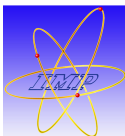


Exploring Pion and Nucleon Structure Through Basis Light Front Quantization

Chandan Mondal



Institute of Modern Physics, CAS

With: S. Xu, J. Lan, S. Nair, X. Zhao (IMP), S. Jia, Y. Li, J. P. Vary (ISU)



Palaiseau, September 18, 2019

Overview

- 1 Basis Light-Front Quantization (BLFQ)
- 2 Application to light mesons
- 3 Application to nucleon
- 4 Conclusions & Outlook

Basis Light-Front Quantization (BLFQ)

Vay et. al. PRC 81 (2010)

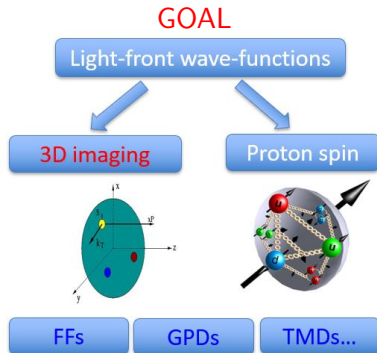
Solve many-body bound state problems in quantum field theories

$$P^-|\beta\rangle = P_\beta^-|\beta\rangle$$

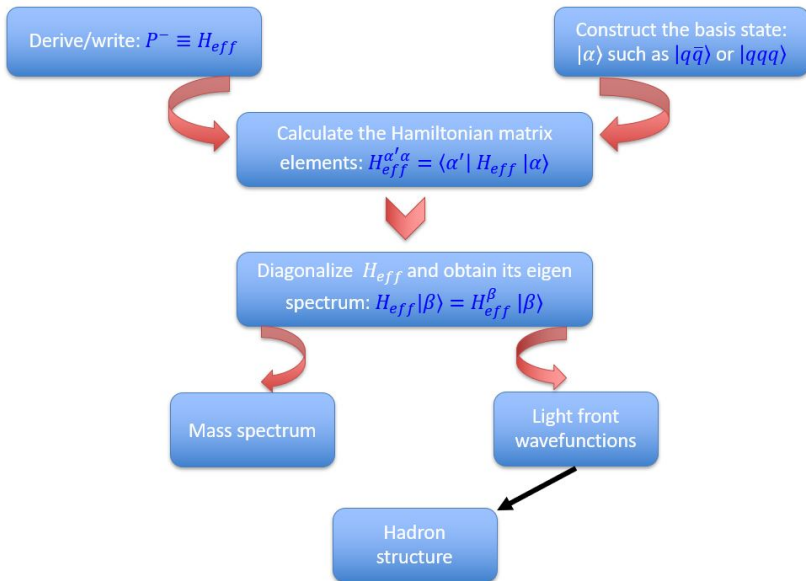
- P^- : light-front Hamiltonian
- $|\beta\rangle$ mass eigenstates
- P_β^- eigenvalue (**light-front energy**) for eigenstate $|\beta\rangle$
- first-principles / effective Hamiltonian as input
- Evaluate observables

$$O \sim \langle\beta|\hat{O}|\beta\rangle$$

- direct access to wave function of bound states



General Procedure for BLFQ



Applications of BLFQ

QCD systems

- **Heavy mesons:** spectrum, decay constant, elastic form factor, radii, radiative transitions, distribution amplitude, GPDs

—Y Li, G Chen, X Zhao, P Maris, J Vary, L Adhikari, M Li, S. Tang, A El-Hady (2016 - 2019)

- **Light mesons:** spectrum, decay constant, elastic form factor, radii, distribution amplitude, PDFs and scale evolution

—S Jia, J Vary, J Lan, CM, X. Zhao (2018 - 2019)

QED systems

- **Electron:** anomalous magnetic moments, GPDs
- **positronium** wave function, spectroscopy, FFs, GPDs

—Zhao, Wiecki, Li, Honkanen, Chakrabarti, Maris, Vary, Brodsky (2013 - 2018)

Talks on BLFQ: Zhao (16:55), Maris (19/9- 11:00), Vary (19/9- 11:55), Meijian (20/9-11:55)

Example: Light Mesons

S. Jia and J. Vary, PRC (2019)

Light front effective Hamiltonian, H_{eff} :

$$\underbrace{\frac{\vec{k}_{\perp}^2 + m_q^2}{x} + \frac{\vec{k}_{\perp}^2 + m_{\bar{q}}^2}{1-x}}_{\text{LF Kinetic energy}} + \underbrace{\kappa^4 x(1-x) \vec{r}_{\perp}^2}_{\text{Transverse}} - \underbrace{\frac{\kappa^4}{(m_q + m_{\bar{q}})^2} \partial_x (x(1-x) \partial_x)}_{\text{Longitudinal}} + H_{\text{NJL}}^{\text{eff}}$$

- quark masses: $[m_{u/d}, m_s] = [337, 445]$ MeV
- Confining strength: $[\kappa_{\pi/\rho}, \kappa_{K/K^*}] = [227, 276]$ MeV
- Coupling constants: $[G_{\pi/\rho}, G_{K/K^*}] = [18.5, 13.6]$ GeV⁻²

Mass	BLFQ-NJL	PDG
m_{π^+}	139.57 MeV	139.57 MeV
m_{ρ^+}	775.23 ± 0.04 MeV	775.26 ± 0.25 MeV
m_{K^+}	493.68 MeV	493.68 ± 0.02 MeV
$m_{K^{*+}}$	891.82 ± 0.06 MeV	891.76 ± 0.25 MeV
$m_{K_0^{*+}}$	858.35 MeV	824 ± 30 MeV

Elastic Form Factor

S. Jia and J. Vary, PRC (2019)

Drell & Yan (PRL, 70); West (PRL, 70)

- The elastic form factors of the pseudoscalar states:

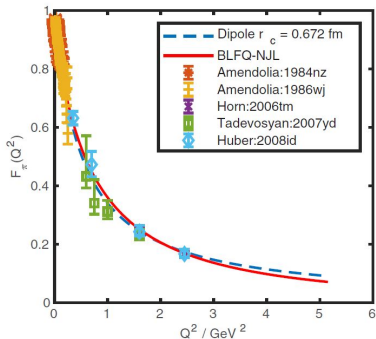
$$F_P(Q^2) = I_{0,0}(Q^2).$$

$$I_{m_J, m_{J'}}(Q^2) = \langle \Psi(P', m_{J'}) | \frac{J^+(0)}{2P^+} | \Psi(P, m_J) \rangle$$

with $q = P' - P$ and $Q^2 = -q^2$.

- The charge radius:

$$\langle r_C^2 \rangle = -6 \lim_{Q^2 \rightarrow 0} \frac{d}{dQ^2} F_P(Q^2).$$

Pion charge radius: $\sqrt{\langle r_C^2 \rangle}_{\pi^+}$ BLFQ: 0.68 ± 0.05 fmPDG: 0.672 ± 0.008 fm

Parton Distribution Amplitude

in preparation, CM *et al.*

Lepage & Brodsky, PRD (1980)

- Distribution amplitude

$$\phi_P(x; \mu) = \frac{2\sqrt{2N_c}}{f_P} \frac{1}{4\pi\sqrt{x(1-x)}} \times \int \frac{d^2k_\perp}{(2\pi)^2} \frac{1}{\sqrt{2}} \left[\psi_{\uparrow\downarrow}(x, k_\perp) - \psi_{\downarrow\uparrow}(x, k_\perp) \right]$$

- Fitting function for pion PDA at $\mu_0 = 442$ MeV

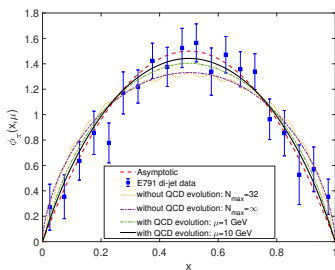
$$\phi(x) = \frac{x^a(1-x)^b}{B(a+1, b+1)},$$

extrapolation: $a = b = 0.60$

Decay constant:

BLFQ: 142.9 MeV

PDG: 130.2 ± 1.7 MeV



- DA evolution: **Gegenbauer basis** Ruiz, *et. al.* PRD 66, (2002)
- Our evolved DA \approx Asymptotic DA

$\pi \rightarrow \gamma(\gamma^*)$ Transition Form Factor

in preparation, CM et al.

Lepage & Brodsky, PRD (1980)

- $\pi \rightarrow \gamma^* \gamma$ TFF:

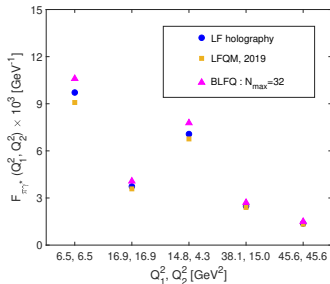
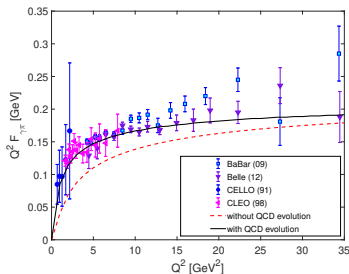
$$F_{\pi\gamma}(Q^2) = \frac{\sqrt{2}}{3} f_P \int_0^1 dx \frac{\phi_\pi(x, xQ)}{Q^2 x}$$

- TFF data prefer QCD evolution of DA

- $\pi \rightarrow \gamma^* \gamma^*$ TFF: $F_{\pi\gamma^*}(Q^2)$

$$\approx \frac{\sqrt{2}}{3} f_P \int_0^1 dx \frac{\phi_\pi(x)}{Q_1^2 x + Q_2^2 (1-x)}$$

LFQM: Choi, Ryu, Ji, PRD 99 (2019)

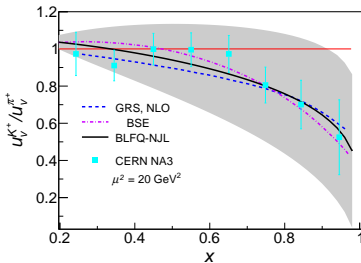
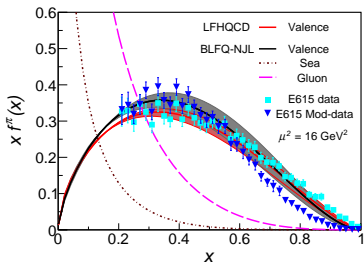


Light Meson PDFs

Lan, CM, Jia, Zhao, Vary: PRL 122 (2019)

$$f(x) = x^a(1-x)^b/B(a+1, b+1),$$

$a = b = 0.5961$ for pion, while $a = 0.6337$ and $b = 0.8546$ for kaon



- LF wavefunctions ► **eigenvectors of effective Hamiltonian.**
PDFs evolution ► based on the **NNLO DGLAP equations.**

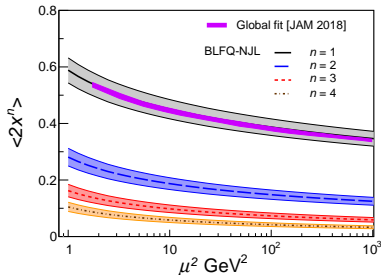
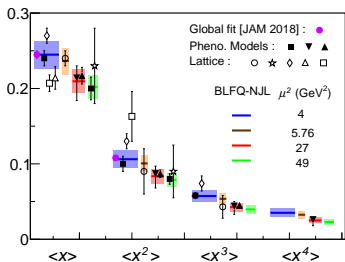
Order	Initial scale of pion	Initial scale of kaon	E-0615 $\chi^2/(d.o.f.)$	NA-003 $\chi^2/(d.o.f.)$
LO	$0.120 \pm 0.012 \text{ GeV}^2$	$0.133 \pm 0.013 \text{ GeV}^2$	6.71	0.88
NLO	$0.205 \pm 0.020 \text{ GeV}^2$	$0.210 \pm 0.021 \text{ GeV}^2$	4.67	0.56
NNLO	$0.240 \pm 0.024 \text{ GeV}^2$	$0.246 \pm 0.024 \text{ GeV}^2$	3.64	0.50

Moments of Pion PDF

Lan, CM, Jia, Zhao, Vary: arXiv: 1907.01509

Moments of the valence quark PDF

$$\langle x^n \rangle = \int_0^1 dx x^n f_v^\pi(x, \mu^2), \quad n = 1, 2, 3, 4.$$



- The robustness of our results motivates the application of analogous effective Hamiltonians to the baryons.

Effective Light-front Hamiltonian for Nucleon

$$P^- = H_{K.E.} + H_{trans} + H_{longi} + H_{OGE}$$

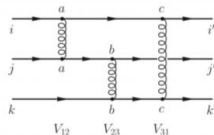
$$H_{K.E.} = \sum_i \frac{p_i^2 + m_q^2}{p_i^+}$$

$$H_{trans} \sim \kappa_T^4 r^2 \quad \text{-- Brodsky, Teramond arXiv: 1203.4025}$$

$$H_{longi} \sim - \sum_{ij} \kappa_L^4 \partial_{x_i} (x_i x_j \partial_{x_j}) \quad \text{---Y Li, X Zhao, P Maris, J Vary, PLB 758(2016)}$$

$$H_{OGE} = - \frac{C_F 4\pi\alpha_s}{Q^2} \sum_{i,j(i<j)} \bar{u}_{s'_i}(k'_i) \gamma^\mu u_{s_i}(k_i) \bar{u}_{s'_j}(k'_j) \gamma_\mu u_{s_j}(k_j) \quad \text{Color factor : } C_F = -\frac{2}{3}$$

$$|P_{baryon}\rangle = |qqq\rangle + |qqqg\rangle + |qqq q\bar{q}\rangle + \dots$$



Three active-quark approach

Basis Construction

- Fock's space expansion:

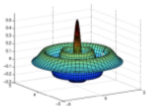
$$|N\rangle_{\text{baryon}} = a|qqq\rangle + b|qqqg\rangle + c|qqqq\bar{q}\rangle + \dots$$

- For each Fock particle:

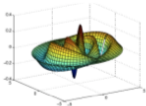
- Transverse:** 2D harmonic oscillator basis: $\Phi_{n,m}^b(\vec{p}_\perp)$

labeled by radial (angular) quantum number n (m); scale parameter b

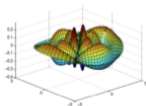
e.g., $n=4$



$m=0$



$m=1$



$m=2$

- Longitudinal:** plane-wave basis, labeled by \mathbf{k}
- Helicity:** labeled by λ

- For the first Fock sector:

$$|qqq\rangle = |n_{q_1}, m_{q_1}, k_{q_1}, \lambda_{q_1}\rangle \otimes |n_{q_2}, m_{q_2}, k_{q_2}, \lambda_{q_2}\rangle \otimes |n_{q_3}, m_{q_3}, k_{q_3}, \lambda_{q_3}\rangle.$$

Basis Truncation Scheme

- Symmetries of Hamiltonian:

- Net fermion number:

$$\sum_i n_i^f = N^f$$

- Total angular momentum projection:

$$\sum_i (m_i + s_i) = J_z$$

- Longitudinal momentum:

$$\sum_i k_i = K$$

- Further truncation:

- Fock sector truncation

- Discretization in longitudinal direction

$$k_i = \begin{cases} 1, 2, 3, \dots & \text{bosons} \\ 0.5, 1.5, 2.5, \dots & \text{fermions} \end{cases}$$

- “ N_{\max} ” truncation in transverse directions

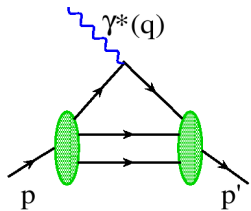
$$\sum_i [2n_i + |m_i| + 1] \leq N_{\max}$$

UV cutoff $\Lambda \sim b\sqrt{N_{\max}}$; IR cutoff $\lambda \sim b/\sqrt{N_{\max}}$

Nucleon Form Factors

In preparation, CM, Siqi Xu, et. al.

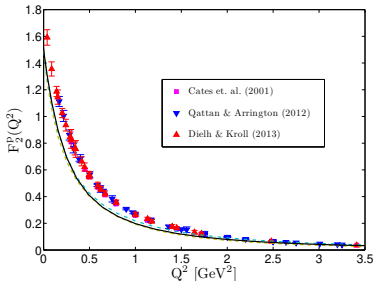
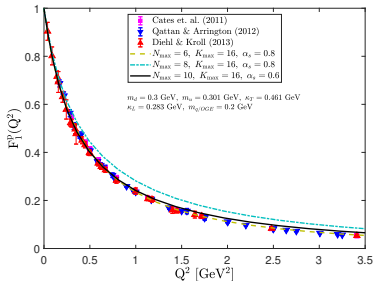
EM form factors in light-front



$$\langle P'; \uparrow | \frac{J^+(0)}{2P^+} | P; \uparrow \rangle = F_1(q^2)$$

$$\langle P'; \uparrow | \frac{J^+(0)}{2P^+} | P; \downarrow \rangle = -(q^1 - iq^2) \frac{F_2(q^2)}{2M}$$

Drell & Yan (PRL, 70); West (PRL, 70)



Nucleon Form Factors

In preparation, CM, Siqi Xu, *et. al.*

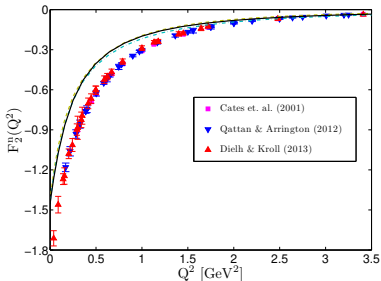
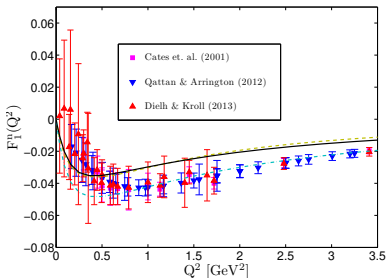
- Only valence quarks contributions
- Missing meson-cloud effects
- $|qqqq\bar{q}\rangle$ has a significant effect on Pauli FF: 30% in proton; 40% in neutron

Sufian *et. al.* PRD 95 (2017)

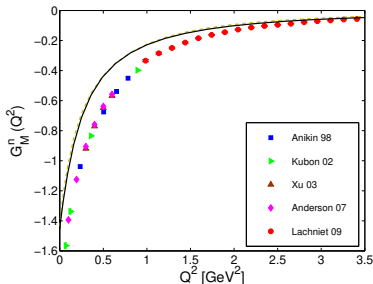
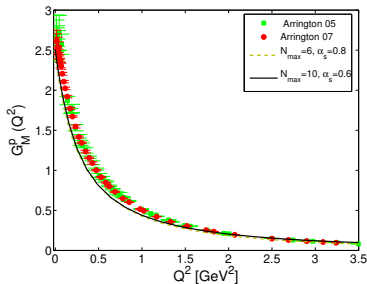
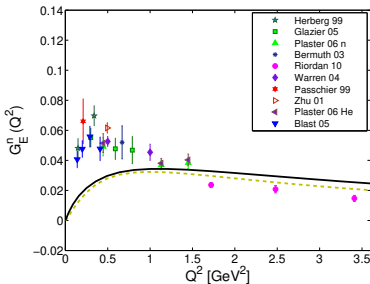
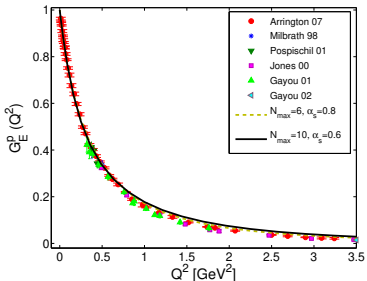
Magnetic moments (μ_N):

Proton: 2.51 (Exp. : 2.79)

Neutron: -1.45 (Exp. : -1.91)



Nucleon Sach's Form Factors

In preparation, CM, Siqi Xu, *et. al.*

Axial Form Factor

In preparation, CM, Siqi Xu, *et. al.*

$$\langle N(p) | A^\mu | N(p') \rangle = \bar{u}(p') \left[\gamma^\mu G_A(t) + \frac{(p' - p)^\mu}{2m} G_P(t) \right] \gamma_5 u(p)$$

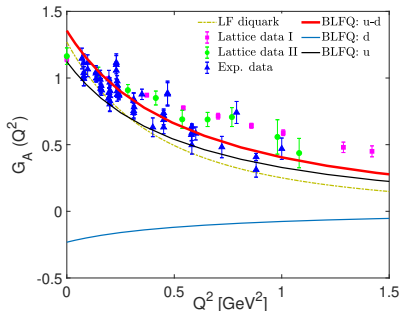
- Axial vector current:

$$A^\mu = \bar{q} \gamma^\mu \gamma_5 q$$

- Axial form factor can be measured by ordinary muon capture (OMC)

$$\mu^-(l) + p(r) \rightarrow \nu_\mu(l') + n(r')$$

- Provide information on spin-isospin distributions



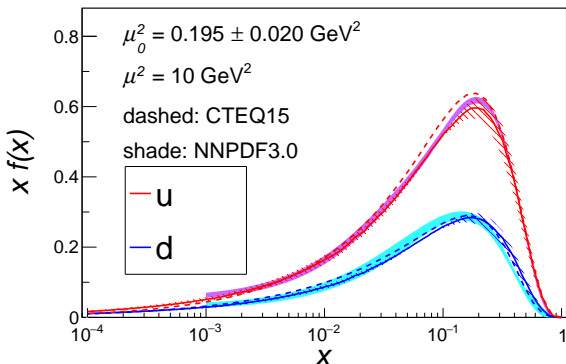
$$G_A(Q^2) = G_u(Q^2) - G_d(Q^2)$$

Unpolarized PDFs

In preparation, CM, Siqi Xu, et. al.

Unpolarized PDF is defined as ($\Gamma \equiv \gamma^+$):

$$\Phi^{\Gamma(\nu)}(x) = \frac{1}{2} \int \frac{dz^-}{2(2\pi)} e^{ip^+z^-/2} \langle P; S | \bar{\psi}(\nu)(0) \Gamma \psi(\nu)(z^-) | P; S \rangle \Big|_{z^+=z_T=0}.$$

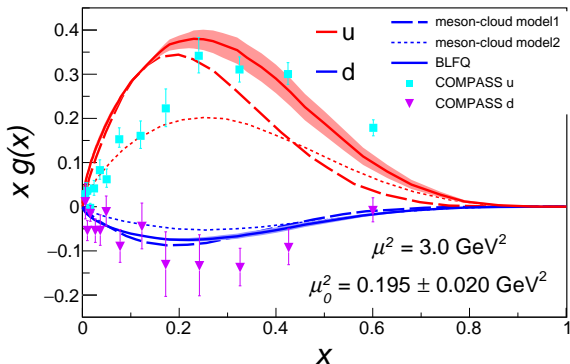


Helicity PDFs

In preparation, CM, Siqi Xu, *et. al.*

Helicity PDF is defined as ($\Gamma \equiv \gamma^+ \gamma_5$):

$$\Phi^{\Gamma(\nu)}(x) = \frac{1}{2} \int \frac{dz^-}{2(2\pi)} e^{ip^+z^-/2} \langle P; S | \bar{\psi}^{(\nu)}(0) \Gamma \psi^{(\nu)}(z^-) | P; S \rangle \Big|_{z^+ = z_T = 0}.$$



model1 & model2: meson-cloud models

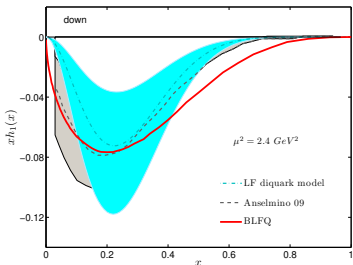
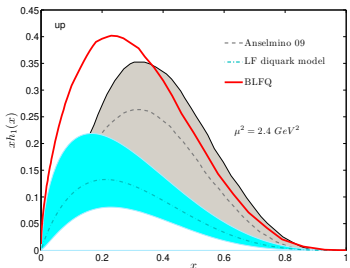
Kofler & Pasquini, PRD 95 (2017)

Transversity PDFs

In preparation, CM, Siqi Xu, *et. al.*

Transversity PDF is defined as ($\Gamma \equiv i\sigma^{J^+}\gamma_5$):

$$\Phi^{\Gamma(\nu)}(x) = \frac{1}{2} \int \frac{dz^-}{2(2\pi)} e^{ip^+z^-/2} \langle P; S | \bar{\psi}(\nu)(0) \Gamma \psi(\nu)(z^-) | P; S \rangle \Big|_{z^+=z_T=0}.$$



Qualitative behavior of BLFQ results are in more or less agreement with other calculations.

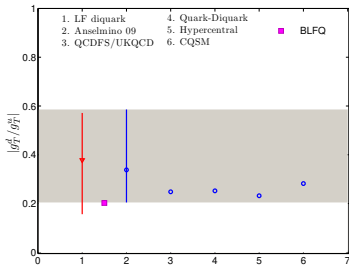
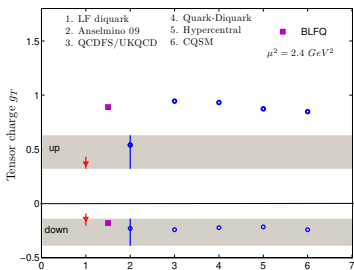
Anselmino, *et al.* Nucl.Phys.Proc.Suppl. 191 (2009); Maji & Chakrabarti, PRD 94 (2016)

Tensor Charge

In preparation, CM, Siqi Xu, et. al.

The tensor charge is defined as

$$g_T^q(\mu^2) = \int_0^1 dx h_1^q(x, \mu^2).$$



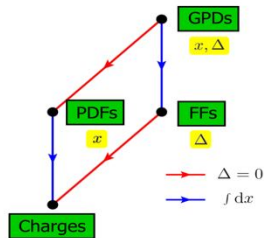
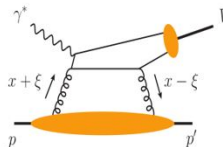
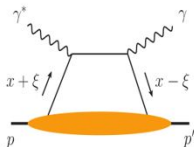
BLFQ results are consistent with other calculations.

Lattice QCD: $g_T^u = 0.3(2)$ and $g_T^d = -0.7(2)$ at $\mu^2 = 2 \text{ GeV}^2$

Lin et. al. PRL, 120 (2018)

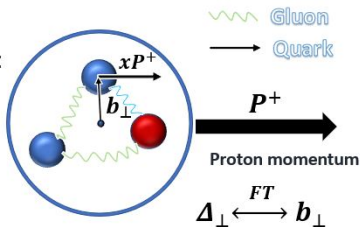
GPDs

➤ Deeply Virtual Compton Scattering (DVCS)/ vector meson productions experiment:



➤ GPDs appear in DVCS processes.

- GPDs are functions of three variables :
 - Longitudinal momentum fraction $x = \frac{k^+}{P^+}$
 - Longitudinal momentum transfer \rightarrow skewness $\xi = \frac{\Delta^+}{P^+} = 0$
 - Square of total mom transfer $t = \Delta^2 = (P' - P)^2$



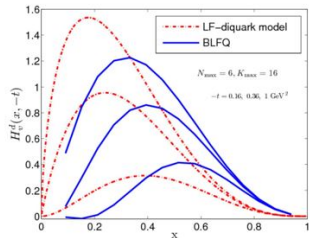
where the \mathbf{b}_\perp is transverse position of parton

Proton unpolarized GPDs in BLFQ

In preparation, CM, Siqi Xu, *et. al.*

Encode information about **three dimensional spatial structure**, as well as the **spin and orbital angular momentum** of the constituents

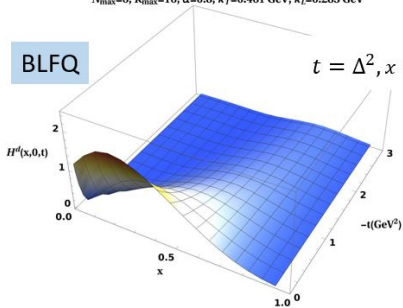
As momentum transfer (t) increases, peaks of distributions shift to larger x ; qualitative behavior of **BLFQ results are consistent** with **LF quark-diquark model results**



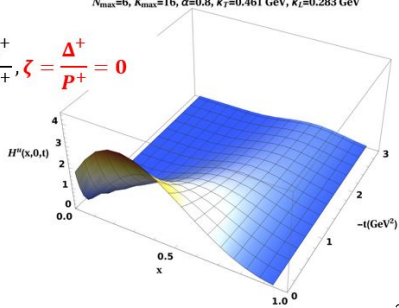
$N_{\max}=6, K_{\max}=16, \alpha=0.8, \kappa_T=0.461 \text{ GeV}, \kappa_L=0.283 \text{ GeV}$

BLFQ

$$t = \Delta^2, x = \frac{k^+}{P^+}, \zeta = \frac{\Delta^+}{P^+} = 0$$



$N_{\max}=6, K_{\max}=16, \alpha=0.8, \kappa_T=0.461 \text{ GeV}, \kappa_L=0.283 \text{ GeV}$



Conclusions & Outlook

- We discussed the structure of pion and proton from the eigenstates of light front effective Hamiltonian
- Consider the leading Fock sectors ($|q\bar{q}\rangle$ for pion, $|qqq\rangle$ for proton).
- BLFQ provides a good description of the experimental data & Global fits for various observables.

Outlook:

- Other distribution functions : spin asymmetries ,GTMD, DPD...
- Study meson-cloud effects and gluon content by including higher Fock sectors.
- Mechanical properties, spin structure of proton, mass decomposition.

Thank You