

Ward-Takahashi Identity on Light-Front

"LC2019-QCD on the Light Cone: from Hadrons to Heavy Ions"

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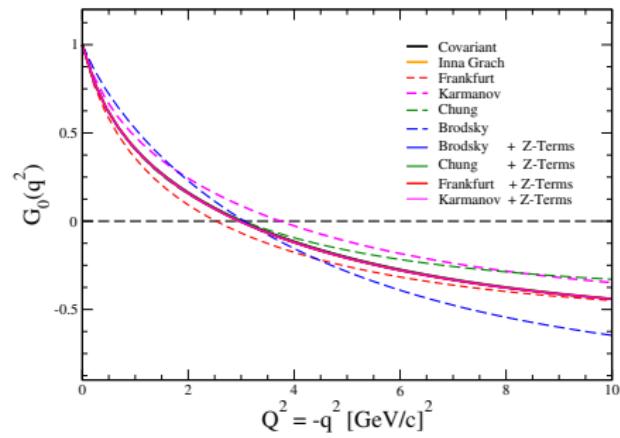
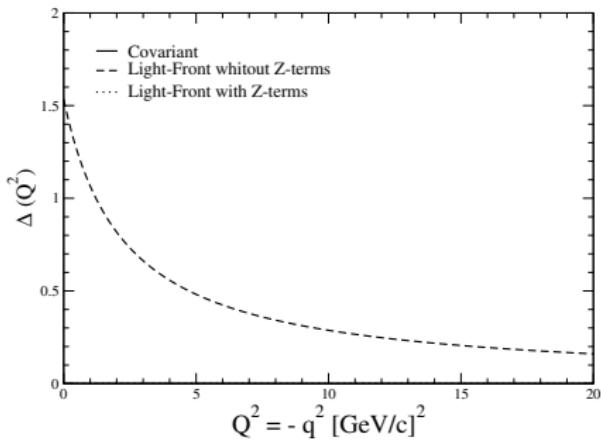
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⇒ Some references:

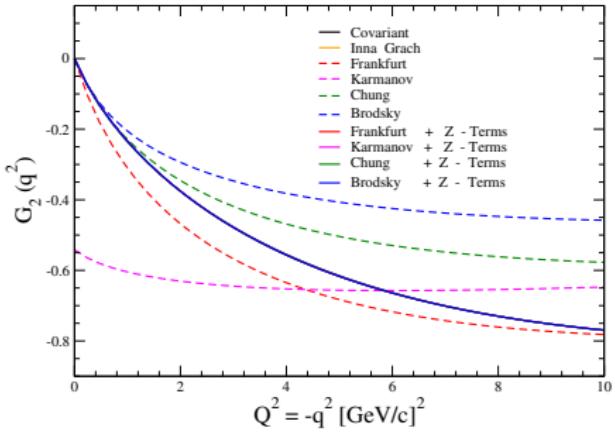
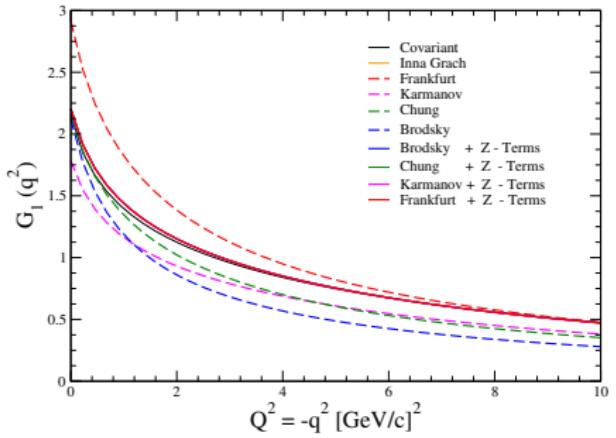
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“Pairs in the light front and covariance,” Nucl. Phys. A 631, 574C (1998)
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Motivation: S=1 / Rho Meson

- J. P. B. C. de Melo and T. Frederico, Phys. Rev. C 55, 2043 (1997)



$$\begin{aligned}\Delta(q^2) &= (1+2\eta)I_{11}^+ + I_{1-1}^+ - \sqrt{8\eta}I_{10}^+ - I_{00}^+ = 0 \\ &= (1+\eta)(J_{yy}^{+V} - J_{zz}^+) = 0\end{aligned}$$



- Eletromagnetic current for Bosons

$$J^\mu = \Gamma^2 \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)^\mu}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)}$$

- Plus Component of the electromagnetic current: J^+

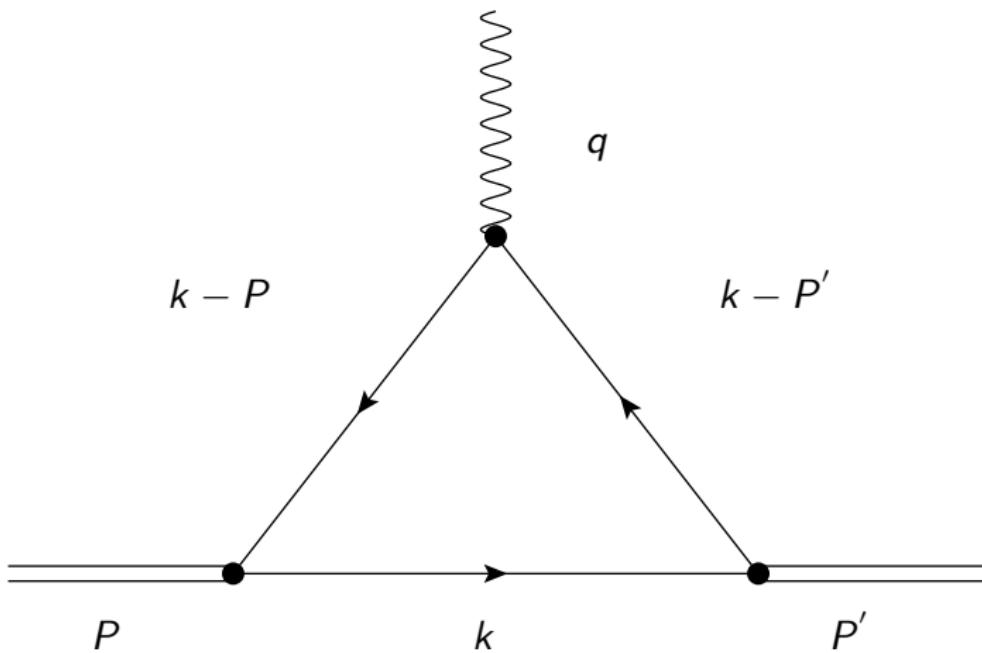
$$J^+ = \Gamma^2 \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)^+}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)}$$

- And with LF coordinates:

$$J^+ = \Gamma^2 \int \frac{d^2 k_\perp dk^+ dk^-}{2(2\pi)^4} \frac{(2k - P' - P)^+}{k^+(P^+ - k^+)(P'^+ - k^+)(k^- - \frac{f_1 - i\epsilon}{k^+})} \\ \frac{1}{(P^- - k^- - \frac{f_2 - i\epsilon}{P^+ - k^+})(P'^- - k^- - \frac{f_3 - i\epsilon}{P'^+ - k^+})}$$

- with the def.:

$$f_1 = k_\perp^2 + m^2, \quad f_2 = (P - k)_\perp^2 + m^2 \quad e \quad f_3 = (P' - k)_\perp^2 + m^2$$



Poles position

$$\begin{aligned}k_1^- &= \frac{f_1 - i\epsilon}{k^+} \\k_2^- &= P^- - \frac{f_2 - i\epsilon}{P^+ - k^+} \\k_3^- &= P'^- - \frac{f_3 - i\epsilon}{P'^+ - k^+}.\end{aligned}$$

i) $k^+ < 0$

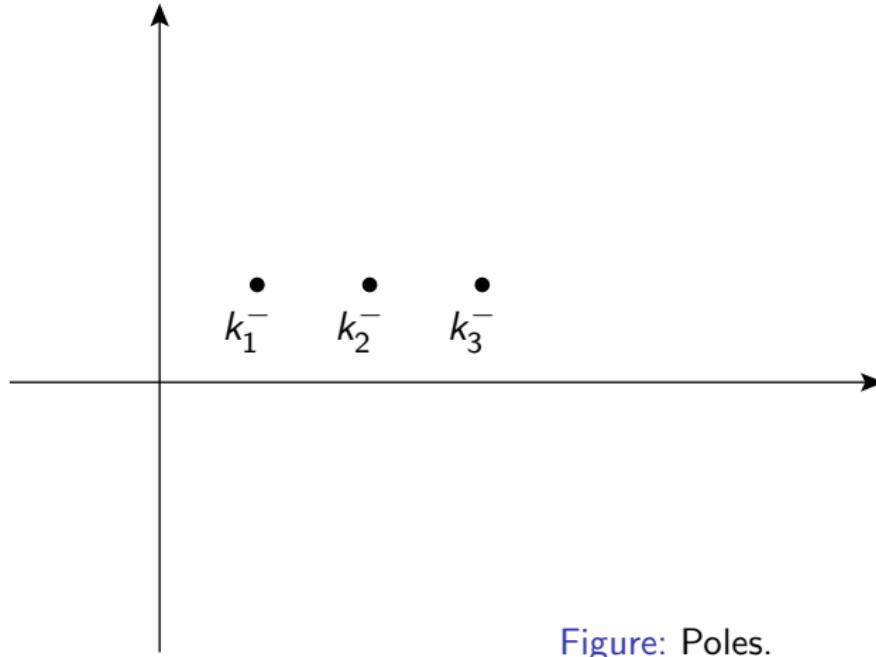


Figure: Poles.

- \Rightarrow **No contribuition to e.m. current!!**

ii) $k^+ > P^+$

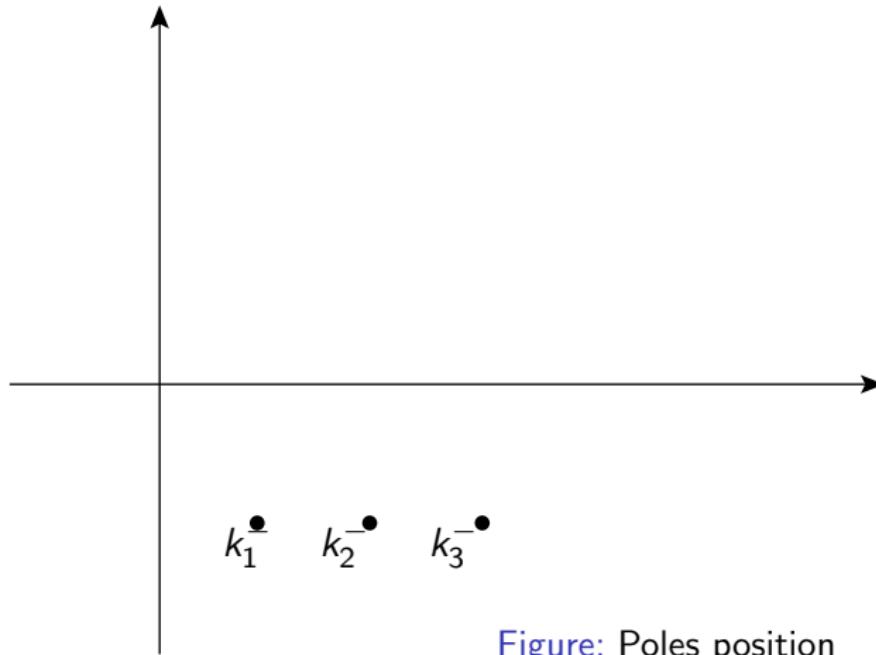


Figure: Poles position

• \Rightarrow No contribuition to e.m. current!!

iii) $0 < k^+ < P^+$

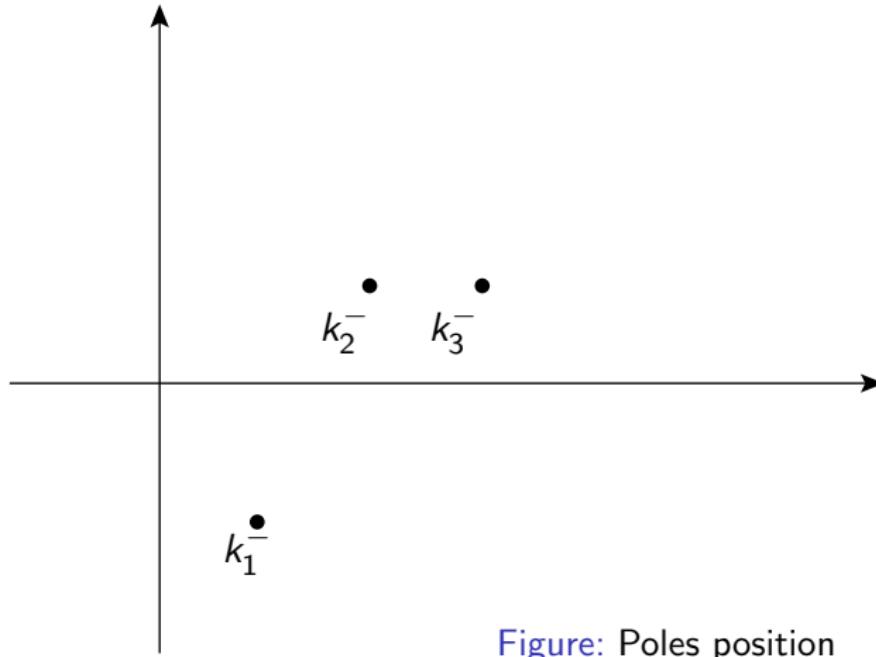


Figure: Poles position

* \implies Contribution to e.m. current!!

- k^- Integration

$$J^+ = \Gamma^2 \int \frac{d^2 k_\perp dk^+}{2(2\pi)^3} \frac{(2k - P' - P)^+ \theta(P^+ - k^+) \theta(k^+)}{k^+(P^+ - k^+)(P'^+ - k^+)} \\ \frac{1}{(P^- - \frac{f_1}{k^+} - \frac{f_2 - i\epsilon}{P^+ - k^+})(P'^- - \frac{f_1}{k^+} - \frac{f_3 - i\epsilon}{P'^+ - k^+})}$$

- $\Rightarrow x = \frac{k^+}{P^+}$, and $0 < x < 1$
- $P^+ P^- - P_\perp^2 = M_b^2 \iff M_b$ Bound state mass (for two bosons)

$$M_0^2 = \frac{f_1}{x} + \frac{f_2}{(1-x)} - P_\perp^2 \quad (2)$$

$$M_0'^2 = \frac{f_1}{x} + \frac{f_3}{(1-x)} - P_\perp'^2$$

$$J^+ = \Gamma^2 \int \frac{d^2 k_\perp dx}{2(2\pi)^3} \frac{(2xP - 2P)^+}{x(1-x)^2} \frac{\theta(x)\theta(1-x)}{(M_b^2 - M_0^2)(M_b^2 - M_0'^2)}$$

- Wave function

$$\Phi(x, \vec{k}_\perp) = \frac{\Gamma}{(M_b^2 - M_0^2)}$$

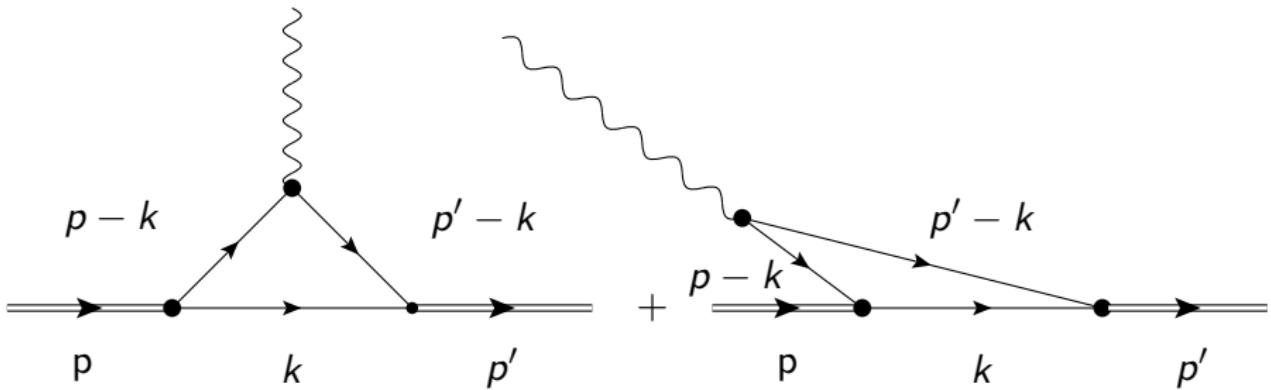
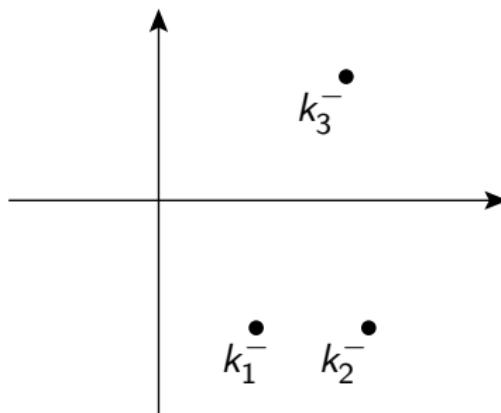


Figure: Feynman diagramms contribuition to the elastic electromagnetic form factor, Left: The valence component contribuition from the electromagnetic current to the electromagnetic form factor. Right panel: Possible non-valence contribuition for the electromagnetic current with frames differents from the Breit frame with the Drell-Yan condition.

Pole dislocation method*

- **Integral region:** $P^+ < k^+ < P'^+$, if $\underline{q^+} \neq 0$
 $\Rightarrow P'^+ = P^+ + q^+$, In the $\lim_{q^+ \rightarrow 0_+} P'^+ = P^+$.



- Integration in the pole k_3^- :

$$J^+ = \Gamma^2 \int \frac{d^2 k_\perp dk^+}{2(2\pi)^4} \frac{(2k - P' - P)^+ \theta(P'^+ - k^+) \theta(k^+ - P^+)}{k^+(P^+ - k^+)(P'^+ - k^+)(P'^- - \frac{f_3}{P'^+ - k^+} - \frac{f_1}{k^+})} \\ \frac{1}{(\frac{f_3}{P'^+ - k^+} - \frac{f_2}{P^+ - k^+})}$$

- The denominator:

$$\left[\frac{f_3}{P'^+ - k^+} - \frac{f_2}{P^+ - k^+} \right]^{-1}$$

- Three virtual particles propagation in $x^+ = 0$
- The denominator

$$\left[P'^- - \frac{f_3}{P'^+ - k^+} - \frac{f_1}{k^+} \right]^{-1}$$

- Is the propagator for two final state particles

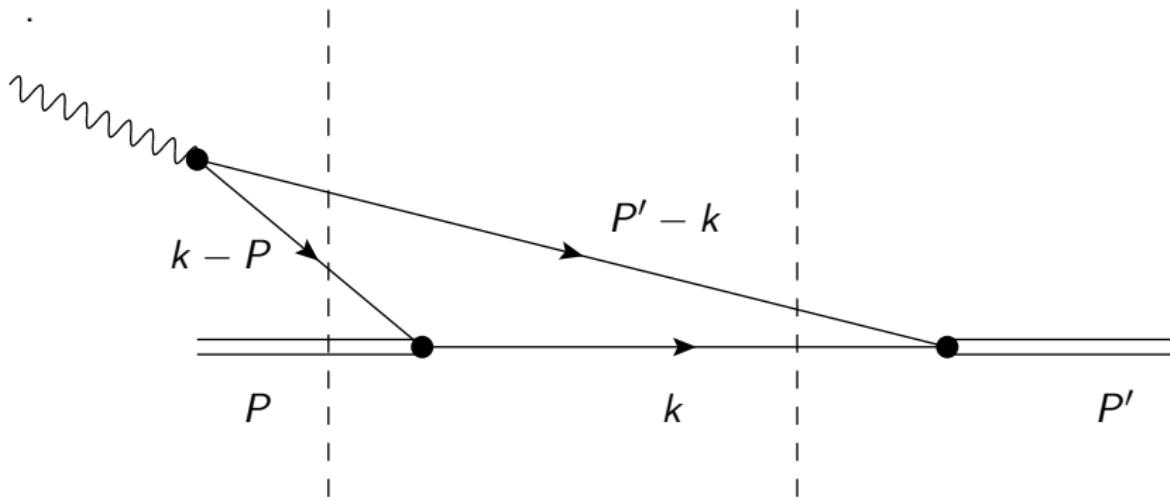


Figure: Pairs diagramm

- Ref. H. Naus, J. de Melo and T. Frederico, Few Body Syst. 24, 99 (1998)
- J. de Melo, J. Sales, T. Frederico and P. U. Sauer, Nucl. Phys. A 631, 574C (1998)

- If use the transformation $x = (k^+ - P^+)/q^+$

$$J^+ = \Gamma^2 \int \frac{d^2 k_\perp dx}{2(2\pi)^3} \frac{q^{+2}(2x-1)\theta(1-x)\theta(x)}{(P^+ + q^+ x)x(1-x)} \frac{1}{(\frac{f_3}{(1-x)} - q^+(\frac{f_1}{(P^+ + q^+ x)} - P'^-) (\frac{f_3}{1-x} + \frac{f_2}{x}))}$$

- Take the limit $q^+ \rightarrow 0_+$, we have $J^+ \propto q^{+2}$

$$\boxed{\lim_{q^+ \rightarrow 0_+} J^+ = 0}$$

Which means physically, for this component of the electromagnetic current the pair creation is suppressed

J_\perp Component of the e.m. current

$$J_\perp = \Gamma^2 \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)_\perp}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)}$$

i) $0 < k^+ < P^+ \implies k^{-(i)} = \frac{f_1 - i\epsilon}{k^+}$

$$\begin{aligned} J_\perp^i &= \Gamma^2 \int \frac{d^2 k_\perp dk^+}{2(2\pi)^3} \frac{(2k - P' - P)_\perp \theta(k^+) \theta(P^+ - k^+)}{k^+ (P^+ - k^+) (P'^+ - k^+)} \\ &\quad \frac{1}{(P^- - \frac{f_1}{k^+} - \frac{f_2 - i\epsilon}{P^+ - k^+})(P'^- - \frac{f_1}{k^+} - \frac{f_3 - i\epsilon}{P'^+ - k^+})} \end{aligned}$$

- Final result

$$J_\perp^i = \Gamma^2 \int \frac{d^2 k_\perp dx}{2(2\pi)^3} \frac{(2k - P' - P)_\perp \theta(x) \theta(1-x)}{x(1-x)^2 (M_b^2 - M_0^2)(M_b^2 - M_0'^2)} \neq 0$$

$$\text{ii) } P^+ < k^+ < P'^+ \implies k_3^- = P'^- - \frac{f_3 - \epsilon}{P'^+ - k^+}$$

- Use $x = (k^+ - P^+)/q^+$, then

$$\begin{aligned} J_\perp^{ii} &= \Gamma^2 \int \frac{d^2 k_\perp dk^+}{2(2\pi)^3} \frac{q^+ (2k - P' - P)_\perp \theta(1-x)\theta(x)}{(P^+ + q^+ x)x(1-x)} \\ &\quad \frac{1}{(-\frac{f_3}{1-x} - q^+ (\frac{f_1}{P^+ + q^+ x} - P'^-))(\frac{f_3}{1-x} + \frac{f_2}{x})} \propto q^+ \end{aligned}$$

$$\boxed{\lim_{q^+ \rightarrow 0_+} J^\perp = 0}$$

Again, the pair creation is suppressed!!

J^- Component of the e.m. current

$$J^- = \Gamma^2 \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)^-}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)}$$

- Integral intervals

i) $0 < k^+ < P^+ \implies k^{-(i)} = \frac{f_1 - i\epsilon}{k^+}$

ii) $P^+ < k^+ < P'^+ \implies k_3^- = P'^- - \frac{f_3 - i\epsilon}{P'^+ - k^+}$

- Results

$$J^{-i} = -i\Gamma^2 \int \frac{d^2 k_\perp dk^+}{2(2\pi)^3} \frac{[2f_1 - k^+(P'^- + P^-)] \theta(P^+ - k^+) \theta(k^+)}{k^{+2}(P^+ - k^+)(P'^+ - k^+)(P^- - \frac{f_1}{k^+} - \frac{f_2}{P^+ - k^+})}$$

$$\frac{1}{(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+ - k^+})} \neq 0$$

$$J^{-ii} = -i\Gamma^2 \int \frac{d^2 k_\perp dk^+}{2(2\pi)^3} \frac{[-2\frac{f_3}{P'^+ - k^+} + P'^- - P^-]}{k^+(P^+ - k^+)(P'^+ - k^+)} \frac{\theta(P'^+ - k^+) \theta(k^+ - P^+)}{(P'^- - \frac{f_3}{P'^+ - k^+} - \frac{f_1}{k^+})(\frac{f_3}{P'^+ - k^+} - \frac{f_2}{P^+ - k^+})}$$

- Final results for interval "(ii)" with $x = \frac{(k^+ - P^+)}{q^+}$

$$J^{-ii} = -\imath \Gamma^2 \int \frac{d^2 k_\perp dx}{2(2\pi)^3} \frac{-2f_3 + q^+(1-x)(P'^- - P^-)\theta(x)\theta(1-x)}{(P^+ + q^+x)x(1-x)^2} \frac{1}{(-\frac{f_3}{1-x} - q^+(\frac{f_1}{P^+ + q^+x} - P'^-))(\frac{f_3}{1-x} + \frac{f_2}{x})}$$

- In the limit $q^+ \rightarrow 0_+$

$$\begin{aligned} J^{-ii} &= -\imath \frac{\Gamma^2}{P^+} \int_0^1 \frac{d^2 k_\perp dx}{(2\pi)^3} \frac{1}{xf_3 + (1-x)f_2} \\ &= \boxed{\imath \frac{\Gamma^2}{P^+} \int \frac{d^2 k_\perp}{(2\pi)^3} \frac{\ln(f_2) - \ln(f_3)}{f_3 - f_2} \neq 0} \end{aligned}$$

- And

$$\boxed{J^+ = J^- = J^{-(i)} + J^{-(ii)}}$$

- Two bosons bound state

$$q_\mu \cdot J^\mu = \Sigma(P') - \Sigma(P)$$

- Electromagnetic current for two bosons bound state

$$q_\mu \cdot J^\mu = \int \frac{d^4 k}{(2\pi)^4} \frac{q_\mu \cdot (2k - P' - P)^\mu}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)}$$

- $\Rightarrow q^\mu = (\mathbf{P}'^\mu - \mathbf{P}^\mu - \mathbf{k}^\mu + \mathbf{k}^\mu)$

Then,

$$\begin{aligned} & -(P'^\mu - k^\mu - P^\mu + k^\mu) \cdot (P_\mu - k_\mu + P'_\mu - k_\mu) \\ &= -(P'^\mu - k^\mu)^2 + (P^\mu - k^\mu)^2 - m^2 + m^2 \end{aligned}$$

$$\int \frac{d^4 k}{(2\pi)^4} \frac{-(P'^{\mu} - k^{\mu})^2 + m^2 + (P^{\mu} - k^{\mu})^2 - m^2}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((P' - k)^2 - m^2 + i\epsilon)} =$$

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{(k^2 - m^2 + i\epsilon)((P' - k)^2 - m^2 + i\epsilon)} -$$

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{(k^2 - m^2 + i\epsilon)((P - k)^2 - m^2 + i\epsilon)}$$

Is the Ward-Takahashi again:

$$q_{\mu} \cdot J^{\mu} = \Sigma(P') - \Sigma(P)$$

- Self-energy

$$\Sigma(P) = \int \frac{d^4 k}{(2\pi)^4} \frac{1}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)}$$

- **Integration with the interval:** $P'^+ > P^+ > 0$

$$\Sigma(P) = \int \frac{d^2 k_\perp dk^+ dk^-}{2(2\pi)^4} \frac{1}{k^+(P^+ - k^+)(k^- - \frac{f_1 - i\epsilon}{k^+})(P^- - k^- - \frac{f_2 - i\epsilon}{P^+ - k^+})}$$

$$\begin{aligned}\Sigma(P) &= 2\pi i \int_0^{P^+} \frac{d^2 k_\perp dk^+}{2(2\pi)^4} \frac{1}{k^+(P^+ - k^+)(P^- - \frac{f_1}{k^+} - \frac{f_2}{P^+ - k^+})} \\ &= \pi i \int_0^1 \frac{d^2 k_\perp dx}{(2\pi)^4} \frac{1}{x(1-x)(P^2 - M_0^2 + i\epsilon)}\end{aligned}$$

- $x = \frac{k^+}{P^+}$, M_0^2 and $M_0'^2$,

$$M_0^2 = \frac{f_1}{x} + \frac{f_2}{(1-x)} - P_\perp^2$$

$$M_0'^2 = \frac{f_1}{x} + \frac{f_3}{(1-x)} - P_\perp'^2$$

$$\Sigma(P') - \Sigma(P) = \pi i \int_0^1 \frac{d^2 k_\perp dx}{(2\pi)^4} \frac{-P'^2 + P^2}{x(1-x)(P'^2 - M_0'^2)(P^2 - M_0^2)}$$

- Or writing as,

$$\Sigma(P') - \Sigma(P) = (P_\mu + P'_\mu)(P^\mu - P'^\mu) I_\Sigma$$

- I_Σ is,

$$I_\Sigma = \pi i \int d^2 k_\perp \int_0^1 \frac{dx}{(2\pi)^4} \frac{1}{x(1-x)(P'^2 - M_0'^2)(P^2 - M_0^2)}$$

- Electromagnetic current in the Light-front

$$q_\mu \cdot J^\mu = \frac{1}{2} (q^+ J^- + q^- J^+) - \vec{q}_\perp \cdot \vec{J}_\perp$$

$$J^- = \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)^-}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)} ,$$

$$J^+ = \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)^+}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)} ,$$

$$J_\perp = \int \frac{d^4 k}{(2\pi)^4} \frac{(2k - P' - P)_\perp}{(k^2 - m^2 + i\epsilon)((k - P)^2 - m^2 + i\epsilon)((k - P')^2 - m^2 + i\epsilon)}$$

Breit frame off-shell: $P^- \neq P'^-, P^+ \neq P'^+$.

$$\begin{aligned}
J^- &= -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P^+} dk^+ \frac{(2\frac{f_1}{k^+} - P'^- - P^-)(P'^- - f_1 - \frac{f_3}{P'^+-k^+})^{-1}}{k^+(P^+ - k^+)(P'^+ - k^+)(P^- - f_1 - \frac{f_2}{P^+-k^+})} \\
&\quad \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_{P^+}^{P'^+} dk^+ \frac{(-2\frac{f_3}{P'^+-k^+} + P'^- - P^-)(P^- - P'^- + \frac{f_3}{P'^+-k^+} - \frac{f_2}{P^+-k^+})^{-1}}{k^+(P^+ - k^+)(P'^+ - k^+)(P'^- - \frac{f_3}{P'^+-k^+} - \frac{f_1}{k^+})} \\
J^+ &= -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P^+} dk^+ \frac{(2k^+ - P'^+ - P^+)(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+-k^+})^{-1}}{k^+(P^+ - k^+)(P'^+ - k^+)(P^- - f_1 - \frac{f_2}{P^+-k^+})} - \\
&\quad \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_{P^+}^{P'^+} dk^+ \frac{(2k^+ + P'^+ - P^+)(P^- - P'^- + \frac{f_3}{P'^+-k^+} - \frac{f_2}{P^+-k^+})^{-1}}{k^+(P^+ - k^+)(P'^+ - k^+)(P'^- - \frac{f_3}{P'^+-k^+} - \frac{f_1}{k^+})} \\
J_\perp &= -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P^+} dk^+ \frac{(2k_\perp - P'_\perp - P_\perp)(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+-k^+})^{-1}}{k^+(P^+ - k^+)(P'^+ - k^+)(P^- - f_1 - \frac{f_2}{P^+-k^+})} - \\
&\quad \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_{P^+}^{P'^+} dk^+ \frac{(2k_\perp - P'_\perp - P_\perp)(P^- - P'^- + \frac{f_3}{P'^+-k^+} - \frac{f_2}{P^+-k^+})^{-1}}{k^+(P^+ - k^+)(P'^+ - k^+)(P'^- - \frac{f_3}{P'^+-k^+} - \frac{f_1}{k^+})}
\end{aligned}$$

- q^μ as, $q^\mu = P'^\mu - P^\mu + \bar{k}^\mu - \bar{k}'^\mu$
and in the interval (i) $\bar{k}^- = \frac{f_1}{k^+}$, $\bar{k}^+ = k_\perp$.

$$[(P - \bar{k})^2 - m^2] - [(P' - \bar{k})^2 - m^2]$$

$$= (P^+ - k^+)[P^- - \frac{f_1}{k^+} - \frac{f_2}{P^+ - k^+}] - (P'^+ - k^+)[P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+ - k^+}]$$

- Final for interval (i)

$$I_i = -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P^+} dk^+ \frac{1}{k^+(P'^+ - k^+)(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+ - k^+})} +$$

$$\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P^+} dk^+ \frac{1}{k^+(P^+ - k^+)(P^- - \frac{f_1}{k^+} - \frac{f_2}{P^+ - k^+})}$$

- **Interval (ii)** $\bar{k}'^- = P'^- - \frac{f_3}{P'^+ - k^+}$, $\bar{k}'^+ = k^+$ and $\bar{k}'_\perp = k_\perp$

$$q^\mu = (P'^\mu - P^\mu) = (P' - \bar{k}')^\mu - (P - \bar{k}')^\mu$$

- And with on-shell condition

$$(P'^\mu - \bar{k}'^\mu)^2 - m^2 = 0$$

- We have below,

$$\begin{aligned} q^\mu \cdot (2\bar{k}'^\mu - P^\mu - P'^\mu) &= \\ -(P'^\mu - \bar{k}'^\mu - P^\mu + \bar{k}'^\mu)(P'^\mu - \bar{k}'^\mu + P^\mu - \bar{k}'^\mu) &= \\ -((P' - \bar{k}')^2 - m^2) + ((P - \bar{k}')^2 - m^2) &= \\ -(P - \bar{k}')^2 + m^2 & \end{aligned}$$

- With Light-front coordinates,

$$(P - \bar{k}')^2 - m^2 = (P^+ - k^+) \left[P^- - \frac{f_2}{P^+ - k^+} - P'^- + \frac{f_3}{P'^+ - k^+} \right]$$

- Final results to interval (ii)

$$I_{ii} = -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_{P^+}^{P'^+} dk^+ \frac{1}{k^+(P'^+ - k^+)(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+ - k^+})}$$

- The sum (i) +(ii) is

$$\begin{aligned} I = I_i + I_{ii} &= -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P'^+} dk^+ \frac{1}{k^+(P'^+ - k^+)(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+ - k^+})} \\ &\quad + \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^{P^+} dk^+ \frac{1}{k^+(P^+ - k^+)(P^- - \frac{f_1}{k^+} - \frac{f_3}{P^+ - k^+})} \\ &\quad - \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_{P^+}^{P'^+} dk^+ \frac{1}{k^+(P'^+ - k^+)(P'^- - \frac{f_1}{k^+} - \frac{f_3}{P'^+ - k^+})} \end{aligned}$$

- Working more, we have the final Ward-Takahashi ,

$$\begin{aligned}
 I_i + I_{ii} = I = & -\pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^1 dx \frac{1}{x(1-x)} \frac{1}{(P'^2 - M_0'^2)} + \\
 & \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^1 dx \frac{1}{x(1-x)} \frac{1}{(P^2 - M_0^2)} \\
 = & \pi i \int \frac{d^2 k_\perp}{(2\pi)^4} \int_0^1 dx \frac{1}{x(1-x)} \frac{(P^2 - P'^2)}{(P'^2 - M_0'^2)(P^2 - M_0^2)}
 \end{aligned}$$

With the conditions, $P'^2 = P^2 = M_b^2$.

If zero-modes, or non-valence contributions, is not included with Light-front calculations, WT is not satisfied!!

- Some works

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