

CSR in Light Sources

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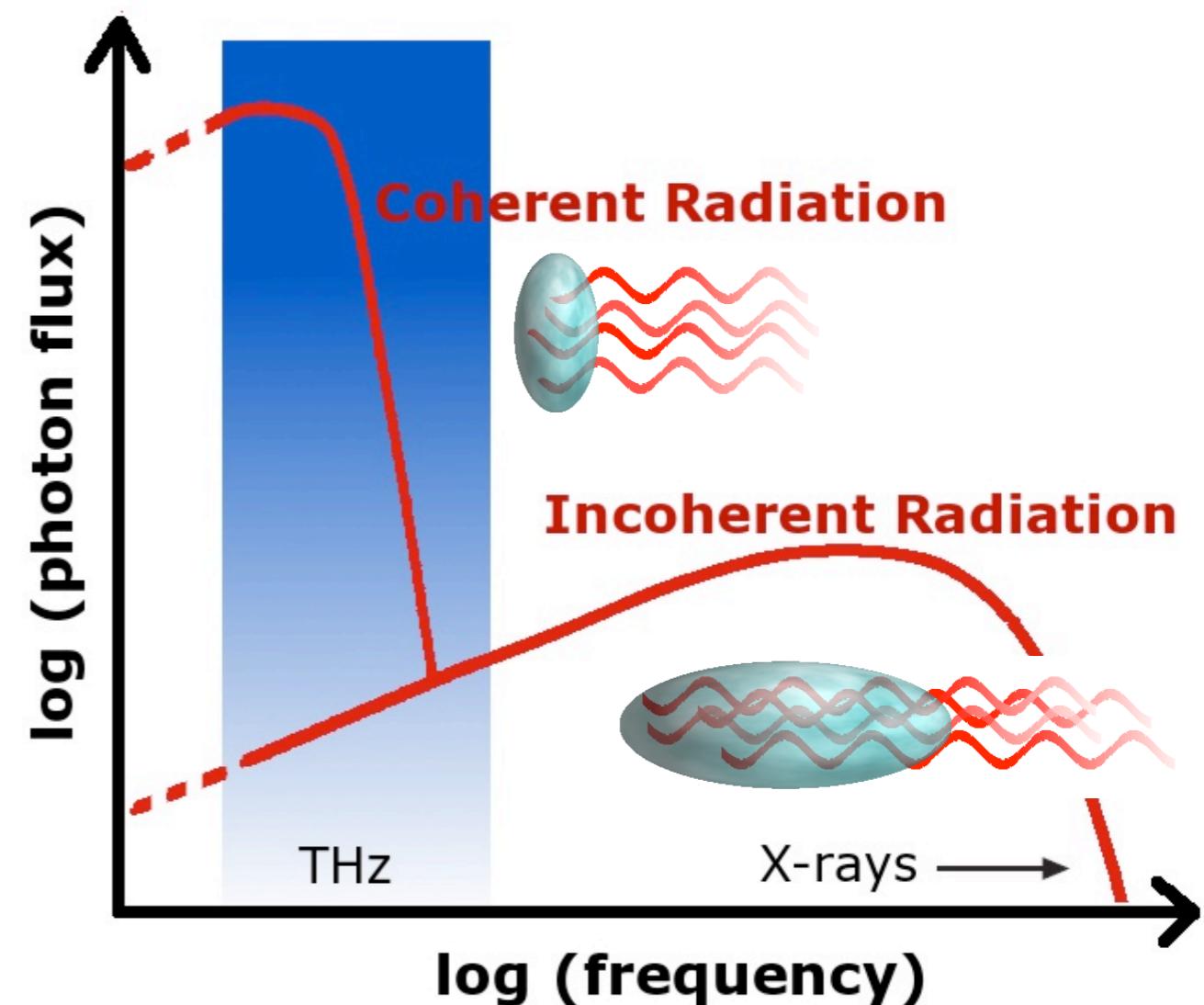
PTB Berlin

R. Müller

- Why CSR?
- Optics considerations for CSR generation in storage rings
- CSR observed in time & space
- CSR and bunch charge distribution
 - bursting/stable emission, bunch deformation
 - influence of other impedances
- Summary

Why Coherent Synchrotron Radiation?

- Enormous increase in power in comparison to incoherent emission
- Extension of successful experimental methods to the low frequency (Terahertz) range
- Coherent synchrotron radiation is emitted from electrons in a deflecting magnetic field for wavelengths equal to or longer than the bunch length
→ short bunches are needed





456. WILHELM UND ELSE HERAEUS SEMINAR

THz RADIATION: GENERATION, DETECTION AND APPLICATIONS

18. - 21. April 2010 Physikzentrum Bad Honnef

■ Lectures on all fields related to THz radiation

- accelerator & laboratory sources
- standard & advanced detection techniques
- applications in life sciences, non destructive testing, metrology, security, astronomy

■ Web page:

http://ankaweb.fzk.de/science_at_anka/ANKA_THz_Group/ANKA_THz_Group/WEH-Seminar/Home.html

■ Venue:

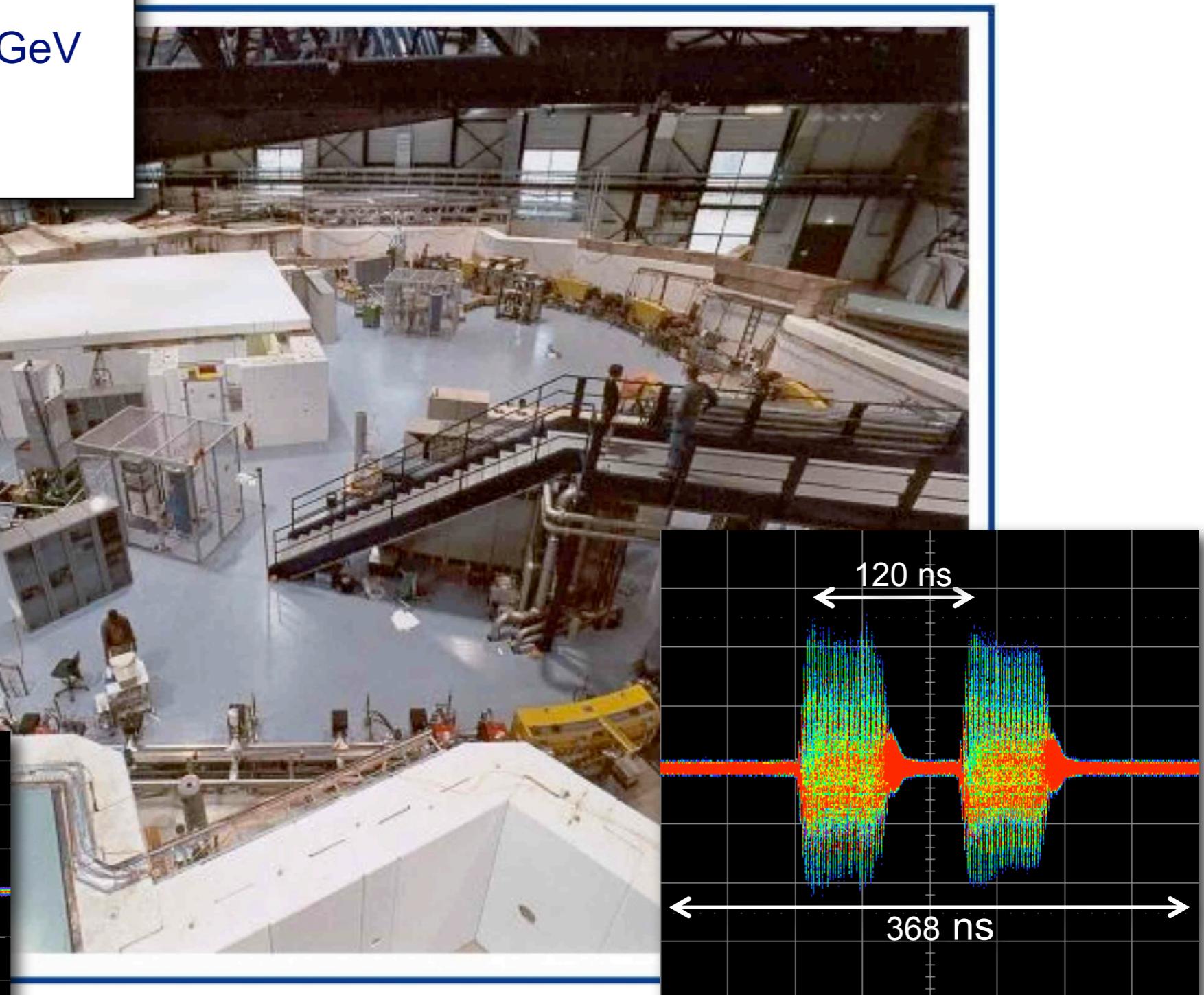
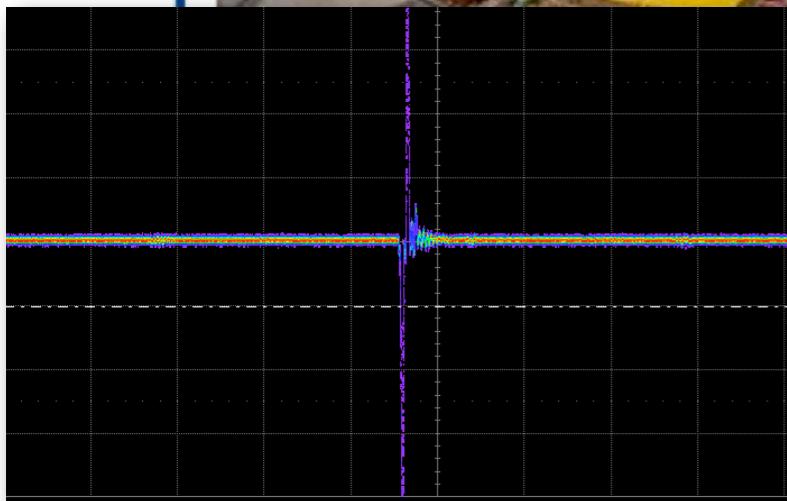


Photographs of the Physikzentrum
Bad Honnef (Courtesy PBH)



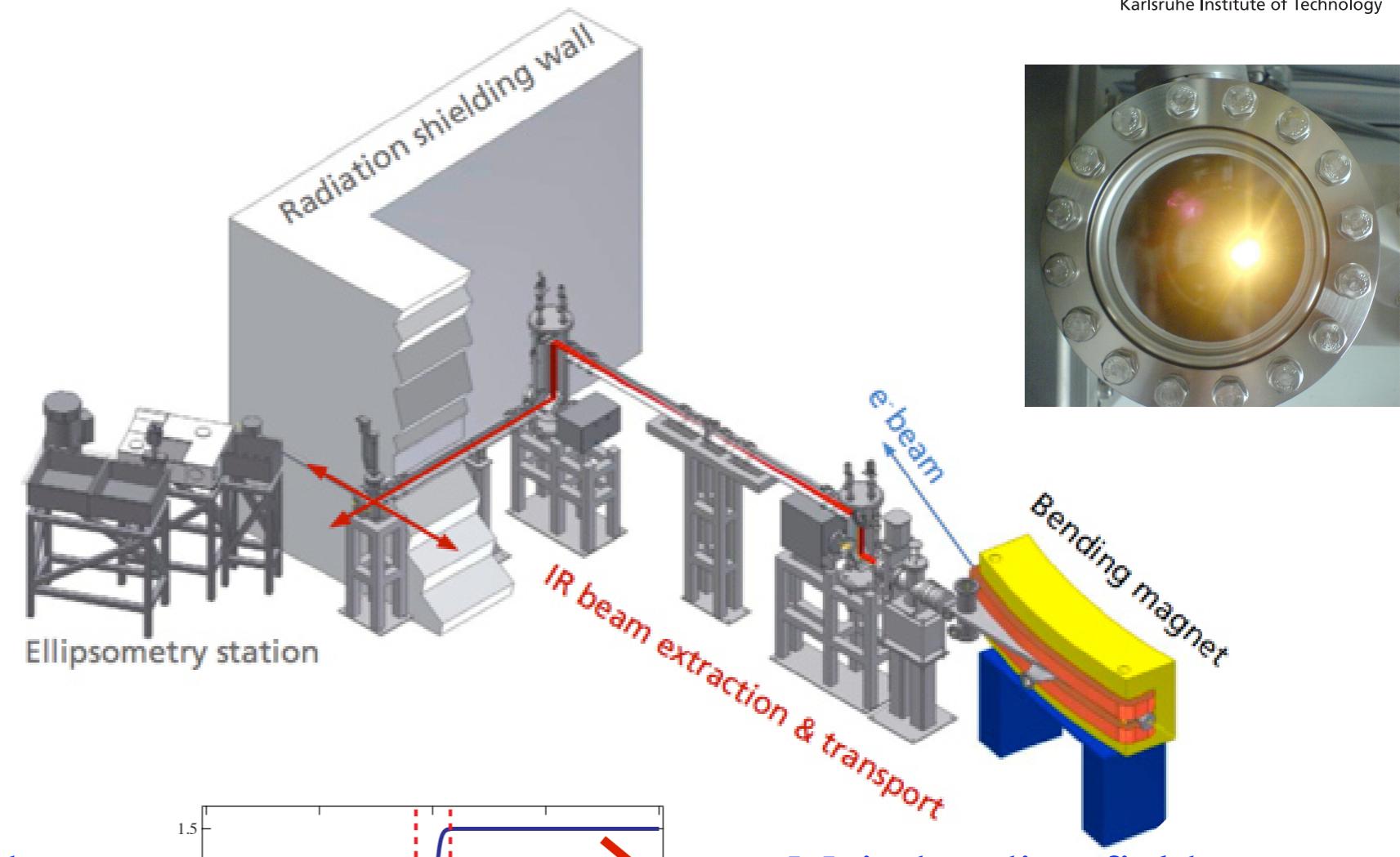
Example: The ANKA Storage Ring

- $C = 110.4 \text{ m}$
- Energy range: 0.5 - 2.5 GeV
- RF frequency 500 MHz
- DBA lattice

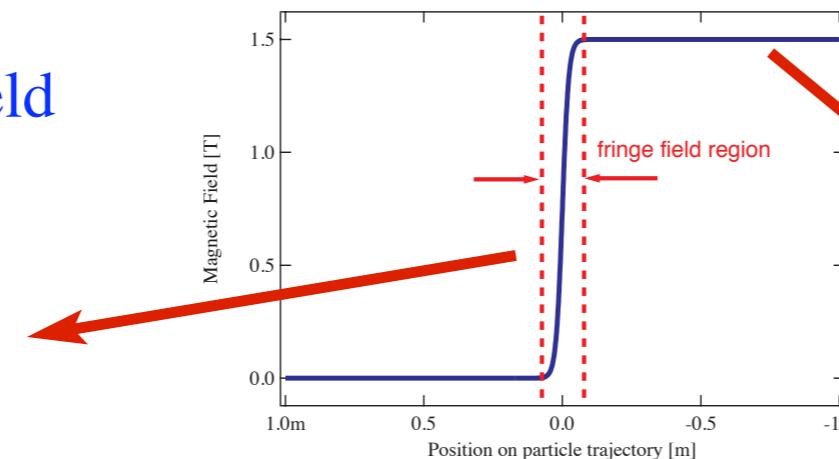
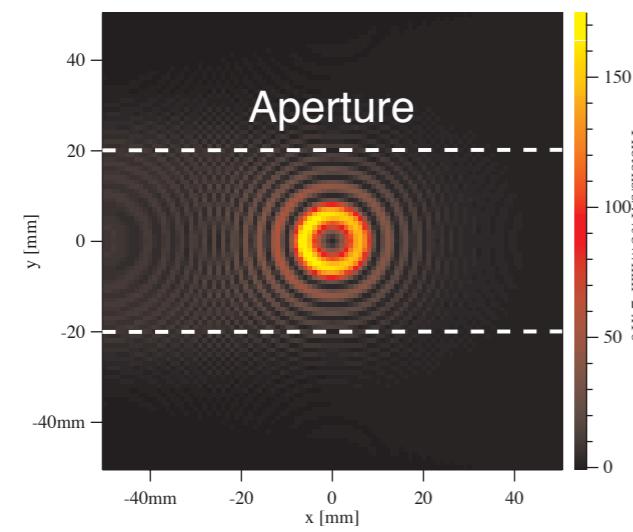


Synchrotron (Edge) Radiation

- CSR is observed as ‘regular’ synchrotron radiation but also as ‘edge’ radiation
- Can be an advantage for a beamline
 - lower frequencies observable for the same aperture

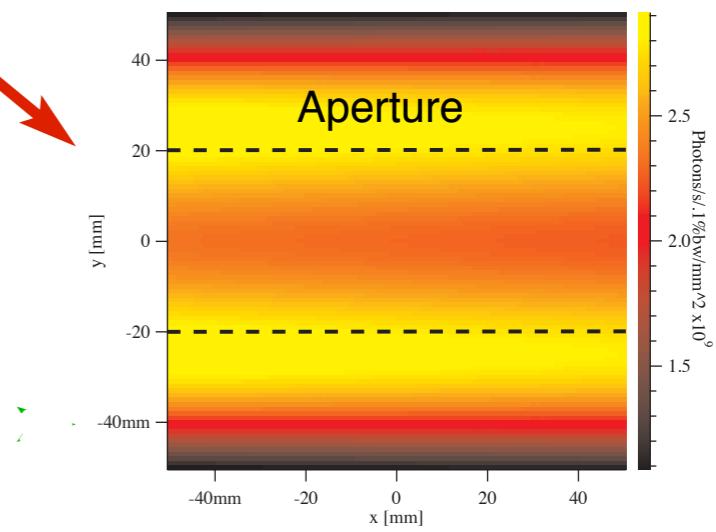


Source at ANKA-IR: fringe field



Courtesy Y.-L.Mathis

Main bending field as source



Low- α_c Optics at ANKA

■ Condition for CSR emission: $\frac{2\pi\sigma_s}{\sqrt{\ln N}} \lesssim \lambda \lesssim 2h\sqrt{\frac{h}{\rho}}$

→ for 100 $\mu\text{A}/\text{bunch}$: $1.4\sigma_s \lesssim \lambda \lesssim 4.9\text{ mm}$

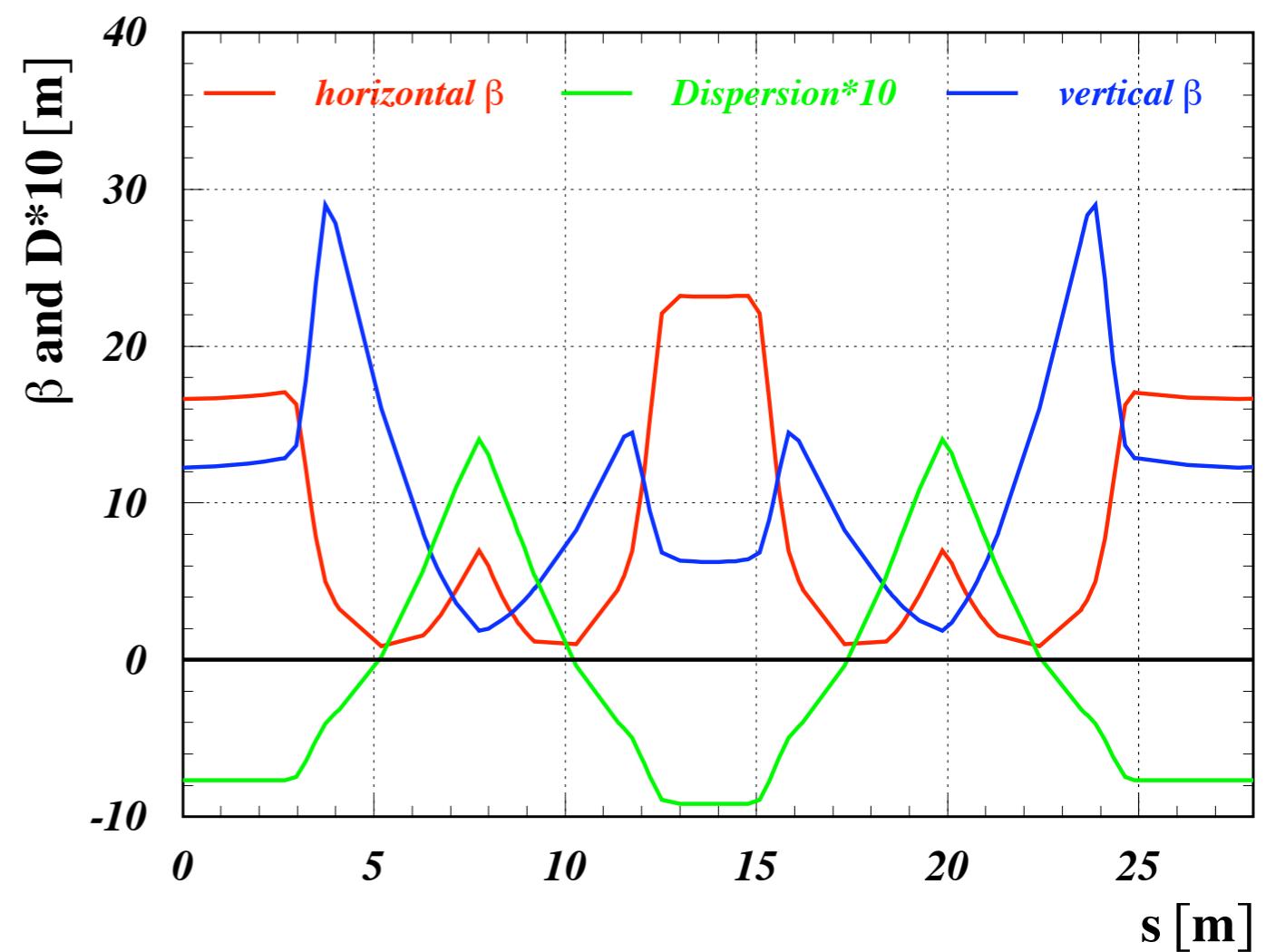
→ further bunch length reduction necessary

$$\alpha_c = \frac{1}{L} \oint ds \frac{D(s)}{\rho(s)}$$

■ Dedicated low- α_c optics with negative dispersion in the long and short straight sections for flexible bunch length tuning following the pioneering work of e.g. BESSY II

■ At ANKA: Observed momentum compaction factor range as extrapolated from Q_s measurements:

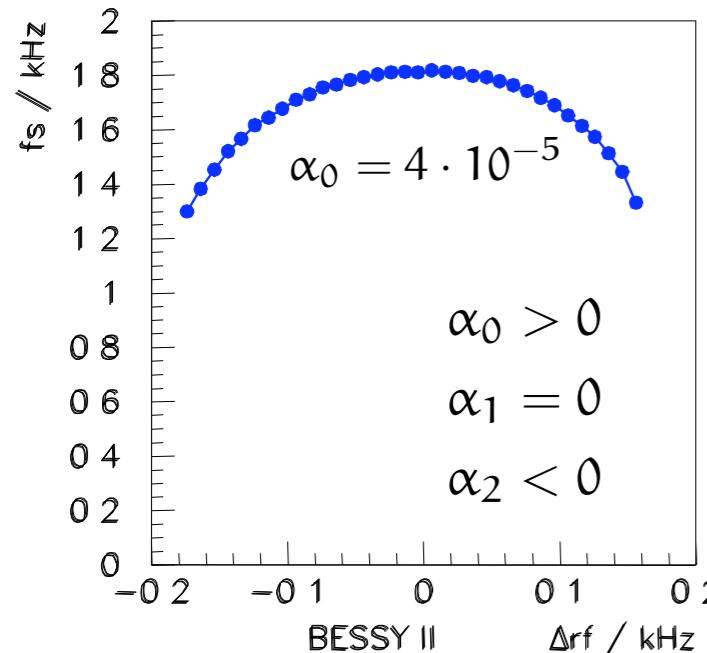
→ from $7.2 \cdot 10^{-3}$ to $1.4 \cdot 10^{-4}$



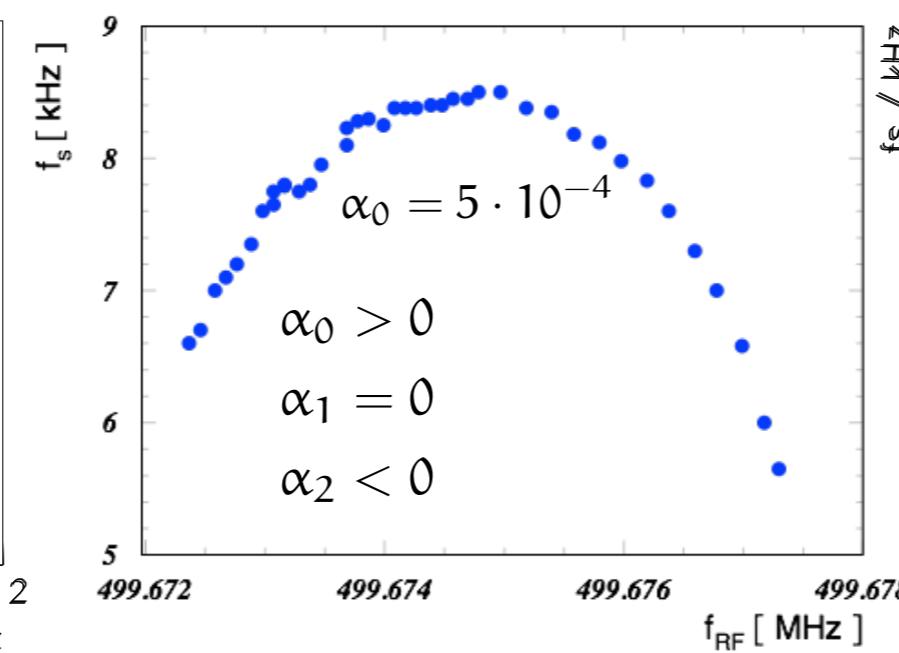
Momentum Compaction Factor

Synchrotron frequency f_s as a function of Δf_{rf} detuning

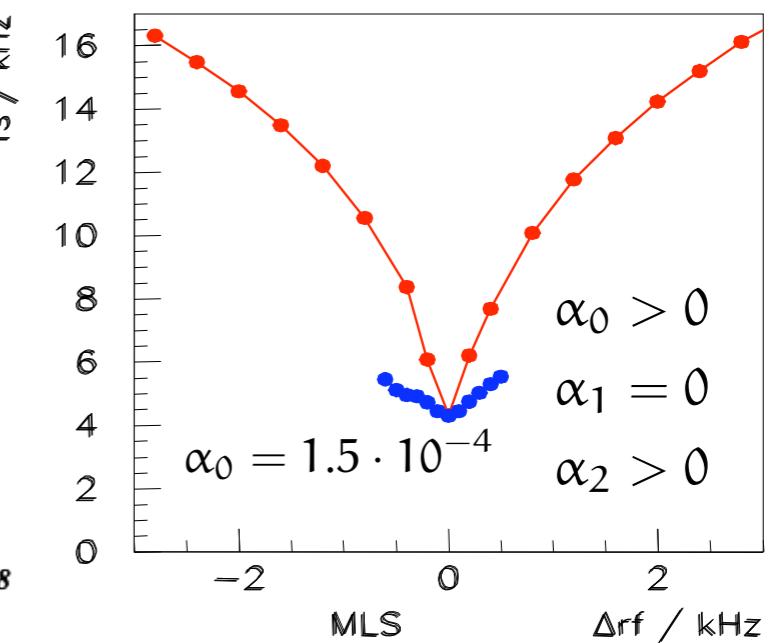
BESSY II, 1.7 GeV
4 chrom. sext. families,
limited flexibility



ANKA, 1.3 GeV
Low alpha optics,
2 chrom. sext. families



MLS, 630 MeV ★
3 chrom. sextupole families,
& octupole family



$\Delta f_{rf} \approx 4$ kHz



MLS: first ring with low alpha correction scheme

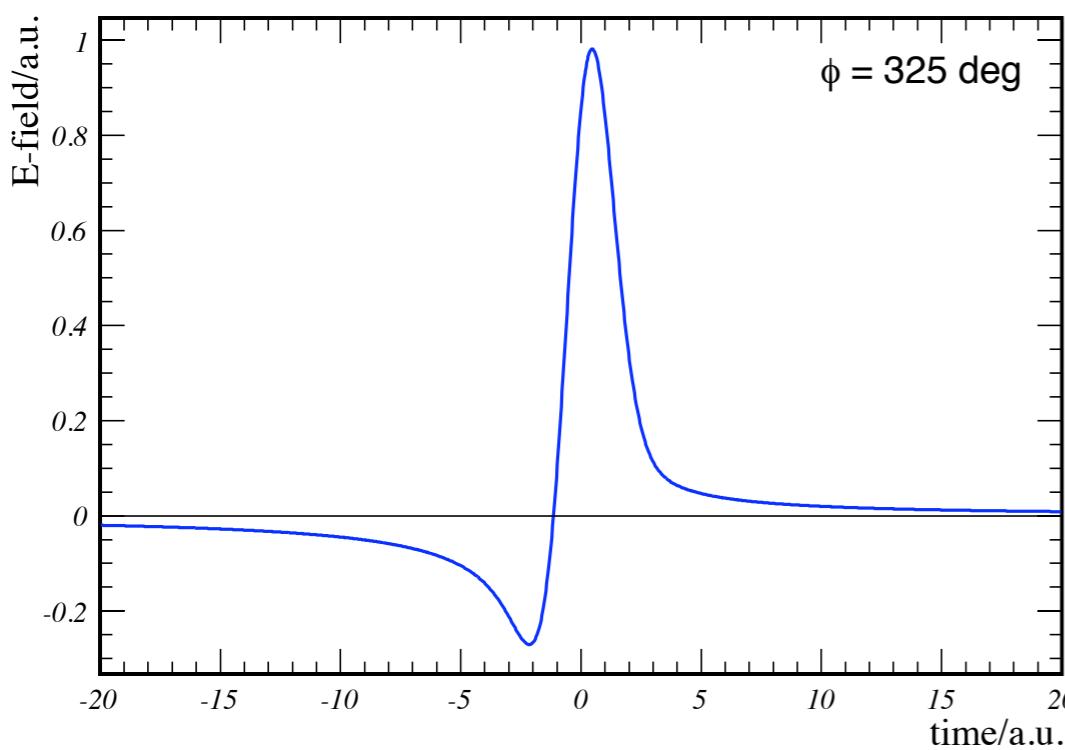
G. Wüstefeld, HZB

THz Pulse in the Time Domain

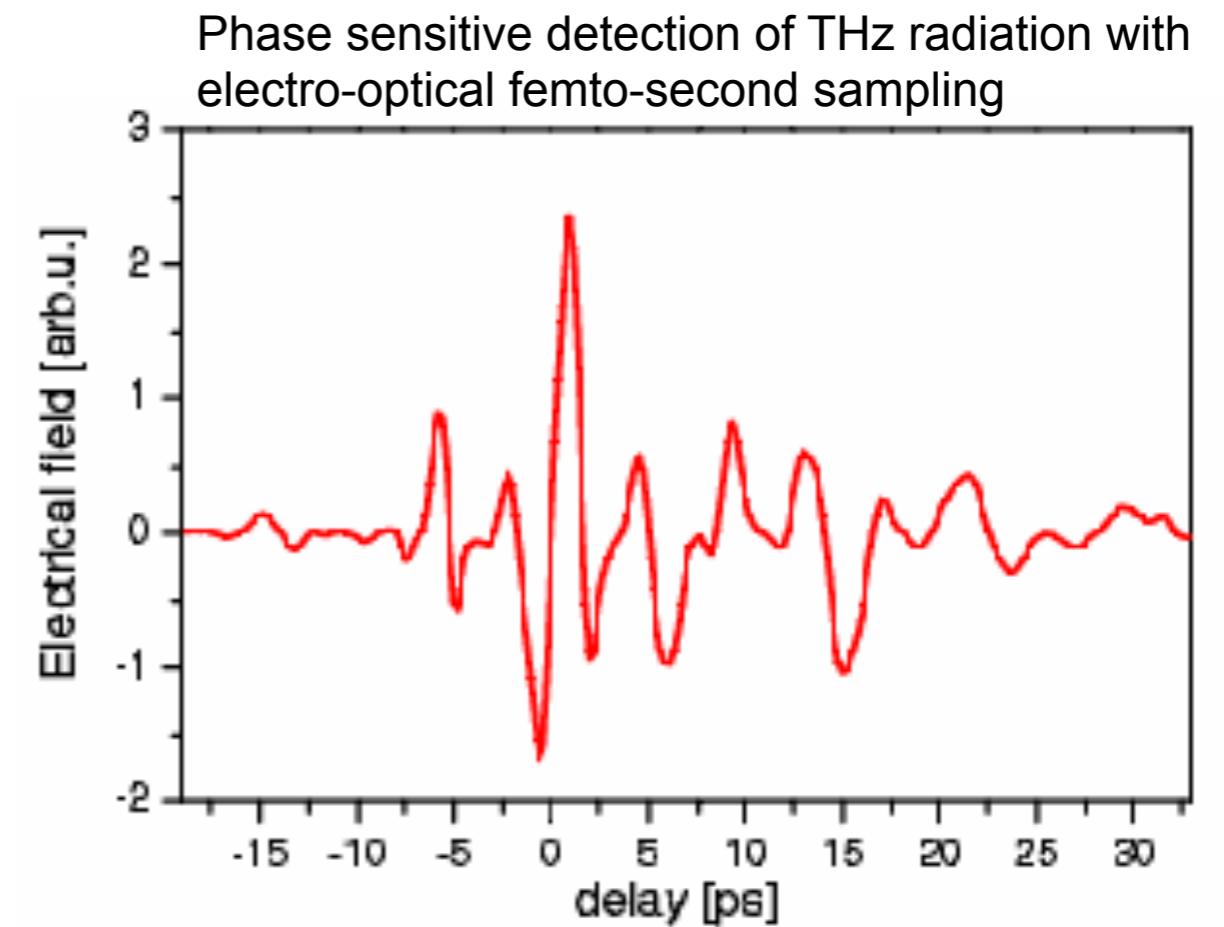
- Superposition of plane waves emitted over the bunch:

$$A(t) = \int_0^\infty d\omega \int_{-\infty}^\infty dx s(\omega) \rho(x) e^{-i\omega(t-x/c)}$$

- Resulting electrical field: $E(t) \sim \cos \phi \cdot \text{Re}A(t) + \sin \phi \cdot \text{Im}A(t)$
(The phase determines the relative weight of the two independent solutions. It is not fixed a priori and given by the ring structure.)



ASM et al.: Modeling the Shape of Coh. THz Pulses
Emitted by Short Bunches in an El. SR, EPAC 2008



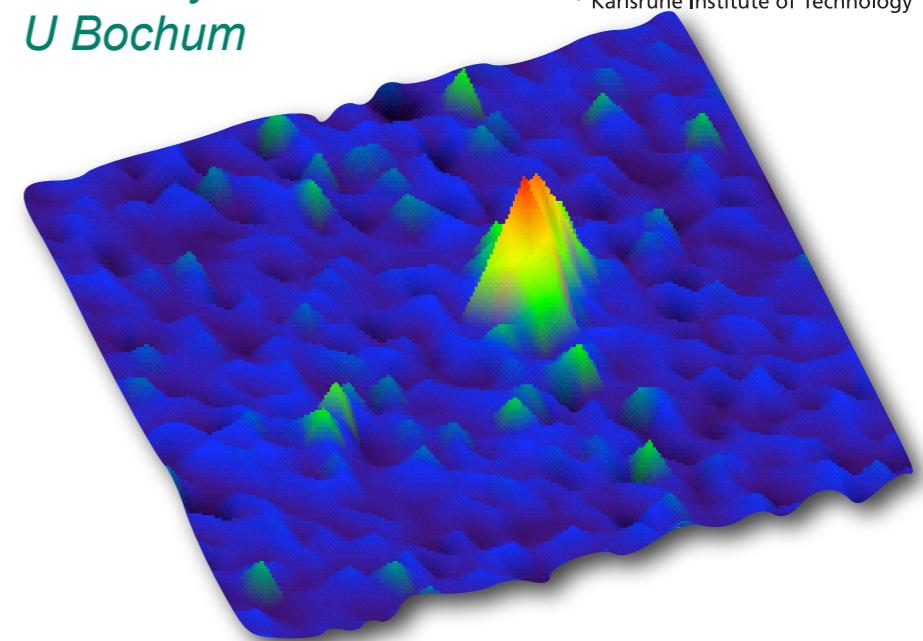
A. Plech et al.: Electro-Optical sampling of Terahertz radiation emitted by short bunches in the ANKA synchrotron, PAC 2009

The THz Beam Profile

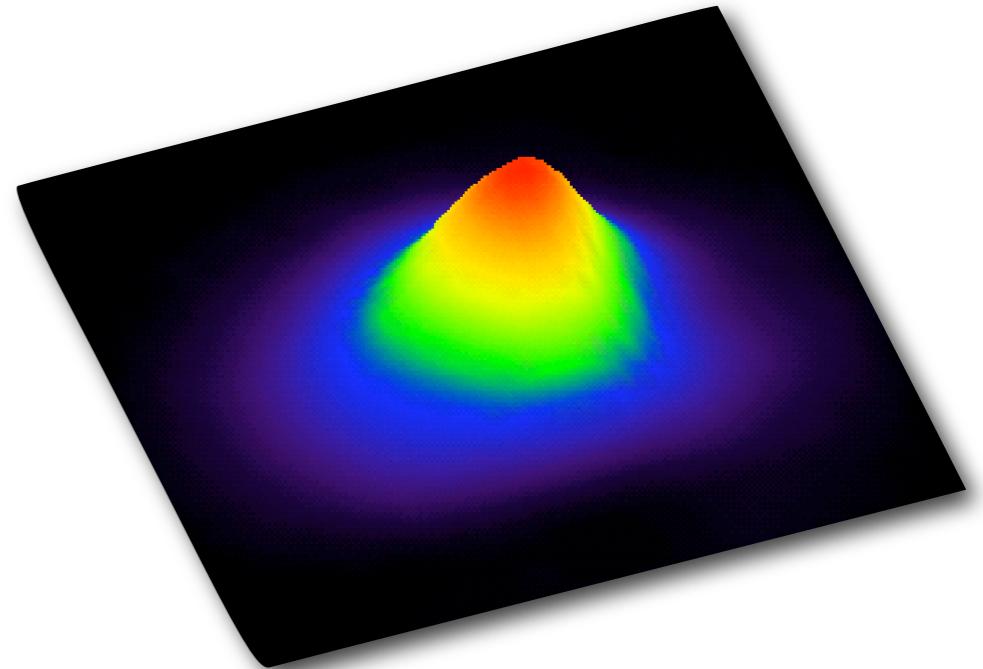
Setup of beam line and detector :

- Measurement behind a Si or CaF₂ vacuum window
- room temperature pneumatic (Golay) detector
- aperture of 6 mm diameter (defined by white high density polyethylene window) plus add. 1.9 mm diameter aperture in front of the detector for better spatial resolution
- two 0.1 mm thick foils of black low density PE to further reduce IR and visible radiation
- setup with 10 Hz chopper and Lock-In amplifier
- detector and aperture are mounted on a x-y imaging stage and scanned vs distance and lateral position relative to the vacuum window

Courtesy E. Bründermann,
U Bochum



incoherent, $A_{\max} \approx 0.1$ mV



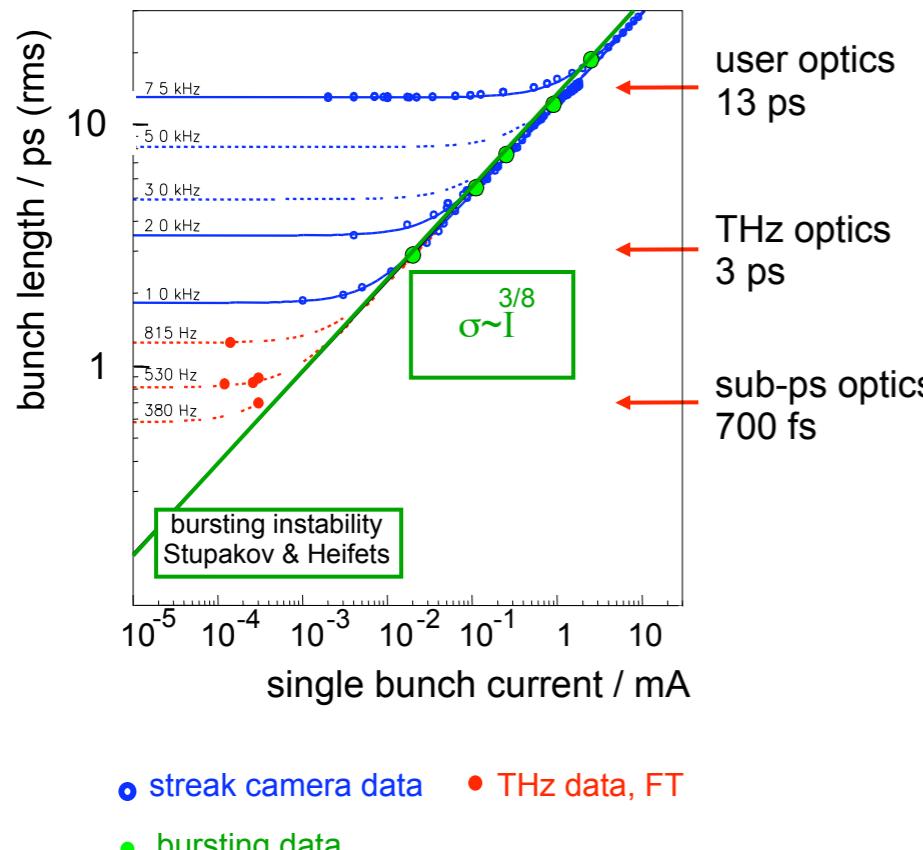
coherent, $A_{\max} \approx 2.9$ mV

Bursting Threshold

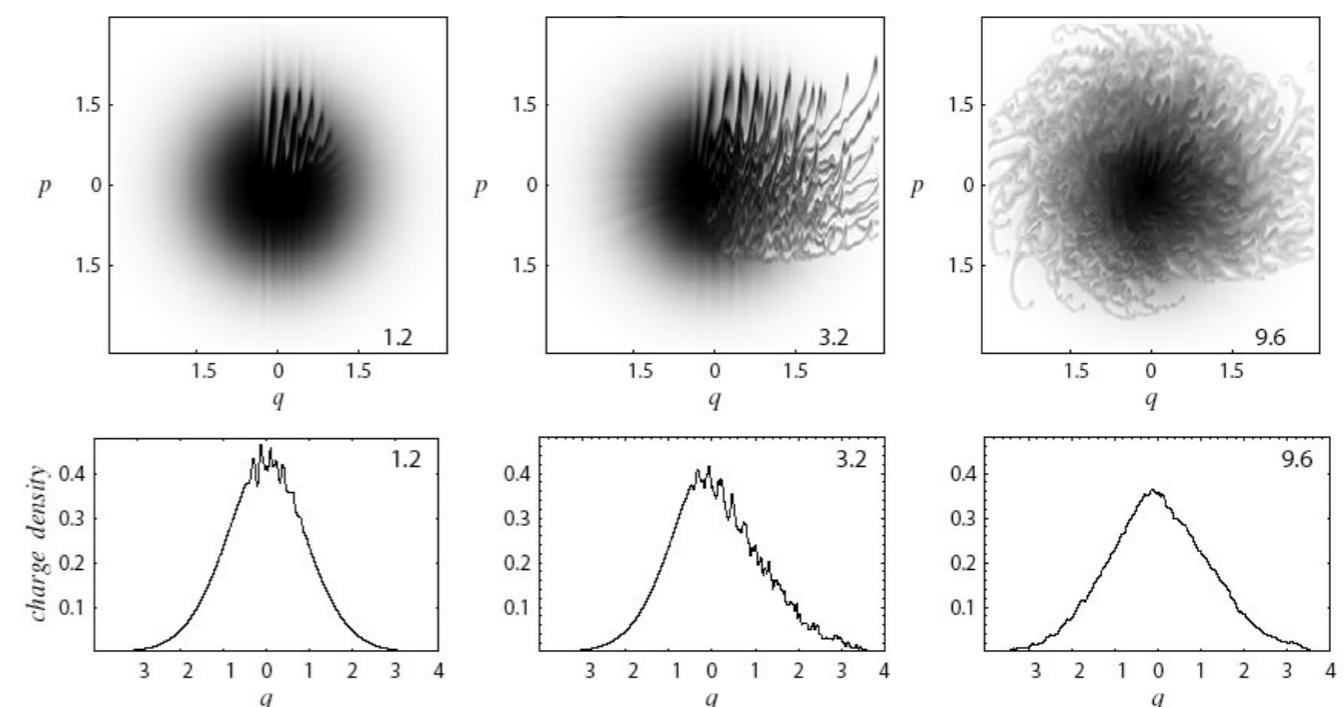
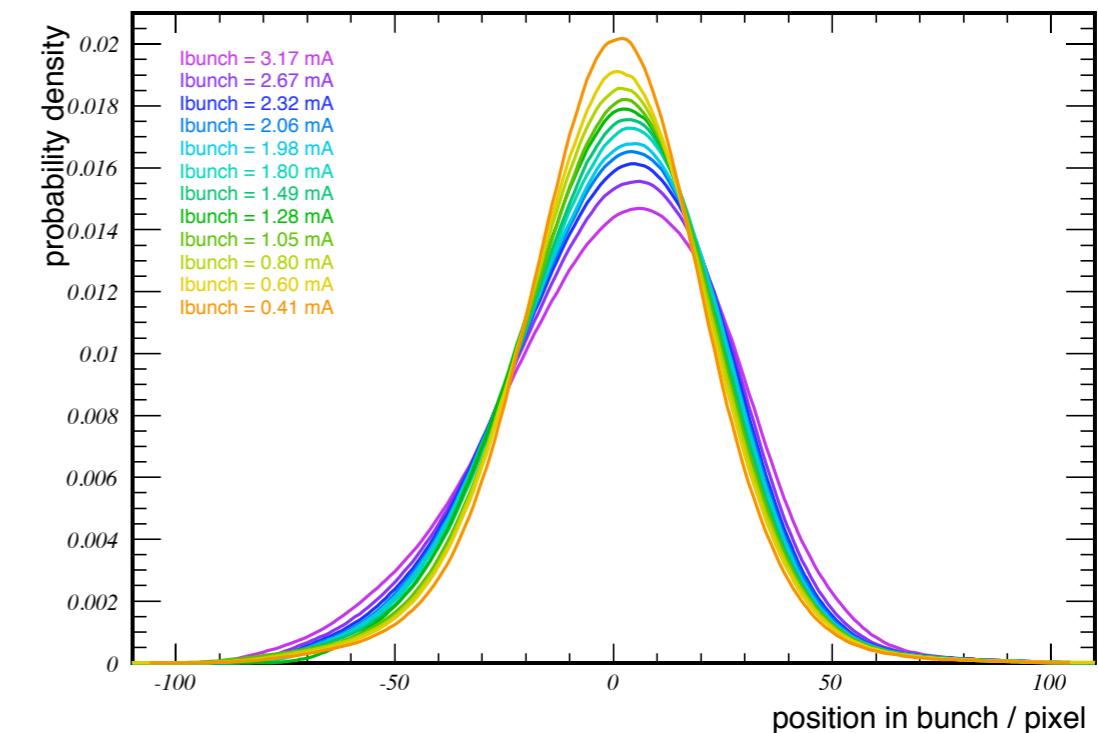
CSR-Bunch interaction:

- deformation with increasing current
- above threshold a microbunching instability results in (periodic) burst of high intensity

BESSY II bunch length - current scaling

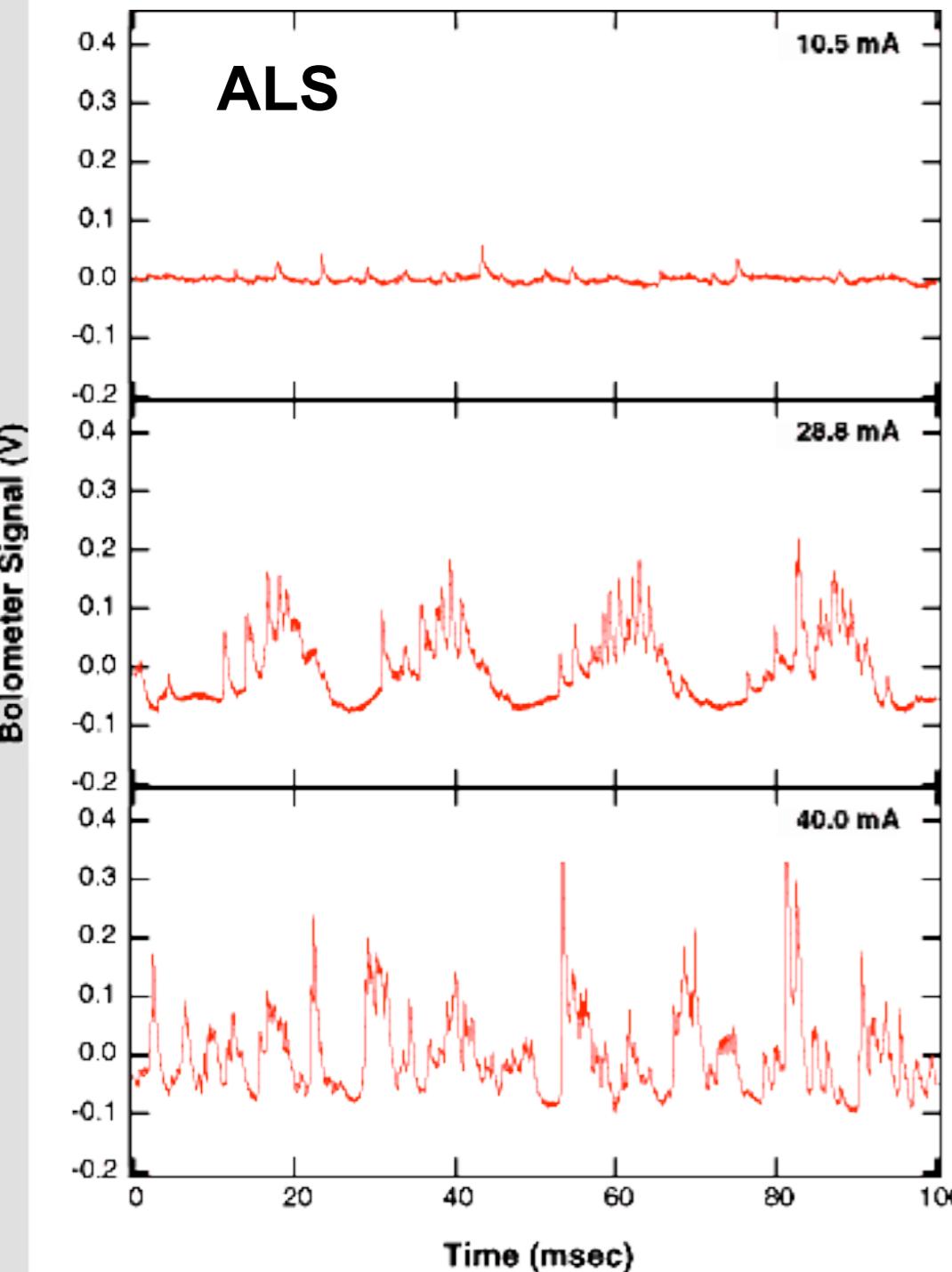


G. Wüstefeld, HZB

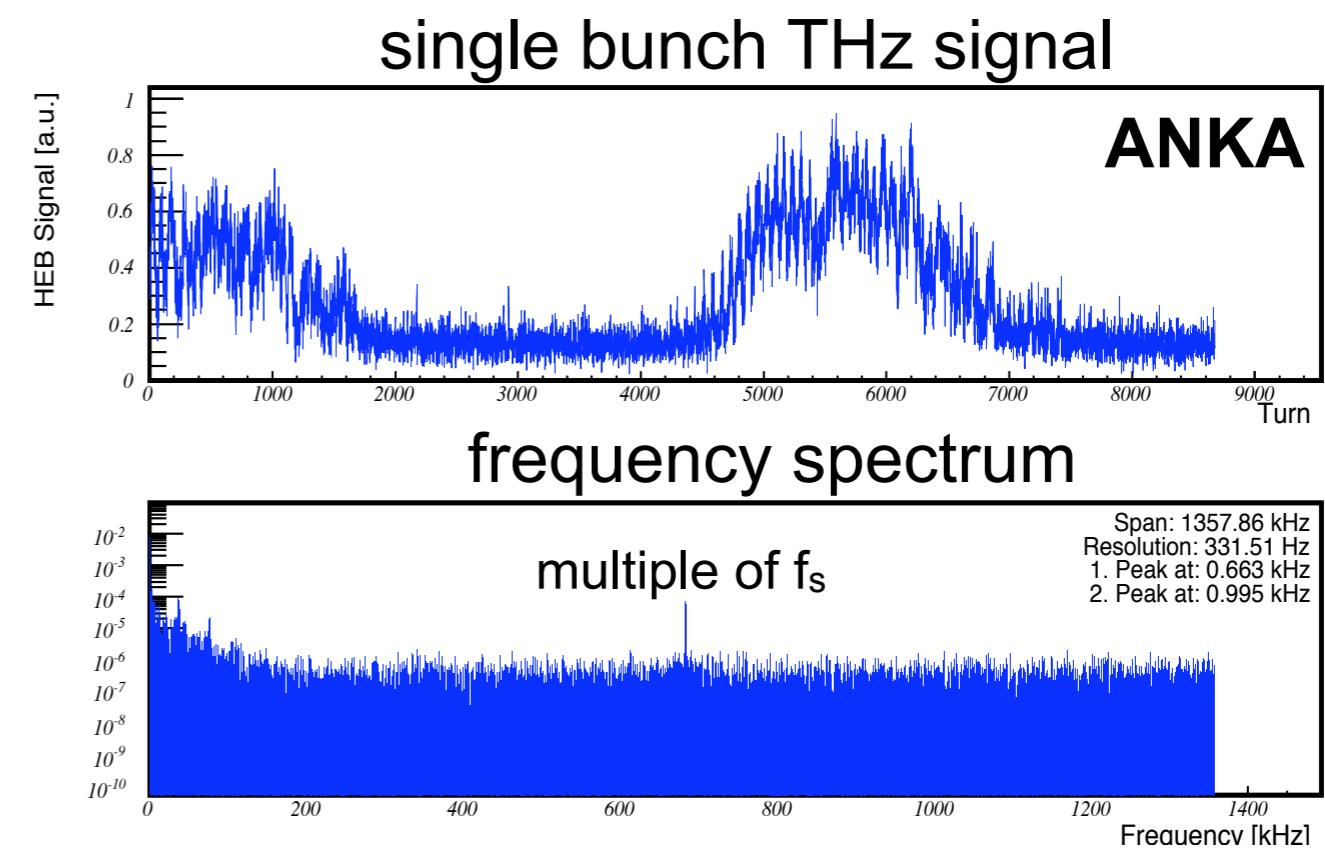


M. Venturini and R. Warnock, PRL 89, 224802 (2002)

Time Evolution of CSR Emission



■ Saturation of the generating instability and subsequent radiation damping leads to a sawtooth-like pattern as a function of time



J. Byrd et al., PRL 89, 224801 (2002)

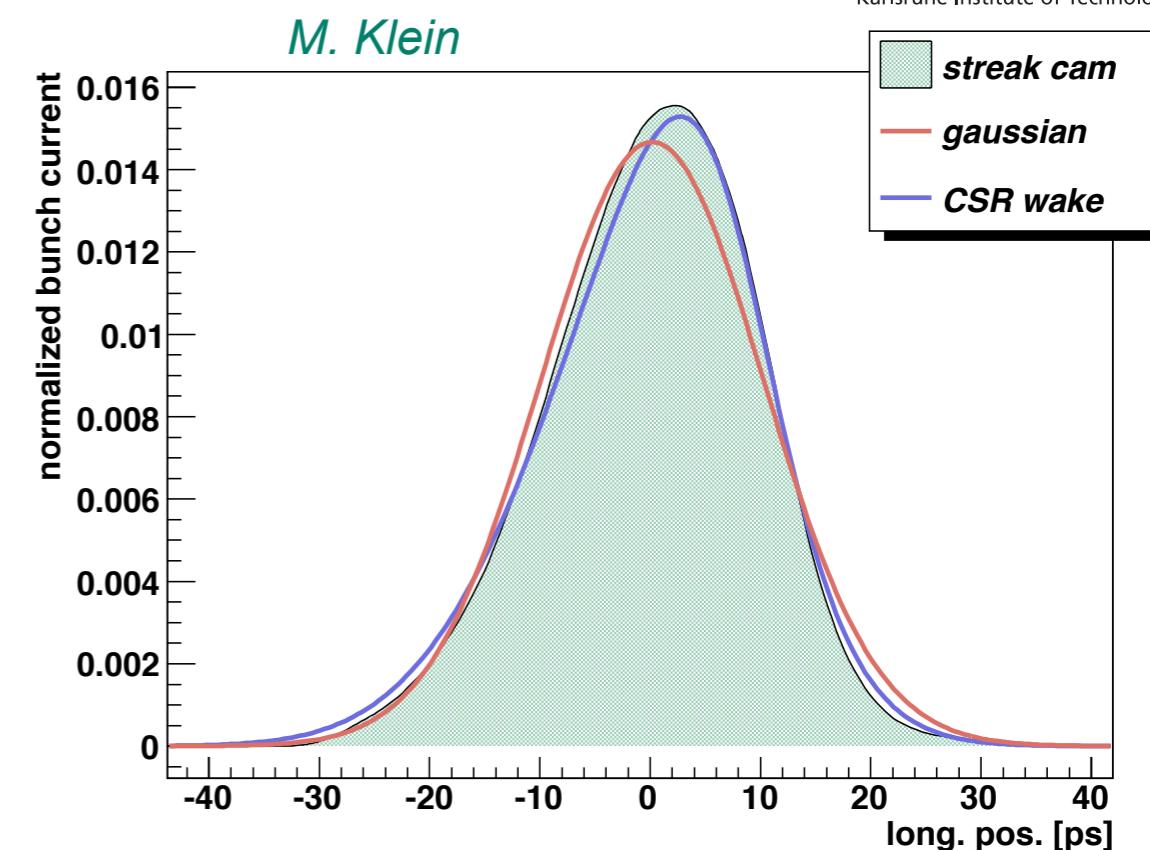
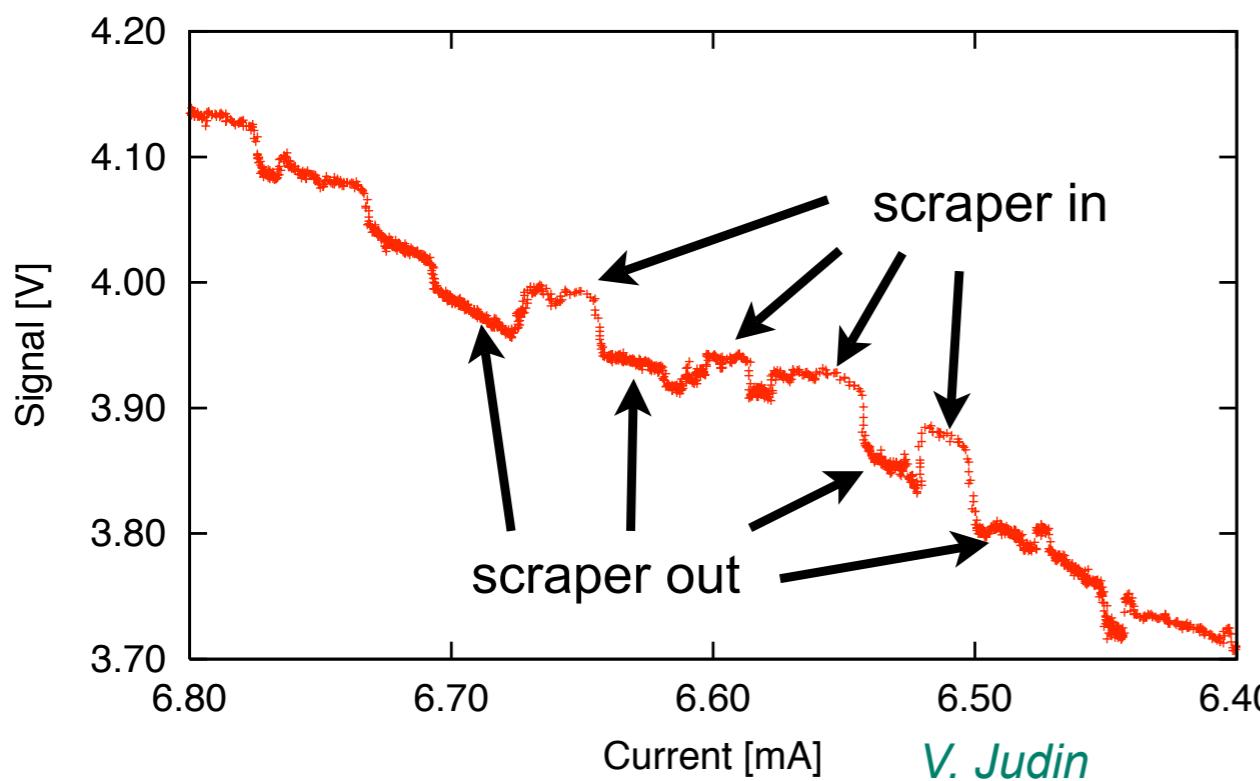
V. Judin

Impedance & CSR Power

- The total power radiated by a bunch of N particles is described by

$$P_{\text{total}} = N P_{\text{incoh}} (1 + N f_\lambda)$$

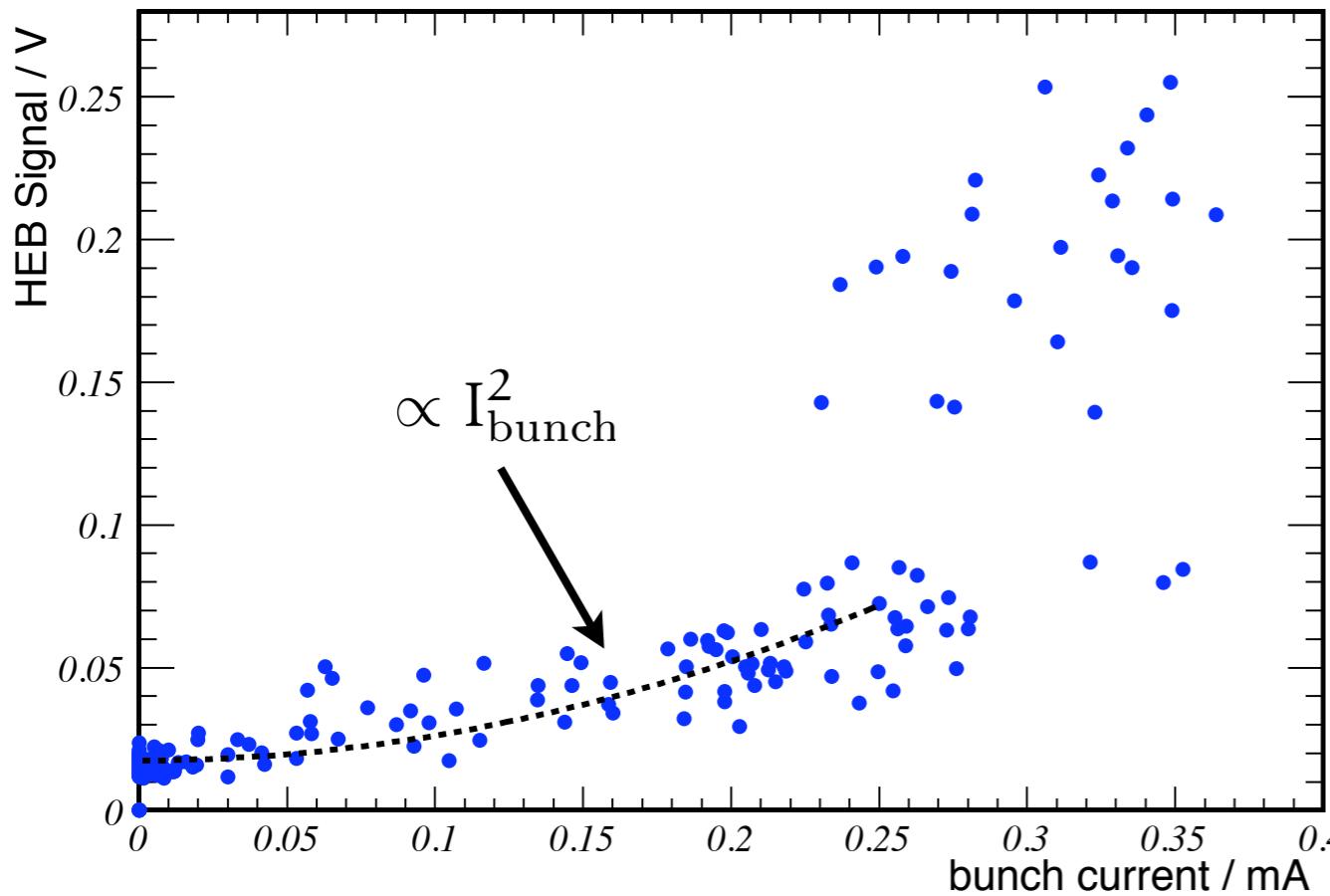
→ change in form factor f_λ is seen on the emitted THz power



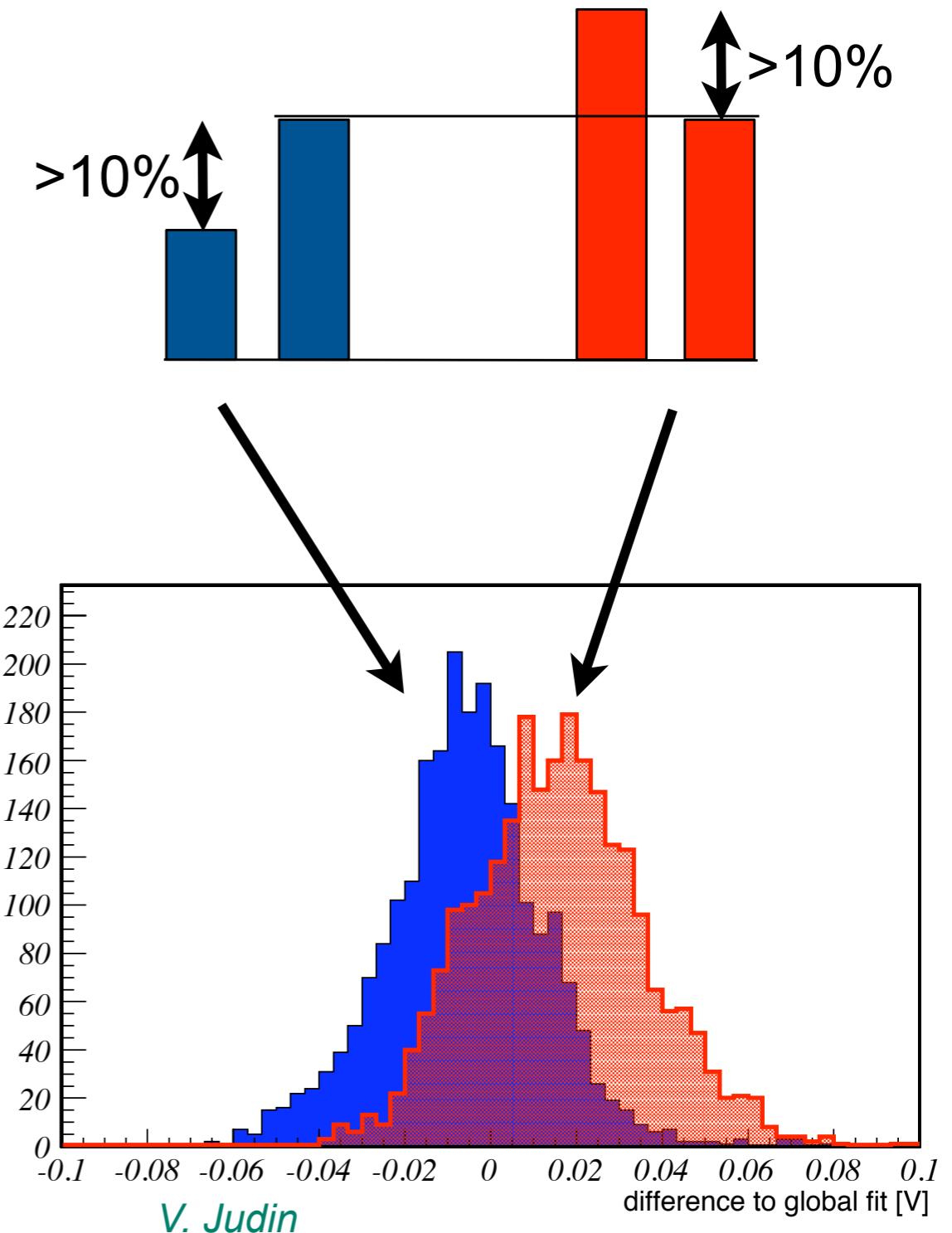
- Controlled change of the impedance by an asymmetric vertical scraper
- clear influence on emitted CSR

Single & Multi-Bunch Effects

- Fast THz detector (HEB) allows to study signals from individual bunches in a multi-bunch environment



→ THz emission depends on filling pattern



Summary

- Unshielded CSR is an important effect for electron rings with short bunches
 - radiation from main and fringe fields
 - mainly in the THz range
- Dedicated low- α_c optics
- CSR emission changes
 - with bunch current (stable → bursting)
 - with shape of charge distribution (CSR & other impedances)



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(CSR & other impedances)

