

# Bonn-Gatchina analysis of the $\eta$ photo-production reactions

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### Bonn-Gatchina Partial Wave Analysis



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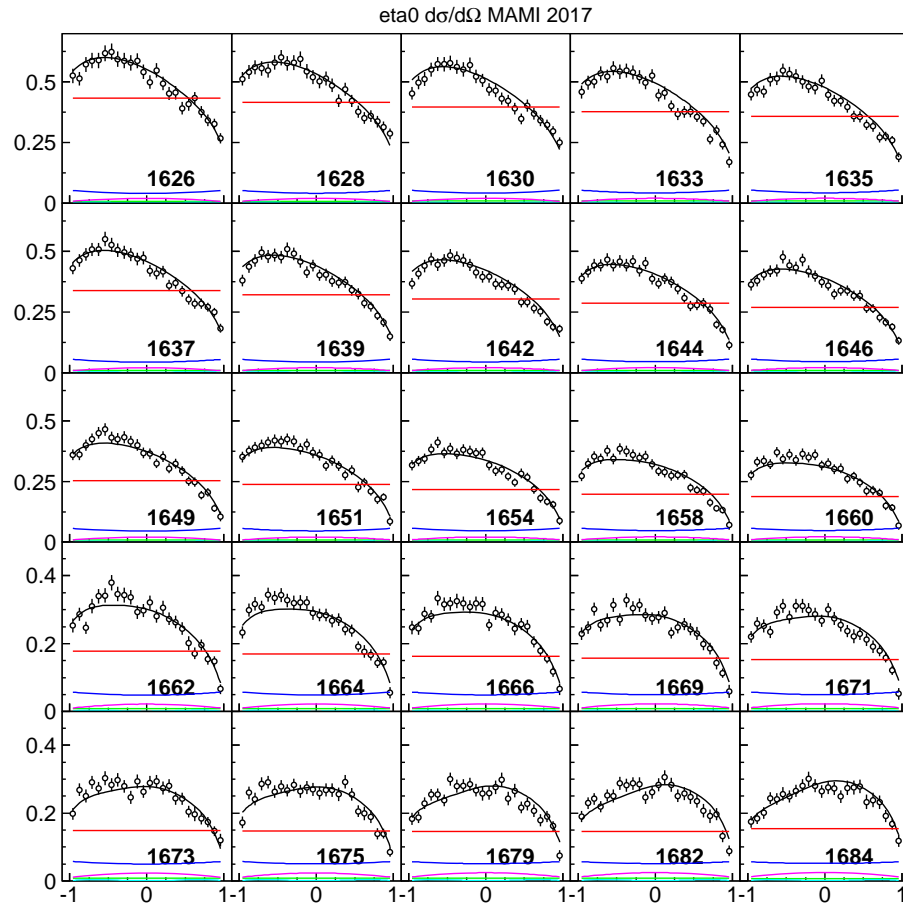
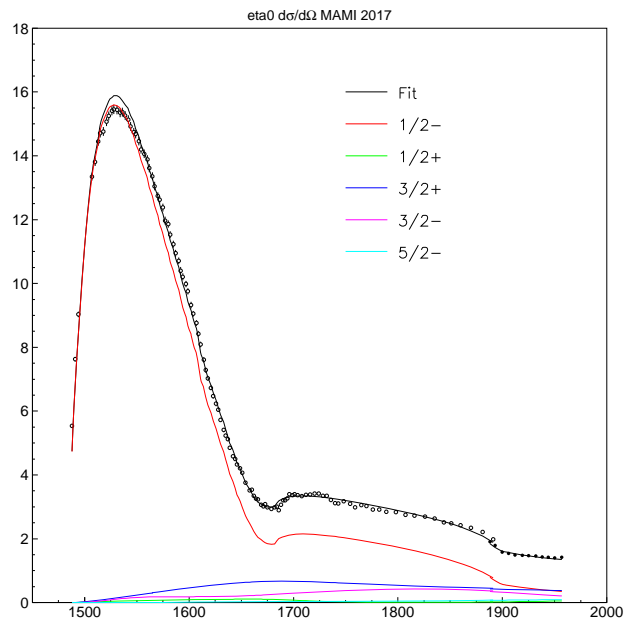
Responsible: Dr. V. Nikonov, E-mail: [nikonov@hiskp.uni-bonn.de](mailto:nikonov@hiskp.uni-bonn.de)  
Last changes: January 26<sup>th</sup>, 2010.

## Recently included meson photoproduction data

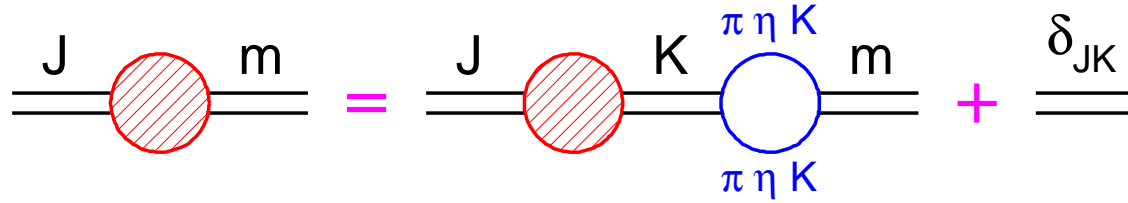
DATA	2011-2016	added in 2016-2018
$\gamma n \rightarrow \Lambda K, \Sigma^- K$		$\frac{d\sigma}{d\Omega}$ (CLAS), E (CLAS)
$\gamma n \rightarrow \pi^- p$	$\frac{d\sigma}{d\Omega}, \Sigma, P$	E, $\Sigma$ (CLAS)
$\gamma n \rightarrow \eta n$ $\gamma p \rightarrow \eta p$	$\frac{d\sigma}{d\Omega}, \Sigma$ $\frac{d\sigma}{d\Omega}, \Sigma$ (GRAAL)	$\frac{d\sigma}{d\Omega}$ (MAMI) $\frac{d\sigma}{d\Omega} (h = \frac{1}{2})$ (CB-ELSA) $\frac{d\sigma}{d\Omega}, F, T$ (MAMI) T, P, H, G, (CB-ELSA) E, $\Sigma$ (CB-ELSA, CLAS)
$\gamma p \rightarrow \eta' p$		$\frac{d\sigma}{d\Omega}, \Sigma$
$\gamma p \rightarrow K^+ \Lambda$ $\gamma p \rightarrow K^+ \Sigma^0$	$\frac{d\sigma}{d\Omega}, \Sigma, P, T, C_x, C_z, O_{x'}, O_{z'}$ $\frac{d\sigma}{d\Omega}, \Sigma, P, C_x, C_z$	$\Sigma, P, T, O_x, O_z$ (CLAS) $\Sigma, P, T, O_x, O_z$ (CLAS)
$\pi^- p \rightarrow \pi^+ \pi^- n$ $\pi^- p \rightarrow \pi^- \pi^0 p$		$d\sigma/d\Omega$ (HADES) $d\sigma/d\Omega$ (HADES)
$\gamma p \rightarrow \pi^0 \pi^0 p$ $\gamma p \rightarrow \pi^+ \pi^- p$	$d\sigma/d\Omega, \Sigma, E, I_c, I_s$	T, P, H, F, $P_x, P_y$ (CB-ELSA) $d\sigma/d\Omega, I_c, I_s$ (CLAS)
$\gamma p \rightarrow \omega p$	$d\sigma/d\Omega, \Sigma, \rho_{ij}^k, E, G$ (CB-ELSA)	$\Sigma$ (CLAS) P, T, F, H (CLAS)
$\gamma p \rightarrow K^* \Lambda$		$d\sigma/d\Omega, \rho_{ij}$

# The analysis of the new $\gamma p \rightarrow \eta p$ data.

## New MAMI data: a strong cusp effect from the $\eta' p$ channel



## N/D based (D-matrix) analysis of the data



$$D_{jm} = D_{jk} \sum_{\alpha} B_{\alpha}^{km}(s) \frac{1}{M_m - s} + \frac{\delta_{jm}}{M_j^2 - s} \quad \hat{D} = \hat{\kappa}(I - \hat{B}\hat{\kappa})^{-1}$$

$$\hat{\kappa} = \text{diag} \left( \frac{1}{M_1^2 - s}, \frac{1}{M_2^2 - s}, \dots, \frac{1}{M_N^2 - s}, R_1, R_2, \dots \right)$$

$$\hat{B}_{ij} = \sum_{\alpha} B_{\alpha}^{ij} = \sum_{\alpha} \int \frac{ds'}{\pi} \frac{g_{\alpha}^{(R)i} \rho_{\alpha}(s', m_{1\alpha}, m_{2\alpha}) g_{\alpha}^{(L)j}}{s' - s - i0}$$

In the present fits we calculate the elements of the  $B_\alpha^{ij}$  using one subtraction taken at the channel threshold  $M_\alpha = (m_{1\alpha} + m_{2\alpha})$ :

$$B_\alpha^{ij}(s) = B_\alpha^{ij}(M_\alpha^2) + (s - M_\alpha^2) \int_{M_\alpha^2}^{\infty} \frac{ds'}{\pi} \frac{g_\alpha^{(R)i} \rho_\alpha(s', m_{1\alpha}, m_{2\alpha}) g_\alpha^{(L)j}}{(s' - s - i0)(s' - M_\alpha^2)}.$$

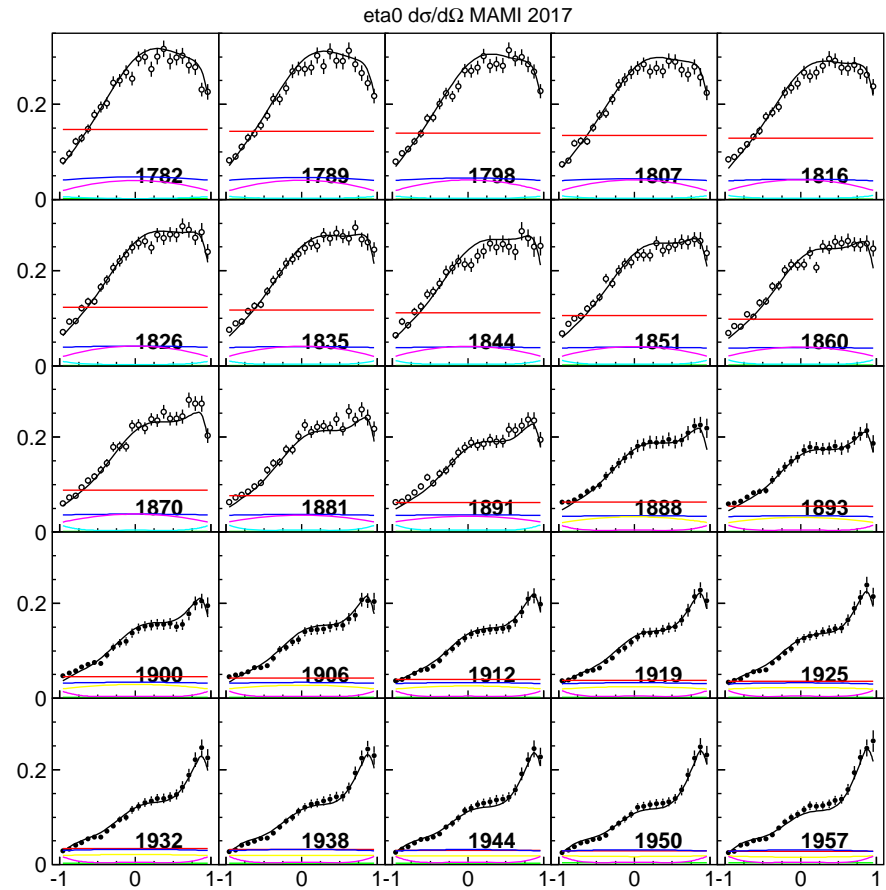
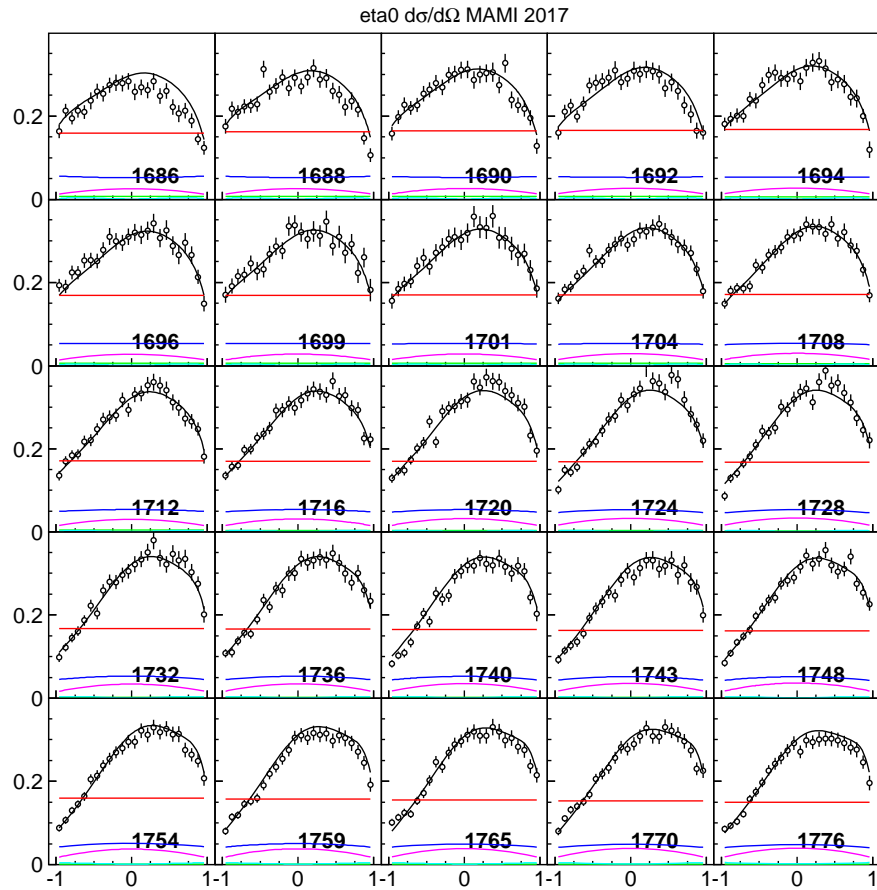
In this case the expression for elements of the  $\hat{B}$  matrix can be rewritten as:

$$B_\alpha^{ij}(s) = g_a^{(R)i} \left( b^\alpha + (s - M_\alpha^2) \int_{M_\alpha^2}^{\infty} \frac{ds'}{\pi} \frac{\rho_\alpha(s', m_{1\alpha}, m_{2\alpha})}{(s' - s - i0)(s' - M_\alpha^2)} \right) g_\beta^{(L)j} = g_a^{(R)i} B_\alpha g_\beta^{(L)j}$$

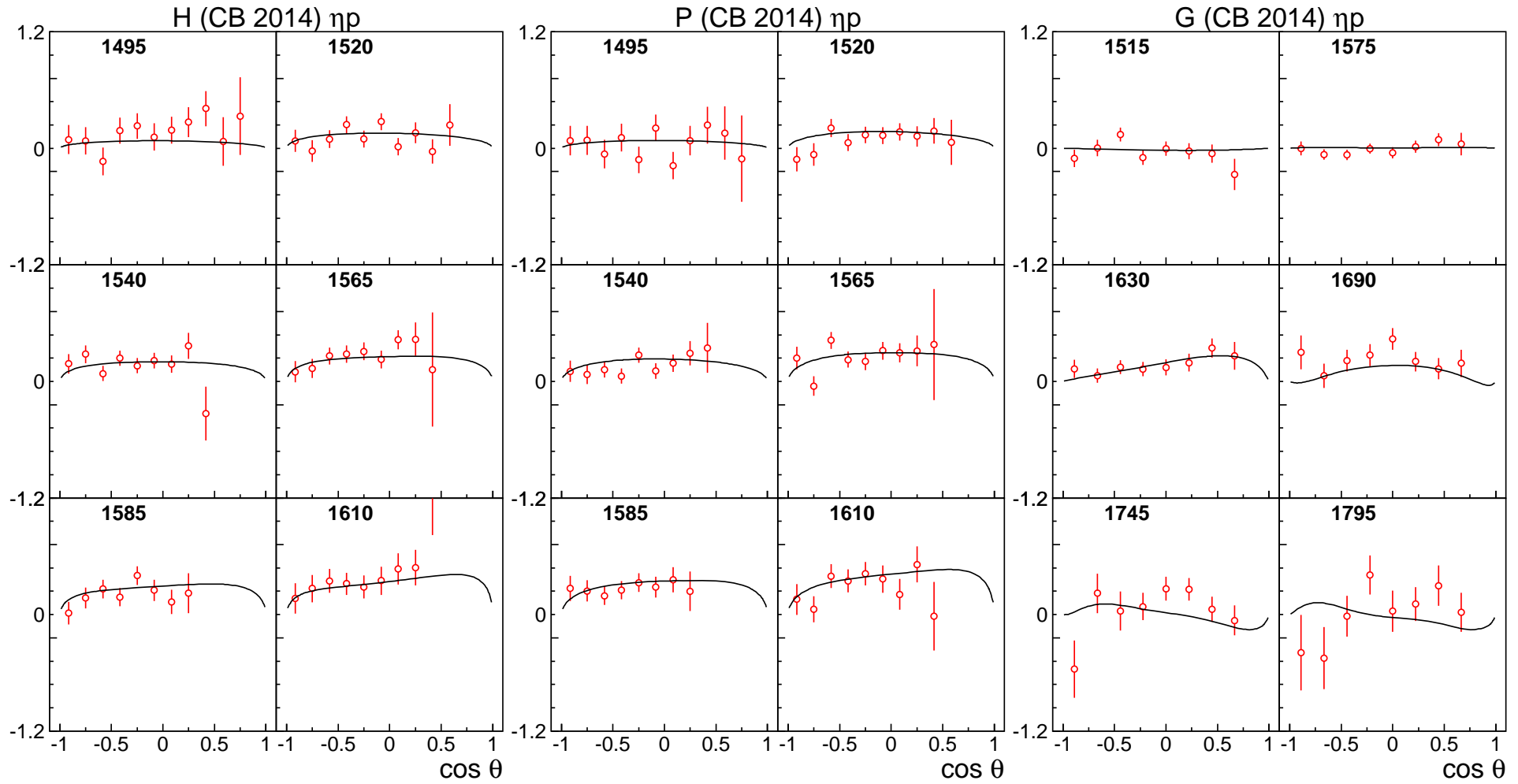
and D-matrix method equivalent to the K-matrix method with loop diagram with real part taken into account:

$$A = \hat{K}(I - \hat{B}\hat{K})^{-1} \quad B_{\alpha\beta} = \delta_{\alpha\beta} B_\alpha$$

# The analysis of the new $\gamma p \rightarrow \eta p$ data. $d\sigma/d\Omega$ (MAMI)

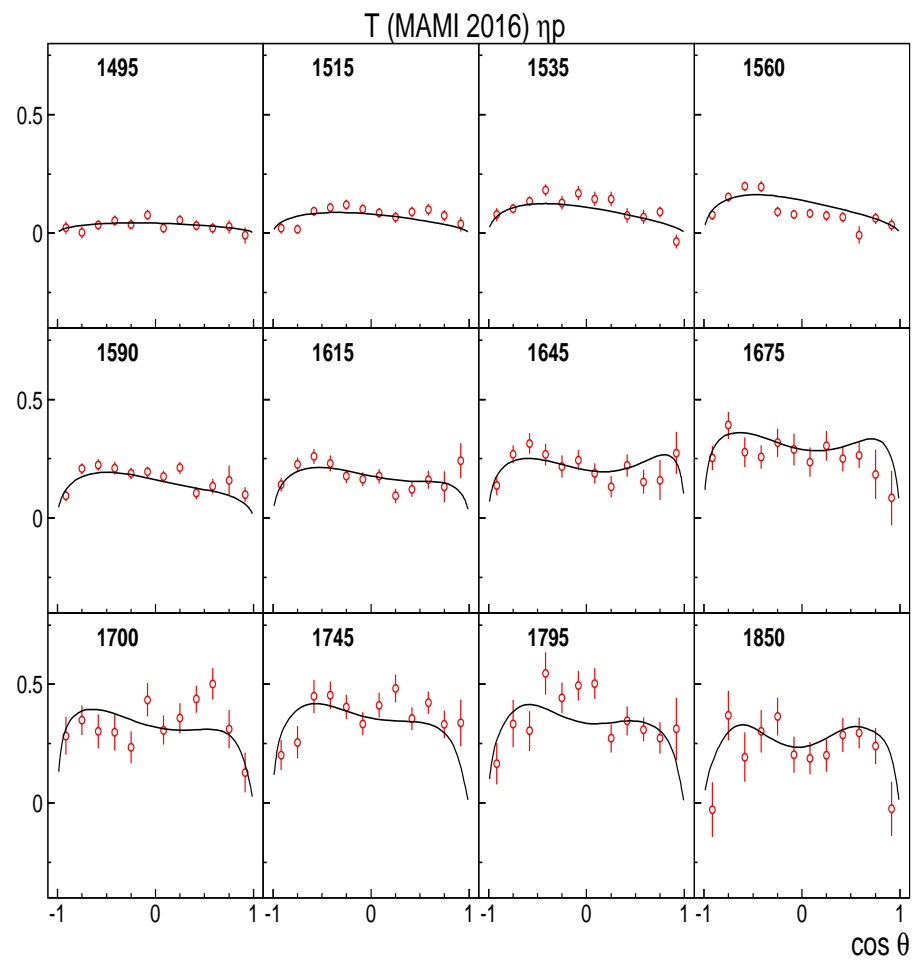
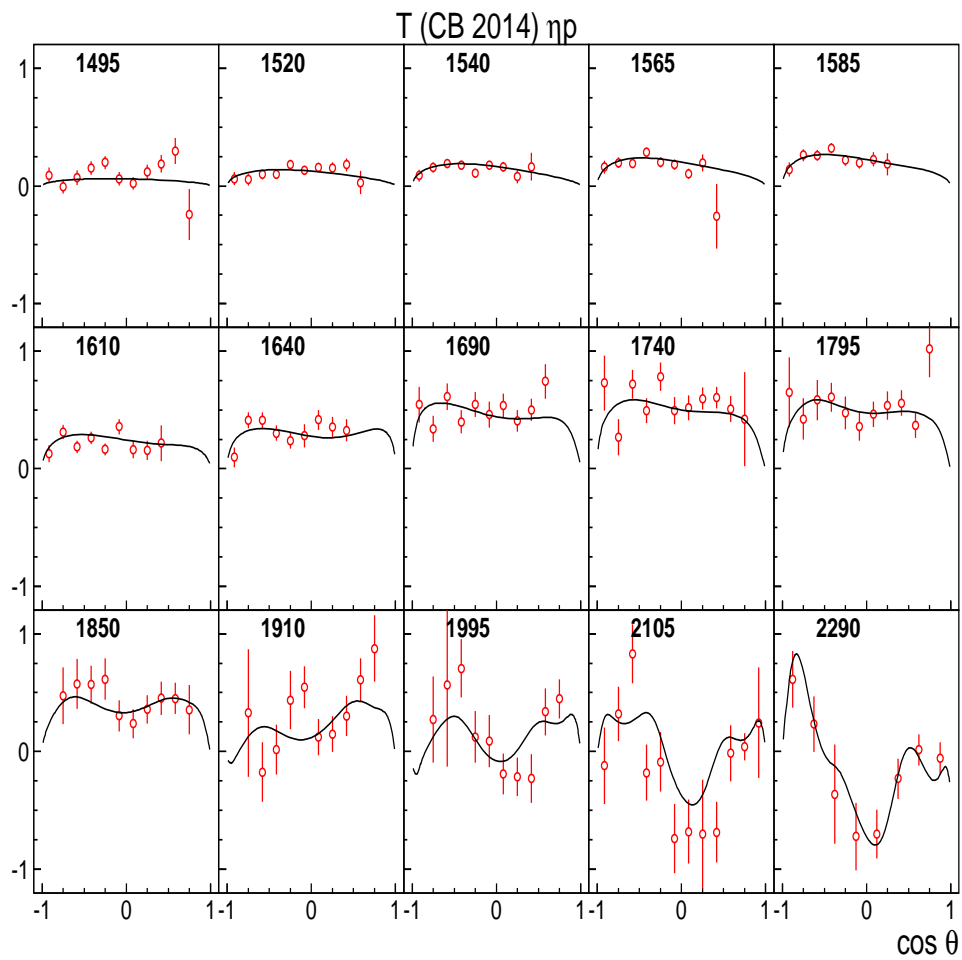


# The analysis of the new $\gamma p \rightarrow \eta p$ data. $H, P, T$ (CB-ELSA, Preliminary)

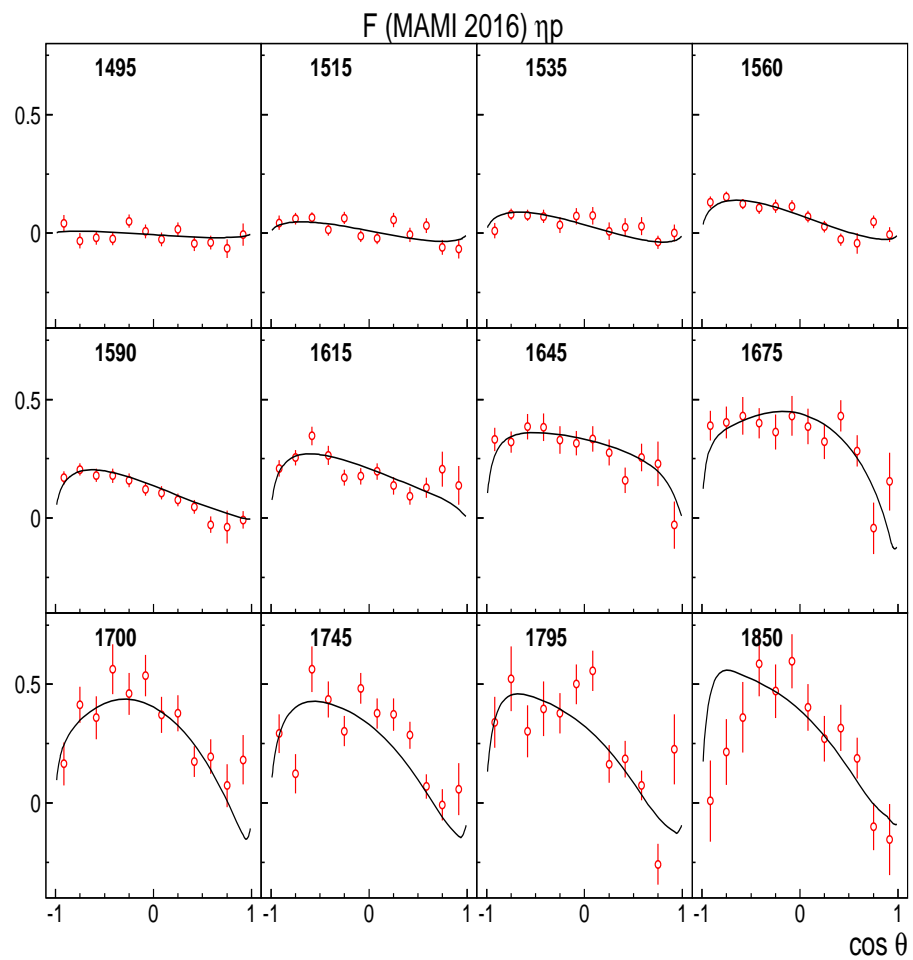
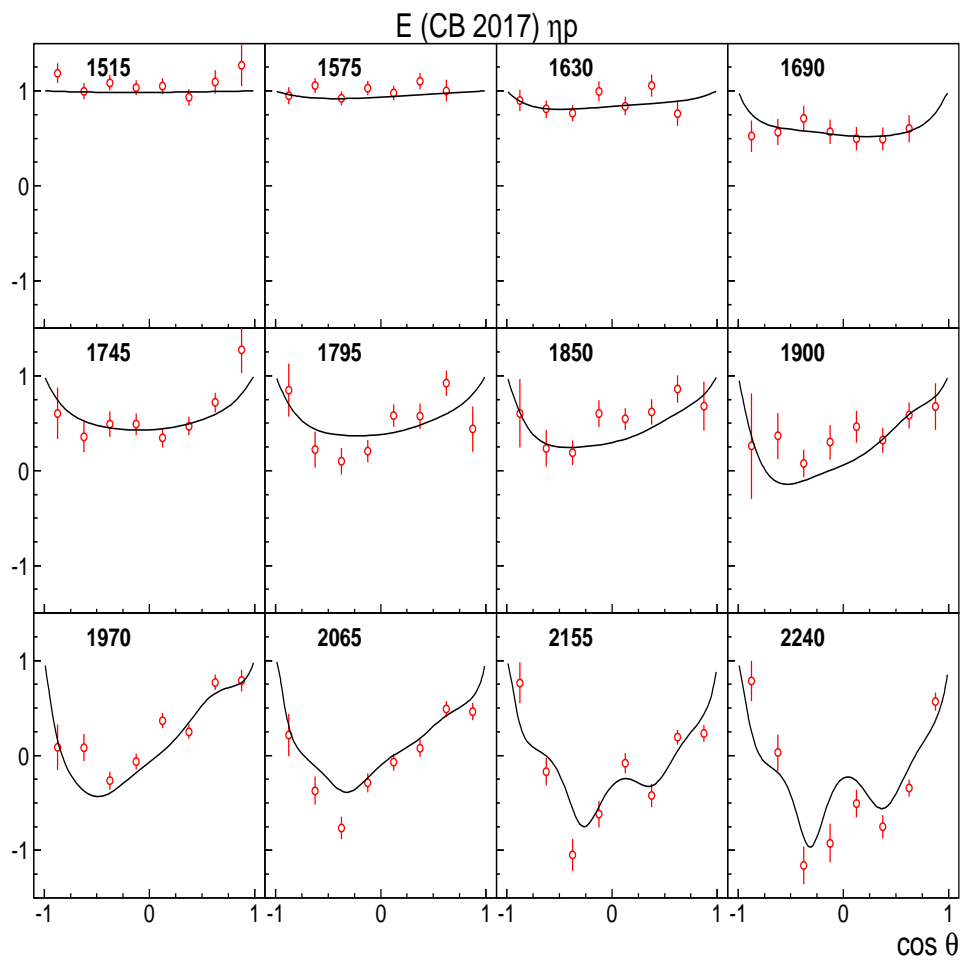




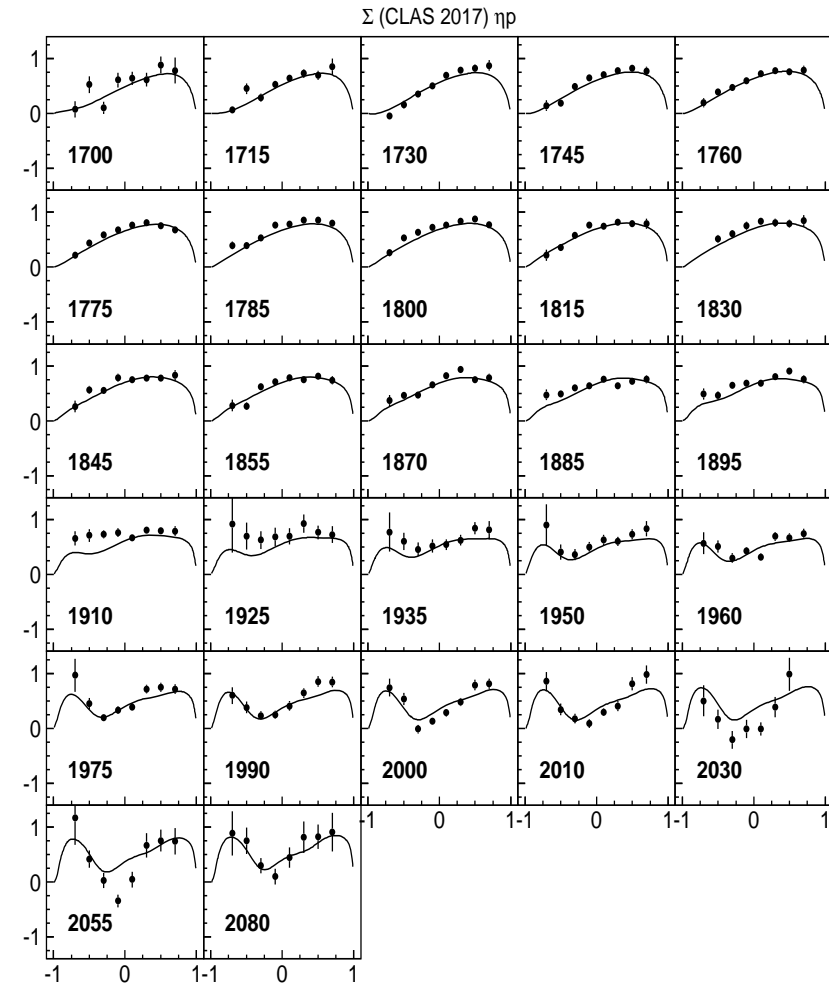
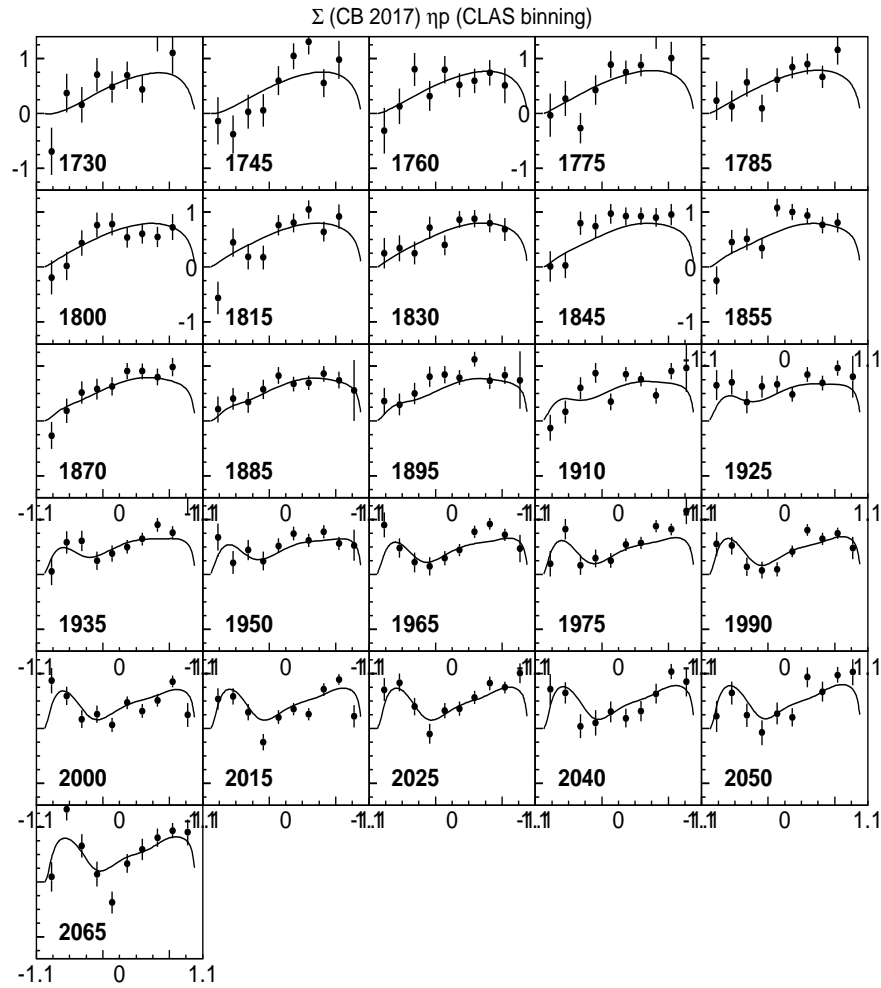
The analysis of the new  $\gamma p \rightarrow \eta p$  data.  $T$  (CB-ELSA, Preliminary), (MAMI scale 1.4)



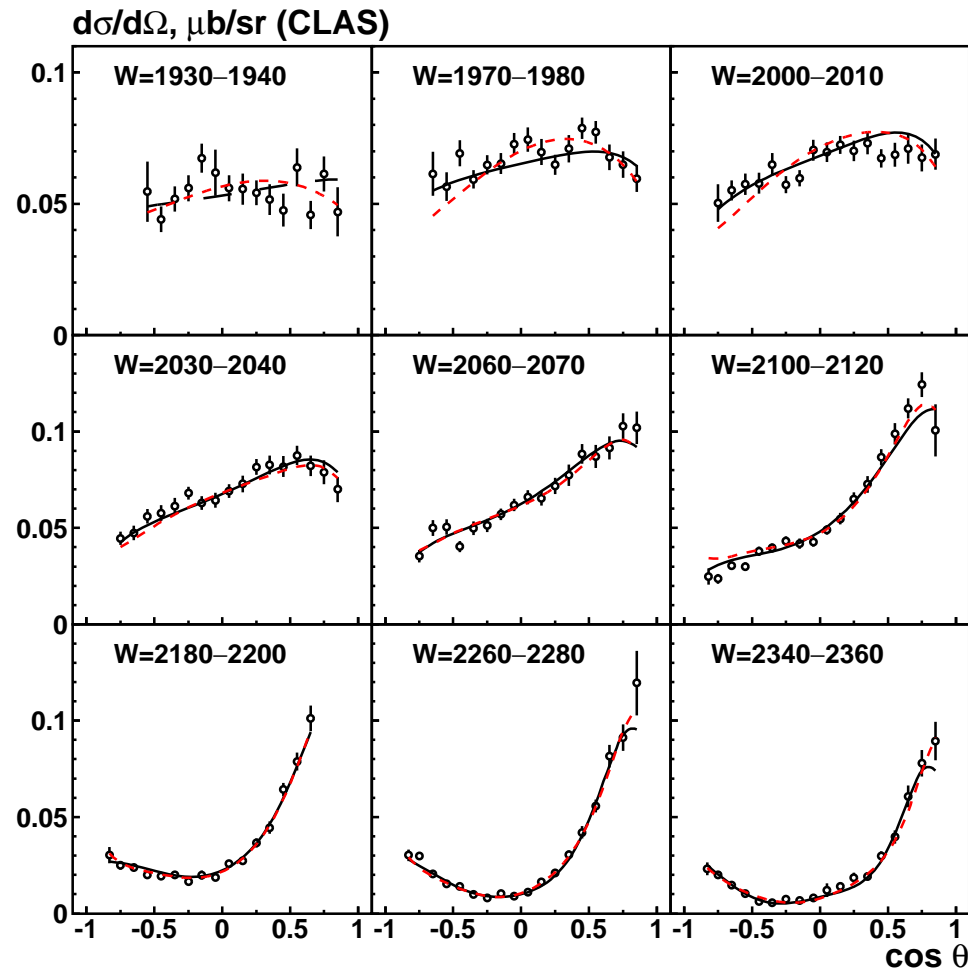
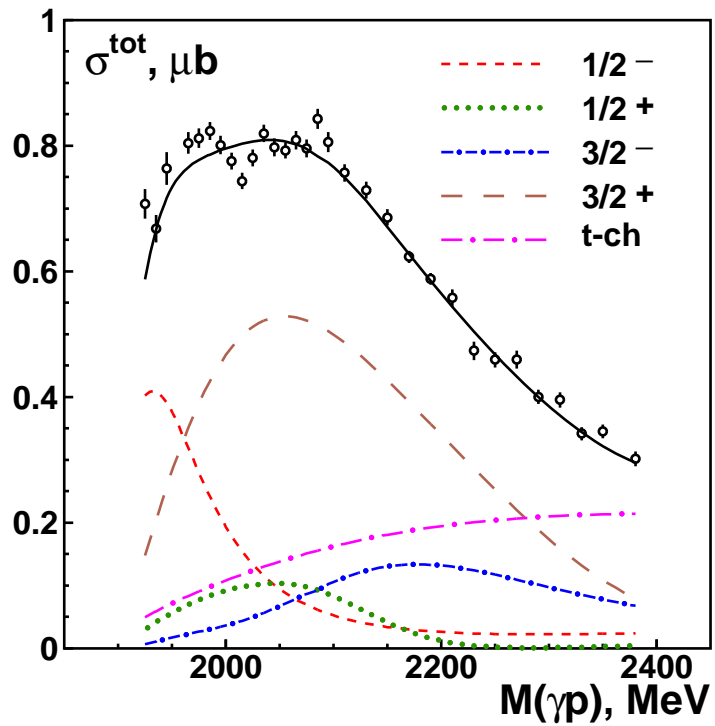
The analysis of the new  $\gamma p \rightarrow \eta p$  data.  $E$  (CB-ELSA, Preliminary),  $F$  (MAMI) (scale 1.4)



# The analysis of the new $\gamma p \rightarrow \eta p$ data. $\Sigma$ (CB-ELSA and CLAS)

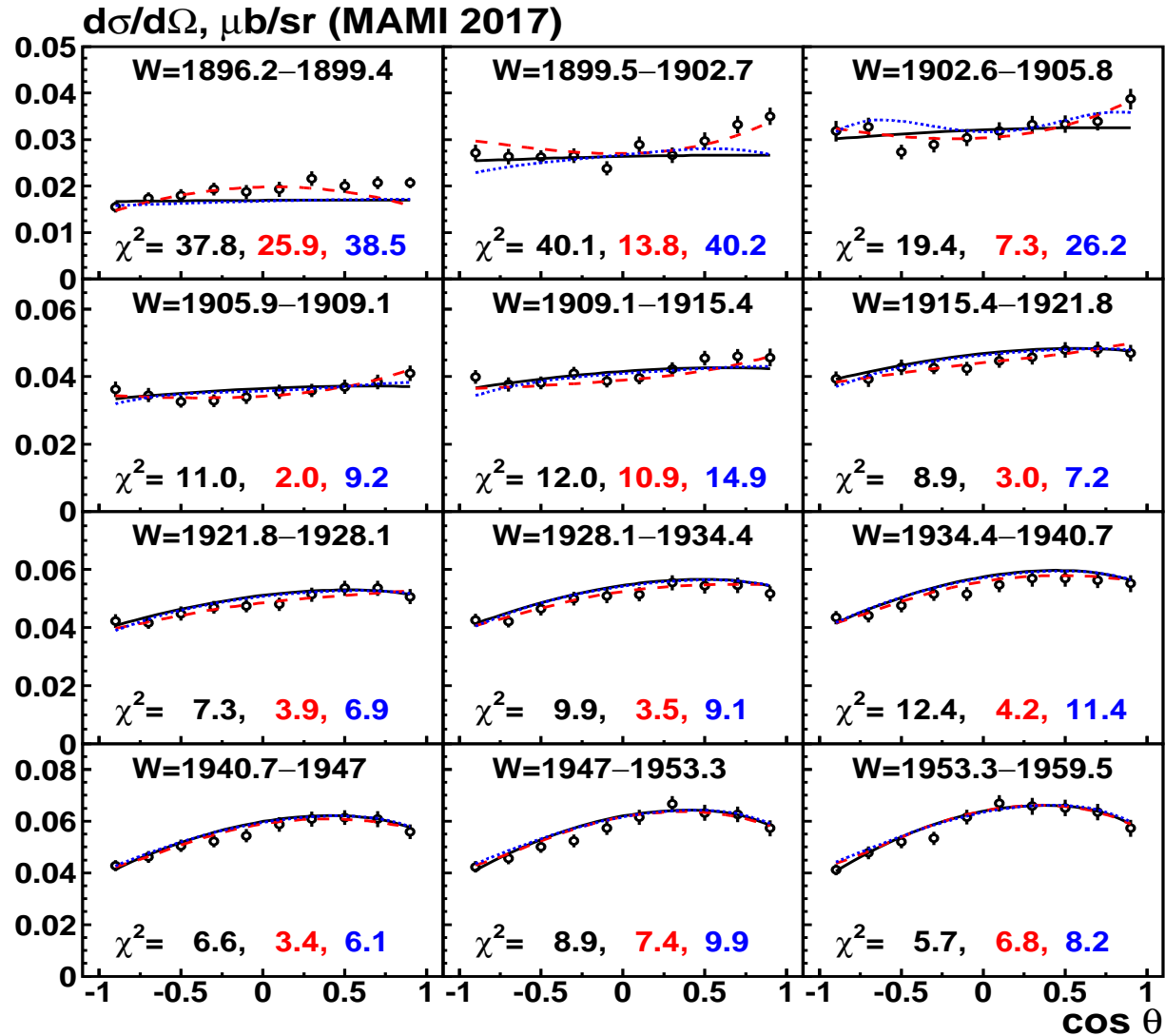


# The analysis of the $\gamma p \rightarrow \eta' p$ data.

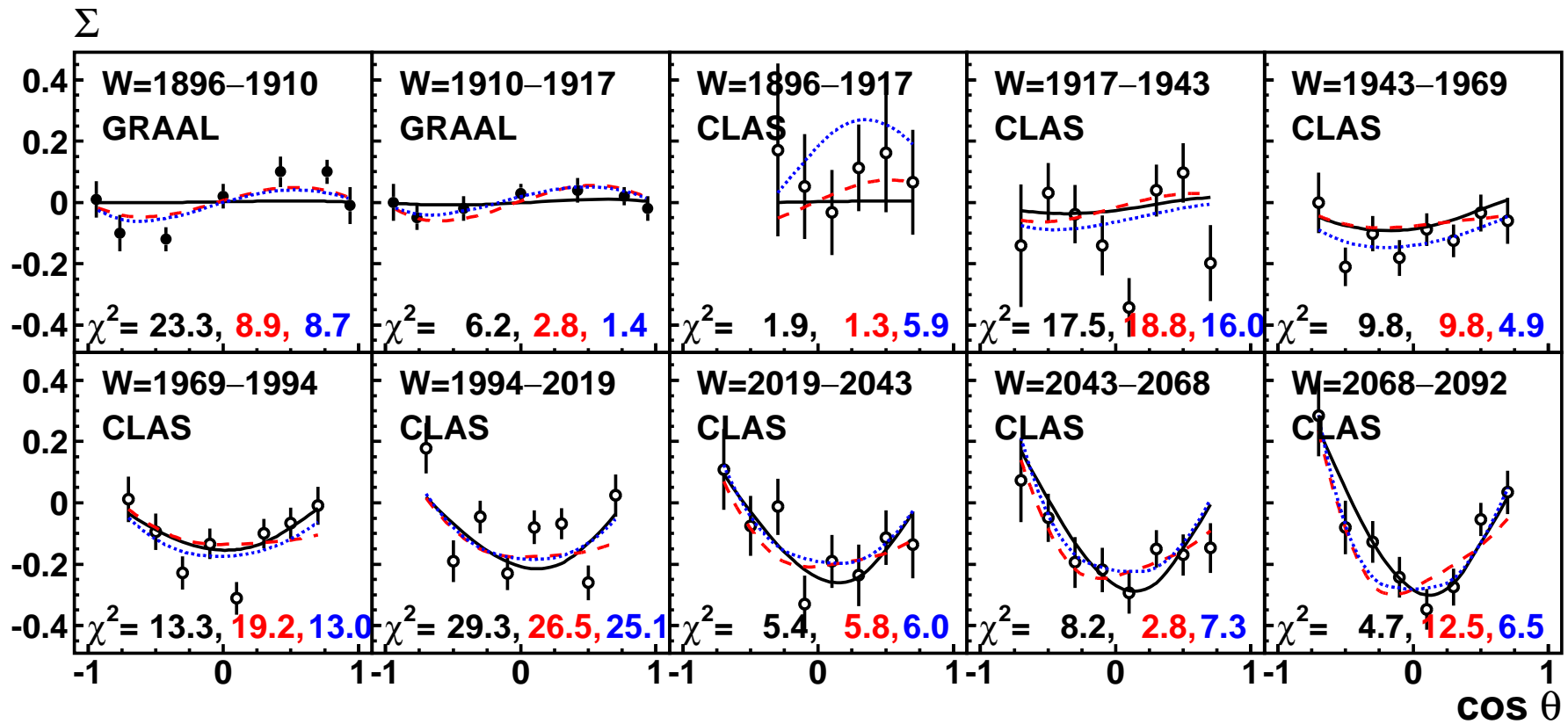


Strong contribution from the  $S_{11}(1895)$ ,  $P_{13}(1900)$ ,  $P_{11}(2100)$  and  $D_{13}(2120)$  states.

# MAMI differential cross section



# The beam asymmetry on $\gamma p \rightarrow \eta' p$



No narrow states

$D_{15}(1903)$

$D_{13}(1900)$

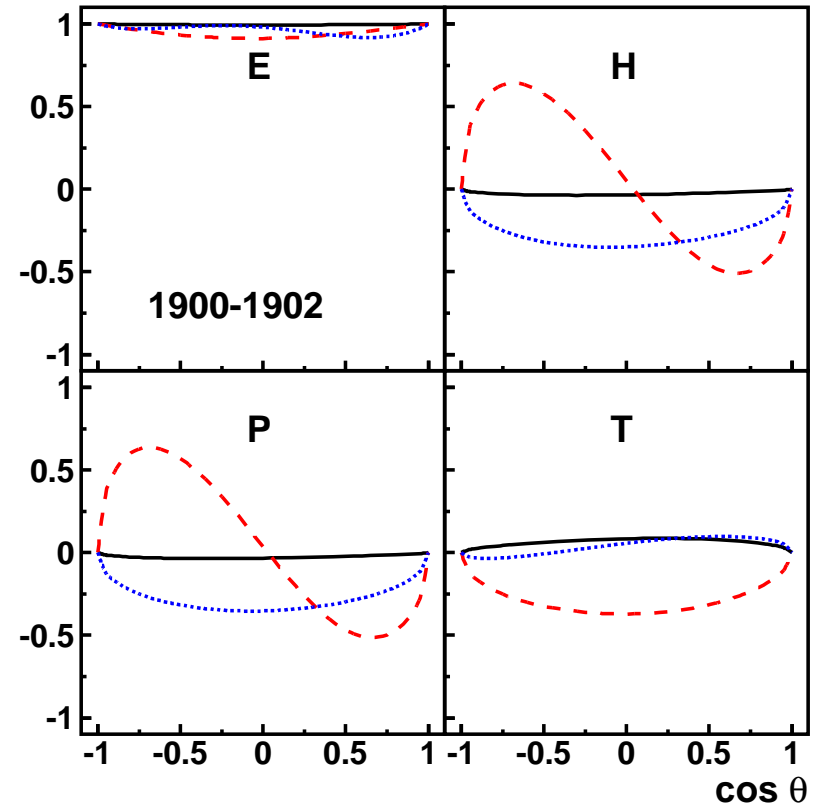
$$BR = \int_{\text{threshold}}^{\infty} \frac{ds}{\pi} \frac{(g_a^\alpha)^2 \rho_a^\alpha(s)}{(M_{\text{BW}}^\alpha)^2 - s)^2 + (\sum_{\alpha} g_a^\alpha)^2 \rho_a^\alpha(s))^2}$$

**Table 1: The masses and widths  $M_{\text{pole}}, \Gamma_{\text{pole}}$  are given in MeV, the moduli of the helicity couplings in  $10^{-3} \text{ GeV}^{-1/2}$ .  $\Delta\pi$  denotes  $\Delta(1232)3/2^+ \pi$  and  $N^* \pi$ ,  $N(1520)3/2^- \pi$ .**

	$M_{\text{pole}}$	$\Gamma_{\text{pole}}$	Branching ratios (in %) for $N^* \rightarrow$						
			$N\eta'$	$N\pi$	$N\eta$	$N\omega$	$\Delta\pi$	$N^* \pi$	$N\sigma$
$N(1895)1/2^-$	<b><math>1907 \pm 10</math></b>	$100_{-10}^{+40}$	<b><math>13 \pm 5</math></b>	<b><math>2.5 \pm 1.5</math></b>	<b><math>10 \pm 5</math></b>	<b><math>28 \pm 12</math></b>	<b><math>7 \pm 4</math></b>	-	-
$N(2100)1/2^+$	<b><math>2100 \pm 30</math></b>	<b><math>280 \pm 35</math></b>	<b><math>8 \pm 3</math></b>	<b><math>13 \pm 5</math></b>	<b><math>25 \pm 10</math></b>	<b><math>15 \pm 10</math></b>	<b><math>10 \pm 4</math></b>	<b><math>30 \pm 4</math></b>	<b><math>20 \pm 6</math></b>
$N(2120)3/2^-$	<b><math>2115 \pm 40</math></b>	<b><math>345 \pm 35</math></b>	<b><math>4 \pm 2</math></b>	<b><math>5 \pm 3</math></b>	<b><math>3 \pm 2</math></b>	<b><math>12 \pm 8</math></b>	<b><math>70 \pm 23</math></b>		<b><math>11 \pm 4</math></b>
$N(1900)3/2^+$	<b><math>1910 \pm 30</math></b>	<b><math>280 \pm 50</math></b>	<b><math>6 \pm 2</math></b>	<b><math>3 \pm 2</math></b>	<b><math>3 \pm 1</math></b>	<b><math>15 \pm 9</math></b>	<b><math>50 \pm 15</math></b>	<b><math>15 \pm 8</math></b>	<b><math>4 \pm 3</math></b>

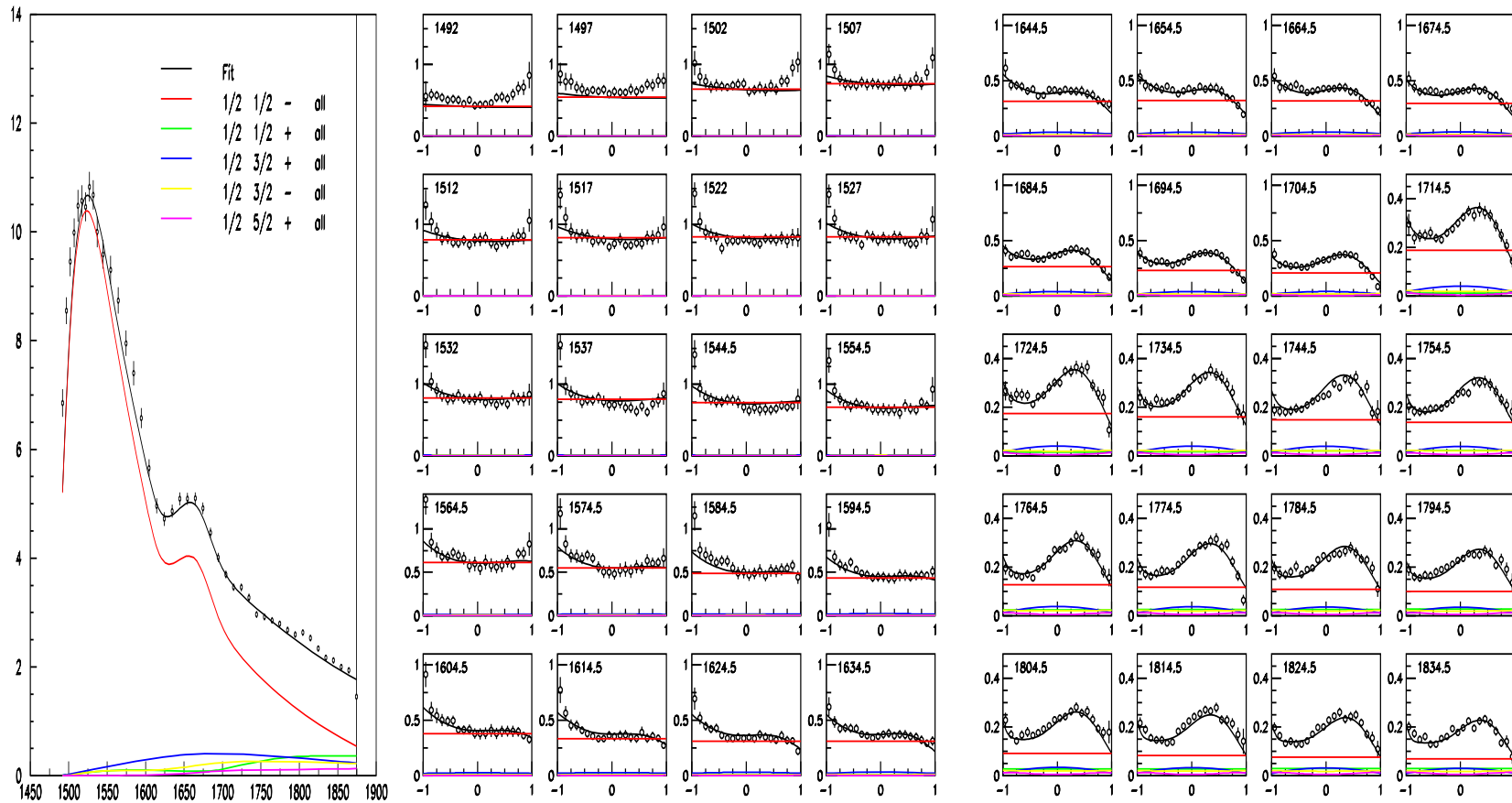
## The description of the data below $W=1917$ MeV and the prediction of other observables

Resonance	N	Basic	$D_{13}$	$D_{15}$
$M$ (MeV)			1900	1903
$\Gamma$ (MeV)			1	1
$\chi^2$ ( $\Sigma$ )	13	29.5	11.7	10.1
$\chi^2$ ( $\frac{d\sigma}{d\Omega}$ )	50	120.3	59.9	129.0

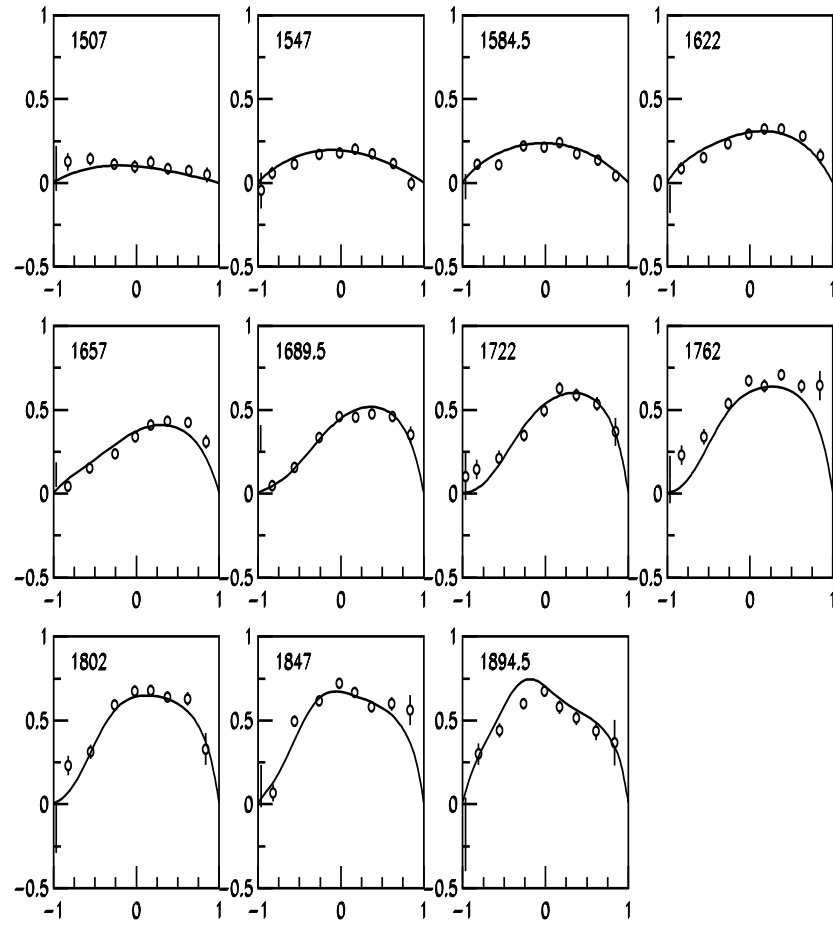
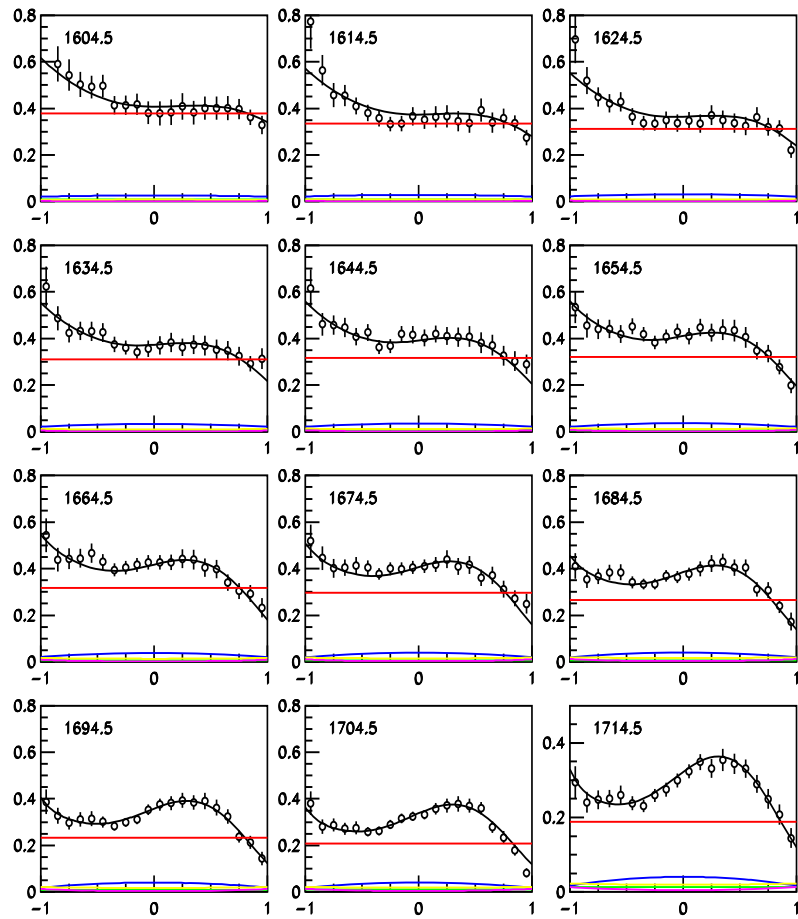




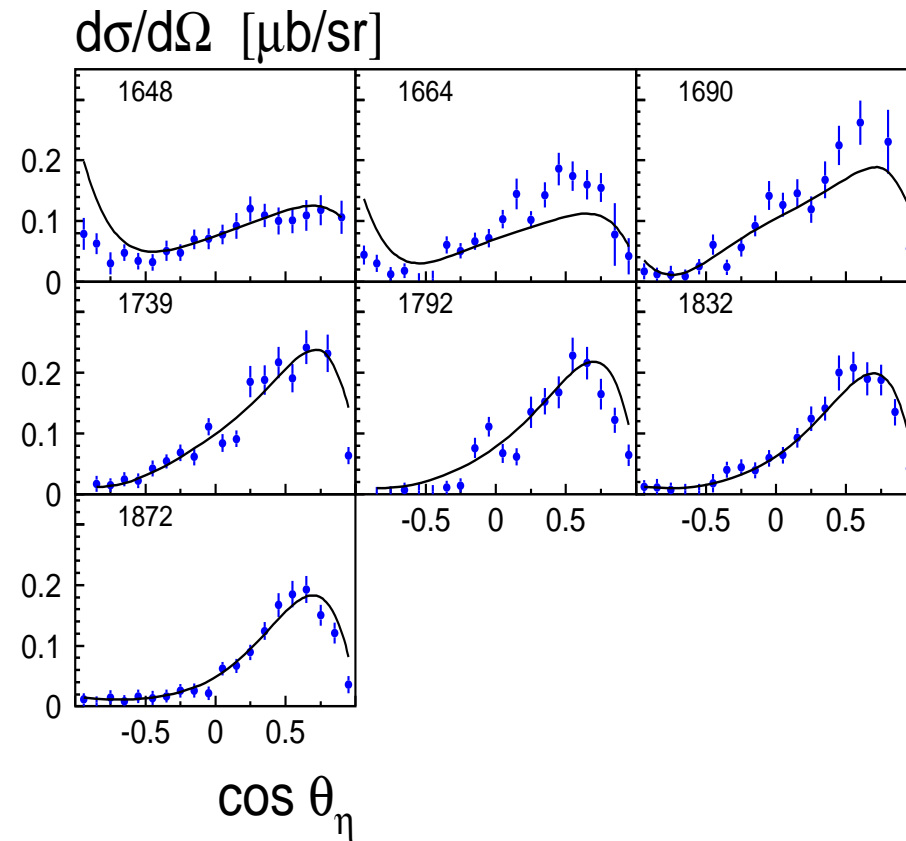
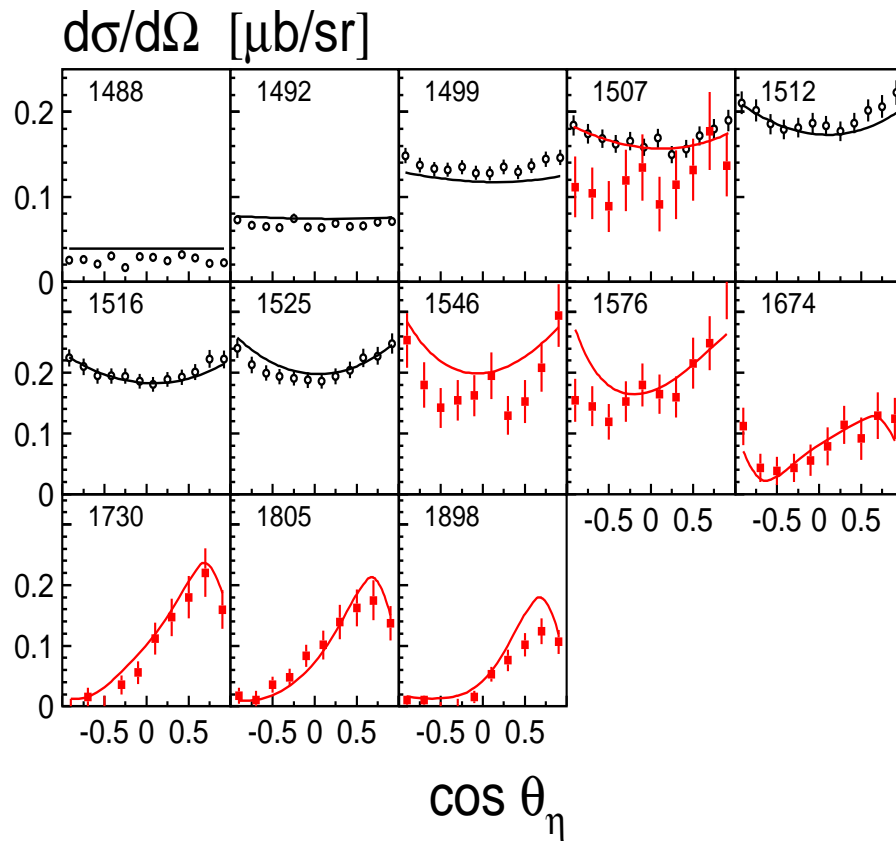
# The description of the $\gamma n \rightarrow \eta n$ data with N/D-based approach



# The description of the differential cross section and beam asymmetry in the selected energy region



## The description of the $\pi p \rightarrow \eta p$ data with N/D-based approach



S. Prakhov et al., Phys. Rev. C72(2005) 015203

W. B. Richards et al., Phys. Rev. D1, 10 (1970)

R.M. Brown et al. Nucl. Phys. B153, 89 (1979)

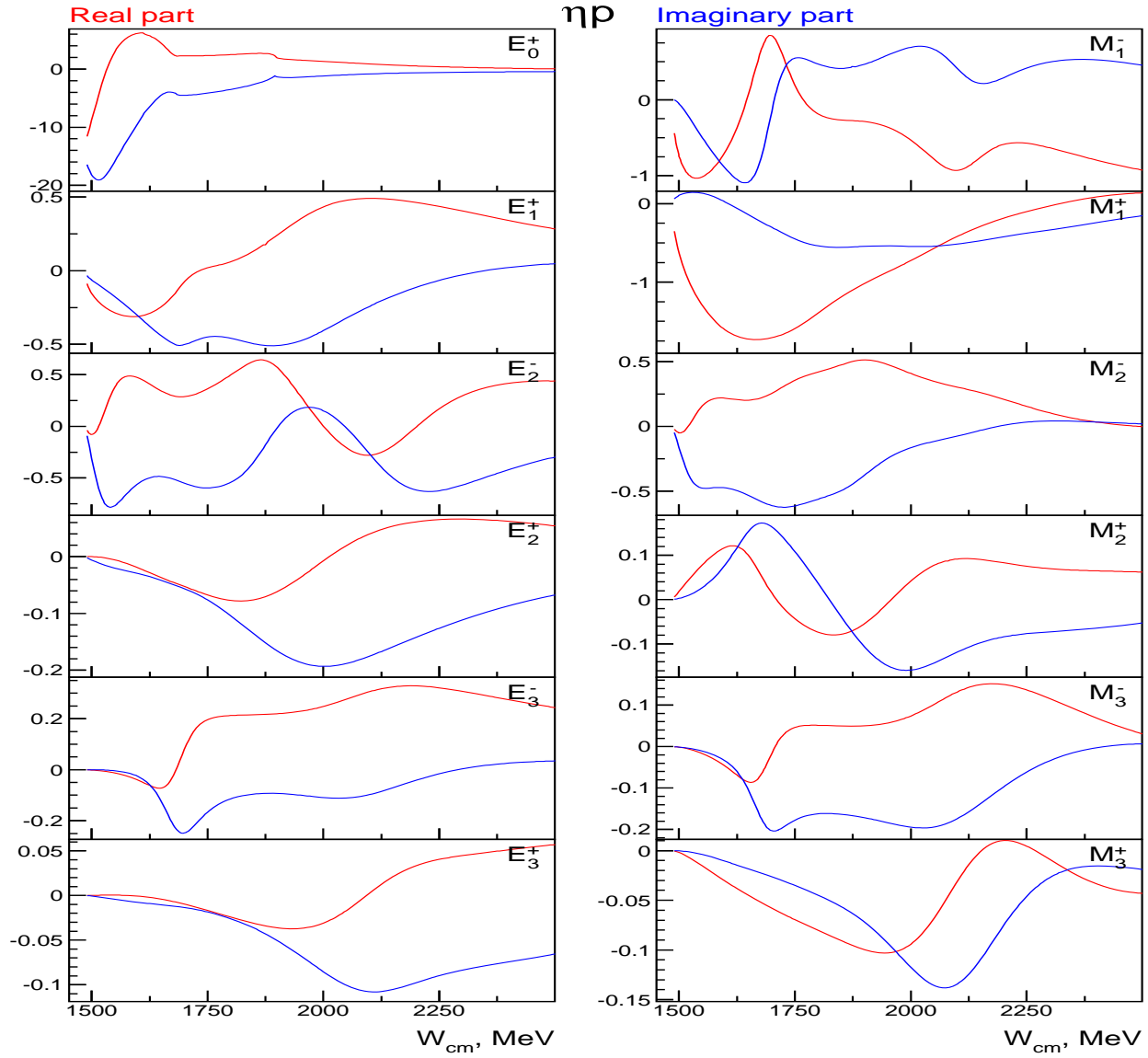
**Residues in the pole for  $\pi N \rightarrow \eta N$  (MeV) and  $\gamma N \rightarrow \eta N$  ( $10^{-3} \text{ GeV}^{-\frac{1}{2}}$ ) reactions amplitudes.**

<b>Res</b>	$N(1535)1/2^-$		$N(1650)1/2^-$		$N(1895)1/2^-$	
$\pi N \rightarrow \eta N$	<b>(29±2)</b>	<b>-(84±3)<sup>o</sup></b>	<b>-(22±3)</b>	<b>(2±12)<sup>o</sup></b>	<b>(5±3)</b>	<b>(25±20)<sup>o</sup></b>
$\gamma p \rightarrow \eta p (\mathbf{E}, -A^{\frac{1}{2}})$	<b>-(19±2)</b>	<b>-(60±3)<sup>o</sup></b>	<b>(4±0.3)</b>	<b>(31±10)<sup>o</sup></b>	<b>(1.7±0.8)</b>	<b>(20±20)<sup>o</sup></b>
<b>Res</b>	$N(1440)1/2^+$		$N(1710)1/2^+$		$N(1880)1/2^+$	
$\pi N \rightarrow \eta N$	<b>-(20±6)</b>	<b>(0±30)<sup>o</sup></b>	<b>(7±2)</b>	<b>(54±15)<sup>o</sup></b>	<b>(12±5)</b>	<b>(60±18)<sup>o</sup></b>
$\gamma p \rightarrow \eta p (\mathbf{M}, A^{\frac{1}{2}})$	<b>(5±2)</b>	<b>(15±30)<sup>o</sup></b>	<b>(2.2±0.7)</b>	<b>-(73±16)<sup>o</sup></b>	<b>(0.9±0.35)</b>	<b>(60±17)<sup>o</sup></b>
<b>Res</b>	$N(1520)3/2^-$		$N(1720)3/2^+$		$N(1900)3/2^+$	
$\pi N \rightarrow \eta N$	<b>(1.7±0.4)</b>	<b>-(88±15)<sup>o</sup></b>	<b>(8.1±2.6)</b>	<b>(45±10)<sup>o</sup></b>	<b>(2.9±0.8)</b>	<b>(55±30)<sup>o</sup></b>
$\gamma p \rightarrow \eta p (A^{\frac{1}{2}})$	<b>(0.22±0.06)</b>	<b>(86±18)<sup>o</sup></b>	<b>(8.7±3.0)</b>	<b>(49±10)<sup>o</sup></b>	<b>-(1.3±0.4)</b>	<b>-(40±25)<sup>o</sup></b>
$\gamma p \rightarrow \eta p (A^{\frac{3}{2}})$	<b>(1.18±0.20)</b>	<b>-(74±17)<sup>o</sup></b>	<b>-(4.0±2.1)</b>	<b>-(53±12)<sup>o</sup></b>	<b>(12±5)</b>	<b>(14±10)<sup>o</sup></b>
$\gamma p \rightarrow \eta p (\mathbf{E})$	<b>-(0.91±0.15)</b>	<b>-(71±16)<sup>o</sup></b>	<b>-(4.3±1.7)</b>	<b>(35±15)<sup>o</sup></b>	<b>(4.0±1.3)</b>	<b>(6±12)<sup>o</sup></b>
$\gamma p \rightarrow \eta p (\mathbf{M})$	<b>-(0.44±0.10)</b>	<b>-(79±14)<sup>o</sup></b>	<b>-(6.1±2.7)</b>	<b>-(98±14)<sup>o</sup></b>	<b>-(10.0±4)</b>	<b>(17±12)<sup>o</sup></b>

## Conclusion

1. The drop of the  $\gamma p \rightarrow \eta p$  total cross section in the mass region around 1900 MeV together with a strong increase of the  $\eta - \text{prime}$  cross section can be explained naturally with the existence of the  $N(1895)1/2^-$  state.
2. The combined analysis allows us to define the parameters of the  $N(1895)1/2^-$  state with a good precision.
3. The obtained description of the  $\eta$  photoproduction data allows us to describe the  $\gamma n \rightarrow \eta n$  data by fitting only the  $\gamma n$  helicity couplings.
4. The beam asymmetry GRAAL data and the MAMI differential cross section on the  $\eta - \text{prime}$  photoproduction show an anomaly near  $\eta' p$  threshold. This anomaly can be explained with a presence of a very narrow state. However the data should be carefully checked and measurements of other polarization observables are greatly welcome.

# $\gamma p \rightarrow \eta p$ multipoles



## Amplitudes squared for the $\pi N \rightarrow \eta N$ and $\gamma p \rightarrow \eta p$ reactions

