

# The spectrum of hyperon resonances from a partial wave analysis of K-p scattering data

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# Quark model

The flavour and spin can be combined in spin-flavour SU(6):

$$6 \otimes 6 \otimes 6 = 56_S \oplus 70_M \oplus 70_M \oplus 20_A$$

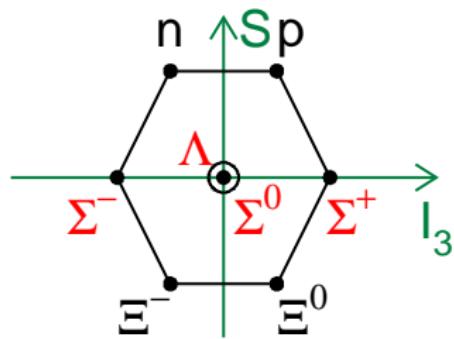
where

$$56 = {}^410 \oplus {}^28$$

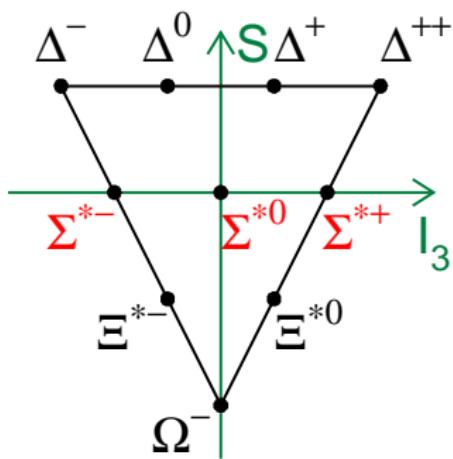
$$70 = {}^210 \oplus {}^48 \oplus {}^28 \oplus {}^21$$

$$20 = {}^28 \oplus {}^41$$

## Octet



## Decuplet



# Quark model classification. Ground and 1st excitation shell

two states on ground shell

2  $\Lambda$  as singlets 6  $\Lambda$  and 6  $\Sigma$  states as octet members on a 1st shell

$(D, L_N^P)$	$S$	$J^P$	Octet members			Singlets
$(56, 0_0^+)$	$\frac{1}{2}$	$\frac{1}{2}^+$	$N(939)$	$\Lambda(1116)^{4*}$	$\Sigma(1193)^{4*}$	-
$(70, 1_1^-)$	$\frac{1}{2}$	$\frac{1}{2}^-$	$N(1535)$	$\Lambda(1670)^{4*}$	$\Sigma(1620)^{2*}$	$\Lambda(1405)^{4*}$
	$\frac{3}{2}$	$\frac{3}{2}^-$	$N(1520)$	$\Lambda(1690)^{4*}$	$\Sigma(1670)^{4*}$	$\Lambda(1520)^{4*}$
		$\frac{1}{2}^-$	$N(1650)$	$\Lambda(1800)^{3*}$	$\Sigma(1750)^{3*}$	-
		$\frac{3}{2}^-$	$N(1700)$			-
		$\frac{5}{2}^-$	$N(1675)$	$\Lambda(1830)^{4*}$	$\Sigma(1775)^{4*}$	-

3  $\Sigma$  states are expected as decuplet members

$(D, L_N^P)$	$S$	$J^P$	Decuplet members	
$(56, 0_0^+)$	$\frac{3}{2}$	$\frac{3}{2}^+$	$\Delta(1232)$	$\Sigma(1385)^{4*}$
$(70, 1_1^-)$	$\frac{1}{2}$	$\frac{1}{2}^-$	$\Delta(1620)$	
	$\frac{1}{2}$	$\frac{3}{2}^-$	$\Delta(1700)$	

# Quark model classification. 2nd excitation shell

$(D, L_N^P)$	$S$	$J^P$	Octet members			Singlets
$(56, 0_2^+)$	$\frac{1}{2}$	$\frac{1}{2}^+$	$N(1440)$	$\Lambda(1600)^3*$	$\Sigma(1660)^3*$	-
$(70, 0_2^+)$	$\frac{1}{2}$	$\frac{1}{2}^+$	$N(1710)$	$\Lambda(1810)^3*$	$\Sigma(1770)^1*$	-
$(56, 2_2^+)$	$\frac{1}{2}$	$\frac{3}{2}^+$	$N(1720)$	$\Lambda(1890)^4*$	$\Sigma(1840)^1*$	-
		$\frac{5}{2}^+$	$N(1620)$	$\Lambda(1820)^4*$	$\Sigma(1915)^4*$	-
$(70, 2_2^+)$	$\frac{1}{2}$	$\frac{3}{2}^+$	$N(1900)$			-
		$\frac{5}{2}^+$	$N(1860)$			-
	$\frac{3}{2}$	$\frac{1}{2}^+$	$N(1880)$			-
		$\frac{3}{2}^+$	$N(1960)$		$\Sigma(2080)^2*$	-
		$\frac{5}{2}^+$	$N(2000)$	$\Lambda(2110)^3*$	$\Sigma(2070)^1*$	-
		$\frac{7}{2}^+$	$N(1990)$	$\Lambda(2020)^1*$	$\Sigma(2030)^4*$	-
$(20, 1_2^+)$	$\frac{1}{2}$	$\frac{1}{2}^+$	$N(2100)$			-
		$\frac{3}{2}^+$	$N(2040)$		$\Sigma(2080)^2*$	-

# Motivation

- ① There are a lot missing states in hyperon spectrum in comparison with  $N$  and  $\Delta$  spectrum in frame of quark model classification
- ② The investigation of hyperon spectrum expands understanding of baryon properties and their classification.
- ③ BnGa approach allow to include near all experimental data in combine analysis, specially to include reaction with three-particle final states

# Scattering amplitude in BnGa approach

- ① Energy dependent approach
- ② Possibility to combine all data in one analysis
- ③ K-matrix satisfies unitarity condition

$$A(s, t) = \sum_{IJN} C_I Q_{JN}(s, t) A_{IJN}(s),$$

- ①  $C_I$  are the Clebsch-Gordan coefficients
- ②  $Q_{JN}$  tensors describe the angular dependent part of the partial wave amplitudes.
- ③  $A_{IJN}(s)$  are partial wave amplitudes (Breit-Wigner function and K-Matrix)

Non-resonance contributions are described by constants in the K-matrix and by amplitudes for t and u-channel exchanges

## Data base of $K^- p$ reactions

Total mass range: (1.46 - 2.3) GeV

Differential cross sections  $d\sigma/d\Omega$  (16316 points)

$K^- p \rightarrow K^- p$	(12, 5170)	$K^- p \rightarrow K^0 n$	(11, 3445)
$K^- p \rightarrow \pi^0 \Lambda$	(11, 2478)	$K^- p \rightarrow \eta \Lambda$	(2, 160)
$K^- p \rightarrow \pi^0 \Sigma^0$	(5, 581)	$K^- p \rightarrow \pi^\mp \Sigma^\pm$	(8, 4177)
$K^- p \rightarrow K^0 \Xi^+$	(5, 305)		

Data on the polarization observable  $P$  (2818 points)

$K^- p \rightarrow K^0 n$	(5, 1180)	$K^- p \rightarrow \pi^0 \Lambda$	(7, 892)
$K^- p \rightarrow \pi^0 \Sigma^0$	(1, 124)	$K^- p \rightarrow \pi^- \Sigma^+$	(5, 593)
$K^- p \rightarrow K^0 \Xi^+$	(1, 29)		

Data three-body final states (3711 points)

$K^- p \rightarrow \omega \Lambda$	(3, 300)	$K^- p \rightarrow \bar{K} \Delta(1232)$	(2, 667)
$K^- p \rightarrow \pi^\mp \Sigma^\pm(1385)$	(2, 899)	$K^- p \rightarrow \pi^0 \Lambda(1520)$	(4, 1011)
$K^- p \rightarrow K^{*0} n$	(2, 371)	$K^- p \rightarrow K^{*-} p$	(2, 463)

Data event-by-event in an event-based likelihood fit

$$K^- p \rightarrow 2\pi^0 \Lambda \quad (1, 26513) \quad | \quad K^- p \rightarrow 2\pi^0 \Sigma \quad (1, 3286)$$

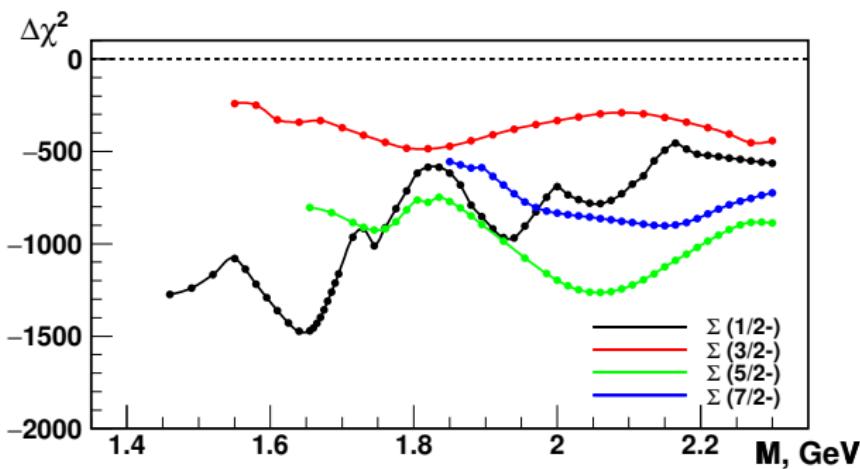
# Step 1: Hyperon set for the primary fit.

	$J^P$	Status	Mass	Width
$\Lambda(1405)$	$1/2^-$	****	$1405^{+1.3}_{-1.0}$	$50.5 \pm 2.0$
$\Lambda(1670)$	$1/2^-$	****	1660 – 1680	25 – 50
$\Lambda(1800)$	$1/2^-$	***	1720 – 1850	200 – 400
$\Lambda(1520)$	$3/2^-$	****	$1519.5 \pm 1.0$	$15.6 \pm 1.0$
$\Lambda(1690)$	$3/2^-$	****	1685 – 1695	50 – 70
$\Lambda(1830)$	$5/2^-$	****	1810 – 1830	60 – 110
$\Lambda(2100)$	$7/2^-$	****	2090 – 2110	100 – 250
$\Lambda(1600)$	$1/2^+$	***	1560 – 1700	50 – 250
$\Lambda(1810)$	$1/2^+$	***	1750 – 1850	50 – 250
$\Lambda(1890)$	$3/2^+$	****	1850 – 1910	60 – 200
$\Lambda(1820)$	$5/2^+$	****	1815 – 1825	70 – 90
$\Lambda(2110)$	$5/2^+$	***	2090 – 2140	150 – 250
	$J^P$	Status	Mass	Width
$\Sigma(1750)$	$1/2^-$	***	1730 – 1800	60 – 160
$\Sigma(1670)$	$3/2^-$	****	1665 – 1685	40 – 80
$\Sigma(1940)$	$3/2^-$	***	1900 – 1950	150 – 300
$\Sigma(1775)$	$5/2^-$	****	1770 – 1780	105 – 135
$\Sigma(1660)$	$1/2^+$	***	1630 – 1690	40 – 200
$\Sigma(1385)$	$3/2^+$	****	$1382.80 \pm 0.35$	$36.0 \pm 0.7$
$\Sigma(1915)$	$5/2^+$	****	1900 – 1935	80 – 160
$\Sigma(2030)$	$7/2^+$	****	2025 – 2040	150 – 200

Primary Breit-Wigner fit  
includes set of hyperons with  
**4 and 3 stars RPP rating  
only**  
Masses and widths are  
allowed to vary **within the  
limits quoted in the RPP**

## Step 2: Mass scan procedure

- to take a primary solution with well establish states
- to add resonance with fix mass to partial wave and to fit
- to repeat it for whole investigated mass region and all partial waves



$$J^P = 1/2^-$$

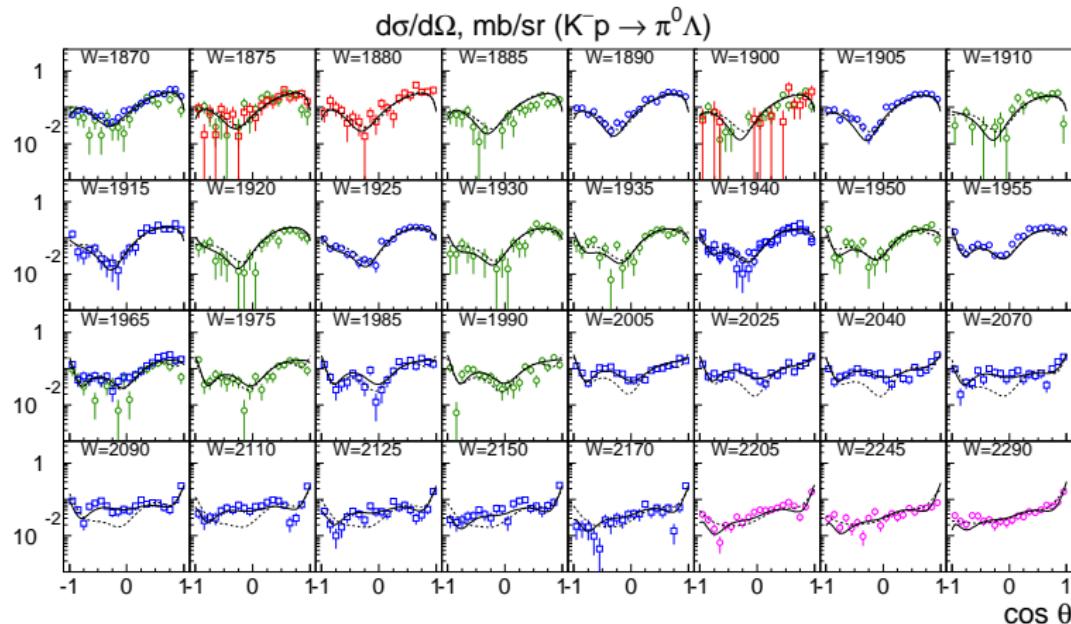
$$J^P = 3/2^-$$

$$J^P = 5/2^-$$

$$J^P = 7/2^-$$

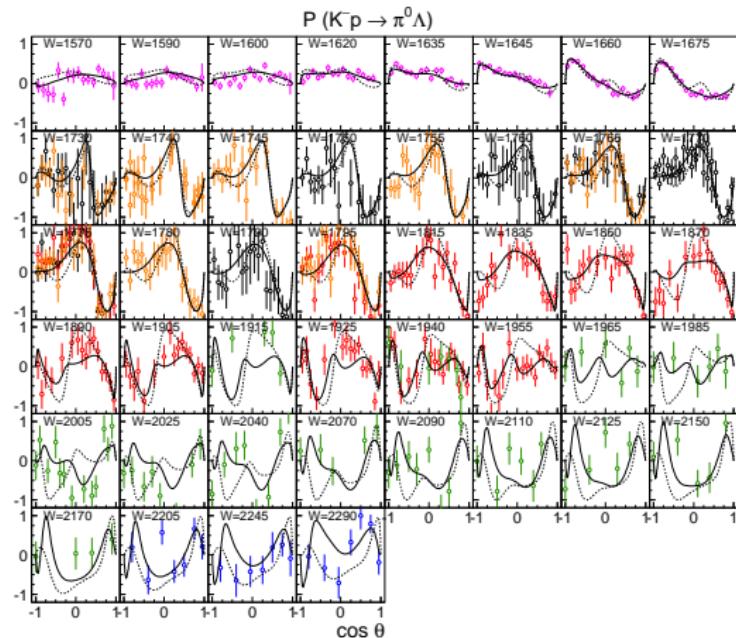
- to establish state which gives essential fit improving
- to repeat mass scan for all partial waves with this additional state
- to continue this procedure up to not significant data description improvement - mass scans are flat

# Differential cross section for $K^- p \rightarrow \pi^0 \Lambda$



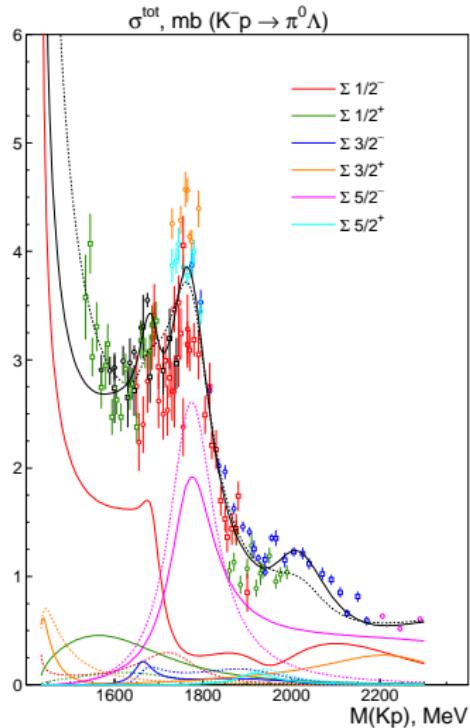
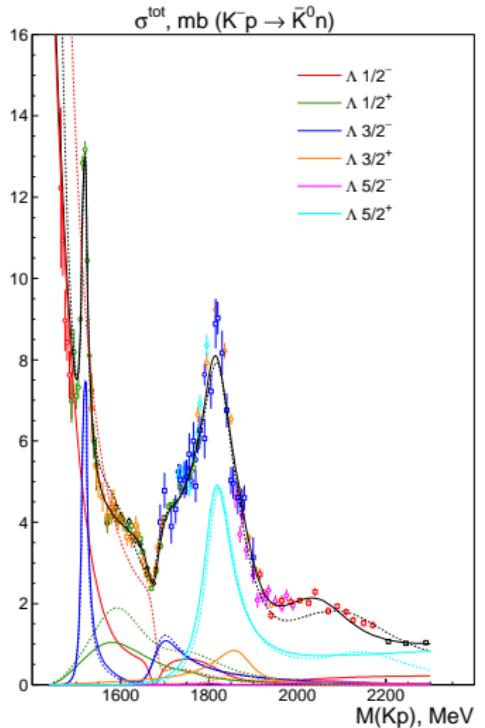
dashed line - primary fit with set of well established hyperons  
solid line - final fit

$$K^- p \rightarrow \pi^0 \Lambda$$



dashed line - primary fit with set of well established hyperons  
 solid line - final fit

# Total cross section for $K^- p \rightarrow \bar{K}^0 n$ and $K^- p \rightarrow \pi^0 \Lambda$



dashed line - primary fit with set of well established hyperons  
solid line - final fit

# $\chi^2$ difference between primary and final fits

## Differential cross sections

$K^- p \rightarrow K^- p$	(2.203, 1.802)	$K^- p \rightarrow K^0 n$	(1.997, 1.545)
$K^- p \rightarrow \pi^0 \Lambda$	(2.135, 1.658)	$K^- p \rightarrow \eta \Lambda$	(2.756, 1.500)
$K^- p \rightarrow \pi^0 \Sigma^0$	(2.169, 1.960)	$K^- p \rightarrow \pi^\mp \Sigma^\pm$	(5.196, 4.190)
$K^- p \rightarrow K^0 \Xi^0$	(3.052, 2.549)		

## Data on the polarization observable $P$

$K^- p \rightarrow K^0 n$	(1.867, 1.406)	$K^- p \rightarrow \pi^0 \Lambda$	(1.901, 1.244)
$K^- p \rightarrow \pi^0 \Sigma^0$	(4.613, 2.411)	$K^- p \rightarrow \pi^- \Sigma^+$	(2.165, 2.094)
$K^- p \rightarrow K^0 \Xi^0$	(2.955, 2.650)		

## Data on the $K^- p$ induced reactions with three-body final states

$K^- p \rightarrow 2\pi^0 \Lambda$	$(\delta\chi^2 = -244^a)$	$K^- p \rightarrow 2\pi^0 \Sigma$	$(\delta\chi^2 = -498^a)$
$K^- p \rightarrow \omega \Lambda$	(1.497, 1.027)	$K^- p \rightarrow \bar{K} \Delta(1232)$	(1.568, 1.301)
$K^- p \rightarrow \pi^\mp \Sigma^\pm(1385)3/2^+$	(3.600, 2.406)	$K^- p \rightarrow \pi^0 \Lambda(1520)3/2^-$	(2.704, 1.390)
$K^- p \rightarrow K^{*0} n$	(2.776, 2.647)	$K^- p \rightarrow K^{*-} p$	(2.218, 1.948)

# The spectrum of $\Lambda$ hyperons with negative parity

	Mass	Width	$\Delta\chi^2$	Status
$\Lambda(1405)1/2^-$	$1420 \pm 3$ $1405.1^{+1.3}_{-1.0}$	$46 \pm 4$ $50.5 \pm 2.0$	4070	****
$\Lambda(1670)1/2^-$	$1677 \pm 2$ 1660 to 1680	$33 \pm 4$ 25 to 50	3610	****
$\Lambda(1800)1/2^-$	$1811 \pm 10$ 1720 to 1850	$209 \pm 18$ 200 to 400	1896	***
$\Lambda(2000)1/2^-$	$2085 \pm 14$ $\approx 2060$	$428 \pm 16$ 100 to 300	845	*
$\Lambda(1520)3/2^-$	$1518.5 \pm 0.5$ $1519.5 \pm 1.0$	$15.7 \pm 1.0$ $15.6 \pm 1.0$	>10 000	****
$\Lambda(1690)3/2^-$	$1689 \pm 3$ 1685 to 1695	$75 \pm 5$ 50 to 70	>10 000	****
$\Lambda(1830)5/2^-$	$1821 \pm 3$ 1810 to 1830	$64 \pm 7$ 60 to 110	1790	***
$\Lambda(2080)5/2^-$	$2082 \pm 13$ -	$181 \pm 29$ -	770	** new
$\Lambda(2100)7/2^-$	$2090 \pm 15$ 2090 to 2110	$290 \pm 30$ 100 to 250	5412	****

# The spectrum of $\Lambda$ hyperons with positive parity

	Mass	Width	$\Delta\chi^2$	Status
$\Lambda(1600)1/2^+$	$1605 \pm 8$ 1560 to 1700	$245 \pm 15$ 50 to 250	$>10\,000$	****
$\Lambda(1810)1/2^+$	$1773 \pm 5$ 1750 to 1850	$36 \pm 6$ 50 to 250	46	*
$\Lambda(1890)3/2^+$	$1873 \pm 5$ 1850 to 1910	$103 \pm 10$ 60 to 200	4480	****
$\Lambda(2070)3/2^+$	$2070 \pm 24$ -	$370 \pm 50$ -	1144	** new
$\Lambda(1820)5/2^+$	$1822 \pm 4$ 1815 to 1825	$80 \pm 8$ 70 to 90	$>10\,000$	****
$\Lambda(2110)5/2^+$	$2086 \pm 12$ 2090 to 2140	$274 \pm 25$ 150 to 250	1418	**
				***

# The spectrum of $\Sigma$ hyperons with negative parity

	Mass	Width	$\Delta\chi^2$	Status
$\Sigma(1620)1/2^-$	$1681 \pm 6$ $\approx 1620$	$40 \pm 12$ 10 to 400	386	*
$\Sigma(1750)1/2^-$	$1692 \pm 11$ 1730 to 1800	$208 \pm 18$ 60 to 160	3032	****
$\Sigma(1900)1/2^-$	$1938 \pm 12$ $1900 \pm 21$	$155 \pm 30$ $191 \pm 47$	1500	***
$\Sigma(2000)1/2^-$	$2165 \pm 23$ $\approx 2000$	$320^{+300}_{-60}$ 100 to 400	1612	**
$\Sigma(1670)3/2^-$	$1665 \pm 3$ 1665 to 1685	$54 \pm 6$ 40 to 80	5894	****
$\Sigma(1860)3/2^-$	$1878 \pm 12$ -	$224 \pm 25$ -	1708	*** new
$\Sigma(1940)3/2^-$	$2005 \pm 14$ 1900 to 1950	$178 \pm 23$ 150 to 300	446	*
$\Sigma(1775)5/2^-$	$1776 \pm 4$ 1770 to 1780	$124 \pm 8$ 105 to 135	>10 000	****
$\Sigma(2100)7/2^-$	$2146 \pm 17$ $\approx 2100$	$260 \pm 40$ 50 to 150	668	*

# The spectrum of $\Sigma$ hyperons with positive parity

	Mass	Width	$\Delta\chi^2$	Status
$\Sigma(1660)1/2^+$	$1665 \pm 20$ 1630 to 1690	$300^{+140}_{-40}$ 40 to 200	1870	**** ***
$\Sigma(1385)3/2^+$	1385 1383.7 $\pm$ 1.0	36 36 $\pm$ 5		****
$\Sigma(2230)3/2^+$	2240 $\pm$ 27	345 $\pm$ 50	1200	** new
$\Sigma(1915)5/2^+$	1918 $\pm$ 6 1900 to 1935	102 $\pm$ 12 80 to 160	2002	****
$\Sigma(2030)7/2^+$	2032 $\pm$ 6 2025 to 2040	177 $\pm$ 12 150 to 200	2856	****

## Non established RPP states

In the mass range below 2200 MeV, we find no any evidence for:

- ① three “bumps”:  $\Sigma(1480)$ ,  $\Sigma(1670)$ , and  $\Sigma(1690)$ ,
- ② eight states with  $1^*$ :  $\Lambda(1710)1/2^+$ ,  $\Lambda(2020)7/2^+$ ,  $\Lambda(2050)3/2^-$   
 $\Sigma(1580)3/2^-$ ,  $\Sigma(1730)3/2^+$ ,  $\Sigma(1770)1/2^+$ ,  $\Sigma(1940)3/2^+$ ,  $\Sigma(2070)5/2^+$
- ③ three states with  $2^*$ :  $\Sigma(1560)$ ,  $\Sigma(1880)1/2^+$ ,  $\Sigma(2080)3/2^+$ .

# Comparison with the Bonn quark model: $\Lambda$ sector

$\Lambda^*$ resonances						
$J^\pi$	RPP	BnGa	$M_{QM}$	$^{21}[70]$	$^{28}[70]$	$^{48}[70]$
$\frac{1}{2}^-$	$1405.1^{+1.3}_{-1.0}$	$1422 \pm 3$	1524	<u>69.4</u>	26.0	0.3
	$1519.5 \pm 1.0$	$1518.5 \pm 0.5$	1508	<u>77.7</u>	18.7	0.1
$\frac{1}{2}^-$	$1670 \pm 10$	$1677 \pm 2$	1630	29.2	<u>61.6</u>	2.1
	$1690 \pm 5$	$1689 \pm 3$	1662	20.1	<u>72.0</u>	2.2
$\frac{1}{2}^-$	$1800^{+50}_{-80}$	$1811 \pm 10$	1816	0.1	3.1	<u>94.9</u>
	-	-	1775	0.4	1.5	<u>96.1</u>
	$1830^{+0}_{-20}$	$1821 \pm 3$	1828	0.0	0.0	<u>99.0</u>
	-	$2082 \pm 13$	$2080$	large		
$\frac{7}{2}^-$	$2100 \pm 10$	$2090 \pm 15$	2090	large		
$J^\pi$	RPP	BnGa	$M_{QM}$	$^{21}[70]$	$^{28}[56]$	$^{28}[70]$
$\frac{1}{2}^+$	$1600^{+100}_{-40}$	$1605 \pm 8$	1677	0.3	<u>88.4</u>	3.0
	$1810^{+40}_{-60}$	$1773 \pm 7$	90% $^{21}[70]$	1747 or 84% $^{28}[70]$	1898	
$\frac{3}{2}^+$	$1890^{+20}_{-40}$	$1872 \pm 5$	1823	9.9	<u>60.0</u>	28.2
	$1820 \pm 5$	$1822 \pm 4$	1834	28.3	<u>57.8</u>	12.2
$\frac{3}{2}^+$	-	$2070 \pm 24$	$1952$	<u>84.0</u>	3.8	7.6
	$2110^{+30}_{-20}$	$2086 \pm 12$	1999	<u>84.1</u>	4.5	8.9

# Comparison with the Bonn quark model: $\Sigma$ sector

$\Sigma^*$ resonances						
$J^\pi$	RPP	BnGa	$M_{QM}$	$^28[70]$	$^48[70]$	$^410[70]$
$\frac{1}{2}^-$ $\sim 1620$	$1681 \pm 6$	$1628$	<u>87.4</u>	2.3	3.4	
$\frac{3}{2}^-$	$1670^{+15}_{-5}$	$1665 \pm 3$	1669	<u>89.0</u>	1.2	3.4
$\frac{1}{2}^-$	$1750^{+50}_{-20}$	$1692 \pm 11$	1771	2.9	<u>94.6</u>	1.1
$\frac{3}{2}^-$	-	-	1728	0.1	<u>82.7</u>	16.0
$\frac{5}{2}^-$	$1775 \pm 5$	$1776 \pm 4$	1770	0.0	<u>99.0</u>	0.0
$\frac{1}{2}^-$	$\sim 1900$	$1938 \pm 12$	1798	2.8	1.7	<u>94.4</u>
$\frac{3}{2}^-$	$1940^{+10}_{-40}$	$2005 \pm 14$	1781	4.4	15.0	<u>79.3</u>
$J^\pi$	RPP	BnGa	$M_{QM}$	$^28[56]$	$^48[70]$	$^410[56]$
$\frac{1}{2}^-$	-	$2165 \pm 23$	2111	large		
$\frac{3}{2}^-$	-	-	2139	large		
$\frac{1}{2}^+$	$1660 \pm 30$	$1665 \pm 20$	1628	<u>87.4</u>	2.3	5.4
$\frac{3}{2}^+$	$\sim 1840$	-	1896	<u>73.9</u>	22.2	0.0
$\frac{5}{2}^+$	$1915^{+20}_{-15}$	$1917 \pm 6$	1956	<u>77.8</u>	18.2	0.0
$\frac{7}{2}^+$	$2030^{+10}_{-5}$	$2032 \pm 6$	2070		29.4	<u>69.6</u>
$\frac{3}{2}^+$		$2240 \pm 27$				

# Conclusion

- ① Combine fit near all available experimental data
- ② 4 new resonances are proposed:  $\Lambda(2080)5/2^-$ (\*\*),  $\Lambda(2070)3/2^+$ (\*\*),  
 $\Sigma(1860)3/2^-$ (\*\*\*) and  $\Sigma(2230)3/2^+$ (\*\*)
- ③ 5 resonances with one and two stars are seen:  $\Lambda(2000)1/2^-$ ,  
 $\Sigma(1620)1/2^-$ (\*),  $\Sigma(1900)1/2^-$ (\*\*\*),  $\Sigma(2000)1/2^-$ (\*),  
 $\Sigma(2100)7/2^-$ (\*)
- ④ We did not find evidence for 3 “bumps” and 11 resonances
- ⑤ Need new experimental data on  $K^-p$  scattering with polarized target.