

# FLOW HARMONIC COEFFICIENTS IN SMALL SYSTEMS AT THE LHC: INITIAL OR FINAL STATE EFFECT?

Alba Soto-Ontoso

+Javier L. Albacete, Hannah Elfner(Petersen) and Harri Niemi

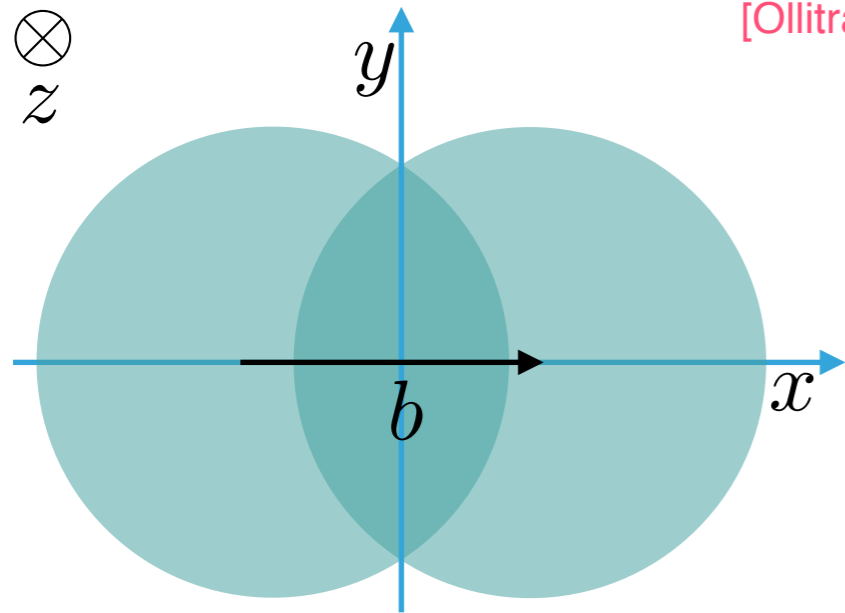
based on Phys. Lett. B770 (2017) 149-153, Phys.Rev. C95 (2017) no. 6, 064909, arXiv:19xx.xxxx

Multiple partonic interactions at the LHC (MPI@LHC)

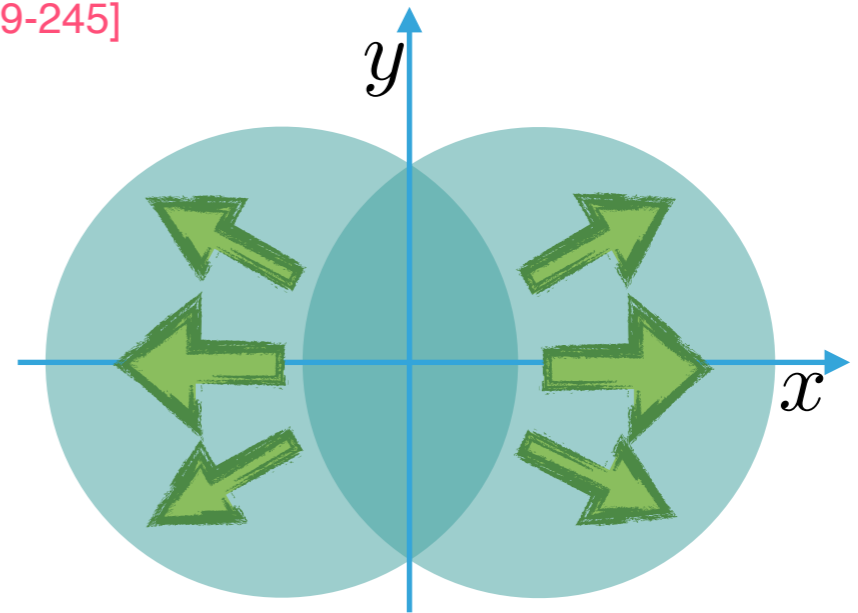
Perugia, 10th December, 2018

# QGP footprints: flow harmonic coefficients

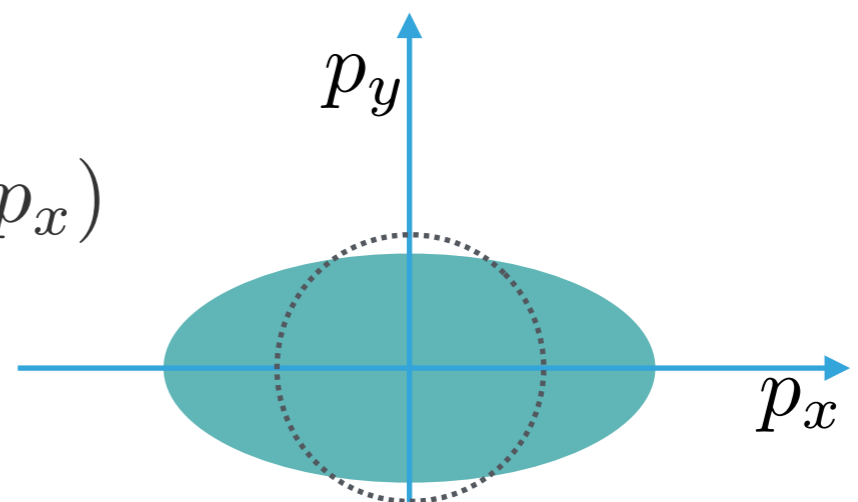
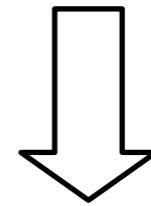
[Ollitrault Phys.Rev. D46 (1992) 229-245]



Initial geometry anisotropy

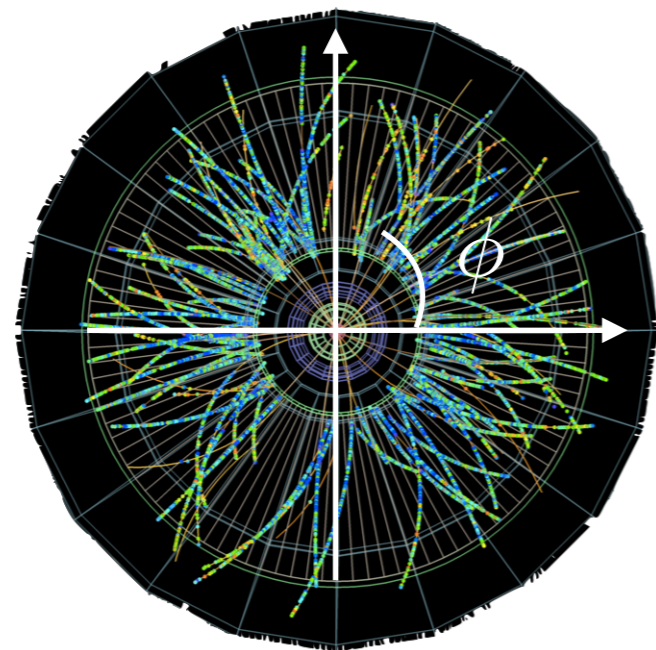
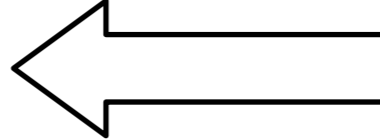


Pressure gradients



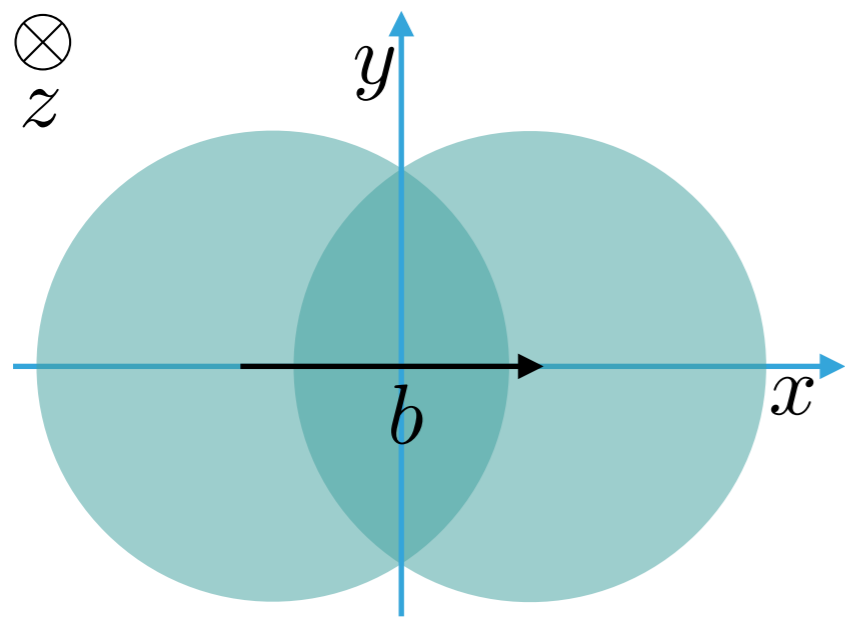
Anisotropy in momentum space

$$\phi = \arctan(p_y/p_x)$$



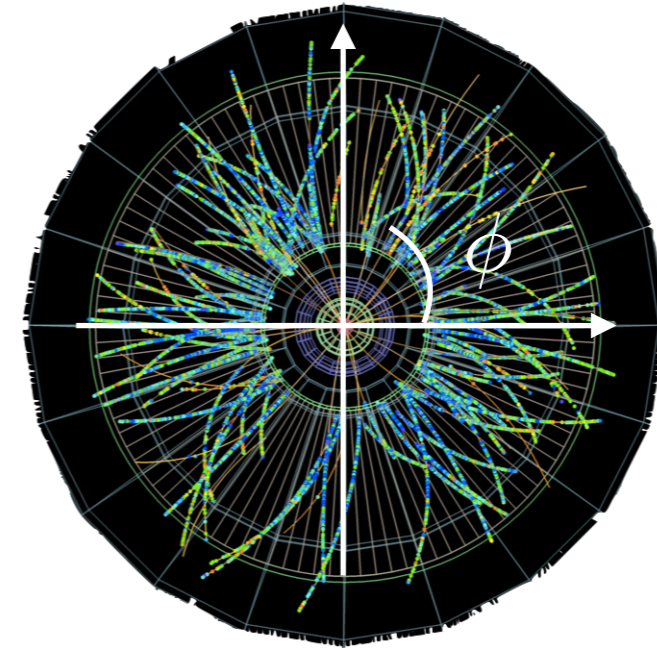
Non-flat azimuthal distribution

# QGP footprints: flow harmonic coefficients



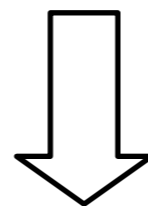
Initial geometry anisotropy

if QGP



Anisotropic azimuthal distribution

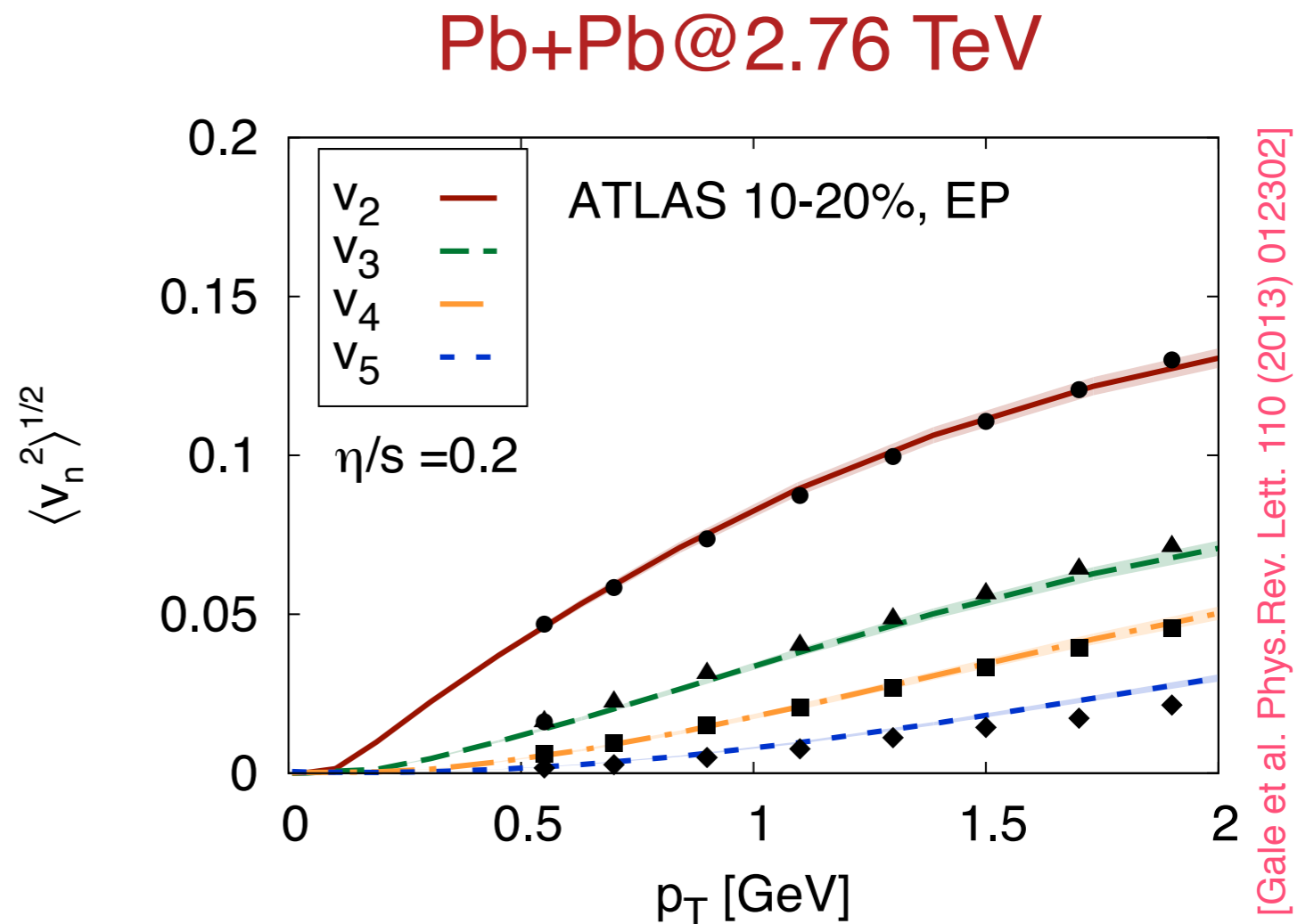
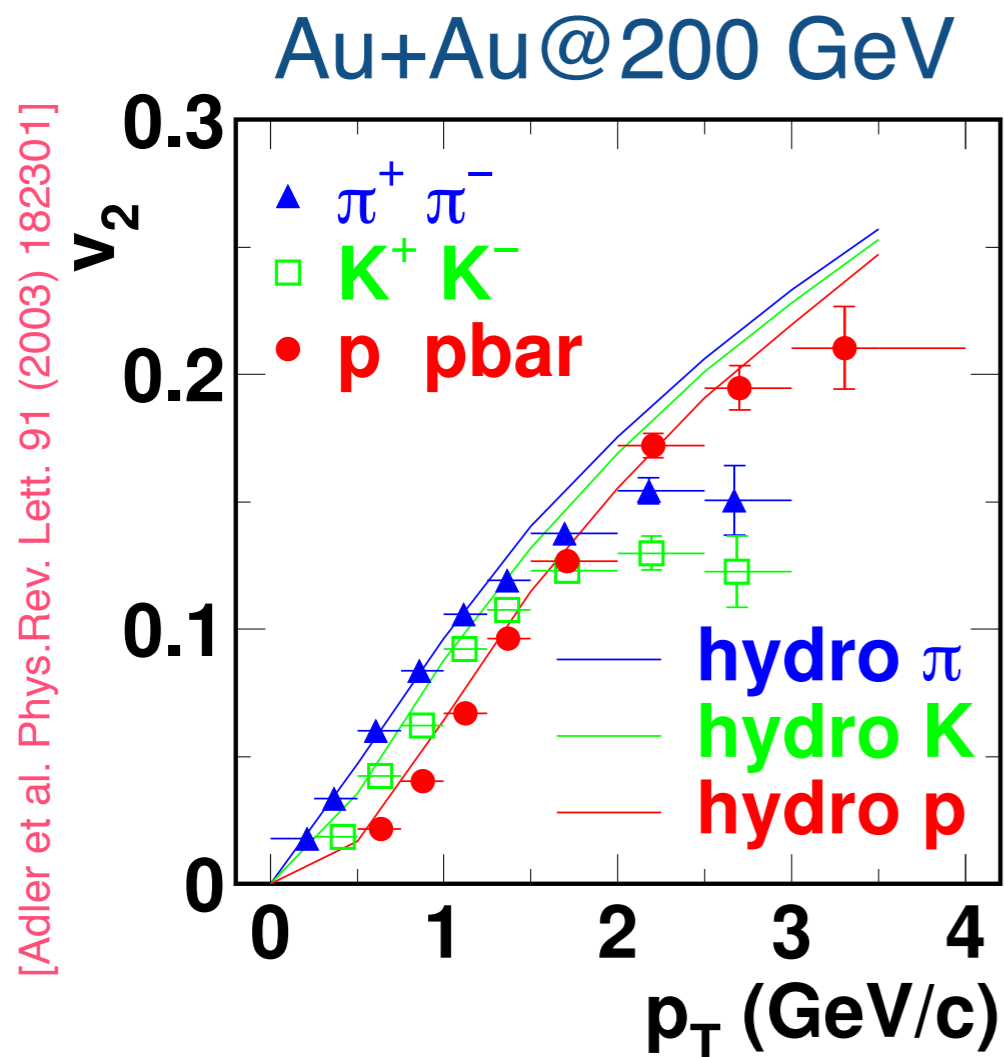
$$E \frac{dN}{d^3p} \propto \frac{dN}{p_T dp_T dy} \left[ 1 + 2 \sum_{n=1}^{\infty} v_n(p_T, y) \cos[n(\phi - \psi_{RP})] \right]$$



$$v_n = \langle \cos(n[\phi_1 - \phi_2]) \rangle = \frac{1}{N(N-1)} \sum_{\substack{i \neq j \\ i, j=1}}^N \cos[n(\phi_i - \phi_j)]$$

\*N: Number of particles in an event

# Flow in heavy ion collisions

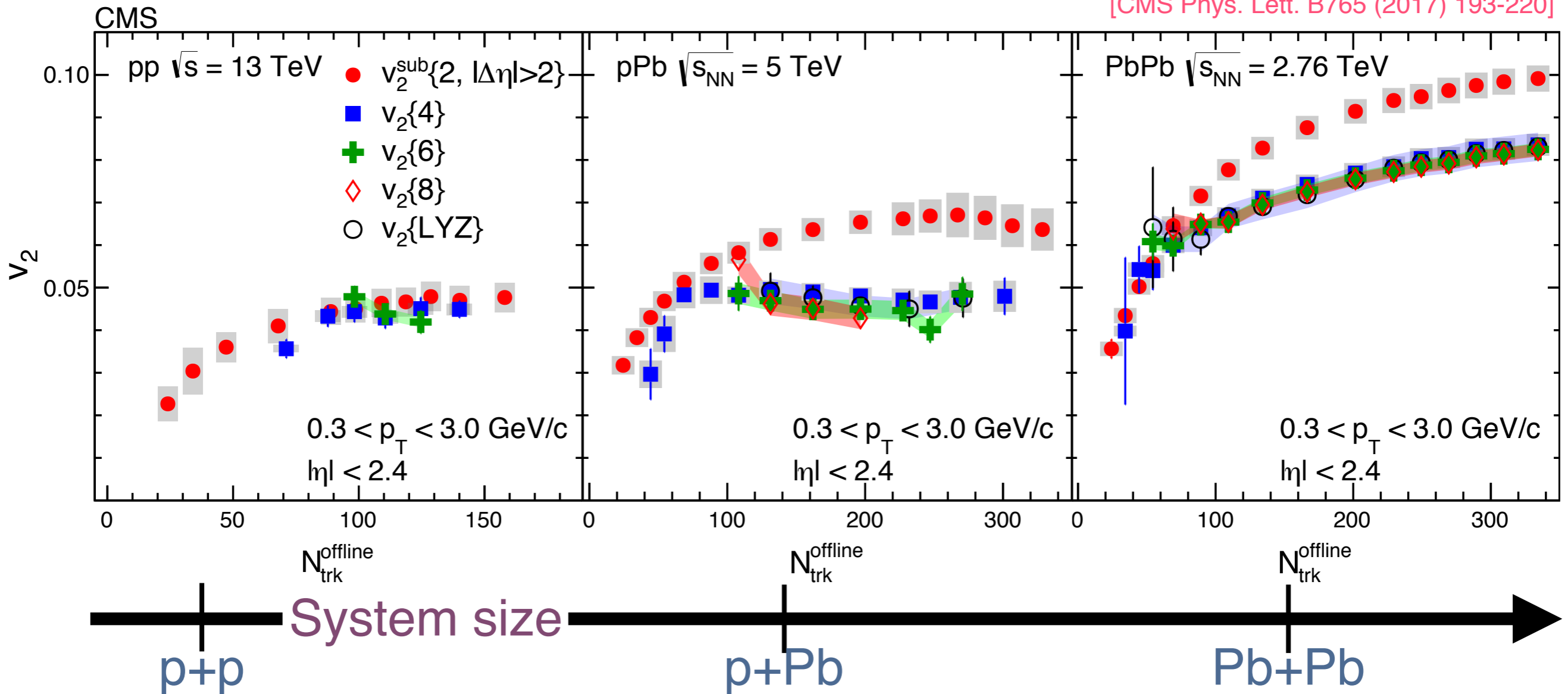


**QGP discovered in ultra relativistic heavy ion collisions\*.  
 Excellent agreement between data and viscous hydro**

\*+ Jet quenching, quarkonia melting, strangeness enhancement...

# Flow in all collision systems@LHC

[CMS Phys. Lett. B765 (2017) 193-220]



**Breaking news from the LHC: non-zero elliptic flow in p+p and p+Pb. “The small system puzzle”**

# Underlying dynamics

CGC

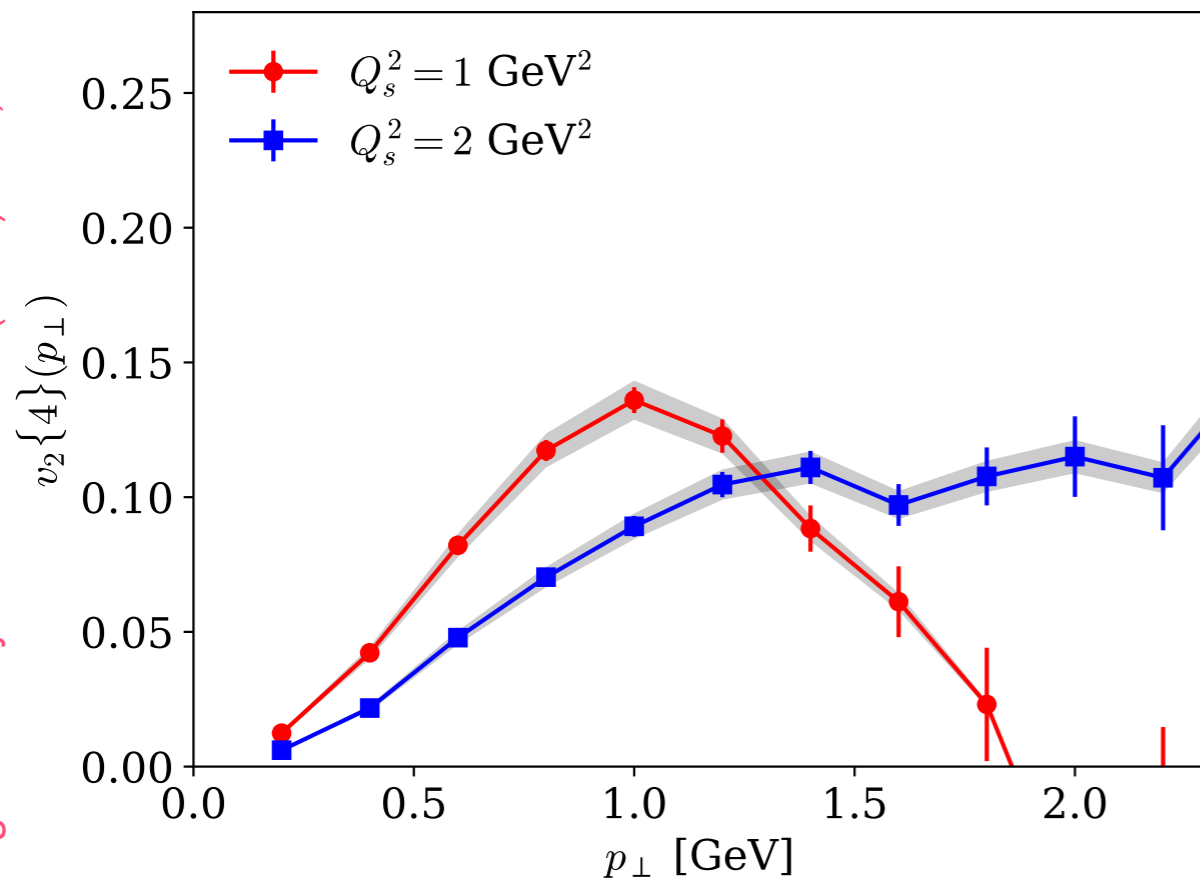
See talk by Skokov

Hydrodynamics

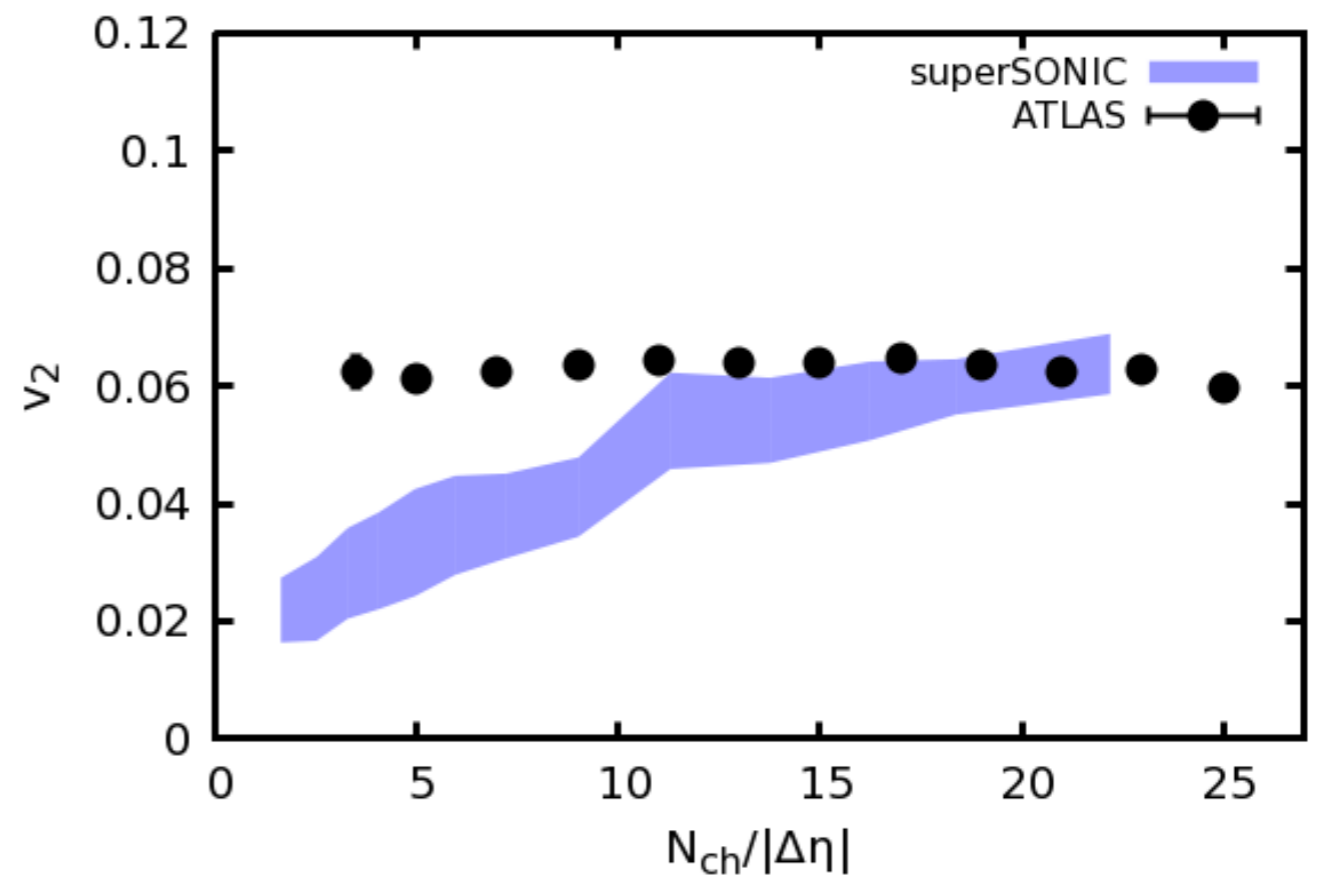
Initial state dynamics; no QGP

Panta Rhei; QGP as in A+A

$v_2\{4\}$



$V_2$

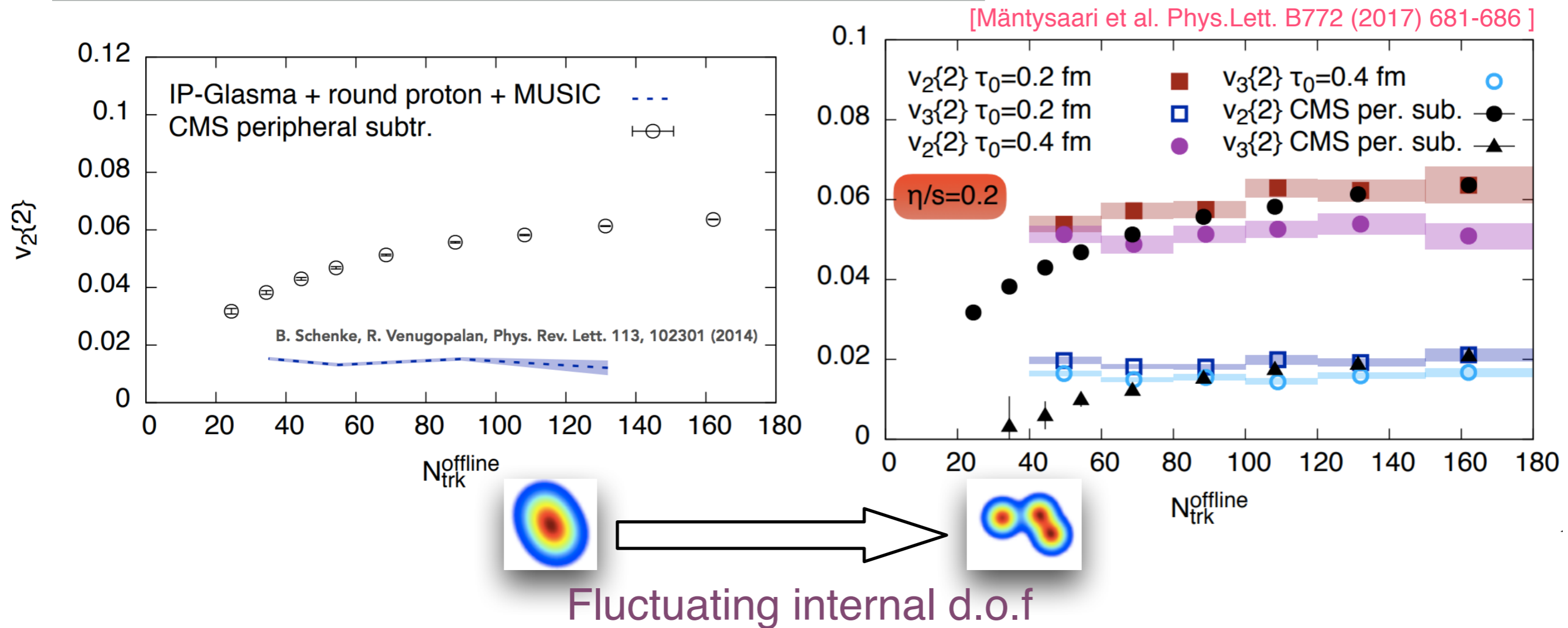


**Common input: geometric structure of the proton**

[Weiler et al. Phys.Lett. B774 (2017) 351-356]

# Role of proton geometry in p+A

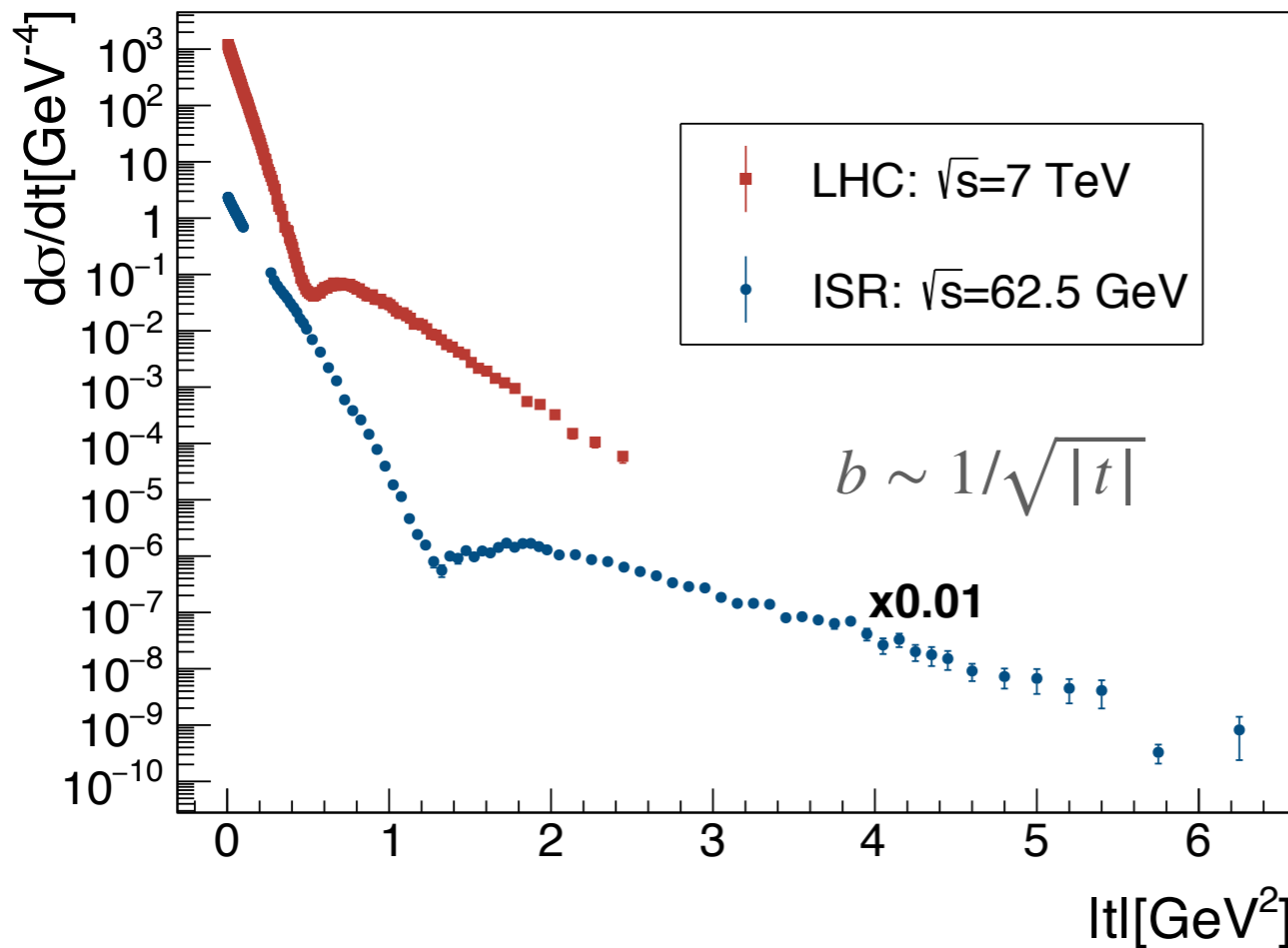
## Initial state dynamics + QGP effects



Highly precise LHC/RHIC data calls for a fine detailed description of the proton structure. How can it be constrained?

# Elastic scattering

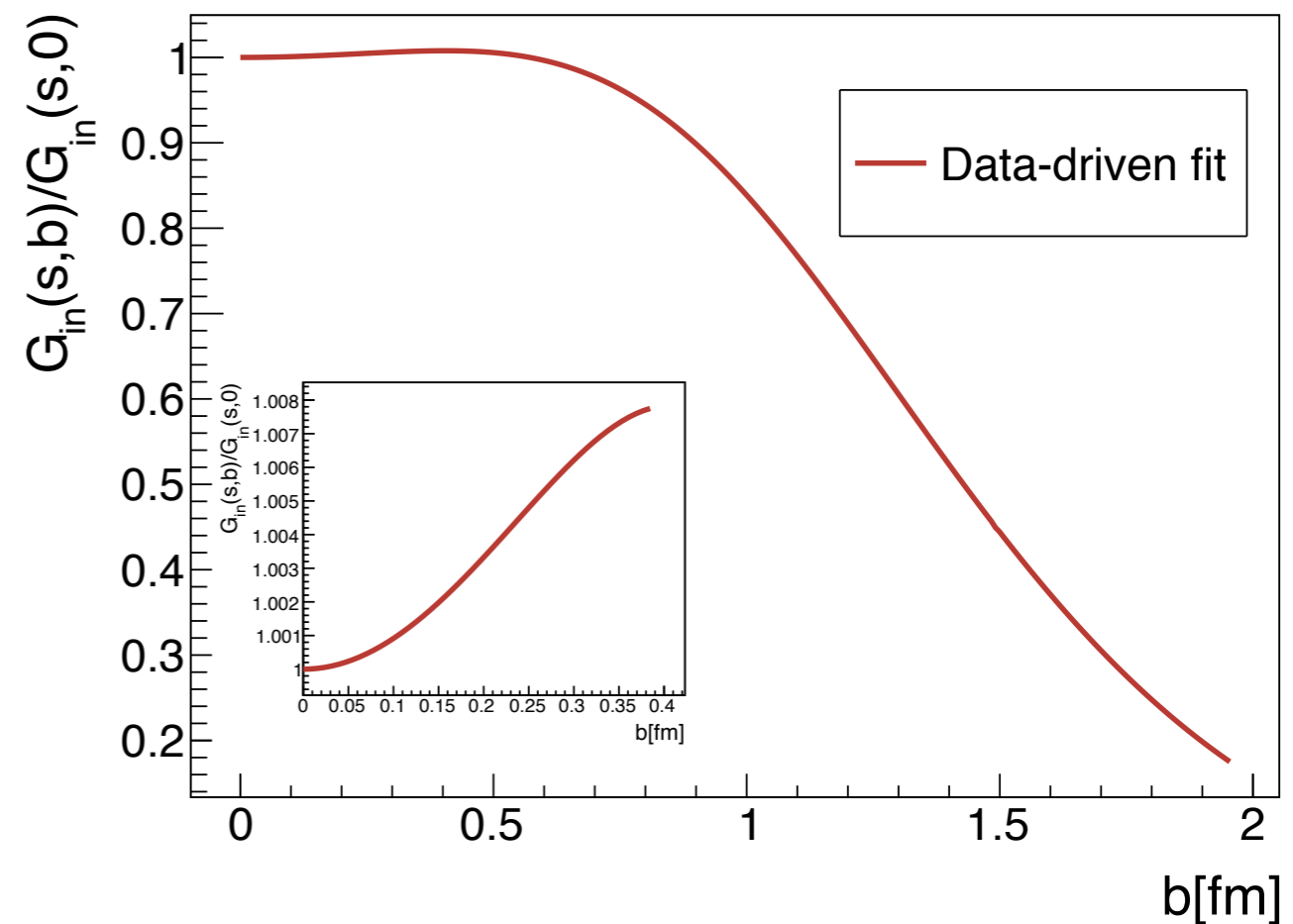
[Albacete, ASO Phys. Lett. B770 (2017) 149-153]



$$\frac{d\sigma_{el}}{dt} \propto |T_{el}(s, t)|^2$$

↓  
Unitarity condition +  
Fourier Transformation

$$G_{in}(s, b) \equiv \frac{d^2\sigma_{inel}}{d^2b} = 2\text{Im}T_{el}(s, b) - |T_{el}(s, b)|^2$$



- **STRIKING** growing behaviour at low impact parameter
- Not observed at **ISR** energies
- Peripheral collisions contribute more to  $\sigma_{inel}$  than central ones

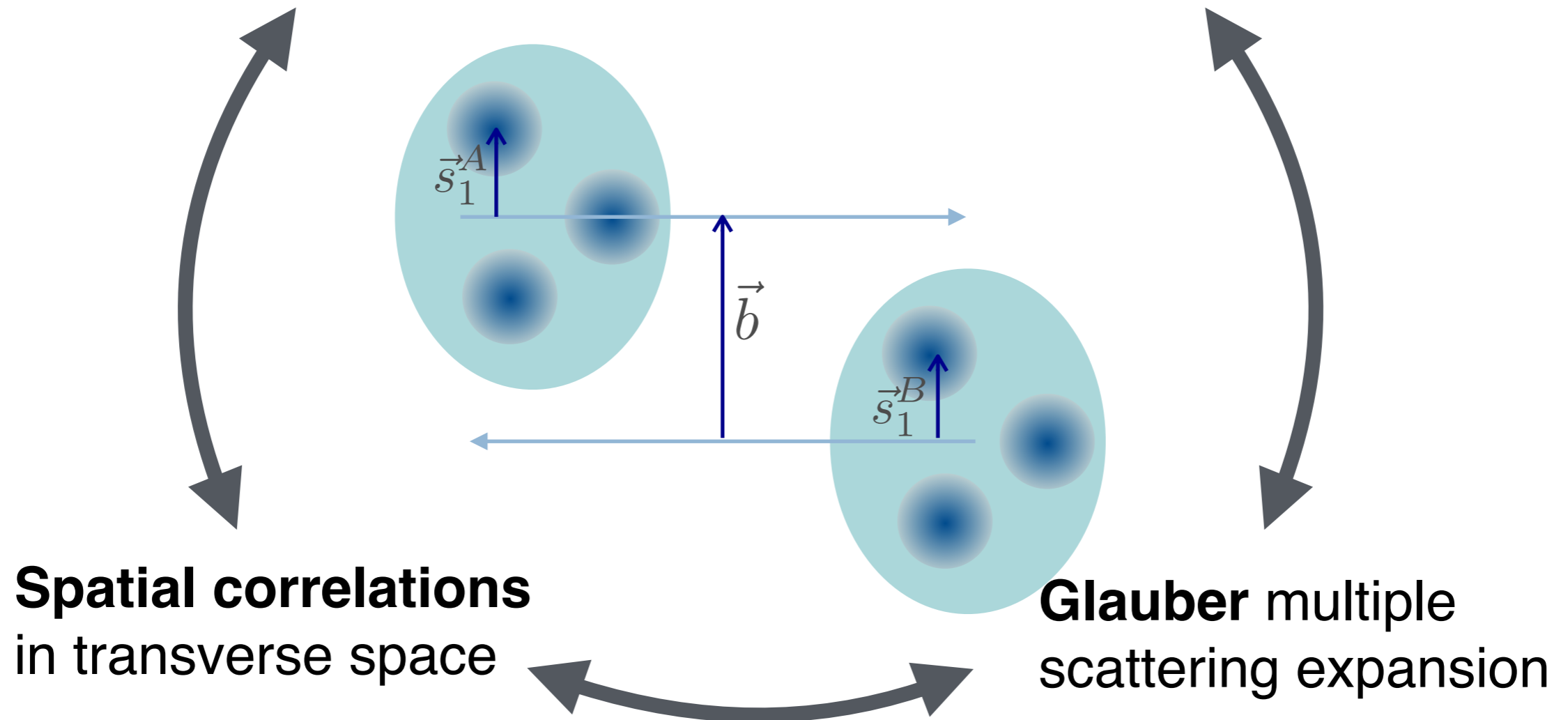
Hollowness effect



# Setup

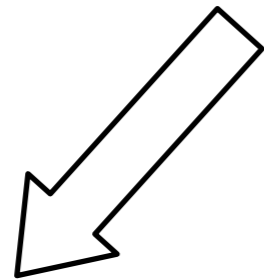
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**Gluonic hot-spots as effective d.o.f**

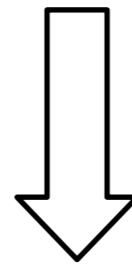


# Spatial distribution of hot spots

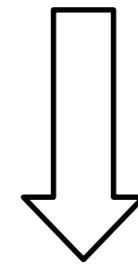
$$D(\vec{s}_1, \vec{s}_2, \vec{s}_3) = C \prod_{i=1}^3 e^{-s_i^2/R^2} \delta^{(2)}(\vec{s}_1 + \vec{s}_2 + \vec{s}_3) \times \prod_{\substack{i < j \\ i, j=1}}^3 \left( 1 - e^{-\mu |\vec{s}_i - \vec{s}_j|^2 / R^2} \right)$$



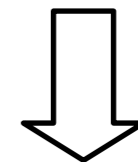
uncorrelated probability distribution characterized by  $R$  (not the proton radius)



fixes the C.o.M of the hot spots system



short range repulsive correlations controlled by  $r_c^2 \equiv R^2 / \mu$

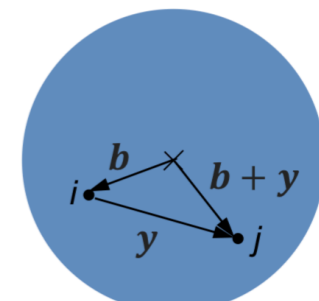


## Simplifying assumptions for DPS cross section

If one ignores correlations between partons in the proton:

$$F^{ij}(x_1, x_2, y) = \int d^2\mathbf{b} D^i(x_1, \mathbf{b}) D^j(x_2, \mathbf{b} + \mathbf{y})$$

Impact parameter dependent PDFs (FT of GPD)



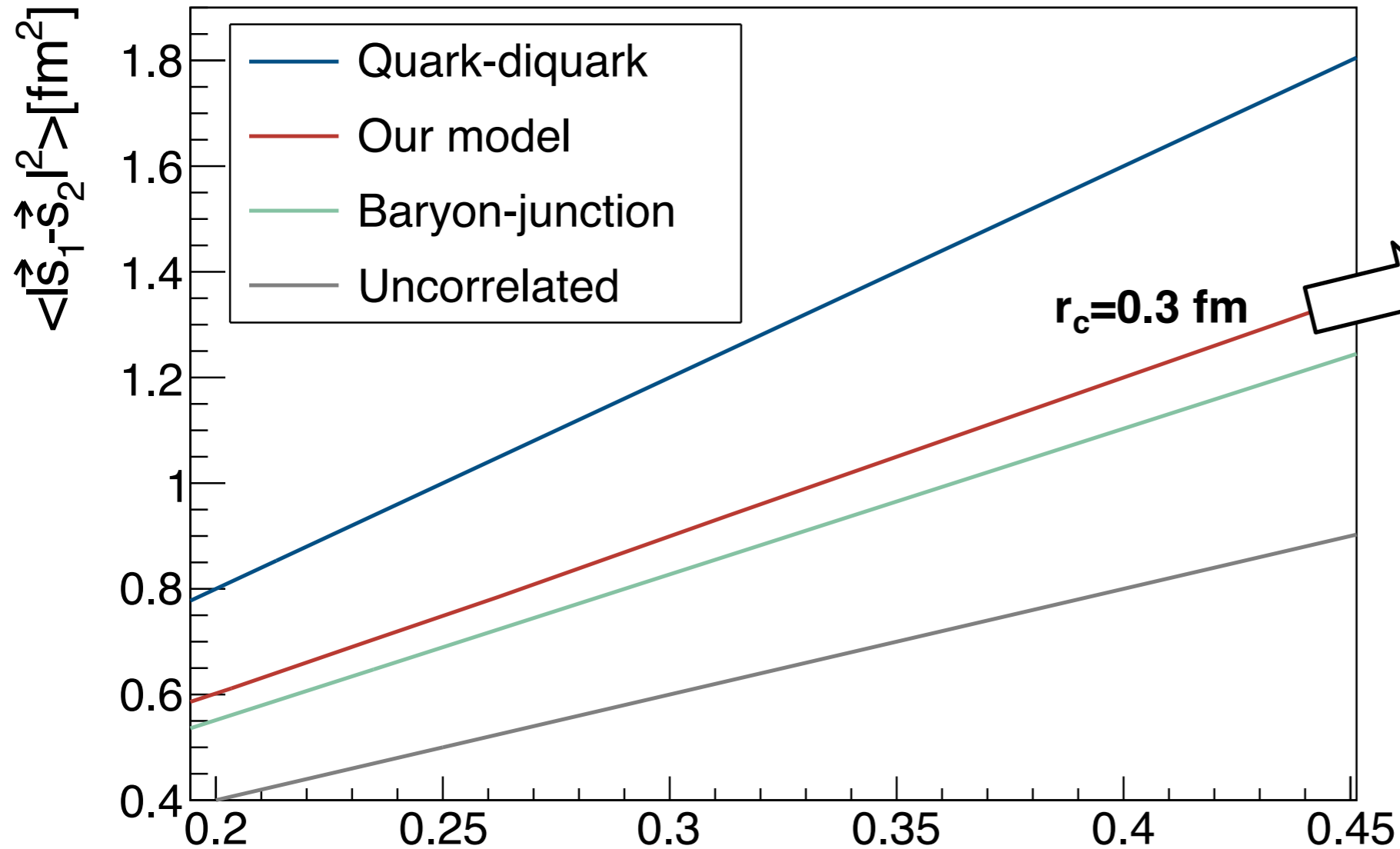
\*C: normalization constant

[Talk by Jonathan Gaunt]

# Spatial distribution of hot spots

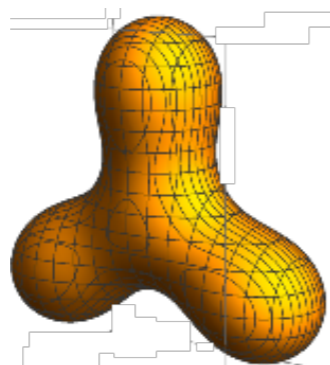
[Albacete, ASO Phys. Lett. B770 (2017) 149-153]

$$\langle X \rangle = \int d\vec{s}_1 d\vec{s}_2 d\vec{s}_3 X D(\vec{s}_1, \vec{s}_2, \vec{s}_3)$$

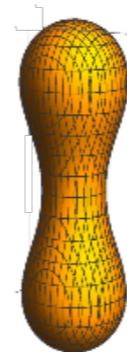


**Interpolates between realistic models of proton substructure**

Baryon-junction



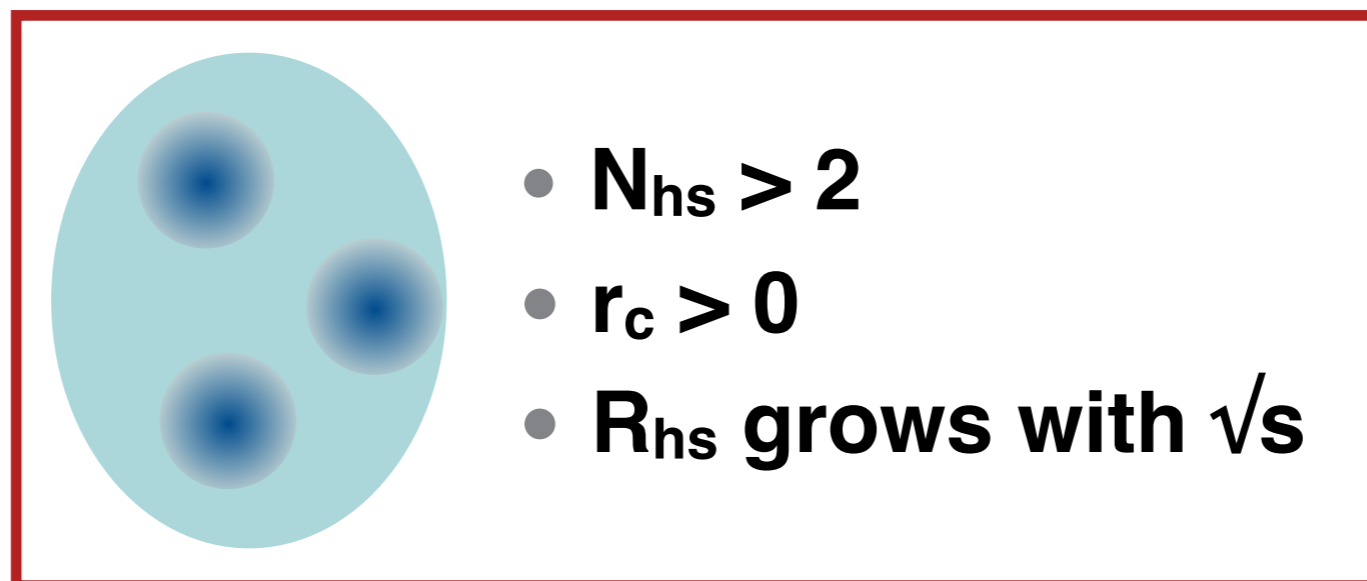
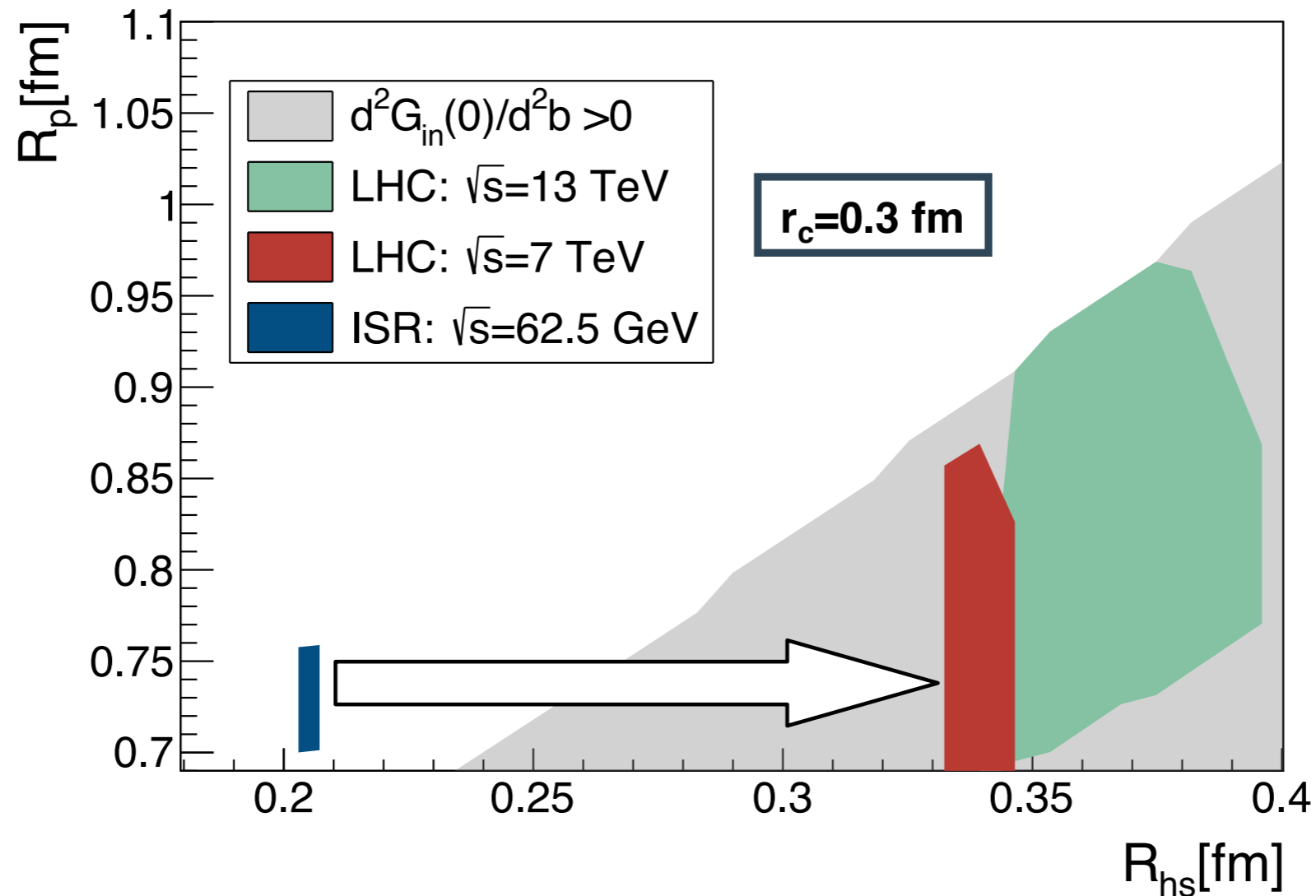
Quark-diquark



$\langle s_1^2 \rangle [\text{fm}^2]$

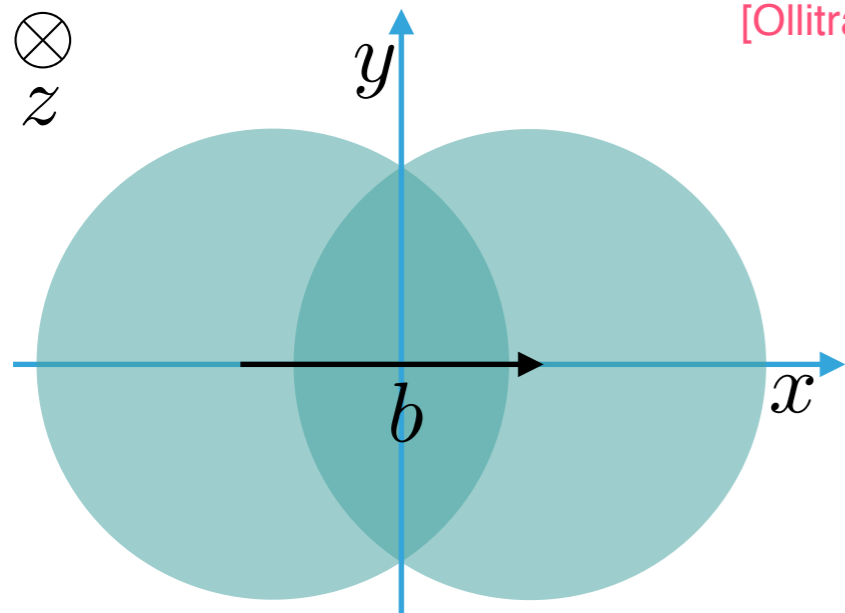
[Kubiczek et al. Lith.J.Phys. 55 (2015) 155]

# Proton transverse structure as a byproduct

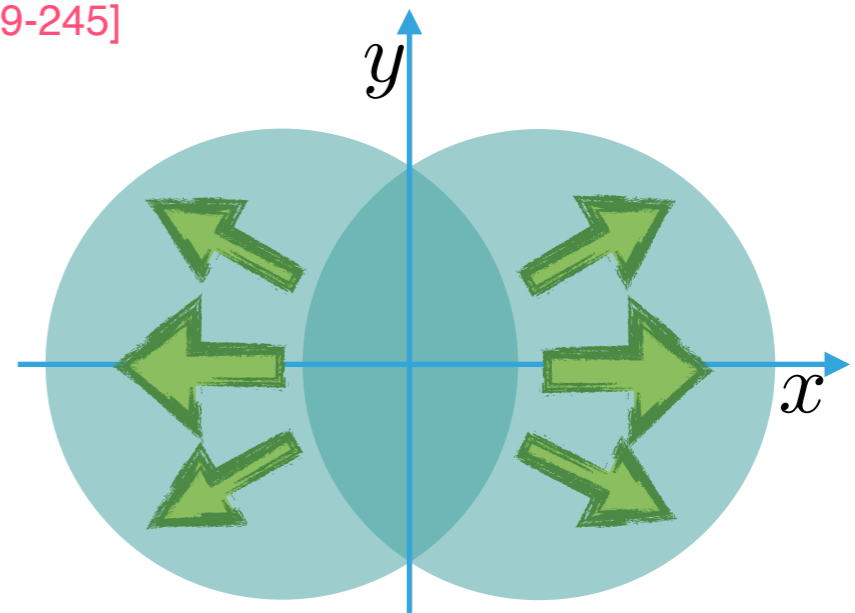


# Test case: hydro responsible for $v_n$ 's in p+p

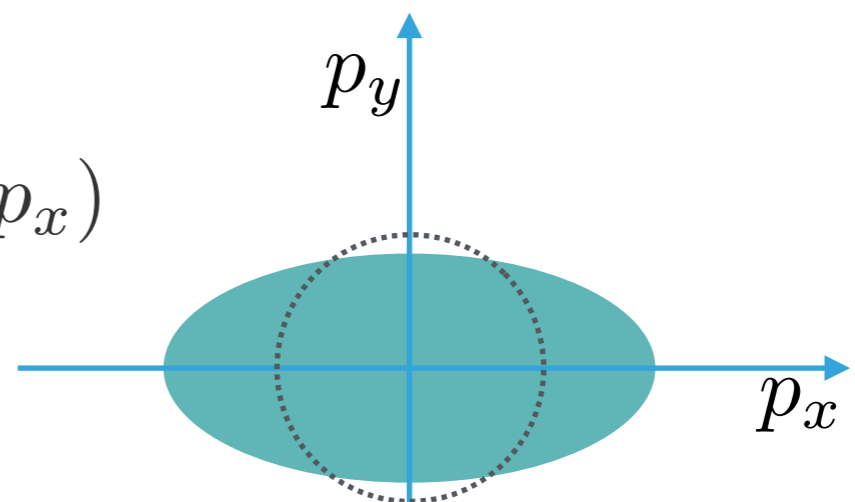
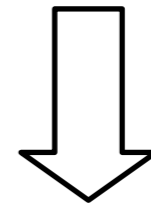
[Ollitrault Phys.Rev. D46 (1992) 229-245]



Initial geometry anisotropy

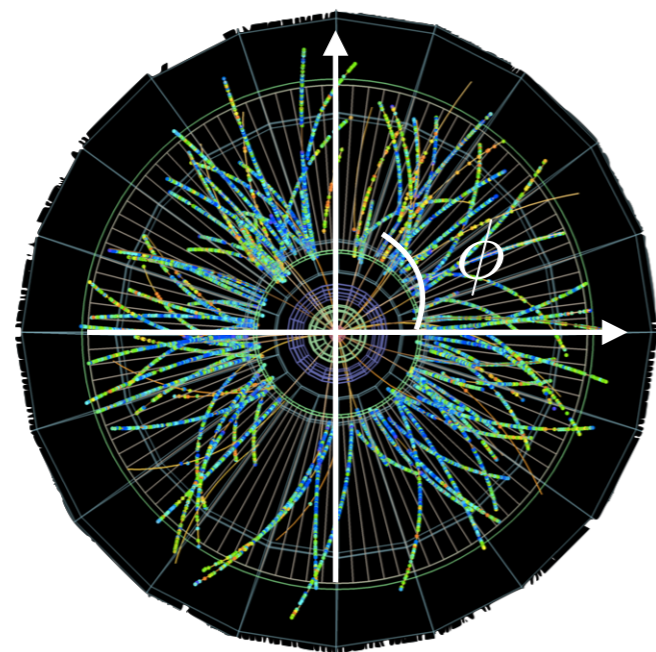
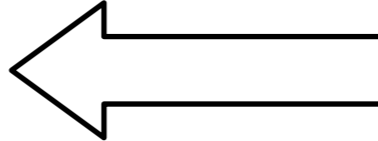


Pressure gradients



Anisotropy in momentum space

$$\phi = \arctan(p_y/p_x)$$



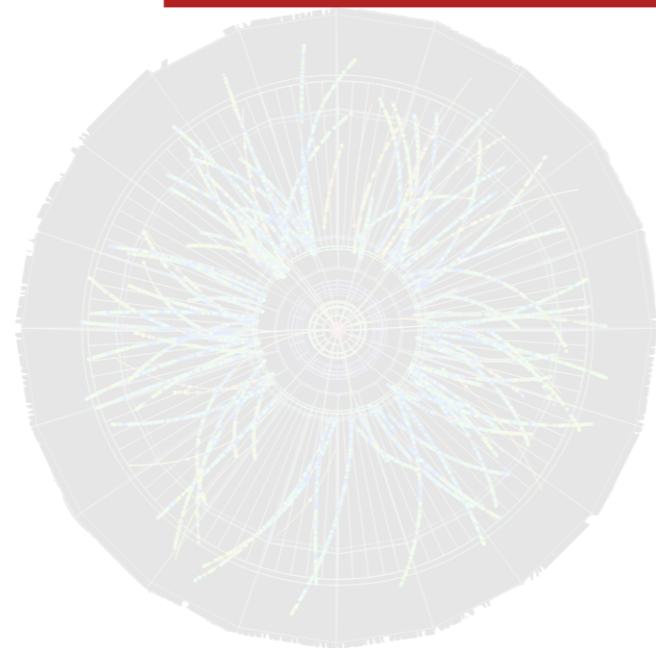
Non-flat azimuthal distribution

# Test case: hydro responsible for $v_n$ 's in p+p

[Ollitrault Phys.Rev. D46 (1992) 229-245]

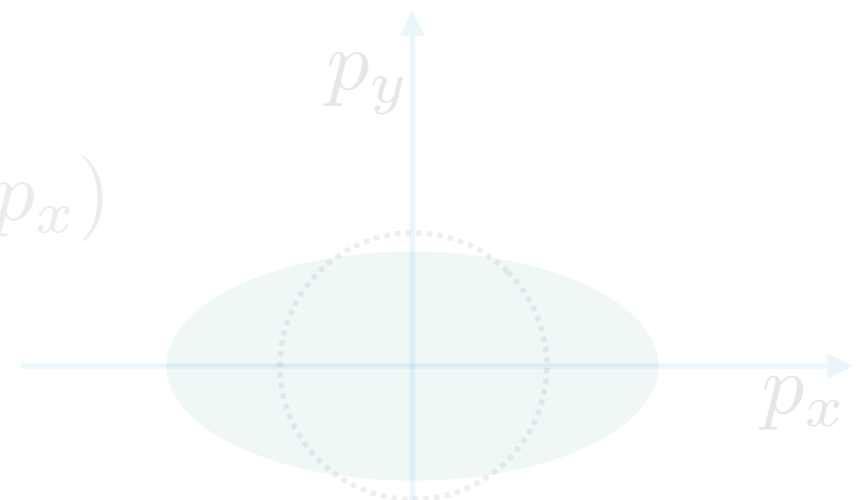


**Goal: gauge the sensitivity of flow harmonic coefficients to **initial geometry** vs. **hydro evolution** parameters in p+p collisions at LHC energies**



Non-flat azimuthal distribution

$$\phi = \arctan(p_y/p_x)$$



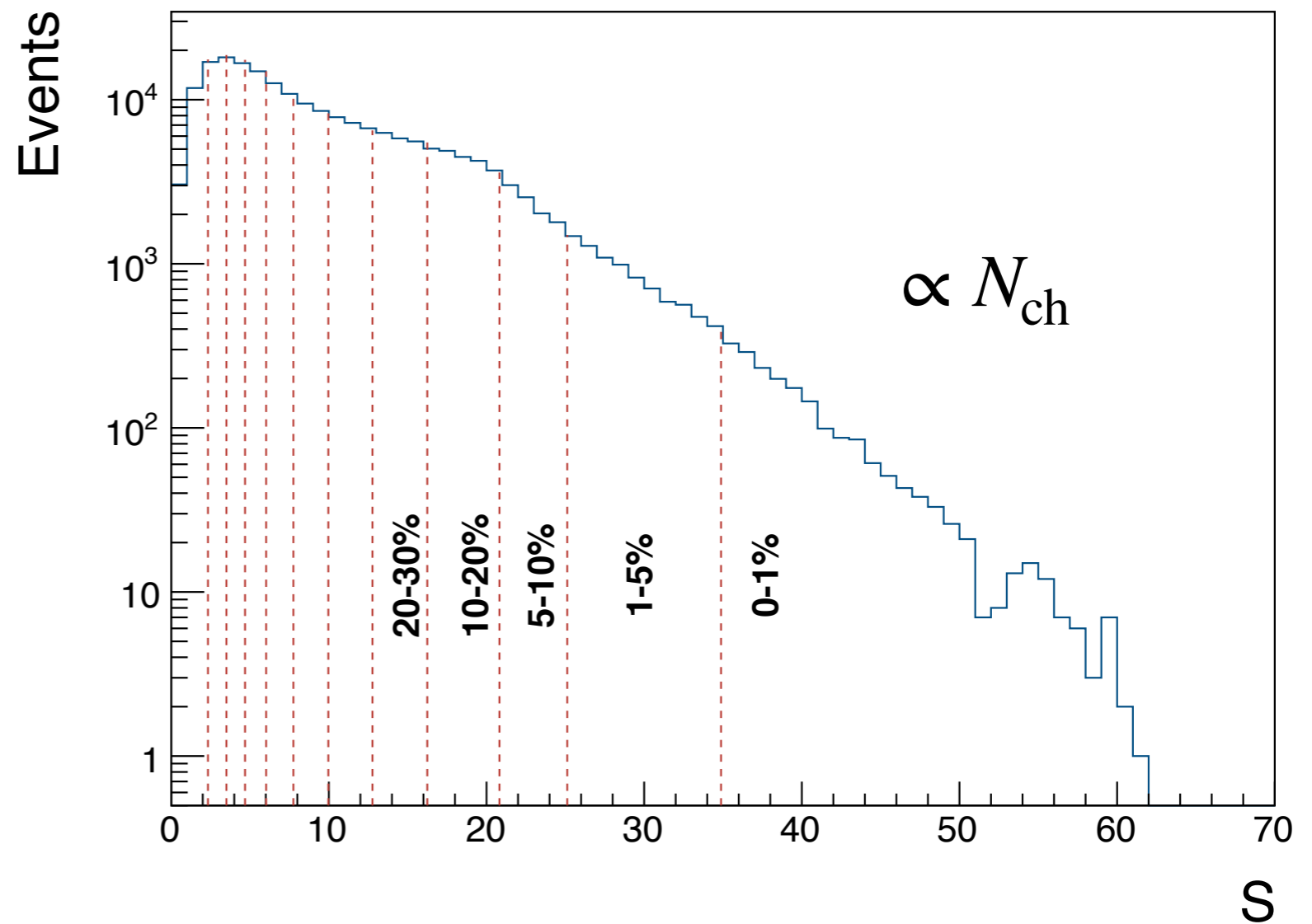
Anisotropy in momentum space

# Initial geometry anisotropy: MC Glauber

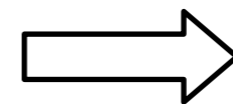
[Albacete, Petersen, ASO Phys. Rev. C95 (6) (2017) 064909]

Analytical  $\Rightarrow$  Monte Carlo: event-by-event fluctuations

- Sampling hot spot positions (correlated/uncorrelated)
- Geometric collision criterion
- Each wounded hot spot deposits entropy in the transverse plane



$$s(x, y) = s_0 \frac{1}{\pi R_{hs}^2} \exp \left( -\frac{(x - x_w)^2 + (y - y_w)^2}{R_{hs}^2} \right)$$



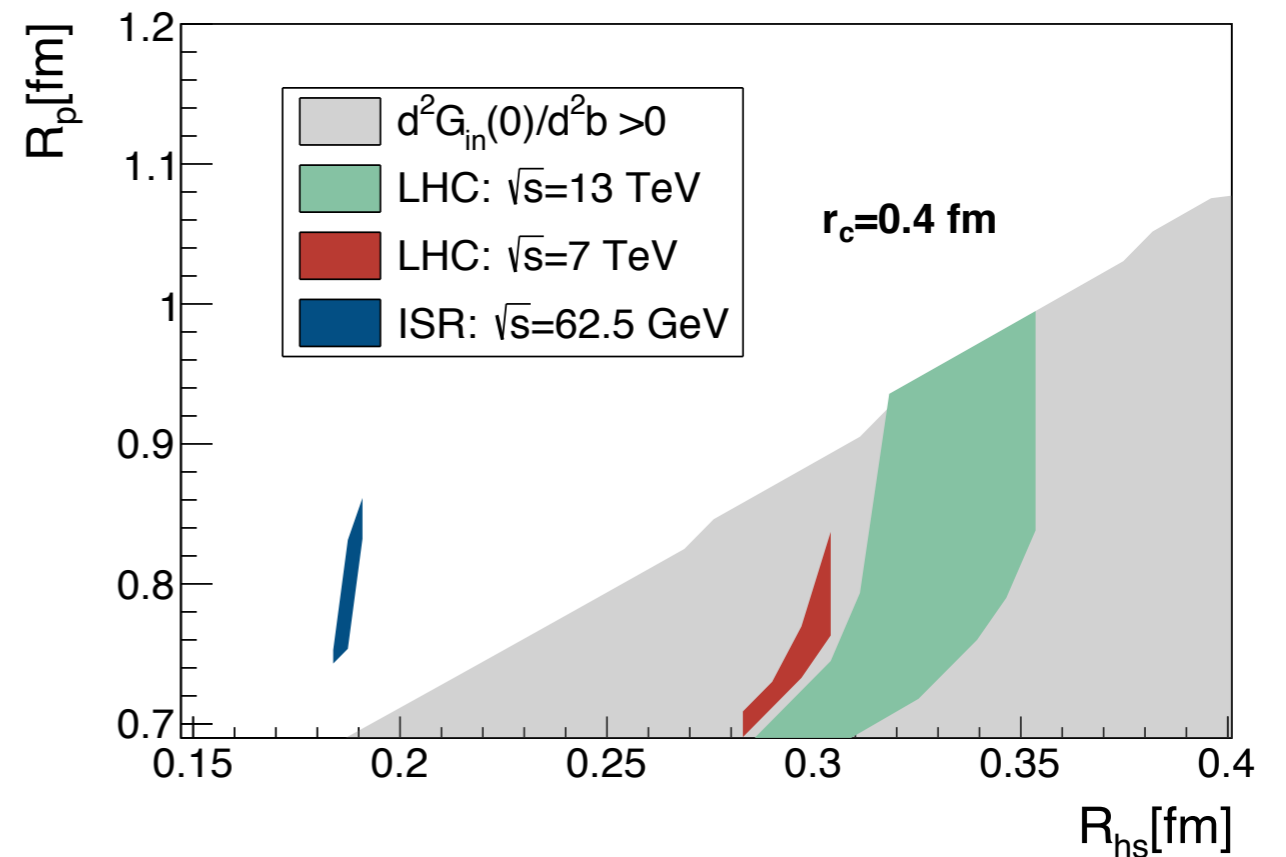
**Input for hydro**

# MC Glauber: parameters

[Albacete, Petersen, ASO Phys. Rev. C95 (6) (2017) 064909]

## ⇒ **Correlated ( $r_c=0.4$ )**

Representative values of the  
hollowness-study phase space  
reproducing



## ⇒ **Uncorrelated**

Problem: Not enough to set  $\mu \rightarrow \infty$ , **swelling effects**

Solution:  $R_{hs}$  as in the **correlated case** and choose  $R$  to have **identical proton size** in both cases

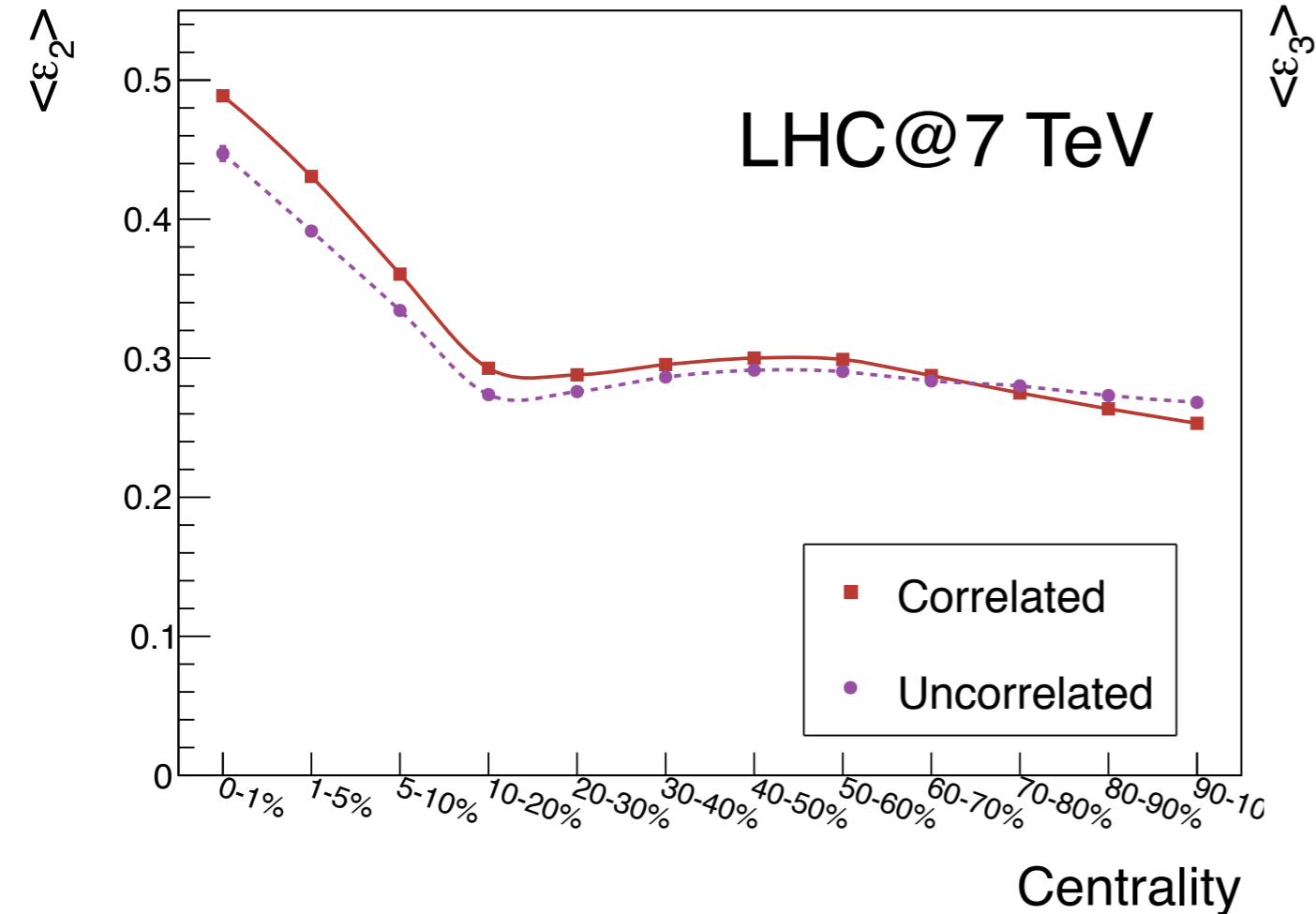
$$\langle s_1 \rangle = \int s_1 d\vec{s}_1 d\vec{s}_2 d\vec{s}_3 D(\vec{s}_1, \vec{s}_2, \vec{s}_3)$$



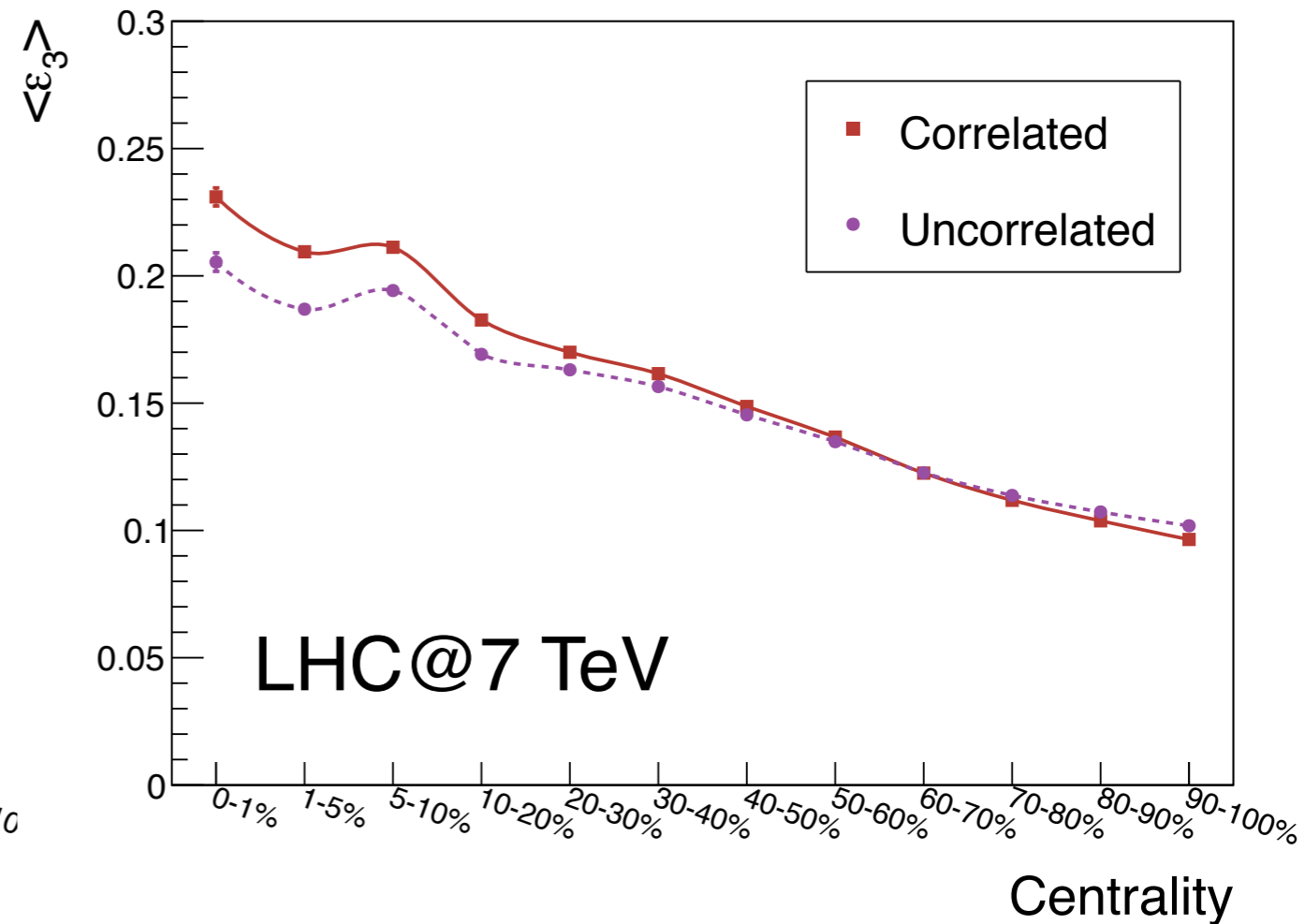
# Eccentricity & triangularity

[Albacete, Petersen, ASO Phys. Rev. C95 (6) (2017) 064909]

## Eccentricity



## Triangularity



**Repulsive correlations make  $\epsilon_{2,3}$  increase in ultra-central collisions**

# 2+1D viscous hydrodynamic evolution

[Albacete, Petersen, ASO Phys. Rev. C95 (6) (2017) 064909]

- Initialized event-by-event with

$$s(x, y, \tau = 0.2 \text{ fm})$$
$$\pi^{\mu\nu} = 0 \quad v_T = 0$$

- Equation of state: s95p-PCE-v1

$$T_{\text{dec}} = 100 \text{ MeV} \quad T_{\text{chem}} = 175 \text{ MeV}$$

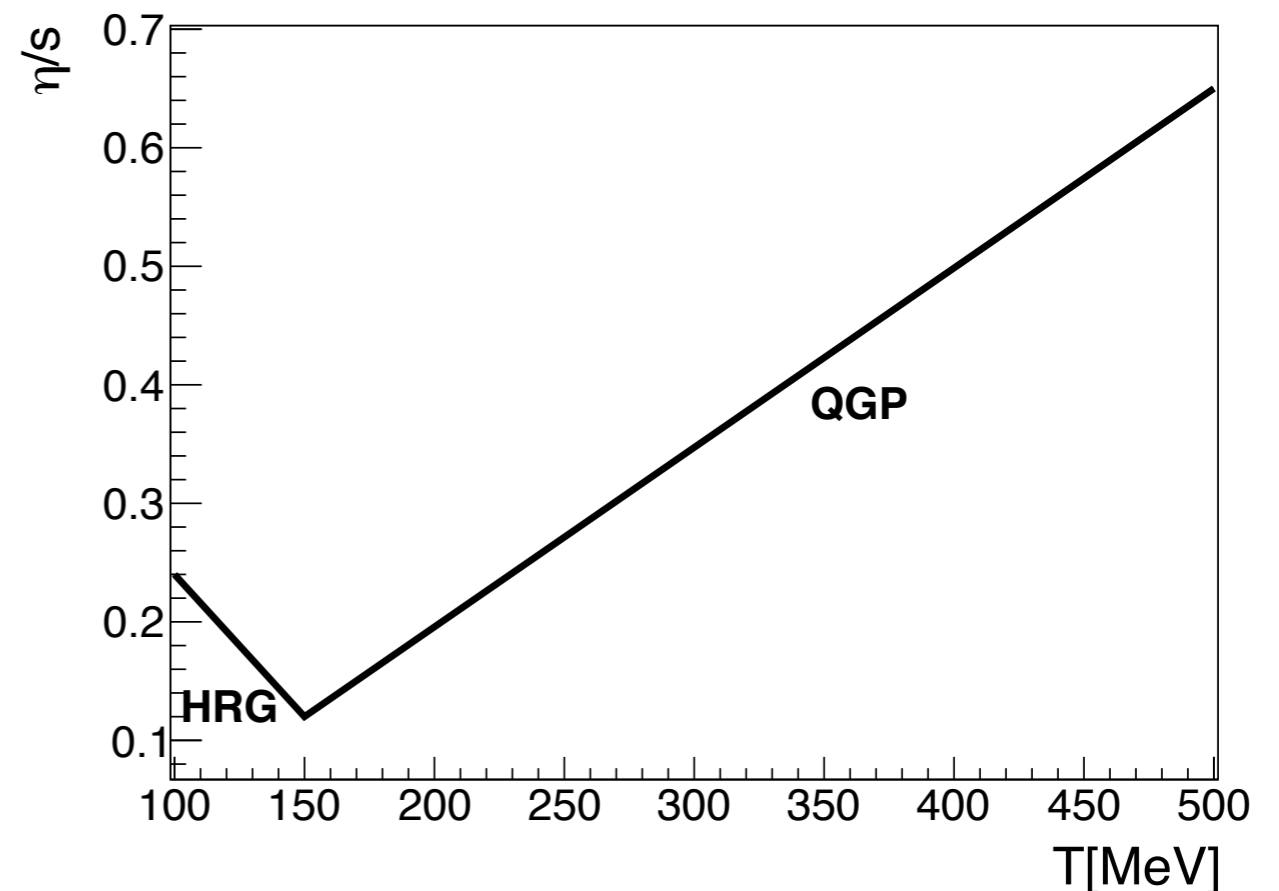
[Huovinen, Petreczky Nucl.Phys. A 837 (2010) 26]

- Resonance decays after freeze out included

- Transport coefficients:  
neglect bulk viscosity and  
heat conductivity. Shear  
viscosity à la ERKT

[Niemi et al., Phys.Rev. C93 (2016) no.2, 024907]

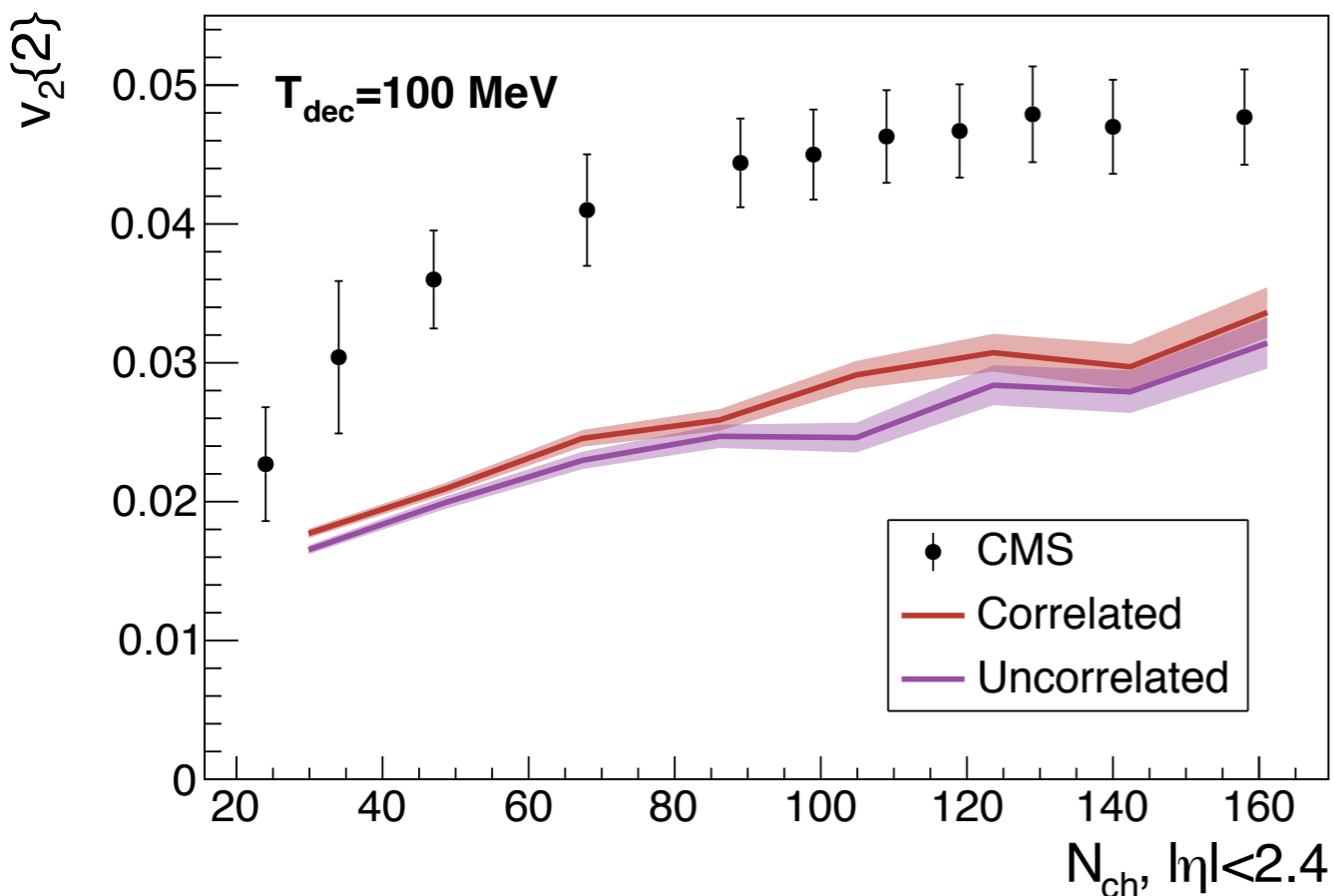
[Eskola et al., Phys.Rev. C93 (2016) no.1, 014912 ]



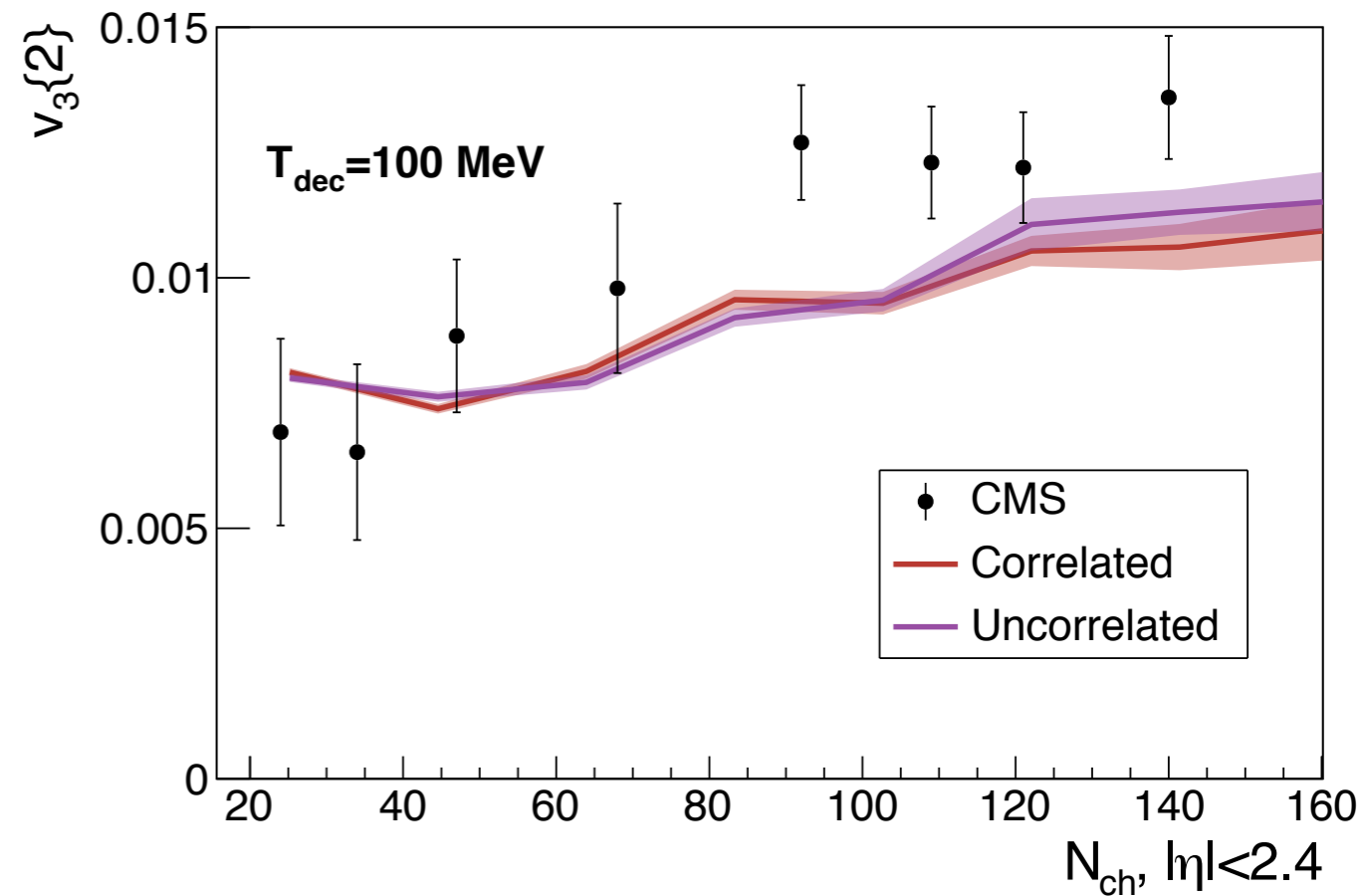
# Impact of spatial correlations in $v_n$ 's

[Albacete, Petersen, Niemi, ASO in preparation]

## Elliptic flow



## Triangular flow

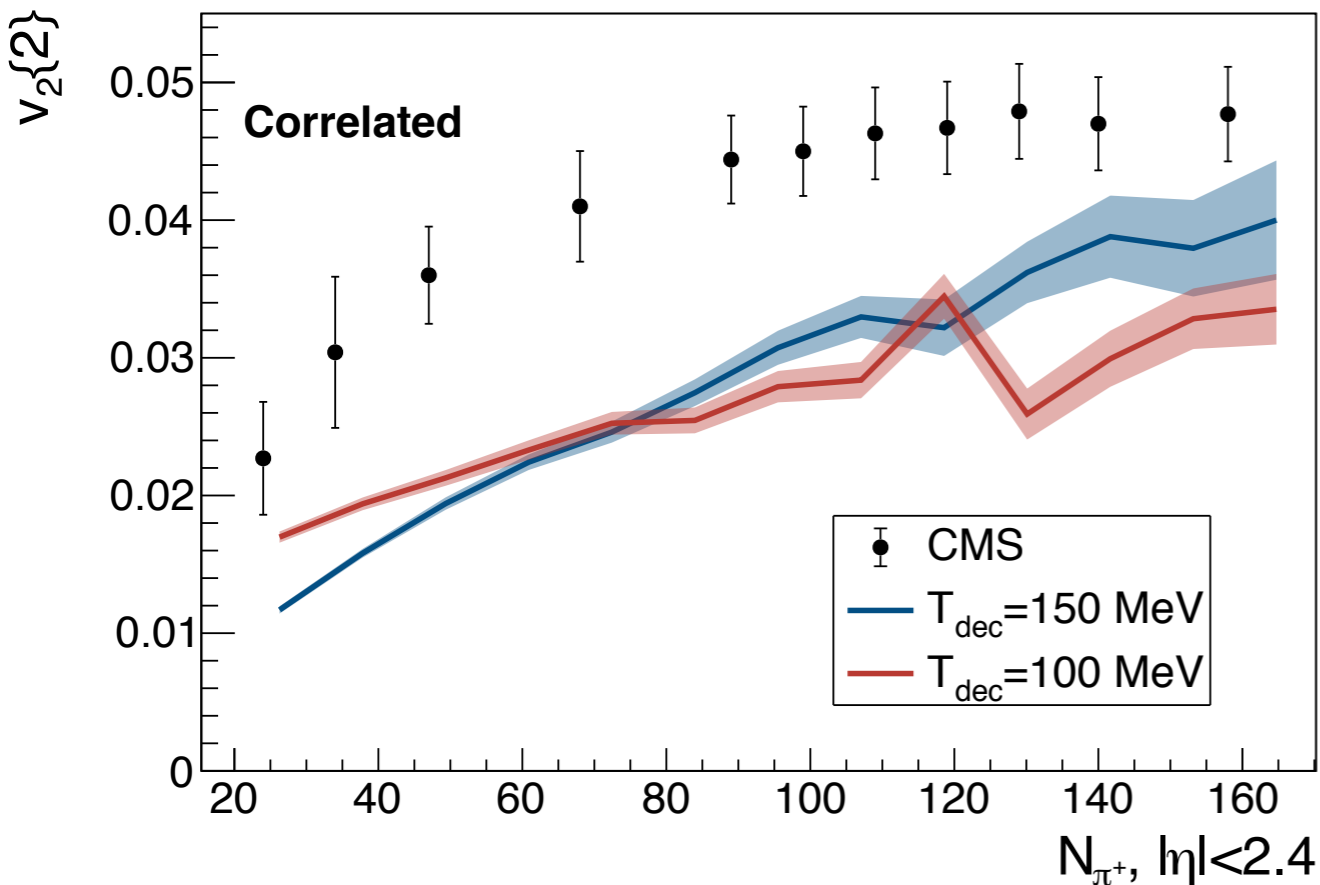


**Repulsive correlations in the initial geometry seem to affect the elliptic flow. No fine tuning. More statistics (and work) needed!!**

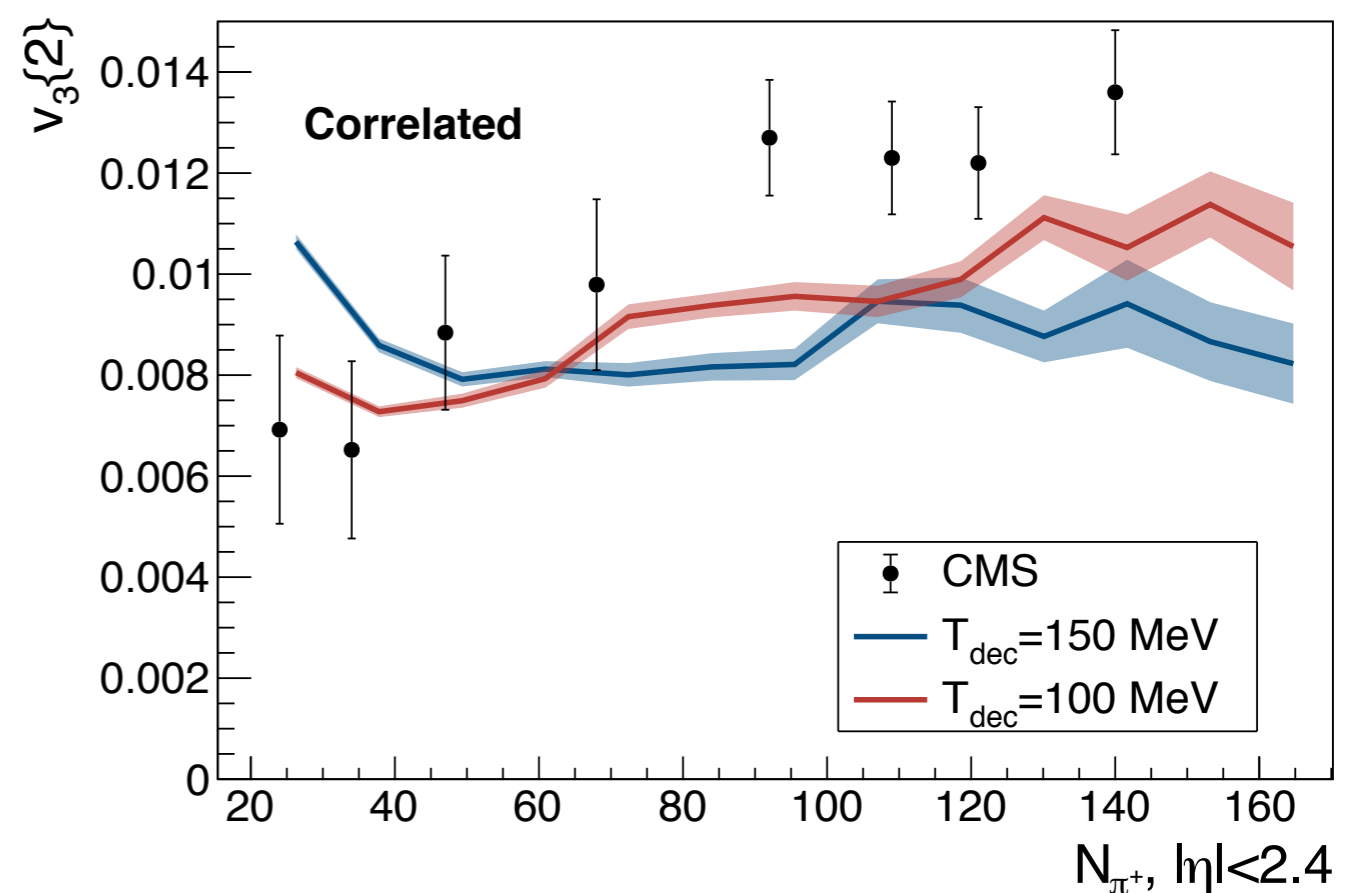
# Impact of $T_{\text{dec}}$ in $v_n$ 's

[Albacete, Petersen, Niemi, ASO in preparation]

## Elliptic flow



## Triangular flow



**Large sensitivity to non-measurable parameters of the hydro evolution observed in the flow coefficients**

# FLOW HARMONIC COEFFICIENTS IN SMALL SYSTEMS AT THE LHC: INITIAL OR FINAL STATE EFFECT?

⇒ Hard to tell!

⇒ Deep and quantitative understanding of details in both paradigms needed before extracting strong conclusions from the experimental data

**Hydro side:** [Albacete, Petersen, Niemi, ASO in preparation] (and others)

**CGC side:** See talk by Skokov

⇒ Constraining power for the proton structure.  
Interesting to the MPI community!