

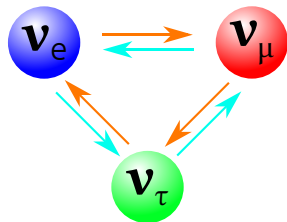
# Signatures from type-I,-II,-III seesaw at LHCb

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Implications of LHCb measurements and future prospects  
October 17, 2018, CERN

# The slide on motivation



Three Generations of Matter (Fermions) spin 1/2

	I	II	III	
mass	2.4 MeV	1.27 GeV	173.2 GeV	0
charge	2/3	2/3	2/3	0
name	u up	c charm	t top	g gluon
	Left Right	Left Right	Left Right	
Quarks				0
	d down	s strange	b bottom	γ photon
	Left Right	Left Right	Left Right	
	0	0	0	91.2 GeV
	ν <sub>e</sub> electron neutrino	ν <sub>μ</sub> muon neutrino	ν <sub>τ</sub> tau neutrino	Z weak force
	Left Right	Left Right	Left Right	
Leptons	0.511 MeV	105.7 MeV	1.777 GeV	126 GeV
	e electron	μ muon	τ tau	H Higgs boson
	Left Right	Left Right	Left Right	spin 0
	-1	-1	-1	0
	W <sup>-</sup> weak force			
	Left Right			

Bosons (Forces) spin 1

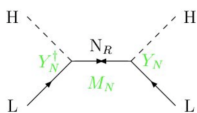
courtesy M. Shaposhnikov

- ▶ Observation of neutrino oscillations  $\Rightarrow$  light neutrinos massive.
  - ▶ No neutrino mass mechanism in the Standard Model (SM).
  - ▶ No mass matrix, no mixing of the neutrino flavour states.
- $\Rightarrow$  Evidence for physics beyond the SM.

# The 3 basic seesaw models

↪ i.e. tree level ways to generate the dim 5 operator

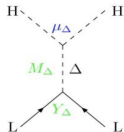
Right-handed singlet:  
(type-I seesaw)



$$m_\nu = Y_N^T \frac{1}{M_N} Y_N v^2$$

Minkowski; Gellman, Ramon, Slansky;  
Yanagida; Glashow; Mohapatra, Senjanovic

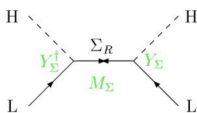
Scalar triplet:  
(type-II seesaw)



$$m_\nu = Y_\Delta \frac{\mu_\Delta}{M_\Delta^2} v^2$$

Magg, Wetterich; Lazarides, Shafi;  
Mohapatra, Senjanovic; Schechter, Valle

Fermion triplet:  
(type-III seesaw)



$$m_\nu = Y_\Sigma^T \frac{1}{M_\Sigma} Y_\Sigma v^2$$

Foot, Lew, He, Joshi; Ma; Ma, Roy; T.H., Lin,  
Notari, Papucci, Strumia; Bajc, Nemevsek,  
Senjanovic; Dorsner, Fileviez-Perez;.....

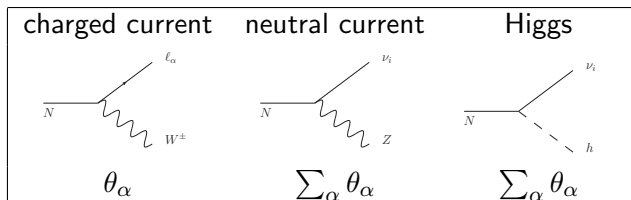
Slide from T. Hambye.

# Type-I seesaw for light neutrino masses

## Lagrangian:

$$\mathcal{L}_N = i \bar{N}_I \not{\partial} N_I - y_{\nu\alpha I} \bar{L}_\alpha \tilde{H} N_I - \frac{M_{N_{IJ}}}{2} \bar{N}_I^c N_J + \text{h.c.},$$

- ▶ Sterile neutrinos  $N_I$ , indices  $I, J = 1, 2, \dots, n$ .
- ▶ Parameters: Mass matrix  $M_N$ , Yukawa matrix  $y_\nu$ .
- ▶ Neutrino oscillations observe  $\Delta m_{\text{atm}}^2, \Delta m_{\text{sol}}^2 \Rightarrow n \geq 2$



# Light neutrino masses

- ▶ Neutrino mass matrix:

$$\begin{pmatrix} 0 & m_D \\ m_D^T & M \end{pmatrix} \begin{array}{l} \text{Majorana mass matrix } M. \\ \text{Dirac mass matrix } m_D = y_\nu v_{EW} \end{array}$$

- ▶ Seesaw  $\Leftrightarrow m_D \ll M$ .
- ▶ Naïve seesaw relation:

$$m_\nu = \frac{1}{2} \frac{m_D^2}{M_N} \Rightarrow |y_\nu|_{1,2} \simeq \sqrt{\frac{2\Delta m_{\text{sol,atm}} M_{N_{1,2}}}{v_{EW}^2}}$$

- ▶ In inverse seesaw  $m_\nu = 0 \Leftrightarrow$  associated symmetry exact.  
 $\Rightarrow$  Heavy neutrinos form degenerate “pseudo Dirac” pairs.  
 $\Rightarrow$  Neutrino Yukawa couplings are not constrained by  $m_\nu$ .

# The meaning of neutrino masses for collider pheno

- ▶ Production rates at LHCb for  $M_N < m_W \simeq 7 \text{ nb} \times |V_{\ell N}|^2$
- ▶ Naïve type I seesaw relation:  $|V_{\ell N}|^2 = \mathcal{O}(10^{-11})$ 
  - ⇒ **no events even in HL phase!**
  - ⇒ Possibly visible at the Z pole run of the FCC-ee

*A. Blondel et al. [FCC-ee study Team], Nucl. Part. Phys. Proc. 273-275 (2016) 1883*

- ▶ Symmetry protected type I seesaw with  $|V_{\ell N}|^2 < 10^{-5}$ .
  - ⇒ Sizable event numbers, **LNV suppressed** by  $m_\nu$ .

*Kersten, Smirnov; Phys. Rev. D 76 (2007) 073005*

- ▶ LNV reintroduced via propagation effects:
  - ⇔ mass splitting between HNL pairs  $>$  decay width.

*Anamiati, Hirsch, Nardi; JHEP 1610 (2016) 010*

- ▶ This effect can be experimentally resolvable.

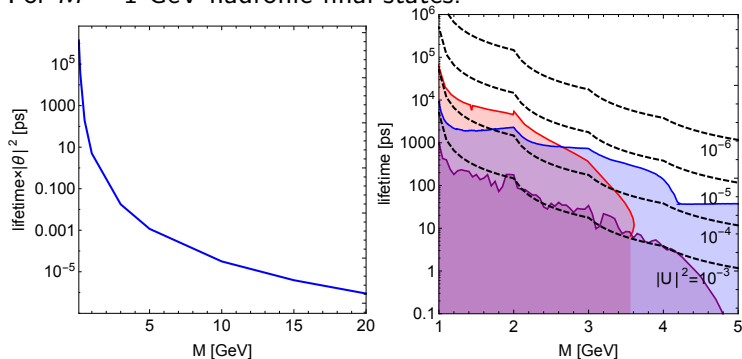
*Antusch, Cazzato, OF; arXiv:1709.03797 [hep-ph].*

Comment:

for  $M > m_W$  LHC cross sections small and backgrounds huge.

# Heavy neutrinos with symmetry protection and $M < m_W$

For  $M \sim 1$  GeV hadronic final states.

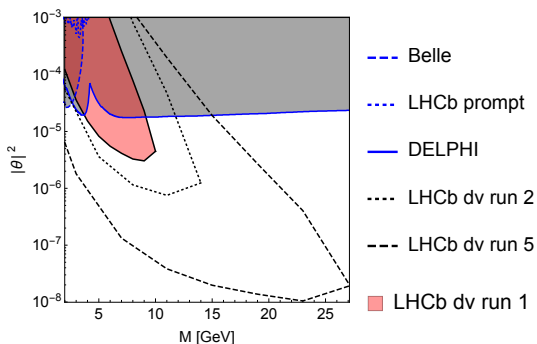


- ▶ Fixing mass and mixing angle also fixes lifetime.
- ▶ Displaced vertex signatures, even for masses  $M = \mathcal{O}(10)$  GeV.
- ▶ Blue: DELPHI, looking for *singly* produced HNL.

P. Abreu *et al.* [DELPHI Collaboration], *Z. Phys. C* **74** (1997) 57



# Forecasts, assuming dedicated analyses

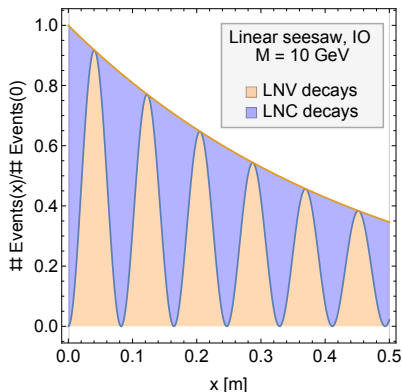
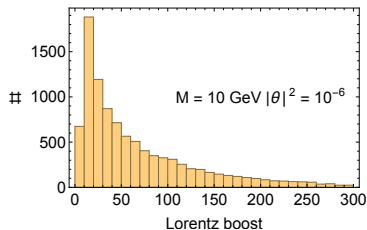
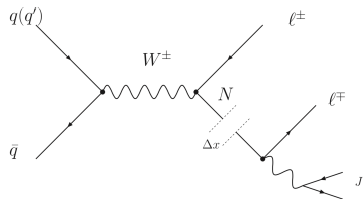


Antusch, Cazzato, OF; Phys. Lett. B 774 (2017) 114

- ▶ Displaced muons plus jets at 13 TeV
- ▶ black dotted: exclusion power for  $5 \text{ fb}^{-1}$  from run 2...
- ▶ black dashed:  $380 \text{ fb}^{-1}$  for the high-luminosity run.
- ▶ All limits for  $|\theta|^2 = |\theta_\mu|^2$  (i.e.  $|\theta_e| = |\theta_\tau| = 0$ ).

Assume we would discover displaced decays from HNL.

# Heavy neutrino oscillations in a nutshell



- ▶ Propagation effect: LNV as function of displacement  $\Delta x$ .
- ▶ Emergent pattern of SS vs OS dileptons in the proper frame.
- ▶ Proper oscillation length:  $\lambda_{\text{osc}}^{\text{lin,IO}} = 3.29 \cdot 10^{-3} \sqrt{\gamma^2 - 1} \text{ m}$
- ▶ In this example, proper lifetime:  $c\tau = 0.0094 \text{ m}$

Dropping some 'standard assumptions' for type-I HNL

Assumption 1:  $|\theta_\mu| \gg |\theta_e|, |\theta_\tau|$

**Instead:** theoretically motivated choice:  $|\theta_\tau| \gg |\theta_e|, |\theta_\mu|$

### Consequences:

- ▶ Precision constraints on  $|\theta_\tau|$  weaker compared to others.  
(They are  $|\theta_\tau|^2$  a few times  $10^{-2}$  for  $M \sim 1$  GeV.)
- ▶ Associated leptons (production) are tau leptons.
- ▶ Associated leptons (decay) are tau leptons if  $M > m_\tau$ .  
⇒ Signatures with two tau leptons important.  
⇒ Example:  $p p \rightarrow \tau^\pm N \rightarrow \tau^\pm (\tau^\mp X)_{\text{displaced}}$
- ▶ Semi leptonic decays suppressed if  $M < m_\tau$ !  
⇒ One prompt tau lepton and displaced vertex.  
⇒ Lepton non-universality from charged current decays.  
⇒ Example:  $B^\pm \rightarrow \tau^\pm N \rightarrow \tau^\pm (X)_{\text{displaced}}, X \supset e^+ e^-, \mu^+ \mu^-$

## Assumption 2: single production

**Notice:** many experimental constraints are based on this assumption, e.g. DELPHI search [P. Abreu et al. \[DELPHI Collaboration\], Z. Phys. C 74 \(1997\) 57.](#)

### HNL pair production via scalars:

- ▶ Scalar mixing between SM Higgs doublet and scalar singlet
- ▶ Production via  $h(125)$  or heavy scalar  $\propto \sin^2 \alpha$  (scalar mixing).  
 $\Rightarrow$  Naïve type-I seesaw relation for  $\theta$  predicts  $h \rightarrow 2 X_{\text{displaced}}$

[Deppisch, Liu, Mitra; JHEP 1808 \(2018\) 181](#)

### HNL pair production via $Z'$ bosons:

- ▶ Production via resonant  $Z'$  possible if  $\text{Br}(Z' \rightarrow NN) \sim 1$ .
- ▶ Production via  $\gamma, Z$  with gauge mixing angle  $< 10^{-2}$ .
- ▶ Production  $\propto g_{Z'}^4 / m_{Z'}^4$ , for  $m_{Z'} > \sqrt{s}$

[E. Accomando et al.; JHEP 1802 \(2018\) 109](#)

Type-III seesaw: neutrino masses from Fermion triplets

# Type-III seesaw for light neutrino masses

## Lagrangian:

- ▶  $\Sigma$ : fermion field, isospin triplet, hypercharge 0.
- ▶  $\mathcal{L}_{\text{kinetic}} = \text{Tr}[\bar{\Sigma} i \not{D} \Sigma]$  Gauge interactions
- ▶  $\mathcal{L}_{\text{Majorana}} = -\frac{1}{2} \text{Tr}[\bar{\Sigma} M_{\Sigma} \Sigma^c + \bar{\Sigma}^c M_{\Sigma}^* \Sigma]$  Majorana mass
- ▶  $\mathcal{L}_{\text{Yukawa}} = \sqrt{2} \tilde{\Phi}^{\dagger} \bar{\Sigma} L$  Yukawa interactions

## Neutrino masses:

$$m_{\nu} = Y_{\Sigma}^T \frac{1}{M_{\Sigma}} Y_{\Sigma} v_{\text{EW}}^2,$$

- ▶ Mass matrix identical to type-I seesaw.
- ▶ The same conclusions apply.



# Some type -III phenomenology

## Leptonic mixing:

- ▶ Neutral component  $\Sigma^0$  mixes with  $\nu_L$
- ▶ Charged components  $\Sigma^\pm$  mix with charged leptons
- ▶ Mass eigenstates:  $\nu_1, \dots, N^0$  and  $e, \mu, \tau, N^\pm \Rightarrow$  Precision measurements of leptons imply mixing  $< 10^{-1.5}$ .
- ▶ Usually naïve seesaw relation applied: mixing  $\sim 10^{-6}$

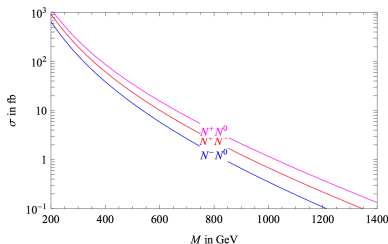
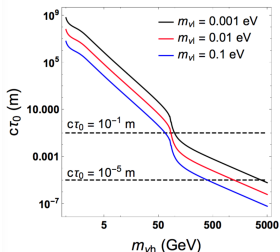
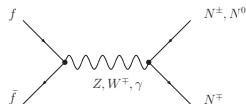
## Interactions:

- ▶ Isospin fixes coupling strength to the gauge bosons.
- ▶ Measurement of  $\Gamma_Z$  at LEP requires  $m_{N^\pm} > m_Z/2$ .

## Lifetime:

- ▶  $N^0$  decays via leptonic mixing  $\Rightarrow$  large  $c\tau$  possible.
- ▶ Radiative corrections:  $m_{N^\pm} = m_{N^0} + 166 \text{ MeV}$   
 $\Rightarrow N^\pm \rightarrow N^0(W^\pm)^*$ , very short lifetime

# Type-III signatures at the LHC



Franceschini, Hambye, Strumia; *Phys. Rev. D* **78** (2008) 033002

- ▶  $N^0$ ,  $N^\pm$  can be pair produced via Drell Yan.
- ▶  $N^\pm$  decay promptly into  $N^0$  plus soft pion or lepton
- ▶ ATLAS/CMS analyses include  $N^0$  decays

The ATLAS collaboration [ATLAS Collaboration], ATLAS-CONF-2018-020.

A. M. Sirunyan *et al.* [CMS Collaboration], *Phys. Rev. Lett.* **119** (2017) no.22, 221802

- ▶ Displacements due to long lifetimes not considered  
 $\Rightarrow$  Long lived  $N^0$  gives neither MET nor prompt tracks!

Type-II seesaw: neutrino masses from scalar triplets

# Neutrino masses via type-II seesaw

## Lagrangian:

- ▶  $\Delta$  scalar field, isospin triplet, hypercharge 2.
- ▶  $\text{Tr}[(D_\mu \Delta)^\dagger (D^\mu \Delta)]$  Gauge interactions
- ▶  $\lambda_1 \Phi^2 \Delta^2 - \lambda_2 \Delta^4 - \kappa \Delta^0 \Phi^2$  Scalar interactions
- ▶  $Y_\Delta \bar{\ell}^c \Delta \ell + \text{h.c}$  Yukawa interactions

## Neutrino masses:

$$m_\nu = Y_\Delta \frac{v_\Delta}{M_\Delta^2} v_{\text{EW}}^2$$

- ▶ Neutral component  $\Delta^0$  with non-zero VEV  $v_\Delta$
- ▶ Small  $m_\nu \Leftrightarrow$  small  $Y_\Delta, v_\Delta, M_\Delta^{-1}$

# Phenomenology

## “Standard” components of the triplet:

- ▶ Neutral mass eigenstate mixes with the SM Higgs doublet  
⇒ two neutral CP even Higgs-like bosons.
- ▶ Singly charged  $\Delta^\pm$ , difficult to search for at LHC.

## Focus on the doubly charged $\Delta^{\pm\pm}$ :

- ▶ Decay into SS dilepton  $\propto Y_\Delta$
- ▶ Decay into SS W boson  $\propto v_\Delta$
- ▶ Searches for prompt SS dileptons:  $m_{\Delta^{\pm\pm}} > 620$  GeV

M. Aaboud *et al.* [ATLAS Collaboration], *Eur. Phys. J. C* **78**, no. 3, 199 (2018)

- ▶ Can be long lived for  $m_\Delta < 2m_W$ , not stable!
- ▶ Searches exist for charged massive stable particles (CMSP)

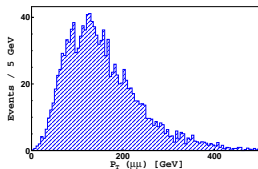
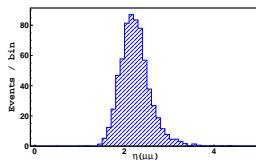
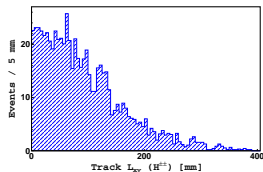
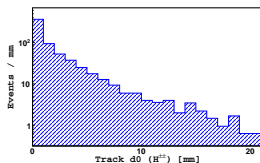
R. Aaij *et al.* [LHCb Collaboration], *Eur. Phys. J. C* **75** (2015) no.12, 595

- ▶ Intermediate lifetimes are subject of study.

Dev, Zhang; arXiv:1808.00943 [hep-ph].

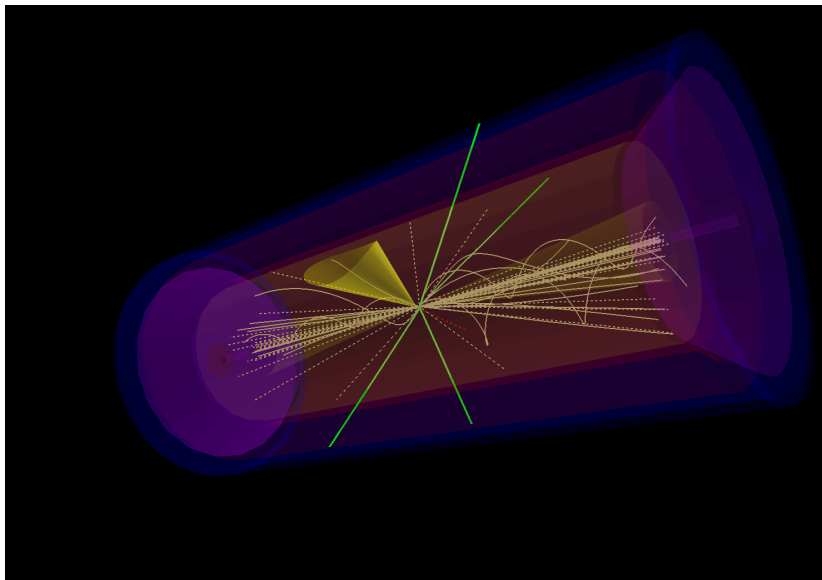
## One example benchmark point

- ▶  $\sqrt{s} = 14$  TeV, luminosity = 3/ab,  $\eta > 2$ ,
- ▶  $m_{\Delta} = 130$  GeV,  $v_{\Delta} = 5 \times 10^{-4}$  GeV
- ▶  $c\tau_{\text{proper}} > 1$  mm, invalidates ATLAS/CMS prompt searches!
- ▶ Compatible with all current constraints.



⇒ Signal: non-prompt dileptons; this is certainly visible?

# A Delphes event display



# Conclusions

## Neutrino masses via seesaw:

- ▶ Type-I: right-handed (or sterile) neutrinos  $\Rightarrow$  HNL.
- ▶ Type-II: scalar isotriplets with hypercharge  
 $\Rightarrow$  neutral scalar mixing, singly/doubly charged scalars, LNV.
- ▶ Type-III: fermion isotriplets, no hypercharge  
 $\Rightarrow$  Gauge interactions fixed, pair production, decay chains, neutral component is HNL

## Signatures related to seesaw:

- ▶ Displaced vertices everywhere in the detector!
- ▶ Drell-Yan produced final states with prompt and displaced  $\tau^\pm$ .
- ▶ Meson decays with prompt  $\tau^\pm$  and displaced dileptons.
- ▶ Pairs of displaced vertices w/ & w/o soft pions/leptons.
- ▶ Displaced same sign dimuons.



Can LHCb detect  $e^+e^-$  or a single photon with  $E < \text{GeV}$  that appear anywhere in the detector?