



# Nucleon resonances in Compton scattering

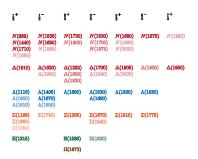
#### **Gernot Eichmann**

IST Lisboa, Portugal

NSTAR 2019 Bonn, Germany

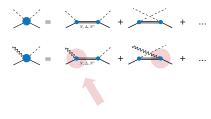
June 12, 2019



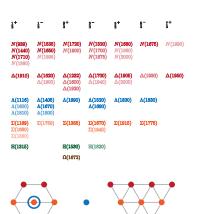




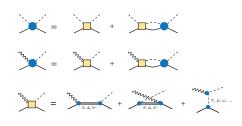
#### Extraction of resonances?



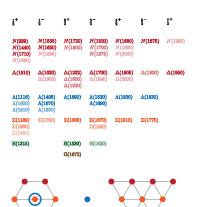
Electromagnetic transition form factors



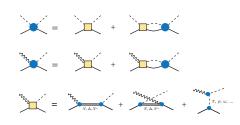
#### Extraction of resonances?



Sato-Lee/EBAC/ANL-Osaka, Dubna-Mainz-Taiwan, Valencia, Jülich-Bonn, GSI. JLab, MAID, SAID, KSU, Giessen, Bonn-Gatchina, JPAC....



#### Extraction of resonances?

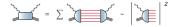


- Lorentz invariance?
- Em. gauge invariance?
- EFT ← QCD?
- What is an "offshell hadron"?

# Compton scattering



#### Structure functions & PDFs in forward limit

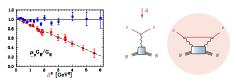


#### Handbag dominance & GPDs in DVCS



#### TPE corrections to form factors

Guichon, Vanderhaeghen, PRL 91 (2003)



#### Proton radius puzzle?

Antonigni et al., 2013, Pohl et al. 2013, Birse, McGovern 2012, Carlson 2015



#### Nucleon polarizabilities

Hagelstein, Miskimen, Pascalutsa, Prog. Part. Nucl. Phys. 88 (2016)



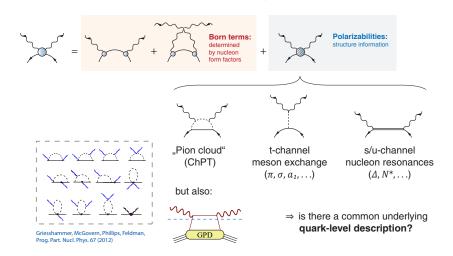
#### Resonances!



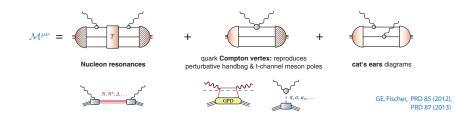


# **Compton scattering**

Compton amplitude = sum of Born terms + 1PI structure part:



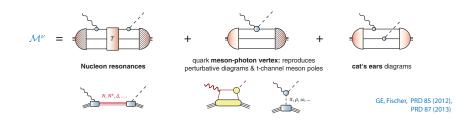
# **Compton scattering**



- Lorentz invariance
- . Crossing symmetry
- Em. gauge invariance
- Chiral symmetry
- Perturbative QCD included
- s, t, u channel poles dynamically generated, no offshell hadrons inside - only quarks & gluons

**But:** consistency important, simple approximations can be dangerous → **challenge** 

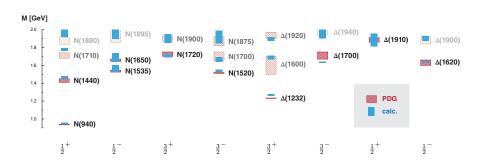
# Meson electroproduction



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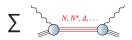
# **Baryons with functional methods**



**Review:** GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 (2016), 1606.09602

- → see **Christian Fischer's** plenary Thursday morning
- → talks by Craig Roberts, José Rodriguez-Quintero, Chen Chen

GE, Ramalho, PRD 98 (2018)



#### Need em. transition FFs

But vertices are half offshell: need 'consistent couplings' Pascalutsa. Timmermans. PRC 60 (1999)

- em gauge invariance:  $Q^{\mu} \Gamma^{\alpha\mu} = 0$
- spin-3/2 gauge invariance:  $k^{\alpha} \Gamma^{\alpha\mu} = 0$
- invariance under point transformations:  $\gamma^{\alpha}\,\Gamma^{\alpha\mu}=0$
- no kinematic dependencies, "minimal" basis

$J^P = \frac{1}{2}^+$	$\frac{3}{2}^{+}$	$\frac{1}{2}$	3 -
N(940)	N(1720)	N(1535)	N(1520)
N(1440)	N(1900)	N(1650)	N(1700)
N(1710)		N(1895)	N(1875)
N(1880)			
$\Delta(1910)$	$\Delta(1232)$	$\Delta(1620)$	$\Delta(1700)$
	$\Delta(1600)$	$\Delta(1900)$	$\Delta(1940)$
	$\Lambda(1920)$		

E.g. Jones-Scadron current cannot be used offshell:

$$\begin{split} \Gamma^{\alpha\mu} &\sim \bar{u}^{\alpha}(k) \left[ m^2 \lambda_{-} (G_M^* - G_E^*) \, \varepsilon_{kQ}^{\alpha\mu} \right. \\ &\left. - G_E^* \, \varepsilon_{kQ}^{\alpha\beta} \, \varepsilon_{kQ}^{\beta\mu} - \frac{1}{2} \, G_C^* \, Q^{\alpha} k^{\beta} t_{QQ}^{\beta\mu} \right] u(k') \end{split}$$

$$t_{AB}^{lphaeta} = A \cdot B \, \delta^{lphaeta} - B^{lpha} \, A^{eta}$$
  $arepsilon_{AB}^{lphaeta} = \gamma_5 \, arepsilon_{lphaeta\gamma\delta}^{lphaeta\gamma\delta} A^{\gamma} B^{\delta}$ 



#### Compton scattering

Tarrach, Nuovo Cim. A28 (1975) GE, Ramalho, PRD 98 (2018)



#### Meson electroproduction

GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 (2016)



#### Nucleon-to-resonance transition currents

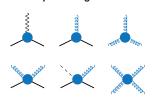
GE, Ramalho, PRD 98 (2018)



#### Light-by-light scattering?

GE, Fischer, Heupel, PRD 92 (2015)

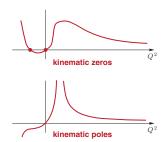
... also n-point functions with quarks & gluons



$$\Gamma^{\mu\nu} \; = \; \sum_{i} \, c_{i} \, K_{i}^{\mu\nu} \; = \; \underbrace{\sum_{i} \, \mathbf{g}_{i} \, G_{i}^{\mu\nu}}_{\mathbf{G}} + \underbrace{\sum_{j} \, \mathbf{f}_{j} \, X_{j}^{\mu\nu}}_{\mathbf{T}} \label{eq:Gamma_problem}$$

**Minimal basis:** neither  $g_i, f_j$  nor  $G_i, X_j$  become singular

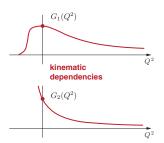
#### Without minimal basis:



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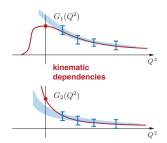
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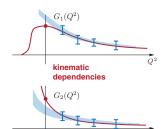
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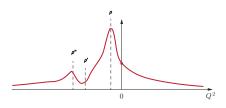
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#### Without minimal basis:



#### With minimal basis:

no kinematic dependencies, only 'physical' poles and cuts!



$$\Gamma^{\mu 
u} = \sum_i c_i K_i^{\mu 
u} = \underbrace{\sum_i g_i G_i^{\mu 
u}}_{\mathbf{G}} + \underbrace{\sum_j f_j X_j^{\mu 
u}}_{\mathbf{T}}$$

**Minimal basis:** neither  $g_i, f_j$  nor  $G_i, X_j$  become singular

Transversality constraints:

$$Q^{\prime\mu} \Gamma^{\mu\nu} = 0$$

$$Q^{\nu} \Gamma^{\mu\nu} = 0$$

$$\Gamma^{\mu 
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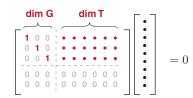
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Transversality constraints:

$$Q^{\prime\mu} \Gamma^{\mu\nu} = 0 \qquad \Rightarrow \qquad Q^{\nu} \Gamma^{\mu\nu} = 0$$



Row-reduced echelon form:



$$\Gamma^{\mu\nu} \; = \; \sum_{i} \, c_{i} \, K_{i}^{\mu\nu} \; = \; \underbrace{\sum_{i} \, g_{i} \, G_{i}^{\mu\nu}}_{\rm G} + \underbrace{\sum_{j} \, f_{j} \, X_{j}^{\mu\nu}}_{\rm T} \label{eq:Gamma_spectrum}$$

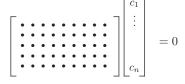
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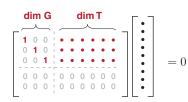
$$\Rightarrow$$

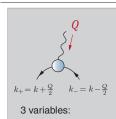


A minimal basis exists, if

- by swapping columns (= renaming basis tensors)
- adding / subtracting rows, multiplying rows with scalars (Gauss-Jordan elimination)

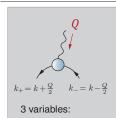
one can find a **row-reduced echelon form**where ::::: is nonsingular in any kinematic limit





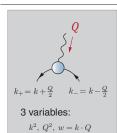
 $k^2$ ,  $Q^2$ ,  $w = k \cdot Q$ 

$$\begin{split} \Gamma^\mu(k,Q) &= c_1\,k^\mu + c_2\,Q^\mu \\ \overline{\Gamma}^\mu(k,Q) &:= \Gamma^\mu(-k,-Q) \stackrel{!}{=} -\Gamma^\mu(k,-Q) \end{split} \quad \text{(charge conjugation)}$$



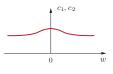
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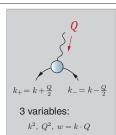
$$\begin{split} \Gamma^{\mu}(k,Q) &= c_1 \, k^{\mu} + c_2 \, (\!\!\! k \cdot Q\!\!\!\! Q) \, Q^{\mu} \\ \overline{\Gamma}^{\mu}(k,Q) &:= \Gamma^{\mu}(-k,-Q) \stackrel{!}{=} -\Gamma^{\mu}(k,-Q) \end{split} \quad \text{(charge conjugation)}$$



$$\Gamma^{\mu}(k,Q) = c_1 k^{\mu} + c_2 (\mathbf{k} \cdot \mathbf{Q}) Q^{\mu}$$

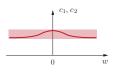
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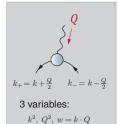




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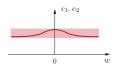


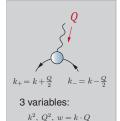


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$$Q^{\mu} \Gamma^{\mu} = c_1 w + c_2 w Q^2 = 0$$



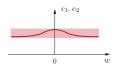


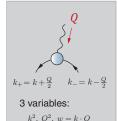
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$$\Rightarrow \left[\begin{array}{cc} w & w Q^2 \end{array}\right] \left[\begin{array}{c} c_1 \\ c_2 \end{array}\right] = 0$$





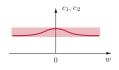
$$\Gamma^{\mu}(k,Q) = c_1 k^{\mu} + c_2 (k \cdot Q) Q^{\mu}$$

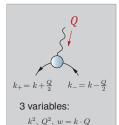
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$$\Rightarrow \left[ \begin{array}{cc} w & w \, Q^2 \end{array} \right] \left[ \begin{array}{c} c_1 \\ c_2 \end{array} \right] = 0$$

$$\Rightarrow \left[ \begin{array}{cc} 1 & Q^2 \end{array} \right] \left[ \begin{array}{c} c_1 \\ c_2 \end{array} \right] = 0$$





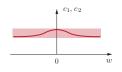
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Transversality:

$$\Rightarrow \begin{bmatrix} w & w Q^2 \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix} = 0$$

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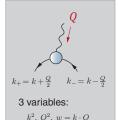
 $Q^{\mu} \Gamma^{\mu} = c_1 w + c_2 w Q^2 = 0$ 



(charge conjugation)

but not

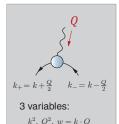
$$\begin{bmatrix} 1 & \frac{1}{Q^2} \end{bmatrix} \begin{bmatrix} c_2 \\ c_1 \end{bmatrix} = 0 \quad !!$$



$$\begin{split} \Gamma^{\mu}(k,Q) &= c_1\,k^{\mu} + c_2\,(\boldsymbol{k}\cdot\boldsymbol{Q})\,Q^{\mu} \\ \overline{\Gamma}^{\mu}(k,Q) &:= \Gamma^{\mu}(-k,-Q) \stackrel{!}{=} -\Gamma^{\mu}(k,-Q) \end{split} \tag{charge conjugation}$$

$$\begin{split} Q^{\mu} \, \Gamma^{\mu} &= c_1 \, w + c_2 \, w \, Q^2 = 0 \\ \\ c_1 &= -c_2 \, Q^2 \quad \Rightarrow \quad \Gamma^{\mu} = -c_2 \, (Q^2 \, k^{\mu} - w \, Q^{\mu}) \\ &= -c_2 \, (Q^2 \, \delta^{\mu\nu} - Q^{\mu} Q^{\nu}) \, k^{\nu} \\ &= -c_2 \, t^{\mu\nu}_{QQ} \, k^{\nu} \\ \\ &\Rightarrow \quad \Gamma^{\mu}(k,Q) = \mathbf{g_1} \, k^{\mu} + \mathbf{f_1} \, t^{\mu\nu}_{QQ} \, k^{\nu} \end{split}$$

$$ightarrow \Gamma^{\mu}(k,Q) = \underbrace{rac{g_1}{g_1}k^{\mu}}_{ extbf{G}} + \underbrace{rac{f_1}{f_2}t^{\mu
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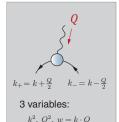
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$$\begin{array}{rcl} c_1 = -c_2\,Q^2 & \Rightarrow & \Gamma^\mu = -c_2\,(Q^2\,k^\mu - w\,Q^\mu) \\ & = -c_2\,(Q^2\,\delta^{\mu\nu} - Q^\mu Q^\nu)\,k^\nu \\ & = -c_2\,t^{\mu\nu}_{OO}\,k^\nu \end{array}$$
 G:

Ward-Takahashi identity only affects G:

$$\begin{split} Q^{\mu} \, \Gamma^{\mu} &= D(k_{+})^{-1} - D(k_{-})^{-1} = g_{1} \, w \\ \Rightarrow & g_{1} = 2 \, \frac{D(k_{+})^{-1} - D(k_{-})^{-1}}{k_{+}^{2} - k_{-}^{2}} = 2 \Delta \end{split} \qquad \qquad \Rightarrow \qquad \begin{split} \Gamma^{\mu}(k,Q) &= \underbrace{g_{1} \, k^{\mu}}_{\mathbf{G}} + \underbrace{f_{1} \, t^{\mu\nu}_{QQ} \, k^{\nu}}_{\mathbf{G}} \end{split}$$



$$\begin{split} \Gamma^{\mu}(k,Q) &= c_1\,k^{\mu} + c_2\,(\boldsymbol{k}\cdot\boldsymbol{Q})\,Q^{\mu} \\ \overline{\Gamma}^{\mu}(k,Q) &:= \Gamma^{\mu}(-k,-Q) \stackrel{!}{=} -\Gamma^{\mu}(k,-Q) \end{split} \tag{charge conjugation}$$

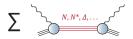
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Transverse-longitudinal separation?

$$\Rightarrow \quad \tilde{g}_1 = \frac{2\Delta}{Q^2} \quad \Rightarrow \quad \tilde{f}_1 = f_1 + \frac{2\Delta}{Q^2} \quad \Rightarrow \quad \text{both kinematically dependent} \\ \text{and singular!}$$



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N(940) N(1440) N(1710) N(1880)	N(1720) N(1900)	N(1535) N(1650) N(1895)	N(1520) N(1700) N(1875)
$\Delta(1910)$	$\Delta$ (1232) $\Delta$ (1600) $\Delta$ (1920)	$\Delta$ (1620) $\Delta$ (1900)	$\Delta$ (1700) $\Delta$ (1940)

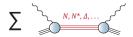
#### Most general offshell vertices satisfying these constraints:

GE. Ramalho, 1806,04579

$$\frac{1}{2}^+ \rightarrow \frac{1}{2}^\pm : \quad \Gamma^\mu = \begin{bmatrix} \mathbf{1} \\ \gamma_5 \end{bmatrix} \sum_{i=1}^8 \mathbf{F_i} T_i^\mu \qquad \begin{cases} \mathbf{f}_{QQ}^\mu \gamma^\nu \\ [\gamma^\mu, Q] \\ \dots \end{cases}$$

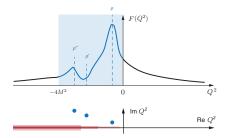
$$\label{eq:continuity} \begin{array}{l} \frac{1}{2}^{+} \rightarrow \frac{3}{2}^{\pm} : \; \Gamma^{\alpha\mu} = \begin{bmatrix} \gamma_{5} \\ 1 \end{bmatrix} \sum_{i=1}^{12} \mathbf{F_{i}} T_{i}^{\alpha\mu} & \begin{cases} \epsilon_{\mathbf{kQ}}^{\alpha\mu} \\ t_{\mathbf{kQ}}^{\alpha} \\ i \epsilon_{\mathbf{kq}}^{\alpha} t_{\mathbf{QQ}}^{\beta\mu} \end{cases}$$

$$\begin{cases} t_{QQ}^{\mu\nu} \gamma^{\nu} \\ [\gamma^{\mu}, Q] \\ \dots \end{cases}$$

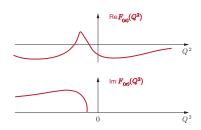


Constraint-free transition FFs: only physical poles and cuts

•  $\rho$  **poles**  $\sim$  monotonous behavior (+ zero crossings for excited states)

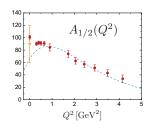


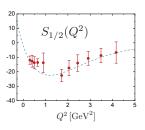
Non-monotonicity at low Q2
 ~ signature for cuts (ρ→ππ, etc.):
 meson cloud





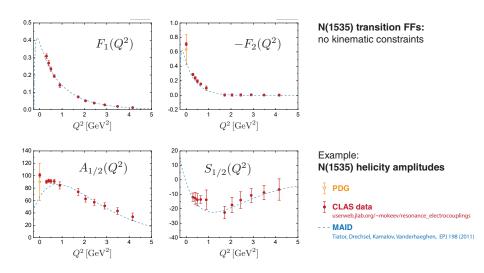
$J^P={1\over 2}^+$	$\frac{3}{2}^{+}$	$\frac{1}{2}$	$\frac{3}{2}$
N(940) N(1440) N(1710) N(1880)	N(1720) N(1900)	N(1535) N(1650) N(1895)	N(1520) N(1700) N(1875)
$\Delta(1910)$	$\Delta$ (1232) $\Delta$ (1600) $\Delta$ (1920)	$\Delta$ (1620) $\Delta$ (1900)	$\Delta$ (1700) $\Delta$ (1940)

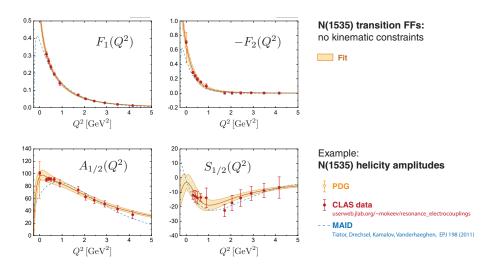




# Example: N(1535) helicity amplitudes

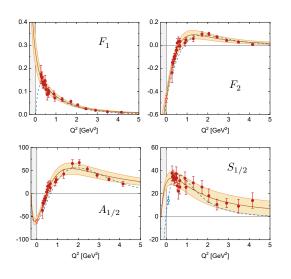
- Y PDG
  - **CLAS** data
  - $userweb.jlab.org/{\sim}mokeev/resonance\_electrocouplings$
- --- MAID
  - Tiator, Drechsel, Kamalov, Vanderhaeghen, EPJ 198 (2011)





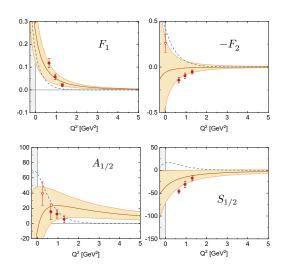


$J^P = \frac{1}{2}^+$	$\frac{3}{2}^{+}$	$\frac{1}{2}$	$\frac{3}{2}^{-}$
N(940) N(1440) N(1710) N(1880)	N(1720) N(1900)	<b>N(1535)</b> <b>N(1650)</b> <i>N</i> (1895)	N(1520) $N(1700)$ $N(1875)$
$\Delta(1910)$	$\Delta(1232)$ $\Delta(1600)$ $\Delta(1920)$	$\Delta$ (1620) $\Delta$ (1900)	$\Delta$ (1700) $\Delta$ (1940)



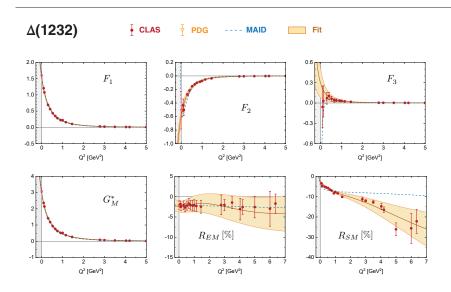
### N(1440)

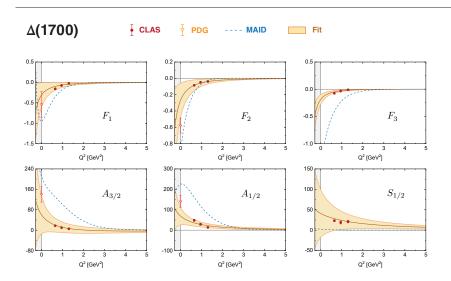


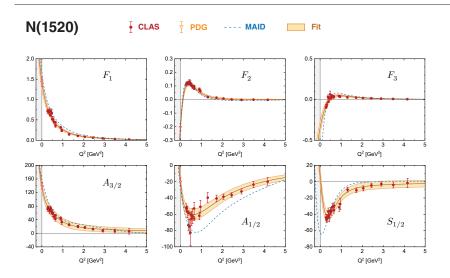


#### Δ(1620)









#### **Kinematics**

$$\sum_{i=1}^{18}$$

$$= \sum c_i(\eta_+, \eta_-, \omega, \lambda) \, \bar{u}(p_f) \, X_i^{\mu\nu}(p, Q, Q') \, u(p_i)$$

#### 18 CFFs

4 kinematic variables:

$$\eta_{+} = rac{Q^{2} + Q'^{2}}{2m^{2}}$$
 $\eta_{-} = rac{Q \cdot Q'}{m^{2}}$ 
 $Q^{2} - Q'^{2}$ 

$$\omega = \frac{Q^2 - Q'^2}{2m^2}$$
$$\lambda = -\frac{p \cdot Q}{m^2}$$

#### 18 Compton tensors, form minimal basis

- · systematic derivation
- similar to Tarrach basis Tarrach, Nuovo Cim. A28 (1975)

$$X_i' = U_{ij} X_j$$
,  $\det U = const$ .

· CFFs free of kinematics

$$X_1^{\mu\nu} = \frac{1}{m^4} t_{Q'p}^{\mu\alpha} t_{pQ}^{\alpha\nu},$$
  
 $X_2^{\mu\nu} = \frac{1}{m^2} t_{Q'Q}^{\mu\nu},$ 

$$X_3^{\mu\nu} = \frac{1}{m^4} t_{Q'Q'}^{\mu\alpha} t_{QQ}^{\alpha\nu},$$

$$X_4^{\mu\nu} = \frac{1}{m^6} t_{Q'Q'}^{\mu\alpha} p^{\alpha} p^{\beta} t_{QQ}^{\beta\nu},$$
  
 $X_5^{\mu\nu} = \frac{\lambda}{m^4} \left( t_{Q'Q'}^{\mu\alpha} t_{pQ}^{\alpha\nu} + t_{Q'p}^{\mu\alpha} t_{QQ}^{\alpha\nu} \right),$ 

$$X_6^{\mu\nu} = \frac{1}{m^2} \epsilon_{Q'Q}^{\mu\nu},$$

$$X_7^{\mu\nu} = \frac{1}{im^3} \left( t_{Q'Q'}^{\mu\alpha} \varepsilon_{\gamma Q}^{\alpha\nu} - \varepsilon_{Q'\gamma}^{\mu\alpha} t_{QQ}^{\alpha\nu} \right),$$

$$X_8^{\mu\nu} = \frac{\omega}{im^3} \left( t_{Q'Q'}^{\mu\alpha} \varepsilon_{\gamma Q}^{\alpha\nu} + \varepsilon_{Q'\gamma}^{\mu\alpha} t_{QQ}^{\alpha\nu} \right),$$



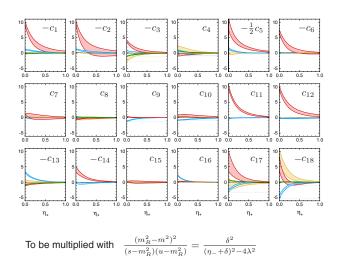






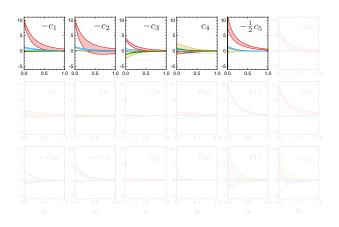


GF. Ramalho. PRD 98 (2018)

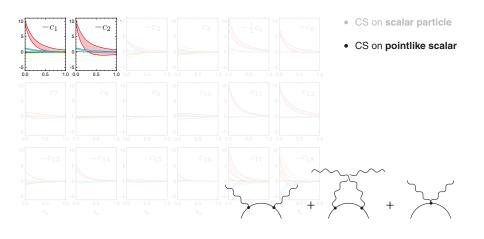


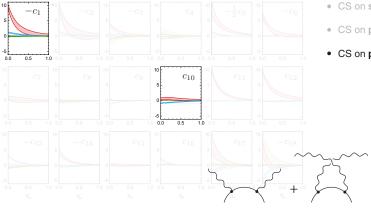




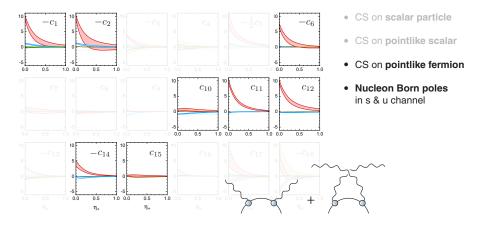


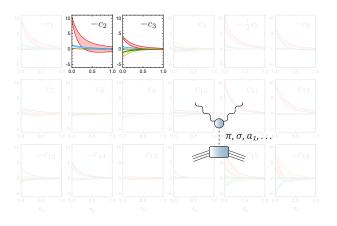
• CS on scalar particle



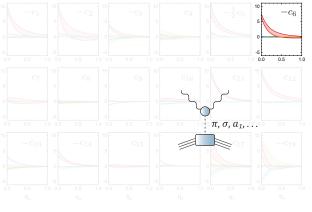


- CS on scalar particle
- CS on pointlike scalar
- CS on pointlike fermion

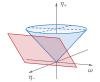




- CS on scalar particle
- CS on pointlike scalar
- CS on pointlike fermion
- Nucleon Born poles in s & u channel
- Scalar pole in t channel

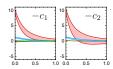


- CS on scalar particle
- CS on pointlike scalar
- CS on pointlike fermion
- Nucleon Born poles in s & u channel
- Scalar pole in t channel
- **Pion pole** in t channel  $(\pi^0 \to \gamma^* \gamma^*)$



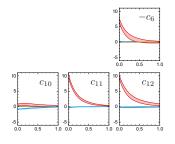
GE, Fischer, Weil, Williams, PLB 774 (2017)

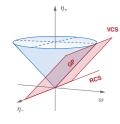
#### **Polarizabilities**



#### Scalar polarizabilities:

$$\left[\begin{array}{c} \alpha+\beta\\ \beta \end{array}\right] = -\frac{\alpha_{\rm em}}{m^3} \left[\begin{array}{c} c_1\\ c_2 \end{array}\right]$$



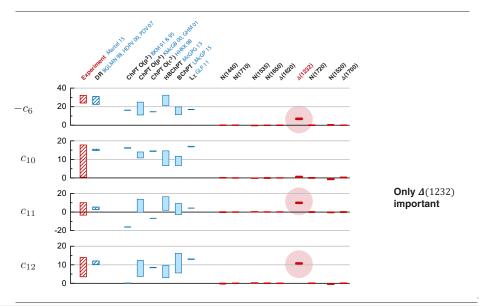


#### Spin polarizabilities:

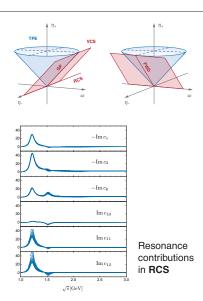
$$\left[ \begin{array}{c} \gamma_{E1E1} \\ \gamma_{M1M1} \\ \gamma_{E1M2} \\ \gamma_{M1E2} \end{array} \right] = \frac{\alpha_{\rm em}}{2m^4} \left[ \begin{array}{c} c_6 + 4c_{11} - 4c_{12} \\ -c_6 - 2c_{10} + 4c_{12} \\ c_6 + 2c_{10} \\ -c_6 \end{array} \right]$$

$$\begin{bmatrix} \gamma_0 \\ \gamma_{\pi} \end{bmatrix} = -\frac{2\alpha_{\rm em}}{m^4} \begin{bmatrix} c_{11} \\ c_6 + c_{10} + c_{11} - 2c_{12} \end{bmatrix}$$

# Spin polarizabilities



### **Summary**

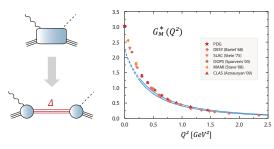


- Compton scattering: minimal basis, constraint-free CFFs
- Electromagnetic transition currents: minimal bases, constraint-free TFFs
- Fits of experimental TFFs
- Nucleon resonances in Compton scattering: derived in general kinematics
- Only Δ(1232) and N(1520) relevant for polarizabilities

GE, Ramalho, PRD 98 (2018)

## **Backup slides**

## Nucleon- $\Delta$ - $\gamma$ transition



- Magnetic dipole transition (G<sup>\*</sup><sub>M</sub>) dominant: quark spin flip (s wave). "Core + 25% pion cloud"
- Electric & Coulomb quadrupole ratios small & negative, encode deformation.
   Reproduced without pion cloud: OAM from p waves!
   GE, Nicmorus, PRD 85 (2012)

