

Search for the QCD Critical Point -

Conserved Charge Fluctuations and Light Nuclei Production at RHIC Beam Energy Scan



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Heavy Flavor Production and Forty Years of Quark-Gluon Plasma

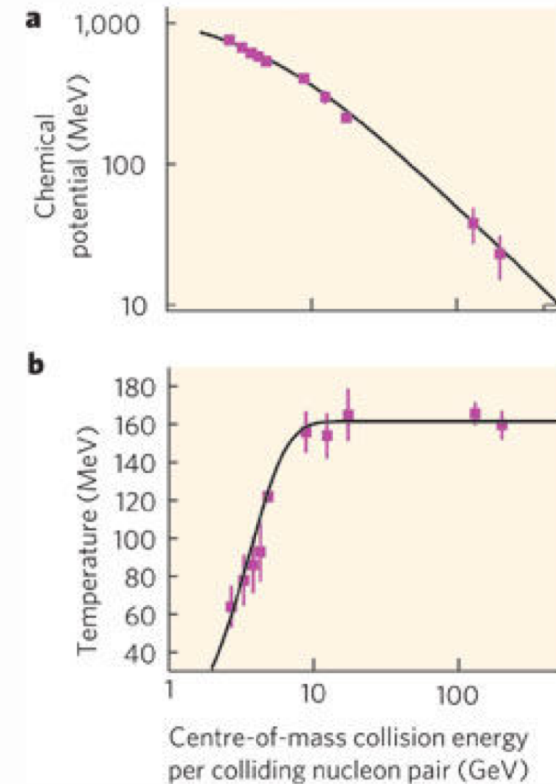
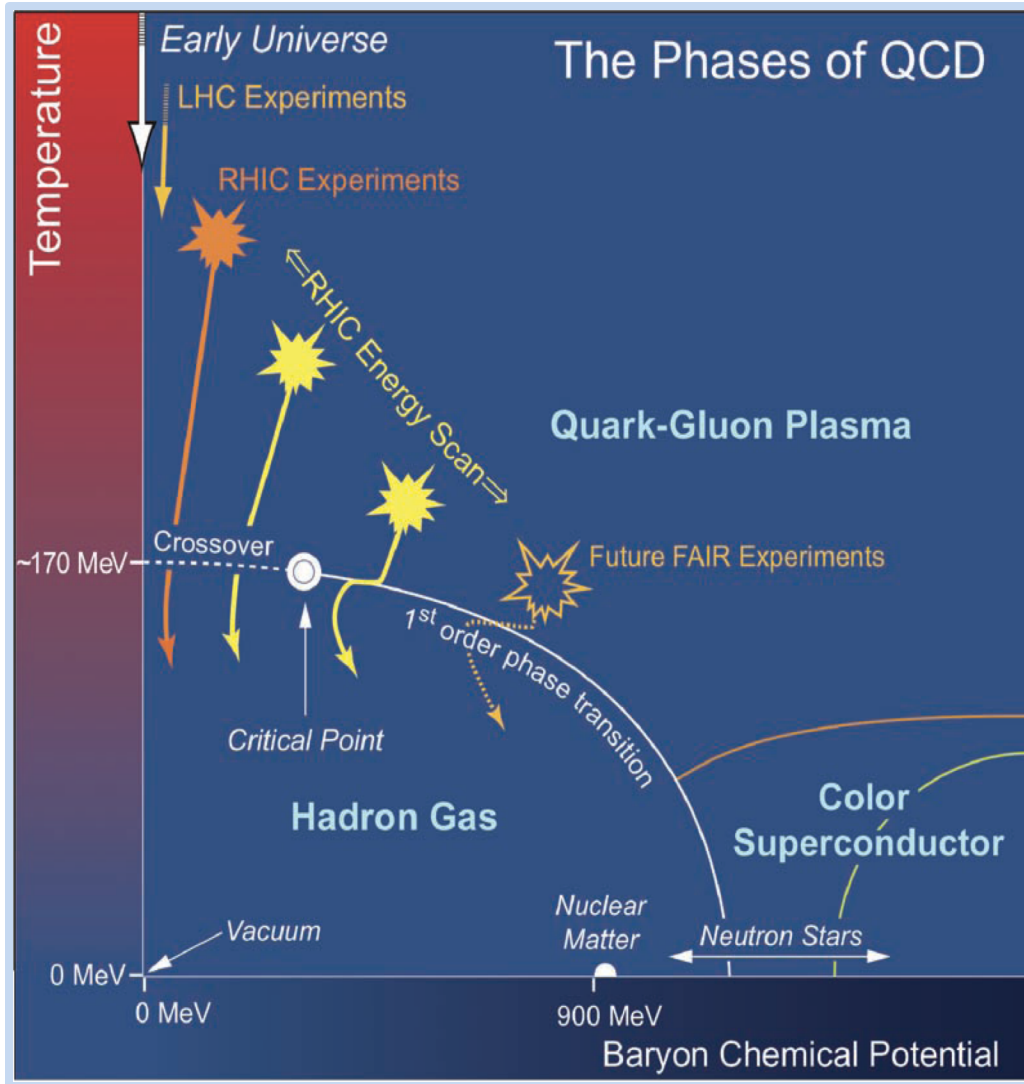
10/08-10/11/2018

Happy Birthday to Edward



After giving seminar at Stony Brook University, April 30, 2015. (photo taken by Jiangyong Jia).

- 1. Introduction**
- 2. Conserved Charge Fluctuations**
- 3. Light Nuclei (d, t) Production**
- 4. Summary**



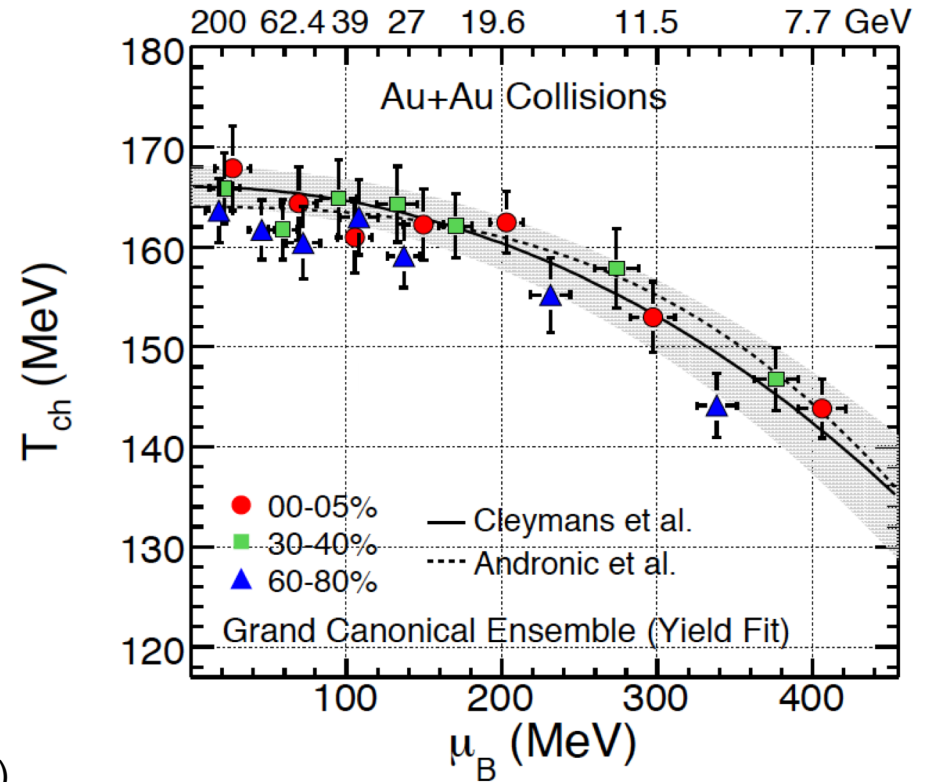
PBM&Johanna, Nature **448**, 302-309 (2007)

Can we find signature of QCD phase transition in heavy-ion collisions ?

- 1. Crossover**
- 2. Critical point**
- 3. 1st order phase transition (P.T.)**

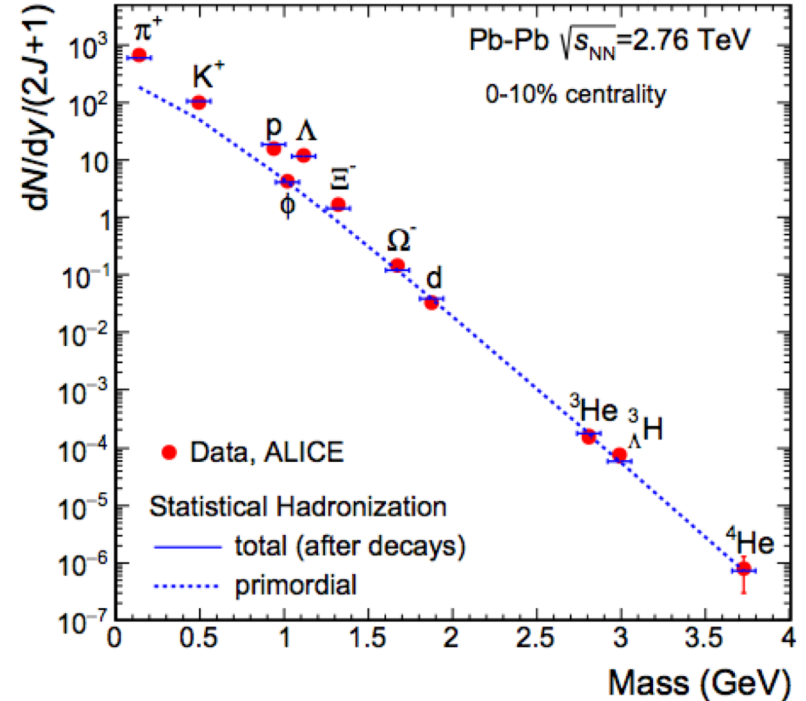
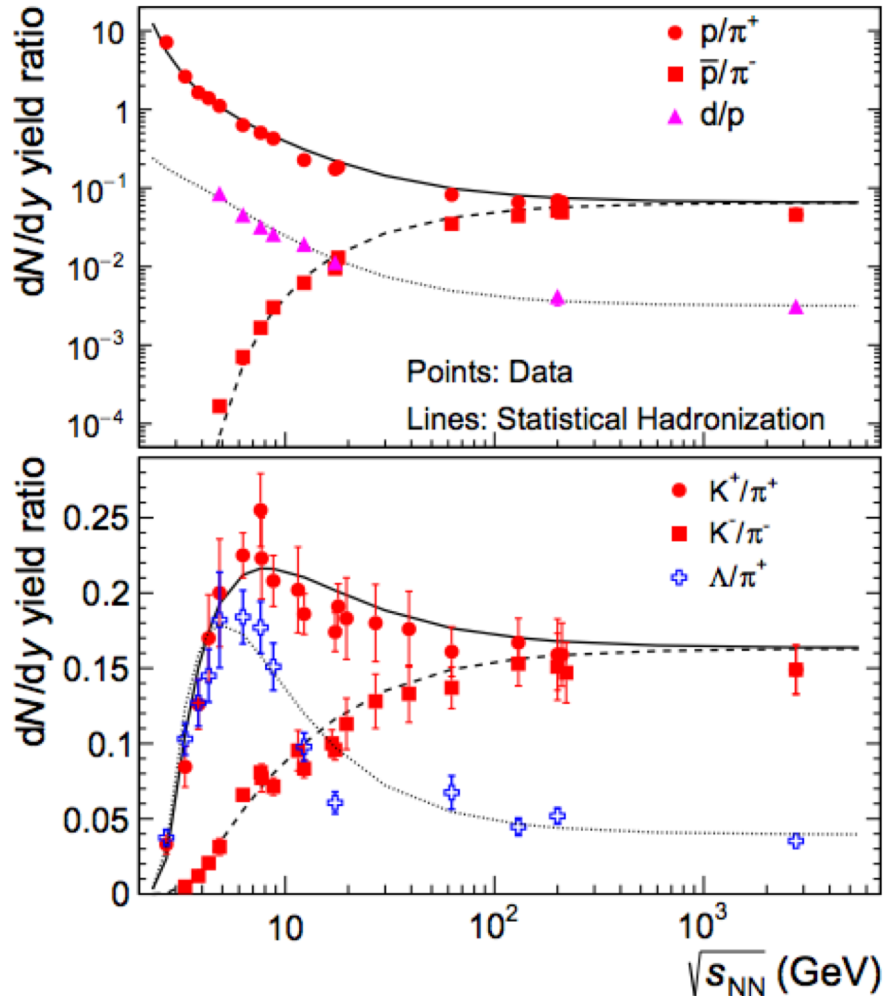
$\sqrt{s_{NN}}$ (GeV)	Events (10^6)	Year	* μ_B (MeV)	* T_{CH} (MeV)
200	350	2010	25	166
62.4	67	2010	73	165
39	39	2010	112	164
27	70	2011	156	162
19.6	36	2011	206	160
14.5	20	2014	264	156
11.5	12	2010	316	152
7.7	4	2010	422	140

*(μ_B, T_{CH}): J. Cleymans et al., PRC 73, 034905 (2006)



STAR: Phys. Rev. C 96, 044904 (2017).

Access the QCD phase diagram: vary collision energies/centralities.



Does this mean the phase structures (or EoS) at high baryon density are the same as small baryon density ?
All crossover ?

How to find 1st order P.T. or critical point signal ?
need to look at: **Fluctuation and Correlation**

- Event-by-Event fluctuations.
- Baryon clustering: Light nuclei production

A. Andronic, P. Braun-Munzinger, K. Redlich, J. Stachel,
Nature 561, 321 (2018).

PHYSICAL REVIEW D, VOLUME 60, 114028

20 years ago...

Signatures of the Tricritical Point in QCD

Event-by-event fluctuations in heavy ion collisions and the QCD critical point

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(Received 23 March 1999; published 10 November 1999)

Citation:939

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³ *Department of Physics and Astronomy, State University of New York, Stony Brook, NY 11794-3800*
(June 1, 1998; ITP-SB-98-39, MIT-CTP-2748, SUNY-NTG-98-17)

Several approaches to QCD with two *massless* quarks at finite temperature T and baryon chemical potential μ suggest the existence of a tricritical point on the boundary of the phase with spontaneously broken chiral symmetry. In QCD with *massive* quarks there is then a critical point at the end of a first order transition line. We discuss possible experimental signatures of this point, which provide information about its location and properties. We propose a combination of event-by-event observables, including suppressed fluctuations in T and μ and, simultaneously, enhanced fluctuations in the multiplicity of soft pions.

Citation:1113

Charged Particle Ratio Fluctuation as a Signal for QGP

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Lawrence Berkeley National Laboratory

Berkeley, CA 94720, USA

letter we argue that the event-by-event fluctuations of the ratio of the positively charged and negatively charged pions provides a signal of quark-gluon plasma. The fact that quarks and gluons are charged is ultimately responsible for this distinct signal.

Citation:515



Proposed experimental observables:

1. Pion multiplicity fluctuations.
2. Mean p_T fluctuations.
3. Particle ratio fluctuations

Scan T - μ Plane: Non-monotonic variation with energy, centrality or rapidity.

1. Higher order cumulants/moments: describe the shape of distributions and quantify fluctuations. (sensitive to the correlation length (ξ))

$$\langle \delta N \rangle = N - \langle N \rangle$$

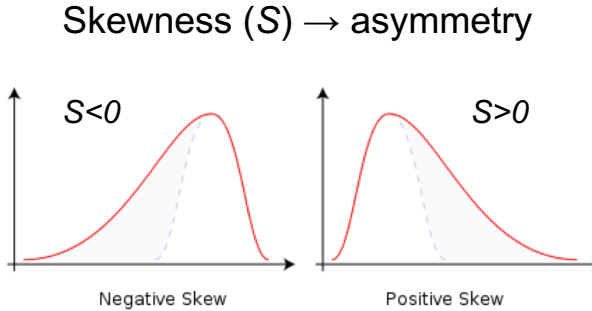
$$C_1 = M = \langle N \rangle$$

$$C_2 = \sigma^2 = \langle (\delta N)^2 \rangle$$

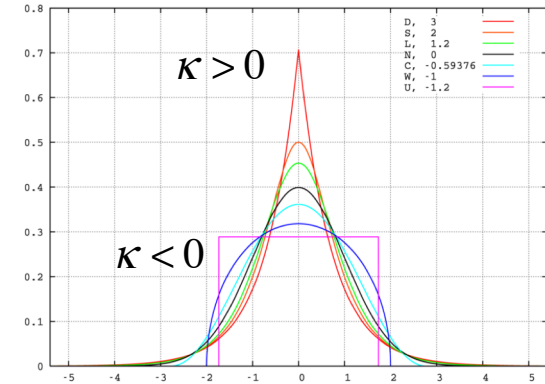
$$C_3 = S\sigma^3 = \langle (\delta N)^3 \rangle$$

$$C_4 = \kappa\sigma^4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2$$

$$\langle (\delta N)^3 \rangle_c \approx \xi^{4.5}, \quad \langle (\delta N)^4 \rangle_c \approx \xi^7$$



Kurtosis (κ) → Sharpness

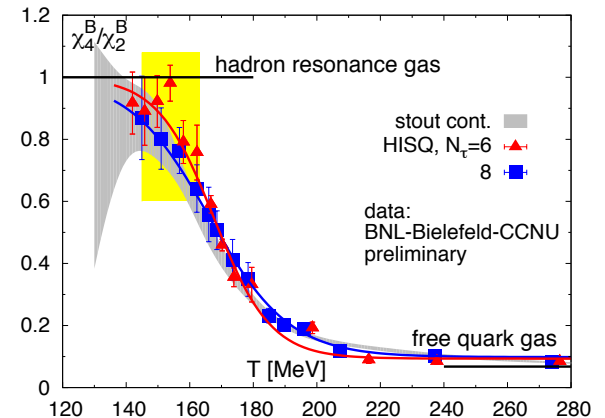


M. A. Stephanov, *Phys. Rev. Lett.* 102, 032301 (2009).
M. Asakawa, S. Ejiri and M. Kitazawa, *Phys. Rev. Lett.* 103, 262301 (2009).
M. A. Stephanov, *Phys. Rev. Lett.* 107, 052301 (2011).

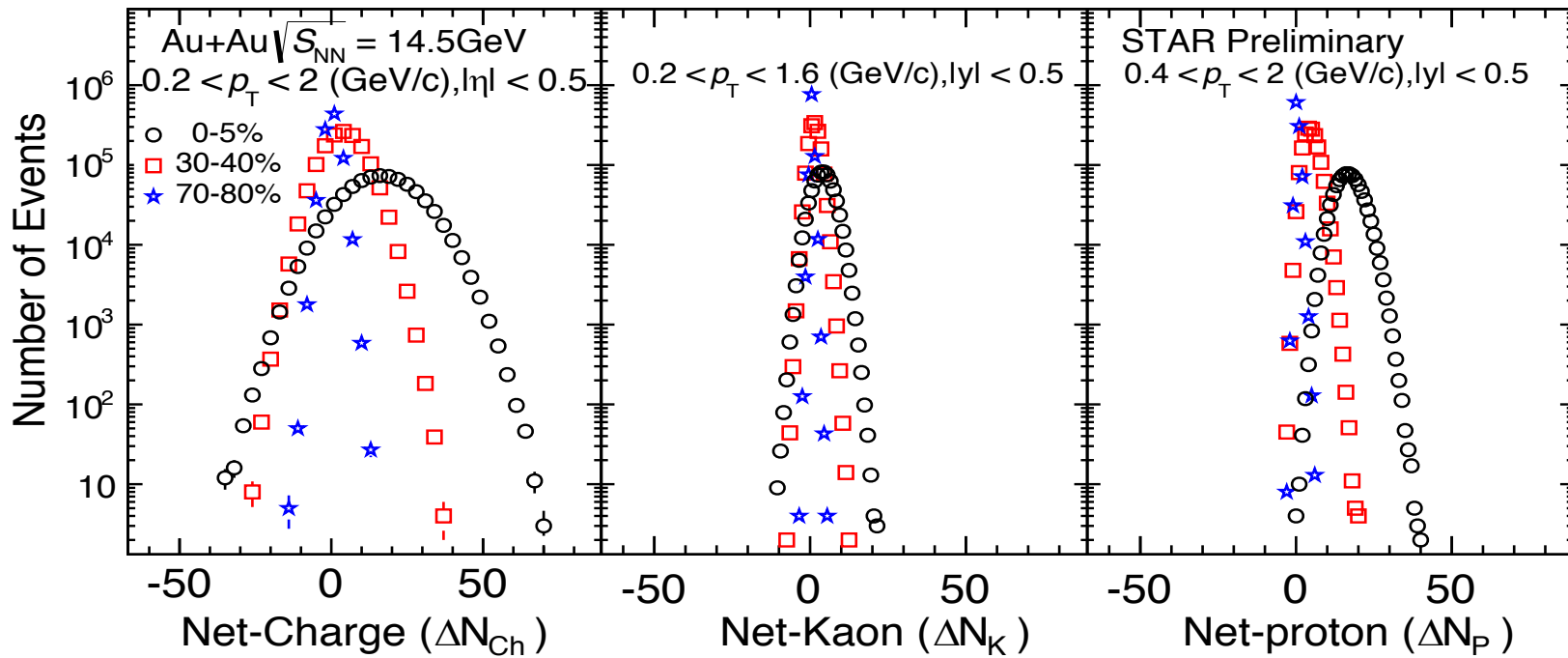
2. Direct connect to the susceptibility of the system.

$$\frac{\chi_q^4}{\chi_q^2} = \kappa\sigma^2 = \frac{C_{4,q}}{C_{2,q}} \quad \frac{\chi_q^3}{\chi_q^2} = S\sigma = \frac{C_{3,q}}{C_{2,q}}$$

$$\chi_q^{(n)} = \frac{1}{VT^3} \times C_{n,q} = \frac{\partial^n (p/T^4)}{\partial (\mu_q)^n}, q = B, Q, S$$



S. Ejiri et al, *Phys. Lett. B* 633 (2006) 275. Cheng et al, *PRD* (2009) 074505. B. Friman et al., *EPJC* 71 (2011) 1694.
F. Karsch and K. Redlich, *PLB* 695, 136 (2011). S. Gupta, et al., *Science*, 332, 1525(2012). A. Bazavov et al., *PRL* 109, 192302(12) // S. Borsanyi et al., *PRL* 111, 062005(13) // P. Alba et al., *arXiv:1403.4903*



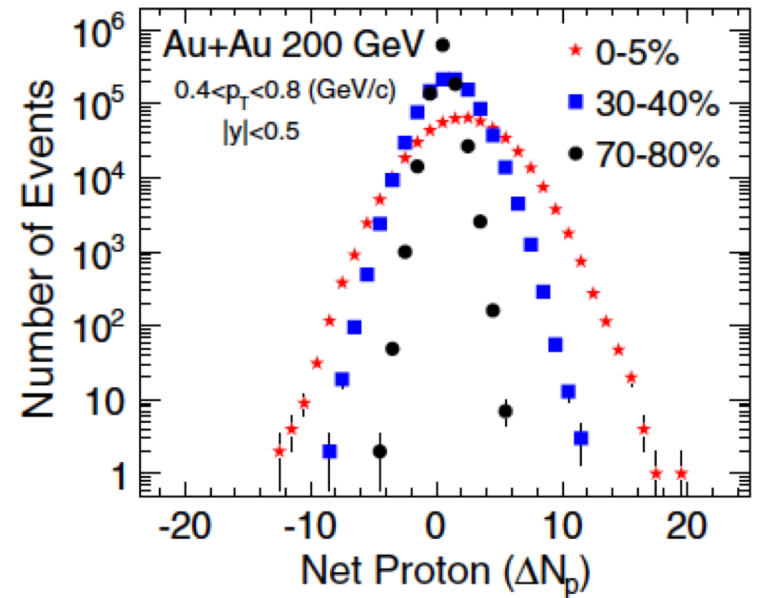
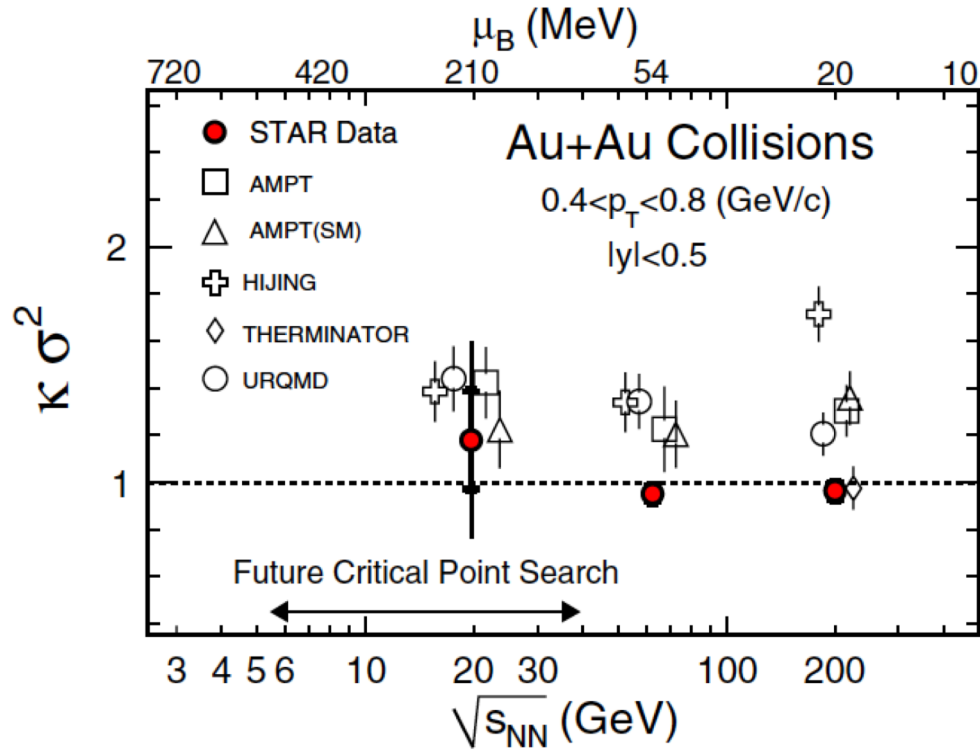
Effects needed to be addressed to get final moments/cumulants:

1. Statistical errors estimation : Delta theorem or bootstrap
2. Avoid auto-correlation effects: New centrality definition.
3. Suppress volume fluctuation: Centrality bin width correction
4. Finite detector efficiency correction (binomial response func.)

X.Luo, et al. J. Phys. G39, 025008 (2012); A. Bzdak and V. Koch, PRC86, 044904 (2012); X.Luo, et al. J. Phys. G40, 105104(2013); X.Luo, Phys. Rev. C 91, 034907 (2015); A. Bzdak and V. Koch, PRC91, 027901 (2015). T. Nonaka et al., PRC95, 064912 (2017). M. Kitazawa and X. Luo, PRC96, 024910 (2017). X. Luo and N. Xu, Nucl. Sci. Tech. 28, 112 (2017).

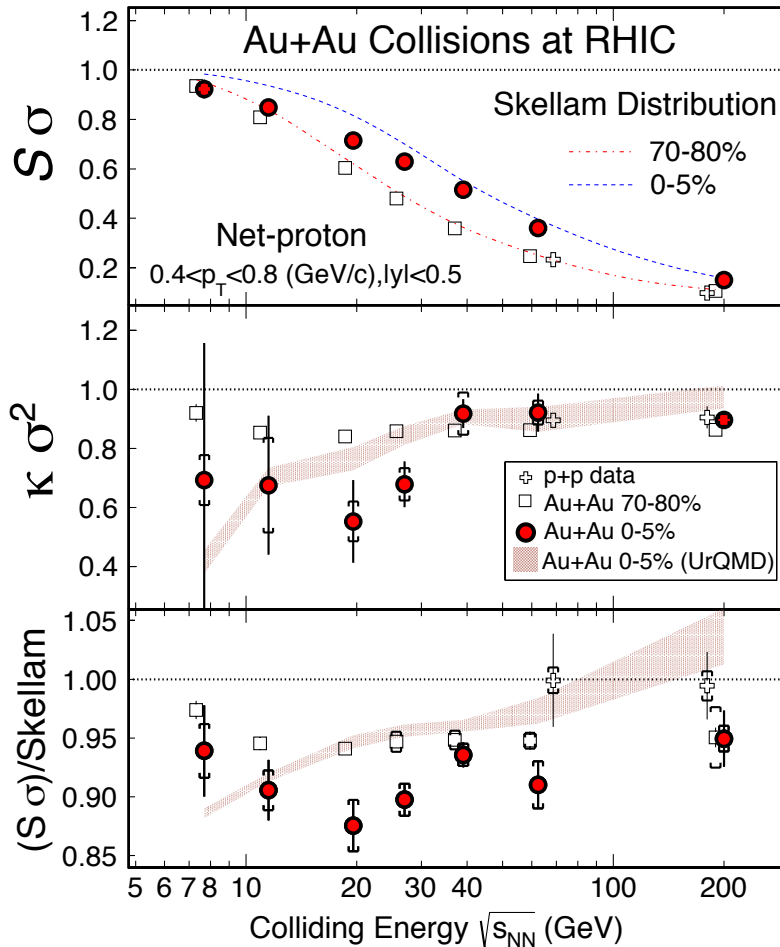
Only TPC used for proton/anti-proton PID.

p_T coverage : $0.4 < p_T < 0.8$ GeV/c



STAR: Phys. Rev. Lett. 105, 022302 (2010).

No critical point signature is observed at $\mu_B < 210$ MeV.



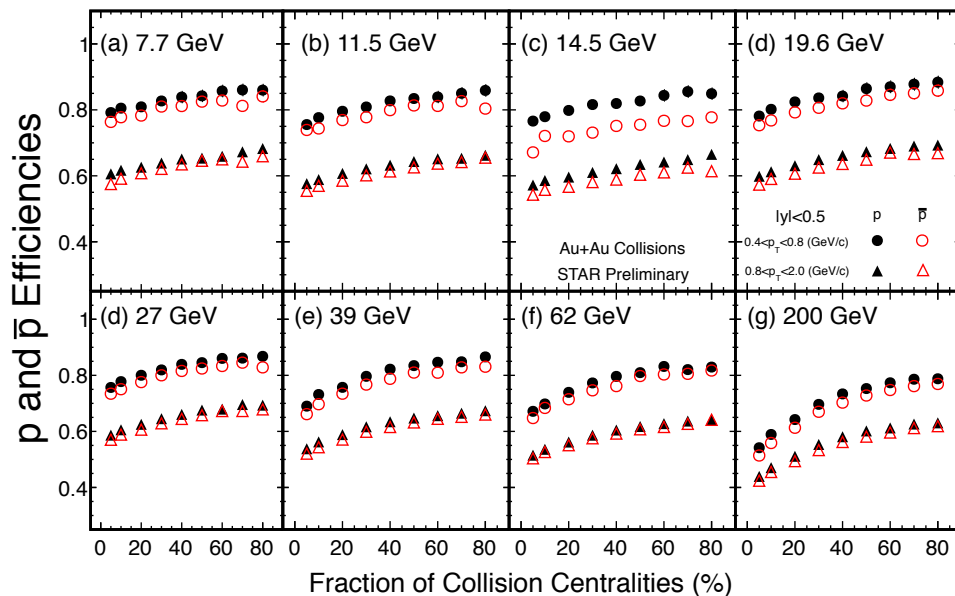
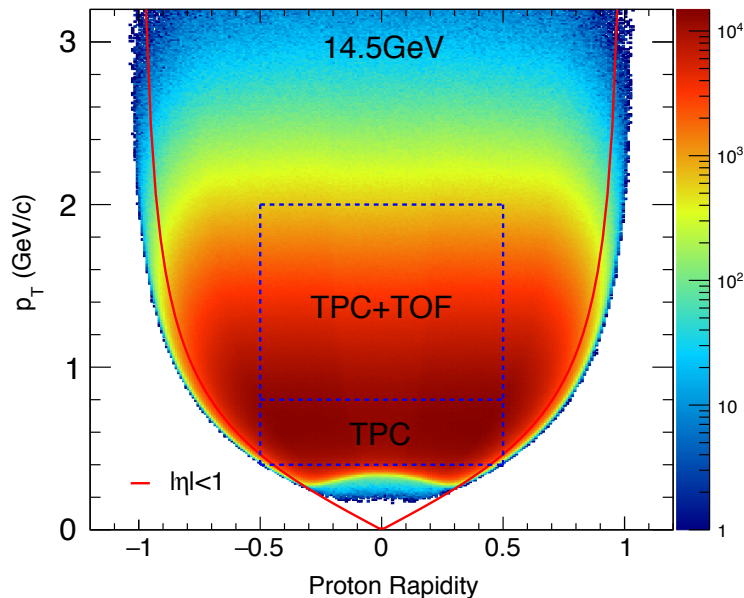
p_T coverage : $0.4 < p_T < 0.8 \text{ GeV/c}$

1. Deviation below Poisson baseline (unity).
2. A minimum shows at 19.6 GeV.

STAR: Phy. Rev. Lett. 112, 032302 (2014).

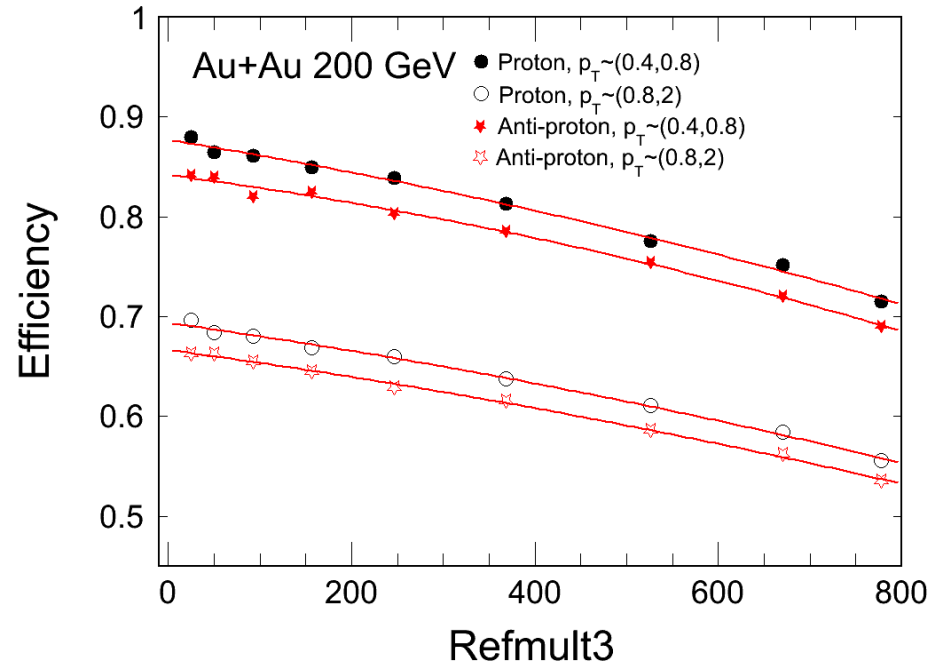
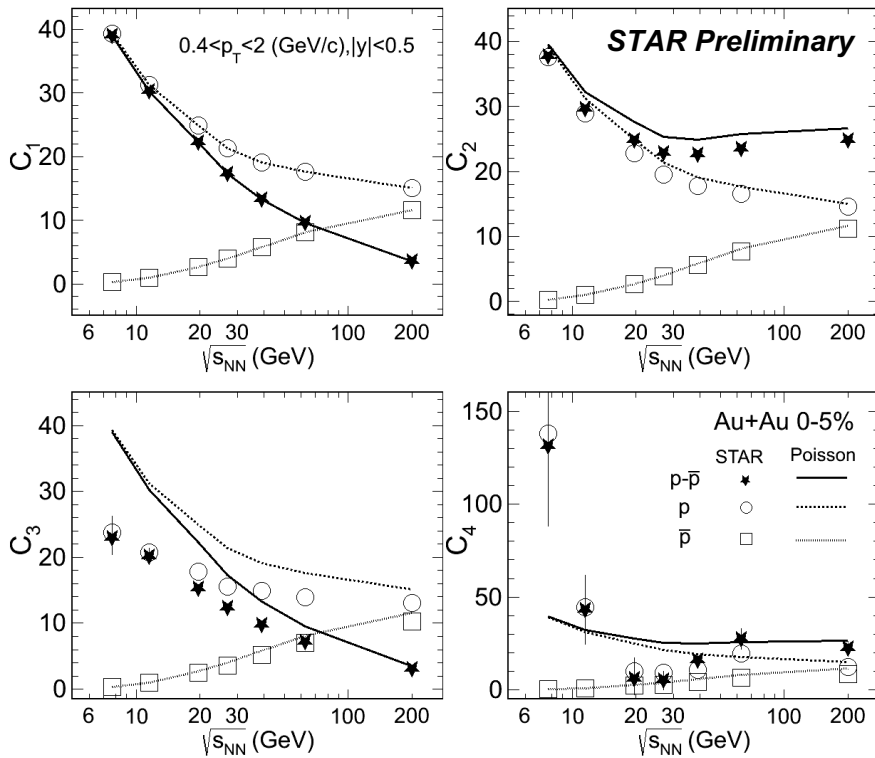
TOF PID can extend the phase space coverage

$|y| < 0.5, 0.4 < p_T < 0.8$ (TPC PID only)
 $0.8 < p_T < 2$ (TPC+TOF PID)



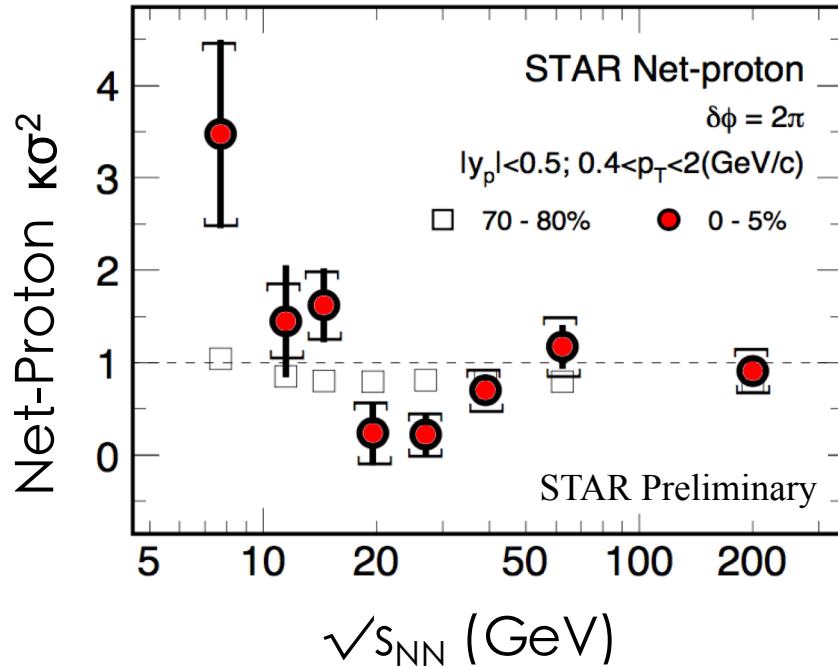
Doubled the accepted number of proton/anti-proton

- Sufficiently large acceptance is important for fluctuation analysis and critical point search.
- Efficiency : Proton > Anti-proton, Low p_T > High p_T , low energy > High Energy, Peripheral > Central



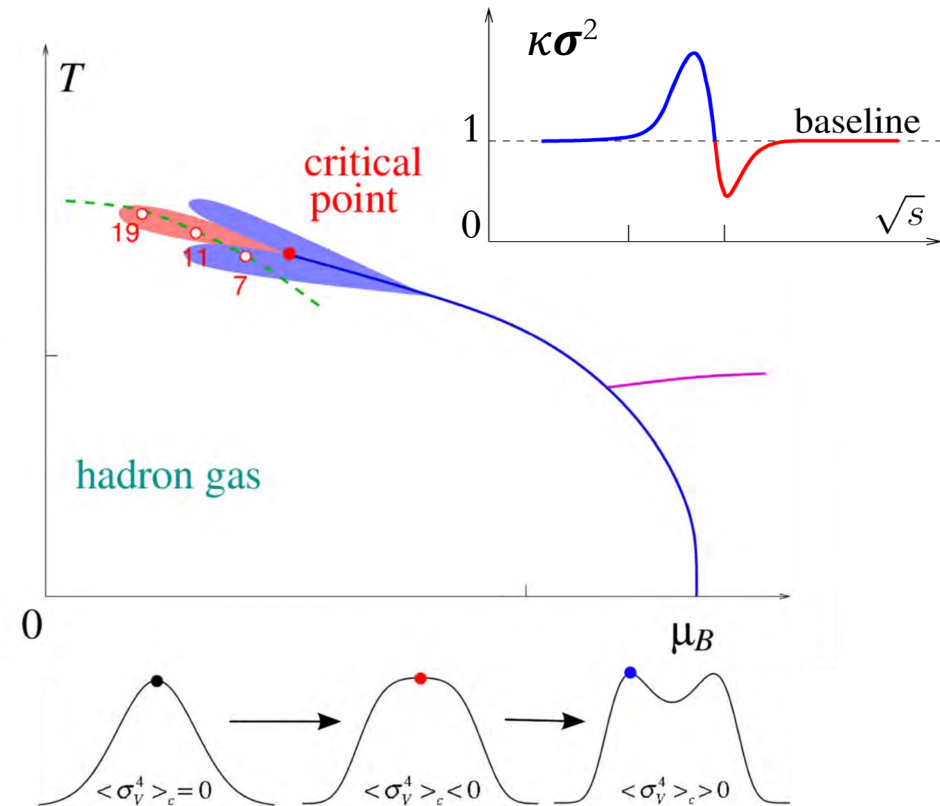
- Net-proton number is dominated by protons at low energies and increases when energy decreases. (Interplay between baryon stopping and pair production)
- The higher the order of cumulants the larger deviations from Poisson expectations for net-proton and proton.

Experimental Measure



STAR: PoS CPOD2014 (2015) 019.

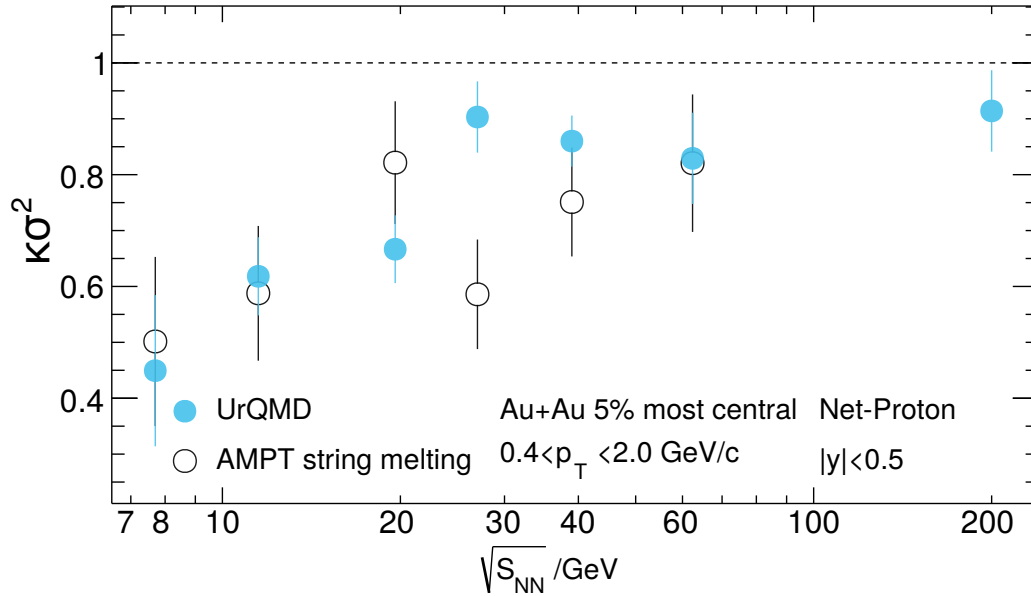
Theoretical calculations



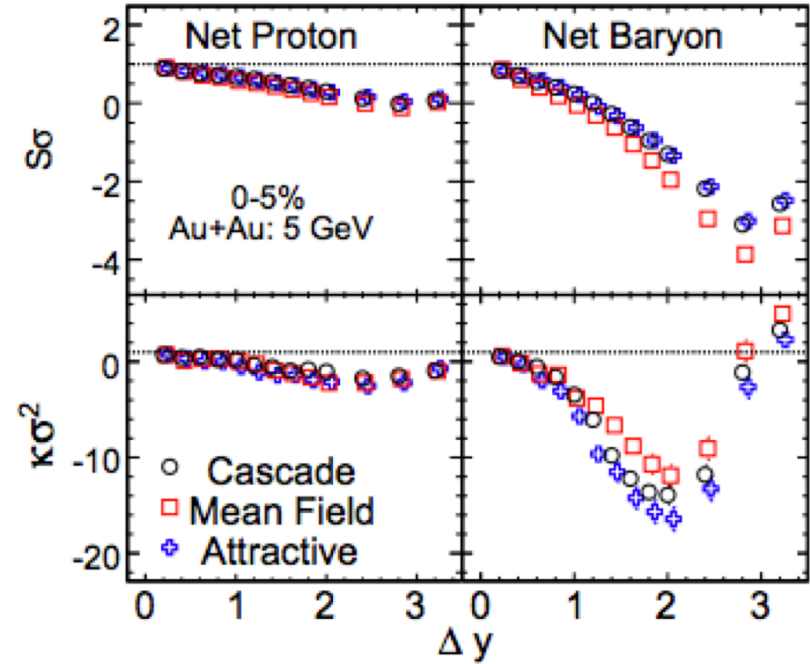
M. Stephanov, PRL107, 052301(2011)
 J. Phys. G: 38, 124147 (2011).

Observe non-monotonic energy dependence in 0-5% most central Au+Au collisions. A hint of entering the critical region.

UrQMD and AMPT models



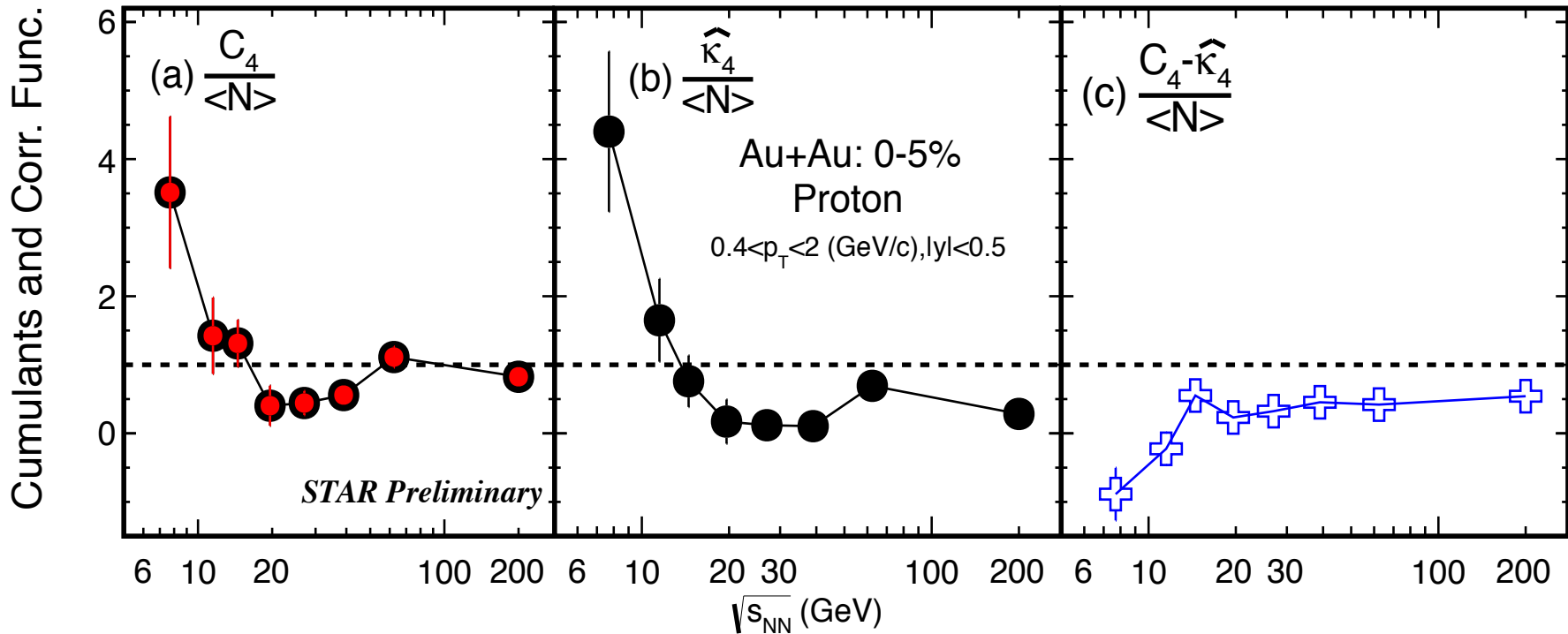
JAM model



At $\sqrt{s_{NN}} \leq 10$ GeV: Data: $\kappa\sigma^2 > 1$ Model: $\kappa\sigma^2 < 1$

➤ Model simulation : *suppress the net-proton fluctuations.*

Z. Feckova, et al., PRC92, 064908(2015). J. Xu, et al., PRC94, 024901(2016). X. Luo et al., NPA931, 808(14), P.K. Netrakanti et al. 1405.4617, NPA947, 248(2016), P. Garg et al. PLB 726, 691(2013). S. He, et al., PLB762, 296 (2016). S. He, X. Luo, PLB 774, 623 (2017).



Four-particle correlation dominated the non-monotonic behavior observed in fourth order net-proton fluctuations.

$$C_2 = \langle N \rangle + \hat{\kappa}_2$$

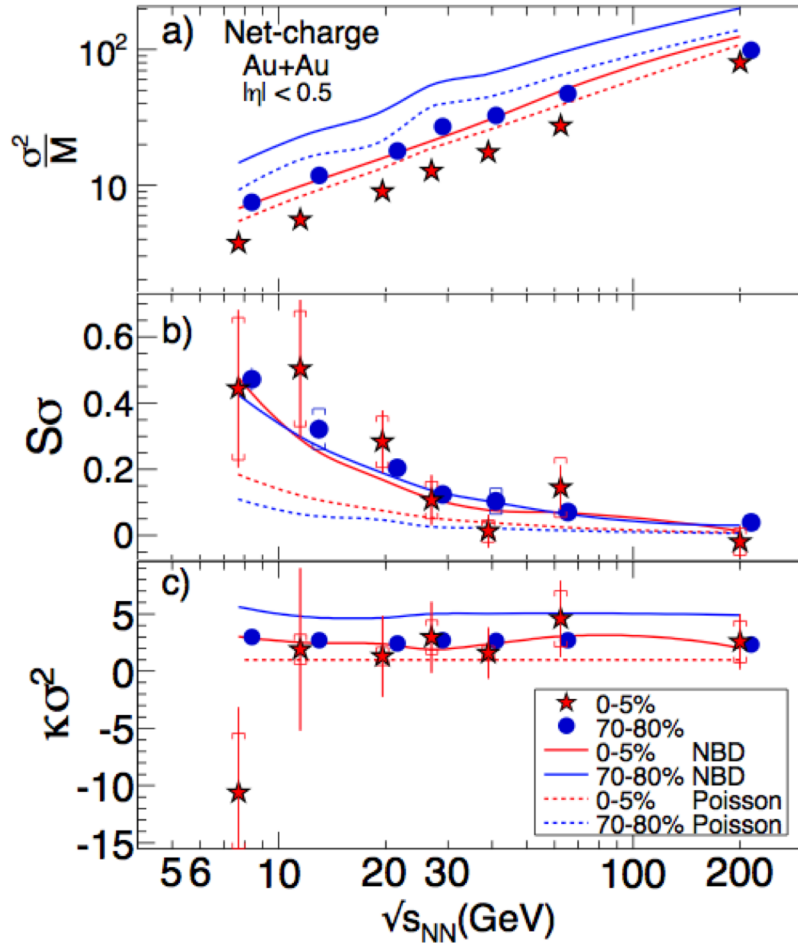
$$C_3 = \langle N \rangle + 3\hat{\kappa}_2 + \hat{\kappa}_3$$

$$C_4 = \langle N \rangle + 7\hat{\kappa}_2 + 6\hat{\kappa}_3 + \hat{\kappa}_4$$

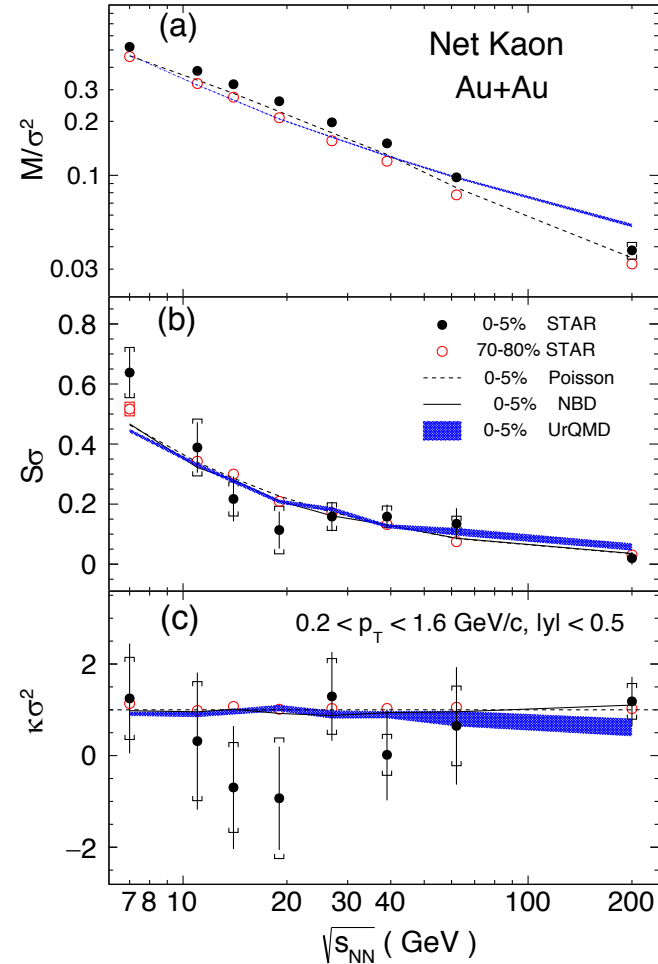
$\hat{\kappa}_2, \hat{\kappa}_3, \hat{\kappa}_4$: 2,3,4-particle correlation function

Large statistical errors, need more data

$$error(\kappa\sigma^2) \propto \frac{\sigma^2}{\varepsilon^2} \frac{1}{\sqrt{N_{evts}}}$$



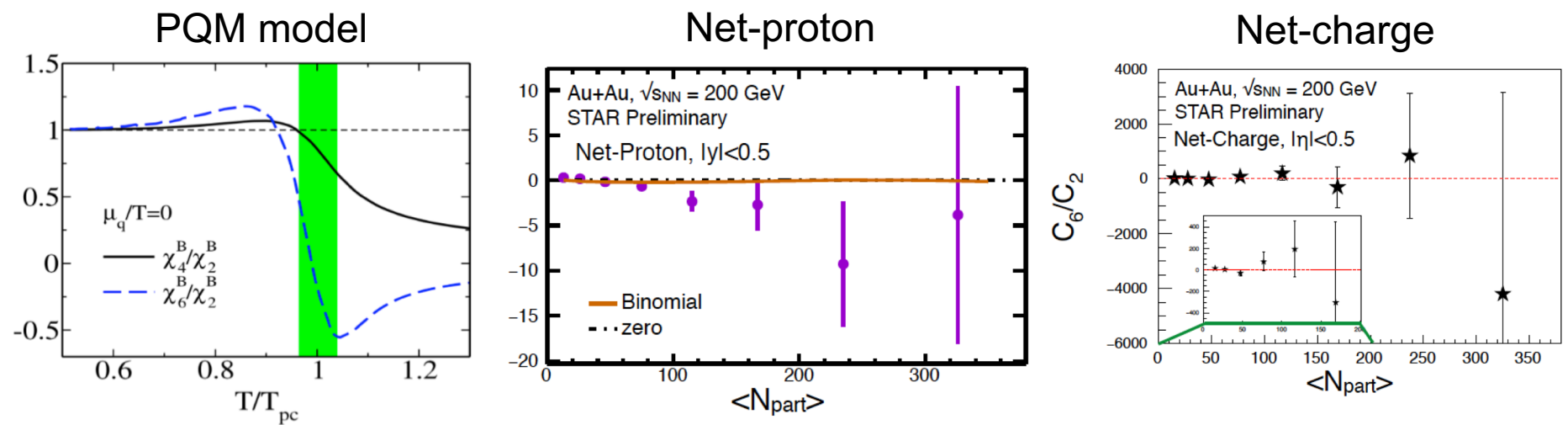
STAR : Phys. Rev. Lett. 113 092301 (2014).



STAR : Phys. Lett. B 785, 551 (2018).

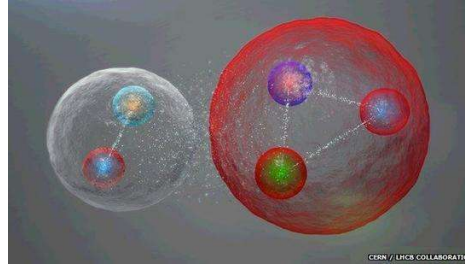
- The sixth-order cumulants of net-charge and net-baryon distributions are predicted to be negative if the chemical freeze-out is close enough to the phase transition.
- In Au+Au collisions at 200 GeV, negative values are observed in net-proton C_6/C_2 systematically from peripheral to central. Results of net-charge C_6/C_2 are consistent with zero within large statistical errors.

$$error(C_n / C_2) \propto \frac{\sigma^{n-2}}{\sqrt{N_{evts}}}$$

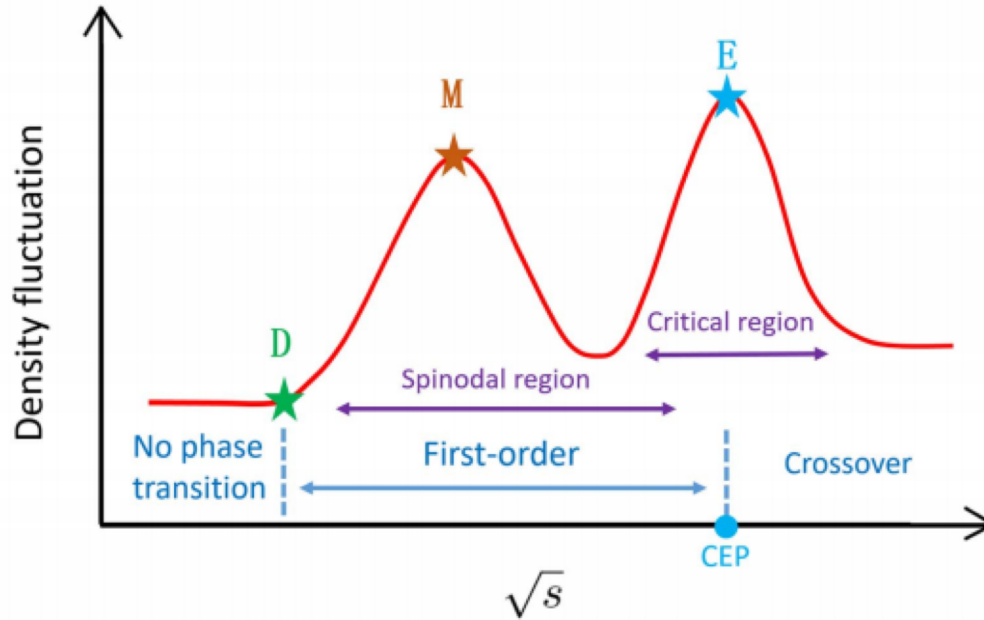


B. Friman et al., Eur. Phys. J. C 71 (2011) 1694

Near CP or 1st order phase transition, baryon density fluctuation become large.



Light nuclei production
(Baryon Clustering)



Coalescence + nucleon density flu.

$$N_d = \frac{3}{2^{1/2}} \left(\frac{2\pi}{m_0 T_{\text{eff}}} \right)^{3/2} N_p \langle n \rangle (1 + \alpha \Delta n),$$

$$N_{3H} = \frac{3^{3/2}}{4} \left(\frac{2\pi}{m_0 T_{\text{eff}}} \right)^3 N_p \langle n \rangle^2 [1 + (1 + 2\alpha) \Delta n],$$

$$N_t \cdot N_p / N_d^2 \approx g(1 + \Delta n)$$

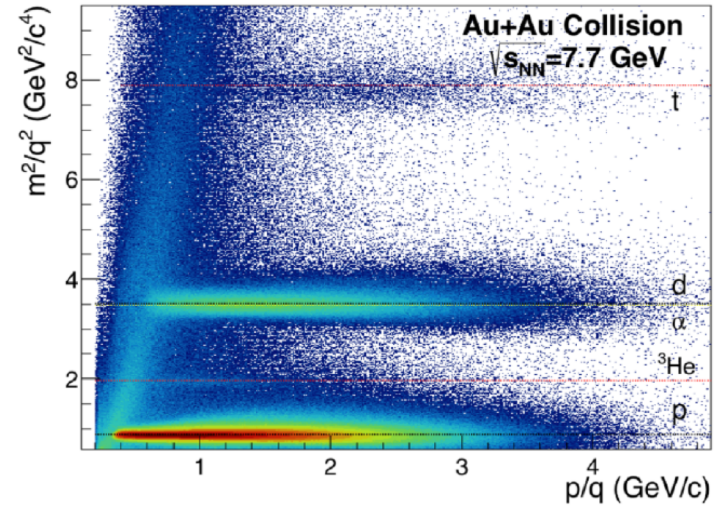
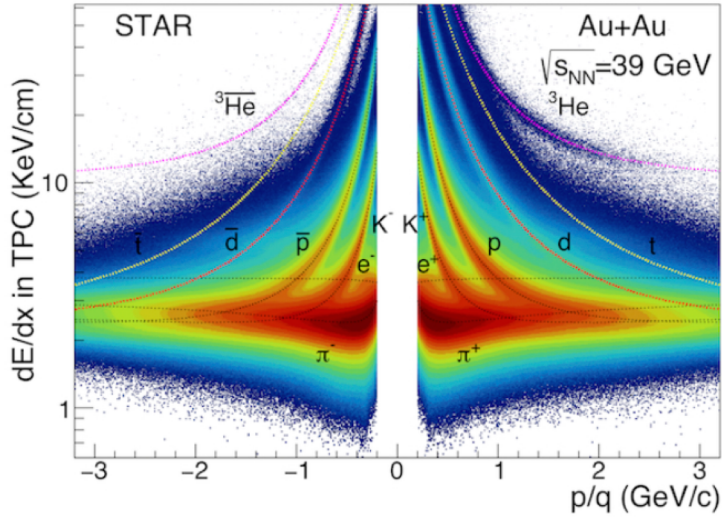
Neutron density flu.

$$\Delta n = \langle (\delta n)^2 \rangle / \langle n \rangle^2$$

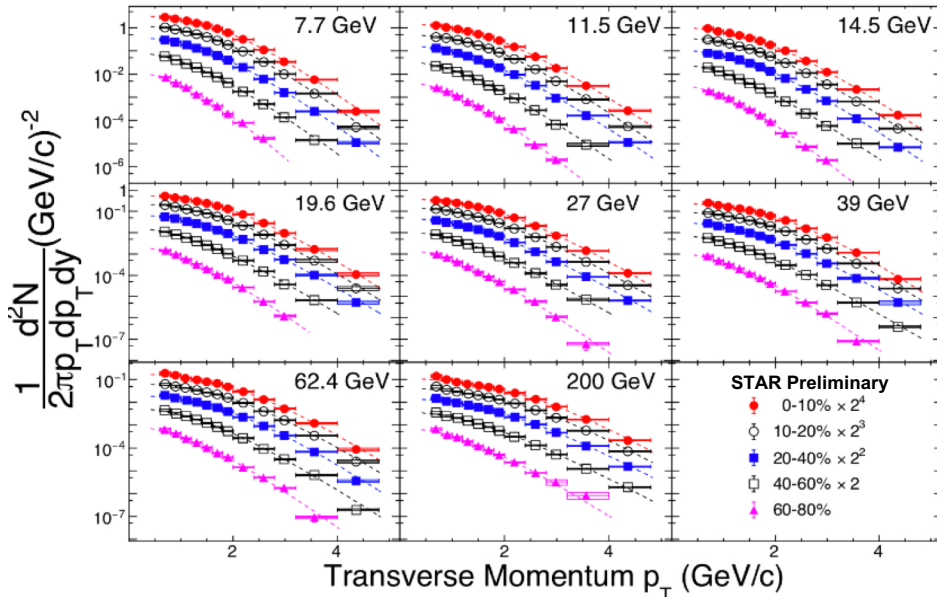
K. J. Sun, L. W. Chen, C. M. Ko, Z. Xu, Phys. Lett. B774, 103 (2017).

K. J. Sun, L. W. Chen, C. M. Ko, J. Pu, Z. Xu, Phys. Lett. B781, 499 (2018).

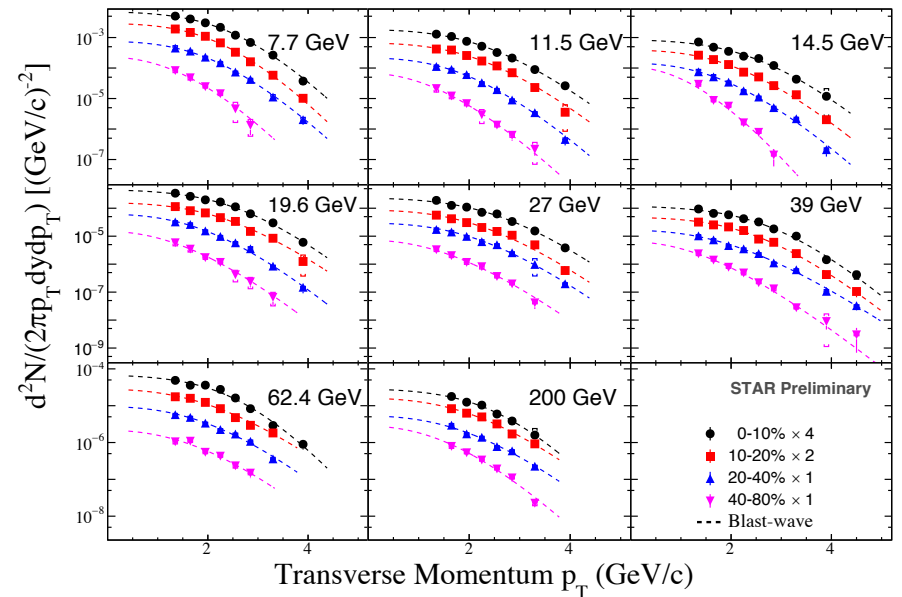
Edward Shuryak and Juan M. Torres-Rincon, arXiv:1805.04444



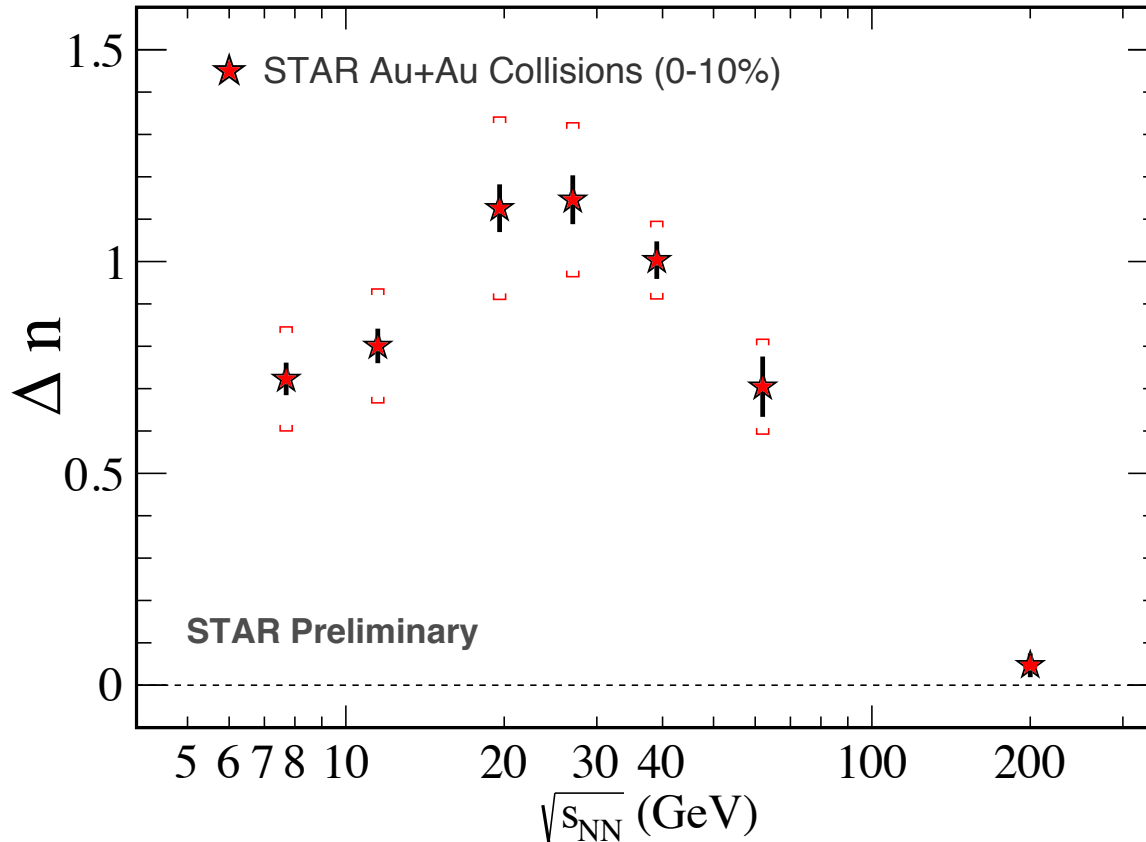
Deuteron from Au+Au Collision



Triton from Au+Au Collision



STAR, QM2018



$$\Delta n = \langle (\delta n)^2 \rangle / \langle n \rangle^2$$

$$N_t \cdot N_p / N_d^2 \approx g(1 + \Delta n)$$

$$g=0.29.$$

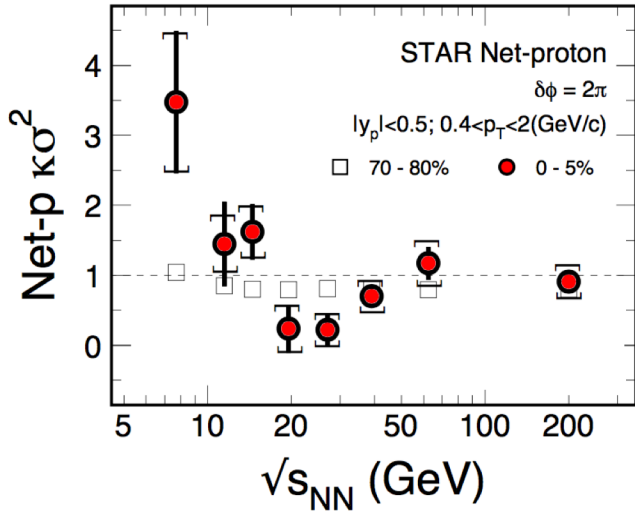
N_t : Triton yield

N_d : Deuteron yield

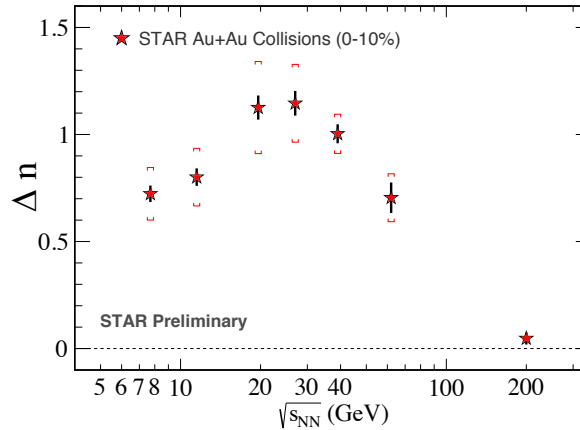
N_p : Proton yield

Observe non-monotonic energy dependence of neutron density fluctuation in 0-10% central Au+Au collisions with a peak around 20-30 GeV.

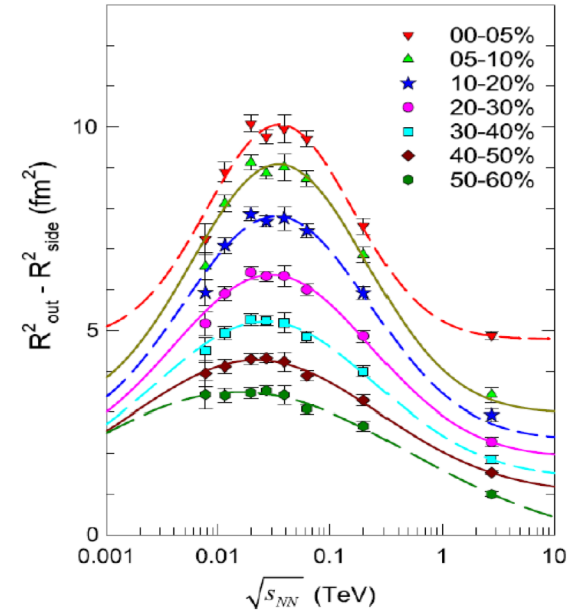
Net-p flu.



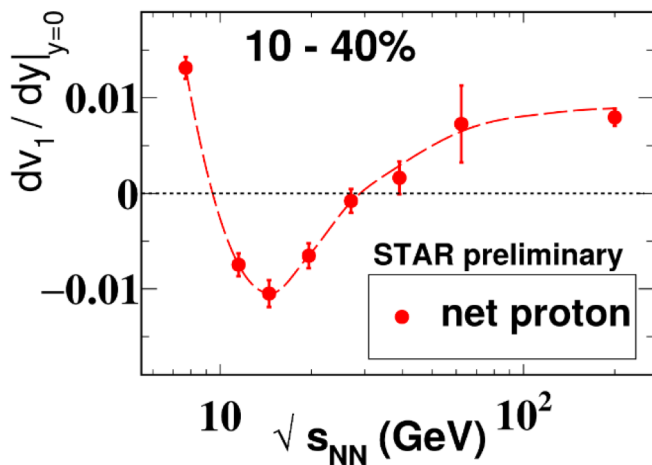
Neutron density flu.



HBT radii



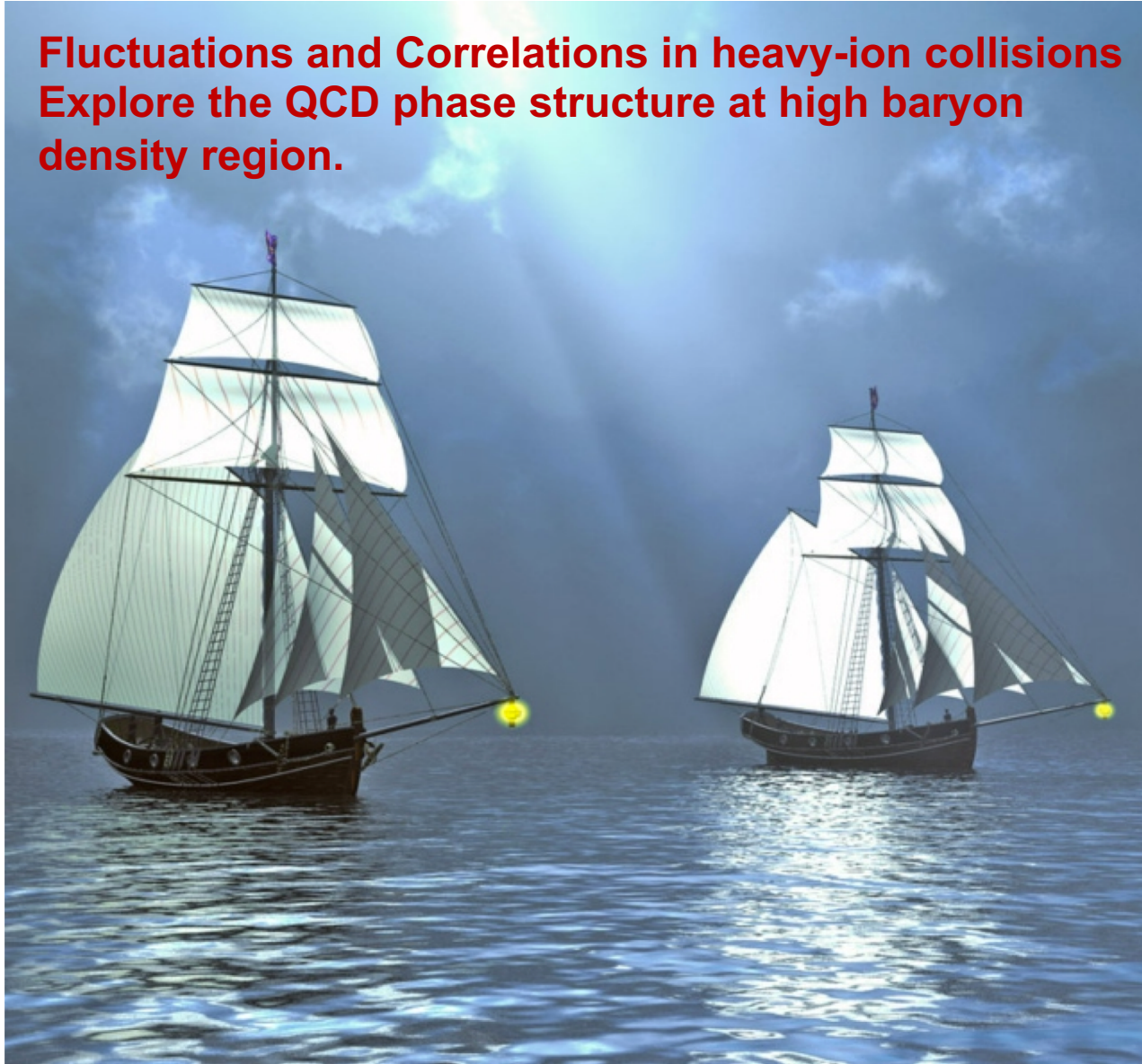
Slope of Directed flow vs y.



Peak and/or dip structures observed at common energy ranges : 20-30 GeV !!

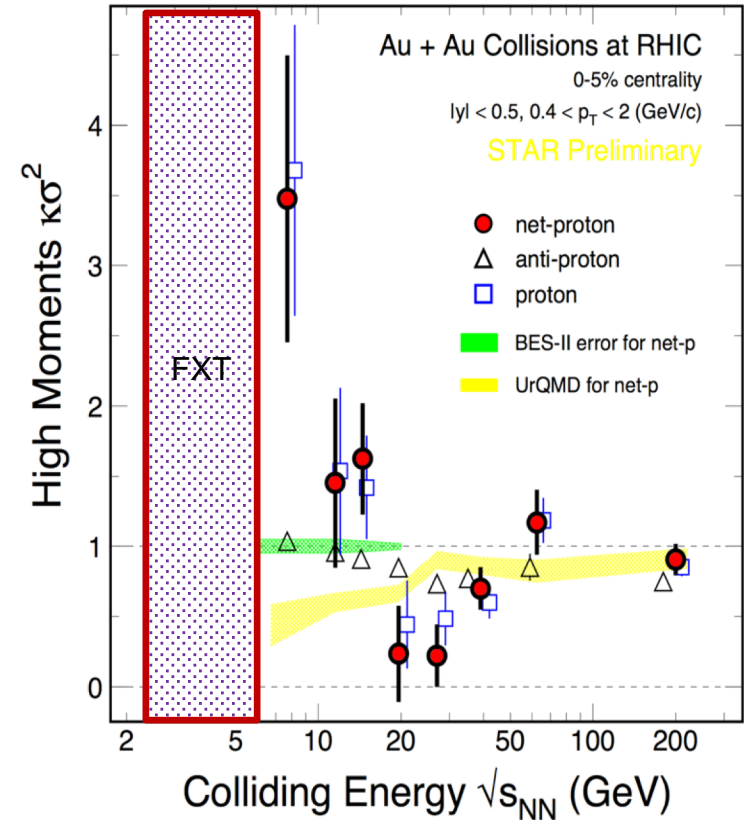
Hard to believe those are driven by different physics.

**Fluctuations and Correlations in heavy-ion collisions
Explore the QCD phase structure at high baryon
density region.**



10-20 times higher statistics than BES-I

$\sqrt{s_{NN}}$ (GeV)	Events (10^6)	BES II / BES I
200	350	2010
62.4	67	2010
54.4	1200	2017
39	39	2010
27	70	2011
19.6	400 / 36	2019-20 / 2011
14.5	300 / 20	2019-20 / 2014
11.5	230 / 12	2019-20 / 2010
9.2	160 / 0.3	2019-20 / 2008
7.7	100 / 4	2019-20 / 2010



- BES-II : Precise mapping the QCD phase diagram **$200 < \mu_B < 420$ MeV**
- FIX-target mode : $\sqrt{s_{NN}} = 3-7.7$ GeV (2018-2020)

- Non-monotonic energy dependence is observed in net-proton C_4/C_2 at most central Au+Au collisions. *A hint of entering the critical region. Need to confirm with more statistics and lower energy data.*
- Within current uncertainties, net-charge and net-kaon fluctuations show flat energy dependent. *Need more statistics.*
- Observe non-monotonic energy dependence of neutron density fluctuations in 0-10% central Au+Au collisions with a peak around 20-30 GeV.
- Study the QCD phase structure with **high precision**: BES-II at RHIC (2019-2020, both collider and fixed-target mode).

Happy Birthday to Edward

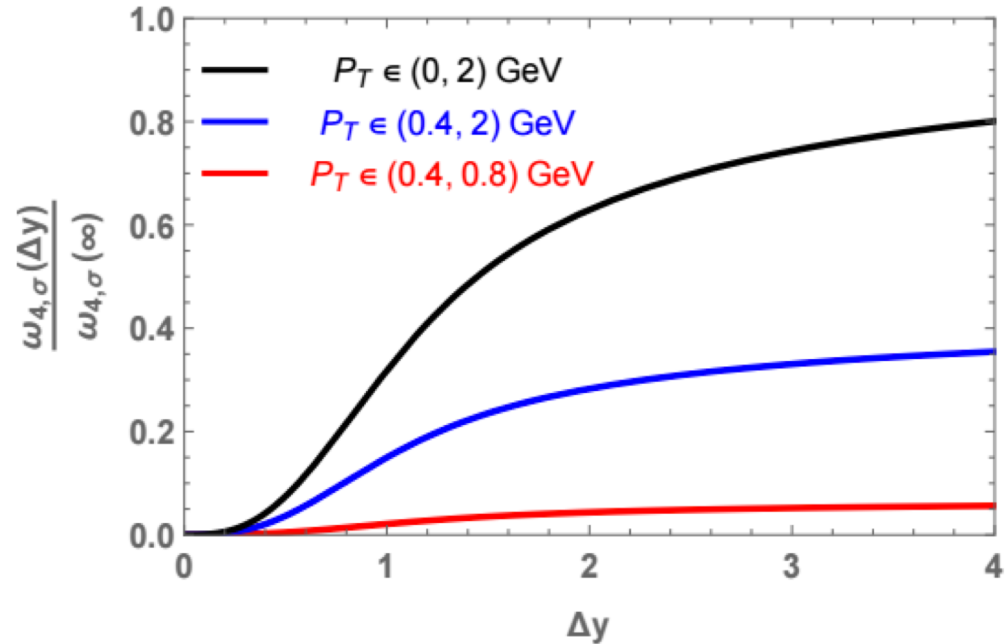
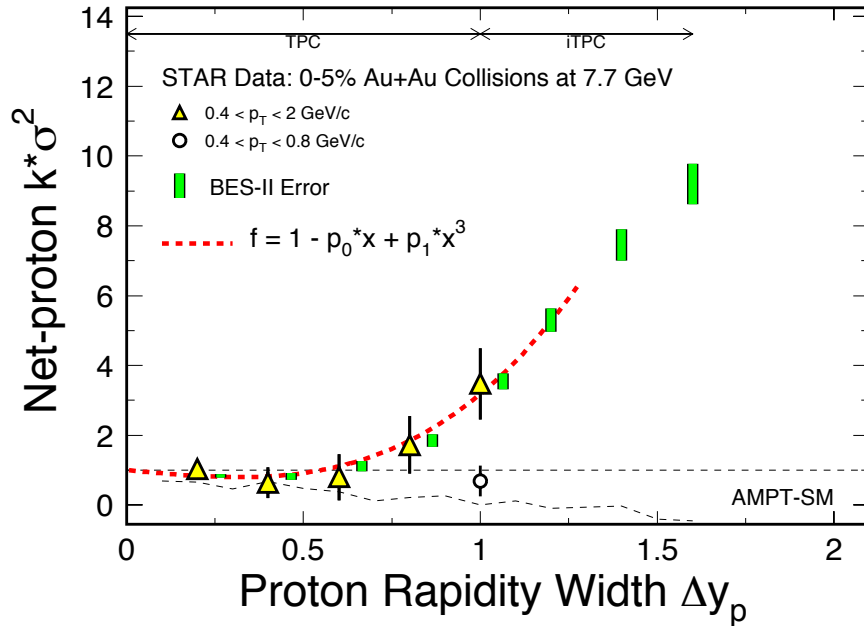
Thank you!

Back up slides

STAR Data

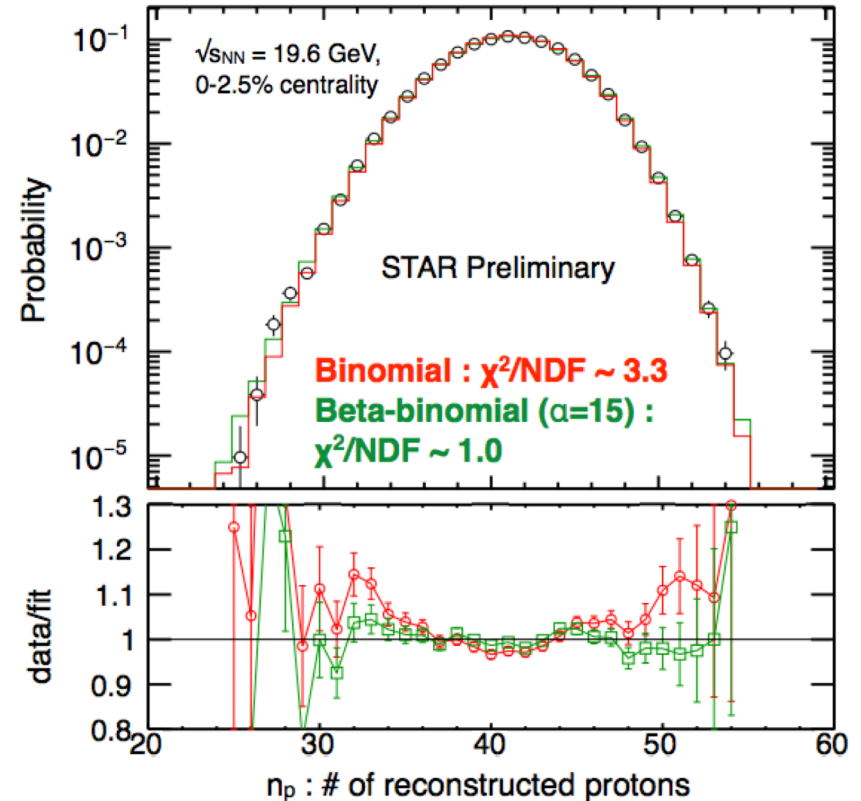
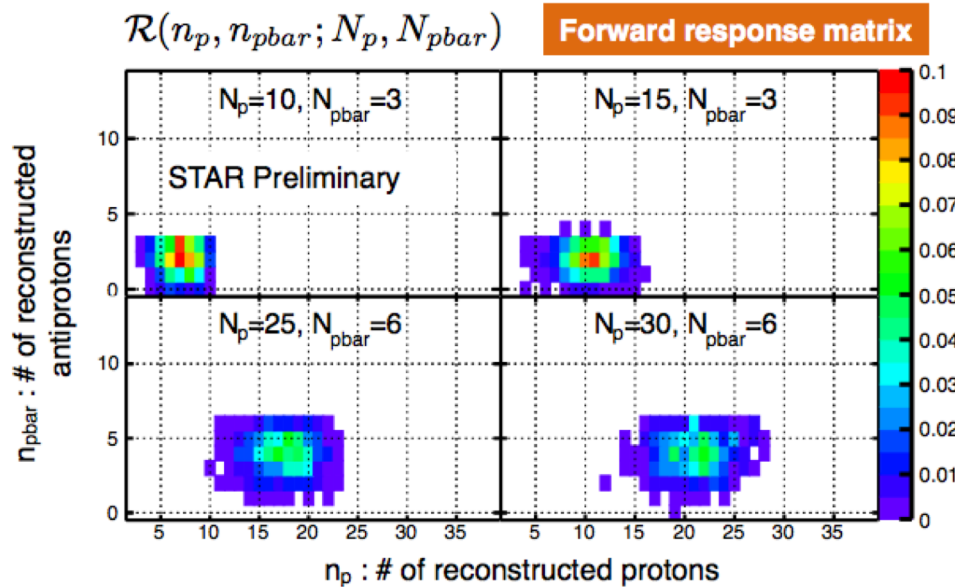
Acceptance dep. near CP

Model

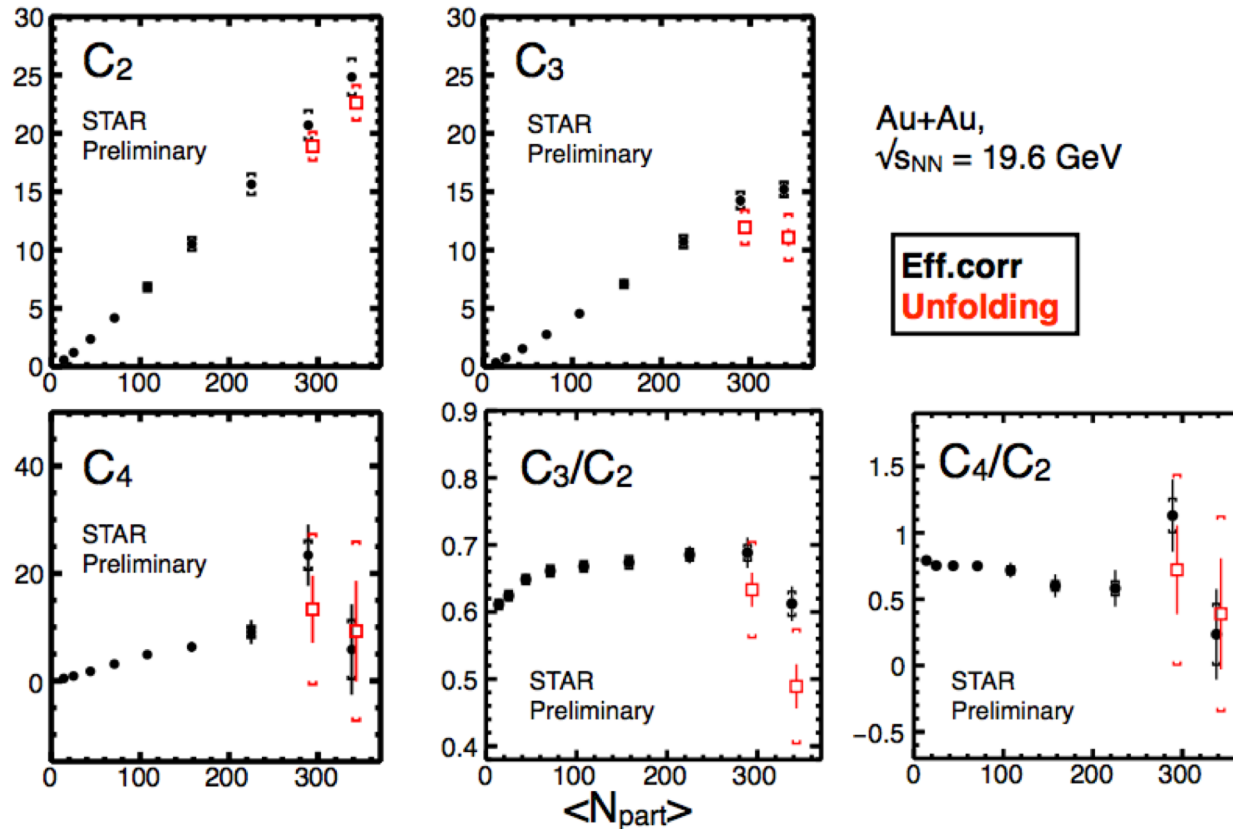


Signals can be enhanced by enlarging the acceptance.

B. Ling, M. Stephanov, *Phys. Rev. C* 93, 034915 (2016).
 A. Bzdak, V. Koch, *Phys. Rev. C* 95, 054906 (2017)
 M. Kitazawa, X. Luo, *PRC* 96, 024910 (2017).



- We perform simulation to construct response matrix by embedding MC tracks of protons and antiprotons into real events, which will be used for unfolding.
- When embed 60 protons (an extreme case), the response matrix is close to beta-binomial, which is wider than binomial (right plot). The deviation from binomial would depend on the # of embedded protons and antiprotons.



- For unfolding, 2.5% centrality width averaging has been done.
- Systematic suppression is observed for C_2 and C_3 with respect to the results of efficiency correction assuming binomial efficiencies.
- C_4 , C_3/C_2 and C_4/C_2 are consistent within large systematic uncertainties limited by embedding samples.