HIGH PRESSURE XENON GAS TPCS WITH SINGLE BARIUM ION IDENTIFICATION

Ben Jones

University of Texas at Arlington

- **Q: "What gives you hope that we could achieve anything like what you guys achieved in constructing the Standard Model today?"**

- **A. "The obvious answer is neutrino masses, which I think clearly take us beyond the standard model, and clearly represent something coming down from a very high energy scale."**

Stephen Weinberg at the SLAC Summer Institute, 2 weeks ago

$$-\frac{1}{2}\partial_\nu g^a_\mu \partial_\nu g^a_\mu - g_s f^{abc}\partial_\mu g^a_\nu g^b_\mu g^c_\nu - \frac{1}{4}g_s^2 f^{abc} f^{ade} g^b_\mu g^c_\nu g^d_\mu g^e_\nu +$$
$$\frac{1}{2}ig_s^2(\bar{q}^\sigma_i \gamma^\mu q^\sigma_j)g^a_\mu + \bar{G}^a\partial^2 G^a + g_s f^{abc}\partial_\mu \bar{G}^a G^b g^c_\mu - \partial_\nu W^+_\mu \partial_\nu W^-_\mu -$$
$$M^2 W^+_\mu W^-_\mu - \frac{1}{2}\partial_\nu Z^0_\mu \partial_\nu Z^0_\mu - \frac{1}{2c_w^2}M^2 Z^0_\mu Z^0_\mu - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H -$$
$$\frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2\phi^+\phi^- - \frac{1}{2}\partial_\mu\phi^0\partial_\mu\phi^0 - \frac{1}{2c_w^2}M\phi^0\phi^0 - \beta_h[\frac{2M^2}{g^2} +$$
$$\frac{2M}{g}H + \frac{1}{2}(H^2 + \phi^0\phi^0 + 2\phi^+\phi^-)] + \frac{2M^4}{g^2}\alpha_h - igc_w[\partial_\nu Z^0_\mu(W^+_\mu W^-_\nu -$$
$$W^+_\nu W^-_\mu) - Z^0_\nu(W^+_\mu \partial_\nu W^-_\mu - W^-_\mu \partial_\nu W^+_\mu) + Z^0_\mu(W^+_\nu \partial_\nu W^-_\mu -$$
$$W^-_\nu \partial_\nu W^+_\mu)] - igs_w[\partial_\nu A_\mu(W^+_\mu W^-_\nu - W^+_\nu W^-_\mu) - A_\nu(W^+_\mu \partial_\nu W^-_\mu -$$
$$W^-_\mu \partial_\nu W^+_\mu) + A_\mu(W^+_\nu \partial_\nu W^-_\mu - W^-_\nu \partial_\nu W^+_\mu)] - \frac{1}{2}g^2 W^+_\mu W^-_\mu W^+_\nu W^-_\nu +$$
$$\frac{1}{2}g^2 W^+_\mu W^-_\nu W^+_\mu W^-_\nu + g^2 c_w^2(Z^0_\mu W^+_\mu Z^0_\nu W^-_\nu - Z^0_\mu Z^0_\mu W^+_\nu W^-_\nu) +$$
$$g^2 s_w^2(A_\mu W^+_\mu A_\nu W^-_\nu - A_\mu A_\mu W^+_\nu W^-_\nu) + g^2 s_w c_w[A_\mu Z^0_\nu(W^+_\mu W^-_\nu -$$
$$W^+_\nu W^-_\mu) - 2A_\mu Z^0_\mu W^+_\nu W^-_\nu] - g\alpha[H^3 + H\phi^0\phi^0 + 2H\phi^+\phi^-] -$$
$$\frac{1}{8}g^2\alpha_h[H^4 + (\phi^0)^4 + 4(\phi^+\phi^-)^2 + 4(\phi^0)^2\phi^+\phi^- + 4H^2\phi^+\phi^- + 2(\phi^0)^2 H^2] -$$
$$gMW^+_\mu W^-_\mu H - \frac{1}{2}g\frac{M}{c_w^2}Z^0_\mu Z^0_\mu H - \frac{1}{2}ig[W^+_\mu(\phi^0\partial_\mu\phi^- - \phi^-\partial_\mu\phi^0) -$$
$$W^-_\mu(\phi^0\partial_\mu\phi^+ - \phi^+\partial_\mu\phi^0)] + \frac{1}{2}g[W^+_\mu(H\partial_\mu\phi^- - \phi^-\partial_\mu H) - W^-_\mu(H\partial_\mu\phi^+ -$$
$$\phi^+\partial_\mu H)] + \frac{1}{2}g\frac{1}{c_w}(Z^0_\mu(H\partial_\mu\phi^0 - \phi^0\partial_\mu H) - ig\frac{s_w^2}{c_w}MZ^0_\mu(W^+_\mu\phi^- - W^-_\mu\phi^+) +$$
$$igs_w MA_\mu(W^+_\mu\phi^- - W^-_\mu\phi^+) - ig\frac{1-2c_w^2}{2c_w}Z^0_\mu(\phi^+\partial_\mu\phi^- - \phi^-\partial_\mu\phi^+) +$$
$$igs_w A_\mu(\phi^+\partial_\mu\phi^- - \phi^-\partial_\mu\phi^+) - \frac{1}{4}g^2 W^+_\mu W^-_\mu[H^2 + (\phi^0)^2 + 2\phi^+\phi^-] -$$
$$\frac{1}{4}g^2\frac{1}{c_w^2}Z^0_\mu Z^0_\mu[H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2\phi^+\phi^-] - \frac{1}{2}g^2\frac{s_w^2}{c_w}Z^0_\mu\phi^0(W^+_\mu\phi^- +$$
$$W^-_\mu\phi^+) - \frac{1}{2}ig^2\frac{s_w^2}{c_w}Z^0_\mu H(W^+_\mu\phi^- - W^-_\mu\phi^+) + \frac{1}{2}g^2 s_w A_\mu\phi^0(W^+_\mu\phi^- +$$
$$W^-_\mu\phi^+) + \frac{1}{2}ig^2 s_w A_\mu H(W^+_\mu\phi^- - W^-_\mu\phi^+) - g^2\frac{s_w}{c_w}(2c_w^2 - 1)Z^0_\mu A_\mu\phi^+\phi^- -$$
$$g^1 s_w^2 A_\mu A_\mu\phi^+\phi^- - \bar{e}^\lambda(\gamma\partial + m_e^\lambda)e^\lambda - \bar{\nu}^\lambda\gamma\partial\nu^\lambda - \bar{u}^\lambda_j(\gamma\partial + m_u^\lambda)u^\lambda_j -$$
$$\bar{d}^\lambda_j(\gamma\partial + m_d^\lambda)d^\lambda_j + igs_w A_\mu[-(\bar{e}^\lambda\gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}^\lambda_j\gamma^\mu u^\lambda_j) - \frac{1}{3}(\bar{d}^\lambda_j\gamma^\mu d^\lambda_j)] +$$
$$\frac{ig}{4c_w}Z^0_\mu[(\bar{\nu}^\lambda\gamma^\mu(1+\gamma^5)\nu^\lambda) + (\bar{e}^\lambda\gamma^\mu(4s_w^2 - 1 - \gamma^5)e^\lambda) + (\bar{u}^\lambda_j\gamma^\mu(\frac{4}{3}s_w^2 -$$
$$1 - \gamma^5)u^\lambda_j) + (\bar{d}^\lambda_j\gamma^\mu(1 - \frac{8}{3}s_w^2 - \gamma^5)d^\lambda_j)] + \frac{ig}{2\sqrt{2}}W^+_\mu[(\bar{\nu}^\lambda\gamma^\mu(1+\gamma^5)e^\lambda) +$$
$$(\bar{u}^\lambda_j\gamma^\mu(1+\gamma^5)C_{\lambda\kappa}d^\kappa_j)] + \frac{ig}{2\sqrt{2}}W^-_\mu[(\bar{e}^\lambda\gamma^\mu(1+\gamma^5)\nu^\lambda) + (\bar{d}^\kappa_j C^\dagger_{\lambda\kappa}\gamma^\mu(1 +$$
$$\gamma^5)u^\lambda_j)] + \frac{ig}{2\sqrt{2}}\frac{m_e^\lambda}{M}[-\phi^+(\bar{\nu}^\lambda(1-\gamma^5)e^\lambda) + \phi^-(\bar{e}^\lambda(1+\gamma^5)\nu^\lambda)] -$$
$$\frac{g}{2}\frac{m_e^\lambda}{M}[H(\bar{e}^\lambda e^\lambda) + i\phi^0(\bar{e}^\lambda\gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}}\phi^+[-m_d^\kappa(\bar{u}^\lambda_j C_{\lambda\kappa}(1-\gamma^5)d^\kappa_j) +$$
$$m_u^\lambda(\bar{u}^\lambda_j C_{\lambda\kappa}(1+\gamma^5)d^\kappa_j] + \frac{ig}{2M\sqrt{2}}\phi^-[m_d^\lambda(\bar{d}^\lambda_j C^\dagger_{\lambda\kappa}(1+\gamma^5)u^\kappa_j) - m_u^\kappa(\bar{d}^\lambda_j C^\dagger_{\lambda\kappa}(1 -$$
$$\gamma^5)u^\kappa_j] - \frac{g}{2}\frac{m_u^\lambda}{M}H(\bar{u}^\lambda_j u^\lambda_j) - \frac{g}{2}\frac{m_d^\lambda}{M}H(\bar{d}^\lambda_j d^\lambda_j) + \frac{ig}{2}\frac{m_u^\lambda}{M}\phi^0(\bar{u}^\lambda_j\gamma^5 u^\lambda_j) -$$
$$\frac{ig}{2}\frac{m_d^\lambda}{M}\phi^0(\bar{d}^\lambda_j\gamma^5 d^\lambda_j) + \bar{X}^+(\partial^2 - M^2)X^+ + \bar{X}^-(\partial^2 - M^2)X^- + \bar{X}^0(\partial^2 -$$
$$\frac{M^2}{c_w^2})X^0 + \bar{Y}\partial^2 Y + igc_w W^+_\mu(\partial_\mu\bar{X}^0 X^- - \partial_\mu\bar{X}^+ X^0) + igs_w W^+_\mu(\partial_\mu\bar{Y}X^- -$$
$$\partial_\mu\bar{X}^+ Y) + igc_w W^-_\mu(\partial_\mu\bar{X}^- X^0 - \partial_\mu\bar{X}^0 X^+) + igs_w W^-_\mu(\partial_\mu\bar{X}^- Y -$$
$$\partial_\mu\bar{Y}X^+) + igc_w Z^0_\mu(\partial_\mu\bar{X}^+ X^+ - \partial_\mu\bar{X}^- X^-) + igs_w A_\mu(\partial_\mu\bar{X}^+ X^+ -$$
$$\partial_\mu\bar{X}^- X^-) - \frac{1}{2}gM[\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w^2}\bar{X}^0 X^0 H] +$$
$$\frac{1-2c_w^2}{2c_w}igM[\bar{X}^+ X^0\phi^+ - \bar{X}^- X^0\phi^-] + \frac{1}{2c_w}igM[\bar{X}^0 X^-\phi^+ - \bar{X}^0 X^+\phi^-] +$$
$$igMs_w[\bar{X}^0 X^-\phi^+ - \bar{X}^0 X^+\phi^-] + \frac{1}{2}igM[\bar{X}^+ X^+\phi^0 - \bar{X}^- X^-\phi^0]$$

← The standard model Lagrangian we all know and love

Write down all terms that are:

1) Consistent with the symmetries groups of the SM

Gauge invariance:

$$SU(3) \times SU(2)_L \times U(1)_Y$$

+ Lorentz invariance

2) Renormalizible

If SM is a low energy effective theory:

$$L = L_{SM} + \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \ldots$$

All that stuff

Dimension-5 terms

Dimension-6 terms

Etc…

E_{new} > 10 TeV — **Seems likely (LHC)**

E_{new} ~ 10^{13} TeV — **Maybe? (GUT scale)**

If SM is a low energy effective theory:

$$L = L_{SM} + \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \dots$$

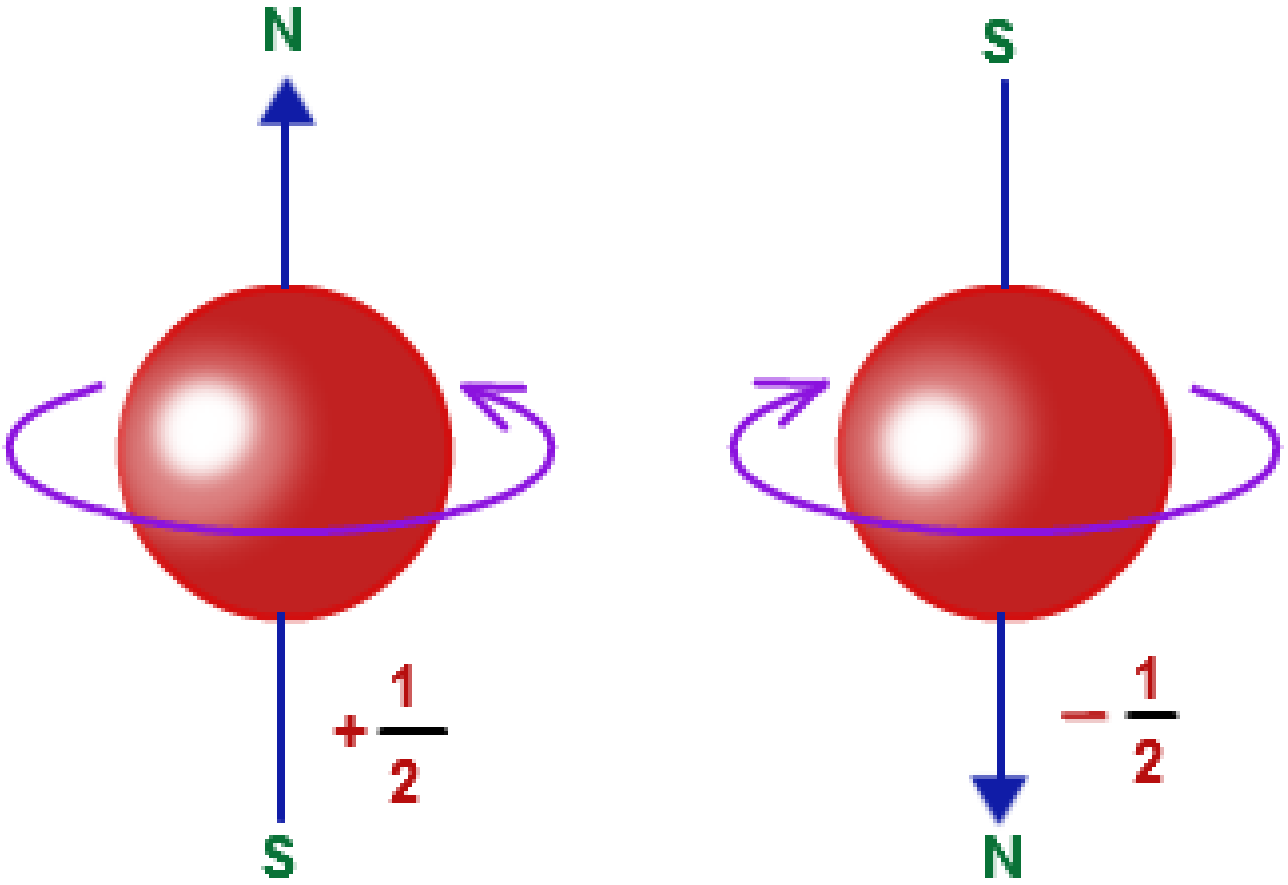
- The only dimension-5 operator one can add obeying SM gauge symmetry:

$$\frac{L_1}{E_{new}} = y_{ij} \frac{\nu^i H \nu^j H}{E_{new}}$$

Weinberg 1979.

- This term does an important thing - it makes neutrinos Majorana particles, with mass suppressed by the new physics scale.
- And it makes the theory non-renormalizible – implying there must be something else at high scale (see also: Seasaw)

The Weirdness of Majorana Fermions



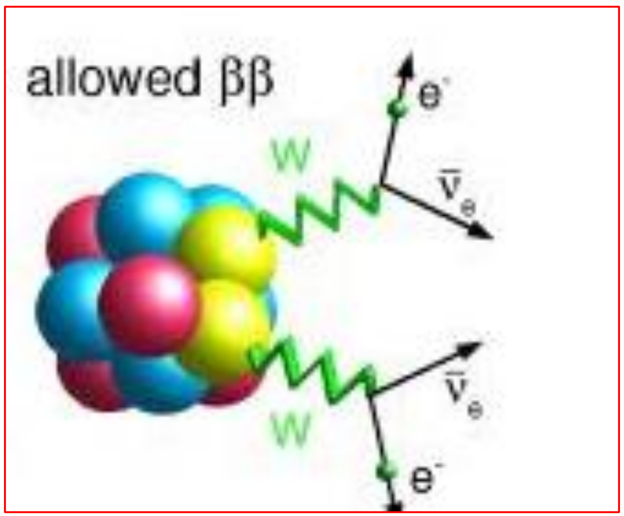
Behaves like matter **Behaves like antimatter**

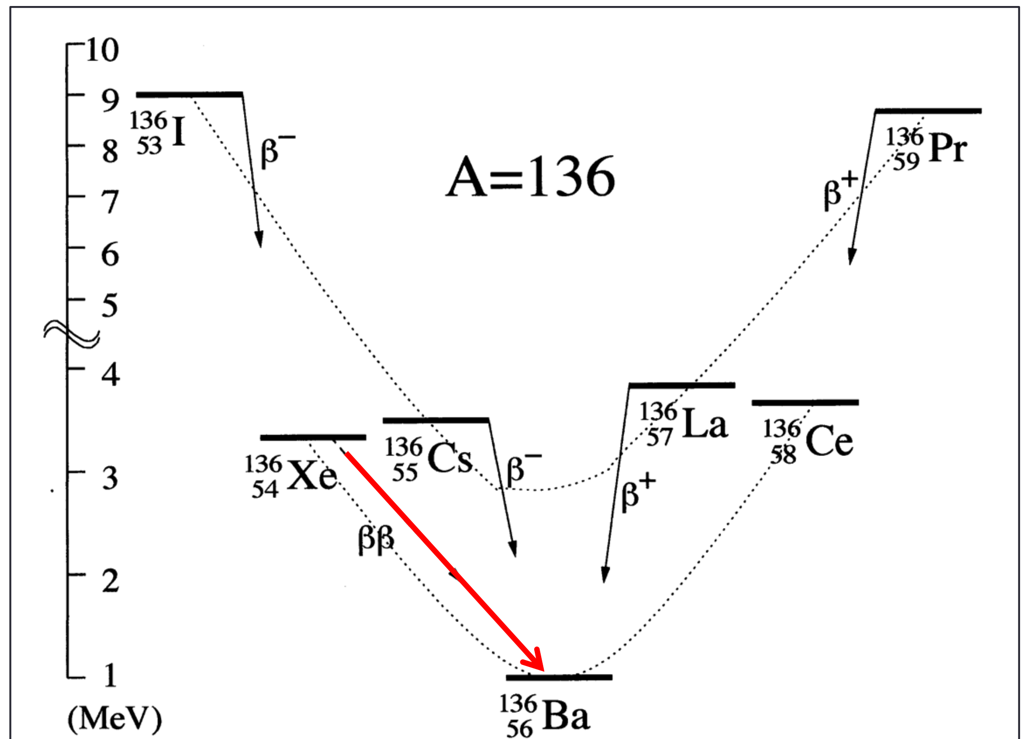
Double-Beta Decay

A rare radioactive process, energetically allowed for some even-even nuclei where $m_{Z+1} > m_Z > m_{Z+2}$

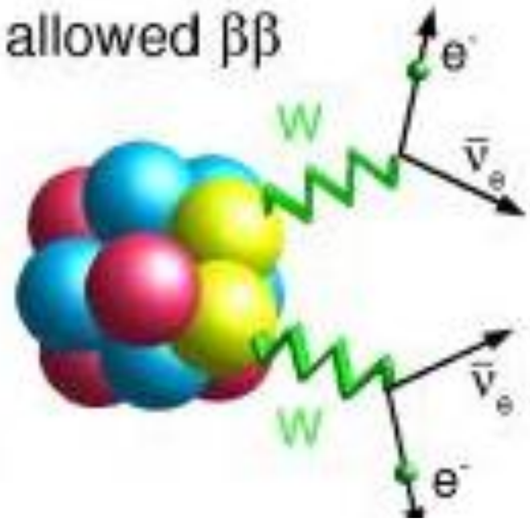
- $(Z,A) \rightarrow (Z+2,A) + e^-_1 + \bar{\nu}_1 + e^-_2 + \bar{\nu}_2$

Discovered in 1987 in ^{82}Se and now seen in multiple isotopes





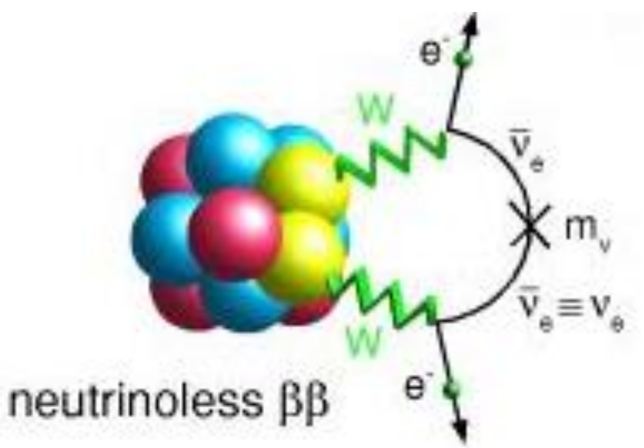
Types of double beta decay



A known standard model process and an important calibration tool

$$T_{\frac{1}{2}} \gg 10^{19-21}\,yrs.$$

Final state: $e^- e^- \nu_e \nu_e$



Observation would prove that the neutrino is a Majorana fermion

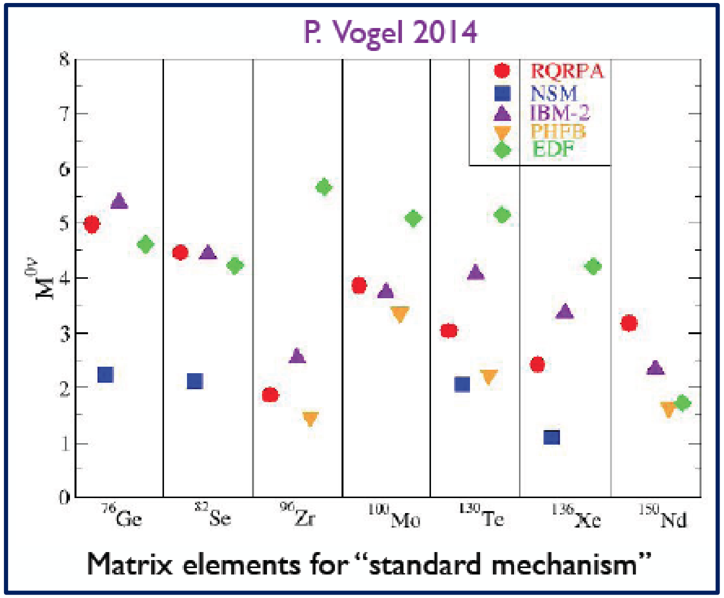
$$\frac{1}{T_{\frac{1}{2}}} = G \times \|\mathbf{M}\|^2 \times m_{\bar{\nu}}^{\;2}$$

Final state: $e^- e^-$

Tale as old
AS TIME.
Song as old
AS RHYME.

1. This formula →

$$\frac{1}{T_{\frac{1}{2}}} = G \times \left\| \mathbf{M} \right\|^2 \times m_{\bar{\nu}}{}^2$$



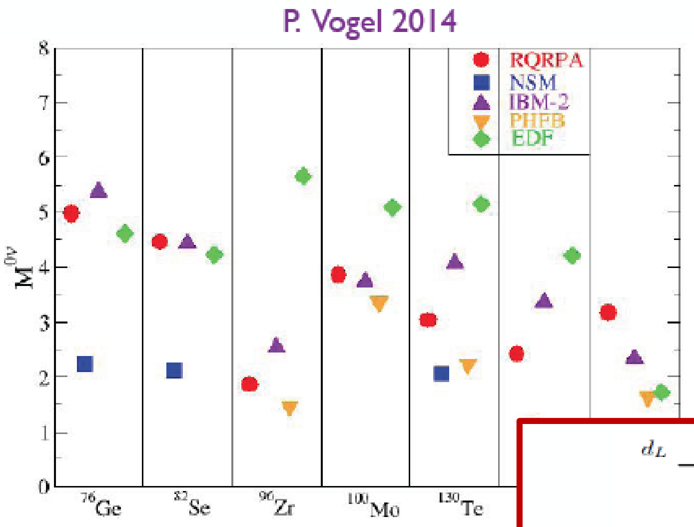
1. This formula →

$$\frac{1}{T_{\frac{1}{2}}} = G \times \|\mathbf{M}\|^2 \times m_{\bar{\nu}}{}^2$$

← 2. Nuclear matrix elements

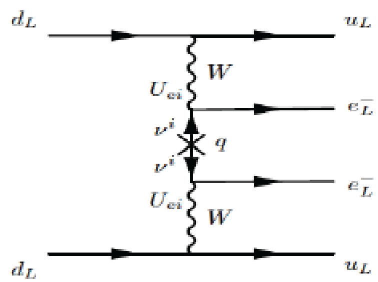
1. This formula →

$$\frac{1}{T_{\frac{1}{2}}} = G \times \|\mathbf{M}\|^2 \times m_{\bar{\nu}}^{\ 2}$$



Matrix elements for "standard me

← 2. Nuclear matrix elements

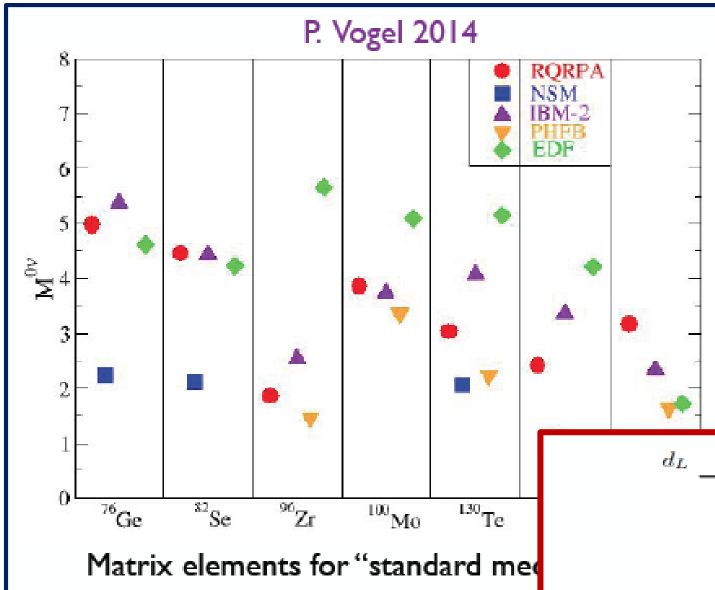


$$\frac{1}{T^{0\nu}_{1/2}} = G_{0\nu} g_A^4 \left| M^{(0\nu)} \right|^2 \langle m_{\beta\beta} \rangle^2$$

$$\langle m_{\beta\beta} \rangle^2 = |\sum U_{ei}^2 m_{\nu i}|^2$$

$$\langle m_{\beta\beta} \rangle^2 = |\cos^2\theta_{12}\cos^2\theta_{13} m_1 + \exp^{2i\lambda_2}\sin^2\theta_{12}\cos^2\theta_{13} m_2 + \exp^{2i(\lambda_3-\delta_{CP})}\sin^2\theta_{13} m_3|^2$$

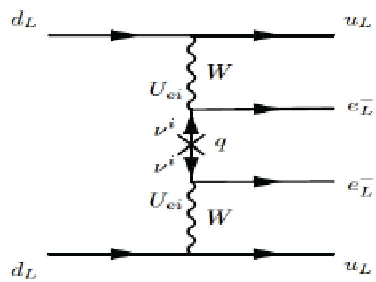
3. Effective nu mass ^



1. This formula →

$$\frac{1}{T_{\frac{1}{2}}} = G \times \|\mathbf{M}\|^2 \times m_{\bar{\nu}}^{\ 2}$$

← 2. Nuclear matrix elements

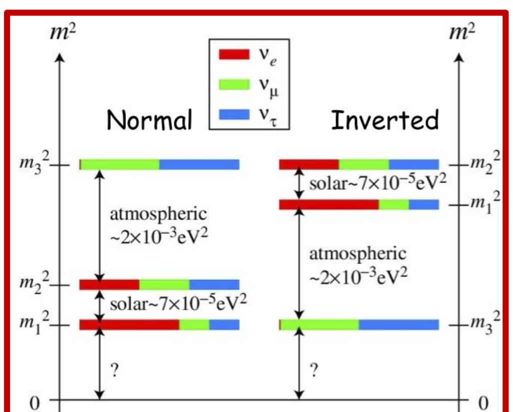


$$\frac{1}{T^{0\nu}_{1/2}} = G_{0\nu} g_A^4 \left| M^{(0\nu)} \right|^2 \langle m_{\beta\beta} \rangle^2$$

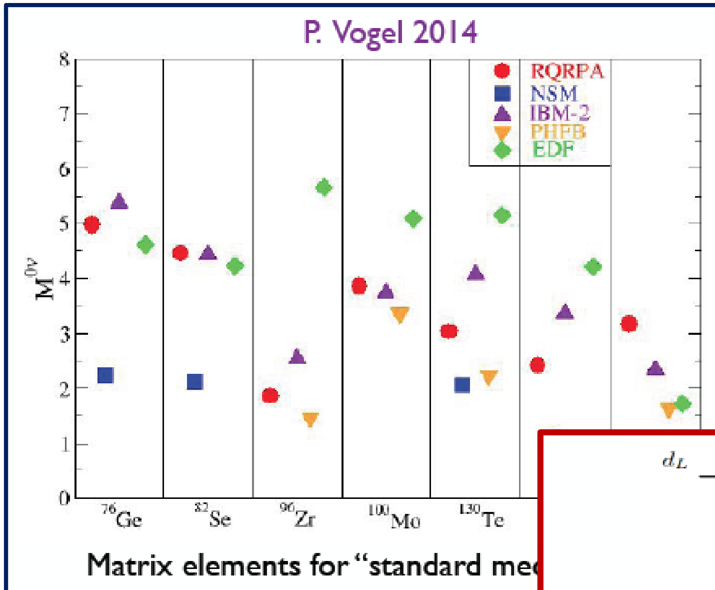
$$\langle m_{\beta\beta} \rangle^2 = |\sum U_{ei}^2 m_{\nu i}|^2$$

$$\langle m_{\beta\beta} \rangle^2 = |\cos^2\theta_{12}\cos^2\theta_{13} m_1 + \exp^{2i\lambda_2}\sin^2\theta_{12}\cos^2\theta_{13} m_2 + \exp^{2i(\lambda_3-\delta_{CP})}\sin^2\theta_{13} m_3|^2$$

3. Effective nu mass ^



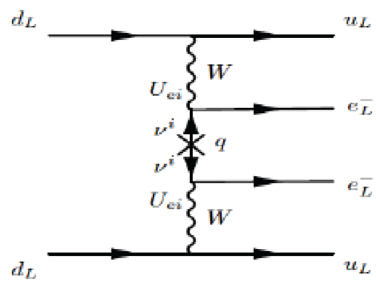
← Special bonus: connect it to this



1. This formula →

$$\frac{1}{T_{\frac{1}{2}}} = G \times \|\mathbf{M}\|^2 \times m_{\bar{\nu}}{}^2$$

← 2. Nuclear matrix elements

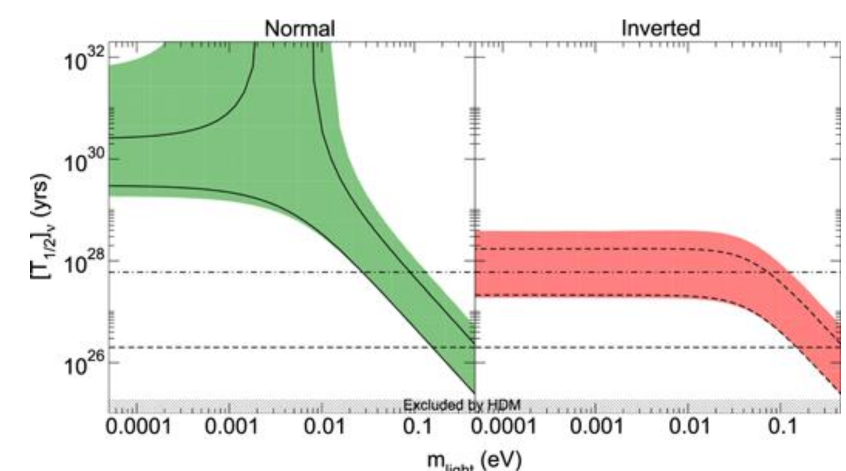


$$\frac{1}{T_{1/2}^{0\nu}} = G_{0\nu} g_A^4 \left|M^{(0\nu)}\right|^2 \langle m_{\beta\beta}\rangle^2$$

$$\langle m_{\beta\beta}\rangle^2 = \left|\sum U_{ei}^2 m_{\nu i}\right|^2$$

$$\langle m_{\beta\beta}\rangle^2 = |\cos^2\theta_{12}\cos^2\theta_{13} m_1 + \exp^{2i\lambda_2}\sin^2\theta_{12}\cos^2\theta_{13} m_2 + \exp^{2i(\lambda_3-\delta_{CP})}\sin^2\theta_{13} m_3|^2$$

3. Effective nu mass ^



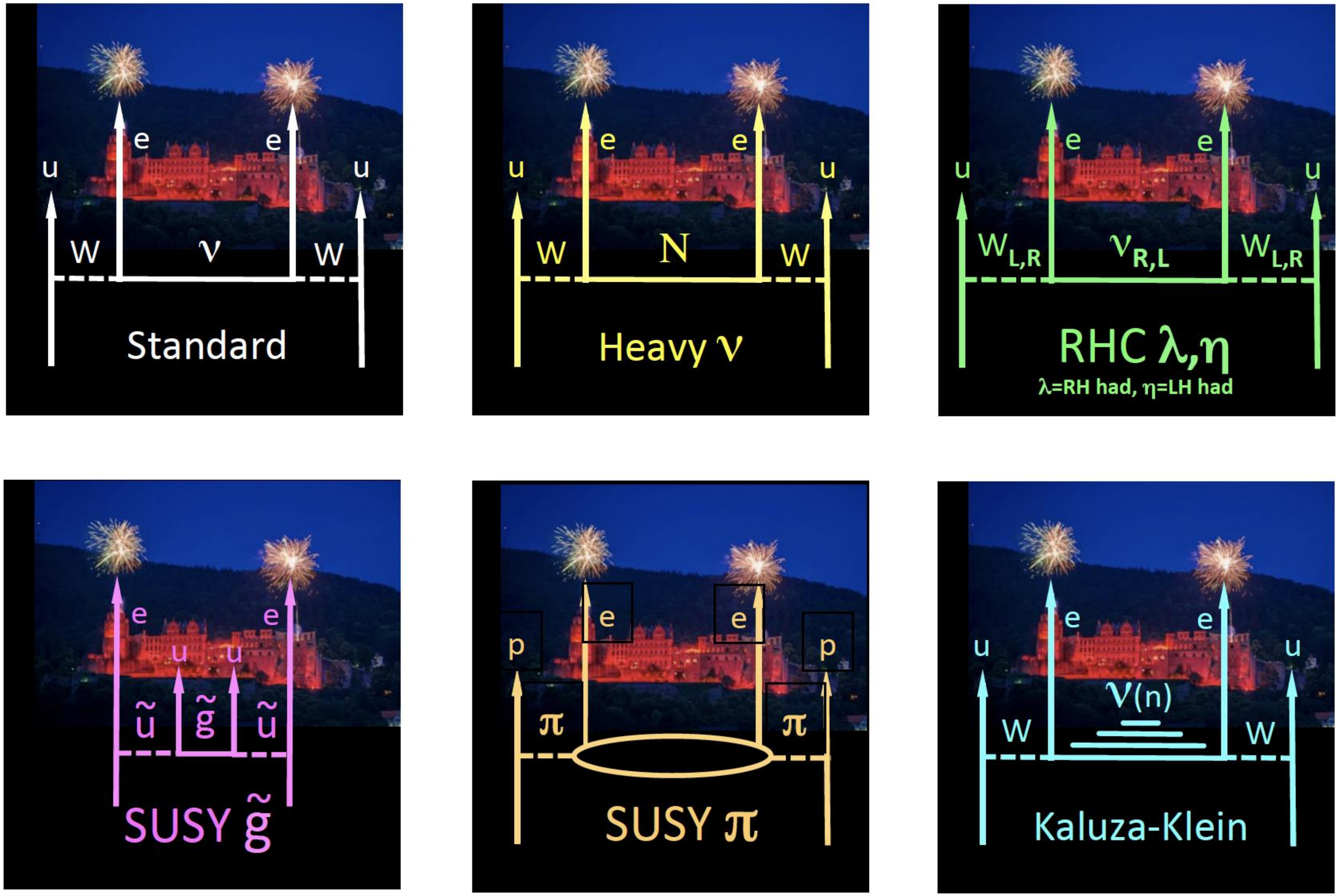
← 4. End up at this plot

Cute story, but…

- Strength of weak interaction is quenched in nuclear medium
  - Rate depends on 4th power of g_A
  - Nobody really knows where the "NH" and "IH" bands are.

- All those plots assume light Majorana neutrino exchange mechanism.
  - *Motivated by seesaw models.*
  - *But Majorana neutrinos imply generically new high-scale physics, and generically new physics can produce generic 0nubb rates.*

Warning: **don't stick to $m_{\beta\beta}$ metric, just go on with $T_{1/2}$!** Variety of $0\nu\beta\beta$ mechanisms:



$0\nu\beta\beta$ from any mechanism → **Majorana nature of ν would be established anyway**

Slide from The Mid and Long Term Future of Neutrinoless Double Beta Decay, Andrea Giuliani, Neutrinno2018, https://doi.org/10.5281/zenodo.1286915

A robust observation of 0nubb would tell us 6 things about nature before breakfast:

1) Lepton number conservation is violated.

2) Massive fermions exist that are neither matter or antimatter but something else (Majorana fermions)

3) The SM with the Majorana term is non-renomalizible → SM is definitely a low energy effective theory.

4) There are other mass generating mechanisms in nature beyond the Higgs mechanism.

5) Like G_F tells us the weak scale and ultimately leads to W and Z mass predictions, m_ν tells us the next scale.

6) Majorana neutrinos are a prediction of the theory of Leptogenesis that may generate observed matter/anti-matter asymmetry of the Universe (given enough CPV – see also: DUNE)

So how about we all agree to:

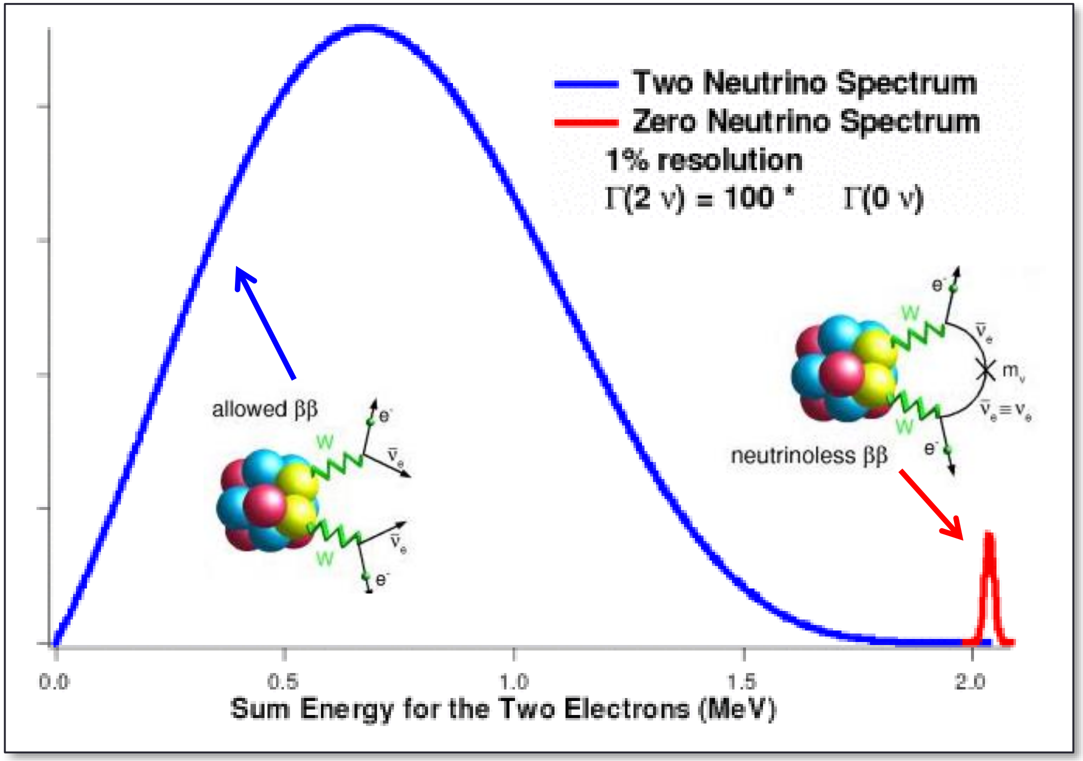
Be as sensitive as possible

(get out logarithmically into unexplored parameter space)

Build experiments that can make discoveries, not just set limits

(well understood detectors with clear, positive signal criteria)

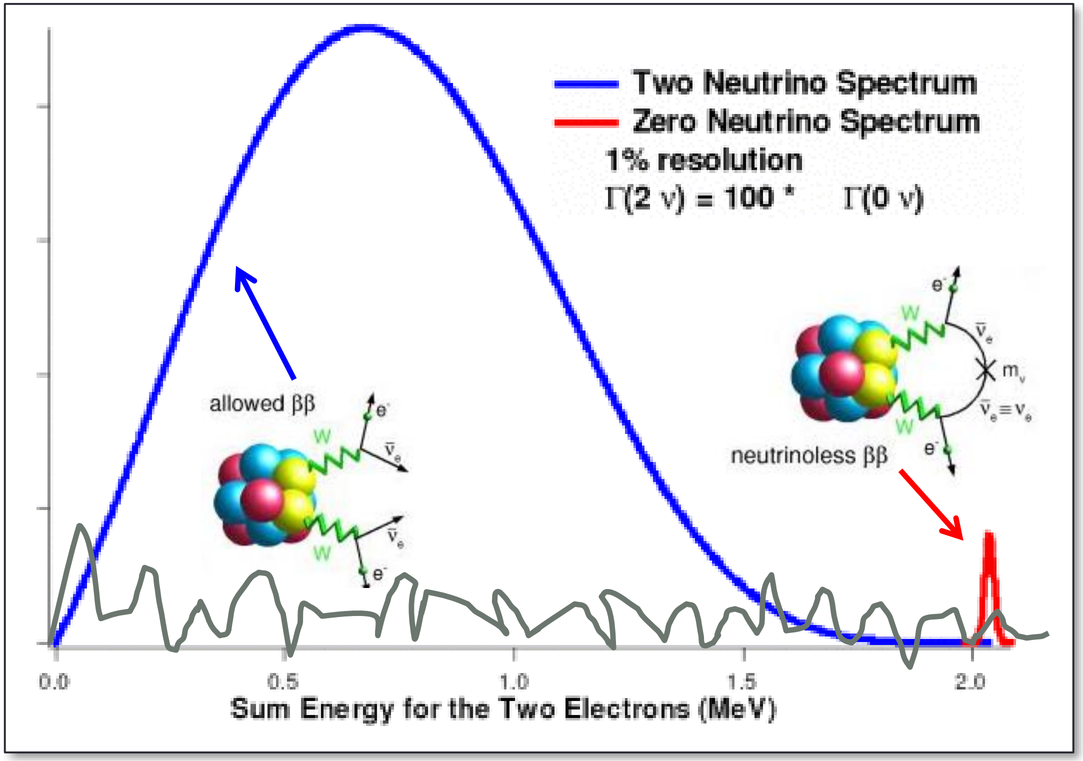
The ideal experiment:



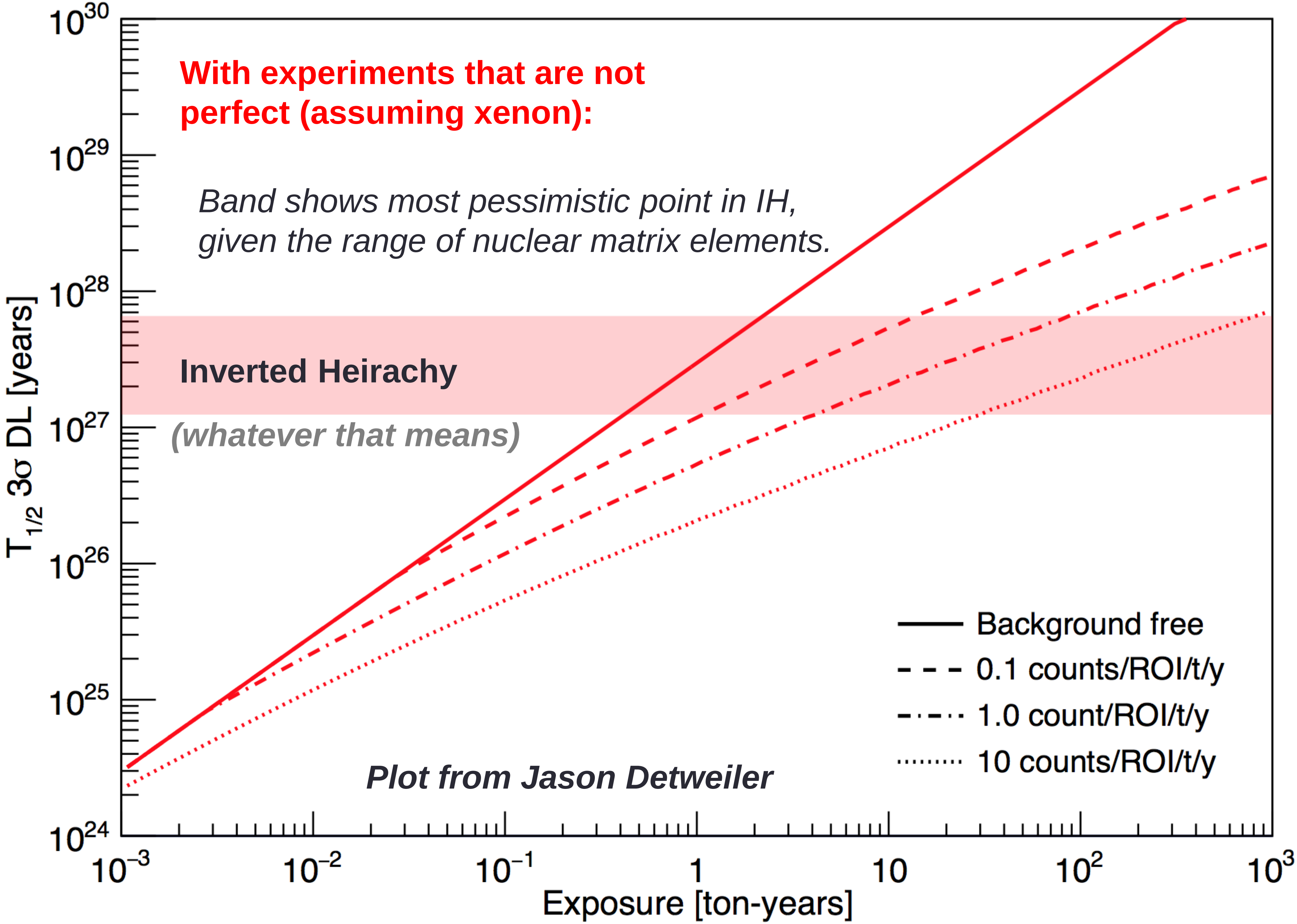
- ***MANDATORY:***
  - Resolution better than ~2% FWHM to fully reject two-neutrino mode

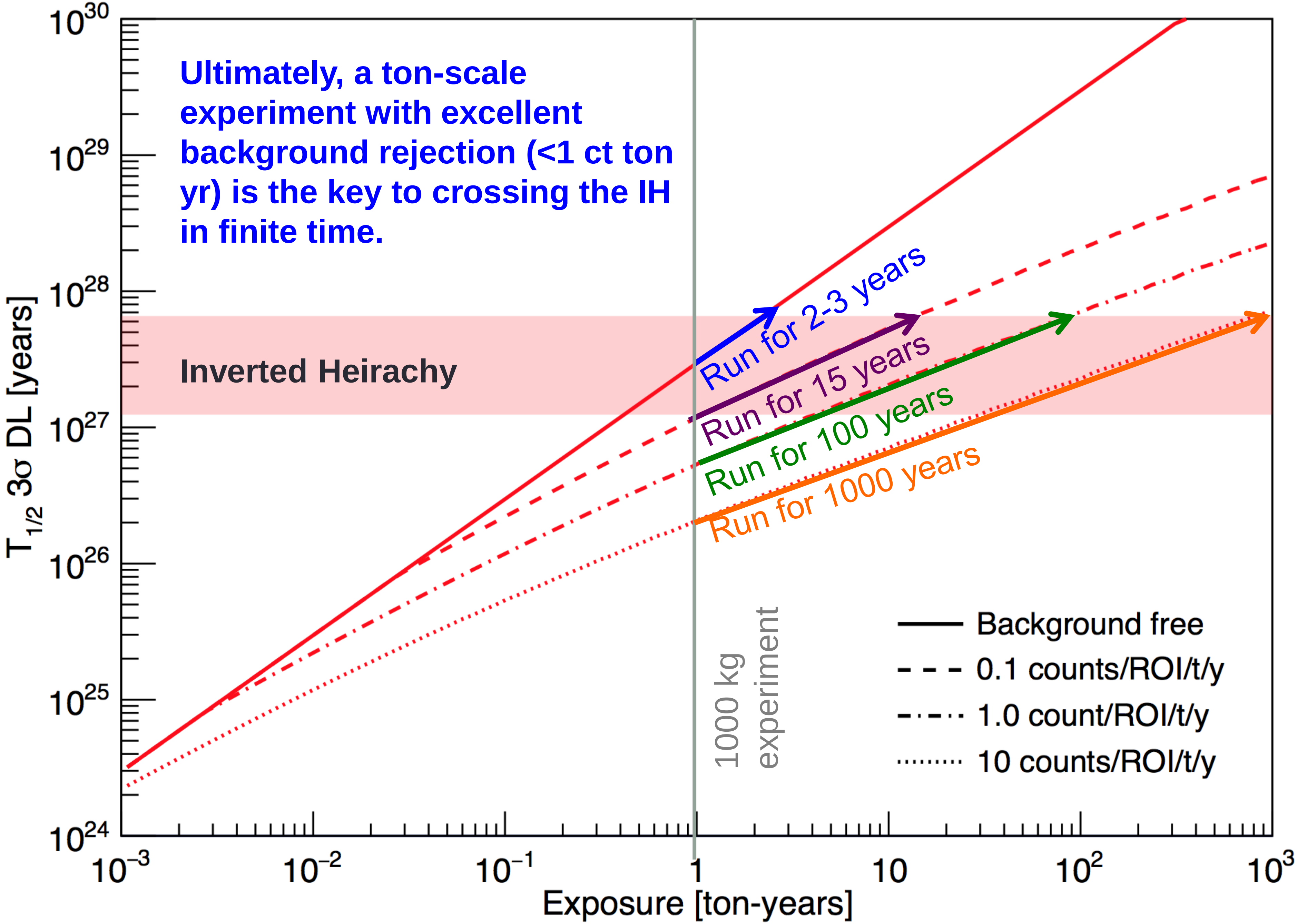
- ***Then just watch and wait…***

The non-ideal experiment:



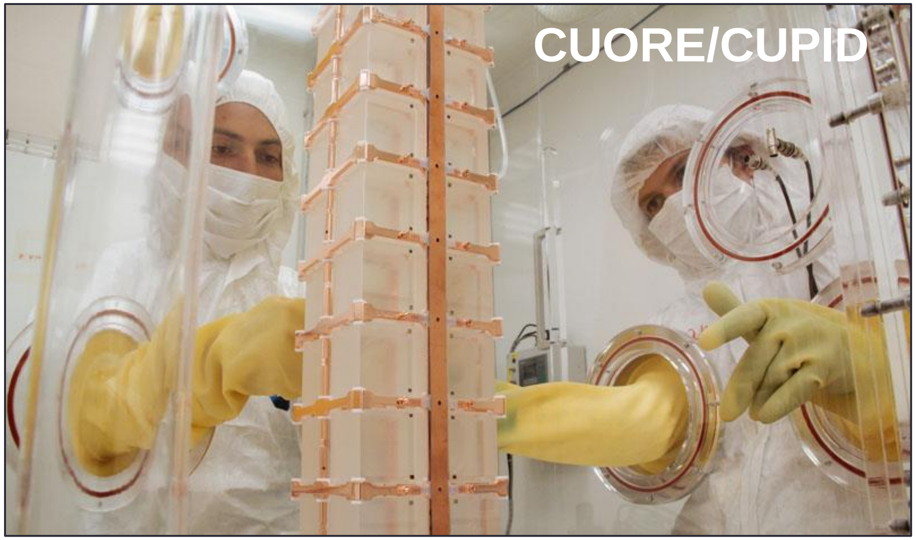
- ***MANDATORY:***
  - Resolution better than ~2% FWHM to fully reject two-neutrino mode
- ***Ideally:***
  - Would have zero other backgrounds
- ***Then just watch and wait…***


With experiments that are not perfect (assuming xenon):
Band shows most pessimistic point in IH, given the range of nuclear matrix elements.
Inverted Heirachy
(whatever that means)
Plot from Jason Detweiler
Background free
0.1 counts/ROI/t/y
1.0 count/ROI/t/y
10 counts/ROI/t/y
$T_{1/2}$ 3σ DL [years]
Exposure [ton-years]
10^{24}
10^{25}
10^{26}
10^{27}
10^{28}
10^{29}
10^{30}
10^{-3}
10^{-2}
10^{-1}
1
10
10^{2}
10^{3}

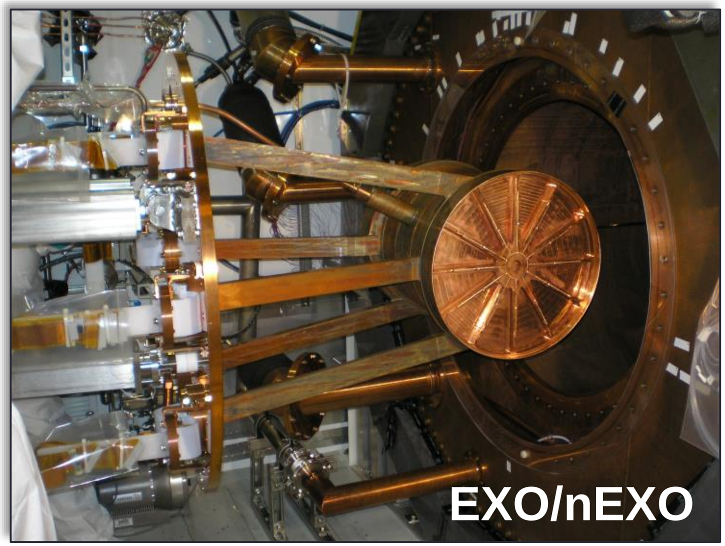


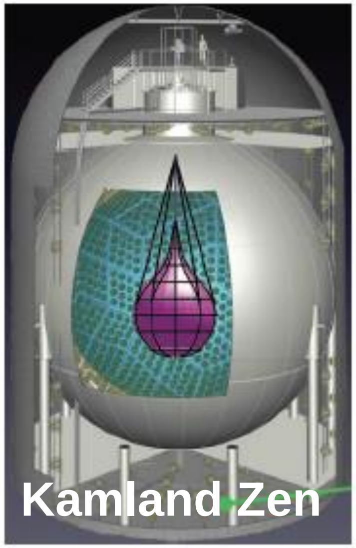
Ultimately, a ton-scale experiment with excellent background rejection (<1 ct ton yr) is the key to crossing the IH in finite time.
Inverted Heirachy
Run for 2-3 years
Run for 15 years
Run for 100 years
Run for 1000 years
1000 kg experiment
Background free
0.1 counts/ROI/t/y
1.0 count/ROI/t/y
10 counts/ROI/t/y
$T_{1/2}$ 3σ DL [years]
Exposure [ton-years]
10^{24}
10^{25}
10^{26}
10^{27}
10^{28}
10^{29}
10^{30}
10^{-3}
10^{-2}
10^{-1}
1
10
10^{2}
10^{3}

Neutrinoless Double Beta Decay Searches

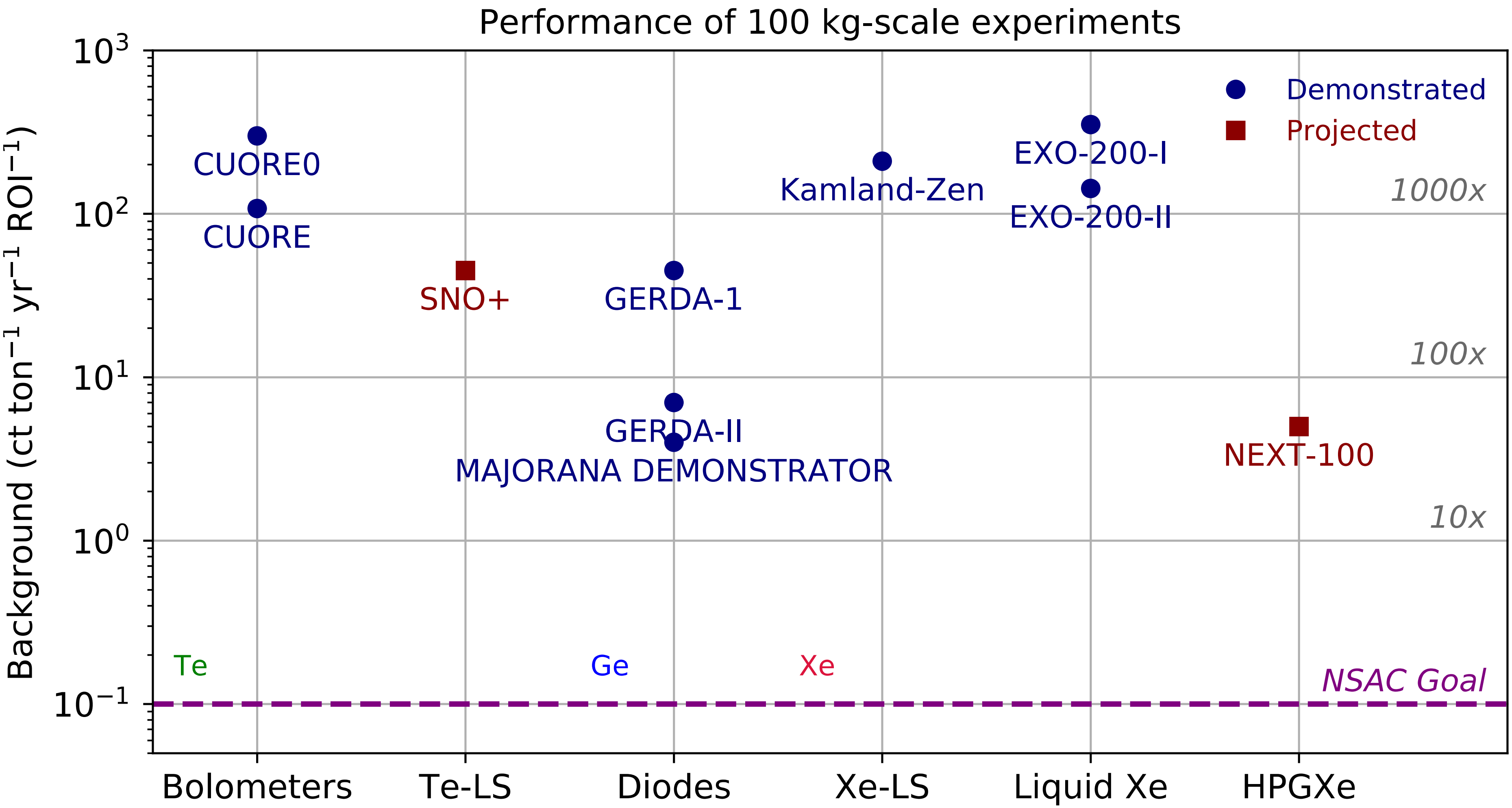


Many completely different technologies…



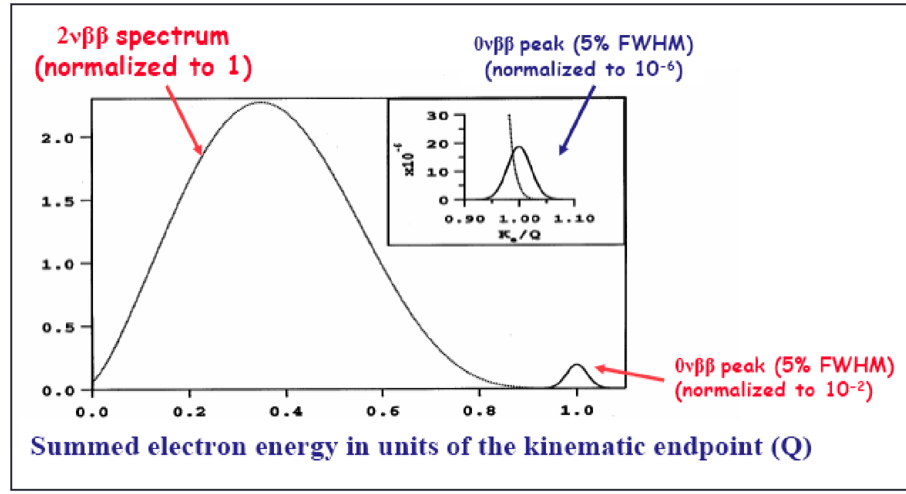


100kg-class experiments:



Methods to achieve lower background

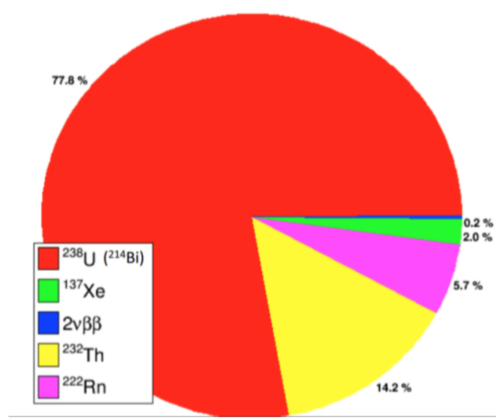
- Reject 2νββ tail:
  - Energy resolution is only handle
    - ~2% FWHM reaches 0.1 ct ton yr.

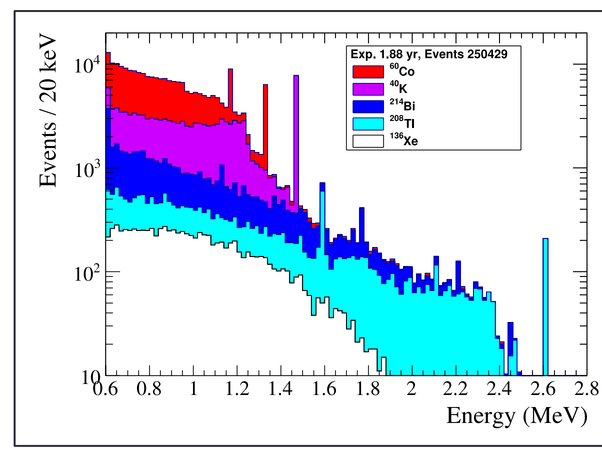


For cartoon purposes only

- Reject radioactive backgrounds

  Typically ^{214}Bi, ^{208}Th from uranium and thorium chains respectively:
  - Shielding / radiopurity
  - Energy resolution
  - Topological signature
  - Daughter tag





nEXO (left) and NEXT-100 (right) background models

The NEXT Program

- **Sequence of HPGXe TPCs, focused on achieving big, very low background xenon 0νββ detector**

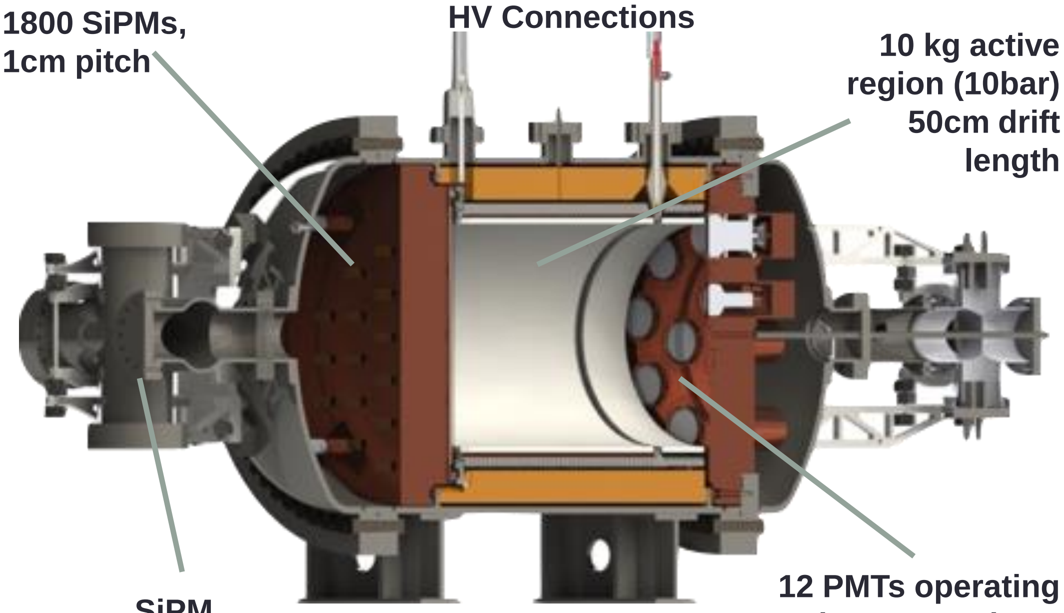
→ NEXT-DBDM
(Berkeley, US)

→ NEXT-DEMO
(Valencia, Spain)

→ **NEXT-White**
(Canfranc, Spain)

→ NEXT-100
(Canfranc, Spain)

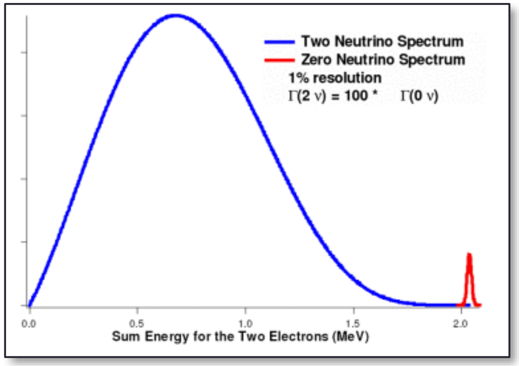
→ NEXT-XXX
???

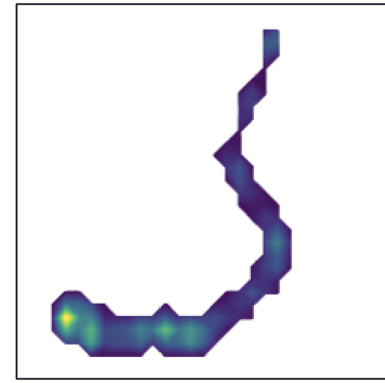


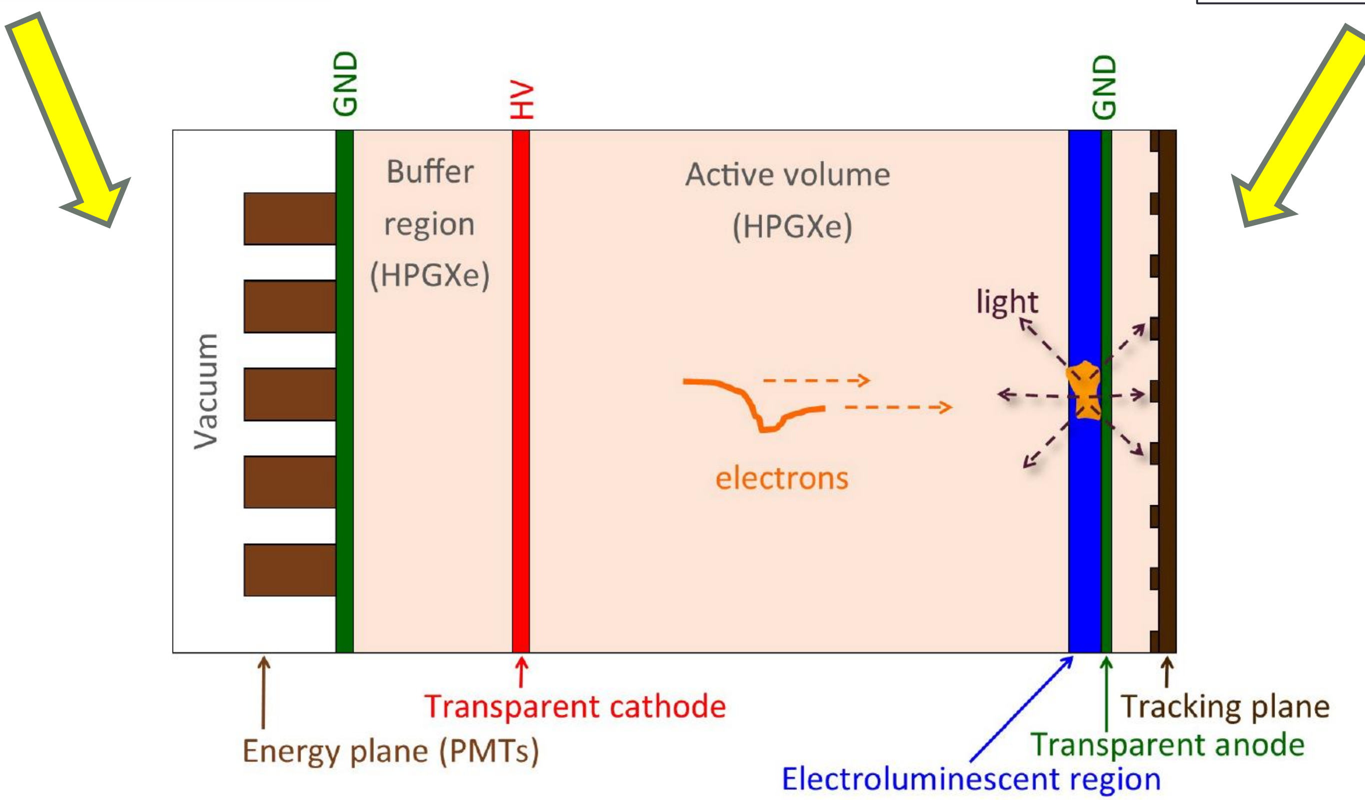
NEXT-White operating now
Full underground technology demonstrator @10kg scale

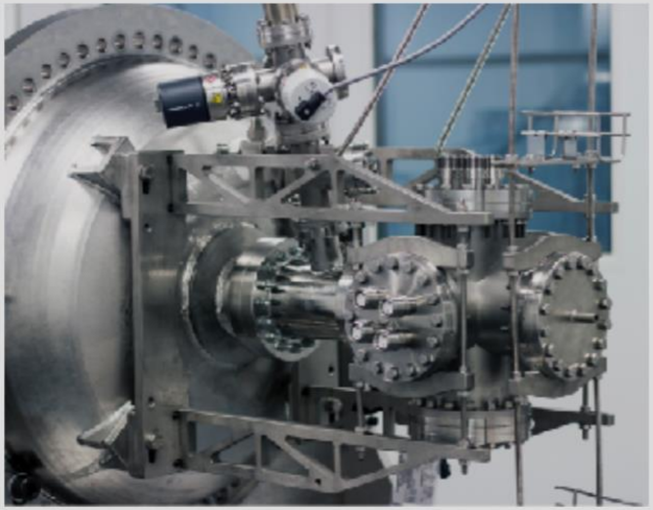
Measure energy this end

Measure topology this end

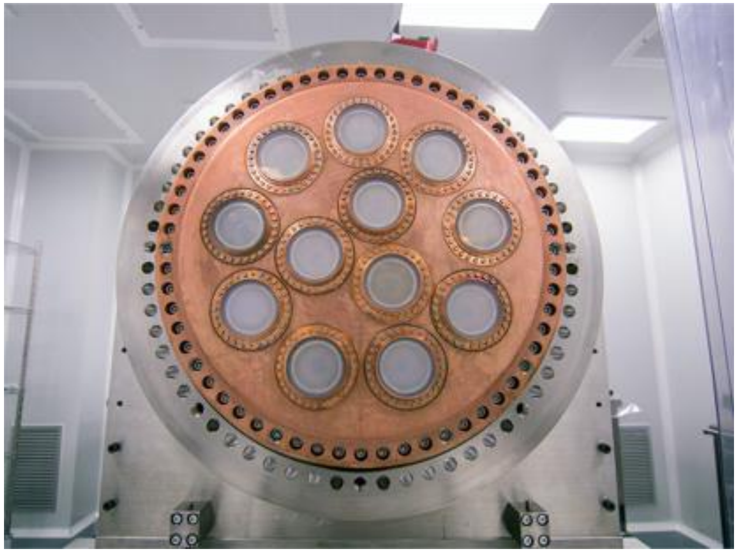






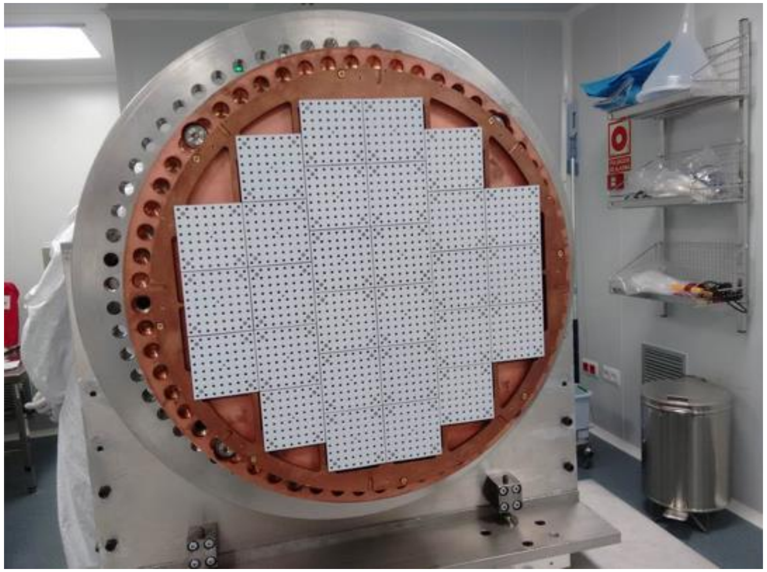


Energy Plane



Field Cage

Tracking Plane



Anode

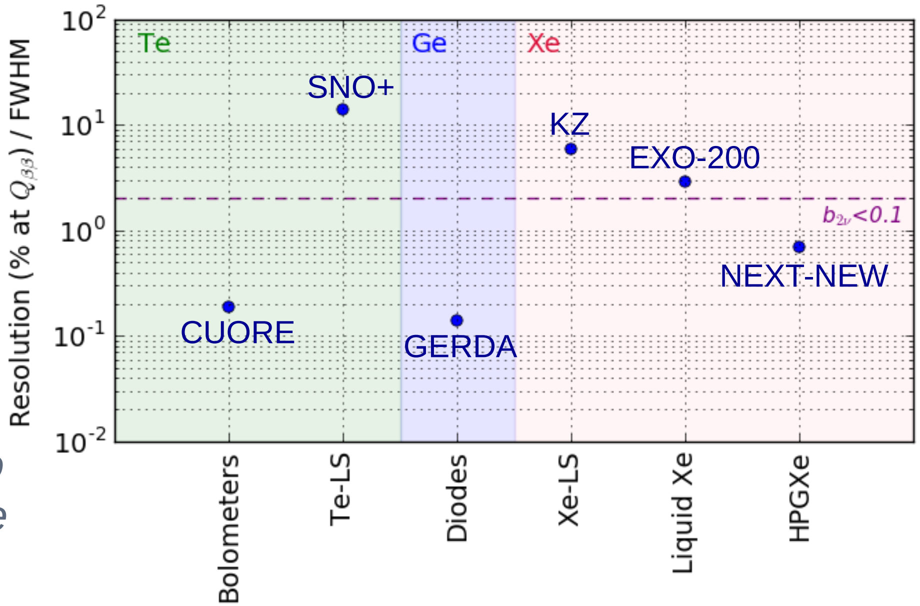
The Next White (NEW) Detector

F. Monrabal et. al. arXiv:1804.02409

Why Xenon Gas?

- 1) Energy resolution
- 2) Topology
- 3) Practicality at ton-scale

Why Xenon Gas?

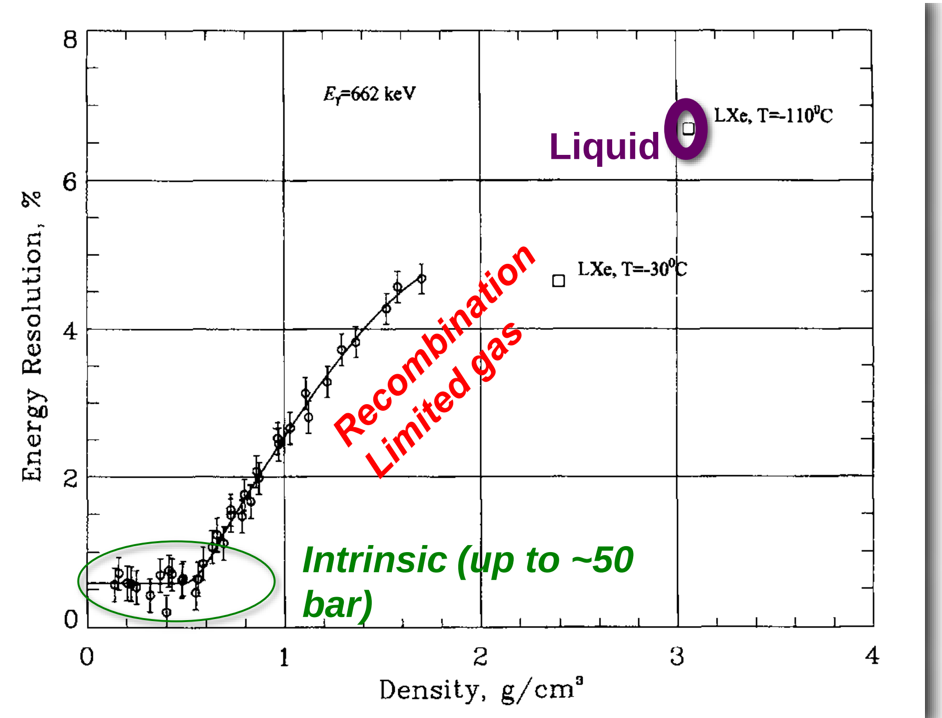


- **1) Energy resolution**

  Fluctuation-less EL gain and low Fano factor produces resolution comparable with solid-state technologies in a monolithic TPC experiment

- 2) Topology

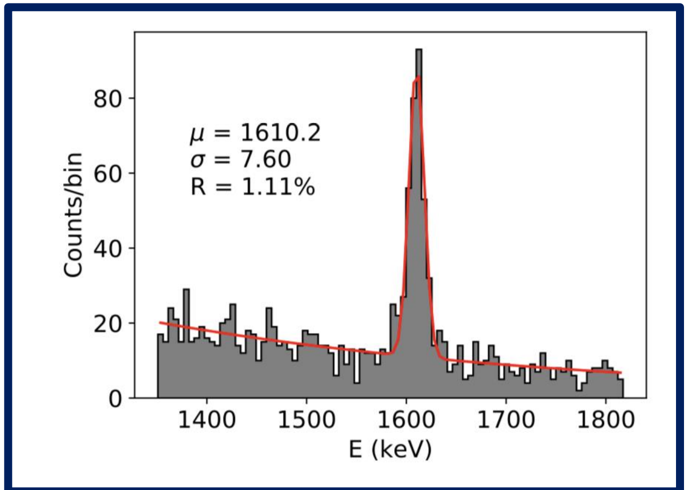
- 3) Practicality at ton-scale

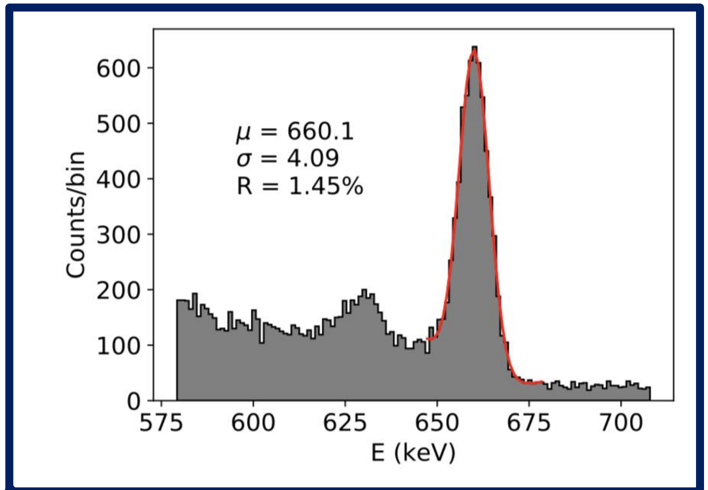


Bolotnikov and Ramsey. "The spectroscopic properties of high-pressure xenon."NIM *A* 396.3 (1997): 360-370

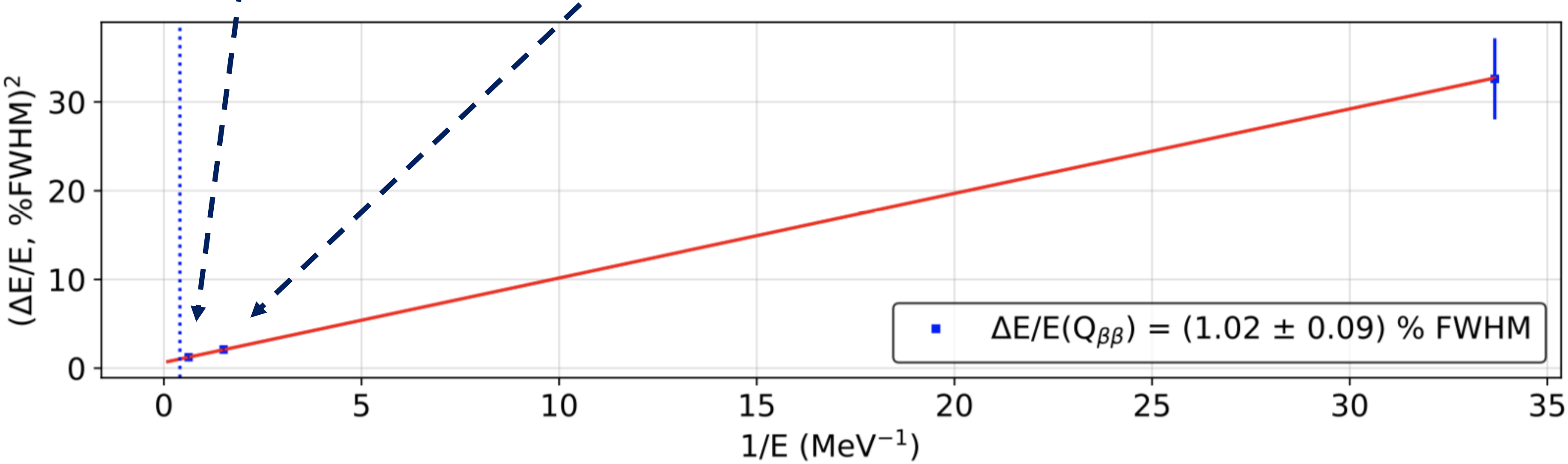
Initial results on energy resolution of the NEXT-White detector

J. Renner et al, arXiv 1808:01804





Energy calibrations and stability still improving: presently sit at ~1% at Q_{bb}



Why Xenon Gas?

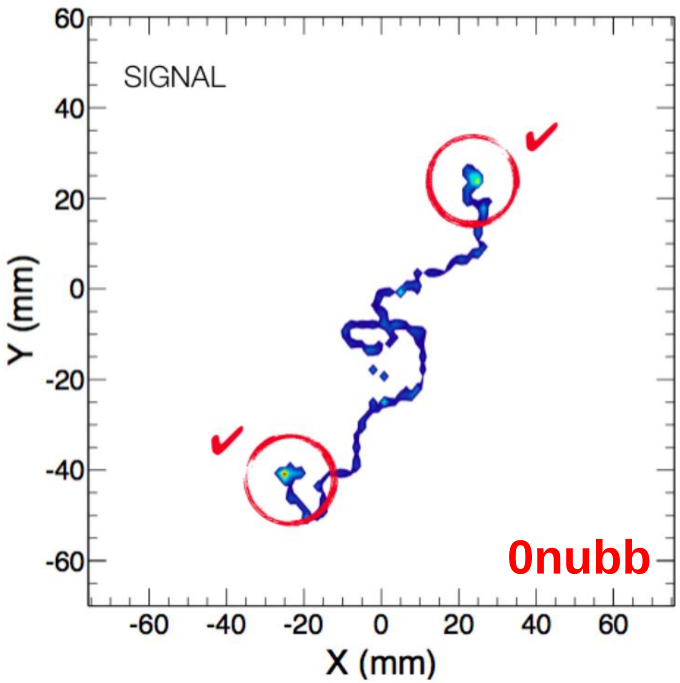
- 1) Energy resolution

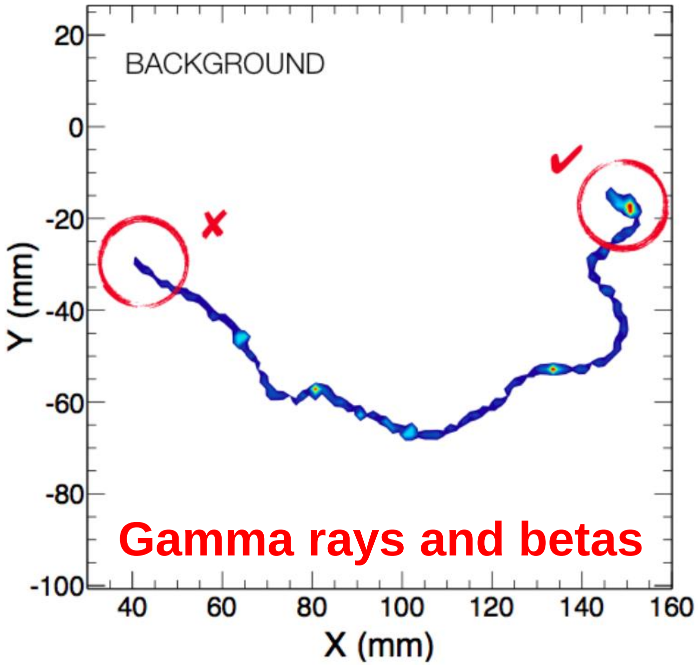
  Fluctuation-less EL gain produces resolution comparable with solid-state technologies in a monolithic TPC experiment

- **2) Topology**

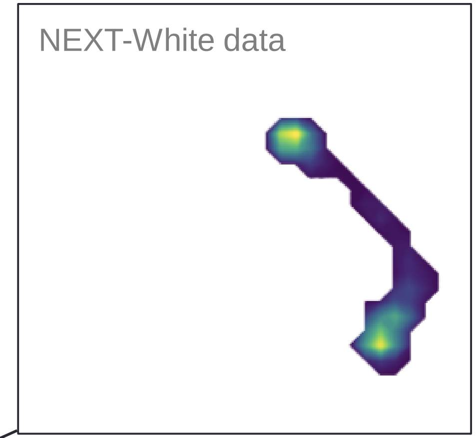
  Lower density allows powerful single-vs-multi electron and single-vs-multi-site topological background rejection

- 3) Practicality at ton-scale

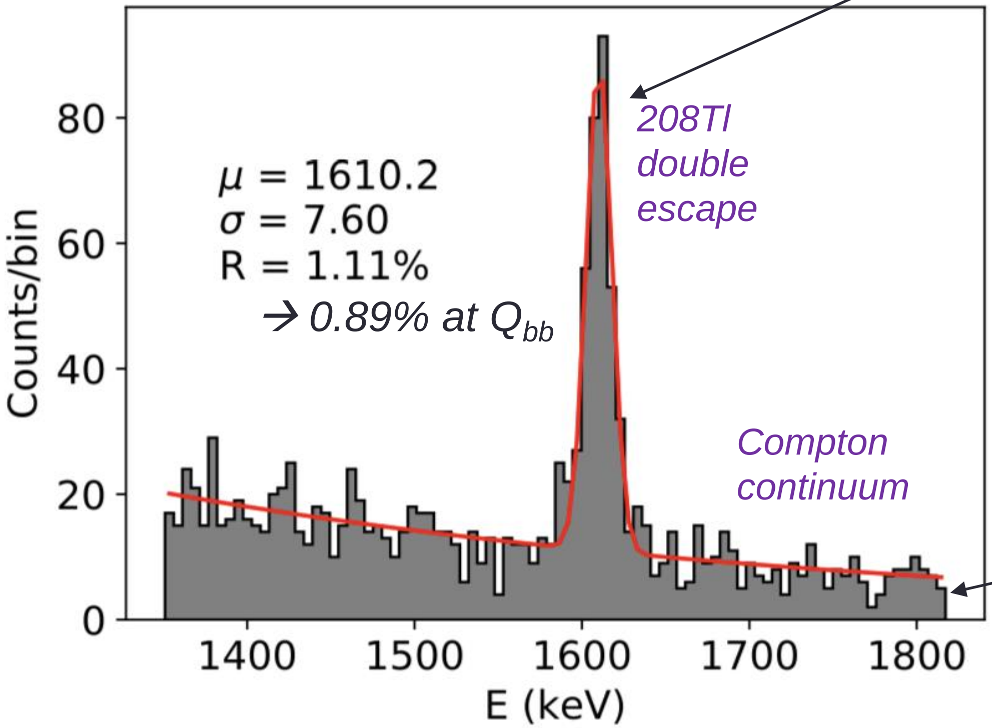


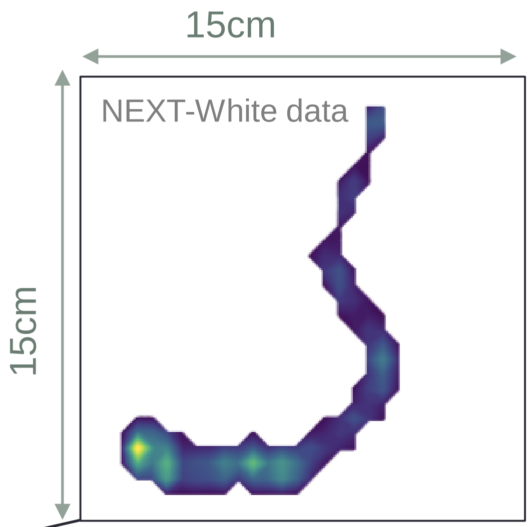


Topological Reco with Double Escape Peaks



On double escape peak

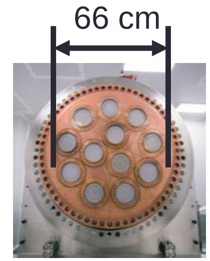




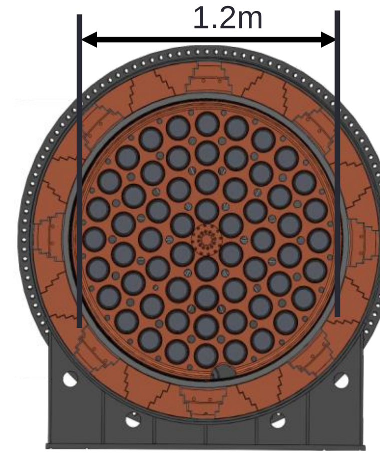
Off double escape peak

Why Xenon Gas?

NEXT-NEW
Running



NEXT-100
2019



- 1) Energy resolution

  Fluctuation-less EL gain produces resolution comparable with solid-state technologies in a monolithic TPC experiment

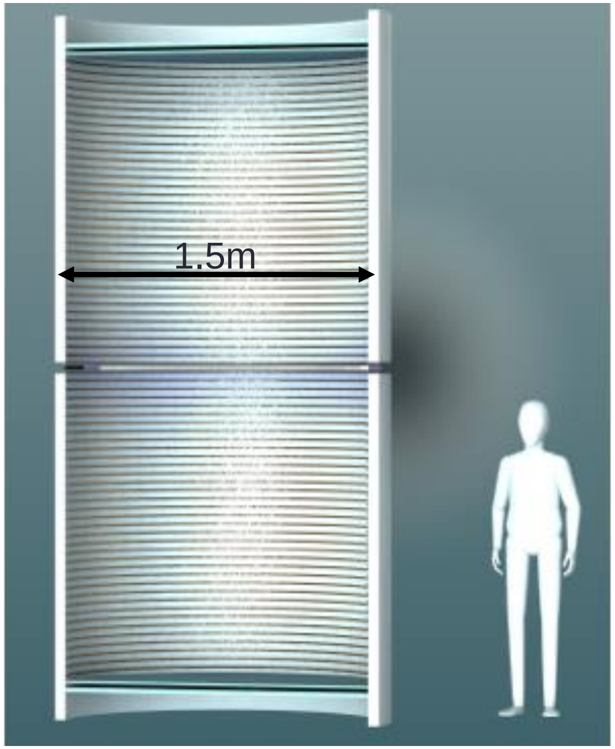
- 2) Topology

  Lower density allows powerful single-vs-multi electron and single-vs-multi-site topological background rejection

- **3) Practicality at ton-scale**

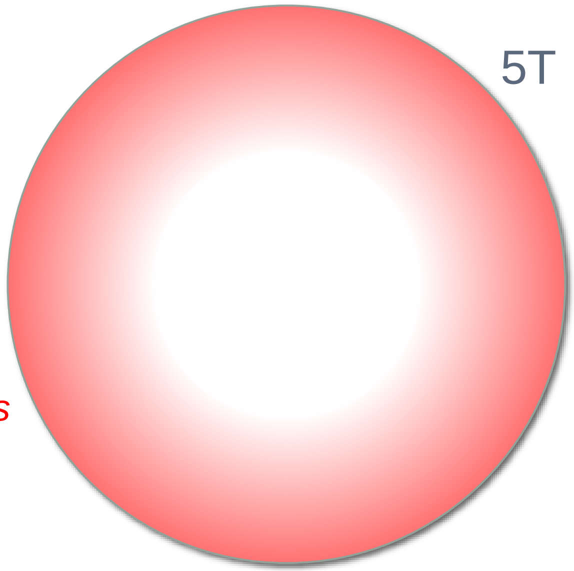
  Reliance on active background rejection rather than self shielding means program uses isotope efficiently and can be phased

NEXT-XXX: *202?*



Self-shielding based
(KamlandZen, nEXO, etc)

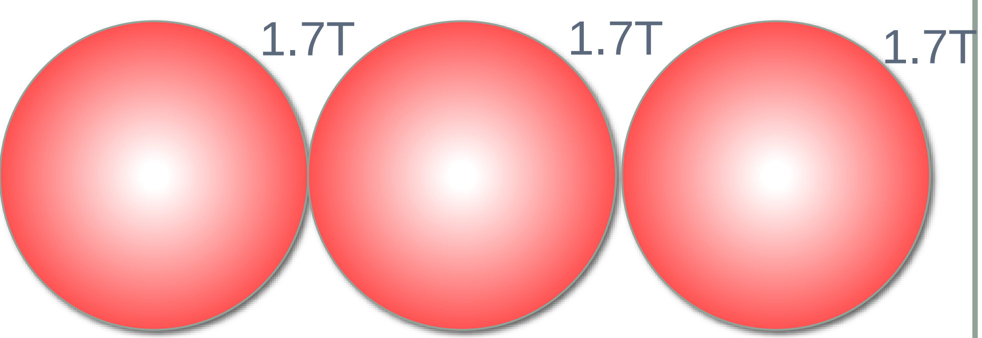
Several tons of isotope for each clean ton



External gamma ray backgrounds absorbed in outer layers

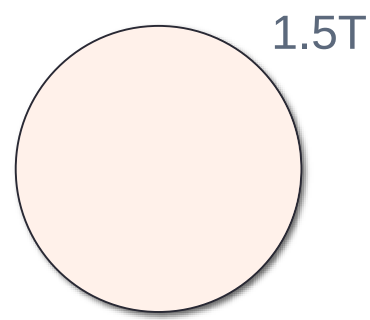
Cannot be phased into smaller modules

have to build it all at once



Active background rejection based
(CUORE, MAJORANA/GERDA, NEXT, etc)

1 ton of isotope for each clean ton.

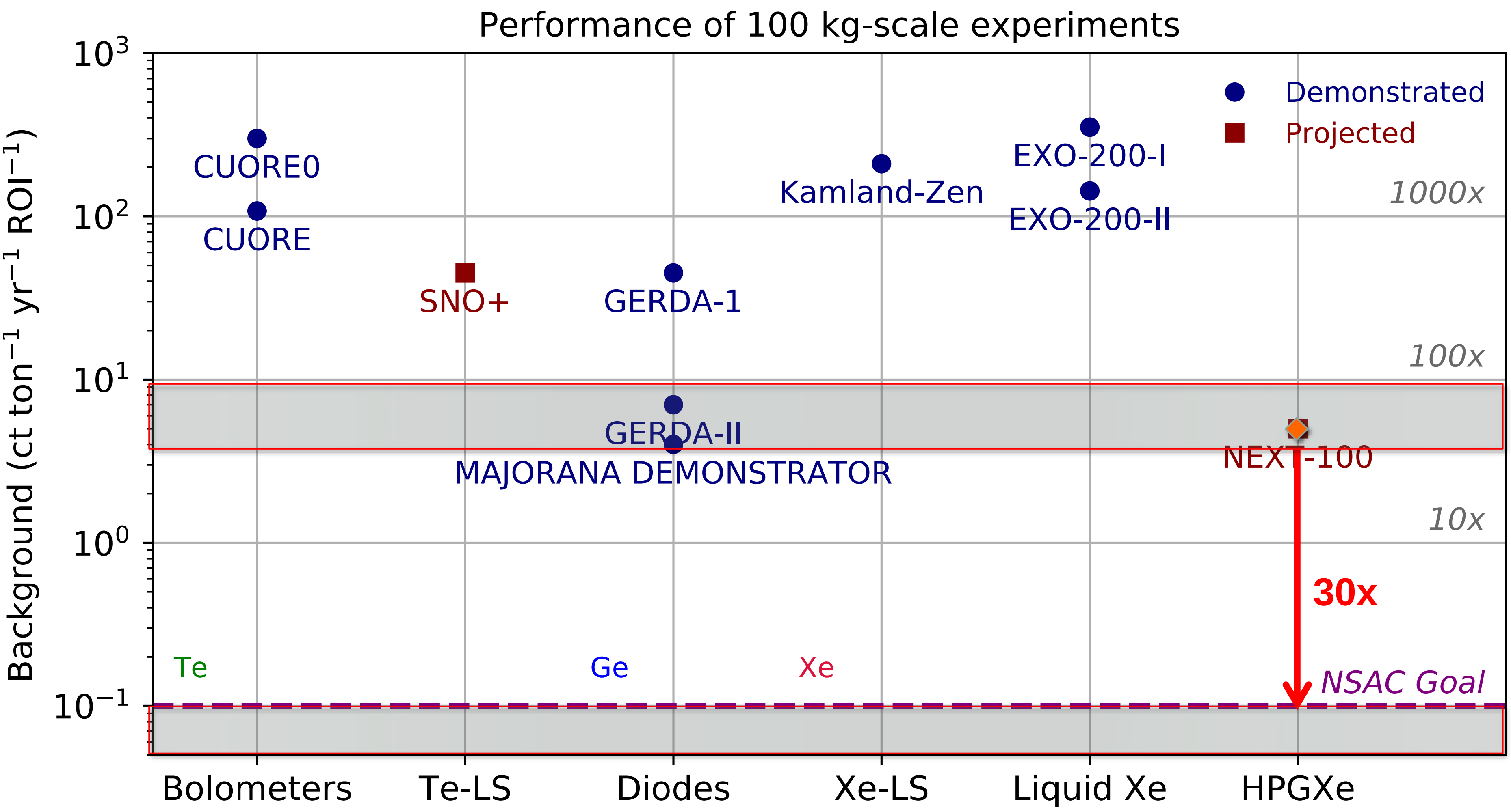


Nearly uniform background throughout

Phase-able deployment possible

Low-risk, reliable, cost effective approach to hitting background targets

NEXT-100 in 2019 will be comfortably world leading in xenon, and competitive with world leaders in other isotopes.

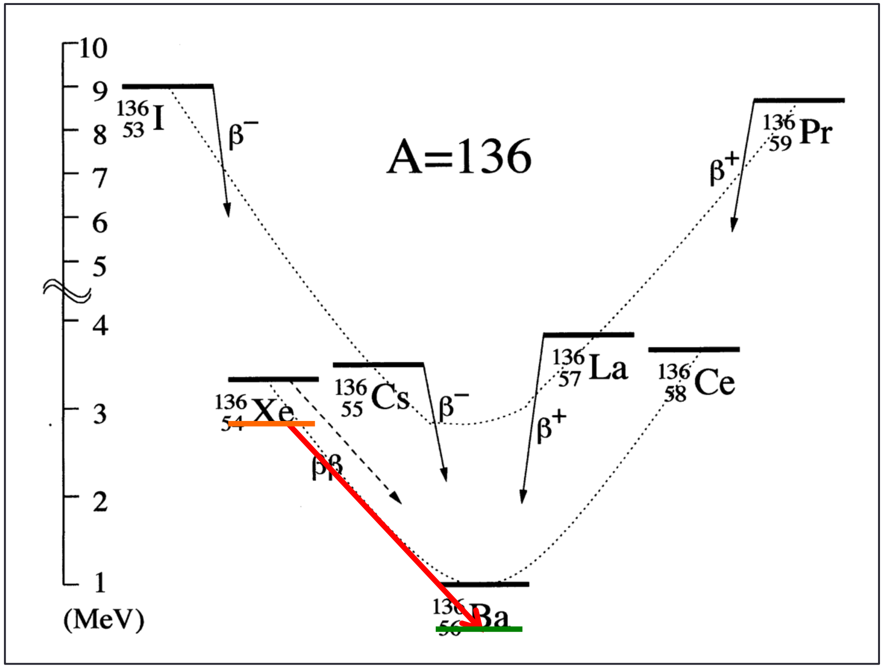


BUT WORLD-LEADING IS NOT ENOUGH!

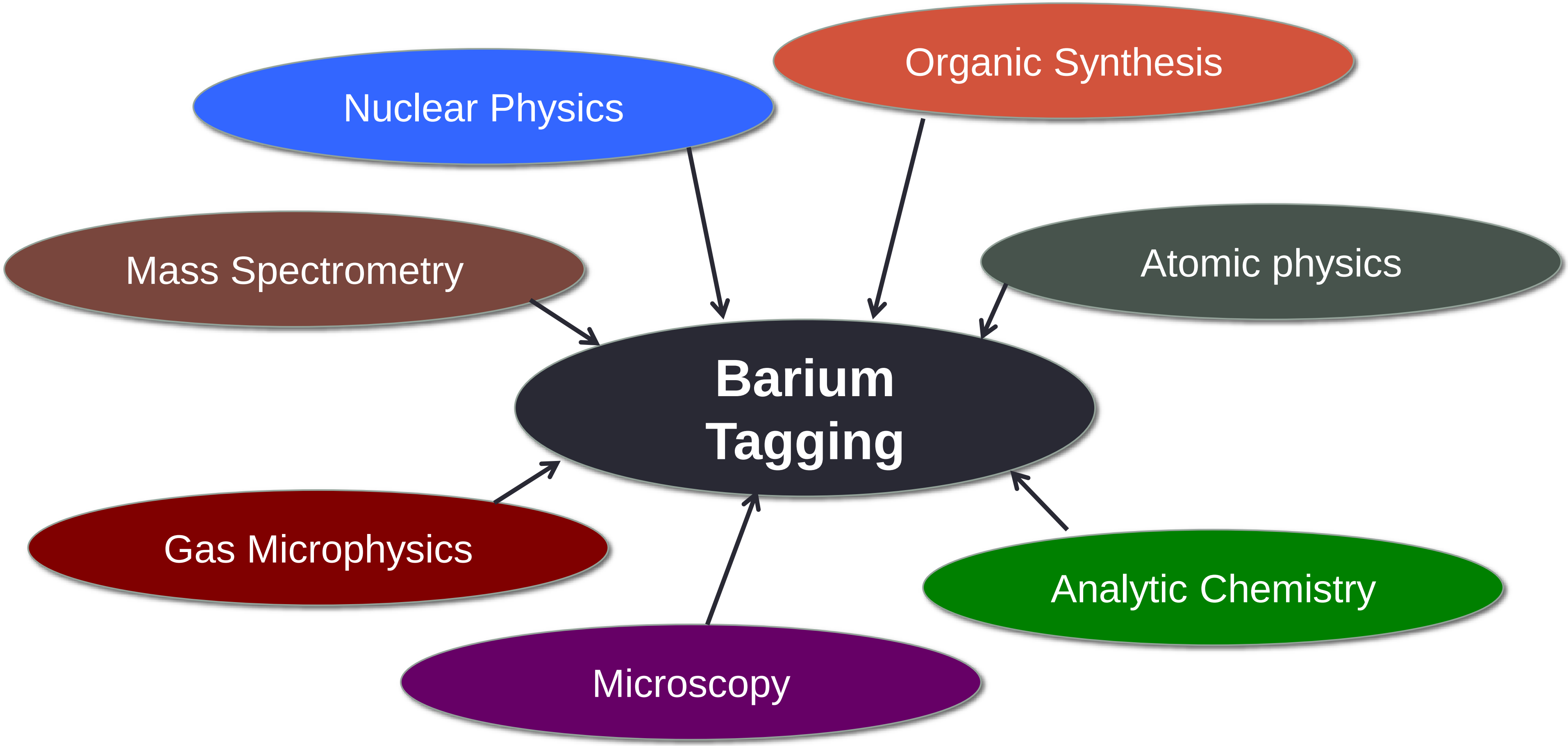
Barium Tagging

- Barium ion is only produced in a true ββ decay, not in any other radioactive event.
- Identification of Ba ion plus ~1% FWHM energy measurement would give a background-free experiment.
- ***Is it plausible to detect an individual barium ion or atom in a ton of material?***

$^{136}Xe \rightarrow {}^{136}Ba + e + e$



Barium tagging for 0nubb has been actively explored in liquid and gas xenon for >15 years, with the holy grail is a scalable single ion sensitive technology.

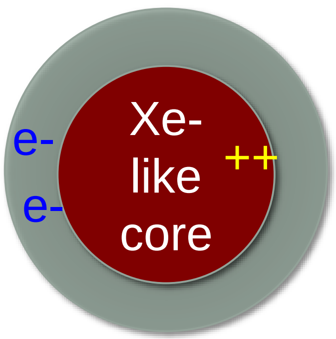

Organic Synthesis
Nuclear Physics
Atomic physics
Mass Spectrometry
Barium Tagging
Gas Microphysics
Analytic Chemistry
Microscopy

Basics - Barium Atoms and Ions

- Barium is born in a high charge state as emerging beta electrons disrupt the atom
- Quickly captures electrons from xenon to reach the Ba^{++} state
- In gas, it ~stops there. In liquid, further recombination happens.

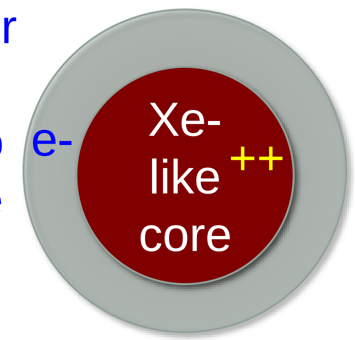
Structure of the **Ba** atom

Two outer electrons bound to +ve core



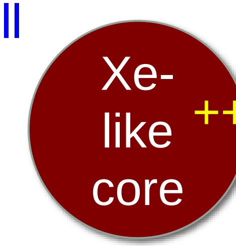
Structure of the **Ba+** ion

One outer electron bound to +ve core



Structure of the **Ba^{++}** ion

Noble-like electron shell



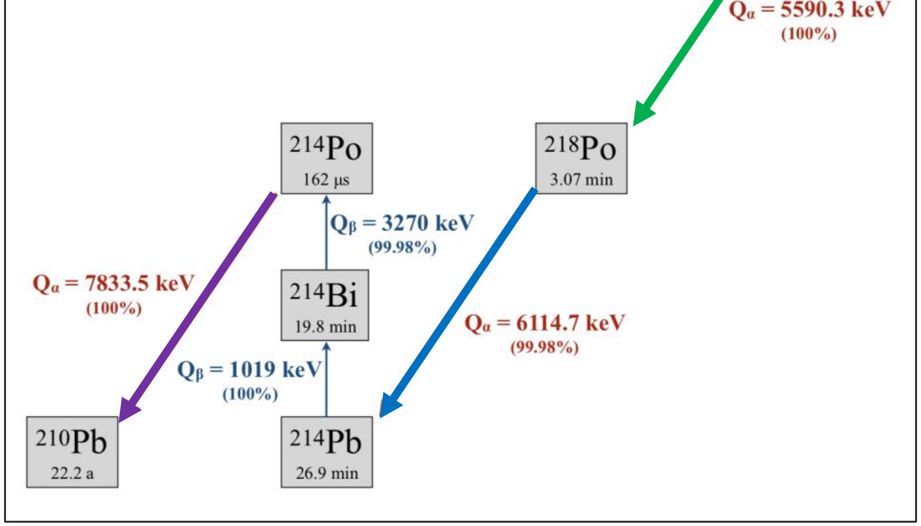
Liquid –
Some distribution

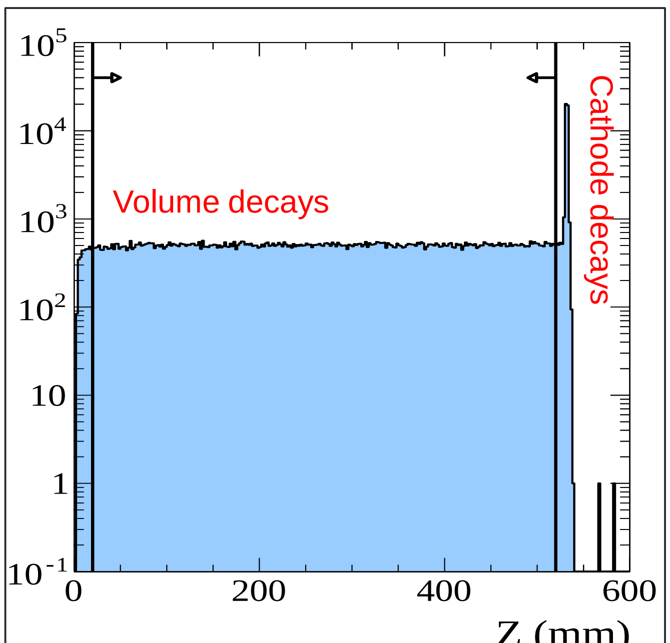
Gas –
Mostly this

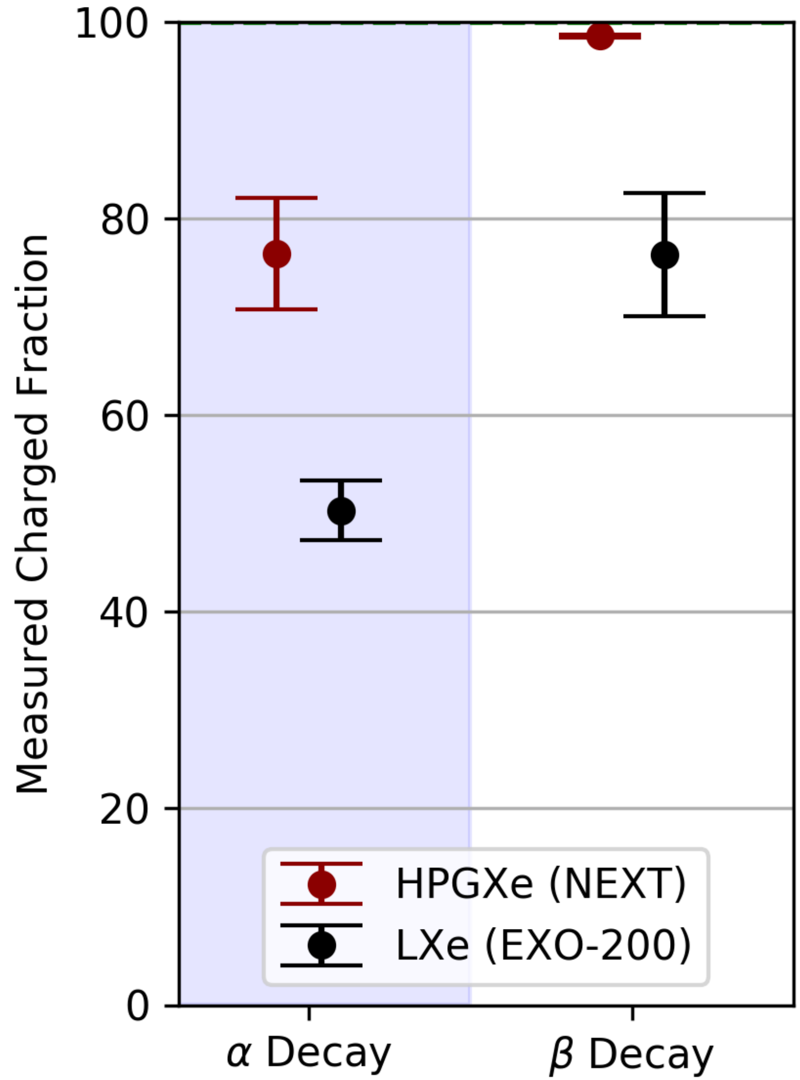
Ionic Charge State in Liquid and Gas

Measurement of radon-induced backgrounds in the NEXT double beta decay experiment

P. Novella et al, arXiv:1804.00471



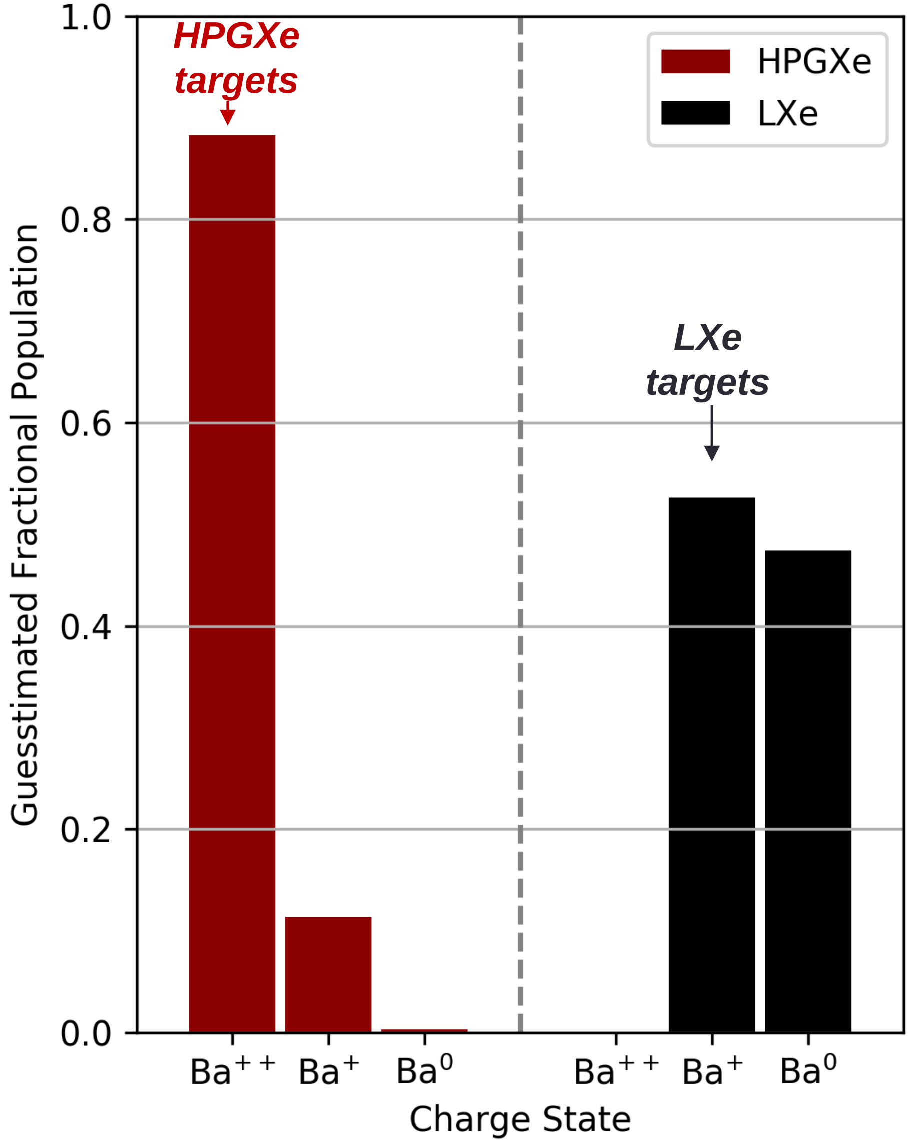




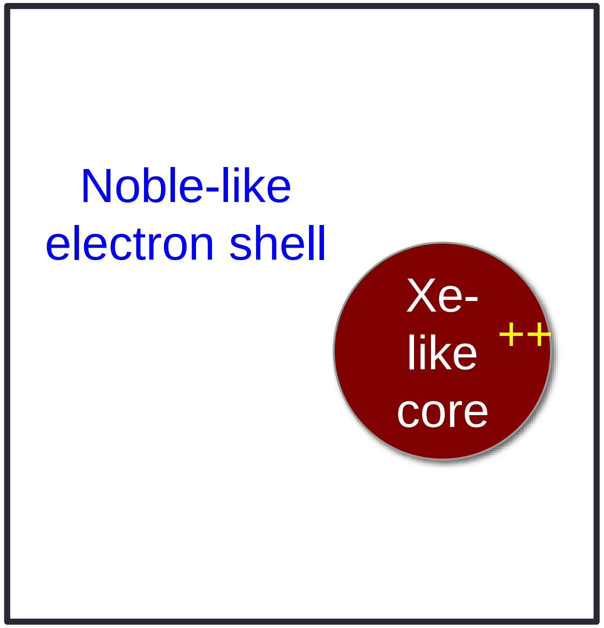
Charge States

- "Best guess" assumptions:
  - Ba++ has 2x larger Onsager radius than Ba+, so expect 4x larger capture cross section for recombination.
  - In double beta decay, two emerging tracks at nucleus, so 2x recombination of single decay.
- Guesstimate Ba charged fractions in double beta decay →

My back-of-envelope extrapolations (your mileage may vary)



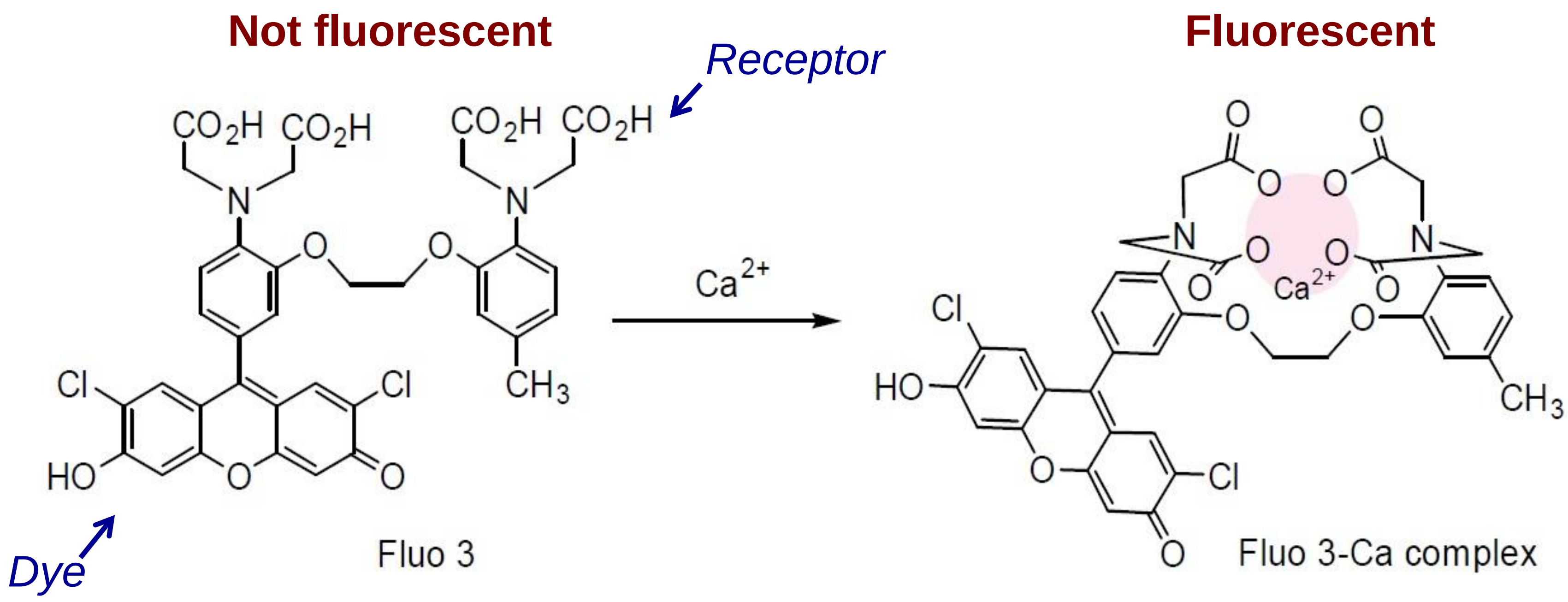
Q: How do you make Ba^{++} shine?



SMFI:

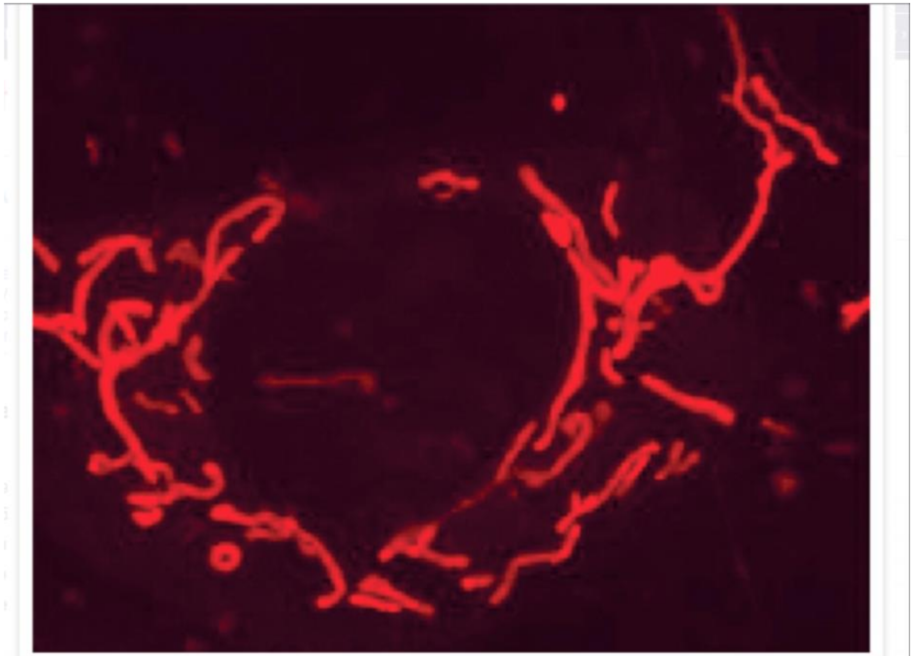
Concept to adapt SMFI for Ba tagging:
***D.R. Nygren*, J.Phys.Conf.Ser. 650 (2015) no.1, 012002**

- A non-fluorescent molecule becomes fluorescent (or vice versa) upon chelation with an incident ion.



Calcium and barium are congeners – many dyes developed for calcium are also expected to respond to barium

SMFI is a technique from biochemistry with demonstrated single-ion resolution.

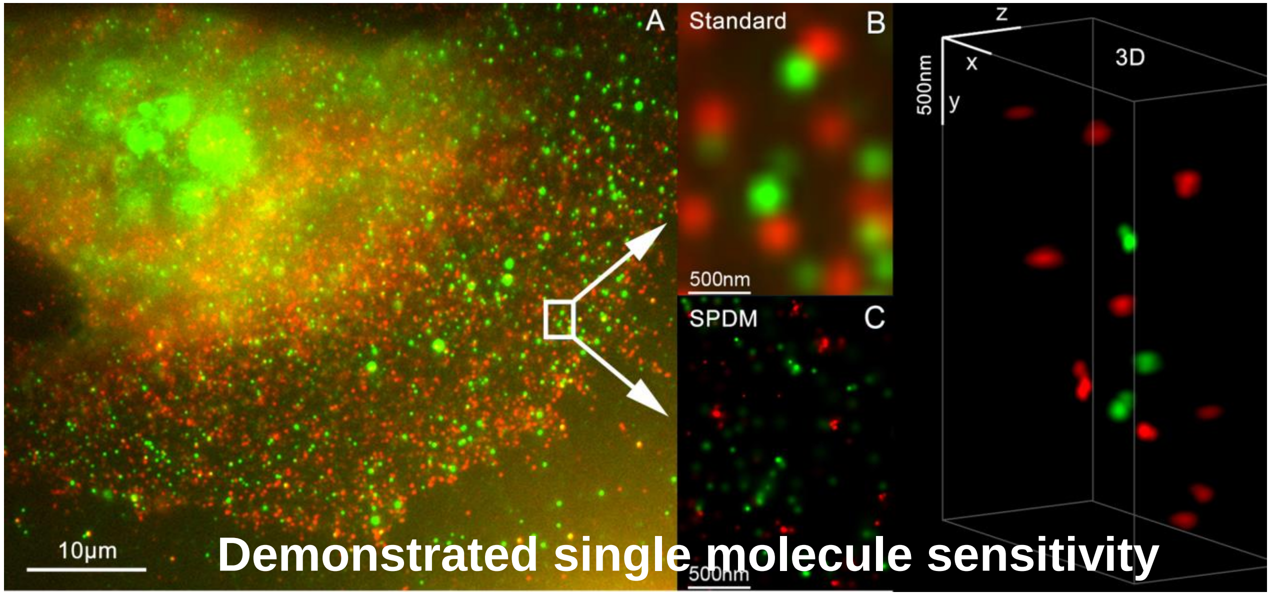


← Rhod-2 sensing Ca^{++} production in rat astrocyte cells

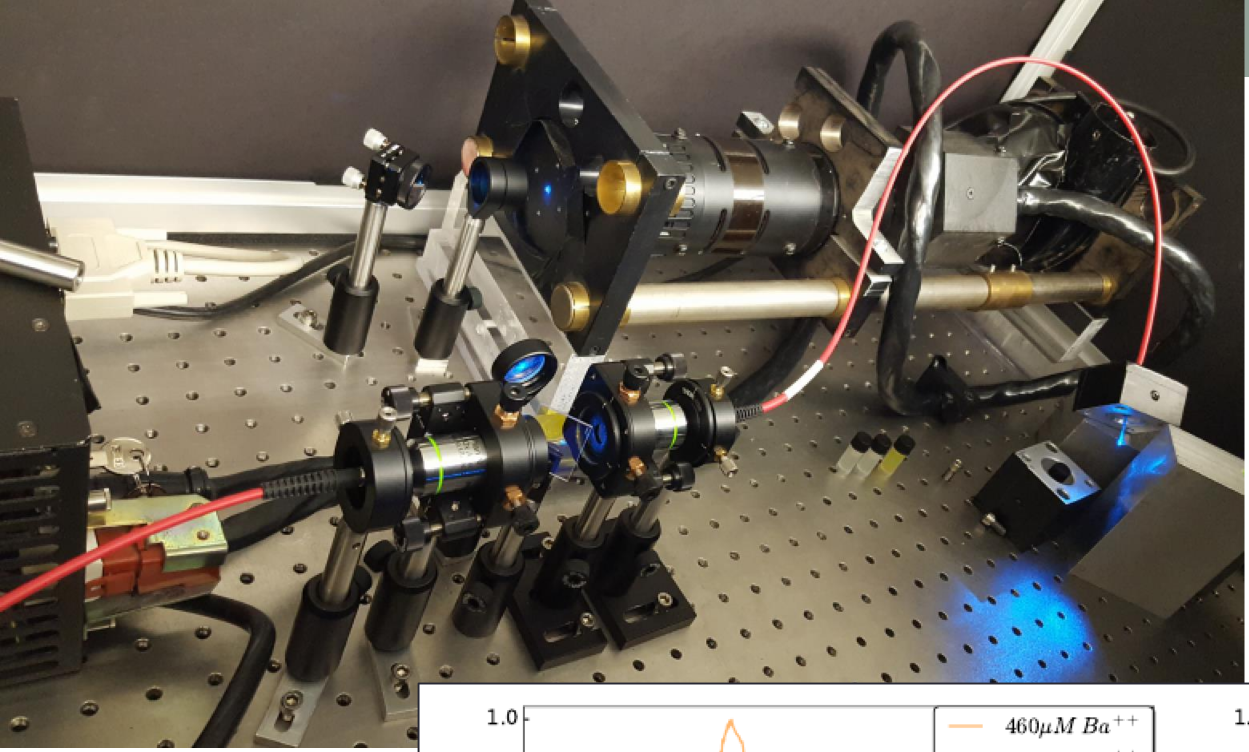
J Microsc. 2011 Apr;242(1):46-54

J Cell Biol 145, 795 (1999).

Single molecule tracking using SMFI is the basis of super-resolution microscopy→



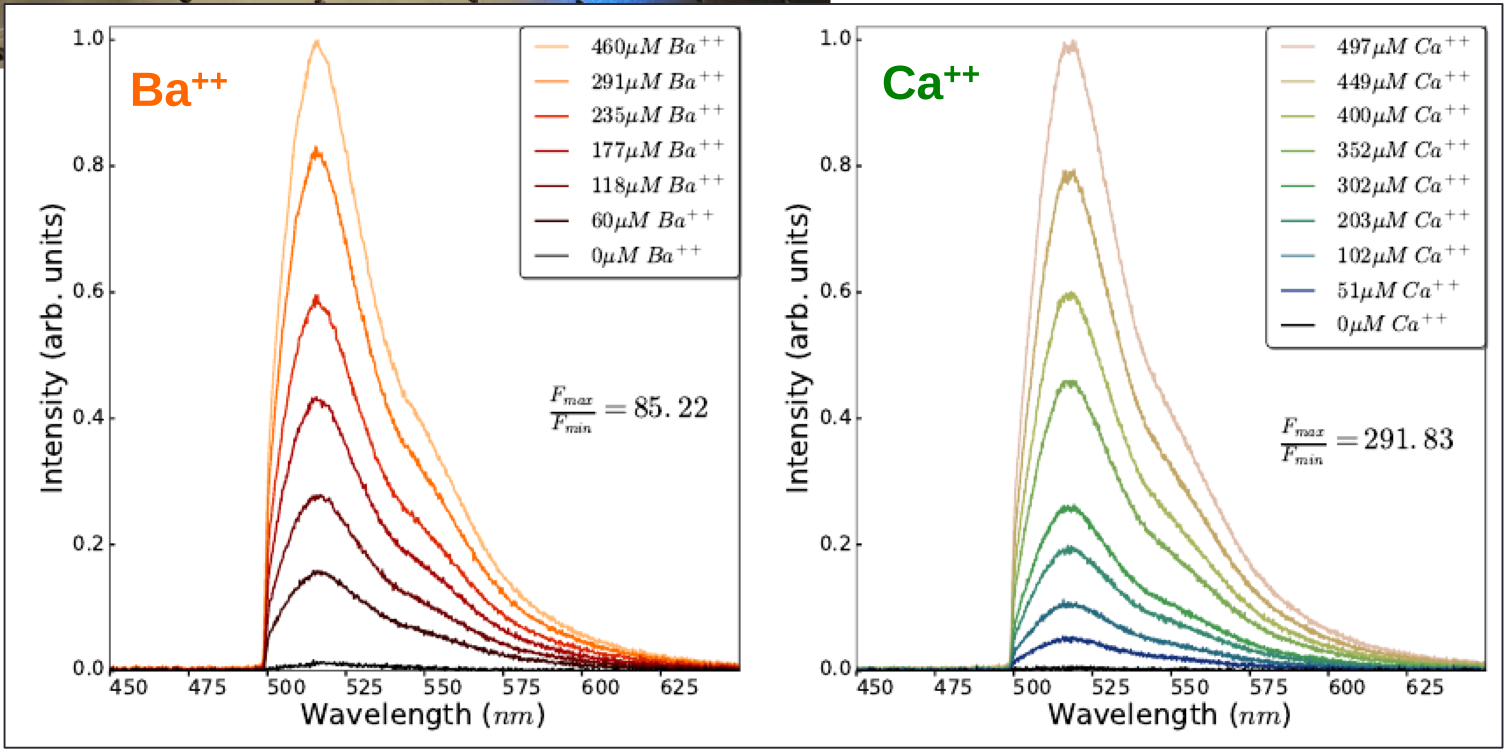
These methods won the Nobel Prize in chemistry in 2014.



← First dabbling - developed a bespoke fluorescence sensor to study barium production at the end of a fiber.

[Single molecule fluorescence imaging as a technique for barium tagging in neutrinoless double beta decay](#)
Jones, McDonald, Nygren, JINST (2016) 11 P12011

We find strong fluorescence from Fluo3 and Fluo4 under chelation with Ba^{++} ions →



Microscope objectives:

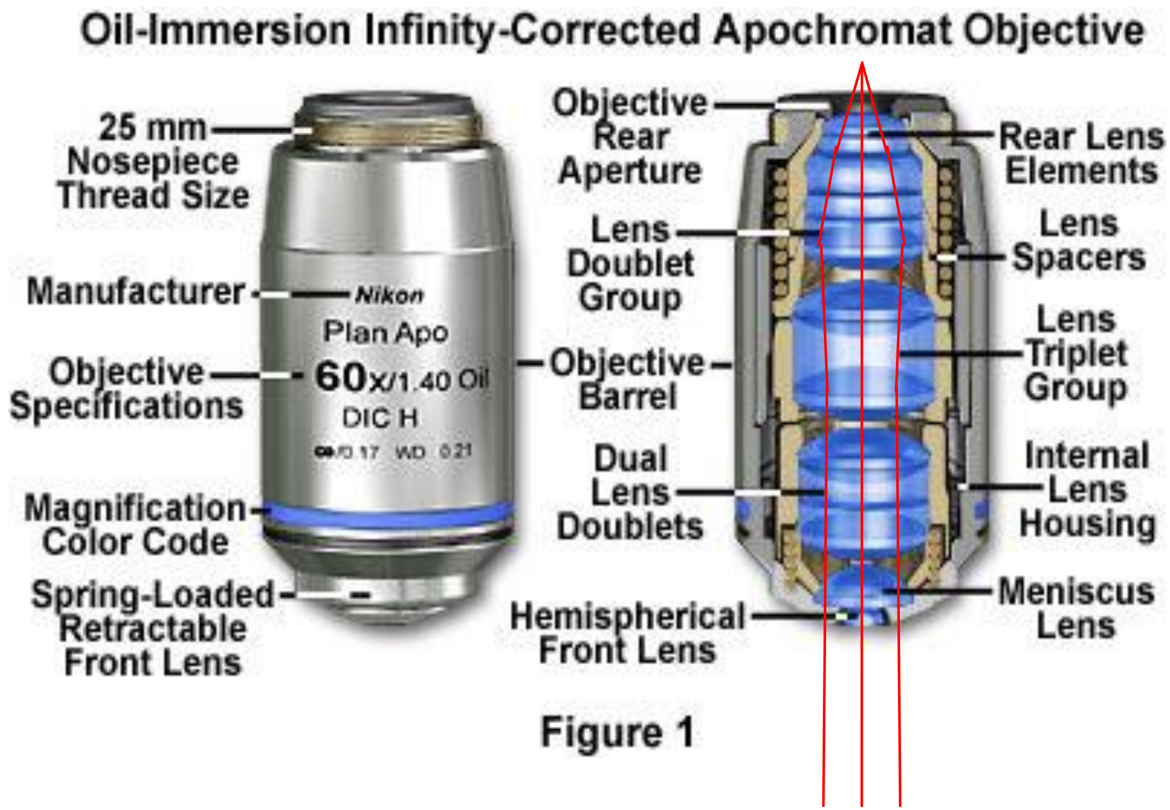
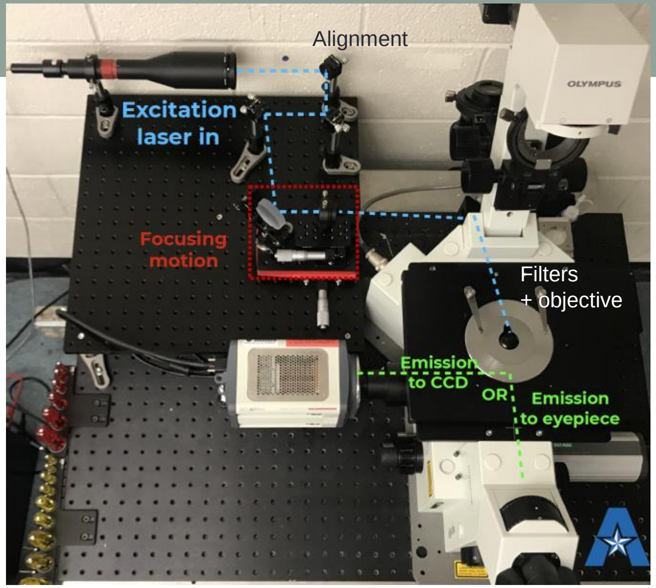
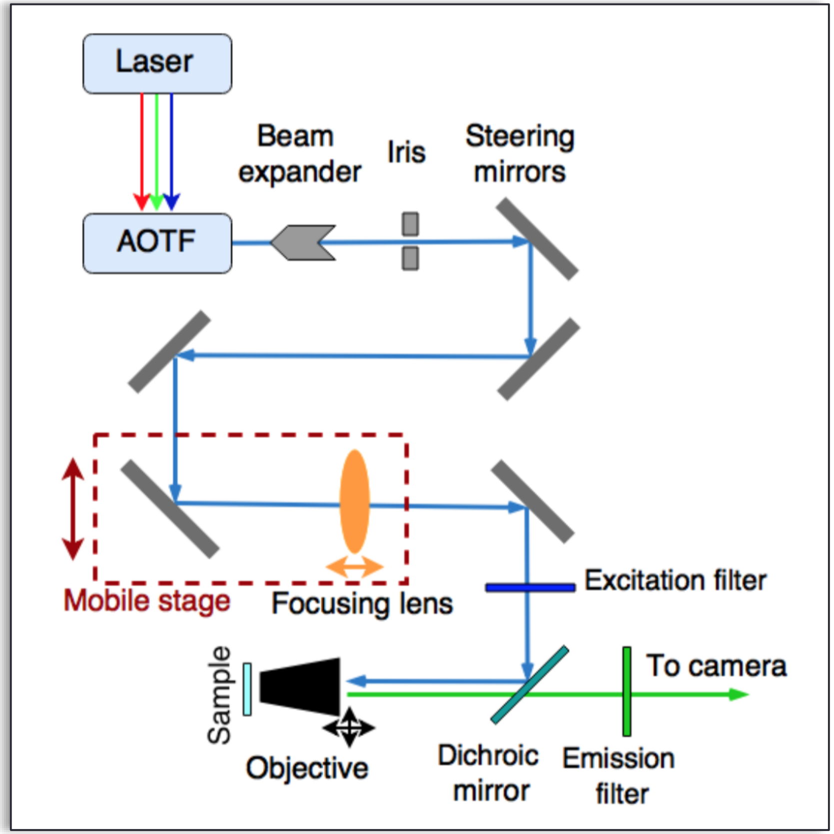
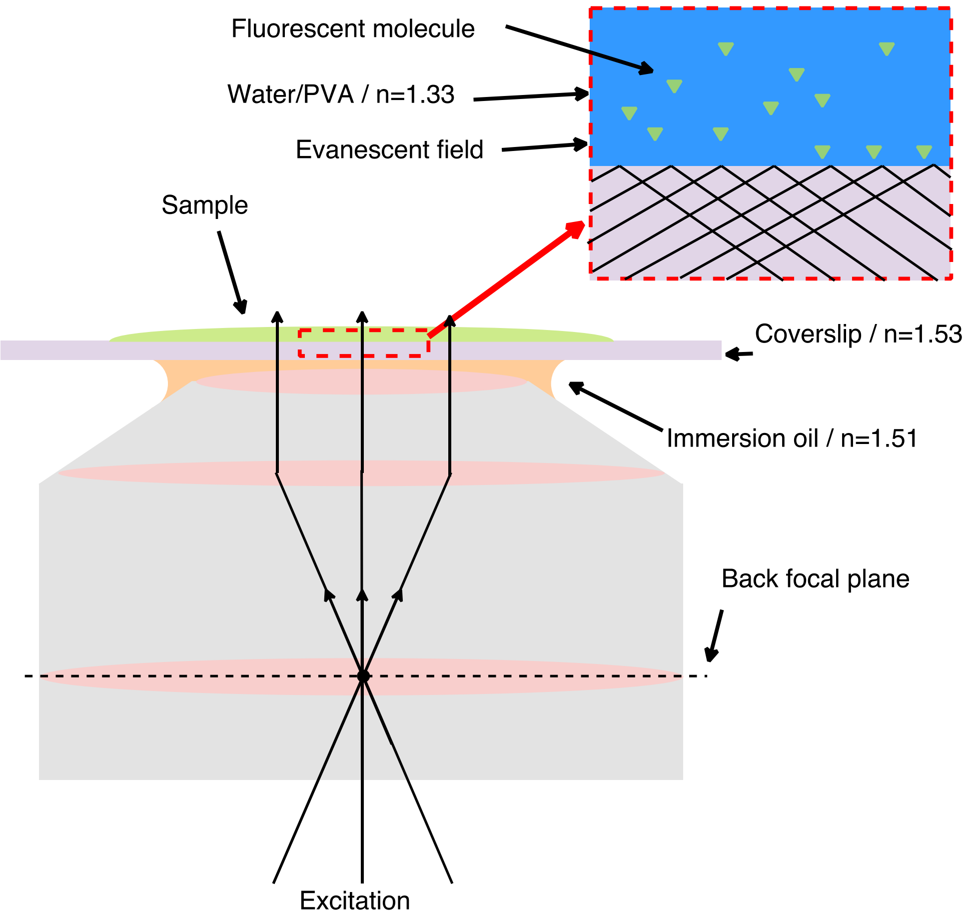


Figure 1

Goal in life: Do its darnedest to turn point sources into parallel rays - and vice versa.

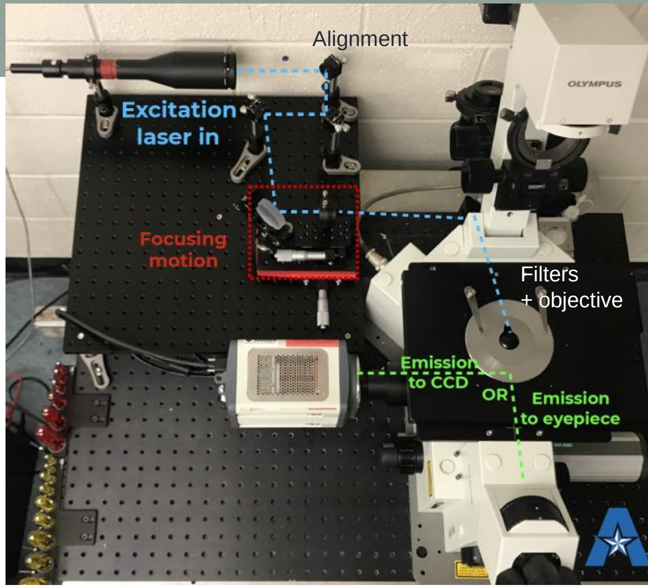


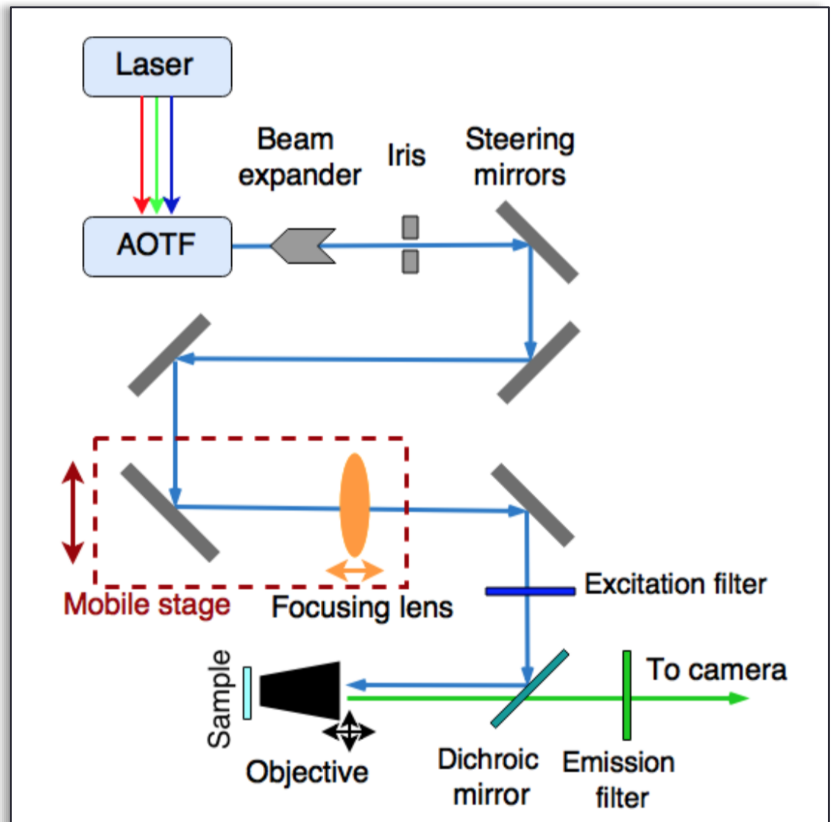


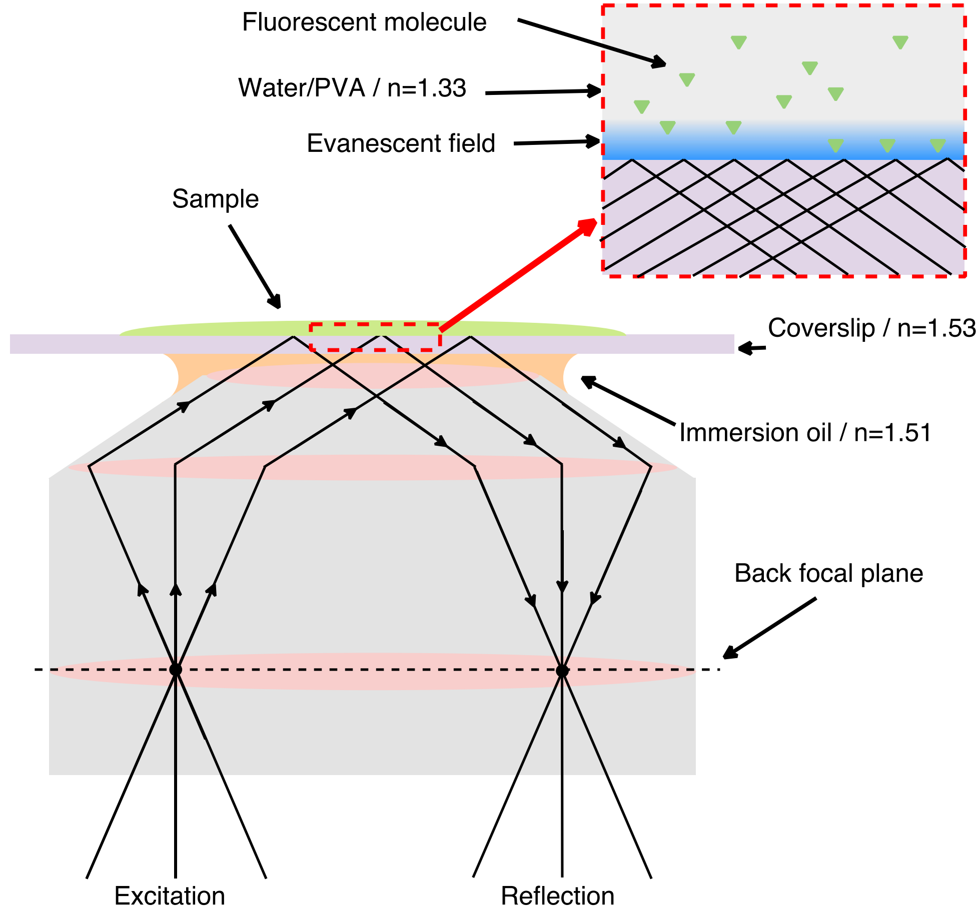


TIRF

Total Internal Reflection Fluorescence microscopy



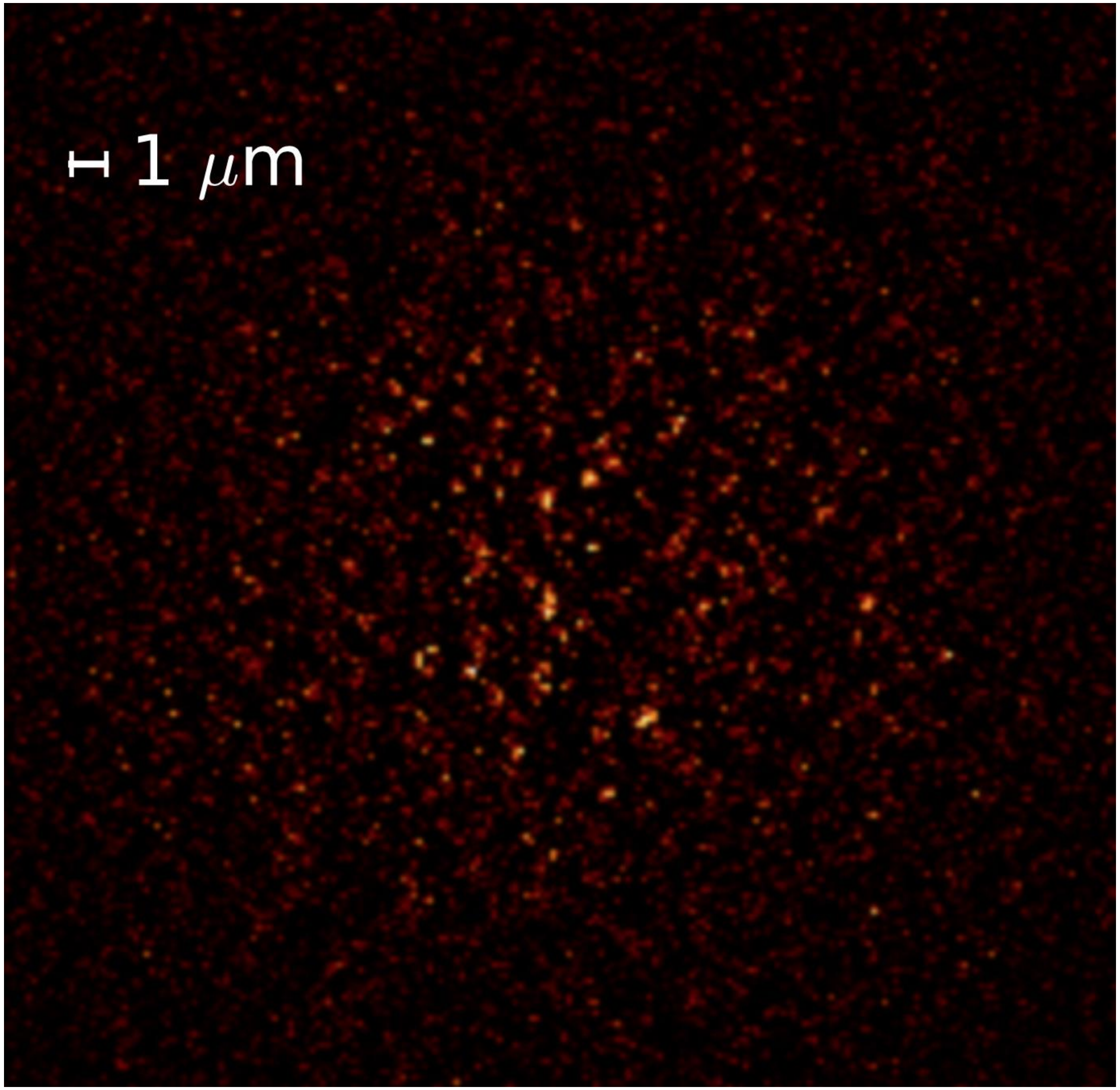




TIRF

Total Internal Reflection Fluorescence microscopy

Single Ba^{++} TIRF images from our lab at UTA

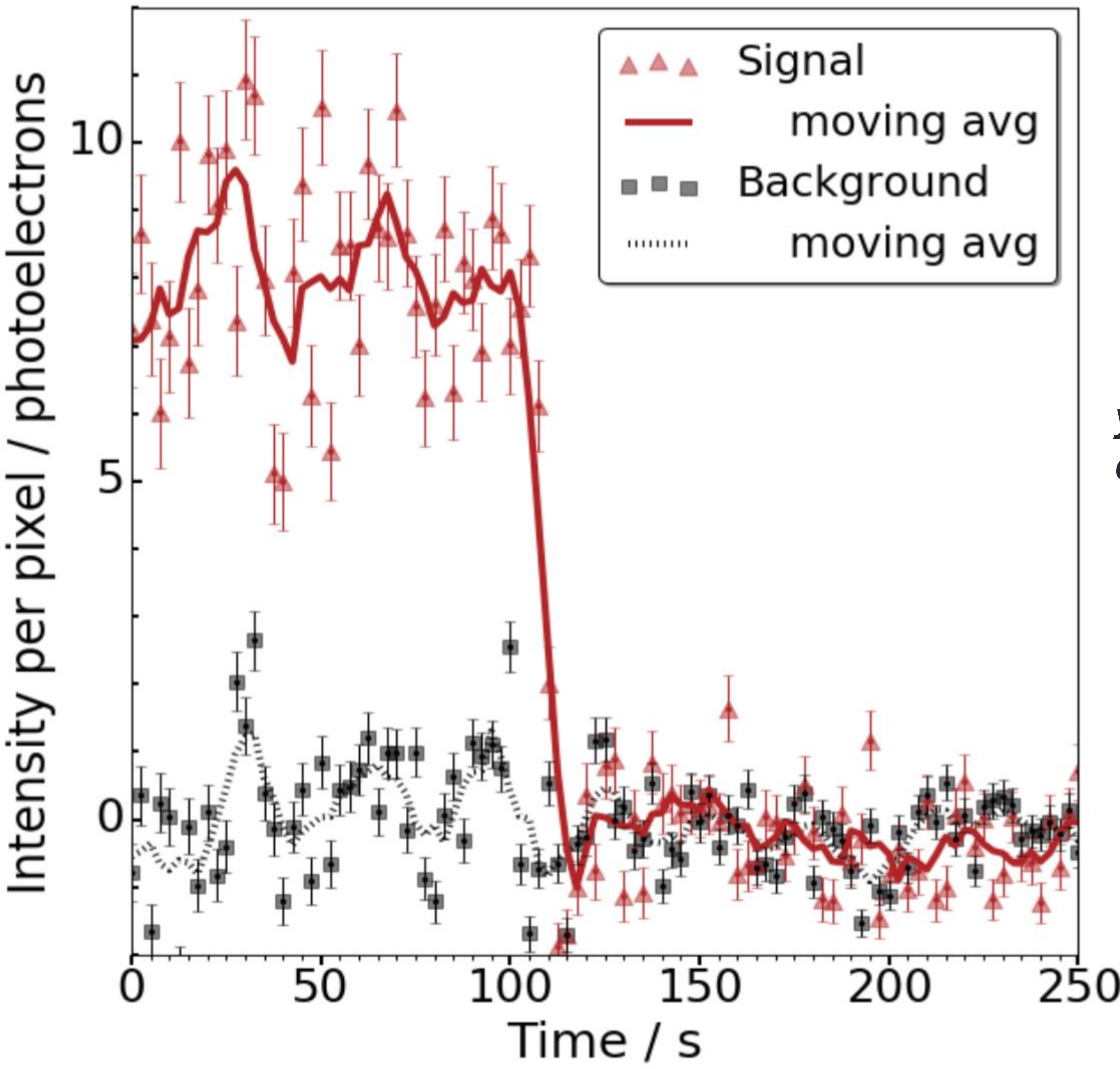


← **This image shows a weak solution of barium perchlorate salt on our sensor.**

Each spot is a **single barium ion.**

Brighter spots are near the TIRF surface, dimmer ones are deeper in the sample.

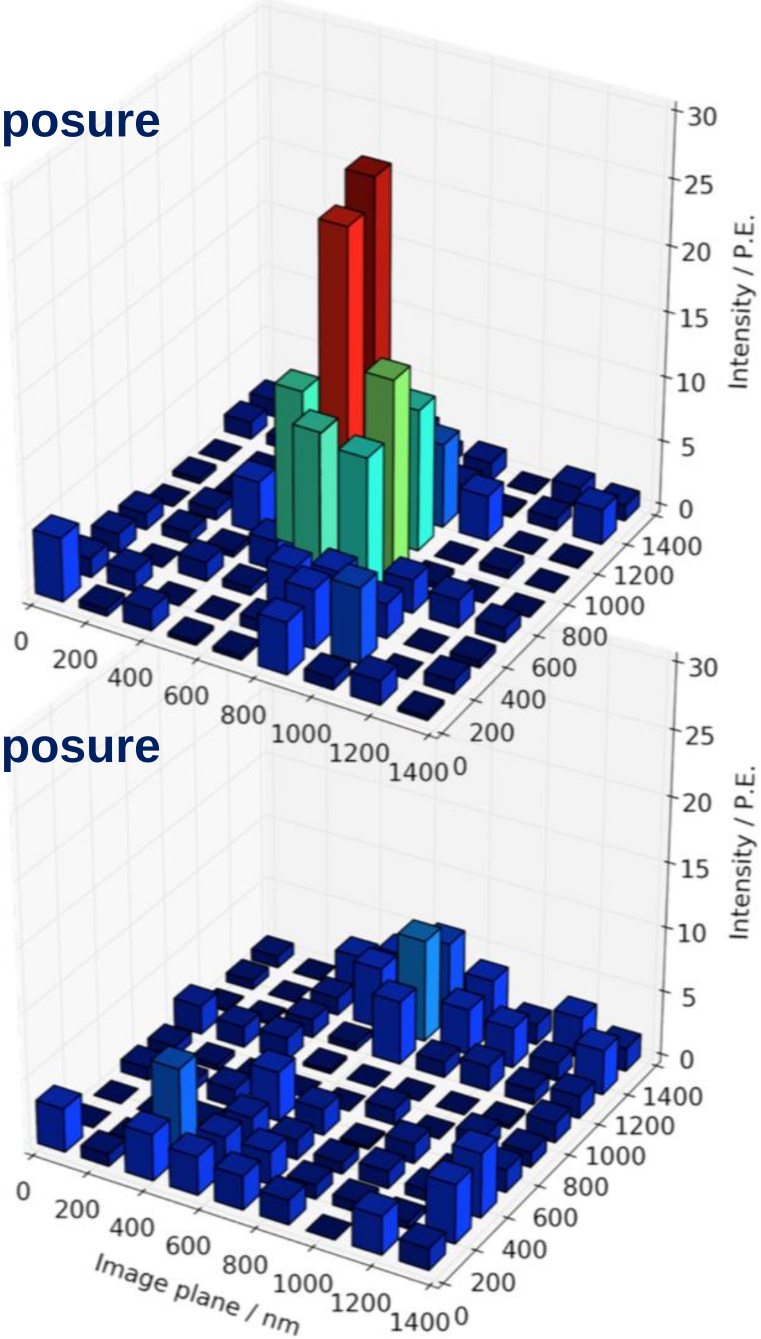
In a xenon detector, dye deposited as a monolayer and only brightest spots at constant depth expected.



This "step" is how you know it is exactly one ion.

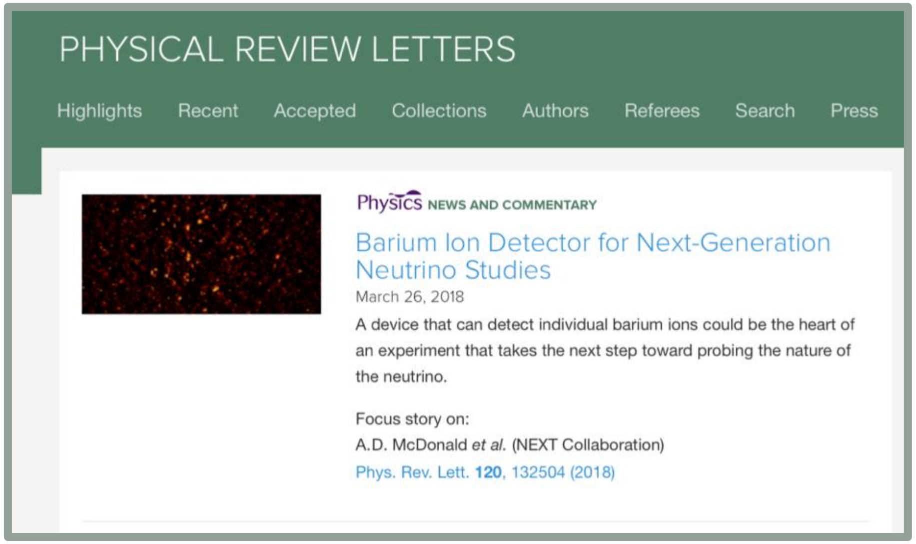
0.5s exposure before:

0.5s exposure after:



Single ions barium resolved with 2nm super-resolution and 12.9 sigma stat. significance.

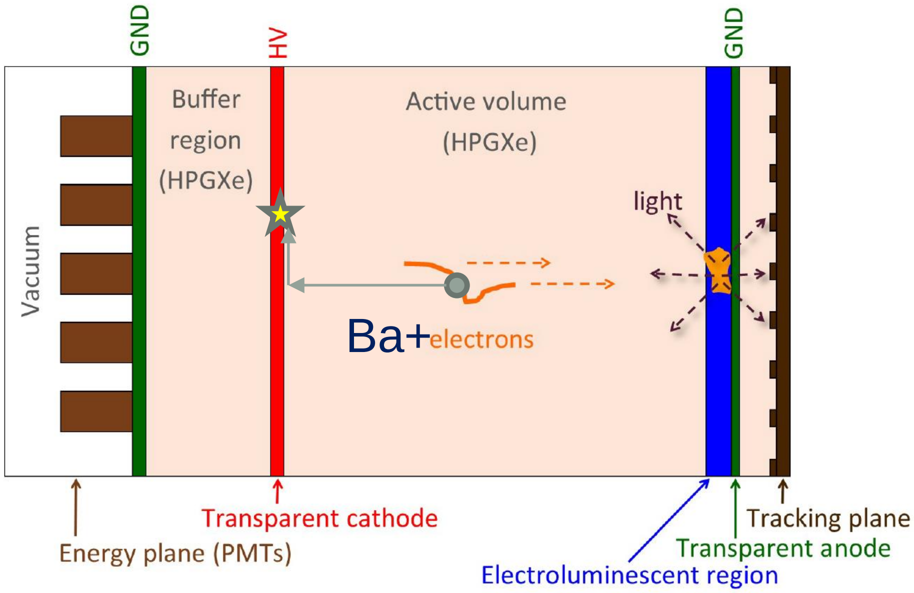
Phys.Rev.Lett. 120 (2018) no.13, 132504



First demonstration of single Ba++ ion resolution.

Next steps: Making it work in gas

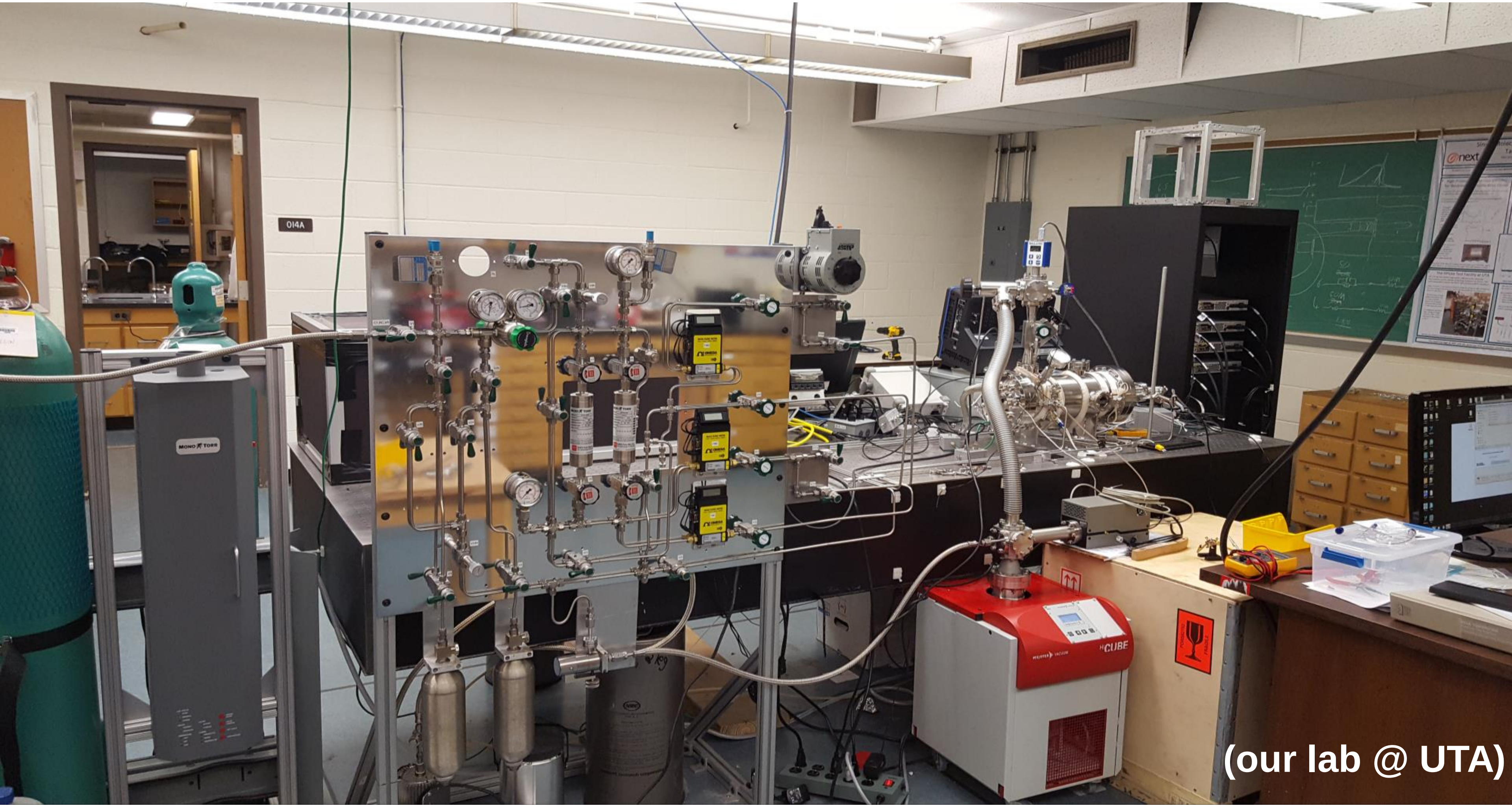
- 1. Barium ion test beam
- 2. Barium drift characterization
- 3. Dry phase microscopy
- 4. Ion concentration to sensors
- 5. Dry SMFI molecule design
- 6. Combine into a working sensor for NEXT prototype

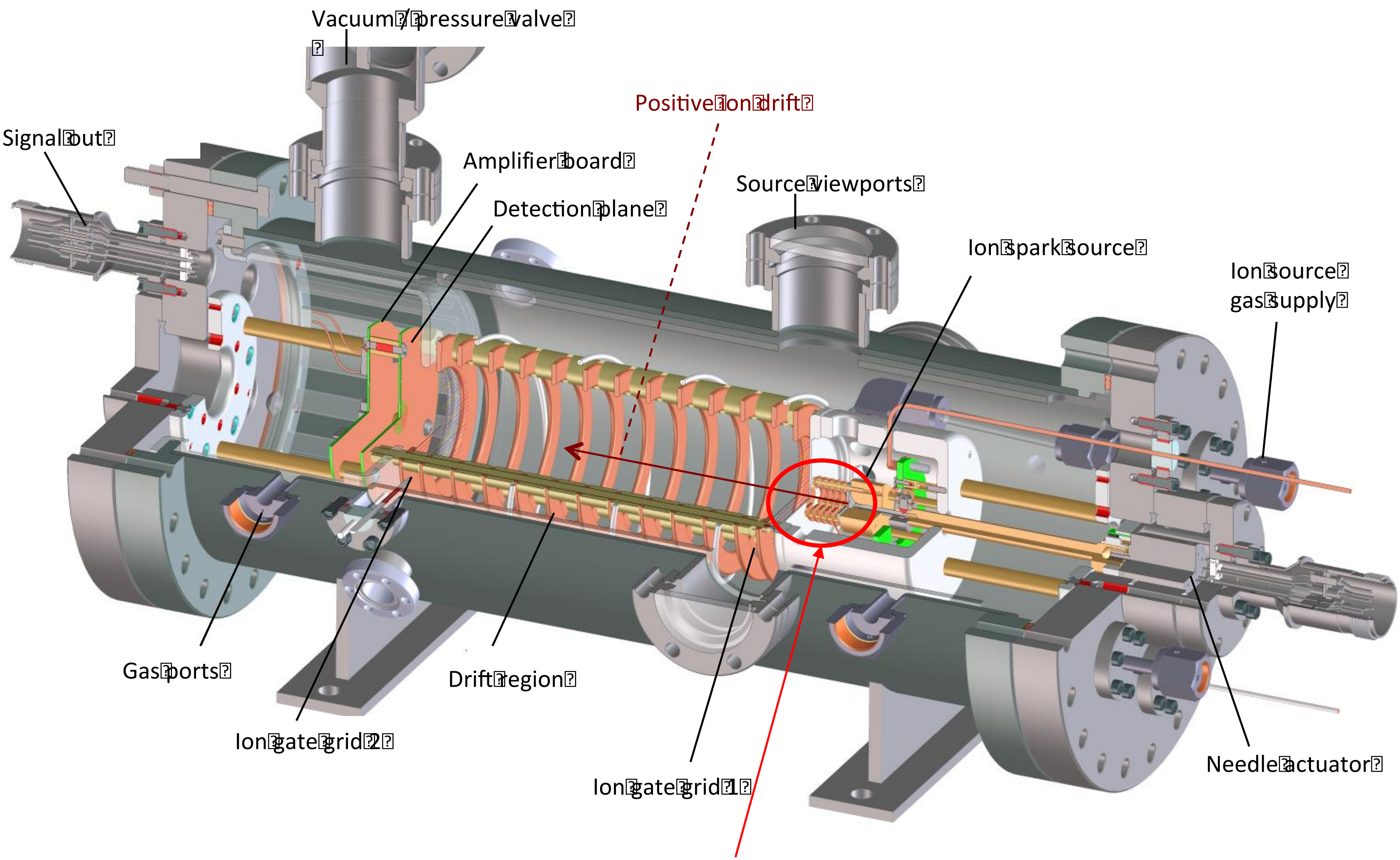


PROUDLY SUPPORTED BY:
U.S. DEPARTMENT OF ENERGY

The barium beam

- The next major step is to test barium sensing dyes in HPGXe environment
- Expect better performance than in solution from both energetics and reactivity considerations

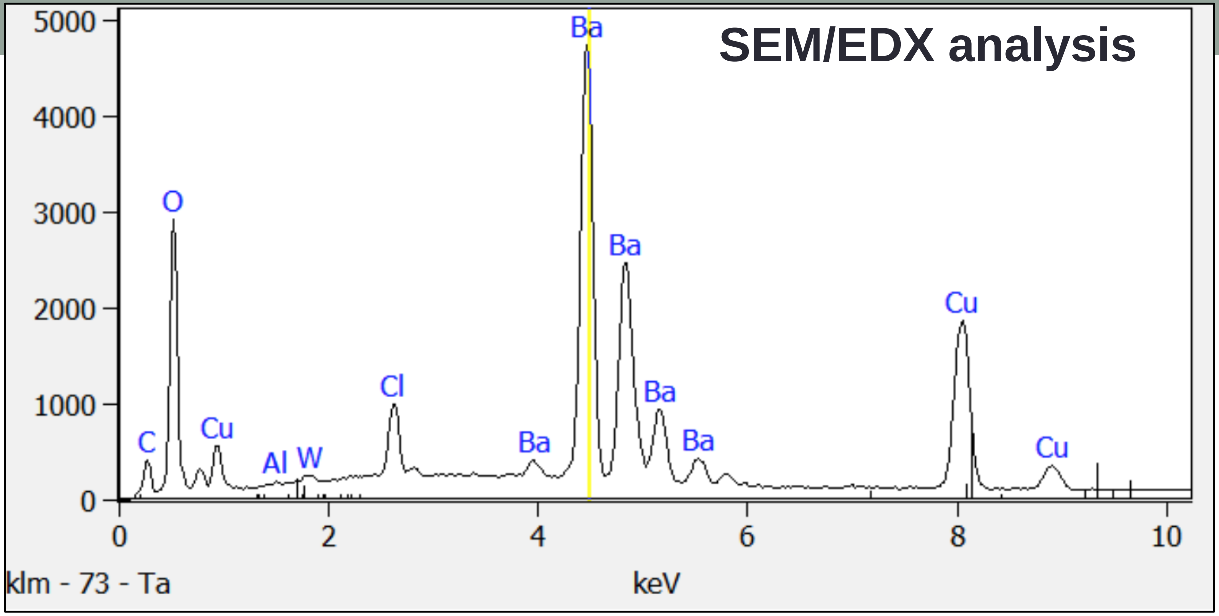




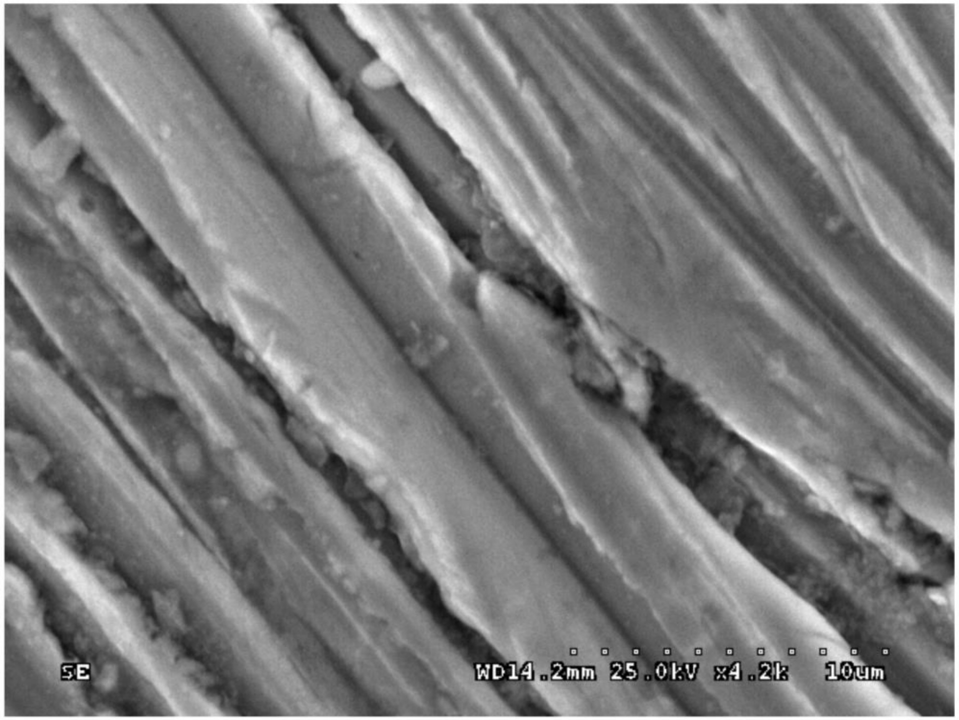
Barium or barium coated needle goes here

We have developed a custom electroplating method to deposit a stable barium-rich coating onto copper for spark source from methanol.

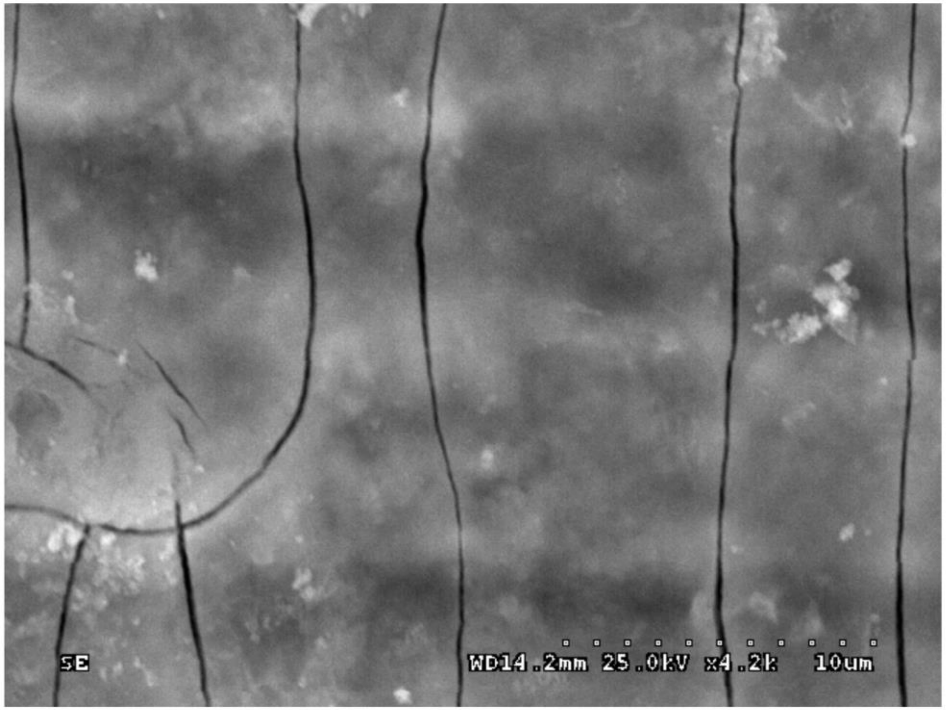
Removes difficulty of barium metal handling.



Bare copper



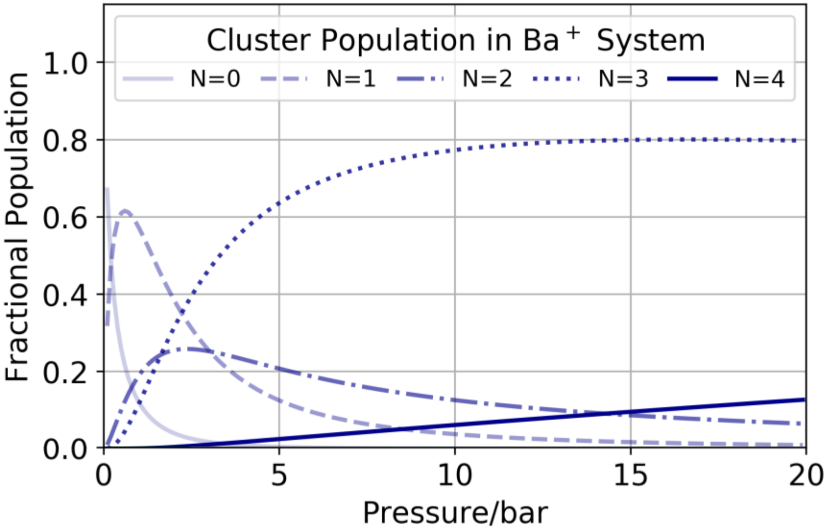
Barium electroplated copper



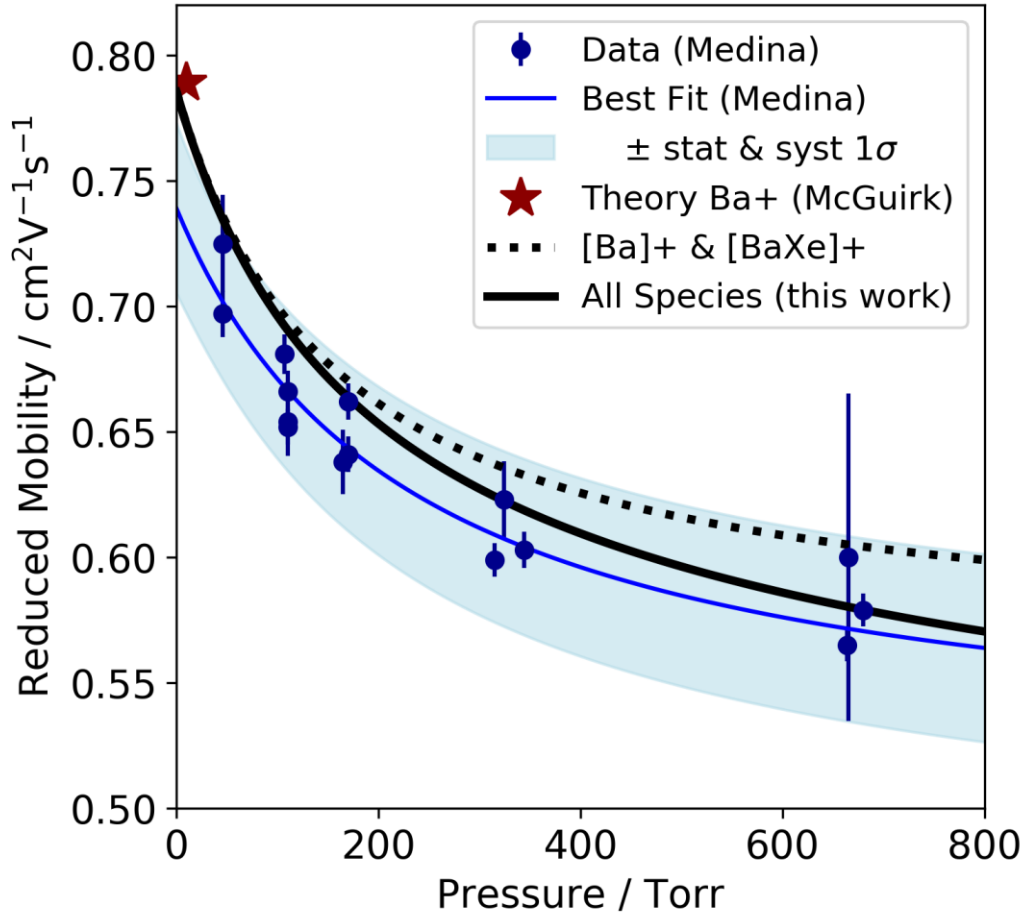
Barium mobility in gas – in theory:

Not a single ion drifting!

$$Ba^+ + Xe + Xe \longleftrightarrow BaXe^+ + Xe$$



Excellent agreement with data for Ba+

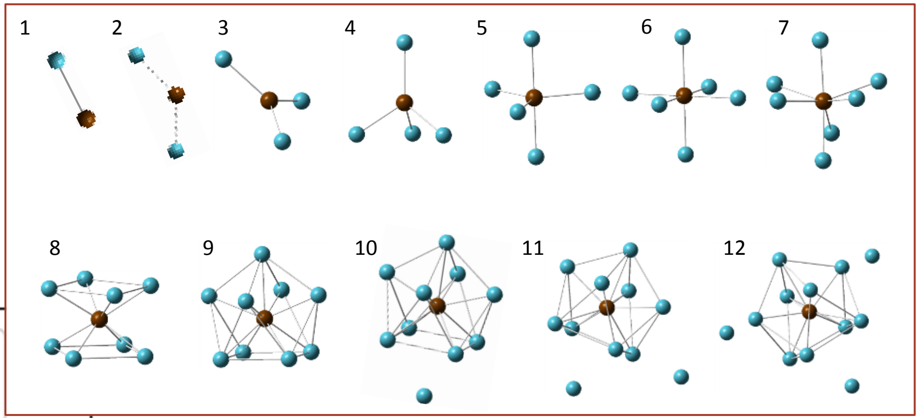


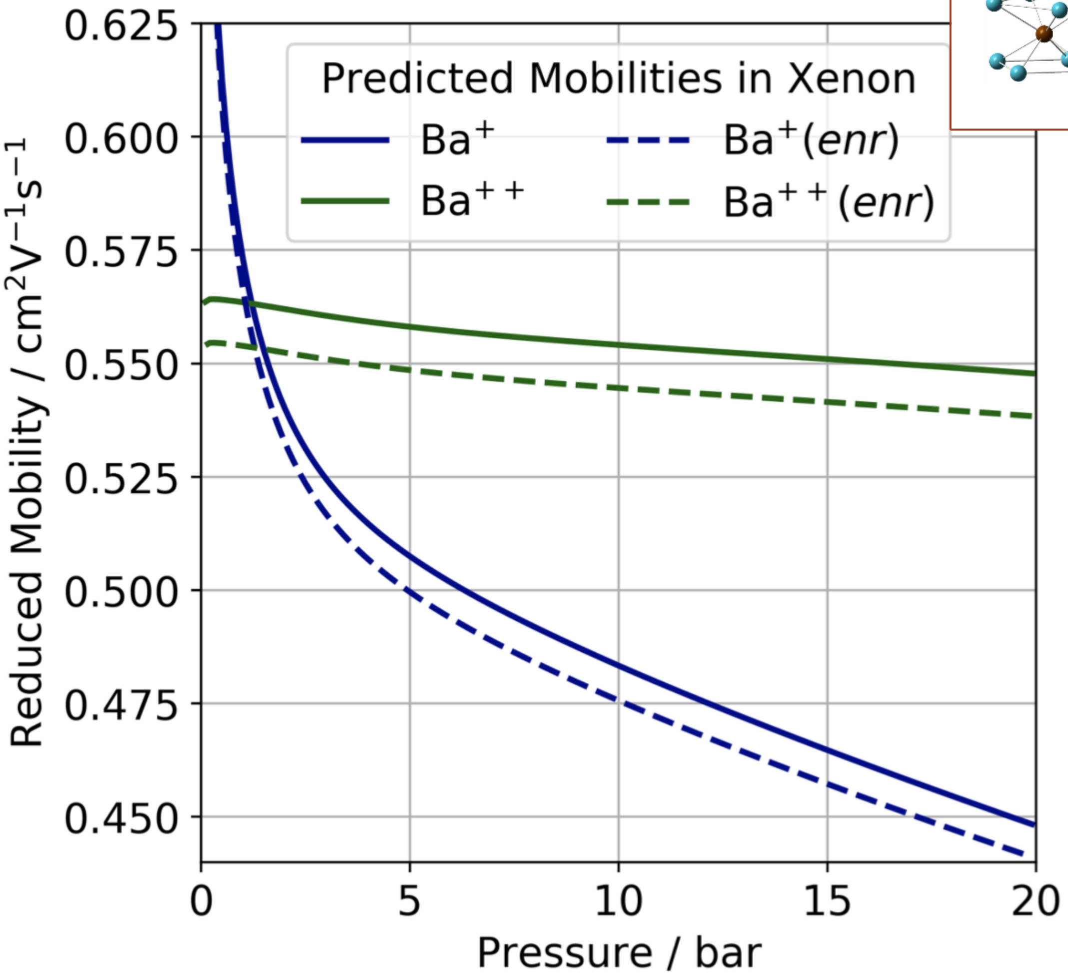
Mobility and Clustering of Barium Ions and Dications in High Pressure Xenon Gas

E. Bainglass, B.J. P. Jones, et. al. Phys.Rev. A97 (2018) no.6, 062509

For Ba++ things get more complicated.

Calculated Ba++ clusters:





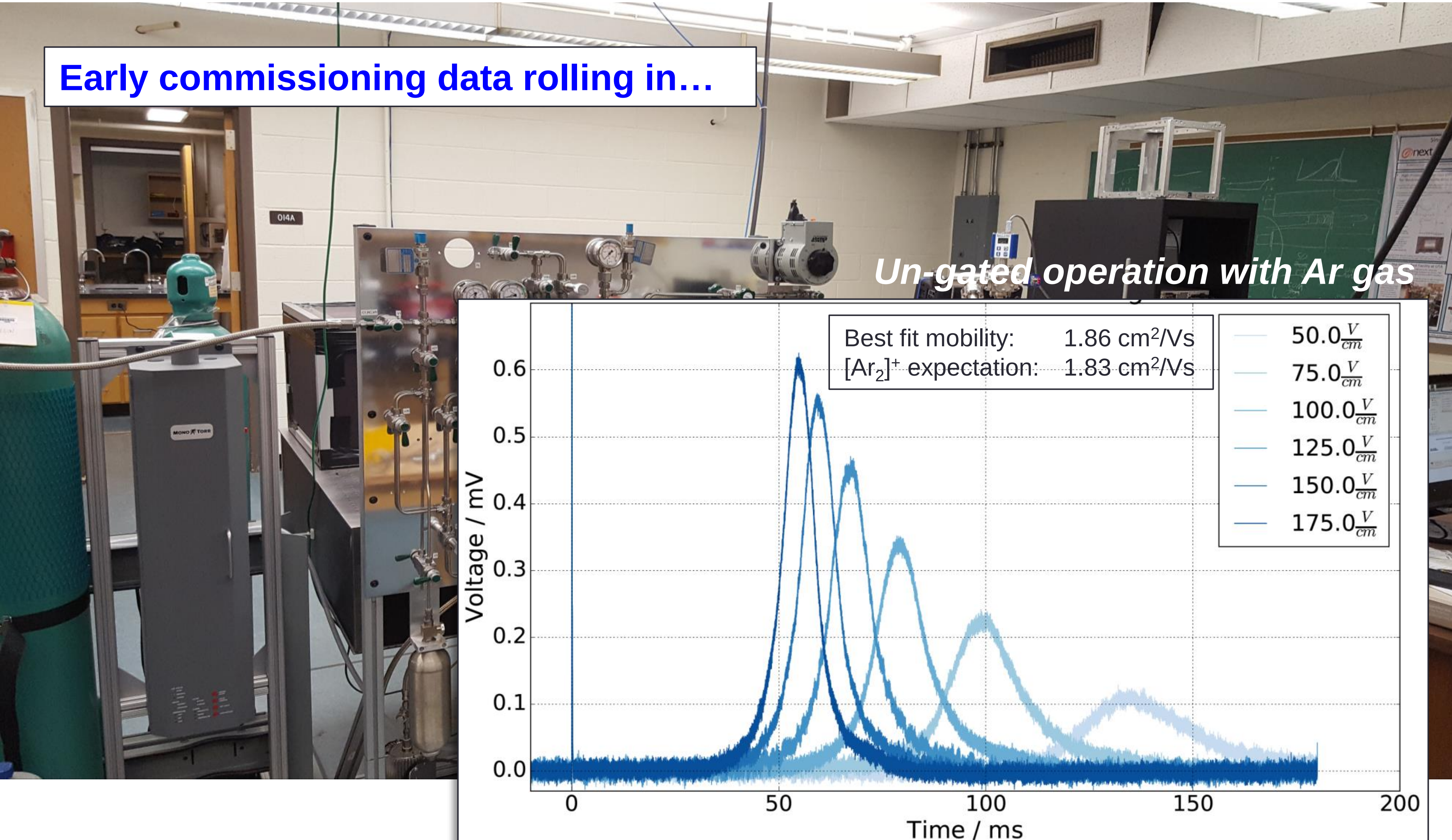
Bigger clusters more similar to each other, so less pressure dependence in Ba++ than Ba+

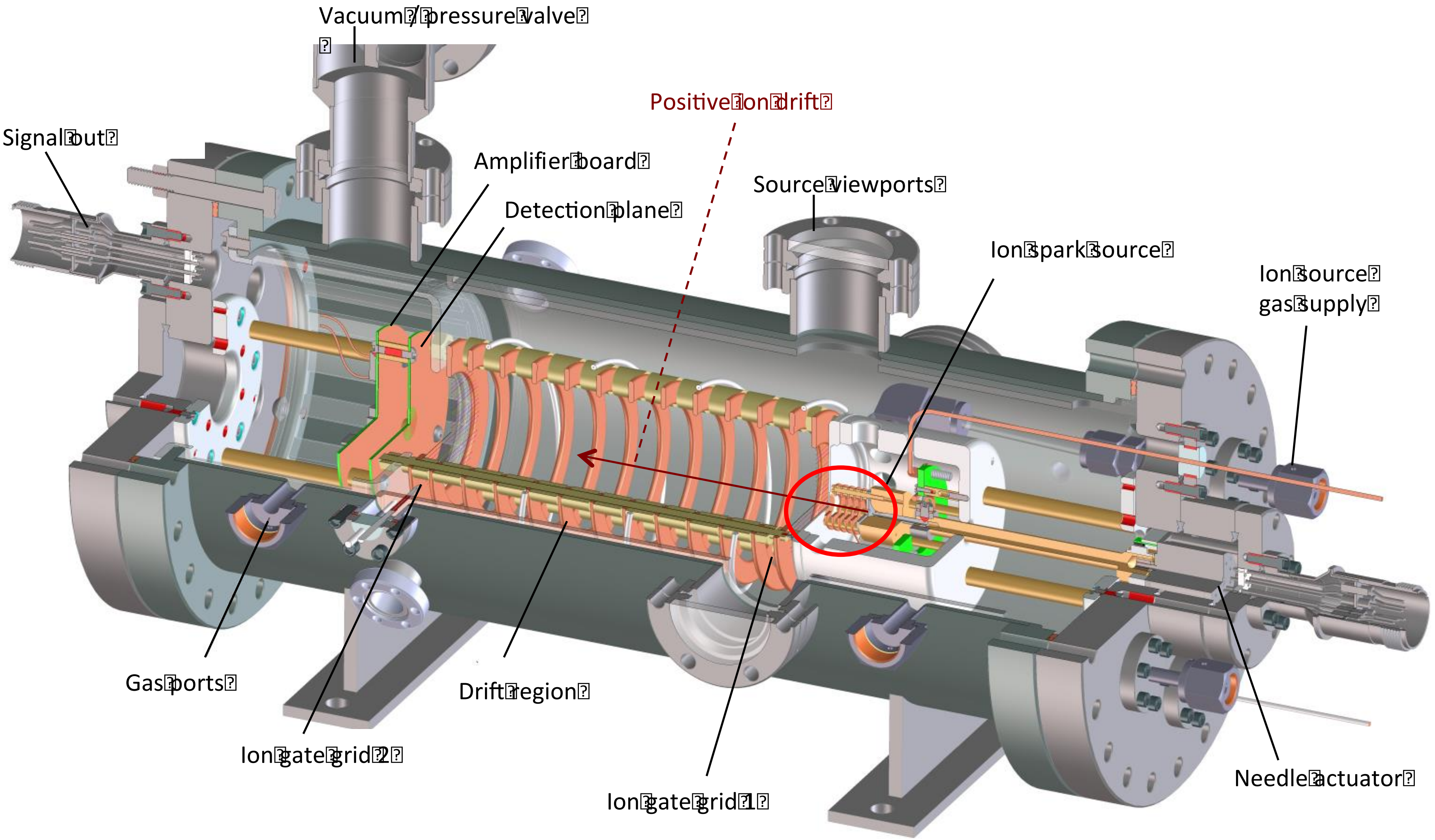
Isotopic composition changes scattering kinematics, so %-level differences with enriched xenon

We will test this with experimental data soon!

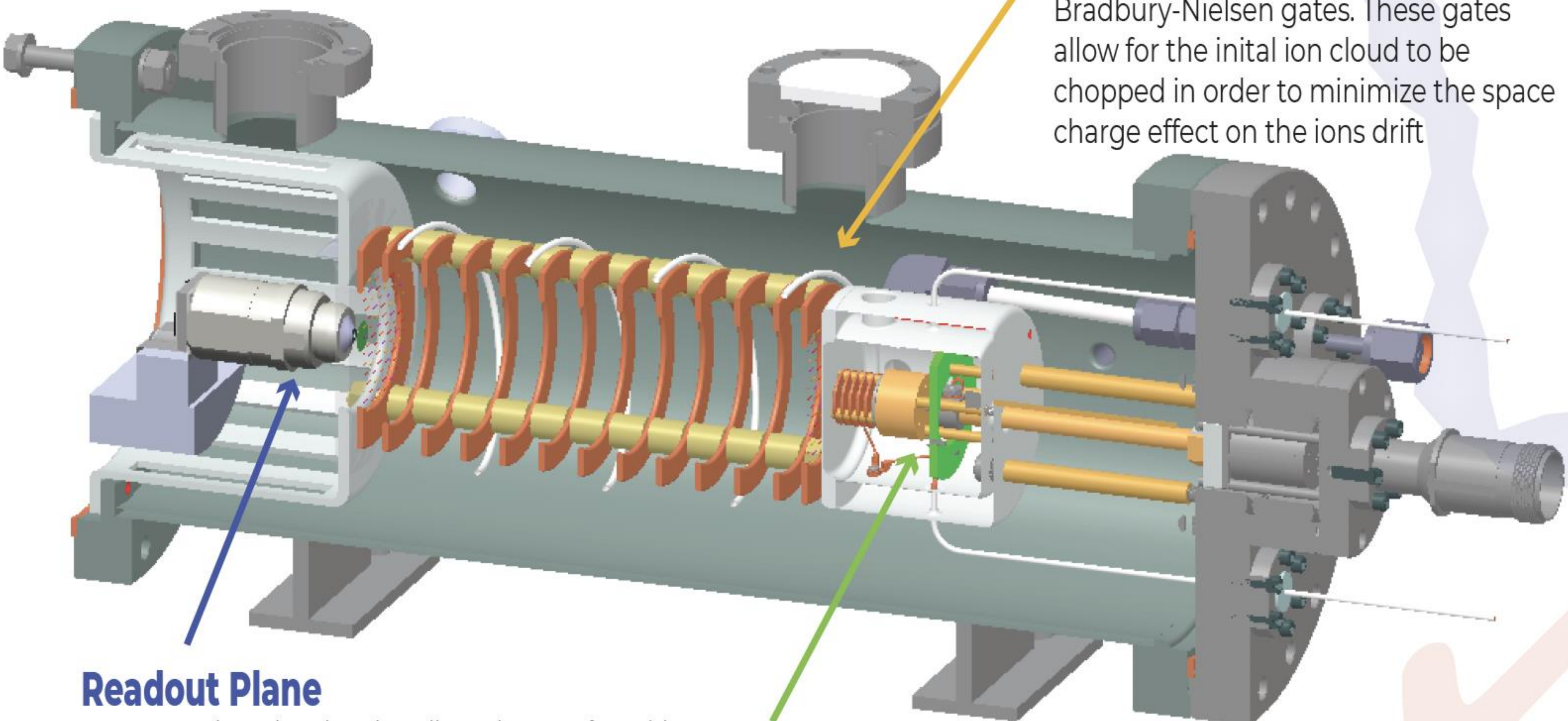
The ion beam is coming…

Early commissioning data rolling in…





Vacuum / pressure valve
Signal out
Amplifier board
Detection plane
Positive ion drift
Source viewports
Ion spark source
Ion source gas supply
Gas ports
Drift region
Ion gate grid 2
Ion gate grid 1
Needle actuator



Gated Drift Region

20cm uniform drift region that utilized Bradbury-Nielsen gates. These gates allow for the inital ion cloud to be chopped in order to minimize the space charge effect on the ions drift

Readout Plane

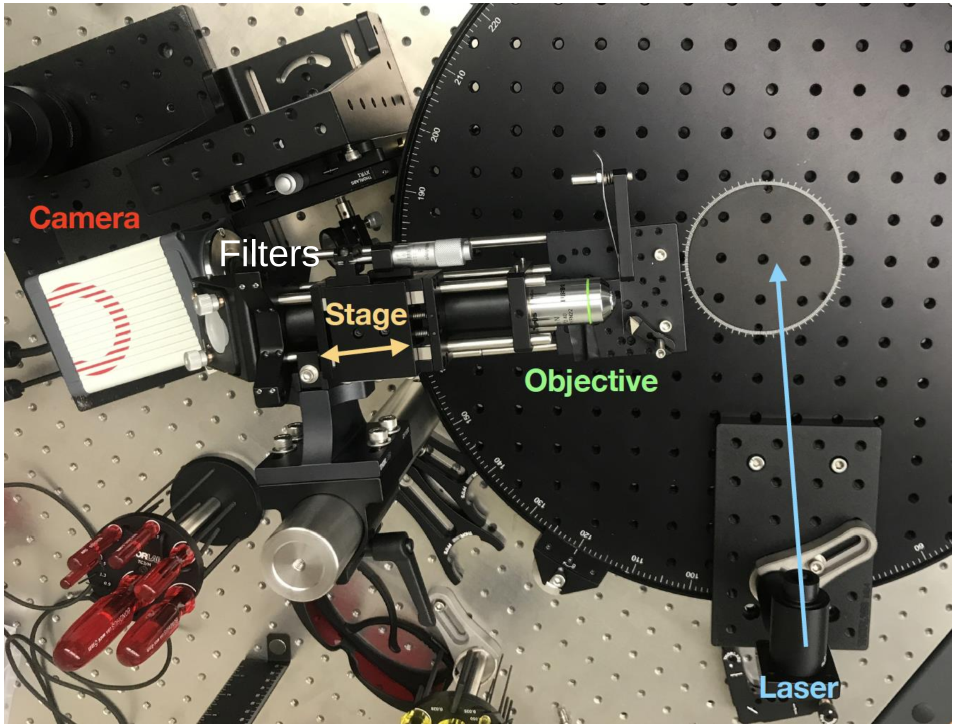
A readout chamber that allows the use of sensitive eletrical amplifier or optical readout. The ions are guided here via an eletric field.

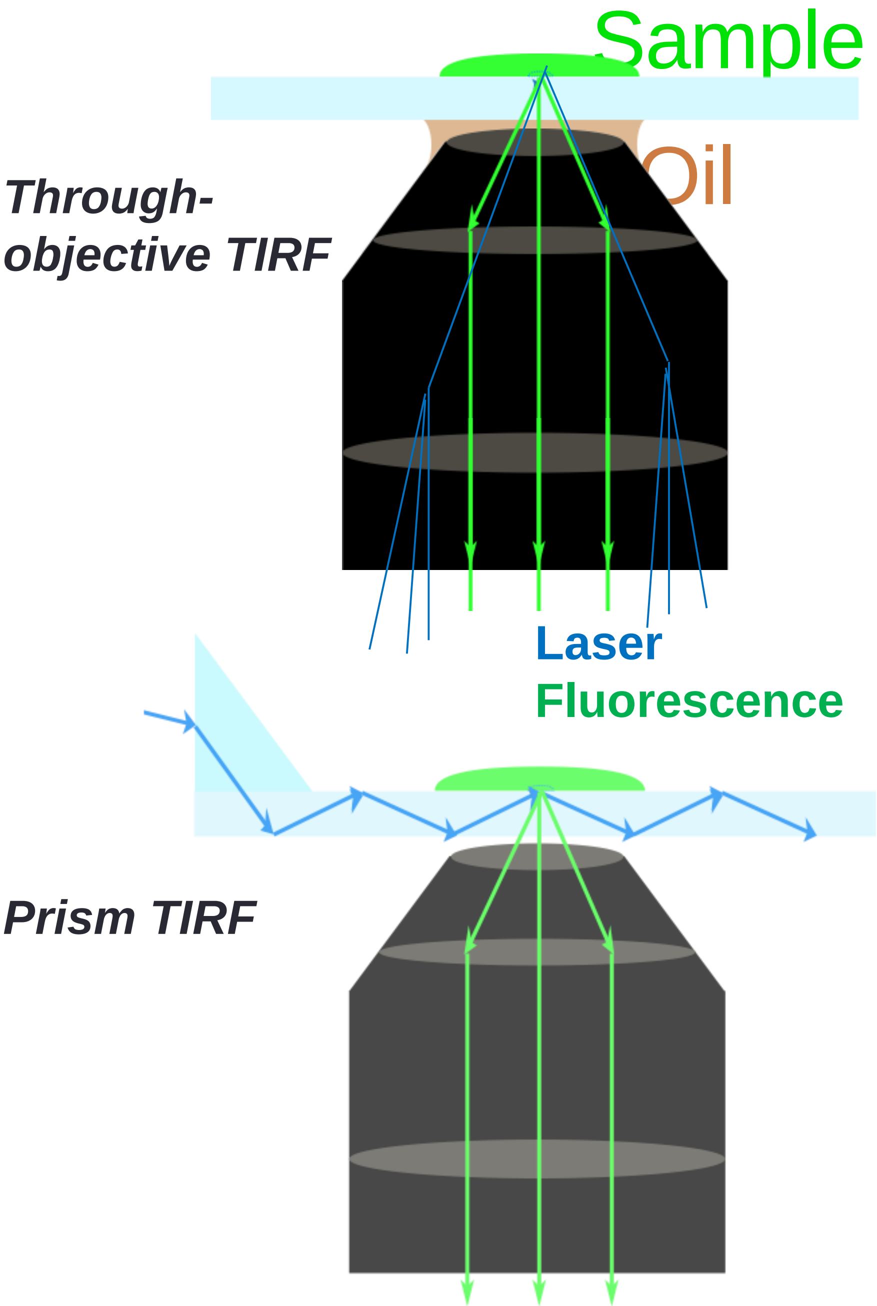
Ion Source

A needle and ring electrode that is adjustable allowing for an adjustable spark gap to acomadate Various pulse and pressure conditions.

Dry Microscopy

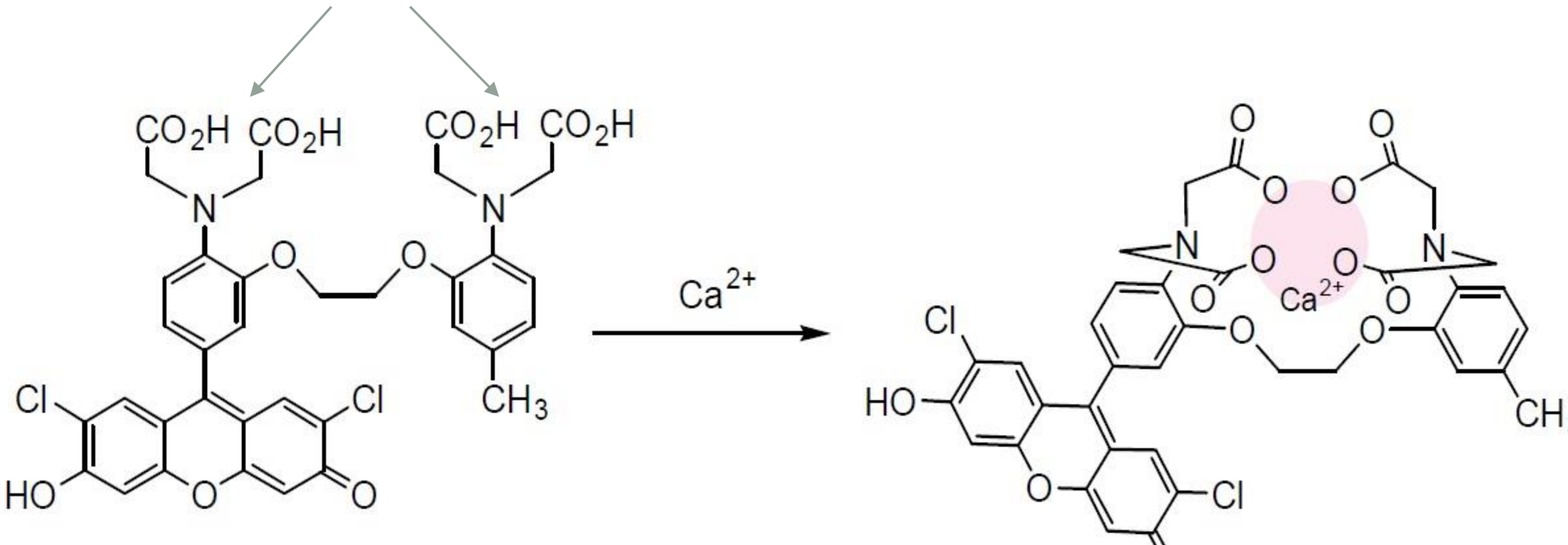
- To remove microscope oil (incompatible with HPGXe) we need to deliver excitation transversely.
- This is achieved through prism TIRF- decouples the motion of the objective and the TIRF source.
- We are building up a prism TIRF array at UTA for operation in gas





Dry Phase Receptor

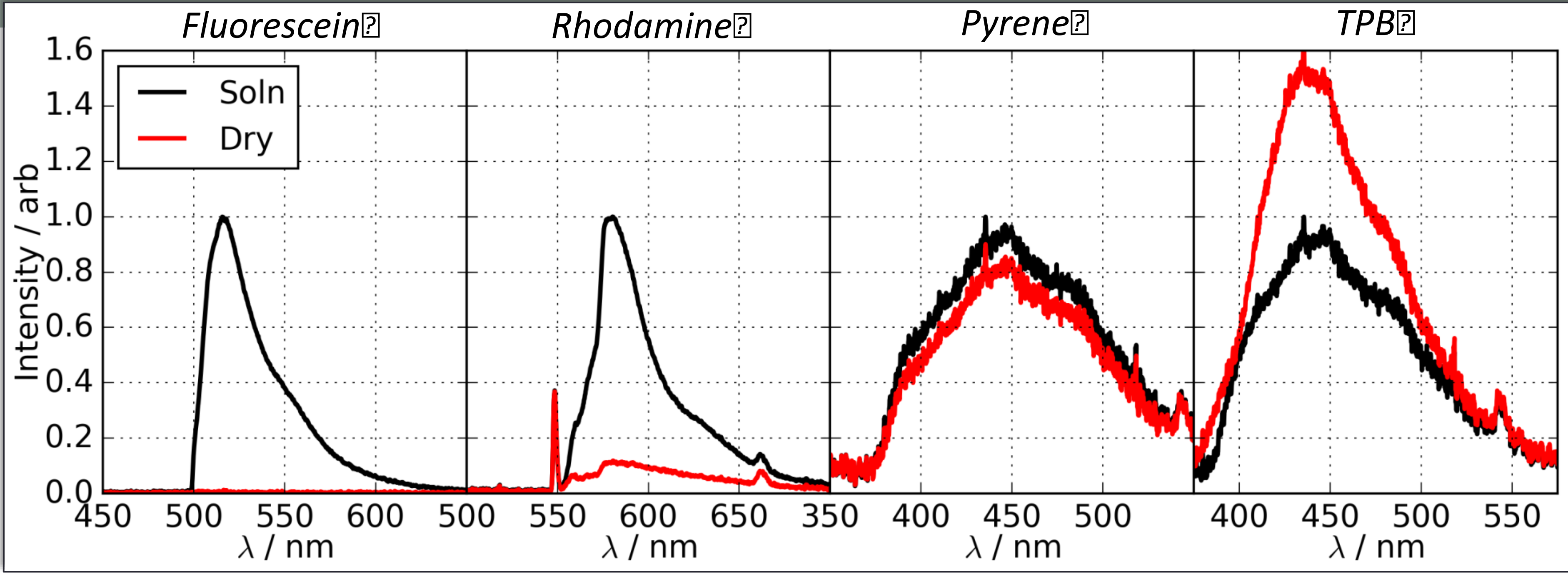
Deprotonation of carboxylic acids is required to accept the ion – we observe the characteristic pH dependence of this in solution



Fluo 3

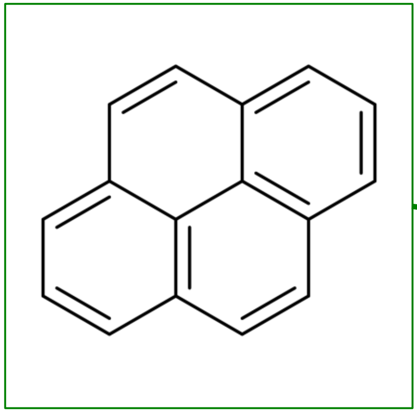
Fluo 3-Ca complex

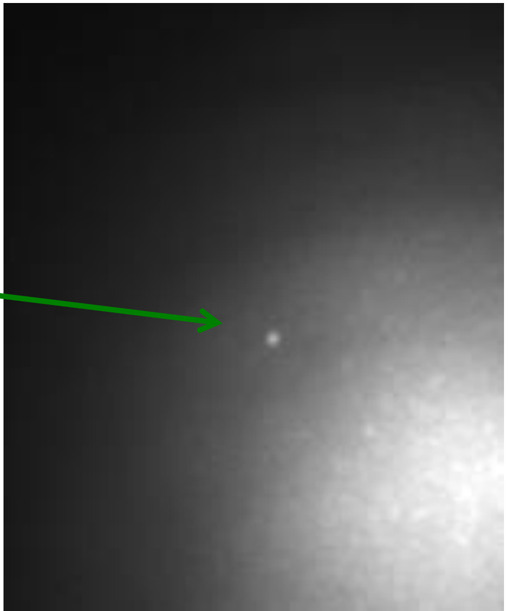
It is unclear if deprotonation will occur in solid phase but we suspect perhaps not. This motivated us to start building “dry-taylored” molecules even before we have a dry system to test in.



Fluo-3 may not be ideal for HPGXe, since fluorescein dye is not bright in dry phase.

But pyrene works. We can resolve single molecules of it too.



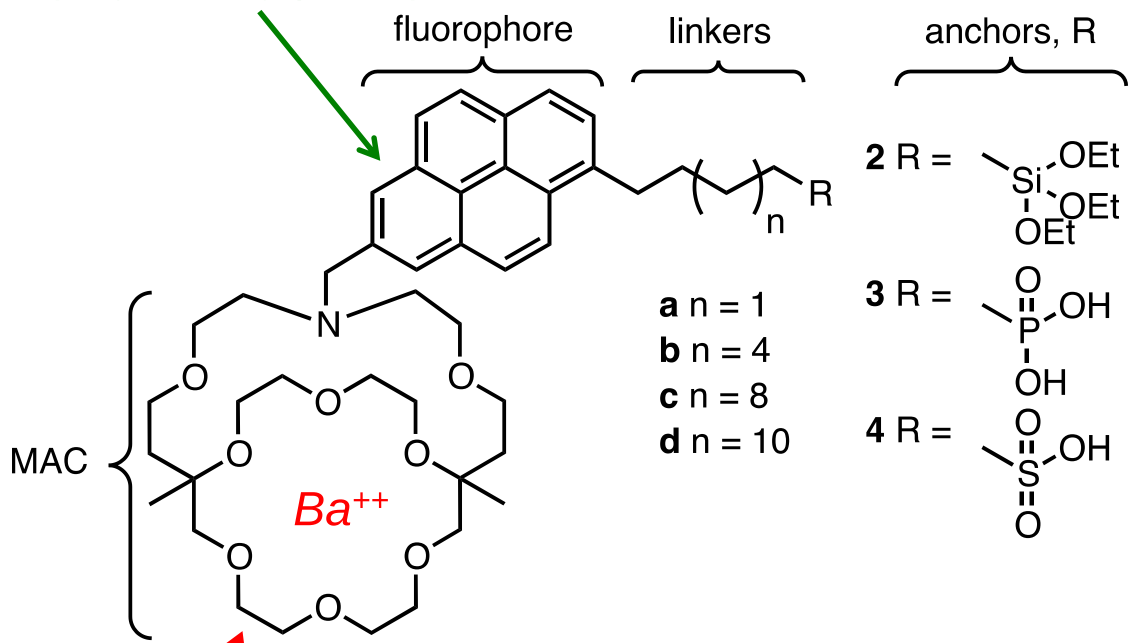


The Dream Molecule?

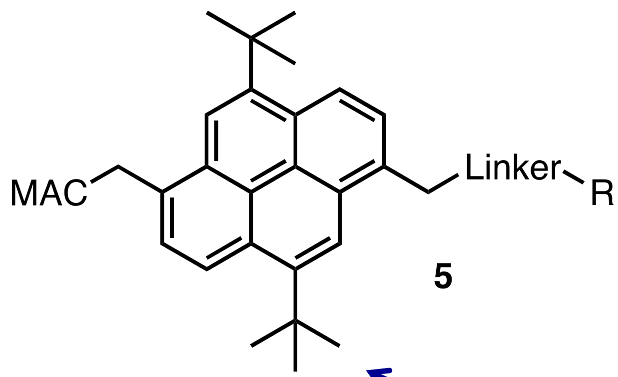
Tethered pyrene substituted monoazocryptand (tPSMA)

Our first batch is being synthesized right now by Foss lab at UTA.

Pyrene
(bright in dry phase)



Monoazocryptand
(hyper Ba^{++} specific)



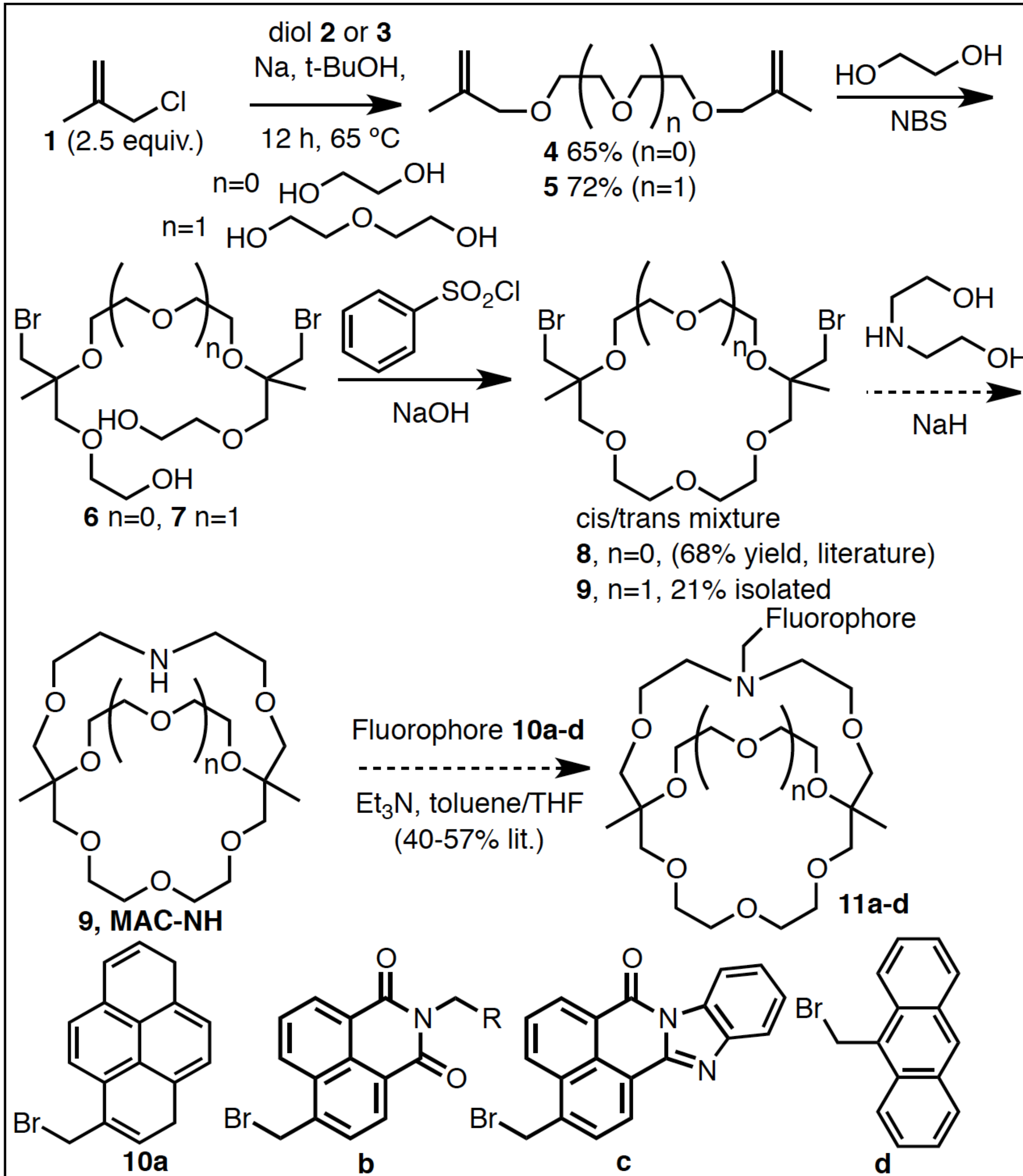
5

Possible "arms" inhibit surface quenching of fluor

Diner, C., Scott, D. E., Tykwinski, R. R., Gray, M. R., & Stryker, J. M. Scalable, Chromatography-Free Synthesis of Alkyl-Tethered Pyrene-Based Materials. Application to First-Generation "Archipelago Model" Asphaltene Compounds. *J. Org. Chem.* **2015**, *80*(3), 1719–1726.

Synthesis of PSMA and its relatives

The story so far…

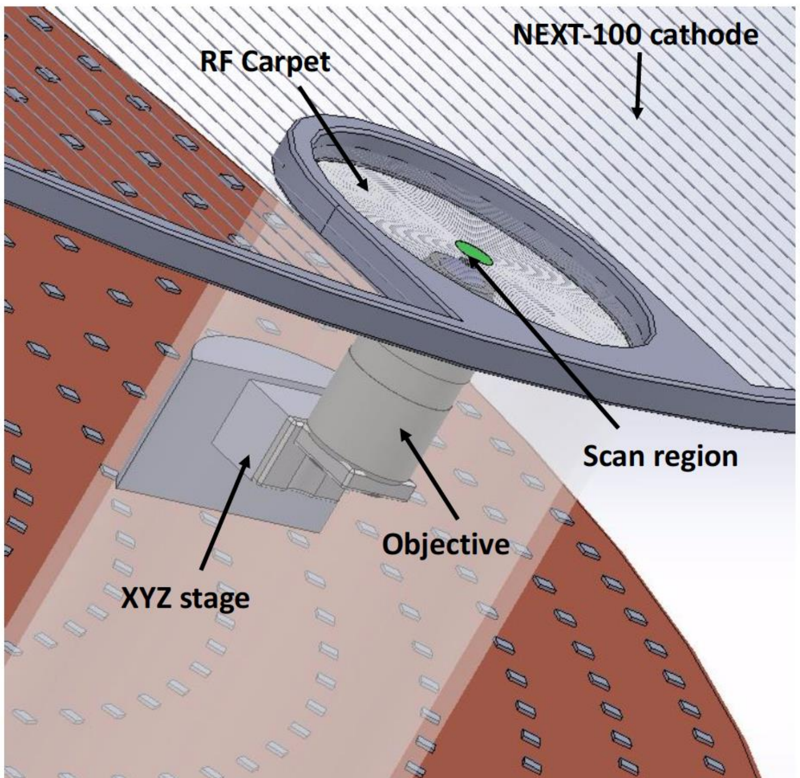


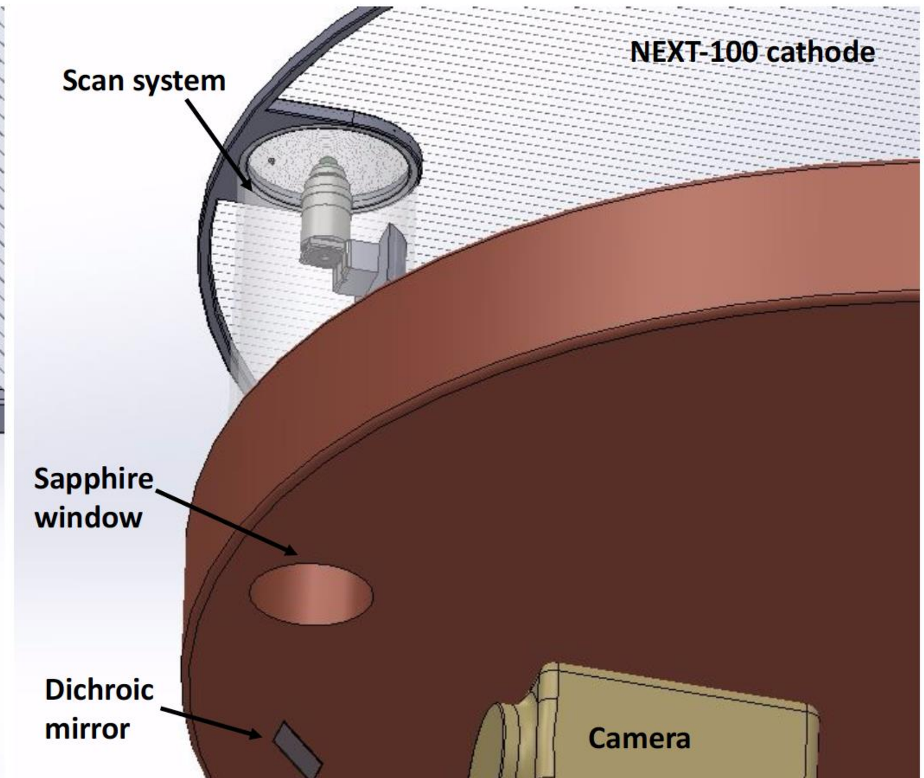
First batch of custom dry fluorophores coming off production line now.

Watch this space!

Implementation in NEXT

- The NEXT collaboration is exploring implementation of a prototype barium tagging sensor in a near-future phase of NEXT.
- Our R&D plan aims to converge on an end-to-end barium tagging scheme by 2022-2023. Optimistic, but looks possible from where we stand now.
- **A ton-scale barium tagging HPGXeTPC may be what we need to uncover the Majorana nature of the neutrino, and open the window to the next energy scale.**





Contacto | A A | english | euskera | Imprimir | Enviar a

→ Inicio → Actividades → Workshops 2018 → Single Atom Ba Tagging (SABAT)

Presentación

Single Atom Ba Tagging (SABAT)

The DIPC-Workshop on **"Single Atom Ba Tagging" (SABAT 2018)** will be held at the "Donostia International Physics Center" (DIPC), Donostia-San Sebastian, on **July 9-10, 2018**. The aim of the workshop is to review the experimental techniques needed to achieved single-atom tagging in Xe-137 double beta decays produced in a gas xenon TPC. Currently, the most promising way to achieve Ba++ tagging is the use of Single Molecule Fluorescence Imaging. The notion, was proposed by Dave Nygren in 2017, and a first proof of concept has been carried out by the NEXT collaboration in 2018. Further progress requires a multi-disciplinary approach involving several fields.

Topics include:
* Applying Single Molecule Fluorescence Imaging to single-Ba++ ion tagging.
* Proof of principle of SMFI technique for Ba++ tagging.
* Developments of dyes for SMFI Ba++ tagging.
* SMFI Ba++ tagging in vacuum. Molecule layers and microscopy.
* SMFI Ba++ tagging in high pressure gas.
* Use of nano-technology to amplify signal.
* Development of Ba++ sources.
* Development of RF ``carpets" for Ba++ collection.
* Implementation in HPXe chambers (e.g, NEXT detector).

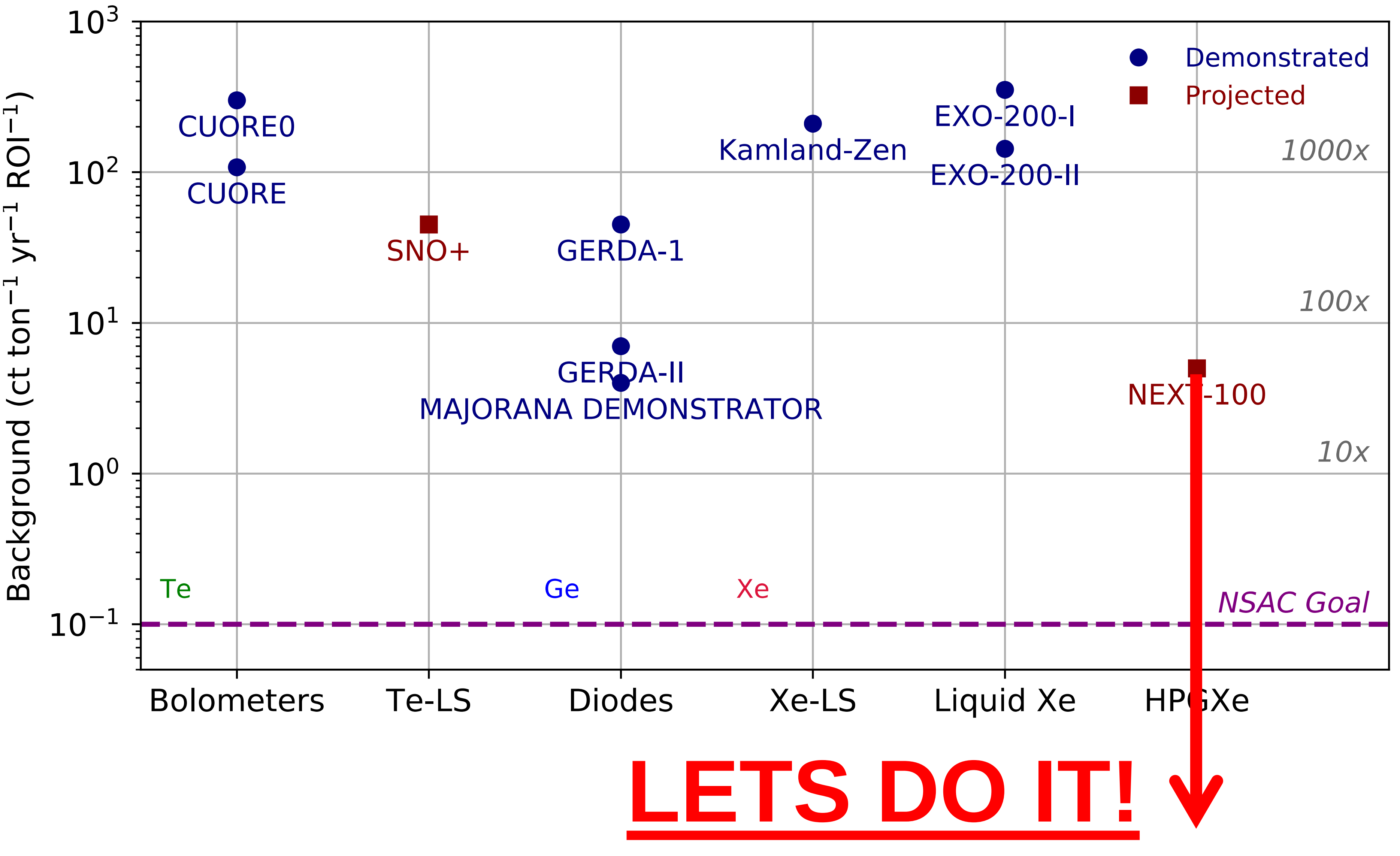
In the spirit of DIPC Workshops, SABAT would like to provide a friendly atmosphere to foster discussions and collaborations between leading scientists working in the various fields involved.

More information is available on the workshop web page: http://sabat2018.dipc.org/

International interest group formed, first meeting at DIPC, San Sebastian, July 2018.

Interested? Get in touch!

THIS QUESTION OF NEUTRINO MASS IS TOO IMPORTANT NOT TO ANSWER, OR TO ANSWER AMBIGUOUSLY.



LETS DO IT!

Thanks!

on behalf of…

The NEXT Collaboration

USA *(co-spokesperson David Nygren)*

University of Texas Arlington

Argonne National Laboratory

Fermilab

Harvard (VE RI TAS)

Iowa State University

A&M

Spain *(co-spokesperson JJ Gomez Cadenas)*

DIPC

IFIC Institut de Física Corpuscular

Universitat de Girona

Universitat Politècnica de València

Universidad de Zaragoza

USC Universidade de Santiago de Compostela

Others

U C

universidade de aveiro

Ben-Gurion University of the Negev

JINR Dubna

UAN Universidad Antonio Nariño

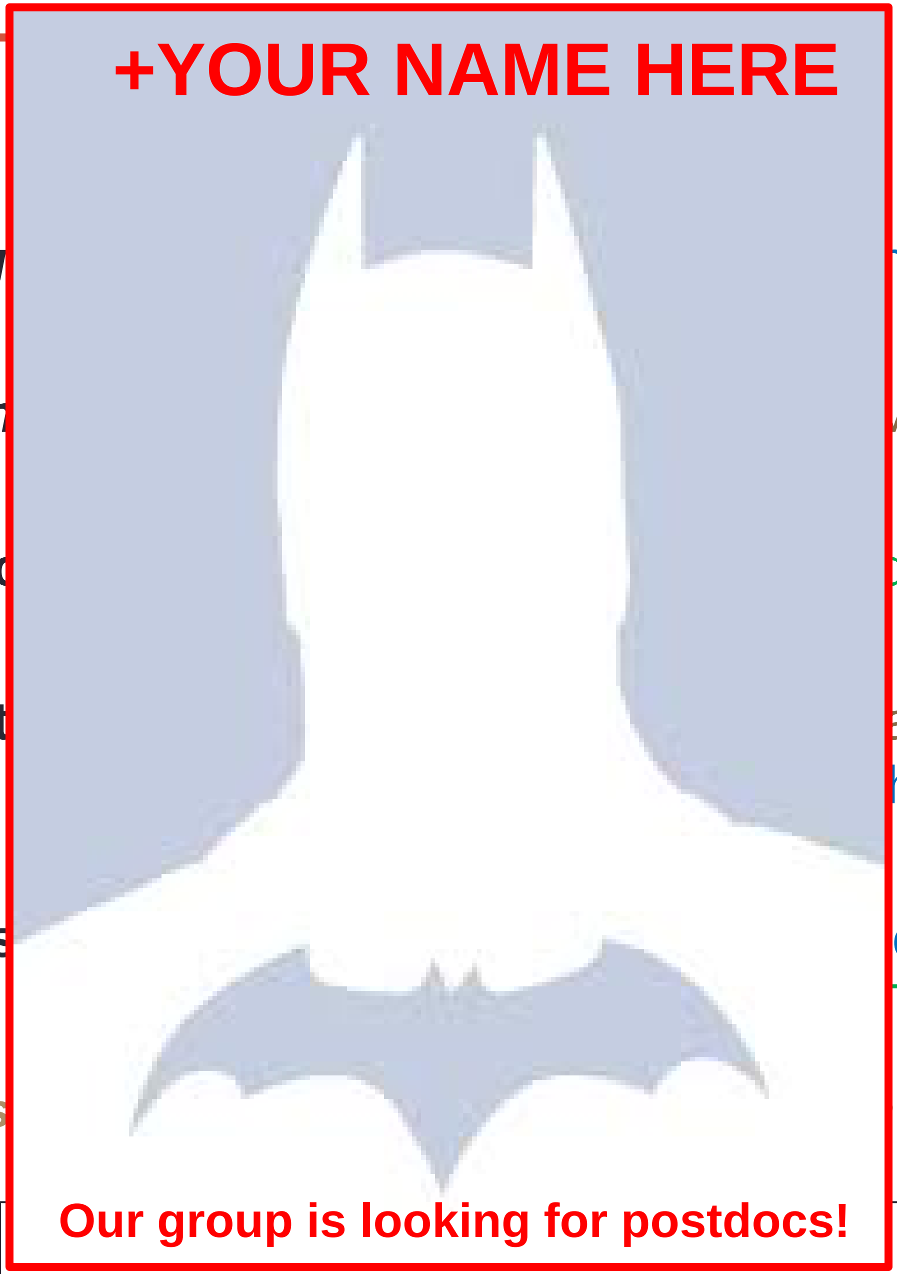
The Ba Tag Team at UTA

- ***SMFI Ba++ Tagging Leaders:*** BJPJ and David Nygren
- ***Barium Beam and SMFI Demonstration:*** Austin McDonald
- **Gas Phase Microscopy R&D:** Fernanda Psihas
- **Condensed Matter Theory:** Edan Bainglass and Mohammed Huda
- **Organic Synthesis of Custom Dyes:** Frank Foss and Pawan Thapa

+ Undergrads Denise Huerta, Zane Miller, Jerry Tram, Ryan Clark, and others…

Prof	***Postdoc***	***Student***

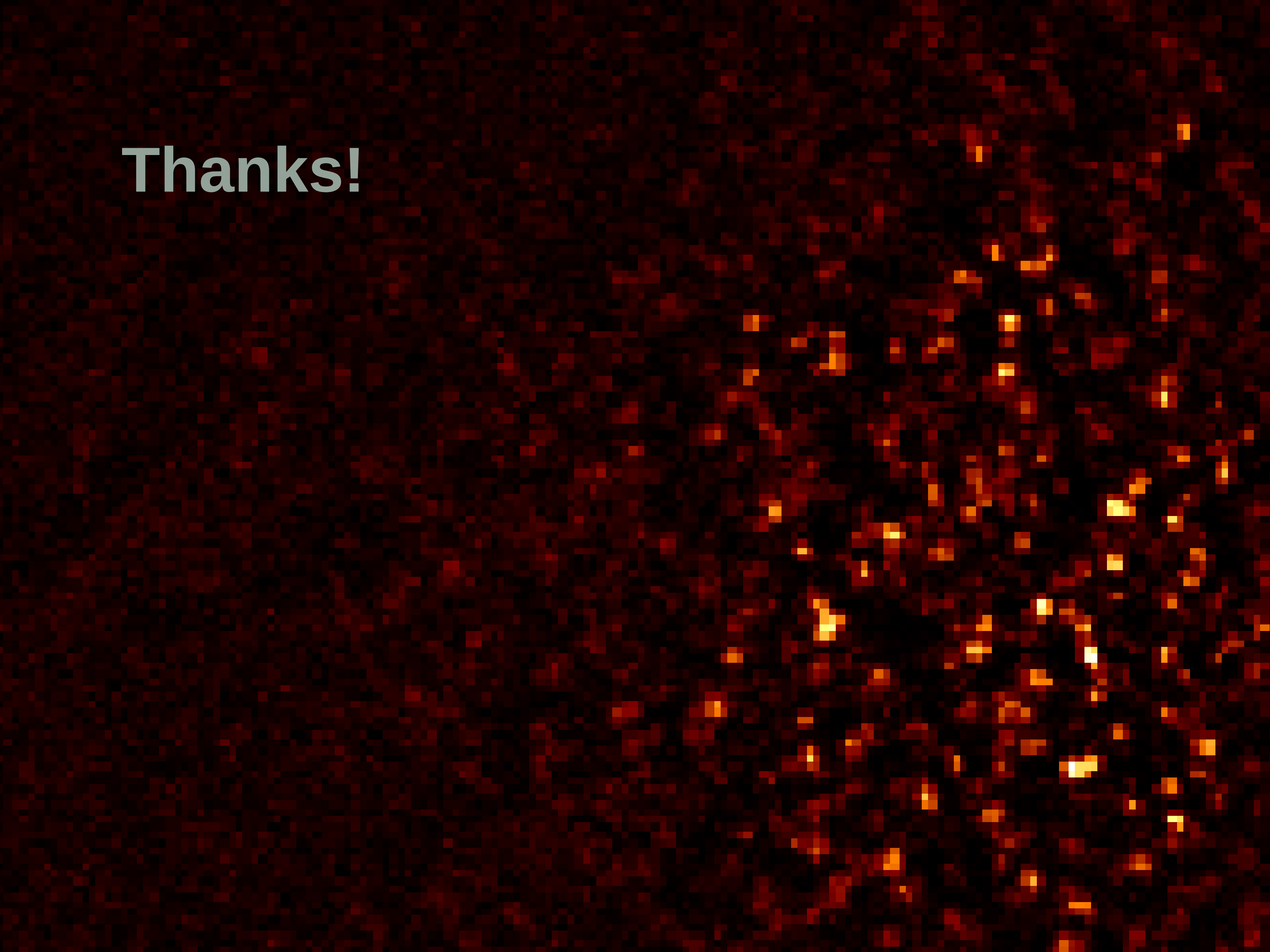
The Ba T



- ***SMFI Ba++ Tagg*** nd David Nygren
- ***Barium Beam an*** McDonald
- **Gas Phase Micr** da Psihas
- **Condensed Matt** ainglass
  hammed Huda
- **Organic Synthes** oss and
  Thapa

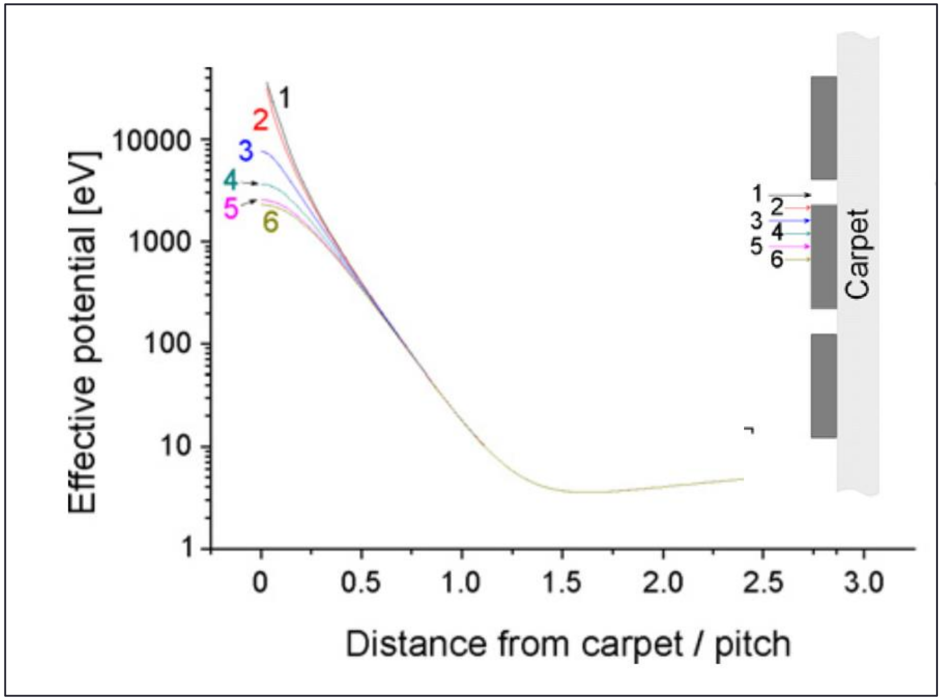
+ Undergrads Denis ***Clark, and others…***

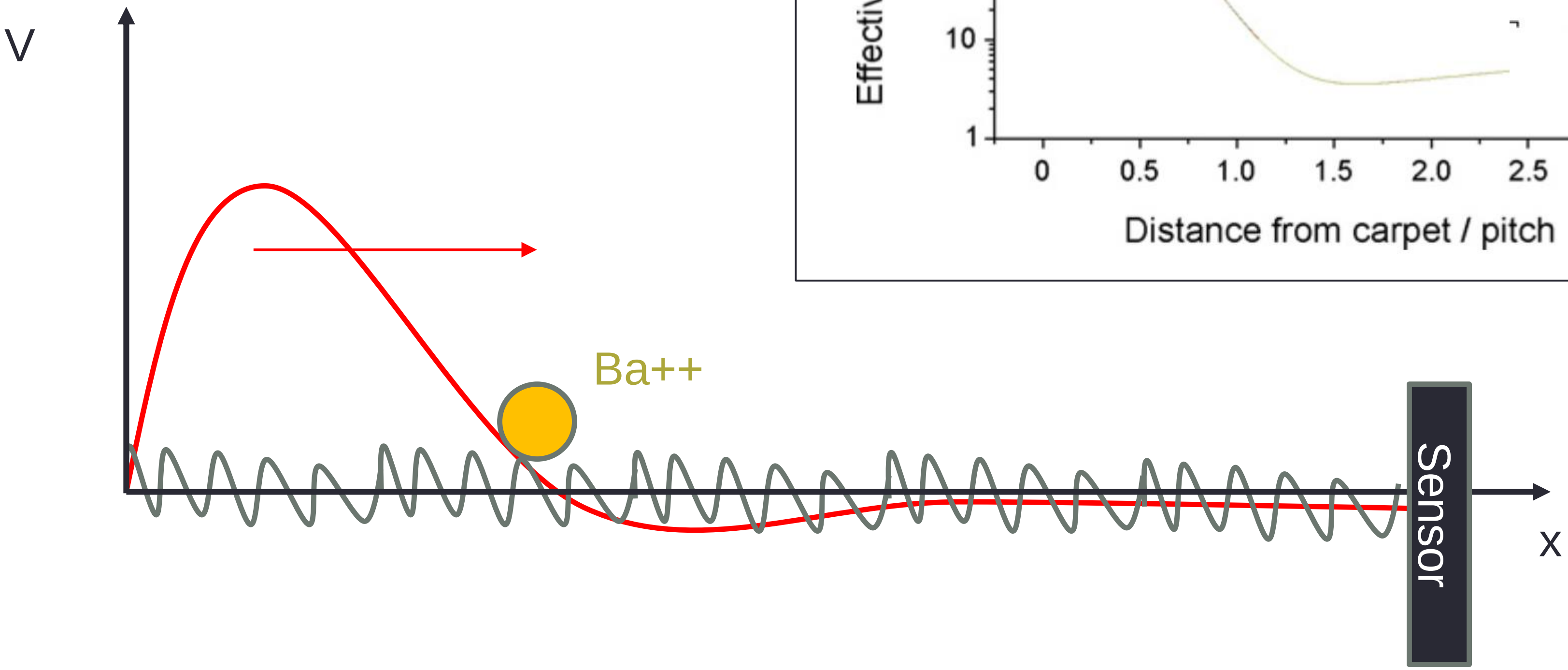
Thanks!



RF Carpet fields

International Journal of Mass Spectrometry 299 (2011) 71–77



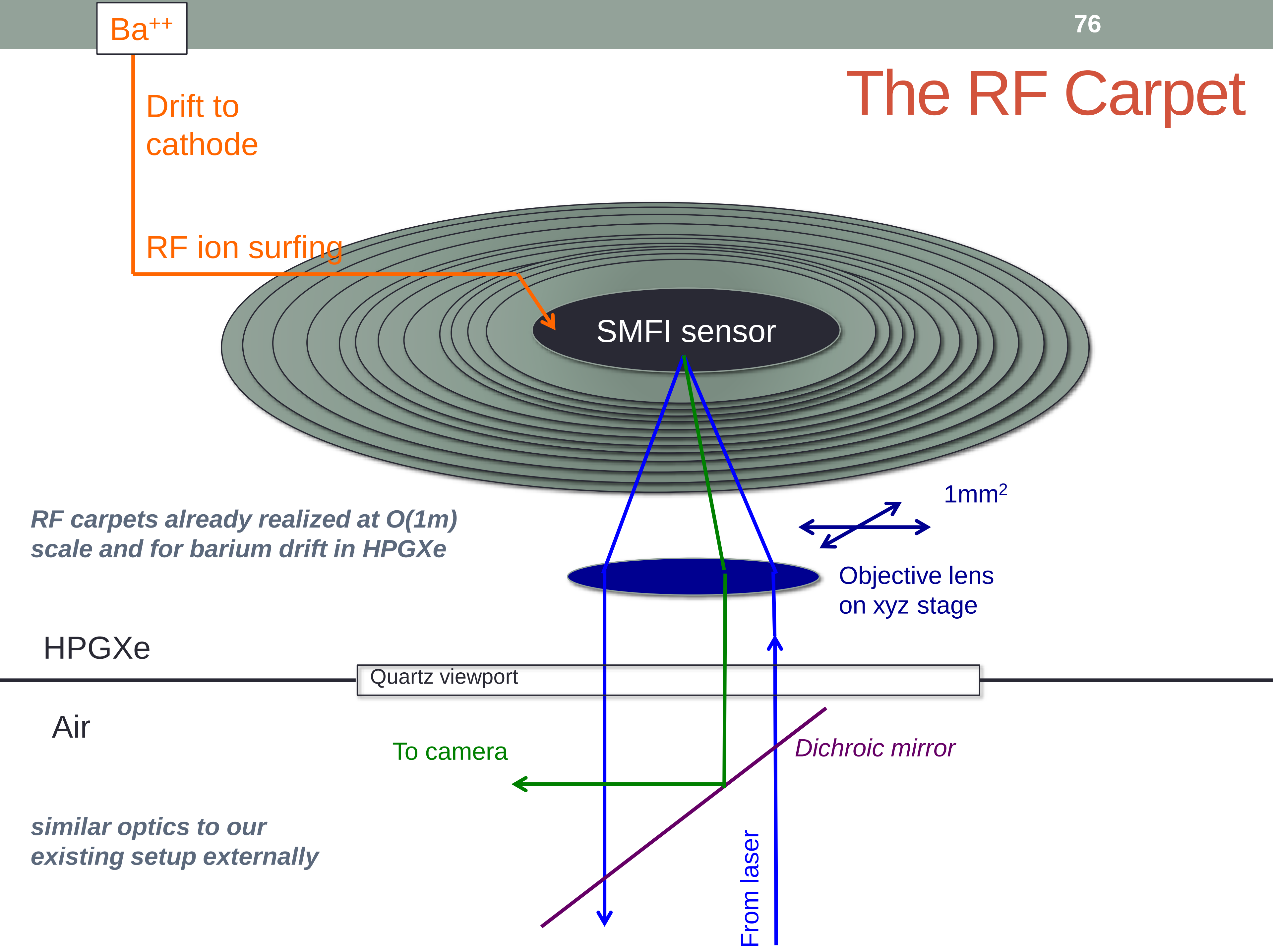


Superposed fields:

RF creates effective potential that levitates ions without neutralization

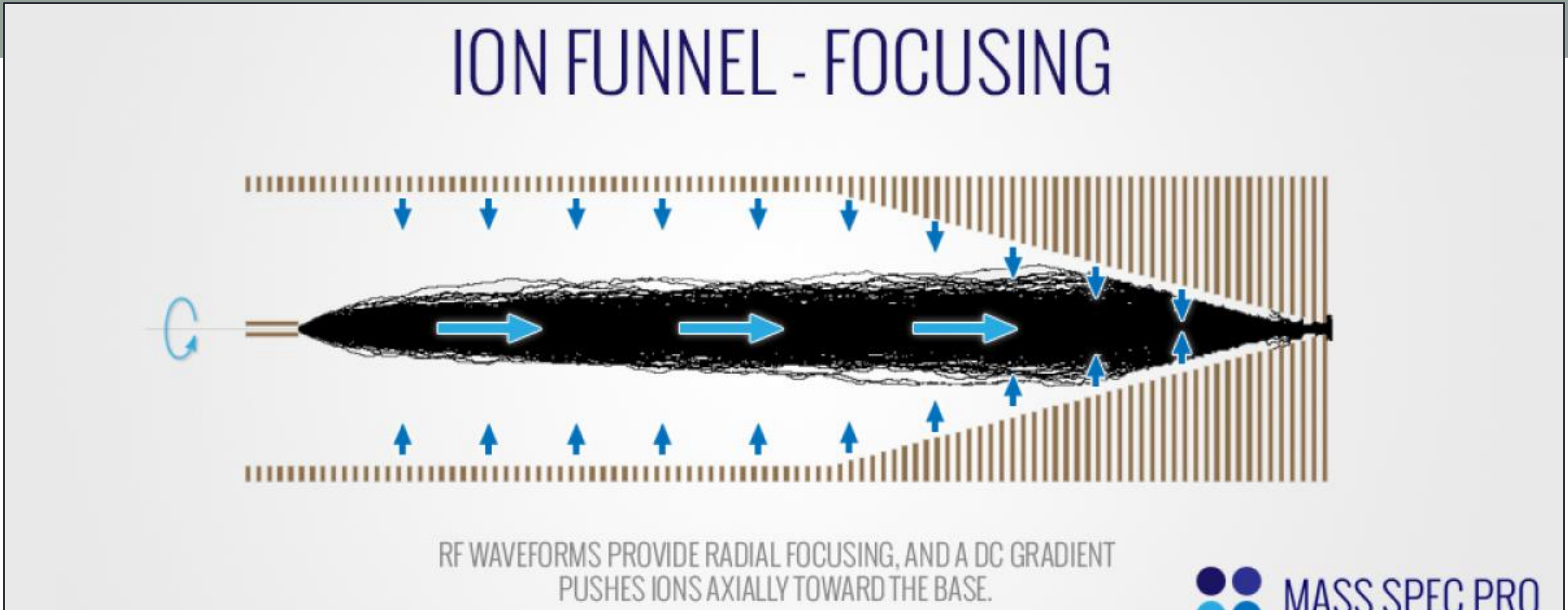
Surfing field sweeps ions to center,

The RF Carpet

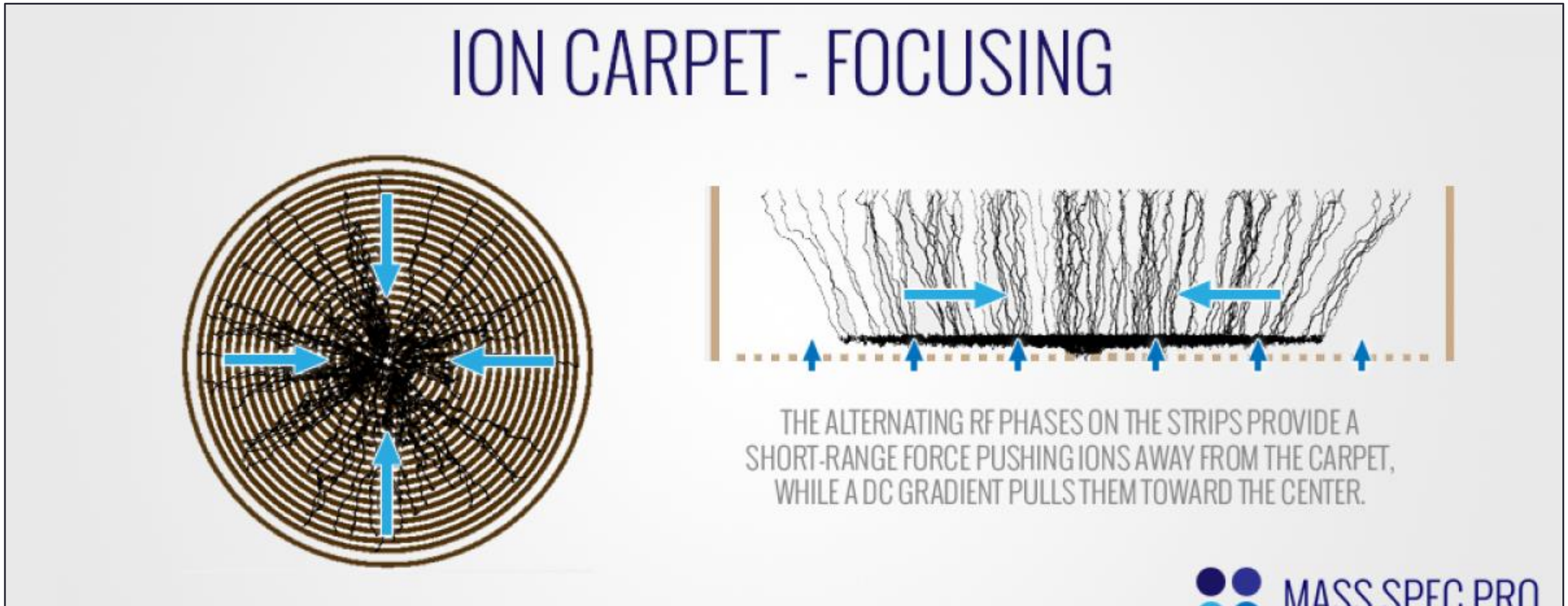


RF carpets already realized at O(1m) scale and for barium drift in HPGXe

similar optics to our existing setup externally

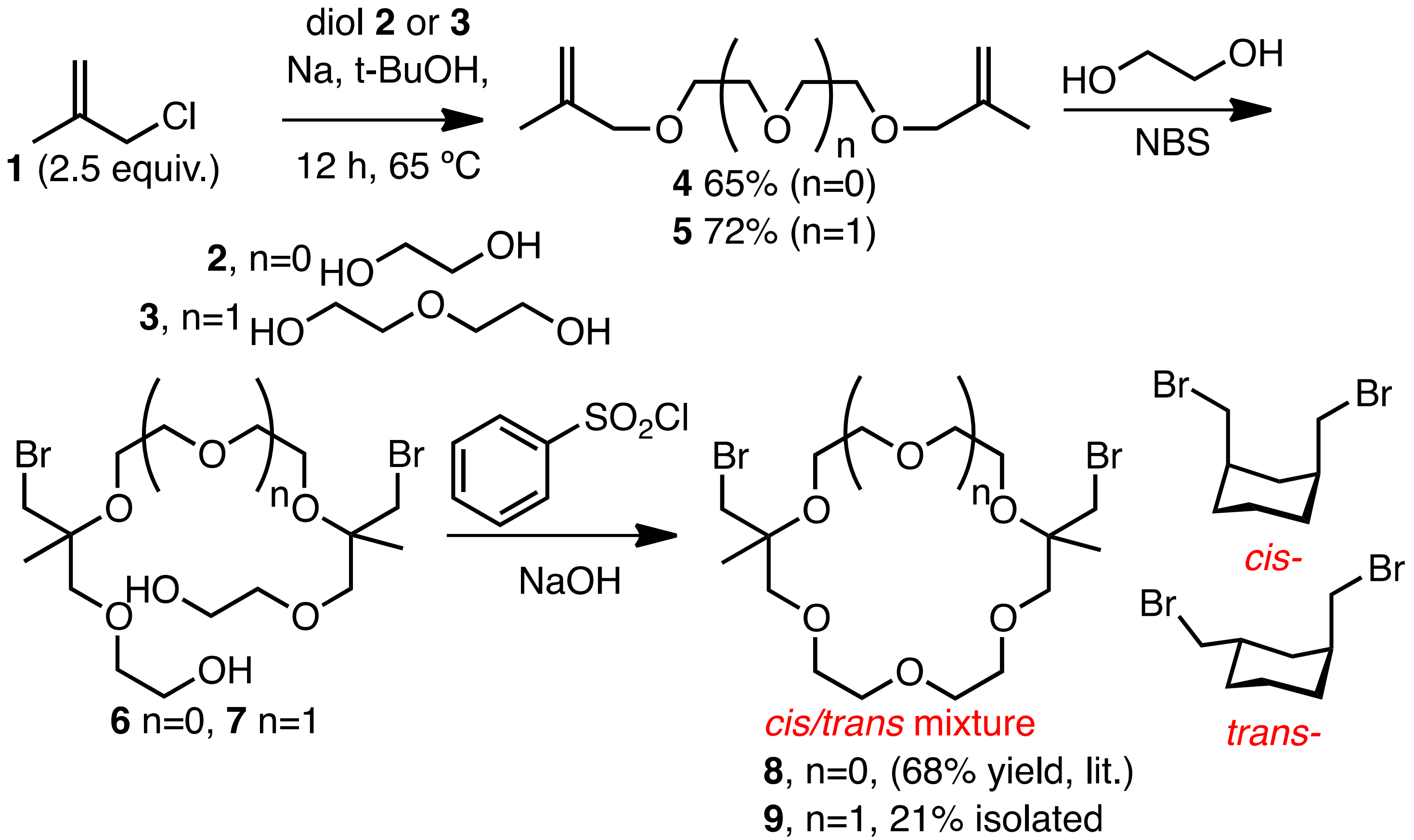


As explored by Brunner / nEXO for extraction to trap

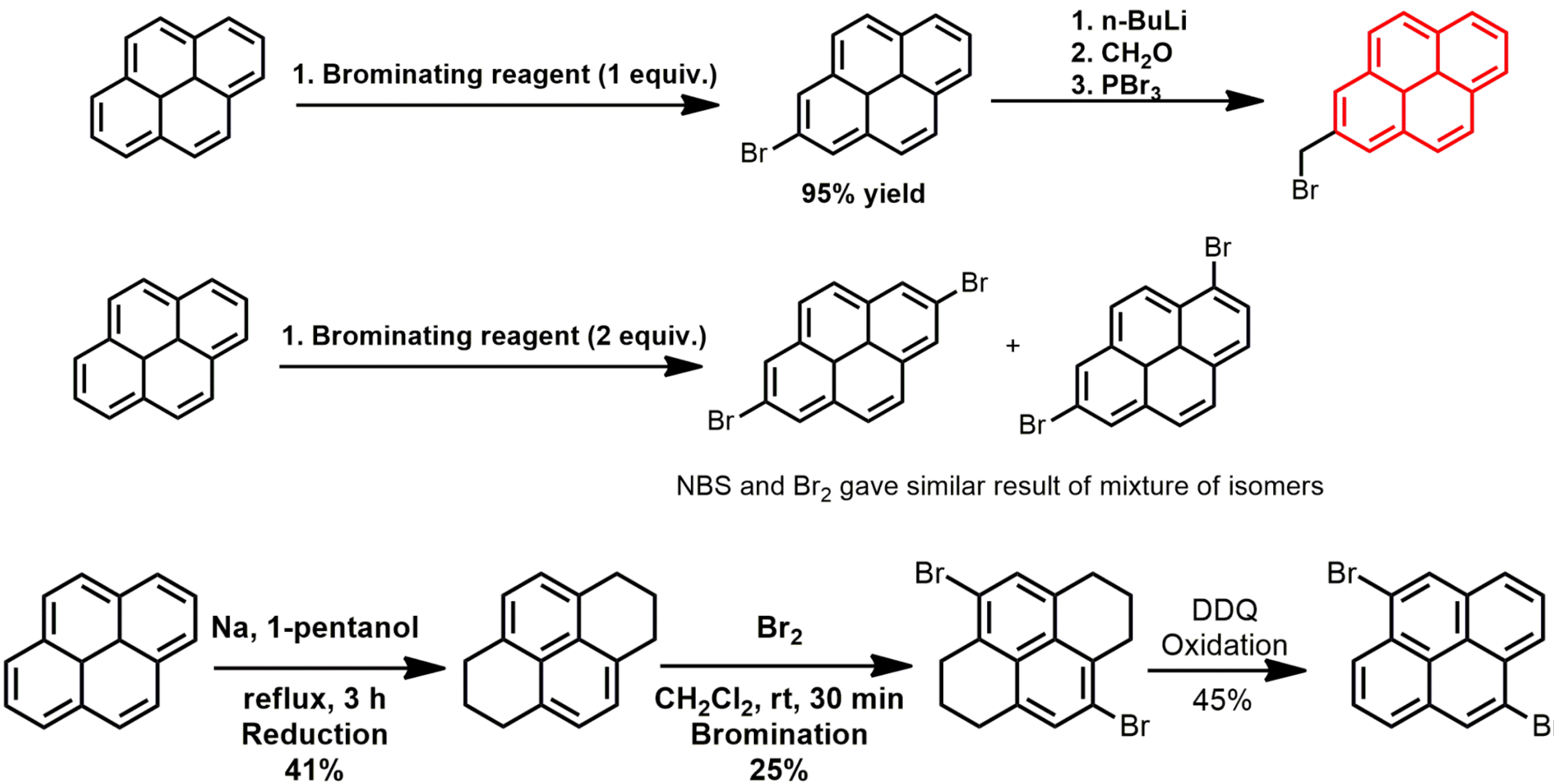


We plan to prototype and demonstrate a small (5cm) RF carpet at UTA

MAC Synthesis: *cis*-crown ether



Orientation of the Pyrene Moiety



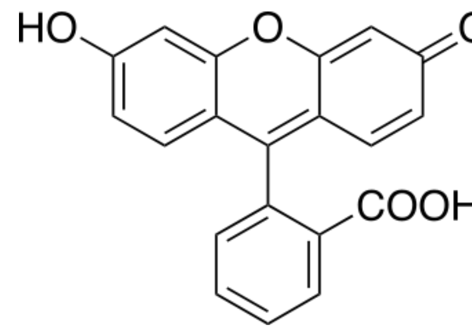
•Nakatsuji, Y.; *Bull. Chem. Soc. Jpn*. **2002**, *75*, 1765-1770.
•*Langmuir*, **2008**, *24*, 5140-5145

168 SINGLE MOLECULE FLUORESCENCE SPECTROSCOPY

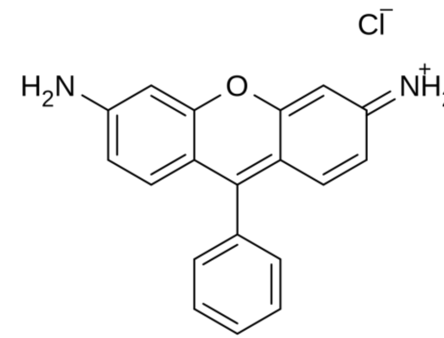
Table 4.1 Photophysical properties of some common dyes with potential for single molecule fluorescence studies

Fluorophore	λ_{ex}^{1p} (nm)	λ_{em}^{1p} (nm)	QY	ε ($cm^{-1} M^{-1}$)	SS (nm)	τ_f (ns)	λ_{ex}^{2p} (nm)	λ_{em}^{2p} (nm)	Reference
FITC	495	520	0.7	73,000	25	—	947	530	[87–89]
FAM	495	520	0.7	83,000	25	—	—	—	[88]
TMR	554	585	0.2–0.5	95,000	31	2.1	849	570	[5,90,91]
R6G	530	556	—	105,000	26	—	—	—	[92]
Cy2	489	506	—	—	17	—	905	520	[87,93]
Cy3	550	570	0.14	150,000	20	~1	1032	578	[3,87,90]
Cy5	650	670	0.15	250,000	20	~1	—	—	[3,5,90]
Cy5.5	675	694	—	250,000	19	—	—	—	[3]
Cy7	743	767	0.02	250,000	24	~0.8	—	—	[3,90,94]
ECFP	458	472	0.4	26,000	14	—	—	—	[94]
EGFP	395,470	509	0.8	30,000	39	3.2	—	—	[5,90,94]
EYFP	514	527	0.6	84,000	13	3.7	—	—	[90,94]
DsRed	532	582	0.29	22,500	50	2.8	—	—	[90]
Bodipy Fl	504	510	—	70,000	6	—	920	526	[87,88]
Bodipy R6G	528	547	—	70,000	19	—	—	—	[88]
AF488	495	520	0.5–0.9	80,000	25	—	985	530	[5,87]
AF546	554	570	—	112,000	16	—	1028	582	[87,95]
AF555	555	565	—	150,000	10	—	—	—	[3]
AF594	590	617	—	92,000	27	—	1074	619	[87,95]
AF633	632	647	—	100,000	15	3.2	—	—	[3,90]

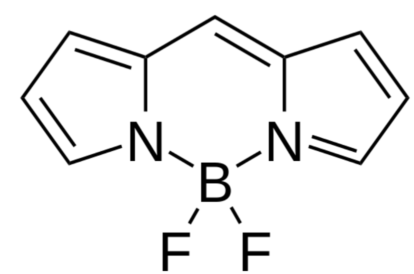
Handbook of Single Molecule Fluorescence, Gell, C. Brockwell, D. Smith, A. OUP Oxford, 2006, New York p. 168-9.



Fluoroscein: FITC + FAM

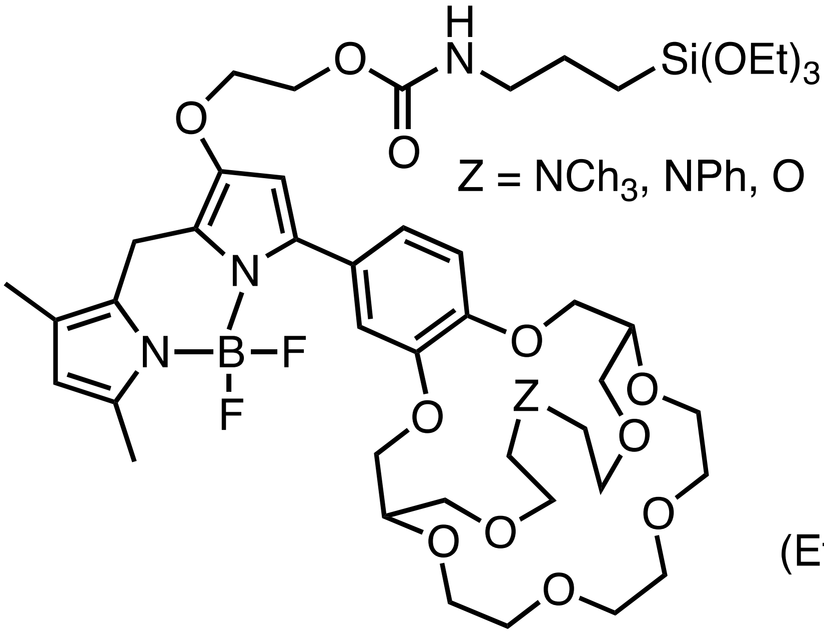


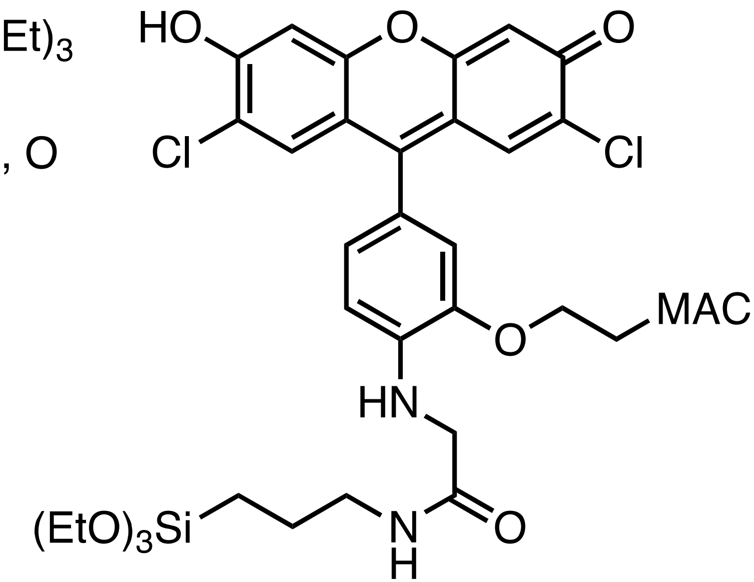
Rhodamine: TMR, R6G

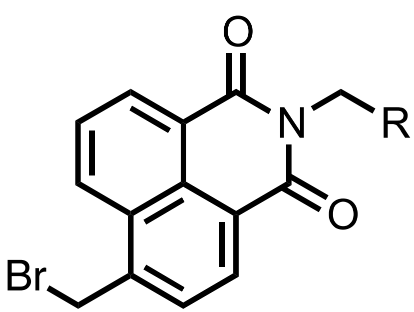


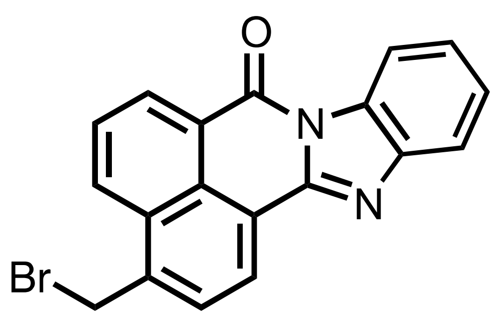
BODIPY

Alternative Fluorophores









Ando, Hiruta, Citterio, Suzuki
Analyst, 2009, 134, 2314-2319

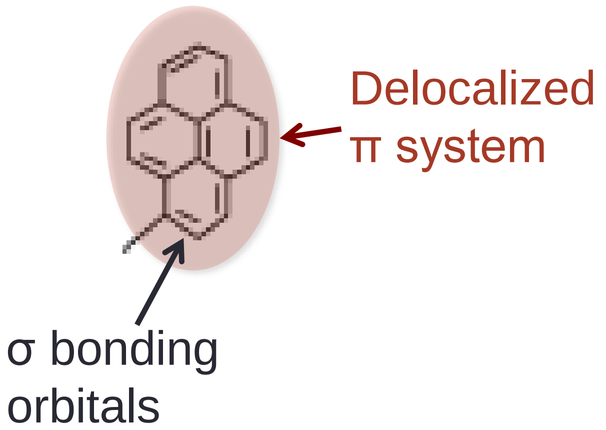
Hiruta, Sato, Takahashi, Kubobuchi, Shichi, Citterio, Suzuki
RSC Adv. 2013, 3, 6499-6506

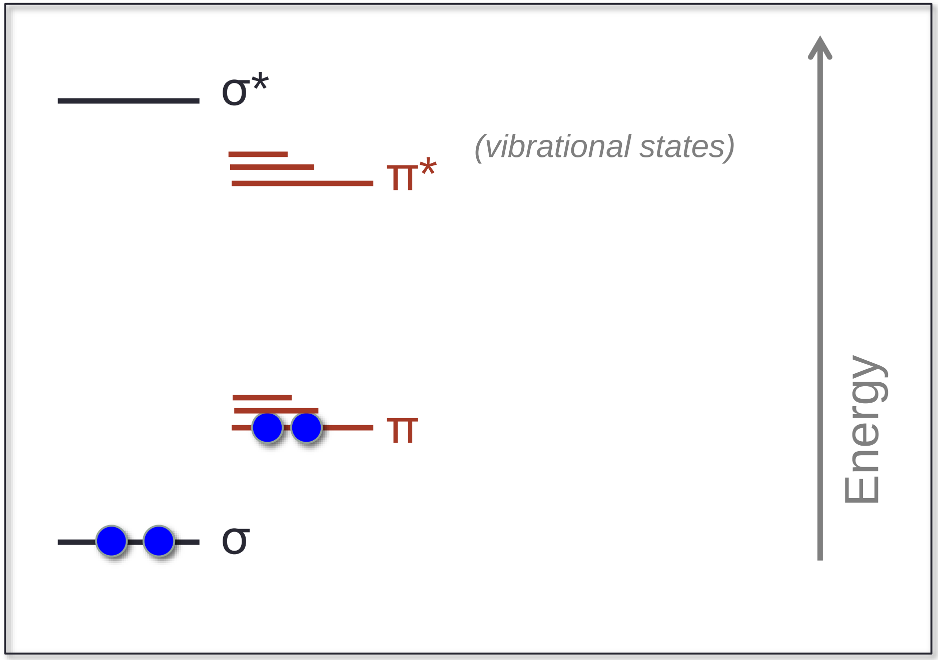
The basic SMFI method:

- A non-fluorescent molecule becomes fluorescent (or vice versa) upon chelation with an incident ion.

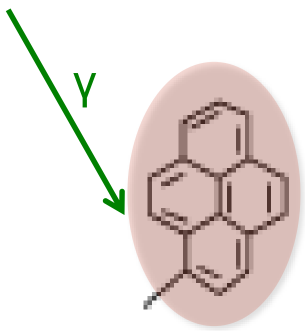
- Calcium and barium are congeners – many dyes developed for calcium are also expected to respond to barium

Fluorescence in pyrene:

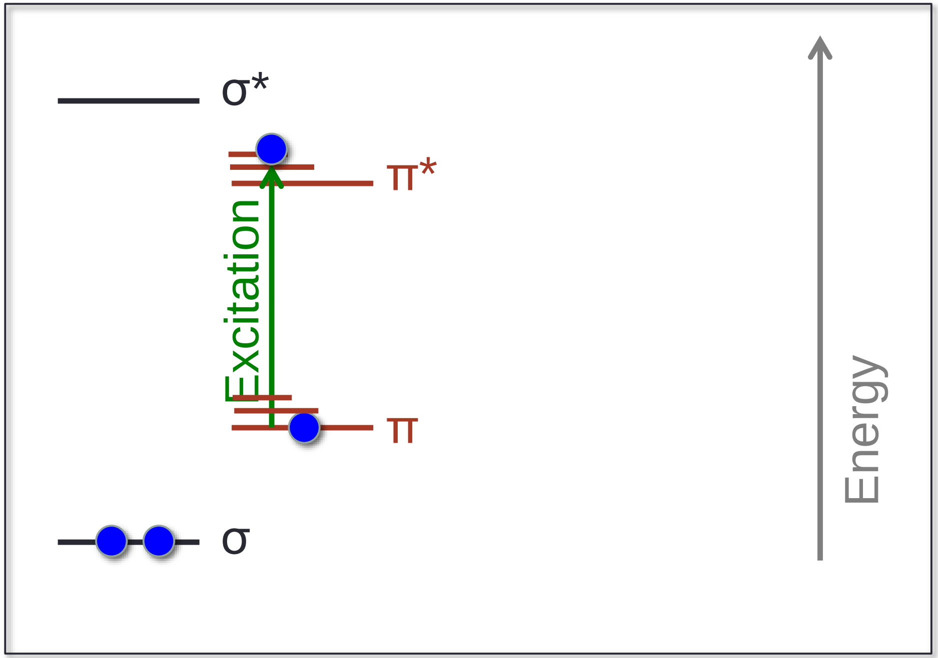




Fluorescence in pyrene:

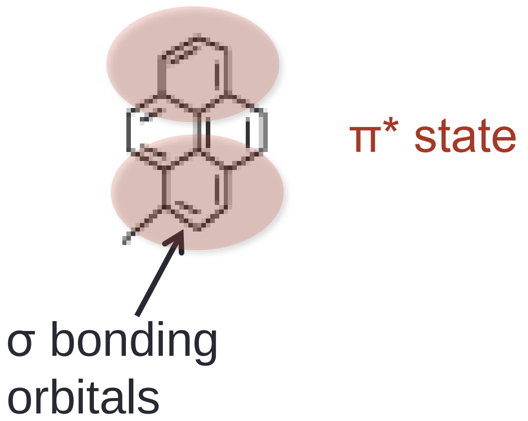


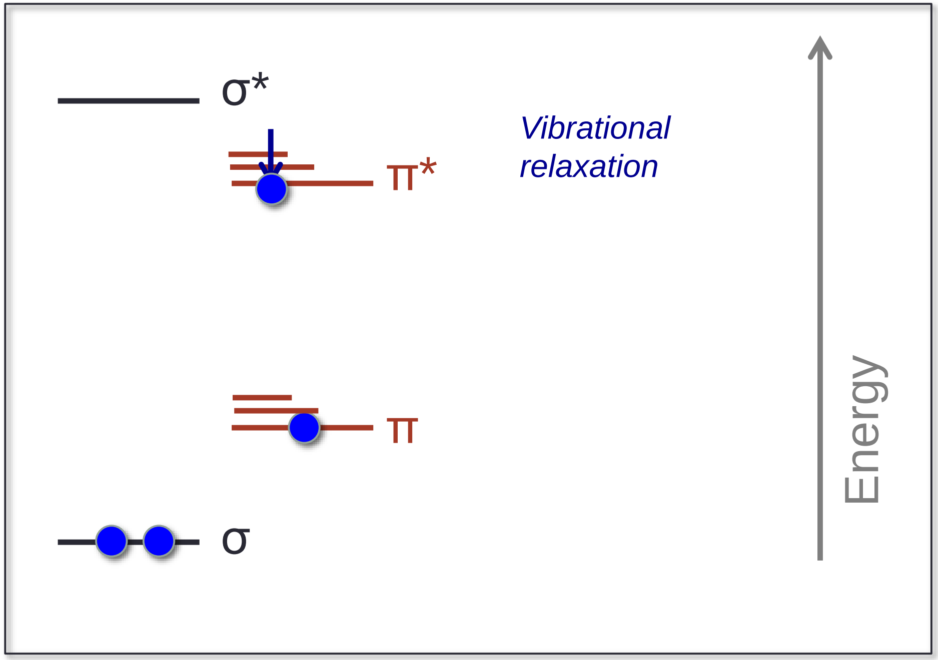
Excitation photon



Fluoresence process

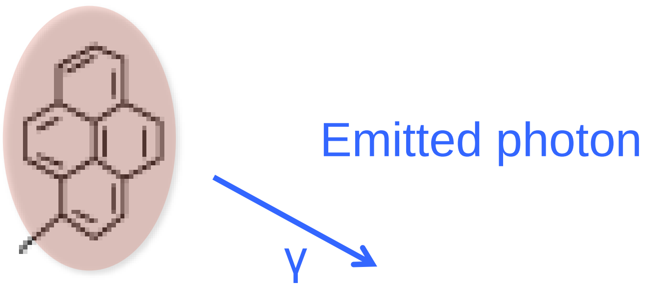
Fluorescence in pyrene:

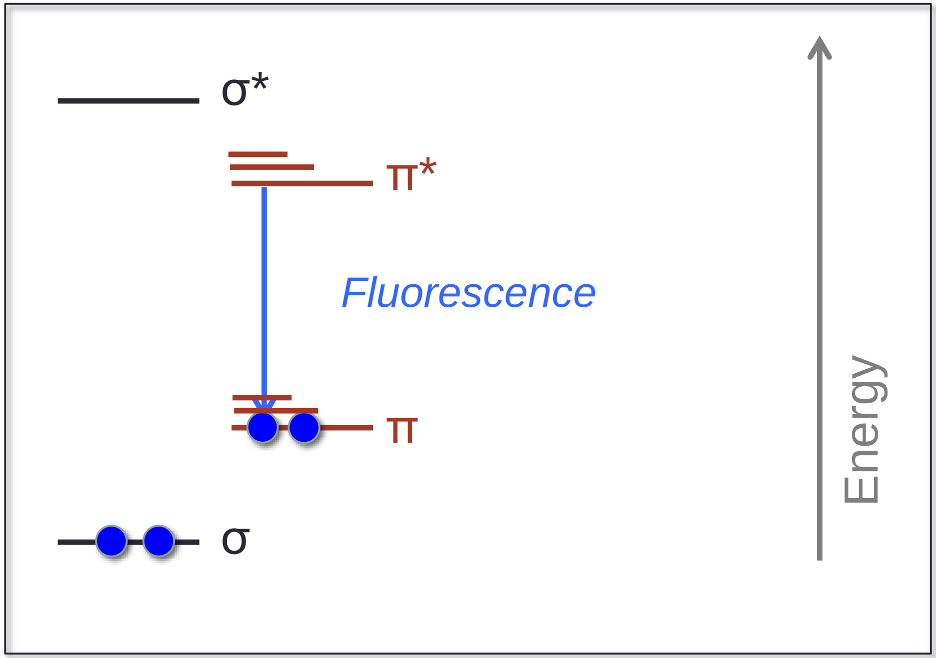




Fluoresence process

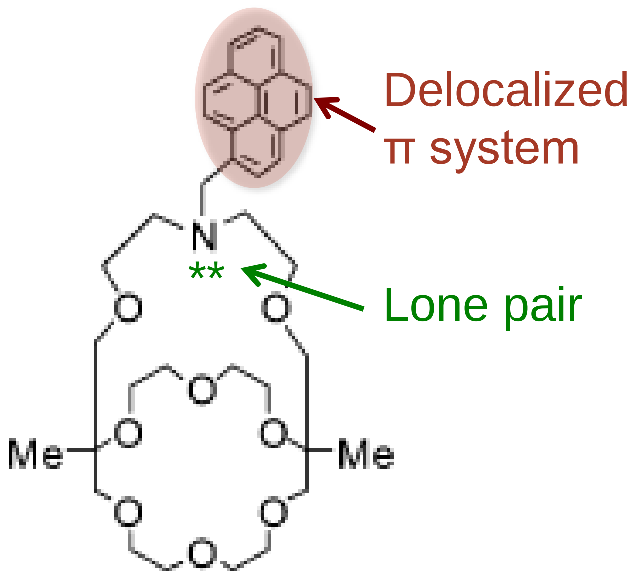
Fluorescence in pyrene:

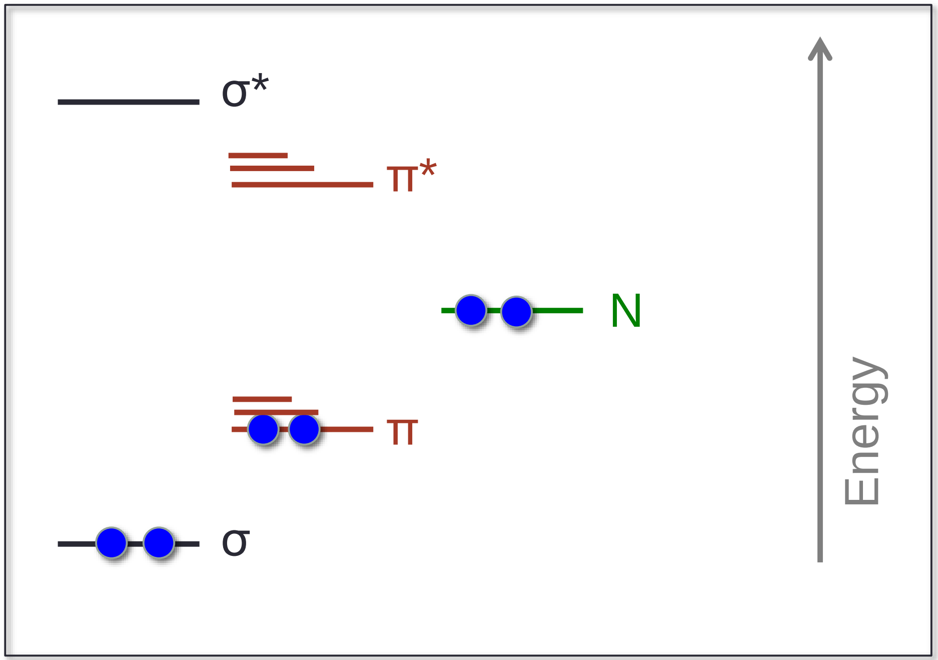




Fluoresence process

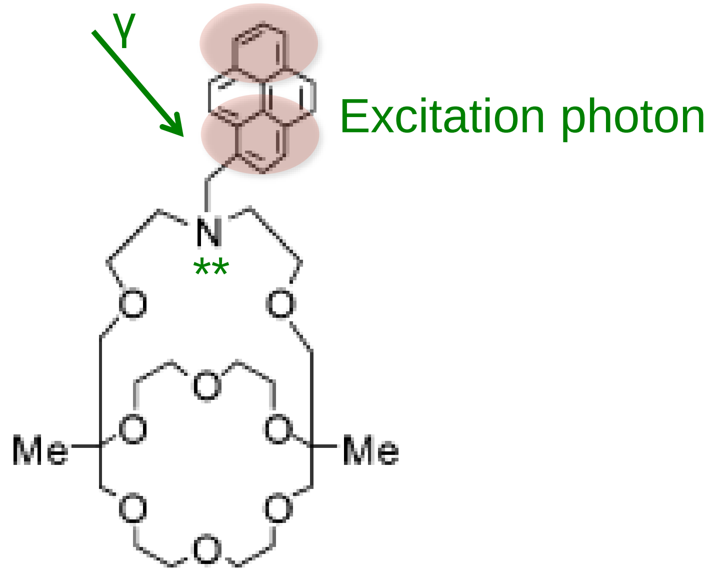
The PSMA magic:

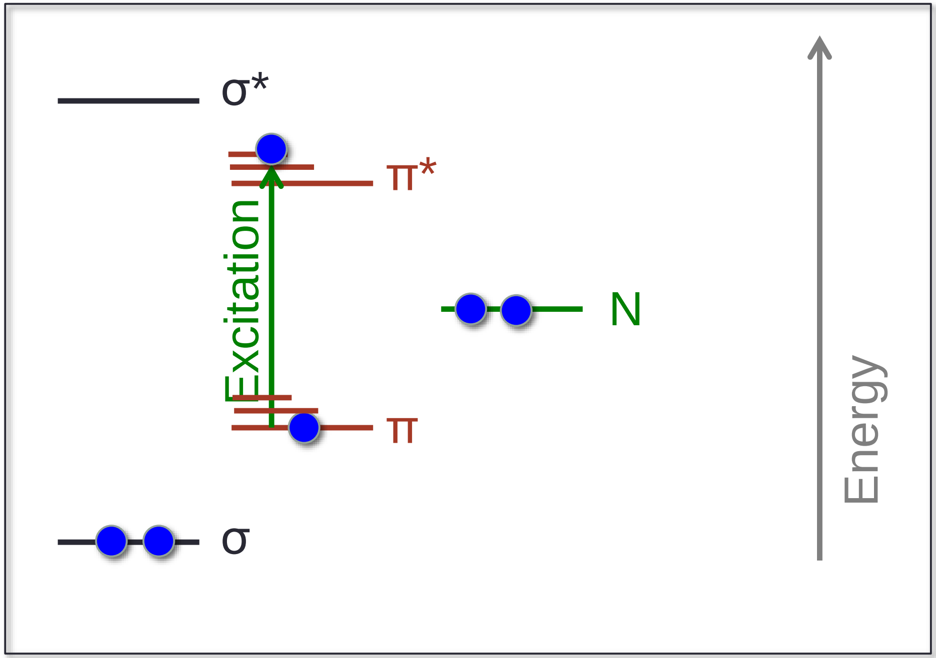




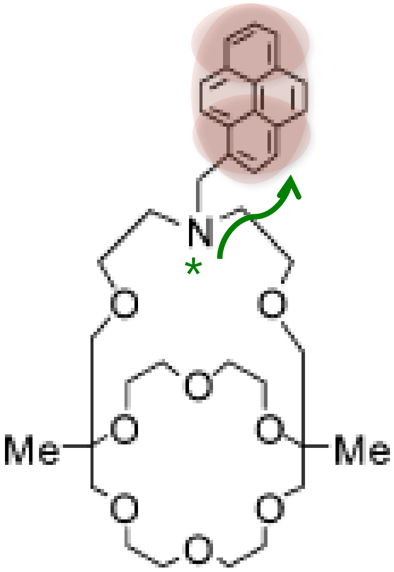
The key to the tagging behaviour of PSMA is a carefully placed lone pair of electrons on the nitrogen atom.

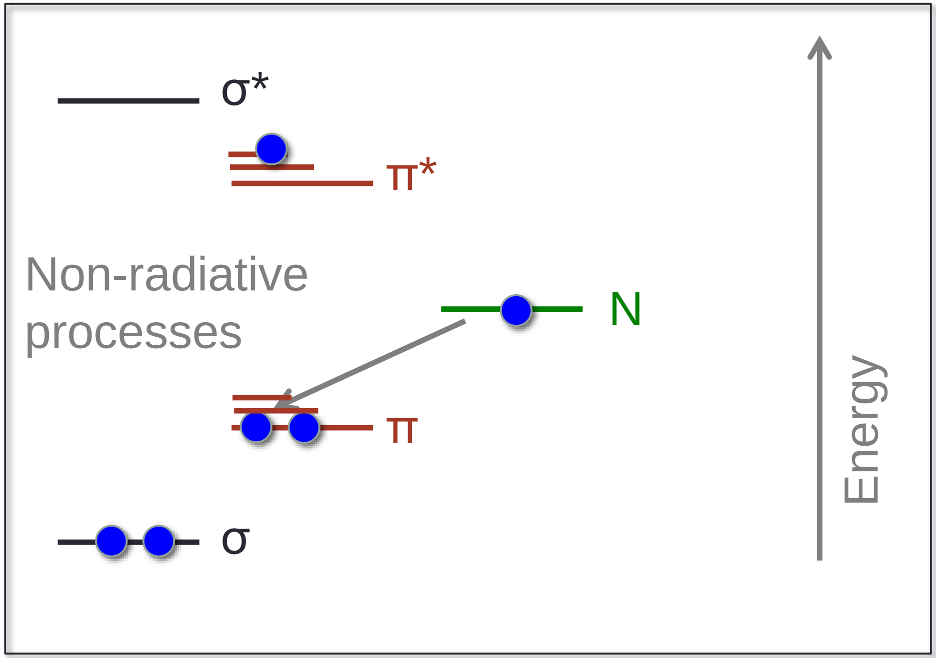
The PSMA magic:



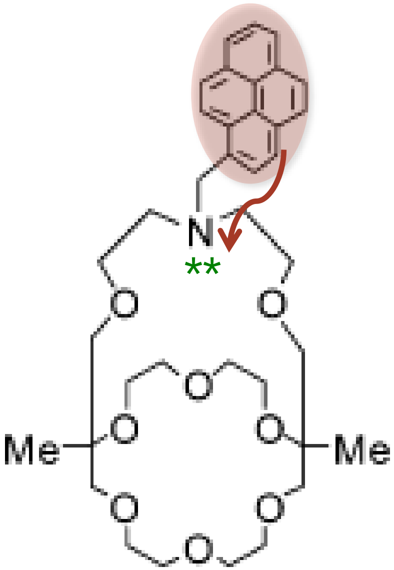


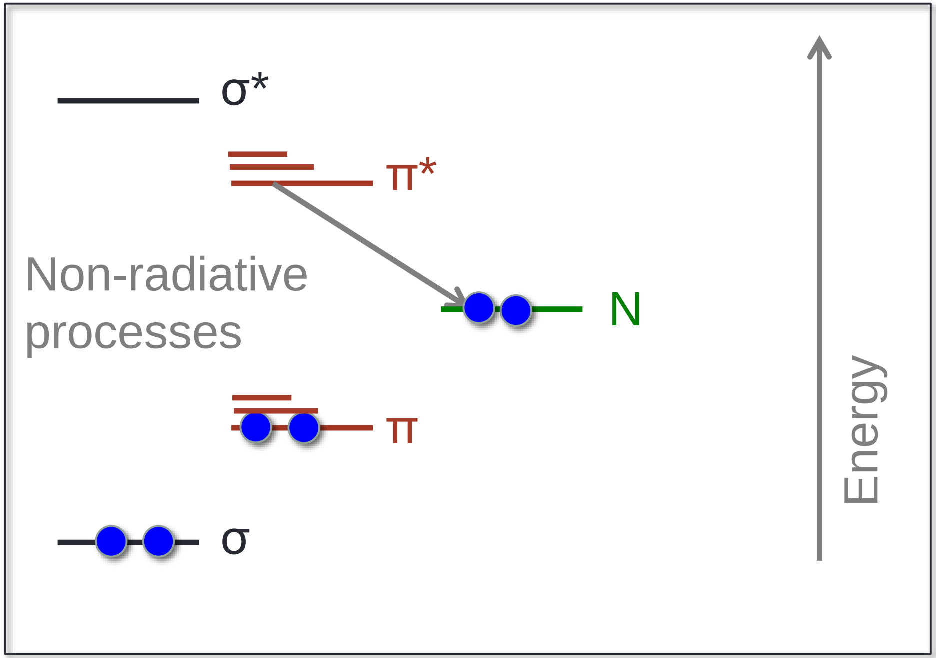
The PSMA magic:





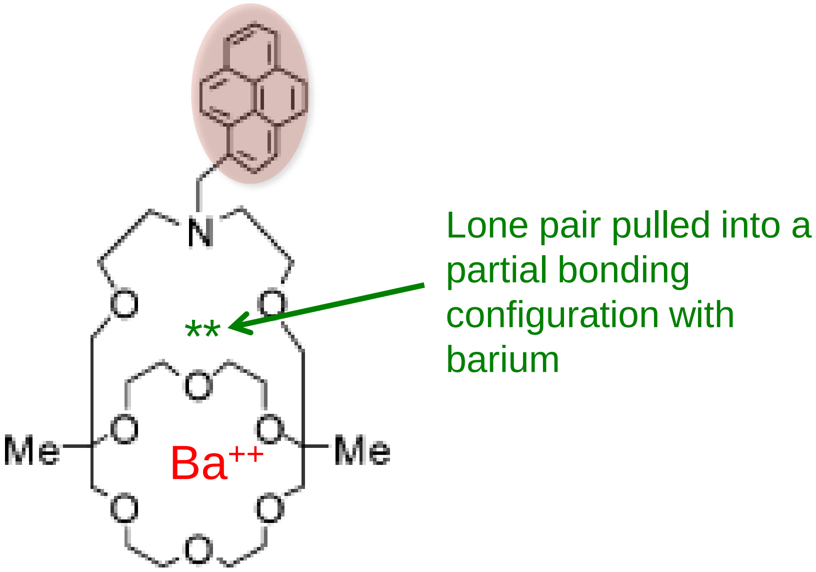
The PSMA magic:

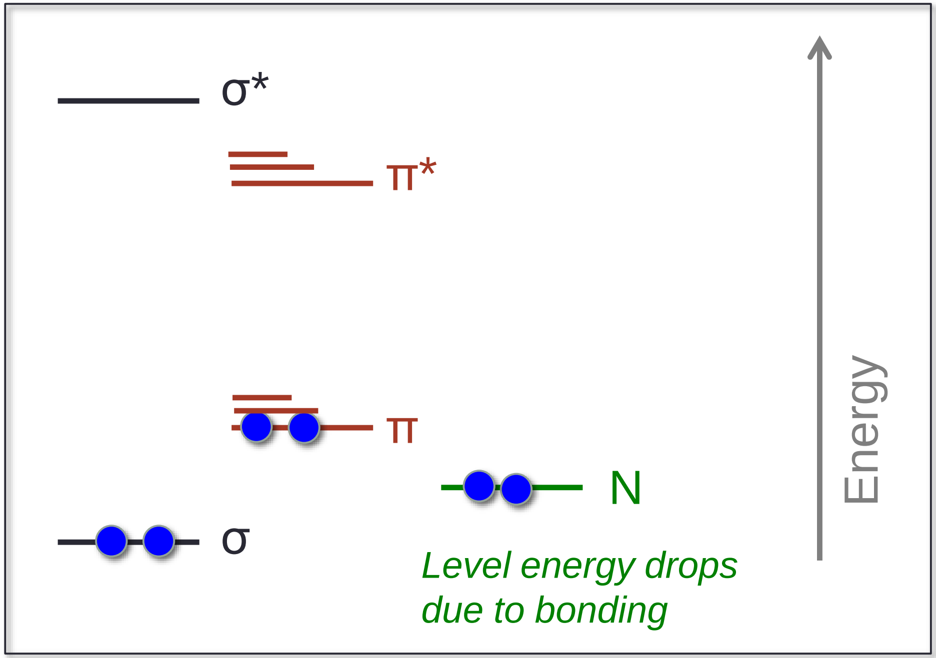




Back in the ground state with no fluorescence!

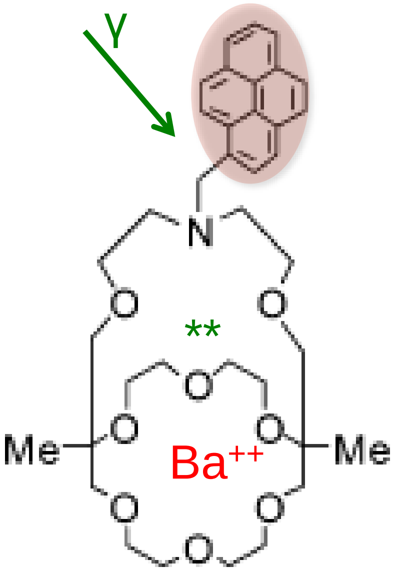
The PSMA magic:

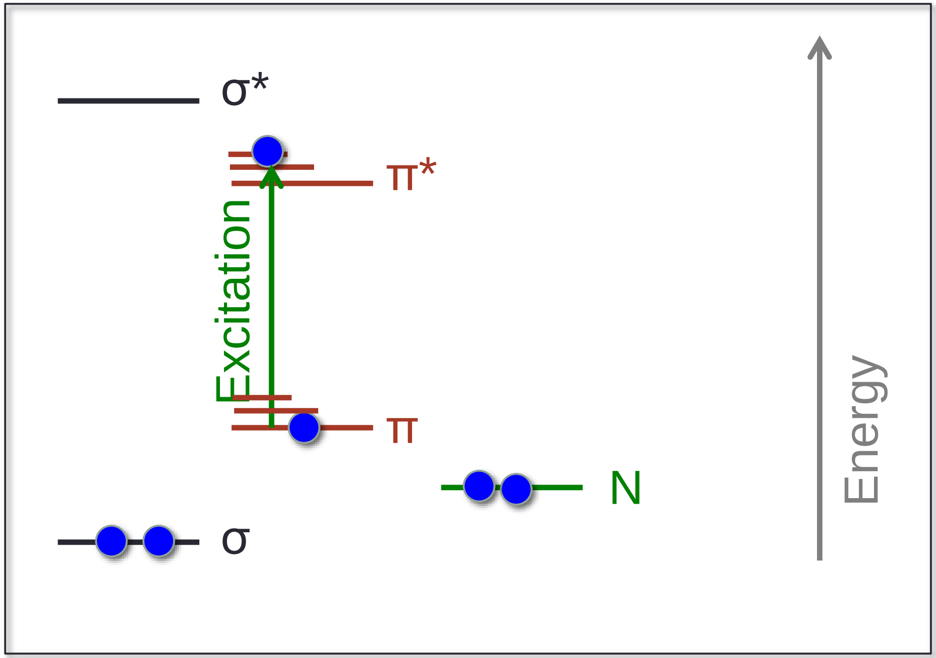




Now when there is a barium dication in the cage…

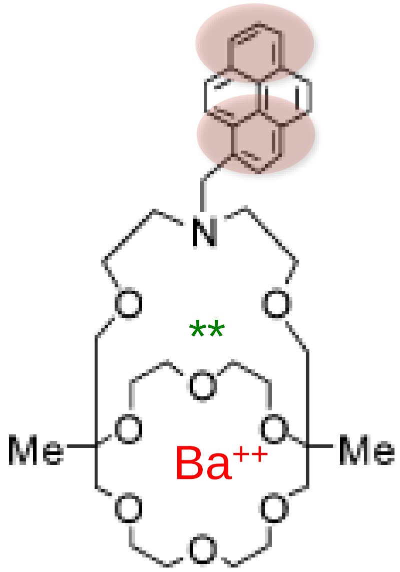
The PSMA magic:

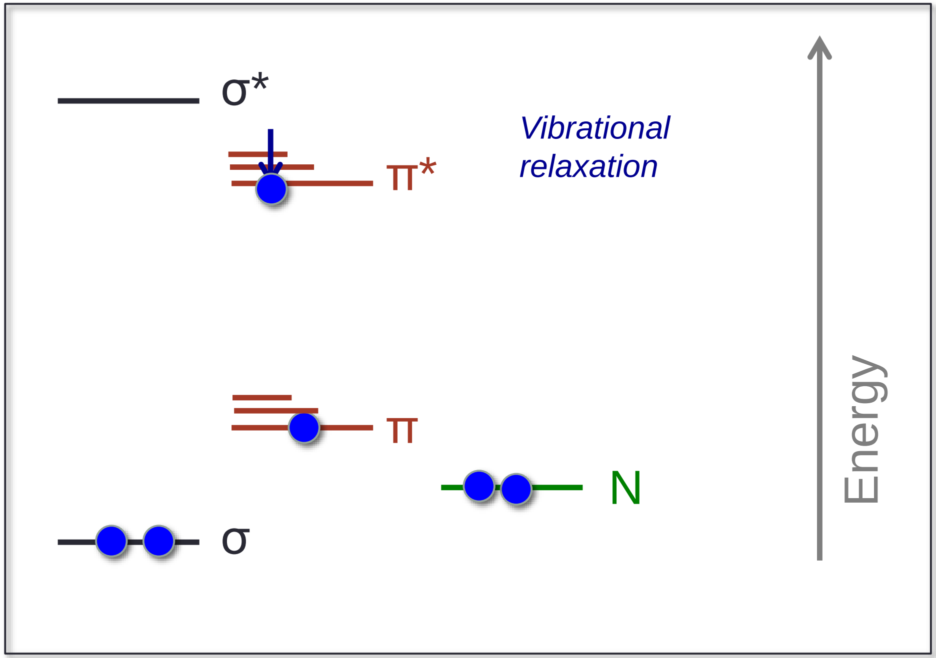




Now when there is a barium dication in the cage…

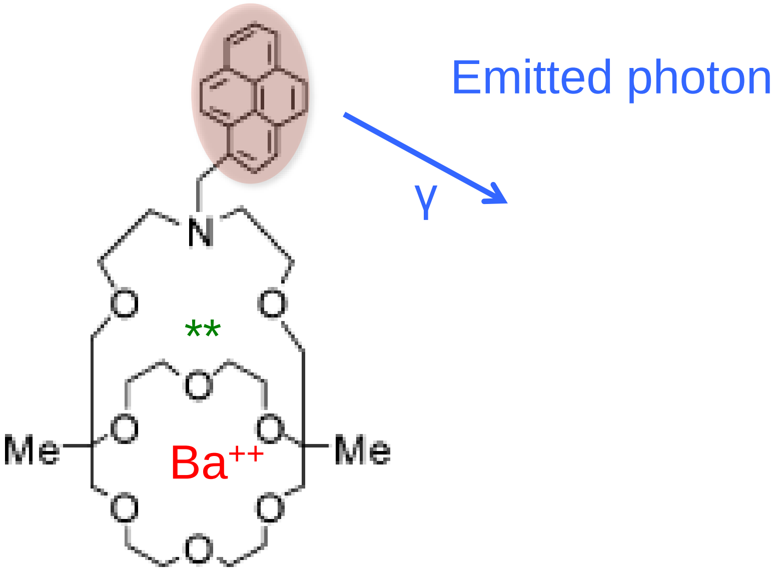
The PSMA magic:

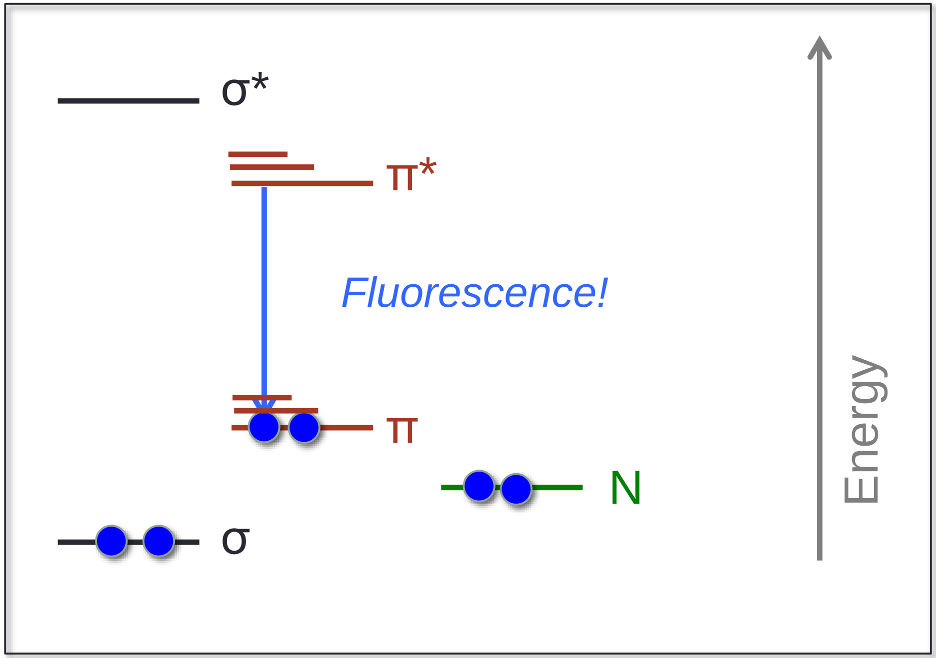




Now when there is a barium dication in the cage…

The PSMA magic:





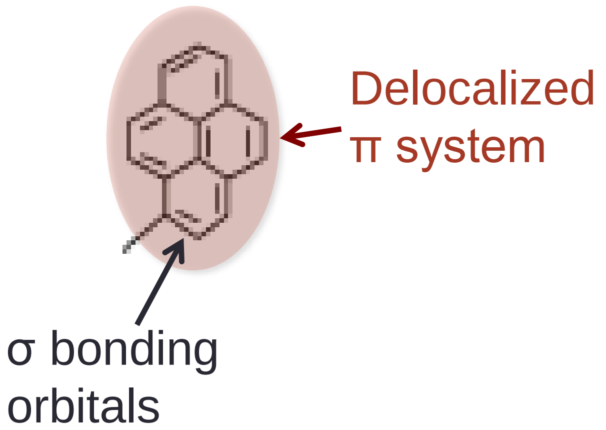
Fluorescence is switched on
only when barium ion is present

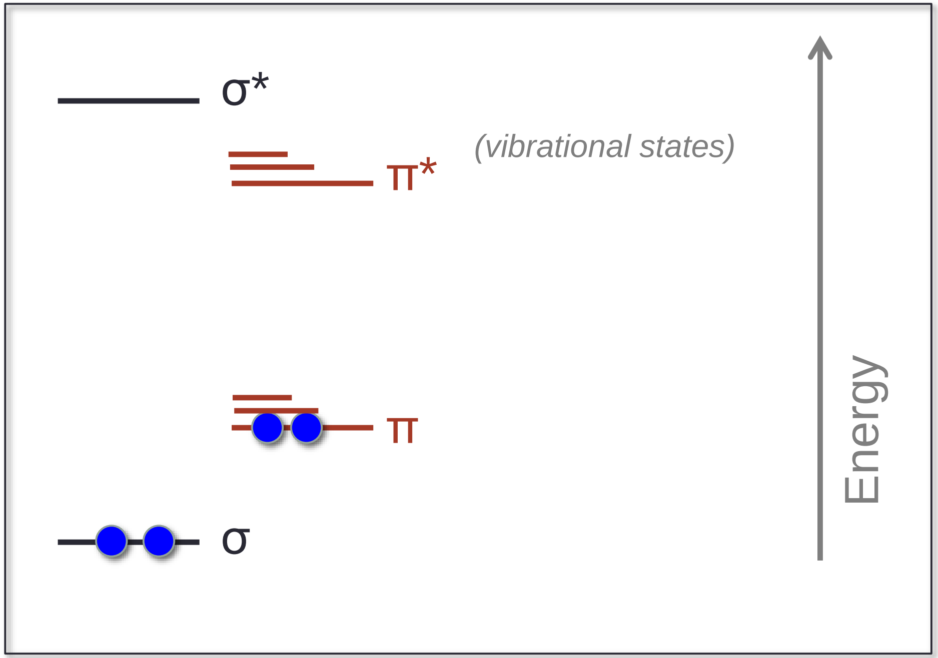
The basic SMFI method:

- A non-fluorescent molecule becomes fluorescent (or vice versa) upon chelation with an incident ion.

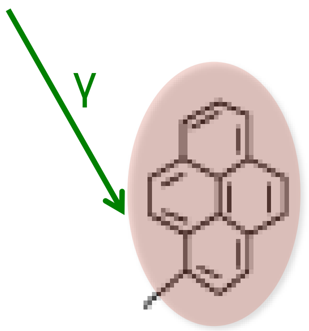
- Calcium and barium are congeners – many dyes developed for calcium are also expected to respond to barium

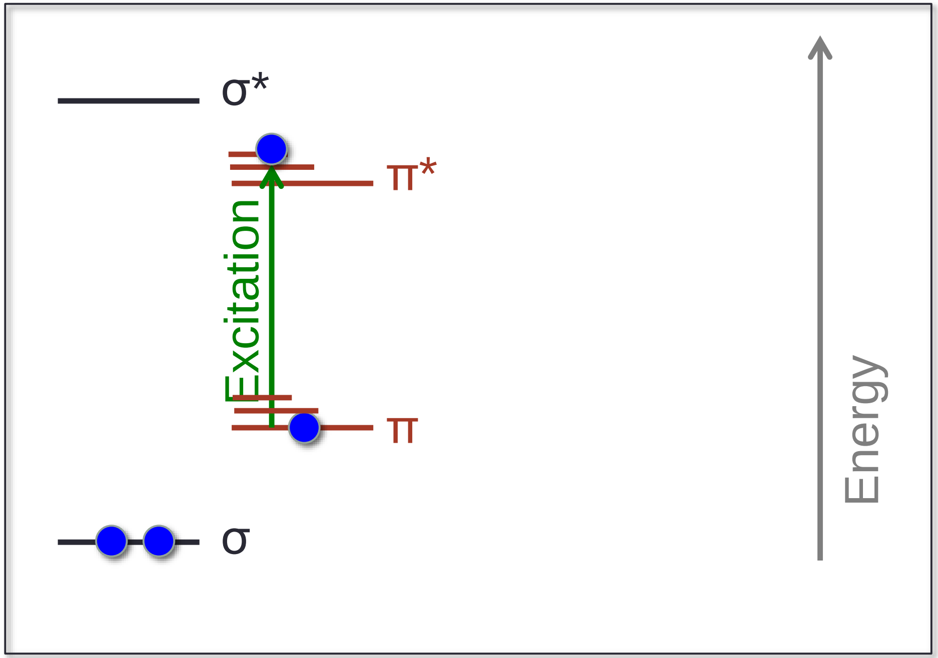
Fluorescence in pyrene:





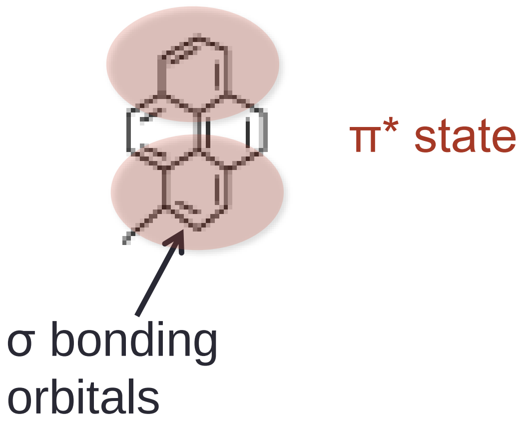
Fluorescence in pyrene:

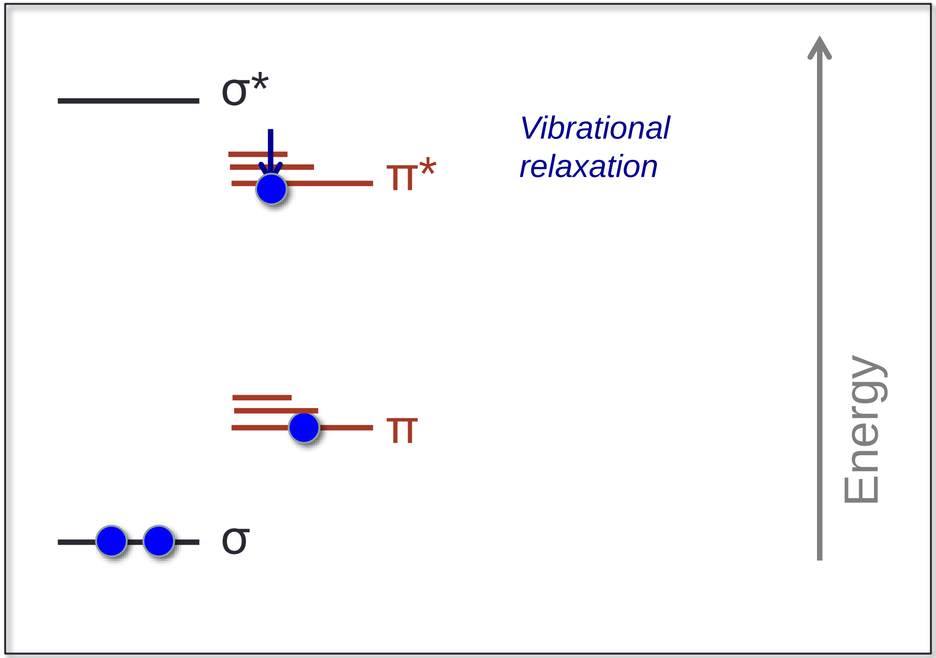




Fluoresence process

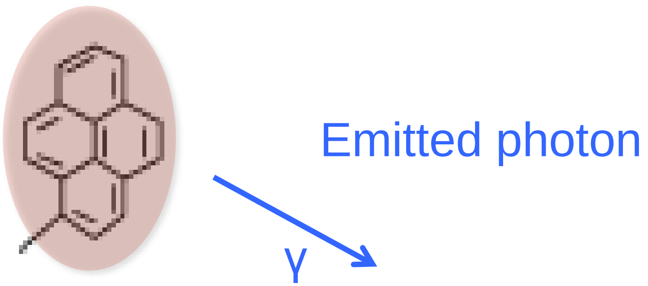
Fluorescence in pyrene:

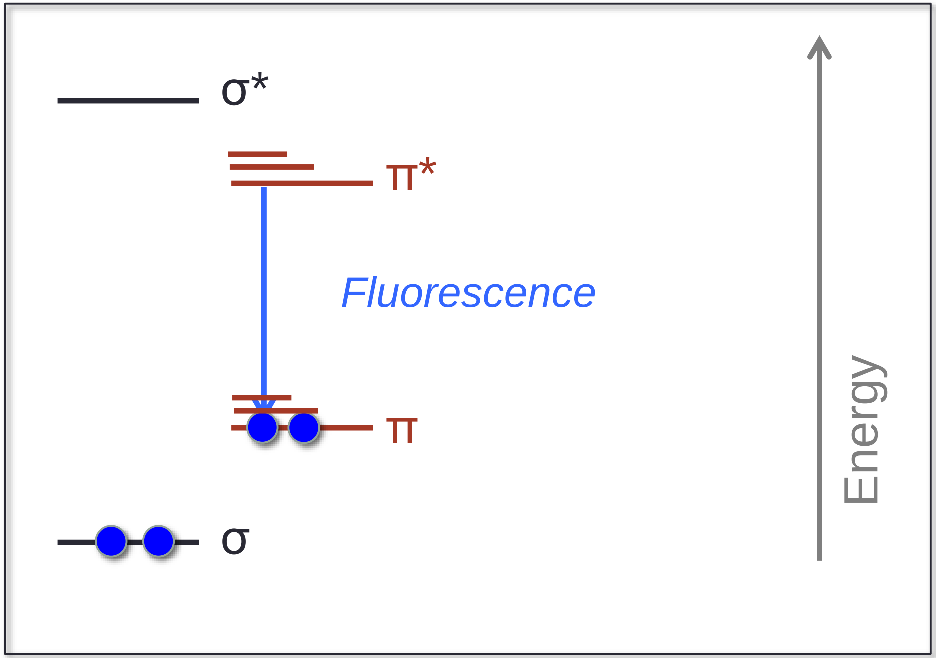




Fluoresence process

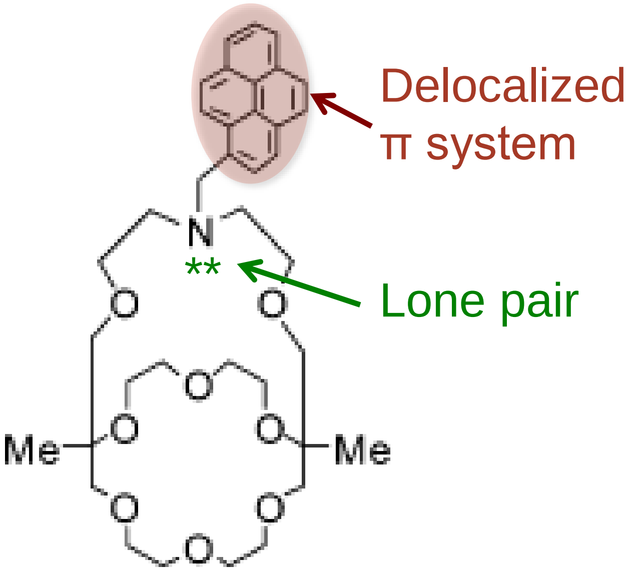
Fluorescence in pyrene:

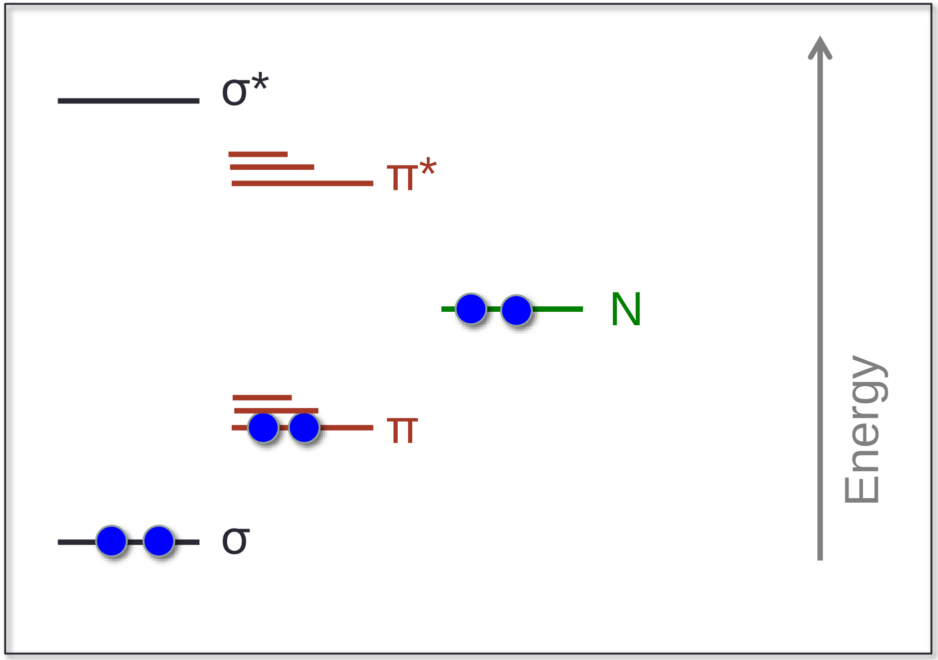




Fluoresence process

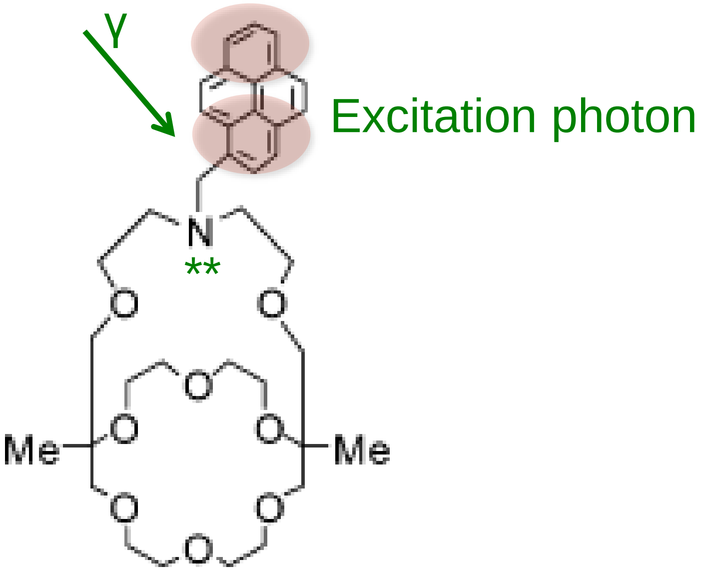
The PSMA magic:

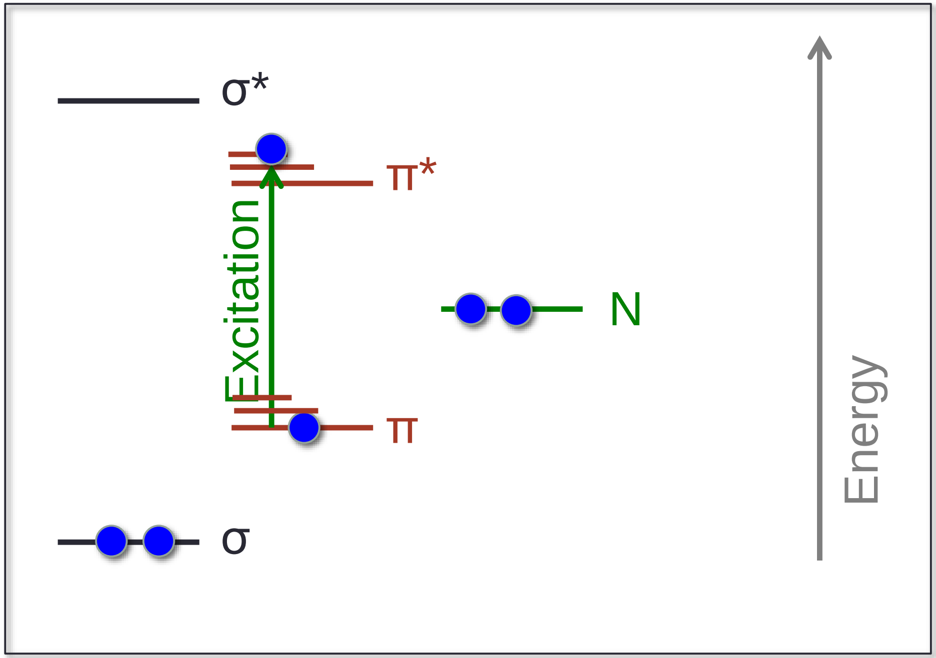




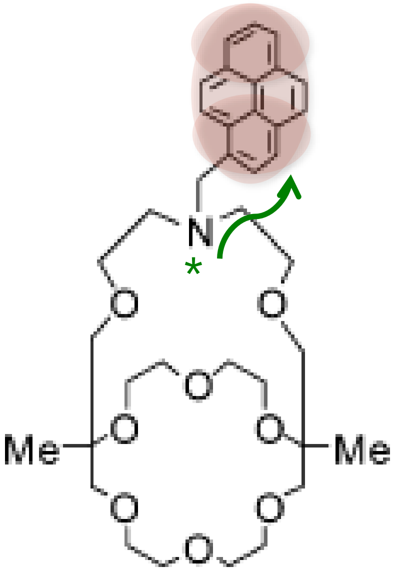
The key to the tagging behaviour of PSMA is a carefully placed lone pair of electrons on the nitrogen atom.

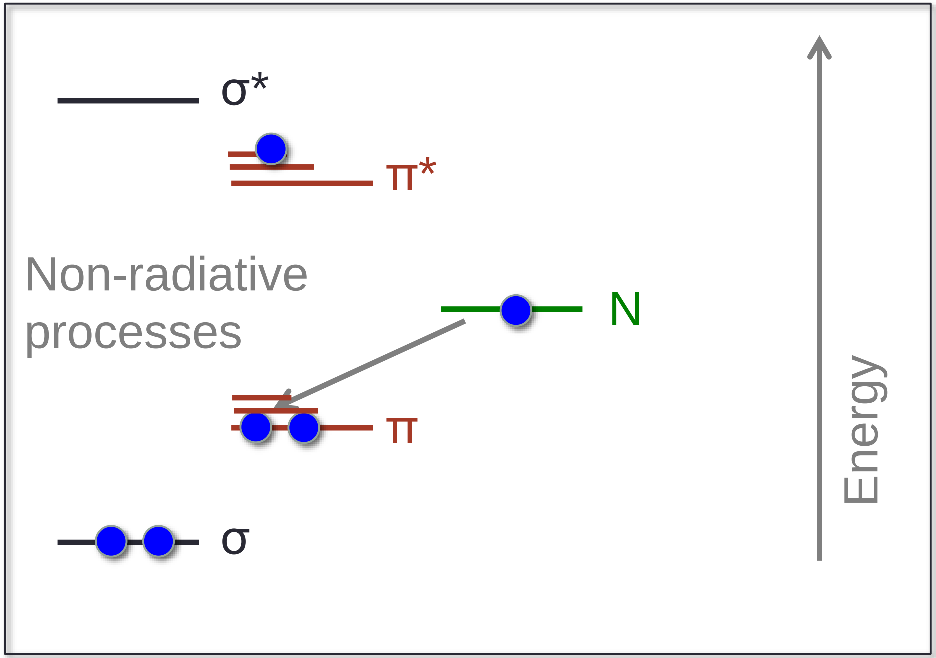
The PSMA magic:



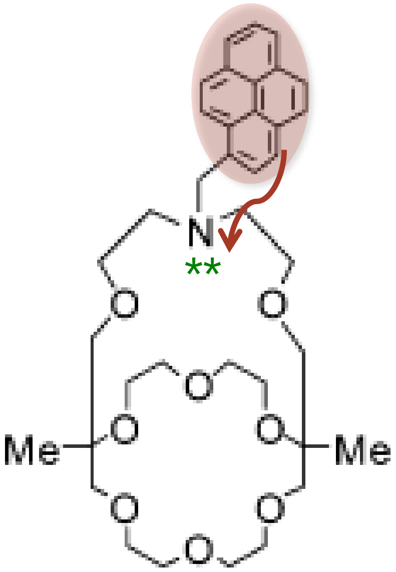


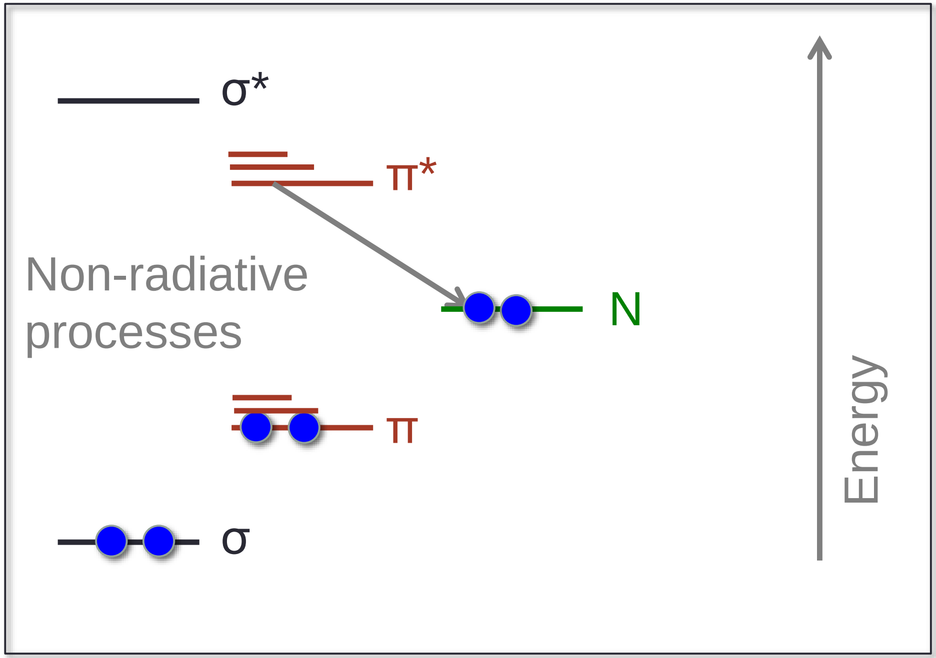
The PSMA magic:





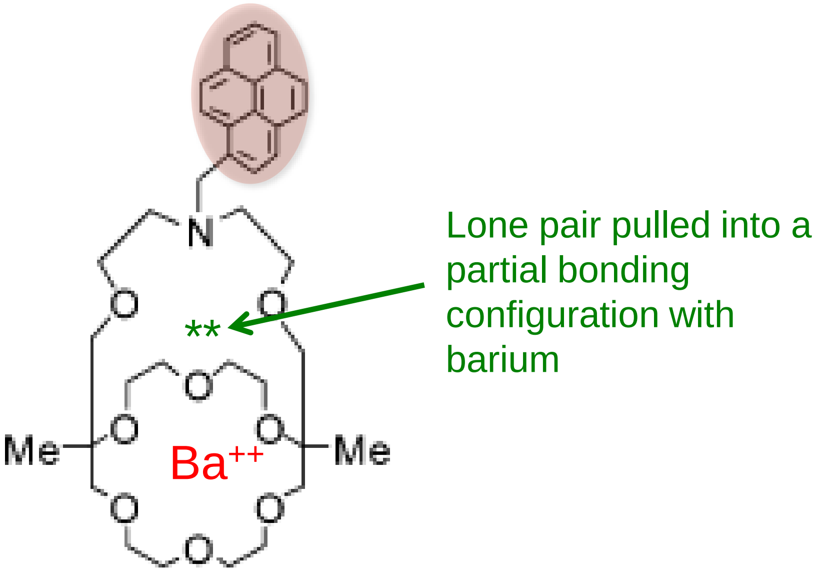
The PSMA magic:

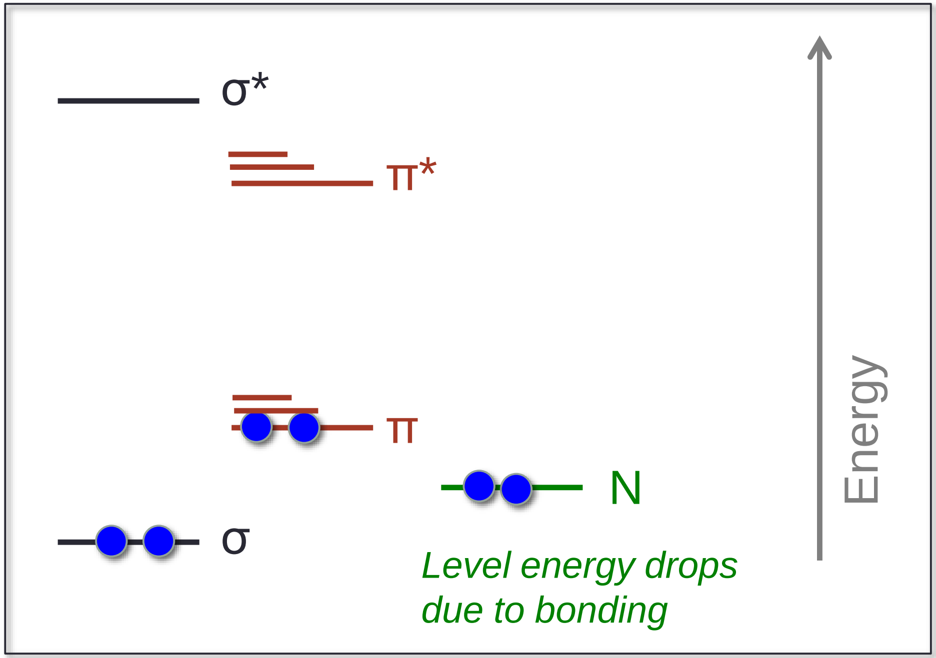




Back in the ground state with no fluorescence!

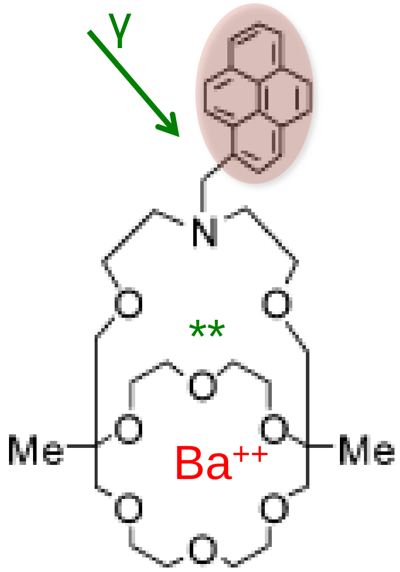
The PSMA magic:

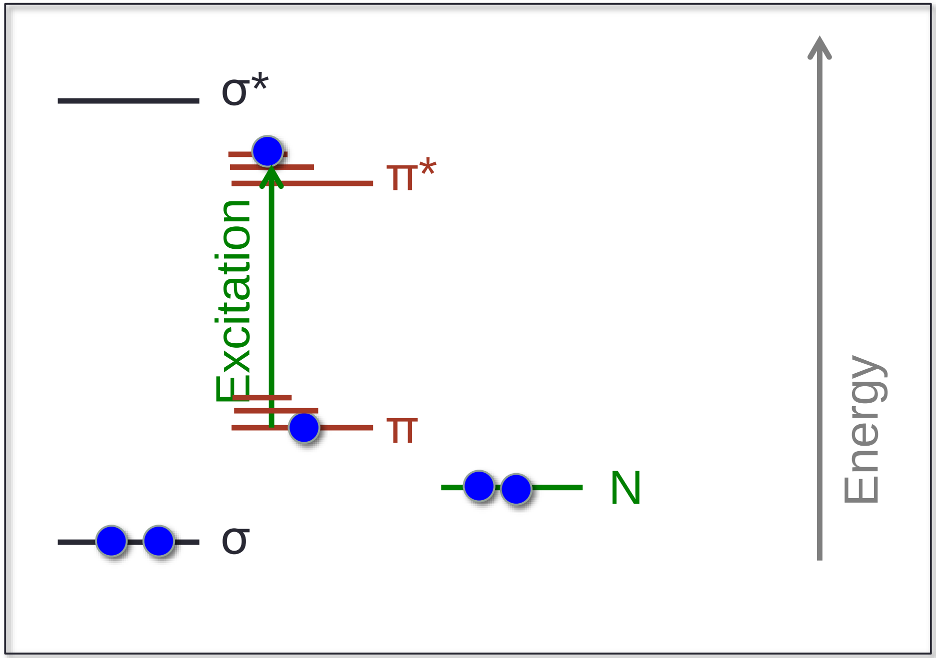




Now when there is a barium dication in the cage…

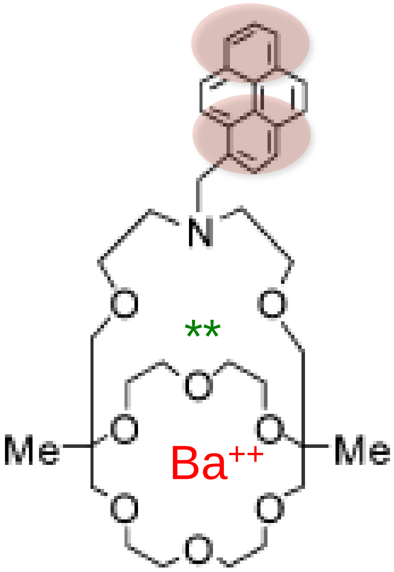
The PSMA magic:

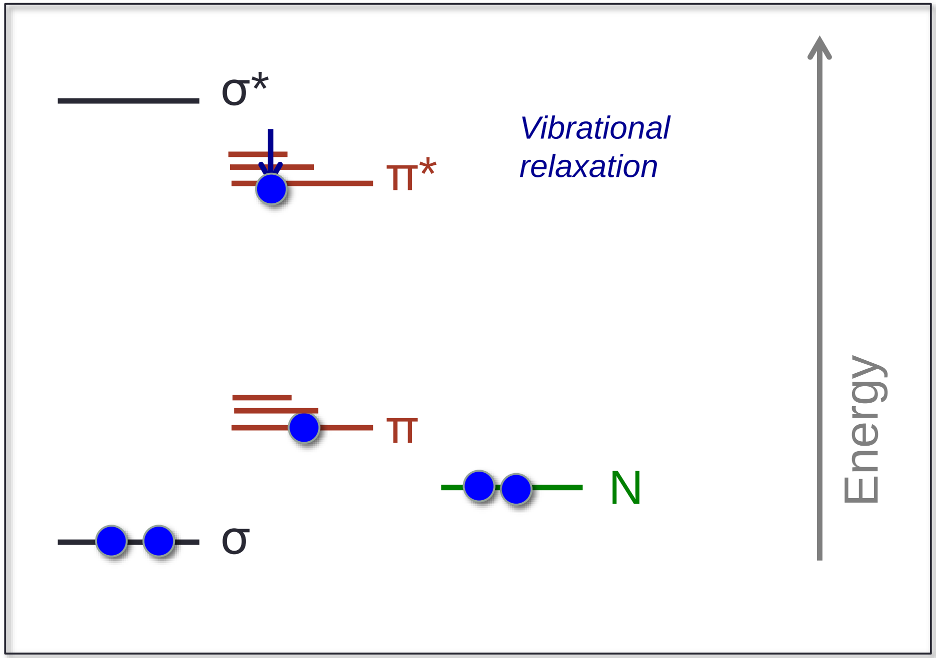




Now when there is a barium dication in the cage…

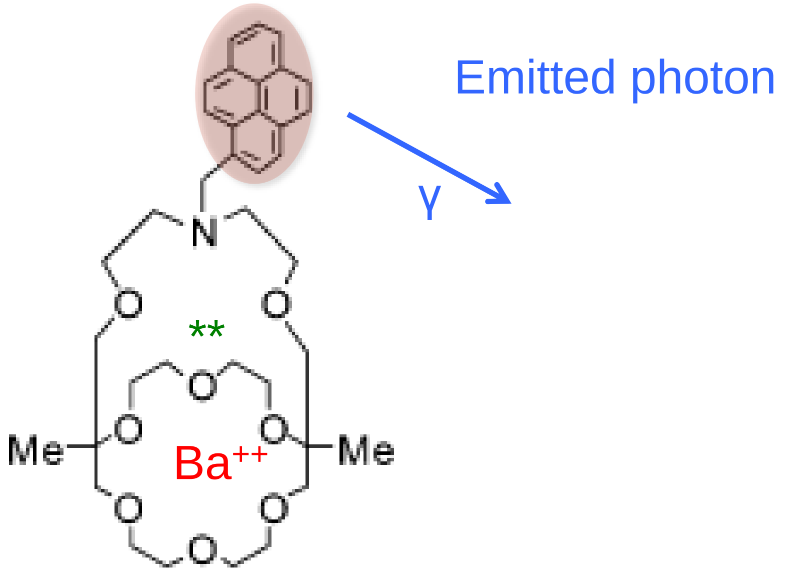
The PSMA magic:

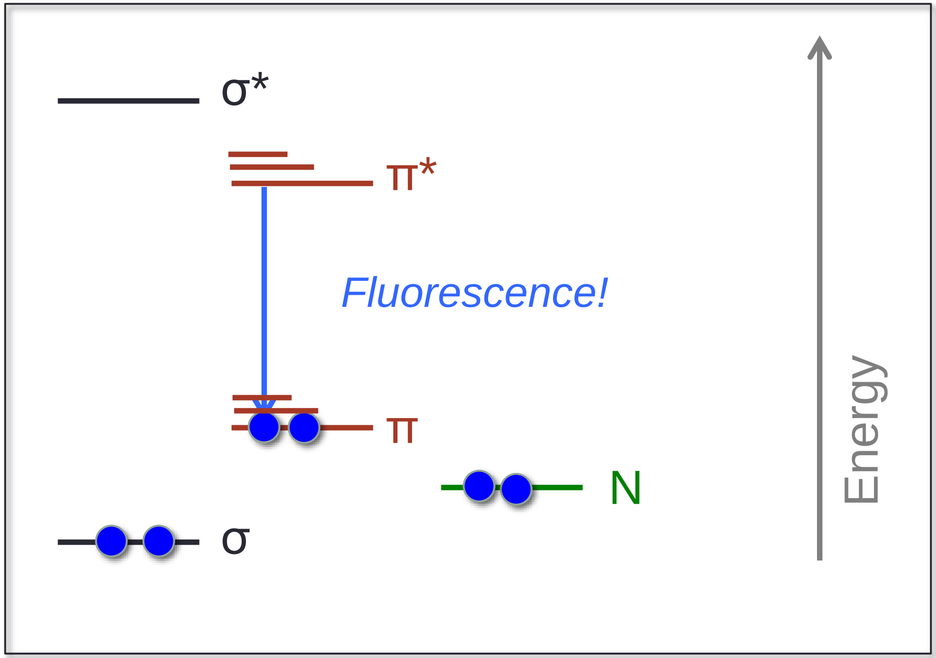




Now when there is a barium dication in the cage…

The PSMA magic:

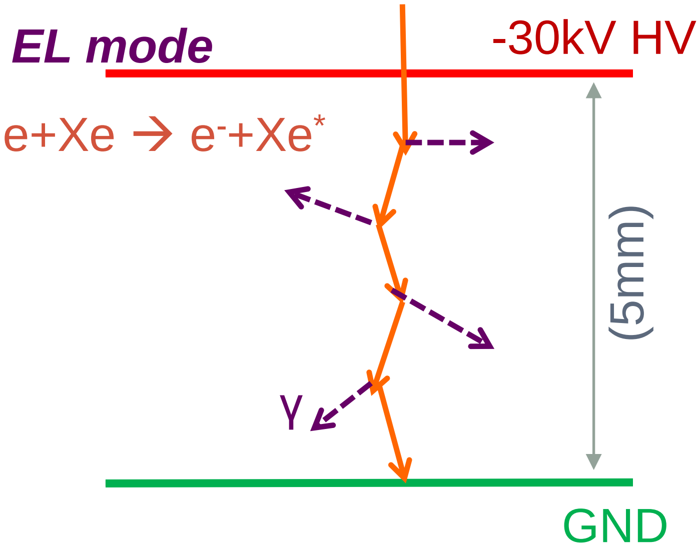




Fluorescence is switched on
only when barium ion is present

A scaling challenge: EL amplification

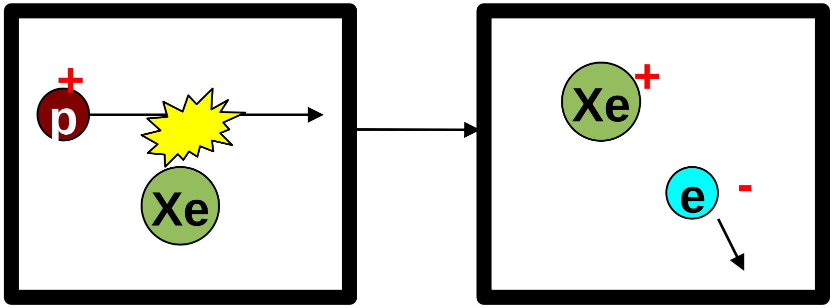
- Ongoing work at UTA to develop large, stable, uniform EL amplification plane for NEXT-100 (D=1.2 m) and NEXT-XXX (D~1.5m)



High Voltage Insulation and Gas Absorption of Polymers in High Pressure Argon and Xenon Gases
arXiv:1804.04116. L.Rogers et. al.

PROUDLY SUPPORTED BY:
U.S. DEPARTMENT OF ENERGY

Energy loss through ionization



Ionization

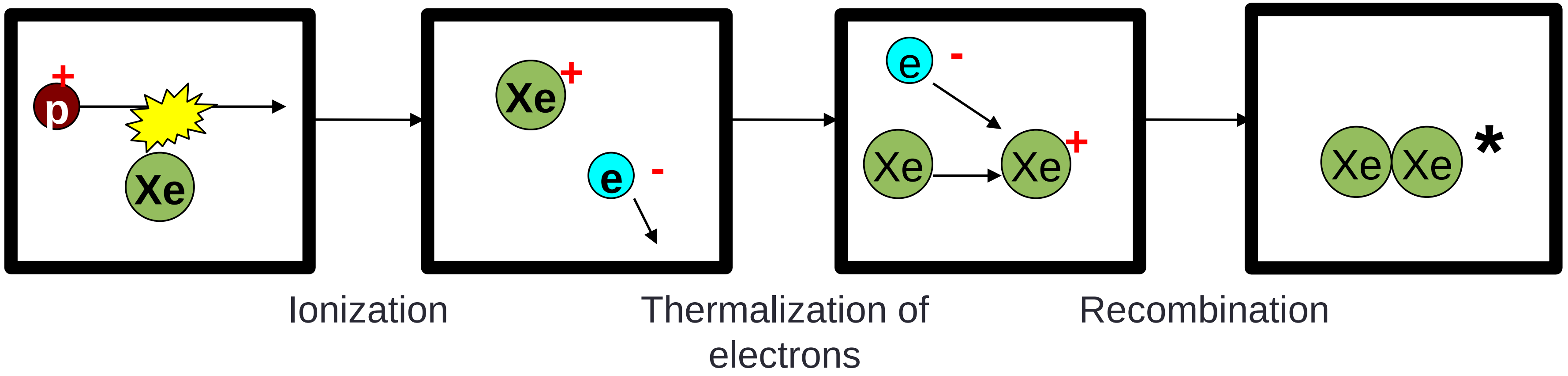
A well defined energy fraction is converted into ionization electrons

“Intrinsic” energy resolution achievable by measuring ionization at 2400MeV is around 0.3% at $Q_{\beta\beta}$

Fano Factor = 0.15

At high densities…

(gas >50 bar and liquid)

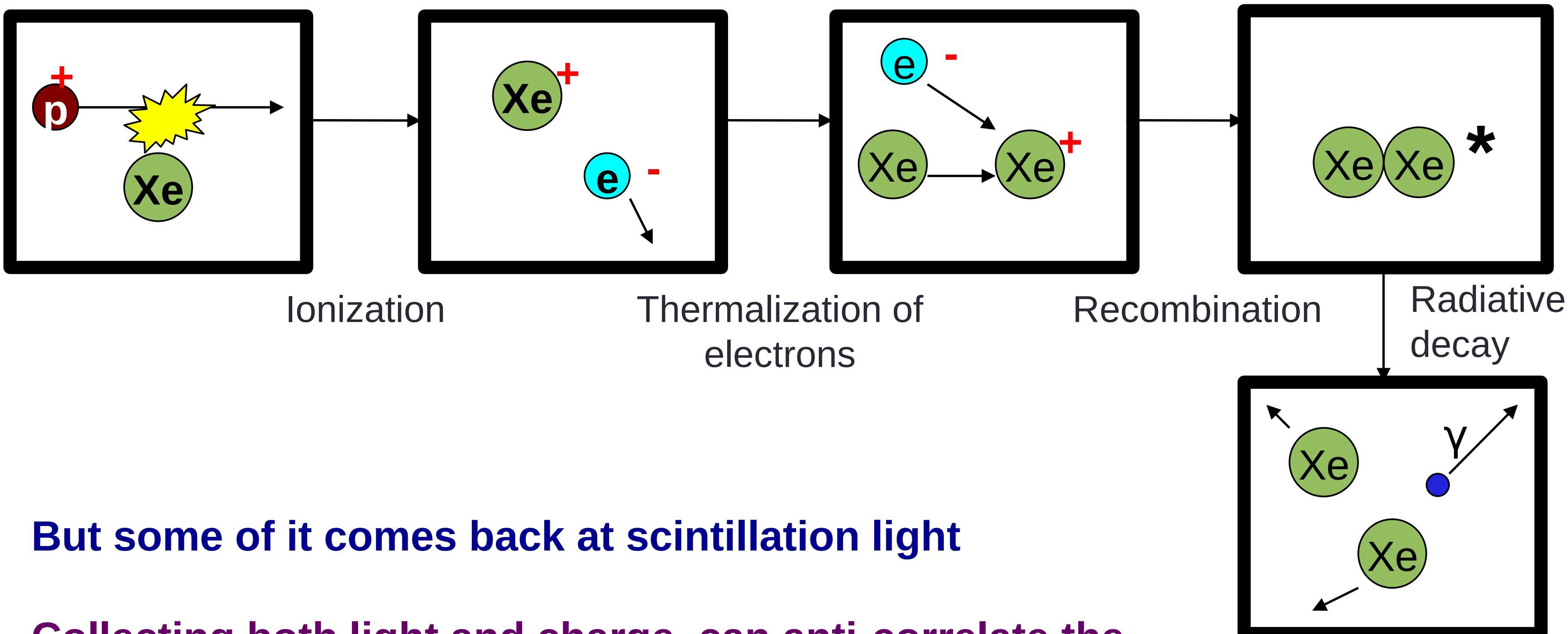


At high densities, ionization charge is lost to recombination

The recombination process is random, and this introduces fluctuations → intrinsic resolution is degraded
(ionization no longer carries a well known fraction of event energy)

At high densities…

(gas >50 bar and liquid)

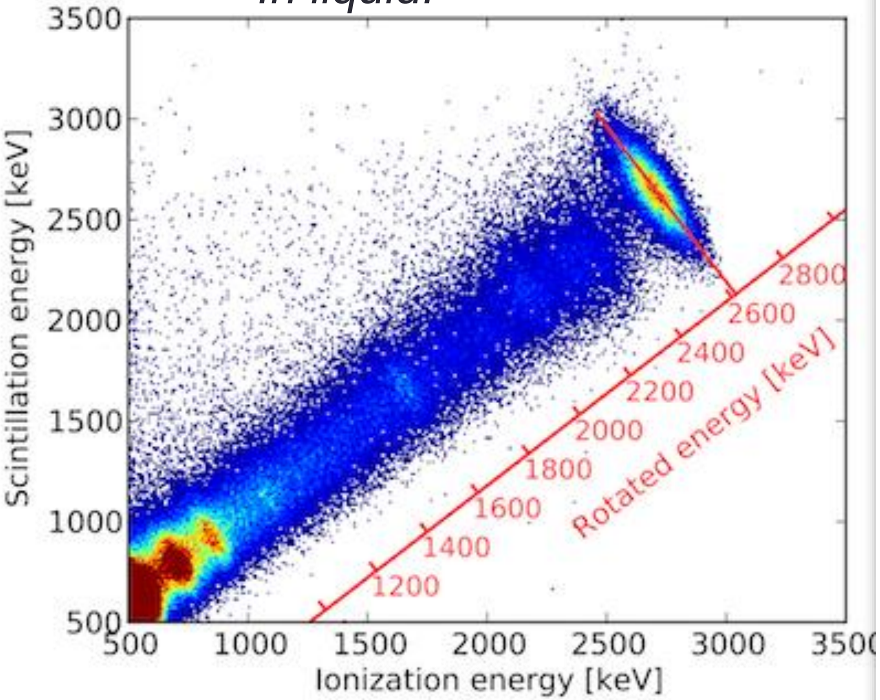


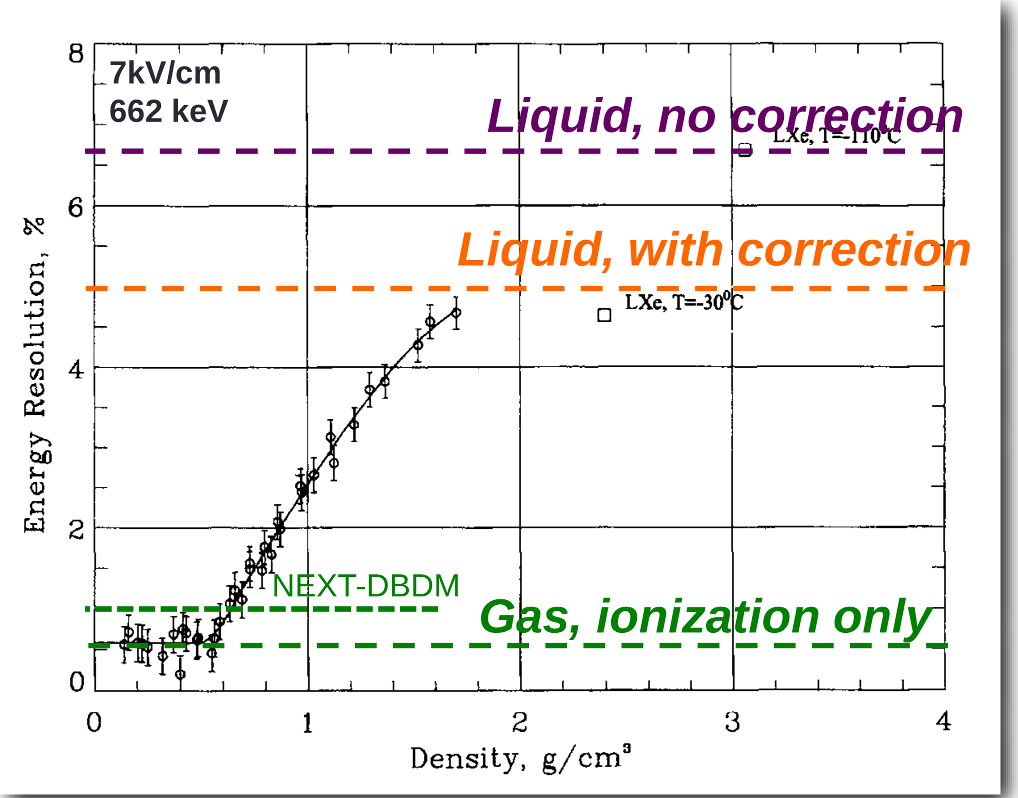
But some of it comes back at scintillation light

Collecting both light and charge, can anti-correlate the fluctuations. This is what liquid experiments do.

In gas @15 bar, no recombination and so no recombination fluctuations.

In liquid:





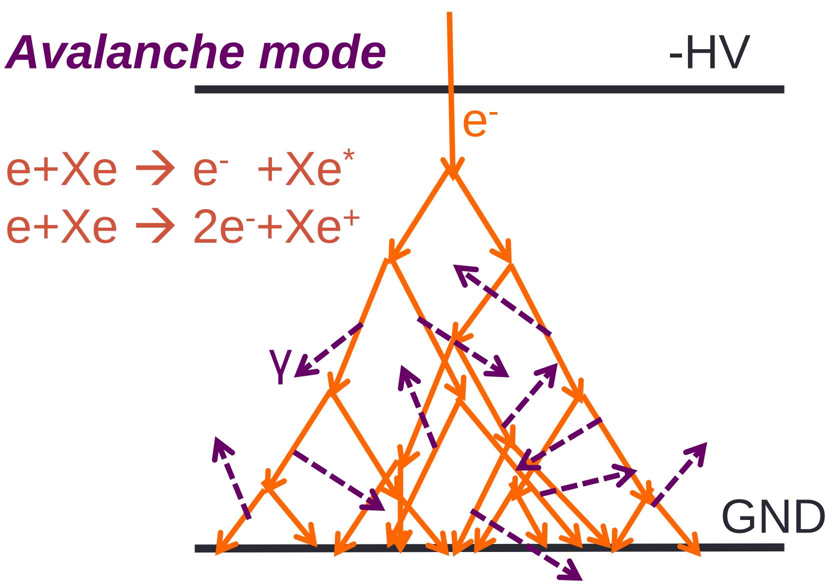
Caveat:

- *All the liquid lines are E-field dependent*
- *Gas has minimal recombination so ~no E-field dependence*

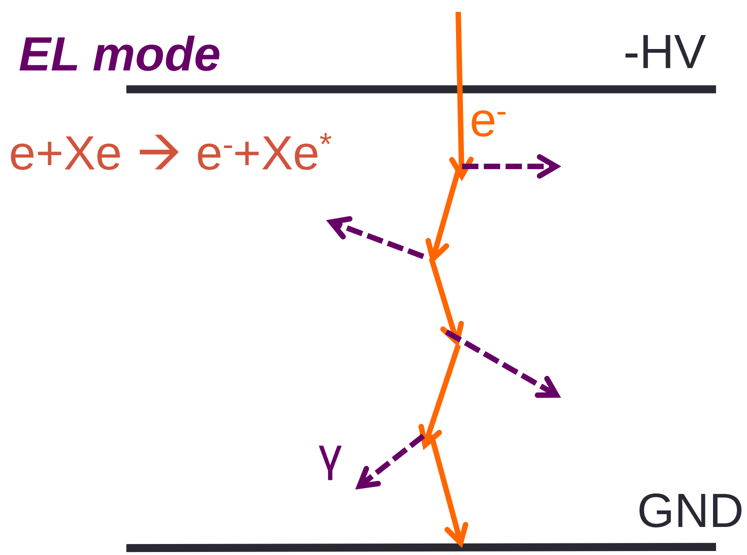
Amplifying the charge

- In high pressure gas, measuring ionization charge **alone** gives excellent energy resolution (0.3%) **in principle.**

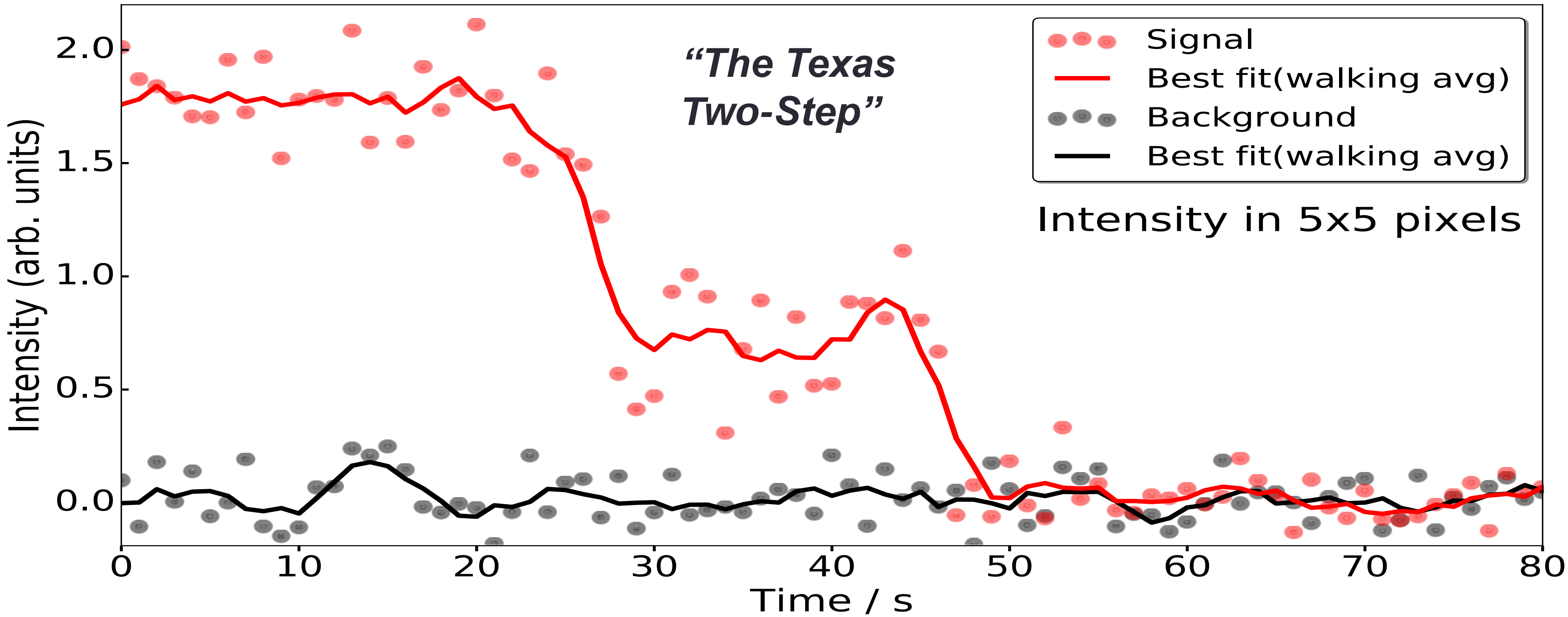
- Now the challenge is to amplify it without introducing fluctuations.
- Use electroluminescent gain mechanism:

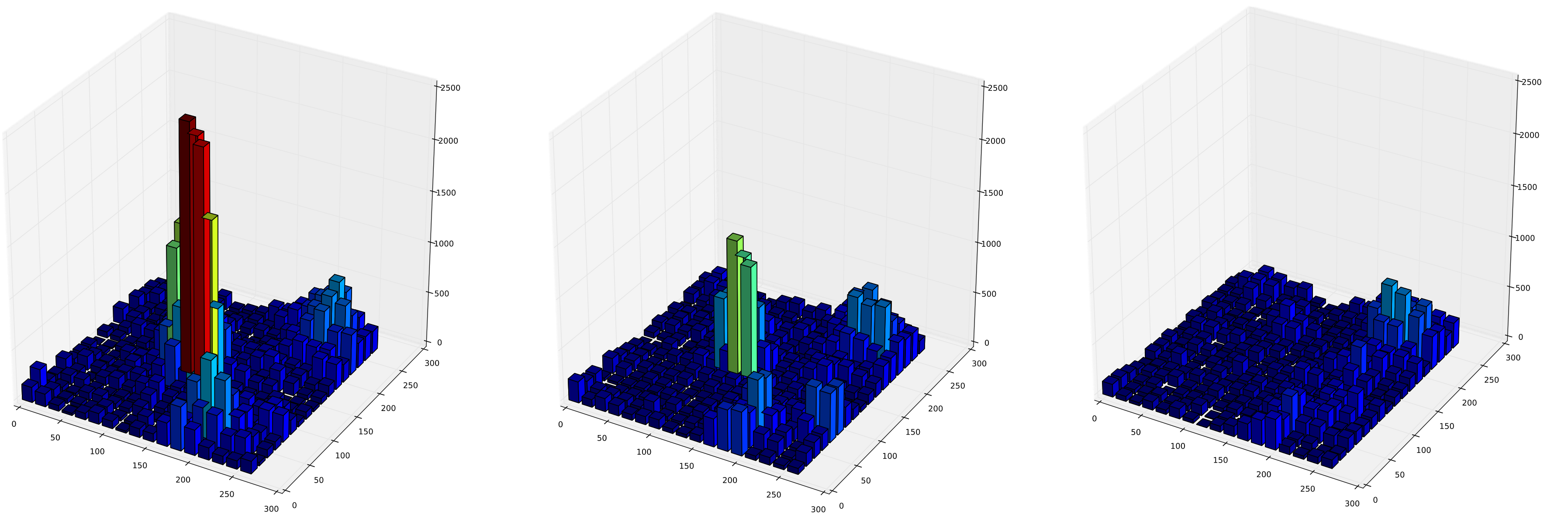


Electrons get enough energy to ionize as well as excite Xe:
→ big fluctuations in photon yield



Only enough energy to excite Xe
→ small fluctuations in photon yield





That this can happen suggests two ions may trap in a single pore in PVA matrix (v. rare)

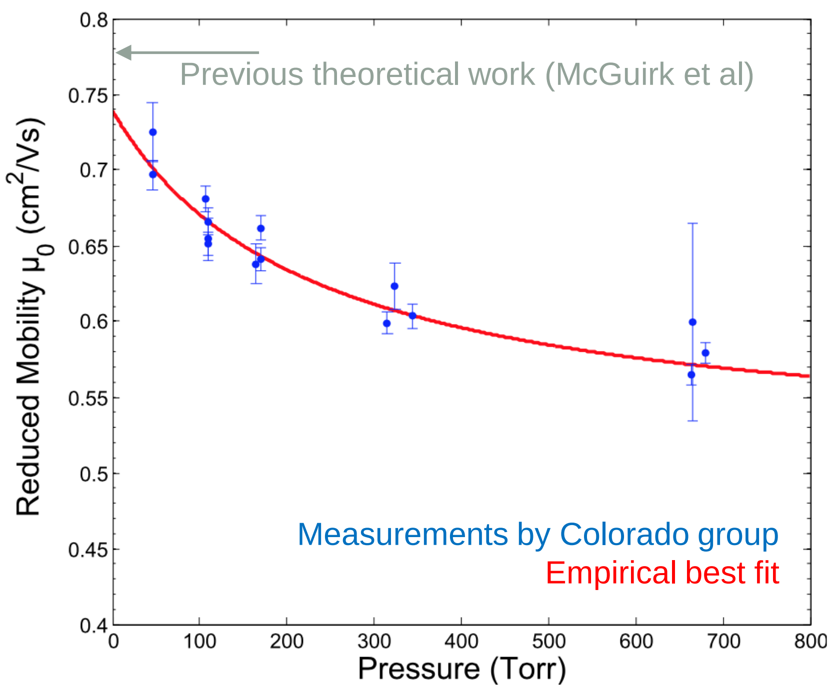
Barium ion mobility in gas

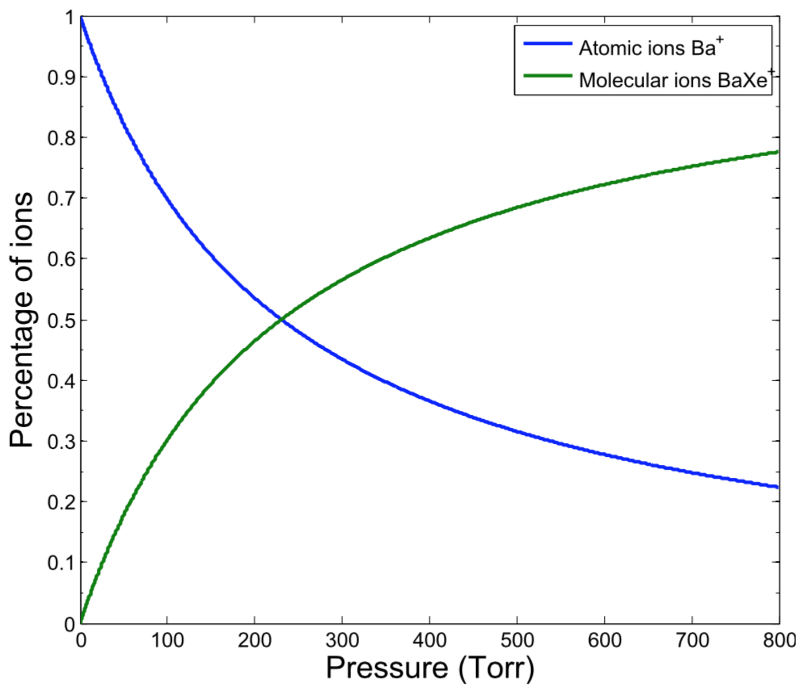
- Measurements from Medina's PhD thesis illustrate the importance of molecular ion formation.

$$Ba^+ + Xe + Xe \longleftrightarrow BaXe^+ + Xe$$

- Never measured in Ba^{++} to our knowledge
- And not measured above 1 bar to our knowledge either

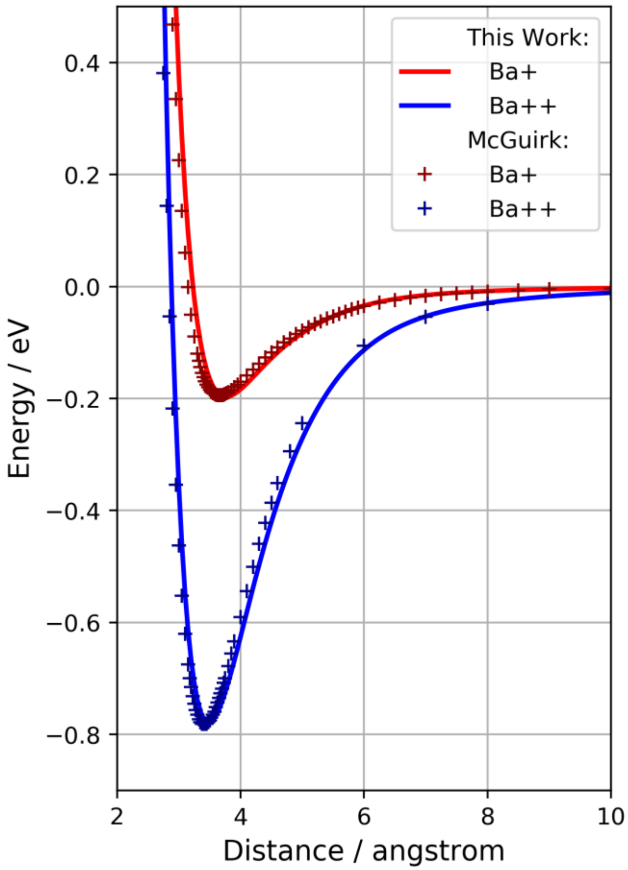
From **"Mobility and fluorescence of barium ions in xenon gas for the EXO experiment" – Julio Cesar Benitez Madina, PhD thesis, Colorado State University**





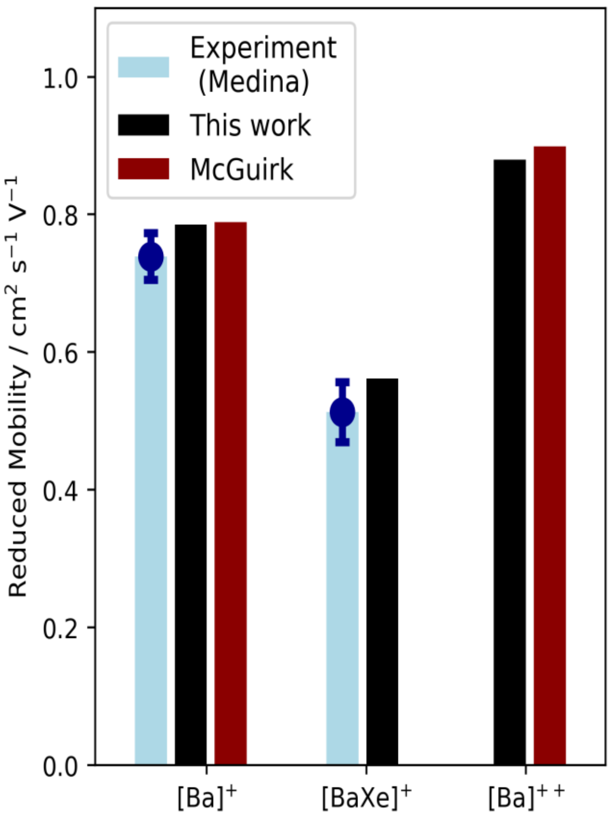
DFT vs coupled clusters

- Past published calcs used coupled cluster theory – good for bare ion ONLY.
- We calculate cluster distributions and cross sections using density functional theory then evaluate mobility.
- DFT is not expected to be as accurate. But too many atoms for coupled cluster method.
- We benchmark our DFT + ion mobility calculations for bare ions against McGuirk, and experimentally extracted data for BaXe system.



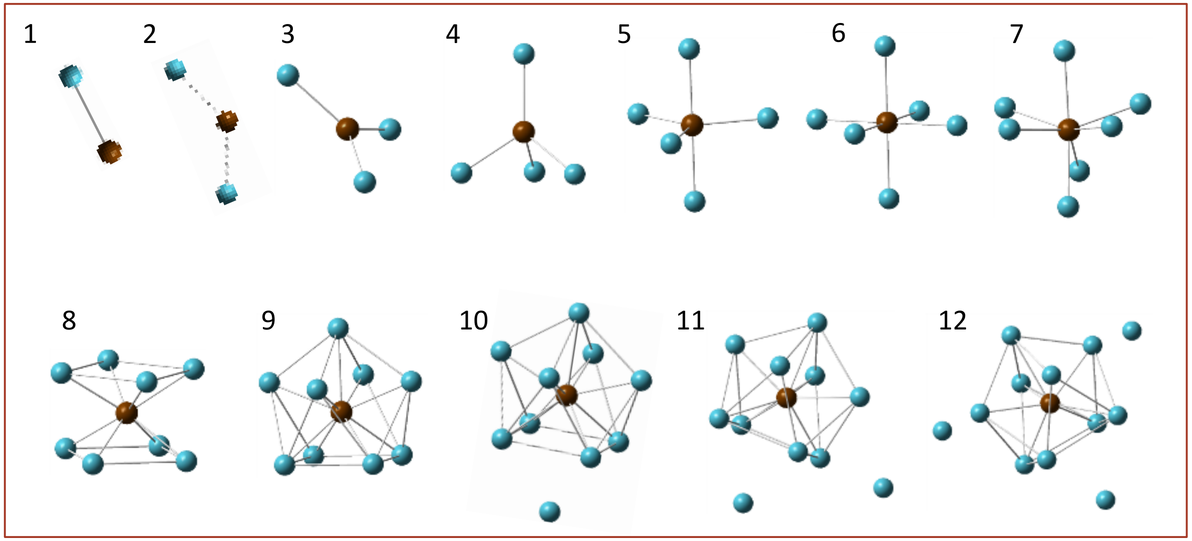
Comparison of potentials

%-level agreement



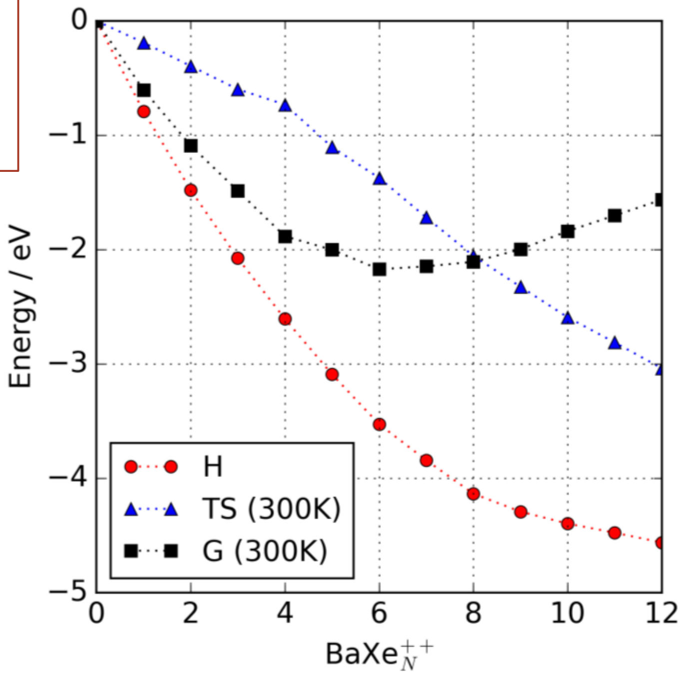
Comparison of mobilities

Ba++ clusters much more



Our calculated DFT clusters

Our calculated thermo quantities

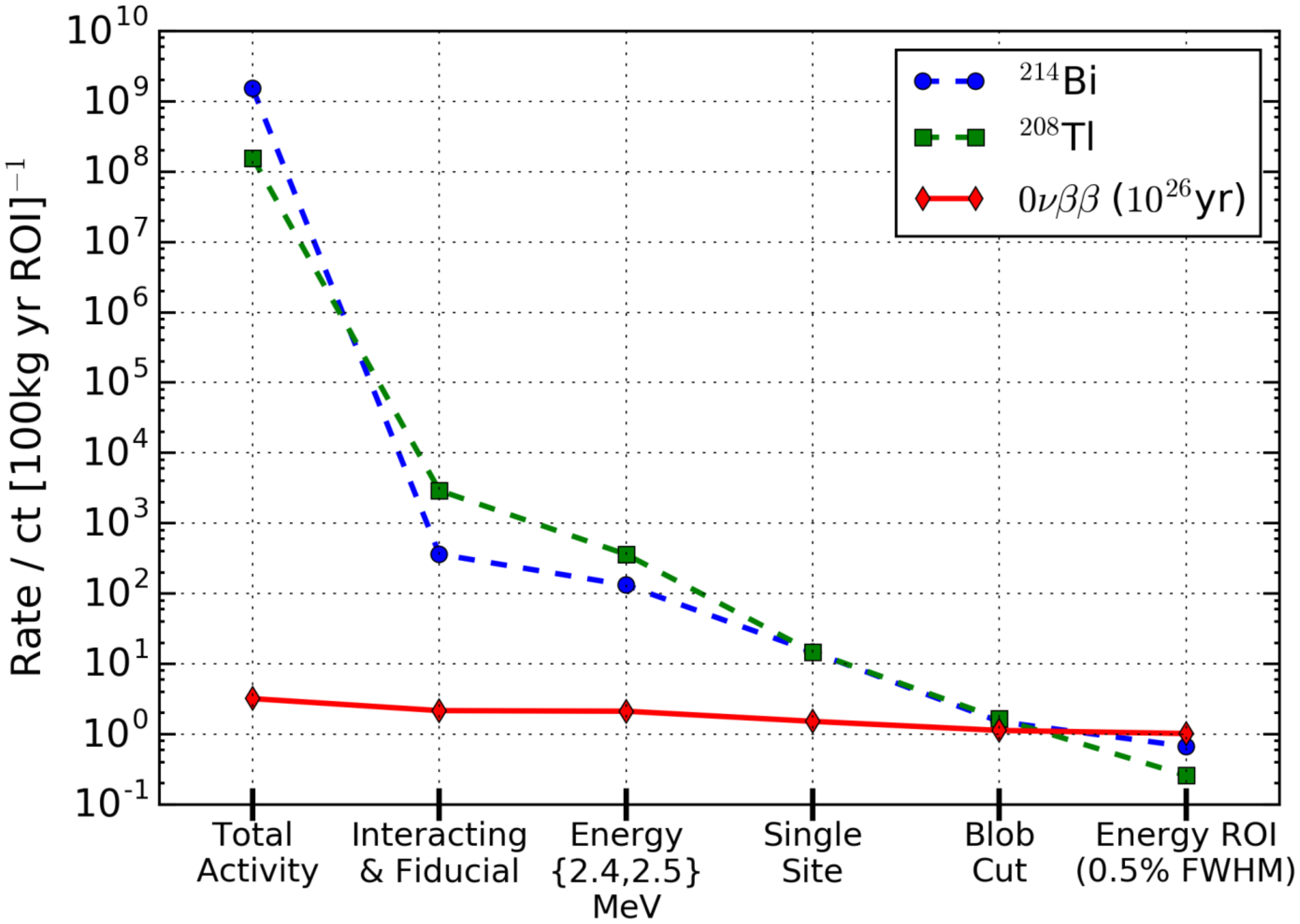


Enthalpy pushes this way

$$[\mathrm{BaXe}_{N-1}]^{++} + \mathrm{Xe} \leftrightarrow [\mathrm{BaXe}_N]^{++}$$

Entropy pushes this way

NEXT-100 projected background index performance



← Approximately Background free @ 100kg-scale

World-leading in xenon by an order of magnitude, but still ~5-7 cts/yr at ton scale

^{22}Na calibration of NEXT-NEW

- ***Preliminary results:***
- **6.75% FWHM (0.74% at $Q_{\beta\beta}$)** resolution for x-rays (29.6 keV).
- **2.35% FWHM (1.07% at $Q_{\beta\beta}$)** resolution for 511 keV γ.
- Continuing improvements from improved analysis methods and detector stability.