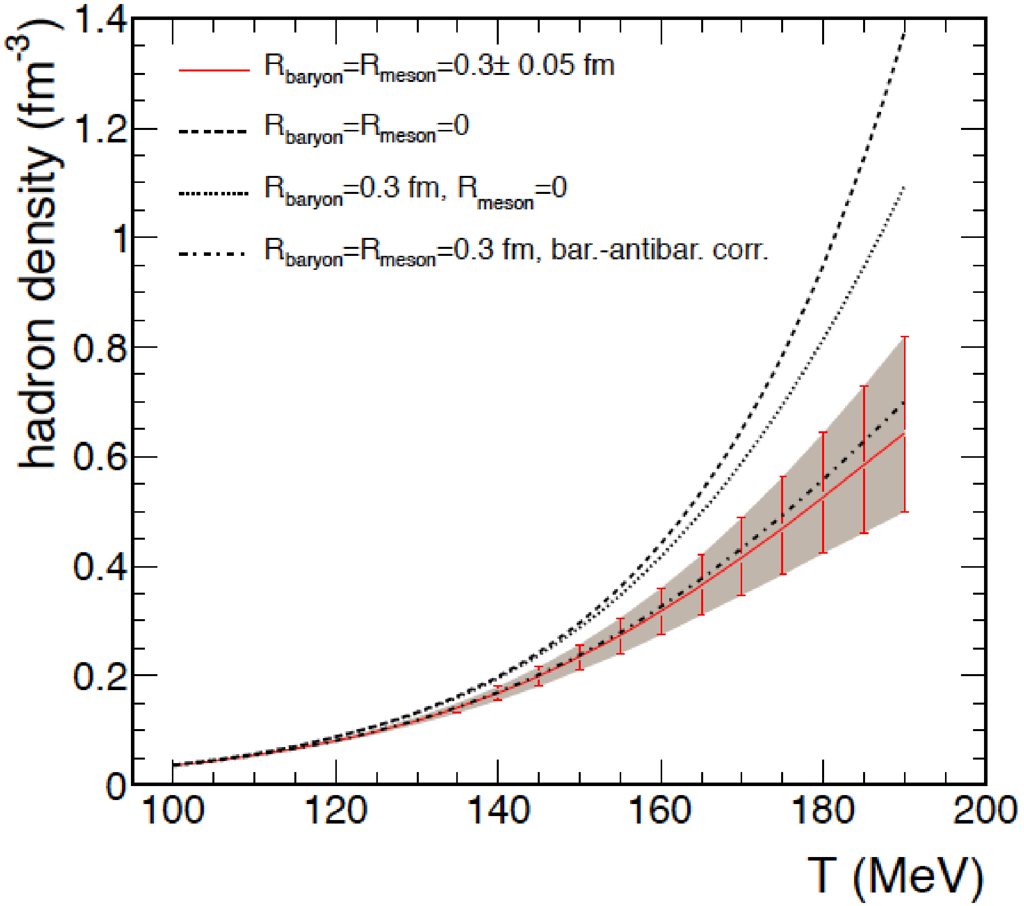
duality between hadrons and quarks/gluons (III)

in the dilute limit T < 165 MeV:

$$\ln Z(T,V,\mu) \approx \sum_{i\in\ mesons} \ln \mathcal{Z}^{M}_{M_i}(T,V,\mu_Q,\mu_S) + \sum_{i\in\ baryons} \ln \mathcal{Z}^{B}_{M_i}(T,V,\mu_b,\mu_Q,\mu_S)$$

where the partition function of the hadron resonance model is expressed in mesonic and baryonic components. The chemical potential μ reflects then the baryonic, charge, and strangeness components $\mu = (\mu_b, \mu_Q, \mu_S)$.

hadron resonance gas and interactions a la 'Dashen-Ma-Bernstein'



note: using S-matrix approach takes care of all aspects of resonance widths as well as repulsive and attractive interactions

no need anymore for theoretically uncontrolled excluded volume approach

within the 2-body framework this is parameter free

for how to do this, see below

for T > 170 MeV the HRG approach breaks down!!!

multi-body collisions become important

previously, we used excluded volume approach only to control the particle density

now not used anymore

hadron production and the QCD phase boundary

part 1: the hadron resonance gas

thermal model of particle production and QCD

partition function Z(T,V) contains sum over the full hadronic mass spectrum and is fully calculable in QCD

low density approach
no interactions

for each particle i, the statistical operator is:

$$\ln Z_i = \frac{V g_i}{2\pi^2} \int_0^\infty \pm p^2 \mathrm{d}p \ln[1 \pm \exp(-(E_i - \mu_i)/T)]$$

particle densities are then calculated according to:

$$n_i = N_i/V = -\frac{T}{V}\frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 \mathrm{d}p}{\exp[(E_i - \mu_i)/T] \pm 1}$$

from analysis of all available nuclear collision data we now know the energy dependence of the parameters T, mu_b, and V over an energy range from threshold to LHC energy and can confidently extrapolate to even higher energies

in practice, we use the full experimental hadronic mass spectrum from the PDG compilation (vacuum masses) to compute the 'primordial yield'

comparison with measured hadron yields needs evaluation of all strong decays

implementation

$$n_i = N_i/V = -\frac{T}{V}\frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2}\int_0^\infty \frac{p^2 \mathrm{d}p}{\exp[(E_i-\mu_i)/T] \pm 1}$$

Latest PDG hadron mass spectrum ...quasi-complete up to m=2 GeV;
our code: 555 species (including light nuclei, charm and bottom hadrons)

for resonances, the width is considered in calculations

Minimize: $\chi^2 = \sum_i \frac{(N_i^{exp} - N_i^{therm})^2}{\sigma_i^2}$

N_i hadron yield, σ_i experimental uncertainty (stat.+syst.)

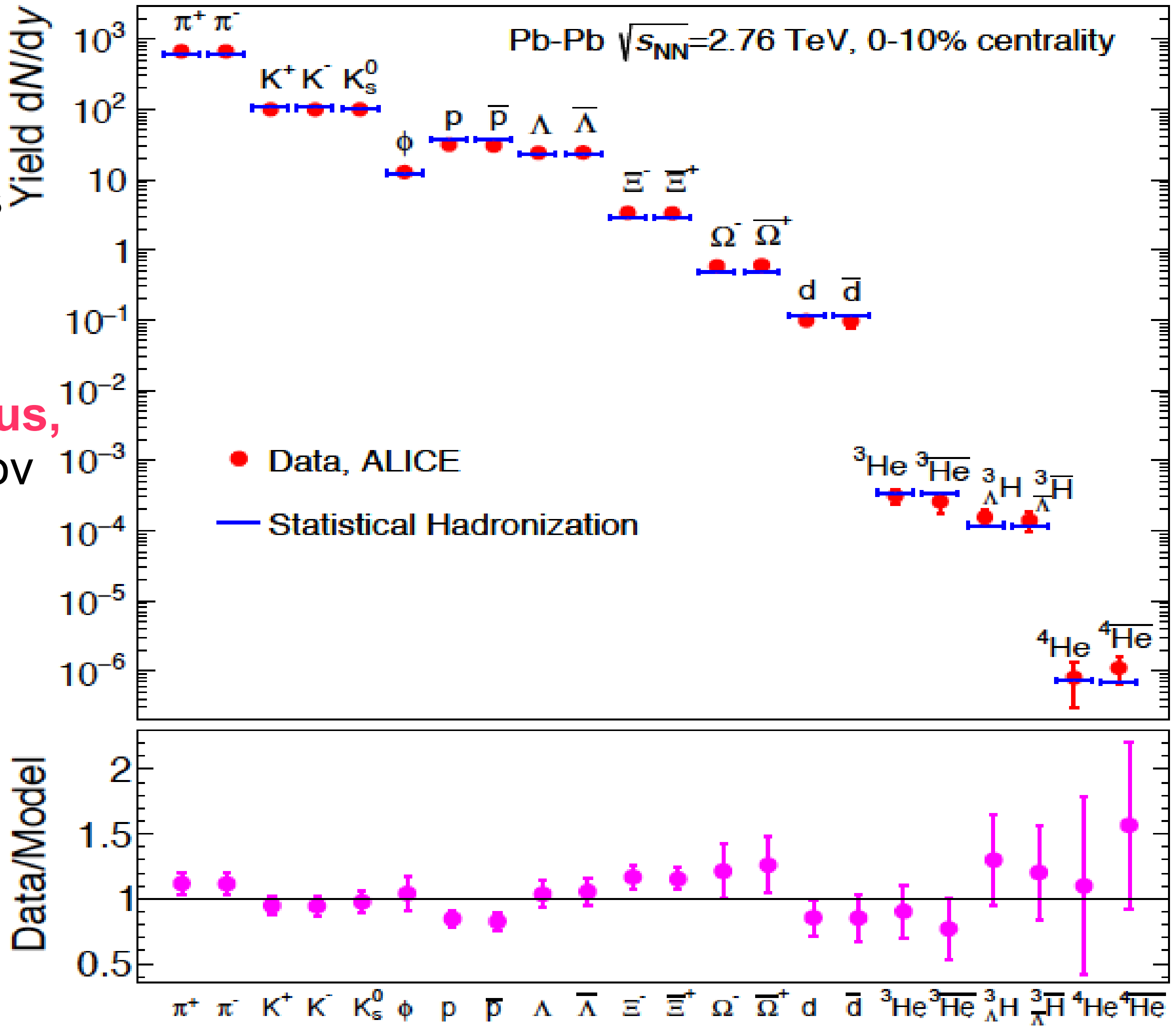
$\Rightarrow (T, \mu_B, V)$

canonical treatment whenever needed (small abundances)

Oct. 2017 update: excellent description of ALICE@LHC data

fit includes loosely bound systems such as deuteron and hypertriton
hypertriton is bound-state of (Λ,p,n),
Λ separation energy about 130 keV
size about 10 fm, the **ultimate halo nucleus,**
produced at T=156 MeV. close to an Efimov state

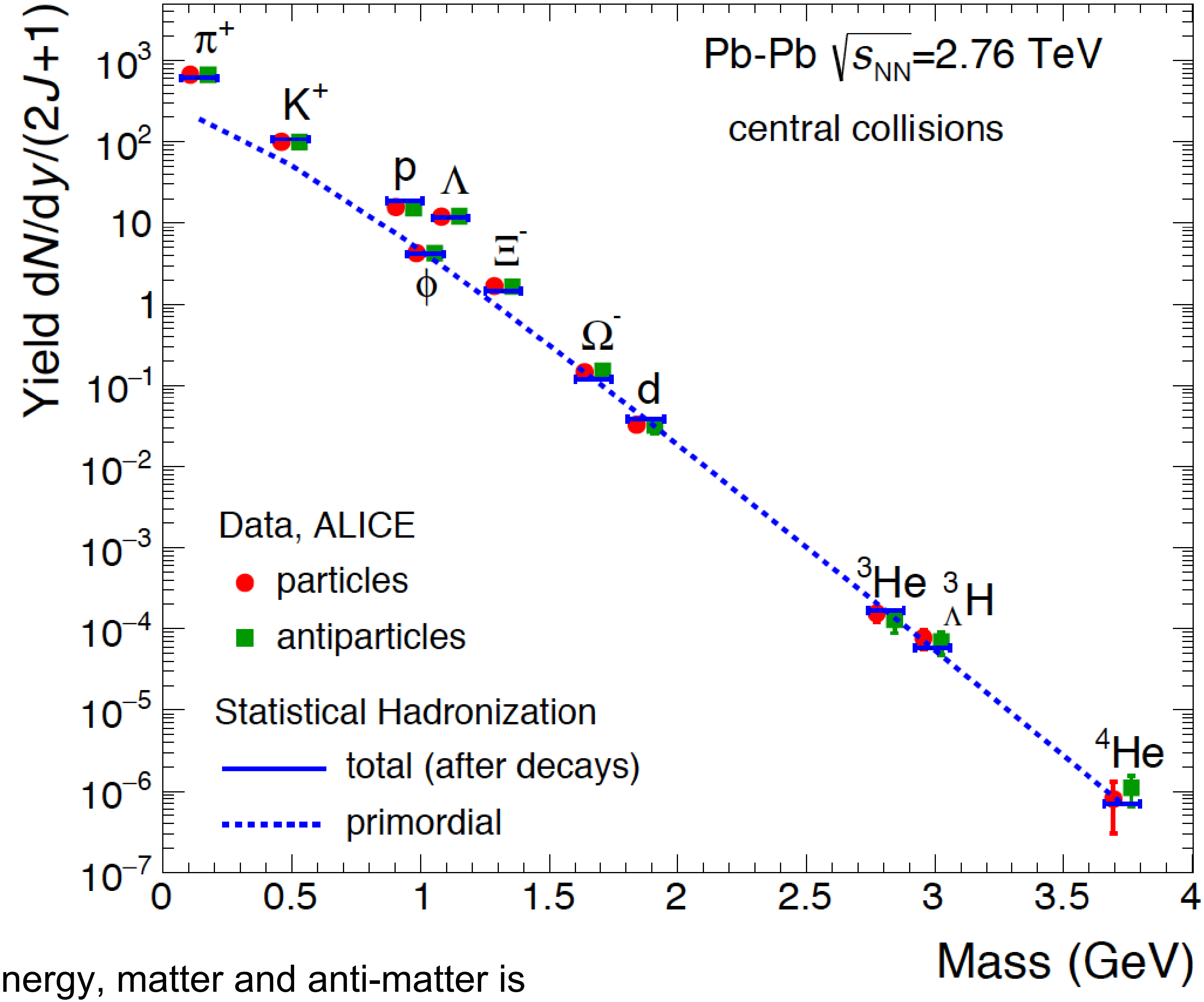
proton discrepancy about 2.8 sigma



Andronic, pbm, Redlich, Stachel, arXiv:1710.09425, Nature (Sep. 20, 2018)

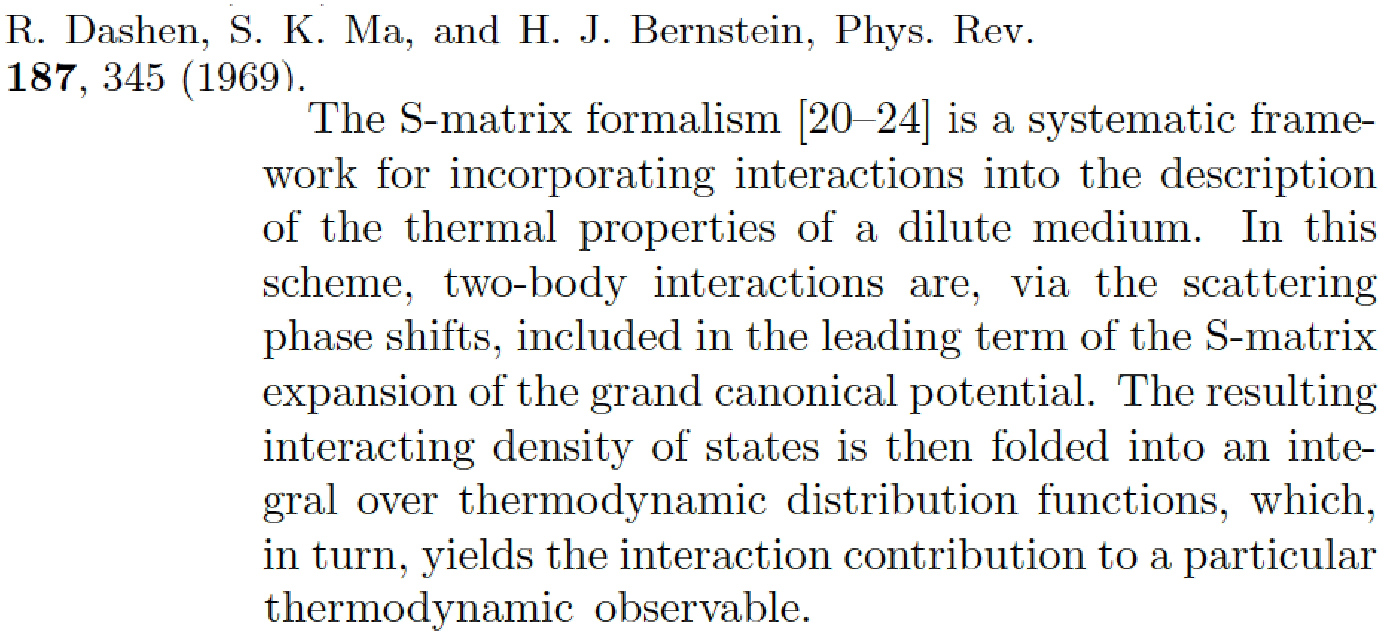
J. Stachel, A. Andronic, P. Braun-Munzinger and K. Redlich, Confronting LHC data with the statistical hadronization model, J.Phys.Conf.Ser.**509** (2014) 012019, arXiv:1311.4662 [nucl-th].

at LHC energy, production of (u,d,s) hadrons is governed
by mass and quantum numbers only
quark content does not matter



at LHC energy, matter and anti-matter is
produced with equal yields

the proton anomaly and the Dashen, Ma, Bernstein S-matrix approach



R. Dashen, S. K. Ma, and H. J. Bernstein, Phys. Rev. **187**, 345 (1969).

The S-matrix formalism [20–24] is a systematic framework for incorporating interactions into the description of the thermal properties of a dilute medium. In this scheme, two-body interactions are, via the scattering phase shifts, included in the leading term of the S-matrix expansion of the grand canonical potential. The resulting interacting density of states is then folded into an integral over thermodynamic distribution functions, which, in turn, yields the interaction contribution to a particular thermodynamic observable.

thermal yield of an (interacting) resonance with mass M, spin J, and isospin I

$$\langle R_{I,J} \rangle = d_J \int_{m_{th}}^{\infty} dM \int \frac{d^3p}{(2\pi)^3} \frac{1}{2\pi} B_{I,J}(M) \times \frac{1}{e^{(\sqrt{p^2+M^2}-\mu)/T}+1},$$

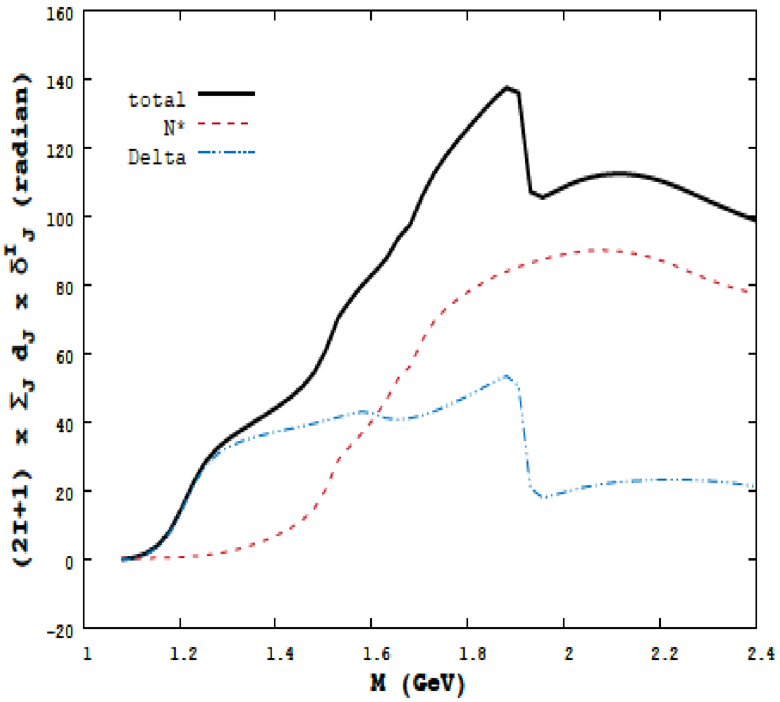
A. Andronic, pbm, B. Friman, P.M. Lo, K. Redlich, J. Stachel, arXiv:1808.03102

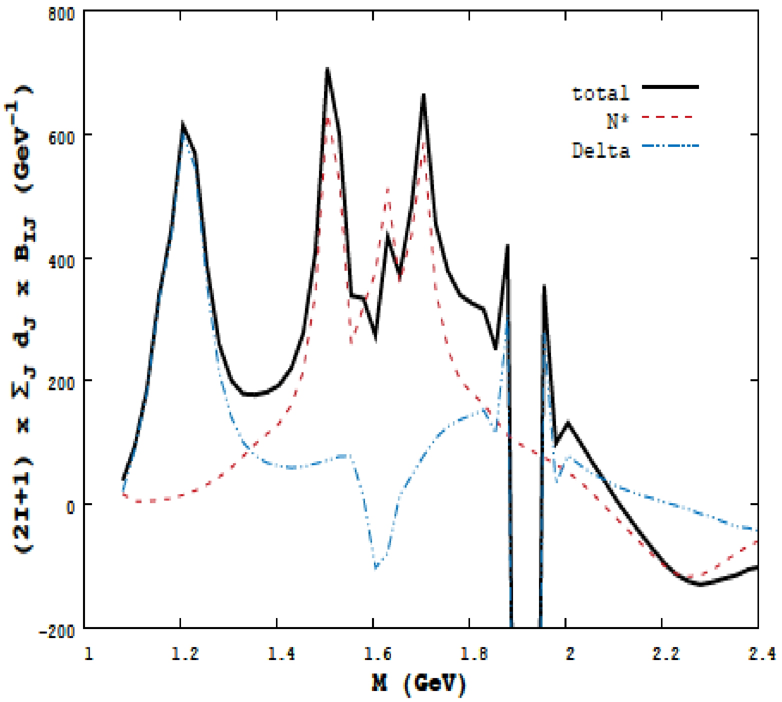
need to know derivatives of phase shifts

$$B_{I,J}(M) = 2\,\frac{d\delta_J^I}{dM}.$$

pion nucleon phase shifts and thermal weights for N* and Δ resonances

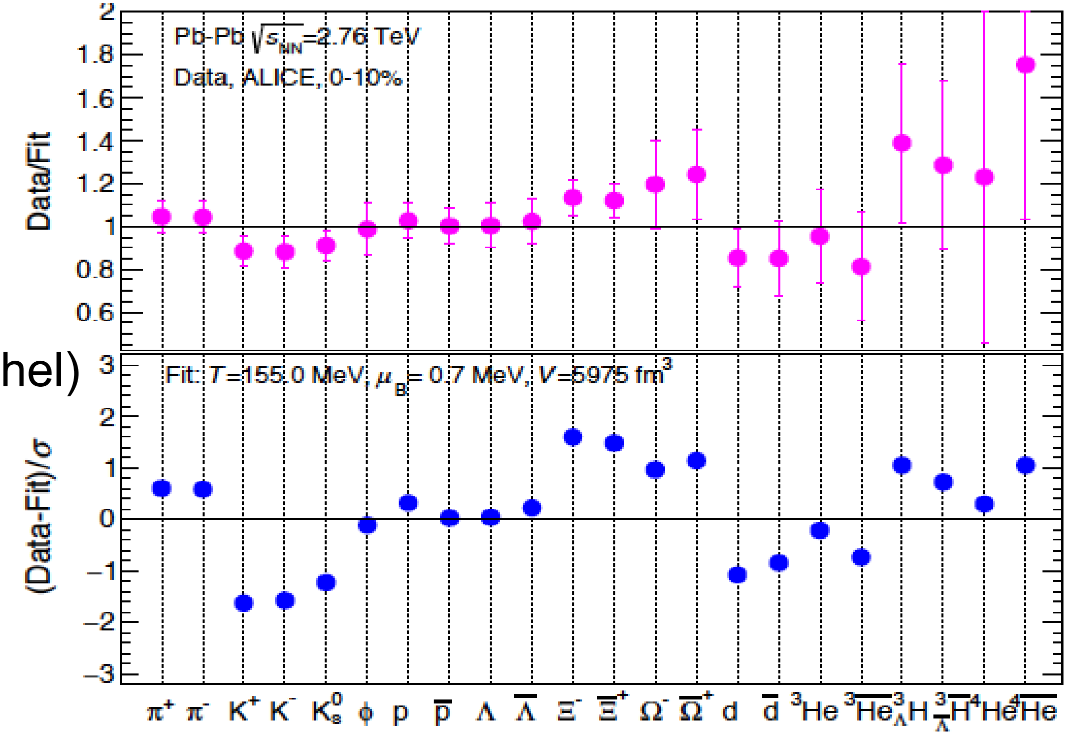
GWU/SAID phase shift analysis, 15 partial waves for each isospin channel





Aug. 2018 update: excellent description of ALICE@LHC data

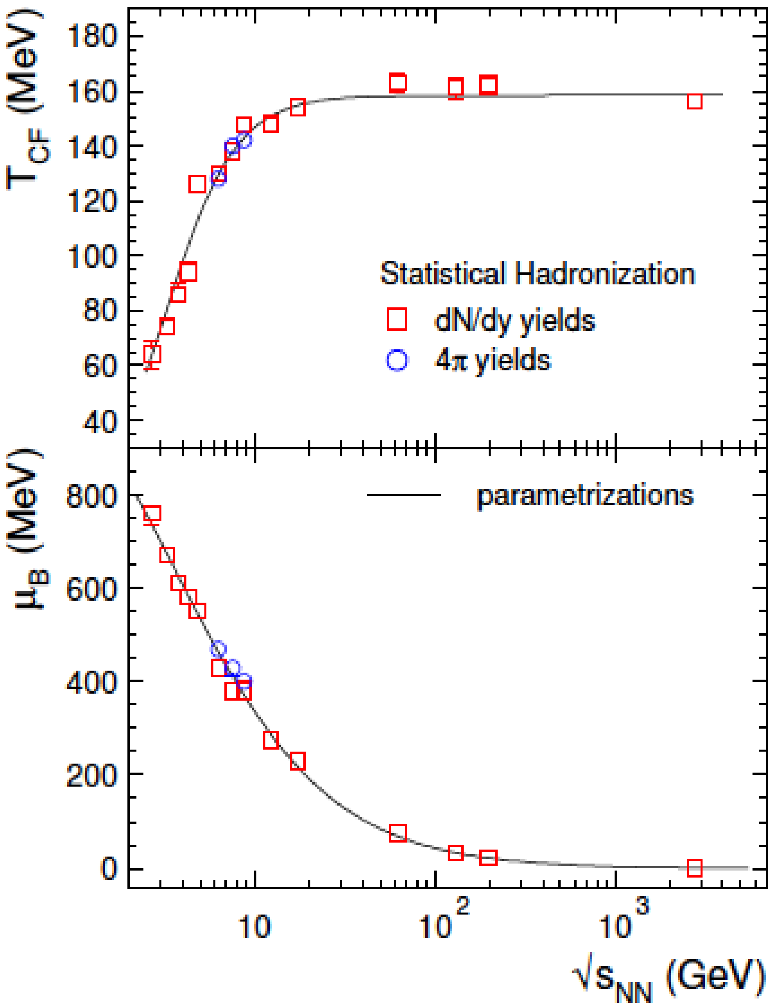
proton discrepancy of 2.8 sigma is now explained in arXiv:1808.03102
explicit phase shift description of baryon resonance region
(Andronic, pbm, Friman, Lo, Redlich, Stachel)

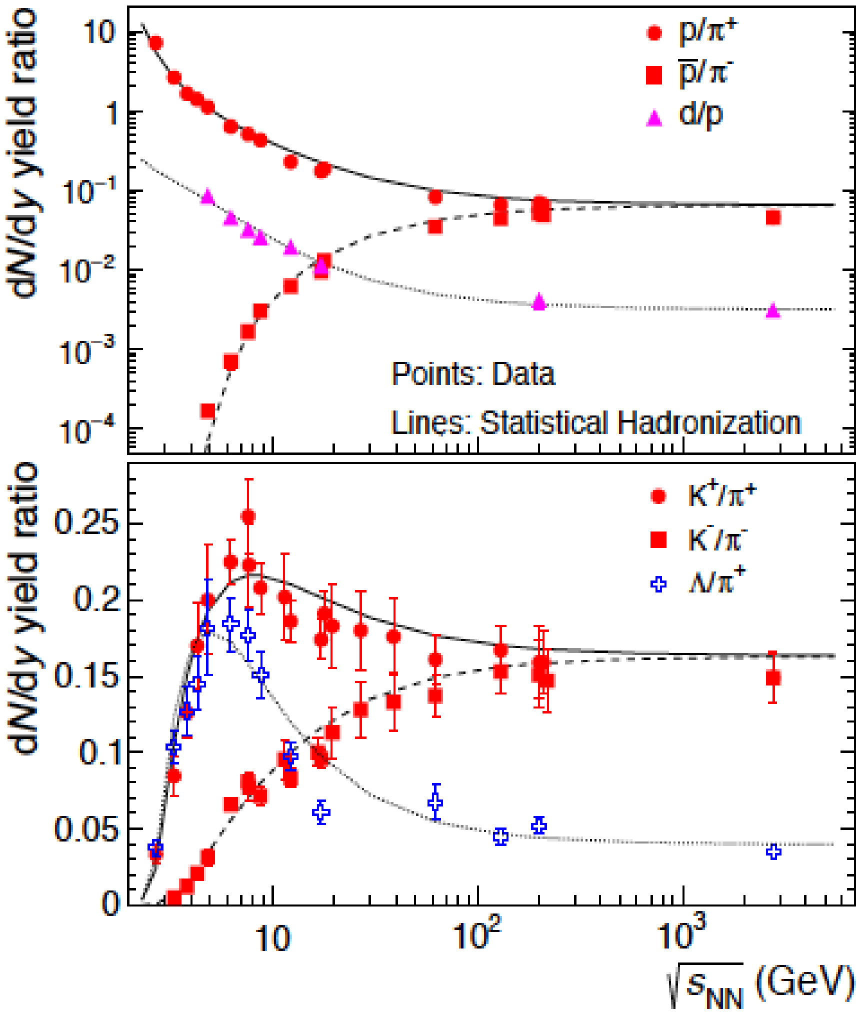


$$\chi^2 = 19.7 \text{ per } 19 \text{ dof}$$

very good fit!

energy dependence of hadron production described quantitatively



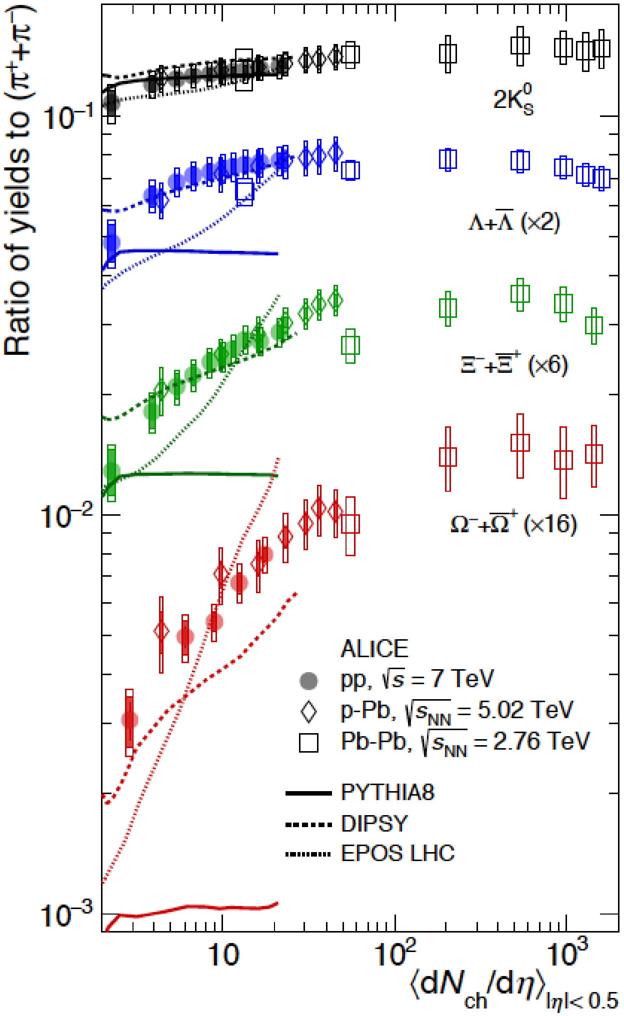


together with known energy dependence of charged hadron production in Pb-Pb collisions we can predict yield of all hadrons at all energies with < 10% accuracy

no new physics needed to describe K+/pi+ ratio including the 'horn'

approach to grand-canonical equilibration for increasing system size

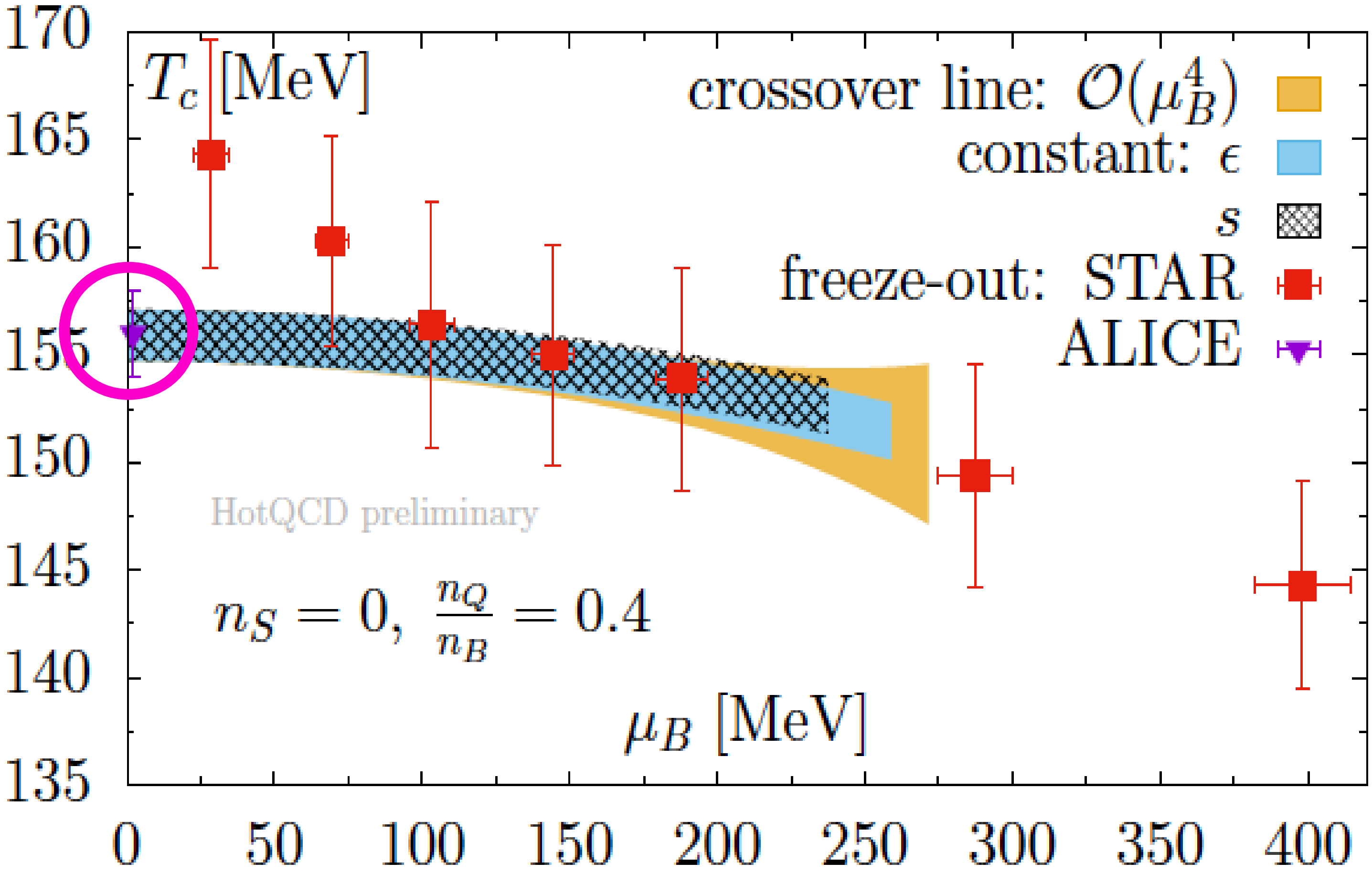
ALICE coll.
Nature Phys. 13 (2017) 535-539



grand-canonical limit
for Omega baryons

chemical freeze-out and the chiral crossover line

ALICE point: 156 ± 1.5 ± 3 (sys) MeV, measured with TPC and Si vertex detector
STAR points: measured with TPC only, feeding from weak decays
lattice: 156 ± 1.5 MeV



lattice: BNL-Bielefeld coll. 1807.05607

a note on the chemical freeze-out temperature

T_{chem} = 156.5 ± 1.5 MeV from fit to all particles

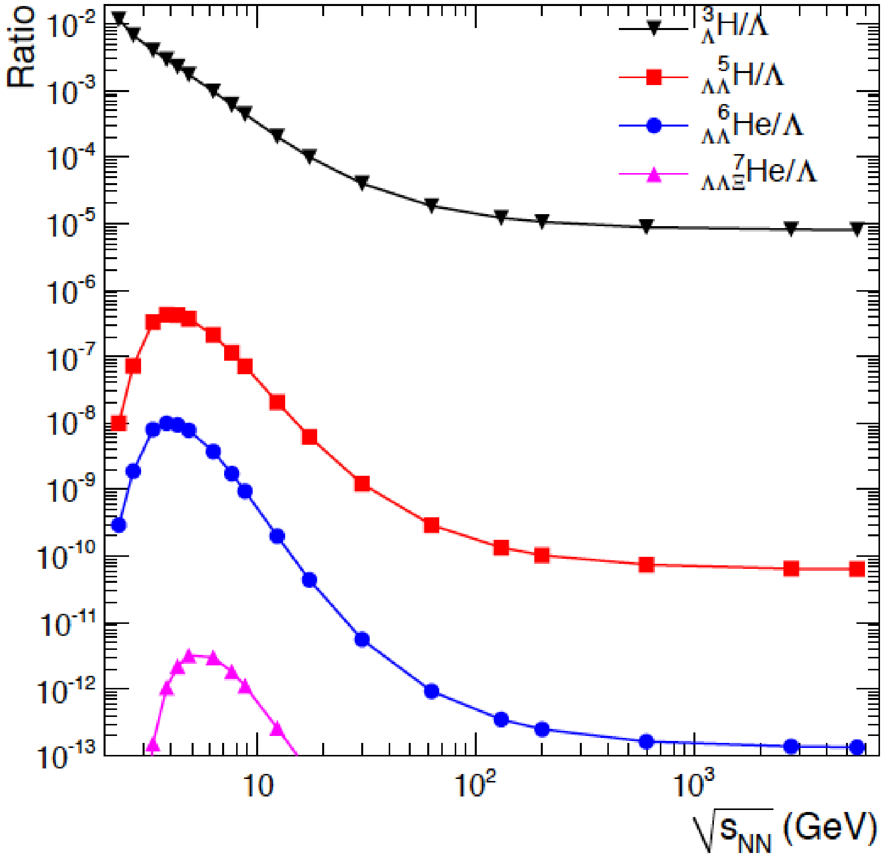
there is an additional uncertainty because of the poorly known hadronic mass spectrum for masses > 2 GeV

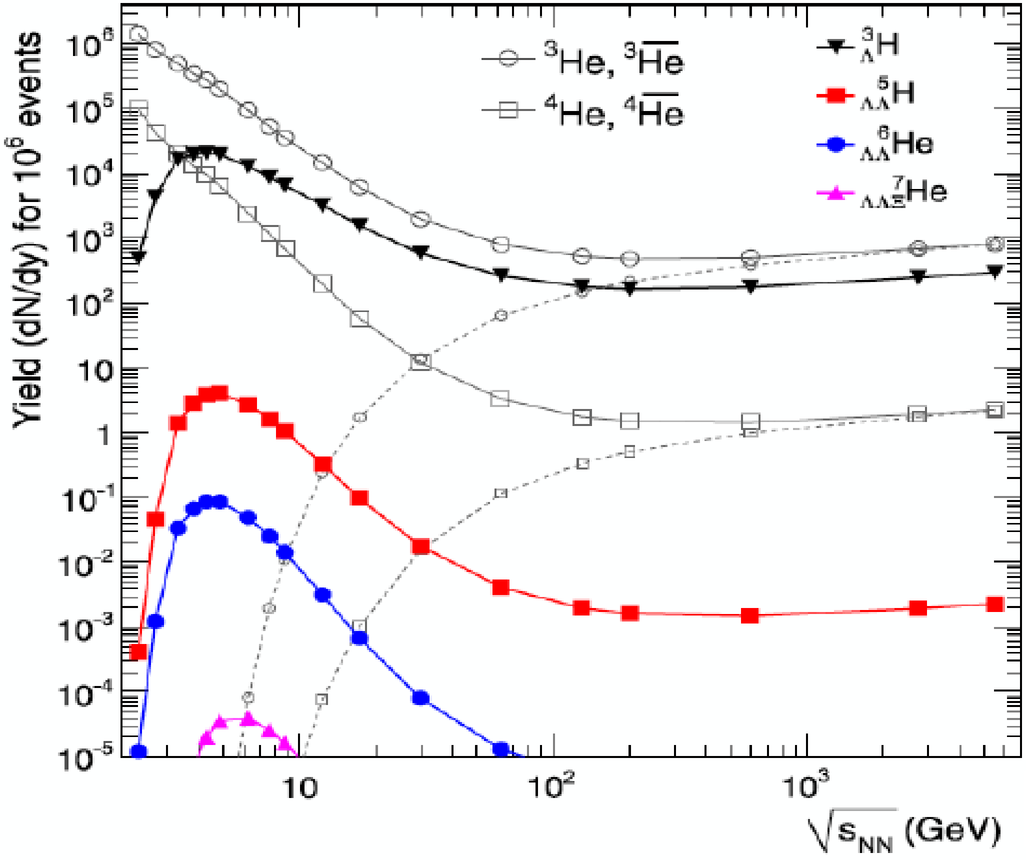
for d, 3He, hypertriton and alpha, there is very little feeding from heavier states and none from high mass states in the hadronic mass spectrum, for these particles the temperature T_{nuc} can be determined 'on the back of an envelope' :

T_{nuc} = 159 ± 5 MeV, independent of hadronic mass spectrum

now loosely bound objects

exciting opportunities for the upcoming accelerator facilities
NICA, FAIR/CBM, J-Parc





Andronic, pbm, Stachel, Stoecker
Phys.Lett. B697 (2011) 203-207

The Hypertriton

mass = 2990 MeV, binding energy = 2.3 MeV

Lambda sep. energy = 0.13 MeV

molecular structure: (p+n) + Lambda

2-body threshold: (p+p+n) + pi- = ^{3}He + pi-

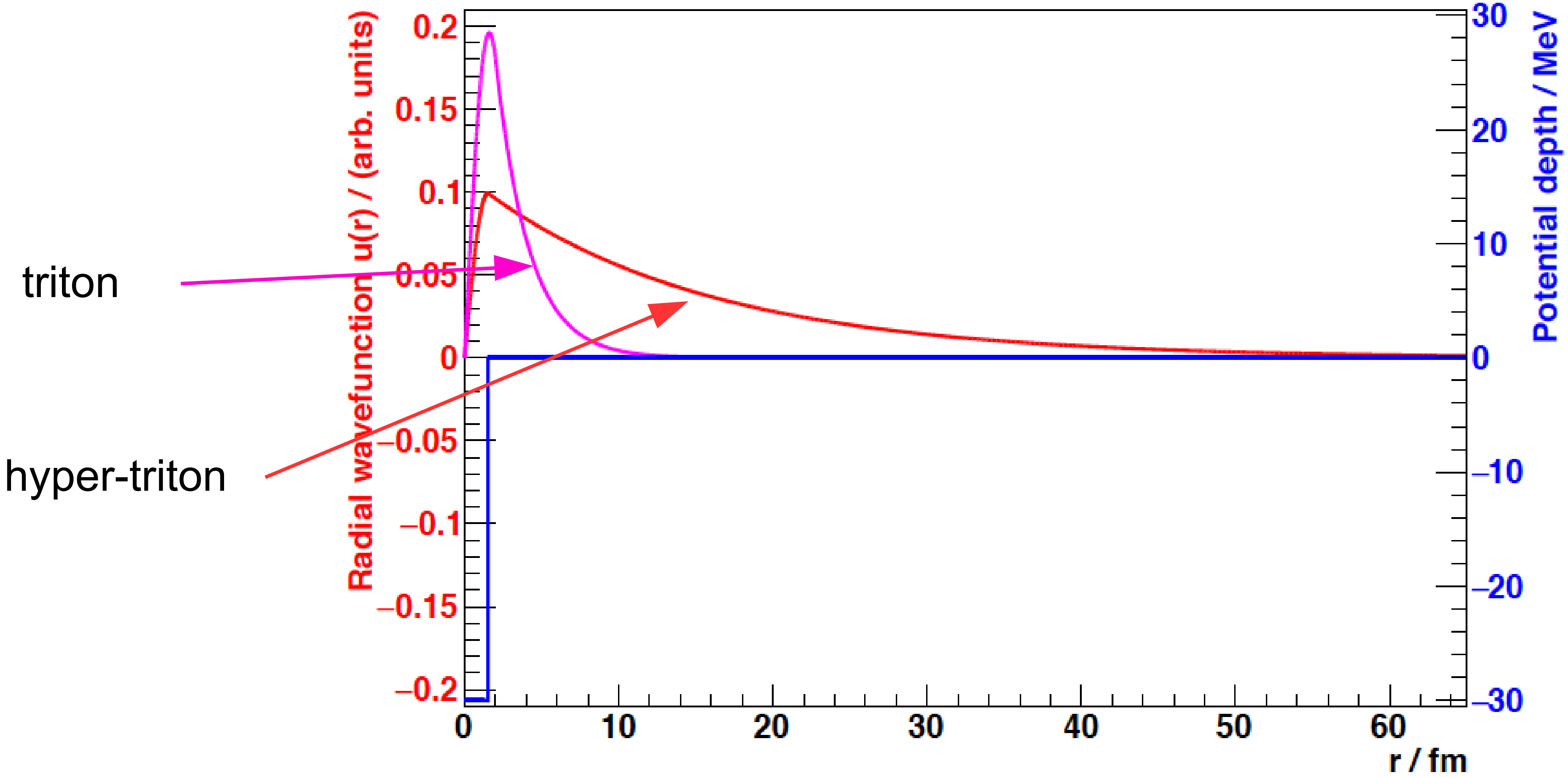
rms radius = $(4\ \text{B.E.}\ M_{red})^{-1/2}$ = 10.3 fm = rms separation between d and Lambda

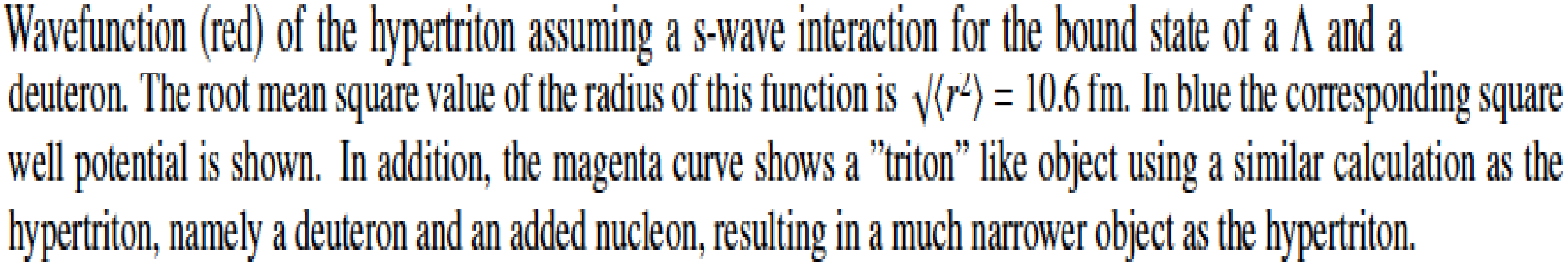
in that sense: hypertriton = (p n Lambda) = (d Lambda) is the ultimate halo state

yet production yield is fixed at 156 MeV temperature (about 1000 x separation energy.)

wave function of the hyper-triton – schematic picture

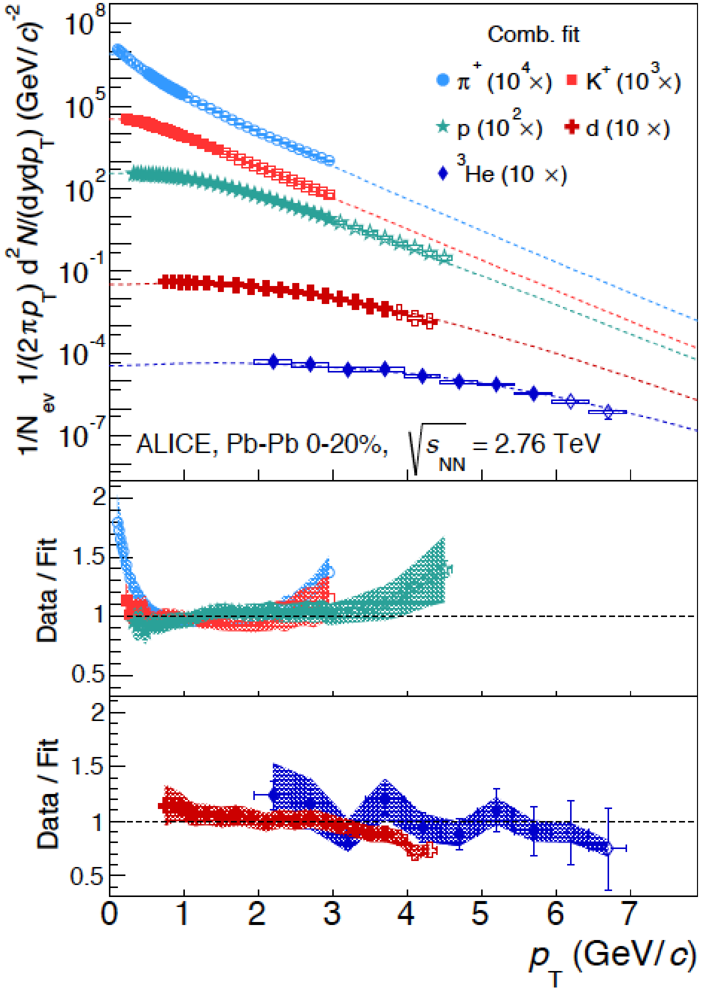
figure by Benjamin Doenigus, August 2017



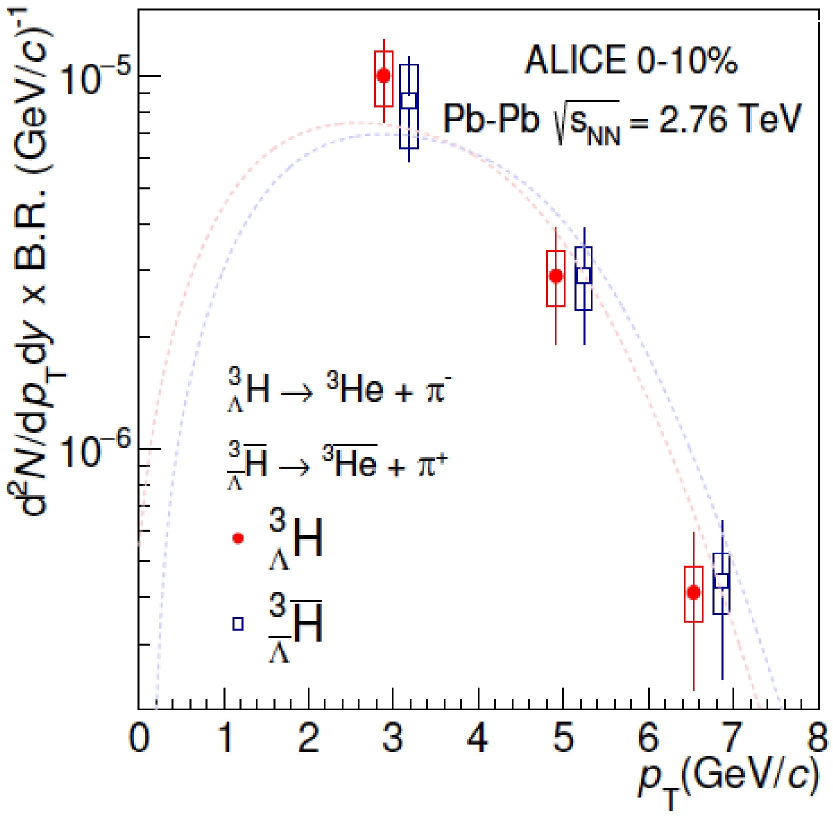


Wavefunction (red) of the hypertriton assuming a s-wave interaction for the bound state of a Λ and a deuteron. The root mean square value of the radius of this function is $\sqrt{\langle r^2\rangle} = 10.6$ fm. In blue the corresponding square well potential is shown. In addition, the magenta curve shows a ”triton” like object using a similar calculation as the hypertriton, namely a deuteron and an added nucleon, resulting in a much narrower object as the hypertriton.

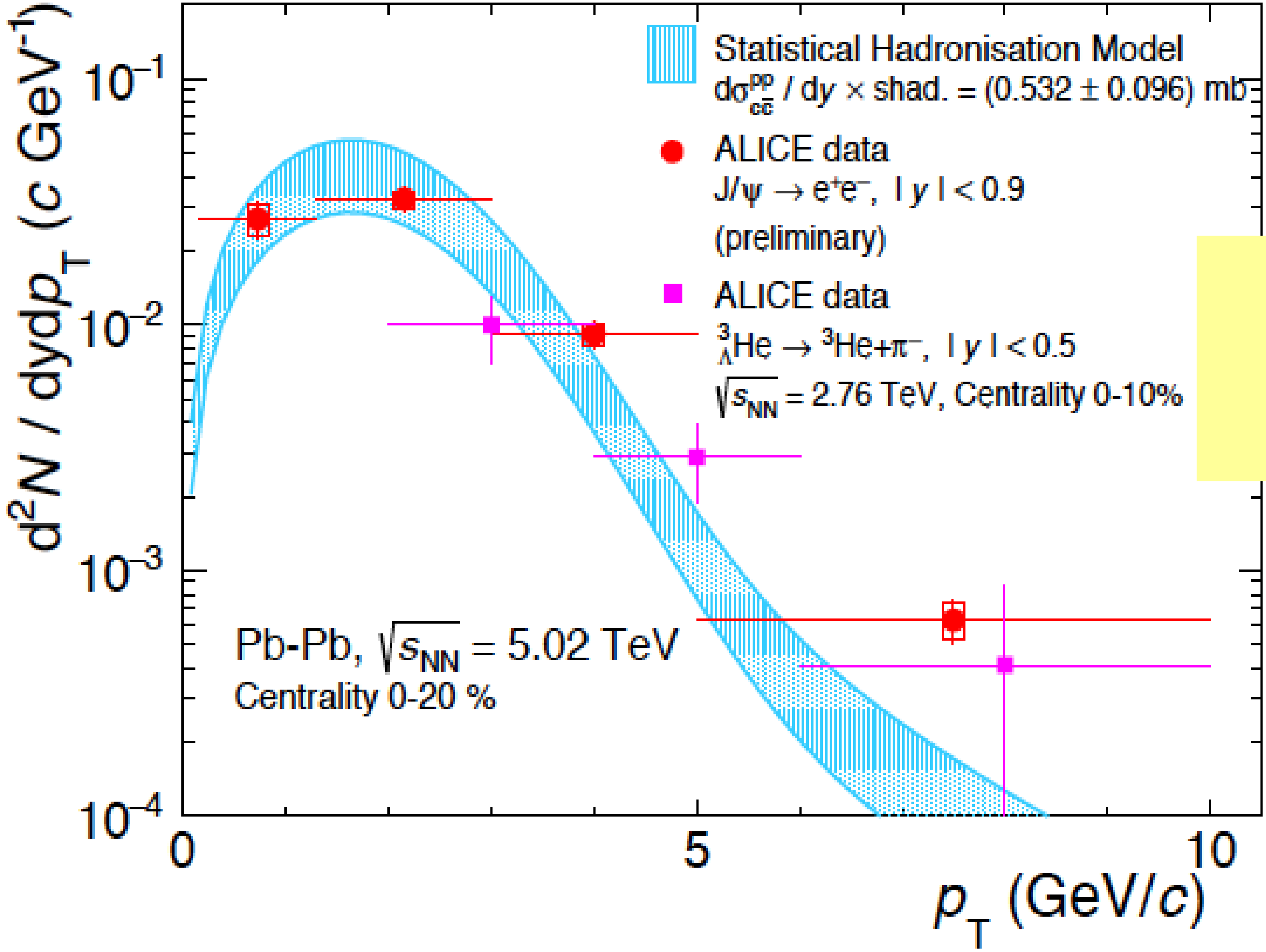
light nuclei flow with same fluid velocity as pions, kaons, and protons



even hyper-triton flows with same common fluid velocity



J/psi and hyper-triton described with the same flow parameters in the statistical hadronization model

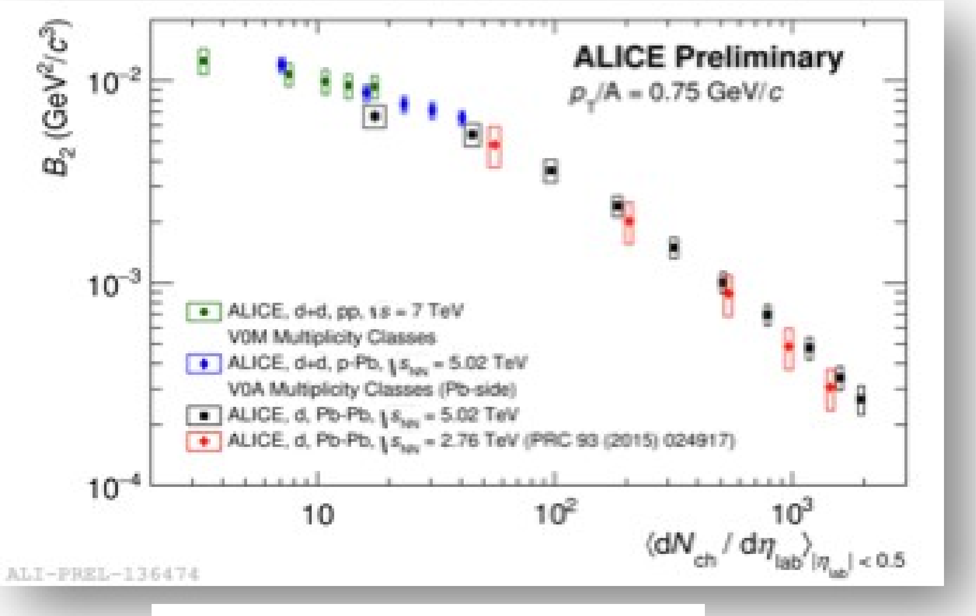


binding energies:
J/psi 600 MeV
hypertriton 2.2 MeV
Lambda S.E. 0.2 MeV

from review: hypernuclei and other loosely bound objects produced in nuclear collisions at the LHC, pbm and Benjamin Doenigus, to appear soon

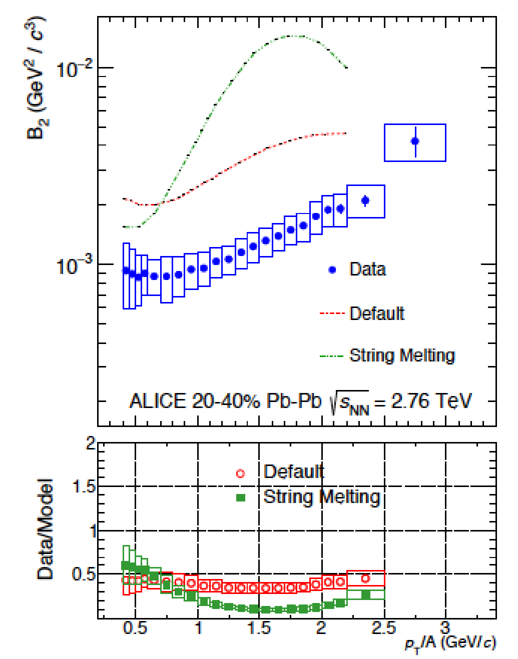
is coalescence approach an alternative?

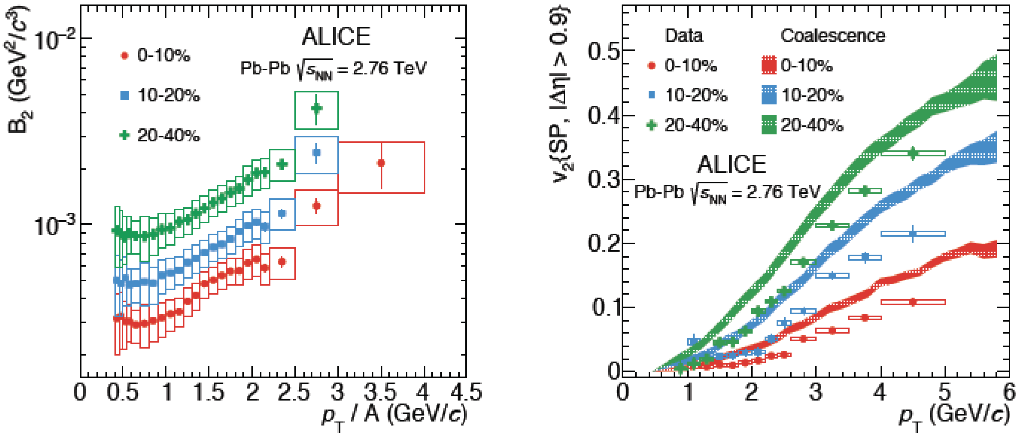
$$E_i \frac{d^3N_i}{dp_i^3} = B_A \left(E_p \frac{d^3N_p}{dp_p^3} \right)^A \qquad B_A = \left(\frac{4\pi}{3} p_0^3 \right)^{A-1} \frac{M}{m^A}$$



centrality and p_T dependence of coalescence parameter not understood and not well reproduced by models such as AMPT

ALICE: arXiv:1707.07304





coalescence approach, general considerations for loosely bound states

- production yields of loosely bound states is entirely determined by mass, quantum numbers and fireball temperature.

- hyper-triton and 3He have very different wave functions but essentially equal production yields.

- energy conservation needs to be taken into account when forming objects with baryon number A from A baryons
- coalescence of off-shell nucleons does not help as density must be << nuclear matter density, see below

- delicate balance between formation and destruction; maximum momentum transfer onto hyper-triton before it breaks up: $\Delta Q_{max} < 20$ MeV/c, typical pion momentum p_pi = 250 MeV/c, typical hadronic momentum tranfer > 100 MeV/c.

- hyper-triton interaction cross section with pions or nucleons at thermal freeze-out is of order $\sigma > 70\ \mathrm{fm}^2$. For the majority of hyper-tritons to survive, the mfp λ has to exceed 15 fm → density of fireball at formation of hyper-triton $n < 1/(\lambda\ \sigma) = 0.001/\mathrm{fm}^3$. Inconsistent with formation at kinetic freeze-out, where $n \approx 0.05/\mathrm{fm}^3$.

is large size of light nuclei and hypernuclei an issue for statistical hadronization model?

note: in thermal approach, the only scale is temperature T
at LHC energy and below, T < 160 MeV

at such a scale, momentum transfer q=T, form factors of hadrons are sampled at q^2 = T^2
this implies that sizes of hadrons < 2 fm cannot be resolved

since $G(q) \sim 1 - q^2 R^2/6$

and since all (rms) radii for nuclei with A = 2, 3, and 4 are smaller than 2 fm, the correction due to the finite size of nuclei will not exceed 35%

the actual change from this on thermal model results should be much less as only the relative change between normal hadrons and light nuclei matters, the overall change only leads to a volume correction, so the correction for nuclei is estimated to be less than 25%

but hyper-triton has much larger radius > 5 fm?
measured yield of hyper-triton and 3He is well compatible with thermal prediction, even though wave function is very different – any wave function correction must be small

the agreement of the baryon number 3 states is also big problem for coalescence model

see also the detailed analysis by Francesca Bellini and Alexander Kalweit, arXiv:1807.05894,
Benjamin Doenigus and Nicole Loeher, GSI-EMMI meeting, Feb. 2018

How can 'thermal production near the phase boundary' i.e. at T ~ 155 MeV be reconciled with binding energies < 5 MeV and large break-up cross sections?

a possible way out

Quark Model Spectroscopy

Why does the quark model work so well?

Why do M and B body plans dominate?

Why don't multibaryons make one big bag?

Frank Wilczek, QM2014 introductory talk

see also the recent review:
Marek Karliner, Jonathan L. Rosner, Tomasz Skwarnicki, arXiv:1711.10626

doorway state hypothesis:
all nuclei and hyper-nuclei, penta-quark and X,Y,Z states are formed as virtual, compact multi-quark states at the phase boundary. Then slow time evolution into hadronic representation. Excitation energy about 20 MeV, time evolution about 10 fm/c

Andronic, pbm, Redlich, Stachel, arXiv :1710.09425

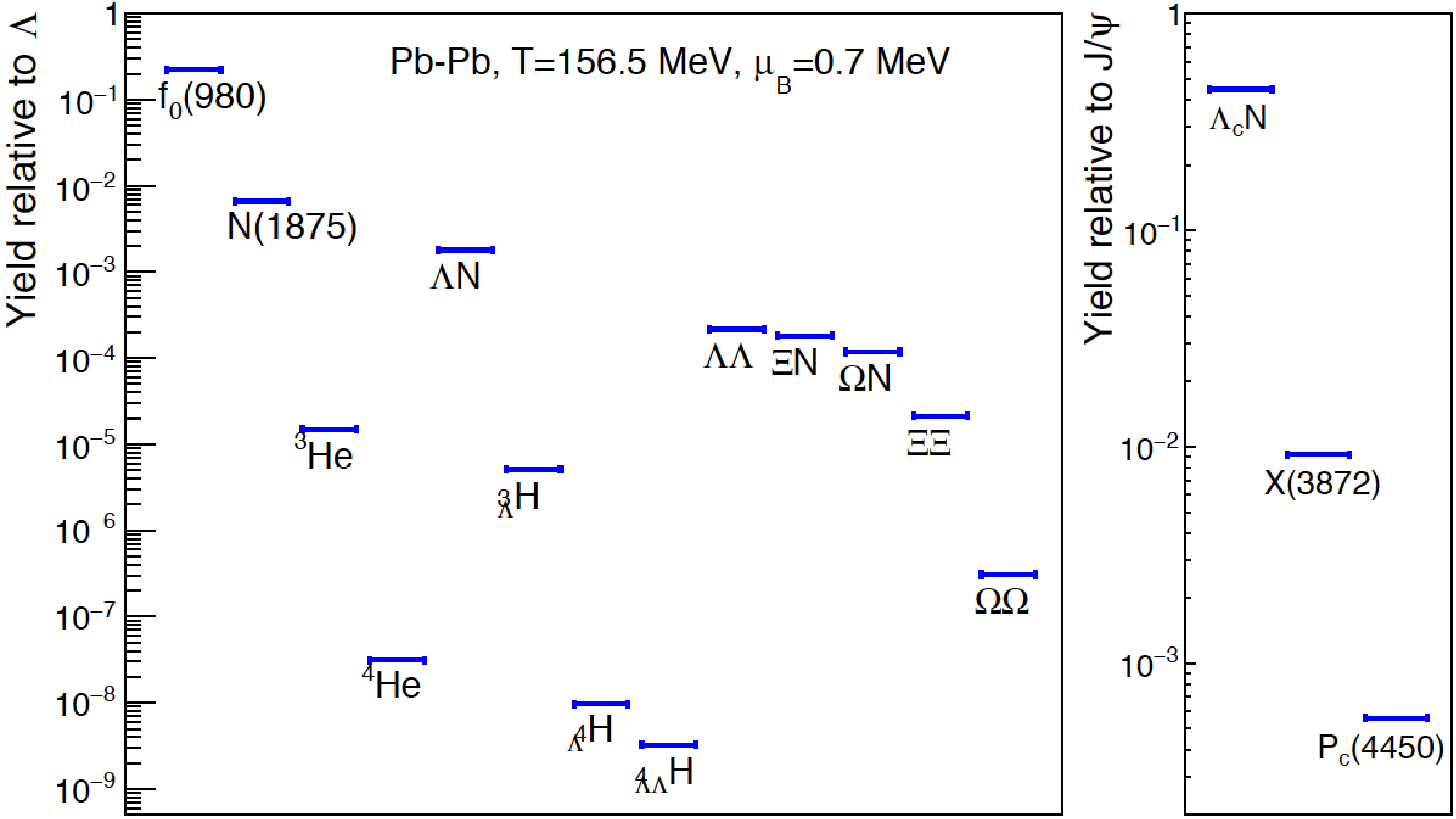
How can this be tested?

precision measurement of spectra and flow pattern for light nuclei and hyper-nuclei, penta-quark and X,Y,Z states from pp via pPb to Pb-Pb

a major new opportunity for ALICE Run3/4
and beyond LS4 for X,Y,Z and penta-quark states

thermal production yields of exotic states in central Pb-Pb collisions at 5 TeV/u

Andronic, pbm, Koehler, Redlich, Stachel
preprint in preparation



example: X(3872)

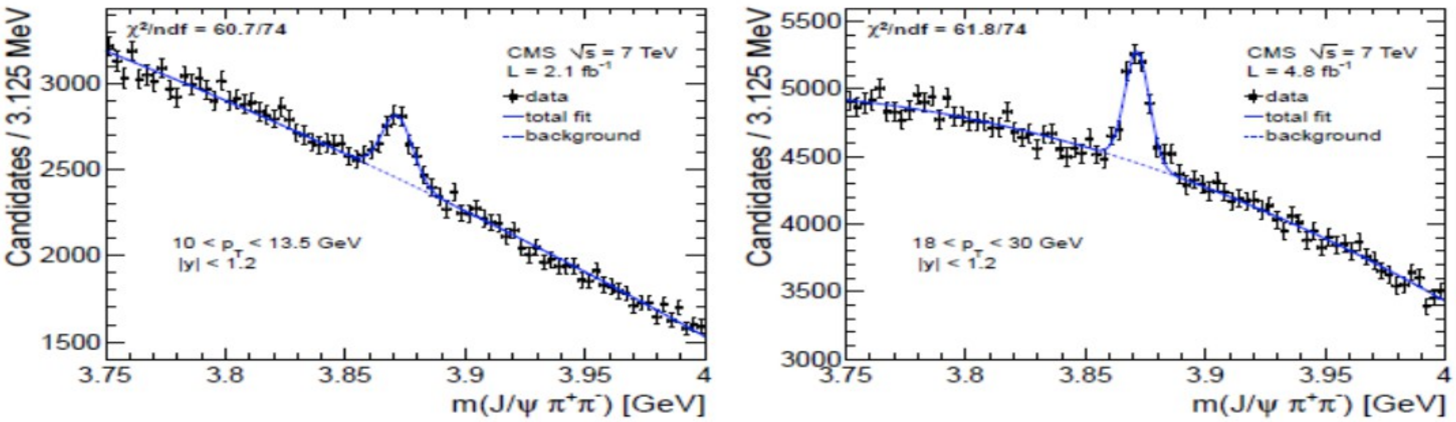
X(3872)

- 2003 -

$$B^{\pm} \rightarrow K^{\pm} \boxed{\pi^{+}\pi^{-} J/\psi}$$

$$M = 3872.0 \pm 0.6 \pm 0.5\,\mathrm{MeV}$$

- 2013 -



***X*(3872)**

$I^{G}(J^{PC}) = 0^{+}(1^{++})$

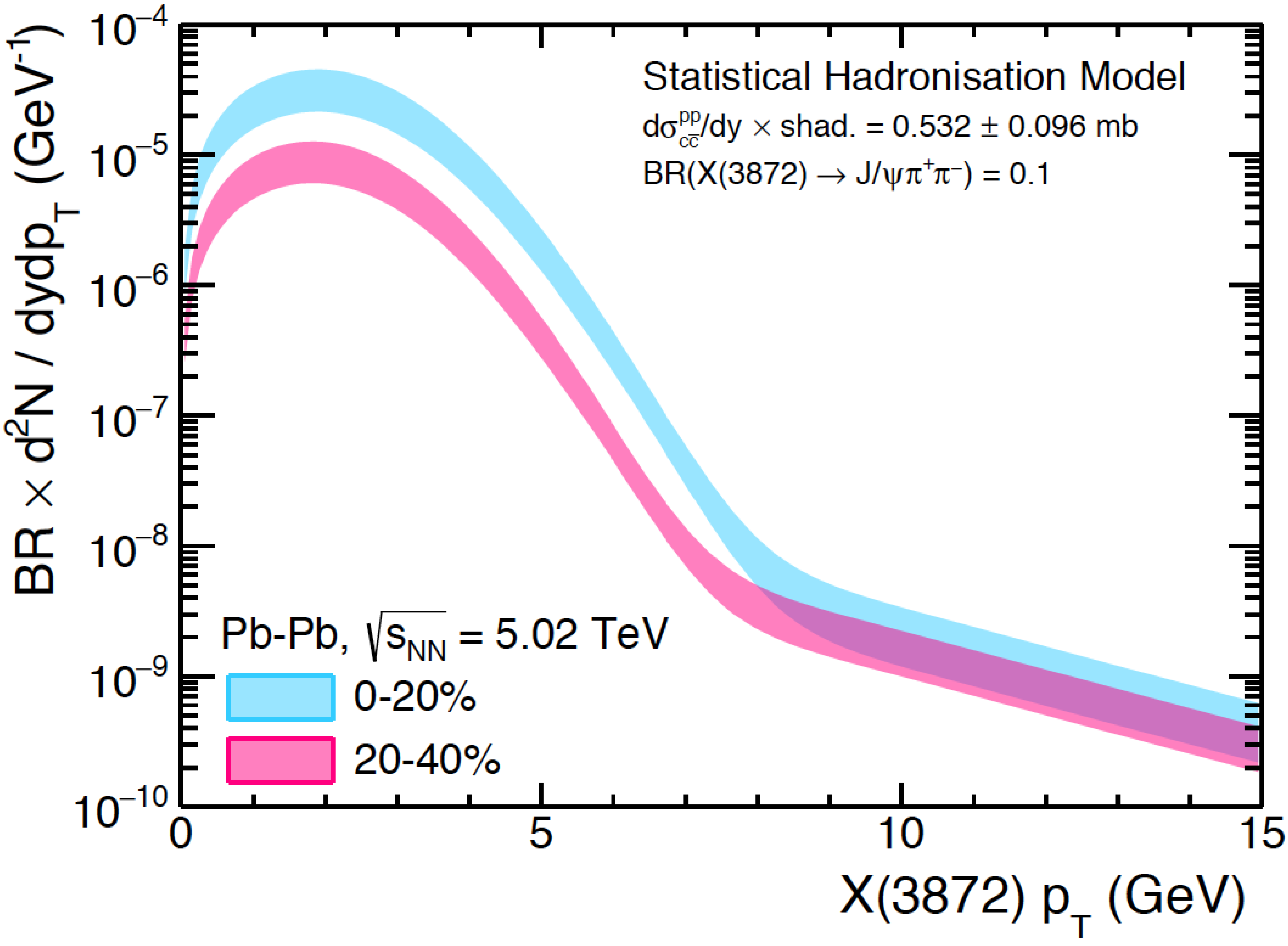
Mass $m = 3871.69 \pm 0.17$ MeV

$m_{X(3872)} - m_{J/\psi} = 775 \pm 4$ MeV

$m_{X(3872)} - m_{\psi(2S)}$

Full width $\Gamma < 1.2$ MeV, CL = 90%

transverse momentum spectrum for X(3872) in the statistical hadronization model Pb-Pb collisions at 5 TeV/u



Summary

- statistical hadronization approach describes well production of hadrons in relativistic nuclear collisions
- at RHIC and LHC energies the resulting chemical freeze-out temperatures agree quantitatively with LQCD predictions
- the 'proton anomaly' is solved by computation of the virial coefficient for the pion-nucleon system
- theoretically ill controllable excluded volume corrections are unnecessary
- even very loosely bound objects are apparently produced near the phase boundary → maybe produced as compact multi-quark objects
- many of the ideas were, over the last decade, discussed in regular workshops at CCNU, Wuhan

additional slides

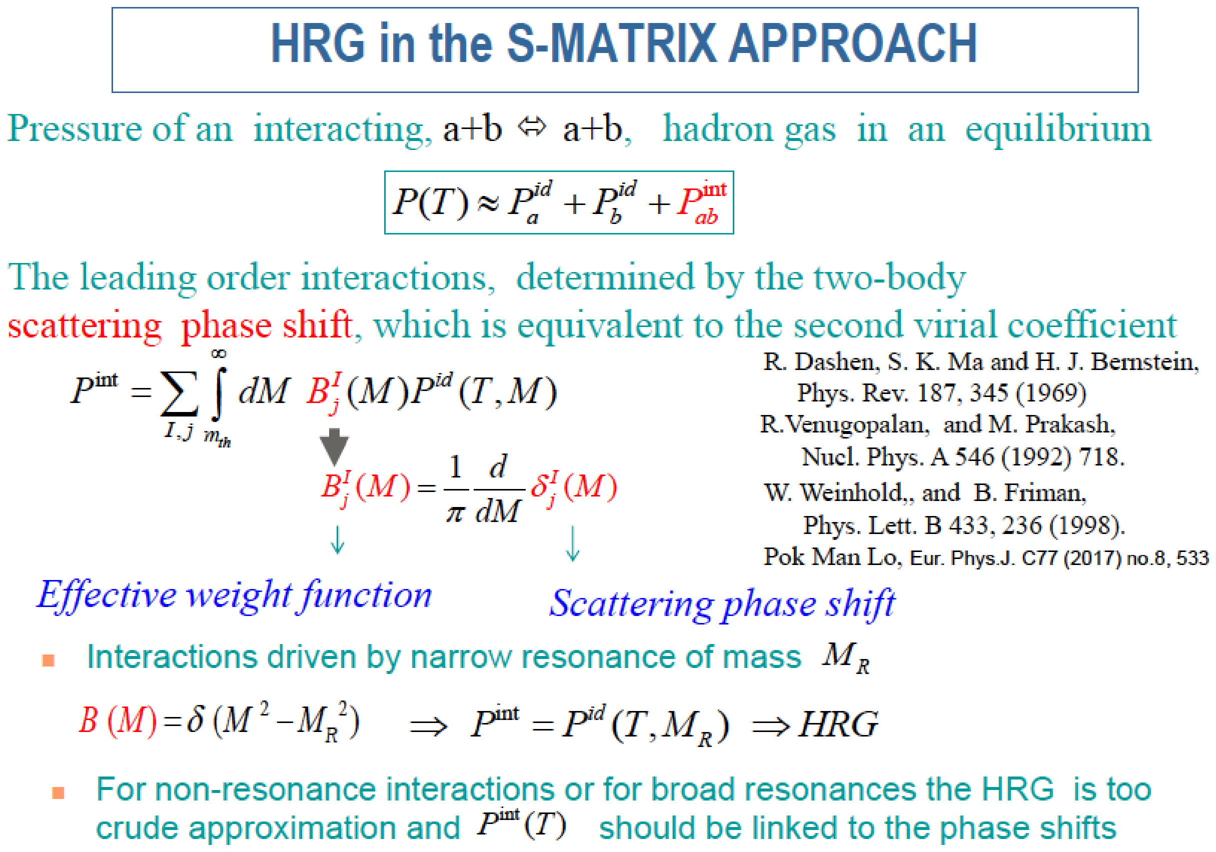
[Braun-Munzinger and Stachel, PLB 490 (2000) 196]
[Andronic, Braun-Munzinger and Stachel, NPA 789 (2007) 334]

- Charm quarks are produced in initial hard scatterings ($m_{c\bar{c}} \gg T_{\mathrm{c}}$) and production can be described by pQCD ($m_{c\bar{c}} \gg \Lambda_{\mathrm{QCD}}$)
- Charm quarks survive and *thermalise* in the QGP
- Full screening before T_{CF}
- Charmonium is formed at phase boundary (together with other hadrons)
- Thermal model input ($T_{\mathrm{CF}}, \mu_{\mathrm{b}} \rightarrow n_X^{\mathrm{th}}$)

$$N_{c\bar{c}}^{\mathrm{dir}} = \underbrace{\frac{1}{2} g_c V \left(\sum_i n_{D_i}^{\mathrm{th}} + n_{\Lambda_i}^{\mathrm{th}} + \cdots \right)}_{\text{Open charm}} + \underbrace{g_c^2 V \left(\sum_i n_{\psi_i}^{\mathrm{th}} + n_{\chi_i}^{\mathrm{th}} + \cdots \right)}_{\text{Charmonia}}$$

- Canonical correction is applied to $n_{\mathrm{oc}}^{\mathrm{th}}$
- Outcome $N_{\mathrm{J}/\psi}, N_D, \ldots$

going beyond the non-interacting HRG – next 3 slides from K. Redlich, QM18



HRG in the S-MATRIX APPROACH

Pressure of an interacting, a+b ⇔ a+b, hadron gas in an equilibrium

$$\boxed{P(T) \approx P_a^{id} + P_b^{id} + P_{ab}^{\text{int}}}$$

The leading order interactions, determined by the two-body scattering phase shift, which is equivalent to the second virial coefficient

$$P^{\text{int}} = \sum_{I,j} \int_{m_{th}}^{\infty} dM \; B_j^I(M) P^{id}(T,M)$$

$$B_j^I(M) = \frac{1}{\pi} \frac{d}{dM} \delta_j^I(M)$$

Effective weight function

Scattering phase shift

R. Dashen, S. K. Ma and H. J. Bernstein,
Phys. Rev. 187, 345 (1969)
R.Venugopalan, and M. Prakash,
Nucl. Phys. A 546 (1992) 718.
W. Weinhold,, and B. Friman,
Phys. Lett. B 433, 236 (1998).
Pok Man Lo, Eur. Phys.J. C77 (2017) no.8, 533

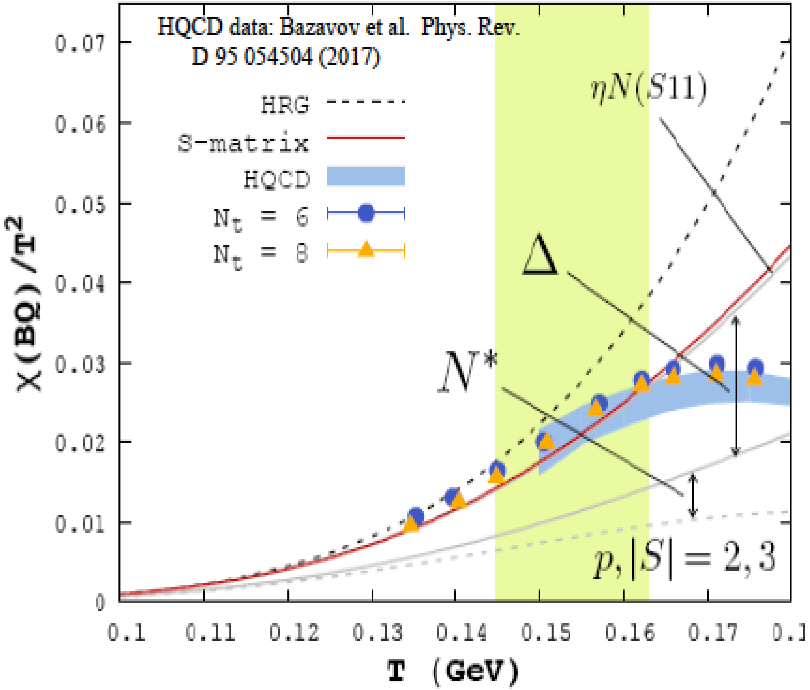
- Interactions driven by narrow resonance of mass M_R

$$B(M) = \delta(M^2 - M_R^2) \quad \Rightarrow \quad P^{\text{int}} = P^{id}(T, M_R) \;\Rightarrow HRG$$

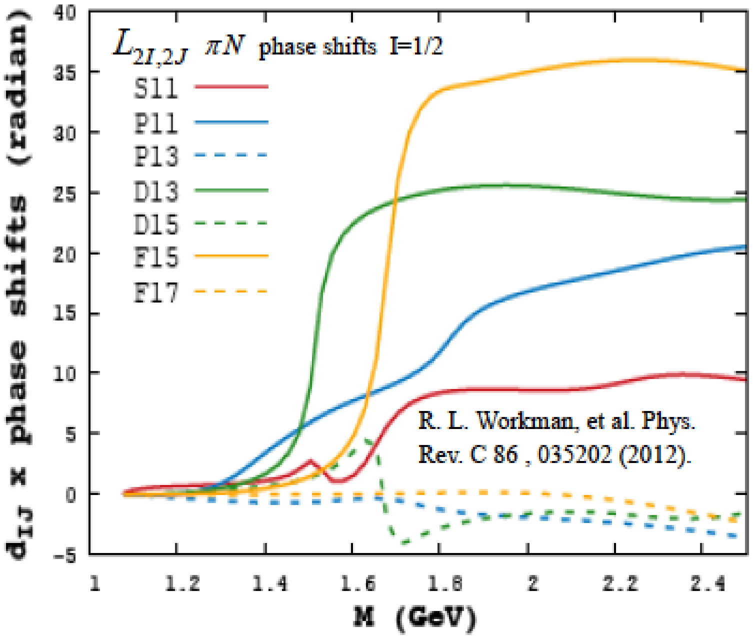
- For non-resonance interactions or for broad resonances the HRG is too crude approximation and $P^{\text{int}}(T)$ should be linked to the phase shifts considering all pion-nucleon phase shifts with isospin 1/2 and 3/2

Probing non-strange baryon sector in πN- system

Pok Man Lo, B. Friman, C. Sasaki & K.R., Phys.Lett. B778 (2018)



$$\chi_{BQ} = (\chi_{BB} - |\chi_{BS}|)/2$$

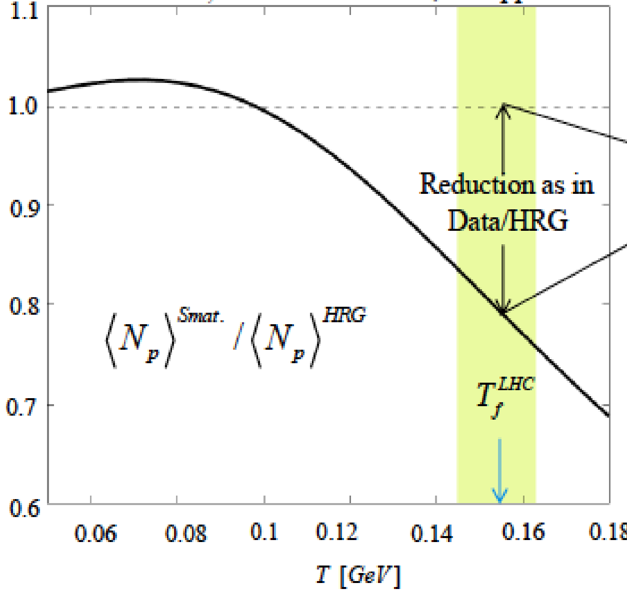


$$\Delta\chi_{BQ} \approx \sum_{I_z,j,B} d_j BQ \int dM \int d^3p \, \frac{1}{T} \frac{d\delta_j^I}{dM} \times e^{-\beta\sqrt{p^2+M^2}} \left(1+e^{-\beta\sqrt{p^2+M^2}}\right)^{-2}$$

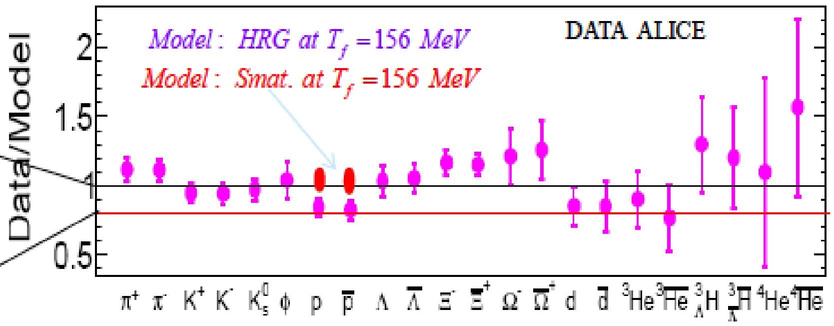
- Considering contributions of all πN $\delta_j^{I=(1/2),\,(3/2)}$ (N^*, Δ^*resonances) to χ_{BQ} within S-matrix approach, reduces the HRG predictions towards the LQCD in the chiral crossover $0.15 < T < 0.16\ GeV$

Phenomenological consequences: proton production yields

A. Andronic, P. Braun-Munzinger, B. Friman, Pok Man Lo, J. Stachel & K.R., too appear.

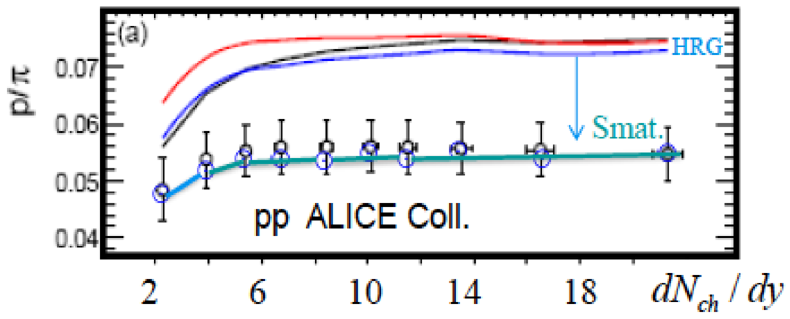


HRG: A. Andronic, P. Braun-Munzinger, J. Stachel & K.R.



- Yields of protons in AA collisions at LHC is consistent with S-matrix result within 1σ

HRG: N. Sharma, J. Cleymans, B. Hippolite, arXiv: 1803.05409



- Yields of protons in the S-matrix is suppressed relative to HRG

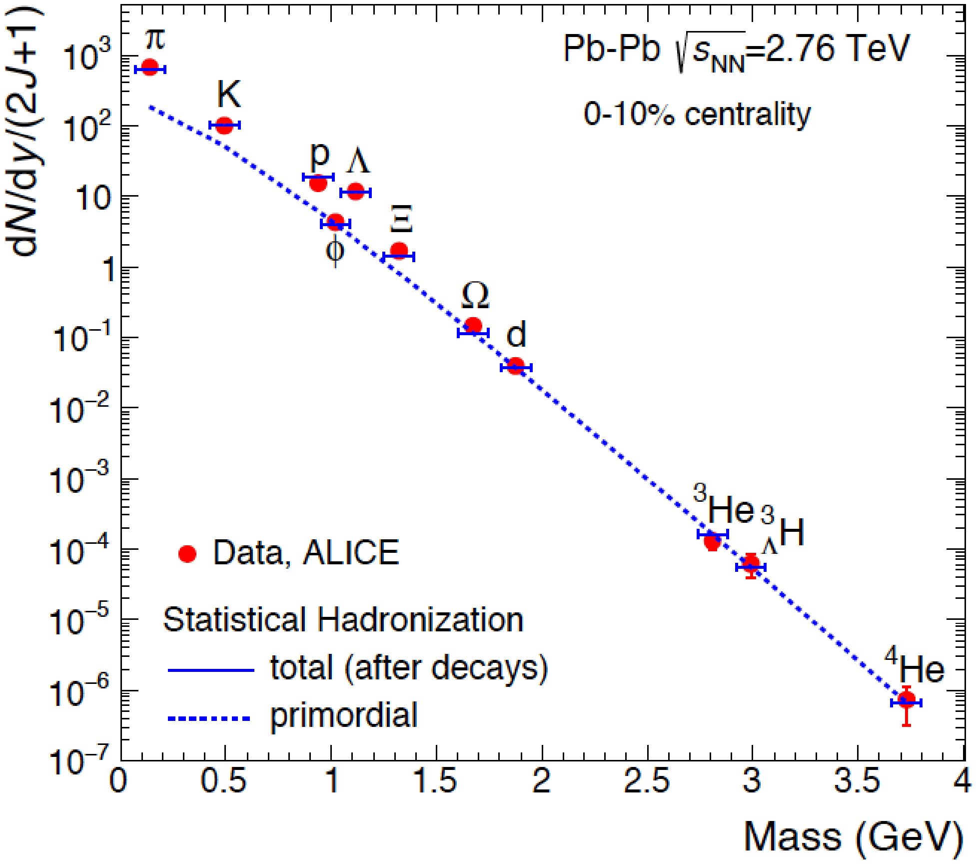
For further consequences of smat. See also: P. Huovinen, P. Petreczky Phys. Lett. B77 (2018) P. Huovinen, poster QM2018

- S-matrix results well consistent with pp data

points a way to explain 'proton puzzle',new description to appear soon

excellent agreement over 9 orders of magnitude

agreement over 9 orders of magnitude with QCD statistical operator prediction



yield of light nuclei predicted in: pbm, J. Stachel, J.Phys. G28 (2002) 1971-1976,
J.Phys. G21 (1995) L17-L20