



# Overview of Flux Trapping at Cornell

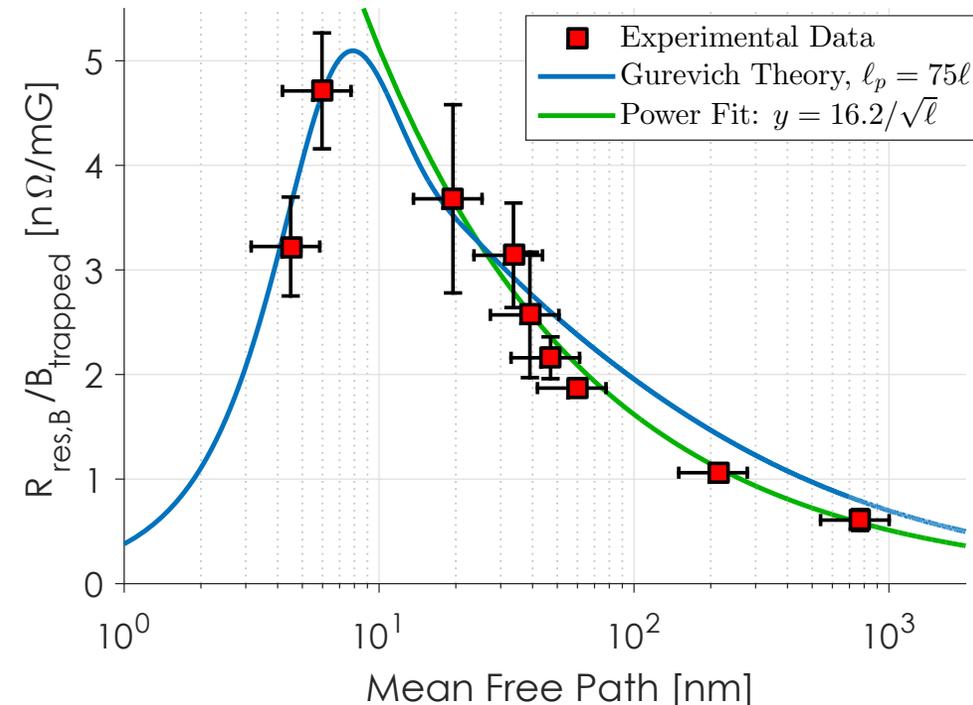
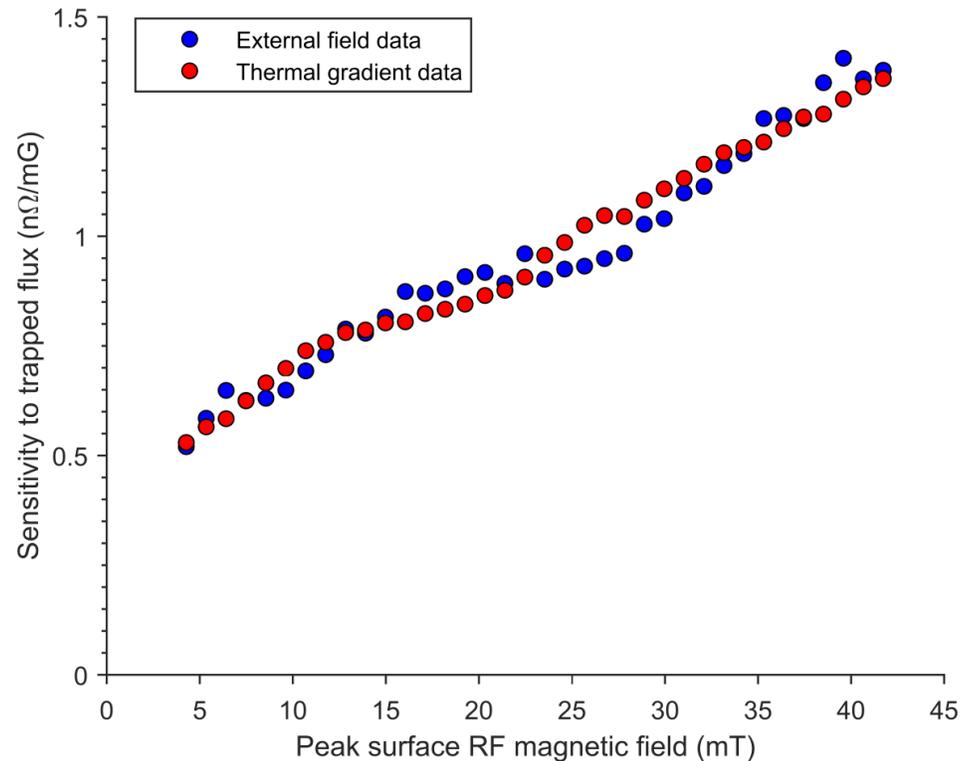
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- In general **linear sensitivity** with an “offset” (y-intercept) parameter and a “slope” parameter. Linear slope predicted by “collective weak pinning” (see D. Liarte’s talk later today!)
- Parameters dependent on **cavity treatment** and **material parameters**, especially the **electron mean free path**, **doping depth**, and **frequency**.
- All Cornell cavities made from the same Nb stock



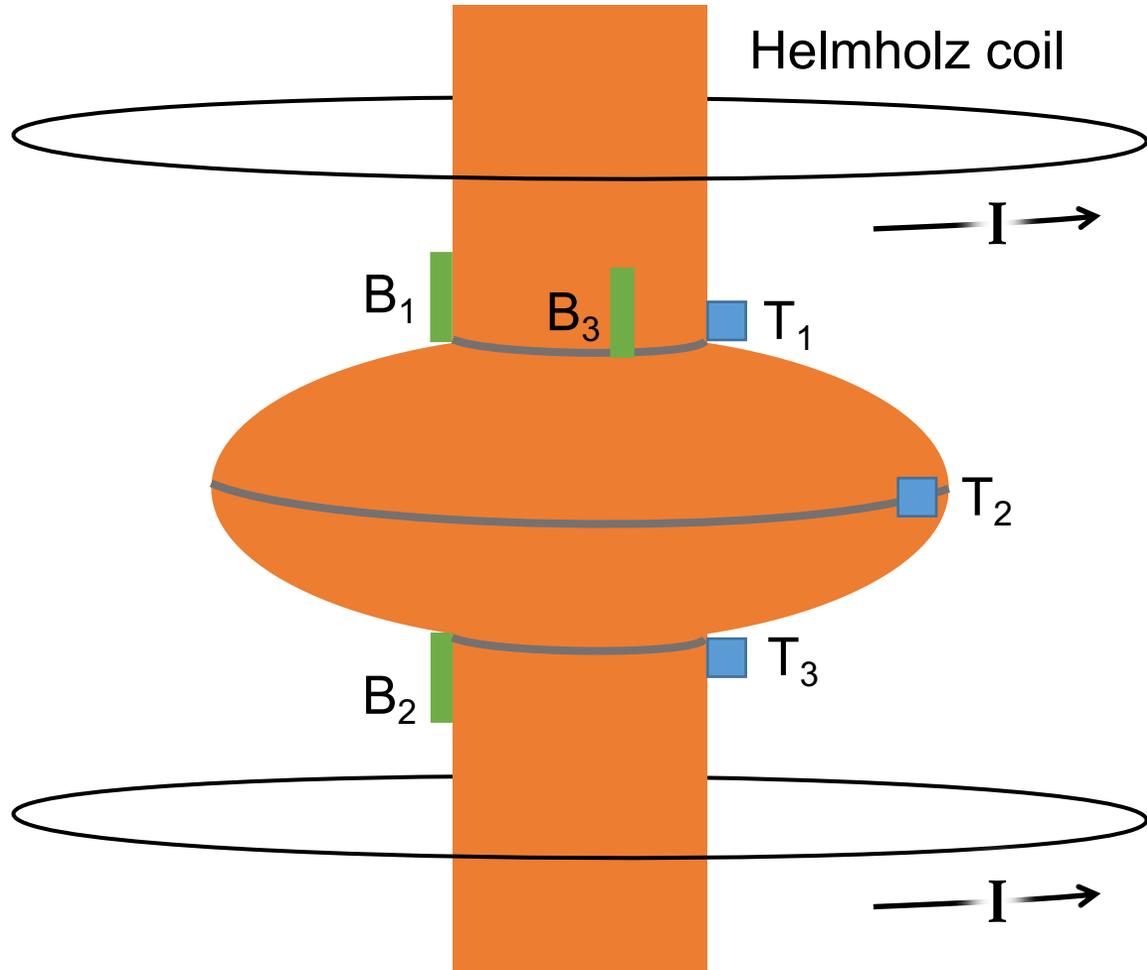


- Overview of flux sensitivity measurement procedure @ Cornell
- Results for doped niobium
  - **1.3 GHz** high-T (**800 °C**) **N-doping** (including “**2/6**” recipe)
  - **1.3 GHz** **160 °C** **impurity doping**
  - **2.6 GHz** “**2/6**” **N-doping**
- Results for Nb<sub>3</sub>Sn
  - **1.3 GHz** cavities





# Flux Sensitivity Test Procedure



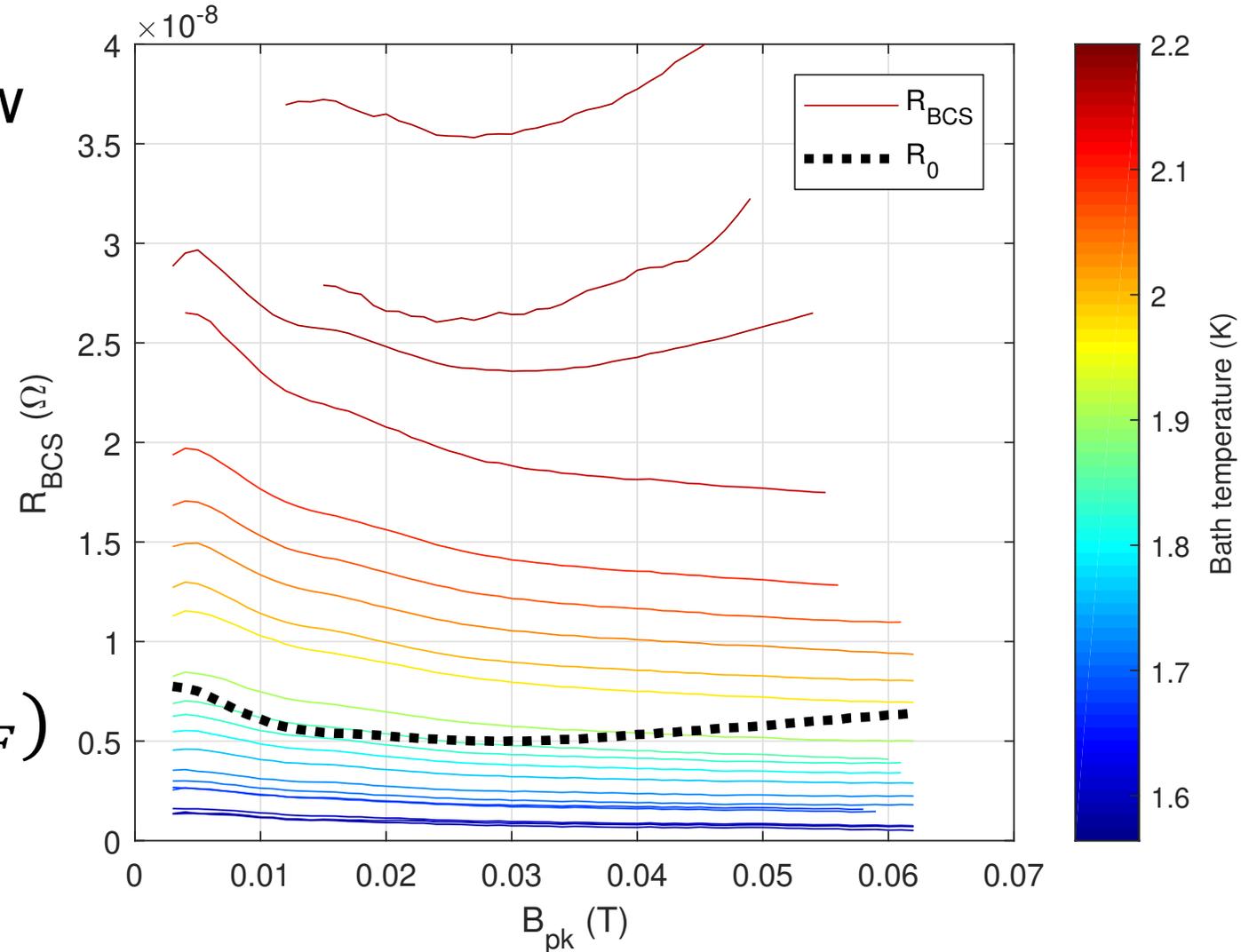
- Fast cool to expel mag. flux;  
slow cool to trap it
- Temperature sensors (Cernox)  
to measure  $\Delta T$  at  $T_c$
- Flux gate magnetometers to  
measure ambient field, applied  
field, expulsion, trapping
- Note:  $dT/dx$  across cell  
matters, not  $dT/dt$ 
  - Though cooling quickly is an  
easy way to get a large  $dT/dx$ .



# Flux Sensitivity Test Procedure

- One test with fast cool
- One or more tests with slow cools, varying trapped flux
- Collect RF test data:  
 $R_s(B_{RF}, T) = G / Q_0$
- Exponential fitting (SRIMP algorithm) to separate  $R_s$ :

$$R_s = R_{BCS}(B_{RF}, T) + R_0(B_{RF})$$





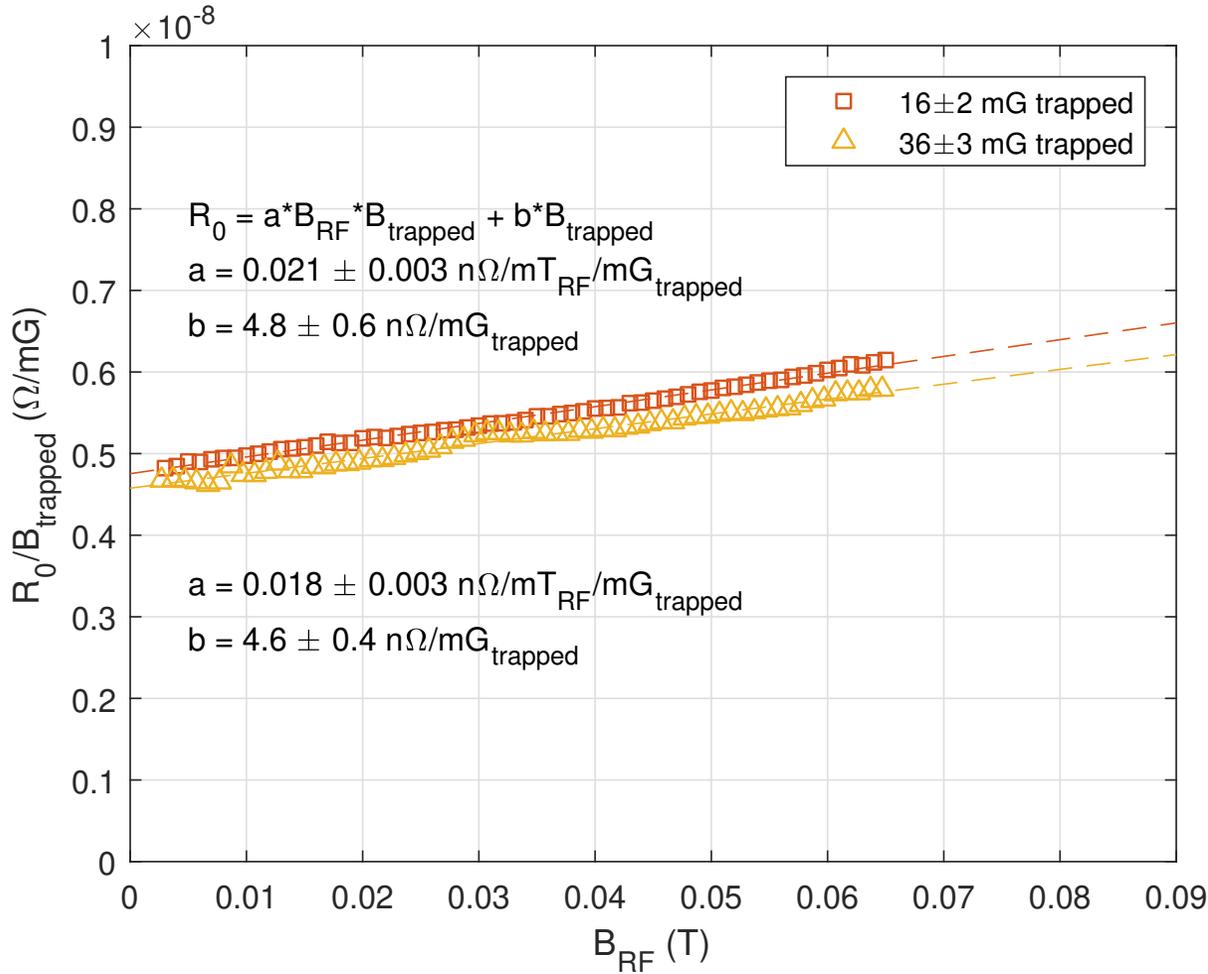
# Flux Sensitivity Test Procedure

- Compare  $R_0$  from different cooldowns to determine sensitivity:

$$R_0(B_{\text{trapped}}) =$$

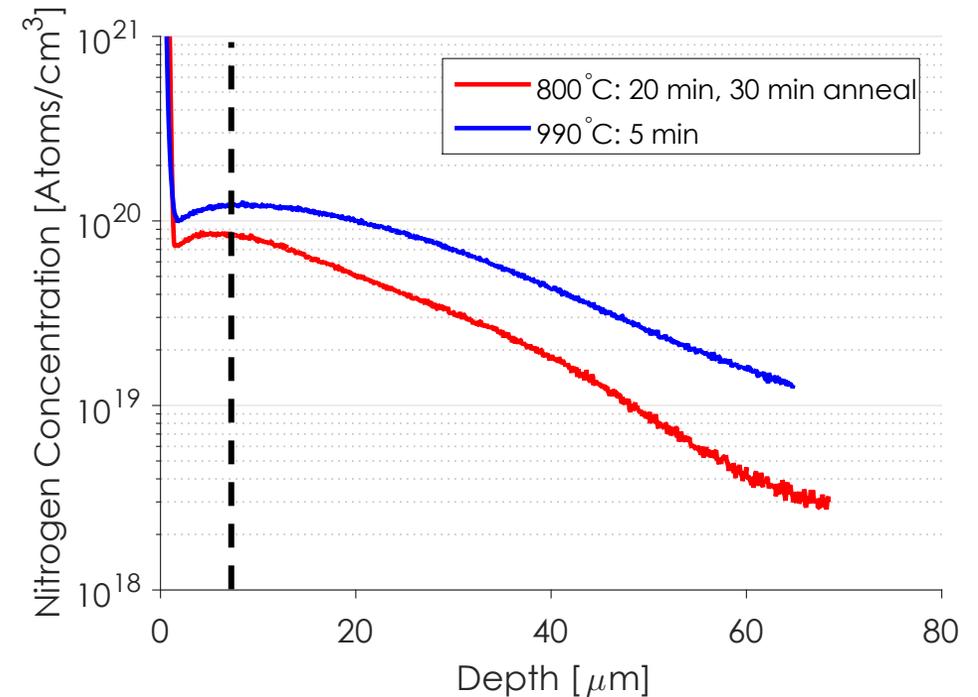
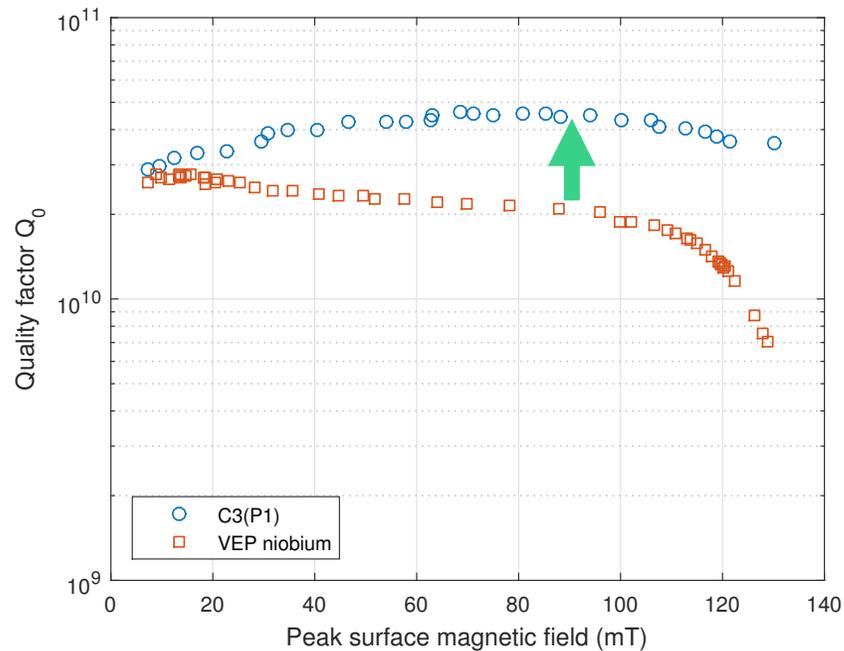
$$R_0(\text{slow cool}) - R_0(\text{fast cool})$$

$$S = R_0(B_{\text{trapped}}) / B_{\text{trapped}} = a \cdot B_{RF} + b$$

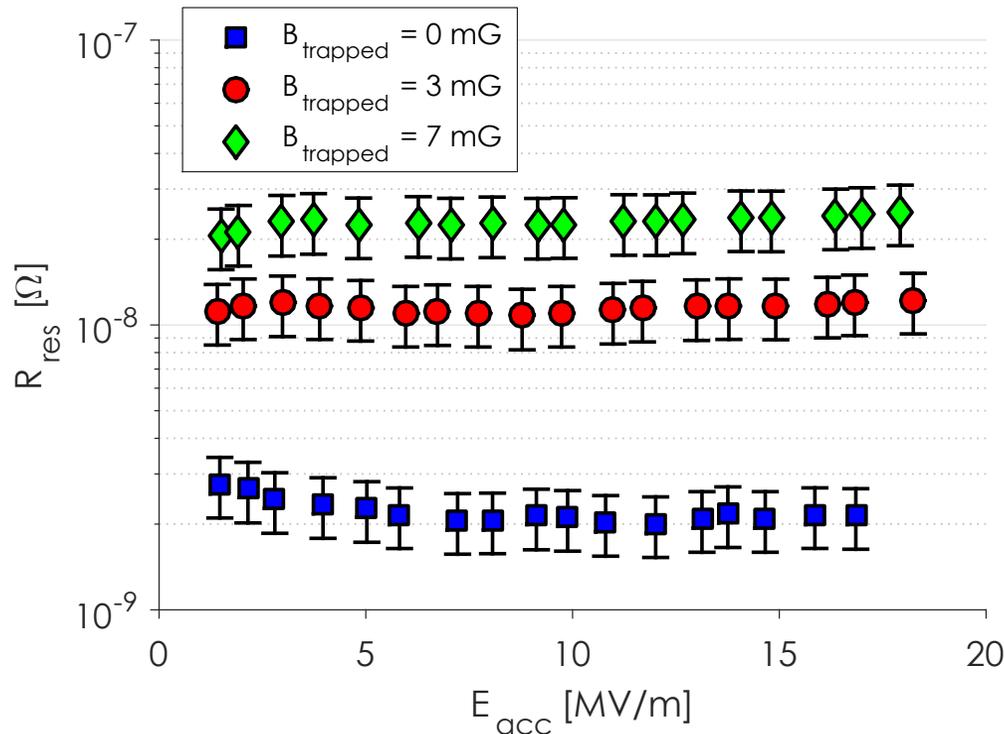


Multiple cooldowns, same sensitivity!

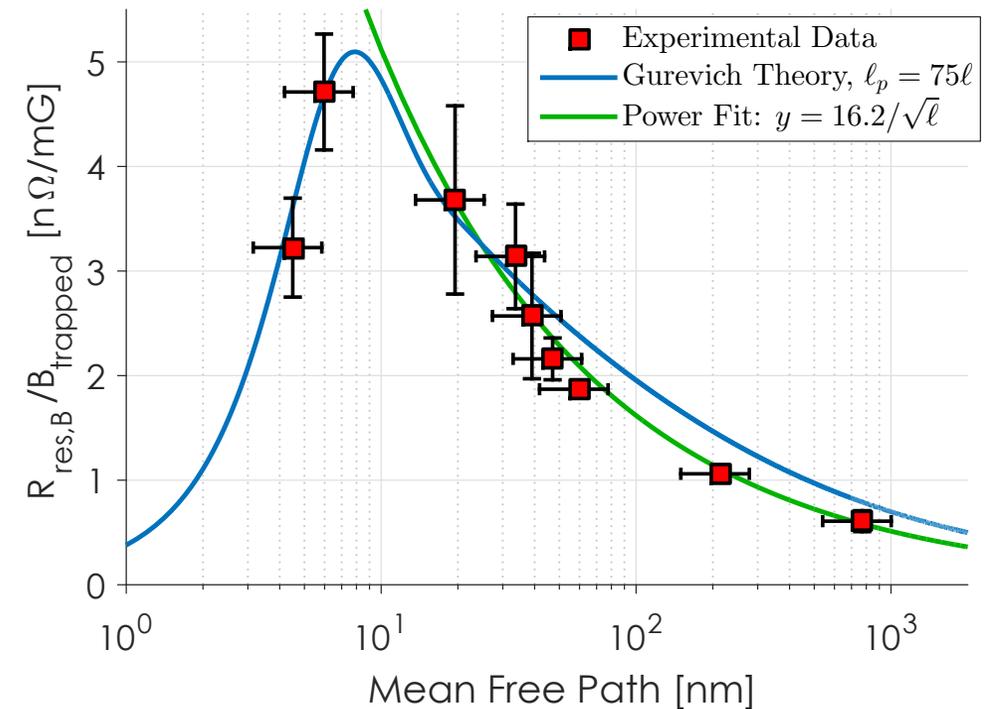
- Treatment protocol:
  - **800 °C UHV** degas bake (5-8 hr)
  - **800 °C 40 mTorr (6 Pa) N<sub>2</sub>** doping bake (2-60 minutes)
  - **800 °C UHV** anneal bake (0-60 minutes)
  - Final VEP (2-20 μm)



Strong “offset”,  
no slope in  $B_{RF}$

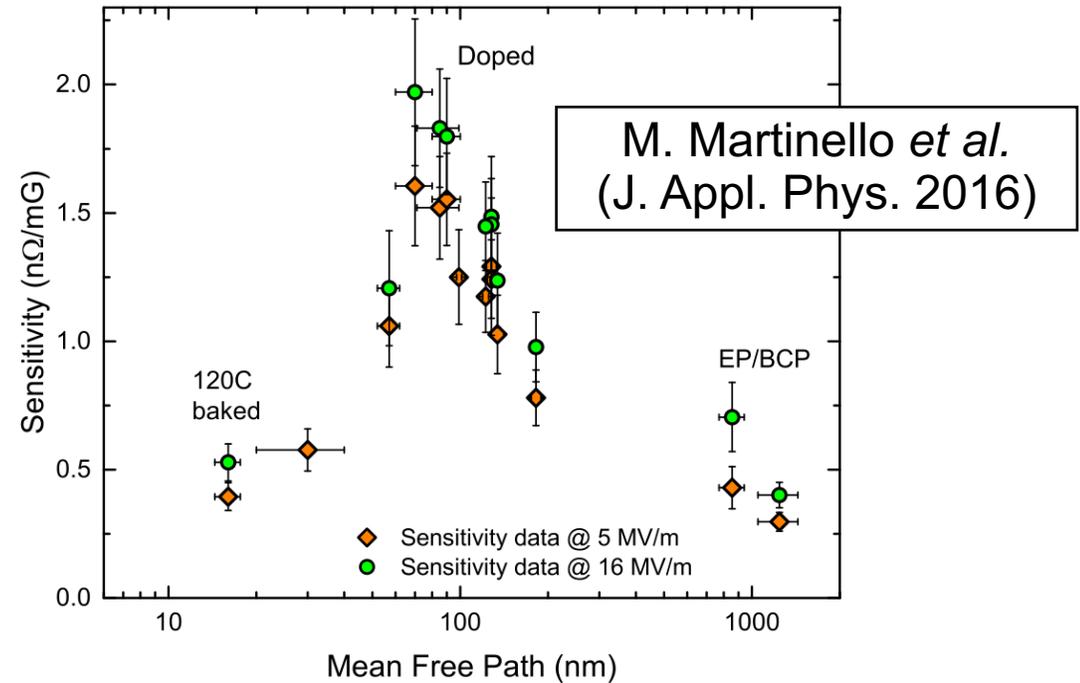
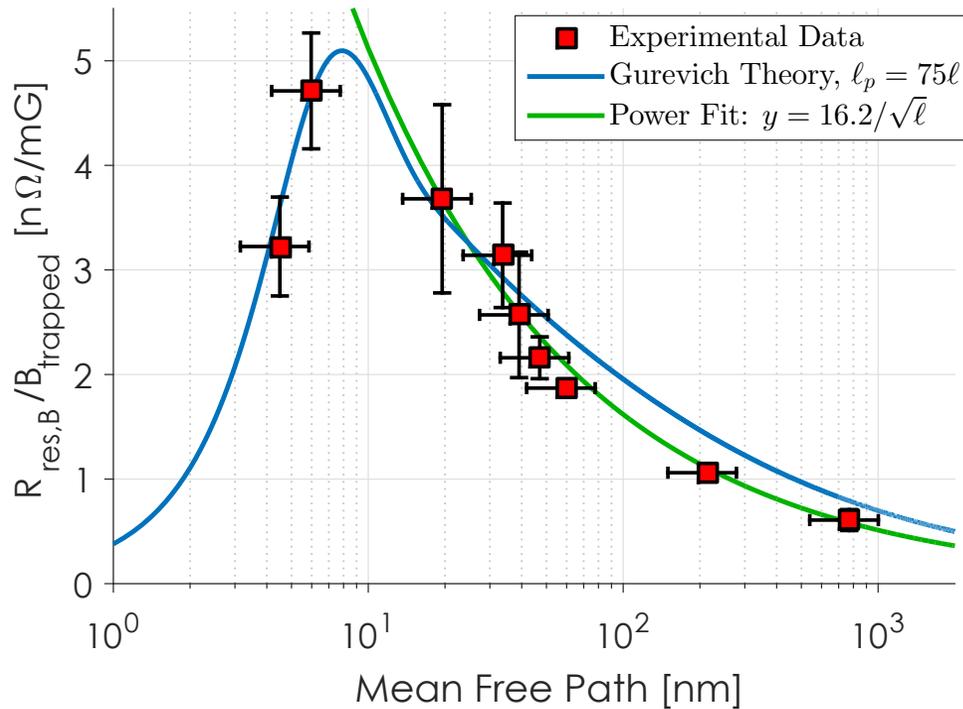


“offset” highly  
sensitive to  $\ell_e$



See D. Gonnella *et al.* (J. Appl. Phys. 2016)  
and J. T. Maniscalco *et al.* (J. Appl. Phys. 2017)

Sensitivity also likely dependent on properties of bulk Nb!  
(different cavity material at different labs)



Flux expulsion also linked to material stock in doped cavities!  
see D. Gonnella *et al.* (NIM-A 2018), S. Posen *et al.* (J. Appl. Phys. 2016)



## Takeaway message:

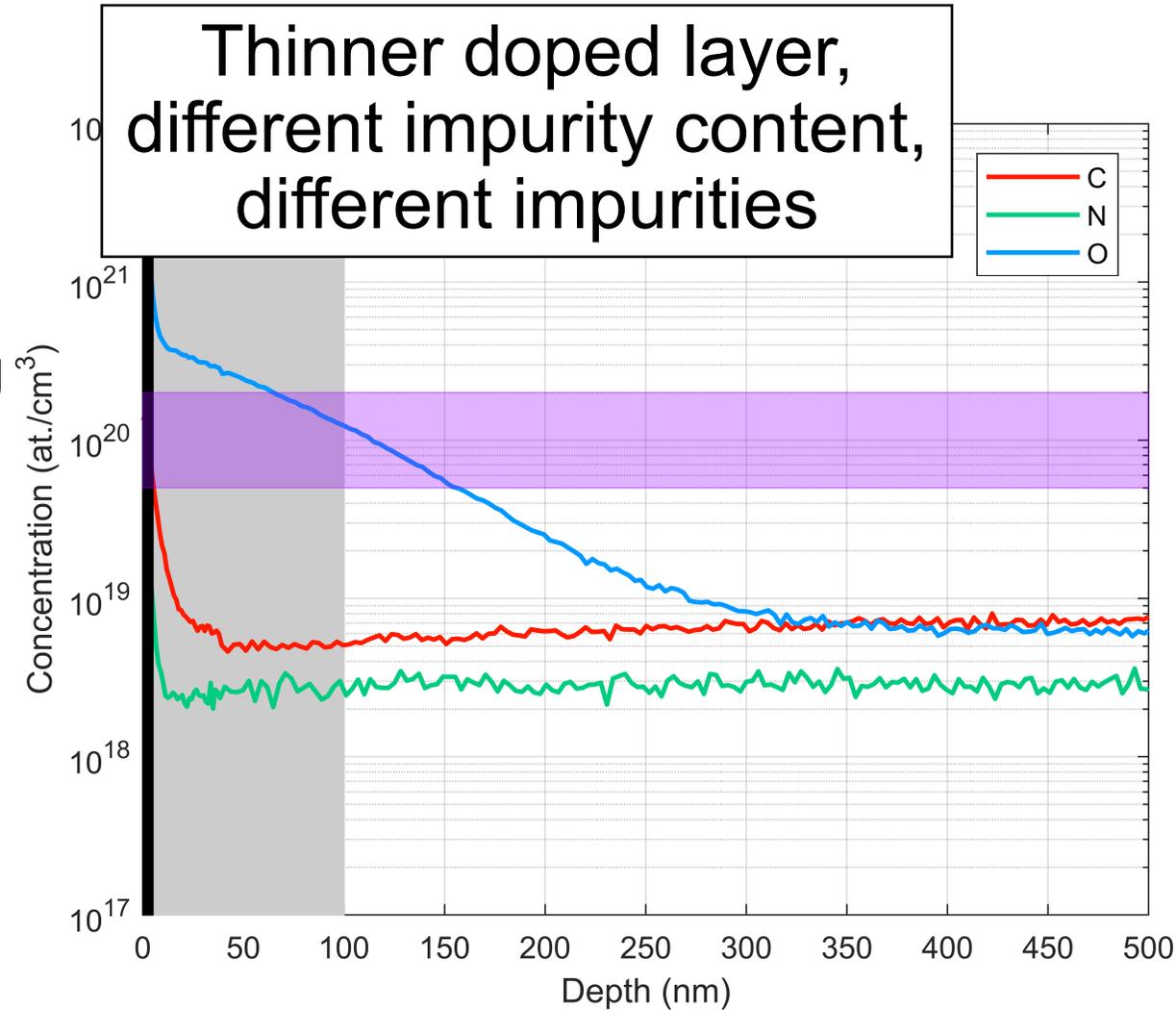
**Doping** strongly impacts trapped flux losses:

- 1) lowers the **mean free path**, which increases losses
- 2) increases **pinning effects** which decrease losses

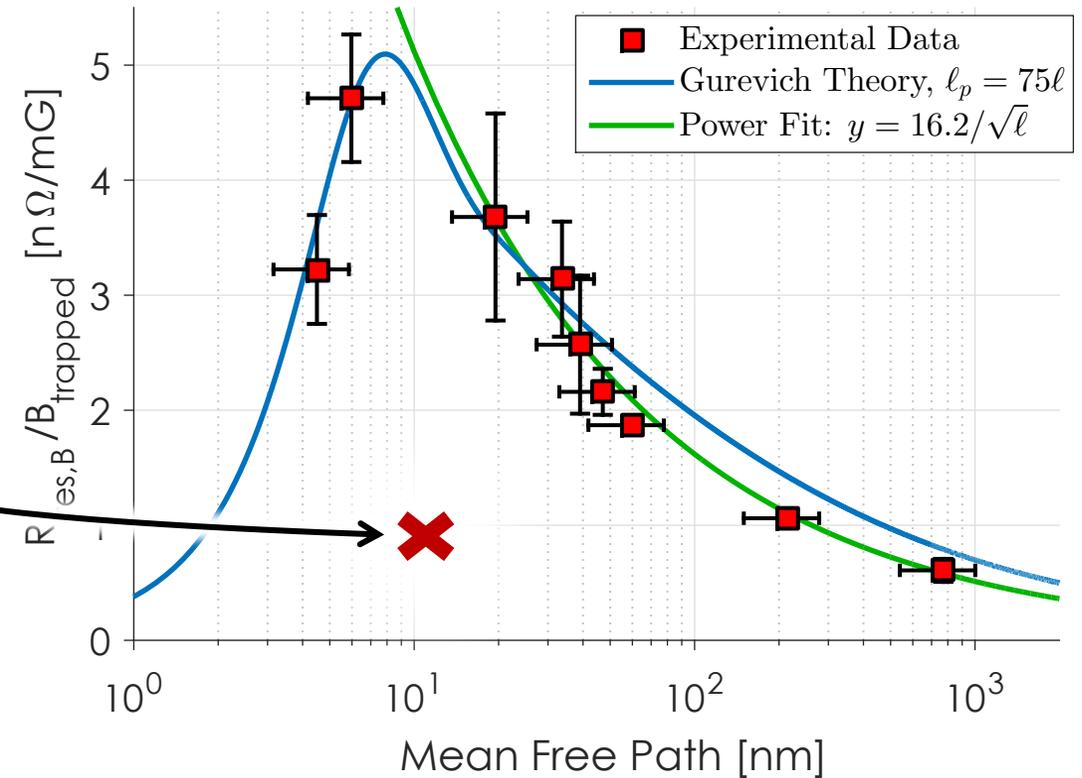
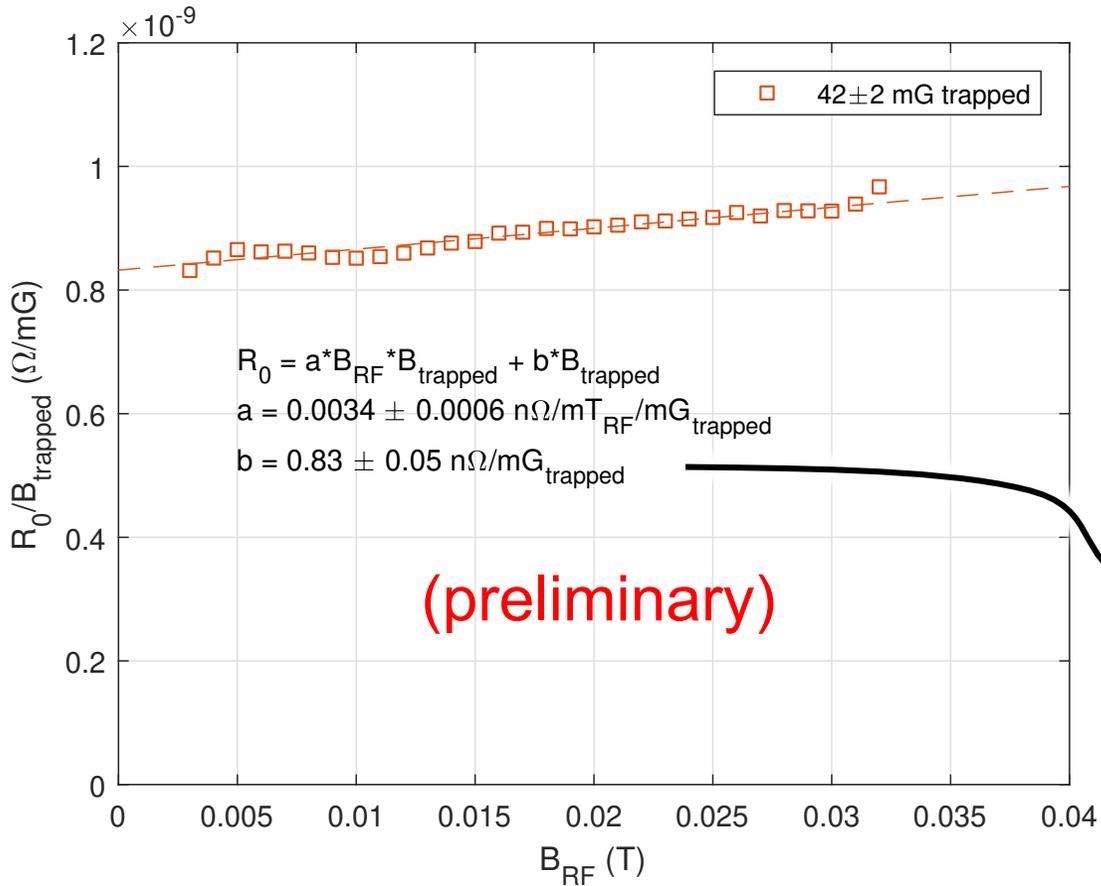
**Nitrogen's** link to pinning is still not well understood.  
Other material parameters might be relevant too!

(see Danilo's work on **collective weak pinning** – how material parameters could affect pinning strength and vortex loss magnitude)

- Treatment protocol:
  - **800 °C UHV** degas bake (5 hr)
  - **UHV** ramp down to **160 °C**
  - **160 °C UHV** rest (3 hr)
  - **160 °C, 40 mTorr (6 Pa) N<sub>2</sub>** doping bake (1 day for test shown here)
  - **160 °C** optional **UHV** anneal bake (not used for tests shown here)
- Little/no post chemistry! Maybe HF rinse or oxypolish for light surface removal  $\ll 1 \mu\text{m}$ .



Lower “offset” sensitivity, higher “slope” sensitivity compared to **800 °C**-doped cavities





Takeaway message:

Shallow **160 °C doping** gives lower sensitivity than **800 °C doping**  
Note: **dopant may be different...**

Results indicate **linear field dependence**  
more prominent for **low-temperature dopings**  
(see again Danilo's work/talk)

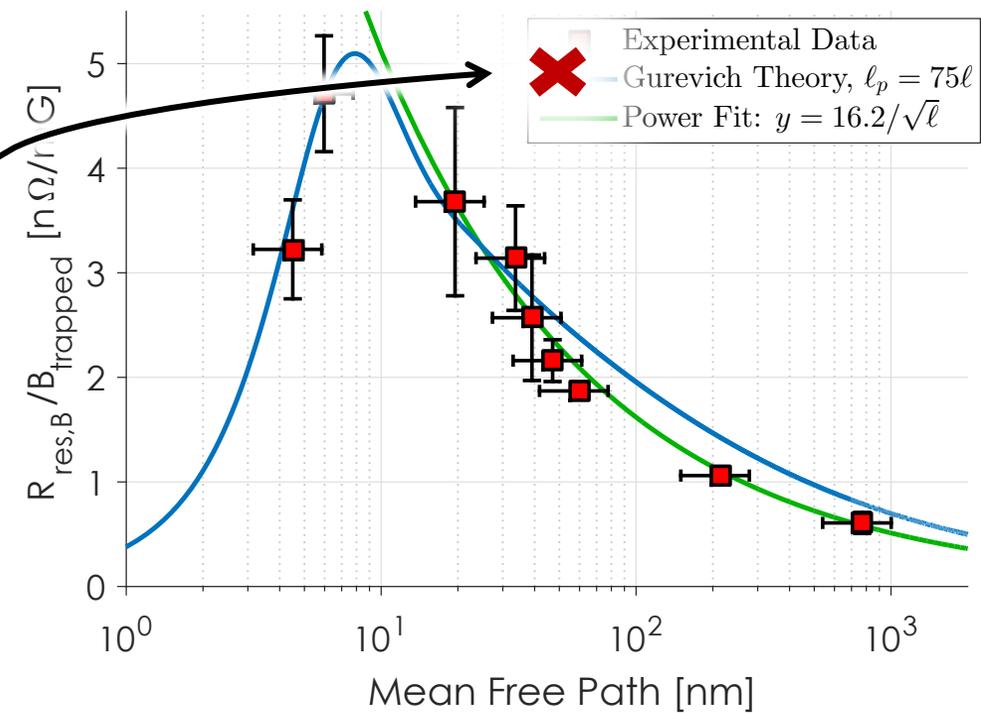
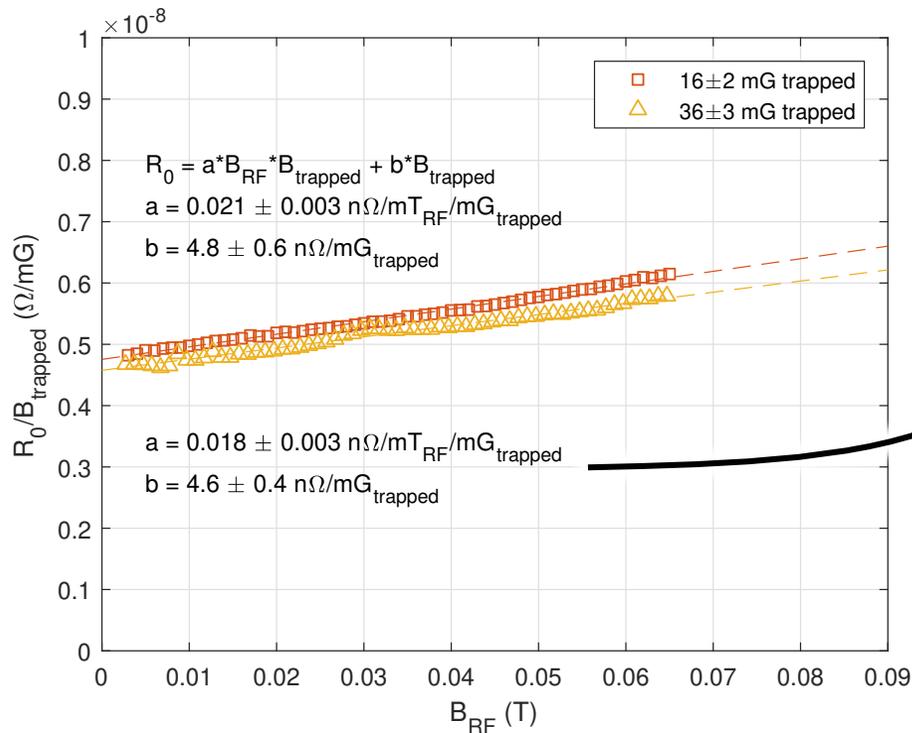


- Treatment protocol:
  - **800 °C UHV** degas bake (3 hr)
  - **800 °C 40 mTorr (6 Pa) N<sub>2</sub>** doping bake (2 minutes)
  - **800 °C UHV** anneal bake (6 minutes)
  - Final VEP (6 μm)



"offset" sensitivity **2x** higher than 1.3 GHz cavities with similar  $\ell_e$

$$R_0/B_{\text{trapped}} \propto \omega?$$



see J. T. Maniscalco *et al.* (LINAC 2018)



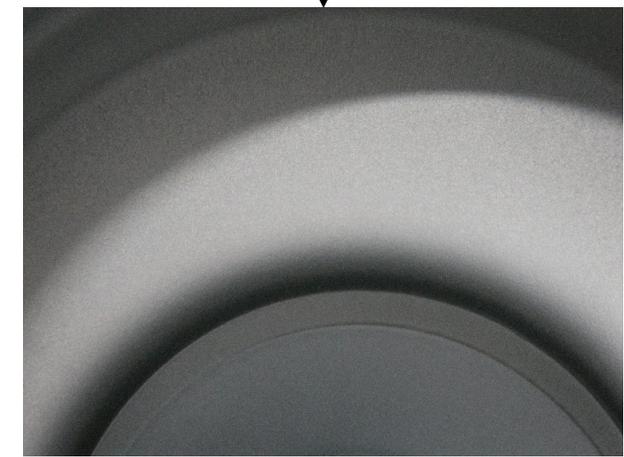
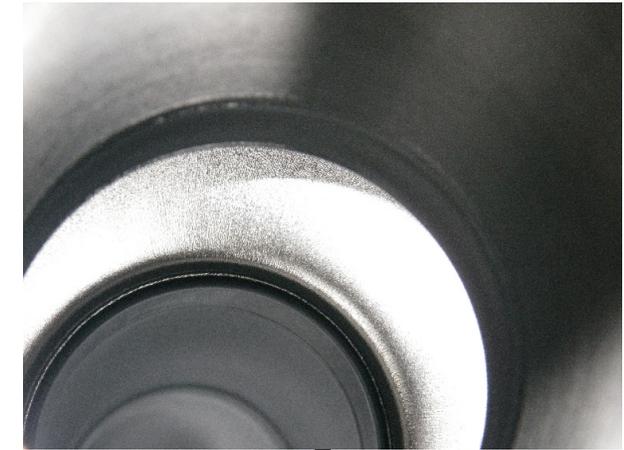
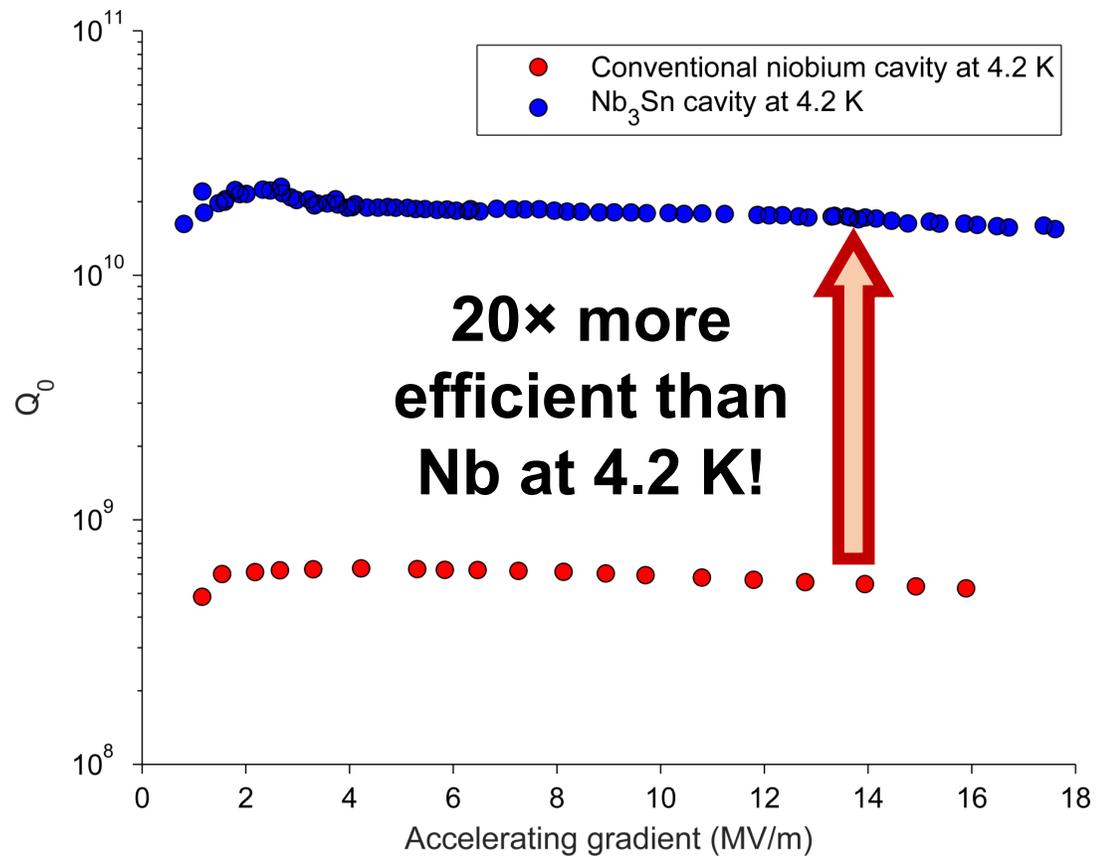
Takeaway message:

**High frequency** → stronger sensitivity

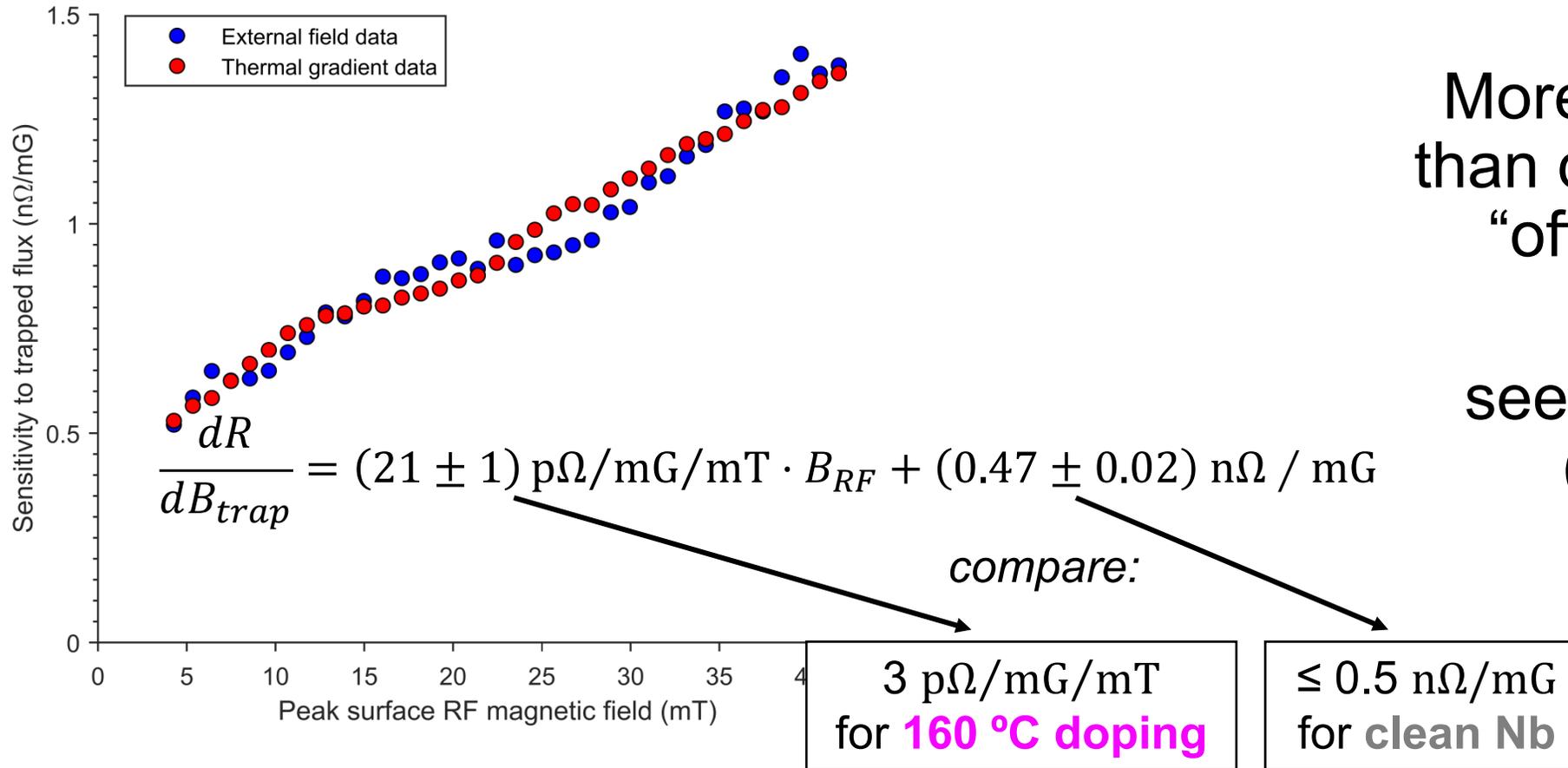
Initial results indicate linear behavior with frequency, *i.e.*

$$R_0/B_{\text{trapped}} \propto \omega$$

- Treatment protocol:
  - Nb<sub>3</sub>Sn grown on Nb cavity substrate by vapor diffusion



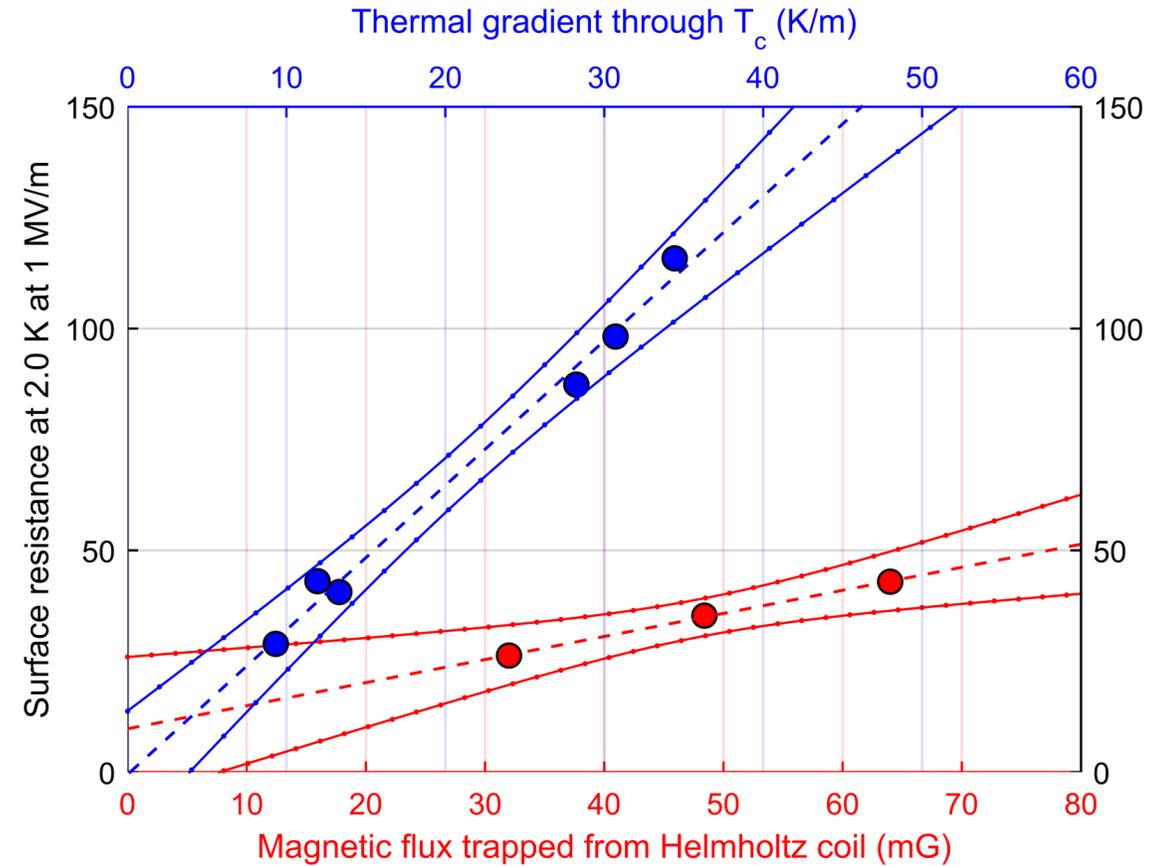
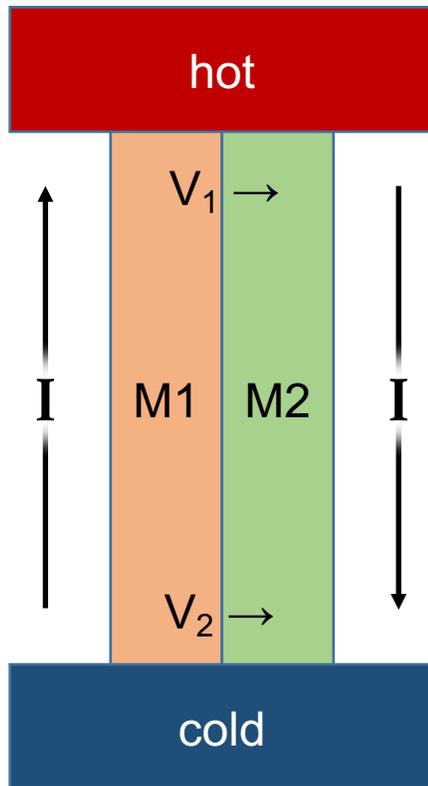
## Typical results: **linear sensitivity**



More sensitive slope than doped Nb, similar “offset” sensitivity

see D. L. Hall *et al.* (IPAC 2017)

Also sensitive to thermoelectric currents (Seebeck effect)



Cavity must be cooled uniformly!  
Need to limit flux sensitivity.



Takeaway message:

Clear **linear behavior**, strong slope

Low “offset” sensitivity – similar to undoped cavities





- Flux sensitivity typically **linear in RF field**
- Coefficients of sensitivity linked to material parameters
  - **Electron mean free path**
  - **Doping depth**
  - **Impurity species**
  - **Frequency**
  - Niobium stock
  - **Nb<sub>3</sub>Sn vs. doped Nb vs. clean Nb**  
(vs. Nb/Cu films not addressed here)
- Nb<sub>3</sub>Sn:  $R_0$  from trapped flux important because  $R_{BCS}$  is so small!
  - Higher “slope” sensitivity than doped Nb
  - Bimetallic interface necessitates slow, uniform cooling and good magnetic shielding/hygiene





# Thanks for your attention!

See also:

D. B. Liarte's talk later today (11h59):  
"Vortex dynamics and hysteretic flux losses due to pinning"

