

# SUSY Breaking in the Klebanov-Strassler Background by Anti-D3 Branes

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based on arxiv:0912.3519

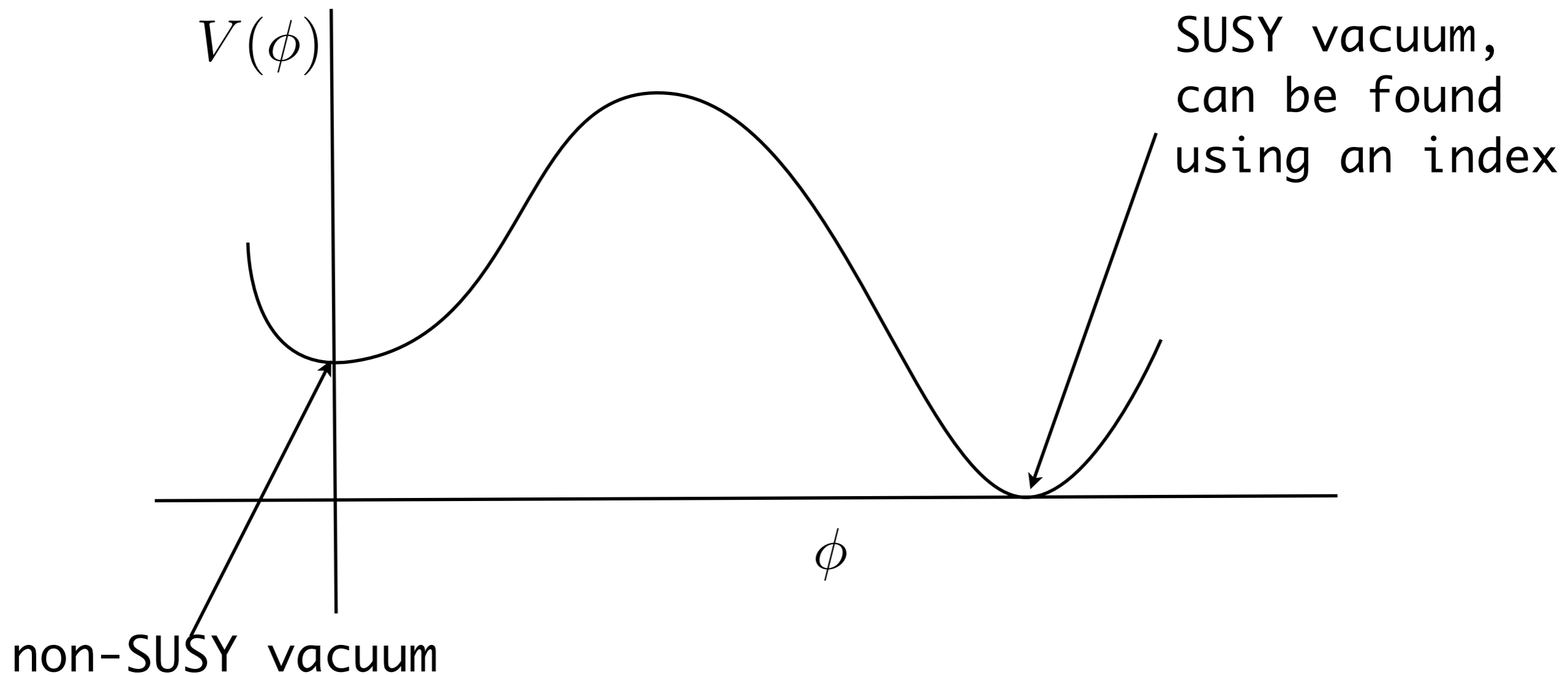
Iosif Bena, Mariana Graña, N.H.

and work in progress...

# Outline

- Review of meta-stable vacua, flux compactifications and the Klebanov-Strassler (KS) Background
- KPV proposal for spontaneous susy breaking in the KS background
- Our computation of the spectrum around KS
- The anti-D3 brane in the Klebanov-Strassler background

# Meta-Stable Vacua in Field Theory



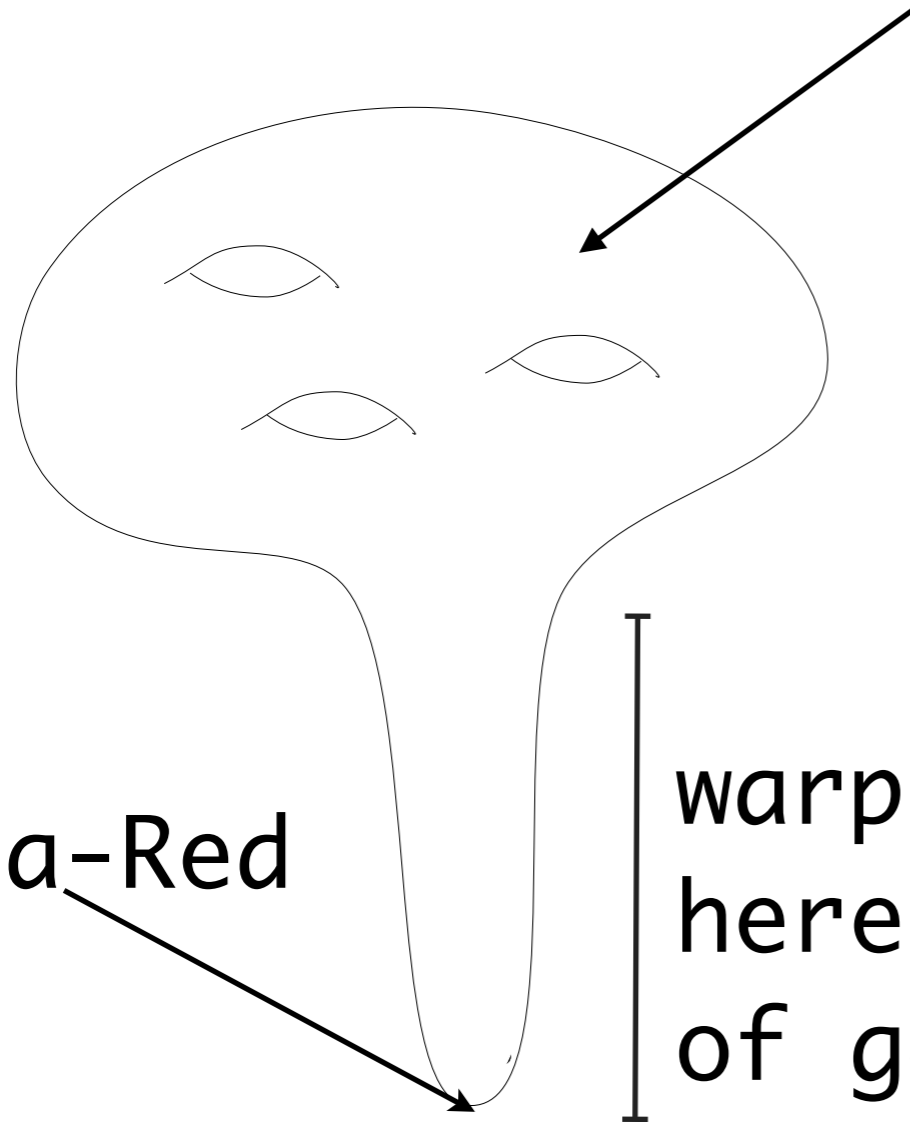
scalar potential  $\xrightarrow{?}$  spontaneous SUSY breaking

- Non-SUSY vacua are difficult to
- identify. In general, strongly coupled QFT  $\rightarrow$  hard to calculate  $V(\phi)$
  - Strategies to study the scalar potential
    - Limit oneself to the holomorphic sector
    - Use a field theory duality
    - Use gauge/gravity duality
    - perhaps perturbation theory in simple examples

# Flux Compactifications in String Theory

six dimensional internal geometry gives the 4d UV completion by gravity

+ moduli stabilization (hard problem)



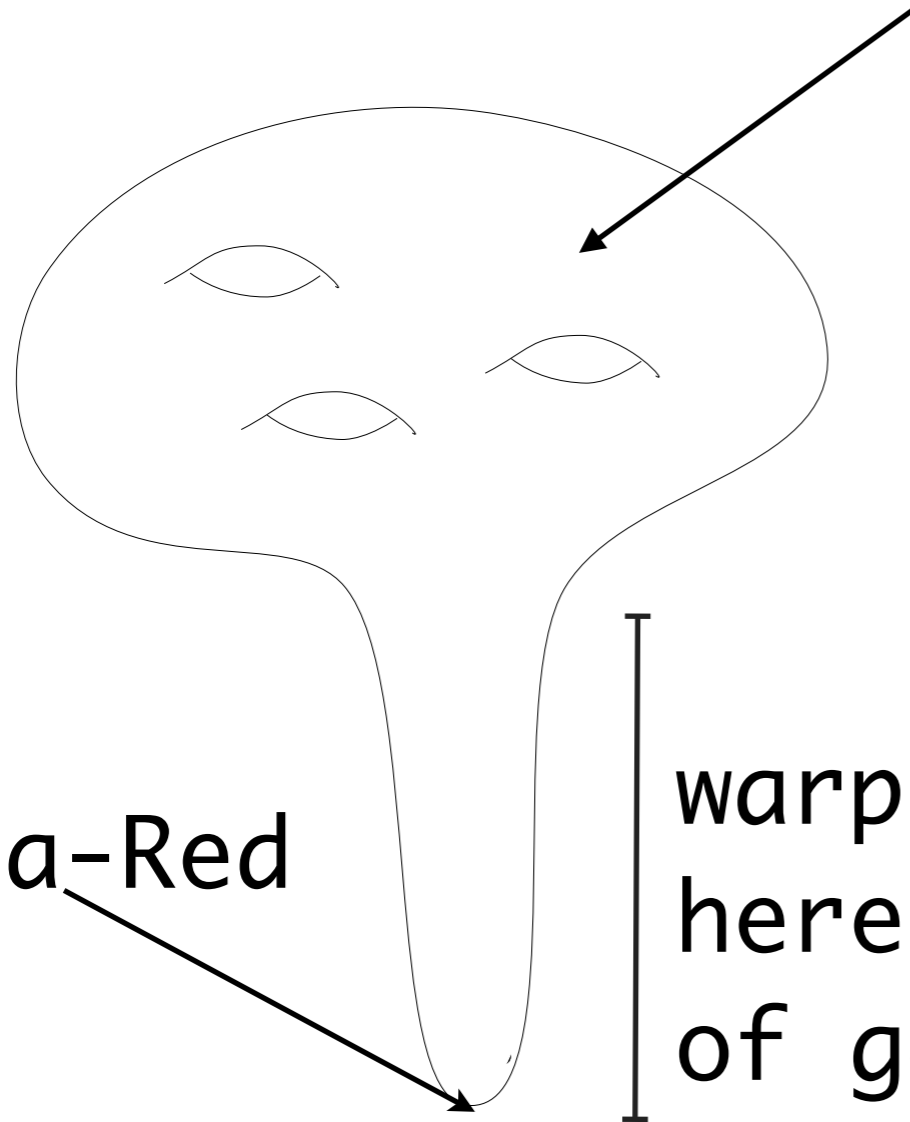
warped throat region,  
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Infra-Red

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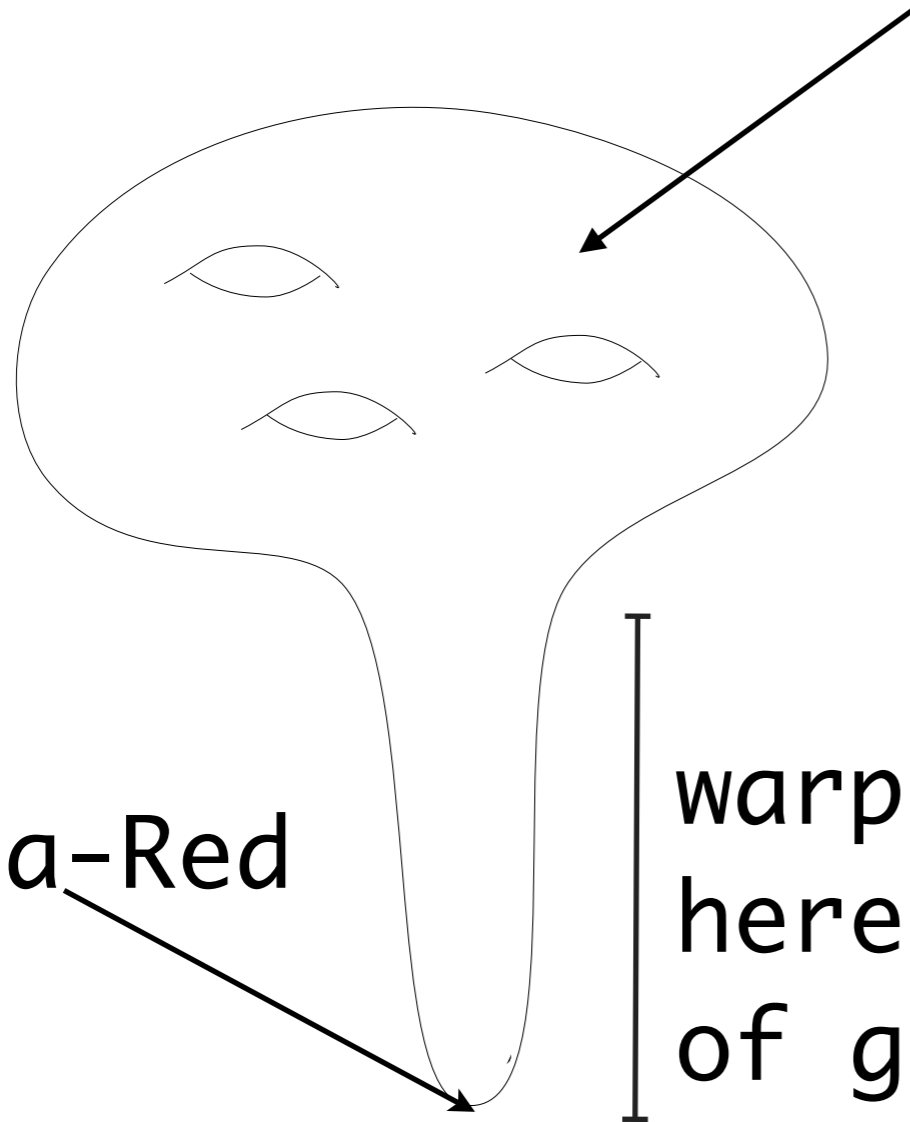
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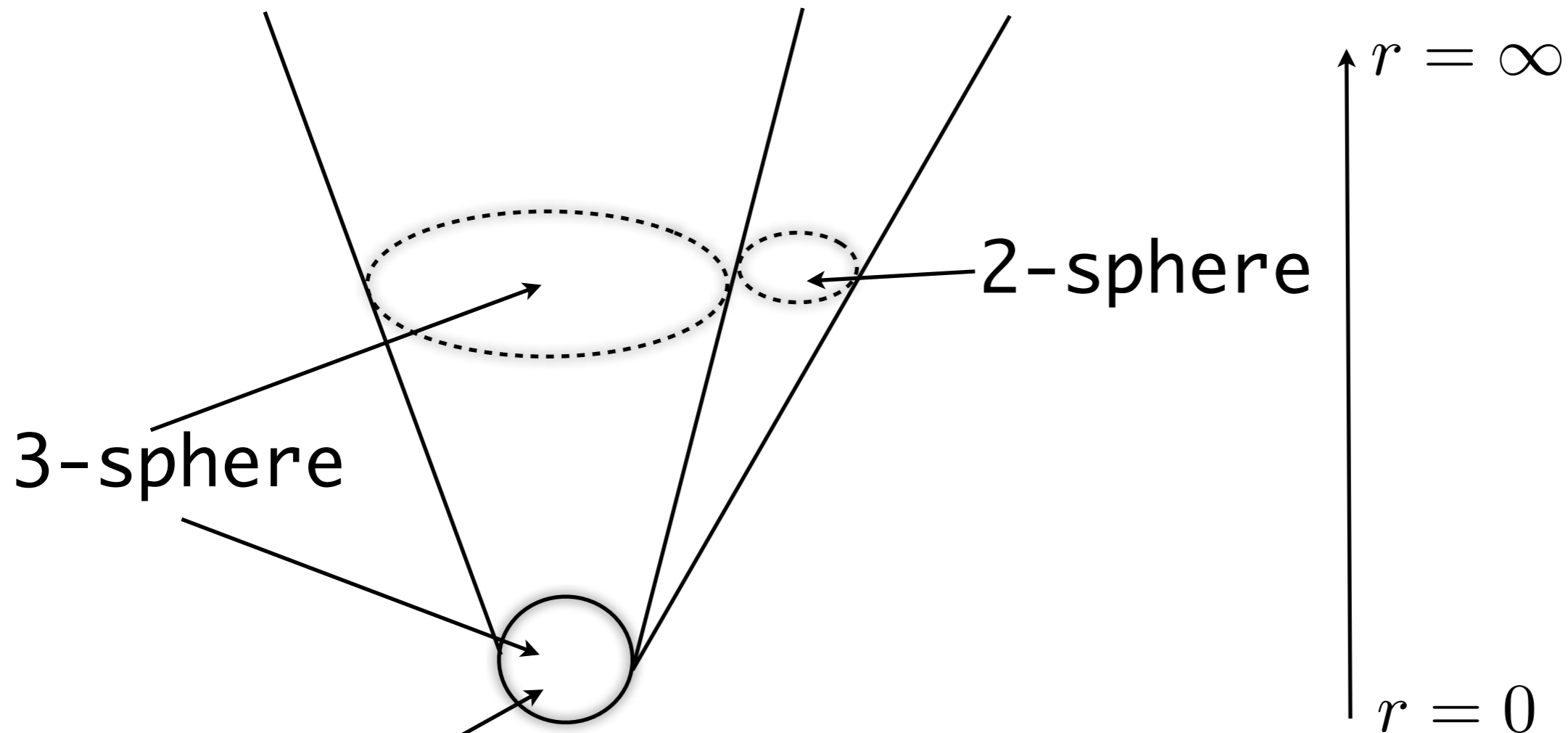
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how is SUSY broken?

at what scale is SUSY broken?

# Warped Throat Region: The Klebanov-Strassler Background

The Deformed Conifold  $R_{ij}^{(6)} = 0$  Klebanov-Strassler '99



$$\frac{1}{4\pi^2 \alpha'} \int_{S^3} F^{(3)} = M$$

despite the branes dissolved in flux, this is a **smooth** sugra solution

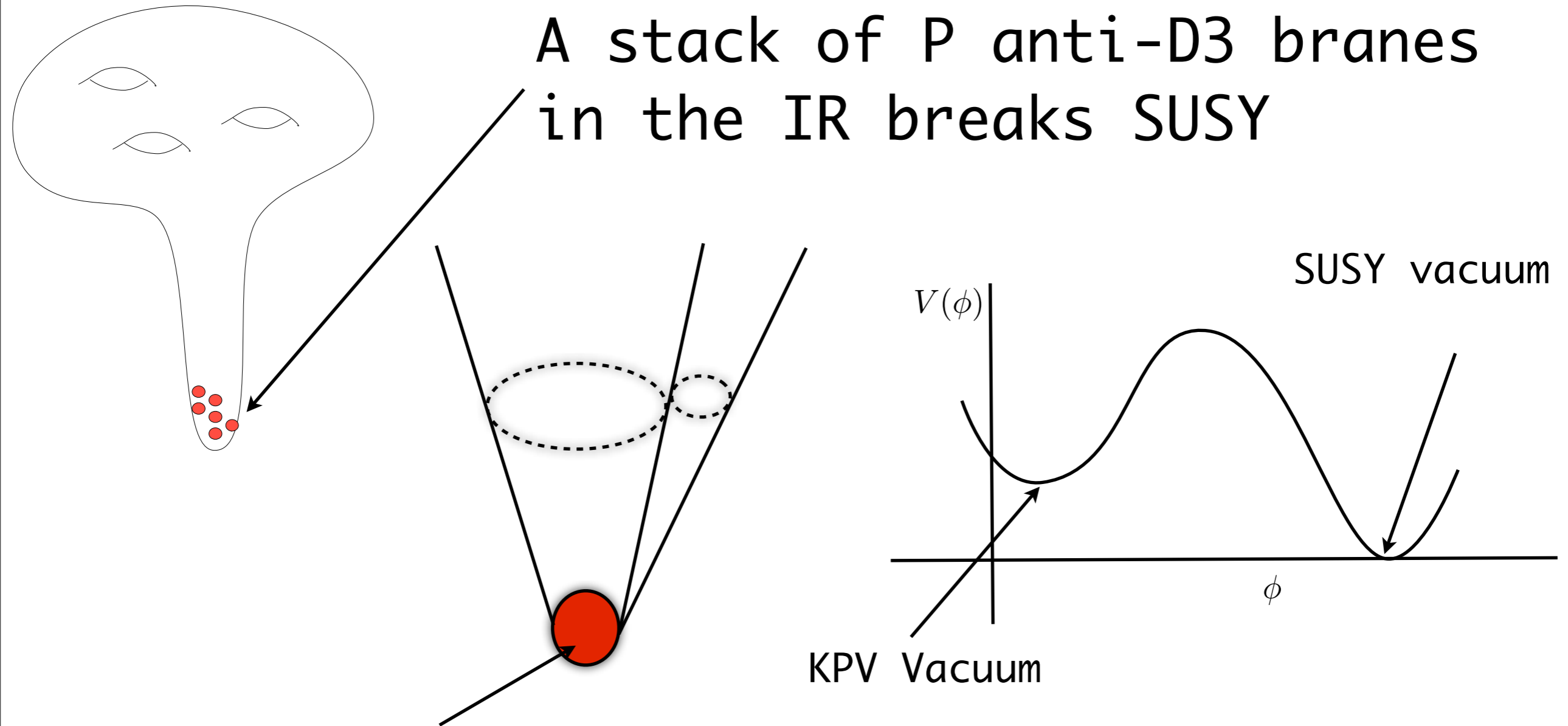


# KPV proposal for spontaneous SUSY breaking

Kachru/Pearson/Verlinde '01

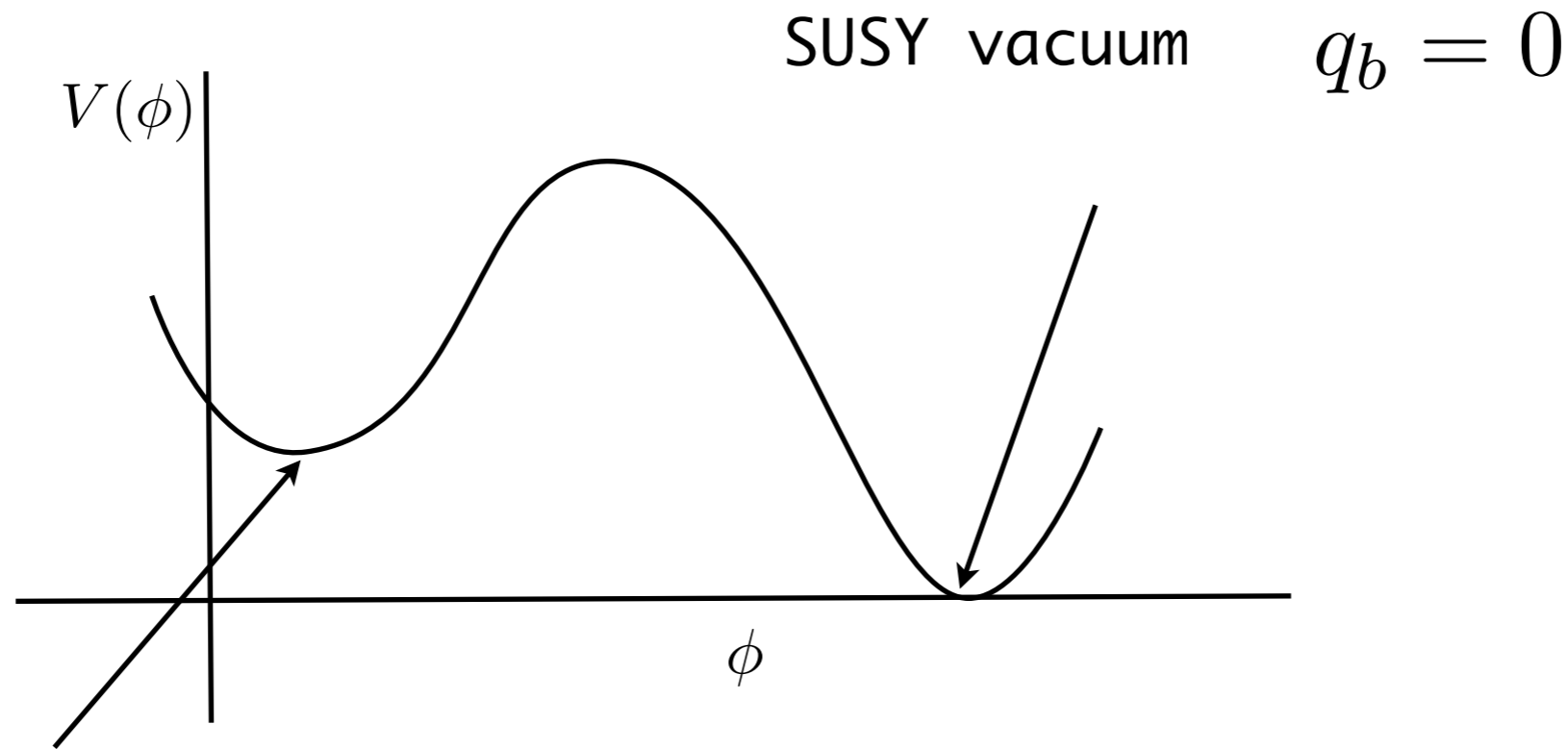
## Probe Approx'

A stack of  $P$  anti-D3 branes  
in the IR breaks SUSY



We will preserve the  
symmetries of the sphere  
by smearing the anti-D3 branes

# KPV proposal for spontaneous SUSY breaking



KPV Vacuum

$$-\bar{N} < q_b < 0$$

$$q_b = \int_{T_0^{11}} F^{(5)}$$

$$\begin{aligned} Q_{D3}^{Max} &= q_b + q_f \\ &= -\bar{N} \end{aligned}$$

$$q_f = \int_{X_6} H^{(3)} \wedge F^{(3)}$$

$$= \left( \int_{T_\infty^{11}} - \int_{T_0^{11}} \right) F^{(5)}$$

Our Goal: find the supergravity solution for the backreacted anti-D3 branes

- This exact supergravity solution would determine unambiguously whether the SUSY breaking by anti-D3 branes is explicit or spontaneous.

- If it is spontaneous, this solution would be very useful for model building, explicit SUSY breaking is less useful

- We work in perturbation theory around the KS background in  $\overline{N}/M$

# AdS/CFT Review: VEV's/Deformations

In AdS space, the wave equation for a scalar field has two possible behaviours near the boundary.

$$\dim(\mathcal{O}) = \Delta$$

$$\phi \sim r^{d-\Delta} \longleftrightarrow$$

Deformation by

$$\delta S = \int d^d x \mathcal{O}(x)$$

$$\phi \sim r^{-\Delta} \longleftrightarrow$$

$$\langle \mathcal{O}(x) \rangle$$

In the UV the KS solution is almost AdS, up to logarithm terms

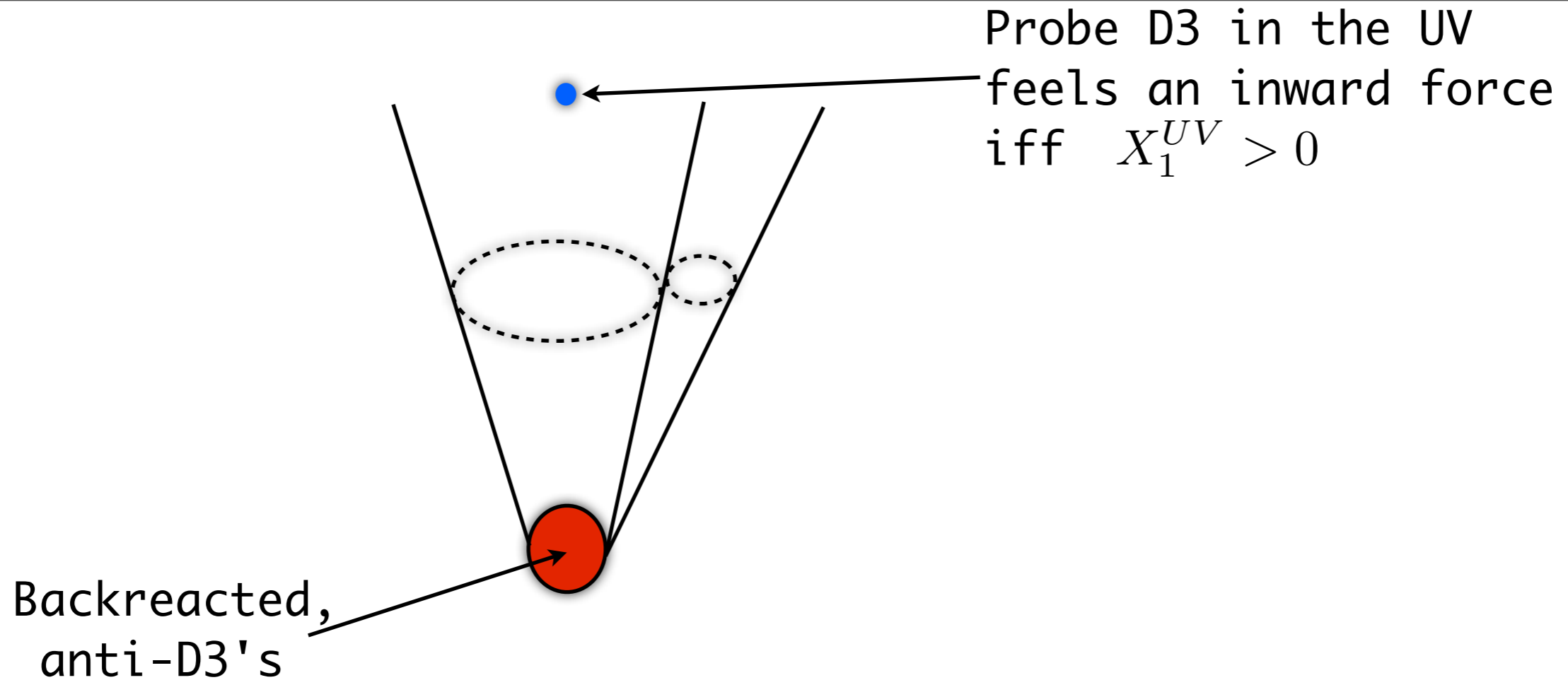
The physical modes in the UV can only be determined by fixing the IR b.c.'s.

# $SU(2) \times SU(2) \times \mathbb{Z}_2$ invariant KS scalar spectrum

## UV Behavior of the spectrum

dim $\Delta$	non-norm/norm	int. constant
8	$r^4 / r^{-8}$	$Y_4 / X_1$
7	$r^3 / r^{-7}$	$Y_5 / X_6$
6	$r^2 / r^{-6}$	$X_3 / Y_3$
5	$r / r^{-5}$	— — —
4	$r^0 / r^{-4}$	$Y_7, Y_8, Y_1 / X_5, X_4, X_8$
3	$r^{-1} / r^{-3}$	$X_2, X_7 / Y_6, Y_2$
2	$r^{-2} / r^{-2}$	— — —

- The  $X_i$  break SUSY
- The  $Y_j$  preserve SUSY



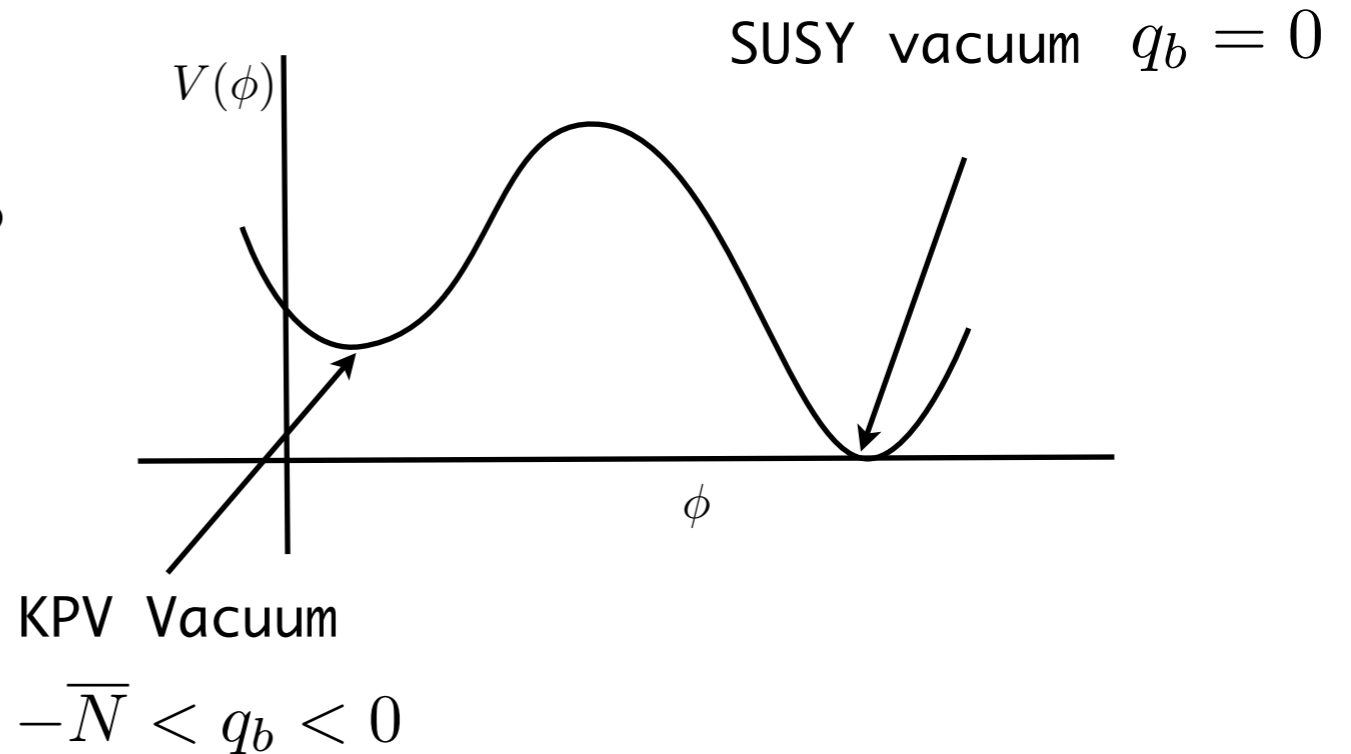
- Universal asymptotic UV falloff for the force on a probe D3 brane in perturbed KS background

$$F \sim \frac{X_1}{r^5} + \mathcal{O}(r^{-11})$$

- $X_1$  is solely responsible for a force on a probe D3 brane!

# Our Strategy

- First enforce UV boundary conditions of the SUSY vacuum



- Examine IR boundary conditions and allow for **physical** singularities
- If no such IR boundary conditions can be found, we would then conclude that the UV boundary conditions must be changed

# IR Singularities

We find various singularities  
in the IR

- warp factor  $\sim \tau^{-1}$   
and thus  $R \sim F_{(5)}^2 \sim \tau^{-4}$

perfectly physical, due to  
the smeared anti-D3 branes



- However, we also find a highly unusual  
singularity

$$H_{(3)}^2 \sim F_{(3)}^2 \sim \tau^{-2}$$

which has no physical interpretation.





- Despite having finite action, this singularity **alone** in the energy density from the three-forms should be considered an unacceptable singularity.

- However, it could be an artifact of the singularity in the energy density of the curvature and five form due to the smearing

# Numerical Analysis

Bena, Graña, N.H.  
work in progress

The superpotential approach is a nice organizing principle for non-SUSY backgrounds which are first order perturbations around a SUSY background

Papadopoulos, Tseytlin 2000  
Borokhov, Gubser 2002

We have two sets of eight coupled first order O.D.E.'s. The bulk mass of these modes is known numerically but we need the actual modes.

Berg, Haack, Muck 2006

# Numerical Analysis...

Numerically integrating the field

- equations will allow us to relate IR and UV integration constants.

$$\begin{array}{l} \text{recall} \\ \text{UV} \rightarrow Q_{D3}^{Max} = q_b + q_f \\ \qquad \qquad \qquad = -\bar{N} \end{array} \quad \begin{array}{l} \text{IR} \\ \swarrow \end{array}$$

- We will compute  $(q_b, q_f)$  in terms of  $\bar{N}$

- We will also compute  $X_1^{UV}$  the coefficient of the force on a probe D3, in terms of  $\bar{N}$

## Summary

- Breaking supersymmetry by adding anti-D3 branes has been common practice and thus understanding the UV behaviour is a vital problem.
- Supergravity is our best tool since the field theory is strongly coupled everywhere.
- The IR singularities in our linear order analysis around KS suggest that the UV boundary conditions cannot be that of KS
  - This seems to imply that the supersymmetry breaking is explicit. However more work needs to be done:

## Future Directions

- Unsmearing the brane
- Beyond perturbation theory
- Examine stability of our solution