



# ALPHA<sub>S</sub> from Lattice QCD

by

J. Rodríguez-Quintero

Dpto. Física Aplicada, Fac. CCEE  
Universidad de Huelva, Spain

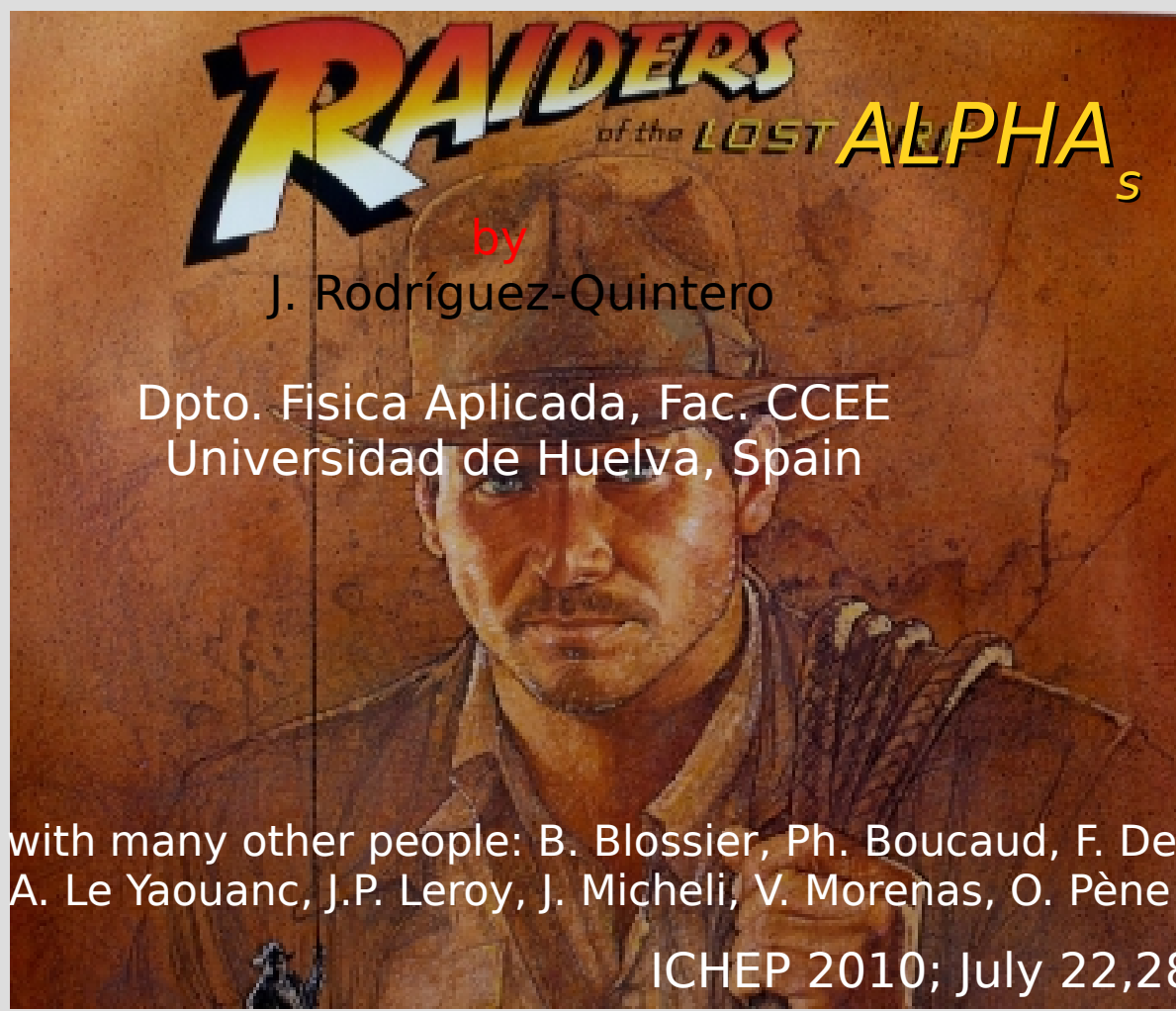


In collaboration with many other people: B. Blossier, Ph. Boucaud, F. De Soto, M. Gravina, A. Le Yaouanc, J.P. Leroy, J. Micheli, V. Morenas, O. Pène

ICHEP 2010; July 22,28; Paris 2010.



## Lattice QCD and



Dpto. Física Aplicada, Fac. CCEE  
Universidad de Huelva, Spain

In collaboration with many other people: B. Blossier, Ph. Boucaud, F. De Soto, M. Gravina, A. Le Yaouanc, J.P. Leroy, J. Micheli, V. Morenas, O. Pène

ICHEP 2010; July 22,28; Paris 2010.

# QCD: The running of $\text{ALPHA}_s$

$$\mathcal{L}_{\text{QCD}} = \sum_q (i\bar{q}\not{D}q - m_q\bar{q}q) - \frac{1}{4}(D \times A)^2$$

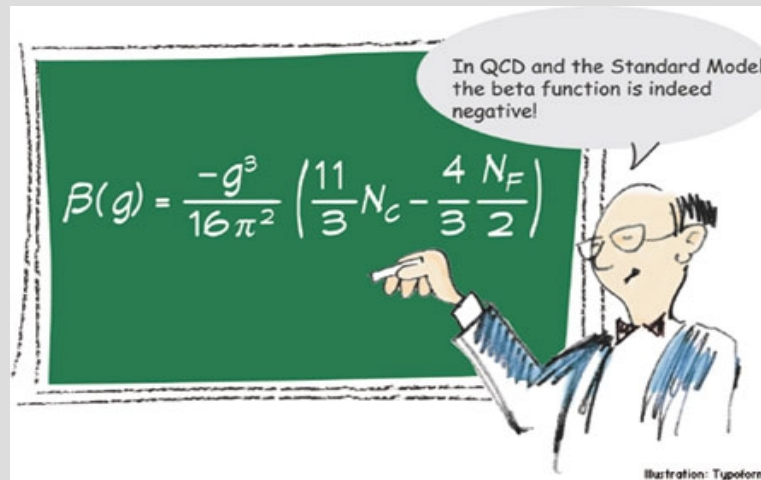
$$(D^\mu \times A^\nu)_a \equiv \partial^\mu A_a^\nu - \partial^\nu A_a^\mu + g \sum f_{abc} A_b^\mu A_c^\nu$$

$$[\lambda^a, \lambda^b] = if_{abc} \lambda^c$$

# QCD: The running of $\text{ALPHA}_s$

$$\mathcal{L}_{\text{QCD}} = \sum_q (i\bar{q}\not{D}q - m_q\bar{q}q) - \frac{1}{4}(D \times A)^2$$

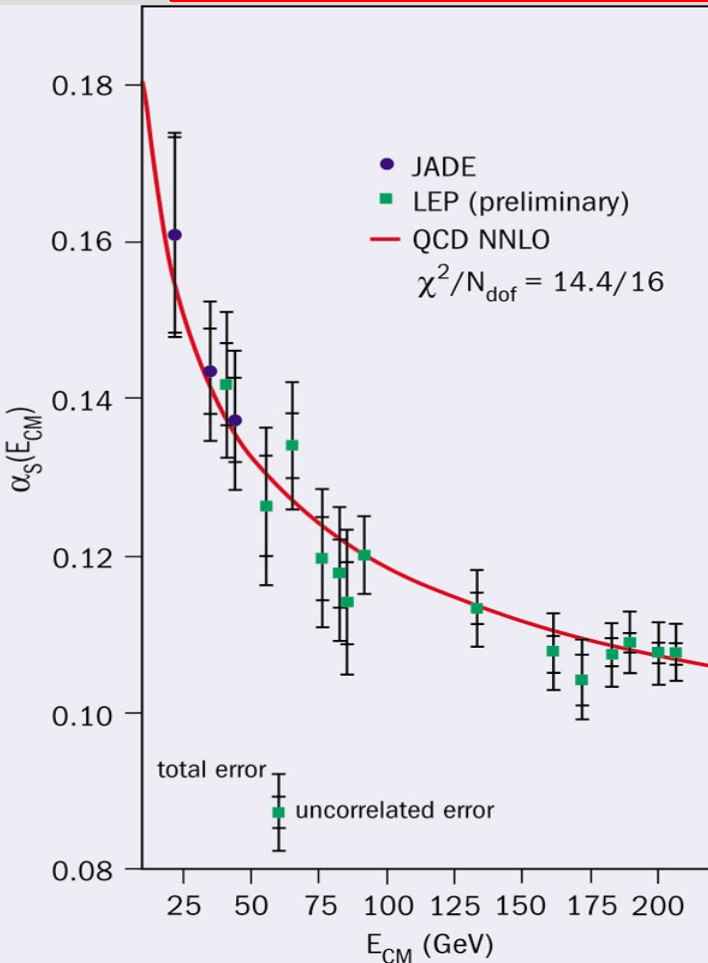
$$(D^\mu \times A^\nu)_a \equiv \partial^\mu A_a^\nu - \partial^\nu A_a^\mu + g \sum f_{abc} A_b^\mu A_c^\nu$$



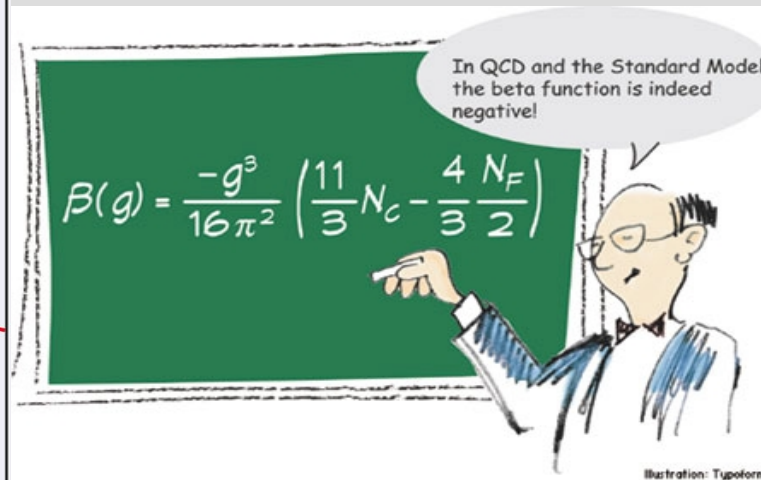
$$[\lambda^a, \lambda^b] = if_{abc} \lambda^c$$

# QCD: The running of $\text{ALPHA}_S$

$$\mathcal{L}_{\text{QCD}} = \sum_q (i\bar{q}\not{D}q - m_q\bar{q}q) - \frac{1}{4}(D \times A)^2$$



$$(D^\mu \times A^\nu)_a \equiv \partial^\mu A_a^\nu - \partial^\nu A_a^\mu + g \sum f_{abc} A_b^\mu A_c^\nu$$

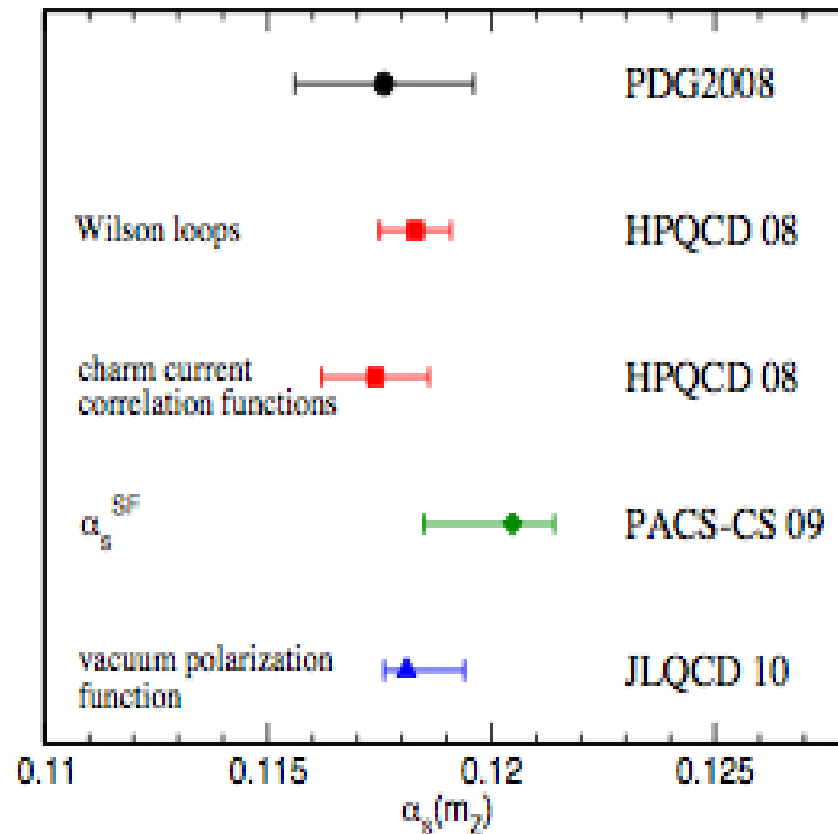


$$[\lambda^a, \lambda^b] = if_{abc} \lambda^c$$

# ALPHA<sub>S</sub> from Lattice QCD:

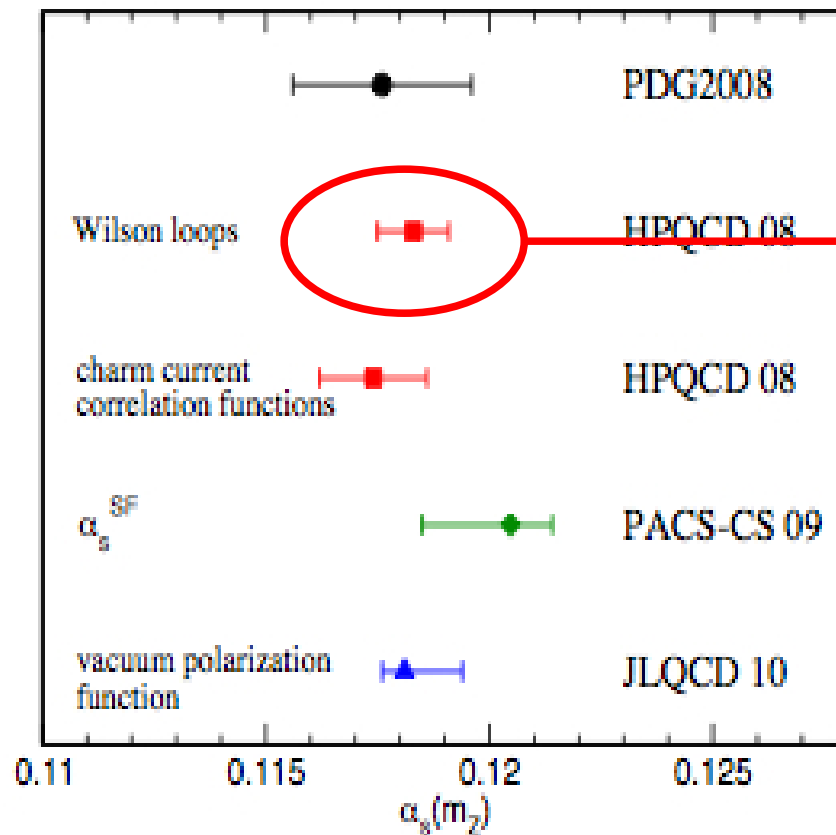
the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:



# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

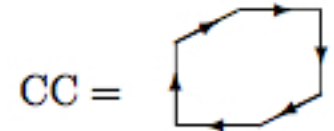
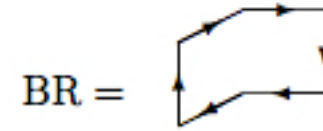
Very recent  $N_f=2+1$  &  $N_f=4$  computations:



C.T.H. Davies et al.;  
Phys. Rev. D78(2008)114507

# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:



C.T.H. Davies et al.;  
Phys. Rev. D78(2008)114507

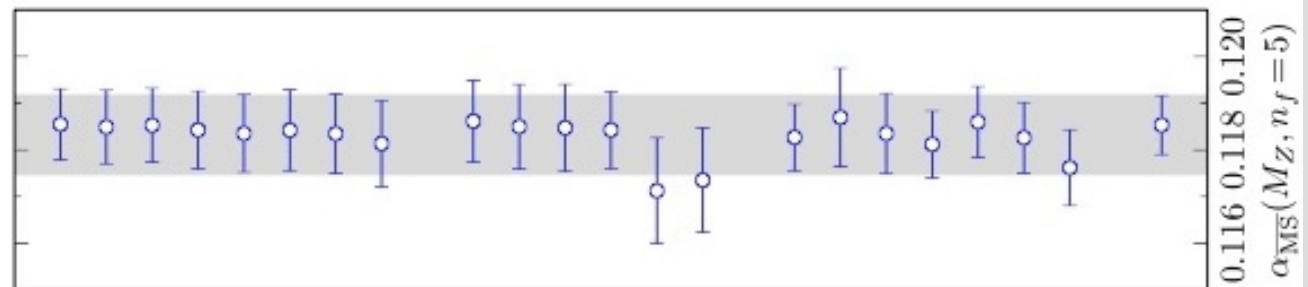
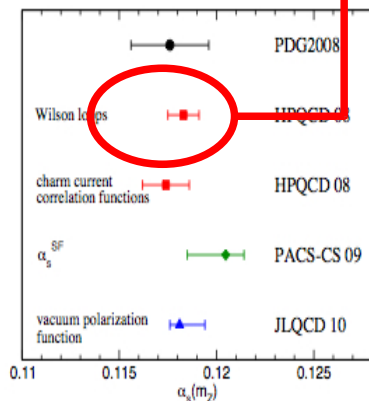
$$W_{mn} \equiv \frac{1}{3} \langle 0 | \text{Re Tr P } e^{-ig \oint_{nm} A \cdot dx} | 0 \rangle,$$

$$q^2 \frac{d\alpha_V(q)}{dq^2} = -\beta_0 \alpha_V^2 - \beta_1 \alpha_V^3 - \beta_2 \alpha_V^4 - \beta_3 \alpha_V^5$$

$$Y = \sum_{n=1}^{\infty} c_n \alpha_V^n(d/a)$$

$$\alpha_0 \equiv \alpha_V(7.5 \text{ GeV}, n_f=3)$$

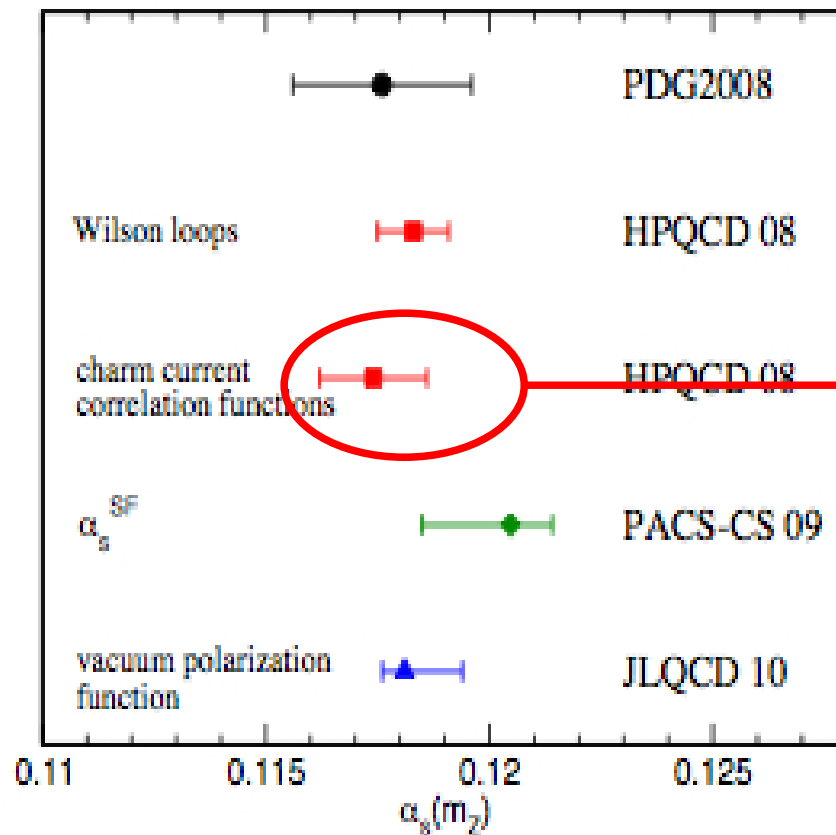
$$\alpha_{\overline{\text{MS}}}(M_Z, n_f=5)$$





# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:



I. Allison et al.;  
Phys. Rev. D78(2008)054513

# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

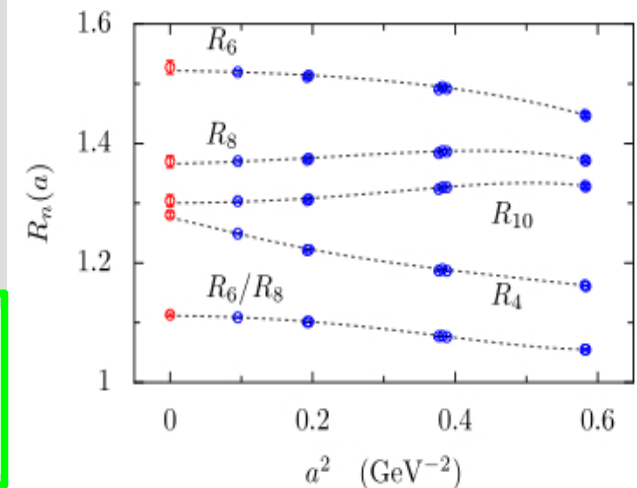
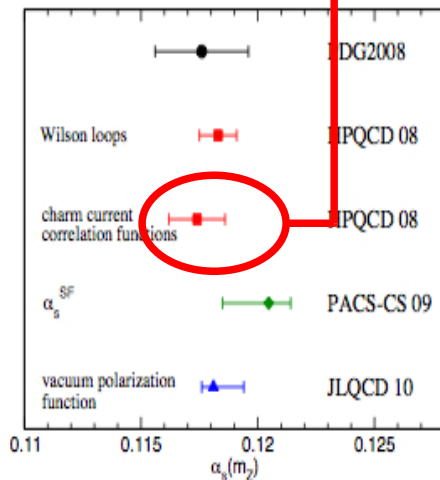
Very recent  $N_f=2+1$  &  $N_f=4$  computations:

I. Allison et al.;  
Phys. Rev. D78(2008)054513

$$G(t) \equiv a^6 \sum_{\mathbf{x}} (am_{0c})^2 \langle 0 | j_5(\mathbf{x}, t) j_5(0, 0) | 0 \rangle$$

$$G_n \equiv \sum_t (t/a)^n G(t),$$

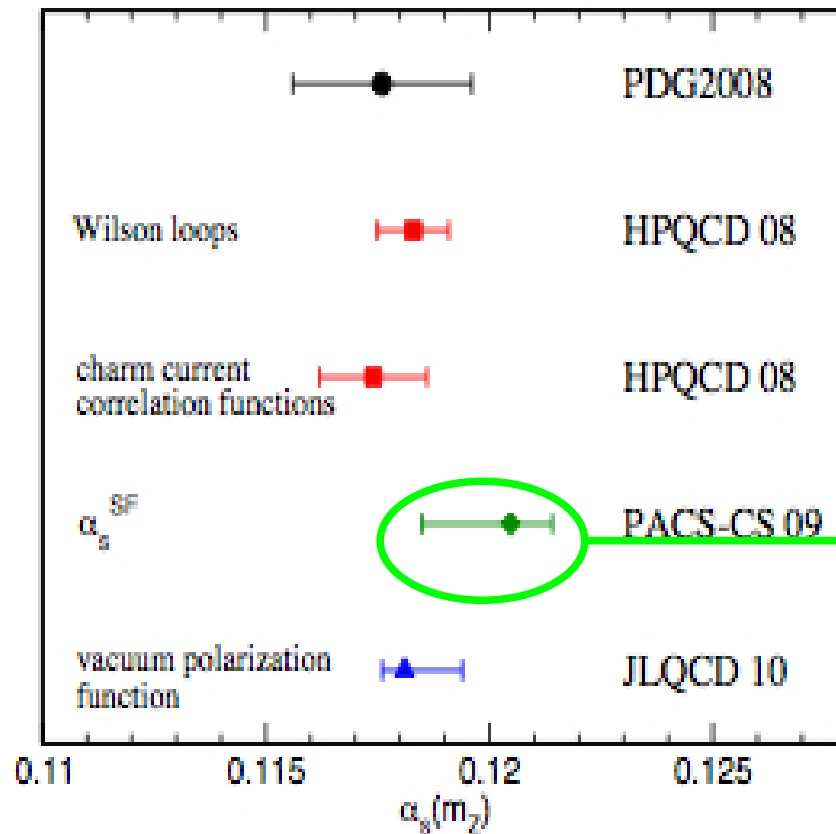
$$R_n \equiv \begin{cases} G_4/G_4^{(0)} & \text{for } n = 4, \\ \frac{am_{\eta_c}}{2am_{\text{pole},c}}^{(0)} \left( G_n/G_n^{(0)} \right)^{1/(n-4)} & \text{for } n \geq 6, \end{cases}$$



$$R_n \equiv \begin{cases} r_4(\alpha_{\overline{\text{MS}}}, \mu/m_c) & \text{for } n = 4, \\ \frac{r_n(\alpha_{\overline{\text{MS}}}, \mu/m_c)}{2m_c(\mu)/m_{\eta_c}} & \text{for } n \geq 6, \end{cases}$$

# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:



S. Aoki et al.;  
JHEP10(2009)053

# ALPHA<sub>S</sub> from Lattice QCD:

the many raiders...

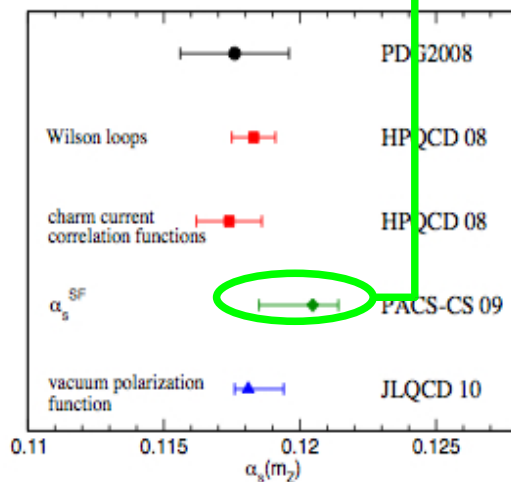
Very recent  $N_f=2+1$  &  $N_f=4$  computations:

S. Aoki et al.;  
JHEP10(2009)053

$$S_g = \frac{\beta}{N} \sum_{C \in S_0} W_0(C, g_0^2) \text{Re tr} (1 - P(C)) + \frac{\beta}{N} \sum_{C \in S_1} W_1(C, g_0^2) \text{Re tr} (1 - R(C)),$$

$$S_f[U, \psi, \bar{\psi}] = a^4 \sum_x \bar{\psi} (D_W + m_0) \psi,$$

$$D_W = \frac{1}{2} \left( \gamma_\mu (\nabla_\mu + \nabla_\mu^*) - a \nabla_\mu^* \nabla_\mu \right) - c_{SW} \frac{1}{4} \sigma_{\mu\nu} P_{\mu\nu}.$$



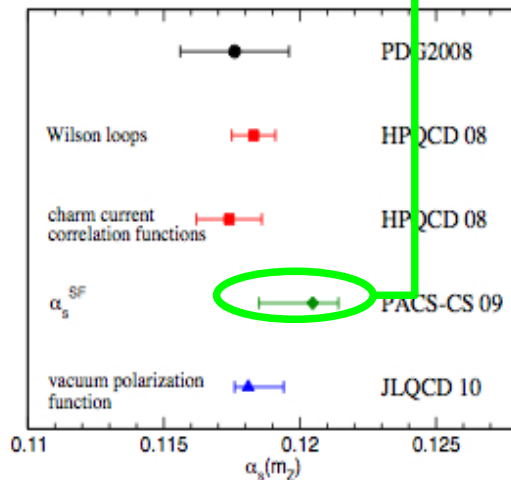
- Iwasaki action for the gauge fields
- Wilson Clover for the fermions

# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:

S. Aoki et al.;  
JHEP10(2009)053

$$\Lambda_{\text{SF}} = \frac{1}{L} (b_0 \bar{g}(L))^{-\frac{b_1}{2b_0^2}} \exp\left(-\frac{1}{2b_0 \bar{g}(L)}\right) \exp\left(-\int_0^{\bar{g}(L)} dg \left(\frac{1}{\beta(g)} + \frac{1}{b_0 g^3} - \frac{b_1}{b_0^2 g}\right)\right),$$



# ALPHA<sub>S</sub> from Lattice QCD:

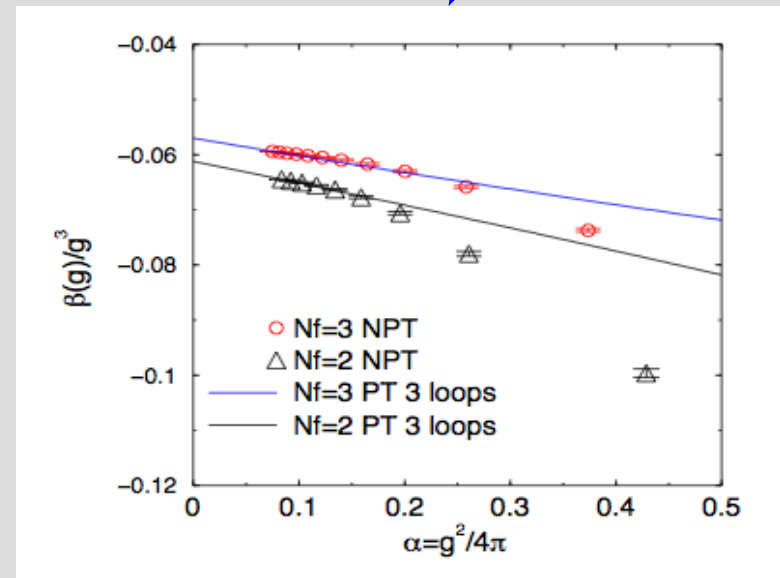
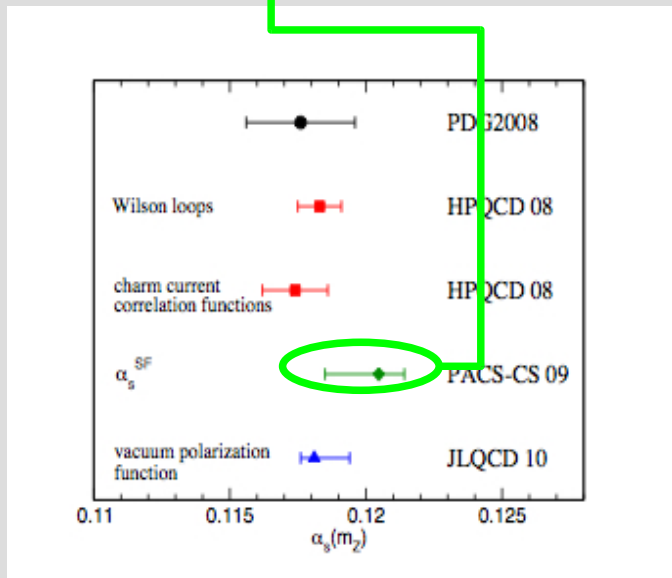
the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:

S. Aoki et al.;  
JHEP10(2009)053

$$\frac{1}{\bar{g}^2(L)} = \frac{1}{k} \left. \frac{\partial \Gamma[V_\mu]}{\partial \eta} \right|_{\eta=0},$$

$$\Lambda_{\text{SF}} = \frac{1}{L} (b_0 \bar{g}(L))^{-\frac{b_1}{2b_0^2}} \exp\left(-\frac{1}{2b_0 \bar{g}(L)}\right) \exp\left(-\int_0^{\bar{g}(L)} dg \left(\beta(g) + \frac{1}{b_0 g^3} - \frac{b_1}{b_0^2 g}\right)\right),$$



# ALPHA<sub>S</sub> from Lattice QCD:

the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:

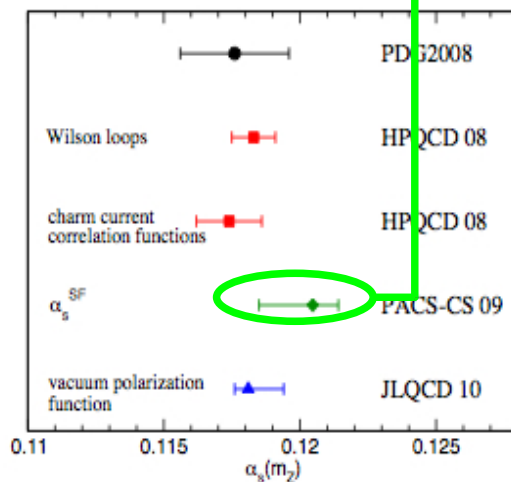
S. Aoki et al.;  
JHEP10(2009)053

$$\frac{1}{\bar{g}^2(L)} = \frac{1}{k} \left. \frac{\partial \Gamma[V_\mu]}{\partial \eta} \right|_{\eta=0},$$

$$\Lambda_{\text{SF}} = \frac{1}{L} (b(\bar{g}(L)))^{-\frac{b_1}{2b_0^2}} \exp\left(-\frac{1}{2b(\bar{g}(L))}\right) \exp\left(-\int_0^{\bar{g}(L)} dg \left(\frac{1}{\beta(g)} + \frac{1}{b_0 g^3} - \frac{b_1}{b_0^2 g}\right)\right),$$

$$\Sigma\left(u, \frac{a}{L}\right) = \bar{g}^2(2L) \Big|_{u=\bar{g}^2(L)}.$$

$$\sigma(u) = \lim_{a/L \rightarrow 0} \Sigma\left(u, \frac{a}{L}\right),$$



# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:

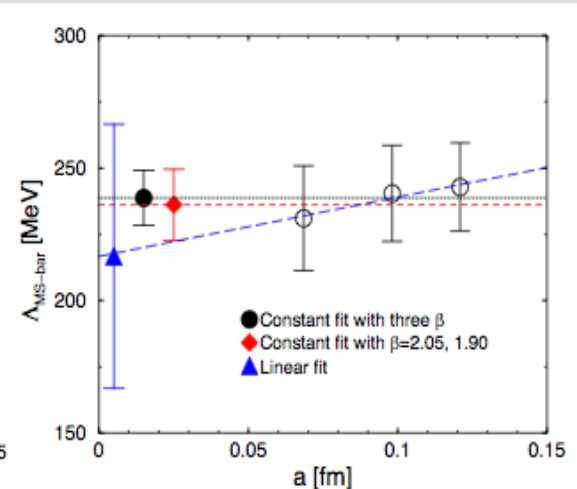
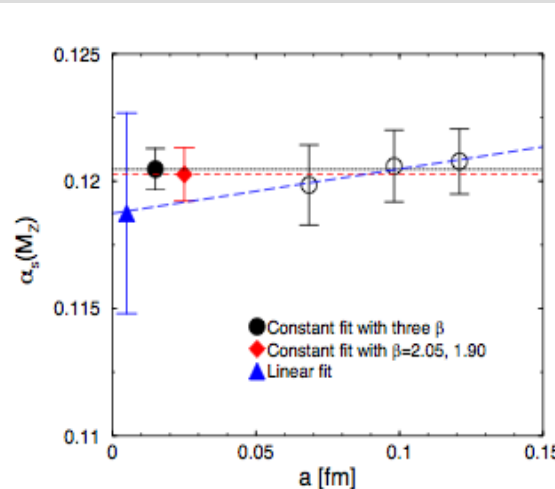
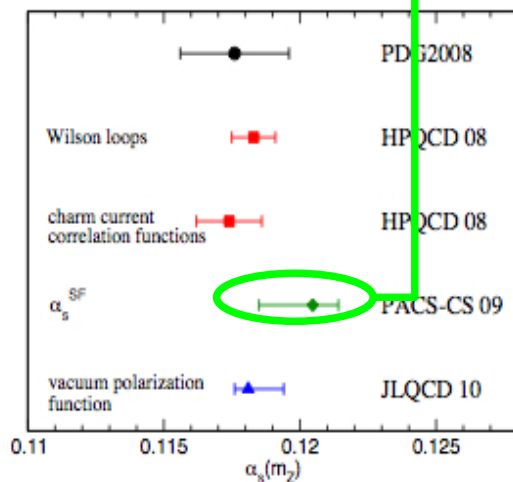
S. Aoki et al.;  
JHEP10(2009)053

$$\frac{1}{\bar{g}^2(L)} = \frac{1}{k} \left. \frac{\partial \Gamma[V_\mu]}{\partial \eta} \right|_{\eta=0},$$

$$\Lambda_{\text{SF}} = \frac{1}{L} (b(\bar{g}(L)))^{-\frac{b_1}{2b_0^2}} \exp\left(-\frac{1}{2b(\bar{g}(L))}\right) \exp\left(-\int_0^{\bar{g}(L)} dg \left(\frac{1}{\beta(g)} + \frac{1}{b_0 g^3} - \frac{b_1}{b_0^2 g}\right)\right),$$

$$\Sigma\left(u, \frac{a}{L}\right) = \bar{g}^2(2L) \Big|_{u=\bar{g}^2(L)}.$$

$$\sigma(u) = \lim_{a/L \rightarrow 0} \Sigma\left(u, \frac{a}{L}\right),$$





# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:

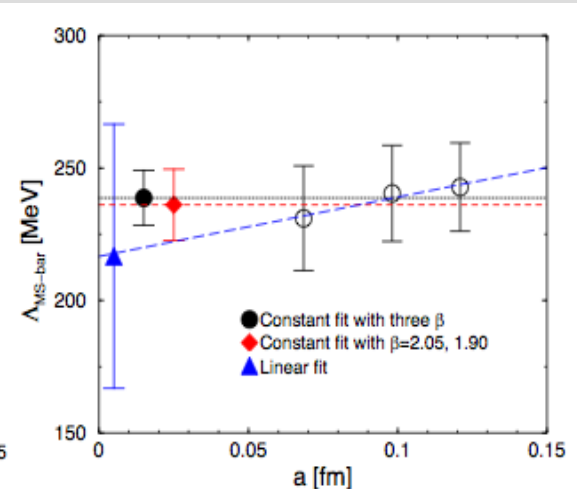
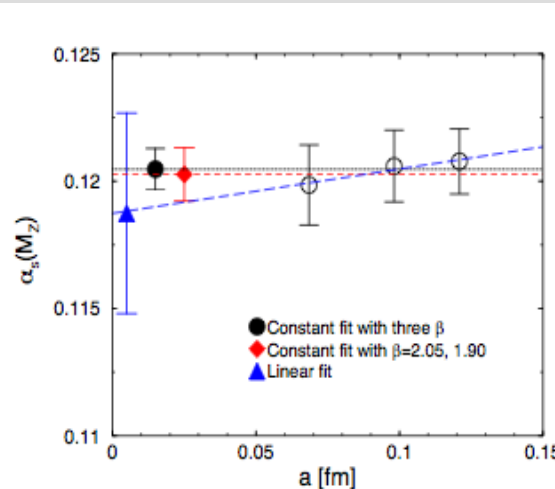
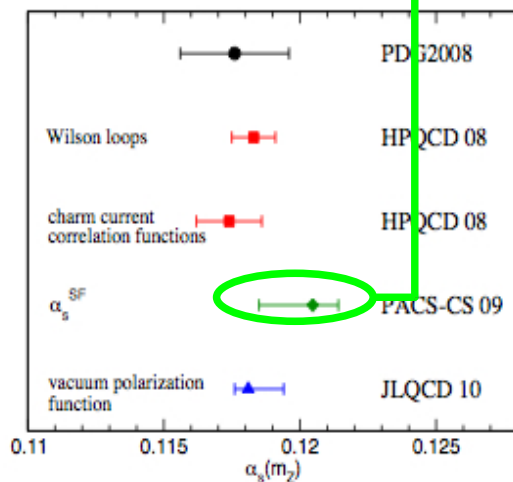
S. Aoki et al.;  
JHEP10(2009)053

$$\frac{1}{\bar{g}^2(L)} = \frac{1}{k} \left. \frac{\partial \Gamma[V_\mu]}{\partial \eta} \right|_{\eta=0},$$

$$\Lambda_{\text{SF}} = \frac{1}{L} (b(\bar{g}(L)))^{-\frac{b_1}{2b_0^2}} \exp\left(-\frac{1}{2b(\bar{g}(L))}\right) \exp\left(-\int_0^{\bar{g}(L)} dg \left(\frac{1}{\beta(g)} + \frac{1}{b_0 g^3} - \frac{b_1}{b_0^2 g}\right)\right),$$

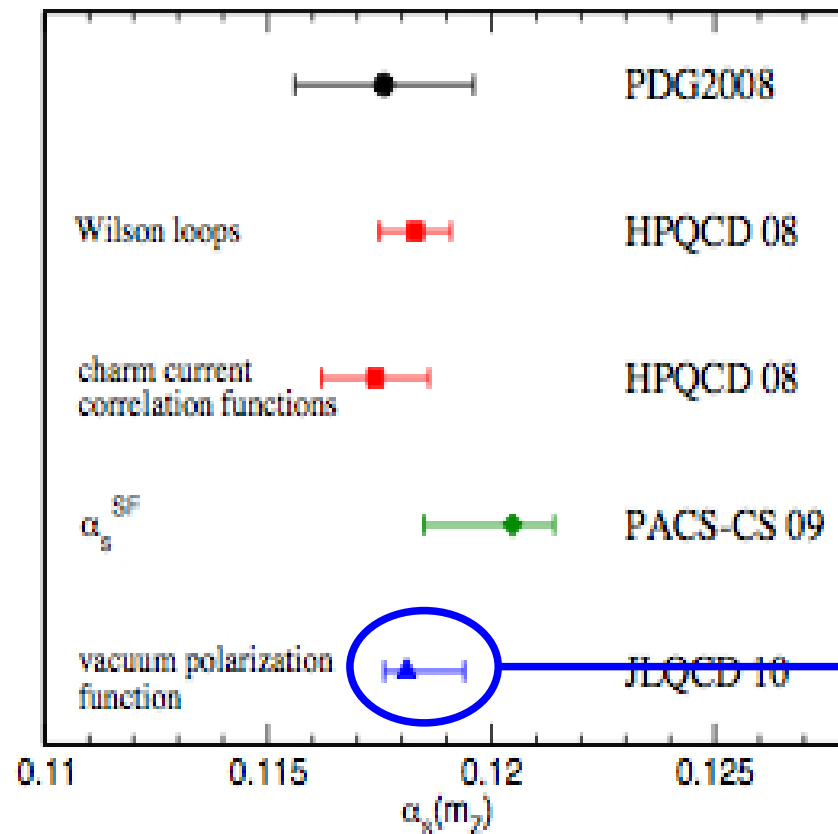
$$\Sigma\left(u, \frac{a}{L}\right) = \bar{g}^2(2L) \Big|_{u=\bar{g}^2(L)}.$$

$$\sigma(u) = \lim_{a/L \rightarrow 0} \Sigma\left(u, \frac{a}{L}\right),$$



# ALPHA<sub>S</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:



E. Shintani et al.;  
arXiv:1002.0371

# ALPHA<sub>s</sub> from Lattice QCD: the many raiders...

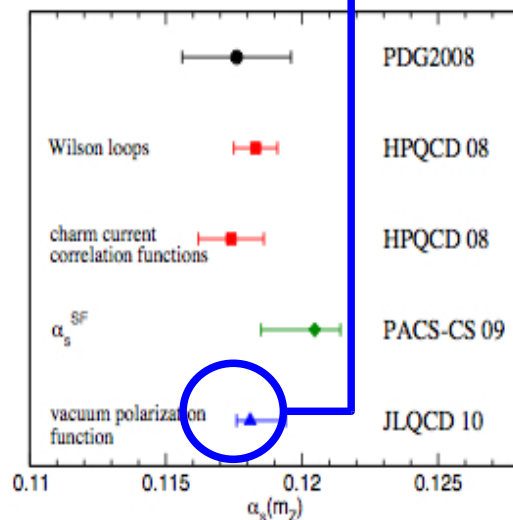
Very recent  $N_f=2+1$  &  $N_f=4$  computations:

E. Shintani et al.;  
arXiv:1002.0371

- Overlap fermions

$$\langle J_\mu^a J_\nu^b \rangle(Q) = \delta^{ab} [(\delta_{\mu\nu} Q^2 - Q_\mu Q_\nu) \Pi_J^{(1)}(Q) - Q_\mu Q_\nu \Pi_J^{(0)}(Q)].$$

$$\begin{aligned} \Pi_{V+A}|_{\text{OPE}}(Q^2, \alpha_s) &= c + C_0(Q^2, \mu^2, \alpha_s) \\ &+ C_m^{V+A}(Q^2, \mu^2, \alpha_s) \frac{\bar{m}^2(Q)}{Q^2} \\ &+ \sum_{q=u,d,s} C_{\bar{q}q}^{V+A}(Q^2, \alpha_s) \frac{\langle m_{\bar{q}q} \rangle}{Q^4} \\ &+ C_{GG}(Q^2, \alpha_s) \frac{\langle (\alpha_s/\pi) GG \rangle}{Q^4} + \mathcal{O}(Q^{-6}). \end{aligned}$$



# ALPHA<sub>s</sub> from Lattice QCD: the many raiders...

Very recent  $N_f=2+1$  &  $N_f=4$  computations:

E. Shintani et al.;  
arXiv:1002.0371

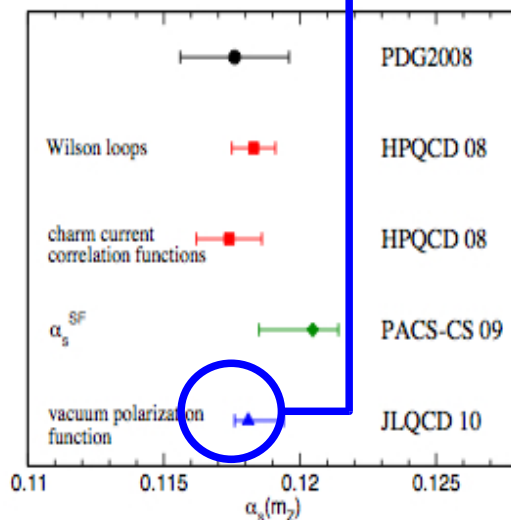
- Overlap fermions

$$\langle J_\mu^a J_\nu^b \rangle(Q) = \delta^{ab} [(\delta_{\mu\nu} Q^2 - Q_\mu Q_\nu) \Pi_J^{(1)}(Q) - Q_\mu Q_\nu \Pi_J^{(0)}(Q)].$$

$$\begin{aligned} \Pi_{V+A}|_{\text{OPE}}(Q^2, \alpha_s) = & c + C_0(Q^2, \mu^2, \alpha_s) \\ & + C_m^{V+A}(Q^2, \mu^2, \alpha_s) \frac{\bar{m}^2(Q)}{Q^2} \\ & + \sum_{q=u,d,s} C_{\bar{q}q}^{V+A}(Q^2, \alpha_s) \frac{\langle m\bar{q}q \rangle}{Q^4} \\ & + C_{GG}(Q^2, \alpha_s) \frac{\langle (\alpha_s/\pi) GG \rangle}{Q^4} + \mathcal{O}(Q^{-6}). \end{aligned}$$

Unknown parameters to be fitted!!

$$\Lambda_{\overline{MS}}^{(3)} = 0.247(5)$$









# ALPHA<sub>s</sub> from Lattice QCD:

## The ghost-gluon coupling

The ghost-gluon vertex:

$$\tilde{\Gamma}_{\nu}^{abc}(-q, k; q - k) = \begin{array}{c} \text{---} \xrightarrow{k=0} \text{---} \\ \text{---} \xrightarrow{q-k} \text{---} \\ \text{---} \xrightarrow{q} \text{---} \end{array} = g_0^{abc} (q_{\nu} H_1(q, k) + (q - k)_{\nu} H_2(q, k))$$

$\tilde{\Gamma}_R = \tilde{Z}_1 \Gamma$

The strong coupling:

$$g_R(\mu^2) = \lim_{\Lambda \rightarrow \infty} Z_g^{-1}(\mu^2, \Lambda^2) g_0(\Lambda^2) = \lim_{\Lambda \rightarrow \infty} \frac{Z_3^{1/2}(\mu^2, \Lambda^2) \tilde{Z}_3(\mu^2, \Lambda^2)}{\tilde{Z}_1(\mu^2, \Lambda^2)} g_0(\Lambda^2)$$

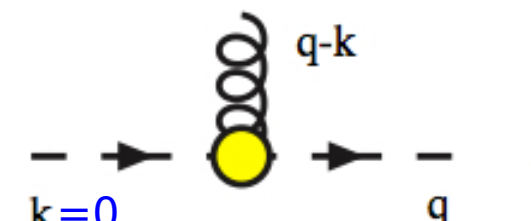
In Taylor scheme



# ALPHA<sub>S</sub> from Lattice QCD:

## The ghost-gluon coupling

The ghost-gluon vertex:

$$\tilde{\Gamma}_{\nu}^{abc}(-q, k; q - k) = \underset{k=0}{\text{---}} \rightarrow \text{---} \rightarrow \text{---} = g_0^{abc} (q_{\nu} H_1(q, k) + (q - k)_{\nu} H_2(q, k))$$


$$\tilde{\Gamma}_R = \tilde{Z}_1 \Gamma$$

The strong coupling:

$$g_R(\mu^2) = \lim_{\Lambda \rightarrow \infty} Z_g^{-1}(\mu^2, \Lambda^2) g_0(\Lambda^2) = \lim_{\Lambda \rightarrow \infty} \frac{Z_3^{1/2}(\mu^2, \Lambda^2) \tilde{Z}_3(\mu^2, \Lambda^2)}{\tilde{Z}_1(\mu^2, \Lambda^2)} g_0(\Lambda^2)$$

In Taylor scheme & Landau gauge



# ALPHA<sub>S</sub> from Lattice QCD:

## The ghost-gluon coupling

The ghost-gluon vertex:

$$\tilde{\Gamma}_\nu^{abc}(-q, k; q - k) = \begin{array}{c} \text{---} \xrightarrow{k=0} \text{---} \\ \text{---} \xrightarrow{q-k} \text{---} \\ \text{---} \xrightarrow{q} \text{---} \end{array} = g_0^{abc} (q_\nu H_1(q, k) + (q - k)_\nu H_2(q, k))$$

$\tilde{\Gamma}_R = \tilde{Z}_1 \Gamma$

The strong coupling:

$$g_R(\mu^2) = \lim_{\Lambda \rightarrow \infty} Z_g^{-1}(\mu^2, \Lambda^2) g_0(\Lambda^2) = \lim_{\Lambda \rightarrow \infty} \frac{Z_3^{1/2}(\mu^2, \Lambda^2) \tilde{Z}_3(\mu^2, \Lambda^2)}{\tilde{Z}_1(\mu^2, \Lambda^2)} g_0(\Lambda^2)$$

In Taylor scheme & Landau gauge

$$\alpha_T(\mu^2) \equiv \frac{g_T^2(\mu^2)}{4\pi} = \lim_{\Lambda \rightarrow \infty} \frac{g_0^2(\Lambda^2)}{4\pi} Z_3(\mu^2, \Lambda^2) \tilde{Z}_3^2(\mu^2, \Lambda^2)$$



# ALPHA<sub>s</sub> from Lattice QCD:

## The ghost-gluon coupling

$$\text{Lattice: } \frac{1}{L^2} \ll p^2 \ll \frac{1}{a^2}$$

Propagators in Landau gauge

$$\langle O \rangle = \frac{1}{Z} \int \mathcal{D}[U] \mathcal{D}[\Phi] O e^{-S[U, \Phi]}$$

$$A_\mu(x + \frac{\hat{\mu}}{2}) = \frac{1}{2iag_0} \left( U_\mu(x) - U_\mu^\dagger(x) \right)$$

- In **Landau** gauge:  $\left[ \text{minimizing } F_U[g] = \text{Re} \sum_x \sum_\mu \left( 1 - \frac{1}{N} g(x) U_\mu(x) g^\dagger(x + \hat{\mu}) \right) \right]$

$$D_{\mu\nu}^{ab}(k) = \int d^4x d^4y e^{ik \cdot (x-y)} \langle A_\mu^a(x) A_\nu^b(y) \rangle_U = \frac{G(k^2)}{k^2} \delta^{ab} \delta_{\mu\nu}^T(k)$$

- In **MOM** renormalization scheme:

# ALPHA<sub>s</sub> from Lattice QCD:

## The ghost-gluon coupling

$$\text{Lattice: } \frac{1}{L^2} \ll p^2 \ll \frac{1}{a^2}$$

Propagators in Landau gauge

$$\langle O \rangle = \frac{1}{Z} \int \mathcal{D}[U] \mathcal{D}[\Phi] O e^{-S[U, \Phi]}$$

$$A_\mu(x + \frac{\hat{\mu}}{2}) = \frac{1}{2iag_0} \left( U_\mu(x) - U_\mu^\dagger(x) \right)$$

- In **Landau** gauge:  $\left[ \text{minimizing } F_U[g] = \text{Re} \sum_x \sum_\mu \left( 1 - \frac{1}{N} g(x) U_\mu(x) g^\dagger(x + \hat{\mu}) \right) \right]$

$$G^{ab}(k) = \int d^4x d^4y e^{ik \cdot (x-y)} (M^{-1})_{xy}^{ab} = -\frac{F(k^2)}{k^2} \delta^{ab}$$

- In **MOM** renormalization scheme:

# ALPHA<sub>s</sub> from Lattice QCD:

## The ghost-gluon coupling

$$\text{Lattice: } \frac{1}{L^2} \ll p^2 \ll \frac{1}{a^2}$$

Propagators in Landau gauge

$$\left(G^{(2)}\right)_{\mu\nu}^{ab}(p^2, \Lambda) = \frac{G(p^2, \Lambda)}{p^2} \delta_{ab} \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right)$$

$$\left(F^{(2)}\right)^{a,b}(p^2, \Lambda) = -\delta_{ab} \frac{F(p^2, \Lambda)}{p^2}$$

Renormalized in non-perturbative MOM-scheme:

$$G_R(p^2, \mu^2) = \lim_{\Lambda \rightarrow \infty} Z_3^{-1}(\mu^2, \Lambda) G(p^2, \Lambda)$$

$$F_R(p^2, \mu^2) = \lim_{\Lambda \rightarrow \infty} \tilde{Z}_3^{-1}(\mu^2, \Lambda) F(p^2, \Lambda)$$

$$G_R(\mu^2, \mu^2) = F_R(\mu^2, \mu^2) = 1$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

4-loops perturbation theory<sup>3</sup>:  $p \gg \Lambda_{QCD}$

$$\alpha_T(\mu^2) = \frac{4\pi}{\beta_0 t} \left( 1 - \frac{\beta_1}{\beta_0^2} \frac{\log(t)}{t} + \frac{\beta_1^2}{\beta_0^4} \frac{1}{t^2} \left( \left( \log(t) - \frac{1}{2} \right)^2 + \frac{\tilde{\beta}_2 \beta_0}{\beta_1^2} - \frac{5}{4} \right) \right) \\ + \frac{1}{(\beta_0 t)^4} \left( \frac{\tilde{\beta}_3}{2\beta_0} + \frac{1}{2} \left( \frac{\beta_1}{\beta_0} \right)^3 \left( -2 \log^3(t) + 5 \log^2(t) + \left( 4 - 6 \frac{\tilde{\beta}_2 \beta_0}{\beta_1^2} \right) \log(t) - 1 \right) \right), \quad t = \ln \frac{\mu^2}{\Lambda_T^2}$$

$$\beta_T(\alpha_T) = \frac{d\alpha_T}{d \ln \mu^2} = -4\pi \sum_{i=0} \tilde{\beta}_i \left( \frac{\alpha_T}{4\pi} \right)^{i+2}, \quad \frac{\Lambda_{\overline{\text{MS}}}}{\Lambda_T} = e^{-\frac{c_1}{2\beta_0}} = e^{-\frac{507 - 40N_f}{792 - 48N_f}} = 0.541449$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

4-loops perturbation theory<sup>3</sup>:  $p \gg \Lambda_{QCD}$

$$\alpha_T(\mu^2) = \frac{4\pi}{\beta_0 t} \left( 1 - \frac{\beta_1 \log(t)}{\beta_0^2 t} + \frac{\beta_1^2}{\beta_0^4} \frac{1}{t^2} \left( \left( \log(t) - \frac{1}{2} \right)^2 + \frac{\tilde{\beta}_2 \beta_0}{\beta_1^2} - \frac{5}{4} \right) \right) \\ + \frac{1}{(\beta_0 t)^4} \left( \frac{\tilde{\beta}_3}{2\beta_0} + \frac{1}{2} \left( \frac{\beta_1}{\beta_0} \right)^3 \left( -2 \log^3(t) + 5 \log^2(t) + \left( 4 - 6 \frac{\tilde{\beta}_2 \beta_0}{\beta_1^2} \right) \log(t) - 1 \right) \right), \quad t = \ln \frac{\mu^2}{\Lambda_T^2}$$

$$\beta_T(\alpha_T) = \frac{d\alpha_T}{d \ln \mu^2} = -4\pi \sum_{i=0}^{\infty} \tilde{\beta}_i \left( \frac{\alpha_T}{4\pi} \right)^{i+2}, \quad \frac{\Lambda_{\overline{\text{MS}}}}{\Lambda_T} = e^{-\frac{c_1}{2\beta_0}} = e^{-\frac{507 - 40N_f}{792 - 48N_f}} = 0.541449$$

K.G. Chetyrkin, Nucl. Phys. B710 (2005) 499

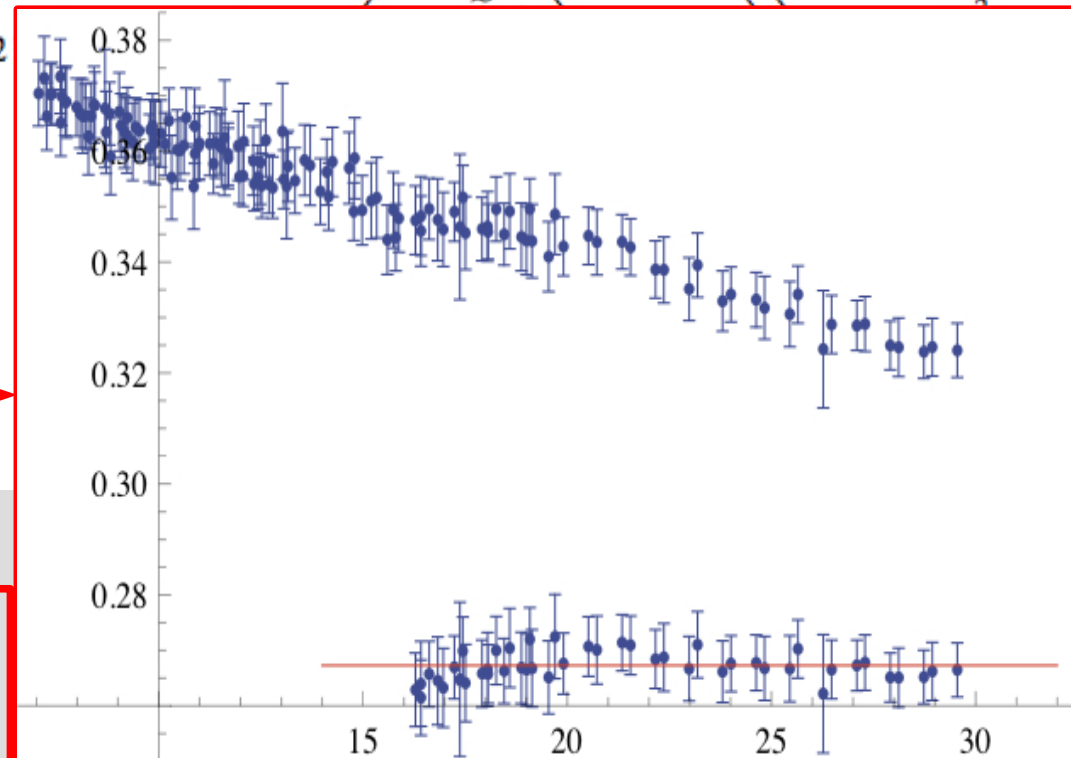
# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

4-loops perturbation theory<sup>3</sup>:  $p \gg \Lambda_{QCD}$

$$\alpha_T(\mu^2) = \frac{4\pi}{\beta_0 t} \left( 1 - \frac{\beta_1 \log(t)}{\beta_0^2 t} + \frac{\beta_1^2}{\beta_0^4} \frac{1}{t^2} \left( \left( \log(t) - \frac{1}{2} \right)^2 + \frac{\tilde{\beta}_2 \beta_0}{\beta_1^2} - \frac{5}{4} \right) \right) + \frac{1}{(\beta_0 t)^4} \left( \frac{\tilde{\beta}_3}{2\beta_0} + \frac{1}{2} \left( \frac{\beta_1}{\beta_0} \right)^3 \right) \left( -2 \right)$$

$$\beta_T(\alpha_T) = \frac{d\alpha_T}{d \ln \mu^2} = -4\pi \sum_{i=0}^{\infty} \tilde{\beta}_i \left( \frac{\alpha_T}{4\pi} \right)^{i+2},$$



Ph. Boucaud et al, PRD79(2009)014508  
Quenched QCD:  $N_f=0$



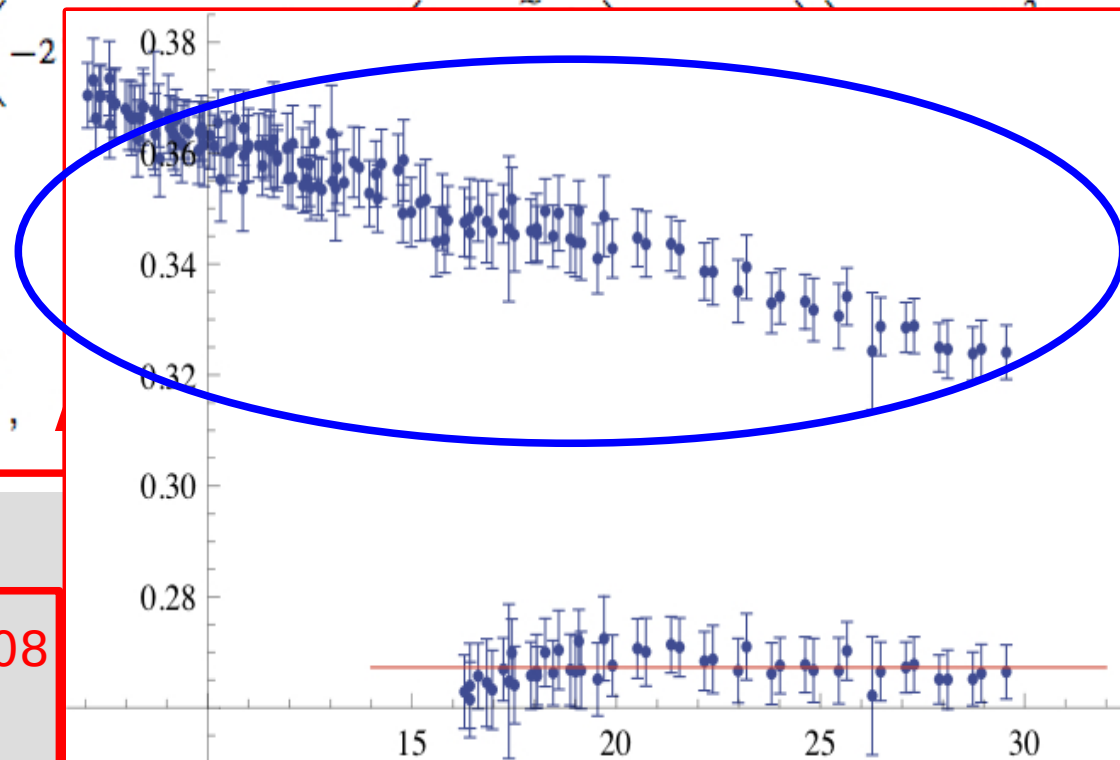
# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

4-loops perturbation theory<sup>3</sup>:  $p \gg \Lambda_{QCD}$

$$\alpha_T(\mu^2) = \frac{4\pi}{\beta_0 t} \left( 1 - \frac{\beta_1 \log(t)}{\beta_0^2 t} + \frac{\beta_1^2}{\beta_0^4} \frac{1}{t^2} \left( \left( \log(t) - \frac{1}{2} \right)^2 + \frac{\tilde{\beta}_2 \beta_0}{\beta_1^2} - \frac{5}{4} \right) \right) + \frac{1}{(\beta_0 t)^4} \left( \frac{\tilde{\beta}_3}{2\beta_0} + \frac{1}{2} \left( \frac{\beta_1}{\beta_0} \right)^3 \right) \left( -2 \right)$$

$$\beta_T(\alpha_T) = \frac{d\alpha_T}{d \ln \mu^2} = -4\pi \sum_{i=0} \tilde{\beta}_i \left( \frac{\alpha_T}{4\pi} \right)^{i+2},$$




Ph. Boucaud et al, PRD79(2009)014508  
Quenched QCD:  $N_f=0$

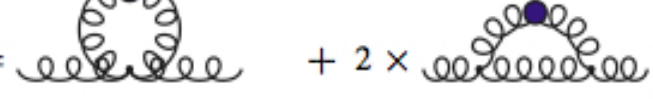


# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

### OPE power corrections

$$(F^{(2)})^{ab}(q^2) = (F_{\text{pert}}^{(2)})^{ab}(q^2) + w^{ab} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab} = 2 \times \text{diagram}$$


$$(G^{(2)})^{ab}_{\mu\nu}(q^2) = (G_{\text{pert}}^{(2)})^{ab}_{\mu\nu}(q^2) + w^{ab}_{\mu\nu} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab}_{\mu\nu} = \text{diagram} + 2 \times \text{diagram}$$


$$F_R(q^2, \mu^2) = F_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R, \mu^2}}{4(N_C^2 - 1)} \right), \quad G_R(q^2, \mu^2) = G_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R, \mu^2}}{4(N_C^2 - 1)} \right)$$


$$\text{Leading logarithm }^4: \alpha_T(\mu^2) = \alpha_T^{\text{pert}}(\mu^2) \left( 1 + \frac{9}{\mu^2} \left( \frac{\alpha_T^{\text{pert}}(\mu^2)}{\alpha_T^{\text{pert}}(q_0^2)} \right)^{1 - \gamma_0^{A^2} / \beta_0} \frac{g_T^2(q_0^2) \langle A^2 \rangle_{R, q_0^2}}{4(N_C^2 - 1)} \right)$$

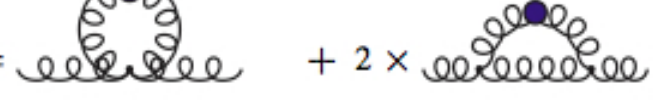
$$1 - \gamma_0^{A^2} / \beta_0 = 1 - \frac{105 - 8N_f}{132 - 8N_f} = \frac{9}{44 - \frac{8}{3}N_f}$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

### OPE power corrections

$$(F^{(2)})^{ab}(q^2) = (F_{\text{pert}}^{(2)})^{ab}(q^2) + w^{ab} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab} = 2 \times \text{diagram}$$


$$(G^{(2)})^{ab}_{\mu\nu}(q^2) = (G_{\text{pert}}^{(2)})^{ab}_{\mu\nu}(q^2) + w^{ab}_{\mu\nu} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab}_{\mu\nu} = \text{diagram} + 2 \times \text{diagram}$$


$$F_R(q^2, \mu^2) = F_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R, \mu^2}}{4(N_C^2 - 1)} \right), \quad G_R(q^2, \mu^2) = G_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R, \mu^2}}{4(N_C^2 - 1)} \right)$$


Leading logarithm <sup>4</sup>:  $\alpha_T(\mu^2) = \alpha_T^{\text{pert}}(\mu^2) \left( 1 + \frac{9}{\mu^2} \left( \frac{\alpha_T^{\text{pert}}(\mu^2)}{\alpha_T^{\text{pert}}(q_0^2)} \right)^{1 - \gamma_0^{A^2} / \beta_0} \frac{g_T^2(q_0^2) \langle A^2 \rangle_{R, q_0^2}}{4(N_C^2 - 1)} \right)$

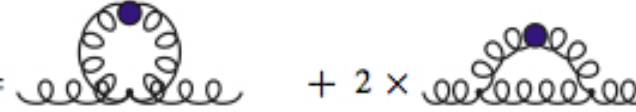
$$1 - \gamma_0^{A^2} / \beta_0 = 1 - \frac{105 - 8N_f}{132 - 8N_f} = \frac{9}{44 - \frac{8}{3}N_f}$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

### OPE power corrections

$$(F^{(2)})^{ab}(q^2) = (F_{\text{pert}}^{(2)})^{ab}(q^2) + w^{ab} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab} = 2 \times \text{diagram}$$


$$(G^{(2)})^{ab}_{\mu\nu}(q^2) = (G_{\text{pert}}^{(2)})^{ab}_{\mu\nu}(q^2) + w^{ab}_{\mu\nu} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab}_{\mu\nu} = \text{diagram} + 2 \times \text{diagram}$$


$$F_R(q^2, \mu^2) = F_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R, \mu^2}}{4(N_C^2 - 1)} \right), \quad G_R(q^2, \mu^2) = G_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R, \mu^2}}{4(N_C^2 - 1)} \right)$$

Leading logarithm <sup>4</sup>:  $\alpha_T(\mu^2) = \alpha_T^{\text{pert}}(\mu^2) \left( 1 + \frac{9}{\mu^2} \left( \frac{\alpha_T^{\text{pert}}(\mu^2)}{\alpha_T^{\text{pert}}(q_0^2)} \right)^{1 - \gamma_0^{A^2} / \beta_0} \frac{g_T^2(q_0^2) \langle A^2 \rangle_{R, q_0^2}}{4(N_C^2 - 1)} \right)$


$$1 - \gamma_0^{A^2} / \beta_0 = 1 - \frac{105 - 8N_f}{132 - 8N_f} = \frac{9}{44 - \frac{8}{3}N_f}$$

Chetyrking & Maier, arXiv:0911.0594  
At the three-loop level !!!

# ALPHA<sub>S</sub> from Lattice QCD:

## Matching Lattice and PTh

### OPE power corrections

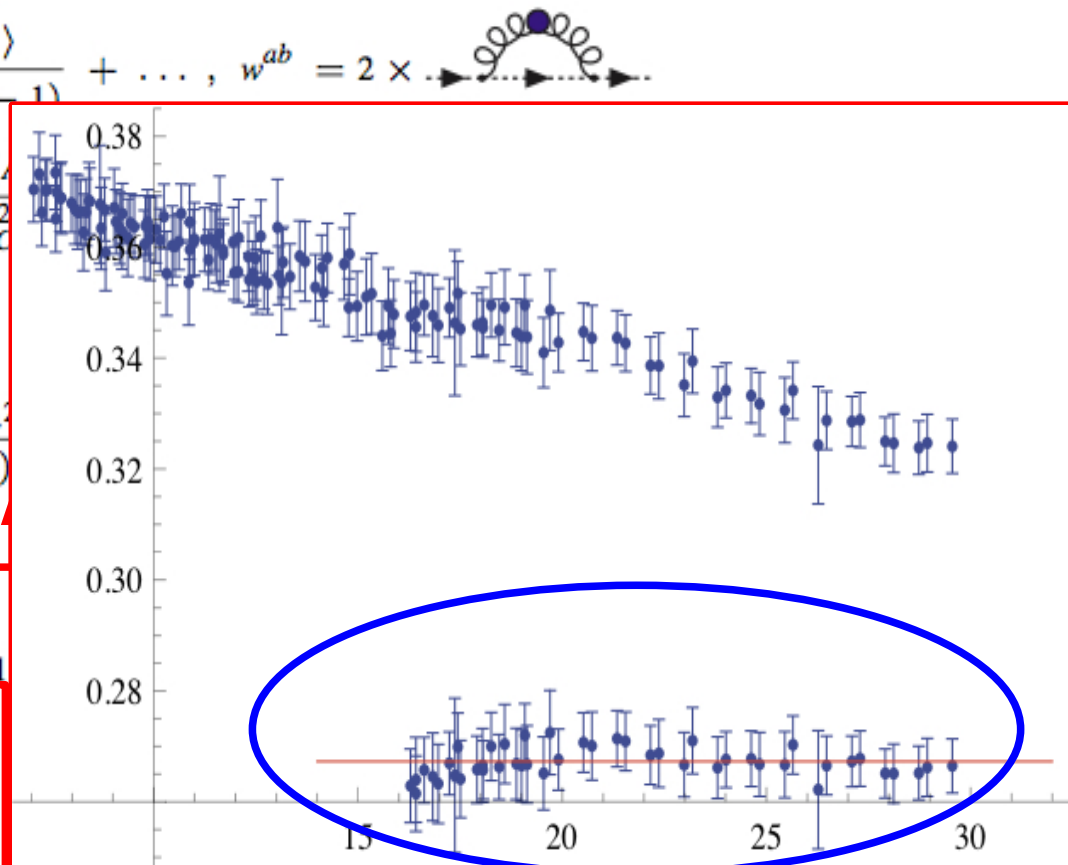
$$(F^{(2)})^{ab}(q^2) = (F_{\text{pert}}^{(2)})^{ab}(q^2) + w^{ab} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)} + \dots, \quad w^{ab} = 2 \times \dots$$


$$(G^{(2)})_{\mu\nu}^{ab}(q^2) = (G_{\text{pert}}^{(2)})_{\mu\nu}^{ab}(q^2) + w_{\mu\nu}^{ab} \frac{\langle A^2 \rangle}{4(N_C^2 - 1)}$$

$$F_R(q^2, \mu^2) = F_{R,\text{pert}}(q^2, \mu^2) \left( 1 + \frac{3}{q^2} \frac{g_R^2 \langle A^2 \rangle_{R,\mu^2}}{4(N_C^2 - 1)} \right)$$

Leading logarithm<sup>4</sup>:  $\alpha_T(\mu^2) = \alpha_T^{\text{pert}}(\mu^2) \left( 1 + \dots \right)$

Ph. Boucaud et al, PRD79(2009)014508  
 Quenched QCD:  $N_f = 0$



# ALPHA<sub>S</sub> from Lattice QCD:

## N<sub>f</sub>=2 twisted-mass QCD

### European Twisted Mass Collaboration

Fermions: twisted-mass action

$$S_{\text{tm}}^{\text{F}} = a^4 \sum_x \left\{ \bar{\chi}_x [D_{\text{W}} + m_0 + i\gamma_5 \tau_3 \mu_q] \chi_x \right\}$$

Gauge fields: tISym action

$$S_g = \frac{\beta}{3} \sum_x \left( b_0 \sum_{\substack{\mu, \nu=1 \\ 1 \leq \mu < \nu}}^4 \{1 - \text{ReTr}(U_{x, \mu, \nu}^{1 \times 1})\} + b_1 \sum_{\substack{\mu, \nu=1 \\ \mu \neq \nu}}^4 \{1 - \text{ReTr}(U_{x, \mu, \nu}^{1 \times 2})\} \right), \quad \beta \equiv 6/g_0^2$$



$(b_0 = 1 - 8b_1, b_1 = -1/12)$  + Maximal twist :  $\mathcal{O}(a^2)$

$$V = 24^3 \times 48 \quad \beta = 3.9 \quad \mu = 0.004, 0.0064, 0.010$$

$$V = 32^3 \times 64 \quad \beta = 4.05 \quad \mu = 0.003, 0.006, 0.008, 0.012$$

$$\beta = 4.2 \quad \mu = 0.0065$$

Artefacts:  $\mathcal{O}(a^2 \Lambda_{\text{QCD}}^2)$ ,  $\mathcal{O}(a^2 p^2)$ ,  $\mathcal{O}(a^2 \mu^2)$

# ALPHA<sub>S</sub> from Lattice QCD:

## N<sub>f</sub>=2 twisted-mass QCD

### Ghost and gluon on the lattice

Landau gauge

$$F_U[g] = \text{Re} \left[ \sum_x \sum_\mu \text{Tr} \left( 1 - \frac{1}{N} g(x) U_\mu(x) g^\dagger(x + \mu) \right) \right]$$

Gluon:

$$A_\mu(x + \hat{\mu}/2) = \frac{U_\mu(x) - U_\mu^\dagger(x)}{2iag_0} - \frac{1}{3} \text{Tr} \left( \frac{U_\mu(x) - U_\mu^\dagger(x)}{2iag_0} \right)$$

$$\text{oooooo} \quad \left( G^{(2)} \right)_{\mu_1 \mu_2}^{a_1 a_2}(p) = \langle A_{\mu_1}^{a_1}(p) A_{\mu_2}^{a_2}(-p) \rangle$$



# ALPHA<sub>S</sub> from Lattice QCD:

## N<sub>f</sub>=2 twisted-mass QCD

### Ghost and gluon on the lattice

Landau gauge

$$F_U[g] = \text{Re} \left[ \sum_x \sum_\mu \text{Tr} \left( 1 - \frac{1}{N} g(x) U_\mu(x) g^\dagger(x + \mu) \right) \right]$$

Ghost:

.....→.....  $(F^{(2)})^{ab}(x-y) \equiv \langle (M^{-1})_{xy}^{ab} \rangle, M(U) = -\frac{1}{N} \nabla \cdot \tilde{D}(U)$

$$\tilde{D}(U)\eta(x) = \frac{1}{2} \left( U_\mu(x)\eta(x+\mu) - \eta(x)U_\mu(x) + \eta(x+\mu)U_\mu^\dagger - U_\mu^\dagger(x)\eta(x) \right)$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

$O(4)$  breaking:  $H(4)$  discretization artefacts<sup>8</sup>

Orbit labeled by  $H(4)$ -invariants:  $p^{[2n]} = \sum_{\mu=1}^4 p_{\mu}^{2n}, n = 1, 2, 3$

Momentum on the lattice:  $\tilde{p}_{\mu} = \frac{1}{a} \sin ap_{\mu}, p_{\mu} = \frac{2\pi n}{Na} \quad n = 0, 1, \dots, N$

$$a^2 \tilde{p}^2 \equiv \sum_{\mu=1}^4 a^2 \tilde{p}_{\mu}^2 = a^2 p^2 + c_1 a^4 p^{[4]} + \dots = a^2 p^2 \left( 1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots \right)$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

$O(4)$  breaking:  $H(4)$  discretization artefacts<sup>8</sup>

Orbit labeled by  $H(4)$ -invariants:  $p^{[2n]} = \sum_{\mu=1}^4 p_{\mu}^{2n}, n = 1, 2, 3$

Momentum on the lattice:  $\tilde{p}_{\mu} = \frac{1}{a} \sin ap_{\mu}, p_{\mu} = \frac{2\pi n}{Na} \quad n = 0, 1, \dots, N$

$$a^2 \tilde{p}^2 \equiv \sum_{\mu=1}^4 a^2 \tilde{p}_{\mu}^2 = a^2 p^2 + c_1 a^4 p^{[4]} + \dots = a^2 p^2 \left( 1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots \right)$$

If  $\epsilon = a^2 p^{[4]} / p^2 \ll 1 \dots$

$$\begin{aligned} Q(a^2 \tilde{p}_{\mu}^2, a^2 \Lambda^2) &\equiv Q \left( a^2 p^2 \left( 1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots \right), a^2 \Lambda^2 \right) \\ &= Q(a^2 p^2, a^2 \Lambda^2) + \left. \frac{dQ}{d\epsilon} \right|_{\epsilon=0} a^2 \frac{p^{[4]}}{p^2} + \dots \end{aligned}$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts


$O(4)$  breaking:  $H(4)$  discretization artefacts<sup>8</sup>

Orbit labeled by  $H(4)$ -invariants:  $p^{[2n]} = \sum_{\mu=1}^4 p_{\mu}^{2n}, n = 1, 2, 3$

Momentum on the lattice:  $\tilde{p}_{\mu} = \frac{1}{a} \sin ap_{\mu}, p_{\mu} = \frac{2\pi n}{Na} \quad n = 0, 1, \dots, N$

$$a^2 \tilde{p}^2 \equiv \sum_{\mu=1}^4 a^2 \tilde{p}_{\mu}^2 = a^2 p^2 + c_1 a^4 p^{[4]} + \dots = a^2 p^2 \left( 1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots \right)$$

If  $\epsilon = a^2 p^{[4]} / p^2 \ll 1 \dots$


$$R = R_0 + R_1 a^2 p^2$$

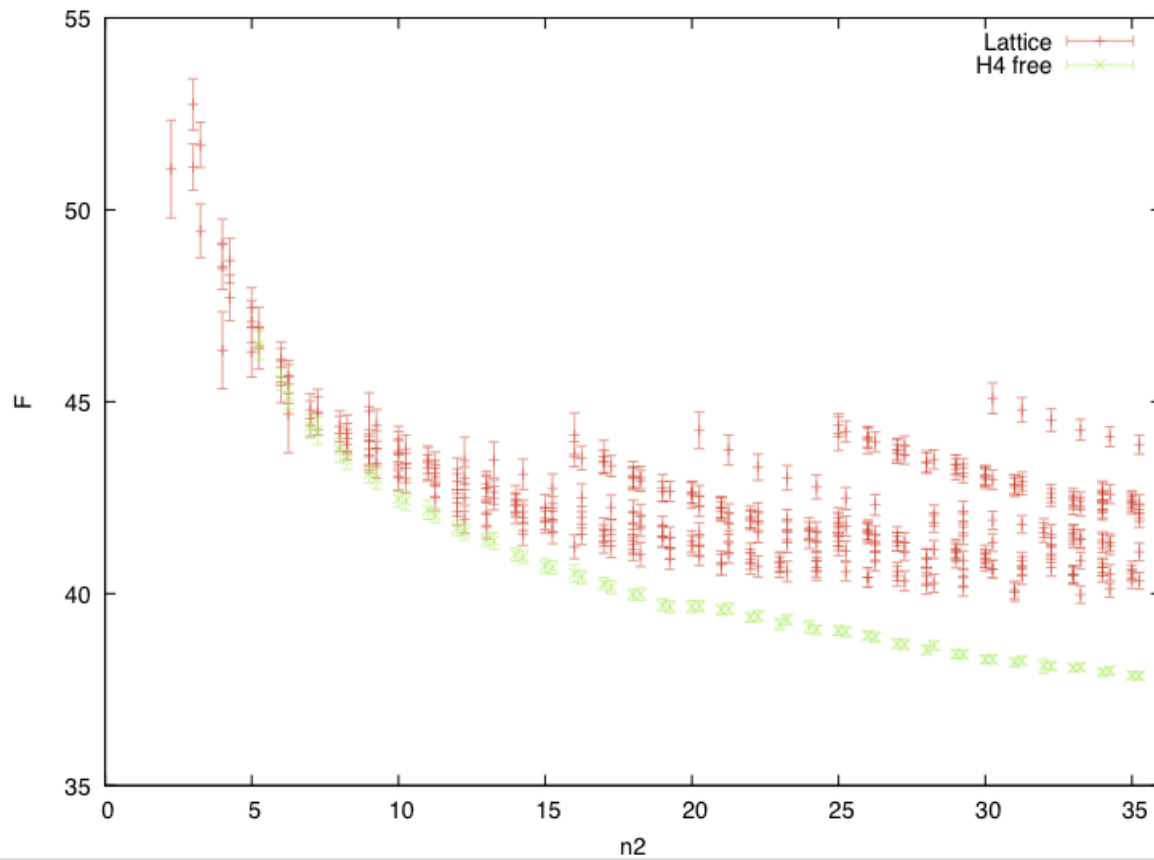
$$Q(a^2 \tilde{p}_{\mu}^2, a^2 \Lambda^2) \equiv Q \left( a^2 p^2 \left( 1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots \right), a^2 \Lambda^2 \right) \\ = Q(a^2 p^2, a^2 \Lambda^2) + \left. \frac{dQ}{d\epsilon} \right|_{\epsilon=0} a^2 \frac{p^{[4]}}{p^2} + \dots$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

$O(4)$  breaking:  $H(4)$  discretization artefacts<sup>8</sup>

Orbit labeled by  $H(4)$ -invariants:  $p^{[2n]} = \sum_{\mu=1}^4 p_{\mu}^{2n}, n = 1, 2, 3$



$n = 0, 1, \dots, N$

$$1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots$$

$p^2 \ll 1 \dots$

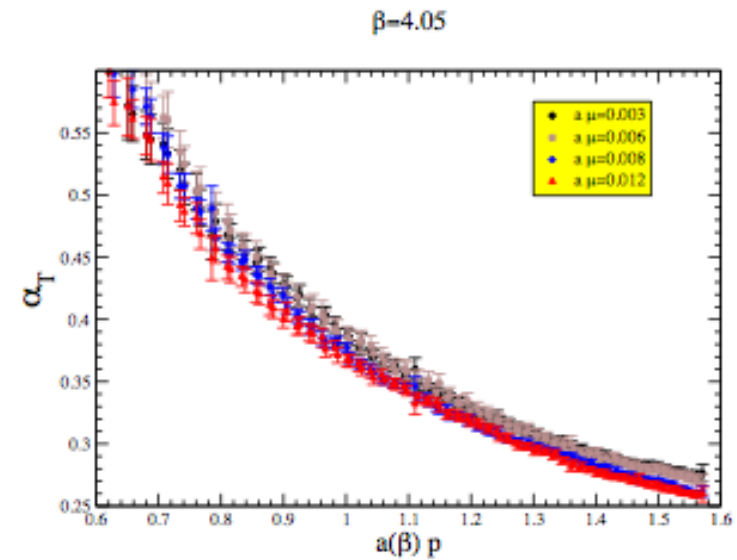
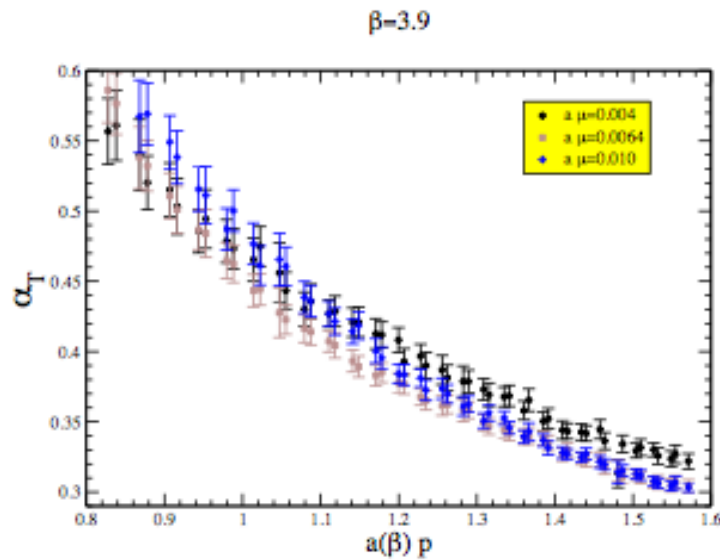
$$\equiv Q \left( a^2 p^2 \left( 1 + c_1 a^2 \frac{p^{[4]}}{p^2} + \dots \right), a^2 \Lambda^2 \right)$$

$$, a^2 \Lambda^2) + \left. \frac{dQ}{d\epsilon} \right|_{\epsilon=0} a^2 \frac{p^{[4]}}{p^2} + \dots$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

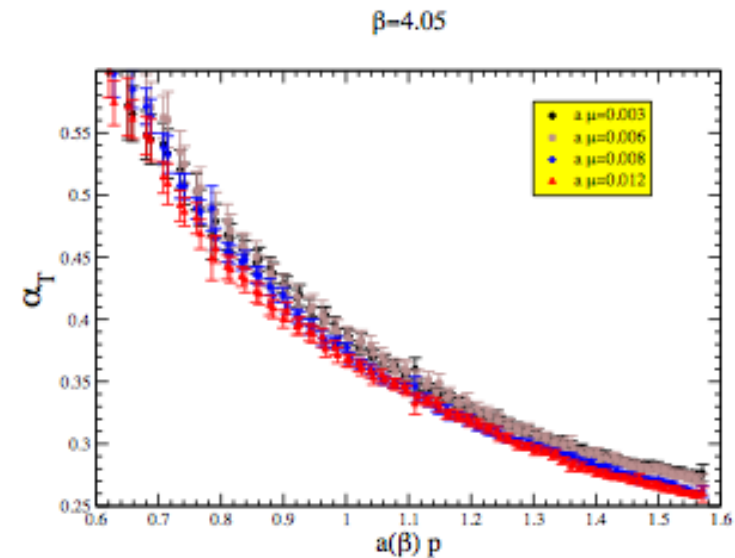
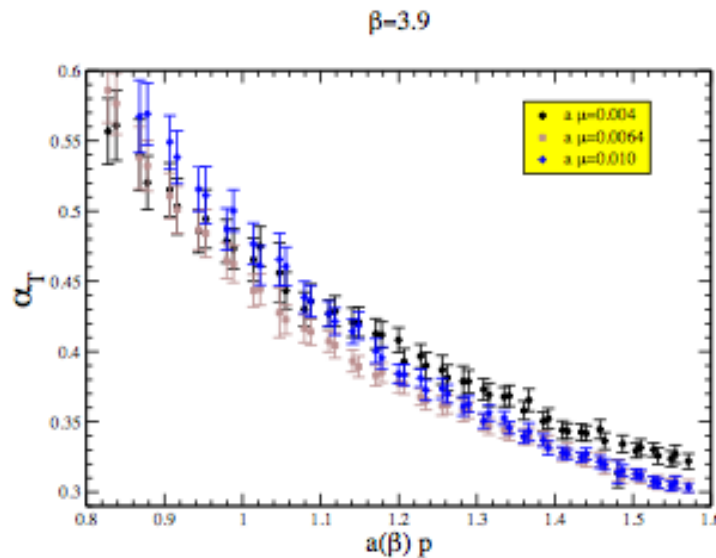
### Quark mass artefacts



# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

### Quark mass artefacts



$\mathcal{O}(a^2 \mu_q^2)$  dependence

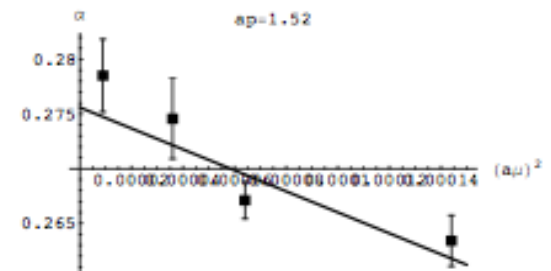
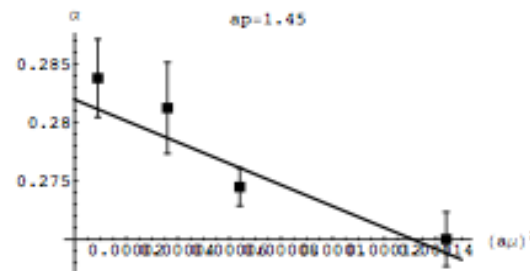
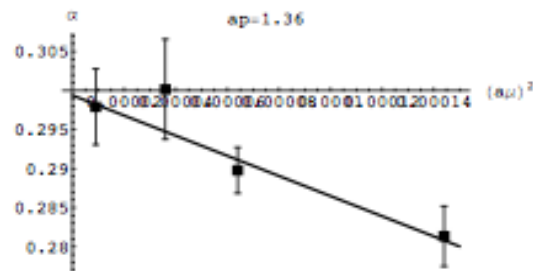
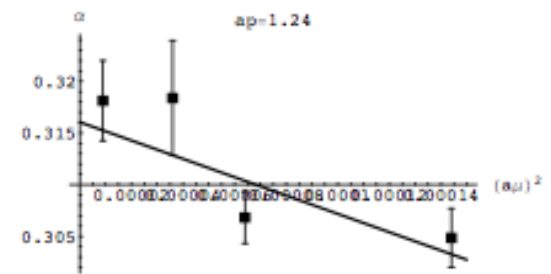
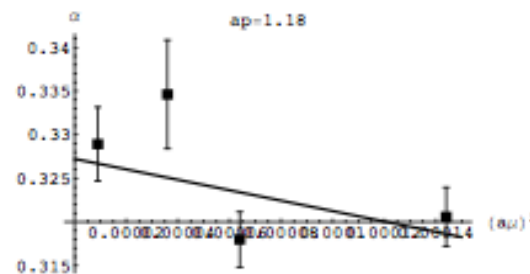
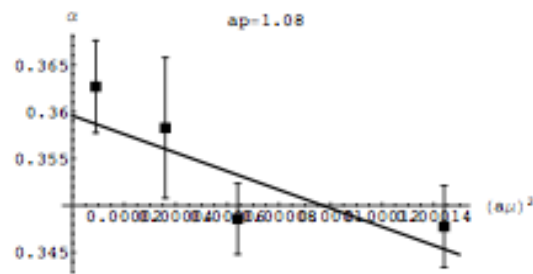
$$\begin{aligned}\widehat{\alpha}_T(a^2 p^2, a^2 \mu_q^2) &= \frac{g_0^2(a^2)}{4\pi} \widehat{G}(a^2 p^2, a^2 \mu_q^2) \widehat{F}^2(a^2 p^2, a^2 \mu_q^2) \\ &= \widehat{\alpha}_T(a^2 p^2, 0) + \frac{\partial \widehat{\alpha}_T}{\partial (a^2 \mu_q^2)}(a^2 p^2) a^2 \mu_q^2 + \dots\end{aligned}$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

### Quark mass artefacts

Example:  $\beta = 4.05$



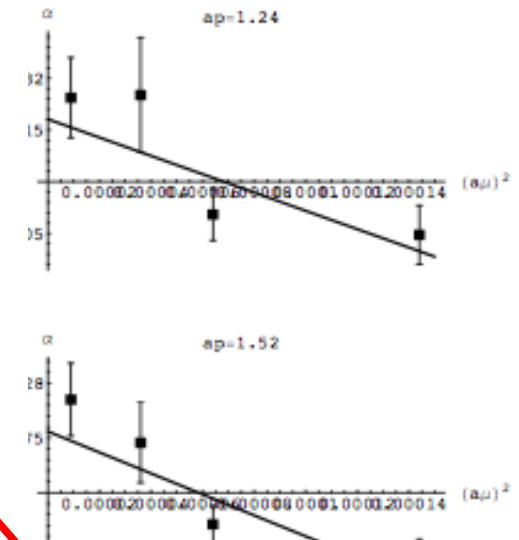
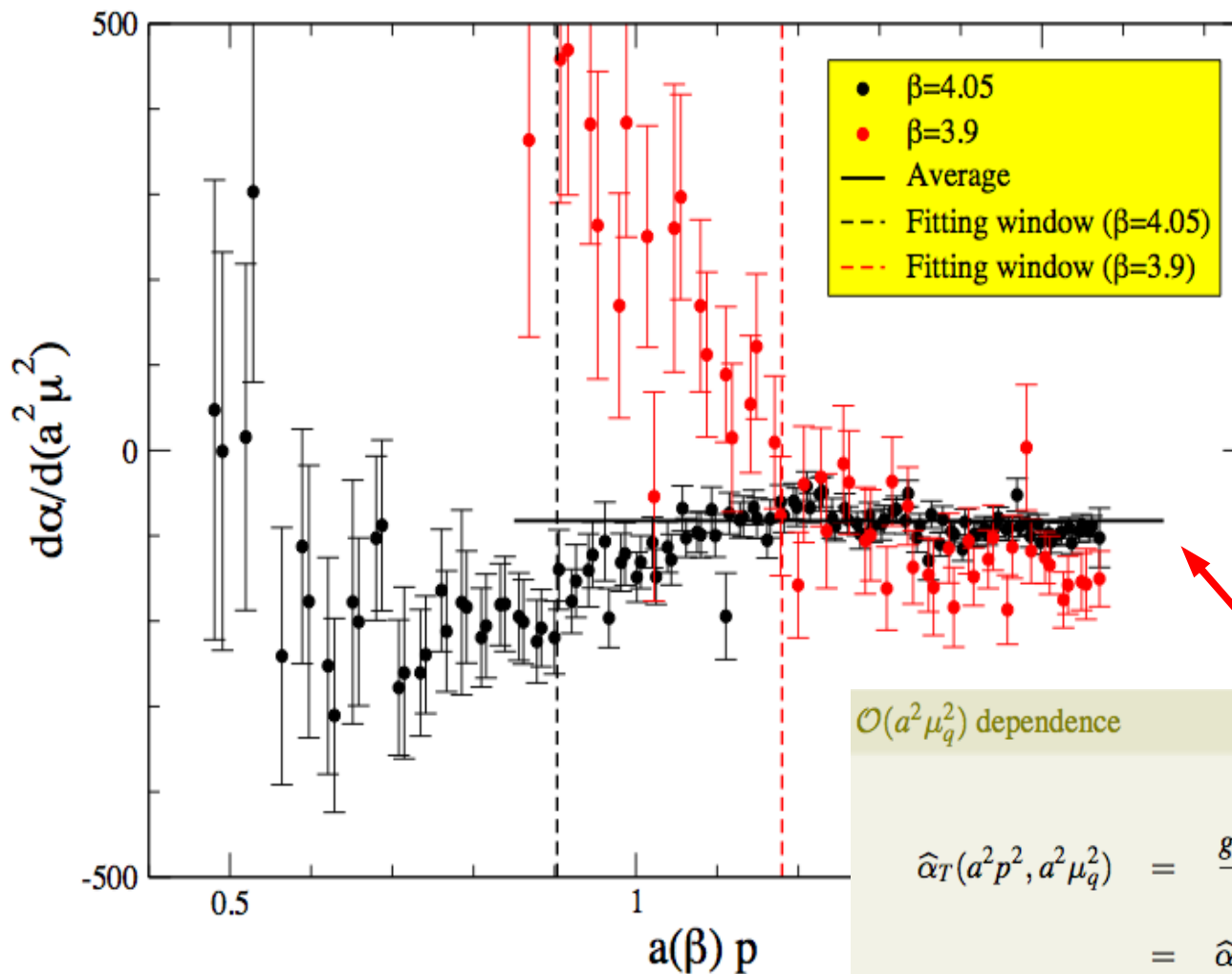
$$\begin{aligned}\widehat{\alpha}_T(a^2 p^2, a^2 \mu_q^2) &= \frac{g_0^2(a^2)}{4\pi} \widehat{G}(a^2 p^2, a^2 \mu_q^2) \widehat{F}^2(a^2 p^2, a^2 \mu_q^2) \\ &= \widehat{\alpha}_T(a^2 p^2, 0) + \frac{\partial \widehat{\alpha}_T}{\partial (a^2 \mu_q^2)} (a^2 p^2) a^2 \mu_q^2 + \dots\end{aligned}$$



# ALPHA<sub>S</sub> from Lattice QCD:

## Lattice artefacts

### Quark mass artefacts

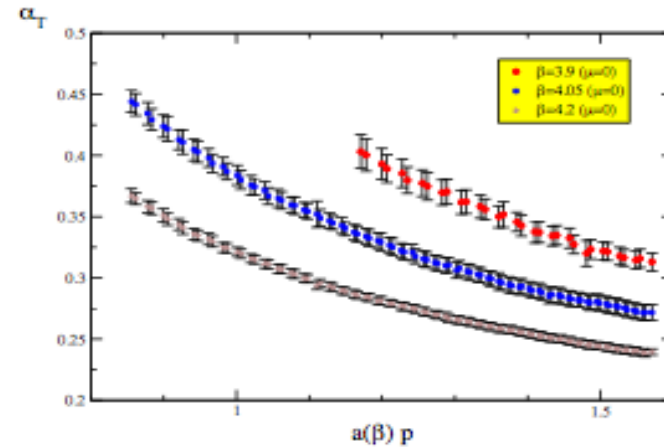
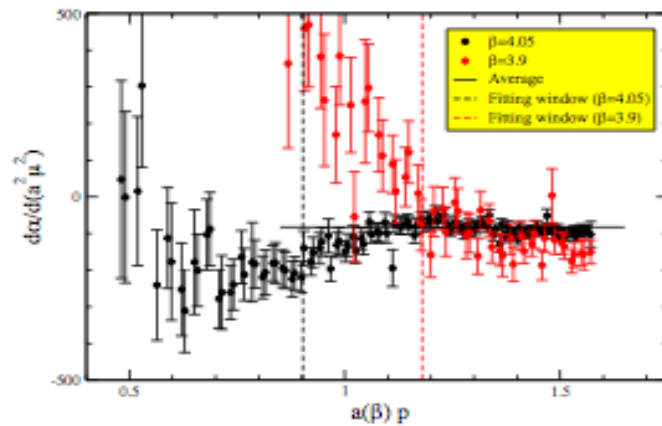


$\mathcal{O}(a^2\mu_q^2)$  dependence

$$\begin{aligned}
 \hat{\alpha}_T(a^2p^2, a^2\mu_q^2) &= \frac{g_0^2(a^2)}{4\pi} \hat{G}(a^2p^2, a^2\mu_q^2) \hat{F}^2(a^2p^2, a^2\mu_q^2) \\
 &= \hat{\alpha}_T(a^2p^2, 0) + \frac{\partial \hat{\alpha}_T}{\partial (a^2\mu_q^2)} (a^2p^2) a^2\mu_q^2 + \dots
 \end{aligned}$$

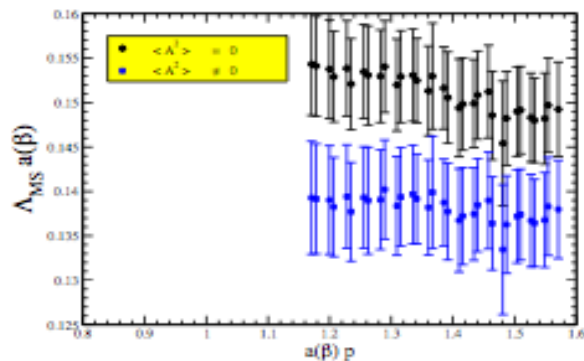
# ALPHA<sub>S</sub> from Lattice QCD:

## After curing lattice artefacts

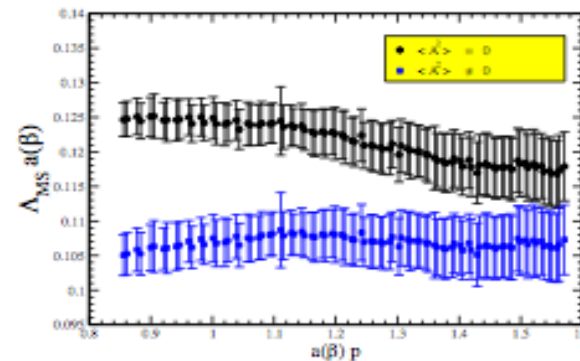


$$R_0 = -92(11), \quad p \geq p_{\min} \simeq 2.8 \text{ GeV} \quad (a(3.9)=0.0801(14) \text{ fm}^9)$$

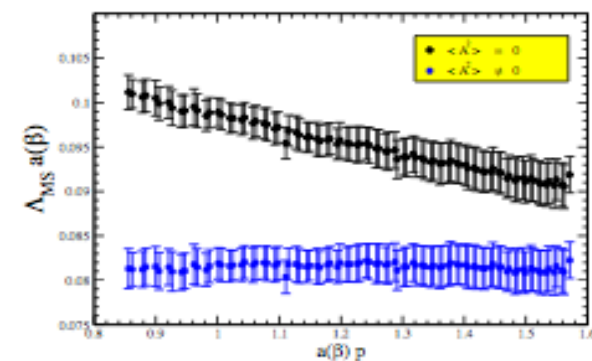
Requiring a plateau for  $\Lambda_{\overline{MS}}$  in the window  $ap \geq p_{\min}$



$$\beta = 3.9$$



$$\beta = 4.05$$



$$\beta = 4.2$$

# ALPHA<sub>s</sub> from Lattice QCD:

**After curing lattice artefacts**

Global fit and calibration of lattice spacing

$$\chi^2 \left( a(\beta_0) \Lambda_{\overline{\text{MS}}}, c, \frac{a(\beta_1)}{a(\beta_0)}, \frac{a(\beta_2)}{a(\beta_0)} \right) = \sum_{j=0}^2 \sum_i \frac{\left( \Lambda_i(\beta_j) - \frac{a(\beta_j)}{a(\beta_0)} a(\beta_0) \Lambda_{\overline{\text{MS}}} \right)^2}{\delta^2(\Lambda_i)}$$

Variables:  $a(\beta_0) \Lambda_{\overline{\text{MS}}}, c, \frac{a(\beta_1)}{a(\beta_0)}, \frac{a(\beta_2)}{a(\beta_0)}$

# ALPHA<sub>S</sub> from Lattice QCD:

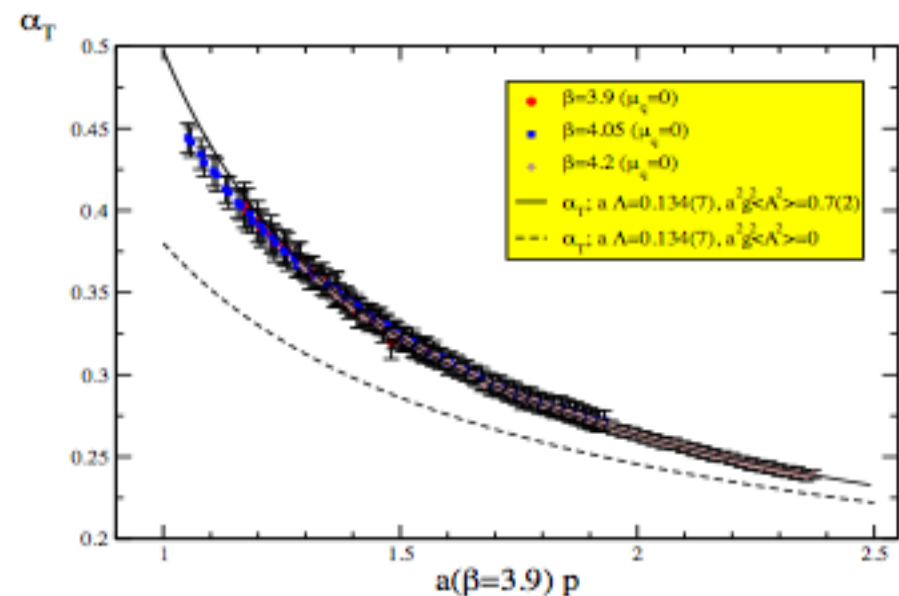
After curing lattice artefacts

Global fit and calibration of lattice spacing

$$\chi^2 \left( a(\beta_0) \Lambda_{\overline{\text{MS}}}, c, \frac{a(\beta_1)}{a(\beta_0)}, \frac{a(\beta_2)}{a(\beta_0)} \right) = \sum_{j=0}^2 \sum_i \frac{\left( \Lambda_i(\beta_j) - \frac{a(\beta_j)}{a(\beta_0)} a(\beta_0) \Lambda_{\overline{\text{MS}}} \right)^2}{\delta^2(\Lambda_i)}$$

Variables:  $a(\beta_0) \Lambda_{\overline{\text{MS}}}, c, \frac{a(\beta_1)}{a(\beta_0)}, \frac{a(\beta_2)}{a(\beta_0)}$

	This paper	String tension <sup>10</sup>
$a(3.9)/a(4.05)$	1.224(23)	1.255(42)
$a(3.9)/a(4.2)$	1.510(32)	1.558(52)
$a(4.05)/a(4.2)$	1.233(25)	1.241(39)
$\Lambda_{\overline{\text{MS}}} a(3.9)$	0.134(7)	
$g^2 \langle A^2 \rangle a^2(3.9)$	0.70(23)	



# ALPHA<sub>S</sub> from Lattice QCD:

After curing lattice artefacts

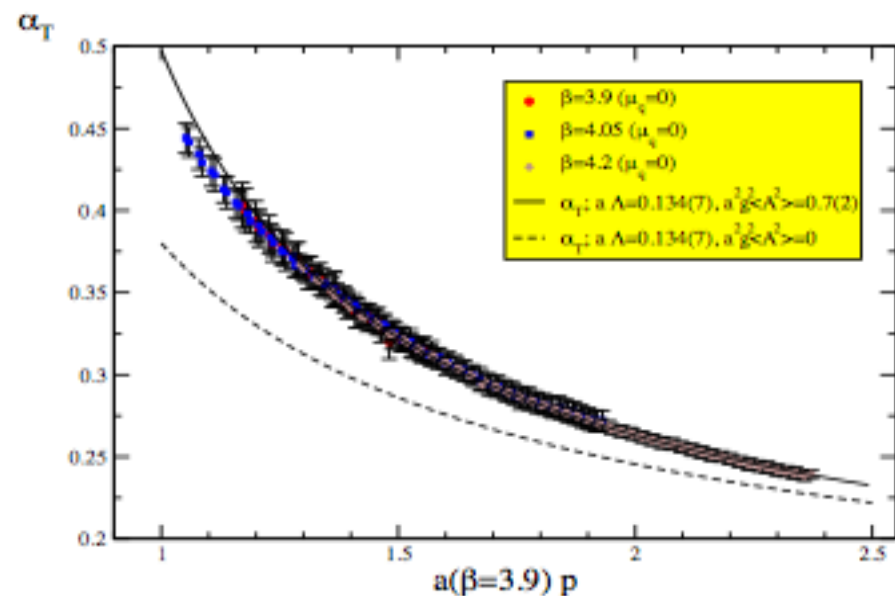
Global fit and calibration of lattice spacing

$$\chi^2 \left( a(\beta_0) \Lambda_{\overline{\text{MS}}}, c, \frac{a(\beta_1)}{a(\beta_0)}, \frac{a(\beta_2)}{a(\beta_0)} \right) = \sum_{j=0}^2 \sum_i \frac{\left( \Lambda_i(\beta_j) - \frac{a(\beta_j)}{a(\beta_0)} a(\beta_0) \Lambda_{\overline{\text{MS}}} \right)^2}{\delta^2(\Lambda_i)}$$

Variables:  $a(\beta_0) \Lambda_{\overline{\text{MS}}}, c, \frac{a(\beta_1)}{a(\beta_0)}, \frac{a(\beta_2)}{a(\beta_0)}$



	This paper	String tension <sup>10</sup>
$a(3.9)/a(4.05)$	1.224(23)	1.255(42)
$a(3.9)/a(4.2)$	1.510(32)	1.558(52)
$a(4.05)/a(4.2)$	1.233(25)	1.241(39)
$\Lambda_{\overline{\text{MS}}} a(3.9)$	0.134(7)	
$g^2 \langle A^2 \rangle a^2(3.9)$	0.70(23)	



# ALPHA<sub>S</sub> from Lattice QCD:

## Systematic deviations

Higher orders for Wilson coefficient...

	One loop	Two loops	Three loops	Four loops
$\Lambda_{\overline{\text{MS}}}a(3.9)$	0.134(7)	0.136(7)	0.137(7)	0.138(7)
$g^2 \langle A^2 \rangle a^2(3.9)$	0.70(23)	0.52(18)	0.44(14)	0.39(14)

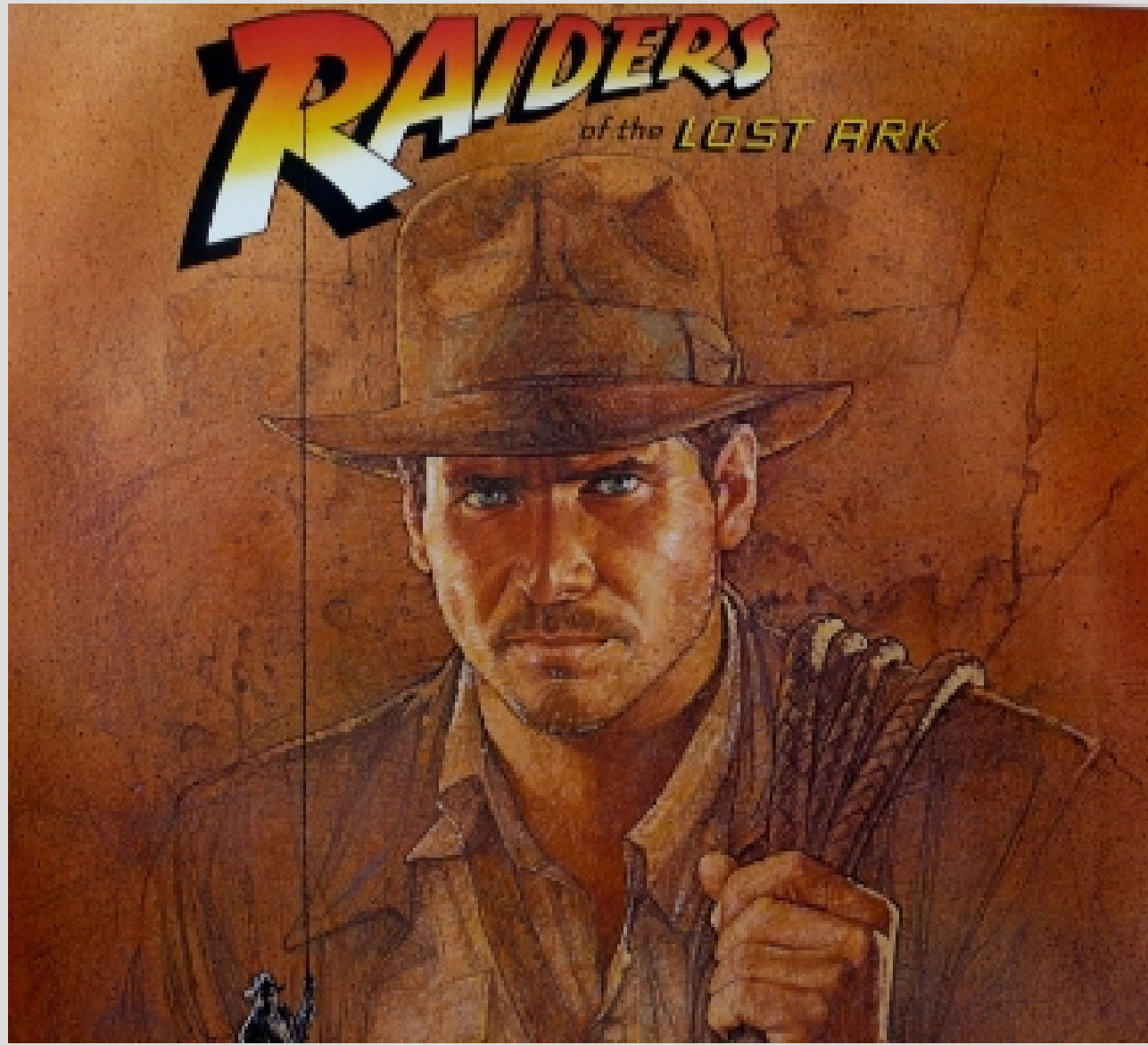
Three-loop versus four-loop perturbative coupling constant...

	Four loops	Three loops
$a(3.9)/a(4.05)$	1.224(23)	1.229(23)
$a(3.9)/a(4.2)$	1.510(32)	1.510(29)
$a(4.05)/a(4.2)$	1.233(26)	1.234(25)
$\Lambda_{\overline{\text{MS}}}a(3.9)$	0.134(7)	0.125(6)
$g^2 \langle A^2 \rangle a^2(3.9)$	0.70(23)	0.80(20)

Higher orders in OPE...unstable!

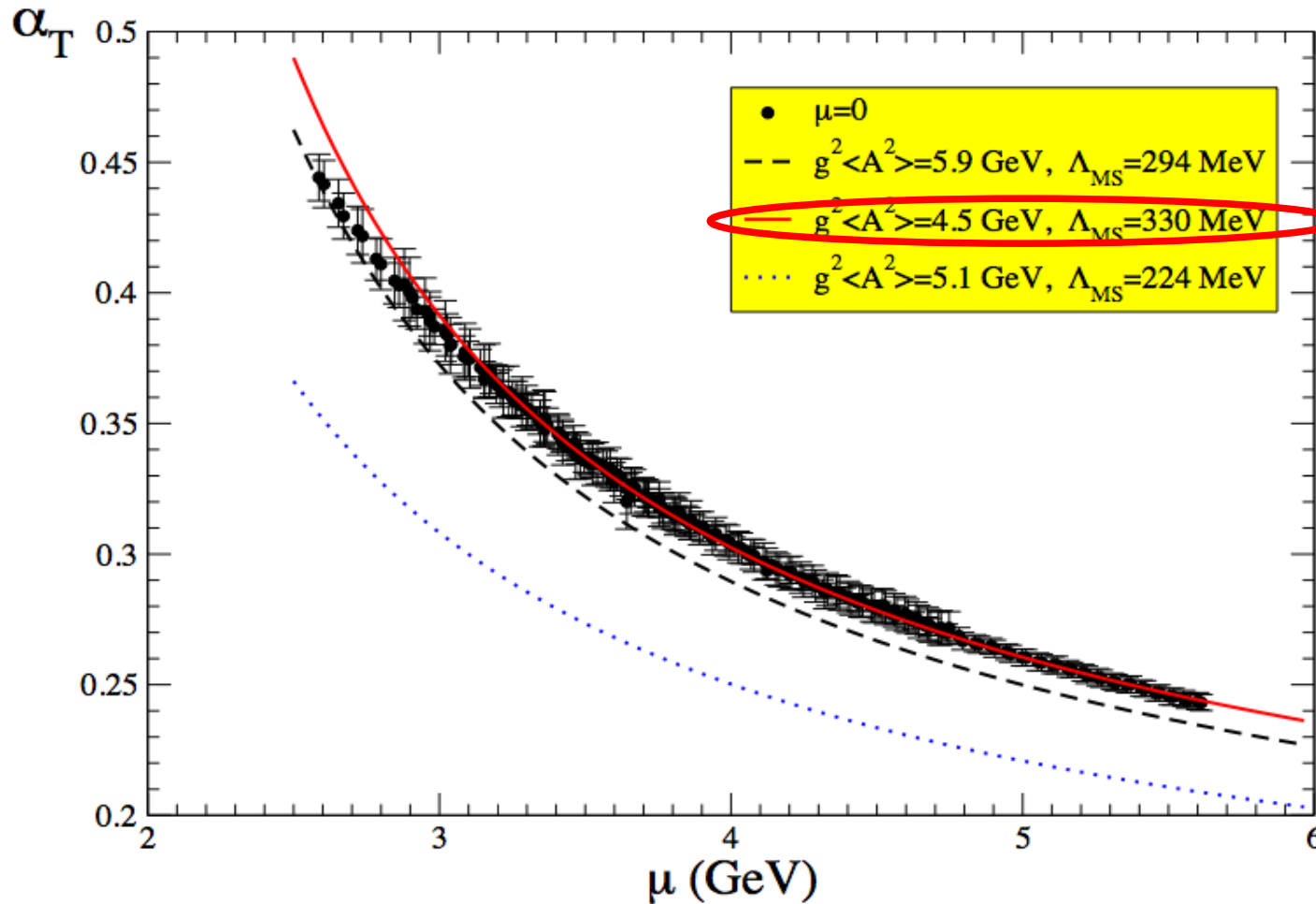
$$\alpha_{T,P4}(\mu^2) = \alpha_T(\mu^2) + \frac{c_4}{\mu^4}$$

# ALPHA<sub>s</sub> from Lattice QCD: Conclusions



# ALPHA<sub>S</sub> from Lattice QCD: Conclusions

RAIDERS





# ALPHA<sub>S</sub> from Lattice QCD: Conclusions

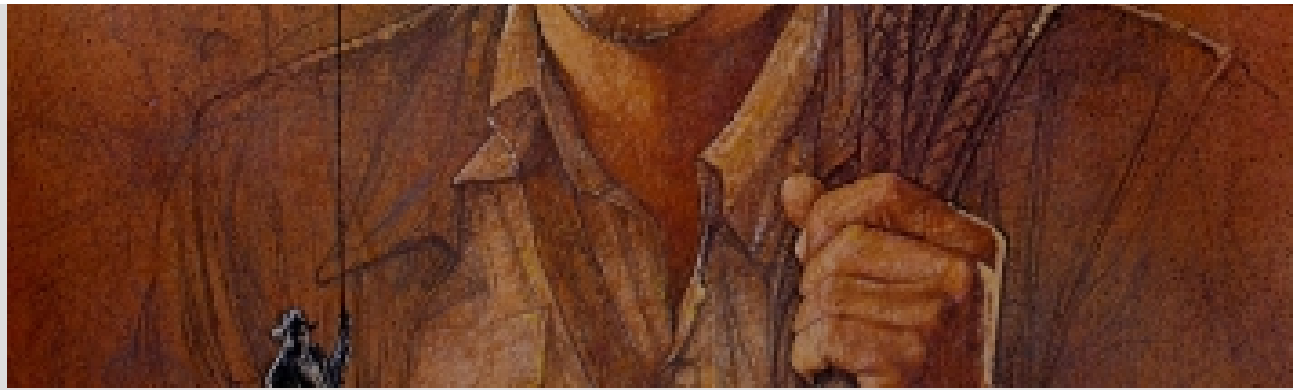
## Conclusions

$N_f = 2$ :

$$\begin{aligned}\Lambda_{\overline{\text{MS}}} &= (330 \pm 23 \pm 22_{-33}) \text{ MeV} \\ g^2(q_0^2) \langle A^2 \rangle_{q_0} &= (4.2 \pm 1.5 \pm 0.7^{+?}) \text{ GeV}^2, \quad q_0 \sim 10 \text{ GeV}\end{aligned}$$

$N_f = 0$ :

$$\begin{aligned}\Lambda_{\overline{\text{MS}}} &= 224_{-5}^{+8} \text{ MeV} \\ g_T^2 \langle A^2 \rangle &= 5.1_{-1.1}^{+0.7} \text{ GeV}^2\end{aligned}$$



# ALPHA<sub>S</sub> from Lattice QCD:

## Conclusions

### Conclusions

$N_f = 2$ :

$$\Lambda_{\overline{\text{MS}}} = (330 \pm 23 \pm 22_{-33}) \text{ MeV}$$
$$g^2(q_0^2) \langle A^2 \rangle_{q_0} = (4.2 \pm 1.5 \pm 0.7^{+?}) \text{ GeV}^2, \quad q_0 \sim 10 \text{ GeV}$$

$N_f = 0$ :

$$\Lambda_{\overline{\text{MS}}} = 224_{-5}^{+8} \text{ MeV}$$
$$g_T^2 \langle A^2 \rangle = 5.1_{-1.1}^{+0.7} \text{ GeV}^2$$

### Outlooks

ETMC new configurations:

$$N_f = 4$$
$$N_f = 2 + 1 + 1$$

# ALPHA<sub>S</sub> from Lattice QCD:

## Conclusions

### Conclusions

$N_f = 2$ :

$$\Lambda_{\overline{\text{MS}}} = (330 \pm 23 \pm 22_{-33}) \text{ MeV}$$

$$g^2(q_0^2) \langle A^2 \rangle_{q_0} = (4.2 \pm 1.5 \pm 0.7^{+?}) \text{ GeV}^2, \quad q_0 \sim 10 \text{ GeV}$$

$N_f = 0$ :

Thank you!!!

### Outlooks

ETMC new configurations:

$$N_f = 4$$

$$N_f = 2 + 1 + 1$$