

# EFT tools & validation

## part I – SMEFT LO

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VILLUM FONDEN



- fundamental assumptions:
- ▶ new physics nearly decoupled:  $\Lambda \gg (v, E)$
  - ▶ at the accessible scale: **SM** fields + symmetries

☛ a Taylor expansion in canonical dimensions ( $v/\Lambda$  or  $E/\Lambda$ ):

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \frac{1}{\Lambda^3} \mathcal{L}_7 + \frac{1}{\Lambda^4} \mathcal{L}_8 + \dots$$

$$\mathcal{L}_n = \sum_i C_i \mathcal{O}_i^{d=n} \quad C_i \text{ free parameters (Wilson coefficients)}$$

$\mathcal{O}_i$  invariant operators that form  
a complete, non redundant basis

# More than a parameterization

by construction the SMEFT is always  
**a valid QFT** and  
**a valid description of Nature**  
(assuming SM sym + fields and at  $E \ll \Lambda$ )

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- ▶ we can do: gauge invariant calculations, loops, RGE, etc.  
systematically improvable with higher orders in loop+EFT expansions

(Alonso), Jenkins, Manohar, Trott 1308.2627, 1310.4838, 1312.2014  
Grojean, Jenkins, Manohar, Trott 1301.2588  
Elias-Miró, Grojean, Gupta, Marzocca 1312.2928  
Alonso, Chang, Jenkins, Manohar, Shotwell 1405.0486  
Ghezzi, Gomez-Ambrosio, Passarino, Uccirati 1505.03706  
Pruna, Signer 1408.3565  
Hartmann, (Shepherd), Trott 1505.02646, 1507.03568, 1611.09879  
Gauld, Pecjak, Scott 1512.02508  
Deutschmann, Duhr, Maltoni, Vryonidou 1708.00460  
Dawson et al 1801.01136, 1807.11504, 1808.05948, 1812.00214  
Dedes, Paraskevas, Rosiek, Suxho, Trifyllis 1805.00302  
...

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- ▶ we can do: gauge invariant calculations, loops, RGE, etc.  
systematically improvable with higher orders in loop+EFT expansions

- ▶ we can set an ambitious & general **goal**:

use it as a self-consistent theory and **measure** its parameters

- a solid, systematic **probe of NP**, able to capture *any* BSM effect
- a truly **universal language** for future data interpretation

# The need for a global approach

goal:

use the SMEFT as a self-consistent theory and  
measure its parameters

- 👉 necessary to retain **all** the operators.  
generally not possible to make a selection *a priori*.  
correspondence to anomalous couplings is *basis dependent*.

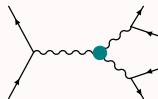
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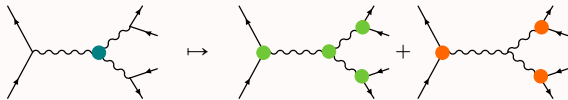
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BSM model  $\rightarrow W_{\mu\nu}^a D^\mu H^\dagger \sigma^a D^\nu H$  affecting



Going to Warsaw basis via EOM:

$W_{\mu\nu}^a D^\mu H^\dagger \sigma^a D^\nu H \mapsto Q_{HW}, Q_{HWB}, Q_{Hq}^{(3)}, Q_{Hl}^{(3)} + \text{Higgs ops.}$



Example

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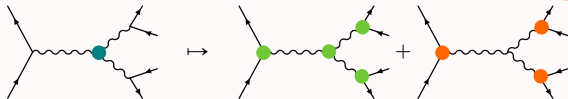
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neglecting operators contributing to  $Vff$  formally spoils the equivalence

↓  
different results in the two bases



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idea: sequential strategy

- ▶ LO, resonance dominated processes
- ▶ + tails  $((\bar{\psi}\psi)(\bar{\psi}\psi))$
- ▶ + NLO ↪ see talk by Eleni

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focus of this talk

# The SMEFTsim package

an UFO & FeynRules model with\*:

Brivio, Jiang, Trott 1709.06492  
feynrules.irmp.ucl.ac.be/wiki/SMEFT

1. the complete B-conserving Warsaw basis for 3 generations , including all complex phases and ~~CP~~ terms
2. automatic field redefinitions to have **canonical kinetic terms**
3. automatic **parameter shifts** due to the choice of an input parameters set

↪ backup

Main scope:

estimate **tree-level**  $|\mathcal{A}_{\text{SM}}\mathcal{A}_{\text{d=6}}^*|$  **interference** terms → theo. accuracy  $\gtrsim$  few %

\* at the moment only LO, unitary gauge implementation

# The SMEFTsim package

6 different implementations available

Brivio, Jiang, Trott 1709.06492

$$\textcircled{3} \text{ flavor structures } \left\{ \begin{array}{l} \text{general} \\ U(3)^5 \text{ symmetric} \\ \text{linear MFV} \end{array} \right. \times \textcircled{2} \text{ input schemes } \left\{ \begin{array}{l} \hat{\alpha}_{em}, \hat{m}_Z, \hat{G}_f \\ \hat{m}_W, \hat{m}_Z, \hat{G}_f \end{array} \right.$$

[feynrules.irmp.ucl.ac.be/wiki/SMEFT](http://feynrules.irmp.ucl.ac.be/wiki/SMEFT)

[viki: SMEFT](#)  
**Standard Model Effective Field Theory – The SMEFTsim package**  
**Authors**  
 Ilaria Brivio, Yun Jiang and Michael Trott  
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 NBIA and Discovery Center, Niels Bohr Institute, University of Copenhagen

Pre-exported UFO files (include restriction cards)

	Set A		Set B	
	$\alpha$ scheme	$m_W$ scheme	$\alpha$ scheme	$m_W$ scheme
Flavor general SMEFT	<a href="#">SMEFTsim_A_general_alphaScheme_UFO.tar.gz</a>	<a href="#">SMEFTsim_A_general_MwScheme_UFO.tar.gz</a>	<a href="#">SMEFT_alpha_UFO.zip</a>	<a href="#">SMEFT_mW_UFO.zip</a>
MFV SMEFT	<a href="#">SMEFTsim_A_MFV_alphaScheme_UFO.tar.gz</a>	<a href="#">SMEFTsim_A_MFV_MwScheme_UFO.tar.gz</a>	<a href="#">SMEFT_alpha_MFV_UFO.zip</a>	<a href="#">SMEFT_mW_MFV_UFO.zip</a>
$U(3)^5$ SMEFT	<a href="#">SMEFTsim_A_U35_alphaScheme_UFO.tar.gz</a>	<a href="#">SMEFTsim_A_U35_MwScheme_UFO.tar.gz</a>	<a href="#">SMEFT_alpha_FLU_UFO.zip</a>	<a href="#">SMEFT_mW_FLU_UFO.zip</a>





# SMEFTsim validation

## 1. Internal validation: 2 independent versions (A, B)

$\sigma(\text{SM}+\text{int}+\text{quadratic})$  for  $C_i = 1, \Lambda = 1 \text{ TeV}$

process	coefficient	general $\alpha$	general Mw	U(3) <sup>5</sup> $\alpha$	U(3) <sup>5</sup> Mw	MFV $\alpha$	MFV Mw
e+ e- > w+ w-	SMLimit	2.6156 0.059793	2.6788 0.061373	2.6156 0.059793	2.6788 0.061373	2.6156 0.059793	2.6788 0.061373
e+ e- > w+ w- NP=1	Hl3	-	-	4.3384 0.10296	4.4249 0.094337	4.3384 0.10296	4.4249 0.094337
e+ e- > w+ w- NP=1	Hl311	4.6686 0.098776	4.7797 0.10282	-	-	-	-
e+ e- > w+ w- NP=1	W	4.9648 0.10804	5.06 0.11063	4.9648 0.10804	5.06 0.11063	4.9648 0.10804	5.06 0.11063
e+ e- > w+ w- NP=1	Wtil	4.9895 0.10855	5.0848 0.1111	4.9895 0.10855	5.0848 0.1111	-	-
e+ e- > z h	SMLimit	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124
e+ e- > z h NP=1	eW	-	-	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124
e+ e- > z h NP=1	eW11	1.9983 0.0050475	0.01302 0.000033124	-	-	-	-
e+ e- > z h NP=1	He	-	-	1.1756 0.0031	1.1838 0.0031194	1.1756 0.0031	1.1838 0.0031194
e+ e- > z h NP=1	He11	1.1756 0.0031	1.1838 0.0031194	-	-	-	-
e+ e- > z h NP=1	HWB	0.040274 0.00009404	0.036476 0.000084148	0.040274 0.00009404	0.036476 0.000084148	0.040274 0.00009404	0.036476 0.000084148
p p > d s-	SMLimit	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.
p p > d s- NP=1	Delta2qd1	-	-	-	-	690 240. 9319.7	690 240. 9319.7
p p > d s- NP=1	DeltadHq3	-	-	-	-	703 760. 9607.7	703 760. 9607.7
p p > d s- NP=1	DeltadW	-	-	-	-	690 240. 9319.7	690 240. 9319.7
p p > d s- NP=1	dW	-	-	690 240. 9319.7	690 240. 9319.7	-	-
p p > d s- NP=1	dW12	692 740. 9950.4	690 240. 9319.7	-	-	-	-
p p > d s- NP=1	Hq312	706 050. 9205.5	706 050. 9205.5	-	-	-	-

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err

MG5 results with set A

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e+ e- > w+ w- NP=1	Wtil	4.9895 0.10855	5.0848 0.1111	4.9895 0.10855	5.0848 0.1111	-	-
e+ e- > z h	SMLimit	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124
e+ e- > z h NP=1	eW	-	-	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124
e+ e- > z h NP=1	eW11	1.9983 0.0050475	0.01302 0.000033124	-	-	-	-
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e+ e- > z h NP=1	HWB	0.040274 0.00009404	0.036476 0.000084148	0.040274 0.00009404	0.036476 0.000084148	0.040274 0.00009404	0.036476 0.000084148
p p > d s-	SMLimit	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.
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p p > d s- NP=1	dW	-	-	690 240. 9319.7	690 240. 9319.7	-	-
p p > d s- NP=1	dW12	692 740. 9950.4	690 240. 9319.7	-	-	-	-
p p > d s- NP=1	dW112	706 050. 9205.5	706 050. 9205.5	-	-	-	-

5–10 coeff.  $\times$   
~ 20 processes

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err

MG5 results with set A

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e+ e- $\rightarrow$ w+ w- NP=1	Hl311	4.6686 0.098776	4.7797 0.10282	-	-	-	-
e+ e- $\rightarrow$ w+ w- NP=1	W	4.9648 0.10804	5.06 0.11063	4.9648 0.10804	5.06 0.11063	4.9648 0.10804	5.06 0.11063
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e+ e- $\rightarrow$ z h	SMLimit	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124
e+ e- $\rightarrow$ z h NP=1	eW	-	-	0.013009 0.000032914	0.01302 0.000033124	0.013009 0.000032914	0.01302 0.000033124
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e+ e- $\rightarrow$ z h NP=1	HWB	0.040274 0.00009404	0.036476 0.000084148	0.040274 0.00009404	0.036476 0.000084148	0.040274 0.00009404	0.036476 0.000084148
p p $\rightarrow$ d s-	SMLimit	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.	688 390. 11 858.
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p p $\rightarrow$ d s- NP=1	W12	692 740. 9950.4	690 240. 9319.7	-	-	-	-
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 $\sim$  20 processes

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2. Validation against dim6top

feynrules.irmp.ucl.ac.be/wiki/dim6top – G.Durieux, C.Zhang

Top WG note: Aguilar-Saavedra et al. 1802.07237

$\sigma(\text{int.})/\sigma(\text{SM})$  for  $C_i$ ,  $\Lambda = 1$  TeV [permille]

	$pp \rightarrow t\bar{t}$	$pp \rightarrow t\bar{t}b\bar{b}$	$pp \rightarrow t\bar{t}t\bar{t}$	$pp \rightarrow t\bar{t}e^+e^-$	$pp \rightarrow t\bar{t}\gamma$	$pp \rightarrow t\bar{t}h$	$pp \rightarrow t\bar{t}j$	$pp \rightarrow t\bar{t}e^+\nu$	$pp \rightarrow t\bar{t}e^+e^-$	$pp \rightarrow t\bar{t}\gamma$	$pp \rightarrow t\bar{t}h$			
SM	as	$5.2 \times 10^7$ pb	1.9 pb	0.0098 pb	0.02 pb	0.016 pb	1.4 pb	0.4 pb	0.4 pb	55 pb	2.5 pb	0.0054 pb	0.39 pb	0.016 pb
$c_{\phi 2}^{(1)}$	-0.25	-1.9	$-1 \times 10^2$	-1.6	-0.87	-0.71								
$c_{\phi 8}^{(1)}$	-0.16	-3.2	-34	-91	-0.5	-0.27								
$c_{\phi 11}^{(1)}$	-0.15	-5.6	$1 \times 10^2$	-0.1	-0.19	-0.55								
$c_{\phi 18}^{(1)}$	-0.053	-1.8	-41	-0.18	-0.095	-0.15								
$c_{\phi 21}^{(1)}$	-0.0095	0.72	-0.052	-0.015	-0.007	-0.026								
$c_{\phi 26}^{(1)}$	0.14	3.9	0.12	0.35	0.16	0.56								
$c_{t 11}^{(1)}$			$-1.8 \times 10^2$											
$c_{t 13}^{(1)}$	-0.0095	0.46	-0.059	-0.02	-0.021	-0.039								
$c_{t 18}^{(1)}$	0.13	3.5	0.11	0.26	0.31	0.56								
$c_{\phi Q 1}^{(1)}$														
$c_{\phi Q 2}^{(1)}$														
$c_{\phi Q 3}^{(1)}$														
$c_{\phi Q 11}^{(1)}$														
$c_{\phi Q 21}^{(1)}$														
$c_{\phi Q 31}^{(1)}$														
$c_{\phi Q 33}^{(1)}$	2.7	-0.11	4.7	-85	-20	15	$4 \times 10^{-14}$	$-6.4 \times 10^{-14}$	$-5.2 \times 10^{-14}$	$-4.1 \times 10^{-14}$				
$c_{\phi 11}^{(2)}$														
$c_{\phi 18}^{(2)}$														
$c_{\phi 26}^{(2)}$														
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$c_{\phi 125}^{(2)}$														
$c_{\phi 126}^{(2)}$														
$c_{\phi 127}^{(2)}$														
$c_{\phi 128}^{(2)}$														





# SMEFTsim validation

1. Internal validation: 2 independent versions (A, B)

2. Validation against dim6top

[feynrules.irmp.ucl.ac.be/wiki/dim6top](https://feynrules.irmp.ucl.ac.be/wiki/dim6top) – G.Durieux, C.Zhang

**Top WG note:** Aguilar-Saavedra et al. 1802.07237

3. Validation against VBFNLO

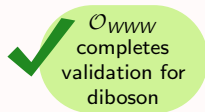
Arnold et al. 0811.4559,1107.4038, Baglio et al 1404.3940

**VBSCan** Thessaloniki Workshop summary. To appear.

**VBFNLO** has hard coded matrix elements for selected EW processes  
uses HISZ basis → could validate  $O_{WWW} = \varepsilon_{ijk} W_\nu^{i\mu} W_\rho^{j\nu} W_\mu^{k\rho}$

checked:  $pp \rightarrow e^+ \nu_e \mu^+ \mu^-$  and  $pp \rightarrow e^+ \nu_e \mu^+ \nu_\mu jj$

LO, compared  $\sigma_{SM} +$  distributions





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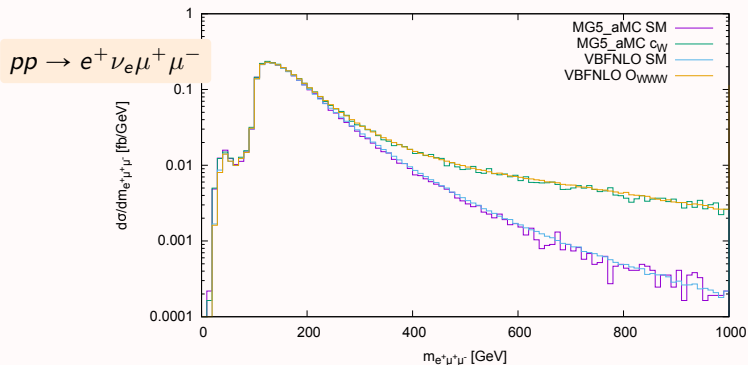
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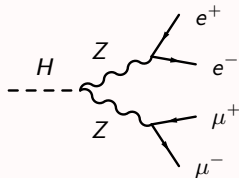
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4. Validation against analytic expressions

Brivio, Trott SMEFT review 1706.08945

Example



	theory	MG interf.	MG full xs
cHW	-0.757133	-0.77948	-0.778724
cHB	-0.217121	-0.223247	-0.223151
cHWB	0.308271	0.295226	0.317418
cHbox	2.	1.99882	2.00469
cHD	0.167224	0.164264	0.170457
cHe	-3.5239	-1.72758	-1.72691
cHl1	4.38291	2.15039	2.14801
cHl3	-1.61513	-3.85776	-3.86201
cll1	2.99835	2.99884	3.00731

$$\sigma(\text{int.})/\sigma(\text{SM}) \text{ for } \bar{C}_i = C_i(v/\Lambda)^2 = 1$$

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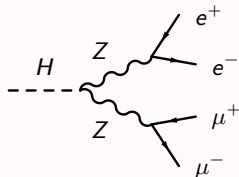
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dominant diag. contribution  
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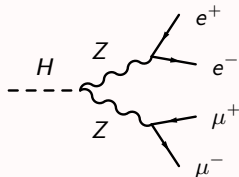
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MG5:  $h \rightarrow e^+ e^- \mu^+ \mu^- / a$

full xsec,  
linearized

pure int.

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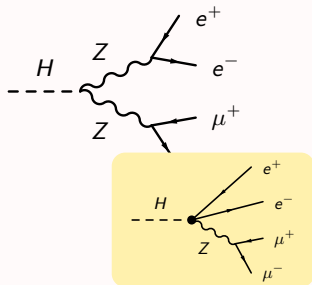
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$z > e^+ e^-$

$w^+ > l^+ \nu_l$

$h > a a$

$h > b b$

$h > e^+ e^- \mu^+ \mu^- / a$

$p p > z h / a$

$p p > w^+ h$

...

$z > u u^{\sim}$

$w^+ > u q d q^{\sim}$

$h > z a$

$h > t a^+ t a^-$

$h > e^+ \nu_e \mu^- \nu_\mu$

$g g > h$

$p p > w^- h$

# SMEFTsim validation

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2. Validation against dim6top
3. Validation against VBFNLO
4. Validation against analytic expressions
5. Further validation still ongoing!

[feynrules.irmp.ucl.ac.be/wiki/dim6top](http://feynrules.irmp.ucl.ac.be/wiki/dim6top) – G.Durieux,C.Zhang

**Top WG note:** Aguilar-Saavedra et al. 1802.07237

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**VBSCan** Thessaloniki Workshop summary. To appear.

Brivio,Trott SMEFT review 1706.08945

# Application: global SMEFT fit

## Key features

Brivio, Hays, Trott, Žemaitytė, in preparation

- ▶ semianalytic LO corrections **up to**  $\mathcal{O}(\Lambda^{-2})$  → fully **analytic**  $\chi^2$  minimization
- ▶ # parameters controlled by
  - ▶ IR symmetries
  - ▶ observables (SM selection rules)



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	total $N_f = 3$	WZH pole obs.
general	2499	~ 46
MFV	~ 108	~ 30
$U(3)^5$	~ 70	~ 24

Brivio,Jiang,Trott 1709.06492

starting point: **resonant-dominated** processes with  $\mathbf{U(3)^5}$

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- ▶ combine different **datasets**.  
so far: EWPD + Higgs + diboson (LEP)  
make it extensible: + top + diboson/EW (LHC) + flavor + ...

combining is **crucial!**

- ▶ needed to break degeneracies
- ▶ large mixings from RGE between measurements at different energies (e.g. LHC vs flavor)

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- ▶ **method**: builds upon
  - Hays,Sanz,Žemaitytė LHCHSWG-INT-2017-001 (Higgs)
  - ATLAS note: ATL-PHYS-PUB-2017-018
  - Brivio,Trott 1701.06424 (EWPD)
  - Berthier,(Bjørn),Trott 1508.05060,1606.06693

see also e.g. Ellis,Murphy,Sanz,You 1803.03252 (with SMEFTsim)  
Almeida,Alves,Rosa-Agostinho,Éboli,Gonzalez-Garcia 1812.01009

## Z,W couplings

$$\begin{aligned}
 Q_{HI}^{(1)} &= (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{l}\gamma^\mu l) \\
 Q_{He} &= (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{e}\gamma^\mu e) \\
 Q_{Hq}^{(1)} &= (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{q}\gamma^\mu q) \\
 Q_{Hq}^{(3)} &= (iH^\dagger \overleftrightarrow{D}_\mu^i H)(\bar{q}\sigma^i\gamma^\mu q) \\
 Q_{Hu} &= (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{u}\gamma^\mu u) \\
 Q_{Hd} &= (iH^\dagger \overleftrightarrow{D}_\mu H)(\bar{d}\gamma^\mu d)
 \end{aligned}$$

$$\begin{aligned}
 Q_{HD} &= (D_\mu H^\dagger H)(H^\dagger D^\mu H) \\
 Q_{HWB} &= (H^\dagger \sigma^i H)W_{\mu\nu}^i B^{\mu\nu} \\
 Q_{HI}^{(3)} &= (iH^\dagger \overleftrightarrow{D}_\mu^i H)(\bar{l}\sigma^i\gamma^\mu l) \\
 Q'_{ll} &= (\bar{l}_p\gamma^\mu l_r)(\bar{l}_r\gamma^\mu l_p)
 \end{aligned}$$

input quantities

$$Q_W = \varepsilon_{ijk} W_\mu^{i\nu} W_\nu^{j\rho} W_\rho^{k\mu}$$

TGC

## Bhabha scattering

$$\begin{aligned}
 Q_{ee} &= (\bar{e}\gamma^\mu e)(\bar{e}\gamma^\mu e) \\
 Q_{le} &= (\bar{l}\gamma^\mu l)(\bar{e}\gamma^\mu e) \\
 Q_{ll} &= (\bar{l}_p\gamma^\mu l_p)(\bar{l}_r\gamma^\mu l_r)
 \end{aligned}$$

$$\begin{aligned}
 Q_{Hbox} &= (H^\dagger H) \square (H^\dagger H) \\
 Q_{HG} &= (H^\dagger H)G_{\mu\nu}^a G^{a\mu\nu} \\
 Q_{HB} &= (H^\dagger H)B_{\mu\nu} B^{\mu\nu} \\
 Q_{HW} &= (H^\dagger H)W_{\mu\nu}^i W^{i\mu\nu} \\
 Q_{uH} &= (H^\dagger H)(\bar{q}Hu) \\
 Q_{dH} &= (H^\dagger H)(\bar{q}Hd) \\
 Q_{eH} &= (H^\dagger H)(\bar{q}He) \\
 Q_G &= \varepsilon_{abc} G_\mu^{a\nu} G_\nu^{b\rho} G_\rho^{c\mu} \\
 Q_{uG} &= (\bar{q}\sigma^{\mu\nu} T^a \tilde{H}u)G_{\mu\nu}^a
 \end{aligned}$$

H processes



# Summary & take-home

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- ▶ It is worth using its full power with a truly **global** analysis
  - ▶ most general BSM characterization (assuming SM sym + fields)
  - ▶ universal parameterization for data interpretation

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- ▶ **SMEFTsim** is a powerful tool specifically designed for this purpose @LO
- ▶ **Validation** of SMEFTsim has been extensive and is still ongoing
- ▶ **Global fits** can and should be done!
  - ▶ starting from flavor symmetric + pole obs: **only** ~ **23** parameters
  - ▶ can be systematically extended to other observables and improved with higher orders

**Backup slides**

$X^3$		$\varphi^6$ and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
$Q_G$	$f^{ABC} G_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu}$	$Q_{\varphi}$	$(\varphi^{\dagger}\varphi)^3$	$Q_{e\varphi}$	$(\varphi^{\dagger}\varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu}$	$Q_{\varphi\Box}$	$(\varphi^{\dagger}\varphi)\Box(\varphi^{\dagger}\varphi)$	$Q_{u\varphi}$	$(\varphi^{\dagger}\varphi)(\bar{q}_p u_r \tilde{\varphi})$
$Q_W$	$\varepsilon^{IJK} W_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$	$Q_{\varphi D}$	$(\varphi^{\dagger} D^{\mu}\varphi)^{\star} (\varphi^{\dagger} D_{\mu}\varphi)$	$Q_{d\varphi}$	$(\varphi^{\dagger}\varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^{\dagger}\varphi G_{\mu\nu}^A G^{A\mu\nu}$	$Q_{eW}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{l}_p \gamma^{\mu} l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^{\dagger}\varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	$Q_{eB}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu}^I \varphi) (\bar{l}_p \tau^I \gamma^{\mu} l_r)$
$Q_{\varphi W}$	$\varphi^{\dagger}\varphi W_{\mu\nu}^I W^{I\mu\nu}$	$Q_{uG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{e}_p \gamma^{\mu} e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^{\dagger}\varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	$Q_{uW}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{q}_p \gamma^{\mu} q_r)$
$Q_{\varphi B}$	$\varphi^{\dagger}\varphi B_{\mu\nu} B^{\mu\nu}$	$Q_{uB}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu}^I \varphi) (\bar{q}_p \tau^I \gamma^{\mu} q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^{\dagger}\varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	$Q_{dG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{u}_p \gamma^{\mu} u_r)$
$Q_{\varphi WB}$	$\varphi^{\dagger} \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	$Q_{dW}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{d}_p \gamma^{\mu} d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^{\dagger} \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$Q_{dB}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^{\dagger} D_{\mu} \varphi) (\bar{u}_p \gamma^{\mu} d_r)$

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
$Q_{ll}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	$Q_{ee}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$Q_{le}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{uu}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{lu}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{dd}$	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{ld}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{eu}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{qe}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{ed}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		$B$ -violating			
$Q_{ledq}$	$(\bar{l}_p^j e_r)(\bar{d}_s^j q_t^j)$	$Q_{duq}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_r^j)^T C l_t^k]$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	$Q_{qqqu}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	$Q_{qqq}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} \varepsilon_{mnn} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^m)^T C l_t^n]$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	$Q_{duu}$	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				

## Gauge bosons

$$\begin{aligned}\mathcal{L}_{\text{SMEFT}} \supset & -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}W_{\mu\nu}^I W^{I\mu\nu} - \frac{1}{4}G_{\mu\nu}^a G^{a\mu\nu} + \\ & + C_{HB}(H^\dagger H)B_{\mu\nu}B^{\mu\nu} + C_{HW}(H^\dagger H)W_{\mu\nu}^I W^{I\mu\nu} + C_{HWB}(H^\dagger \sigma^I H)W_{\mu\nu}^I B^{\mu\nu} \\ & + C_{HG}(H^\dagger H)G_{\mu\nu}^a G^{a\mu\nu}\end{aligned}$$

to have **canonically normalized kinetic terms** we need to

1. redefine fields and couplings keeping  $(gV_\mu)$  unchanged:

$$\begin{aligned}B_\mu &\rightarrow B_\mu(1 + C_{HB}v^2) & g_1 &\rightarrow g_1(1 - C_{HB}v^2) \\ \mathcal{W}_\mu^I &\rightarrow W_\mu^I(1 + C_{HW}v^2) & g_2 &\rightarrow g_2(1 - C_{HW}v^2) \\ G_\mu^a &\rightarrow G_\mu^a(1 + C_{HG}v^2) & g_s &\rightarrow g_s(1 - C_{HG}v^2)\end{aligned}$$

2. correct the rotation to mass eigenstates:

$$\begin{pmatrix} \mathcal{W}_\mu^3 \\ B_\mu \end{pmatrix} = \begin{pmatrix} 1 & -v^2 C_{HWB}/2 \\ -v^2 C_{HWB}/2 & 1 \end{pmatrix} \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} Z_\mu \\ A_\mu \end{pmatrix}$$

(equivalent to a shift of the Weinberg angle)

Grinstein, Wise Phys. Lett. B265(1991)326  
Alonso, Jenkins, Manohar, Trott 1312.2014

## Higgs

$$\mathcal{L}_{\text{SMEFT}} \supset \frac{1}{2} D_\mu H^\dagger D^\mu H + C_{H\Box} (H^\dagger H) (\Box H) + C_{HD} (H^\dagger D_\mu H)^* (\Box H)$$

to have a canonically normalized kinetic term, in unitary gauge, we need to replace

$$h \rightarrow h \left( 1 + v^2 C_{H\Box} - \frac{v^2}{4} C_{HD} \right)$$

Alonso, Jenkins, Manohar, Trott 1312.2014

# Shifts from input parameters

SM case.

Parameters in the canonically normalized Lagrangian :  $\bar{v}, \bar{g}_1, \bar{g}_2, s_{\bar{\theta}}$

The values can be inferred from the measurements e.g. of  $\{\alpha_{\text{em}}, m_Z, G_f\}$  :

$$\alpha_{\text{em}} = \frac{\bar{g}_1 \bar{g}_2}{\bar{g}_1^2 + \bar{g}_2^2}$$

$$m_Z = \frac{\bar{g}_2 \bar{v}}{2c_{\bar{\theta}}}$$

$$G_f = \frac{1}{\sqrt{2}\bar{v}^2}$$

→

$$\hat{v}^2 = \frac{1}{\sqrt{2}G_f}$$

$$\sin^2 \hat{\theta} = \frac{1}{2} \left( 1 - \sqrt{1 - \frac{4\pi\alpha_{\text{em}}}{\sqrt{2}G_f m_Z^2}} \right)$$

$$\hat{g}_1 = \frac{\sqrt{4\pi\alpha_{\text{em}}}}{\cos \hat{\theta}}$$

$$\hat{g}_2 = \frac{\sqrt{4\pi\alpha_{\text{em}}}}{\sin \hat{\theta}}$$

in the SM at tree-level  $\bar{\kappa} = \hat{\kappa}$



# Shifts from input parameters

SMEFT case.

Parameters in the canonically normalized Lagrangian :  $\bar{v}, \bar{g}_1, \bar{g}_2, s_{\bar{\theta}}$

The values can be inferred from the measurements e.g. of  $\{\alpha_{\text{em}}, m_Z, G_f\}$  :

$$\begin{aligned}\alpha_{\text{em}} &= \frac{\bar{g}_1 \bar{g}_2}{\bar{g}_1^2 + \bar{g}_2^2} \left[ 1 + \bar{v}^2 C_{HWB} \frac{\bar{g}_2^3 / \bar{g}_1}{\bar{g}_1^2 + \bar{g}_2^2} \right] & \hat{v}^2 &= \frac{1}{\sqrt{2} G_f} \\ m_Z &= \frac{\bar{g}_2 \bar{v}}{2 c_{\bar{\theta}}} + \delta m_Z(C_i) & \sin \hat{\theta}^2 &= \frac{1}{2} \left( 1 - \sqrt{1 - \frac{4\pi \alpha_{\text{em}}}{\sqrt{2} G_f m_Z^2}} \right) \\ G_f &= \frac{1}{\sqrt{2} \bar{v}^2} + \delta G_f(C_i) & \hat{g}_1 &= \frac{\sqrt{4\pi \alpha_{\text{em}}}}{\cos \hat{\theta}} \\ & & \hat{g}_2 &= \frac{\sqrt{4\pi \alpha_{\text{em}}}}{\sin \hat{\theta}}\end{aligned}$$

in the SM at tree-level  $\bar{\kappa} = \hat{\kappa}$

in the SMEFT  $\bar{\kappa} = \hat{\kappa} + \delta \kappa(C_i)$

# Shifts from input parameters

To have numerical predictions it is necessary to replace  $\bar{\kappa} \rightarrow \hat{\kappa} + \delta\kappa(C_i)$  for all the parameters in the Lagrangian.

---

$\{\alpha_{\text{em}}, m_Z, G_f\}$  scheme

$$\delta m_Z^2 = m_Z^2 \hat{v}^2 \left( \frac{c_{HD}}{2} + 2c_{\hat{\theta}} s_{\hat{\theta}} c_{HWB} \right)$$

$$\delta G_f = \frac{\hat{v}^2}{\sqrt{2}} \left( (c_{HI}^{(3)})_{11} + (c_{HI}^{(3)})_{22} - (c_{II})_{1221} \right)$$

$$\delta g_1 = \frac{s_{\hat{\theta}}^2}{2(1-2s_{\hat{\theta}}^2)} \left( \sqrt{2}\delta G_f + \delta m_Z^2/m_Z^2 + 2\frac{c_{\hat{\theta}}^3}{s_{\hat{\theta}}} c_{HWB} \hat{v}^2 \right)$$

$$\delta g_2 = -\frac{c_{\hat{\theta}}^2}{2(1-2s_{\hat{\theta}}^2)} \left( \sqrt{2}\delta G_f + \delta m_Z^2/m_Z^2 + 2\frac{s_{\hat{\theta}}^3}{c_{\hat{\theta}}} c_{HWB} \hat{v}^2 \right)$$

$$\delta s_{\hat{\theta}}^2 = 2c_{\hat{\theta}}^2 s_{\hat{\theta}}^2 (\delta g_1 - \delta g_2) + c_{\hat{\theta}} s_{\hat{\theta}} (1 - 2s_{\hat{\theta}}^2) c_{HWB} \hat{v}^2$$

$$\delta m_h^2 = m_h^2 \hat{v}^2 \left( 2c_{H\Box} - \frac{c_{HD}}{2} - \frac{3c_H}{2lam} \right)$$

# Shifts from input parameters

To have numerical predictions it is necessary to replace  $\bar{\kappa} \rightarrow \hat{\kappa} + \delta\kappa(C_i)$   
for all the parameters in the Lagrangian.

---

$\{m_W, m_Z, G_f\}$  scheme

$$\delta m_Z^2 = m_Z^2 \hat{v}^2 \left( \frac{c_{HD}}{2} + 2c_{\hat{\theta}} s_{\hat{\theta}} c_{HWB} \right)$$

$$\delta G_f = \frac{\hat{v}^2}{\sqrt{2}} \left( (c_{HI}^{(3)})_{11} + (c_{HI}^{(3)})_{22} - (c_{II})_{1221} \right)$$

$$\delta g_1 = -\frac{1}{2} \left( \sqrt{2} \delta G_f + \frac{1}{s_{\hat{\theta}}^2} \frac{\delta m_Z^2}{m_Z^2} \right)$$

$$\delta g_2 = -\frac{1}{\sqrt{2}} \delta G_f$$

$$\delta s_{\hat{\theta}}^2 = 2c_{\hat{\theta}}^2 s_{\hat{\theta}}^2 (\delta g_1 - \delta g_2) + c_{\hat{\theta}} s_{\hat{\theta}} (1 - 2s_{\hat{\theta}}^2) c_{HWB} \hat{v}^2$$

$$\delta m_h^2 = m_h^2 \hat{v}^2 \left( 2c_{CH} - \frac{c_{HD}}{2} - \frac{3c_H}{2\lambda m} \right)$$

120 observables included so far

- ▶ 10 near-Z-pole EWPO:  $\Gamma_Z$ ,  $R_{\ell,c,b}^0$ ,  $A_{FB}^{\ell,c,b,\mu,\tau}$ ,  $\sigma_h^0$  LEPI combination hep-ex/0509008
- ▶ 21 distribution bins for bhabha scattering at LEP II LEPII combination 1302.3415
- ▶ 74 dist. bins for  $W^+W^-$  production at LEP II L3: hep-ex/0409016  
OPAL: 0708.1311  
ALEPH: Eur.Phys.J. C38 (2004) 147  
differential combined: 1302.3415
- ▶ 15 inclusive obs. for Higgs measurements in  $H \rightarrow \gamma\gamma$  and  $H \rightarrow 4\ell$  at LHC
  - ▶ ATLAS ( $36 \text{ fb}^{-1}$ ) ATLAS-CONF-2017-047
  - ▶ CMS ( $36 \text{ fb}^{-1}$ ) CMS PAS HIG-17-031

## Example: dependence for $gg \rightarrow h \rightarrow 4\ell$

correction to the inclusive rate, relative to the SM  
obtained automatically with SMEFTsim

$$\begin{aligned} & 1 + 0.0185 * CG + 0.000425 * CH_{\text{box}} - 0.0001062 * CHD + 22.3 * CHG \\ & - 0.000422 * CuHAbs - 0.000425 * CHL3 + 0.000212 * Cll1 + 0.1212 * \\ & CH_{\text{box}} + 0.1193 * CHD + 0.04691 * CHW + 0.01345 * CHB + 0.1284 * CHWB \\ & + 0.1279 * CHl1 + 0.01765 * CHl3 + 0.003545 * CHe + 0.0925 * Cl1 + \\ & 0.1819 * Cl11 - 0.000491 * CHWB + 0.0001946 * CHl1 + 0.001461 * CHl3 \\ & + 0.0001942 * CHe - 0.0004985 * CHq1 - 0.001724 * CHq3 - 0.000259 * \\ & CHu + 0.0001917 * CHd - 0.00107 * Cl11 - (0.1166 * CH_{\text{box}} + 0.000747 * \\ & CHD + 1.445 * CHG + 0.01088 * CHW + 0.0001615 * CHB + 0.04346 * CHWB \\ & + 0.0001276 * CHl1 + 0.000786 * CHl3 + 0.000598 * CHq1 + 0.01186 \\ & * CHq3 + 0.0002017 * CHu + 0.0729 * Cl11 + 0.01098 * CHD - 0.0706 * \\ & CHWB + 0.0001807 * CdWAbs + 0.02797 * CHl1 + 0.2101 * CHl3 + 0.02792 \\ & * CHe - 0.0717 * CHq1 - 0.2479 * CHq3 - 0.03722 * CHu + 0.02755 * CHd \\ & - 0.1537 * Cl11 + 0.0002095 * CeWAbs + 0.0003167 * CuWAbs + 2.622 * \\ & CHl3 - 2.551 * CHq3 - 1.965 * Cl11) \end{aligned}$$

- ▶ all the relevant operators are included
- ▶ only **interference** is kept  $\rightarrow$  simple **linear** expressions

## Basics of the fit method

Likelihood:

$$L(C_i) = \frac{1}{\sqrt{(2\pi)^n \det V}} \exp\left(-\frac{1}{2} (\hat{O} - \bar{O})^T V^{-1} (\hat{O} - \bar{O})\right)$$



$$\chi^2 = -2 \log L(C_i)$$



extract **best-fit values** on each  $C_i$   
after profiling the  $\chi^2$  over the others

# Global fit to EW precision data - method

Likelihood:

$$L(C_i) = \frac{1}{\sqrt{(2\pi)^n \det V}} \exp \left( -\frac{1}{2} (\hat{O} - \bar{O})^T V^{-1} (\hat{O} - \bar{O}) \right)$$

# observables

SMEFT prediction ( $C_i$ )

exp. measurement

covariance matrix

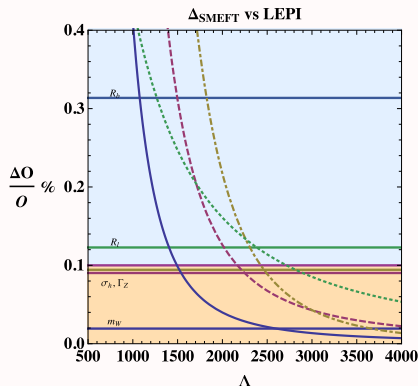
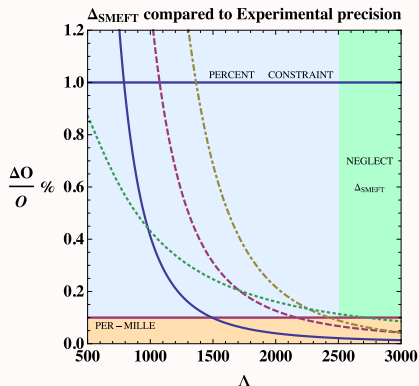
$$V_{i,j} = \Delta_i^{\text{exp}} \rho_{ij}^{\text{exp}} \Delta_j^{\text{exp}} + \Delta_i^{\text{th}} \rho_{ij}^{\text{th}} \Delta_j^{\text{th}}$$

error on  $O_i$

correlation mat.

$$\Delta_i^{\text{th}} = \sqrt{\Delta_{i,\text{SM}}^2 + \Delta_{\text{SMEFT}}^2 \bar{O}_i^2}$$

- SMEFT uncertainty:  $\rightarrow$  impact of  $d \geq 8$  operators + radiative corrections  
 $\rightarrow$  initial/final state radiation  
 $\rightarrow$  ...



Berthier, Trott 1508.05060  
 Hays, Martin, Sanz, Setford 1808.00442

in the fit: taken to be a fixed flat relative uncertainty  $0 \leq \Delta_{\text{SMEFT}} \leq 1\%$



# Constructing convenient observables

looking for an optimal set of observables

only a **few** operators contributing significantly  
many observables **share the same** relevant ops.  
sufficient experimental **sensitivity**

Working assumption:

the dominant effect is the **tree-level interference**  $|\mathcal{A}_{SM}\mathcal{A}_{d=6}^*| \sim \frac{C_i}{\Lambda^2}$ .

whenever this is **suppressed**, the coefficient  $C_i$  can be neglected even if  $C_i \neq 0$

- ▶ in specific kinematic regions. e.g. for  $\psi^4$  ops. close to **W, Z, h poles**

# Constructing convenient observables

Example – close to a pole

Brivio, Jiang, Trott 1709.06492

most  $\psi^4$  operators give diagrams with less resonances

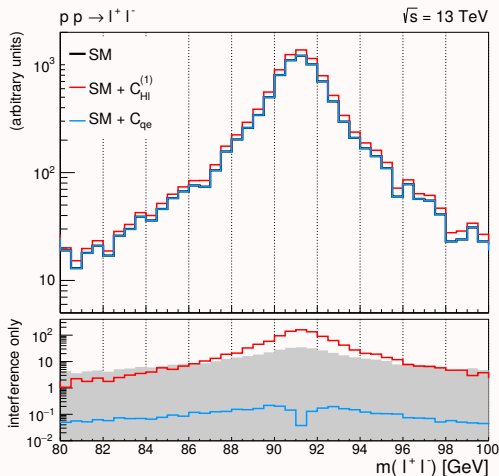
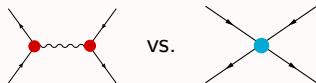
expected to be **suppressed**  
wrt. “pole operators” by

$$\left(\frac{\Gamma_B m_B}{v^2}\right)^n \sim \begin{cases} 1/300 & (Z,W) \\ 1/10^6 & (h) \end{cases}$$

$$B = \{Z, W, h\}$$

$n = \#$  missing resonances

Drell-Yan via Z resonance  $\rightarrow$



# Constructing convenient observables

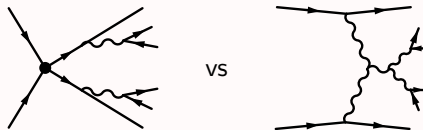
Example – close to a pole

Brivio, Jiang, Trott 1709.06492

most  $\psi^4$  operators give diagrams with less resonances

! Not *always* the case. The impact must be checked case by case

E.g. VBS



the 4-fermion diagram is not removed by poles selection.

# Constructing convenient observables

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whenever this is **suppressed**, the coefficient  $C_i$  can be neglected even if  $C_i \neq 0$

- ▶ in specific kinematic regions. e.g. for  $\psi^4$  ops. close to **W, Z, h poles**
- ▶ for operators with interference  $\propto m_f$

Example: **dipole operators** can be neglected for  $f \neq t, b$



# Constructing convenient observables

looking for an optimal set of observables

only a **few** operators contributing significantly  
many observables **share the same** relevant ops.  
sufficient experimental **sensitivity**

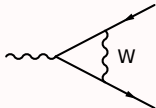
Working assumption:

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whenever this is **suppressed**, the coefficient  $C_i$  can be neglected even if  $C_i \neq 0$

- ▶ in specific kinematic regions. e.g. for  $\psi^4$  ops. close to **W, Z, h poles**
- ▶ for operators with interference  $\propto m_f$
- ▶ for operators inducing FCNC

$\mathcal{A}_{SM}$  is very suppressed:


$$\sim \frac{m_j^2 V_{jk}^* V_{ji}}{32\pi^2 m_W^2}$$

# Constructing convenient observables

looking for an optimal set of observables

- only a **few** operators contributing significantly
- many observables **share the same** relevant ops.
- sufficient experimental **sensitivity**

Working assumption:

the dominant effect is the **tree-level interference**  $|\mathcal{A}_{SM}\mathcal{A}_{d=6}^*| \sim \frac{C_i}{\Lambda^2}$ .

whenever this is **suppressed**, the coefficient  $C_i$  can be neglected even if  $C_i \neq 0$

- ▶ in specific kinematic regions. e.g. for  $\psi^4$  ops. close to **W, Z, h poles**
- ▶ for operators with interference  $\propto m_f$
- ▶ for operators inducing FCNC
- ▶ ...

Brivio, Jiang, Trott 1709.06492

	total $N_f = 3$	WZH poles
general	2499	$\sim 46$
MFV	$\sim 108$	$\sim 30$
$U(3)^5$	$\sim 70$	$\sim 24$

**The counts reduce significantly!**