

# New developments in EPOS

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# 1 Multiple scattering in EPOS

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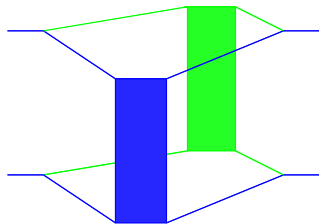
[Parton based Gribov-Regge theory](#). By H.J. Drescher, M. Hladik, S. Ostapchenko, T. Pierog, K. Werner. hep-ph/0007198. Published in Phys.Rept. 350 (2001) 93-289.

Be  $T$  the elastic (pp,pA,AA) scattering T-matrix =>

$$2s \sigma_{\text{tot}} = \frac{1}{i} \text{disc } T$$

**Basic assumption :**  
**Multiple “Pomerons”**

$$iT = \sum_k \frac{1}{k!} \{ iT_{\text{Pom}} \times \dots \times iT_{\text{Pom}} \}$$

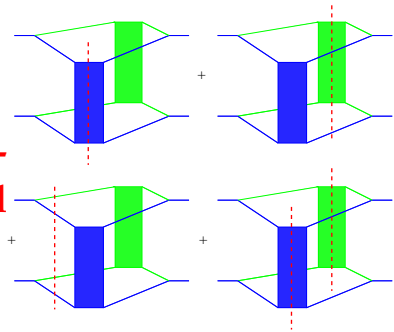


Evaluate

$$\frac{1}{i} \text{disc} \{ iT_{\text{Pom}} \times \dots \times iT_{\text{Pom}} \}$$

using “cutting rules” :

A “cut” multi-Pomeron diagram amounts to the sum of all possible cuts

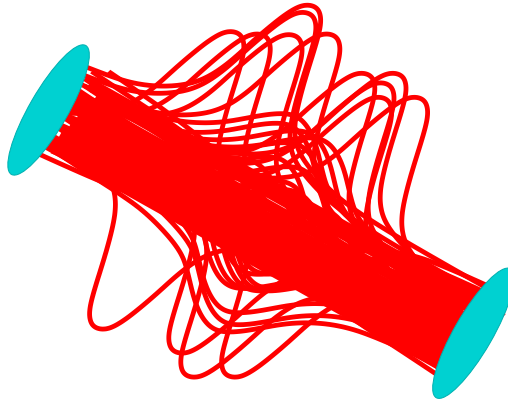


Summing uncut contributions

=> MuScatt weights, high-dim probab. distributions

=> Sampled via Markov chains

**High multiplicity pp or AA: Many cut Pomerons  
=> Many kinky strings**



**=> core + corona**

**core => hydro evolution => statistical decay**

## 2 News

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### Current activities:

### Unification of EPOS LHC & EPOS3

- Implementation of HF
- Saturation and factorization
- **Statistical hadronization**

- **Core-corona separation (pp,pA,AA)**
  
- **EPOS 3:**
  - **Hydrodynamic expansion of core**
  - **Statistical decay of fluid**  
(Grand canonical, big systems)
  
- **EPOS LHC:**
  - **Effective flow, droplet decay**  
(like resonance decay, small systems, small hadron list)
  
- **“Unification”:** **Microcanonical decay**  
for small and big systems, works for pp and PbPb

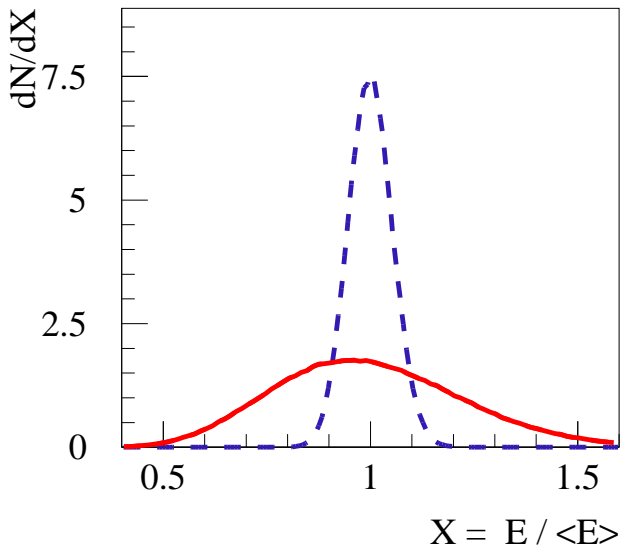
## **2.1 Microcanonical hadronization of plasma droplets**

- No need to match dynamical part of hydro evolution**
- Energy and flavor conservation for small systems**
- Needed to “unify” EPOSLHC and EPOS3**



## Grand canonical decay, $T = 130$ MeV

$V=50 \text{ fm}^3$ ;  $V=1000 \text{ fm}^3$



## Microcanonic decay

of given volume in its CMS into  $n$  hadrons

$$dP = C_{\text{vol}} C_{\text{deg}} C_{\text{ident}} \times \delta(E - \Sigma E_i) \delta(\Sigma \vec{p}_i) \prod_A \delta_{Q_A, \Sigma q_{Ai}} \prod_{i=1}^n d^3 p_i$$

$$C_{\text{vol}} = \frac{V^n}{(2\pi\hbar)^{3n}}, \quad C_{\text{deg}} = \prod_{i=1}^n g_i, \quad C_{\text{ident}} = \prod_{\alpha \in \mathcal{S}} \frac{1}{n_\alpha!}$$

( $n_\alpha$  is the number of particles of species  $\alpha$ ,  $\mathcal{S}$  is the set of particle species)

**Different from decay rate of a massive particle (using LIPS), where asymptotic states are defined over an infinitely large volume**

**(see Becattini et al, EPJC35:243-258,2004). But  $E_i = \sqrt{p_i^2 + m_i^2}$**

## Microcanonical decay

$$dP \propto d\Phi_{\text{NRPS}} = \delta(M - \Sigma E_i) \delta(\Sigma \vec{p}_i) \prod_{i=1}^n d^3 p_i$$

- Hagedorn 1958 methods to compute  $\Phi_{\text{NRPS}}$
- Lorentz invariant phase space (LIPS) (James 1968)
- Hagedorn methods used for decaying QGP droplets (Werner, Aichelin, 1994, Becattini 2003)
- 2012 (Bignamini, Becattini, Piccinini) compute  $\Phi_{\text{NRPS}}$  via the Lorentz invariant phase space (LIPS)

- **LIPS method very fast for small  $n$ , gets time consuming at large  $n$**
- **Hagedorn integral method can be made very efficient at large  $n$  (new), while it is VERY time consuming at small  $n$**
- **around  $n \approx 30 - 40$  both methods work (=> checks)**

## Hagedorn integral method

The phase-space integral:

$$\begin{aligned} \phi_{\text{NRPS}}(M, m_1, \dots, m_n) \\ = (4\pi)^n \int \prod_{i=1}^n p_i^2 \delta(E - \sum_{i=1}^n E_i) W(p_1, \dots, p_n) \prod_{i=1}^n dp_i, \end{aligned}$$

with the “random walk function”  $W$  (angular integral)

$$W(p_1, \dots, p_n) := \frac{1}{(4\pi)^n} \int \delta\left(\sum_{i=1}^n p_i \times \vec{u}_i\right) \prod_{i=1}^n d\Omega_i$$

**New:** Very efficient procedures to compute  $W$  for large  $n$  with high precision

We obtain (Werner, Aichelin 94)

$$\phi(M, m_1, \dots, m_n) = \int_0^1 dr_1 \dots \int_0^1 dr_{n-1} \psi(r_1, \dots, r_{n-1})$$

$$\psi = \frac{(4\pi)^n T^{n-1}}{(n-1)!} \prod_{i=1}^n p_i E_i W(p_1, \dots, p_n),$$

with  $z_i = r_i^{1/i}$ ,  $x_i = z_i x_{i+1}$ ,  $s_i = x_i T$ ,  $t_i = s_i - s_{i-1}$ ,  
 $E_i = t_i + m_i$ ,  $T = M - \sum_{i=1}^n m_i$

**Suitable for MC provided  $W$  is known**

Sampling hadron configurations  $K = \{h_1, \dots, h_n; \vec{p}_1, \dots, \vec{p}_n\}$   
via Markov chains

We construct sequences of random configurations

$$K_1, K_2, K_3, \dots, K_t, \dots$$

such that  $f_t(K_t)$  converges towards  $f(K)$  for  $t \rightarrow \infty$

with  $f$  = microcanonical probability distribution

**The law changes step by step (  $f_t \rightarrow f_{t+1}$  ) :**

$$f_{t+1}(K) = \sum_{K'} f_t(K') p(K' \rightarrow K)$$

**with  $p(K \rightarrow K')$  of the form**

$$w(K \rightarrow K') \times \min \left( 1, \frac{f(K')}{f(K)} \frac{w(K' \rightarrow K)}{w(K \rightarrow K')} \right)$$



## 2.2 Grand canonical limit

For very large  $M$  we should recover the “grand canonical limit” for single particle spectra:

$$f_k = \frac{g_k V}{(2\pi\hbar)^3} \exp\left(-\frac{E_k}{T}\right),$$

The average energy is

$$\bar{E} = \frac{g_k V}{(2\pi\hbar)^3} \sum_k \int_0^\infty E_k \exp\left(-\frac{E_k}{T}\right) 4\pi p^2 dp$$

Changing variables via  $E_k dE_k = p dp$ , and using  $K_1(z) = z \int_1^\infty \exp(-zx) \sqrt{x^2 - 1} dx$ , and  $3K_2(z) = z^2 \int_1^\infty \exp(-zx) \sqrt{x^2 - 1}^3 dx$

=>

$$\bar{E} = \frac{4\pi g_k V}{(2\pi\hbar)^3} m^2 T \left( 3TK_2\left(\frac{m}{T}\right) + mK_1\left(\frac{m}{T}\right) \right).$$

The microcanonical decay of an object of mass  $M$  and volume  $V$  should converge (for  $M \rightarrow \infty$ ) to the GC single particle spectra

**with  $T$  obtained from  $M = \bar{E}$ .**

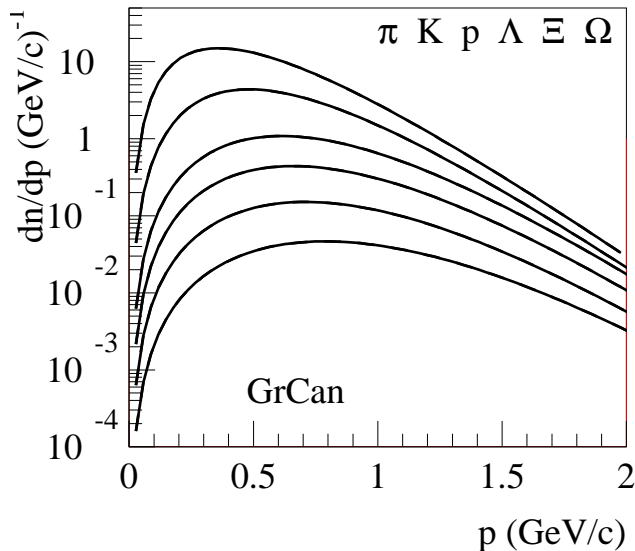
## 2.3 Comparing GC et MiC decay

We consider a complete (?) set of hadrons  
( $\approx 400$ , PDG list)

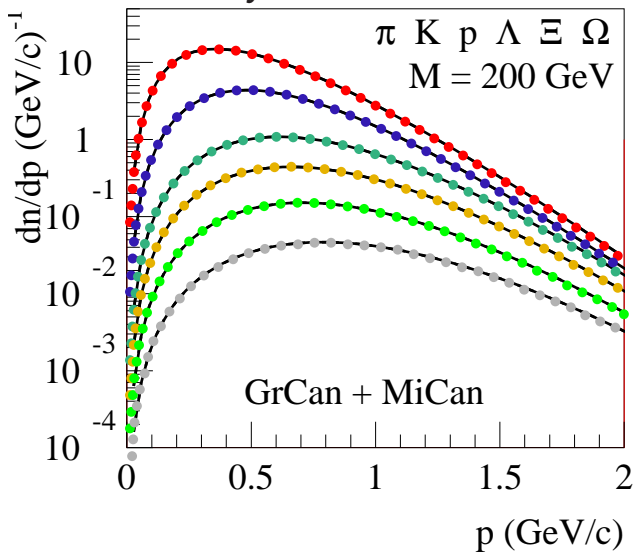
We check the effect of

- energy conservation
- flavor conservation

## GC decay, $E/V = 0.333 \text{ GeV}/\text{fm}^3$ $T = 164 \text{ MeV}$



**GC+MiC decay,  $E/V = 0.333 \text{ GeV}/\text{fm}^3$   $M = 200 \text{ GeV}$**

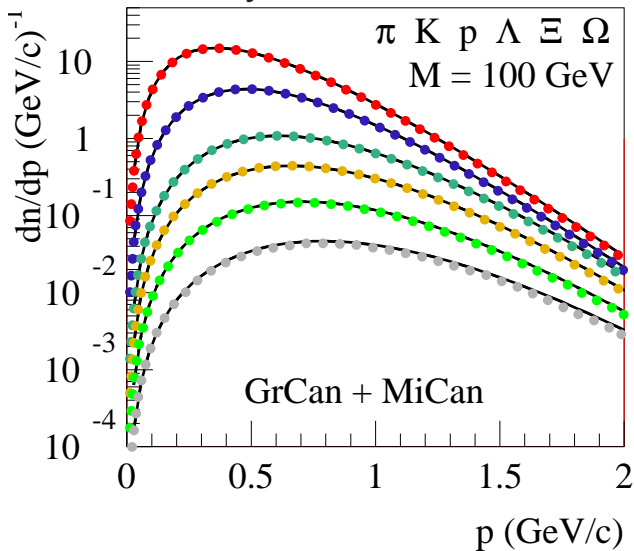


$V = 600 \text{ fm}^3$

$\times \frac{1}{4}$

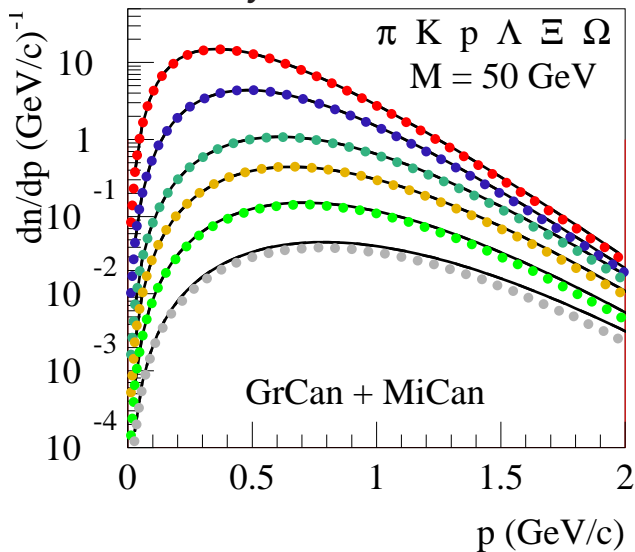
good test for  
Metropolis proposal

**GC+MiC decay,  $E/V = 0.333 \text{ GeV}/\text{fm}^3$   $M = 100 \text{ GeV}$**



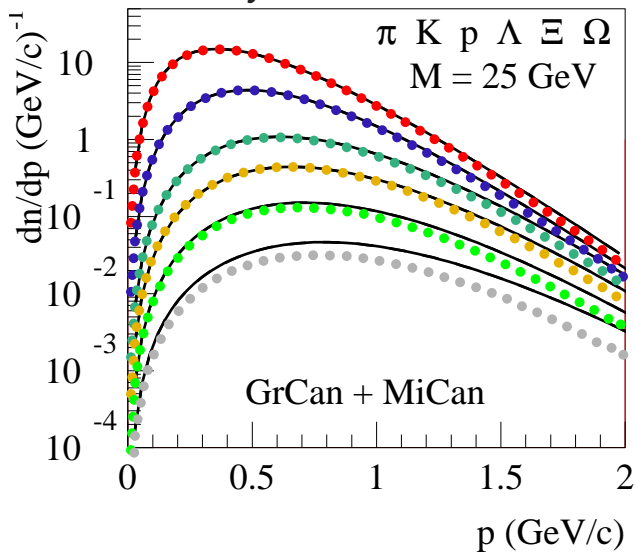
$V = 300 \text{ fm}^3$   
 $\times \frac{1}{2}$

**GC+MiC decay,  $E/V = 0.333 \text{ GeV}/\text{fm}^3$   $M = 50 \text{ GeV}$**



$V = 150 \text{ fm}^3$   
 $\times 1$

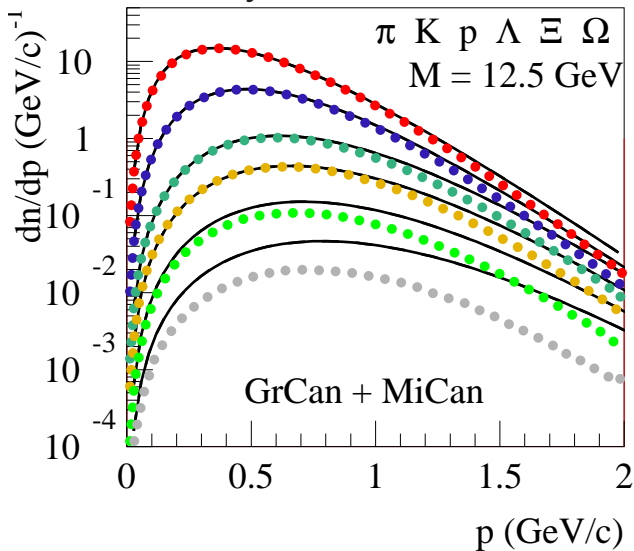
**GC+MiC decay,  $E/V = 0.333 \text{ GeV}/\text{fm}^3$   $M = 25 \text{ GeV}$**



$V = 75 \text{ fm}^3$   
 $\times 2$

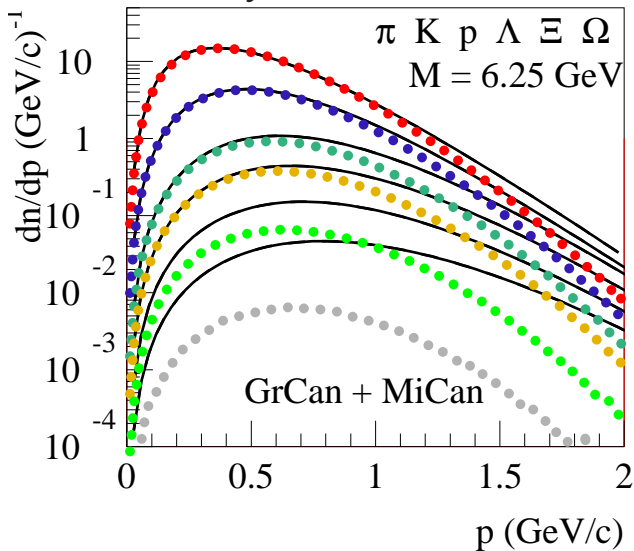


GC+MiC decay,  $E/V = 0.333 \text{ GeV/fm}^3$   $M = 12.5 \text{ GeV}$



$V = 37.5 \text{ fm}^3$   
 $\times 4$

**GC+MiC decay,  $E/V = 0.333 \text{ GeV}/\text{fm}^3$   $M = 6.25 \text{ GeV}$**



$V = 18.75 \text{ fm}^3$

$\times 8$

### 3 Summary

- **New microcanonical hadronization procedure:**
  - **Very efficient, possible for the first time to simulate “big” systems**
  - **Works for “complete” hadron set (PDG)**
  - **Coincides with GC results for big systems**
  - **Unique procedure, for big and small systems**
  
- **Todo: Incorporate flow**