

Rare decays, dark matter and baryon asymmetry

Mikhail Shaposhnikov

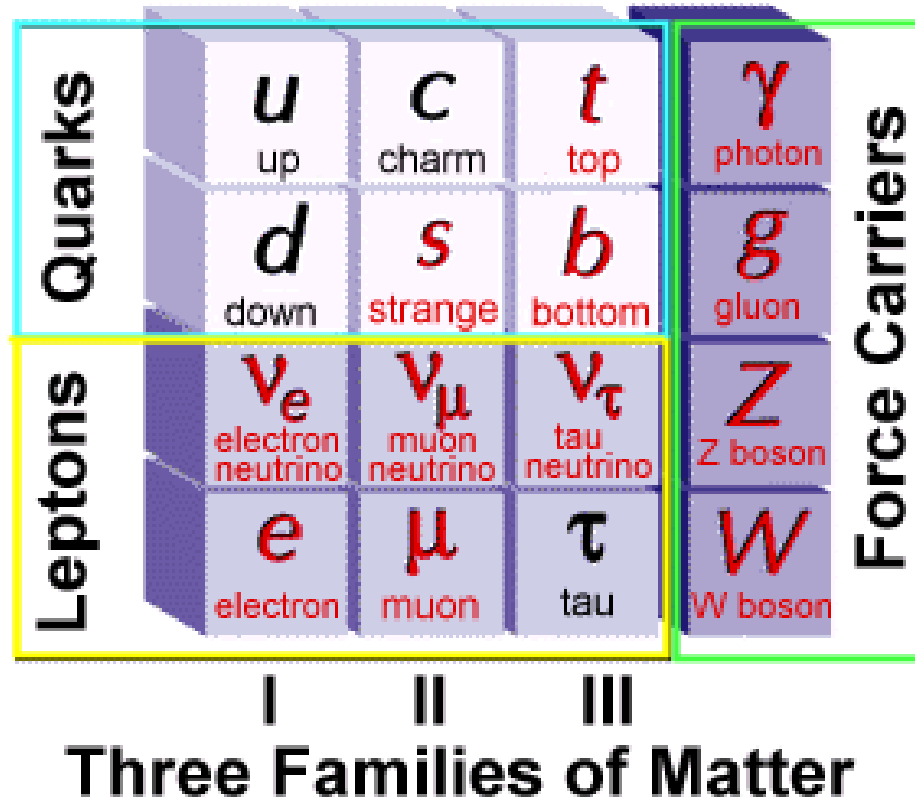


**The New, the Rare and the
Beautiful**

- Motivation
- Neutrino masses
- Dark Matter
- Baryon Asymmetry
- Rare decays
- Conclusions

Standard Model of particle interactions is in great shape: it agrees with all **accelerator** experiments

Elementary Particles



The only missing particle - the Higgs boson. It will be searched at the LHC

Still, the Standard Model cannot accommodate a number of cosmological observations and discoveries in neutrino physics, it also has a number of “fine tuning” problems from theory side.

A strategy

- Select the most important problems to solve (may be subjective).
- Use Ockham’s razor principle: “entities must not be multiplied beyond necessity”. For particle physics: entities = new hypothetical particles.

NATURALLY LIGHT DIRAC NEUTRINOS IN GAUGE THEORIES

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Received 23 September 1983

In order to have naturally small Dirac masses, extra symmetries (to be broken spontaneously) are clearly needed, and along with them, also new fields[‡].

[‡] This is true for a significantly better understanding of many phenomena; simple book-keeping arguments (fewest fields) should be mistrusted.

Experimental evidence in favour of new physics

- 0.511 MeV annihilation line in the direction of the Galaxy center
- Baryon asymmetry of the Universe
- Cosmological inflation
- DAMA annual modulations
- Dark energy
- Dark matter in the Universe
- Discrepancy between experiment and the theory prediction of anomalous magnetic moment of muon
- Egret gamma-ray excess
- Heidelberg neutrinoless double β decays
- LSND anomaly (light sterile neutrino?)
- MiniBooNE anomaly (excess of low energy events)
- Neutrino masses and oscillations
- Origin of high energy cosmic rays
- Pamela positron excess
- Pulsar-kick velocities
- ...

Confirmed experimental evidence in favour of new physics

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Fine-tuning problems

- Cosmological constant problem: Why $\epsilon_{vac}/M_{Pl}^4 \lll 1$?
- Hierarchy problem: Why $M_W/M_{Pl} \ll 1$?
- Stability of the Higgs mass against radiative corrections.
- Strong CP-problem: Why $\theta_{QCD} \ll 1$?
- Fermion mass matrix: Why $m_e \ll m_t$?
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Will concentrate on:

- Neutrino masses and oscillations
- Dark matter
- Baryon asymmetry of the Universe
- Inflation

**New dedicated experiments
in flavour physics**

The ν MSM

Neutrino masses and oscillations, Dark matter, and Baryon asymmetry of the Universe can be explained if the Standard Model is extended in a minimal way by adding 3 right-handed neutrinos (other names: Majorana fermions, singlet leptons, neutral fermions)

SM \Rightarrow ν MSM

ν MSM = Neutrino Minimal Standard Model

SM fermions

quarks

left	u	d	c	s	t	b
right	u	d	c	s	t	b
left	ν_e	e	ν_μ	μ	ν_τ	τ
right		e		μ		τ

leptons

ν MSSM fermions

quarks

left	u	d	c	s	t	b
right	u	d	c	s	t	b
left	ν_e	e	ν_μ	μ	ν_τ	τ
right	N_e	e	N_μ	μ	N_τ	τ

leptons

Role of N_e with mass in keV region: dark matter

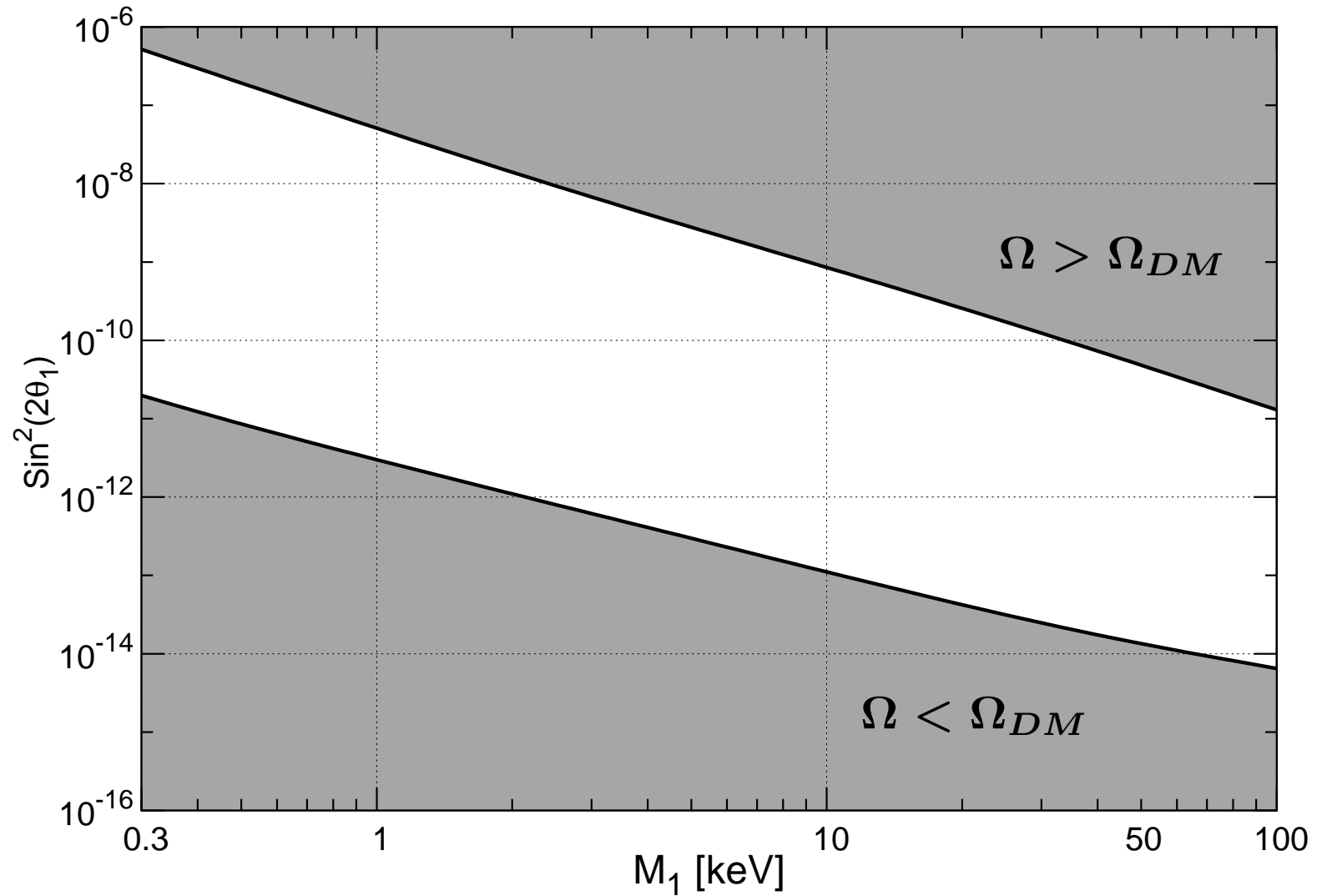
Role of N_μ , N_τ with mass in 100 MeV – GeV region: “give” masses to neutrinos and produce baryon asymmetry of the Universe

Automatic electric charge quantisation without Grand Unification: from requirement of gravitational and gauge anomalies cancellation.

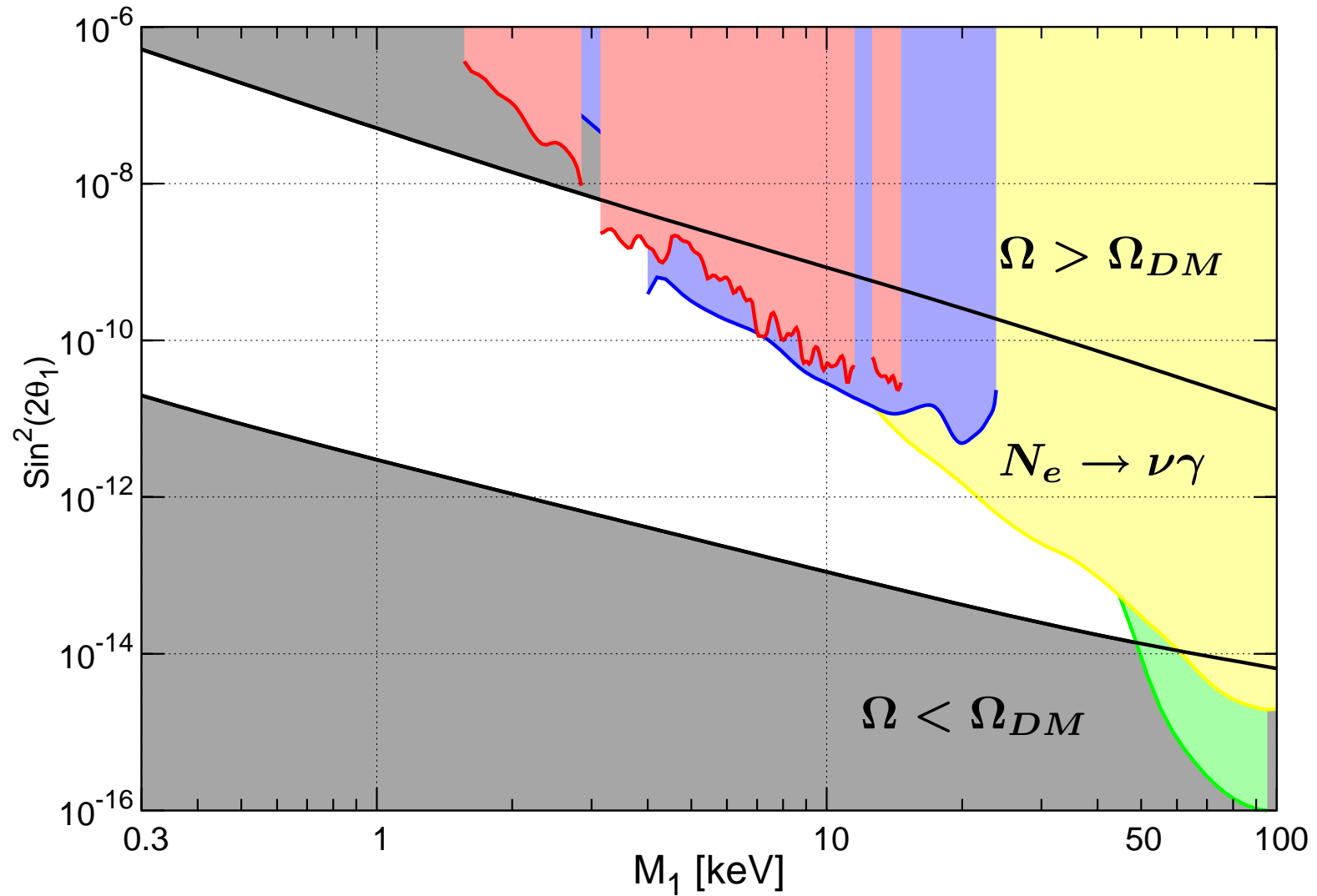
Constraints on DM sterile neutrino

- **Production.** N_e are created in the early Universe in reactions $l\bar{l} \rightarrow \nu N_e$, $q\bar{q} \rightarrow \nu N_e$ etc. We should get correct DM abundance.
- **X-rays.** N_e decays radiatively, $N_e \rightarrow \gamma\nu$, producing a narrow line which can be detected. This line has not been seen (yet).
- **Structure formation.** If N_e is too light it may have considerable free streaming length and erase fluctuations on small scales. This can be checked by the study of Lyman- α forest spectra of distant quasars.

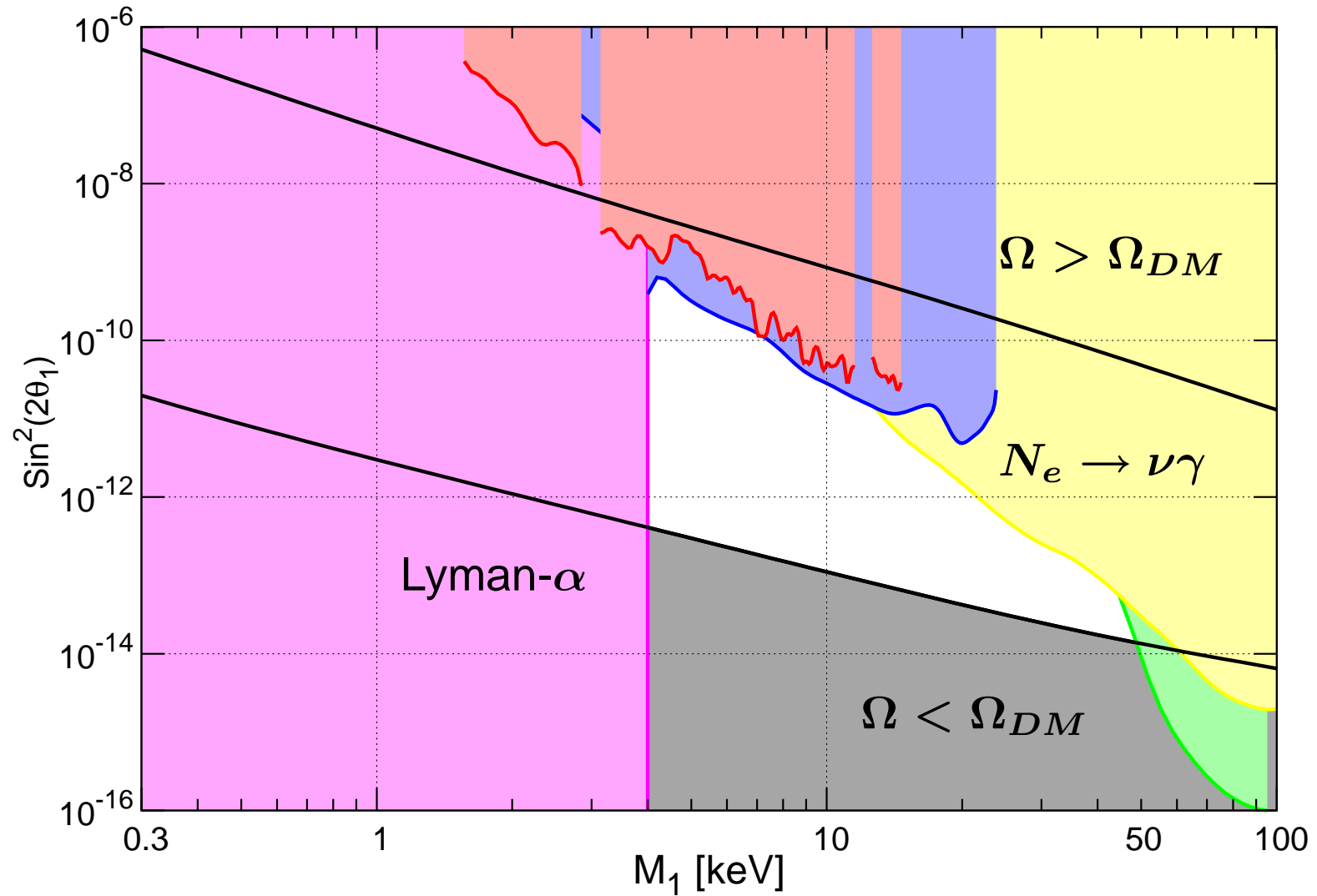
DM: production



DM: production + X-ray constraints



DM: production + X-ray constraints + Lyman- α bounds



Direct experimental test: astrophysics

Over the last years restrictions on sterile neutrino parameters were improved **by several orders of magnitude**.

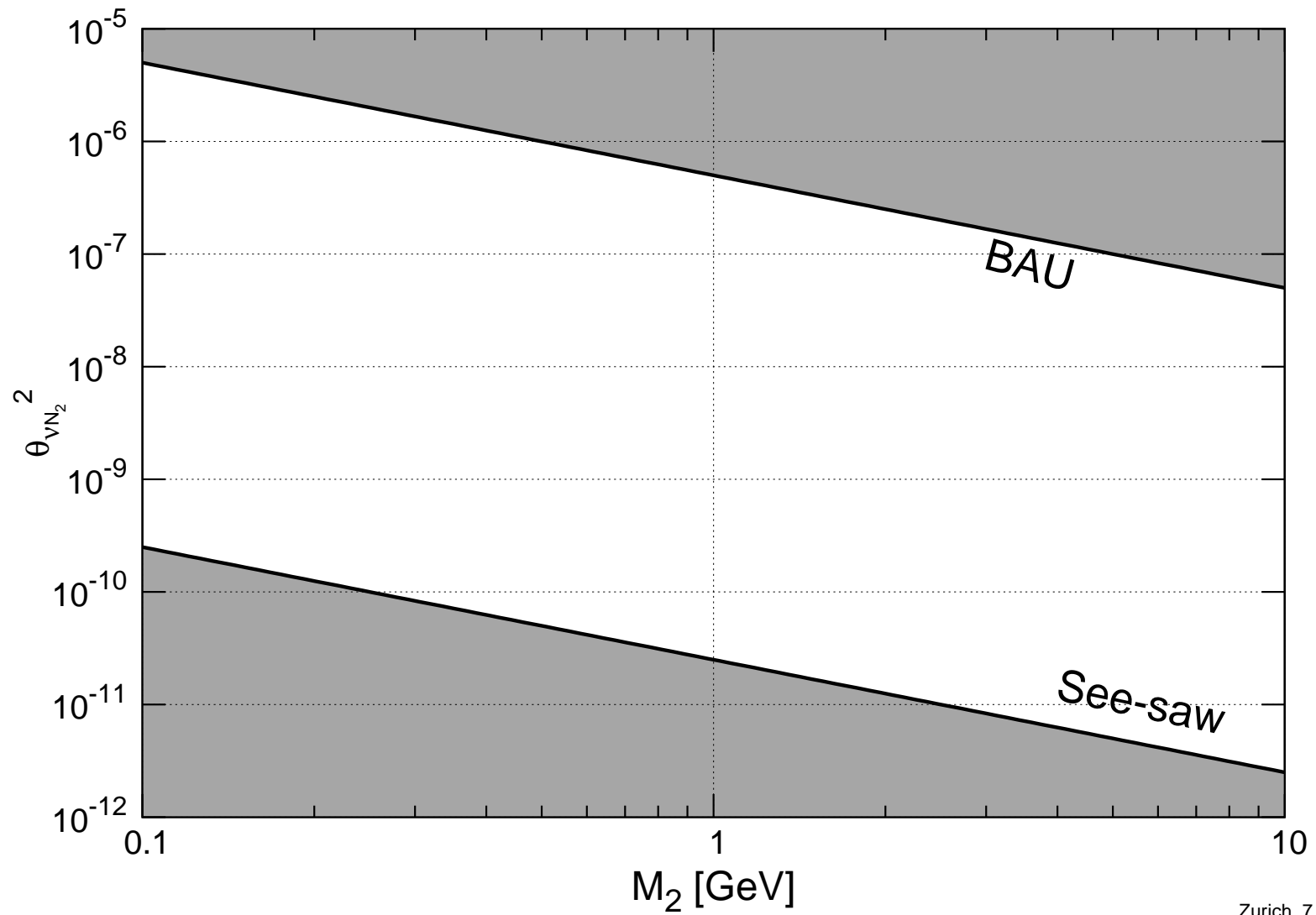
The new data from *Chandra* and *XMM-Newton* can hardly improve constraints by more than a factor 10. One needs:

- Improvement of spectral resolution up to the natural line width ($\Delta E/E \sim 10^{-3}$).
- $\text{FoV} \sim 1^\circ$ (size of a dSph).
- Wide energy scan, from $\mathcal{O}(100)$ eV to $\mathcal{O}(10)$ MeV.

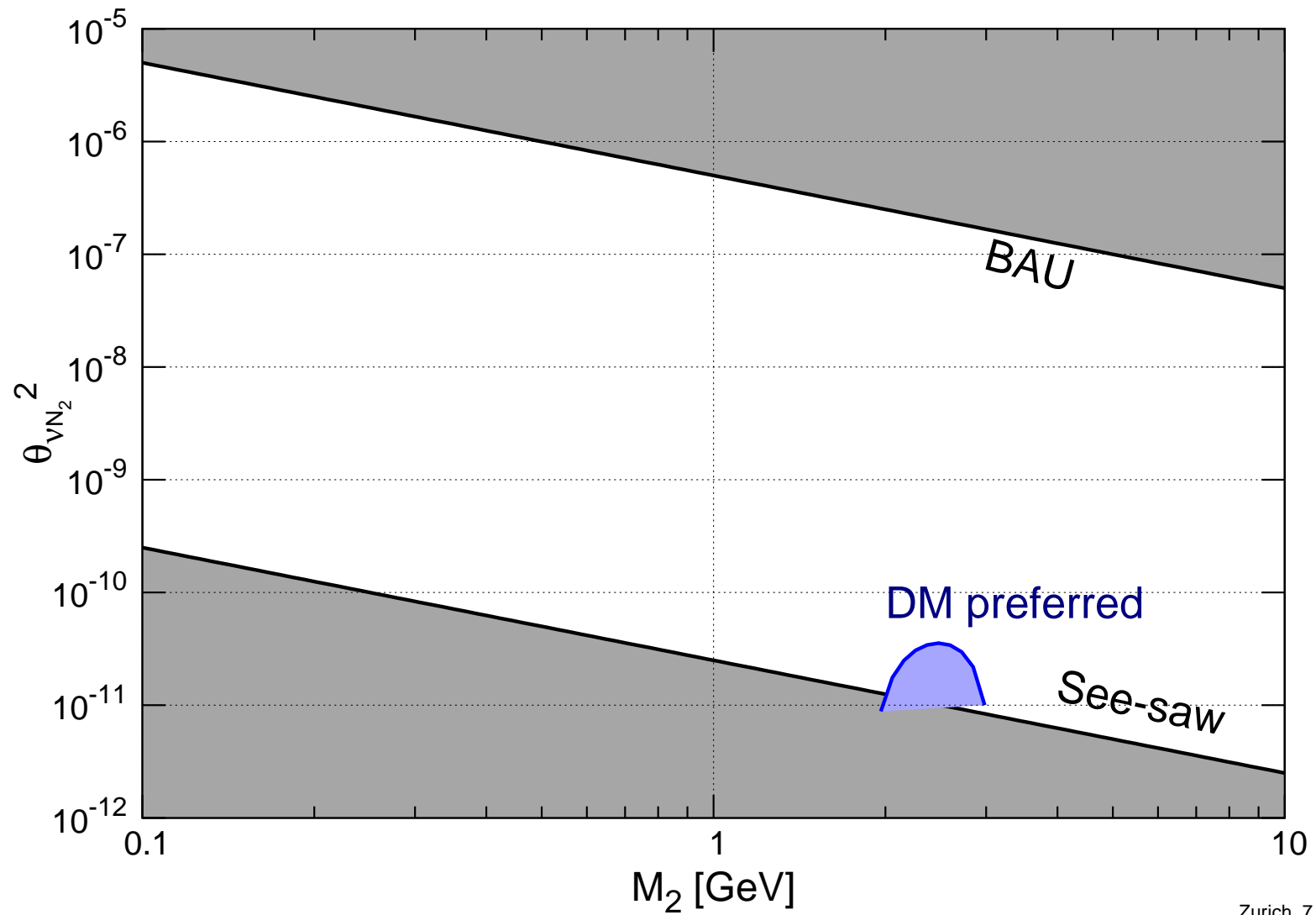
Constraints on BAU Majorana fermions

- **BAU generation** requires out of equilibrium: mixing angle of $N_{\mu,\tau}$ to active neutrinos cannot be too large
- **Neutrino masses.** Mixing angle of $N_{\mu,\tau}$ to active neutrinos cannot be too small
- **Dark matter and BAU.** Concentration of DM sterile neutrinos must be much larger than concentration of baryons
- **BBN.** Decays of $N_{\mu,\tau}$ must not spoil Big Bang Nucleosynthesis
- **Experiment.** $N_{\mu,\tau}$ have not been seen (yet).

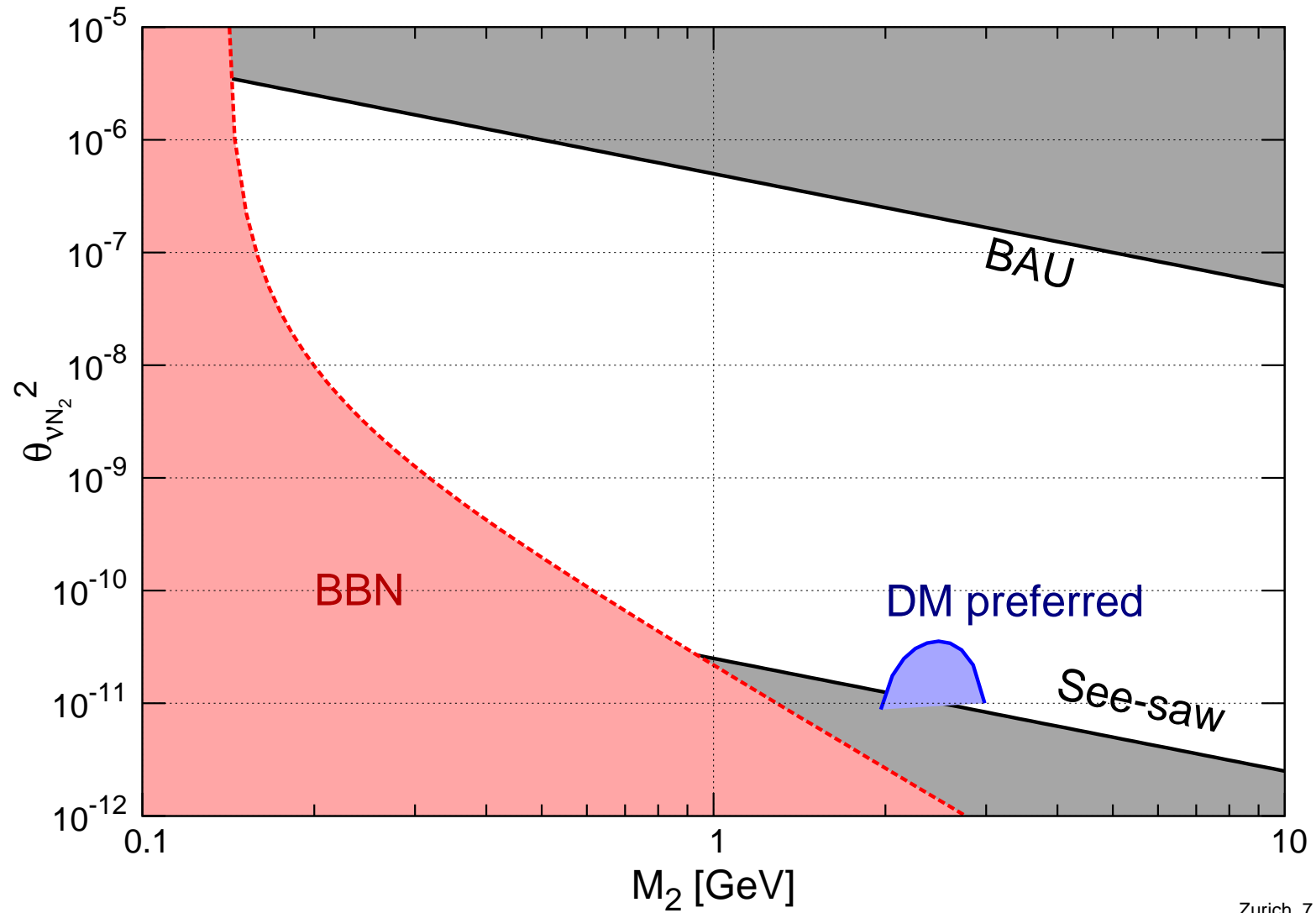
$N_{\mu,\tau}$: BAU



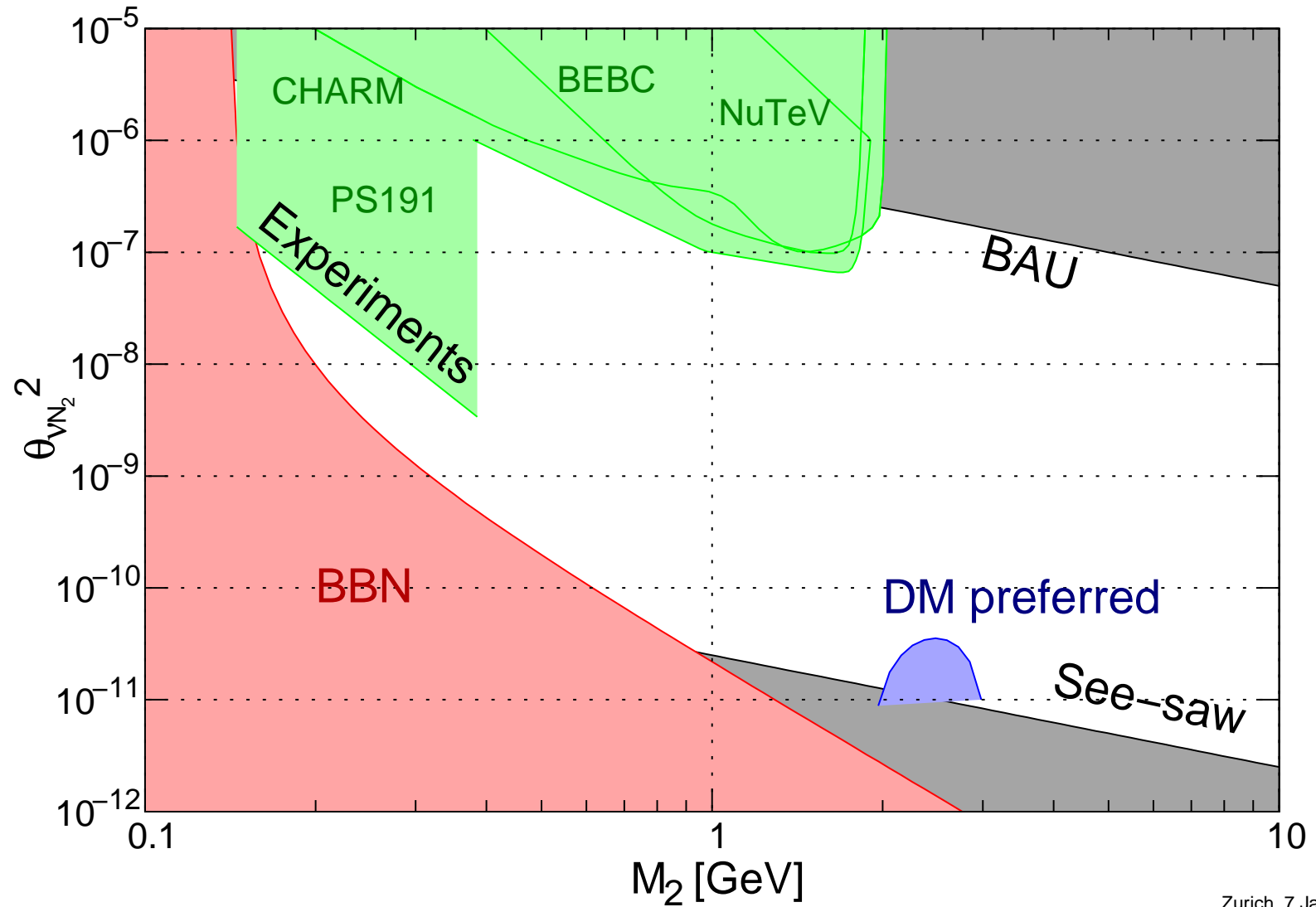
$N_{\mu,\tau}$: BAU + DM



$N_{\mu,\tau}$: BAU + DM + BBN



$N_{\mu,\tau}$: BAU + DM + BBN + Experiment

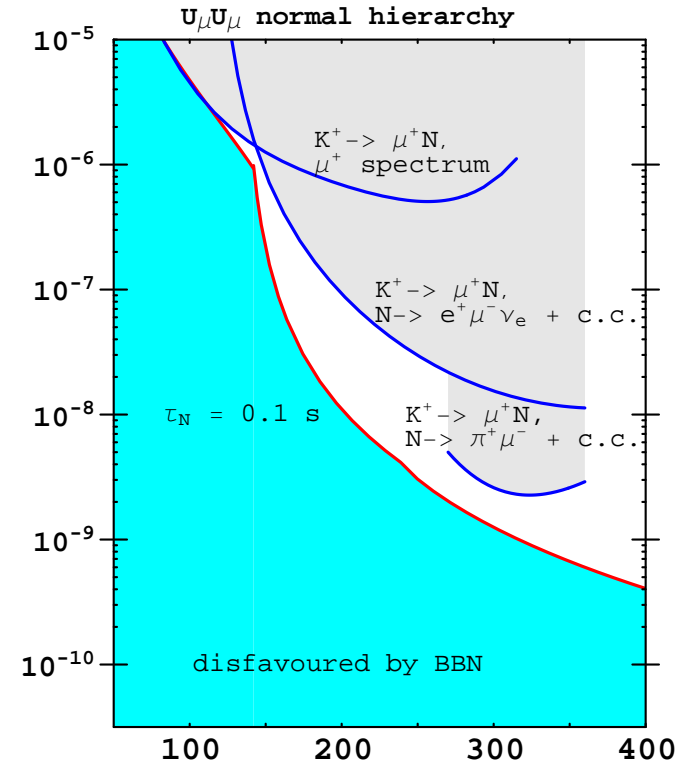
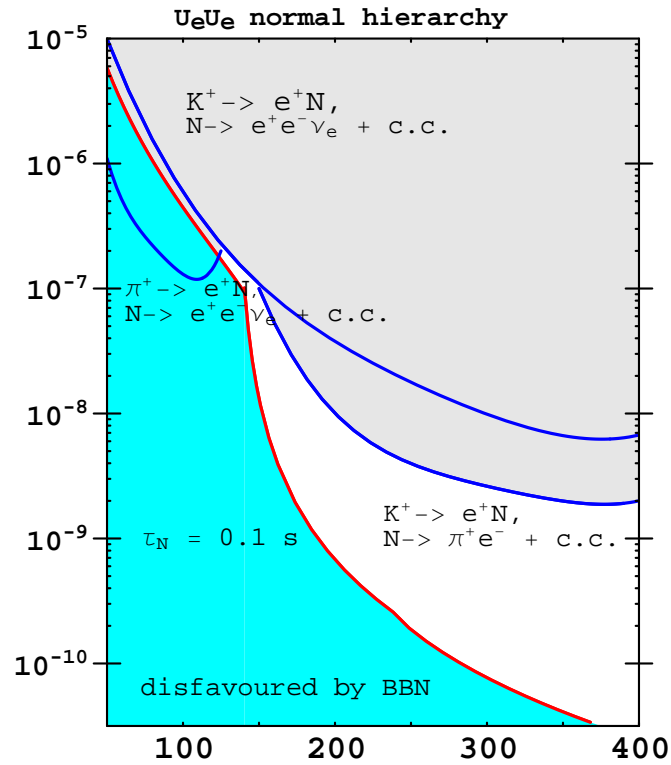
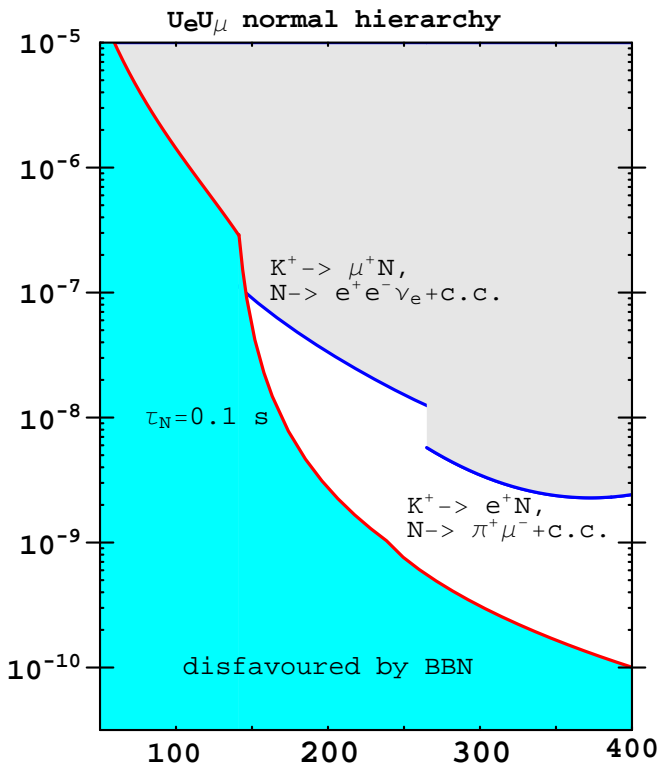


Direct experimental tests: rare decays

Previous searches at CERN

- A. M. Cooper-Sarkar *et al.* [WA66 Collaboration] “Search For Heavy Neutrino Decays In The Bebc Beam Dump Experiment”, 1985
- J. Dorenbosch *et al.* [CHARM Collaboration] “A search for decays of heavy neutrinos in the mass range 0.5-GeV to 2.8-GeV”, 1985
- G. Bernardi *et al.* [PS191 Collaboration], “Search For Neutrino Decay”, 1986;
“Further Limits On Heavy Neutrino Couplings”, 1988
- P. Astier *et al.* [NOMAD Collaboration], “Search for heavy neutrinos mixing with tau neutrinos”, 2001
- P. Achard *et al.* [L3 Collaboration], “Search for heavy neutral and charged leptons in e^+e^- annihilation at LEP”, 2001

CERN PS191 experiment, 1988



Conclusion: $M_{2,3} > 140 \text{ MeV}$

Experimental signatures 1

Challenge - from baryon asymmetry: $\theta^2 \lesssim 5 \times 10^{-7} \left(\frac{\text{GeV}}{M}\right)$

- Peak from 2-body decay and missing energy signal from 3-body decays of K , D and B mesons (sensitivity θ^2)

Example:

$$K^+ \rightarrow \mu^+ N, \quad M_N^2 = (p_K - p_\mu)^2 \neq 0$$

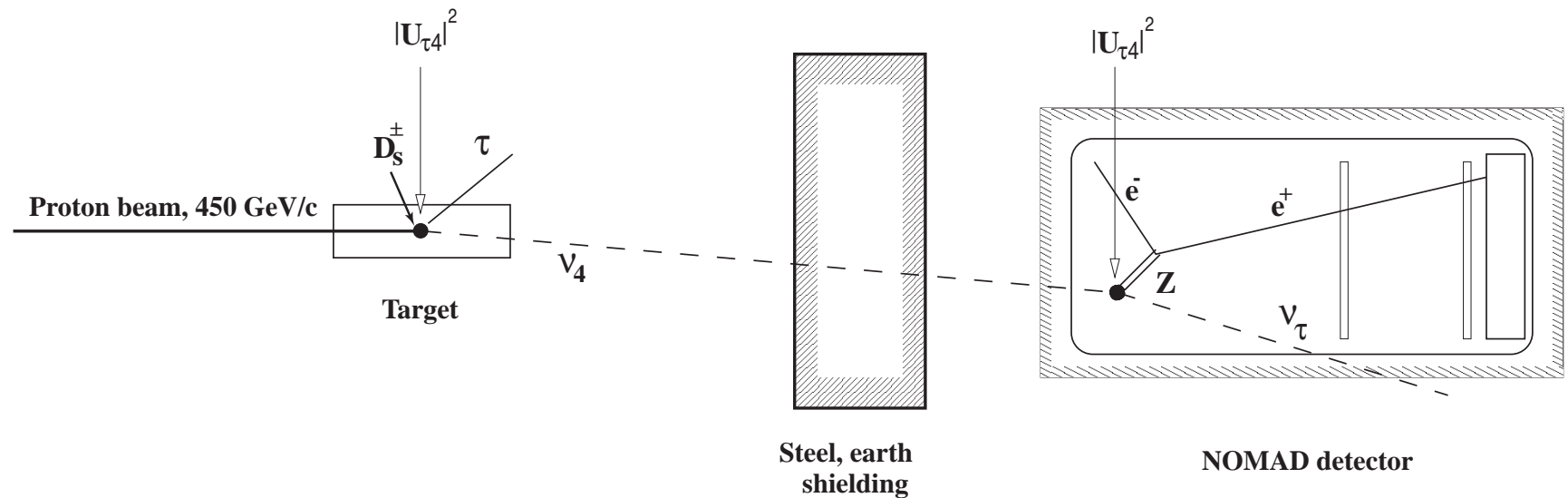
Similar for charm and beauty.

- $M_N < M_K$: KLOE, NA62, E787
- $M_K < M_N < M_D$: charm and τ factories, CLEO
- $M_N < M_B$: B-factories (planned luminosity is not enough to get into cosmologically interesting region)

Experimental signatures 2

- Two charged tracks from a common vertex, decay processes $N \rightarrow \mu^+ \mu^- \nu$, etc. (sensitivity $\theta^4 = \theta^2 \times \theta^2$)
First step: proton beam dump, creation of N in decays of K , D or B mesons: θ^2
Second step: search for decays of N in a near detector, to collect all N s: θ^2
 - $M_N < M_K$: Any intense source of K-mesons (e.g. from proton targets of PS.)
 - $M_N < M_D$: SPS or PS2 beam + near detector
 - $M_N < M_B$: Project X (?) + near detector
 - $M_N > M_B$: extremely difficult

$N_{2,3}$ production and decays



Type on neutrino mass hierarchy - from branching ratios of $N_{2,3}$ decays to e, μ, τ .

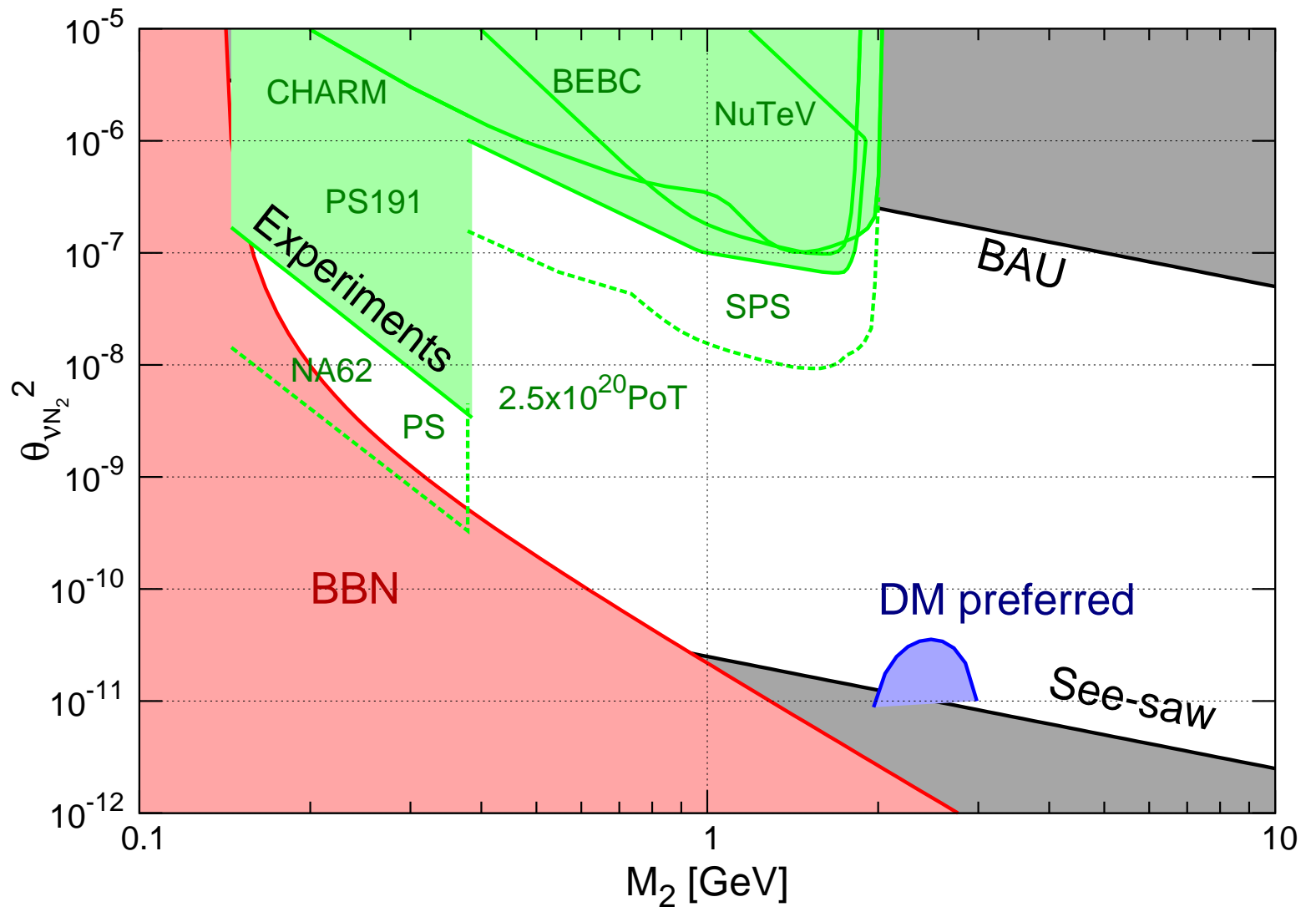
CP asymmetry can be as large as 1% - from BAU and DM

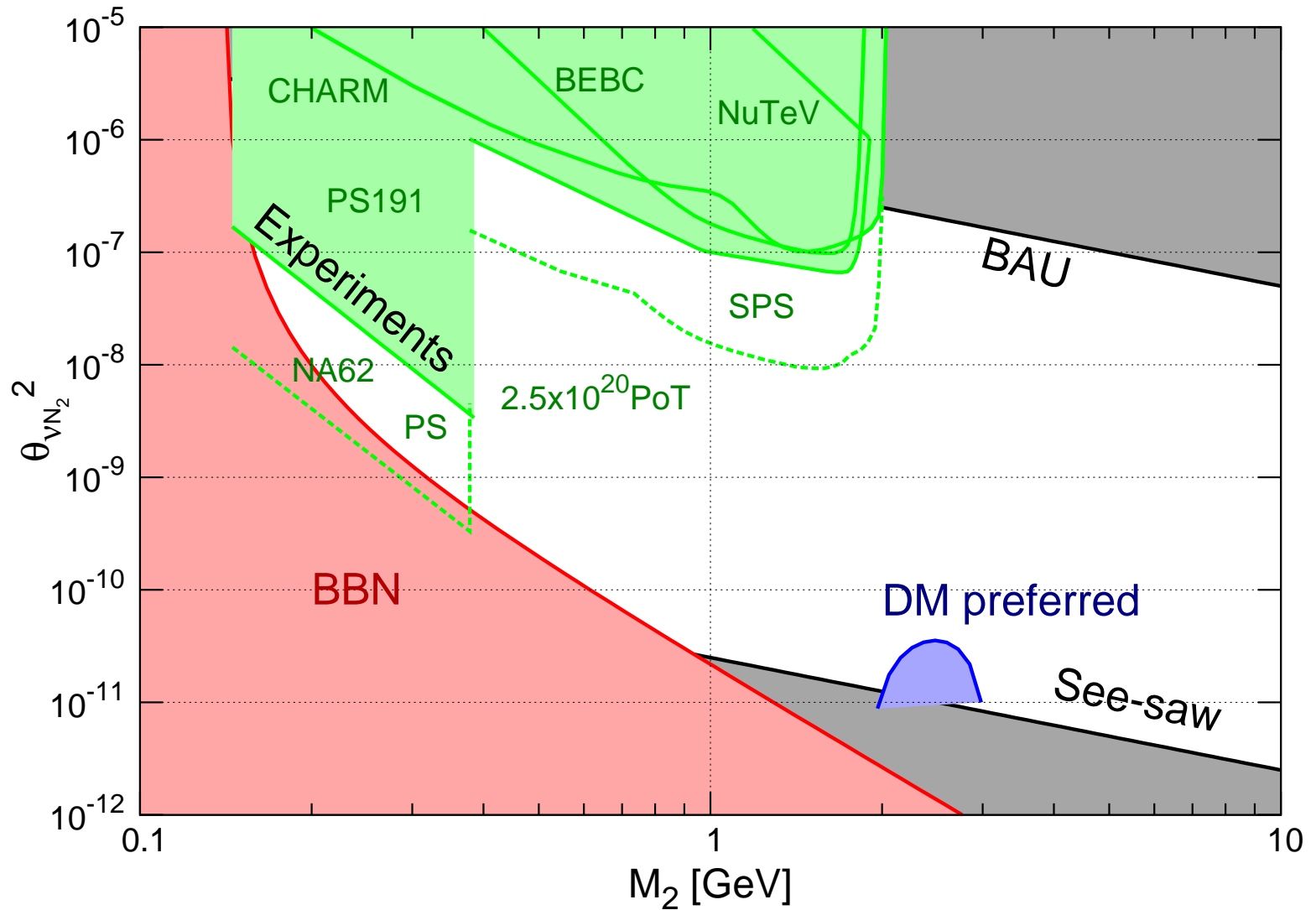
Conclusions

- New physics, responsible for **neutrino masses and mixings**, for **dark matter**, and for **baryon asymmetry of the Universe** may hide itself **below** the EW scale

Conclusions

- New physics, responsible for **neutrino masses and mixings**, for **dark matter**, and for **baryon asymmetry of the Universe** may hide itself **below** the EW scale
- The estimates based on PS191, BEBC and CHARM results show that the experiments with a sensitivity in mixing strength θ^2 improved by a factor more than 10 are feasible. The coverage of significant amount of parameter space would require from CERN a delivery of more than 10^{20} POTs





The efforts for performing the dedicated experiments, pioneered by CERN almost 25 year ago, are well justified, and, hopefully, may lead these searches to the happy end.

ν MSM prediction for LHC

Cosmological inflation from the Higgs of the Standard Model

(Bezrukov, MS):

Nothing but the SM Higgs in the mass interval $M_H \in [m_{min}, m_{max}]$

$$m_{min} = [126.1 + \frac{m_t - 171.2}{2.1} \times 4.1 - \frac{\alpha_s - 0.1176}{0.002} \times 0.6] \text{ GeV} ,$$

$$m_{max} = [193.9 + \frac{m_t - 171.2}{2.1} \times 0.6 - \frac{\alpha_s - 0.1176}{0.002} \times 0.1] \text{ GeV} .$$

Theoretical uncertainty ± 2.2 GeV.

If gravity is asymptotically safe, then (MS, Wetterich): $M_H = m_{min}$