

Bottomonium in Heavy Ion Collisions: AA case

Quarkonia as Tools

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16/Jan/2019



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- Bottomonia measurements: “go heavy or go home” in AA.
- Bottomonium.

- what can beauty (+ charm) bound states tell us about **Hot QCD**
- How **hot** does it get?
- Is **color** deconfined?

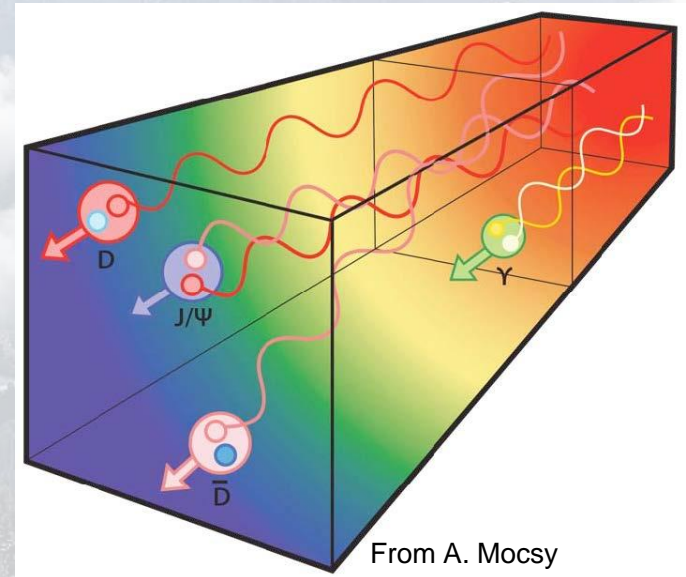
– Bottomonium measurements:

- Double ratios, RAA, centrality, pT, y, \sqrt{s} .
 - STAR: pp, dAu, AuAu, $\sqrt{s}=200$ GeV 1S, 2S+3S, UU $\sqrt{s}=193$ GeV.
 - » PLB 735 (2014) 127, PRC 94 (2016) 64904
 - PHENIX: pp, AuAu $\sqrt{s}=200$ GeV 1S+2S+3S, PRC 91, 024913 (2015)
 - ALICE: PbPb 2.76 TeV, 5.02 TeV 1S, 2S
 - » PLB 738 (2014) 361, arXiv:1805.04387 (to appear in PLB)
 - CMS, PbPb 2.76 TeV, 5.02 TeV 1S, 2S, 3S.
 - » PRL 109, 222301 (2012), PRL 120 (2018) 142301, arXiv:1805.09215 (to appear in PLB)



Bottomonium states in AA

- States are massive, produced early
 - pQCD can estimate production
- Sensitive to temperature and deconfined color fields: input from Lattice QCD
 - Debye screening, Landau damping
 - Re and Im $V(r, T)$
 - Different states have different sizes/binding energy
 - Sequential suppression
- Cold-nuclear matter
 - Initial state effects: e.g. nPDF, energy loss
 - Final state: absorption/co-mover interaction
- Regeneration
 - Uncorrelated heavy-quarks can pair up
- Bottomonium: a cleaner probe than charmonium...
 - 3 states are accessible experimentally
 - Differentiate between initial and final state effects
 - expect some small CNM effects (shadowing/nPDF)
 - expect small regeneration effects



From A. Mocsy



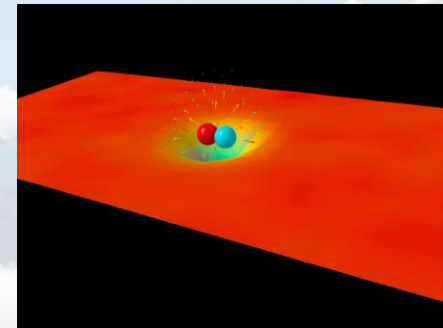
Heavy Quark Potential at High T

- Quarkonium suppression: longstanding QGP signature
 - Original idea: High T leads to QCD Debye screening

QED:
$$V \sim -\frac{a_{\text{eff}}}{r} \quad V \sim -\frac{a_{\text{eff}}}{r} e^{-\frac{r}{r_D(T)}} \\ T = 0 \quad T > T_c$$

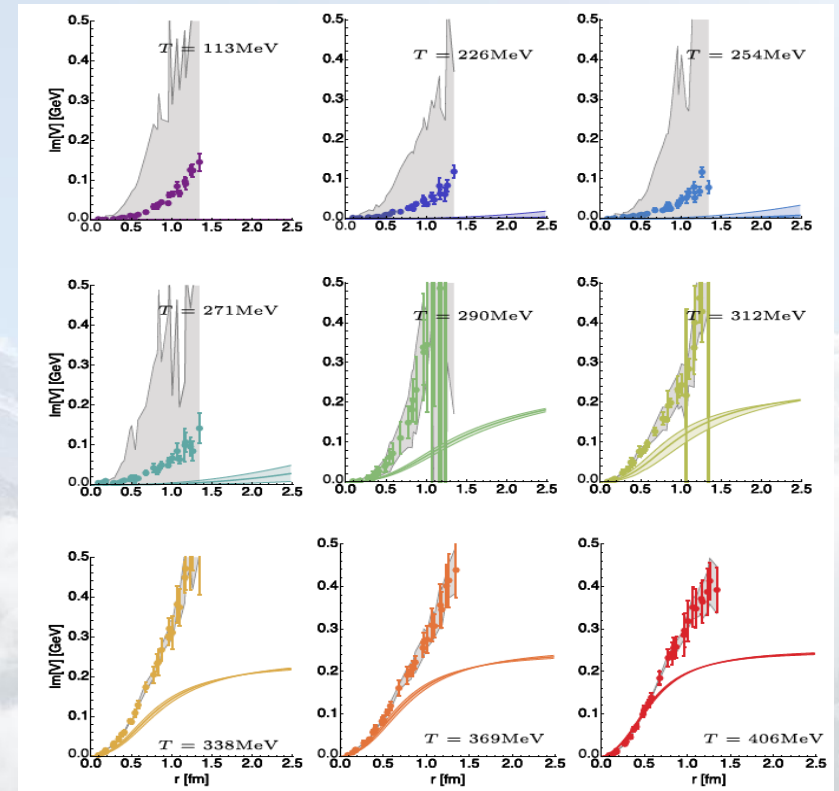
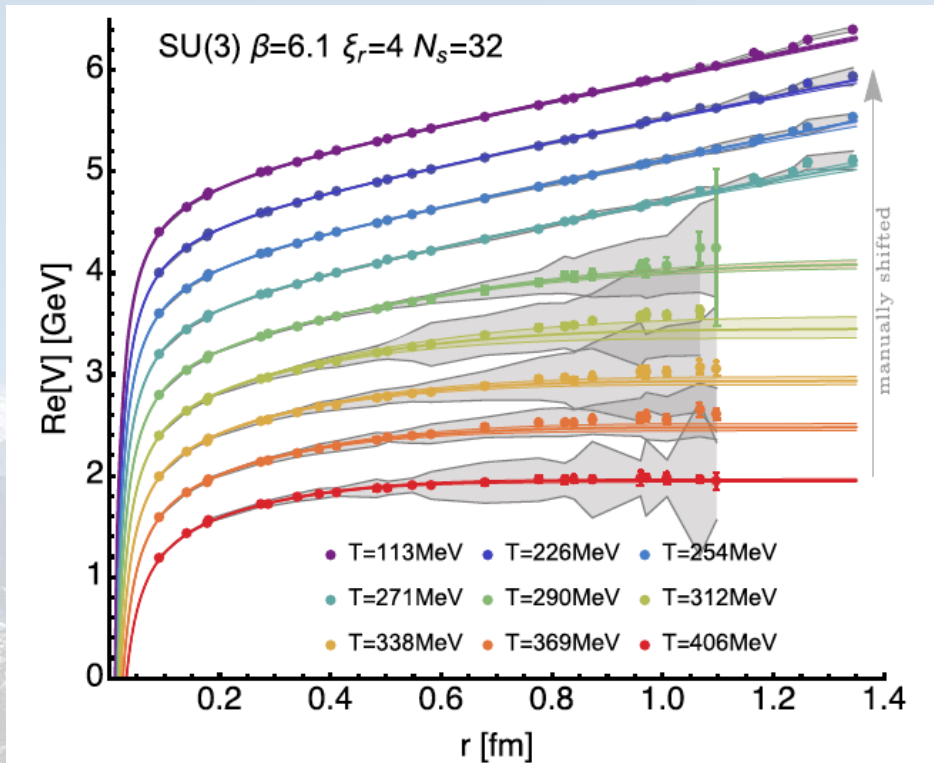
QCD:

- T=0, Cornell potential: $V = -\frac{\alpha}{r} + \sigma r + c$
- Presence of Debye mass in Cornell potential, screening of potential at large r .
- Screening prevents heavy quark bound states from forming!
- **Original idea of J/ ψ suppression:**
 - Matsui and Satz, *Phys. Lett. B* **178** (1986) 416
- What do the latest calculations in lattice QCD say?





Lattice QCD Heavy Quark Potential



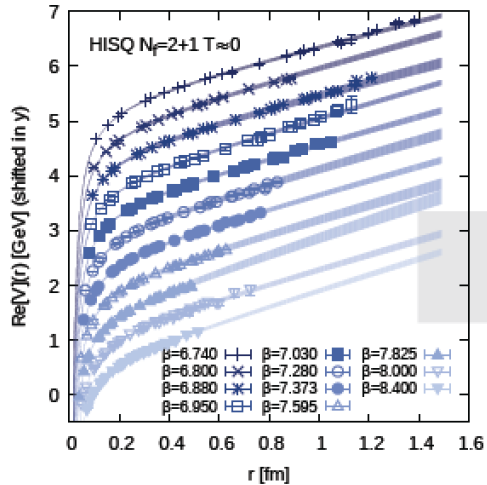
Y. Burnier & A. Rothkopf
PRD 95 054511 (2017)

- Lattice calculations confirm screening effects
 - **Screening: Re V**
 - **Landau damping, gluodissociation: Im V**
 - Both contribute to Quarkonium Suppression

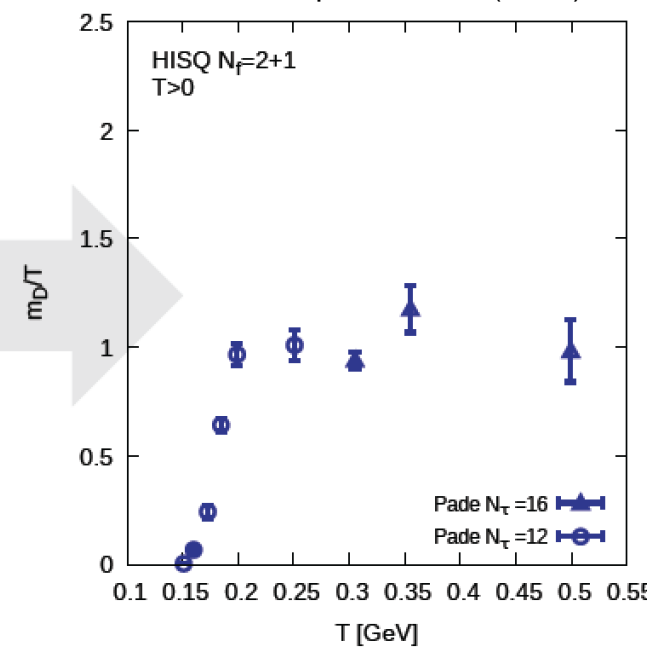
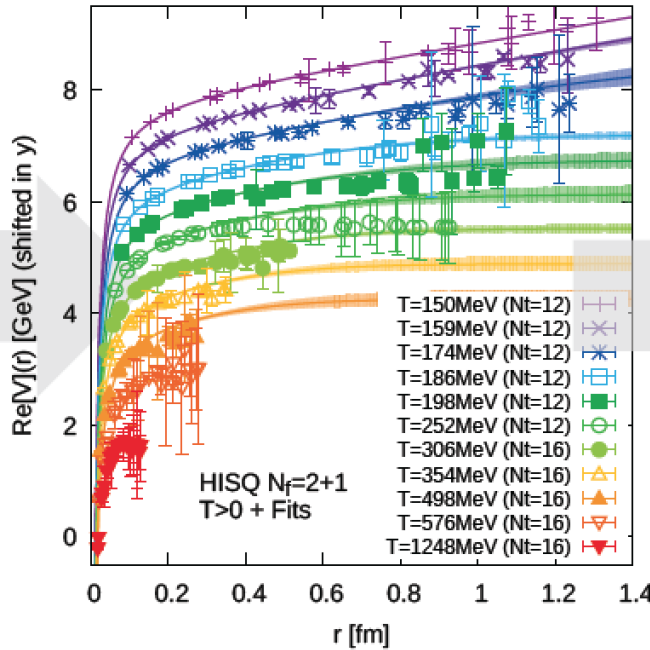


Interpreting the $T > 0$ potential

Burnier & Rothkopf, PLB 753 (2016) 232



fit $T=0$ Cornell parameters
 $V_{\text{Cornell}}(r) = -\alpha_s/r + \sigma r + c$

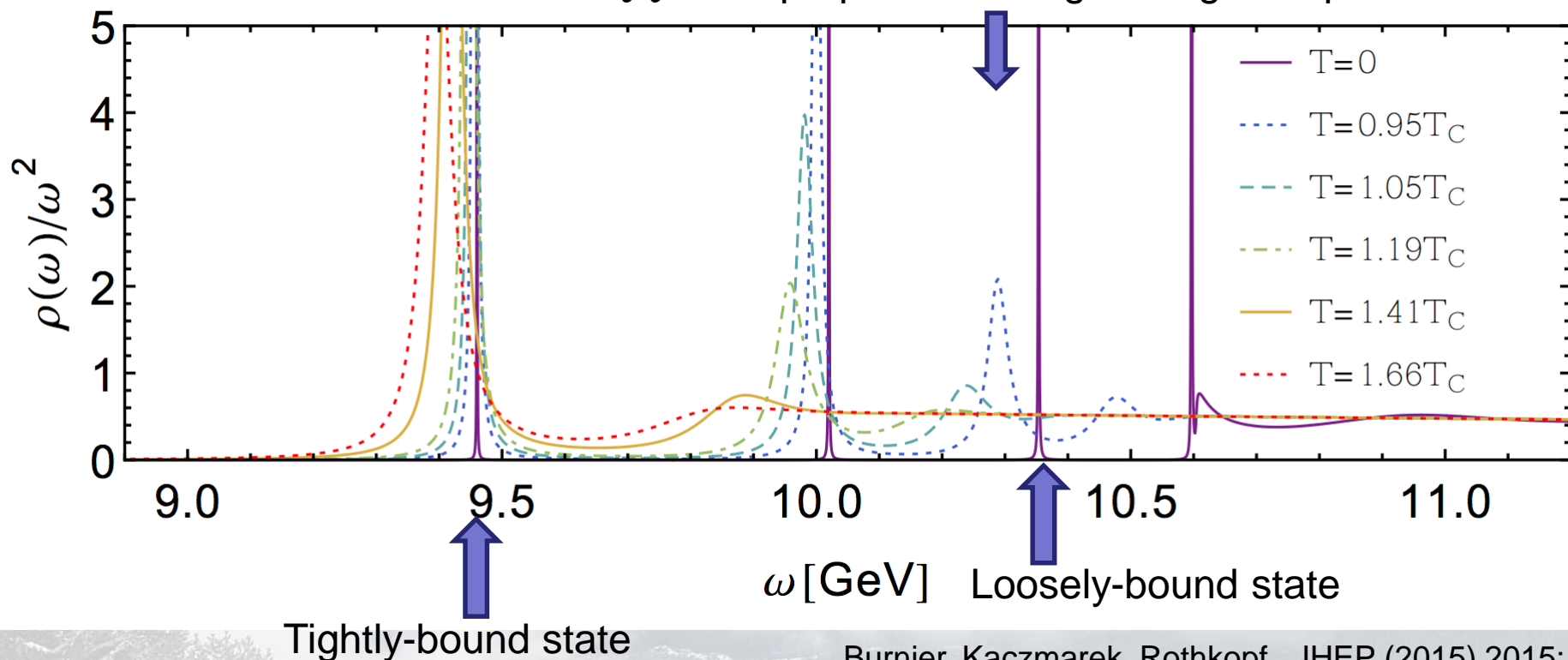


- Combine non-perturbative $T=0$ Cornell potential and perturbative medium:
 - $\text{Re}[V](r,T)$ and $\text{Im}[V](r,T)$ from one T -dependent parameter: Debye mass, m_D .
 - Lattice results from $\text{Re}[V]$ are well described by tuning m_D , smooth onset for $T > T_c$.
 - $\text{Im}[V]$ predictions based on m_D are slightly lower.



Sequential Melting of Bottomonia

$Q\bar{Q}$ state properties change at high temperature



Burnier, Kaczmarek, Rothkopf, JHEP (2015) 2015: 1

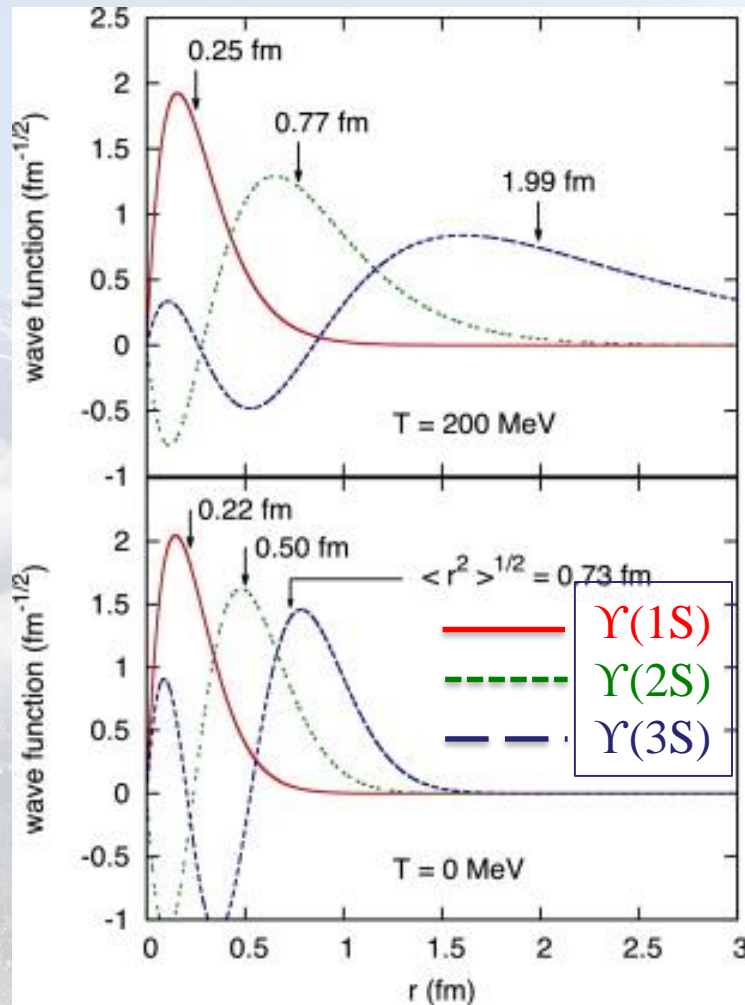
- As temperature increases:
 - The peaks broaden and their masses shift to lower values.
 - Highest states broaden and shift first, followed sequentially by lower states.
 - Peaks eventually disappear completely \rightarrow States melt.
 - $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$ melt at $2.66T_C$, $1.25T_C$ and $1.01T_C$.
- Sequential melting of Υ states is sensitive to color deconfinement!



Bottomonia spectroscopy as a tool

- $\Pi \rightarrow \mu\mu$: 3 states to study
 - Hot Nuclear Matter Effects
 - Different sizes, binding energies.
 - Sequential melting.
 - Expectation:
 - $\Upsilon(1S)$ almost no melting
 - $\Upsilon(2S)$ likely to melt
 - $\Upsilon(3S)$ fully melted?
 - Feed-down is important.
- Recombination/coalescence
 - $\Pi(1S)$ recombination \rightarrow much smaller effect than for J/ψ .

– Ko et al. PRC 85, 014902 (2012)

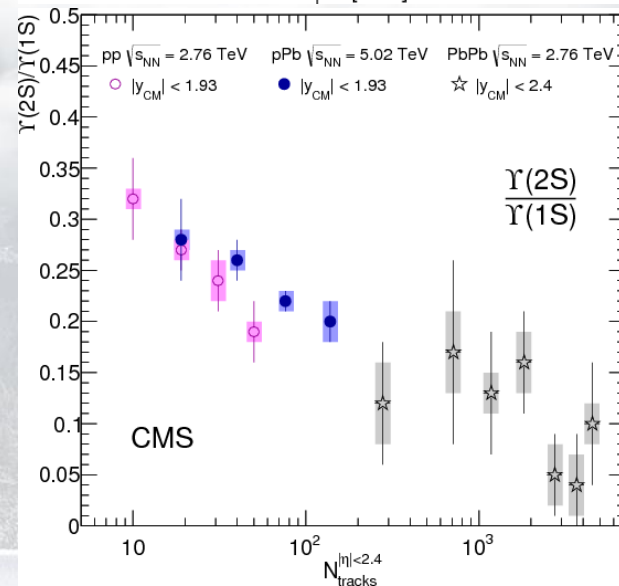
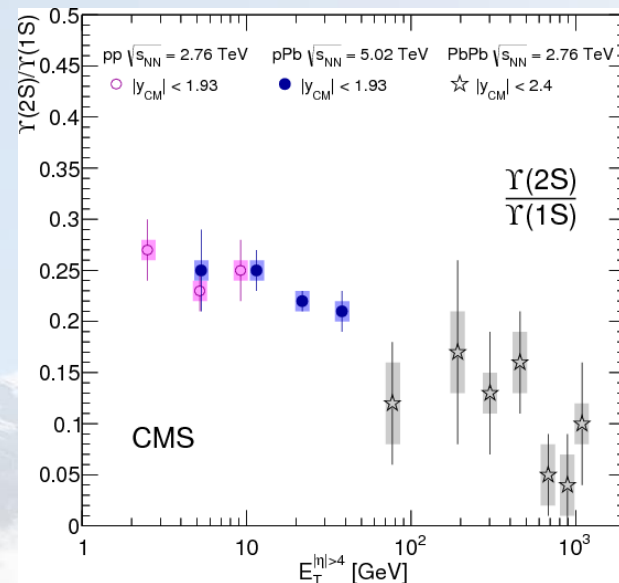


Brezinski & Wolschin PLB 707 (2012) 534



Bottomonia: from pA to AA

- pA measurements are crucial to study CNM effects.
 - Necessary step before discussing possible Hot Nuclear Matter effects.
 - e.g. ratios vs. activity in pp, pA, AA.
 - Extrapolation from pA to AA not trivial.
- Bottomonia:
- Cold Nuclear Matter Effects
 - Shadowing small near $y \sim 0$
 - Small hadronic absorption for $\Upsilon(1S)$.
 - Lin & Ko, PLB 503 (2001) 104
 - Nuclear breakup \sim before formation time.
 - Comover Interaction Model: AA and pA?
 - Ferreiro and Lasberg: JHEP 10 (2018) 094





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Bottomonia in AA



Υ Relative and Absolute Modification

- Observables:

- Double Ratio of excited to ground state yields

- Relative modification of excited states compared to ground state

- Cancellation of efficiency and acceptance corrections

- Cancellation of initial state effects, e.g. nuclear shadowing of PDF's

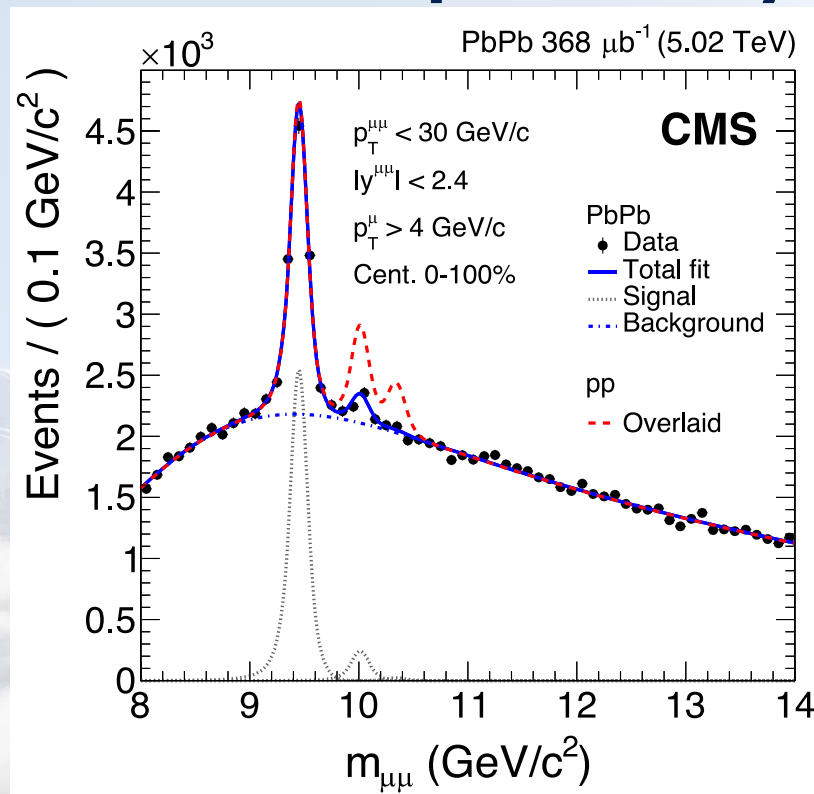
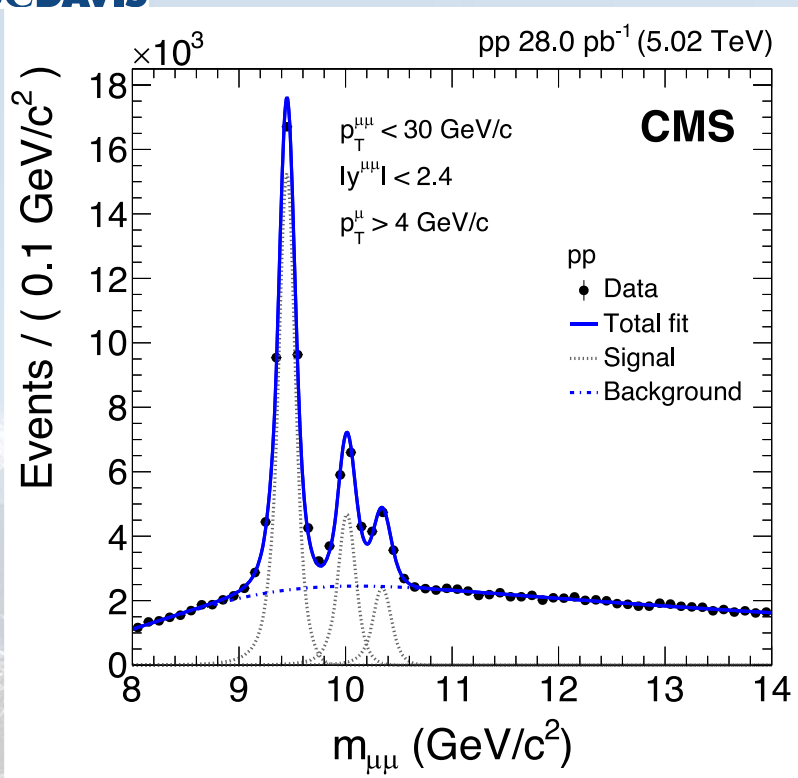
- $$\frac{\frac{Y(nS)}{Y(1S)_{PbPb}}}{\frac{Y(nS)}{Y(nS)_{pp}}} = \frac{R_{AA}(Y(nS))}{R_{AA}(Y(1S))}$$

- Nuclear Modification Factor, R_{AA}
$$R_{AA} = \frac{1}{N_{coll}} \frac{N(Y)_{PbPb}}{N(Y)_{pp}}$$

- Ratio of invariant yields (or cross section) from PbPb to pp
- Scaled by the number of nucleon-nucleon collisions
- Absolute modification from pp to AA



□ Double ratio Graphically

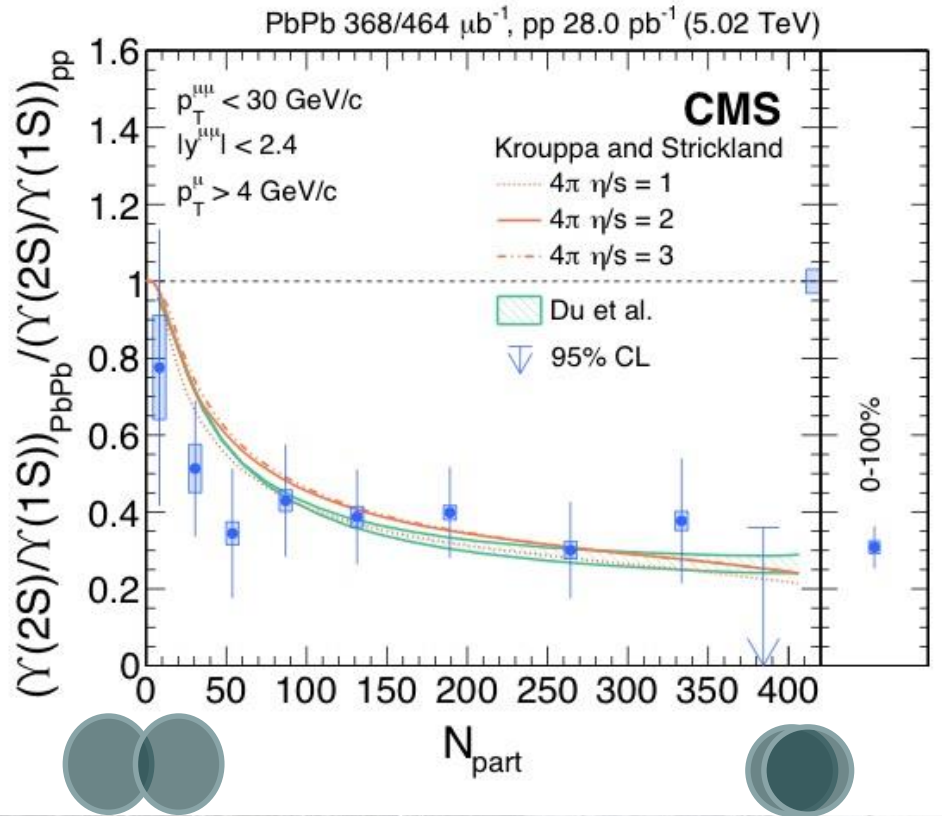


CMS: PRL 120 (2018) 142301

- Visual representation of Double Ratio
- pp shapes scaled to 1S in PbPb and overlaid on PbPb data:
 - Relative suppression of excited states.
 - Strong relative suppression of 3S state!



$\Upsilon(2S)$ Double Ratio vs. Centrality



Peripheral collisions

Central collisions

CMS: PRL 120 (2018) 142301

- Larger suppression toward more central events
- Consistent with unity in most peripheral bin

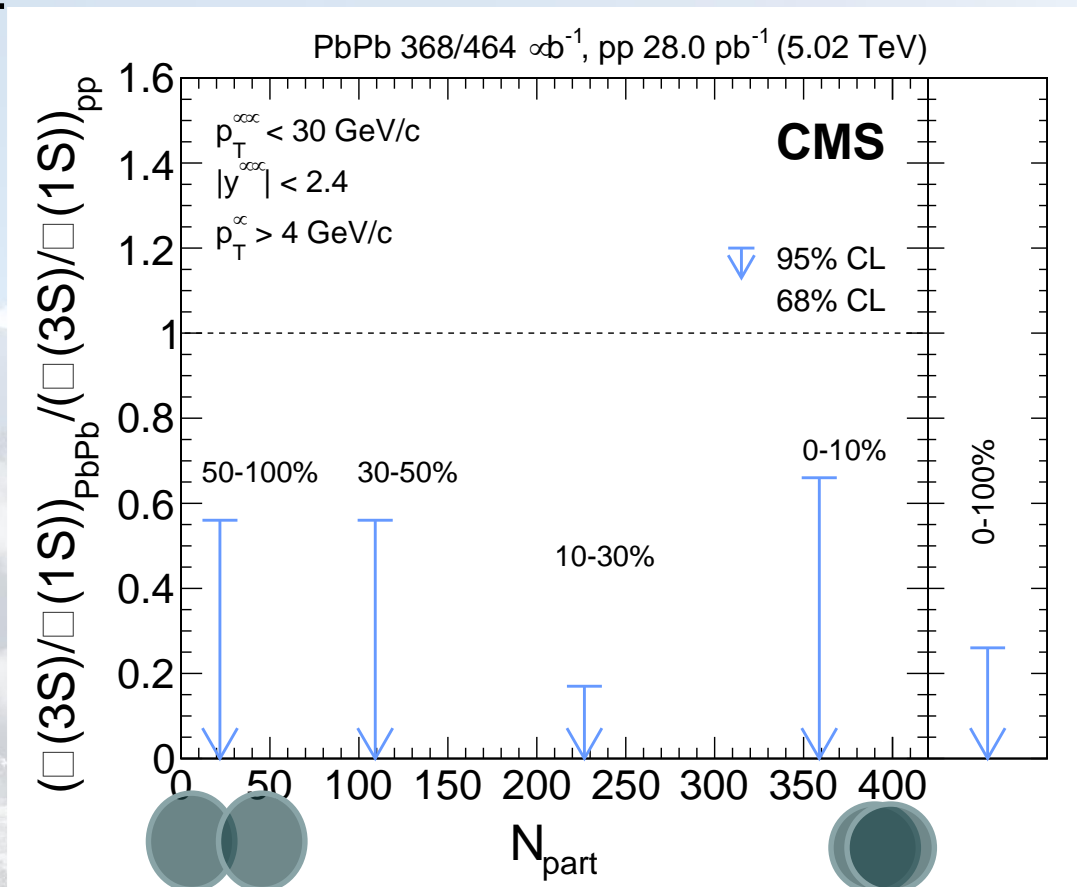
• Comparison to theory:

- Model: Strickland et al.
- Containing bottomonia evolved using anisotropic hydrodynamics
 - Curves:
 - $\frac{4\pi\eta}{s} = \{1,2,3\}, T_0 = \{641, 632, 629\} \text{MeV}$
- Consistent with our data
 - Similar results by Rapp et al.



$\Upsilon(3S)$: Strong suppression!

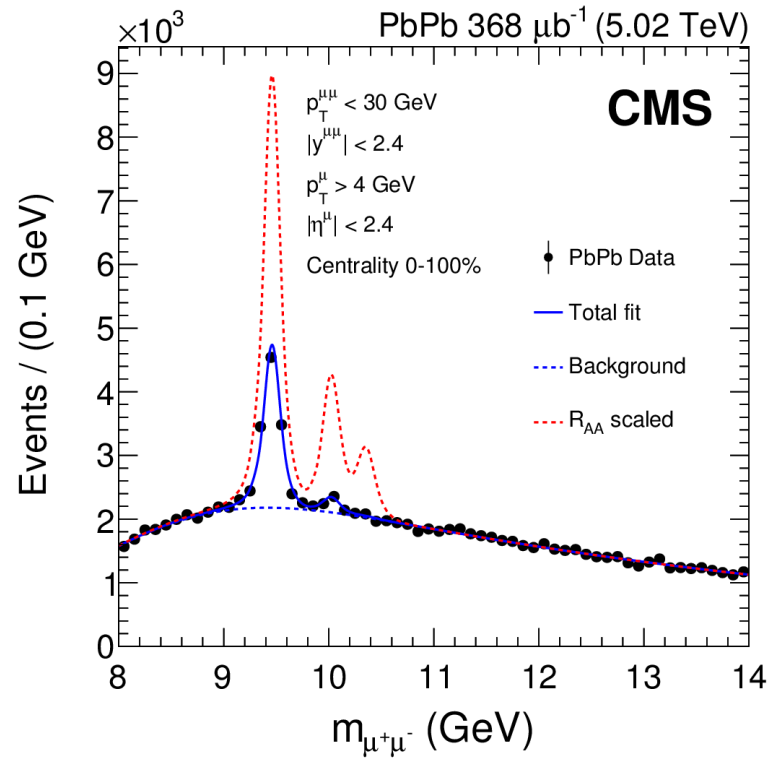
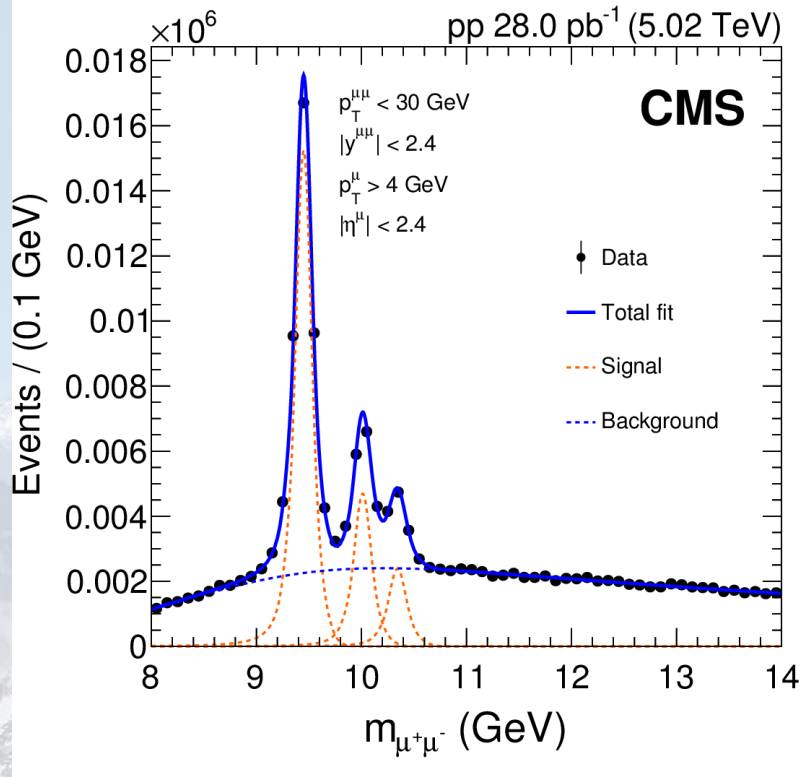
- Strong suppression of $\Upsilon(3S)$ relative to the $1S$ in all centralities
- Upper limits calculated in all cases
- $\Upsilon(3S)$ has smallest binding energy in Υ family
 - Sequential suppression of Υ states
 - Supports picture of melting in a color-deconfined QGP



CMS: PRL 120 (2018) 142301



Υ Absolute modification: Striking result



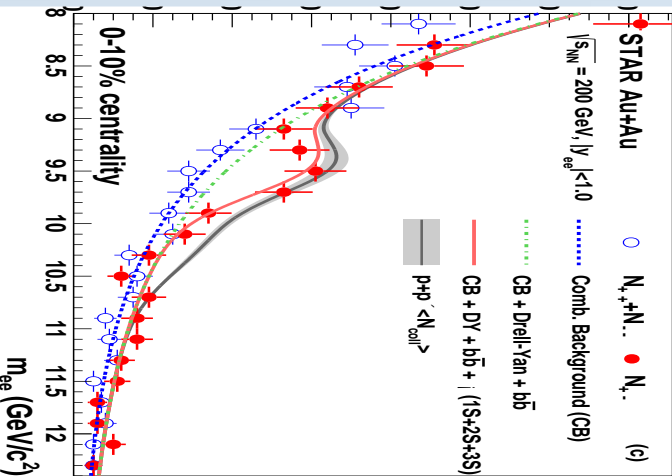
CMS [arXiv:1805.09215](https://arxiv.org/abs/1805.09215)
PLB in press

- Measure production in pp.
- Scale by number of binary nucleon-nucleon collisions, compare to AA measurement.
- **Suppression observed!**

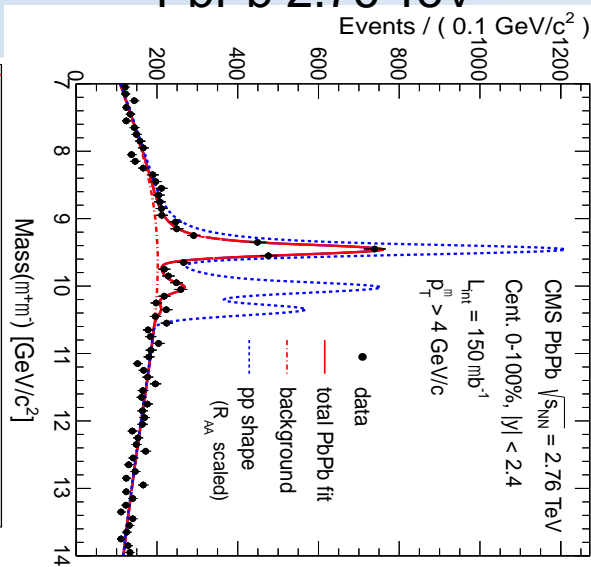


R_{AA} Graphically, RHIC and LHC

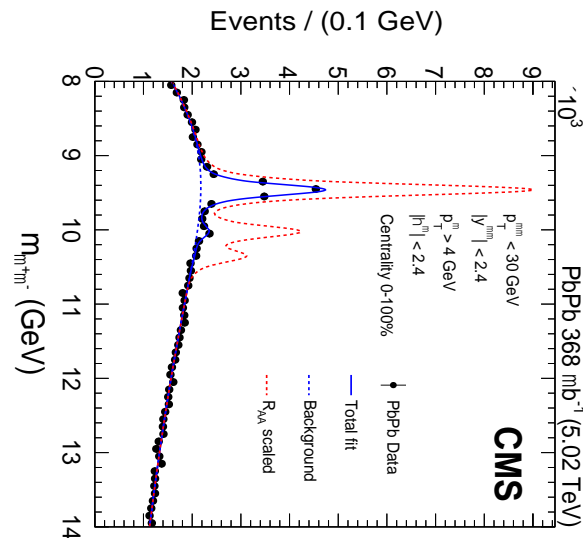
AuAu 200 GeV



PbPb 2.76 TeV



PbPb 5.02 TeV



STAR: PLB 735 (2014) 127

CMS: PLB 770 (2017) 357

CMS: arXiv:1805.09215

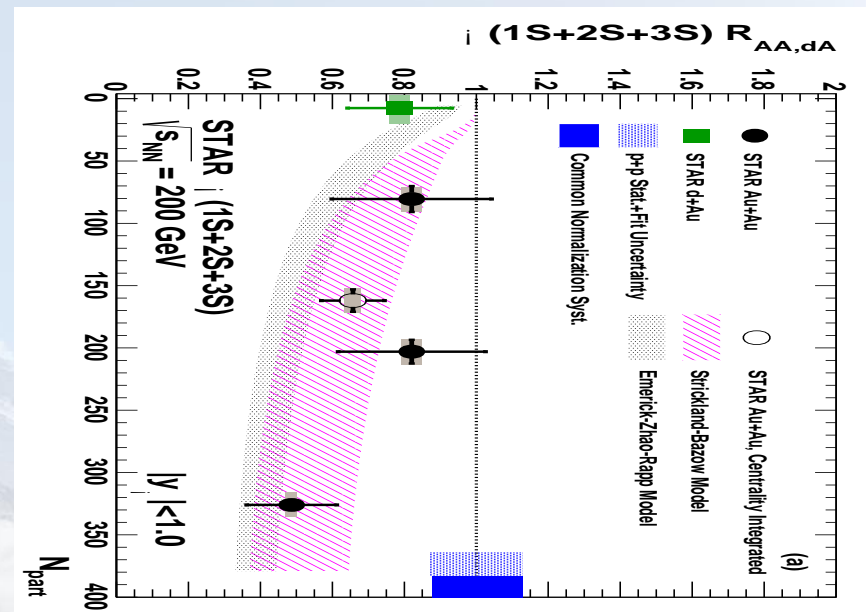
- Heavy ion data: dilepton invariant mass in \otimes region
- pp reference shape: overlay after scaling by number of binary collisions, N_{coll}
- Difference between Data and pp reference: evidence for \otimes suppression, $R_{AA} < 1$.



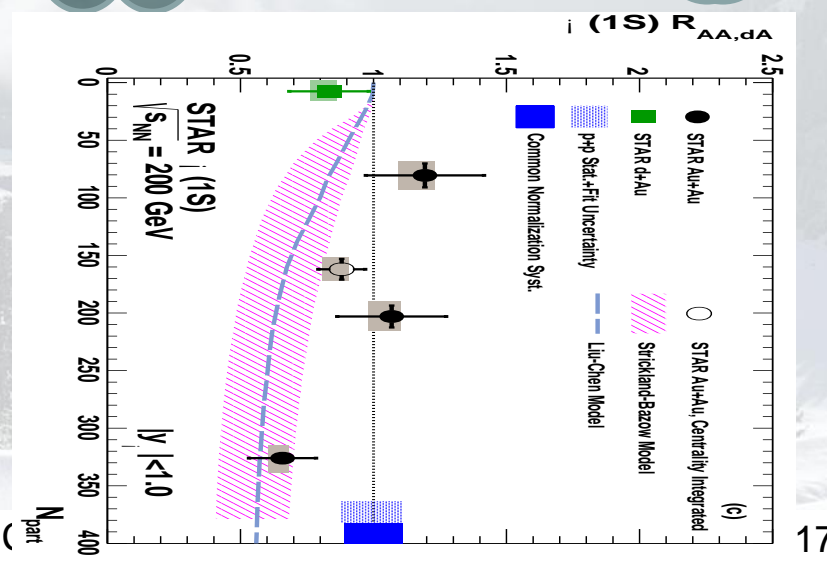
Υ Nuclear Modification in STAR

- All data in STAR acceptance $|y| < 1$
 - dAu, and two most peripheral Au+Au bins: consistent with no suppression
 - Suppression in the most central Au+Au bin:
 - Consistent with expectations for hot & cold nuclear matter
- Calculations:
 - Strickland & Bazow (Nucl. Phys. A 879 (2012) 25):
 - Includes estimate of heavy quarkonium potential, Re and Im.
 - Evolution through anisotropic hydro.,
 - T: 428 – 442 MeV (Depending on η/s , to match $dN/d\eta$)
 - Emerick, Zhao & Rapp:
 - attempt to include both Hot & Cold nuclear effects.
 - $T_0 = 330$ MeV.

- First combined dAu, AuAu analysis:
 - Peripheral R_{AA} similar to R_{dAu} .
 - Hint of stronger suppression for more central collisions, but not enough data to conclude difference with dAu



STAR: PLB 735 (2014) 127





$\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$ R_{AA} at 2.76 TeV

- 2011 150 / μb

- $\Upsilon(1S)$ R_{AA} , 7 centrality bins

- First results on $\Upsilon(2S)$ R_{AA}

- Clear suppression of $\Upsilon(2S)$

- $\Upsilon(1S)$ suppression

- Excited state suppression removes feeddown contribution.

- 60-70% $\Upsilon(1S)$ directly produced.

- Is direct $\Upsilon(1S)$ suppressed?

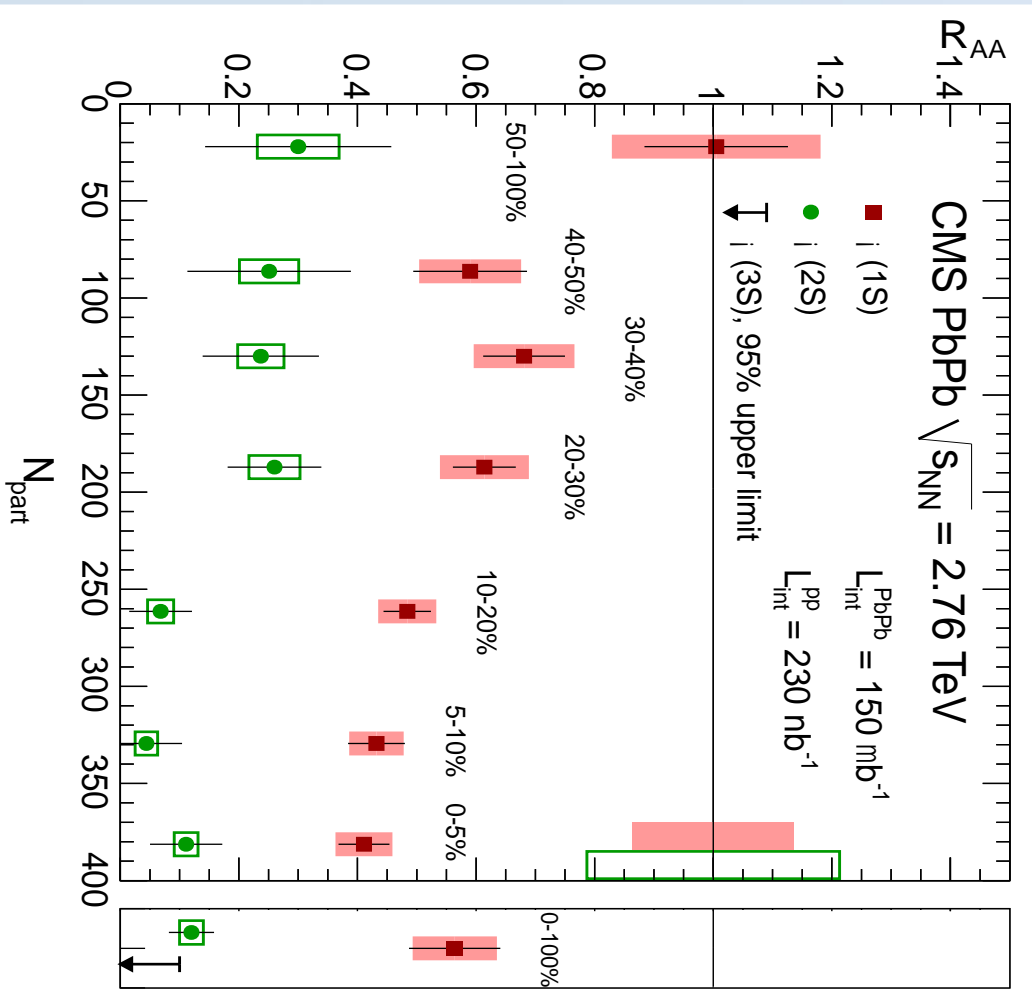
- Centrality integrated

$$R_{AA}(\Upsilon(1S)) = 0.56 \pm 0.08 \text{ (stat.)} \pm 0.07 \text{ (syst.)}$$

$$R_{AA}(\Upsilon(2S)) = 0.12 \pm 0.04 \text{ (stat.)} \pm 0.02 \text{ (syst.)}$$

$$R_{AA}(\Upsilon(3S)) < 0.1 \text{ (at 95\% C.L.)}$$

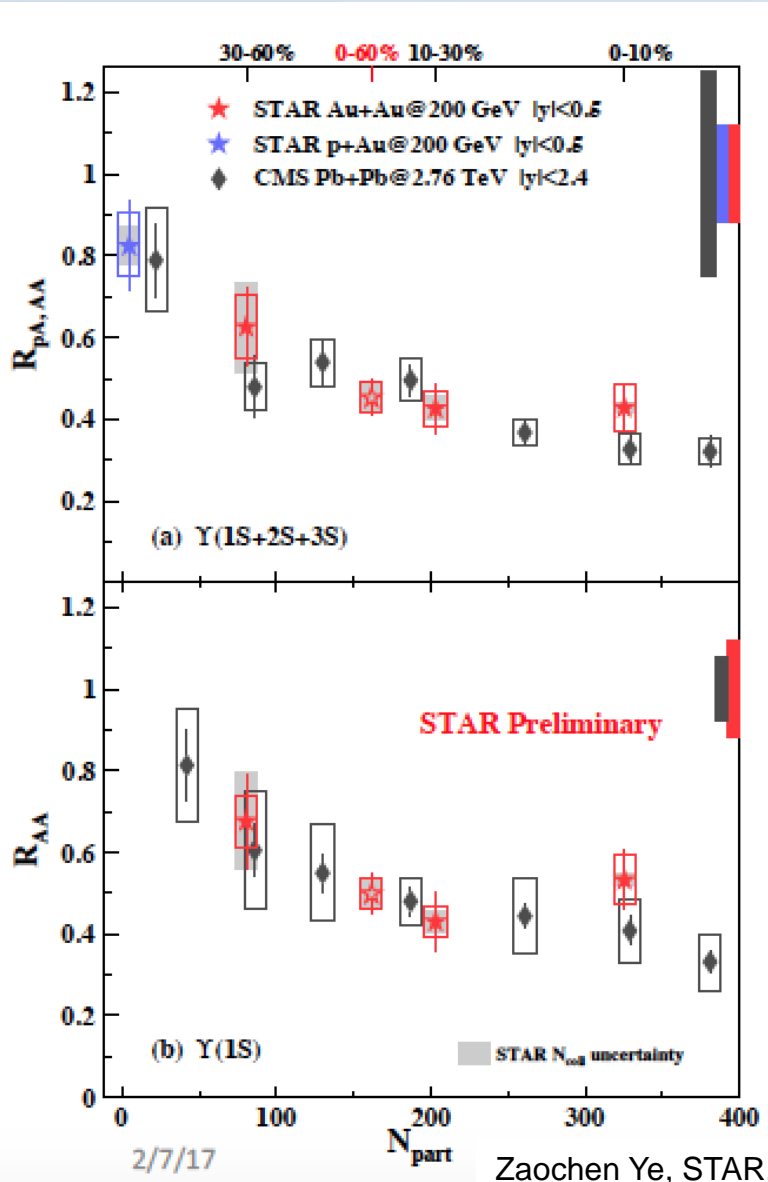
- **Observation of sequential suppression of Υ in order of their binding energy**



PRL 109, 222301 (2012)



RHIC and LHC midrapidity comparison



- 1S+2S+3S: via dielectrons and dimuons.

- R_{pA} : slightly smaller than 1
- Centrality dependence: more suppression in central events
- R_{AA} in central events significantly smaller than R_{pA}

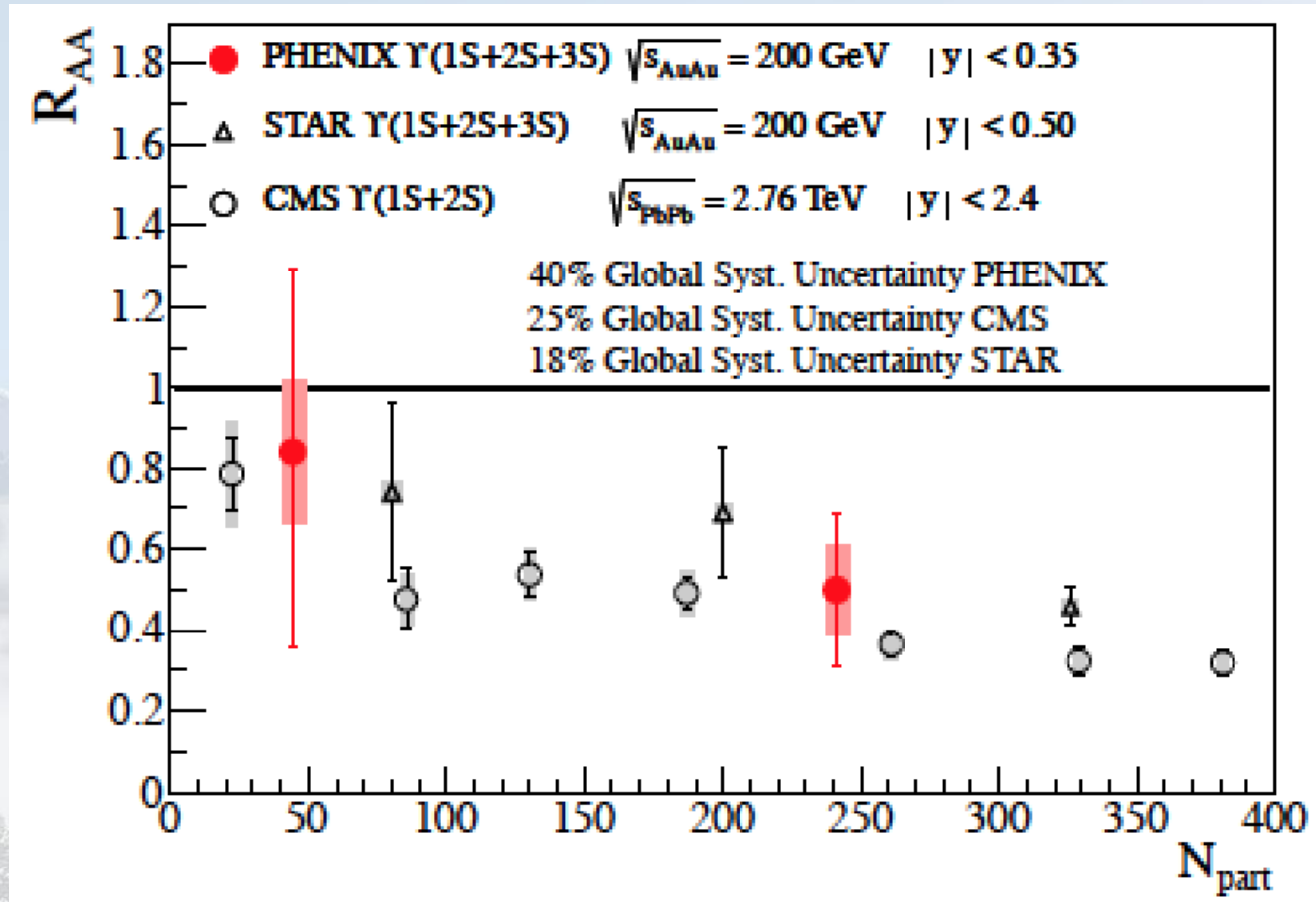
- Separation of 1S (bottom panel)

- Centrality dependence RHIC vs LHC:

- Similar suppression of $Y(1S)$!
- Feeddown dominated?



PHENIX: Υ results

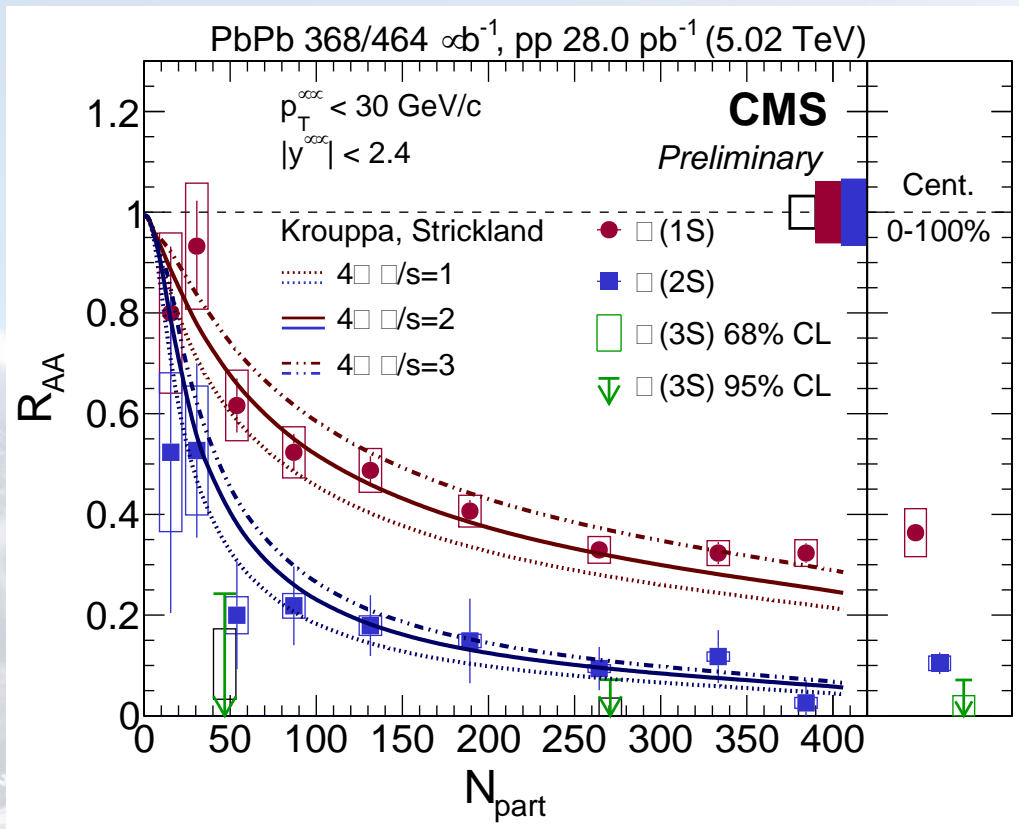


- Suppression in central events of similar magnitude as STAR and CMS.



$\Upsilon(nS) R_{AA}$ vs. Centrality at 5.02 TeV

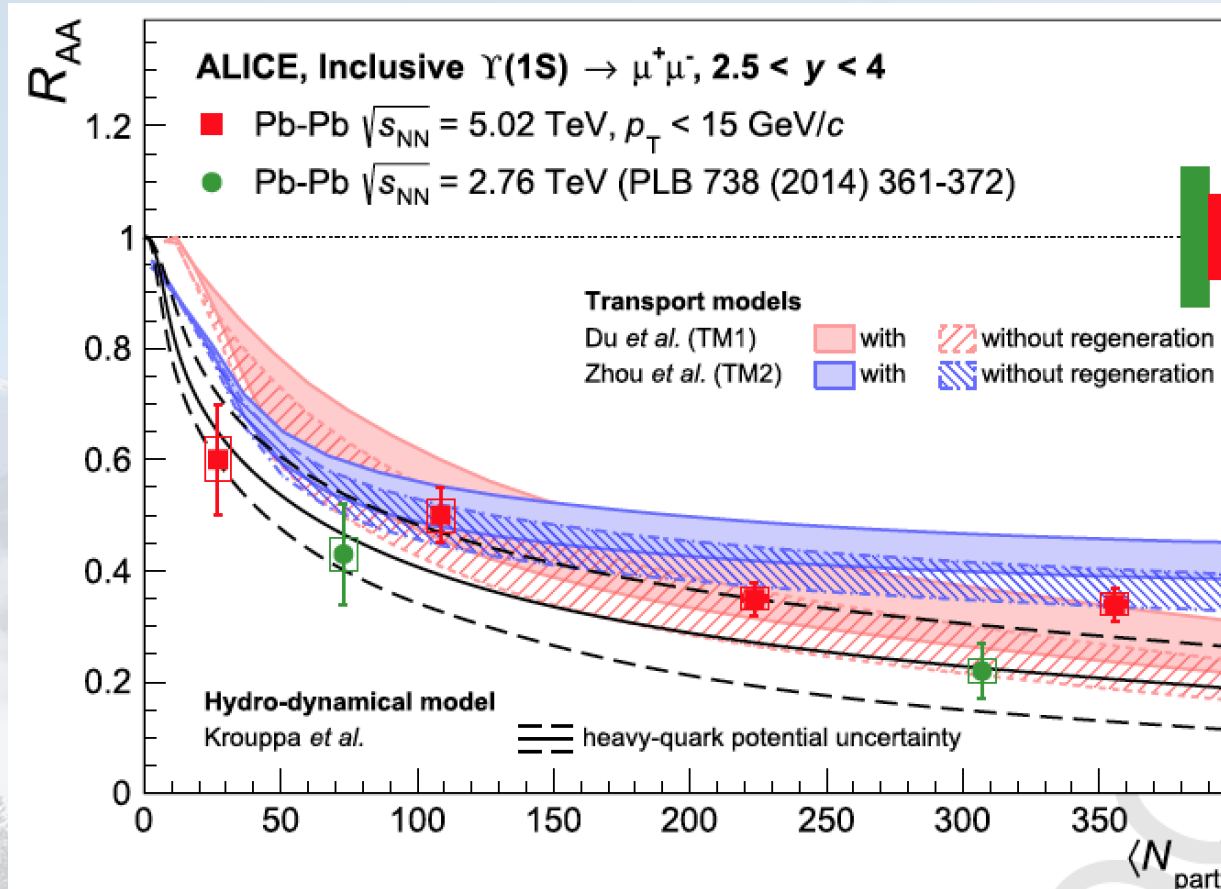
- Nuclear modification for the 3 S-states:
 - Sequential melting!
- Suppression of 1S and 2S:
 - Increasing for more central events
- R_{AA} Integrated results (0-100%)
 - 1S: $0.364 \pm 0.014 \pm 0.048$
 - 2S: $0.104 \pm 0.021 \pm 0.014$
 - 3S: 0.071 at 95% CL



[CMS PAS HIN-16-023](#)
arXiv:1805.09215



$\Upsilon(1S)$ R_{AA} in ALICE

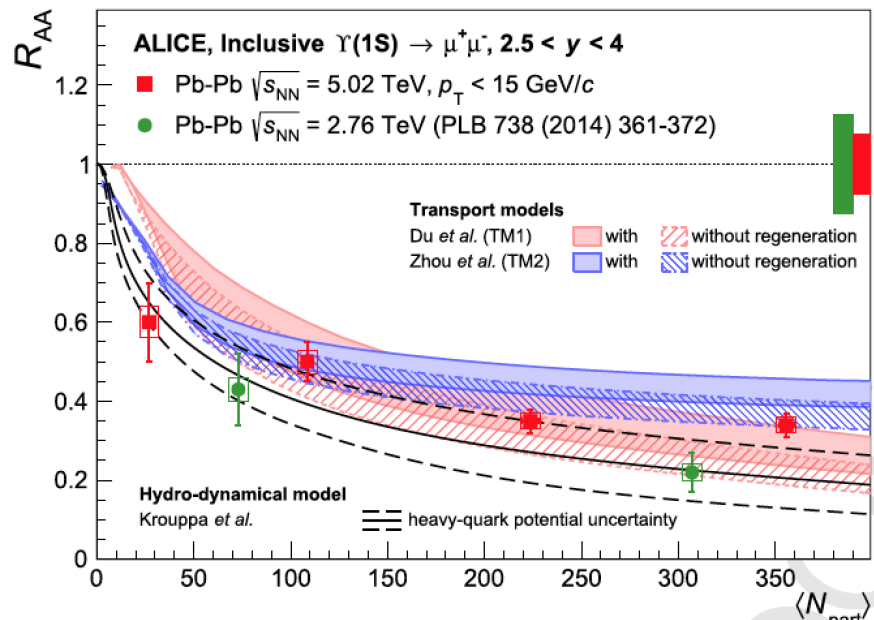
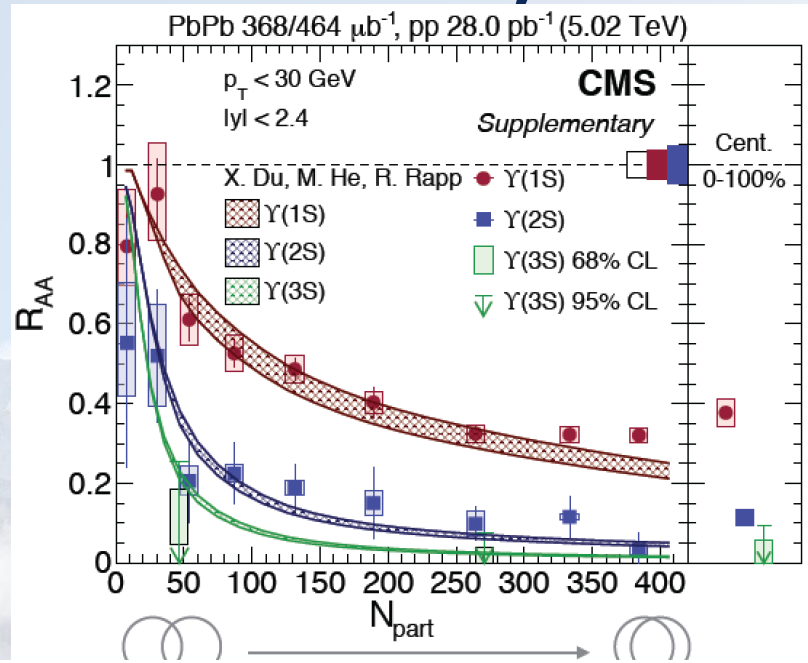


- Suppression increases for central events in ALICE kinematic region.
- Same trends seen in 2.76 and 5 TeV.
 - No energy difference within the $\sim 20\%$ uncertainties.



Comparison to Model, TAMU

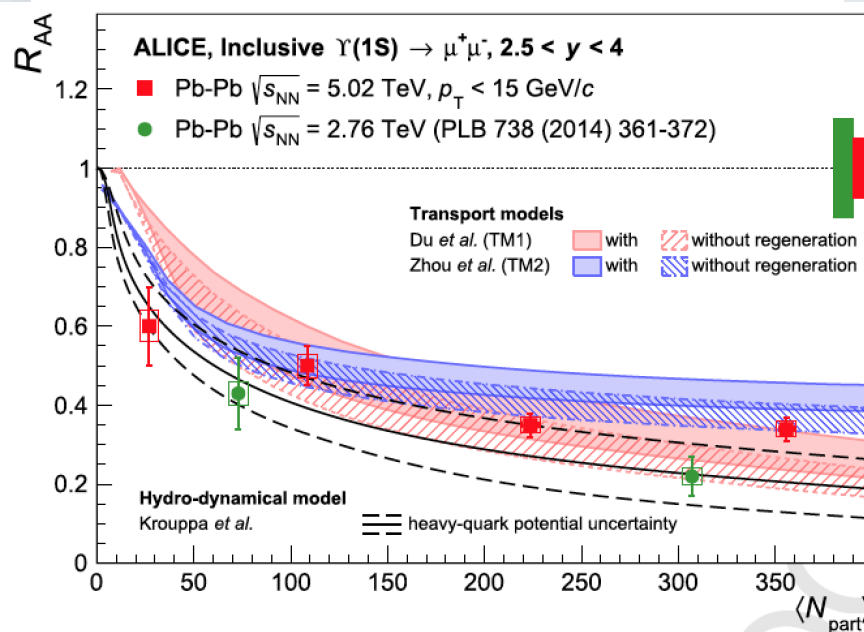
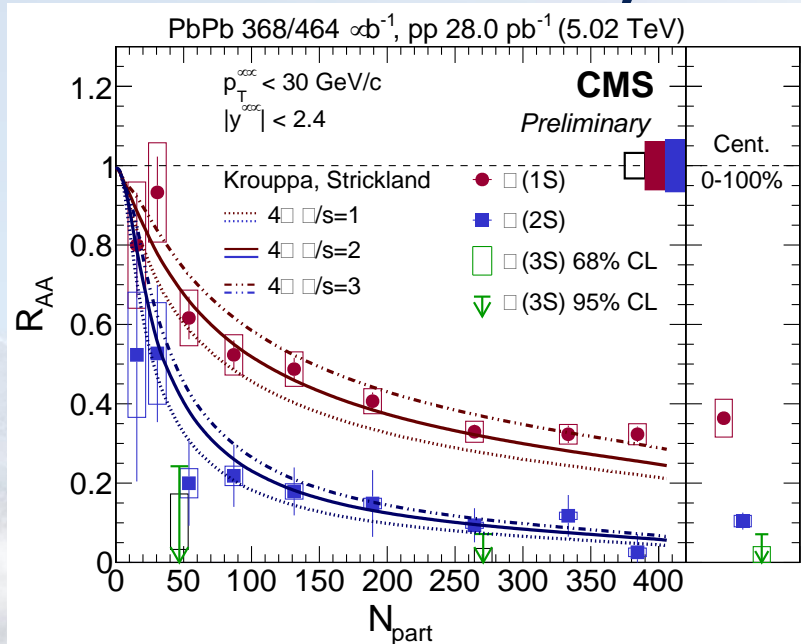
- Transport Model by R. Rapp
 - PRC 96 (2017) 054901
 - $Re[V]$ and $Im[V]$
 - In-medium binding energies
 - Internal energy potentials
 - Feed down from excited states
 - Directly produced $\Upsilon(1S) \sim 67\%$
 - Model QGP via Kinetic rate equation
 - Regeneration
- Melting temperatures:
 - $\Upsilon(1S, 2S, 3S)$: 500, 240, 190 MeV
- Initial temperature:
 - 2.76 TeV: 520 - 750 MeV
 - 5.02 TeV: 7% increase
 - $\Upsilon(1S)$ RAA: only slight decrease
 - Direct $\Upsilon(1S)$ melts!





Comparison to Model, KSU

- Hydro Model by M. Strickland
 - Universe 2 (2016) 16
 - $Re[V]$ and $Im[V]$
 - In-medium binding energies
 - Internal energy potentials
 - Feed down from excited states
 - Directly produced $\Upsilon(1S)$ ~67%
 - Model QGP via Hydrodynamics
 - Momentum-space anisotropy
 - No regeneration
- Melting temperatures:
 - $\Upsilon(1S, 2S, 3S)$: 600, 230, 170 MeV
- Initial temperature:
 - 2.76 TeV: 544 - 552 MeV
 - 5.02 TeV: 629 - 641 MeV (16% increase)
 - $\Upsilon(1S)$ RAA: ~25% decrease
 - Direct $\Upsilon(1S)$ melts!



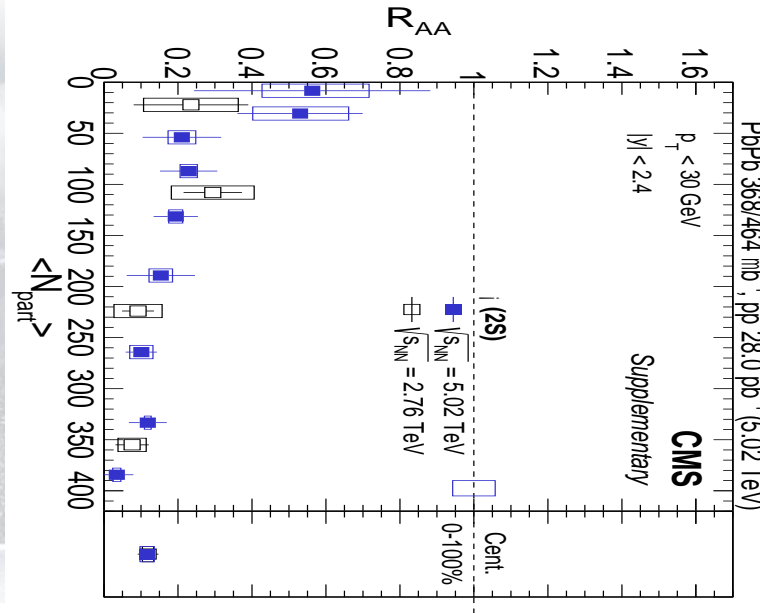
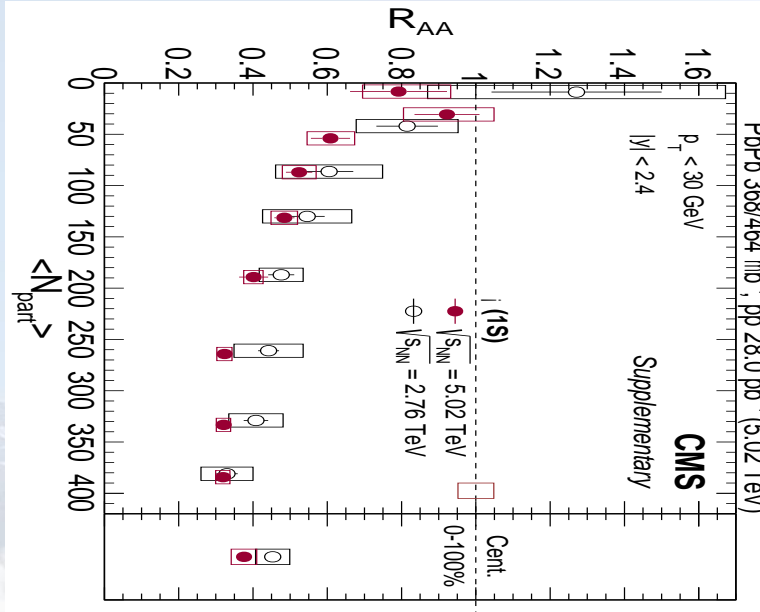


Comparison to 2.76 TeV Results

- $\Upsilon(1S)$:
 - Models expect increased suppression with larger \sqrt{s}
 - Larger suppression at higher energy? Can't really say:
 - Factor of 1.2 ± 0.15 (syst).
 - Need to reduce uncertainties below 10%.

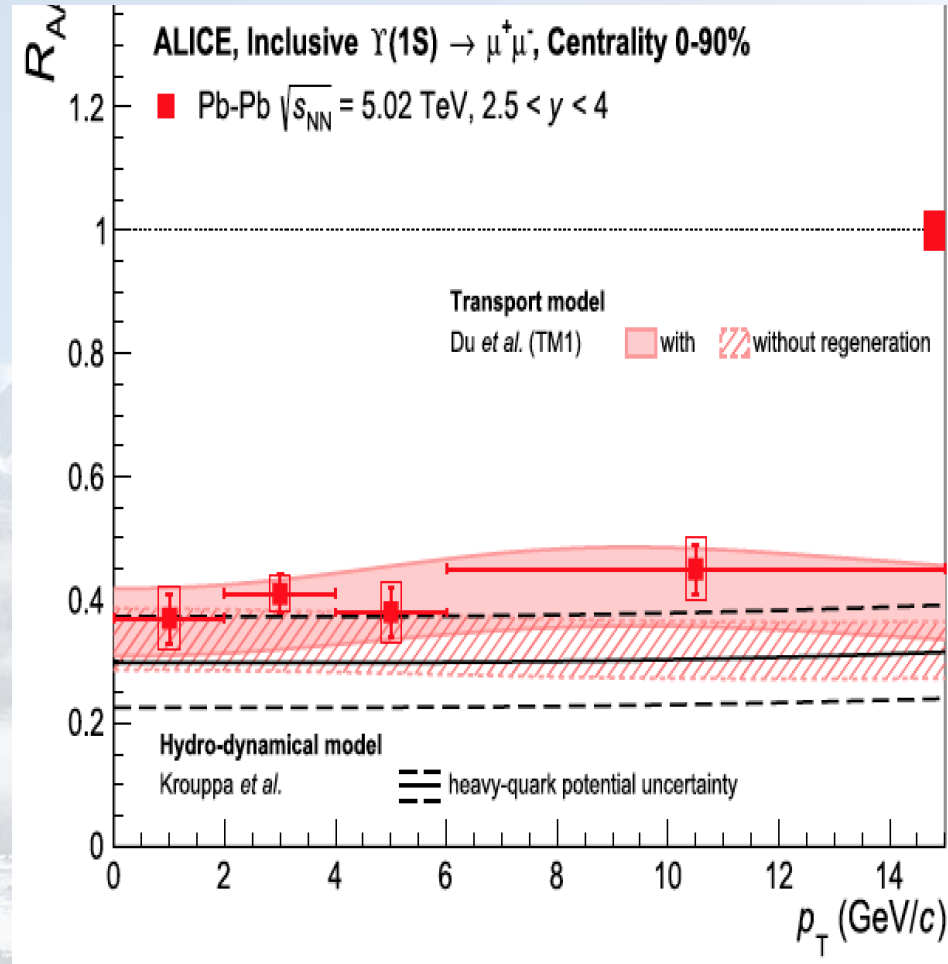
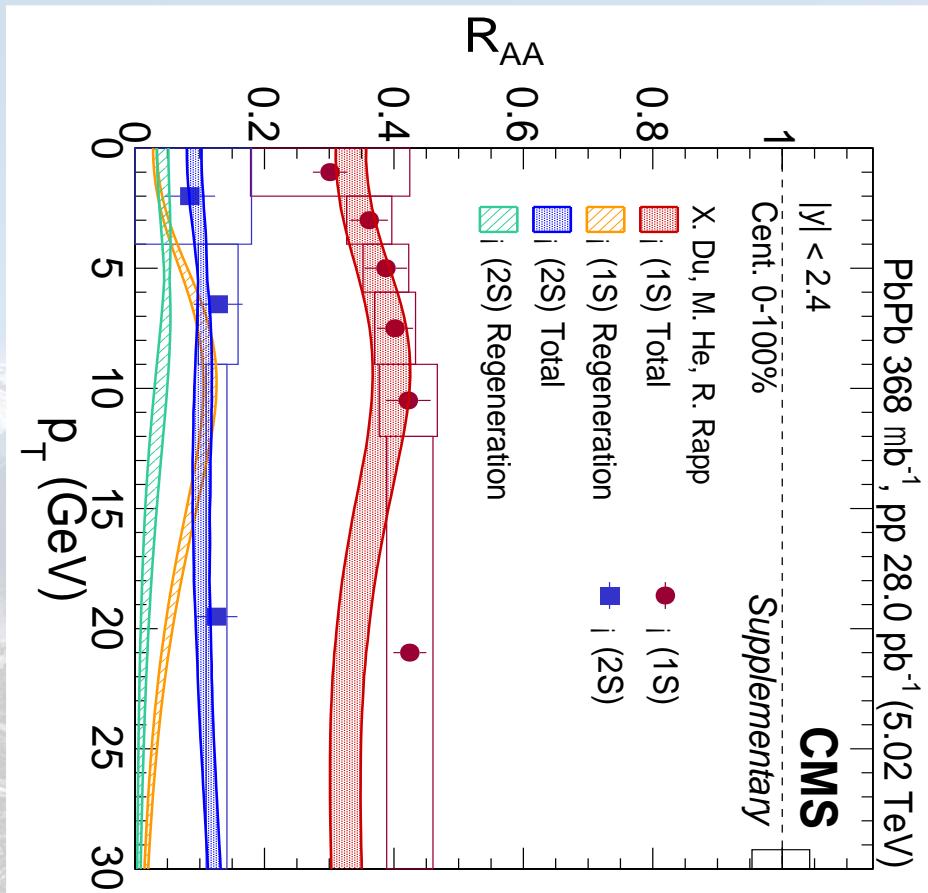
- $\Upsilon(2S)$:
 - Suppression level is consistent between the two energies
 - Larger role of regeneration for 2S state?
 - Rapp et al. arXiv:1706.08670

CMS 5.02 TeV R_{AA} : [CMS PAS HIN-16-023](#) arXiv:1805.09215
 CMS 2.76 TeV R_{AA} : [PLB 770 \(2017\) 357](#)





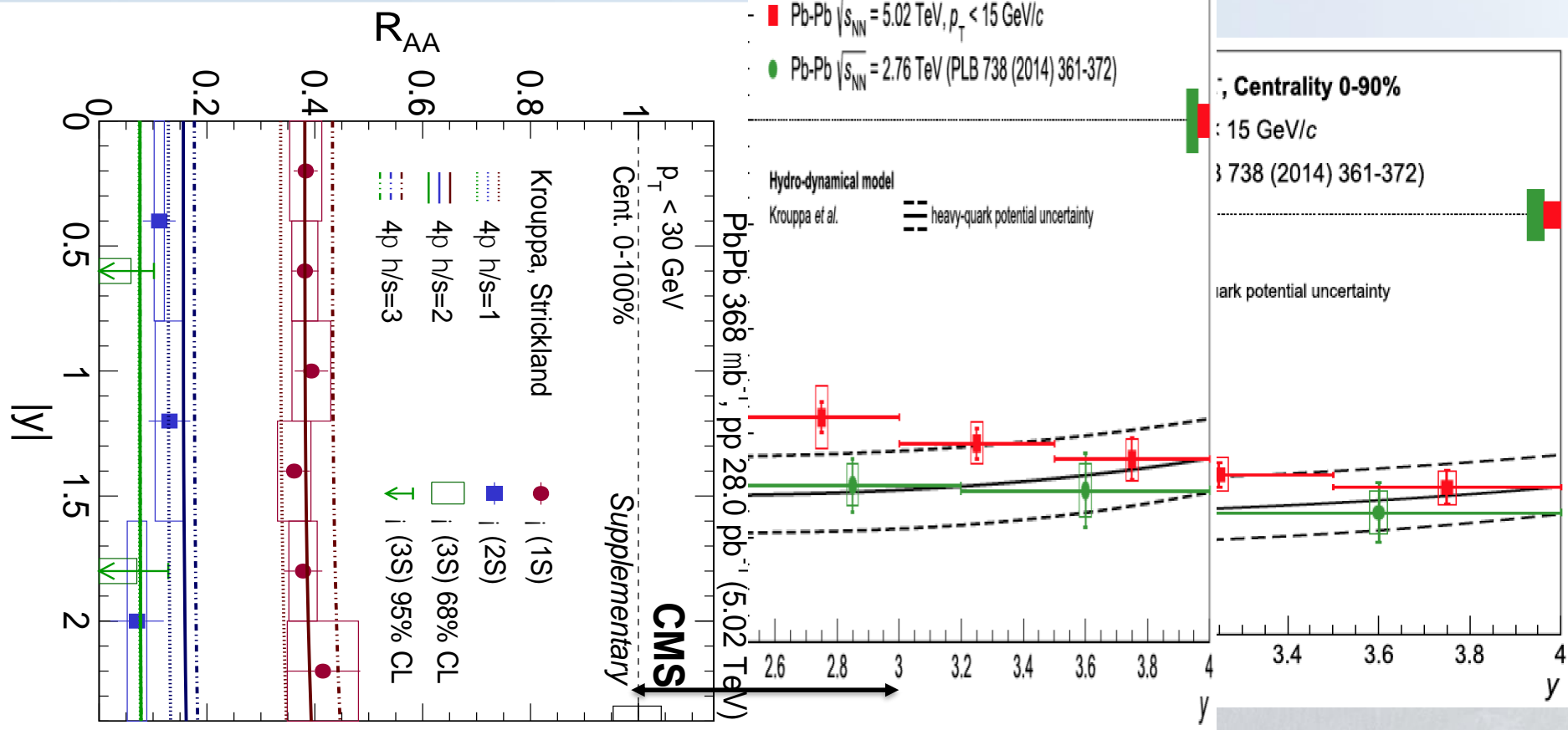
Υ R_{AA} p_T dependence



- $\Upsilon(nS)$: No significant dependence of R_{AA} on p_T ,
- p_T , Rapp Model: Contributions from regeneration have different shapes.
 - Not visible within our uncertainties.



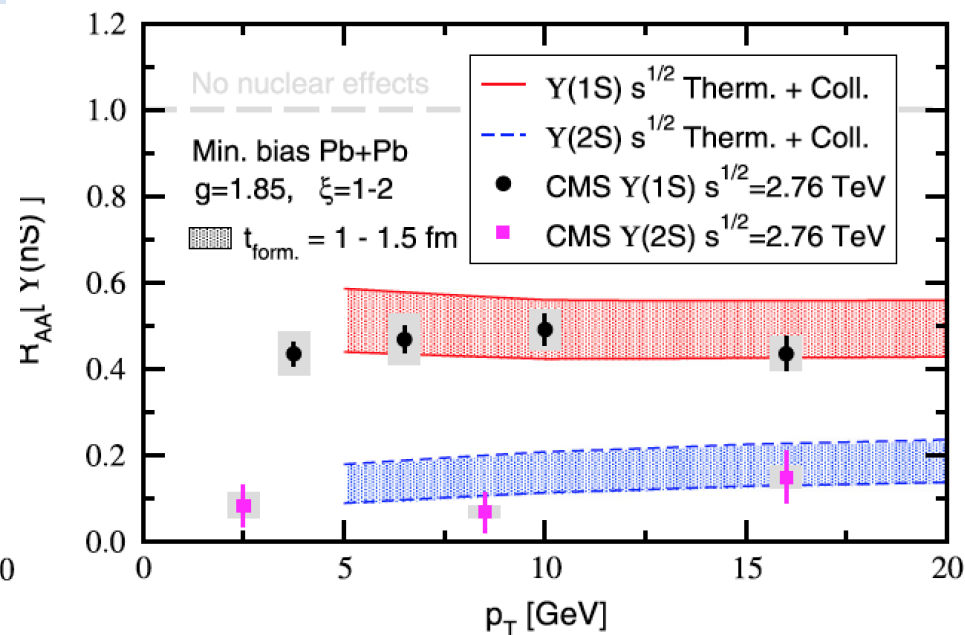
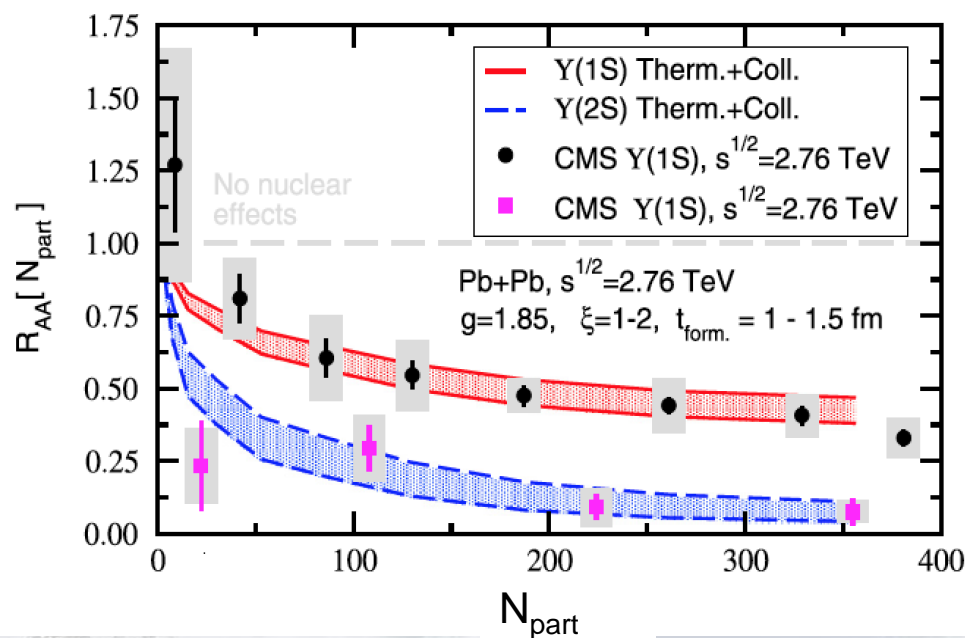
Υ R_{AA} y dependence



- $\Upsilon(nS)$: No significant dependence of R_{AA} on rapidity.
- y , Strickland Model: Modest increase of R_{AA} at higher rapidity. Consistent with data.



Υ Collisional + Thermal dissociation model



Aronson, Borrás, Odegard, Sharma, Vitev,
PLB 778 (2018) 384

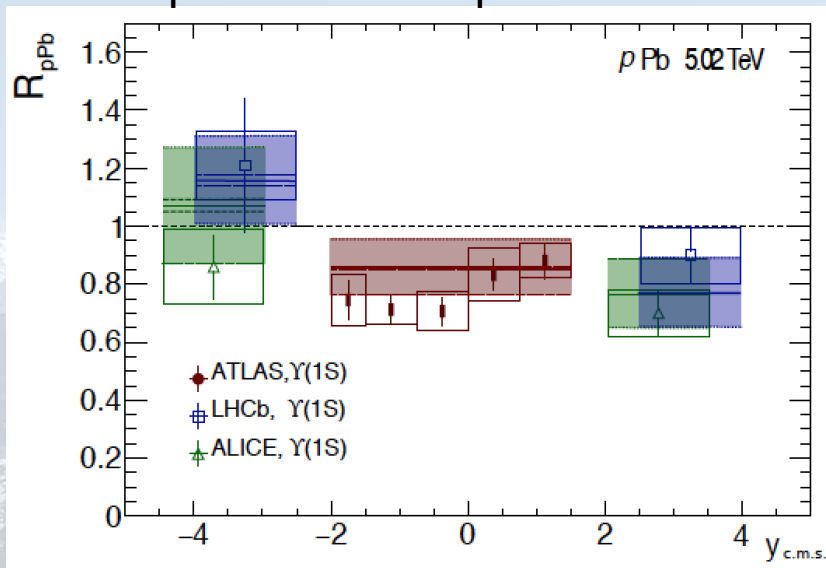
- Model:

- $Q\bar{Q}$ proto-quarkonia state with formation time ~ 1 fm interacts via Debye screened potential.
 - Thermal screening dominates the suppression of excited states.
- Broadening due to multiple scattering in QGP, controlled by ξ parameter.
 - Ground state is also dissociated due to collisional interactions.
- Reproduce data at 2.76 TeV

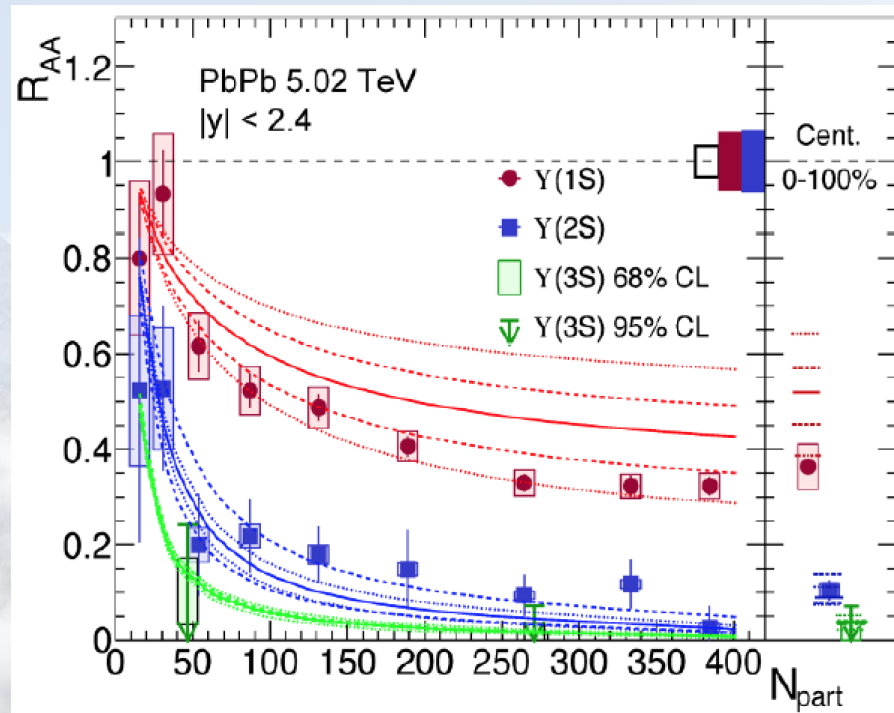


Υ in Co-mover Interaction Model

Fit dissociation cross section parameters in pA



Extrapolate to AA using same parameters

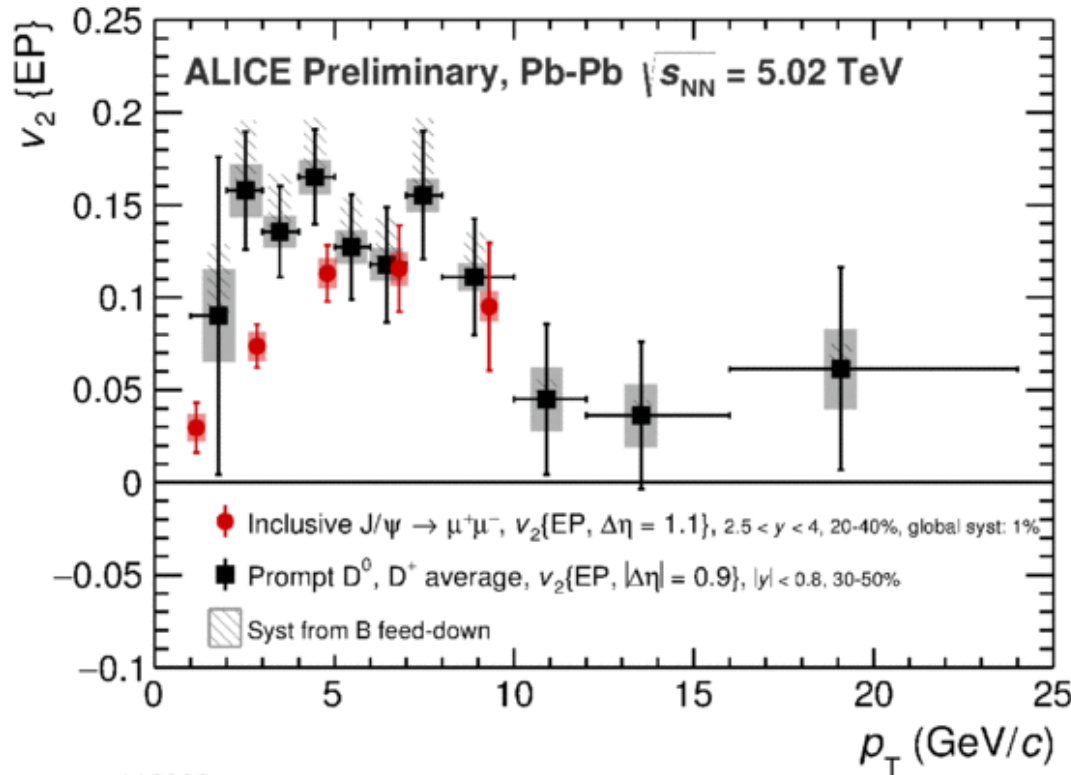


- Same physics in pA and AA?
- Hadrons in pA vs gluons in AA? Same parameter for scattering strength in pp and AA...Coincidence?
- Do we agree these are gluons in all cases? QGP drop in pA? Is this a crude way to deal with gluodissociation?
- Note uncertainties from PDF: with reduced uncertainties, will this still work?

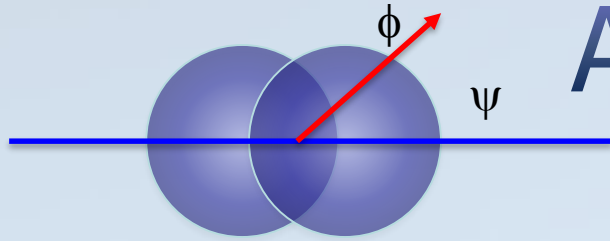


Next Υ measurements...

- Do heavy quarks flow?

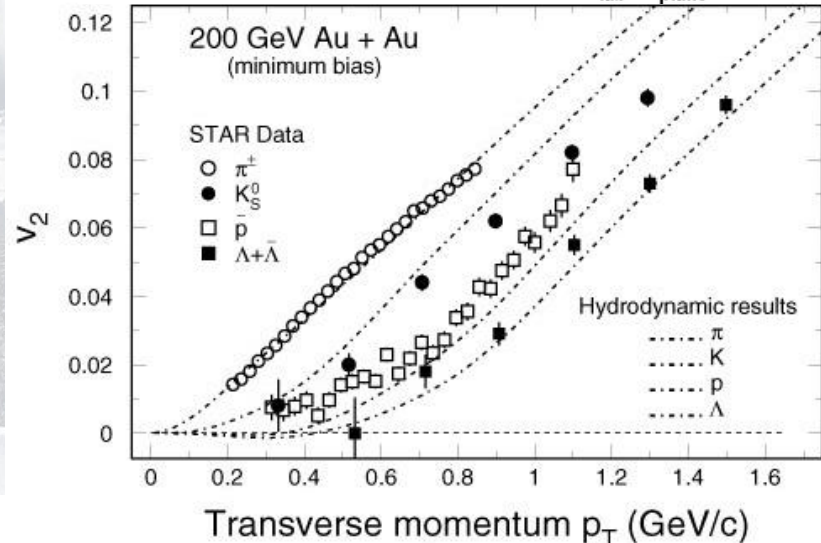
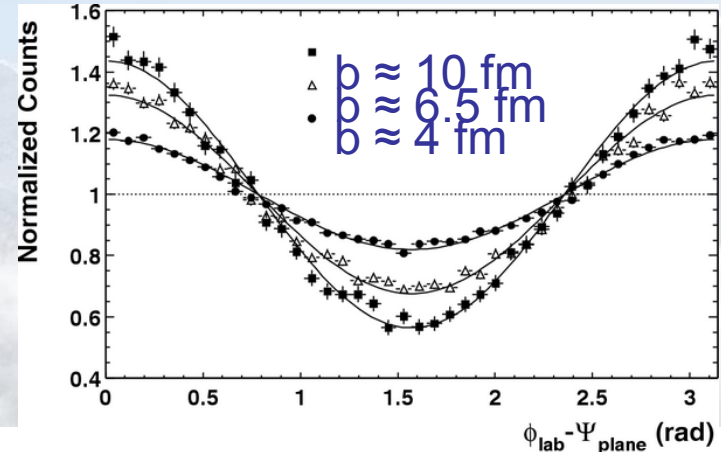
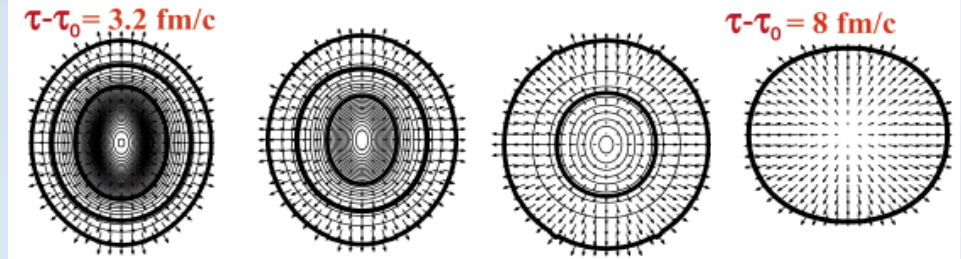


- J/ψ v_2 is large, and puzzling. Should heavy quarks flow? What will Υ do?



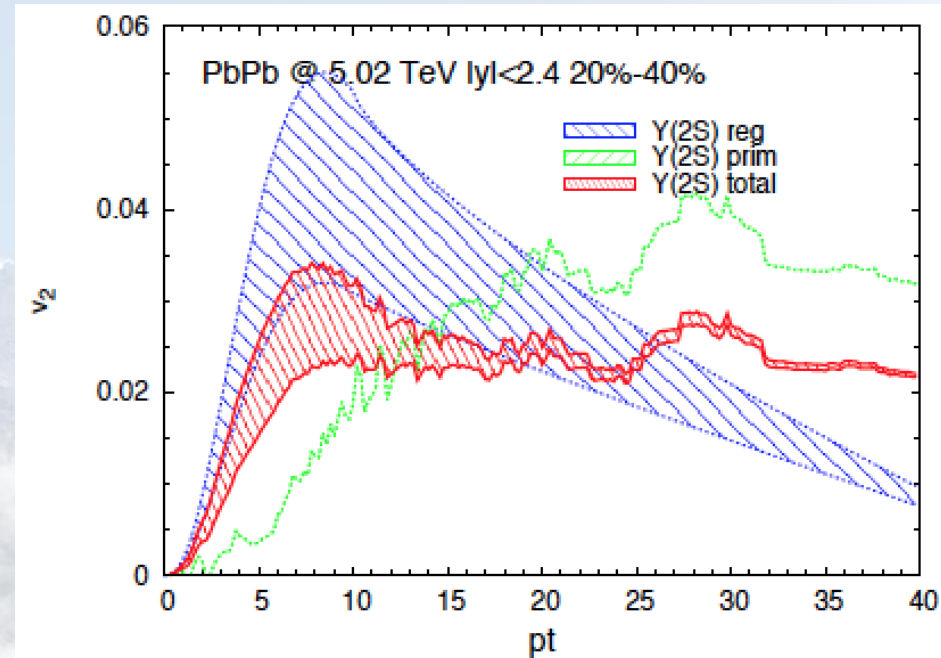
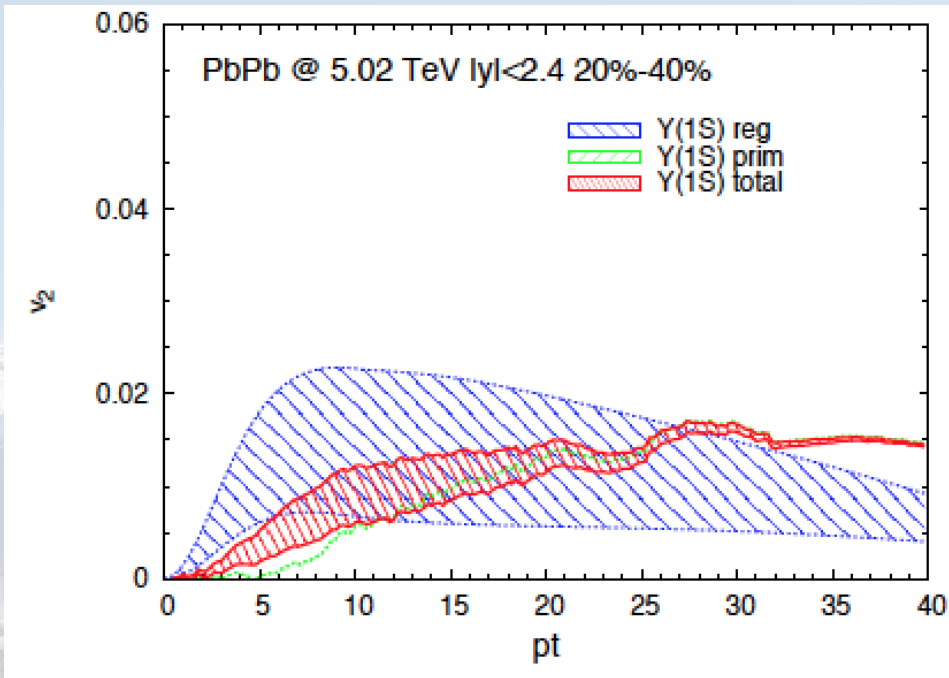
Azimuthal anisotropy

- Non-central collisions produce elliptical shape in transverse plane.
- Azimuthal distribution of particles around the reaction plane:
 - driven by large anisotropic pressure gradients in relativistic QGP fluid.
 - in-plane vs out-of plane difference
 - Characterized by Fourier series. 2nd coefficient: $v_2 \sim \langle 2\cos(2(\phi-\psi)) \rangle$ called “elliptic flow”.
- Low pt azimuthal anisotropy:
- Characteristic mass ordering vs p_T in Hydro
 - Same velocity, but different mass.





Υ v_2 predictions, Transport Model

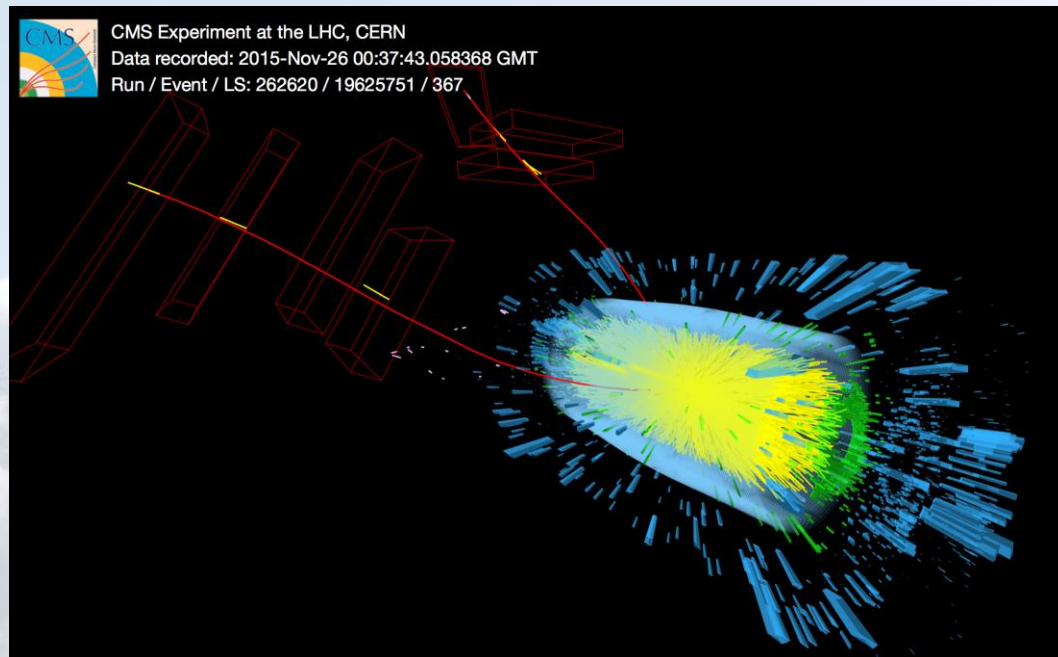
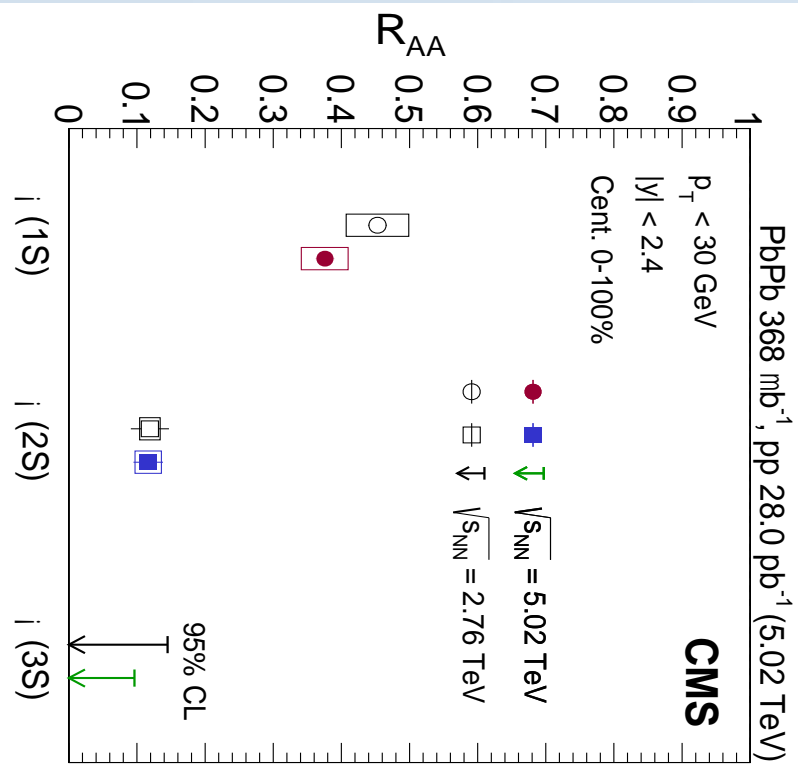


Du, He, Rapp, PRC96 (2017) 054901

- Expect small ($< 2\%$) v_2 for 1S (with reg). Otherwise ~ 0 up to 10 GeV/c.
 - Also expected small v_2 for J/ψ ... didn't happen. Attempt Υ v_2 with 2018 data.
- Test of regeneration: $v_2(2S) \sim 2 v_2(1S)$
 - Will need higher luminosity, more data for $v_2(2S)$.
- Toward correlating Υ with activity in azimuth in AA.



Summary



Bottomonia: essential to quantitatively explore the temperature and color deconfinement properties of the Quark Gluon Plasma.

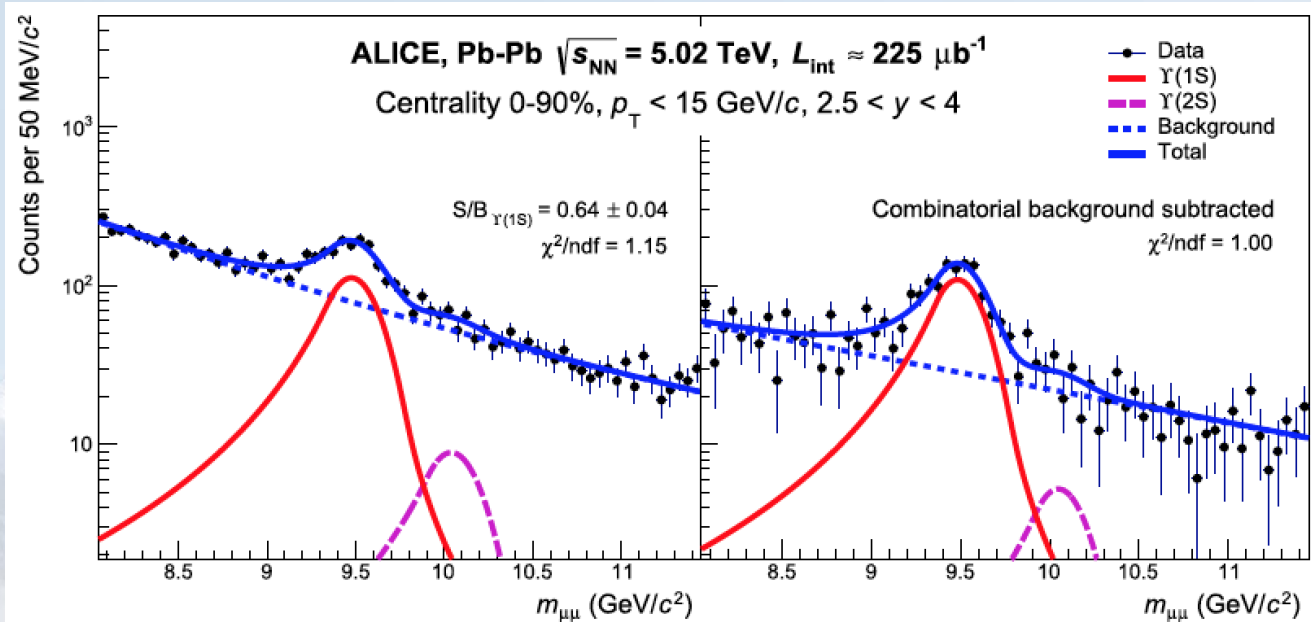
1S: Direct ~67%, RAA ~0.4, how much suppression from CNM?

2S: Dominated by regeneration? Not all models include it.

3S: Fully melted?!

Open questions: Feed-down, suppression mechanisms, energy dependence, pA to AA...

Backup Material

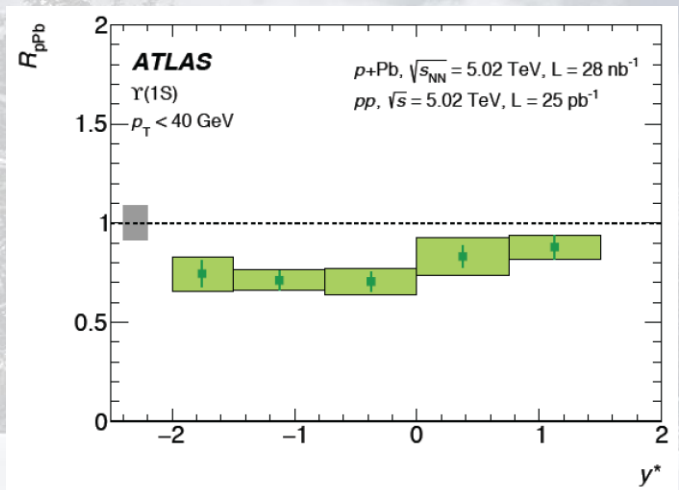
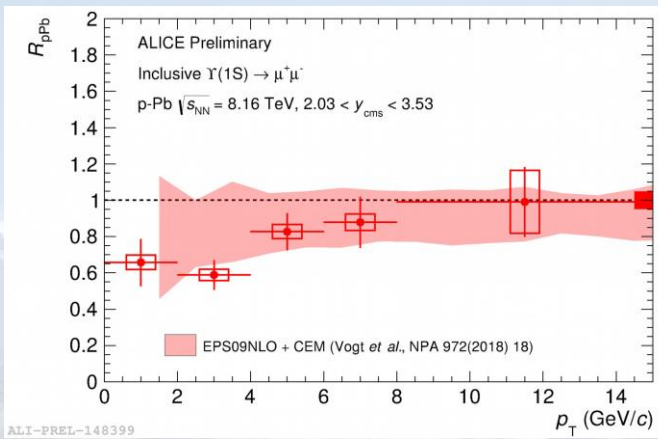
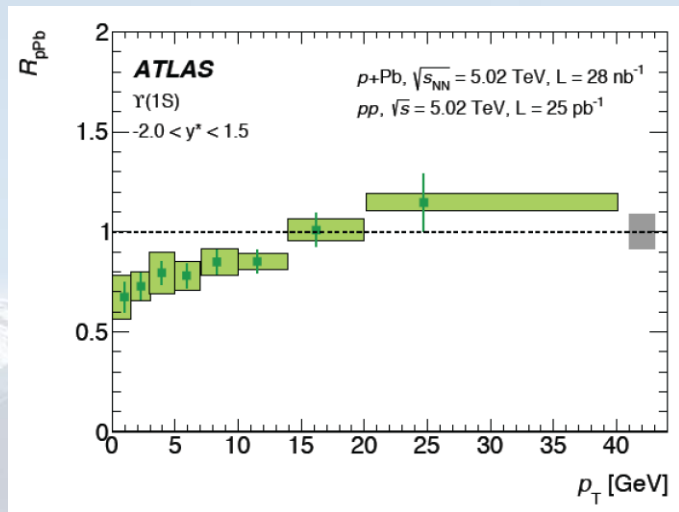


- 1S and 2S observed. 3S consistent with zero.



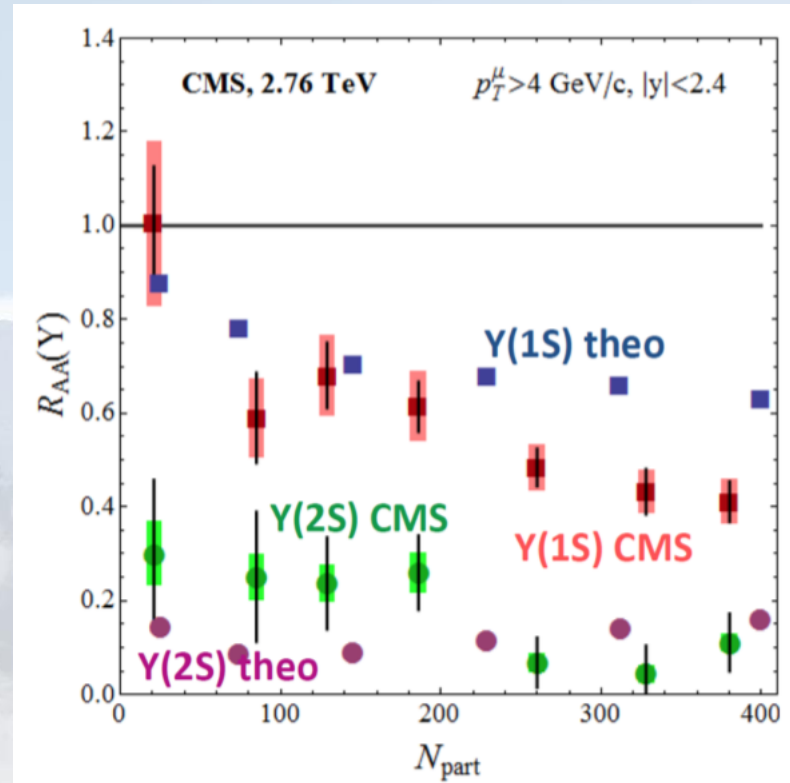
LHC pPb Results

- ATLAS, ALICE, LHCb





Schrödinger-Langevin Approach

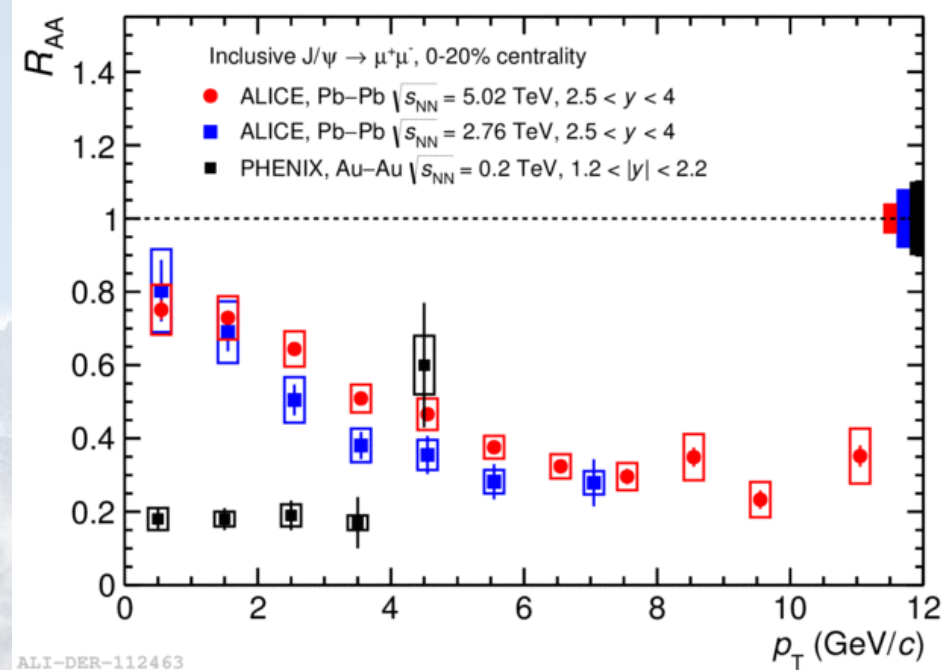
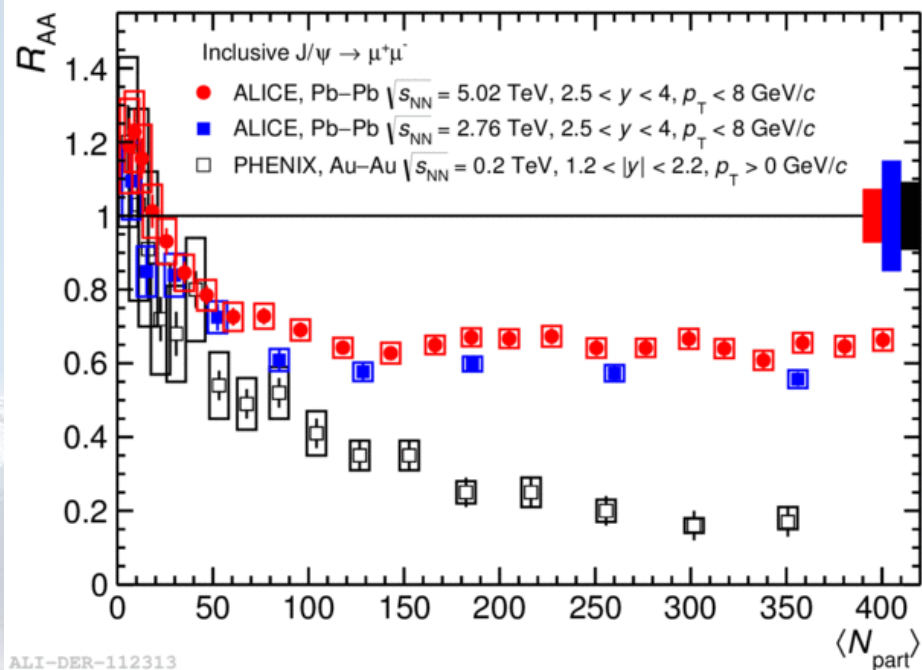


Gossiaux and Katz
[arXiv:1611.06499](https://arxiv.org/abs/1611.06499)

- Central collisions: 1S states survive more in Model.
 - P \rightarrow S transitions, not fully compensated by suppression of feed-down from higher excited states.
 - Need to include CNM effects in model.



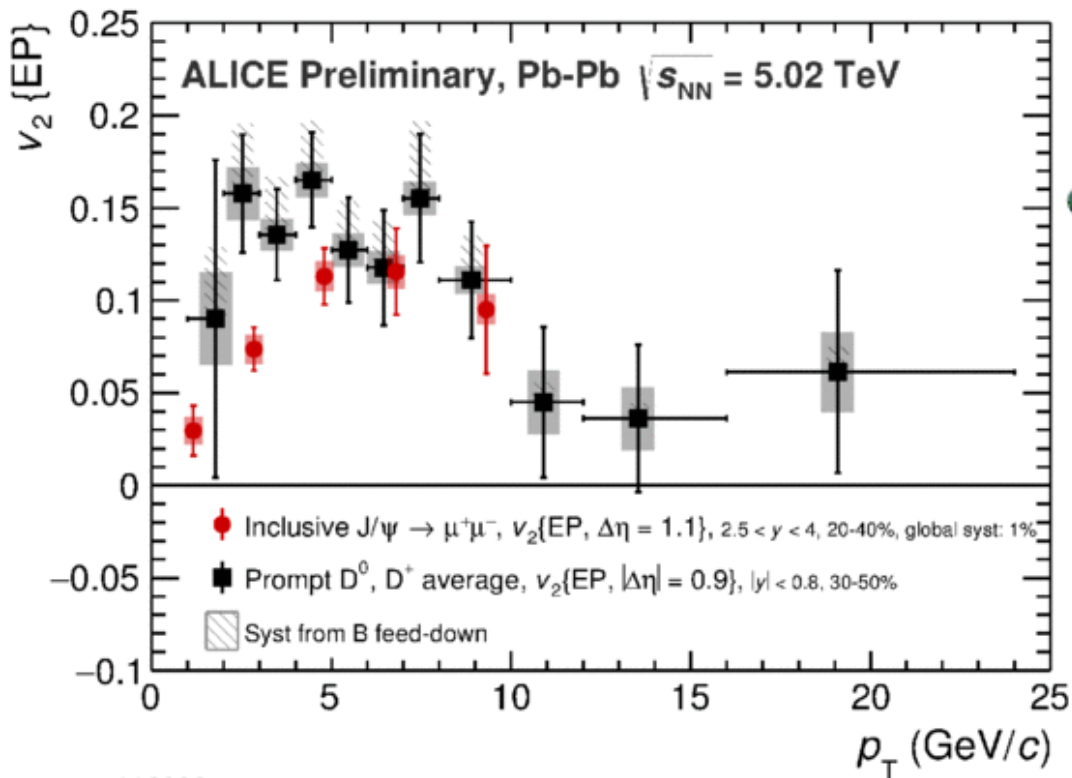
J/ψ Forward, ALICE & PHENIX



- p_T -integrated R_{AA} : systematically larger at LHC than at RHIC.
 - 5.02 TeV RAA > than at 2.76 TeV
- Excess seen at LHC from low- p_T J/ψ .
 - Theoretical interpretation: screening + regeneration



J/ψ v₂: recombination signal



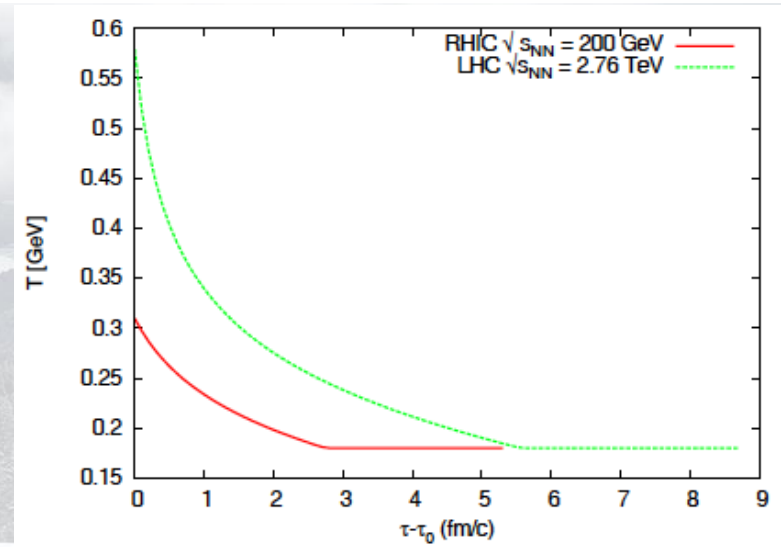
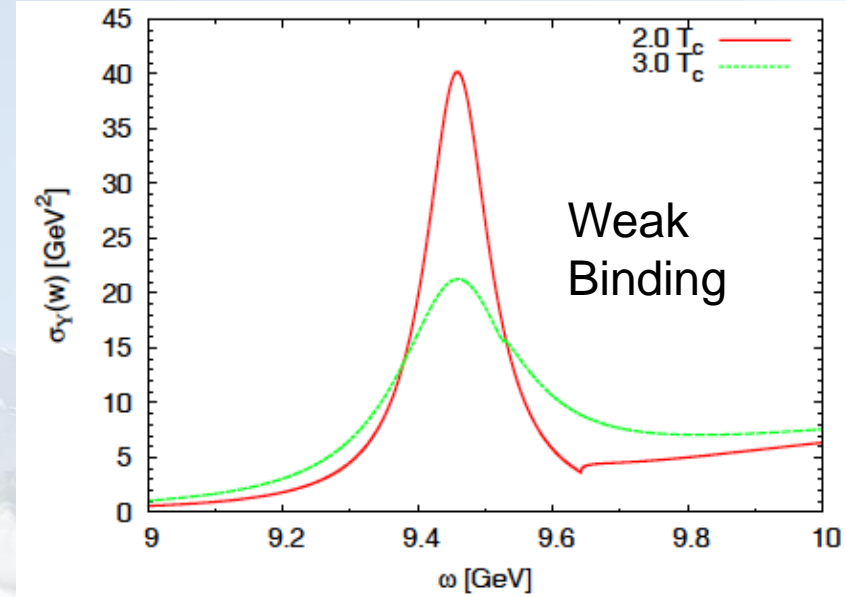
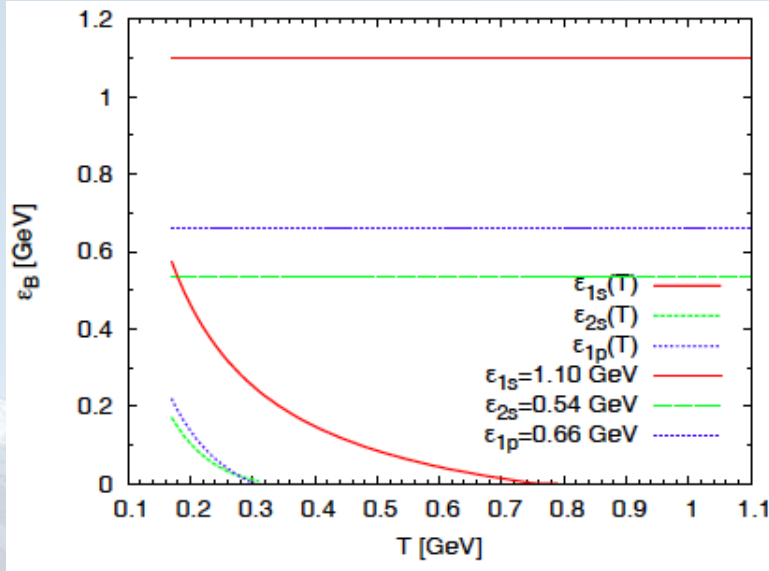
- Centrality 20-40%
- Non-zero v₂ observed!!
- Effect is 5-7σ for some p_T bins.

ALI-PREL-119009

- Supports the idea that charm quarks thermalize in the QGP
 - If so, expect D meson v₂ to be similar, but maybe a little higher.



Υ R_{AA} Model II



- Weak vs. Strong Binding
 - Narrower spectral functions for “Strong” case
 - Ratios of correlators compared to Lattice: favor “Strong” binding case
- Kinetic Theory Model
 - Rate Equation: dissociation + regeneration
 - Fireball model: T evolution. $T_0 \sim 300$ MeV