

# Electron EDM and dark sectors

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# Plan

1. Intro: Why EDMs?
2. Recent progress with electron EDMs
3. **Can “dark” sectors** (for this talk = *renormalizable extension of SM, consisting only of neutral particles*) **induce observable levels of electron EDMs?**
4. Models with heavy Dirac neutrinos and Higgs-portal singlets.
5. Sensitivity of electron-EDM-like observables to the hadronic CP-violation.
6. *Conclusions*

Purcell and Ramsey (1949) (“How do we know that strong interactions conserve parity?”  $\longrightarrow |d_n| < 3 \times 10^{-18} \text{ ecm.}$ )

$$H = -\mu \mathbf{B} \cdot \frac{\mathbf{S}}{S} - d \mathbf{E} \cdot \frac{\mathbf{S}}{S}$$

$d \neq 0$  means that both P and T are broken. If CPT holds then CP is broken as well.

CPT is based on locality, Lorentz invariance and spin-statistics = very safe assumption.

*search for EDM = search for CP violation, if CPT holds*

Relativistic generalization

$$H_{\text{T,P-odd}} = -d \mathbf{E} \cdot \frac{\mathbf{S}}{S} \rightarrow \mathcal{L}_{\text{CP-odd}} = -d \frac{i}{2} \bar{\psi} \sigma^{\mu\nu} \gamma_5 \psi F_{\mu\nu},$$

corresponds to dimension five effective operator and naively suggests  $1/M_{\text{new physics}}$  scaling. Due to  $SU(2) \times U(1)$  invariance, however, it scales as  $m_f/M^2$ .

## Why bother with EDMs?

Is the accuracy sufficient to probe TeV scale and beyond?

Typical energy resolution in modern EDM experiments

$$\Delta\text{Energy} \sim 10^{-6}\text{Hz} \sim 10^{-21}\text{eV}$$

translates to limits on EDMs

$$|d| < \frac{\Delta\text{Energy}}{\text{Electric field}} \sim 10^{-25}\text{e} \times \text{cm}$$

Comparing with theoretically inferred scaling,

$$d \sim 10^{-2} \times \frac{1 \text{ MeV}}{\Lambda_{CP}^2},$$

we get **sensitivity to**

$$\Lambda_{CP} \sim 1 \text{ TeV}$$

**Comparable with the LHC reach! EDMs are one of the very few low-energy measurements sensitive to the fundamental particle physics.**

## Current Experimental Limits

”paramagnetic EDM”, Berkeley experiment

$$|d_{\text{Tl}}| < 9 \times 10^{-25} e \text{ cm} \quad \text{Interpreted } |d_e| < 1.6 \times 10^{-27}$$

”diamagnetic EDM”, U of Washington experiment

$$|d_{\text{Hg}}| < 2 \times 10^{-28} e \text{ cm}$$

factor of 7 improvement in 2009!

And another factor of 4 in 2016

$$|d_{\text{Hg}}| < 3 \times 10^{-29} e \text{ cm} \quad 7.4 \times 10^{-30} e \text{ cm}$$

neutron EDM, ILL experiment

$$|d_n| < 3 \times 10^{-26} e \text{ cm}$$

Notice that Thallium EDM is usually quoted as  $d_e < 1.6 \times 10^{-27} e \text{ cm}$

bound. It was modestly improved by YbF results.  $|d_e| < 1.1 \times 10^{-29}$

2013 ThO result by Harvard-Yale collaboration:  $|d_e| < 8.7 \times 10^{-29}$

”Confirmed” using different techniques at JILA,  $|d_e| < 1.3 \times 10^{-28}$  <sup>5</sup>

# “Paramagnetic” EDMs:

- Paramagnetic EDM (EDM carried by electron spin) can be induced not only by a purely leptonic operator

$$d_e \times \frac{-i}{2} \bar{\psi} \sigma_{\mu\nu} \gamma_5 F_{\mu\nu} \psi$$

but by semileptonic operators as well:

$$C_S \times \frac{G_F}{\sqrt{2}} \bar{N} N \bar{\psi} i \gamma_5 \psi$$

- Only a linear combination is limited in any single experiment.  
ThO 2018 ACME result is:

$$|d_e| < 1.1 \times 10^{-29} \text{ e cm} \quad \text{at } C_S = 0$$

$$|C_S^{\text{singlet}}| < 7.3 \times 10^{-10} \quad \text{at } d_e = 0$$

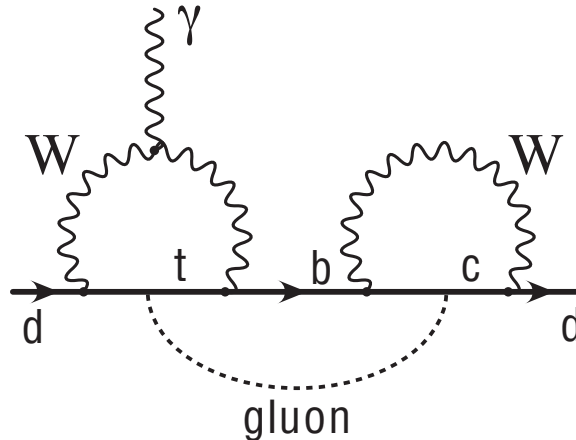
- In the SM the answer is very small, and  $C_S$  is more important, both for  $\Theta_{\text{QCD}}$  and  $\delta_{\text{CKM}}$ .

# Recent progress is very significant

- From  $10^{-27}$  cm to  $10^{-29}$  cm is a factor of 1/100, and sensitivity to  $\Lambda$  is increased by an order of magnitude.
- In terms of probing  $\Lambda$ , progress in electron EDM is similar to the transition from Tevatron to the LHC [but of course requires flavour-diagonal CP-violation]
- E.g. EDMs indirectly probe contact CP-odd Higgs-gamma-gamma coupling with accuracy far greater than usual  $h\text{-}\gamma\text{-}\gamma$ .
- With some luck extends sensitivity to super-partners to a multi-10-TeV/100 TeV regime.
- More progress with  $d_e$  could be anticipated.

(1806.06774 suggests a possibility of going down to  $10^{-34}$  e cm)

# EDMs from CKM



CKM phase generates tiny EDMs:

$$d_d \sim \text{Im}(V_{tb}V_{td}^*V_{cb}V_{cd}^*)\alpha_s m_d G_F^2 m_c^2 \times \text{loop suppression} \\ < 10^{-33} \text{ ecm}$$

Direct quark EDMs identically vanish at 1 and 2 loop levels

(**Shabalin**, 1981). 3-loop EDMs are calculated by **Khriplovich**.

$d_e$  first appears at 4 loops (**Khriplovich, MP**, 1991)  $< 10^{-38}$  e cm



# Two questions for today

- Q1: If in the next round of improvements we see a non-zero electron EDM, does it imply new physics above the EW scale with new charged particles (SUSY, ED, LR models, multi-Higgs models etc)? Or electron EDM can emerge as a consequence of new low-energy UV complete sectors?
- Q2: given rapid progress in ThO measurements, what is the current sensitivity to purely hadronic CP violation?

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(S. Okawa, MP, A. Ritz, *to appear*)

- Q2: given rapid progress in ThO measurements, what is the current sensitivity to purely hadronic CP violation?

(V. Flambaum, MP, A. Ritz, Y. Stadnik, *to appear*)

# Neutral “portals” to the SM

Let us *classify* possible connections between Dark sector and SM

$H^+H$  ( $\lambda S^2 + A S$ ) Higgs-singlet scalar interactions (scalar portal)

$B_{\mu\nu} V_{\mu\nu}$  “Kinetic mixing” with additional U(1)’ group

(becomes a specific example of  $J_\mu^i A_\mu$  extension)

$LHN$  neutrino Yukawa coupling,  $N$  – RH neutrino

$J_\mu^i A_\mu$  requires gauge invariance and anomaly cancellation

I’ll stick to renormalizable UV complete extensions for this talk (i.e. will not touch axions or multi-Higgs theories...)

A systematic search for light particles coupled through vector, Higgs, neutrino and axion portal has intensified in recent years.

(See recent “Physics Beyond Colliders” working group summary, arxiv [1901.09966](https://arxiv.org/abs/1901.09966), for the up-to-date limits)

# Precision frontier: UV physics or IR?

**Typical approach:** we measure an observable (e.g.  $\mu \rightarrow e \gamma$ , EDM, rare meson decays etc), we perform calculation of the same quantity in the SM, take a difference, and whatever is left is interpreted in terms of physics at a TeV, 10 TeV, XXX TeV scales – *all of them being UV scales.*

**More correct approach:** Assume that New Physics consist of UV pieces, Assume IR pieces or both,

$$\mathcal{L}_{\text{NP}} = \mathcal{L}_{\text{UV}} + \mathcal{L}_{\text{IR}}.$$

$$\mathcal{L}_{\text{UV}} = \sum_{d \geq 5} \frac{1}{\Lambda_{\text{UV}}^{d-4}} \mathcal{O}_d. \quad \mathcal{L}_{\text{IR}} = \kappa B^{\mu\nu} V_{\mu\nu} - H^\dagger H (AS + \lambda S^2) - Y_N L H N + \mathcal{L}_{\text{hid}}$$

If result for NP is consistent with 0, we can set constraints on both. If it is non-zero: then *more work is required in deciding IR or UV*

# UV physics or IR: examples of NP that we know

**Neutrino oscillations:** We know that new phenomenon exists, and if interpreted as neutrino masses and mixing, is it coming from deep UV, via e. .g Weinberg's operator

$$\mathcal{L}_{\text{NP}} \propto (HL)(HL)/\Lambda_{\text{UV}} \text{ with } \Lambda_{\text{UV}} \gg \langle H \rangle$$

or it is generated by *new IR field*, such as RH component of Dirac neutrinos? *New dedicated experimental efforts are directed in trying to decide between these possibilities.*

**Dark matter:** 25% of Universe's energy balance is in dark matter: we can set constraints on both. If it is embedded in particle physics, then e.g. neutralinos or axions imply new UV scales.

However, *there are models of DM where NP is completely localized in the UV, and no new scales are necessary.*

New efforts underway both in the UV and IR category.

[Some people at this meeting believe that] *flavour anomalies may already be pointing to NP. It necessarily has states in the UV. Discrepancy in  $(g-2)_\mu$ , on the other hand, if real, does not unambiguously point to either IR or UV.* 13

# Classification of all precision observables

In 2015, **LeDall, MP, Ritz** classified all precision observables from that point of view, IR or UV or both. [*Among EDMs, for example, neutron EDM can be a consequence of UV physics, or IR physics, theta term in QCD*]

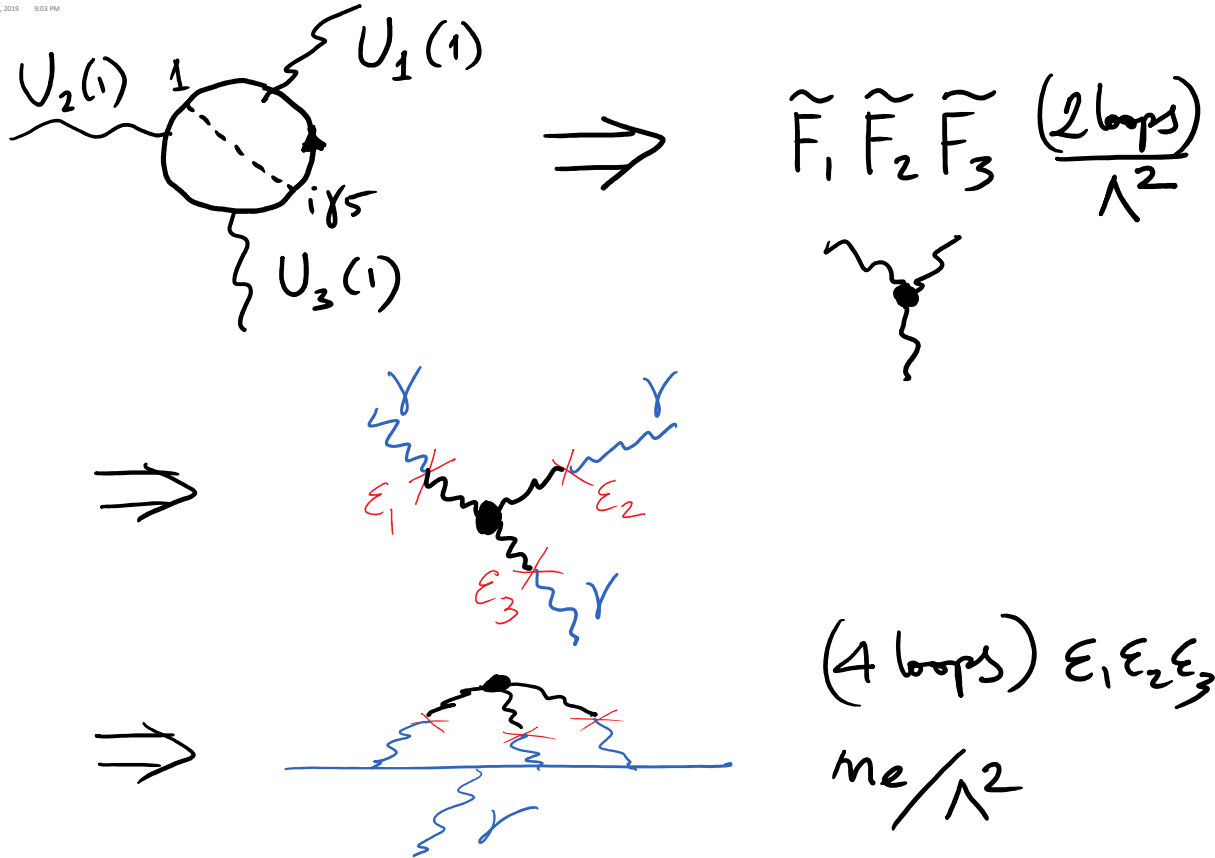
Electron EDM close to current sensitivity – no EW or below EW models were found.

- *This question could be important if one sees a nonzero  $d_e$ . Would it then imply new physics at high scale and “justify” new collider, or IR models can be found?*

$$m_e \ll \Lambda ; m_W \rightarrow \text{infinity}$$

In this limit, only kinetic mixing [and some new gauge symmetries] remain as viable option. You need 3 new U(1) gauge groups and CP violation in the dark sector.

March 26, 2019 9:03 PM

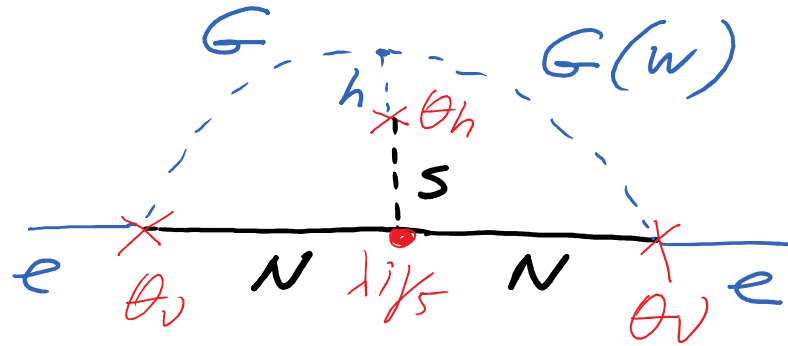


Minimum is 4 loops and 3 kinetic mixings. For  $\Lambda$  below 10 GeV, the result is too small

assuming  
 $m_2, m_1, m_3 \sim \Lambda$   
 $m_1 \neq m_2 \neq m_3$

If it is only 1 extra U(1), then interm. operator is FFFFdual, and  $d_e \sim \epsilon^4 \times 5 \text{ loops}$

# Model with heavy Dirac Neutrinos and Higgs portal singlets



Edm of electron  
is generated  
at 2 loops  
and 3 mixing  
angles

$$HLN; (H^\dagger H)S; S\bar{N}i\gamma_5 N \text{ portals}$$

CP violation in the dark sector

Mixing angles  $\theta_h$  and  $\theta_v$  do not have to be small if S, N are EW scale<sup>16</sup>



# Two-loop calculation

We have performed calculation in  $\alpha_W \rightarrow 0$  limit, where only Goldstone loops matter (and for  $m_N > m_W$ ).

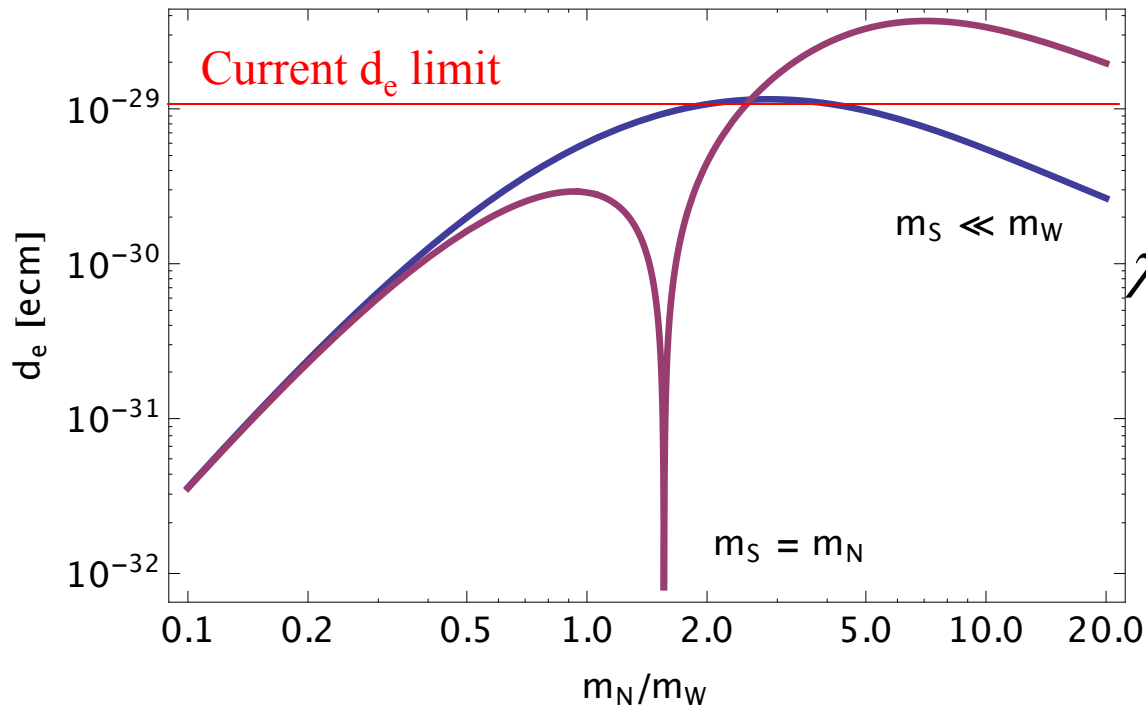
**Answer =  $d_e^{\text{scale}} \times 2\text{loop function}$**

$$d_e^{\text{scale}} = e\lambda_N\theta_\nu^2 \times \frac{2}{(16\pi^2)^2} \frac{Am_em_N}{v^2m_h^2} = e\lambda_N\theta_\nu^2\theta_h \times \frac{2}{(16\pi^2)^2} \frac{m_em_N}{v^3}$$

$$d_e^{\text{scale}} = 4.4 \times 10^{-29} e \text{ cm} \times \frac{\lambda_N\theta_\nu^2\theta_h}{10^{-2}} \times \frac{m_N}{m_W}.$$

*2-loop function* is calculated numerically, and analytic expressions are obtained in various hierarchical limits (e.g.  $m_W \ll m_S \ll m_N$ )

# Two-loop calculation results



$$\lambda_N \theta_h \theta_\nu^2 = 10^{-2}.$$

$m_N \gg m_W$  has a chance of creating electron EDM above  $10^{-29}$  cm.

- **Answer to Q1:** two-loop  $d_e$  via Higgs and  $\nu$  portals. Example of model inducing  $d_e$  using only neutral BSM particles around EW scale. For the *first example of this type*, see [A. Abada, T. Toma, 2015](#), in a model with two heavy RH neutrinos.
- **Best chance at probing EW scale N's – high-luminosity LHC**
- A good chance for EW baryogenesis – N carry lepton number

# Two-photon exchange induced $C_S$

- Th used by ACME collaboration is a spin-less nucleus.
- ThO is mostly sensitive to CP violation in the lepton sector. If CP is broken in the strong interaction sector, *two photon exchange* can communicate it to the electron shells.
- Cutting across the two photons, the intermediate result can be phrased via CP-odd nuclear polarizability,  $\mathbf{EB} \delta(\mathbf{r})$ , where E and B are created by an electron.
- Good scale separation is possible,  $m_p \gg p_F$ ,  $m_\pi \gg m_e \sim Z\alpha m_e$
- Nuclear uncertainties could be under control if the result is driven by “bulk” [as opposed to valence] nucleons.

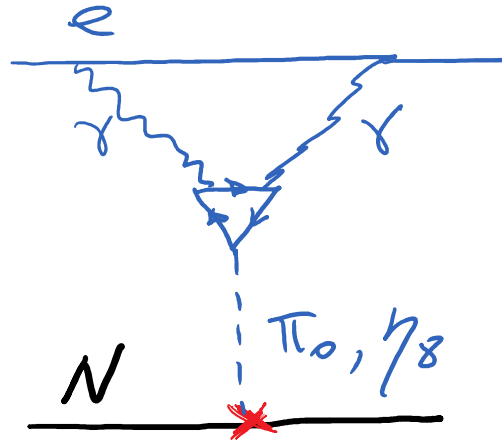
# Two questions for today

$(\bar{e} \gamma_5 e) \bar{N} N$  operators

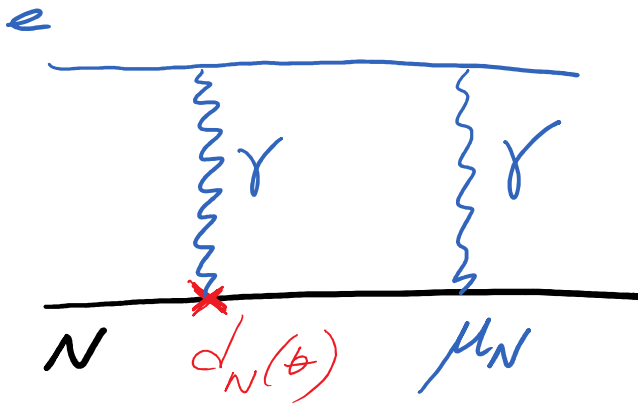
$m_q$  counting:

$$\theta m_q / m_\pi^2$$

$$\sim O(m_q^0)$$



Almost complete  
cancellation of  
π<sub>0</sub> and η<sub>8</sub>  
contributions



+ cross diagrams

$m_q$  counting:

$$\sim O(m_q \log(m_q))$$

# Unexpected complication:

- T-channel pion exchange gives

$$\begin{aligned}\mathcal{L} &= \theta \times \frac{1}{m_\pi^2} \times 0.017 \times 3.5 \times 10^{-7} (\bar{e}i\gamma_5 e)(\bar{n}n - \bar{p}p) \\ &= (\bar{e}i\gamma_5 e)(\bar{n}n - \bar{p}p) \times \frac{3.2 \times 10^{-13}\theta}{\text{MeV}^2}.\end{aligned}$$

implying  $|\theta| < 8.4 \times 10^{-8}$  sensitivity. However, adding exchange of  $\eta_8$ ,

$$1 \rightarrow 1 - \frac{1}{3} \frac{f_\pi^2 m_\pi^2}{f_\eta^2 m_\eta^2} \times \frac{m_d - m_u}{m_d + m_u} \times \frac{A \times \sigma_N}{\frac{m_d - m_u}{2} \langle p | \bar{u}u - \bar{d}d | p \rangle \times (N - Z)}$$

$$1 \rightarrow 1 - 0.88 \simeq 0.12.$$

The effect can completely cancel within error bars on nucleon sigma term  $\sigma_N$ .

# Photon box diagrams:

- Diagrams are IR divergent but regularized by Fermi momentum in the Fermi gas picture of a nucleus (intermediate N is above Fermi surface).

$$\mathcal{L} = \bar{e}i\gamma_5 e \bar{N}N \times \frac{2m_e \times 4\alpha \times \bar{d}\mu \times 6.2}{\pi p_F} = \bar{e}i\gamma_5 e \bar{N}N \times 2.4 \times 10^{-4} \times \bar{d}\mu$$

$$\bar{d}\mu = \frac{Z}{A}\mu_p d_p + \frac{A-Z}{A}\mu_n d_n = \frac{e}{2m_p} \times (1.08d_p - 1.16d_n)$$

- Nucleon EDM (theta) is very much a triplet,  $d_p \simeq -d_n \simeq 1.6 \times 10^{-3} \text{efm}\theta$

Constructive interference. [Preliminary] answer:

$$C_S^s \sim 0.025 \times \theta \quad \text{or} \quad |\theta| < 3 \times 10^{-8}.$$

Bonus: limit on protons EDM

$$C_S^s \sim \frac{\sqrt{2}}{G_F} \times 2.4 \times 10^{-4} \times \frac{4\pi\alpha}{2m_p} \times 1.08 \times \frac{d_p}{e} < 7.3 \times 10^{-10} \implies d_p < 10^{-23} \text{ecm}$$

# Conclusions

- Electron EDM/ semileptonic  $C_S$  operator are limited rather strongly by the recent results of ACME collaboration.
- New levels of sensitivity,  $10^{-29}$  e cm, is probing not only UV models of new physics with new SM-charged states, but can also probe “dark sectors” where new physics is neutral [and UV complete].
- Models with heavy [singlet] Dirac Neutrinos and [singlet] scalar mixed with the Higgs can induce  $d_e$  at two loops and three portal insertions. The answer was demonstrated to be as big as  $10^{-29}$  e cm.
- Electron EDM experiments are sensitive to the hadronic CP violation via a two-photon exchange diagrams.
- Simple evaluations of sensitivity results in  $|\theta| < 3 \times 10^{-8}$ . Further progress by O(100) for  $d_e$  type of experiments will bring the sensitivity to hadronic CP violation on par with current  $d_n$  limits.