

A gauged horizontal $SU(2)$ symmetry and R_κ

Diego Guadagnoli
CNRS, LAPTh Annecy

Based on work with M. Reboud and O. Sumensari

Several references on horizontal symmetries for B anomalies

E.g.

[Crivellin, D'Ambrosio, Heeck, PRD2015]

[Alonso, Cox, Han, Yanagida, PRD2017]

[Cline, Martin Camalich, PRD2017]

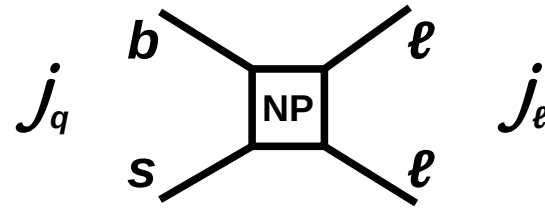
However, theory arguments quite distant from the one pursued here

$b \rightarrow s$ anomalies' basic challenge

- $R_K \approx 0.75$



$O(15-25\%)$ effects in

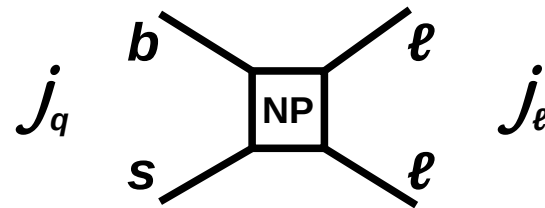


$b \rightarrow s$ anomalies' basic challenge

- $R_K \approx 0.75$

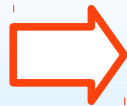


$O(15-25\%)$ effects in

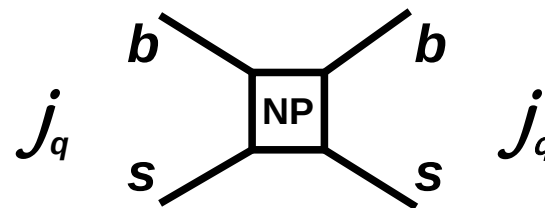


At the same time:

- $\Delta M_s \approx (\Delta M_s)_{SM}$



small corrections to

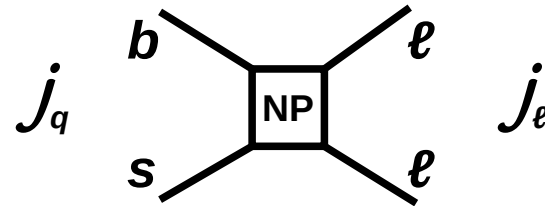


$b \rightarrow s$ anomalies' basic challenge

- $R_K \approx 0.75$



$O(15-25\%)$ effects in

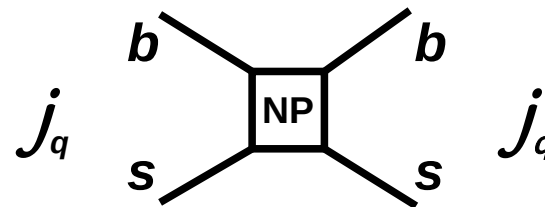


At the same time:

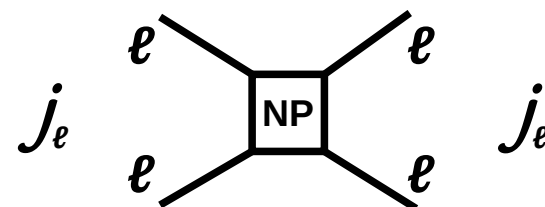
- $\Delta M_s \approx (\Delta M_s)_{SM}$



small corrections to



and small corrections to



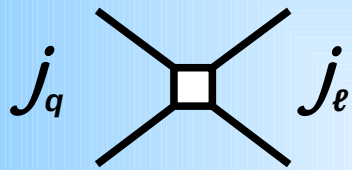
- $l \rightarrow l' + X$
< current limits



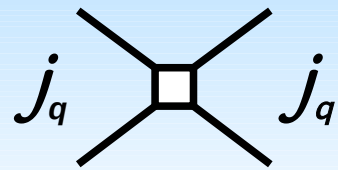
Z' - like NP

- *The challenge in short*

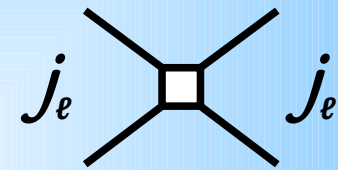
large enough



yet



&&

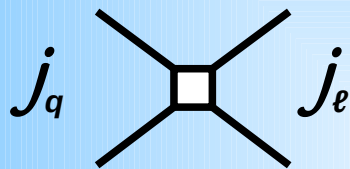


small enough

Z' - like NP

- *The challenge in short*

large enough

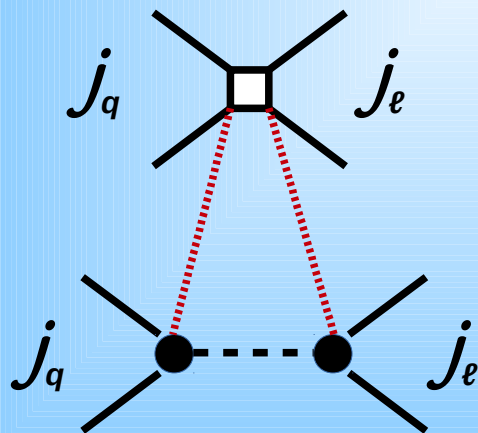


yet

small enough



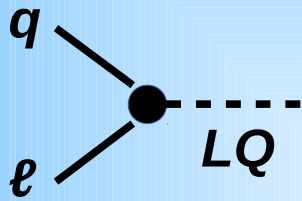
- *This is potentially a problem when*



*i.e. when the semi-lep. 4-f structure
arises from Z'-like NP*

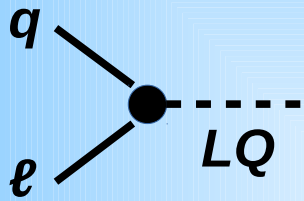
Leptoquark-like NP

Take



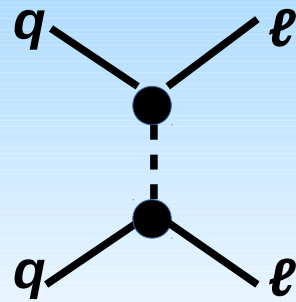
Leptoquark-like NP

Take



then

J_q

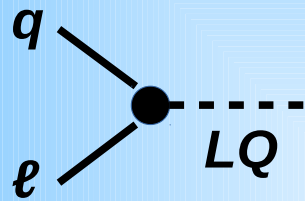


J_e

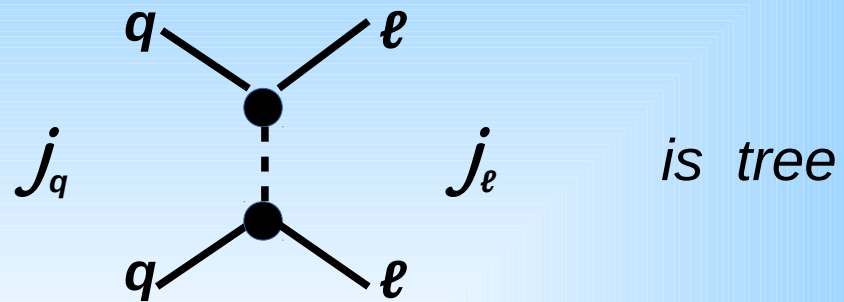
is tree

Leptoquark-like NP

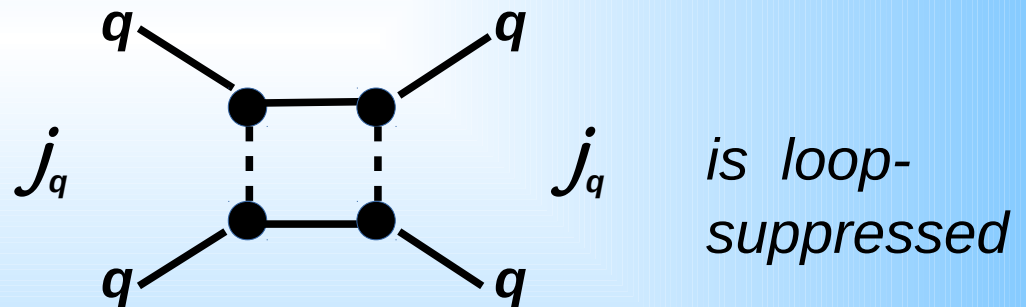
Take



then



but



(at least for “genuine” LQs [Dorsner et al., LQ review])

**Can one accomplish
a mechanism for suppressing
flavour-changing $j_q \otimes j_q$ && $j_e \otimes j_e$
within gauge extensions?**

A gauged horizontal $SU(2)$

- Place the two heavier generations of each fermion

in a doublet $\mathcal{F} \equiv \begin{pmatrix} f_2 \\ f_3 \end{pmatrix}$ w/ $f = u_L, d_L, \ell_L, \nu_L,$
or RH counterparts

A gauged horizontal $SU(2)$

- Place the two heavier generations of each fermion

in a doublet $\mathcal{F} \equiv \begin{pmatrix} f_2 \\ f_3 \end{pmatrix}$ w/ $f = u_L, d_L, \ell_L, \nu_L$,
or RH counterparts

- Consider a new $SU(2)$ interaction for each such doublet

$$\delta \mathcal{L} = g \sum_{\mathcal{F}} \bar{\mathcal{F}}_L \vec{\tau} \cdot \vec{G} \mathcal{F}_L + \text{RH counterpart}$$

A gauged horizontal $SU(2)$

- Place the two heavier generations of each fermion

in a doublet $\mathcal{F} \equiv \begin{pmatrix} f_2 \\ f_3 \end{pmatrix}$ w/ $f = u_L, d_L, \ell_L, \nu_L,$
or RH counterparts

- Consider a new $SU(2)$ interaction for each such doublet


$$\delta \mathcal{L} = g \sum_{\mathcal{F}} \bar{\mathcal{F}}_L \vec{\tau} \cdot \vec{\mathcal{G}} \mathcal{F}_L + \text{RH counterpart}$$

- Integrate out horizontal bosons

$$\delta \mathcal{L}_{\text{eff}} = - \sum_{\mathcal{F}, \mathcal{F}', a} \frac{g_L^2}{2M_{G_a}^2} \left(\bar{\mathcal{F}}_L \gamma^\mu \tau^a \mathcal{F}_L \right) \left(\bar{\mathcal{F}}'_L \gamma^\mu \tau^a \mathcal{F}'_L \right)$$

Basic argument

- Doublets $\mathcal{F} \equiv \begin{pmatrix} f_2 \\ f_3 \end{pmatrix}$ aren't yet in the mass basis.

Rotate as: $\mathcal{F} = U_{\mathcal{F}} \hat{\mathcal{F}}$  mass eigenbasis

Basic argument

- Doublets $\mathcal{F} \equiv \begin{pmatrix} f_2 \\ f_3 \end{pmatrix}$ aren't yet in the mass basis.

Rotate as: $\mathcal{F} = U_{\mathcal{F}} \hat{\mathcal{F}}$  mass eigenbasis

- How does \mathcal{L}_{eff} change?

$$\delta \mathcal{L}_{\text{eff}} \propto \frac{1}{2M_{G_a}^2} \left(\bar{\hat{\mathcal{F}}}_L U_{\mathcal{F}}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}} \hat{\mathcal{F}}_L \right) \left(\bar{\hat{\mathcal{F}}}'_L U_{\mathcal{F}'}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}'} \hat{\mathcal{F}}'_L \right)$$

Basic argument

$$\delta \mathcal{L}_{\text{eff}} \propto \frac{1}{2 M_{G_a}^2} \left(\bar{\hat{\mathcal{F}}}_L U_{\mathcal{F}}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}} \hat{\mathcal{F}}_L \right) \left(\bar{\hat{\mathcal{F}}}'_L U_{\mathcal{F}'}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}'} \hat{\mathcal{F}}'_L \right)$$

- Consider terms with $\mathcal{F} = \mathcal{F}'$

Basic argument

$$\delta \mathcal{L}_{\text{eff}} \propto \frac{1}{2 M_{G_a}^2} \left(\bar{\hat{\mathcal{F}}}_L U_{\mathcal{F}}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}} \hat{\mathcal{F}}_L \right) \left(\bar{\hat{\mathcal{F}}}'_L U_{\mathcal{F}'}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}'} \hat{\mathcal{F}}'_L \right)$$

- Consider terms with $\mathcal{F} = \mathcal{F}'$

If G_a
degenerate



Rotations can be traded
for G_a basis redefinition



$\mathcal{F} = \mathcal{F}'$ terms
flavour-diag. in all generality

Basic argument

$$\delta \mathcal{L}_{\text{eff}} \propto \frac{1}{2 M_{G_a}^2} \left(\bar{\hat{\mathcal{F}}}_L U_{\mathcal{F}}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}} \hat{\mathcal{F}}_L \right) \left(\bar{\hat{\mathcal{F}}}'_L U_{\mathcal{F}'}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}'} \hat{\mathcal{F}}'_L \right)$$

- Consider terms with $\mathcal{F} = \mathcal{F}'$

If G_a
degenerate



Rotations can be traded
for G_a basis redefinition



$\mathcal{F} = \mathcal{F}'$ terms
flavour-diag. in all generality

- Our currents of interest: $j_q \otimes j_q$, $j_\ell \otimes j_\ell$
would be flavour-diagonal in all generality

Basic argument

$$\delta \mathcal{L}_{\text{eff}} \propto \frac{1}{2 M_{G_a}^2} \left(\bar{\hat{\mathcal{F}}}_L U_{\mathcal{F}}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}} \hat{\mathcal{F}}_L \right) \left(\bar{\hat{\mathcal{F}}}'_L U_{\mathcal{F}'}^\dagger \gamma^\mu \tau^a U_{\mathcal{F}'} \hat{\mathcal{F}}'_L \right)$$

- Consider terms with $\mathcal{F} = \mathcal{F}'$

If G_a
degenerate



Rotations can be traded
for G_a basis redefinition



$\mathcal{F} = \mathcal{F}'$ terms
flavour-diag. in all generality

- Our currents of interest: $j_q \otimes j_q$, $j_\ell \otimes j_\ell$
would be flavour-diagonal in all generality

For the original argument
(in unrelated context) see:
Cahn, Harari, NPB1980

But

- *Mixing beneath the EWSB scale has to involve all generations*



Contributions to meson mixings & leptonic decays not exactly zero

But

- *Mixing beneath the EWSB scale has to involve all generations*



Contributions to meson mixings & leptonic decays not exactly zero

- *Contributions suppressed by powers of 1st – (2nd or 3rd) mixing*
So they are “small”

But

- *Mixing beneath the EWSB scale has to involve all generations*



Contributions to meson mixings & leptonic decays not exactly zero

- *Contributions suppressed by powers of 1st – (2nd or 3rd) mixing*

So they are “small”

- *But processes like $K^0 - \bar{K}^0$ mixing and $\mu \rightarrow 3e$ very constraining*

Is “small” small enough?

Scenario 0: degenerate G_a masses


Scenario 0: degenerate G_a masses

- Need to generalize our 2-generation relation $\mathcal{F} = U_{\mathcal{F}} \hat{\mathcal{F}}$ to 3 generations

- It is these $U_{3 \times 3}$ that are unitary

Then $CKM = (U_{UL})^\dagger U_{DL}$

Scenario 0: degenerate G_a masses

- Need to generalize our 2-generation relation $\mathcal{F} = U_{\mathcal{F}} \hat{\mathcal{F}}$ to 3 generations
- It is these $U_{3 \times 3}$ that are unitary
Then $CKM = (U_{UL})^\dagger U_{DL}$
- Taking $U_{DL} = 1$  $K^0 - \bar{K}^0$ constraint disappears


Scenario 0: degenerate G_a masses

- Need to generalize our 2-generation relation $\mathcal{F} = U_{\mathcal{F}} \hat{\mathcal{F}}$ to 3 generations

- It is these $U_{3 \times 3}$ that are unitary

Then $CKM = (U_{UL})^\dagger U_{DL}$

- Taking $U_{DL} = 1$  $K^0 - \bar{K}^0$ constraint disappears

but $U_{UL} = CKM^\dagger$  $D^0 - \bar{D}^0$ mixing $100\times$ exp limit

Scenario 0: degenerate G_a masses

- Need to generalize our 2-generation relation $\mathcal{F} = U_{\mathcal{F}} \hat{\mathcal{F}}$ to 3 generations

- It is these $U_{3 \times 3}$ that are unitary

Then $CKM = (U_{UL})^\dagger U_{DL}$

- Taking $U_{DL} = 1$ \Rightarrow $K^0 - \bar{K}^0$ constraint disappears

but $U_{UL} = CKM^\dagger$ \Rightarrow $D^0 - \bar{D}^0$ mixing $100\times$ exp limit

- Still exploring whether, with different $U_{UL,DL}$ assumptions, scenario 0 fulfils all main constraints:

CKM

small
 $K^0 - \bar{K}^0$

small
 $D^0 - \bar{D}^0$

R_K
as measured

Besides, several ways to generalize the idea that will fulfil all constraints.

Examples:

(i) non-degenerate G_a masses

(ii) non-zero (but small) 1^{st} – (2^{nd} & 3^{rd}) gen. mixing terms

Scenario 1: split G_a masses

- *Take one mass split from the other two, e.g.:*

$$m_{G_1} = m_{G_2} \ll m_{G_3}$$

Scenario 1: split G_a masses

- Take one mass split from the other two, e.g.:

$$m_{G_1} = m_{G_2} \ll m_{G_3}$$

- **Minus:**
With such pattern, one has to forsake the initial symmetry

Scenario 1: split G_a masses

- Take one mass split from the other two, e.g.:

$$m_{G_1} = m_{G_2} \ll m_{G_3}$$

- **Minus:**
With such pattern, one has to forsake the initial symmetry
- **Plus:** All data explained at one stroke
 - R_K & Co.

Scenario 1: split G_a masses

- Take one mass split from the other two, e.g.:

$$m_{G_1} = m_{G_2} \ll m_{G_3}$$

- **Minus:**

With such pattern, one has to forsake the initial symmetry


- **Plus:** *All data explained at one stroke*


- R_K & Co.
- ΔM_s ok, if somewhat $< SM$
- $B \rightarrow K \nu \bar{\nu}$ shift small, due to underlying $SU(2)$ sym.
- Small shifts to $\tau \rightarrow \ell \nu \nu$ & $D^0 \rightarrow \mu \mu$
- Small effects in di-muon tails [Greljo, Marzocca, EPJC2017]

Scenario 1: predictions

1 $\delta C_{9,10}^{\tau\tau} = -\delta C_{9,10}^{\mu\mu}$

Scenario 1: predictions

① $\delta C_{9,10}^{\tau\tau} = -\delta C_{9,10}^{\mu\mu}$ 

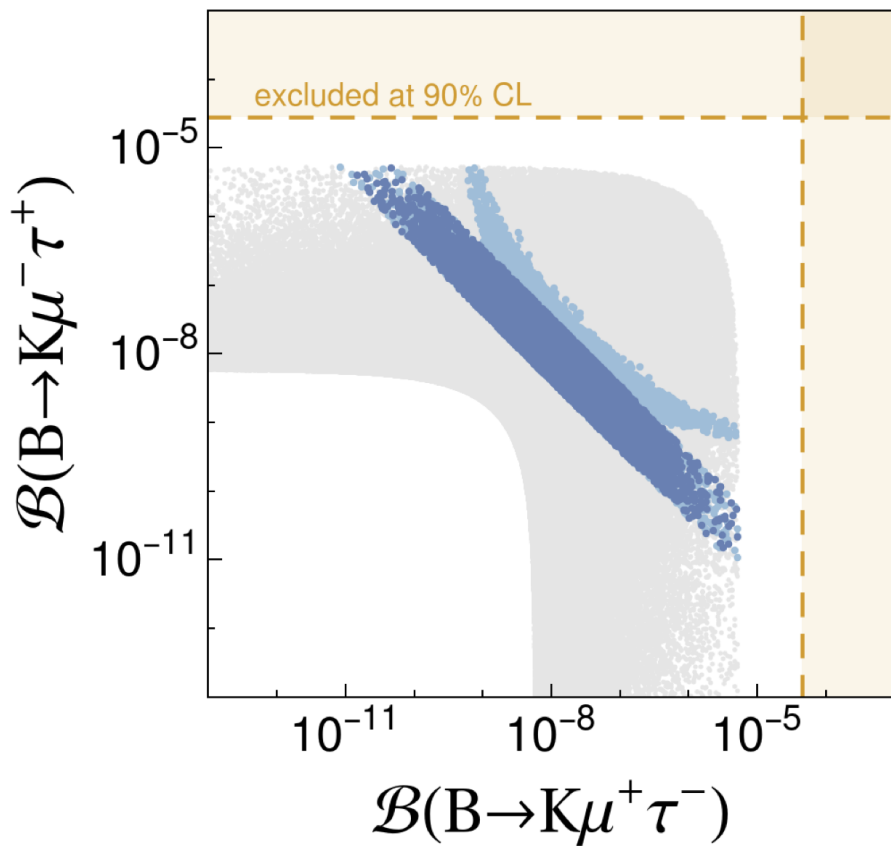
Shifts to $BR(B_{(s)} \rightarrow (K) \tau\tau)$  R_K shift

Scenario 1: predictions

1 $\delta C_{9,10}^{\tau\tau} = -\delta C_{9,10}^{\mu\mu}$ \Rightarrow

Shifts to $BR(B_{(s)} \rightarrow (K) \tau\tau)$ $\Leftrightarrow R_K$ shift

2 LFV-mode correlations

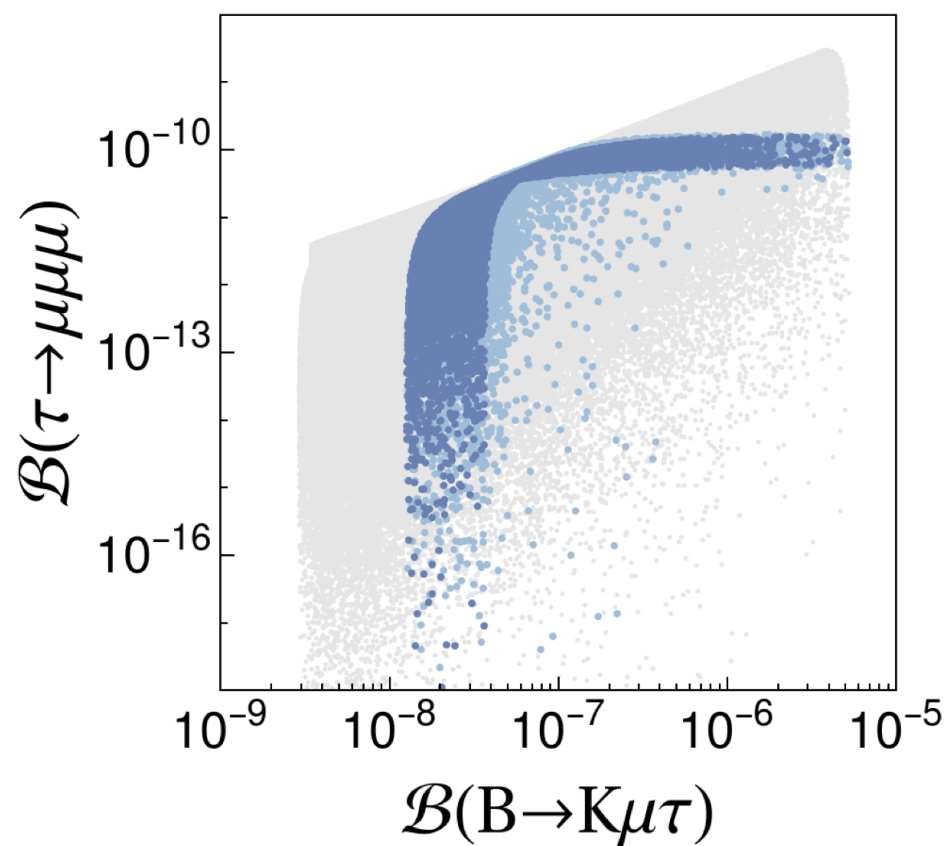
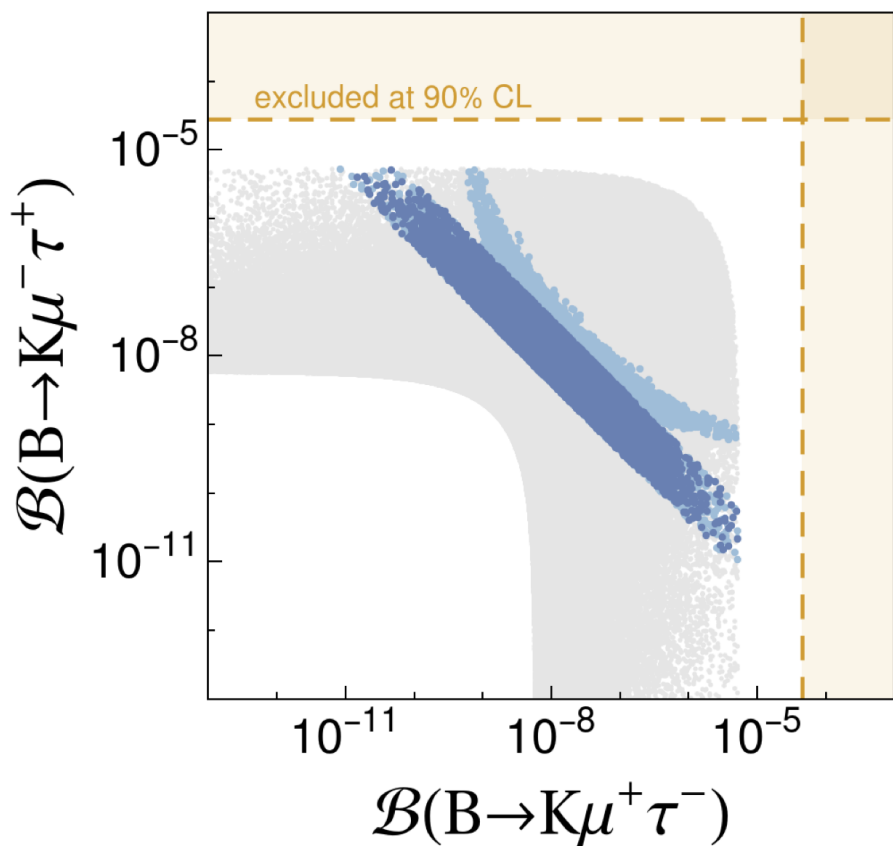


Scenario 1: predictions

1 $\delta C_{9,10}^{\tau\tau} = -\delta C_{9,10}^{\mu\mu}$ \Rightarrow

Shifts to $BR(B_{(s)} \rightarrow (K) \tau\tau)$ $\Leftrightarrow R_K$ shift

2 LFV-mode correlations

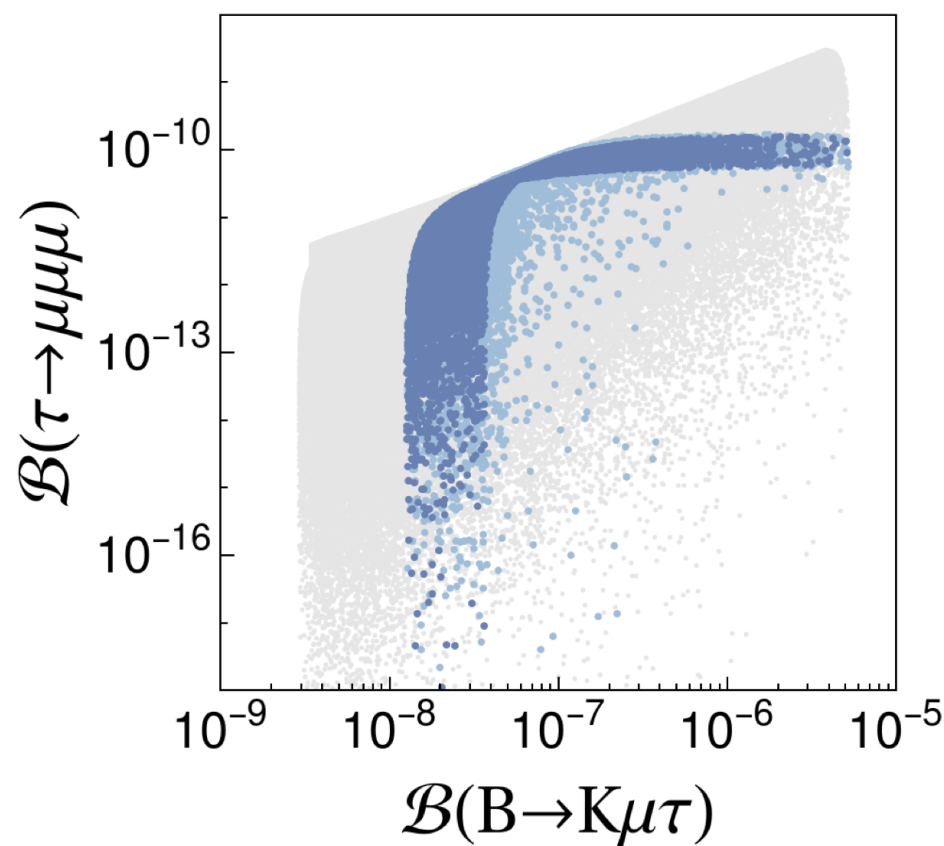
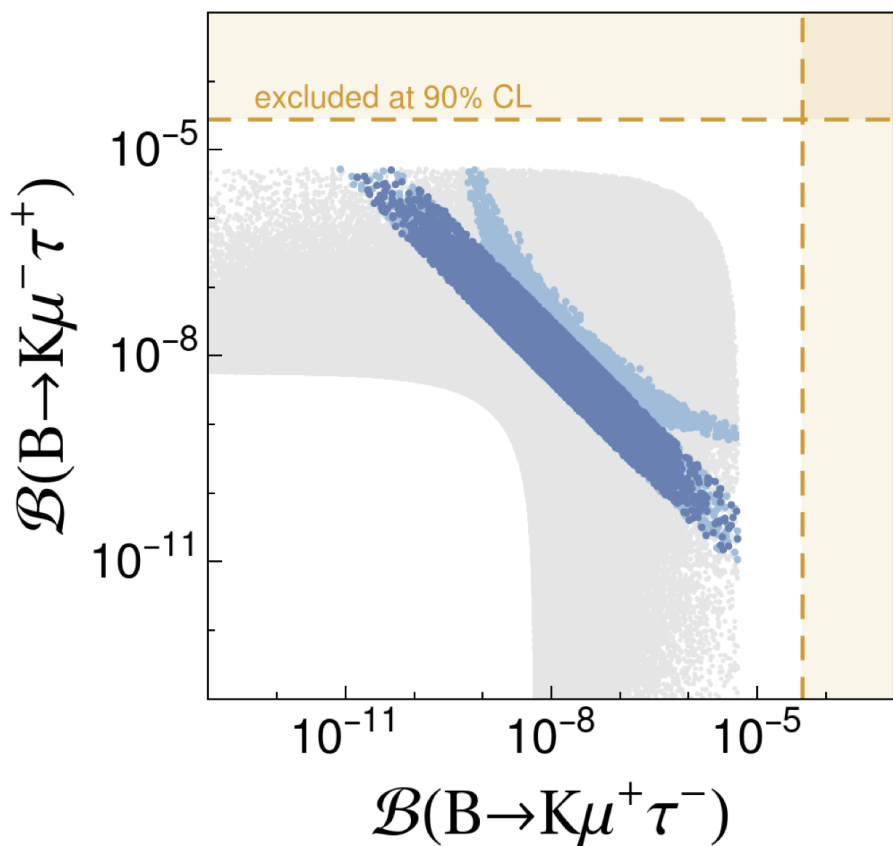


Scenario 1: predictions

① $\delta C_{9,10}^{\tau\tau} = -\delta C_{9,10}^{\mu\mu}$ \Rightarrow

Shifts to $BR(B_{(s)} \rightarrow (K) \tau\tau)$ $\Leftrightarrow R_K$ shift

② LFV-mode correlations



Note: $1.3 \times 10^{-8} \lesssim \mathcal{B}(B \rightarrow K\mu^+\tau^-) + \mathcal{B}(B \rightarrow K\mu^-\tau^+) \lesssim 5.2 \times 10^{-6}$

Conclusions / Outlook

- I discussed a Z' -like setup that accomplishes:

large

$$j_q \otimes j_e$$

(for R_K)

yet

small

$$j_q \otimes j_q \quad \&\& \quad j_e \otimes j_e$$

(ΔM_s && leptonic LFV constraints)

Conclusions / Outlook

- I discussed a Z' -like setup that accomplishes:

large

$$j_q \otimes j_e$$

(for R_K)

yet

small

$$j_q \otimes j_q \quad \&\& \quad j_e \otimes j_e$$

(ΔM_s && leptonic LFV constraints)

through a symmetry

Conclusions / Outlook

- I discussed a Z' -like setup that accomplishes:

large

$$j_q \otimes j_e$$

(for R_K)

yet

small

$$j_q \otimes j_q \quad \&\& \quad j_e \otimes j_e$$

(ΔM_s && leptonic LFV constraints)

through a symmetry

- The $SU(2)$ case implies automatically no gauge anomalies

Conclusions / Outlook

- I discussed a Z' -like setup that accomplishes:

large

$$j_q \otimes j_e$$

(for R_κ)

yet

small

$$j_q \otimes j_q \quad \&\& \quad j_e \otimes j_e$$

(ΔM_s && leptonic LFV constraints)

through a symmetry

- The $SU(2)$ case implies automatically no gauge anomalies



A larger group would require extra matter

But maybe it would help solve extra problems?