Color Confinement and Supersymmetric Features of Hadron Physics from Light-Front Holography and Superconformal Algebra





with Guy de Tèramond, Hans Günter Dosch, Marina Nielsen, F. Navarra, Tianbo Llu, Liping Zou, S. Groote, S. Koshkarev, C. Lorcè, R. S. Sufian, A. Deur



LC2019 - QCD ON THE LIGHT CONE: FROM HADRONS TO HEAVY IONS

Ecole Polytechnique, Palaiseau, France



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ABORATORY



The leading Regge trajectory: Δ resonances with maximal J in a given mass range. Also shown is the Regge trajectory for mesons with J = L+S.

E. Klempt and B. Ch. Metsch



Profound Questions for Hadron Physics

- Color Confinement
- Origin of the QCD Mass Scale
- Meson and Baryon Spectroscopy
- Exotic States: Tetraquarks, Pentaquarks, Gluonium,
- Universal Regge Slopes: n, L, Mesons and Baryons
- Massless Pion: Quark Anti-Quark Bound State
- Hadron Structure and Dynamics: QCD Coupling at all Scales

Need a First Approximation to QCD

Comparable in simplicity to Schrödinger Theory in Atomic Physics

Relativistic, Frame-Independent, Color-Confining

Origin of hadronic mass scale

AdS/QCD Líght-Front Holography Superconformal Algebra

No parameters except for quark masses!



P.A.M Dirac, Rev. Mod. Phys. 21, 392 (1949)

Dírac's Amazing Idea: The "Front Form"

Evolve in light-front time!



Bound States in Relativistic Quantum Field Theory:

Light-Front Wavefunctions Dirac's Front Form: Fixed $\tau = t + z/c$

Fixed
$$\tau = t + z/c$$

 $\psi(x_i, \vec{k}_{\perp i}, \lambda_i)$
 $x = \frac{k^+}{P^+} = \frac{k^0 + k^3}{P^0 + P^3}$

Invariant under boosts. Independent of P^{μ}

$$\mathbf{H}_{LF}^{QCD}|\psi\rangle = M^2|\psi\rangle$$

Direct connection to QCD Lagrangian

LF Wavefunction: off-shell in invariant mass

Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space



Measurements of hadron LF wavefunction are at fixed LF time

Like a flash photograph

Fixed
$$\tau = t + z/c$$

$$x_{bj} = x = \frac{k^+}{P^+}$$

Invariant under boosts! Independent of P^{μ}

Light-Front Wavefunctions: rigorous representation of composite systems in quantum field theory



Causal, Frame-independent. Creation Operators on Simple Vacuum, Current Matrix Elements are Overlaps of LFWFS

 $= 2p^+F(q^2)$

Front Form



Drell, sjb

Exact LF Formula for Paulí Form Factor

$$\begin{split} \frac{F_2(q^2)}{2M} &= \sum_a \int [\mathrm{d}x] [\mathrm{d}^2 \mathbf{k}_{\perp}] \sum_j e_j \; \frac{1}{2} \; \times & \text{Drell, sjb} \\ \left[\; -\frac{1}{q^L} \psi_a^{\uparrow *}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \; \psi_a^{\downarrow}(x_i, \mathbf{k}_{\perp i}, \lambda_i) + \frac{1}{q^R} \psi_a^{\downarrow *}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \; \psi_a^{\uparrow}(x_i, \mathbf{k}_{\perp i}, \lambda_i) \right] \\ \mathbf{k}'_{\perp i} &= \mathbf{k}_{\perp i} - x_i \mathbf{q}_{\perp} & \mathbf{k}'_{\perp j} = \mathbf{k}_{\perp j} + (1 - x_j) \mathbf{q}_{\perp} \\ \mathbf{z}_{\mathbf{q}}^{\mathbf{q}} \mathbf{1} & q_{R,L} = q^x \pm i q^y \end{split}$$



Must have $\Delta \ell_z = \pm 1$ to have nonzero $F_2(q^2)$

Nonzero Proton Anomalous Moment --> Nonzero orbítal quark angular momentum Dae Sung Hwang, Bo-Qiang Ma, Ivan Schmidt, sjb

Terayev, Okun: B(0) Must vanish because of Equivalence Theorem



Vanishing Anomalous gravitomagnetic moment B(0)



Must include vacuum-induced currents to compute form factors and other current matrix elements in instant form

Boost are dynamical in instant form

Advantages of the Dirac's Front Form for Hadron Physics Poincare' Invariant

Physics Independent of Observer's Motion

- Measurements are made at fixed τ
- Causality is automatic
- Structure Functions are squares of LFWFs
- Form Factors are overlap of LFWFs
- LFWFs are frame-independent: no boosts, no pancakes!

Penrose, Terrell, Weisskopf

- Same structure function measured at an e p collider and the proton rest frame
- No dependence of hadron structure on observer's frame
- LF Holography: Dual to AdS space
- LF Vacuum trivial up to zero modes
- Implications for Cosmological Constant

Roberts, Shrock, Tandy, sjb





Light-Front QCD

Physical gauge: $A^+ = 0$

Exact frame-independent formulation of nonperturbative QCD!

$$\begin{split} L^{QCD} &\to H_{LF}^{QCD} \\ H_{LF}^{QCD} &= \sum_{i} [\frac{m^{2} + k_{\perp}^{2}}{x}]_{i} + H_{LF}^{int} \\ H_{LF}^{int}: \text{ Matrix in Fock Space} \\ H_{LF}^{QCD} |\Psi_{h} \rangle &= \mathcal{M}_{h}^{2} |\Psi_{h} \rangle \\ |p, J_{z} \rangle &= \sum_{n=3} \psi_{n}(x_{i}, \vec{k}_{\perp i}, \lambda_{i}) |n; x_{i}, \vec{k}_{\perp i}, \lambda_{i} \rangle \end{split}$$

Eigenvalues and Eigensolutions give Hadronic Spectrum and Light-Front wavefunctions

LFWFs: Off-shell in P- and invariant mass



Light-Front Perturbation Theory for pQCD

$$T = H_I + H_I \frac{1}{\mathcal{M}_{initial}^2 - \mathcal{M}_{intermediate}^2 + i\epsilon} H_I + \cdots$$

- "History": Compute any subgraph only once since the LFPth numerator does not depend on the process — only the denominator changes!
- Wick Theorem applies, but few amplitudes since all $k^+ > 0$.

- Unitarity is explicit
- Loop Integrals are 3-dimensional

• hadronization: coalesce comoving quarks and gluons to hadrons using light-front wavefunctions
$$\ \Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

K. Chiu, Lorcé, sjb
$$\int_0^1 dx \int d^2 k_\perp$$

$$\sum_{initial} S^z - \sum_{final} S_z \mid \leq n \text{ at order } g^n$$

$$\begin{array}{c} \text{Light-Front QCD} \\ \mathcal{L}_{QCD} \longrightarrow H_{QCD}^{LF} \\ (H_{LF}^{0} + H_{LF}^{I}) |\Psi \rangle = M^{2} |\Psi \rangle \\ (H_{LF}^{0} + H_{LF}^{I}) |\Psi \rangle = M^{2} |\Psi \rangle \\ [\frac{\vec{k}_{\perp}^{2} + m^{2}}{x(1 - x)} + V_{\text{eff}}^{LF}] \psi_{LF}(x, \vec{k}_{\perp}) = M^{2} \psi_{LF}(x, \vec{k}_{\perp}) \\ (-\frac{d^{2}}{d\zeta^{2}} + \frac{1 - 4L^{2}}{4\zeta^{2}} + U(\zeta)] \psi(\zeta) = \mathcal{M}^{2} \psi(\zeta) \\ \hline \text{AdS/QCD:} \\ (U(\zeta) = \kappa^{4} \zeta^{2} + 2\kappa^{2}(L + S - 1)) \end{array}$$

Semiclassical first approximation to QCD

Fixed $\tau = t + z/c$



Coupled Fock states

Elímínate hígher Fock states and retarded ínteractíons

Effective two-particle equation

Azímuthal Basís
$$\zeta, \phi$$

Single variable Equation $m_q = 0$

Confining AdS/QCD potential!

Sums an infinite # diagrams

Maldacena





 \bullet Isomorphism of SO(4,2) of conformal QCD with the group of isometries of AdS space

$$s^2 = \frac{R^2}{z^2} (\eta_{\mu\nu} dx^{\mu} dx^{\nu} - dz^2),$$

 $x^{\mu} \rightarrow \lambda x^{\mu}, \ z \rightarrow \lambda z$, maps scale transformations into the holographic coordinate z.

• AdS mode in z is the extension of the hadron wf into the fifth dimension.

d

• Different values of z correspond to different scales at which the hadron is examined.

$$x^2 \to \lambda^2 x^2, \quad z \to \lambda z.$$

 $x^2 = x_\mu x^\mu$: invariant separation between quarks

• The AdS boundary at $z \to 0$ correspond to the $Q \to \infty$, UV zero separation limit.

AdS/CFT

Dílaton-Modífied Ads

$$ds^{2} = e^{\varphi(z)} \frac{R^{2}}{z^{2}} (\eta_{\mu\nu} x^{\mu} x^{\nu} - dz^{2})$$



- Soft-wall dilaton profile breaks conformal invariance $e^{\varphi(z)} = e^{+\kappa^2 z^2}$
- Color Confinement in z
- •Introduces confinement scale к
- Uses AdS₅ as template for conformal theory



Supersymmetric Features of Hadron Physics from Superconformal Algebra and Light-Front Holography



$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$

Ads Soft-Wall Schrödinger Equation for bound state of two scalar constituents:

$$\left[-\frac{d^2}{dz^2} - \frac{1 - 4L^2}{4z^2} + U(z)\right]\Phi(z) = \mathcal{M}^2\Phi(z)$$

$$U(z) = \kappa^4 z^2 + 2\kappa^2 (L + S - 1)$$

Derived from variation of Action for Dilaton-Modified AdS₅ **Identical to Single-Variable Light-Front Bound State Equation in** ζ !



Light-Front Holography: Unique mapping derived from equality of LF and AdS formula for EM and gravitational current matrix elements and identical equations of motion

de Tèramond, Dosch, sjb

AdS/QCD Soft-Wall Model

 $e^{\varphi(z)} = e^{+\kappa^2 z^2}$



Líght-Front Holography

$$\left[-\frac{d^2}{d\zeta^2} - \frac{1 - 4L^2}{4\zeta^2} + U(\zeta)\right]\psi(\zeta) = M^2\psi(\zeta)$$



Light-Front Schrödinger Equation

 $U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$ Single variable ζ Unique Confinement Potential!

Conformal Symmetry of the action

Confinement scale:

the $\kappa \simeq 0.5 \ GeV$

de Alfaro, Fubini, Furlan:Fubini, Rabinovici:

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

GeV units external to QCD: Only Ratios of Masses Determined

Meson Spectrum in Soft Wall Model

$$m_{\pi} = 0$$
 if $m_q = 0$

Pion: Negative term for J=0 cancels positive terms from LFKE and potential

Massless pion!

- Effective potential: $U(\zeta^2) = \kappa^4 \zeta^2 + 2\kappa^2 (J-1)$
- LF WE

$$\left(-rac{d^2}{d\zeta^2}-rac{1-4L^2}{4\zeta^2}+\kappa^4\zeta^2+2\kappa^2(J-1)
ight)\phi_J(\zeta)=M^2\phi_J(\zeta)$$

• Normalized eigenfunctions $\ \langle \phi | \phi
angle = \int d\zeta \, \phi^2(z)^2 = 1$

$$\phi_{n,L}(\zeta) = \kappa^{1+L} \sqrt{rac{2n!}{(n+L)!}} \, \zeta^{1/2+L} e^{-\kappa^2 \zeta^2/2} L_n^L(\kappa^2 \zeta^2)$$

Eigenvalues

$$\mathcal{M}_{n,J,L}^2 = 4\kappa^2\left(n+rac{J+L}{2}
ight)$$

$$\vec{\zeta}^2 = \vec{b}_\perp^2 x (1-x)$$

G. de Teramond, H. G. Dosch, sjb



Light meson orbital (a) and radial (b) spectrum for $\kappa=0.6$ GeV.





Effective mass from $m(p^2)$

Roberts, et al.

Prediction from AdS/QCD: Meson LFWF



Holographic Mapping of AdS Modes to QCD LFWFs

Integrate Soper formula over angles:

Drell-Yan-West: Form Factors are Convolution of LFWFs

$$F(q^2) = 2\pi \int_0^1 dx \, \frac{(1-x)}{x} \int \zeta d\zeta J_0\left(\zeta q \sqrt{\frac{1-x}{x}}\right) \tilde{\rho}(x,\zeta),$$

with $\widetilde{\rho}(x,\zeta)$ QCD effective transverse charge density.

• Transversality variable

$$\zeta = \sqrt{x(1-x)\vec{b}_{\perp}^2}$$

• Compare AdS and QCD expressions of FFs for arbitrary Q using identity:

$$\int_0^1 dx J_0\left(\zeta Q\sqrt{\frac{1-x}{x}}\right) = \zeta Q K_1(\zeta Q),$$

the solution for $J(Q,\zeta) = \zeta Q K_1(\zeta Q)$!

de Teramond, sjb

Identical to Polchinski-Strassler Convolution of AdS Amplitudes



Timelike Pion Form Factor from AdS/QCD and Light-Front Holography



week ending 24 AUGUST 2012



AdS/QCD Holographic Wave Function for the ρ Meson and Diffractive ρ Meson Electroproduction

• Light Front Wavefunctions: $\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$ off-shell in P^- and invariant mass $\mathcal{M}_{q\bar{q}}^2$



Boost-invariant LFWF connects confined quarks and gluons to hadrons



Comparison for xq(x) in the pion from LFHQCD (red band) with the NLO fits [82,83] (gray band and green curve) and the LO extraction [84]. NNLO results are also included (light blue band). LFHQCD results are evolved from the initial scale $\mu_0 = 1.1\pm0.2$ GeV at NLO and the initial scale $\mu_0 = 1.06\pm0.15$ GeV at NLO.

Universality of Generalized Parton Distributions in Light-Front Holographic QCD

Guy F. de Te´ramond, Tianbo Liu, Raza Sabbir Sufian, Hans Günter Dosch, Stanley J. Brodsky, and Alexandre Deur

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Connection to the Linear Instant-Form Potential





Harmonic Oscillator $U(\zeta) = \kappa^4 \zeta^2$ LF Potential for relativistic light quarks

A.P. Trawinski, S.D. Glazek, H. D. Dosch, G. de Teramond, sjb

QCD Lagrangian

$$\mathcal{L}_{QCD} = -\frac{1}{4} Tr(G^{\mu\nu}G_{\mu\nu}) + \sum_{f=1}^{n_f} i\bar{\Psi}_f D_{\mu}\gamma^{\mu}\Psi_f + \sum_{f=1}^{n_f} i_f\bar{\Psi}_f\Psi_f$$

$$iD^{\mu} = i\partial^{\mu} - gA^{\mu} \qquad G^{\mu\nu} = \partial^{\mu}A^{\mu} - \partial^{\nu}A^{\mu} - g[A^{\mu}, A^{\nu}]$$

Classical Chiral Lagrangian is Conformally Invariant

Where does the QCD Mass Scale come from?

QCD does not know what MeV units mean! Only Ratios of Masses Determined

ode Alfaro, Fubini, Furlan:

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

Unique confinement potential!
• de Alfaro, Fubini, Furlan (dAFF)

$$\begin{aligned} G|\psi(\tau) > &= i\frac{\partial}{\partial\tau}|\psi(\tau) > \\ G &= uH + vD + wK \\ G &= H_{\tau} = \frac{1}{2} \left(-\frac{d^2}{dx^2} + \frac{g}{x^2} + \frac{4uw - v^2}{4}x^2 \right) \end{aligned}$$

Retains conformal invariance of action despite mass scale! $4uw-v^2=\kappa^4=[M]^4$

Identical to LF Hamiltonian with unique potential and dilaton!

$$\begin{bmatrix} -\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \end{bmatrix} \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$
$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L+S-1)$$

Dosch, de Teramond, sjb



- Identify with difference of LF time $\Delta x^+/P^+$ between constituents
- Finite range
- Measure in Double-Parton Processes

Retains conformal invariance of action despite mass scale!

Remarkable Features of Líght-Front Schrödínger Equation

- Relativistic, frame-independent
- •QCD scale appears unique LF potential
- Reproduces spectroscopy and dynamics of light-quark hadrons with one parameter
- Zero-mass pion for zero mass quarks!
- Regge slope same for n and L -- not usual HO
- Splitting in L persists to high mass -- contradicts conventional wisdom based on breakdown of chiral symmetry
- Phenomenology: LFWFs, Form factors, electroproduction
- Extension to heavy quarks

Stan Brodsky

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

Supersymmetric Features of Hadron Physics from Superconformal Algebra and Light-Front Holography



Dynamics + Spectroscopy!

Light-Front Holography: First Approximation to QCD

- Color Confinement, Analytic form of confinement potential
- Retains underlying conformal properties of QCD despite mass scale (DeAlfaro-Fubini-Furlan Principle)
- Massless quark-antiquark pion bound state in chiral limit, GMOR
- QCD coupling at all scales
- Connection of perturbative and nonperturbative mass scales
- Poincarè Invariant
- Hadron Spectroscopy-Regge Trajectories with universal slopes in n, L
- Supersymmetric 4-Plet: Meson-Baryon -Tetraquark Symmetry
- Light-Front Wavefunctions
- Form Factors, Structure Functions, Hadronic Observables
- OPE: Constituent Counting Rules
- Hadronization at the Amplitude Level: Many Phenomenological Tests
- Systematically improvable: Basis LF Quantization (BLFQ)

Stan Brodsky

Supersymmetric Features of Hadron Physics from Superconformal Algebra and Light-Front Holography



Haag, Lopuszanski, Sohnius (1974)

Superconformal Quantum Mechanics $\{\psi,\psi^+\} = 1$ $B = \frac{1}{2}[\psi^+,\psi] = \frac{1}{2}\sigma_3$ $\psi = \frac{1}{2}(\sigma_1 - i\sigma_2), \quad \psi^+ = \frac{1}{2}(\sigma_1 + i\sigma_2)$ $Q = \psi^{+}[-\partial_{x} + \frac{f}{x}], \quad Q^{+} = \psi[\partial_{x} + \frac{f}{x}], \quad S = \psi^{+}x, \quad S^{+} = \psi x$ $\{Q, Q^+\} = 2H, \{S, S^+\} = 2K$ $\{Q, S^+\} = f - B + 2iD, \ \{Q^+, S\} = f - B - 2iD$ generates conformal algebra [H,D] = i H, [H, K] = 2 i D, [K, D] = - i K $Q \simeq \sqrt{H}, S \simeq \sqrt{K}$

Superconformal Quantum Mechanics

Baryon Equation $Q \simeq \sqrt{H}, S \simeq \sqrt{K}$

Consider
$$R_w = Q + wS;$$

w: dimensions of mass squared

$$G = \{R_w, R_w^+\} = 2H + 2w^2K + 2wfI - 2wB \qquad 2B = \sigma_3$$

Retains Conformal Invariance of Action

Fubini and Rabinovici

New Extended Hamíltonían G ís díagonal:

$$G_{11} = \left(-\partial_x^2 + w^2 x^2 + 2wf - w + \frac{4(f + \frac{1}{2})^2 - 1}{4x^2}\right)$$

$$G_{22} = \left(-\partial_x^2 + w^2 x^2 + 2wf + w + \frac{4(f - \frac{1}{2})^2 - 1}{4x^2}\right)$$

Identify $f - \frac{1}{2} = L_B$, $w = \kappa^2$
Firmulas of C : $M^2(n, I) = 4w^2(n + I - 1)$

Eigenvalue of G: $M^2(n, L) = 4\kappa^2(n + L_B + 1)$

de Téramond, Dosch, Lorcé, sjb LF Holography Ba

Baryon Equation

Superconformal Quantum Mechanics

S=0, P=+

$$\begin{split} \left(-\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}(L_{B}+1) + \frac{4L_{B}^{2}-1}{4\zeta^{2}} \right)\psi_{J}^{+} &= M^{2}\psi_{J}^{+} \\ \left(-\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}L_{B} + \frac{4(L_{B}+1)^{2}-1}{4\zeta^{2}} \right)\psi_{J}^{-} &= M^{2}\psi_{J}^{-} \\ M^{2}(n, L_{B}) &= 4\kappa^{2}(n + L_{B}+1) \qquad \text{S=I/2, P=+} \\ Meson \ Equation \qquad \lambda &= \kappa^{2} \\ \left(-\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}(J-1) + \frac{4L_{M}^{2}-1}{4\zeta^{2}} \right)\phi_{J} &= M^{2}\phi_{J} \end{split}$$

 $M^2(n, L_M) = 4\kappa^2(n + L_M)$ Same κ ! S=0, I=I Meson is superpartner of S=1/2, I=I Baryon Meson-Baryon Degeneracy for L_M=L_B+1

LF Holography



Superconformal Quantum Mechanics

• Nucleon LF modes

$$\psi_{+}(\zeta)_{n,L} = \kappa^{2+L} \sqrt{\frac{2n!}{(n+L)!}} \zeta^{3/2+L} e^{-\kappa^{2}\zeta^{2}/2} L_{n}^{L+1} \left(\kappa^{2}\zeta^{2}\right)$$
$$\psi_{-}(\zeta)_{n,L} = \kappa^{3+L} \frac{1}{\sqrt{n+L+2}} \sqrt{\frac{2n!}{(n+L)!}} \zeta^{5/2+L} e^{-\kappa^{2}\zeta^{2}/2} L_{n}^{L+2} \left(\kappa^{2}\zeta^{2}\right)$$

Normalization

$$\int d\zeta \,\psi_+^2(\zeta) = \int d\zeta \,\psi_-^2(\zeta) = 1$$

Eigenvalues $\int_{0}^{\infty} d\zeta \int_{0}^{1} dx \psi_{+}^{2}(\zeta^{2}, x) = \int_{0}^{\infty} d\zeta \int_{0}^{1} dx \psi_{-}^{2}(\zeta^{2}, x) = \frac{1}{2}$ *Quark Chiral Symmetry of Eigenstate!*

Nucleon: Equal Probability for L=0, I $J^{z} = +1/2$: $\frac{1}{\sqrt{2}}[|S_{q}^{z}| = +1/2, L^{z}| = 0 > + |S_{q}^{z}| = -1/2, L^{z}| = +1 >]$

Nucleon spin carried by quark orbital angular momentum









de Téramond, Dosch, Lorcé, sjb





Fit to the slope of Regge trajectories, including radial excitations

Same Regge Slope for Meson, Baryons: Supersymmetric feature of hadron physics

Superconformal Algebra

2X2 Hadronic Multiplets

Bosons, Fermions with Equal Mass!



Superconformal Algebra 4-Plet



New Organization of the Hadron Spectrum

	Meson			Baryon			Tetraquark			
	q-cont	$J^{P(C)}$	Name	q-cont	J^p	Name	q-cont	$J^{P(C)}$	Name	
	$\bar{q}q$	0-+	$\pi(140)$	_	_		_		_	
	$\bar{q}q$	1+-	$b_1(1235)$	[ud]q	$(1/2)^+$	N(940)	$[ud][\overline{u}\overline{d}]$	0++	$f_0(980)$	
	$\bar{q}q$	2^{-+}	$\pi_2(1670)$	[ud]q	$(1/2)^{-}$	$N_{\frac{1}{2}}(1535)$	$[ud][\overline{u}d]$	1-+	$\pi_1(1400)$	
					$(3/2)^{-}$	$N_{\frac{3}{2}}(1520)$			$\pi_1(1600)$	
	āq	1	$\rho(770), \omega(780)$							
	$\bar{q}q$	2++	$a_2(1320), f_2(1270)$	[qq]q	$(3/2)^+$	$\Delta(1232)$	$[qq][\bar{u}\bar{d}]$	1++	$a_1(1260)$	
	$\bar{q}q$	3	$\rho_3(1690), \ \omega_3(1670)$	[qq]q	$(1/2)^{-}$	$\Delta_{\frac{1}{2}}(1620)$	$[qq][\bar{u}d]$	2	$\rho_2 (\sim 1700)?$	
					$(3/2)^{-}$	$\Delta_{a-}(1700)$				
	$\bar{q}q$	4++	$a_4(2040), f_4(2050)$	[qq]q	$(7/2)^+$	$\Delta_{\frac{7}{8}^+}(1950)$	$[qq][\bar{u}\bar{d}]$	3++	$a_3 (\sim 2070)?$	
	\bar{qs}	0-(+)	K(495)			_				
	$\bar{q}s$	1+(-)	$\bar{K}_{1}(1270)$	[ud]s	$(1/2)^+$	Λ(1115)	$[ud][\bar{s}\bar{q}]$	0+(+)	$K_0^*(1430)$	
	$\bar{q}s$	$2^{-(+)}$	$K_2(1770)$	[ud]s	$(1/2)^{-}$	Λ(1405)	$[ud][\bar{s}\bar{q}]$	1-(+)	$K_1^* (\sim 1700)?$	
					$(3/2)^{-}$	$\Lambda(1520)$				
	$\bar{s}q$	0-(+)	K(495)	_	_					
	$\bar{s}q$	1+(-)	$K_1(1270)$	[sq]q	$(1/2)^+$	$\Sigma(1190)$	$[sq][\bar{s}\bar{q}]$	0++	$a_0(980)$	
	_		te: (000)						$f_0(980)$	
(są	1-(-)	<u>K*(890)</u>		(0.(0))		(1()	41(1)		
C	sq	2+(+)	$K_{2}^{*}(1430)$	[sq]q	$(3/2)^+$	Σ(1385) Σ(1650)	sq qq	1+(+)	$K_1(1400)$	
	sq -	3 (-)	$K_{3}(1780)$	[<i>sq</i>] <i>q</i>	$(3/2)^{-}$	Σ(1070) Σ(2020)	[<i>sq</i>][<i>qq</i>]	2 (-)	$K_2(\sim 1700)?$	
	sq -	4	R ₄ (2045)	[sq]q	$(1/2)^{+}$	2(2030)	[<i>sq</i>][<i>qq</i>]	3.(1)	$K_{3}(\sim 2070)$?	
	88	1+-	$\eta(550)$	[an]a	(1.(9)+		[][==]	0++	£ (1970)	
	88	1.	$n_1(1170)$	[sq]s	$(1/2)^{-1}$	2(1320)	[sq][sq]	0	$J_0(1370)$ $q_{-}(1450)$	
	āe	2-+	m(1645)	[20]2	$(7)^{?}$	豆(1690)	[90][90]	1-+	$\Phi'(1750)?$	
		1	Φ(1020)	[04]0	(.)	=======================================	[94][94]	_	* (1150).	
	38	2++	$f'_{2}(1525)$	[sq]s	$(3/2)^+$	E*(1530)	[sq][sq]	1++	$f_1(1420)$	
	38	3	$\Phi_{a}(1850)$	[sq]s	$(3/2)^{-}$	三(1820)	$[sq][\bar{s}\bar{q}]$	2	$\Phi_2(\sim 1800)?$	
	ŝs	2++	f2(1950)	[88]8	(3/2)+	Ω(1672)	$[ss][\bar{s}\bar{q}]$	1+(+)	$K_1(\sim 1700)?$	
	M	esc	n	Barvon			Tetraquark			
					• / ~					

M. Níelsen, sjb Universal Hadronic Decomposition

$$\frac{\mathcal{M}_{H}^{2}}{\kappa^{2}} = (1 + 2n + L) + (1 + 2n + L) + (2L + 4S + 2B - 2)$$
• Universal quark light-front kinetic energy
Equal:
Virial
Virial
Heorem
• Universal quark light-front potential energy
$$\Delta \mathcal{M}_{LFFE}^{2} = \kappa^{2}(1 + 2n + L)$$
• Universal quark light-front potential energy
$$\Delta \mathcal{M}_{LFPE}^{2} = \kappa^{2}(1 + 2n + L)$$
• Universal Constant Contribution from AdS
and Superconformal Quantum Mechanics
$$\Delta \mathcal{M}_{spin}^{2} = 2\kappa^{2}(L + 2S + B - 1)$$

hyperfine spin-spin

Using SU(6) flavor symmetry and normalization to static quantities







Comparison for xq(x) in the proton from LFHQCD (red bands) and global fits: MMHT2014 (blue bands) [5], CT14 [6] (cyan bands), and NNPDF3.0 (gray bands) [77]. LFHQCD results are evolved from the initial scale $\mu_0 = 1.06 \pm 0.15$ GeV.

Universality of Generalized Parton Distributions in Light-Front Holographic QCD

Guy F. de Te´ramond, Tianbo Liu, Raza Sabbir Sufian, Hans Günter Dosch, Stanley J. Brodsky, and Alexandre Deur PHYSICAL REVIEW LETTERS 120, 182001 (2018)

de Téramond, Dosch, Lorcé, sjb

Supersymmetry across the light and heavy-light spectrum



de Téramond, Dosch, Lorcé, sjb

Supersymmetry across the light and heavy-light spectrum



Heavy charm quark mass does not break supersymmetry

Superpartners for states with one c quark

	Me	eson		Bar	yon	Tetraquark			
q-cont	$J^{P(C)}$	Name	$q ext{-cont}$	J^P	Name	q-cont	$J^{P(C)}$	Name	
$\bar{q}c$	0-	D(1870)							
$\bar{q}c$	1+	$D_1(2420)$	[ud]c	$(1/2)^+$	$\Lambda_c(2290)$	$[ud][\bar{c}\bar{q}]$	0^{+}	$\bar{D}_{0}^{*}(2400)$	
$\bar{q}c$	2^{-}	$D_J(2600)$	[ud]c	$(3/2)^{-}$	$\Lambda_c(2625)$	$[ud][\bar{c}\bar{q}]$	1-		
$\bar{c}q$	0-	$\bar{D}(1870)$							
$\bar{c}q$	1+	$D_1(2420)$	[cq]q	$(1/2)^+$	$\Sigma_{c}(2455)$	$[cq][\bar{u}\bar{d}]$	0^{+}	$D_0^*(2400)$	
$\bar{q}c$	1-	$D^{*}(2010)$			_ \				
$\bar{q}c$	2^{+}	$D_2^*(2460)$	(qq)c	$(3/2)^+$	$\Sigma_{c}^{*}(2520)$	$(qq)[\bar{c}\bar{q}]$	1+	D(2550)	
$\bar{q}c$	3^{-}	$D_3^*(2750)$	(qq)c	$(3/2)^{-}$	$\Sigma_{c}(2800)$	$(qq)[\bar{c}\bar{q}]$			
$\bar{s}c$	0-	$D_s(1968)$							
$\bar{s}c$	1+	$D_{s1}(2460)$	[qs]c	$(1/2)^+$	$\Xi_c(2470)$	$[qs][ar{c}ar{q}]$	0^{+}	$\bar{D}_{s0}^{*}(2317)$	
$\bar{s}c$	2^{-}	$Q_{s2}(\sim 2860)?$	[qs]c	$(3/2)^{-}$	$\Xi_c(2815)$	$[sq][\bar{c}\bar{q}]$	1-		
$\bar{s}c$	1-	$D_{s}^{*}(2110)$	$\backslash -$						
$\overline{s}c$	2^{+}	$D_{s2}^{*}(2573)$	(sq)c	$(3/2)^+$	$\Xi_{c}^{*}(2645)$	$(sq)[\bar{c}\bar{q}]$	1+	$D_{s1}(2536)$	
$\bar{c}s$	1+	$Q_{s1}(\sim 2700)?$	[cs]s	$(1/2)^+$	$\Omega_c(2695)$	$[cs][\bar{s}\bar{q}]$	0^{+}	??	
$\bar{s}c$	2^{+}	$D_{s2}^* (\sim 2750)?$	(ss)c	$(3/2)^+$	$\Omega_c(2770)$	$(ss)[\bar{c}\bar{s}]$	1+	??	
M. Níelsen, sjb				pr	edictions	beautiful agreement! 58			

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Supersymmetry across the light and heavy-light spectrum



Heavy bottom quark mass does not break supersymmetry

Heavy-light and heavy-heavy hadronic sectors

• Extension to the heavy-light hadronic sector

[H. G. Dosch, GdT, S. J. Brodsky, PRD 92, 074010 (2015), PRD 95, 034016 (2017)]

• Extension to the double-heavy hadronic sector

[M. Nielsen and S. J. Brodsky, PRD, 114001 (2018)]

[M. Nielsen, S. J. Brodsky, GdT, H. G. Dosch, F. S. Navarra, L. Zou, PRD 98, 034002 (2018)]

• Extension to the isoscalar hadronic sector

[L. Zou, H. G. Dosch, GdT,S. J. Brodsky, arXiv:1901.11205 [hep-ph]]



Supersymmetry in QCD

- A hidden symmetry of Color SU(3)c in hadron physics
- QCD: No squarks or gluinos!
- Emerges from Light-Front Holography and Super-Conformal Algebra
- Color Confinement
- Massless Pion in Chiral Limit

de Téramond, Dosch, Lorcé, sjb



Supersymmetric Features of Hadron Physics from Superconformal Algebra and Light-Front Holography



Superconformal Algebra

2X2 Hadronic Multiplets

Bosons, Fermions with Equal Mass!





Use counting rules to identify composite structure

Lebed, sjb

Underlying Principles

- Polncarè Invariance: Independent of the observer's Lorentz frame: Quantization at Fixed Light-Front Time τ
- Causality: Information within causal horizon: Light-Front
- Light-Front Holography: AdS₅ = LF (3+1)

 $z \leftrightarrow \zeta$ where $\zeta^2 = b_{\perp}^2 x(1-x)$

- Introduce mass scale к while retaining the Conformal Invariance of the Action (dAFF)
 "Fmerce
- Unique Dilaton in AdS₅: $e^{+\kappa^2 z^2}$

Stan Brodsky

- Unique color-confining LF Potential $~U(\zeta^2)=\kappa^4\zeta^2$
- Superconformal Algebra: Mass Degenerate 4-Plet:

Meson $q\bar{q} \leftrightarrow \text{Baryon } q[qq] \leftrightarrow \text{Tetraquark } [qq][\bar{q}\bar{q}]$

Supersymmetric Features of Hadron Physics from Superconformal Algebra and Light-Front Holography





19 Sept 2019

NAL ACCELERATOR LABORATOR



Running Coupling from Modified AdS/QCD Deur, de Teramond, sjb

Consider five-dim gauge fields propagating in AdS₅ space in dilaton background $arphi(z) = \kappa^2 z^2$

$$e^{\phi(z)} = e^{+\kappa^2 z^2} \qquad S = -\frac{1}{4} \int d^4 x \, dz \, \sqrt{g} \, e^{\varphi(z)} \, \frac{1}{g_5^2} \, G^2$$

Flow equation

$$\frac{1}{g_5^2(z)} = e^{\varphi(z)} \frac{1}{g_5^2(0)} \quad \text{or} \quad g_5^2(z) = e^{-\kappa^2 z^2} g_5^2(0)$$

where the coupling $g_5(z)$ incorporates the non-conformal dynamics of confinement

- YM coupling $\alpha_s(\zeta) = g_{YM}^2(\zeta)/4\pi$ is the five dim coupling up to a factor: $g_5(z) \to g_{YM}(\zeta)$
- $\bullet\,$ Coupling measured at momentum scale Q

$$\alpha_s^{AdS}(Q) \sim \int_0^\infty \zeta d\zeta J_0(\zeta Q) \,\alpha_s^{AdS}(\zeta)$$

Solution

$$\alpha_s^{AdS}(Q^2) = \alpha_s^{AdS}(0) e^{-Q^2/4\kappa^2}.$$

where the coupling α_s^{AdS} incorporates the non-conformal dynamics of confinement

$${}^{1} dx[g_{1}^{ep}(x,Q^{2}) - g_{1}^{en}(x,Q^{2})] \equiv \frac{g_{a}}{6}[1 - \frac{\alpha_{g1}(Q^{2})}{\pi}]$$

 $\alpha_{q1}(Q^2)$

•Can be used as standard QCD coupling

• Well measured

- Asymptotic freedom at large Q²
- Computable at large Q² in any pQCD scheme
- Universal β_0 , β_1



Analytic, defined at all scales, IR Fixed Point

AdS/QCD dilaton captures the higher twist corrections to effective charges for Q < 1 GeV

$$e^{\varphi} = e^{+\kappa^2 z^2}$$

Deur, de Teramond, sjb



Poincarè Invariance: Independent of the observer's Lorentz frame: Quantization at Fixed Light-Front Time τ

Underlying Principles

- Causality: Information within causal horizon: Light-Front
- Light-Front Holography: AdS₅ = LF (3+1) $z \leftrightarrow \zeta$ where $\zeta^2 = b_{\perp}^2 x(1-x)$



- Introduce Mass Scale κ while retaining the Conformal Invariance of the Action (dAFF)
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Stan Brodsky

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Supersymmetric Features of Hadron Physics from Superconformal Algebra and Light-Front Holography





uniquely identify the ß terms

Renormalization scale depends on the thrust

Not constant !



1**-**T

T. Gehrmann, N. H'afliger, P. F. Monni

S.-Q. Wang, L. Di Giustino, X.-G. Wu, sjb



Principle of Maximum Conformality (PMC)




Features of BLM/PMC

- Predictions are scheme-independent at every order
- Matches conformal series
- No n! Renormalon growth of pQCD series
- New scale appears at each order; n_F determined at each order matches virtuality of quark loops
- Multiple Physical Scales Incorporated (Hoang, Kuhn, Tuebner, sjb)
- Rigorous: Satisfies all Renormalization Group Principles
- Realistic Estimate of Higher-Order Terms
- Reduces to standard QED scale $N_C \rightarrow 0$
- GUT: Must use the same scale setting procedure for QED, QCD
- Eliminates unnecessary theory error
- Maximal sensitivity to new physics
- Commensurate Scale Relations between observables: Generalized Crewther Relation (Kataev, Lu, Rathsman, sjb)
- PMC Reduces to BLM at NLO: Example: BFKL intercept (Fadin, Kim, Lipatov, Pivovarov, sjb)





A-1

Front-Face Nucleon N1 struck Front-Face Nucleon N₁ not struck One-Step / Two-Step Interference Study Double Virtual Compton Scattering $\gamma^* A \to \gamma^* A$

Cannot reduce to matrix element Liuti, Schmidt sjb of local operator! No Sum Rules!

- Unlike shadowing, anti-shadowing from Reggeon exchange is flavor specific;
- Each quark and anti-quark will have distinctly different constructive interference patterns
- The flavor dependence of antishadowing explains why anti- shadowing is different for electron (neutral electro- magnetic current) vs. neutrino (charged weak current) DIS reactions.
- Test of the explanation of antishadowing: Bjorken-scaling leading-twist charge exchange DDIS reaction $\gamma^*p \rightarrow nX^+$ with a rapidity gap due to I=1 Reggeon exchange
- The finite path length due to the on-shell propagation of V⁰ between N₁ and N₂ contributes a finite distance $(\Delta z)^2$ between the two virtual photons in the DVCS amplitude.

The usual "handbag" diagram where the two J $\mu(x)$ and J $\nu(0)$ currents acting on an uninterrupted quark propagator are replaced by a local operator T $\mu\nu(0)$ as Q² $\rightarrow \infty$, is inapplicable in deeply virtual Compton scattering from a nucleus since the currents act on different nucleons.

$$\Delta z^2$$
 does not vanish as $\frac{1}{Q^2}$.

OPE and Sum Rules invalid for nuclear pdfs

Invariance Principles of Quantum Field Theory

- Polncarè Invariance: Physical predictions must be independent of the observer's Lorentz frame: Front Form
- Causality: Information within causal horizon: Front Form
- Gauge Invariance: Physical predictions of gauge theories must be independent of the choice of gauge
- Scheme-Independence: Physical predictions of renormalizable theories must be independent of the choice of the renormalization scheme — Principle of Maximum Conformality (PMC)
- Mass-Scale Invariance: Conformal Invariance of the Action (DAFF)

Stan Brodsky





Color Confinement and Supersymmetric Features of Hadron Physics from Light-Front Holography and Superconformal Algebra





with Guy de Tèramond, Hans Günter Dosch, Marina Nielsen, F. Navarra, Tianbo Llu, Liping Zou, S. Groote, S. Koshkarev, C. Lorcè, R. S. Sufian, A. Deur



LC2019 - QCD ON THE LIGHT CONE: FROM HADRONS TO HEAVY IONS



Ecole Polytechnique, Palaiseau, France

September 19, 2019

ABORATORY