

Role of the triangle singularity
in $\Lambda(1405)$ production in
 $\pi^-p \rightarrow K^0\pi\Sigma$ and $pp \rightarrow pK^+\pi\Sigma$ reactions

M. Bayar, R. Pavao, S. Sakai, and E. Oset

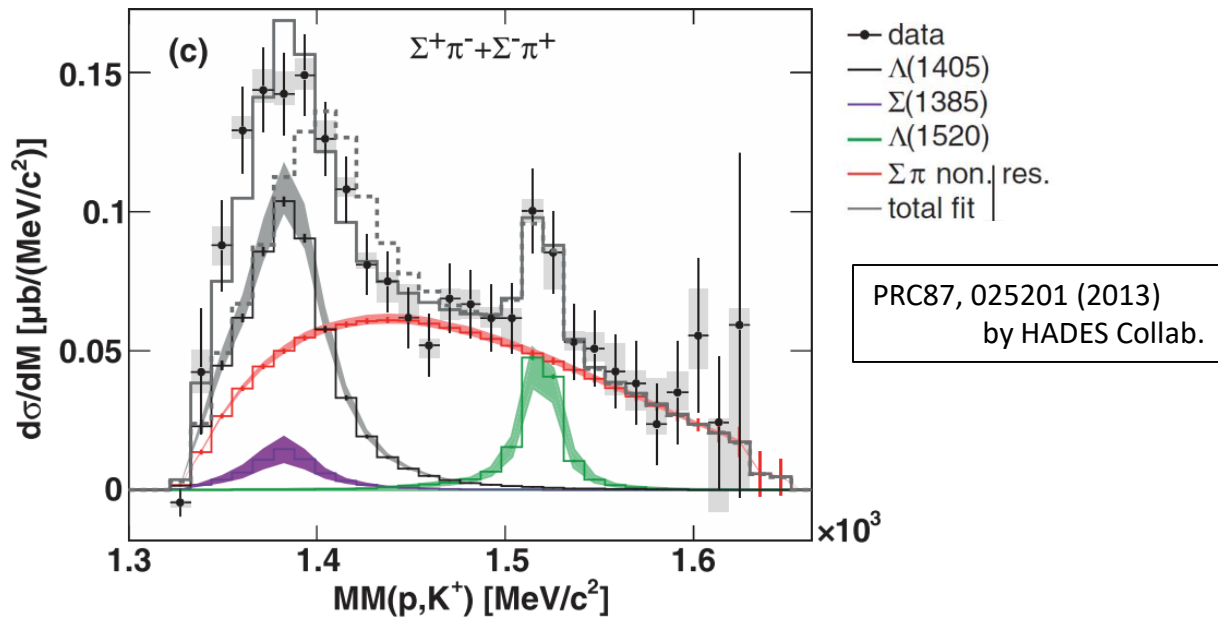
Phys. Rev. C 97, 035203 (2018)

Shuntaro Sakai [Institute of Theoretical Physics, CAS (China)]

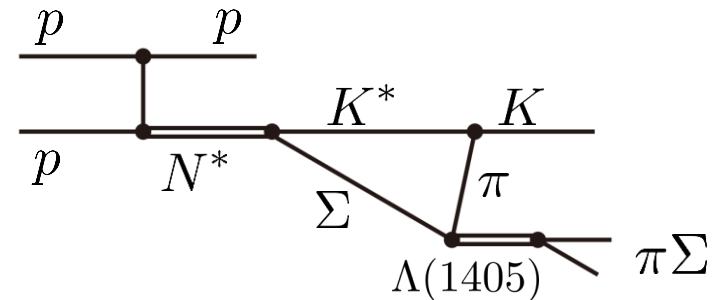
HADES Experiment

$$pp \rightarrow \Sigma^\pm + \pi^\mp + K^+ + p$$

with 3.5 GeV kinetic energy beam

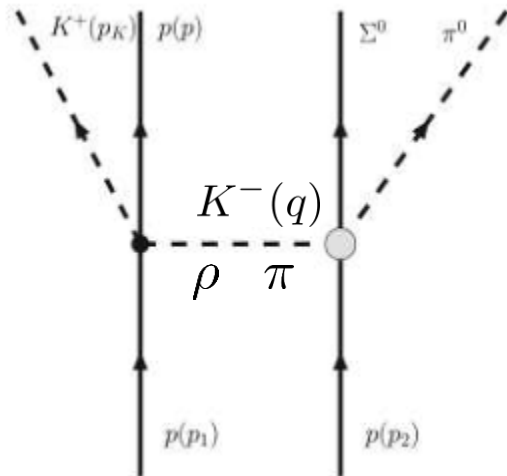


Possible description with triangle mechanism



Theoretical study of $pp \rightarrow pK^+\pi\Sigma$

Geng-Oset (2007)



■ π and ρ exchange

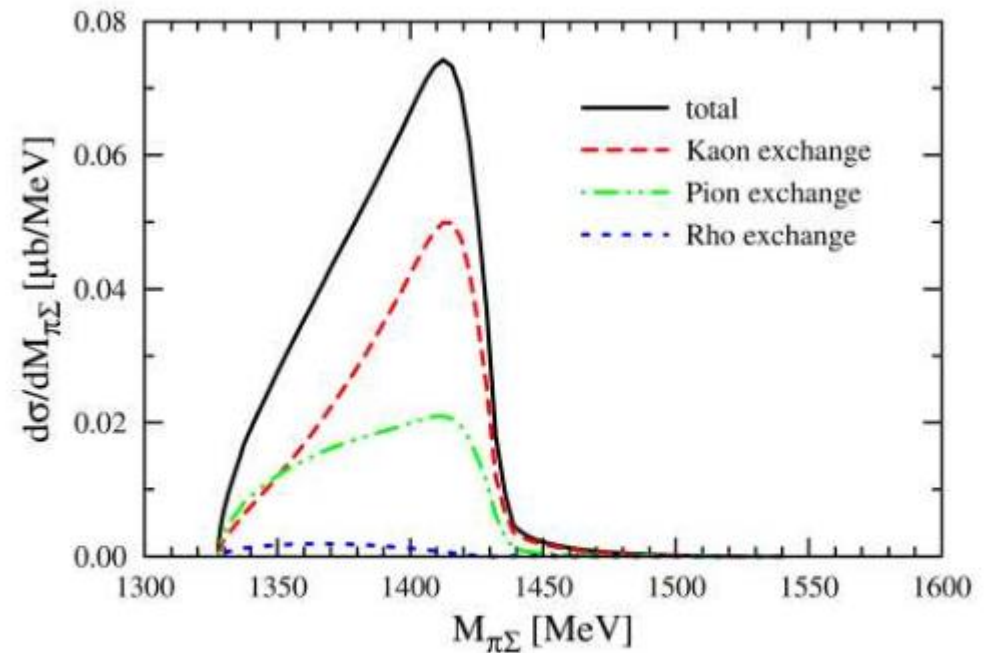
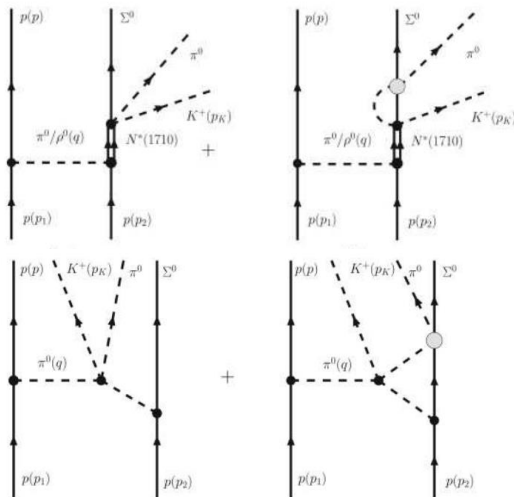
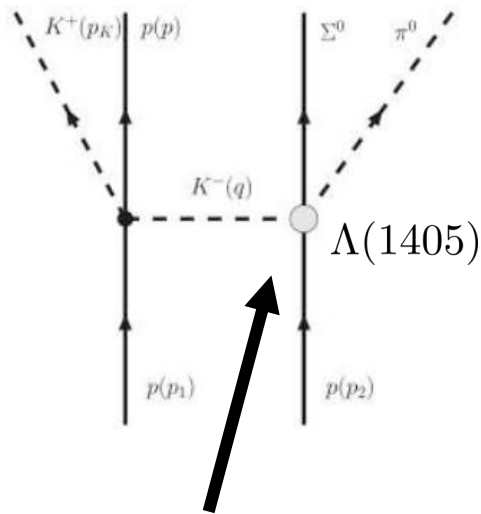


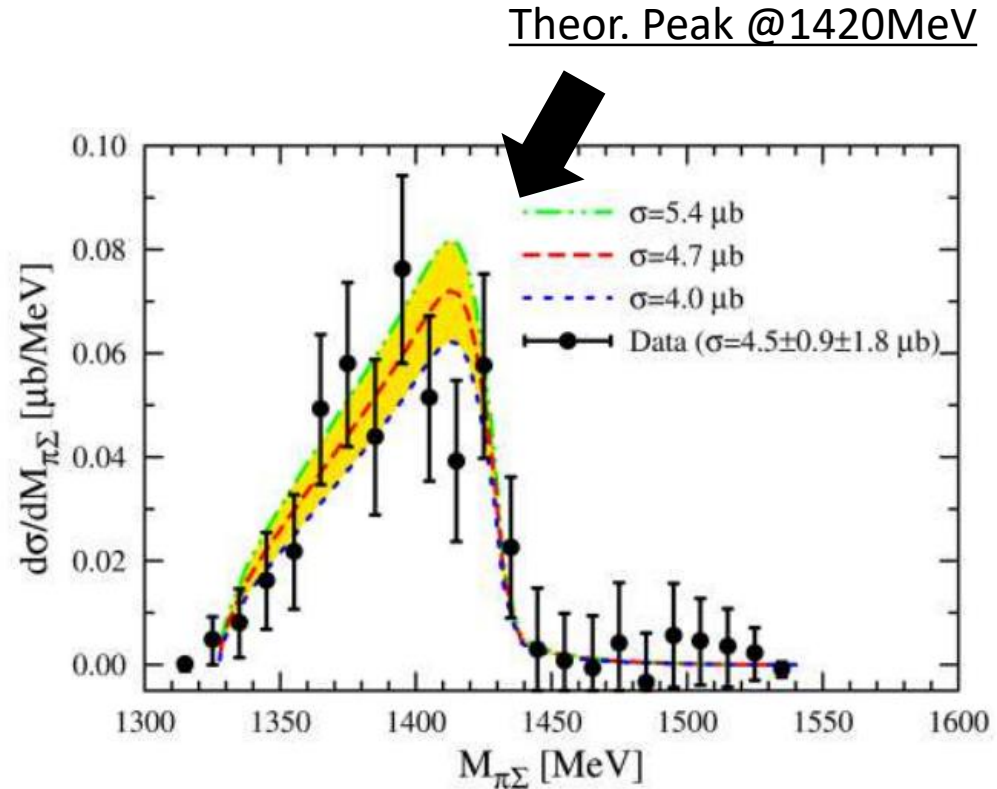
Fig. 6. The contribution of the three different mechanisms with $G_{\rho NN^*} = -0.62$.

Theoretical study of $pp \rightarrow pK^+ \pi \Sigma$

Geng-Oset (2007)



$\Lambda(1405)$
in $K^- p \rightarrow \pi^0 \Sigma^0$

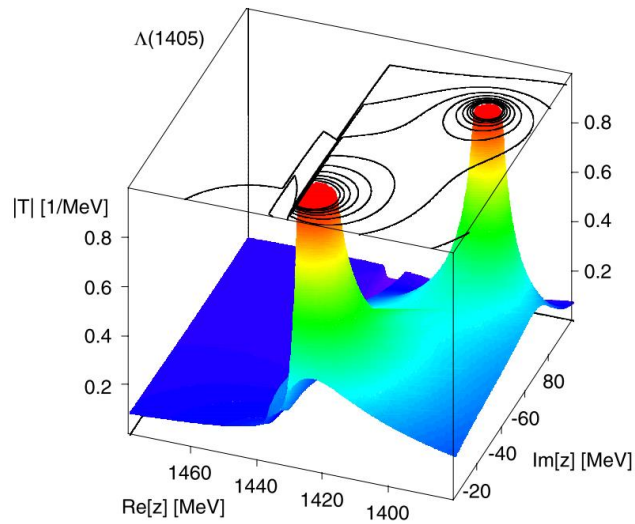


data: COSY-Julich [PLB660, 167 (2008)]

✓ Pole structure of $\Lambda(1405)$

[based on chiral SU(3)]

Oset-Ramos (1998), Oller-Meissner (2001),...



Hyodo-Jido (2012)

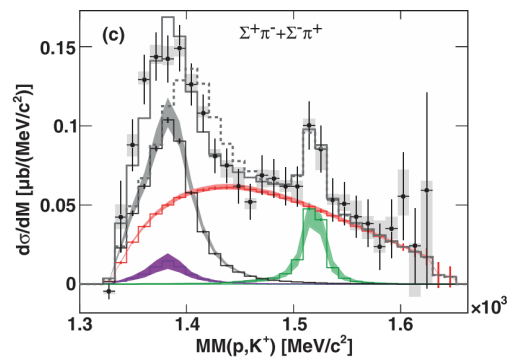
$\Lambda(1405)$

1390 – 66i

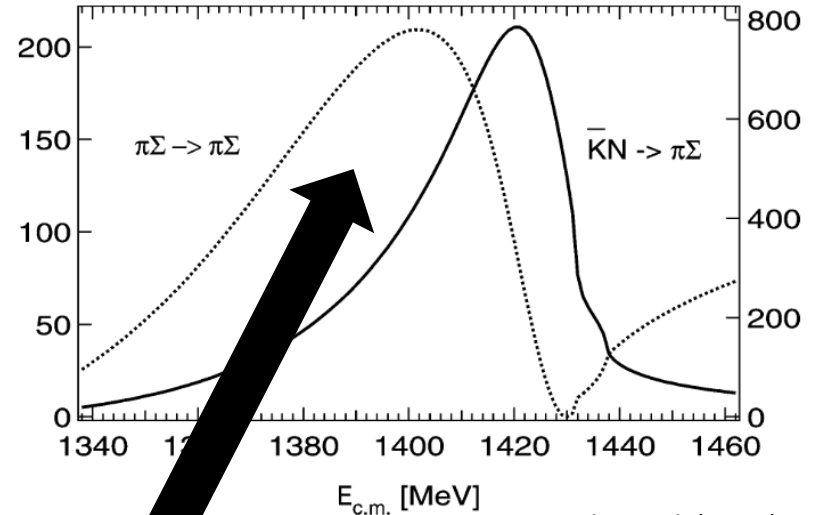
1426 – 16i

$\rightarrow \pi\Sigma$

$\rightarrow \bar{K}N$

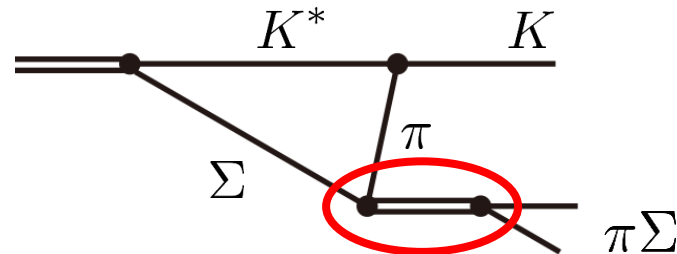


✓ Mass distribution of $\Lambda(1405) \rightarrow \pi\Sigma$



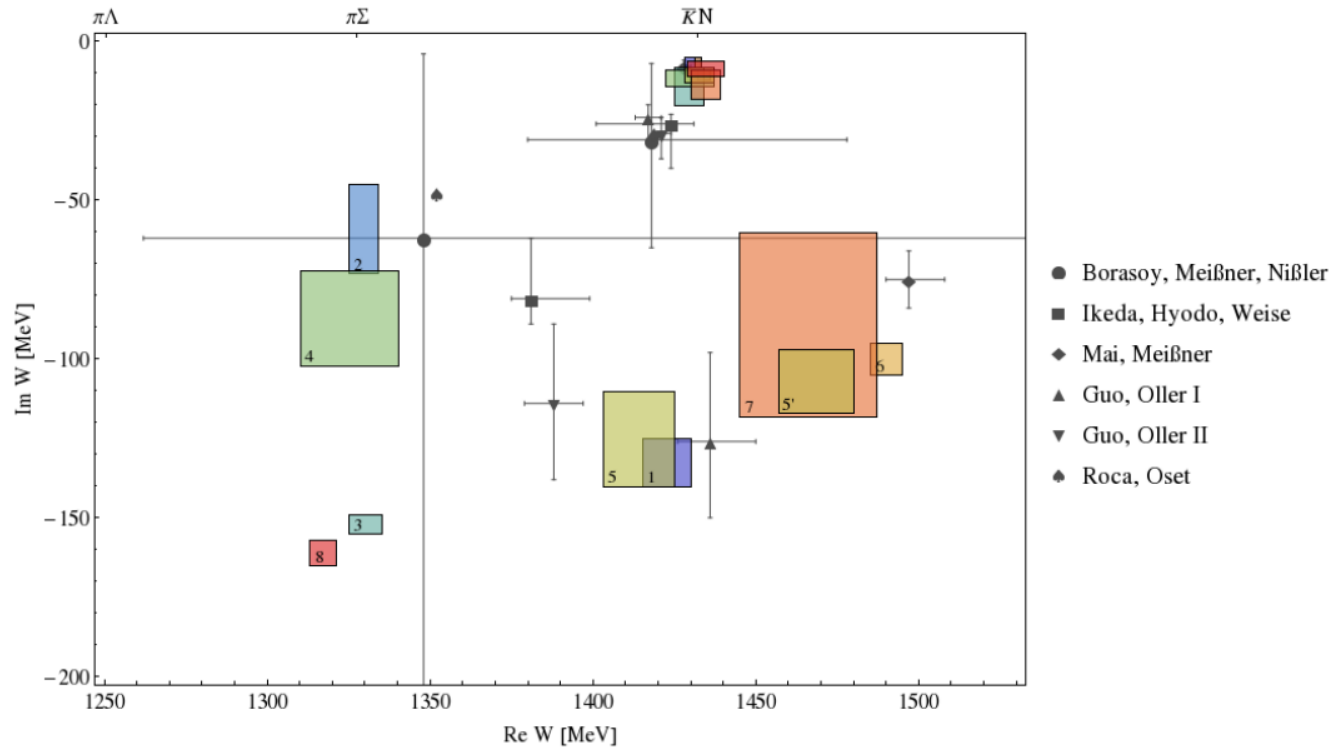
Jido et al. (2003)

**Large contribution of lower pole
for $\pi\Sigma \rightarrow \pi\Sigma$ peak**



$\pi\Sigma \rightarrow \Lambda(1405) \rightarrow \pi\Sigma$

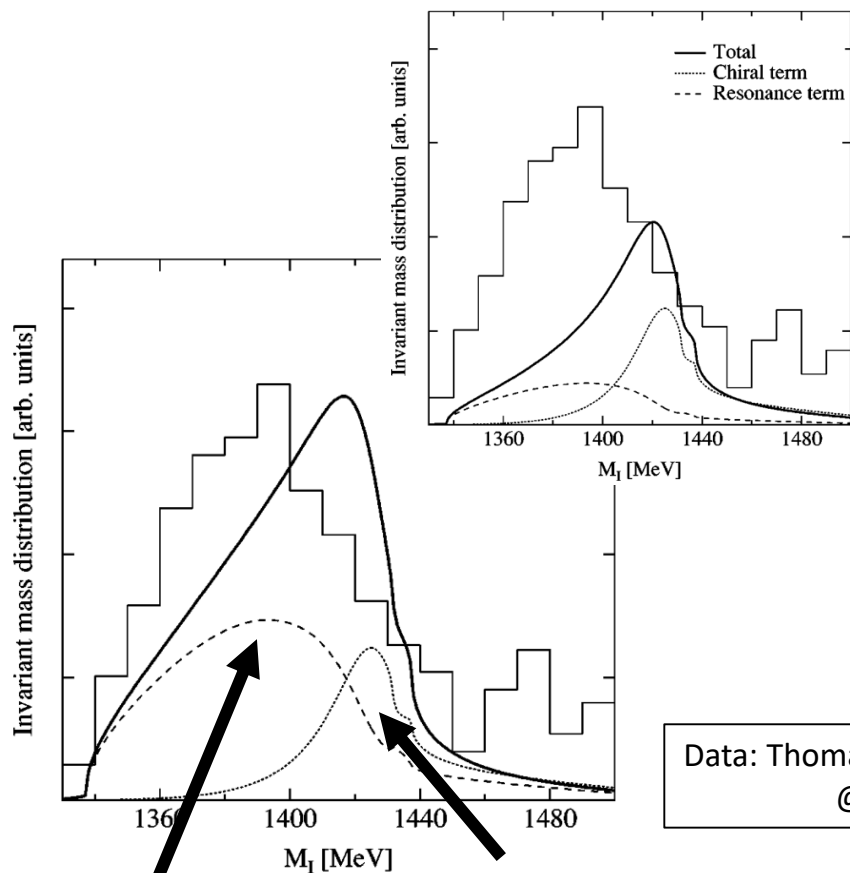
Analyses



Mai-Meissner (2015)

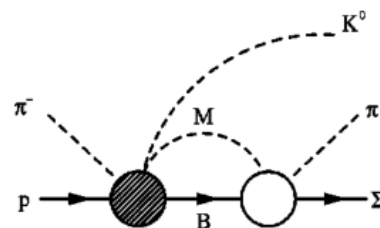
Theoretical study of $\pi^- p \rightarrow K^0 \pi \Sigma$

Hyodo *et al.* (2003)

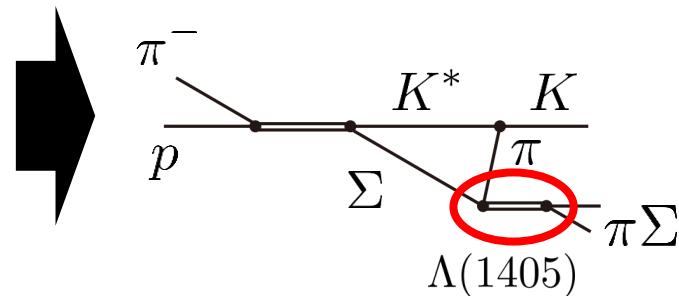
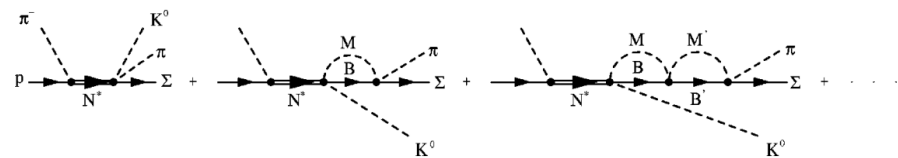


Data: Thomas *et al.* (1973)
 @ $\sqrt{s} \sim 2020$ MeV

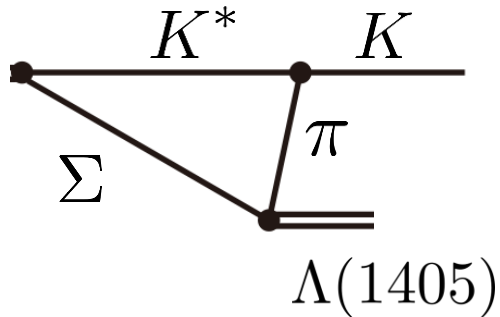
from $\bar{K} N \rightarrow \pi \Sigma$
 from intermediate $N^*(1710)$
 [induce $\pi \Sigma \rightarrow \pi \Sigma$]



with



Triangle mechanism from $K^*\Sigma\pi$ loop



: singularity@2140 MeV in $M_{K\Lambda(1405)}$

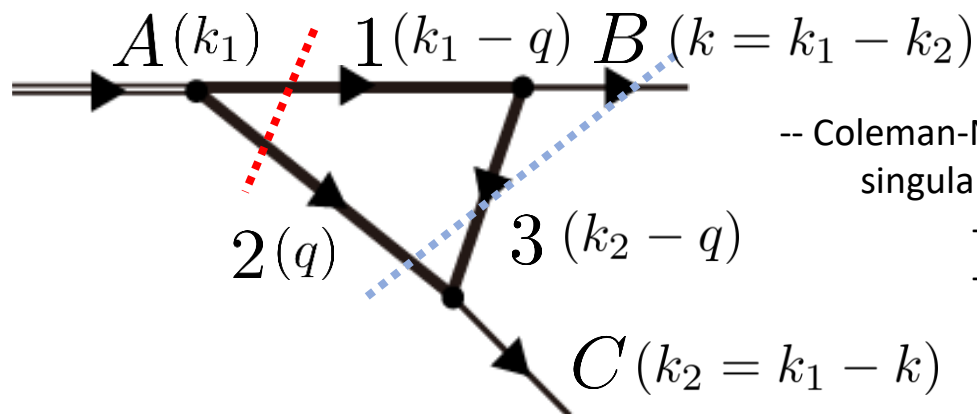


Large contribution in the process

@ $M_{K\Lambda(1405)}=2140\text{MeV}$

Triangle singularity

Landau (1959),....



-- Coleman-Norton theorem (1965):

singularity \sim process occurs classically

- internal particle: onshell

- energy-mom. of each vertex: conserved

$$i \int \frac{d^4 q}{(2\pi)^4} \frac{1}{[q^2 - m_2^2 + i\epsilon][(k_1 + q)^2 - m_1^2 + i\epsilon][(k_2 + q)^2 - m_3^2 + i\epsilon]}$$

$$\sim \int dq d\cos\theta \frac{1}{[E_A - \omega_1(q) - \omega_2(q) + i\epsilon][E_{23} - \omega_2(q) - \omega_3(\vec{q} + \vec{k}) + i\epsilon]}$$

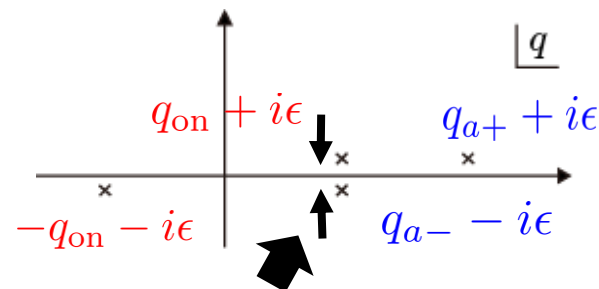
$$q = \pm q_{\text{on}} \pm i\epsilon$$

$$q = q_{a\pm} \pm i\epsilon$$

$$\text{with } \cos\theta = -1$$

$$[\omega_i(\vec{q}) \sqrt{m_i^2 + |\vec{q}|^2}]$$

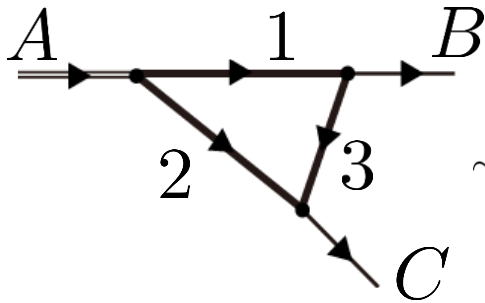
$$E_{23} = E_A - E_B$$



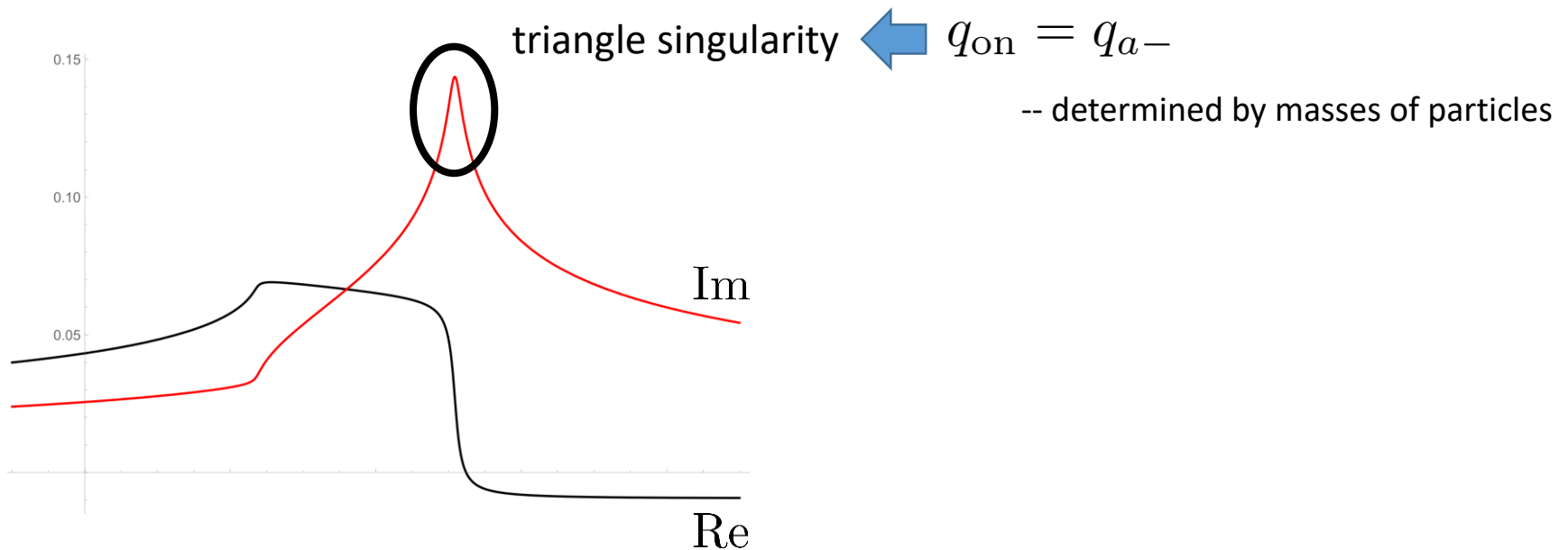
Pinching

$$\text{@ } q_{\text{on}} = q_{a-}$$

Singularity



$$\sim i \int \frac{d^4 q}{(2\pi)^4} \frac{1}{[q^2 - m_1^2 + i\epsilon][(k_1 + q)^2 - m_2^2 + i\epsilon][(k_2 + q)^2 - m_3^2 + i\epsilon]}$$



-- finite width of particles: smear the singularity

-- momentum or angle dependence of vertices: modify the structure

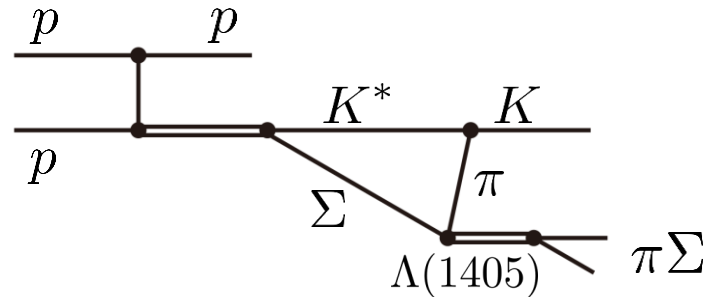
➡ **Necessity of phenomenological study of triangle mechanism**

Role of triangle mechanism

- Enhancement of Isospin-violating $\pi^0 f_0(980)$ production in $\eta(1405/1475)$ decay
 - >> Wu *et al.* (2012,2013), Aceti *et al.* (2012),...
- “a1(1420)” peak in p-wave $\pi^+ f_0(980)$
 - >> Mikhasenko-Ketzer-Sarantsev (2015), Aceti-Dai-Oset (2016),...
- $\gamma p \rightarrow K\Lambda(1405)$ process
 - >> Wang *et al.* (2017)
- Near-threshold $\pi^0 N(1535)$ photoproduction
 - >> Debastiani-Sakai-Oset (2017)
- P_c peak in $J/\psi N$ distribution by LHCb
 - >> Mikhasenko (2015), Guo *et al.* (2015),...
- B_c or B decays
 - >> Liu-Meissner (2017), Nakamura (2019),...
- and so on...

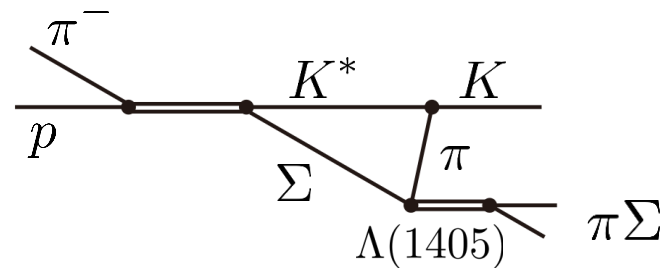
This work

- $pp \rightarrow pK^+\pi\Sigma$



No exp. data ~ 2.14 GeV (πN ene.)

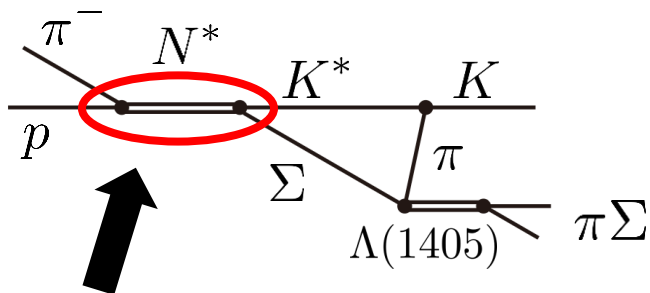
- $\pi^- p \rightarrow K^0 \pi \Sigma$



Formalism

- $\pi^- p \rightarrow K^0 \pi \Sigma$

$\checkmark \pi N \rightarrow K^* \Sigma$

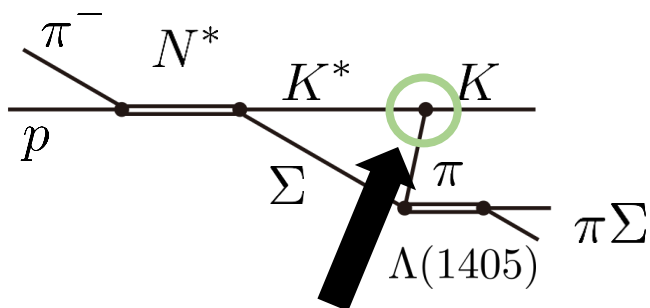


$K^* \Sigma$: dominantly produced by $N^*(2030)$ in s wave

- large coupling: Oset-Ramos(2010,2013) ← from $\gamma N \rightarrow K \Sigma$
- used in the $\gamma p \rightarrow K \Lambda(1405)$ study (Wang *et al.* (2017))

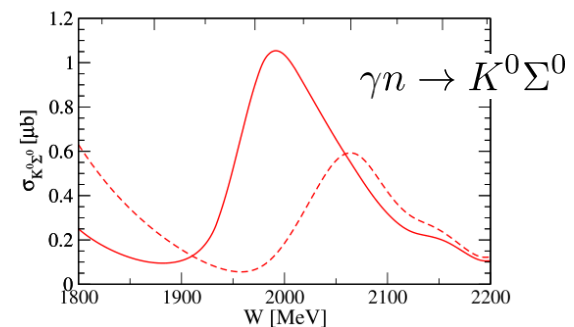
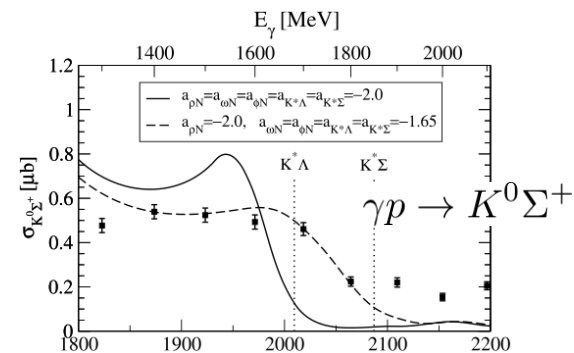
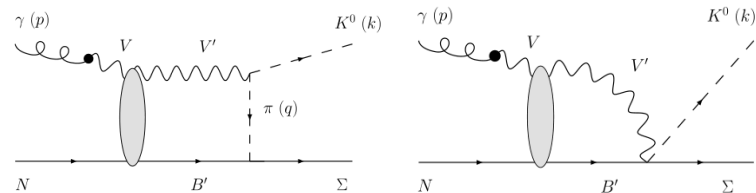
⊗ No information of $\pi N - N^*$ coupling:
estimation with typical value of $N^* \rightarrow \pi N$ width ($\sim 70 \text{ MeV}$)

$\checkmark K^* \rightarrow \pi K$

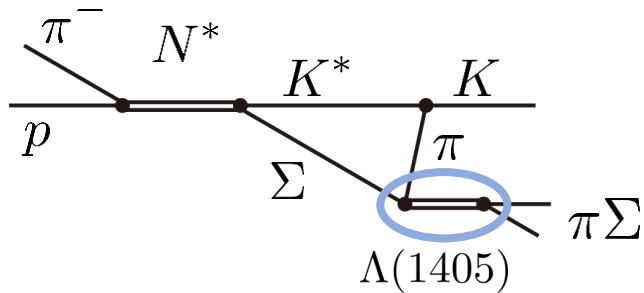


$K^* \pi K$: effective Lagrangian

$$\mathcal{L} = -ig \text{tr} ([P, \partial^\mu P] V_\mu)$$



$$\checkmark \pi\Sigma \rightarrow \Lambda(1405) \rightarrow \pi\Sigma$$



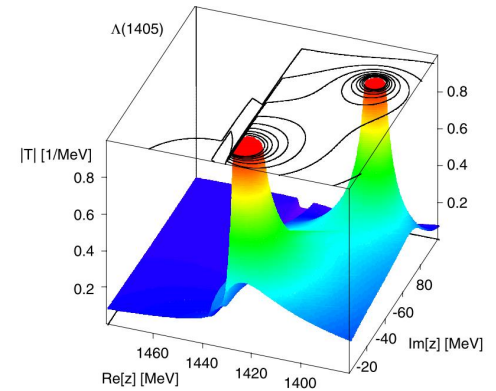
-- Chiral unitary approach

$$t_{ij} = ([1 - vg]^{-1}v)_{ij} \quad \leftarrow \text{elastic unitarity}$$

$$(i, j : \pi\Sigma, \bar{K}N, \eta\Lambda, \pi\Lambda, \eta\Sigma, K\Xi)$$

v : interaction kernel

\leftarrow from LO SU(3) chiral Lagrangian



Two poles:

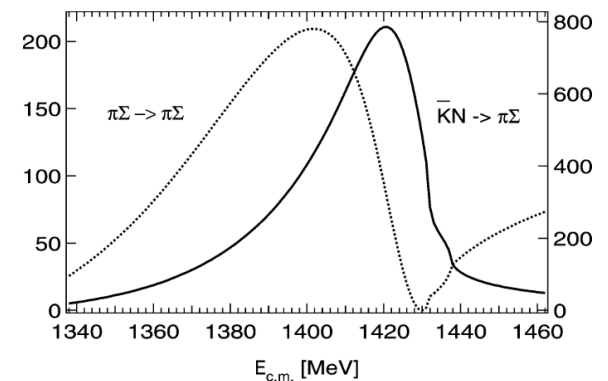
Higher pole(~ 1420 MeV):

large coupling to $\bar{K}N$

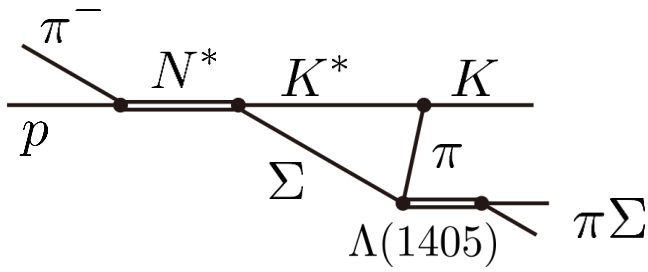
Lower pole(~ 1390 MeV):

large coupling to $\pi\Sigma$

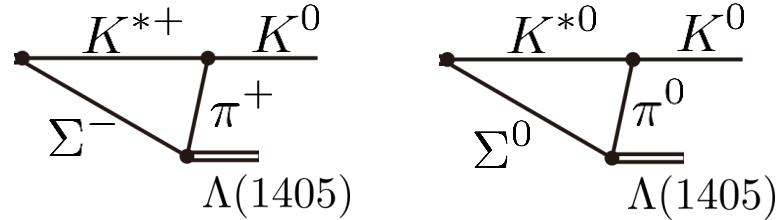
Hyodo-Jido (2012)



Jido et al. (2003)



$$t = -\frac{2M_\Sigma}{3\sqrt{3}} \frac{g_{\pi N, N^*}^{I=1/2} g_{N^*, K^* \Sigma}^{I=1/2}}{\sqrt{s} - M_{N^*} + i\Gamma_{N^*}/2} \vec{\sigma} \cdot \vec{k} t_T \left(t_{\Sigma^- \pi^+, \Sigma \pi} + \frac{1}{2} t_{\Sigma^0 \pi^0, \Sigma \pi} \right)$$



t_T : triangle loop

$$t_T = i \int \frac{d^4 q}{(2\pi)^4} \left(2 + \frac{\vec{k} \cdot \vec{q}}{k^2} \right) \frac{1}{[q^2 - m_\Sigma^2 + i\epsilon][(P - q)^2 - m_{K^*}^2 + i\epsilon][(P - q - k)^2 - m_\pi^2 + i\epsilon]}$$

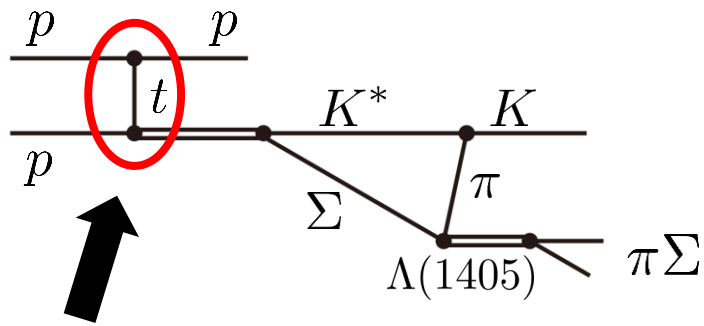
$$\sim \int \frac{d^3 q}{(2\pi)^3} \frac{\left(2 + \frac{\vec{k} \cdot \vec{q}}{k^2} \right)}{8\omega\omega'\omega^*} \frac{2P^0\omega + 2k^0\omega' - 2[\omega + \omega'][\omega + \omega' + \omega^*]}{[k^0 - \omega' - \omega^* + i\epsilon][P^0 + \omega + \omega' - k^0][P^0 - \omega - \omega' - k^0 + i\epsilon][P^0 - \omega^* - \omega + i\epsilon]}$$

$$\frac{d\sigma_{K^0 \pi \Sigma}}{dm_{\pi \Sigma}} = \frac{M_p M_\Sigma |\vec{k}| |\vec{p}_\pi|}{2(2\pi)^3 s |\vec{p}_\pi|} \sum \sum |t_{\pi^- p \rightarrow K^0 \pi \Sigma}|^2$$

Formalism

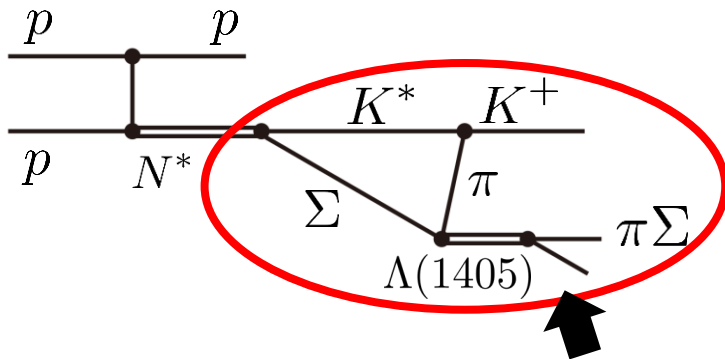
- $pp \rightarrow pK^+ \pi \Sigma$

$$\checkmark \underline{pp \rightarrow pN^*}$$



$C : t(\cos\theta)$ dependence

← averaged

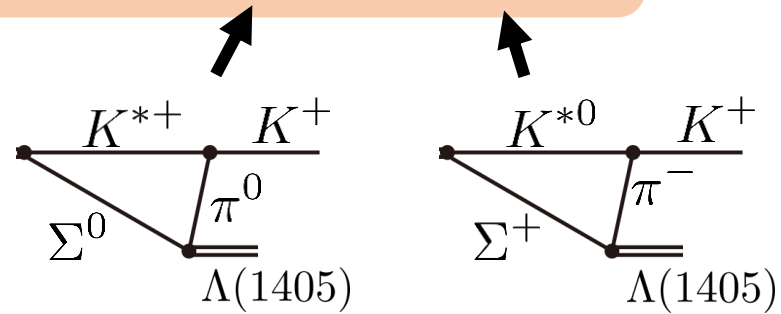


Same as $\pi N \rightarrow K \Lambda(1405)$

$K^*\Sigma$: large coupling to N^* : Oset-Ramos(2010)

$pp \rightarrow K\Lambda(1405)$ amplitude

$$t = -\frac{2M_\Sigma}{3\sqrt{3}} \frac{Cgg_{N^*,K^*\Sigma}^{I=1/2}}{M_{\text{inv}} - M_{N^*} + i\Gamma_{N^*}/2} \vec{\sigma} \cdot \vec{k} t_T \left(t_{\Sigma^-\pi^+,\Sigma\pi} + \frac{1}{2} t_{\Sigma^0\pi^0,\Sigma\pi} \right)$$



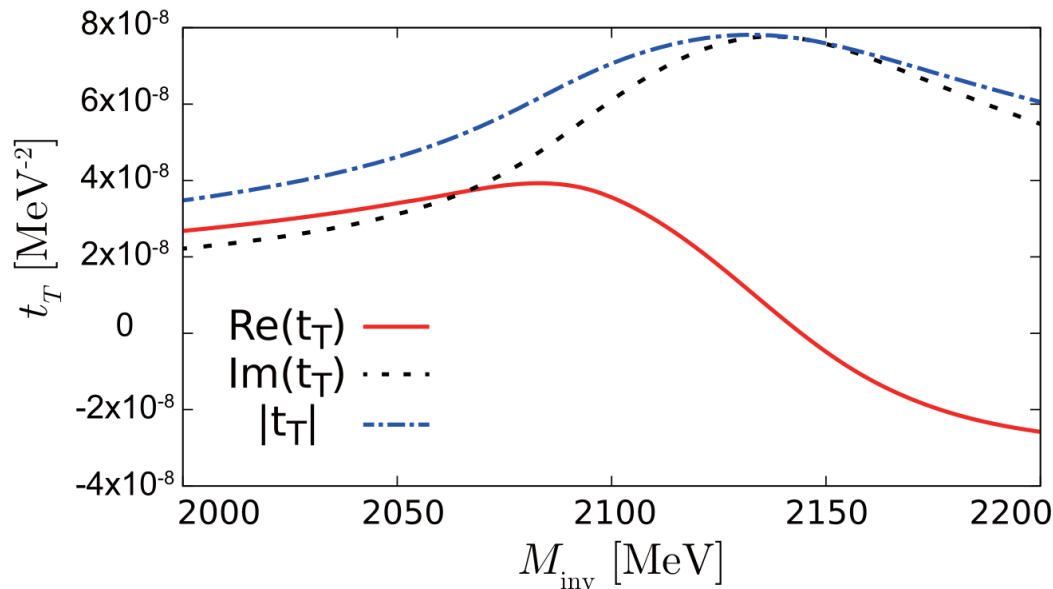
$$\frac{d^2\sigma}{dM_{K\Lambda^*} dm_{\pi\Sigma}} = \frac{Cp_1p_a}{|M_{\text{inv}} - M_{N^*} + i\Gamma_{N^*}/2|^2} \times |\tilde{p}_\pi k^3 |t_T|^2 |t_{\Sigma^+\pi^-, \Sigma\pi} + \frac{1}{2} t_{\Sigma^0\pi^0, \Sigma\pi}|^2$$

○ pp energy: coincide with experiment

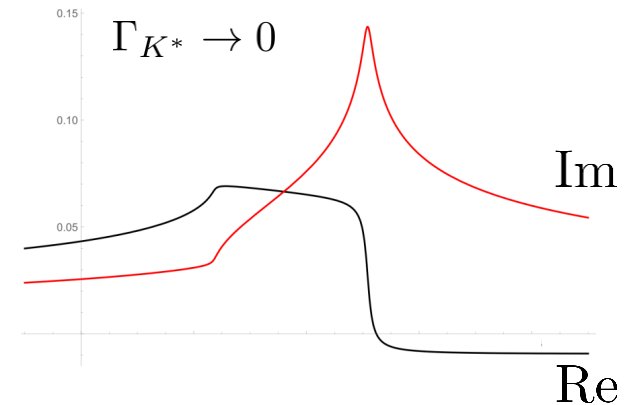
Results

Triangle amplitude

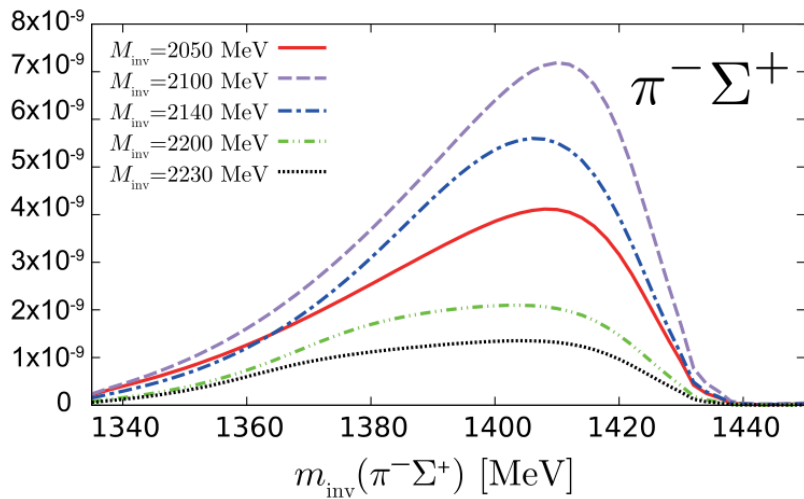
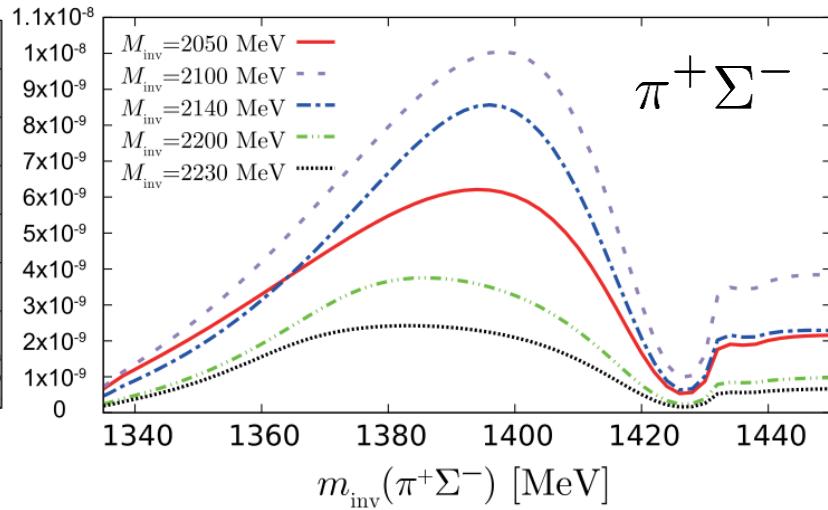
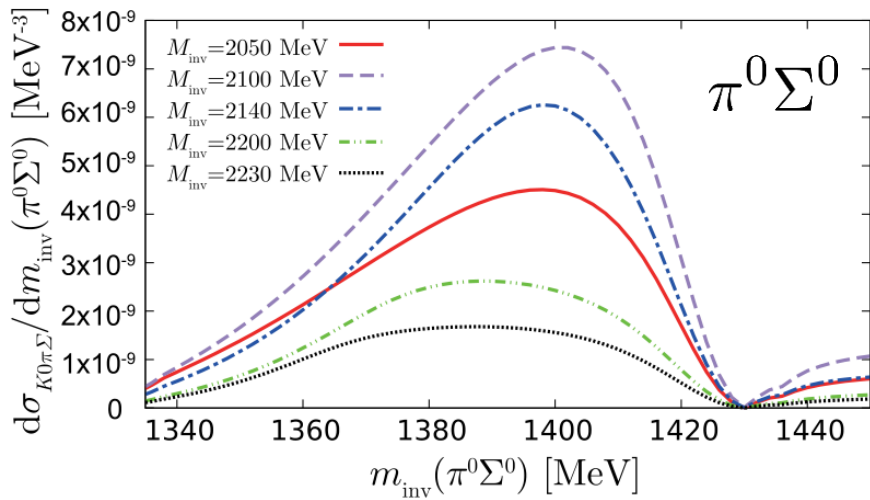
$$t_T = \int \frac{d^3q}{(2\pi)^3} \frac{\left(2 + \frac{\vec{k} \cdot \vec{q}}{k^2}\right)}{8\omega\omega'\omega^*} \frac{2P^0\omega + 2k^0\omega' - 2[\omega + \omega'][\omega + \omega' + \omega^*]}{[k^0 - \omega' - \omega^* + i\epsilon][P^0 + \omega + \omega' - k^0][P^0 - \omega - \omega' - k^0 + i\epsilon][P^0 - \omega^* - \omega + i\epsilon]}$$



Peak @2140MeV
from triangle singularity



$$\frac{d\sigma_{\pi^- p \rightarrow K^0 \pi \Sigma}}{dm_{\pi \Sigma}}$$



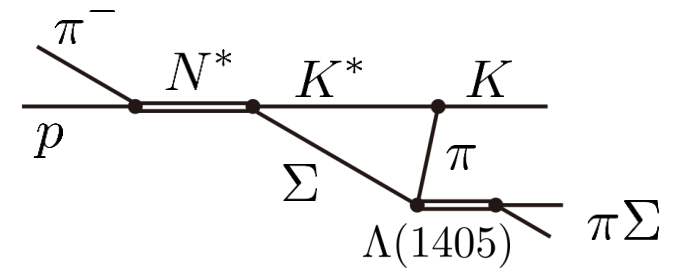
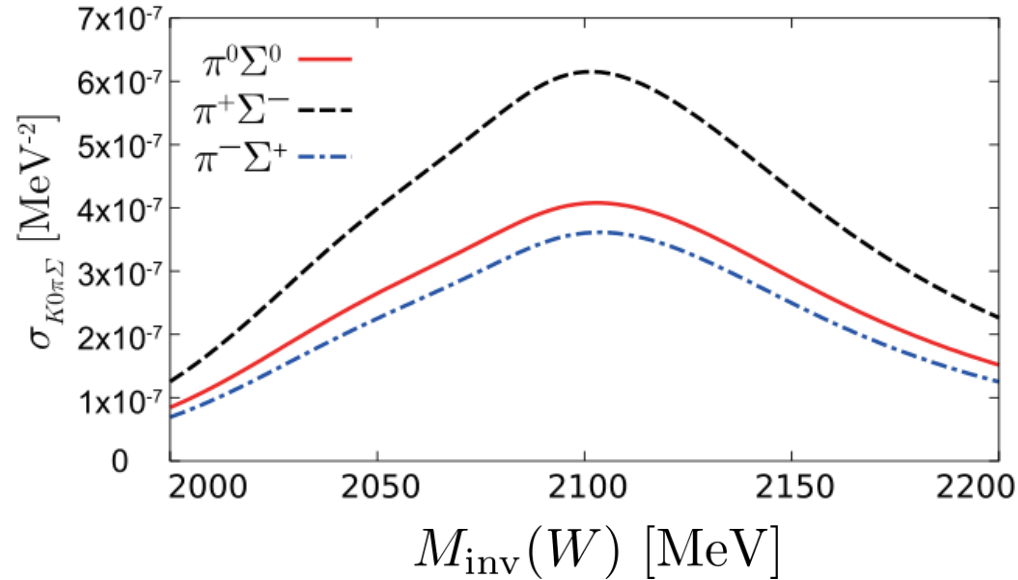
Peaks around @ $m_{\text{inv}} \sim 1380-1410 \text{ MeV}$

-- Role of lower pole of $\Lambda(1405)$

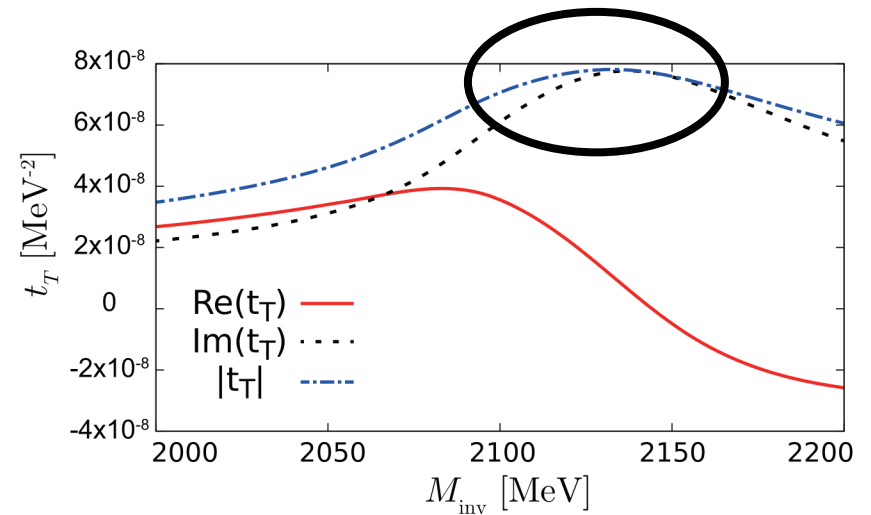
-- Peak position: lower than 1420 MeV

Maximum @ $W \sim 2100 \text{ MeV}$

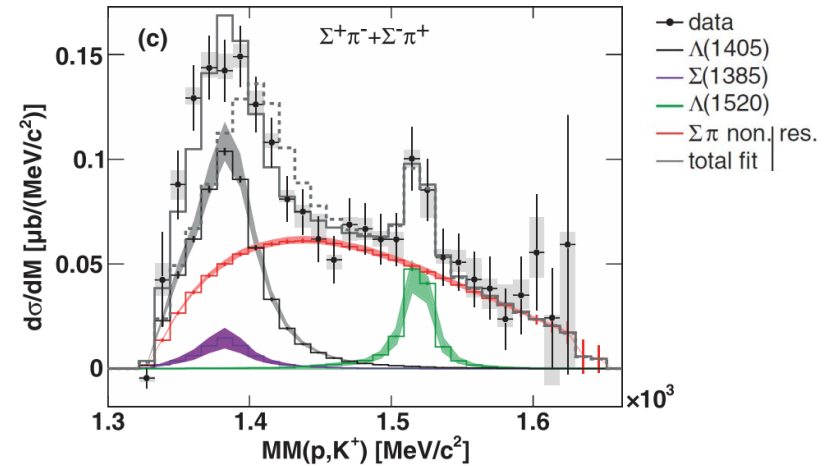
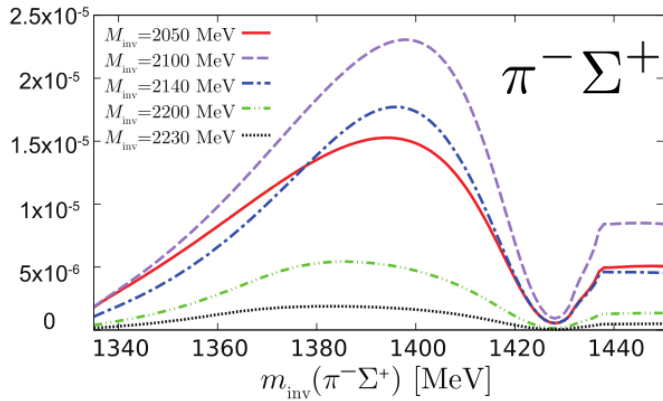
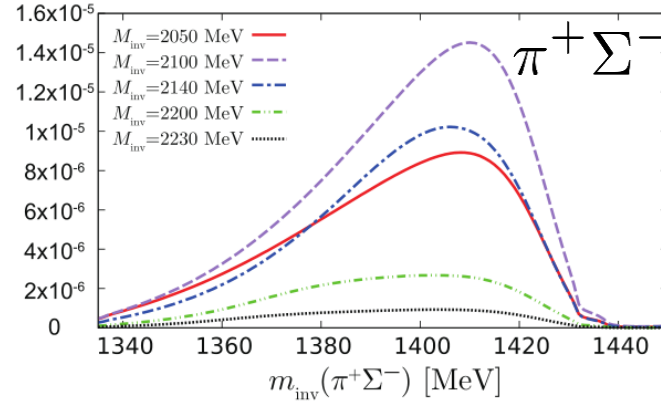
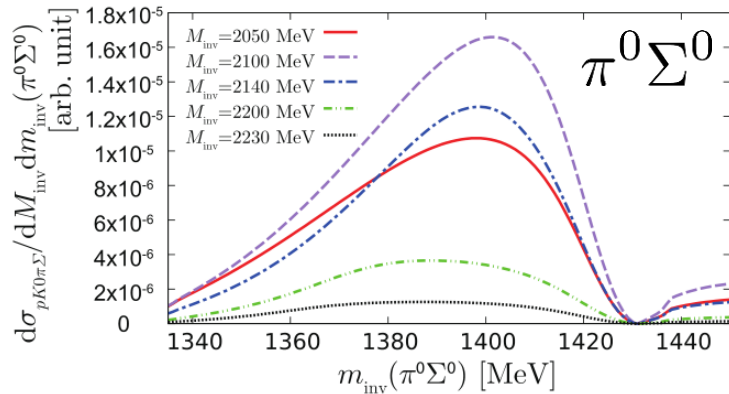
$$\sigma_{\pi^- p \rightarrow K^0 \pi \Sigma}(W)$$



- Peak @2100MeV
from triangle mechanism
- [lowered by the resonance amplitude
@2030MeV]



$$\frac{d^2 \sigma_{pp \rightarrow pK^+ \pi \Sigma}}{dM_{K\Lambda^*} dm_{\pi \Sigma}}$$



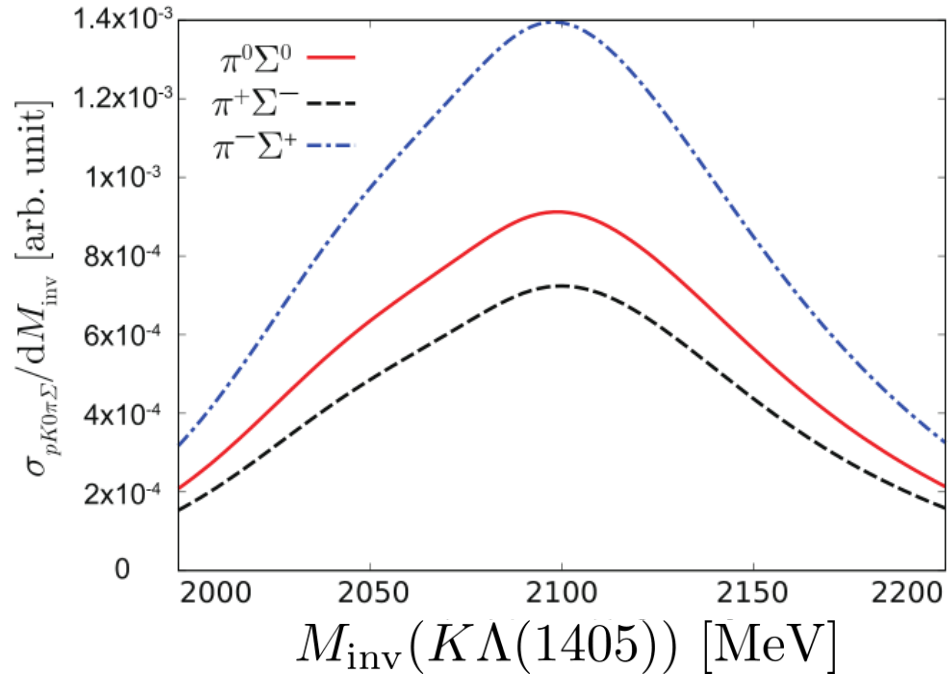
Peaks around @ $m_{\text{inv}} \sim 1380-1410 \text{ MeV}$

- Role of lower pole of $\Lambda(1405)$
- Peak position: lower than 1420 MeV

Maximum @ $M_{K\Lambda^*} \sim 2100 \text{ MeV}$

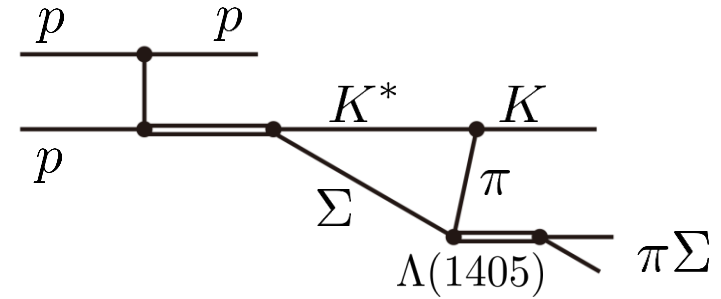
- Role of triangle mechanism and N^* resonance

$$\frac{d\sigma_{pp \rightarrow pK^+\pi\Sigma}}{dM_{\text{inv}}(K\Lambda(1405))}$$



■ Peak @2100MeV
from triangle mechanism

[lowered by the resonance amplitude @2030MeV]

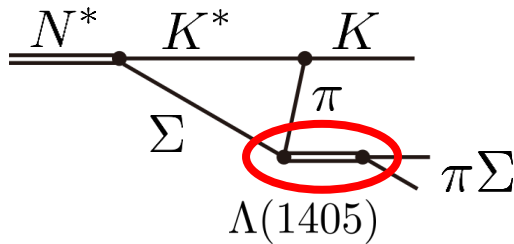


Summary

$$\begin{cases} \text{-- } \pi^- p \rightarrow K^0 \pi \Sigma \\ \text{-- } pp \rightarrow pK^+ \pi \Sigma \end{cases}$$

with triangle mechanism

○ $K^* \Sigma \pi$ loop: peak @2140MeV



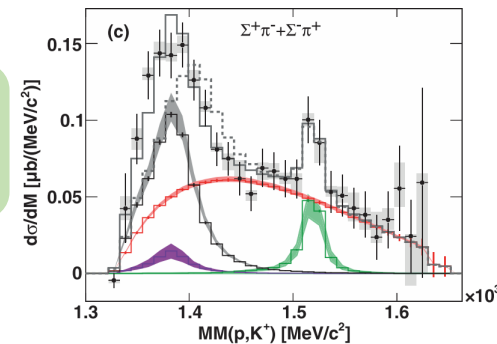
← $\pi \Sigma \rightarrow \pi \Sigma$ amplitude

: larger weight of **lower $\Lambda(1405)$ pole**



$$\frac{d\sigma_{\pi^- p \rightarrow K^0 \pi \Sigma}}{dm_{\pi \Sigma}} \quad \text{and} \quad \frac{d^2\sigma_{pp \rightarrow pK^+ \pi \Sigma}}{dM_{K\Lambda^*} dm_{\pi \Sigma}} : \text{peak @1380-1410MeV}$$

-- explain experimental peak position



No exp. data

@W=2140MeV

$$\sigma_{\pi^- p \rightarrow K^0 \pi \Sigma} \quad \text{and} \quad \frac{d\sigma_{pp \rightarrow pK^+ \pi \Sigma}}{dM_{K\Lambda^*}} : \text{peak @2100MeV}$$

← from triangle mechanism and resonance amplitude